A space-like determination of a_{μ}^{HLO} via μ -e scattering data

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Muon-electron scattering: Theory kickoff workshop

4-5 September 2017

Padova Europe/Rome timezone

Muon g-2: summary of the present status

- E821 experiment at BNL has generated enormous interest: $a_{\mu}^{E821} = 11659208.9(6.3) \times 10^{-10}$ (0.54 ppm)
- Tantalizing $\sim 3\sigma$ deviation with SM (persistent since >10 years):

 $a_{\mu}^{SM} = 11659180.2(4.9) \times 10^{-10} (DHMZ)$

M. Davier, A. Hoecker, B. Malaescu and Z. Zhang, Eur. Phys. J. C71 (2011)

$$a_{\mu}^{E821} - a_{\mu}^{SM} \sim (28 \pm 8) \times 10^{-10}$$

- Current discrepancy limited by:
 - Experimental uncertainty → New experiments at FNAL and J-PARC x4 accuracy
 - Theoretical uncertanty → limited by hadronic effects



$(g-2)_{\mu}$: a new experiment at FNAL (E989)

- New experiment at FNAL (E989) at magic momentum, consolidated method. 20 x stat. w.r.t. E821. First result with BNL accuracy (0.54 ppm) expected in 2018-19.
 - → Ultimate precision: $\delta a_{\mu} x4$ improvement (0.14ppm)



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 - → Ultimate precision: $\delta a_{\mu} x4$ improvement (0.14ppm)

If the central value remains the same \Rightarrow 5-8 σ from SM* (enough to claim discovery of New Physics!)

*Depending on the progress on Theory BNL-E821 04 ave.



Complementary proposal at J-PARC in progress





(g-2)_u: First wiggle plot at FNAL (E989)

Number of high energy positrons as a function of time



Complementary proposal at J-PARC in progress

 $\frac{\alpha}{2\pi}$

П

89

230

Three Recent papers relevant for g-2!

20 years effort!

25 April 2017

High-precision calculation of the 4-loop contribution to the electron q-2 in QED

Stefano Laporta*

1100 digits!

Dipartimento di Fisica, Università di Bologna, Istituto Nazionale Fisica Nucleare, Sezione di Bologna, Via Irnerio 46, I-40126 Bologna, Italy

Abstract

I have evaluated up to 1100 digits of precision the contribution of the 891 4-loop Feynman diagrams contributing to the electron g-2 in QED. The total 4-loop contribution is

 $a_e = -1.912245764926445574152647167439830054060873390658725345\dots \left(\frac{\alpha}{-}\right)^4$

I have fit a semi-analytical expression to the numerical value. The expression contains harmonic polylogarithms of argument $e^{\frac{i\pi}{3}}$, $e^{\frac{i\pi}{3}}$, $e^{\frac{i\pi}{2}}$, one-dimensional integrals of products of complete elliptic integrals and six finite parts of master integrals, evaluated up to 4800 digits.

Eur. Phys. J. C (2017) 77:139 DOI 10.1140/epjc/s10052-017-4633-z

 $\delta a_{\mu}^{HLO}/a_{\mu}^{HLO} \rightarrow 0.3\%_{stat}$

Regular Article - Experimental Physics

Measuring the leading hadronic contribution to the muon g-2 via μe scattering

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Received: 17 October 2016 / Accepted: 17 January 2017 / Published online: 1 March 2017 © The Author(s) 2017. This article is published with open access at Springerlink.com The hadronic vacuum polarization contribution to the muon g - 2 from lattice QCD

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Abstract

We present a calculation of the hadronic vacuum polarization contribution to the muon anomalous magnetic moment, $a_{\mu}^{\rm hvp}$, in lattice QCD employing dynamical up and down quarks. We focus on controlling the infrared regime of the vacuum polarization function. To this end we employ several complementary approaches, including Padé fits, time moments and the time-momentum representation. We correct our results for finite-volume effects by combining the Gounaris-Sakurai parameterization of the timelike pion form factor with the Lüscher formalism. On a subset of our ensembles we have derived an upper bound on the magnitude of quark-disconnected diagrams and found that they decrease the estimate for $a_{\mu}^{\rm hvp}$ by at most 2%. Our final result is $a_{\mu}^{\rm hvp} = (654 \pm 32 \substack{+21 \\ -23}) \cdot 10^{-10}$, where the first error is statistical, and the second denotes the combined systematic uncertainty. Based on our findings we discuss the prospects for determining $a_{\mu}^{\rm hvp}$ with sub-percent precision.



(see talks from U. Marconi, S. Laporta and M. Marinkovic)

a_{μ}^{HLO} calculation, traditional way: time-like data

[C. Bouchiat, L. Michel, '61; N. Cabibbo, R. Gatto 61; L. Durand '62-'63; M. Gourdin, E. De Rafael, '69; S. Eidelman F. Jegerlehner '95,....]

• Optical theorem and analyticity:

$$\sigma(s)_{(e^+e^- \to had)} = \frac{4\pi}{s} \operatorname{Im} \Pi_{hadron}(s)$$

$$a_{\mu}^{HLO} = \frac{1}{4\pi^3} \int_{4m_{\pi}^2}^{\infty} ds \, K(s) \cdot \sigma(s)_{(e^+e^- \to had)}$$

• The main contribution is in the highly fluctuating low energy region.

$$K(s) = \int_0^1 dx \, \frac{x^2(1-x)}{x^2 + (1-x)(s/m^2)} \sim \frac{1}{s}$$

The enhancement at low energy implies that the $\rho \rightarrow \pi^+\pi^-$ resonance is dominating the dispersion integral (~ 75 %). Current precision at 0.6% \rightarrow need to be reduced by a factor ~2

G. Venanzoni, μ -e Theory Workshop, Padova, 4 September 2017



The high-energy tail of the integral is calculated using pQCD 7

Timelike data aiming at 0.2% on a_{μ}^{HLO} ?

- Not an easy task!
 - >30 channels to keep under control (at (sub)percent level)
 - local discrepancies in main channels (2π (KLOE/Babar), K⁺K⁻ CMD2/Babar)
 - Isospin corrections for not measured channels
 - Treatment of narrow resonances? (See F. Jegerlehner, ArXiv:1511.04473)



M. Davier, TAU16 WS

An independent/complementary approach is highly desirable!

Lattice-QCD progress on a_{μ}^{HVP}



- Can calculate nonperturbative vacuum polarization function П(Q²) directly in lattice QCD from simple 2-point correlation function of EM quark current [Blum, PRL 91 (2003) 052001]
- Several ongoing lattice efforts yielding new results since ICHEP 2014 including:
- (1) First calculation of quark-disconnected contribution [RBC/UKQCD, PRL116, 232002 (2016)]
- (2) Second complete calculation of leading-order a_{μ}^{HVP} [HPQCD, arXiv:1601.03071]
 - First to reach precision needed to observe significant deviation from experiment
 - ~1% total uncertainty by 2018 possible
 - Sub-percent precision will require inclusion of isospin breaking & QED, and hence take longer

$\alpha_{_{em}}$ running and the Vacuum Polarization

- Due to Vacuum Polarization effects α_{em}(q²) is a running parameter from its value at vanishing momentum transfer to the effective q².
- The "Vacuum Polarization" function Π(q²) can be "absorbed" in a redefinition of an effective charge:

$$e^{2} \rightarrow e^{2}(q^{2}) = \frac{e^{2}}{1 + (\Pi(q^{2}) - \Pi(0))} \qquad \alpha(q^{2}) = \frac{\alpha(0)}{1 - \Delta\alpha}; \quad \Delta\alpha = -\Re e \left(\Pi(q^{2}) - \Pi(0)\right)$$
$$\Delta\alpha = \Delta\alpha_{|} + \Delta\alpha_{|}^{(5)} + \Delta\alpha_{|}^{(5)} + \Delta\alpha_{|}^{(5)}$$

> Δa takes a contribution by non perturbative hadronic effects ($\Delta a^{(5)}_{had}$) which exibits a different behaviour in time-like and space-like region







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$$\Delta \alpha = \Delta \alpha_{\rm l} + \Delta \alpha^{(5)}_{\rm had} + \Delta \alpha_{\rm top}$$

➤ ∆a takes a contribution by non perturbative hadronic effects (∆a⁽⁵⁾_{had}) which exibits a different behaviour in time-like and space-like region







Running of α_{em}



Direct measurement of α running



Measurement of the running of $\alpha(s)$ time-like at KLOE ($e^+e^- \rightarrow \mu^+\mu^-\gamma$)



 $\Delta \alpha_{had}(s) = - \left(\frac{\alpha(0)s}{3\pi}\right) Re \int_{m_{\pi}^2}^{\infty} ds' \frac{R(s')}{s'(s'-s-i\epsilon)} \qquad R(s) = \frac{\sigma_{tot}(e^+e^- \rightarrow \gamma \ast \rightarrow hadrons)}{\sigma_{tot}(e^+e^- \rightarrow \gamma \ast \rightarrow \mu^+\mu^-)}$

Alternative approach: a_{μ}^{HLO} from space-like region

$$a_{\mu}^{HLO} = -\frac{\alpha}{\pi} \int_{0}^{1} (1-x) \Delta \alpha_{had} \left(-\frac{x^{2}}{1-x} m_{\mu}^{2}\right) dx$$

$$\begin{split} t &= \frac{x^2 m_{\mu}^2}{x-1} \quad 0 \leq -t < +\infty \\ x &= \frac{t}{2m_{\mu}^2} (1 - \sqrt{1 - \frac{4m_{\mu}^2}{t}}); \quad 0 \leq x < 1; \end{split}$$

- a_μ^{HLO} is given by the integral of the curve (smooth behaviour)
- It requires a measurement of the hadronic contribution to the effective electromagnetic coupling in the space-like region Δα_{had}(t) (t=q²<o)
- It enhances the contribution from low q² region (below 0.11 GeV²)
- Its precision is determined by the uncertainty on $\Delta \alpha_{had}$ (t) in this region





Measurement of $\Delta \alpha_{had}$ (t) spacelike at LEP

- $\Delta \alpha_{\rm had}$ (t) (t<o) has been measured at LEP using small angle Bhabha scattering

$$f(t) = \frac{N_{\text{data}}(t)}{N_{\text{MC}}^0(t)} \propto \left(\frac{1}{1 - \Delta\alpha(t)}\right)^2.$$

Accuracy at per mill level was achieved!

 For low t values (≤0.11 GeV²) and higher precision (~10⁻⁵) as in our case a different approach is needed!







Experimental approach:

Use of a 150 GeV μ beam on Be target at CERN (elastic scattering $\mu e \rightarrow \mu e$)



Why measuring $\Delta \alpha_{had}$ (t) with a 150 GeV μ beam on e⁻ target ?

It looks an ideal process!

- $\mu e \rightarrow \mu e$ is pure t-channel (at LO)
- Simple kinematics (2 body process, t=-2m_eE_e<o) allows to span the region o<-t<0.143 GeV² (o<x<0.93); 87% of total a_{μ}^{HLO} (the rest can be computed by pQCD/time-like data)
- Angular measurement: high boosted system gives access to all angles (t) in the cms region $\begin{array}{l} \theta_e^{\ LAB} < 32 \ mrad \ (E_e > 1 \ GeV) \\ \theta_\mu^{\ LAB} < 5 \ mrad \end{array}$
- It allows using the same detector for signal and normalization (x<0.3, $\Delta \alpha_{had}(t) < 10^{-5}$) \rightarrow cancellation of detector effects at first order



Detector considerations

- Modular apparatus: 20 layers of 3 cm Be (target), each coupled to 1 m distant Si (0.3 mm) planes. It provides a 0.02 mrad resolution on the scattering angle
- The t=q² <0 of the interaction is determined by the electron (or muon) scattering angle (a` la NA7)
- ECAL and μ Detector located downstream to solve PID ambiguity below 5 mrad. Above that, angular measurement gives correct PID
- It provides uniform full acceptance, with the potential to keep the systematic errors at 10⁻⁵
- Statistical considerations show that a 0.3% error can be achieved on a_{μ}^{HLO} in 2 years of data taking with 1.3x10⁷ μ /s (available at CERN)

More details in Umberto's talk



Muon beam M2 at CERN

"Forty years ago, on 7 May 1977, CERN inaugurated the world's largest accelerator at the time – the Super Proton Synchrotron". COMPASS



$I_{beam} > 10^7 \, muon/s, E_{\mu} = 150 \, GeV$

Systematics

Affordable by means of

GEANT4 based simulations

- 1. Acceptance
- 2. Tracking
- 3. Trigger
- 4. PID
- 5. Effects of E_e energy cut
- Signal/Background:
 It requires a dedicated event generator.
- Uncertainty in the location of interaction vertices: Segmented/ active target to resolve the vertex position
- 8. Uncertainty in the muon beam momentum: Scattering kinematics to determine the beam momentum
- Effects of Multiple Scattering (must be known at ≤1%): It requires dedicated work on simulation and measurements (test beam).
- 10. Theoretical uncertainty on the mu-e cross section (see later)

All the systematic effects must be known to ensure an errror on the cross section < 10ppm

Activity on the theory side

- Development of dedicated NNLO MC event generator for mu-e scattering. Ultimate precisione ≤10ppm. Many ingredients...see this workshop!!!
- 2. Theory workshop next year in Mainz (19-24 February 2018). You are all invited!

Plans

2017 - 2019

- Detector optimization studies
- Test beams (first on 27 Sep-3 October 2017 at CERN)
- Set up a collaboration
- Theoretical studies
- Letter of Intent to the SPSC

2020

- Detector construction and installation
- (a staged version of the detector may be) LHC roadmap, according to MTP 2016-2020* 2021 2024
- - Ure HLO Long Shutdown (LS) – Sta (not necessarily the ultimate precision)



Conclusion

- Proposal part of the CERN "Physics Beyond Collider Study Group" at CERN (<u>http://pbc.web.cern.ch/</u>)
- If approved (by CERN SPSC) first results in the same period of the g-2 measurements at Fermilab and J-Parc (first results around 2021-2024)
- Many experimental and theoretical challenges: very exciting!!
- We need your help!!



Thanks!

Reference papers

A new approach to evaluate the leading hadronic corrections to the muon g-2

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Measuring the leading hadronic contribution to the muon g-2 via μe scattering

G. Abbiendi¹, C. M. Carloni Calame², U. Marconi¹, C. Matteuzzi³, G. Montagna^{4,2},
O. Nicrosini², M. Passera⁵, F. Piccinini², R. Tenchini⁶, L. Trentadue^{7,3}, and G. Venanzoni⁸ ¹INFN, Sezione di Bologna, Bologna, Italy ²INFN, Sezione di Pavia, Pavia, Italy ³INFN, Sezione di Milano Bicocca, Milano, Italy ⁴Dipartimento di Fisica, Università di Pavia, Pavia, Italy ⁵INFN, Sezione di Padova, Padova, Italy ⁶INFN, Sezione di Pisa, Pisa, Italy ⁷Dipartimento di Fisica e Scienze della Terra "M. Melloni", Università di Parma, Parma, Italy ⁸INFN, Laboratori Nazionali di Frascati, Frascati, Italy

Spare

Previous tests of the hadronic contribution to VP

1) '73: $\phi(1020)$ contribution to VP at ACO (Orsay e+e-) in the e⁺e⁻ $\rightarrow \mu + \mu$ - process: evidence at 3 σ in the region ±5 MeV around the second the

2) 70's: g-2 experiment at CERN: evidence for hadronic contribution to g-2 at 6σ
[Phys. Lett. 67B (1977) 225; Phys. Lett. 68B
(1977) 191]

3) 2006: OPAL at LEP: evidence for hadronic contributon Δa_{had} (t) (t<0) at 3σ in Bhabha scattering at small angle [Eur.Phys.J. C45 (2006) 1-21]
 G. Venanzoni, PHIPSI17, Mainz, 26 June 2017



0.995

-t (GeV²)