Muon-electron Scattering @NNLO in QED

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> In collaboration with: - Passera, Peraro, Primo, Ossola, Schubert, Torres-Bobadilla







Outline

Motivations

Adaptive Integrand Decomposition

Improved reduction @ 1- and 2-loop

Automated two-loop corrections for generic processes

Feynman Integrals in Dimensional Regularization
Integration-by-parts identities, Master Integrals & Differential Equations
Magnus Exponential Matrix and Canonical Forms

Results

Conclusions/Outlook

Motivations



Muon-electron scattering

Abbiendi, Carloni Calame, Marconi, Matteuzzi, Montagna, Nicrosini, MP, Piccinini, Tenchini, Trentadue, Venanzoni EPJC 2017 - arXiv:1609.08987

>> see Trentadue & Marconi's talk

What we know :: Anatomy of NLO



What we need :: Anatomy of NNLO



What we need :: Anatomy of NNLO



Amplitudes Decomposition:



 $a_x = a.i$

 $a_y = a.$

 $a_z = a_k$



Projections :: On-Shell Cut-Conditions



y

Completeness Relations: cutting "1"

• the richness of factorization

$$i(-i) = 1$$

$$\sum_{n} |\psi_{n}\rangle \langle \psi_{n}| = 1$$

$$(p^2 - m^2) = (\not p - m)(\not p + m)$$

$$\varepsilon^{\mu\nu} = \varepsilon^{\mu}\varepsilon^{\nu}$$

Completeness Relations: cutting "1"

• the richness of factorization —> ideas for workshop organisation









Amplitudes Decomposition: *the algebraic way*



Multi-loop Integrand Decomposition



graph

Multi-loop Integrand Decomposition

Polynomial Division

Ossola & **P.M.** (2011); Zhang (2012); Badger Frellesvig Zhang (2012) Mirabella, Ossola, Peraro, & **P.M.** (2012)

$$\mathcal{N}_{i_1\dots i_n} = \sum_{\kappa=1}^n \mathcal{N}_{i_1\dots i_{\kappa-1}i_{\kappa+1}\dots i_n} D_{i_\kappa} + \Delta_{i_1\dots i_n}$$

$$\frac{\mathcal{N}_{i_1\dots i_n}}{D_{i_1}\cdots D_{i_n}} = \sum_{\kappa=1}^n \frac{\mathcal{N}_{i_1\dots i_{\kappa-1}i_{\kappa+1}\dots i_n} D_{i_\kappa}}{D_{i_1}\cdots D_{i_{\kappa-1}}D_{i_\kappa}D_{i_{\kappa+1}}\cdots D_{i_n}} + \frac{\Delta_{i_1\dots i_n}}{D_{i_1}\cdots D_{i_n}}$$

Multi-loop Integrand Decomposition

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$$\frac{\mathcal{N}_{i_1\dots i_n}}{D_{i_1}\cdots D_{i_n}} = \sum_{\kappa=1}^n \frac{\mathcal{N}_{i_1\dots i_{\kappa-1}i_{\kappa+1}\dots i_n} D_{i_\kappa}}{D_{i_1}\cdots D_{i_{\kappa-1}}D_{i_\kappa}D_{i_{\kappa+1}}\cdots D_{i_n}} + \frac{\Delta_{i_1\dots i_n}}{D_{i_1}\cdots D_{i_n}}$$



Multi-Loop Integrand Recurrence

Ossola & **P.M.** (2011); Zhang (2012); Badger Frellesvig Zhang (2012) Mirabella, Ossola, Peraro, & **P.M.** (2012)

el-Loop Recurrence Relation





Unique Methodology: Multiple-cuts as Projectors



Unique Methodology: Multiple-cuts as Projectors

What about higher orders?



Longitudinal and Transverse Space

Peraro Primo & P.M. (2016)

Dimensional Regularization

$$d = 4 - 2\epsilon$$



Denominators do not depend on "the angular variables" of the Transverse Space

Mumerators depend on "all" loop variables

$$egin{aligned} d = 4 - 2\epsilon \ q^lpha = q^lpha_{[4]} + \mu^lpha, & q^lpha_{[4]} = \sum_{i=1}^4 x_i e^lpha_i, & q^2 = q^2_{[4]} + \mu^2. \ \mathcal{D}_i = \left(q_{[4]} + \sum_{j=0}^i p_j\right)^2 + \mu^2 + m^2_i, \end{aligned}$$

$$d=d_{//}+d_{\perp}$$

$$q^lpha = q^lpha_{[k]} + \lambda^lpha, \qquad q^lpha_{[k]} = \sum_{j=1}^k x_j e^lpha_j, \qquad q^2 = q$$

$$q^2 = q_{[k]}^2 + \lambda^2,$$

k-dimensional the space spanned by the external momenta

$$\lambda^{\alpha} = \sum_{j=k+1}^{4} x_j e_j^{\alpha} + \mu^{\alpha}, \qquad \lambda^2 = \sum_{j=k+1}^{4} x_j^2 + \mu^2,$$

$$(d-k)$$
-dimensional orthogonal subspace.

$$\mathcal{D}_i = \left(q_{[k]} + \sum_{j=0}^i p_j\right)^2 + \lambda^2 + m_i^2.$$

Adaptive Unitarity @ 2-loop

Novel Integrand red'n



Arbitrary (external and internal) kinematics!

8 and 7 legs

$\mathcal{I}_{i_1 \cdots i_n}$	$\Delta_{i_1 \cdots i_n}$	$\mathcal{I}_{i_1 \cdots i_n}$	$\Delta_{i_1 \cdots i_n}$
τ^{P}	1	τP \rightarrow	6
	{1}		$\{1, x_{41}\}$
\mathcal{T}^{NP1}	1	$\mathcal{T}_{12}^{\text{NP1}}$	10
2123456789101.	{1}		$\{1, x_{42}\}$
Inp2	1	$\mathcal{I}_{122456801011}^{\text{NP1}}$	6
-12345078910117	{1}		$\{1, x_{42}\}$
$\mathcal{I}^{\mathrm{P}}_{234567801011}$	6	$\mathcal{I}_{124567801011}^{NP2}$	10
	$\{1, x_{41}\}$	124507851011 Y	$\{1, x_{42}\}$
$\mathcal{I}_{224567801011}^{\text{NP1}}$	10	$\mathcal{I}_{24567801011}^{\text{NP1}}$	15
254507891011	$\{1, x_{42}\}$	24507891011	$\{1, x_{31}, x_{41}\}$
$\mathcal{I}^{\mathrm{NP2}}_{123457891011} \qquad $	6	τ^{NP2}	33
	$\{1, x_{42}\}$		$\{1, x_{41}, x_{42}\}$
$ au^{NP2}$	10	$\mathcal{I}^{\mathrm{NP1}}_{12456891011} \qquad $	39
	$\{1, x_{42}\}$		$\{1, x_{41}, x_{42}\}$
TP T	15	$\mathcal{I}_{12245691011}^{\text{NP1}}$	15
	$\{1, x_{31}, x_{41}\}$		$\{1, x_{32}, x_{42}\}$
TP	33	τ^{NP2}	45
-23457891011	$\{1, x_{41}, x_{42}\}$		$\{1, x_{41}, x_{42}\}$
7NP1	39	$\mathcal{I}_{247891011}^{\rm NP1}$	20
-23457891011	$\{1, x_{41}, x_{42}\}$		$\{1, x_{21}, x_{31}, x_{41}\}$
$\mathcal{T}_{1224}^{\text{NP1}}$	15	$\tau^{\rm NP1}_{\rm ext}$	76
	$\{1, x_{32}, x_{42}\}$		$\{1, x_{31}, x_{41}, x_{42}\}$
Inp2	45	$\mathcal{I}_{2457801011}^{\text{NP1}}$	116
-2340/891011	$\{1, x_{41}, x_{42}\}$		$\{1, x_{41}, x_{32}, x_{42}\}$
		$\mathcal{T}_{1047701011}^{\text{NP1}}$	80
			$\{1, x_{31}, x_{41}, x_{42}\}$

6 and 5 legs

$\mathcal{I}_{i_1 \cdots i_n}$	$\Delta_{i_1 \cdots i_n}$	$\mathcal{I}_{i_1 \cdots i_n}$	$\Delta_{i_1 \cdots i_n}$
$\mathcal{I}_{13567891011}^{P}$	15	$\mathcal{I}^{\rm P}_{1567891011}$	20
	$\{1, x_{31}, x_{41}\}$		$\{1, x_{21}, x_{31}, x_{41}\}$
$\mathcal{I}_{12456791011}^{P}$	$\begin{cases} 62\\ 1 & \pi \\ 1 & \pi \\ 2 & \pi $	$\mathcal{I}^{\mathrm{P}}_{1356791011}$	76
	$\frac{1, x_{41}, x_{42}}{30}$		$\frac{1, x_{31}, x_{41}, x_{42}}{80}$
$\mathcal{I}_{2356891011}^{\rm NP1}$	$\{1, x_{41}, x_{42}\}$	$\mathcal{I}_{1567891011}^{\rm NP1}$	$\{1, x_{31}, x_{41}, x_{42}\}$
τ NP1 $-$	15	τ^{p}	15
	$\{1, x_{32}, x_{42}\}$		$\{1, x_{11}, x_{21}, x_{31}, x_{41}\}$
τ^{NP2}	45	$\tau^{\rm NP1}$	116
	$\{1, x_{41}, x_{42}\}$		$\{1, x_{31}, x_{32}, x_{42}\}$
$\mathcal{T}^{\mathrm{P}}_{T}$	20	$\mathcal{T}^{\mathrm{P}}_{\mathrm{P}}$	94
22567891011	$\{1, x_{21}, x_{31}, x_{41}\}\$		$\{1, x_{21}, x_{31}, x_{41}, x_{42}\}$
$\tau^{\rm P}$	76	τ^{P}	66
	$\{1, x_{31}, x_{41}, x_{42}\}$	21678911	$\{1, x_{11}, x_{21}, x_{31}, x_{41}, x_{42}\}$
τ^{NP1}	80	$\mathcal{T}^{\mathrm{P}}_{\mathrm{P}}$	160
	$\{1, x_{31}, x_{41}, x_{42}\}\$		$\{1, x_{31}, x_{41}, y_{32}, x_{42}\}\$
τ NP1	116	$\tau^{\rm NP1}$	185
22456891011	$\{1, x_{41}, x_{32}, x_{42}\}$		$\{1, x_{31}, x_{41}, x_{32}, x_{42}\}$
τ^{P}	15	τ^{P}	180
2367891011	$\{1, x_{11}, x_{21}, x_{31}, x_{41}\}$	21256911	$\{1, x_{11}, x_{31}, x_{41}, x_{32}, x_{42}\}$
$\tau^{\rm P}$	94	$\tau^{\rm NP1}_{\rm NP1}$	246
	$\{1, x_{21}, x_{31}, x_{41}, x_{42}\}\$		$\{1, x_{31}, x_{41}, x_{22}, x_{32}, x_{42}\}$
7P	160		
-235791011	$\{1, x_{31}, x_{41}, x_{32}, x_{42}\}\$		
$\mathcal{T}^{\text{NP1}}_{\text{NP1}}$	185		
	$\{1, x_{31}, x_{41}, x_{32}, x_{42}\}\$		

• 4 legs: divide-integra-divide



$\mathcal{I}_{i_1 \cdots i_n}$	$\Delta_{i_1 \cdots i_n}$	$\Delta^{\mathrm{int}}_{i_1\cdots i_n}$	$\Delta'_{i_1\cdots i_n}$
τ^{P}	94	53	10
L156791011	$\{1, x_{21}, x_{31}, x_{41}, x_{42}\}\$	$\{1, x_{21}, x_{31}, \lambda_{11}, \lambda_{22}, \lambda_{12}\}$	$\{1, x_{21}, x_{31}\}$
τ^{P}	160	93	22
	$\{1, x_{31}, x_{41}, x_{32}, x_{42}\}$	$\{1, x_{31}, x_{32}, \lambda_{11}, \lambda_{22}, \lambda_{12}\}$	$\{1, x_{31}, x_{32}\}$
τ^{NP1}	184	105	25
	$\{1, x_{31}, x_{42}, x_{32}, x_{42}\}$	$\{1, x_{31}, x_{32}, \lambda_{11}, \lambda_{22}, \lambda_{12}\}$	$\{1, x_{31}, x_{32}\}$
τ^{P}	180	101	39
L ₁₃₅₆₈₁₁	$\{1, x_{31}, x_{41}, x_{22}, x_{32}, x_{42}\}$	$\{1, x_{31}, x_{22}, x_{32}, \lambda_{11}, \lambda_{22}, \lambda_{12}\}$	$\{1, x_{31}, x_{22}, y_{32}\}$
$\tau^{\rm P}$	66	35	10
	$\{1, x_{11}, x_{21}, x_{31}, x_{41}, x_{42}\}$	$\{1, x_{11}, x_{21}, x_{31}, \lambda_{11}, \lambda_{22}, \lambda_{12}\}$	$\{1, x_{11}, x_{21}, x_{31}\}$
τ^{NP1}	245	137	55
	$\{1, x_{31}, x_{41}, x_{21}, x_{32}, x_{42}\}$	$\{1, x_{31}, x_{22}, x_{32}, \lambda_{11}, \lambda_{22}, \lambda_{12}\}$	$\{1, x_{31}, x_{22}, y_{32}\}$
$\mathcal{I}_{3681011}^{\mathrm{P}}$	115	66	35
	$\{1, x_{31}, x_{41}, x_{12}, x_{22}, x_{32}, x_{42}\}\$	$\{1, x_{31}, x_{12}, x_{22}, x_{32}, \lambda_{11}, \lambda_{22}, \lambda_{12}\}$	$\{1, x_{31}, x_{12}, x_{22}, x_{32}\}$
τ^{P}	180	103	60
	$\{1, x_{11}, x_{31}, x_{41}, x_{22}, x_{32}, x_{42}\}$	$\{1, x_{11}, x_{31}, x_{22}, x_{32}, \lambda_{11}, \lambda_{22}, \lambda_{12}\}$	$\{1, x_{11}, x_{31}, x_{22}, x_{32}\}$

• 3, 2, 1 legs: divide-integra-divide



$\mathcal{I}_{i_1 \cdots i_n}$	$\Delta_{i_1 \cdots i_n}$	$\Delta^{\mathrm{int}}_{i_1\cdots i_k}$	$\Delta'_{i_1\cdots i_k}$
τ^{P}	180	22	4
L ₁₃₅₆₉₁₁	$\{1, x_{31}, x_{41}, x_{22}, x_{32}, x_{42}\}$	$\{1, x_{22}, \lambda_{11}, \lambda_{22}, \lambda_{12}\}$	$\{1, x_{22}\}$
τ^{NP1} \rightarrow	240	30	6
L15691011	$\{1, x_{31}, x_{41}, x_{22}, x_{32}, x_{42}\}$	$\{1, x_{22}, \lambda_{11}, \lambda_{22}, \lambda_{12}\}$	$\{1, x_{22}\}$
<i>I</i> ^P ₁₅₇₁₀₁₁	180	33	13
	$\{1, x_{21}, x_{31}, x_{41}, x_{12}, x_{32}, x_{42}\}$	$\{1, x_{21}, x_{12}, \lambda_{11}, \lambda_{22}, \lambda_{12}\}$	$\{1, x_{21}, x_{12}\}$
<i>I</i> ^P ₁₆₉₁₀₁₁	115	20	6
	$\{1, x_{31}, x_{41}, x_{12}, x_{22}, x_{32}, x_{42}\}$	$\{1, x_{11}, x_{22}\lambda_{11}, \lambda_{22}, \lambda_{12}\}$	$\{1, x_{12}, x_{22}\}$
τ^{P}	100	26	16
2361011	$\left\{1, x_{11}, x_{21}, x_{31}, x_{41}, x_{22}, x_{32}, x_{42}\right\}$	$\{1, x_{11}, x_{21}, x_{22}, \lambda_{11}, \lambda_{22}, \lambda_{12}\}$	$\{x_{11}, x_{21}, x_{22}\}$

$\mathcal{I}_{i_1 \cdots i_n}$	$\Delta_{i_1 \cdots i_n}$	$\Delta_{i_1\cdots i_n}^{\mathrm{int}}$	$\Delta'_{i_1\cdots i_n}$
τP	180	8	1
	$\{1, x_{21}, x_{31}, x_{41}, x_{22}, x_{32}, x_{42}\}$	$\{1,\lambda_{11},\lambda_{22},\lambda_{12}\}$	{1}
$\tau^{\rm P}$ $-$	100	8	3
L161011	$\{1, x_{11}, x_{21}, x_{31}, x_4, x_{22}, y_3, x_{42}\}$	$\{1, x_{11}, \lambda_{11}, \lambda_{22}, \lambda_{12}\}$	$\{1, x_{11}\}$
$\mathcal{I}^{\mathrm{P}}_{131011}$ ~~	100	26	16
	$\{1, x_{11}, x_{21}, x_{31}, x_{41}, x_{12}, x_{32}, x_{42}\}$	$\{1, x_{11}, x_{21}, x_{12}, \lambda_{11}, \lambda_{22}, \lambda_{12}\}$	$\{1, x_{11}, x_{21}, x_{12}\}$
τ^{P} \frown	45	9	6
	$\{1, x_{11}, x_{21}, x_{31}, x_{41}, x_{12}, x_{22}, x_{32}, x_{42}\}\$	$\{1, x_{11}, x_{12}, \lambda_{11}, \lambda_{22}, \lambda_{12}\}$	$\{1, x_{11}, x_{12}\}$
$\tau^{\rm P}$	45	18	15
	$\{1, x_{11}, x_{21}, x_{31}, x_{41}, x_{12}, x_{22}, x_{32}, x_{42}\}\$	$\{1, x_{11}, x_{21}, x_{12}, x_{22}, \lambda_{11}, \lambda_{22}, \lambda_{12}\}$	$\{1, x_{11}, x_{22}, x_{21}, x_{22}\}$

\mathcal{I}_{i_1}	$_1 \cdots i_n$	$\Delta_{i_1 \cdots i_n}$	$\Delta_{i_1\cdots i_n}^{\rm int}$	$\Delta'_{i_1\cdots i_n}$
$ au^{\mathrm{P}}$	\bigcirc	45	4	1
L ₁₁₀ 11		$\{1, x_{11}, x_{21}, x_{31}, x_{41}, x_{12}, x_{22}, x_{32}, x_{42}\}\$	$\{1, \lambda_{11}, \lambda_{22}, \lambda_{12}\}$	{1}

Adaptive Integrand Decomposition

Peraro Primo & **P.M.** (2016)

Integrand reduction beyond multivariate polynomial division

$$d = d_{//} + d_{\perp}$$

idea n.1Integrating over Transverse Spaceidea n.2Cutting in the Longitudinal Space

1&2-loop Automation :: AIDA Peraro Primo TorresBobadilla & P.M. (w.i.p.)

Analytic 1- and 2-loop decomposition

Application to Mu-e scattering

Ossola Peraro Primo TorresBobadilla Schubert & P.M. (w.i.p.)

Dimensionally Regulated Integrals

Graph Topology & Integrals



 $e = # \text{ legs :: } p_i, \quad (i = 1, ..., e);$ $\ell = # \text{ loops :: } q_i \quad (i = 1, ..., \ell);$ $n = # \text{ denominators :: } D_i \quad (i = 1, ..., n);$

N = # scalar products (of types $q_i \cdot p_j$ and $q_i \cdot q_j$) $N = \ell(e-1) + \frac{\ell(\ell+1)}{2}$

n = # reducible scalar products (expressed in terms of denominators);

m = # irreducible scalar products $= N - n :: S_i \quad (i = 1, ..., m)$

Graph Topology & Integrals



$$e = \# \text{ legs } :: p_i, \quad (i = 1, \dots, e);$$

$$\ell = \# \text{ loops } :: q_i \quad (i = 1, \dots, \ell);$$

$$n = \# \text{ denominators } :: D_i \quad (i = 1, \dots, n);$$

 $N=\texttt{\texttt{\#}}$ scalar products (of types $q_i\cdot p_j$ and $q_i\cdot q_j$)

$$N = \ell(e - 1) + \frac{\ell(\ell + 1)}{2}$$

n = # reducible scalar products (expressed in terms of denominators);

m = # irreducible scalar products $= N - n :: S_i \quad (i = 1, ..., m)$

Graph Topology & Integrals



$$e = \# \text{ legs } :: p_i, \quad (i = 1, ..., e);$$

 $\ell = \# \text{ loops } :: q_i \quad (i = 1, ..., \ell);$
 $n = \# \text{ denominators } :: D_i \quad (i = 1, ..., n),$

$$N = \#$$
 scalar products (of types $q_i \cdot p_j$ and $q_i \cdot q_j$) $N = \ell(e-1) + \frac{\ell(\ell+1)}{2}$

n = # reducible scalar products (expressed in terms of denominators);

$$m = \#$$
 irreducible scalar products $= N - n :: S_i \quad (i = 1, ..., m)$

Associated Integrals ::

$$F_{n,m}^{[d]}(\mathbf{x}, \mathbf{y}) \equiv \int_{q_1 \dots q_\ell} f_{n,m}(\mathbf{x}, \mathbf{y}) , \qquad \int_{q_1 \dots q_\ell} \equiv \int \frac{\mathrm{d}^d q_1}{(2\pi)^d} \dots \frac{\mathrm{d}^d q_\ell}{(2\pi)^d}$$
$$f_{n,m}(\mathbf{x}, \mathbf{y}) = \frac{S_1^{y_1} \dots S_m^{y_m}}{D_1^{x_1} \dots D_n^{x_n}} \longleftarrow$$

Integration-by-parts Identities (IBPs)

Tkachov; Chetyrkin Tkachov; Laporta;

$$\int_{q_1\dots q_\ell} \frac{\partial}{\partial q_i^{\mu}} \Big(v^{\mu} f_{n,m}(\mathbf{x}, \mathbf{y}) \Big) = 0 , \qquad v = q_1, \dots, q_\ell, \ p_1, \dots, p_{\ell-1}.$$

 $\forall (n,m), N_{\text{IBP}} = \# \text{ of IBP relations} = \ell(\ell + e - 1)$

Relations between integrals associated to the same topology (or subtopologies)

$$c_0 \ F_{n,m}^{[d]}(\mathbf{x}, \mathbf{y}) + \sum_{i,j} c_{i,j} \ F_{n,m}^{[d]}(\mathbf{x}_i, \mathbf{y}_j) = 0 ,$$
$$\mathbf{x}_i = \{x_1, \dots, x_i \pm 1, \dots, x_n\}$$

 $\mathbf{y}_{\mathbf{i}} = \{y_1, \dots, y_j \pm 1, \dots, y_n\}$

public codes :: AIR; Reduze2; FIRE; LiteRed; private codes :: ... many authors ... Laporta, Sturm ...

Master Integrals (MIs)

Independent set of integrals $M_i^{[d]}$,

$$M_i^{[d]} \equiv \int_{q_1...q_\ell} m_i(\bar{\mathbf{x}}, \bar{\mathbf{y}}) ,$$

with a definite set of powers $\bar{\mathbf{x}}, \bar{\mathbf{y}}$ such that

$$F_{n,m}^{[d]}(\mathbf{x}, \mathbf{y}) \stackrel{\text{IBP}}{=} \sum_{k} c_k M_k^{[d]}, \quad \forall (n, m)$$

They form a *basis* for the integrals of the corresponding topology.

Two special cases

Two types of integrals generated from the master integrands

• Polynomial insertion:

$$\int_{q_1\dots q_\ell} P(q_i \cdot p_j, q_i \cdot q_j) \ m_i(\bar{\mathbf{x}}, \bar{\mathbf{y}}) = \sum_{n,m} \alpha_{n,m} \ F_{n,m}^{[d]}(\mathbf{x}, \mathbf{y}) \stackrel{\text{IBP}}{=} \sum_i c_i \ M_i^{[d]}$$

• External-leg derivatives:

$$p_i^{\mu} \frac{\partial}{\partial p_j^{\mu}} M_k^{[d]} = \int_{q_1 \dots q_\ell} p_i^{\mu} \frac{\partial}{\partial p_j^{\mu}} \ m_k(\bar{\mathbf{x}}, \bar{\mathbf{y}}) = \sum_{n,m} \beta_{n,m} \ F_{n,m}^{[d]}(\mathbf{x}, \mathbf{y}) \stackrel{\text{IBP}}{=} \sum_i c_i \ M_i^{[d]}$$



Differential Equations for Master Integrals



Differential Equations for Master Integrals

$$p^{2}\frac{\partial}{\partial p^{2}}\left\{p-p\right\} = \frac{1}{2}p_{\mu}\frac{\partial}{\partial p_{\mu}}\left\{p-p\right\}$$

Kotikov; Remiddi; Gehrmann Remiddi Argeri Bonciani Ferroglia Remiddi **P.M**.

Henn; Henn Smirnov Lee; Papadopoulos; Argeri diVita Mirabella Schlenk Schubert Tancredi **P.M**. diVita Schubert Yundin **P.M**. Remiddi Tancredi; Primo Tancredi Papadopoulos Frellesvig Zheng

$$P^{2}\frac{\partial}{\partial P^{2}}\left\{ \int_{p_{2}}^{p_{1}} -p_{3} \right\} = \left[A\left(p_{1,\mu}\frac{\partial}{\partial p_{1,\mu}} + p_{2,\mu}\frac{\partial}{\partial p_{2,\mu}} \right) + B\left(p_{1,\mu}\frac{\partial}{\partial p_{2,\mu}} + p_{2,\mu}\frac{\partial}{\partial p_{1,\mu}} \right) \right] \left\{ \int_{p_{2}}^{p_{1}} -p_{3} \right\}$$

$$P = p_{1} + p_{2},$$

$$P^{2}\frac{\partial}{\partial P^{2}}\left\{ \underbrace{p_{1}}_{p_{2}} \underbrace{p_{3}}_{p_{4}} \right\} = \left[C\left(p_{1,\mu}\frac{\partial}{\partial p_{1,\mu}} - p_{3,\mu}\frac{\partial}{\partial p_{3,\mu}} \right) + Dp_{2,\mu}\frac{\partial}{\partial p_{2,\mu}} + E(p_{1,\mu} + p_{3,\mu})\left(\frac{\partial}{\partial p_{3,\mu}} - \frac{\partial}{\partial p_{1,\mu}} + \frac{\partial}{\partial p_{2,\mu}} \right) \right] \left\{ \underbrace{p_{1}}_{p_{2}} \underbrace{p_{3}}_{p_{4}} \right\}$$

.

In general, *n* MIs obey a system of 1st ODE

 $\partial_z \mathbf{M}^{[d]} = \mathbb{A}(d, z) \ \mathbf{M}^{[d]}$

Two-Loop Integrals for Mu-E Scattering





Planar Integrals :: Family-1



Planar Integrals :: Family-1

Planar Integrals :: Family-2



Planar Integrals :: Family-1

Planar Integrals :: Family-2

Non-Planar Integrals





30 MIs for non-planar diagrams

Primo Schubert & P.M. (w.i.p.)



Quantum Mechanics

Schroedinger Eq'n (eps-linear Hamiltonian)

 $i\hbar \partial_t |\Psi(t)\rangle = H(\epsilon, t) |\Psi(t)\rangle$, $H(\epsilon, t) = H_0(t) + \epsilon H_1(t)$

Interaction Picture

 $H_{i,I}(t) = B^{\dagger}(t) \ H_i(t) \ B(t)$

Search Matrix Transform

$$i\hbar \partial_t B(t) = H_0(t)B(t) \qquad B(t) = e^{-\frac{i}{\hbar}\int_{t_0}^t d\tau H_0(\tau)}$$

Schroedinger Eq'n (canonical form)

 $i\hbar \partial_t |\Psi_I(t)\rangle = \epsilon H_{1,I}(t) |\Psi_I(t)\rangle,$

Magnus Expansion

System of 1st ODE

 $Y(x) = Y_0 +$

 $\partial_x Y(x) = A(x)Y(x)$, $Y(x_0) = Y_0$. A(x) non-commutative

Series :: Matrix Exponential

$$Y(x) = e^{\Omega(x,x_0)} Y(x_0) \equiv e^{\Omega(x)} Y_0,$$

Argeri, Di Vita, Mirabella, Schlenk, Schubert, Tancredi, P.M. (2014)

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$$\Omega(x) = \sum_{n=1}^{\infty} \Omega_n(x) .$$

$$\Omega(x) = \frac{1}{2} \int_{x_0}^x d\tau_1 \int_{x_0}^{\tau_1} d\tau_2 [A(\tau_1), A(\tau_2)] ,$$

$$\Omega_3(x) = \frac{1}{6} \int_{x_0}^t d\tau_1 \int_{x_0}^{\tau_1} d\tau_2 \int_{x_0}^{\tau_2} d\tau_3 [A(\tau_1), [A(\tau_2), A(\tau_3)]] + [A(\tau_3), [A(\tau_2), A(\tau_1)]] .$$

$$(generalized polylogs)$$

$$Y(x) = Y_0 + \sum_{n=1}^{\infty} Y_n(x) ,$$

$$Y_n(x) = \int_{x_0}^x d\tau_1 \dots \int_{x_0}^{\tau_{n-1}} d\tau_n A(\tau_1) A(\tau_2) \dots A(\tau_n)$$

$$BCH-formula$$

$$BCH-formula$$

$$BCH-formula$$

Argeri, Di Vita, Mirabella, Schlenk, Schubert, Tancredi, **P.M**. (2014)

Quantum Mechanics

- Time-evolution in Perturbation Theory
- ^φ perturbation parameter: ε
- ^ω Linear Hamiltonian in ε
- Unitary transform
- Schroedinger Equation in the interaction picture (ε-factorization)
- Solution: Dyson series

• Feynman Integrals

- Kinematic-evolution in Dimensional Regularization
- space-time dimensional parameter: $\varepsilon = (4-d)/2$
- $\stackrel{\scriptstyle{\smile}}{\scriptstyle{\leftarrow}}$ Linear system in ϵ
- non-Unitary Magnus transform
- System of Differential Equations in canonical form (ε-factorization) Henn (2013)
- Solution: Dyson/Magnus series

boundary term (simpler integral)

Feynman integrals can be determined from differential equations that looks like gauge transformations

 $\mathrm{e}^{\Omega(d,x)}$

Argeri, Di Vita, Mirabella, Schlenk, Schubert, Tancredi, **P.M**. (2014)

• pre-Canonical form Linear-eps Matrix $\partial_x f(\epsilon, x) = A(\epsilon, x) f(\epsilon, x) ,$ $A(\epsilon, x) = A_0(x) + \epsilon A_1(x) ,$ change of basis :: Magnus #1 $B_0(x) \equiv e^{\Omega[A_0](x,x_0)} , \qquad \partial_x B_0(x) = A_0(x) B_0(x) ,$ $f(\epsilon, x) = B_0(x) q(\epsilon, x)$, • Canonical form Henn (2013) $\partial_x g(\epsilon, x) = \epsilon \hat{A}_1(x) g(\epsilon, x)$ $\hat{A}_1(x) = B_0^{-1}(x)A_1(x)B_0(x)$ (*) Solution :: Magnus #2 (or Dyson) $g(\epsilon, x) = B_1(\epsilon, x)g_0(\epsilon)$, $B_1(\epsilon, x) = e^{\Omega[\epsilon \hat{A}_1](x, x_0)}$

Feynman integrals can be determined from differential equations that looks like gauge transformations

• Kinematic variables

$$-\frac{s}{m^2} = x, \quad -\frac{t}{m^2} = \frac{(1-y)^2}{y}$$

• eps-linear basis

$$\partial_x f(x, y, \epsilon) = \left(A_{10}(x, y) + \epsilon A_{11}(x, y) \right) f(x, y, \epsilon)$$
$$\partial_y f(x, y, \epsilon) = \left(A_{20}(x, y) + \epsilon A_{21}(x, y) \right) f(x, y, \epsilon)$$

canonical form: Magnus #1

$$\partial_x g(x, y, \epsilon) = \epsilon \hat{A}_1(x, y) g(x, y, \epsilon)$$
$$\partial_y g(x, y, \epsilon) = \epsilon \hat{A}_2(x, y) g(x, y, \epsilon)$$



Canonical systems and Iterated Integrals Henn (2013)

- Canonical system of DE $d\mathbf{I} = \epsilon \, d\mathbb{A} \, \mathbf{I}$ $dA = \sum_{i=1}^{n} \mathbb{M}_i \, d\log \eta_i$ $\vec{x} = (x_1, x_2, \dots, x_n)$
- Solution as path-ordered exponential

$$\mathbf{I}(\epsilon, \vec{x}) = \mathcal{P} \exp\left\{\epsilon \int_{\gamma} dA\right\} \mathbf{I}(\epsilon, \vec{x}_0), \qquad \qquad \mathcal{P} \exp\left\{\epsilon \int_{\gamma} dA\right\} = \mathbb{1} + \epsilon \int_{\gamma} dA + \epsilon^2 \int_{\gamma} dA \, dA + \epsilon^3 \int_{\gamma} dA \, dA \, dA \, dA \dots,$$

• Path invariance $\gamma : [0,1] \ni t \mapsto \gamma(t) = (\gamma_1(t), \gamma_2(t), \dots, \gamma_n(t)), \qquad \gamma(0) = \vec{x}_0 \text{ and } \gamma(1) = \vec{x}.$

Taylor expansion and Dyson/Magnus series

$$= \mathbf{I}^{(0)}(\vec{x}) + \epsilon \mathbf{I}^{(1)}(\vec{x}) + \epsilon^{2} \mathbf{I}^{(2)}(\vec{x}) + \dots$$

$$I^{(0)}(\vec{x}) = \mathbf{I}^{(0)}(\vec{x}_{0}),$$

$$\mathbf{I}^{(1)}(\vec{x}) = \mathbf{I}^{(1)}(\vec{x}_{0}) + \int_{\gamma} dA \, \mathbf{I}^{(0)}(\vec{x}_{0}),$$

$$\mathbf{I}^{(2)}(\vec{x}) = \mathbf{I}^{(2)}(\vec{x}_{0}) + \int_{\gamma} dA \, \mathbf{I}^{(1)}(\vec{x}_{0}) + \int_{\gamma} dA \, dA \, \mathbf{I}^{(0)}(\vec{x}_{0}),$$

$$\mathbf{I}^{(3)}(\vec{x}) = \mathbf{I}^{(3)}(\vec{x}_{0}) + \int_{\gamma} dA \, \mathbf{I}^{(2)}(\vec{x}_{0}), + \int_{\gamma} dA \, dA \, \mathbf{I}^{(1)}(\vec{x}_{0}) + \int_{\gamma} dA \, dA \, \mathbf{I}^{(0)}(\vec{x}_{0}),$$

$$\mathbf{I}^{(4)}(\vec{x}) = \mathbf{I}^{(4)}(\vec{x}_{0}) + \int_{\gamma} dA \, \mathbf{I}^{(3)}(\vec{x}_{0}) + \int_{\gamma} dA \, dA \, \mathbf{I}^{(1)}(\vec{x}_{0}) + \int_{\gamma} dA \, dA \, \mathbf{I}^{(1)}(\vec{x}_{0}) + \int_{\gamma} dA \, dA \, \mathbf{I}^{(0)}(\vec{x}_{0}).$$

Chen's Iterated integral

 $I(\epsilon, \vec{x})$

$$\int_{\gamma} \underbrace{dA \dots dA}_{k \text{ times}}, \longleftrightarrow \quad \mathcal{C}_{i_k, \dots, i_1}^{[\gamma]} \equiv \int_{\gamma} d\log \eta_{i_k} \dots d\log \eta_{i_1} \equiv \int_{0 \le t_1 \le \dots \le t_k \le 1} g_{i_k}^{\gamma}(t_k) \dots g_{i_1}^{\gamma}(t_1) dt_1 \dots dt_k, \qquad g_i^{\gamma}(t) = \frac{d}{dt} \log \eta_i(\gamma(t))$$

• a special case :: Goncharov's polylogs

$$G(\vec{w}_n; x) \equiv G(w_1, \vec{w}_{n-1}; x) \equiv \int_0^x dt \frac{1}{t - w_1} G(\vec{w}_{n-1}; t), \qquad G(\vec{0}_n; x) \equiv \frac{1}{n!} \log^n(x), \qquad \partial_x G(\vec{w}_n; x) = \partial_x G(w_1, \vec{w}_{n-1}; x) = \frac{1}{x - w_1} G(\vec{w}_{n-1}; x).$$





124 x¹ :: B^pund²4^p corditions



 $p_3 p_2$

 $p_4 p_1$

p₃

p₄

massless leg

regular @ s = 0
\$
\$
trivial conditions !



p₃ **p**₂

 $p_4 p_1$



p₃ p₂

 $p_4 p_1$



PAR1 :: Brandar y conditions



Fregular @ s = 0

trivial conditions !



 $p_4 p_1$

 $p_3 p_2 (k_2 + p_1)^2 p_3 p_2$

*p*₄**p*₁

 $p_3 p_2$

 $p_4 p_1$

vegular @ $p_1^2 = 0$ **vegular** @ $p_1^2 = 0$ **vegular** b.c. conditions !



 $\int I(0) =$

p4

$$P_{A} P_{A} P_{A$$

$$\begin{array}{l} p_{A} p_{1} & p_{A} p_{A} & p_{A} &$$

Summary ...

Feasibility of Two-Loop QED Corrections *analytically*

Master Integrals via Differential Equations + Magnus Series

Amplitude decomposition via Adaptive Integrand Decomposition (AID) *Mu-e scattering :: a first example of 2-loop automation for massive amplitudes*

...Outlook

Building the 2-loop amplitude (Form Factors and AID)
 Analytic continuation and Numerical Evaluation of 2-loop MIs
 The 1-Loop amplitude and 2-loop renorm. counterterms (GoSam, AID)
 Implementing a Subtraction Scheme for NNLO (*hinc sunt leones*)
 MonteCarlo Integration >> see Piccinini's talk

Progress in Mu-e Scattering @ 2-loop QED
 => benefit for e+ e- -> Mu+ Mu- @ 2-loop QED
 => benefit for p p -> t T @ 2-loop QCD



Muon-electron scattering: Theory kickoff workshop 4-5 September 2017 Padova

https://agenda.infn.it/internalPage.py?pageId=0&confId=13774



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