

Higher-rank discrete symmetries: Octahedral shapes in *sdg*-IBM

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Motivation

Quadrupole and hexadecapole shapes

The *sdg*-IBM

Octahedral shapes in *sdg*-IBM

Conclusion

Motivation

How to realize higher-rank discrete symmetries in
the context of an algebraic model.

Derive the necessary conditions on the parameters
of the Hamiltonian.

Example: Octahedral shape in *sdg*-IBM.

4-pole and 16-pole shapes

Nuclear surface with quadrupole and hexadecapole deformation:

$$\begin{aligned} R(\theta, \phi) = & R_0 \left(1 + a_{20} Y_{20}(\theta, \phi) + a_{22} [Y_{2-2}(\theta, \phi) + Y_{2+2}(\theta, \phi)] \right. \\ & + a_{40} Y_{40}(\theta, \phi) + a_{42} [Y_{4-2}(\theta, \phi) + Y_{4+2}(\theta, \phi)] \\ & \left. + a_{44} [Y_{4-4}(\theta, \phi) + Y_{4+4}(\theta, \phi)] \right) \end{aligned}$$

4-pole and 16-pole shapes

4-pole variables ($\beta_2 \geq 0, 0 \leq \gamma_2 \leq \pi/3$):

$$a_{20} = \beta_2 \cos \gamma_2$$

$$a_{22} = \beta_2 \sin \gamma_2$$

16-pole variables ($\beta_4 \geq 0, 0 \leq \gamma_4 \leq \pi/3, 0 \leq \delta_4 \leq \pi$):

$$a_{40} = \beta_4 \left(\sqrt{\frac{7}{12}} \cos \delta_4 + \sqrt{\frac{5}{12}} \sin \delta_4 \cos \gamma_4 \right)$$

$$a_{42} = -\sqrt{\frac{1}{2}} \beta_4 \sin \delta_4 \sin \gamma_4$$

$$a_{44} = \beta_4 \left(\sqrt{\frac{5}{24}} \cos \delta_4 - \sqrt{\frac{7}{24}} \sin \delta_4 \cos \gamma_4 \right)$$

Octahedral symmetry

A shape with octahedral symmetry is obtained for

$$a_{20} = a_{22} = 0, \quad a_{42} = 0, \quad a_{44} / a_{40} = \pm\sqrt{5/14}$$

Three cases in the alternative parameterisation:

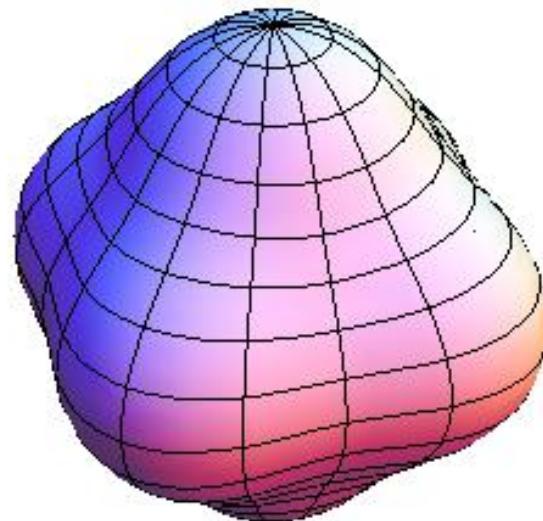
(a) octahedron : $\beta_2 = 0, \gamma_2$ free, $\beta_4 > 0, \gamma_4$ free, $\delta_4 = 0$

(b) cube : $\beta_2 = 0, \gamma_2$ free, $\beta_4 > 0, \gamma_4$ free, $\delta_4 = \pi$

(c) octahedron : $\beta_2 = 0, \gamma_2$ free, $\beta_4 > 0, \gamma_4 = 0, \delta_4 = \arccos\left(\frac{1}{6}\right)$

Octahedral symmetry

(a) octahedron

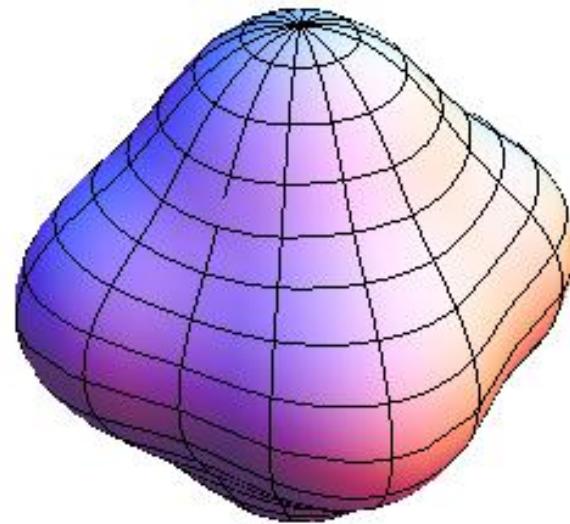


$$a_{40} = +0.2, \quad a_{42} = 0, \quad a_{44} / a_{40} = +\sqrt{5/14}$$

$$\beta_4 = 0.2, \quad \delta_4 = 0$$

Octahedral symmetry

(c) octahedron

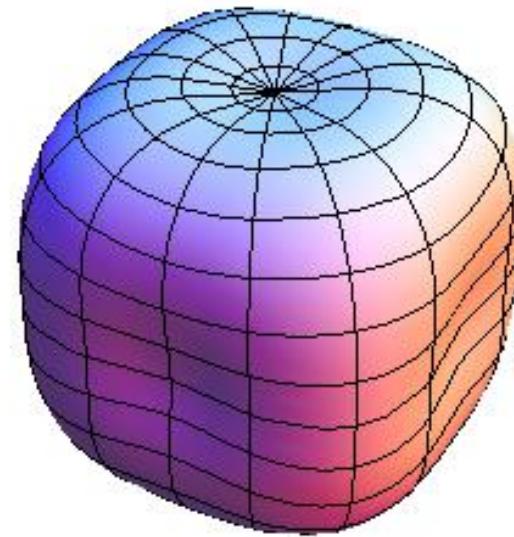


$$a_{40} = +0.2, \quad a_{42} = 0, \quad a_{44} / a_{40} = -\sqrt{5/14}$$

$$\beta_4 = 0.2, \quad \delta_4 = \arccos(1/6)$$

Octahedral symmetry

(b) cube

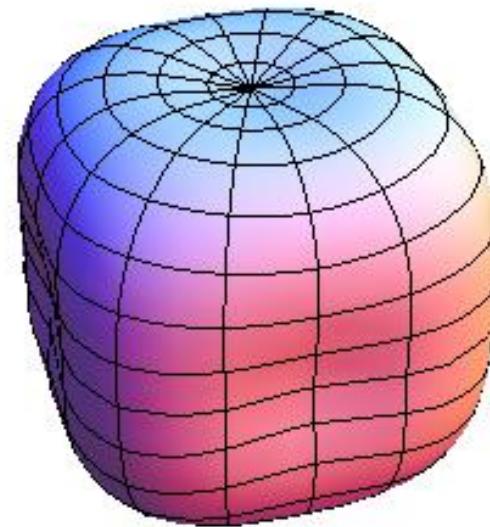


$$a_{40} = -0.2, \quad a_{42} = 0, \quad a_{44} / a_{40} = +\sqrt{5/14}$$

$$\beta_4 = 0.2, \quad \delta_4 = \pi$$

Octahedral symmetry

(d) cube



$$a_{40} = -0.2, \quad a_{42} = 0, \quad a_{44} / a_{40} = -\sqrt{5/14}$$

not possible

Hamiltonian of the sdg -IBM

Hamiltonian with up to 2-body interactions:

$$\hat{H}_{sdg} = \varepsilon_s \hat{n}_s + \varepsilon_d \hat{n}_d + \varepsilon_g \hat{n}_g + \sum_{l_1 \leq l_2, l'_1 \leq l'_2, L} \frac{(-)^L v_{l_1 l_2 l'_1 l'_2}^L}{\sqrt{(1 + \delta_{l_1 l_2})(1 + \delta_{l'_1 l'_2})}} (b_{l_1}^+ \times b_{l_2}^+)^{(L)} \cdot (\tilde{b}_{l'_1} \times \tilde{b}_{l'_2})^{(L)}$$

H_{sdg} contains 3 single-boson energies and 32 two-body boson-boson interactions

$$\varepsilon_l = \langle l | \hat{H}_{sdg} | l \rangle$$

$$v_{l_1 l_2 l'_1 l'_2}^L = \langle l_1 l_2; L | \hat{H}_{sdg} | l'_1 l'_2; L \rangle$$

Classical limit of the *sdg*-IBM

Two quadrupole (β_2, γ_2) and three hexadecapole $(\beta_4, \gamma_4, \delta_4)$ shape parameters.

Coherent state:

$$\begin{aligned} |N; \beta_2, \gamma_2, \beta_4, \gamma_4, \delta_4\rangle \propto \\ \left\{ s^+ + \beta_2 \left[\cos \gamma_2 d_0^+ + \sqrt{\frac{1}{2}} \sin \gamma_2 (d_{-2}^+ + d_{+2}^+) \right] \right. \\ + \beta_4 \left[\left(\sqrt{\frac{7}{12}} \cos \delta_4 + \sqrt{\frac{5}{12}} \sin \delta_4 \cos \gamma_4 \right) g_0^+ \right. \\ - \sqrt{\frac{1}{2}} \sin \delta_4 \sin \gamma_4 (g_{-2}^+ + g_{+2}^+) \\ \left. \left. + \left(\sqrt{\frac{5}{24}} \cos \delta_4 - \sqrt{\frac{7}{24}} \sin \delta_4 \cos \gamma_4 \right) (g_{-4}^+ + g_{+4}^+) \right] \right\}^N |\text{o}\rangle \end{aligned}$$

Classical limit of the sdg -IBM

Classical limit of H_{sdg}

$$\langle \hat{H}_{sdg} \rangle = E(\beta_2, \gamma_2, \beta_4, \gamma_4, \delta_4; c)$$

$$= N \frac{\varepsilon_s + \varepsilon_d \beta_2^2 + \varepsilon_g \beta_4^2}{1 + \beta_2^2 + \beta_4^2}$$

$$+ \frac{N(N-1)}{(1 + \beta_2^2 + \beta_4^2)^2} \sum_{kl} \beta_2^k \beta_4^l \left[c_{kl} + \sum_{ij} c_{kl}^{ij} \cos(i\gamma_2 + j\gamma_4) \phi_{kl}^{ij}(\delta_4) \right]$$

in terms of trigonometric functions $\Phi(\delta_4)$ and with
22 coefficients c depending on interactions U .

Catastrophe analysis

Energy surface has extremum at p^* if

$$\frac{\partial E}{\partial \beta_2} \Big|_{p^*} = \frac{\partial E}{\partial \gamma_2} \Big|_{p^*} = \frac{\partial E}{\partial \beta_4} \Big|_{p^*} = \frac{\partial E}{\partial \gamma_4} \Big|_{p^*} = \frac{\partial E}{\partial \delta_4} \Big|_{p^*} = 0$$

Shape of extremum is octahedral [(a), upper sign]
or cubic [(b), lower sign] if

$$\begin{aligned} \beta_4 & \left[-4c'_{00} + 2c'_{02} \pm 21c^{00}_{03}\beta_4 \right. \\ & \left. - (2c'_{02} - 4c'_{04} - 76c^{00}_{04})\beta_4^2 \mp 7c^{00}_{03}\beta_4^3 \right] = 0 \end{aligned}$$

A real solution $\beta_4 > 0$ always exists except for

$$c^{00}_{03} = \frac{2}{\sqrt{429}} v_{sg\cdot gg}^4 = 0$$

Catastrophe analysis

Shape of extremum is octahedral (c) if

$$\beta_4^2 \left[-3(\sqrt{35}c_{12}^{1-1} + 35c_{12}^{12}) + (\sqrt{35}c_{13}^{1-1} - 5c_{13}^{12})\beta_4 \right] = 0$$

$$\begin{aligned} & \beta_4 \left[324(-2c_{00}' + c_{02}') + 63(4c_{03}^{00} + 5\sqrt{35}c_{03}^{03}) \right. \\ & + 4(-81c_{02}' + 162c_{04}' + 1853c_{04}^{00} + 35\sqrt{35}c_{04}^{03})\beta_4^2 \\ & \left. - 21(4c_{03}^{00} + 5\sqrt{35}c_{03}^{03})\beta_4^3 \right] = 0 \end{aligned}$$

$$\beta_4^3 \left[9(-2\sqrt{35}c_{03}^{00} + 7c_{03}^{03}) + 224(\sqrt{35}c_{04}^{00} - c_{04}^{03})\beta_4 \right] = 0$$

First and third conditions are identically satisfied;
second reduces to previous one (upper sign).

Catastrophe analysis

Sufficient conditions that the *extremum* with octahedral shape is a *minimum*

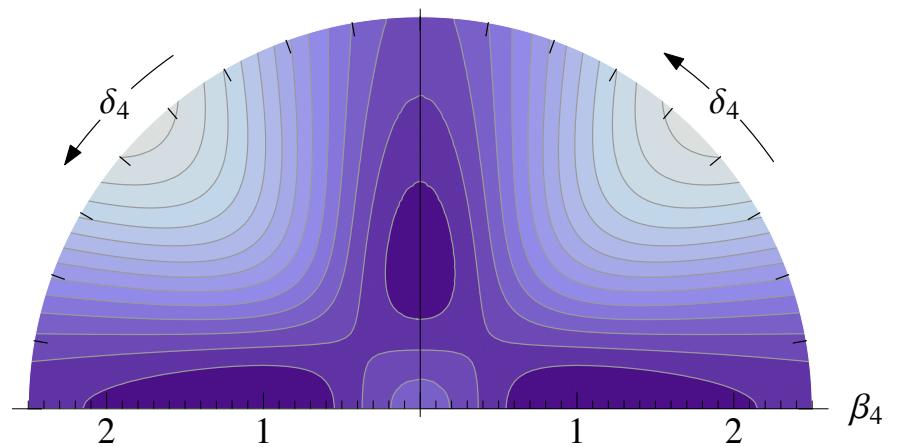
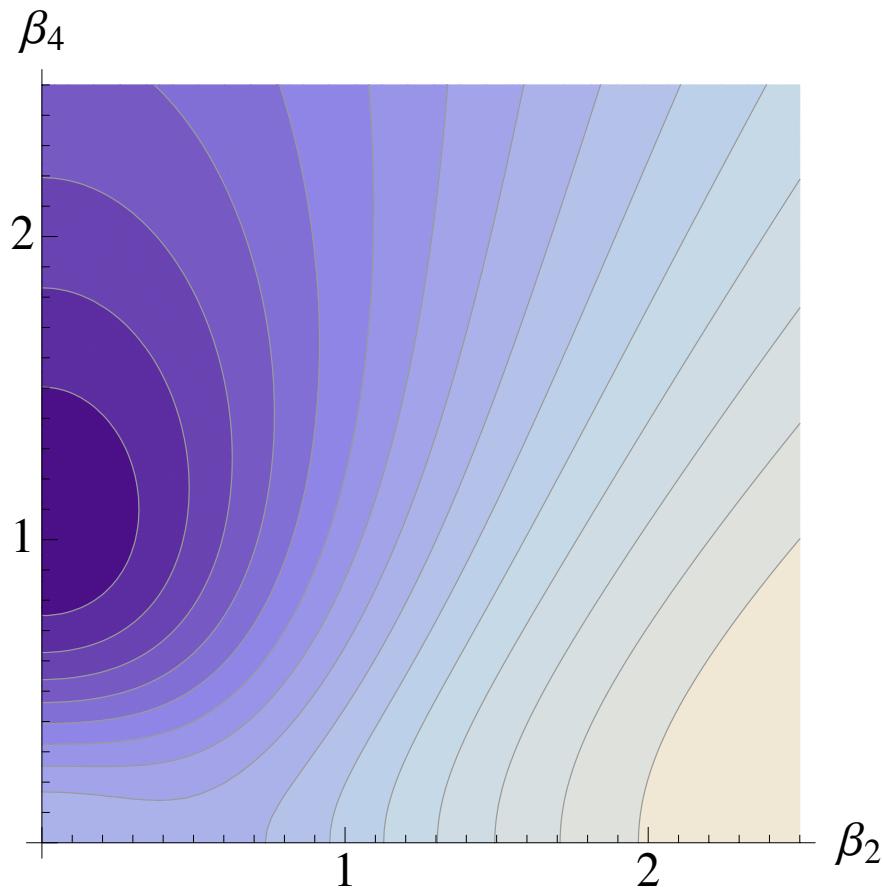
$$c_{02} + \frac{\varepsilon_g}{N-1} \leq 2c_{00} \leq c_{20} + \frac{\varepsilon_d}{N-1}$$

$$c_{02} \leq 2c_{04} + 38c_{04}^{00} + \frac{\varepsilon_g}{N-1} \leq c_{22} + \frac{\varepsilon_d}{N-1}$$

Reduce to conditions on interactions, e.g.

$$\frac{1}{3}\boldsymbol{\upsilon}_{ss\cdot gg}^0 + \boldsymbol{\upsilon}_{sgsg}^4 + \frac{\varepsilon_g}{N-1} \leq \boldsymbol{\upsilon}_{ssss}^0 \leq \sqrt{\frac{1}{5}}\boldsymbol{\upsilon}_{ss\cdot dd}^0 + \boldsymbol{\upsilon}_{sdsd}^2 + \frac{\varepsilon_d}{N-1}$$

Example



Conclusion

A procedure to treat higher-rank discrete symmetries in an algebraic framework.

Analysis completed for octahedral symmetry in the *sdg*-IBM:

- *s-g* mixing strong enough compared to $\varepsilon_g - \varepsilon_s$ to develop a $\beta_4 > 0$ minimum;
- *s-d* mixing weak enough compared to $\varepsilon_d - \varepsilon_s$ not to develop a $\beta_2 > 0$ minimum.

Analysis of tetrahedral case can be done similarly.