



50 Years of the Veneziano Model



**Heterotic- and Bosonic-String
Amplitudes from Field Theories**

Oliver Schlotterer (Perimeter Institute & AEI Potsdam)

based on arXiv:1106.2645 with C. Mafra and S. Stieberger

arXiv:1803.05452 with T. Azevedo, M. Chiodaroli and H. Johansson

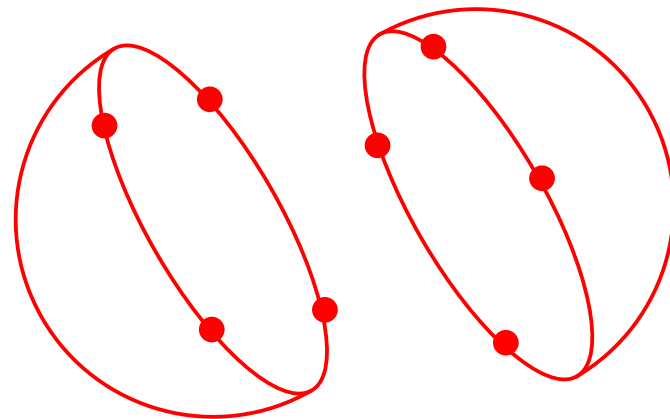
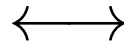
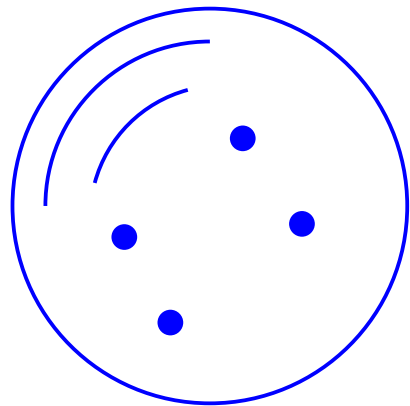
15.05.2018

Introduction: Double copy – from string to field theory

Birth of double copy: **KLT relations** among string amplitudes at tree-level

$$\mathcal{M}_{\text{closed}}^{4\text{pt}}(\alpha') = \bar{\mathcal{A}}_{\text{open}}(1, 2, 4, 3; \alpha') \sin\left(\frac{\pi\alpha'}{2} k_1 \cdot k_2\right) \mathcal{A}_{\text{open}}(1, 2, 3, 4; \alpha') .$$

closed
strings:
sphere



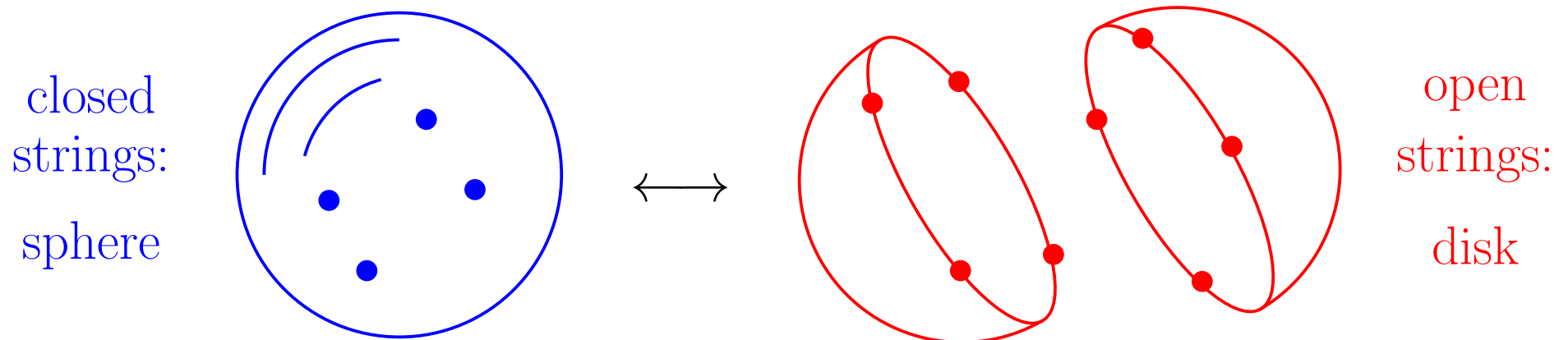
open
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disk

[Kawai, Lewellen, Tye 1986; see Zvi Bern's talk]

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Field-theory limit $\alpha' \rightarrow 0$: relate **gravity** to **double copy of gauge theories**:

$$M_{\text{SUGRA}}^{4\text{pt}} = \bar{A}_{\text{SYM}}(1, 2, 4, 3) k_1 \cdot k_2 A_{\text{SYM}}(1, 2, 3, 4) \equiv \bar{A}_{\text{SYM}} \otimes_{\text{KLT}} A_{\text{SYM}} .$$

Will refer to operation \otimes_{KLT} at $\alpha' \rightarrow 0$ as **field-theory double copy**.


Introduction: Double copy – from string to field theory

At n points, more combinatorics and $(n-3)!$ -element BCJ bases

$$\{A_{\text{SYM}}(1, \rho(2, 3, \dots, n-2), n-1, n), \quad \text{permutation } \rho \in S_{n-3}\}$$

[Bern, Carrasco, Johansson 0805.3993; see Zvi Bern's talk]

$$M_{\text{SUGRA}}^{n \text{ pt}} = \sum_{\rho, \tau \in S_{n-3}} \bar{A}_{\text{SYM}}(1, \rho, n, n-1) S(\rho|\tau)_1 A_{\text{SYM}}(1, \tau, n-1, n)$$



 $(n-3)! \times (n-3)! \text{ KLT matrix, entries are } \sim (k_i \cdot k_j)^{n-3}$

[Bern, Dixon, Perelstein, Rozowsky 1998]

[Bjerrum-Bohr, Damgaard, Feng, Sondergaard, Vanhove 2010]

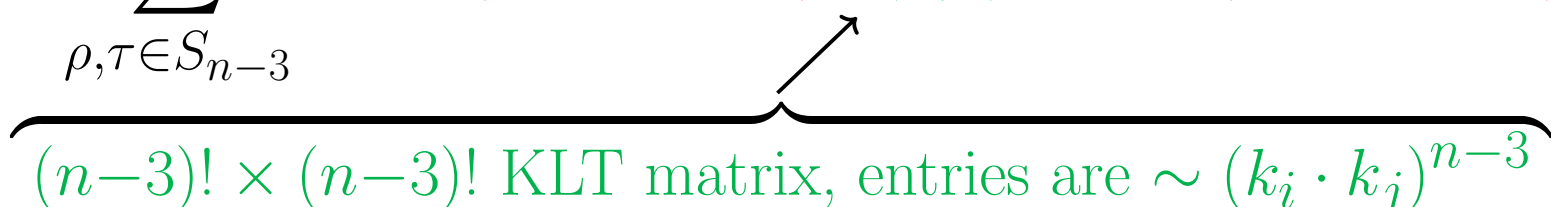
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e.g. 2×2 terms

at 5 points with

$$S(\rho(2, 3)|\tau(2, 3))_1 = \begin{pmatrix} (k_1 \cdot k_2)(k_{1+2} \cdot k_3) & (k_1 \cdot k_2)(k_1 \cdot k_3) \\ (k_1 \cdot k_2)(k_1 \cdot k_3) & (k_1 \cdot k_3)(k_{1+3} \cdot k_2) \end{pmatrix}$$

Shorthand for KLT formulae:

$$M_{\text{SUGRA}} = \bar{A}_{\text{SYM}} \otimes_{\text{KLT}} A_{\text{SYM}}$$

Key result: Massless tree amplitudes in various string theories

Field-theory double copy $\otimes_{\text{KLT}} \Rightarrow$ web of relations for string amplitudes

\rightarrow representations of the flavour (field theory) \otimes_{KLT} (stringy building block)

\otimes_{KLT}	SYM		
Z-theory	open superstring		

- “Z-theory” \leftrightarrow α' -dependent disk integrals (over moduli space $\mathcal{M}_{0,n}$)

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sv (open superstring)	closed superstring	heterotic string (grav)	
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Outline

I. Open superstrings as a field-theory double copy

[Mafra, OS, Stieberger 1106.2645; Broedel, OS, Stieberger 1304.7267]

II. Closed superstrings from single-valued open superstrings

[OS, Stieberger 1205.1516; Stieberger 1310.3259]

III. Bosonic strings from $(DF)^2 + \text{YM}$ field theory

[Azevedo, Chiodaroli, Johansson, OS 1803.05452]

IV. Heterotic strings and $(DF)^2 + \text{YM} + \phi^3$ field theory

[Azevedo, Chiodaroli, Johansson, OS 1803.05452]

V. Conclusions & Outlook

I. Open superstrings as a field-theory double copy

Color ordered n -point trees of open superstring (with $\sigma = \sigma(2, 3, \dots, n-2)$)

$$\mathcal{A}_{\text{super}}^{\text{open}}(1, \sigma, n-1, n; \alpha') = \sum_{\tau \in \mathcal{S}_{n-3}} F_{\sigma}^{\tau}(\alpha') A_{\text{SYM}}(1, \tau, n-1, n)$$

[Mafra, OS, Stieberger 1106.2645, 1106.2646]

- all polarizations in BCJ basis of 10-dim SYM amplitudes $A_{\text{SYM}}(\dots)$

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- all polarizations in BCJ basis of 10-dim SYM amplitudes $A_{\text{SYM}}(\dots)$
- all the α' dependence in $(n-3)! \times (n-3)!$ basis of disk integrals $F_{\sigma}^{\tau}(\alpha')$

$$F_{\sigma}^{\tau}(\alpha') = \int_{0 < z_{\sigma(2)} < z_{\sigma(3)} < \dots < z_{\sigma(n-2)} < 1} dz_2 \dots dz_{n-2} \prod_{i < j}^{n-1} |z_i - z_j|^{2\alpha' k_i \cdot k_j} \tau \left\{ \prod_{l=2}^{n-2} \sum_{m=1}^{l-1} \frac{2\alpha' k_l \cdot k_m}{z_l - z_m} \right\}$$
- $\tau \in S_{n-3}$ acts on k_j and z_j enclosed in $\{\dots\}$ with $j = 2, 3, \dots, n-2$
- recover $\mathcal{A}_{\text{super}}^{\text{open}} \rightarrow A_{\text{SYM}}$ from field-theory limit $F_{\sigma}^{\tau}(\alpha') = \delta_{\sigma}^{\tau} + \mathcal{O}(\alpha'^2)$

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A KLT formula in disguise involving disk integrals (with $z_{ij} \equiv z_i - z_j$)

$$Z_{\sigma}(\rho(1, 2, \dots, n)) \equiv (2\alpha')^{n-3} \int_{z_{\sigma(i)} < z_{\sigma(i+1)}} \frac{dz_1 \dots dz_n}{\text{vol SL}_2(\mathbb{R})} \frac{\prod_{i < j}^n |z_{ij}|^{2\alpha' k_i \cdot k_j}}{\rho(z_{12} z_{23} \dots z_{n-1, n} z_{n, 1})}$$

Permutation $\rho = \rho(1, 2, \dots, n)$ acts on cyclic denominator $(z_{12} z_{23} \dots z_{n, 1})^{-1}$

... and integrand of F_{σ}^{τ} is $\sum_{\rho} S(\rho|\tau) \rho(z_{12} z_{23} \dots z_{n, 1})^{-1}$ w. KLT matrix

[Brödel, OS, Stieberger 1304.7267]

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$$\mathcal{A}_{\text{super}}^{\text{open}}(1, \sigma, n-1, n; \alpha') = \sum_{\rho, \tau \in S_{n-3}} Z_{\sigma}(1, \rho, n, n-1) S(\rho|\tau)_1 A_{\text{SYM}}(1, \tau, n-1, n)$$

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\implies field-theory double copy

$$\mathcal{A}_{\text{super}}^{\text{open}}(\sigma) = Z_{\sigma} \otimes_{\text{KLT}} A_{\text{SYM}}$$

Integrals Z_{σ} dubbed “Z-theory amplitudes” [Carrasco, Mafra, OS 1608.02569]

II. Closed superstrings from single-valued open superstrings

α' -expansion of F_σ^τ & $\mathcal{A}_{\text{super}}^{\text{open}}$ involves multiple zeta values (MZVs)

$$\zeta_{n_1, n_2, \dots, n_r} \equiv \sum_{0 < k_1 < k_2 < \dots < k_r}^{\infty} k_1^{-n_1} k_2^{-n_2} \dots k_r^{-n_r}, \quad n_r \geq 2$$

[Terasoma 2002 & Brown 2006]

Schematically,

$$\mathcal{A}_{\text{super}}^{\text{open}}(\sigma) = \overbrace{\left(\mathbb{1} + \zeta_2(\alpha' k_i \cdot k_j)^2 + \zeta_3(\alpha' k_i \cdot k_j)^3 + \mathcal{O}(\alpha'^4) \right)}^{F_\sigma^\tau(\alpha')} \sigma^\tau A_{\text{SYM}}(\tau)$$

Polynomial structure in $\alpha' k_i \cdot k_j$ at n points can be determined to any order

[Brödel, OS, Stieberger, Terasoma 1304.7304 & Mafra, OS 1609.07078]

Explicit results at $n \leq 7$ points available for download

<http://wwth.mpp.mpg.de/members/stieberg/mzv/index.html>

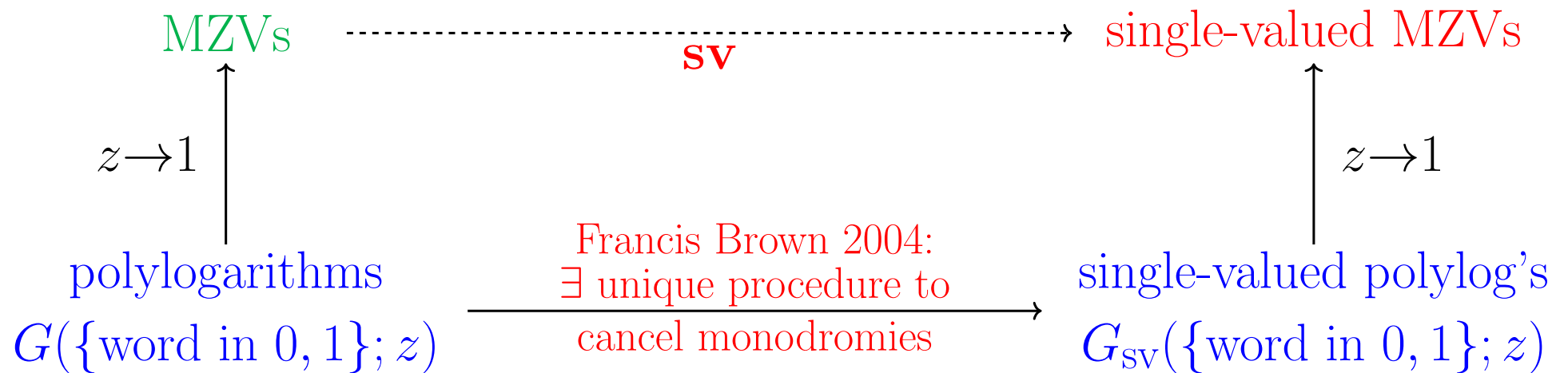
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Define single-valued projection **sv** of MZVs via their polylogarithm origin

[Schnetz 1302.6445 & Brown 1309.5309]



e.g. $G(1; z) = \log(1-z) \longrightarrow G_{\text{sv}}(1; z) = \log |1-z|^2$

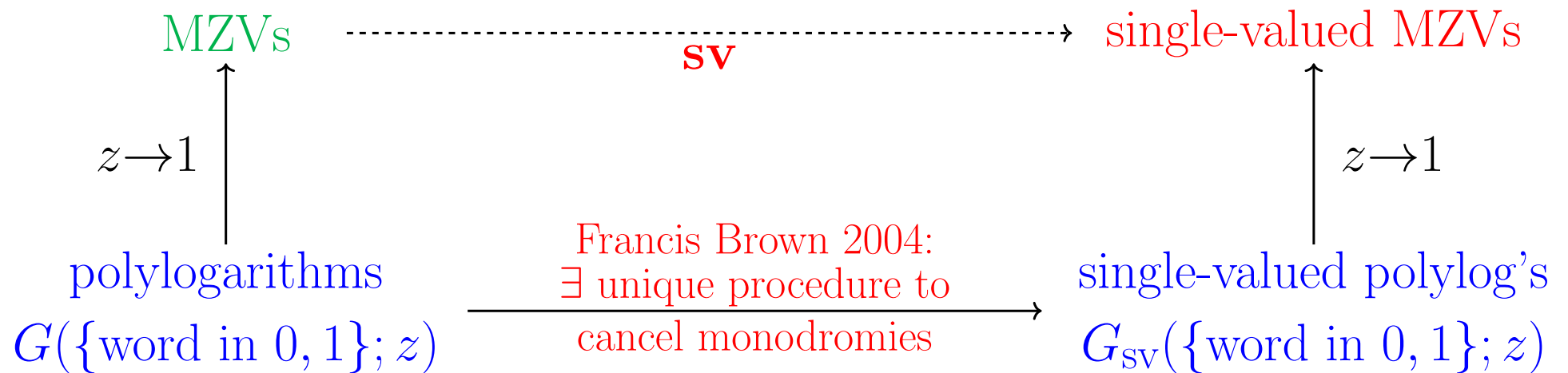
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$$\mathbf{sv}(\zeta_{2k}) = 0, \quad \mathbf{sv}(\zeta_{2k+1}) = 2 \zeta_{2k+1}, \quad \mathbf{sv}(\zeta_{3,5}) = -10 \zeta_3 \zeta_5, \quad \text{etc.}$$

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In principle, [Kawai, Lewellen, Tye 1986] determine closed-superstring trees as

$$\mathcal{M}_{\text{super}}^{\text{closed}}(\alpha') \sim \bar{\mathcal{A}}_{\text{super}}^{\text{open}}(\dots; \alpha') \prod \sin\left(\frac{\pi\alpha'}{2} k_i \cdot k_j\right) \mathcal{A}_{\text{super}}^{\text{open}}(\dots; \alpha') ,$$

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Simplifies to field-theory double copy @ $\sin\left(\frac{\pi\alpha'}{2} k_i \cdot k_j\right) \rightarrow k_i \cdot k_j$

$$\mathcal{M}_{\text{super}}^{\text{closed}}(\alpha') = \bar{\mathcal{A}}_{\text{SYM}} \otimes_{\text{KLT}} \mathbf{sv} \mathcal{A}_{\text{super}}^{\text{open}}(\alpha')$$

[OS, Stieberger 1205.1516 & Stieberger 1310.3259]

Schematically, by $\mathbf{sv}(\zeta_{2k}) = 0$ and $\mathbf{sv}(\zeta_{2k+1}) = 2\zeta_{2k+1}$,

$$\mathcal{A}_{\text{super}}^{\text{open}}(\sigma) = (\mathbf{1} + \zeta_2(\alpha' k_i \cdot k_j)^2 + \zeta_3(\alpha' k_i \cdot k_j)^3 + \mathcal{O}(\alpha'^4))_{\sigma}^{\tau} A_{\text{SYM}}(\tau)$$

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- emergence of $\mathbf{sv} \mathcal{A}_{\text{super}}^{\text{open}}(\alpha')$ still conjectural (tested to high orders in α')
- \exists first one-loop echos of \mathbf{sv} -map from open to closed strings

[D'Hoker, Green, Gürdogan, Vanhove 1512.06779 & Brödel, OS, Zerbini 1803.00527]

III. Bosonic strings from $(DF)^2 + \text{YM field theory}$

Open **bosonic** string: can still expand n -point trees via integrals F_σ^τ

$$\mathcal{A}_{\text{super}}^{\text{open}}(1, \sigma, n-1, n; \alpha') = \sum_{\tau \in S_{n-3}} F_\sigma^\tau(\alpha') A_{\text{SYM}}(1, \tau, n-1, n)$$

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α' -dependent **kinematic factors** $B(\dots; \alpha')$, e.g.

$$B(1, 2, 3; \alpha') = A_{\text{YM}}(1, 2, 3) - 4\alpha' (e_1 \cdot k_2)(e_2 \cdot k_3)(e_3 \cdot k_1) \quad \text{tachyon pole}$$

$$B(1, 2, 3, 4; \alpha') = A_{\text{YM}}(1, 2, 3, 4) - 4\alpha' s_{13} \left\{ \left[\frac{f_{12} f_{34}}{s_{12}^2 (1 - 2\alpha' s_{12})} + \text{cyc}(2, 3, 4) \right] - \frac{g_1 g_2 g_3 g_4}{s_{12}^2 s_{13}^2 s_{23}^2} \right\}$$

with $s_{ij} \equiv k_i \cdot k_j$ and $f_{ij} \equiv s_{ij}(e_i \cdot e_j) - (k_i \cdot e_j)(k_j \cdot e_i)$ and $g_i \equiv (k_{i-1} \cdot e_i)s_{i,i+1} - (k_{i+1} \cdot e_i)s_{i-1,i}$.

[Huang, OS, Wen 1602.01674]

Is there a **field-theory interpretation** of these $B(\dots; \alpha')$ with BCJ rel's?

Open **bosonic** string: can still expand n -point trees via integrals F_σ^τ

$$\mathcal{A}_{\text{bos}}^{\text{open}}(1, \sigma, n-1, n; \alpha') = \frac{\alpha'}{2} \sum_{\tau \in S_{n-3}} F_\sigma^\tau(\alpha') A_{(DF)^2+\text{YM}}(1, \tau, n-1, n; \alpha')$$

“(DF)² + YM” gauge theory \implies kin. factors $B \rightarrow \frac{\alpha'}{2} A_{(DF)^2+\text{YM}}$

[Azevedo, Chiodaroli, Johansson, OS 1803.05452]

$$\begin{aligned} \mathcal{L}_{(DF)^2+\text{YM}} \equiv & \frac{1}{2} (D_\mu F^{a\mu\nu})^2 - \frac{1}{4} m^2 (F_{\mu\nu}^a)^2 - \frac{1}{3} f^{abc} F_\mu^{a\nu} F_\nu^{b\lambda} F_\lambda^{c\mu} \\ & + \frac{1}{2} (D_\mu \varphi^\alpha)^2 - \frac{1}{2} m^2 (\varphi^\alpha)^2 + \frac{1}{2} C^{\alpha ab} \varphi^\alpha F_{\mu\nu}^a F^{b\mu\nu} + \frac{1}{3!} d^{\alpha\beta\gamma} \varphi^\alpha \varphi^\beta \varphi^\gamma \Big|_{m^2 = -\frac{1}{\alpha'}} \end{aligned}$$

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- kin. operator $(\partial^4 - m^2 \partial^2) A^2 \implies$ 2 gluon modes: $\left(\begin{smallmatrix} \text{massless} \\ \text{physical} \end{smallmatrix} \right) \oplus \left(\begin{smallmatrix} \text{massive} \\ \text{ghost} \end{smallmatrix} \right)$
- massive-ghost scalar φ^α : index $\alpha \leftrightarrow$ real representation of gauge group

$$\implies m^2 = -\frac{1}{\alpha'} \text{ gives required tachyon poles } A_{(DF)^2+\text{YM}} \sim \frac{1}{1 - 2\alpha' k_i \cdot k_j}$$

Open **bosonic** string: can still expand n -point trees via **integrals** F_σ^τ

$$\mathcal{A}_{\text{bos}}^{\text{open}}(1, \sigma, n-1, n; \alpha') = \frac{\alpha'}{2} \sum_{\tau \in S_{n-3}} F_\sigma^\tau(\alpha') A_{(DF)^2+\text{YM}}(1, \tau, n-1, n; \alpha')$$

“(DF)² + YM” gauge theory \implies kin. factors $B \rightarrow \frac{\alpha'}{2} A_{(DF)^2+\text{YM}}$
 [Azevedo, Chiodaroli, Johansson, OS 1803.05452]

$$\begin{aligned} \mathcal{L}_{(DF)^2+\text{YM}} \equiv & \frac{1}{2} (D_\mu F^{a\mu\nu})^2 - \frac{1}{4} m^2 (F_{\mu\nu}^a)^2 - \frac{1}{3} f^{abc} F_\mu^{a\nu} F_\nu^{b\lambda} F_\lambda^{c\mu} \\ & + \frac{1}{2} (D_\mu \varphi^\alpha)^2 - \frac{1}{2} m^2 (\varphi^\alpha)^2 + \frac{1}{2} C^{\alpha ab} \varphi^\alpha F_{\mu\nu}^a F^{b\mu\nu} + \frac{1}{3!} d^{\alpha\beta\gamma} \varphi^\alpha \varphi^\beta \varphi^\gamma \Big|_{m^2 = -\frac{1}{\alpha'}} \end{aligned}$$

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- massive-ghost scalar φ^α : index $\alpha \leftrightarrow$ real representation of gauge group
- for external gluons, Clebsch Gordan $C^{\alpha ab}$ & $d^{\alpha\beta\gamma}$ conspire to $\prod f^{abc}$
- initial construction of $\mathcal{L}_{(DF)^2+\text{YM}}$ guided by imposing BCJ relations
 [Johansson, Nohle 1707.02965]

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Since $F_\sigma^\tau = \sum_\rho Z_\sigma(\rho) S(\rho|\tau)$ signal KLT formula in disguise,

\implies field-theory double copy

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Corollaries by recycling the underlying open-string CFT correlator:

- via **sv** projection of MZVs:

$$\mathcal{M}_{\text{bos}}^{\text{closed}} = \mathbf{sv} \bar{\mathcal{A}}_{\text{bos}}^{\text{open}} \otimes_{\text{KLT}} A_{(DF)^2+\text{YM}}$$

- heterotic strings (gravity): can put SUSY on either side of double copy

$$\mathcal{M}_{\text{grav}}^{\text{het}} = \begin{cases} \mathbf{sv} \bar{\mathcal{A}}_{\text{super}}^{\text{open}} \otimes_{\text{KLT}} A_{(DF)^2+\text{YM}} & : \text{SUSY on } \mathbf{sv}(\text{string}) \text{ side} \\ \mathbf{sv} \bar{\mathcal{A}}_{\text{bos}}^{\text{open}} \otimes_{\text{KLT}} A_{\text{SYM}} & : \text{SUSY on field-theory side} \end{cases}$$

[Azevedo, Chiodaroli, Johansson, OS 1803.05452]

IV. Heterotic strings and $(DF)^2 + \text{YM} + \phi^3$ field theory

Now incorporate gauge sector & gauge/gravity couplings of heterotic string

- single-trace gluon amplitudes: $\mathcal{A}_{\text{s.tr.}}^{\text{het}}(1, \dots, n) = \mathbf{sv} \mathcal{A}_{\text{super}}^{\text{open}}(1, \dots, n)$
[Stieberger, Taylor 1401.1218]
- various mixed gauge/gravity & double-trace amplitudes reduced to $\mathcal{A}_{\text{s.tr.}}^{\text{het}}$
[OS 1608.00130]

Suggests double-copy structure in color-dressed het. amplitudes $\mathcal{M}_{\text{gauge} \oplus \text{grav}}^{\text{het}}$,

$$\mathcal{M}_{\text{gauge} \oplus \text{grav}}^{\text{het}} \sim A_{??} \otimes_{\text{KLT}} \mathbf{sv} \mathcal{A}_{\text{super}}^{\text{open}}$$

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 \mathcal{M}_{\oplus\text{grav}}^{\text{het gauge}} & \sim & A_{??} \otimes_{\text{KLT}} \mathbf{sv} \mathcal{A}_{\text{super}}^{\text{open}} \\
 \begin{array}{l} \nearrow \\ n \text{ gravitons} \\ k \text{ gluons} \end{array} & & \begin{array}{l} \uparrow \\ n \text{ gluons} \\ k \text{ scalars} \end{array} \\
 & & \begin{array}{l} \nwarrow \\ \text{adds spin one and} \\ \text{carries the SUSY} \end{array}
 \end{array}$$

As an amplitude $A_{??} \supset$ colored scalars “s” & uncolored vectors “g” e.g.

$$A_{??}(1_s, 2_s, 3_g) \sim \delta^{A_1 A_2} (e_3 \cdot k_1) .$$

Further examples of $A_{??}$ (with colored scalar “s” & uncolored vectors “g”)

$$A_{??}(1_s, 2_s, 3_s, 4_s) = \overbrace{\frac{\bar{f}^{A_1 A_2 B} \bar{f}^{B A_3 A_4}}{2s_{12}} + \frac{\bar{f}^{A_2 A_3 B} \bar{f}^{B A_4 A_1}}{2s_{23}}}^{\text{single trace}} + \overbrace{s_{13} \left\{ \frac{2\alpha' \delta^{A_1 A_2} \delta^{A_3 A_4}}{s_{12}(1 - 2\alpha' s_{12})} + \text{cyc}(2, 3, 4) \right\}}^{\text{double trace}}$$

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 \end{aligned}$$

suggest extension of $(DF)^2 + \text{YM}$ by bi-adjoint scalars $\phi = \sum_{a,A} \phi^{aA} t^a \otimes \bar{t}^A$

$$\begin{aligned}
 \mathcal{L}_{(DF)^2 + \text{YM} + \phi^3} &\equiv \frac{1}{2} (D_\mu F^{a\mu\nu})^2 - \frac{1}{4} m^2 (F_{\mu\nu}^a)^2 - \frac{1}{3} f^{abc} F_\mu^{a\nu} F_\nu^{b\lambda} F_\lambda^{c\mu} \\
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 \end{aligned}$$

- only A_μ^a, φ^α have tachyonic mode @ $m^2 = -\alpha'^{-1}$, the ϕ^{aA} are massless
- only color order w.r.t. $\text{Tr}(t^a t^b \dots)$, not w.r.t. $\text{Tr}(\bar{t}^A \bar{t}^B \dots)$

$$A_{(DF)^2 + \text{YM} + \phi^3}(1, 2, \dots, n) \equiv M_{(DF)^2 + \text{YM} + \phi^3} \Big|_{\text{Tr}(t^1 t^2 \dots t^n)}$$

All mixed gluon/graviton amplitudes of heterotic string from $(DF)^2 + \text{YM} + \phi^3$,

$$\begin{array}{ccc}
 \mathcal{M}_{\text{gauge} \oplus \text{grav}}^{\text{het}} & \sim & \mathcal{A}_{(DF)^2 + \text{YM} + \phi^3} \otimes_{\text{KLT}} \mathbf{sv} \mathcal{A}_{\text{super}}^{\text{open}} \\
 \uparrow & & \uparrow \\
 n \text{ gravitons} & & n \text{ gluons} \\
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V. Conclusions & Outlook

- generated all massless tree amplitudes of superstrings, bosonic strings and heterotic strings from field-theory double-copy “ \otimes_{KLT} ”

$$\left(\begin{array}{c} \text{string} \\ \text{amplitudes} \end{array} \right) = \left(\begin{array}{c} \text{gauge theory: possibly} \\ \text{with } (DF)^2/\phi^3 \text{ extension} \end{array} \right) \otimes_{\text{KLT}} \left(\begin{array}{c} \text{disk integrals or} \\ \text{sv(open strings)} \end{array} \right)$$

\otimes_{KLT}	SYM	$(DF)^2 + \text{YM}$	$(DF)^2 + \text{YM} + \phi^3$
Z-theory	open superstring	open bosonic string	comp(open bos. string)
sv(open superstring)	closed superstring	heterotic string (grav)	het. string (gauge/grav)
sv(open bos. string)	heterotic string (grav)	closed bosonic string	comp(closed bos. string)

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- also at one loop: evidence for double-copy structure of open superstrings

[Mafra, OS 1711.09104 & in progress]

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Thank you for your attention !