Topological Superfluid Phase with Repulsive Fermionic Atoms

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QPTn-9 — May 24th 2018

Is Landau theory the only possible framework for (Quantum or Thermal) Phase Transitions?

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"There is life beyond the Landau paradigm"

— Anonymous

- Anonymous

Landau versus Topological orders

Landau orders

Topological orders

(Local) Order parameter

String (non-local) Order /
Topological Invariants

Broken (global) Symmetry

Robustness against local perturbations/
Bulk-boundary correspondence

The Nobel Prize in Physics 2016

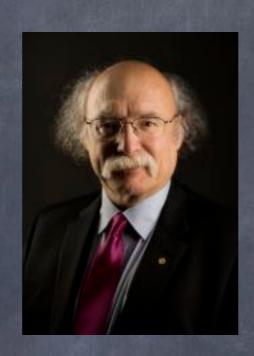


The Nobel Prize in Physics 2016





David J. Thouless Prize share: 1/2



F. Duncan M. Haldane Prize share: 1/4



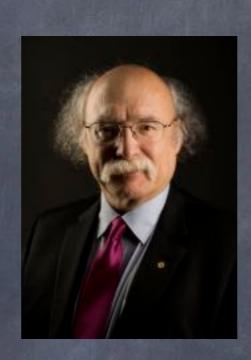
J. Michael Kosterlitz Prize share: 1/4

The Nobel Prize in Physics 2016





David J. Thouless
Prize share: 1/2

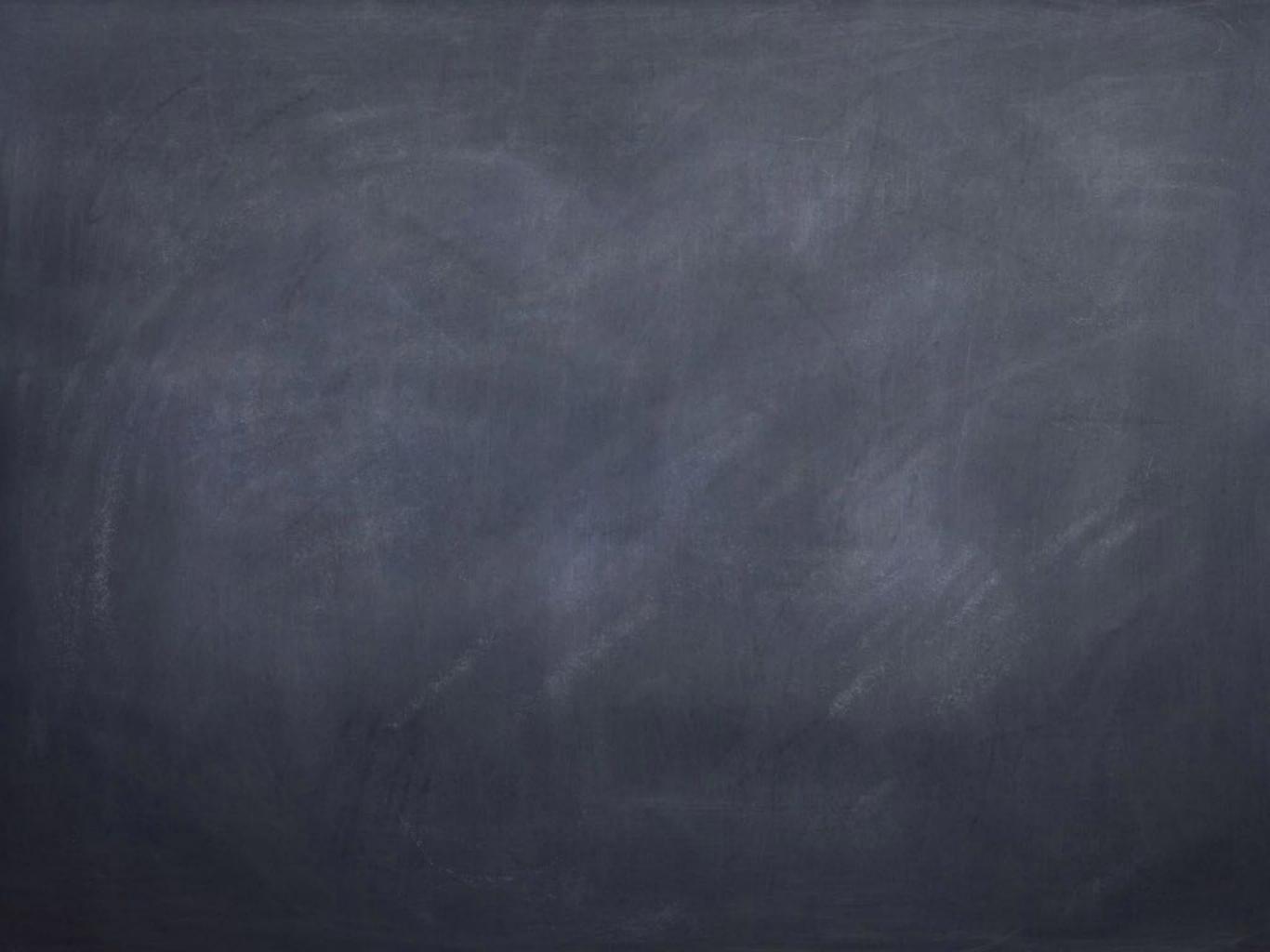


F. Duncan M. Haldane Prize share: 1/4



J. Michael Kosterlitz Prize share: 1/4

"For theoretical discoveries of topological phase transitions and topological phases of matter"



- Most of the work on Topological Quantum Matter is based on the Mean-Field quasi-particle picture
 - Paired Superfluids are particle-number non-conserving
 - Zero-energy modes are Majorana fermions (MFs) by default
 - MFs are blueprints for topological quantum computation

- Most of the work on Topological Quantum Matter is based on the Mean-Field quasi-particle picture
 - Paired Superfluids are particle-number non-conserving
 - Zero-energy modes are Majorana fermions (MFs) by default
 - MFs are blueprints for topological quantum computation

- But real "closed systems" are interacting particle-number conserving (pnc) systems
 - What is a topological superfluid in a pnc system?
 - What is a Majorana fermion in a pnc system?
 - How one detects Majorana fermions?
 - Can one braid MFs in pnc systems?

PRL 113, 267002 (2014); Ann. Phys. 372, 357 (2016); PRB 95, 201114(RC) (2017)

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What is a Particle-Conserving Topological Superfluid?

Characterization of Topological Superfluidity in generic interacting many-body systems: Fermion Parity Switches

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Characterization of Topological Superfluidity in generic interacting many-body systems: Fermion Parity Switches

$$\lim_{N,V\to\infty}\frac{N}{V}=\rho=\mathrm{constant}$$

$$E \hspace{-0.2cm} \uparrow \hspace{-0.2cm} \text{broken } U(1) \hspace{-0.2cm} \downarrow \hspace{-0.2cm} E \hspace{-0.2cm} \uparrow \hspace{-0.2cm} \text{broken } U(1) \wedge \mathbb{Z}_2$$

$$N-1 \hspace{-0.2cm} N+1 \hspace{-0.2cm} N+3 \hspace{-0.2cm} \downarrow N-1 \hspace{-0.2cm} N+1 \hspace{-0.2cm} N+3 \hspace{-0.2cm} \downarrow N-2 \hspace{-0.2cm} N \hspace{-0.2cm} N+2 \hspace{-0.2cm} \uparrow \hspace{-0.2cm} \hat{N}$$

$$N-2 \hspace{-0.2cm} N \hspace{-0.2cm} N+2 \hspace{-0.2cm} Topo-Trivial Topological$$

PRL 113, 267002 (2014); Ann. Phys. 372, 357 (2016); PRB 95, 201114(RC) (2017)

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What is the Fate of Majoranas Beyond Mean-Field?

Meaning to emergent many-body Majorana zero-energy modes

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What is the Fate of Majoranas Beyond Mean-Field?

- Meaning to emergent many-body Majorana zero-energy modes
 - Coherent superpositions of states with different # of particles
 - lacksquare $\Gamma_{1,2}$ modes anti-commute with fermionic parity
 - Non-number conserving in number conserving systems

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Can one prepare/manipulate/braid coherent superpositions of states with a different number of particles?

INTERVIEW 8 May 2013

Nothing to see: The man who made a Majorana particle

Majorana in the News

By Lisa Grossman

Physicist **Leo Kouwenhoven** ended a 75-year hunt for the tricky Majorana fermion – a particle that is its own antiparticle – by creating one on a chip

The New York Times

http://nyti.ms/1mfRwAV

TECHNOLOGY

Microsoft Makes Bet Quantum Computing Is Next Breakthrough

By JOHN MARKOFF JUNE 23, 2014

That may change soon. The company has been spending heavily and is contributing to 10 of the roughly 20 academic research groups exploring a long-hypothesized class of subatomic particles known as Majorana fermions. Beyond being a scientific advance, proving the existence of the Majorana would mean that it was likely they could be used to form qubits for this new form of quantum computing.

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News > Science

'Angel particle' which is both matter and anti-matter discovered in 'landmark' quantum physics breakthrough

Scientists say they have found the first evidence that 'Majorana fermions' exist, 80 years after they were first suggested

Ian Johnston Science Correspondent | @montaukian | Thursday 20 July 2017 18:41 BST | ₱98 comments





Two angels painted in Byzantium between 395 and 1453AD Public domain image/St Catherine's Monastery, Mount Sinai

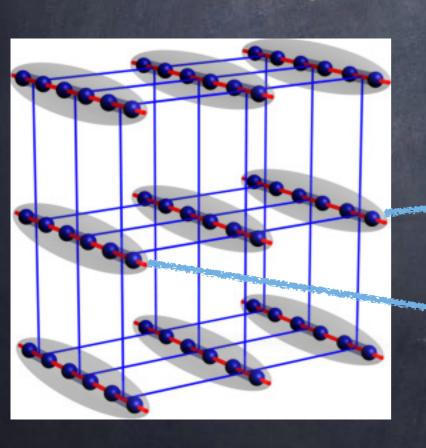
This Talk is about

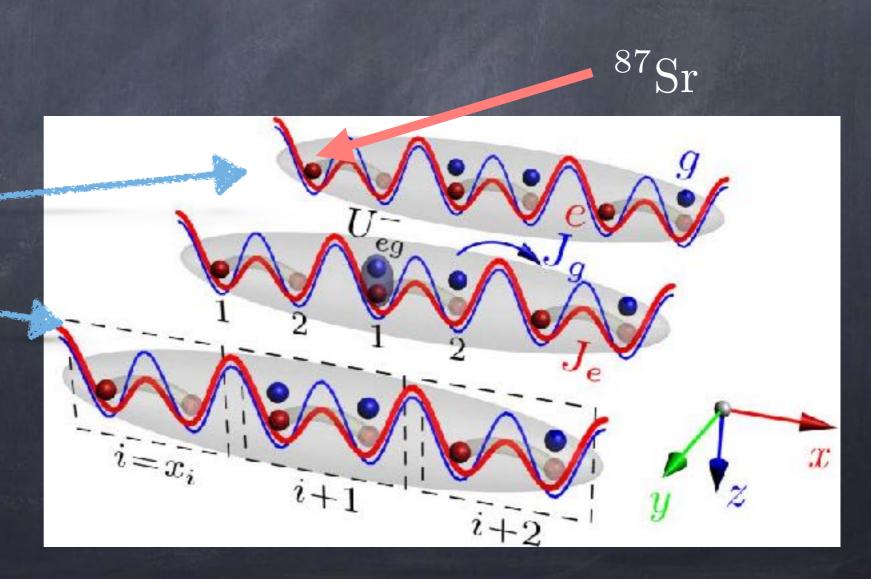
This Talk is about

- Engineering a number-conserving Topological Superfluid (TS) in Optical Lattices with repulsive fermionic atoms
 - New physical mechanism for emergent superfluidity
 - New magneto-electric phenomenon to detect TSs

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Collaborators:

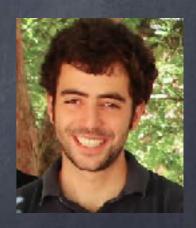
Leonid Isaev: JILA, University of Colorado - Boulder



Ana Maria Rey: JILA, University of Colorado - Boulder



Adam Kaufman: JILA, University of Colorado - Boulder



Mostly: PRL 113, 267002 (2014)

Ann. Phys. 372, 357 (2016)

and arXiv:1710.02768

Outline

Motivation

What and Why Topological Superfluidity?

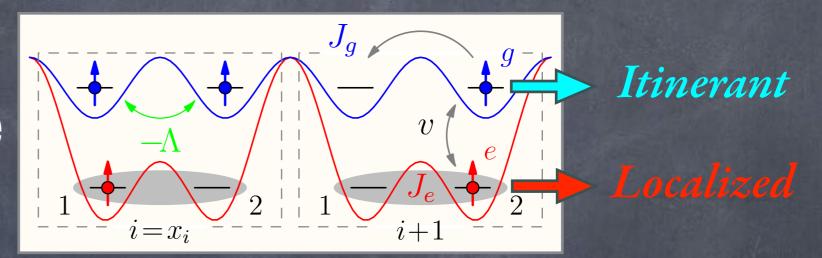
Repulsive Superfluidity in Optical Lattices

- **■** Fermion Pairing from Repulsion
- Attraction from Local Fluctuations: New Mechanism
- Topological Superfluid State in a quasi-1D Lattice
- Topological Superfluidity in 2-D

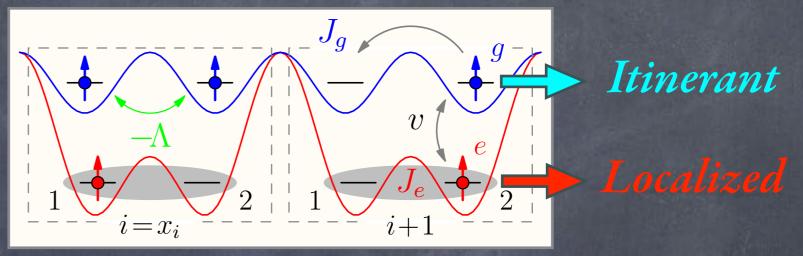
Developing Probes in AMO Experiments

Probing Topological Superfluidity

- ► Main ingredients
 - Optical superlattice
 - Spin-orbit coupling



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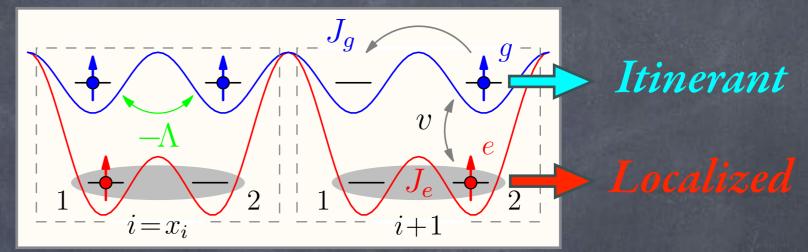


- ► Alcaline-earth atoms ¹⁷¹Yb, ¹⁷³Yb, ⁸⁷Sr
 - Electronic states: 1S_0 , 3P_0 with J=0
 - Many nuclear spin states
 - Nuclear spins decoupled from electrons

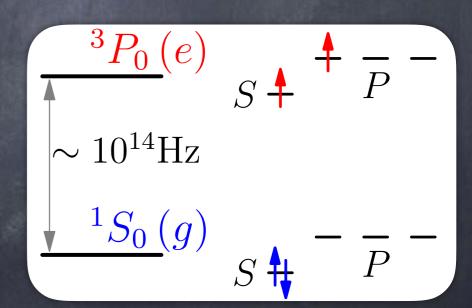
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\end{array}$$

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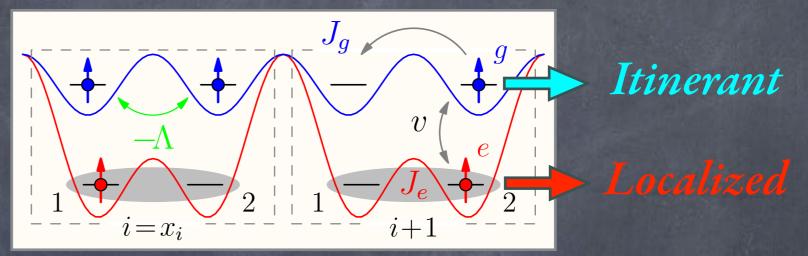
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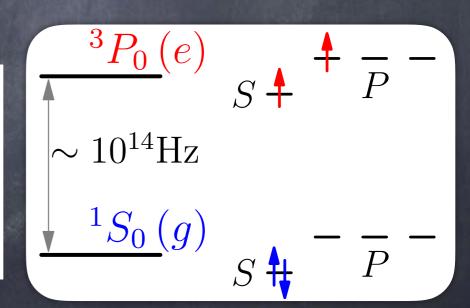
- ► New Mechanism of Pairing: Engineer Subsystem Mediator
 - Interplay between atomic repulsion and kinetic energy

(Clock-State-Dependent Optical Superlattice)

- ► Main ingredients
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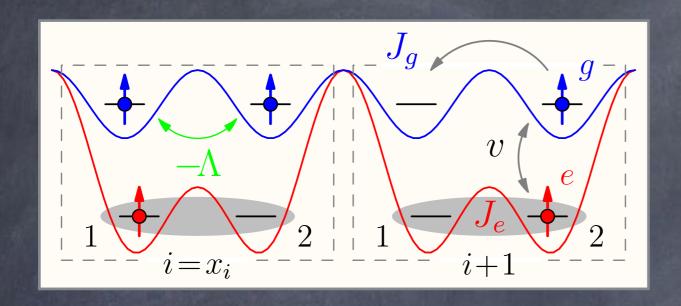


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Superfluid State in a quasi-1D Lattice



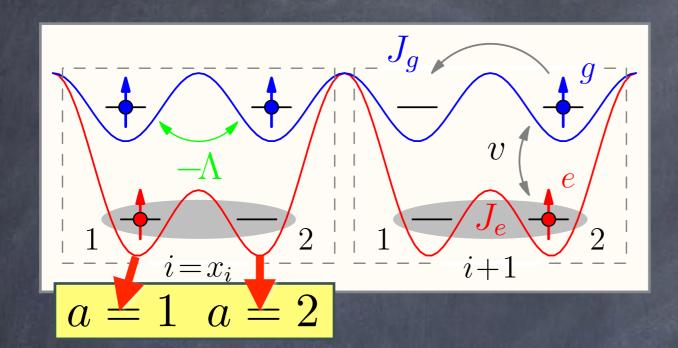
- Mobile g-atoms
- e-atoms: 1 atom/dimer

•
$$i=1\dots N_d, \quad v \sim a_{eg}^-$$

• m, n = Nuclear spins

$$H = -J_{e} \sum_{in} (e_{i,1n}^{\dagger} e_{i,2n} + \text{h.c.}) + U_{gg} \sum_{ia} n_{ia}^{g} (n_{ia}^{g} - 1) - J_{g} \sum_{in} (g_{i,1n}^{\dagger} g_{i,2n} + g_{i+1,1n}^{\dagger} g_{i,2n} + \text{h.c.}) + J_{eg} \sum_{ia} n_{ia}^{e} n_{ia}^{g} + V_{ex} \sum_{iamn} e_{i,an}^{\dagger} e_{i,an} g_{i,am}^{\dagger} g_{i,an}$$

Superfluid State in a quasi-1D Lattice



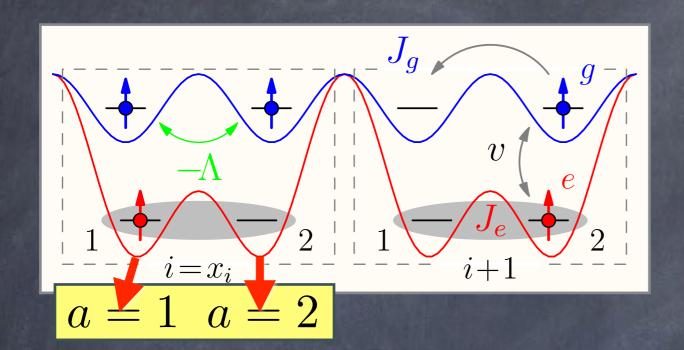
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Superfluid State in a quasi-1D Lattice

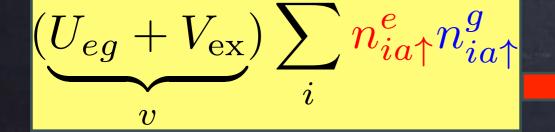


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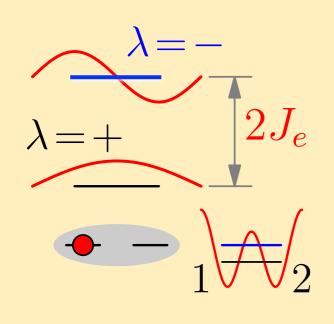
$$-J_{g} \sum_{in} (g_{i,1n}^{\dagger} g_{i,2n} + g_{i+1,1n}^{\dagger} g_{i,2n} + \text{h.c.}) +$$

$$+U_{eg} \sum_{ia} n_{ia}^{e} n_{ia}^{g} + V_{ex} \sum_{iamn} e_{i,am}^{\dagger} g_{i,am}^{\dagger} g_{i,am}^{\dagger} g_{i,an}$$

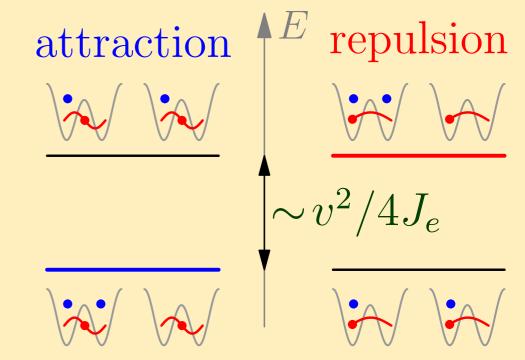


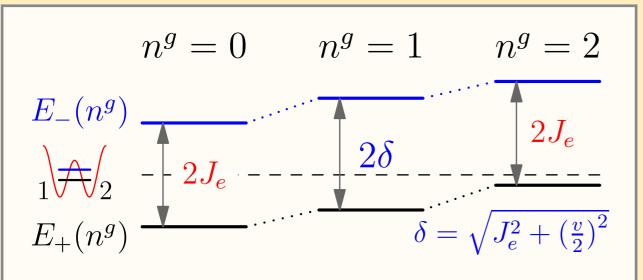
Nuclear spin-polarized

Emergent Fermion pairing



- e-atom quantum fluctuations
 - ullet Consider $J_e\gg J_g,v$
 - Preparing e-atoms in a given motional state
- ▶ Local Pairing: $\Delta_{\pm} \approx \pm v^2/4J_e = \pm \Lambda$
 - sym \longrightarrow repulsion of g-atoms antisym \longrightarrow attraction of g-atoms
 - e-g interaction: $H_{eg} = v(n_1^e n_1^g + n_2^e n_2^g)$



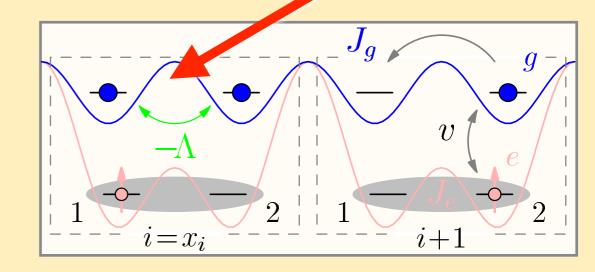


- $\Delta_+ = E_+(2) + E_+(0) 2E_+(1) \approx v^2/4J_e$
- $\Delta_{-} = E_{-}(2) + E_{-}(0) 2E_{-}(1) \approx -v^{2}/4J_{e}$

$$H_{\text{ef}} = -J_g \sum_{k} (g_{k1}^{\dagger} g_{k2}^{\dagger}) \begin{pmatrix} 0 & 1 + e^{-ik} \\ 1 + e^{ik} & 0 \end{pmatrix} \begin{pmatrix} g_{k1} \\ g_{k2} \end{pmatrix} - \frac{v^2}{4J_e} \sum_{i} n_{i,1}^g n_{i,2}^g$$

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Attraction is intra-dimer

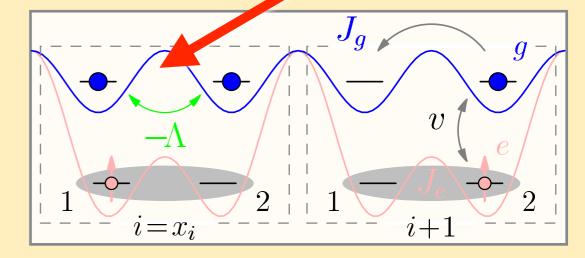


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$$H_{\text{SOI}} = (1 + \cos k)\sigma^x + \sin k \sigma^y$$

- Attraction is intra-dimer
- **Effective spin-orbit interaction**

due to doubling of unit cell

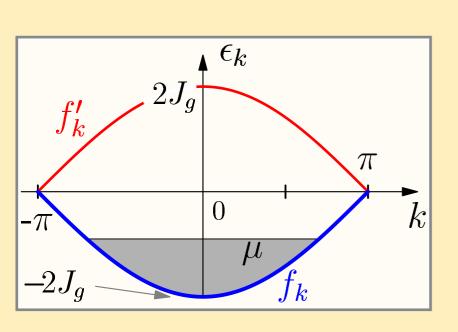


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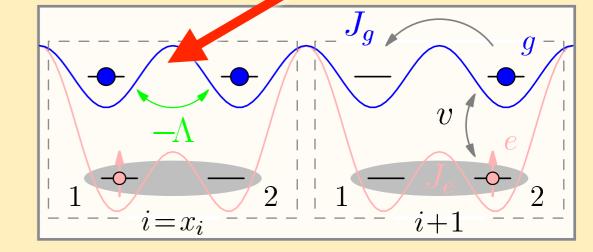
$$\underbrace{H_{\text{SOI}} = (1 + \cos k)\sigma^x + \sin k \sigma^y}_{\text{Homogeneous equation}}$$

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p-wave symmetry



- ► Non-interacting band-structure
 - Quasiparticle states

$$\begin{pmatrix} g_{k1} \\ g_{k2} \end{pmatrix} = \frac{1}{\sqrt{2}} \left(\frac{f_k + f'_k}{\frac{1 + e^{ik}}{\sqrt{2(1 + \cos k)}}} [f_k - f'_k] \right)$$

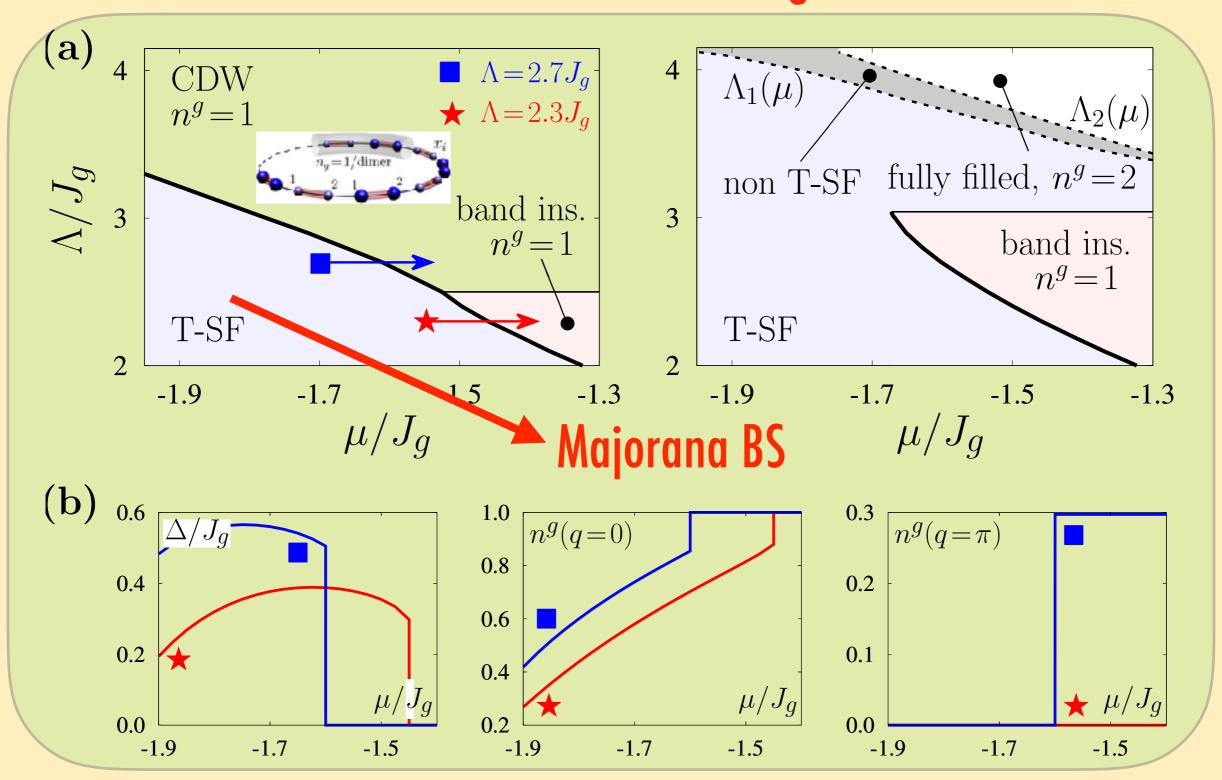
Energies:

$$E_k = \pm J_g \sqrt{2(1+\cos k)}$$

Phase Diagram in quasi-1D

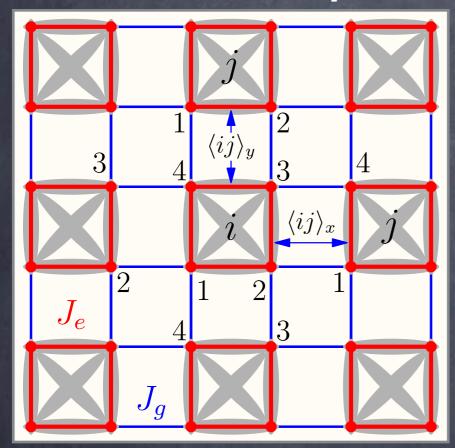


Single-site unit cell



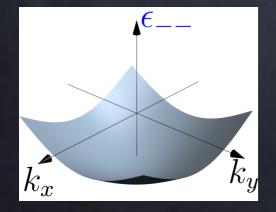
Topological Superfluidity in two-dimensions

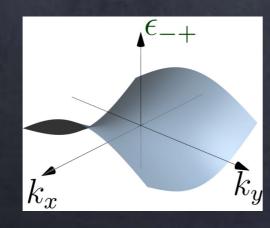
 $(p_x + ip_y)$ -wave symmetry

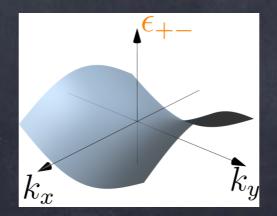


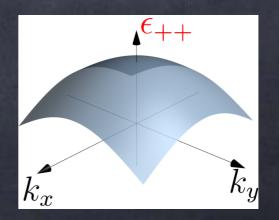
$$H = -J_e \sum_{i, \square_{ab}} (e_{ia}^{\dagger} e_{ib} + \text{h.c.}) + v \sum_{ia} n_{ia}^e n_{ia}^g - J_g \sum_{i, \square_{ab}} (g_{ia}^{\dagger} g_{ib} + \text{h.c.}) - J_g \sum_{i, \square_{ab}} (g_{i,2}^{\dagger} g_{j,1} + g_{i,3}^{\dagger} g_{j,4} + \text{h.c.}) - J_g \sum_{\langle ij \rangle_y} (g_{i,4}^{\dagger} g_{j,1} + g_{i,3}^{\dagger} g_{j,2} + \text{h.c.})$$

g-atom dispersion: $\epsilon_{n_x n_y} = 2J_g \left(n_x \cos \frac{k_x}{2} + n_y \cos \frac{k_y}{2}\right), \, n_{x,y} = \pm 1$









Outline

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What and Why Topological Superfluidity?

Repulsive Superfluidity in Optical Lattices

- Fermion Pairing from Repulsion
- Attraction from Local Fluctuations: New Mechanism
- Topological Superfluid State in a quasi-1D Lattice
- Topological Superfluidity in 2-D

Developing Probes in AMO Experiments (^{87}Sr)

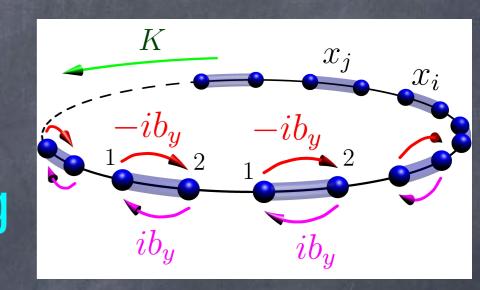
Probing Topological Superfluidity

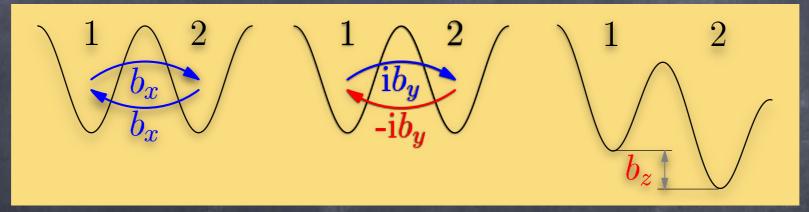
Synthetic Magneto-Electric Effect (Spin-Galvanic Effect)

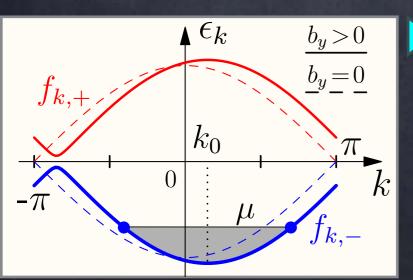
► Laser-induced Magnetic Field

$$\delta H_M = -\sum_{i,a,b} (\boldsymbol{b} \cdot \boldsymbol{\sigma}_{ab}) g_{ia}^{\dagger} g_{ib}$$

- $b_x/b_y \leftrightarrow \text{intra-dimer real/complex hopping}$
- $ullet b_z = ext{relative energy shift of dimer sites}$



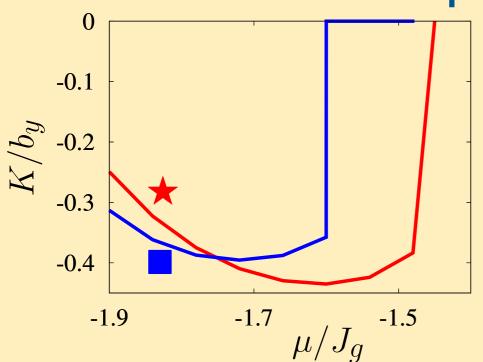


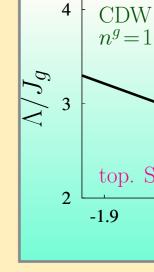


- ► Time-reversal Symmetry *T*
 - $g_{k,a} o g_{-k,a}$ and a c-number $c o c^*$
 - ullet Only b_y breaks ${\mathcal T}$
 - We consider $b = (0, b_y, 0)$

<u>Unveiling the Superfluid State</u>

► Linear Mass Current Response





 $\Lambda = 2.7 J_c$

 \star $\Lambda = 2.3 J_c$

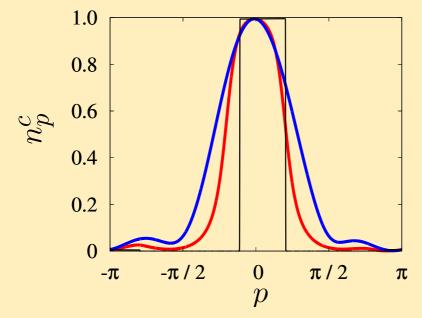
band ins.

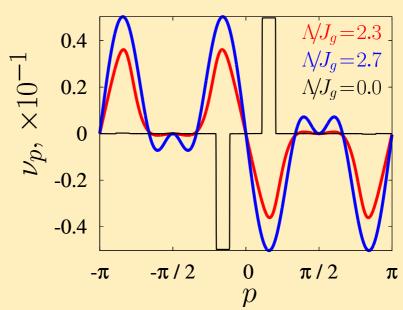
-1.5

 μ/J_q

- ullet $b_y \ll |\Delta|$ to avoid FFLO states
- ullet Supercurrent K is a fingerprint of a SF

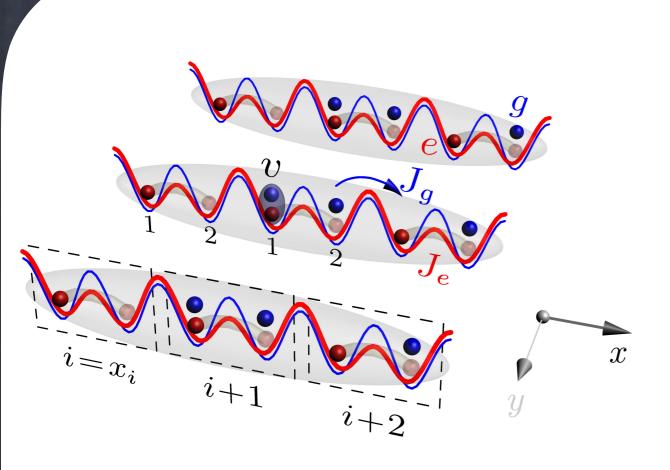
Momentum Distribution Asymmetry $[\nu_p = J_g(n_p^c - n_{-p}^c)/b_y]$





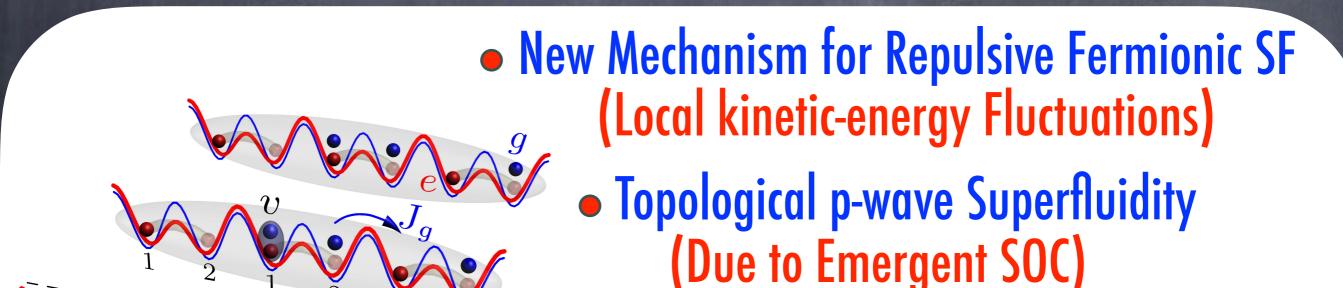
Time-of-flight measurement

Fingerprint of a SF



New Mechanism for Repulsive Fermionic SF (Local kinetic-energy Fluctuations)
 Topological p-wave Superfluidity (Due to Emergent SOC)

New Magneto-Electric Effect
 and Experimental Probes



- New Magneto-Electric Effect
 and Experimental Probes
- Use of realistic exp-controllable interactions (no heating)
- Superfluidity without Feshbach Resonance (no losses)



(Due to Emergent SOC)

New Magneto-Electric Effect
 and Experimental Probes

- Use of realistic exp-controllable interactions (no heating)
- Superfluidity without Feshbach Resonance (no losses)
 - Open Question: Can we detect/manipulate Majorana modes?

Chi l'ha visto?



ettore Majorana, ordinario di fisica teorica all' Università di Napoli, è misteriosamente scomparso dagli ultimi di marzo. Di anni 31, alto metri 1,70, snello, con capelli neri, occhi scuri, una lunga cicatrice sul dorso di una mano. Chi ne sapesse qualcosa è pregato di scrivere al R. P. E. Maria-

necci, Viale Regina Margherita 66 -

Use of realistic

New Mechanism for Repulsive Fermionic SF (Local kinetic-energy Flynns)

• Topological p-v So Judity (Due to Sol)

(D

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and Experimental Probes

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