

Quantum Phase Transitions within the Shell Model

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**9th international workshop: Quantum Phase Transitions
in Nuclei and Many-body Systems , 05/22/2018**

- Quantum phase transition induced by changing the strength of the two-body matrix elements of the shell model Hamiltonian
- Quantum phase transition demonstrates itself in many nuclear observables, including level densities

Quantum Phase Transition in atomic nuclei

Algebraic models are used, possible phases of a system obtained by considering possible breakings of the algebra g into its subalgebras, each subalgebra having its own **dynamical symmetry**, structural shape.

$$g \supset g_1 \supset \dots$$

$$g \supset g_2 \supset \dots$$

...

$$g \supset g_\nu \supset \dots$$

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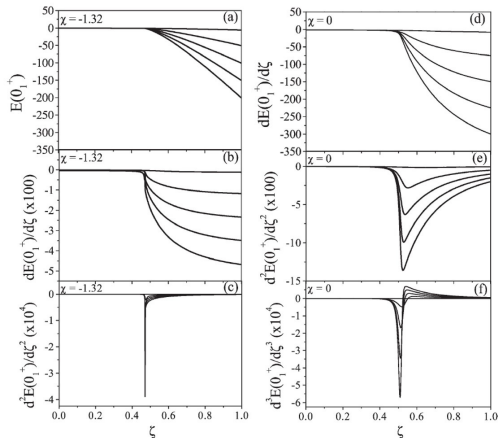
$$H = (1 - \lambda)H_1 + \lambda H_2,$$

$$\lambda = 0 \rightarrow H_1, \lambda = 1 \rightarrow H_2$$

Order of Quantum Phase Transition

Ehrenfest criterion:

ground state energy functional and its derivatives, n th order if its n derivative with respect to λ is discontinuous



Lincoln D. Carr, *Understanding QPTs* ($N_B = 10, 100, 200, 300, 400$)

Shell model Hamiltonian

Two-Body Matrix Elements:

The interaction scatters 2 nucleons into a different pair of orbits compared to the pair of orbits before the scattering

fall into 3 categories

V_1 : one particle transfer matrix elements,

V_2 : the rest matrix elements

V_1 : deformation

H. Hori *et al.*, PRC **81**, 034396 (2010)

sd shell: 63 matrix elements

pf shell: 195 matrix elements

one-particle transfer

$$V_1: \left\langle \begin{array}{c} l_3 j_3 \\ l_2 j_2 \\ l_1 j_1 \end{array} \left| V \right| \begin{array}{c} l_3 j_3 \\ l_2 j_2 \\ l_1 j_1 \end{array} \right\rangle_{JT}$$

two-particle transfer

$$\left\langle \begin{array}{c} l_3 j_3 \\ l_2 j_2 \\ l_1 j_1 \end{array} \left| V \right| \begin{array}{c} l_3 j_3 \\ l_2 j_2 \\ l_1 j_1 \end{array} \right\rangle_{JT}$$

V_2 :

no-particle transfer

$$\left\langle \begin{array}{c} l_3 j_3 \\ l_2 j_2 \\ l_1 j_1 \end{array} \left| V \right| \begin{array}{c} l_3 j_3 \\ l_2 j_2 \\ l_1 j_1 \end{array} \right\rangle_{JT}$$

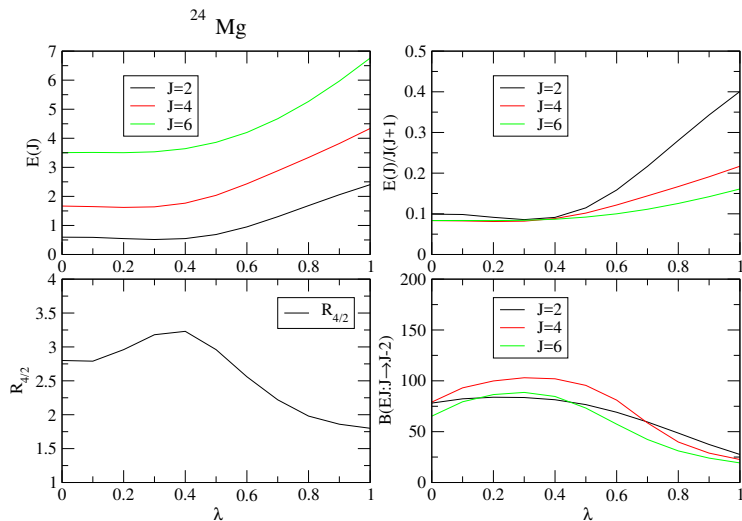
Quantum Phase Transition

$$H = (1 - \lambda)V_1 + \lambda V_2$$

- ✓ Varying the strength of the V_1 and V_2 shell model matrix elements in the sd and pf shells using a control parameter λ going from 0 to 1

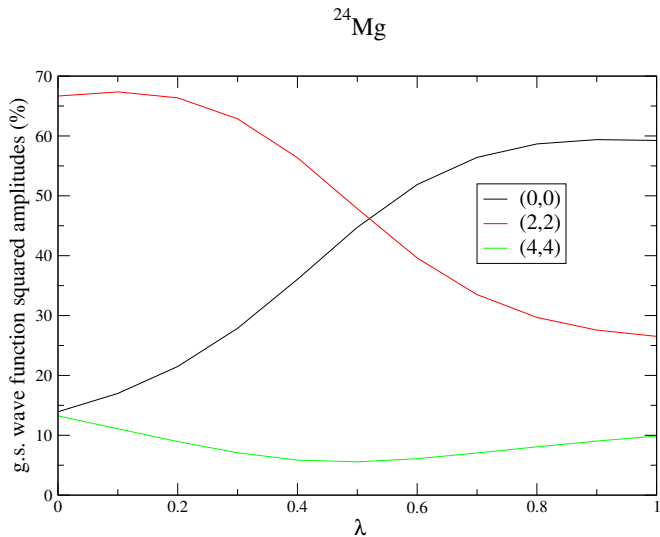
²⁴ Mg	2_1^+	4_1^+	6_1^+	$R_{4/2}$	$Q(2_1^+)$	$B(E2; 2_1^+ \rightarrow 0_1^+)$	$B(E2; 4_1^+ \rightarrow 2_1^+)$	$B(E2; 6_1^+ \rightarrow 4_1^+)$
0.0	0.596	1.667	3.507	2.80	-16.32	7.81E+01	7.91E+01	6.53E+01
0.1	0.590	1.649	3.512	2.79	-18.02	8.23E+01	9.31E+01	7.94E+01
0.2	0.548	1.620	3.504	2.96	-18.65	8.39E+01	9.99E+01	8.65E+01
0.3	0.515	1.640	3.533	3.18	-18.87	8.36E+01	1.03E+02	8.86E+01
0.4	0.547	1.766	3.642	3.23	-18.68	8.14E+01	1.02E+02	8.46E+01
0.5	0.688	2.036	3.860	2.96	-17.97	7.66E+01	9.55E+01	7.32E+01
...	0.951	2.433	4.201	2.56	-16.65	6.91E+01	8.10E+01	5.74E+01
1.0	2.404	4.337	6.761	1.80	-7.29	2.76E+01	2.26E+01	1.93E+01

Signs of a QPT in even-even nuclei

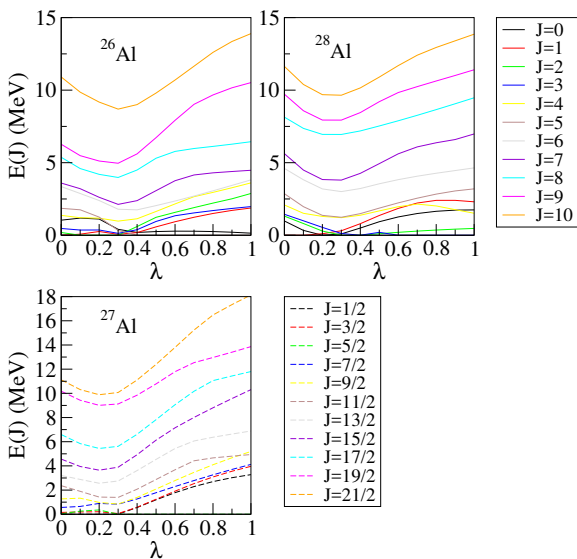


S. Karampagia, V. Zelevinsky, PRC **94** 014321 (2016)

Angular momentum wavefunction decomposition

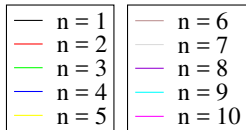
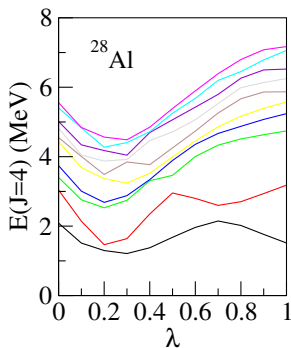
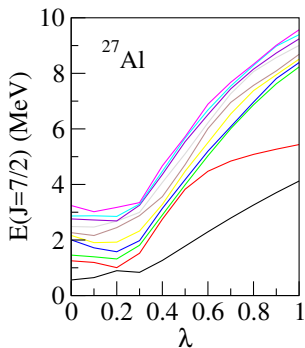
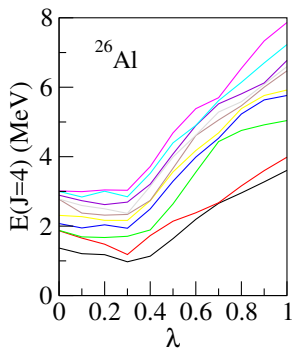


Signs of a QPT in odd–even, odd–odd nuclei

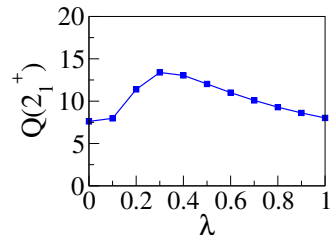
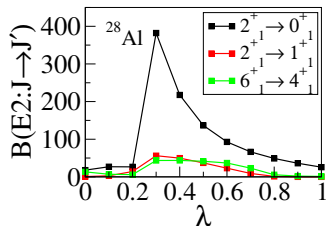
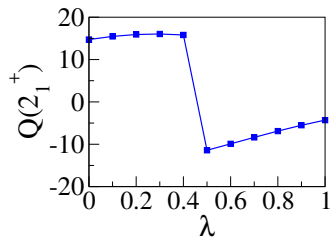
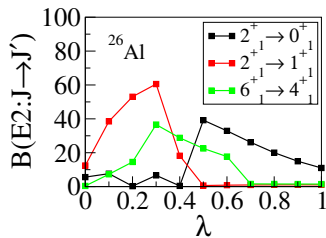


S. Karampagia, A. Renzaglia, V. Zelevinsky, NPA **962** 46 (2017)

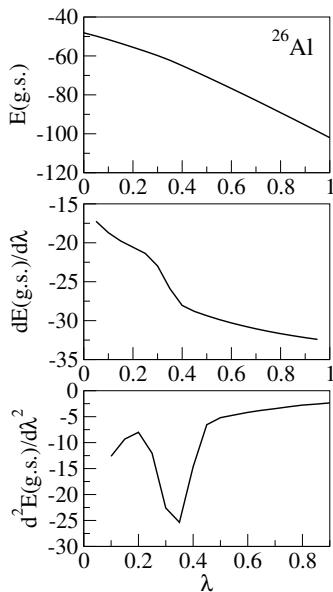
Excited states in odd–even, odd–odd nuclei



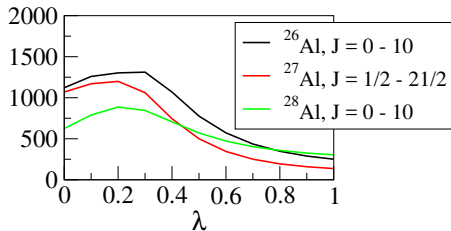
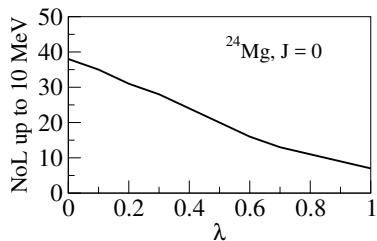
B(E2)s, Q(2₁⁺) in odd–even, odd–odd nuclei



Order of QPT



Level Densities



Conclusions

- Found a QPT induced by changing the strength of the single particle transfer matrix elements with respect to the rest matrix elements
- QPT in even-even (rotational-spherical), odd-even, odd-odd nuclei
- Single particle transfer matrix elements carriers of deformation, provide rotational observables, strong transition probabilities, quadrupole moments and enhanced level densities
- Even in a small system collective motion, phase transitions, thermodynamics are present

Additional Slides

Definition of Quantum Phase Transition (QPT)

Quantum Phase Transitions (QPT) refer to phase transitions, which occur at zero temperature ($T = 0$), when a non-thermal parameter, λ , is changed (pressure, magnetic field). This parameter(s) is called **control parameter** and the value for which the phase transition takes place is the critical value, $\lambda = \lambda_c$, for a system described by

$$H(\lambda) = \epsilon(H_1 + \lambda H_2)$$

Ising Model in a transverse field (ferromagnetic \rightarrow paramagnetic)

$$H = -J \sum_{ij} S_i^z S_j^z - h \sum_i S_i^x$$

Signs of Collective Enhancement in even–even nuclei

Collective enhancement: Noticeable change of the low lying nuclear level density due to collective effects

Many cases with signs of collective enhancement, selected just a few...

J=0 (up to 10 MeV)	Deformed		Spherical		Deformed	Spherical
	$R_{4/2}$	λ	$R_{4/2}$	λ	NoL	NoL
^{28}Si	3.01	0.2	2.15	0.5	45	15
^{24}Mg	3.23	0.4	1.98	0.8	24	11
^{52}Fe	3.10	0.2	2.25	1.0	412	30

Calculation of level density—Moments Method

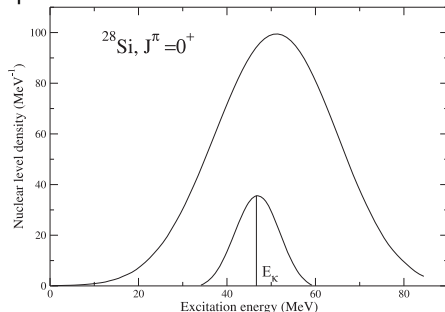
- Many different models (Fermi Gas, Hartree-Fock-Bogolyubov model, ...)
- Shell model, Moments Method (Statistical description)

$$\rho(E; \alpha) = \sum_p D_{\alpha p} G(E - E_{\alpha p} + E_{g.s.}; \sigma_{\alpha p})$$

$D_{\alpha p}$ – dimension of the class of states

α – quantum numbers of the class of states

p – partition, e.g. 6 particles in sd shell make 15 partitions



p	$d_{5/2}$	$s_{1/2}$	$d_{3/2}$
1	6	0	0
2	5	1	0
3	5	0	1
4	4	2	0
...
15	0	2	4

Moments of H for each partition p :

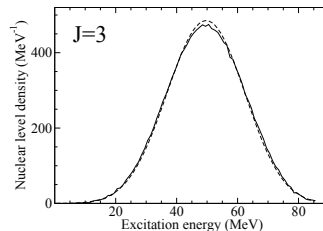
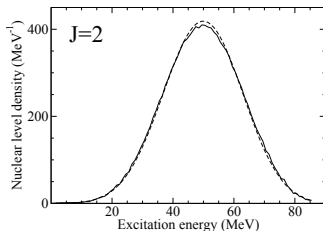
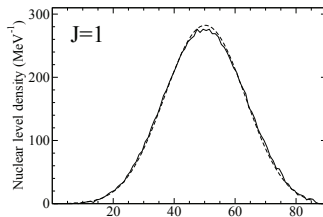
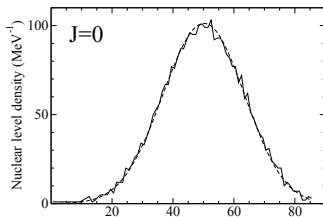
$$E_{\alpha p} = \frac{1}{D_{\alpha p}} \text{Tr}^{\alpha p} H$$

$$\sigma_{\alpha p}^2 = \frac{1}{D_{\alpha p}} \text{Tr}^{\alpha p} H^2 - E_{\alpha p}^2$$

R. Sen'kov, V. Zelevinsky, PRC **93** 064304

Moments method vs exact shell model

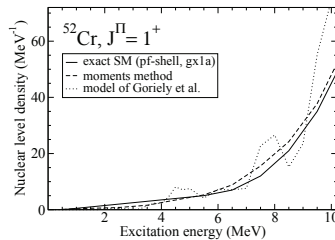
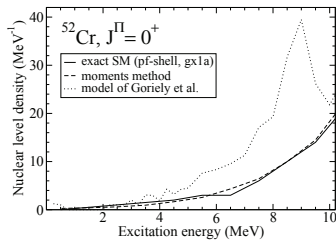
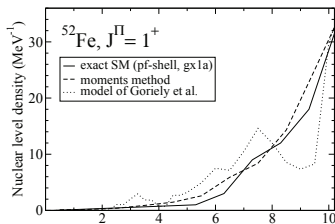
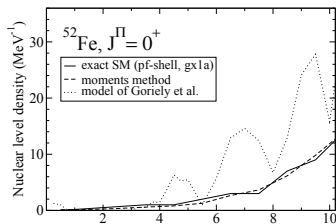
^{28}Si , $J = 0, 1, 2, 3$, SD shell, USDB interaction



R. Sen'kov, V. Zelevinsky, arXiv:1508.03683

Moments method vs exact shell model

^{52}Fe , ^{52}Cr , $J = 0, 1$, PF shell, GX1A interaction

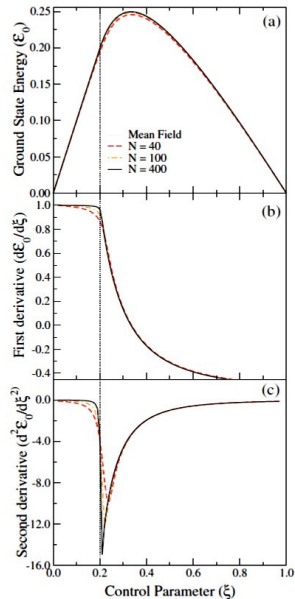


Order of Quantum Phase Transition

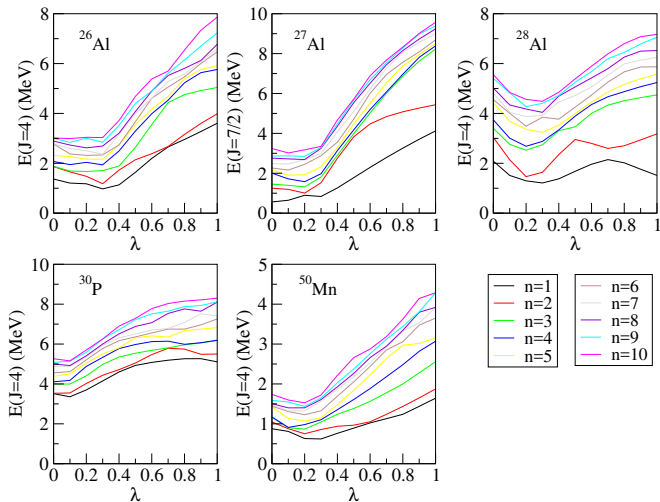
Ehrenfest criterion: ground state energy functional and its derivatives, n th order if its n derivative with respect to λ is discontinuous

image: second order QPT

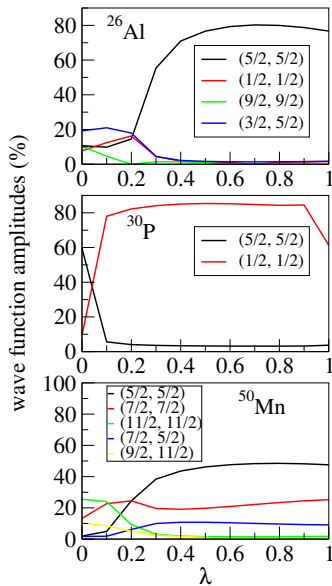
F. Pérex – Bernal *et al.*, PRA **77**, 032115 (2008)



Excited states in odd–even, odd–odd nuclei



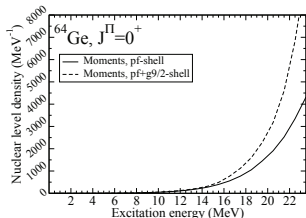
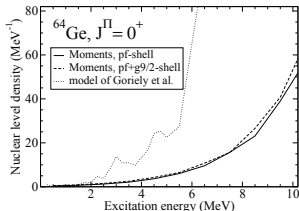
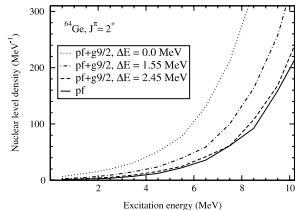
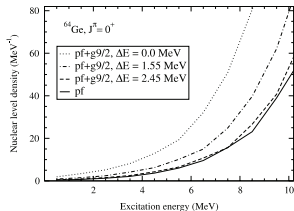
Angular momentum wavefunction decomposition, 1_1^+



Expansion of model space

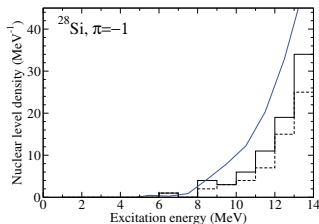
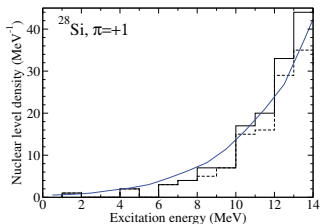
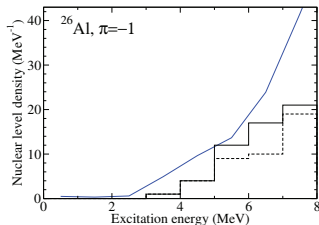
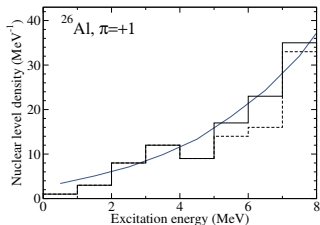
Ways of finding the g.s. energy in larger spaces?

- Exponential method, M. Horoi, A. Volya, V. Zelevinsky, PRL82 8
- Adjusting the level density to the previous space



Moments method vs experimental data

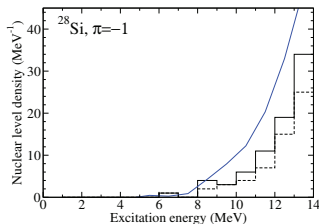
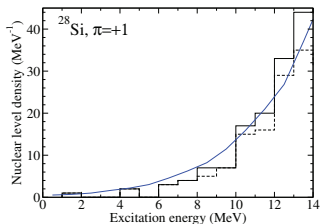
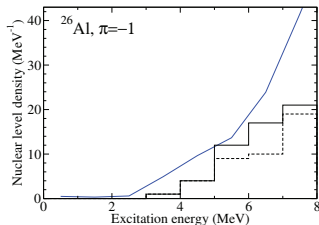
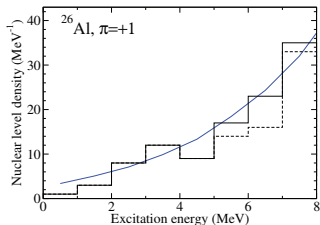
^{26}Al , ^{28}Si all J , s-p-sd-pf model space, WBT interaction



R. Sen'kov, V. Zelevinsky, arXiv:1508.03683

Moments method vs experimental data

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