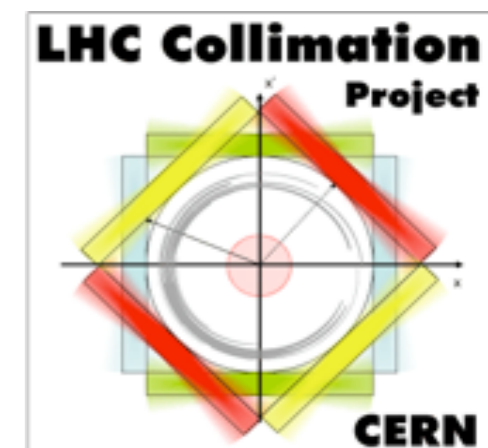


The LHC machine protection and beam collimation

Part 2

Stefano Redaelli

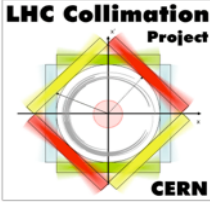
CERN Beams Department
Accelerator Physics Group



- **Introduction to the LHC**
 - Recap. of basic accelerator physics
 - CERN accelerator complex
 - LHC parameters and detailed layouts
- **Machine protection and collimation**
 - Machine protection and collimation system
 - Design of beam halo collimation
 - The LHC beam collimation system
- **Advanced beam collimation**
 - Collimation in practice: LHC operation
 - Simulations and measurements
 - HL-LHC upgrade
 - Advanced concepts: crystals, hollow lenses



Outline - 2nd lecture

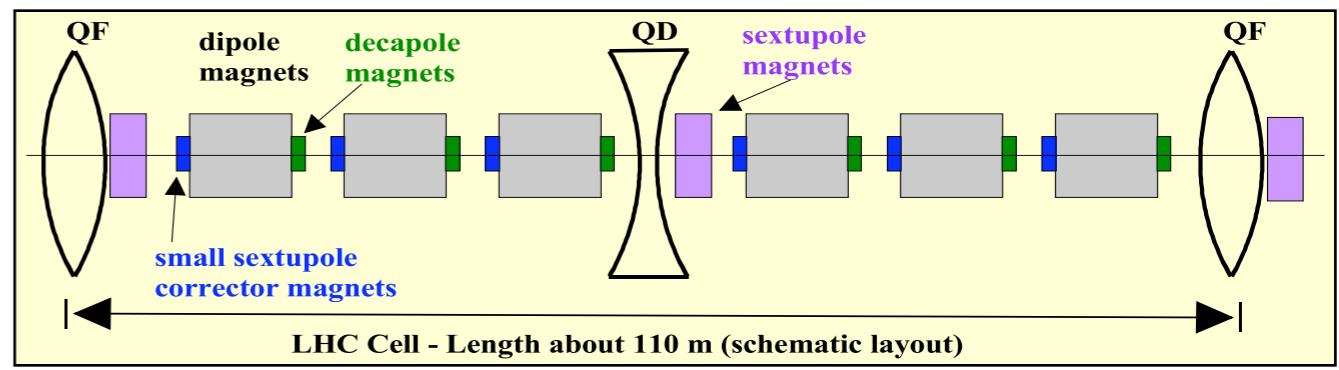


- **Main points from 1st lecture**
- **Machine protection and collimation**
 - Concepts and LHC implementation
 - Case study: 2008 event
- **Beam losses and collimation**
 - Roles of beam collimation systems
 - Beam losses mechanisms
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 - Detailed layouts

High-intensity
circular hadron
accelerators

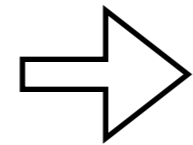
Beam rigidity:

$$B\rho = \frac{p}{e}$$



Equation of motion and its solution:

$$x'' + K(s)x = \frac{1}{\rho} \frac{\Delta p}{p_0}$$



$$x(s) = A\sqrt{\beta_x(s)} \cos[\phi(s) + \phi_0] + D(s) \times \frac{\Delta p}{p}$$

Beta function *Dispersion*

Betatron tune and chromaticity:

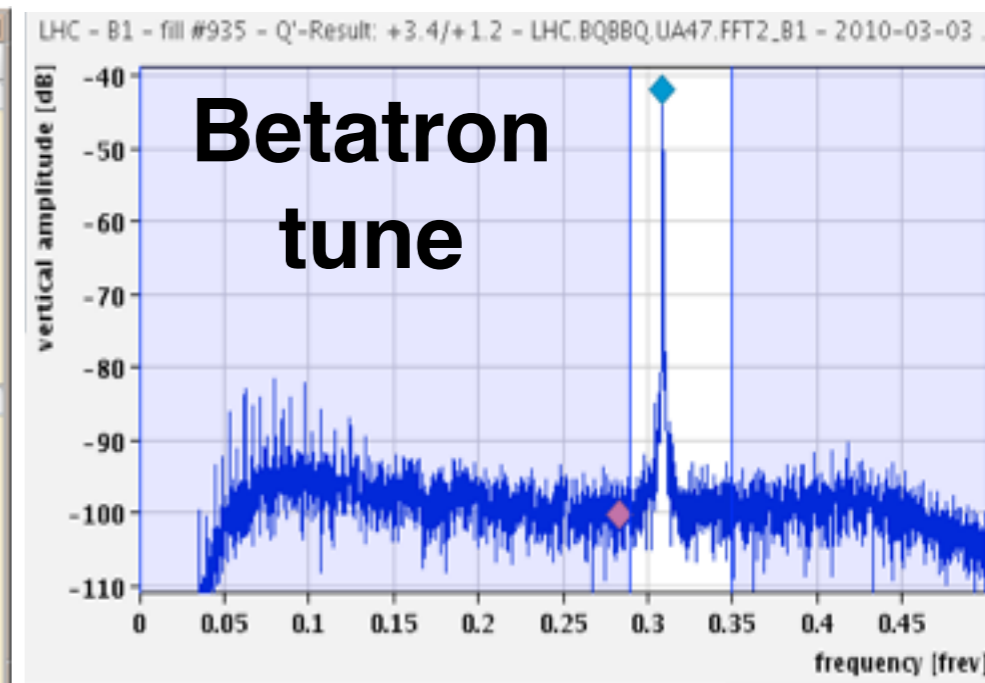
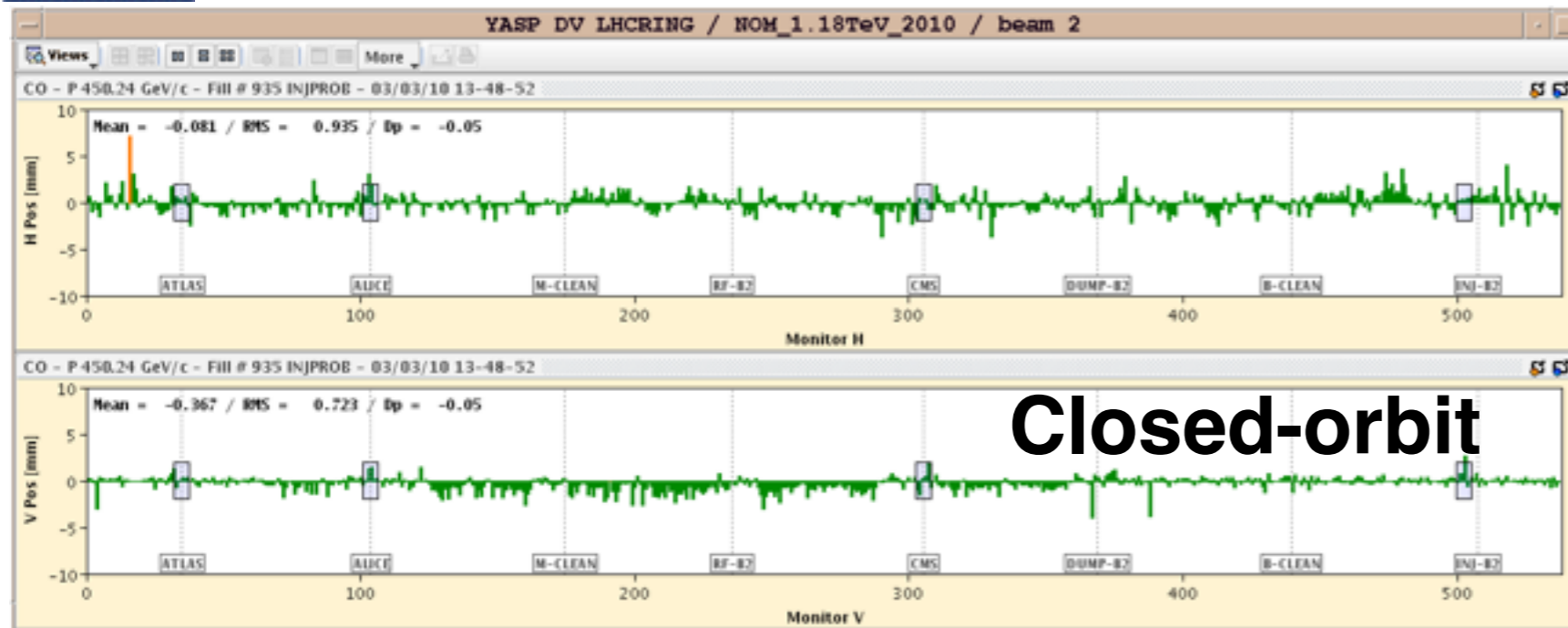
$$Q = \frac{1}{2\pi} \int \frac{ds}{\beta(s)}$$

$$Q' = \frac{\Delta Q}{\Delta p/p}$$

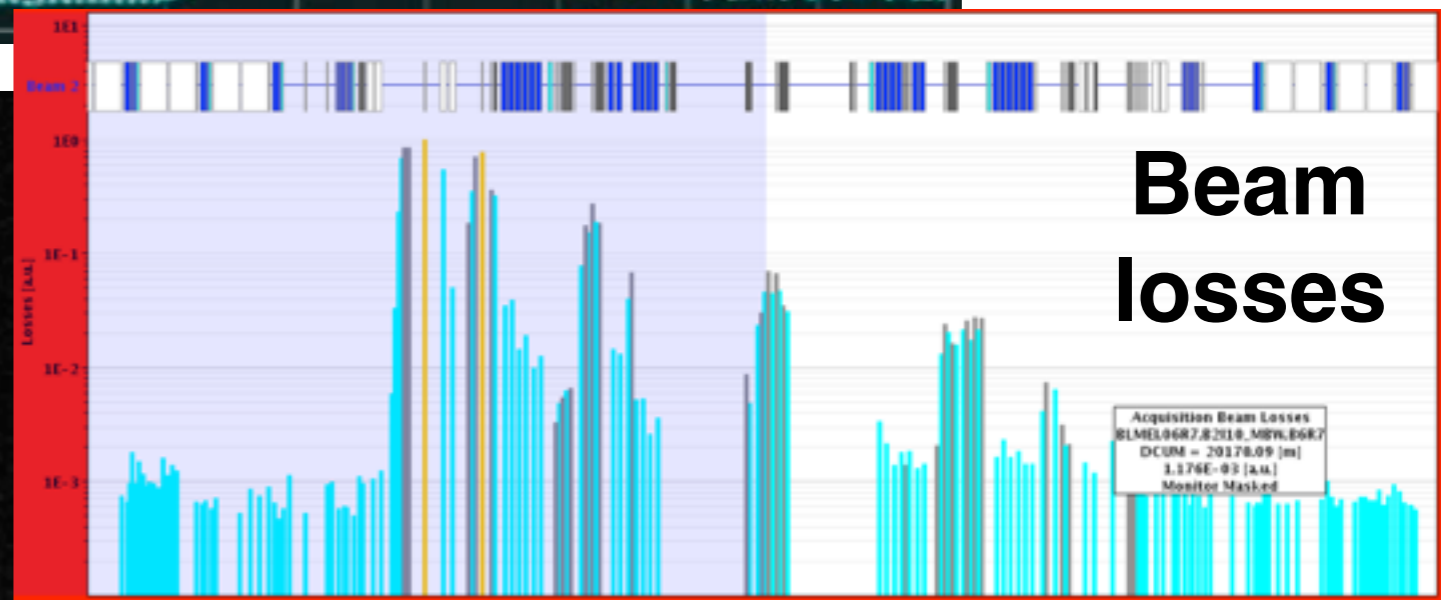
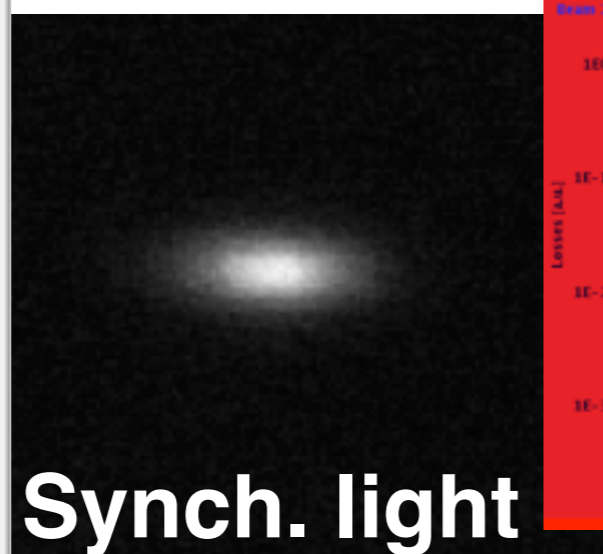
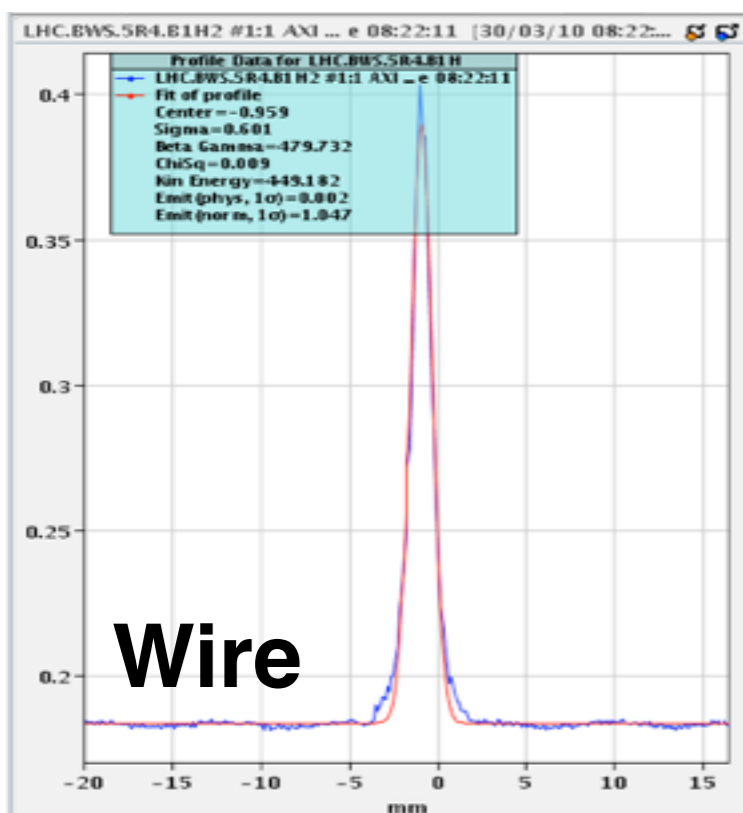
Emittance and beam size:

$$\sigma_x(s) = \sqrt{\epsilon\beta_x(s) + [D_x(s)\delta]^2}$$

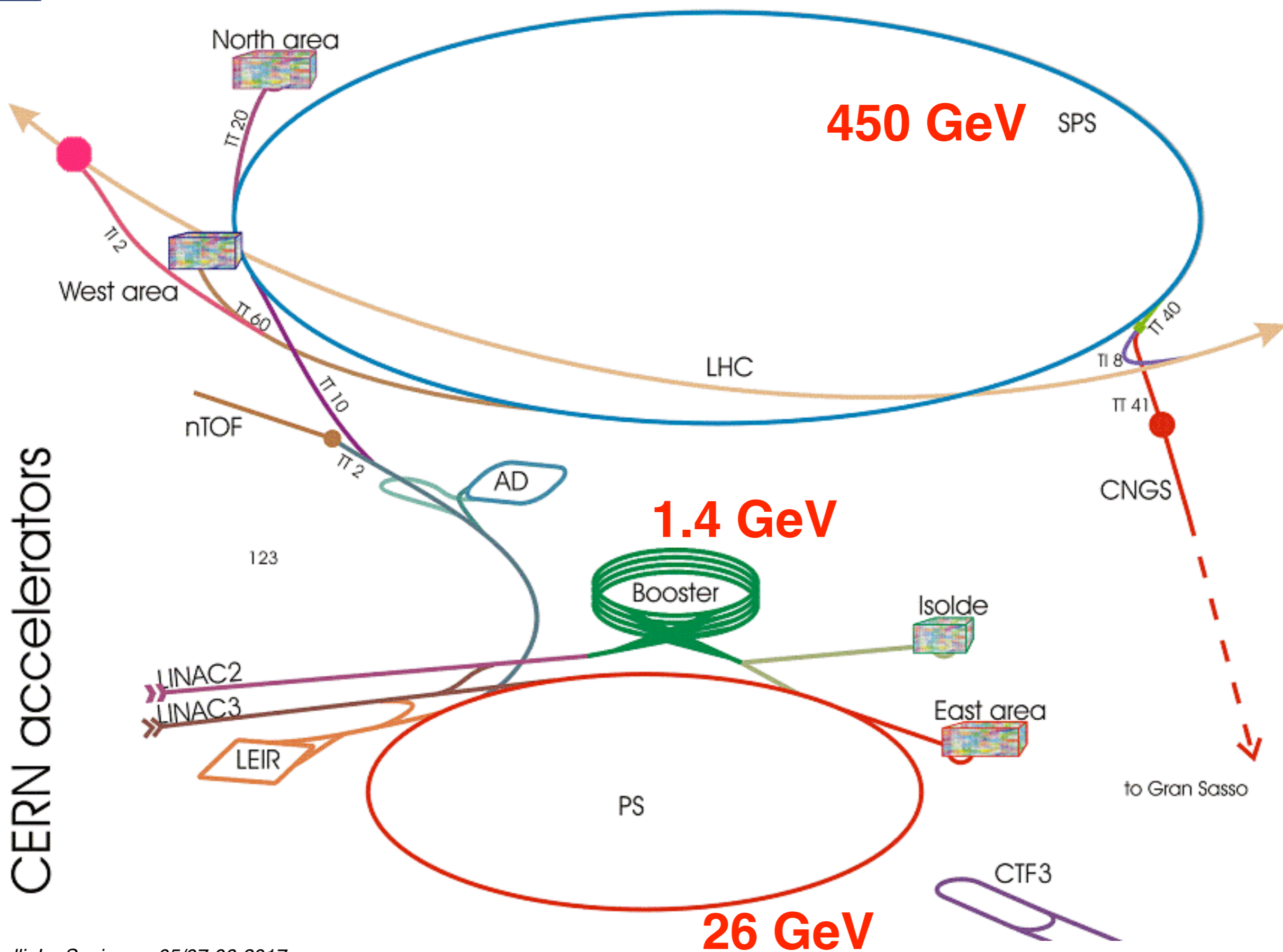
Beam measurements



Transverse size



LHC injector complex

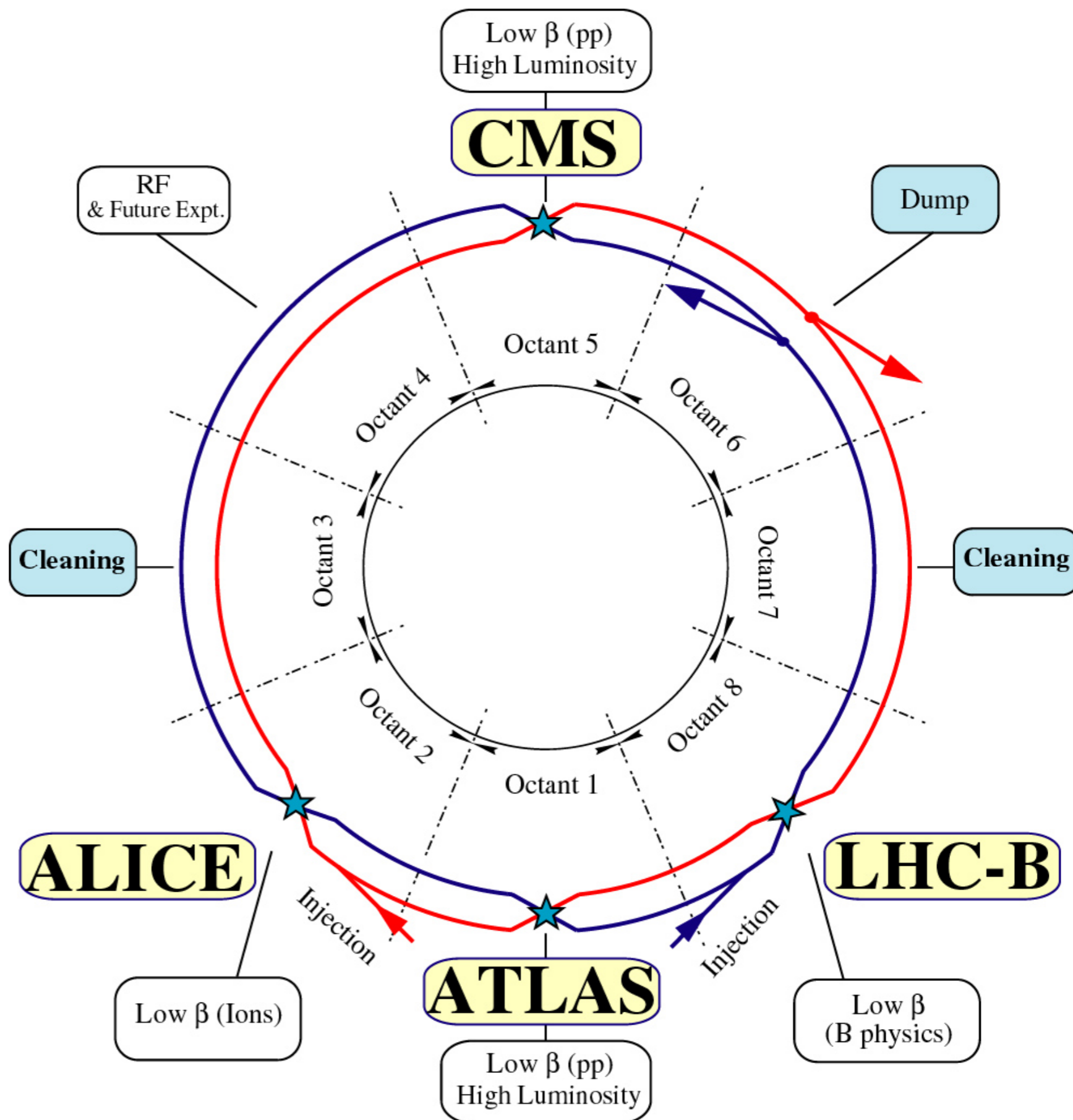


CERN accelerators

The Large Hadron Collider

Eight arcs and eight straight sessions:

- Point 1: **Atlas**, **LHCf**
- Point 2: **Alice**, injection
- Point 3: Momentum cleaning
- Point 4: RF
- Point 5: **CMS**, **TOTEM**
- Point 6: Beam Dumps
- Point 7: Betatron cleaning
- Point 8: **LHCb**, injection



Outline

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Machine protection

Why do we have to care??

Energy stored in the superconducting magnet

10.4 GJ

Energy stored in the 7 TeV beams

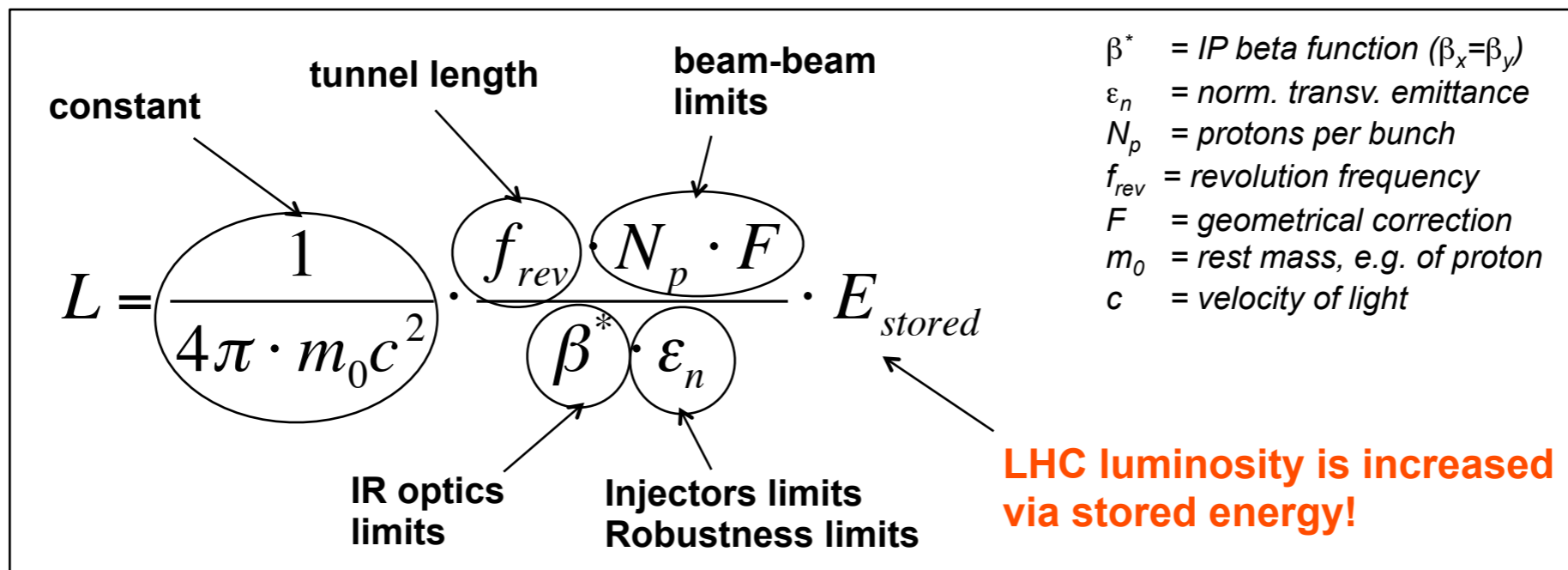
362 MJ

We've seen what damage a 2.5 MJ beam, or even 1 bunch at 7 TeV, can do!

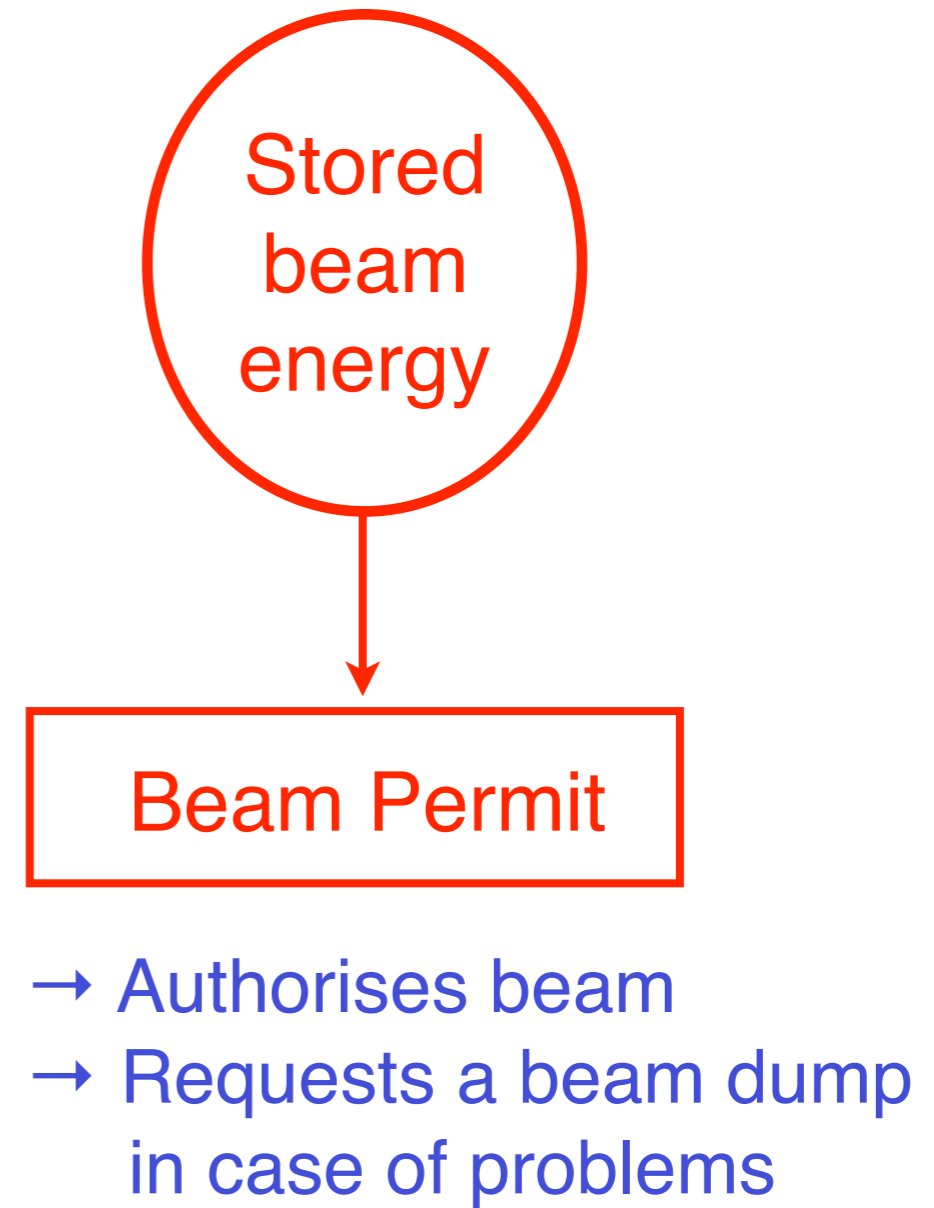
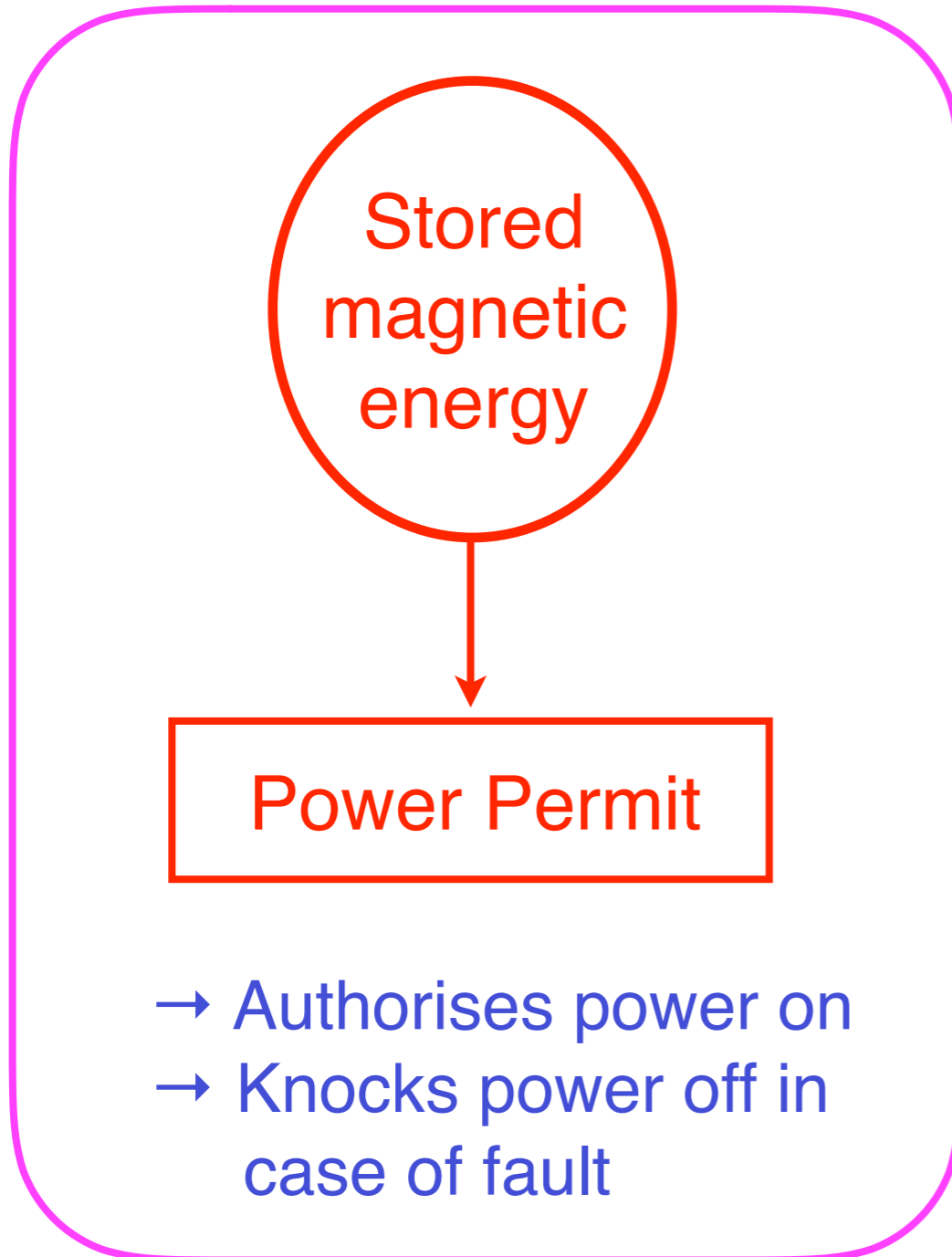
Why do we need so much?

Magnet energy is driven by the high-field requirement.

Beam stored energy is driven by luminosity increase!

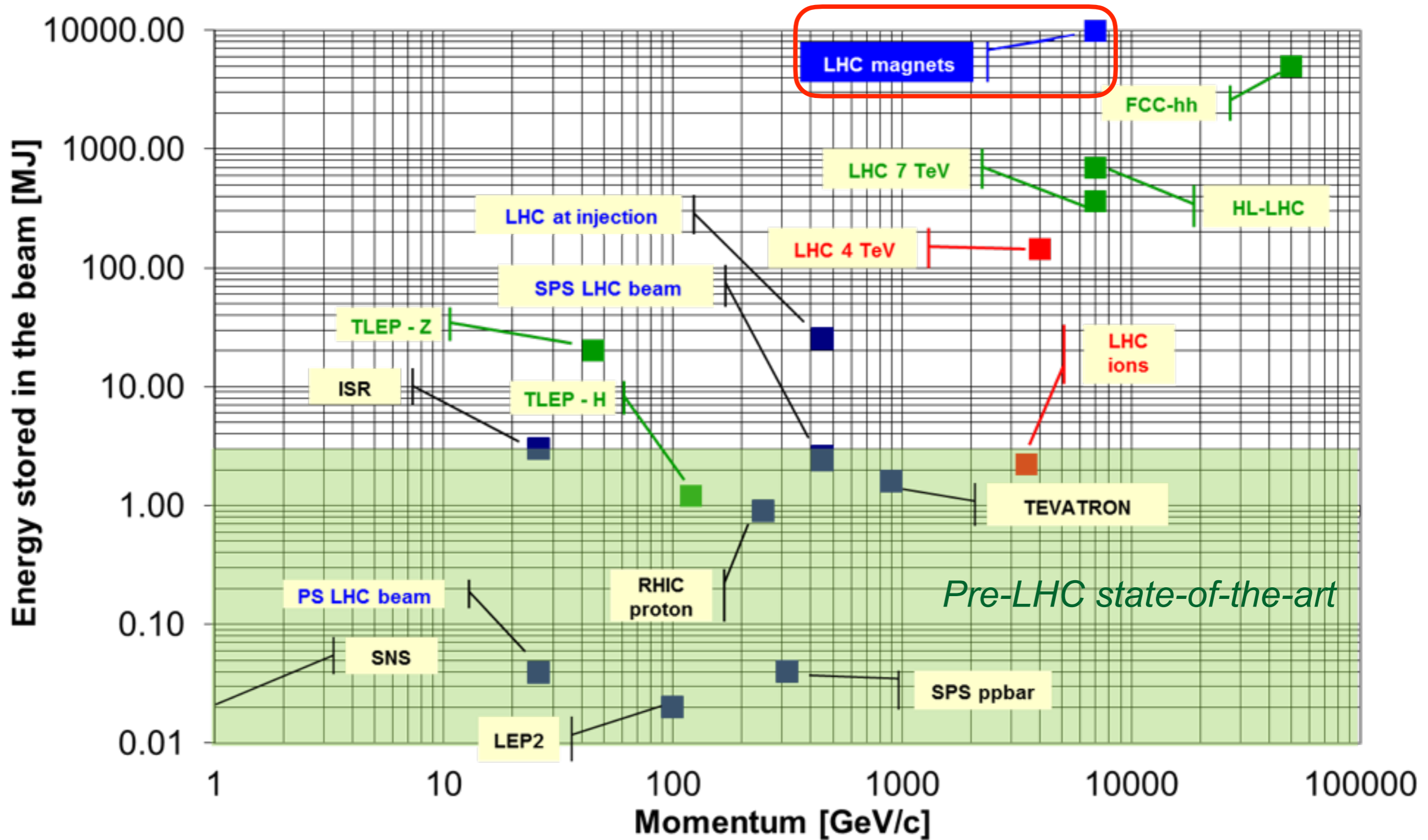


Two sides of machine protection

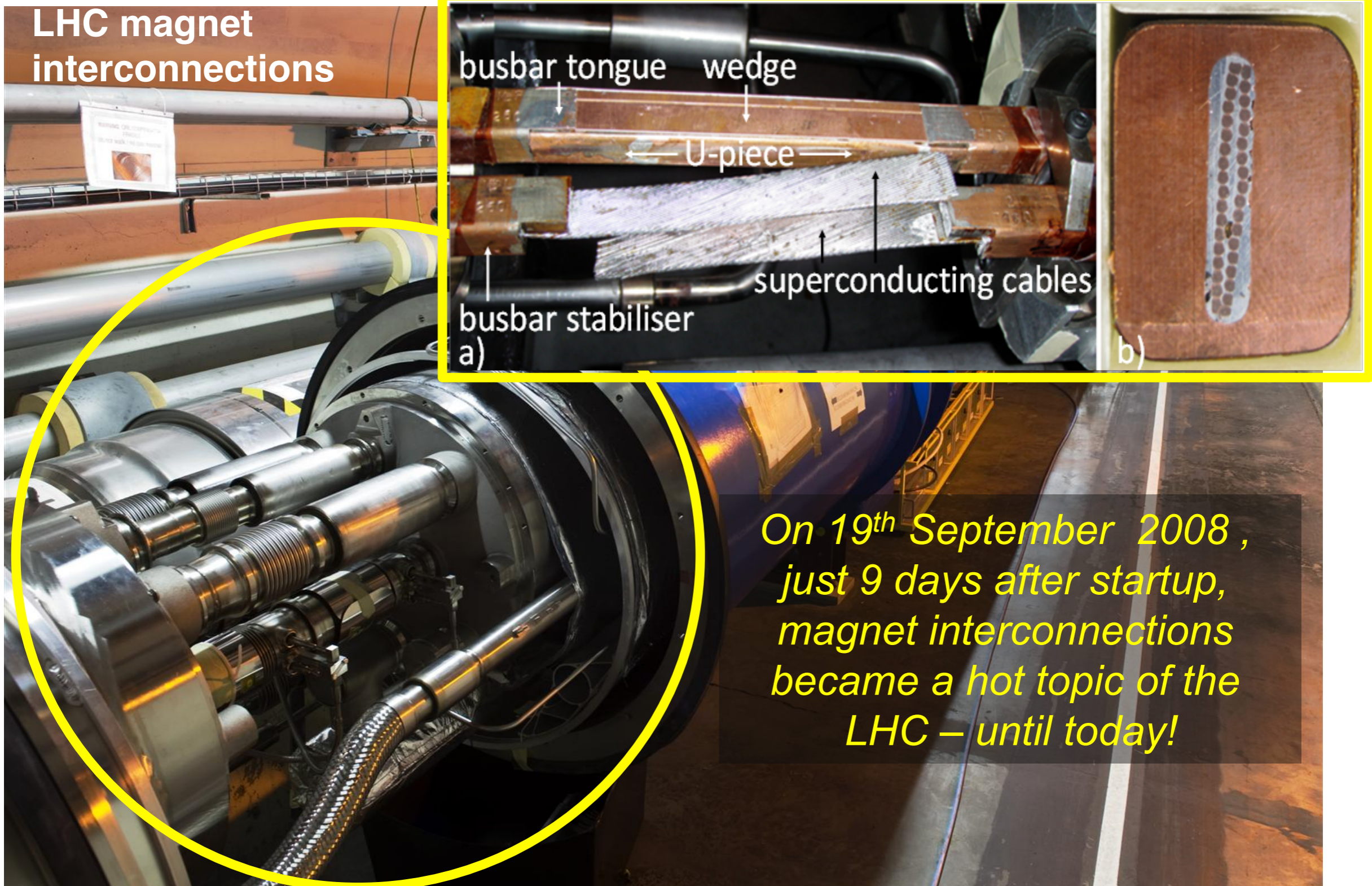


Remark: aspects relates to people safety and environment protection (legal obligations!) are not treated here.

The stored energy challenge



2008 incident on LHC magnet system



List of events

LHC incident on September 19th 2008

- Last commissioning step of one out of the 8 main dipole electrical circuit in sector 34 : **ramp to 9.3kA (5.5 TeV)**.
- At 8.7kA an electrical fault developed in the **dipole bus bar** located in the interconnection between quadrupole Q24.R3 and the neighboring dipole.

Later correlated to a local resistance of $\sim 220 \text{ n}\Omega$ – nominal value $0.35 \text{ n}\Omega$.

- An electrical arc developed which punctured the helium enclosure.

Secondary arcs developed along the arc.

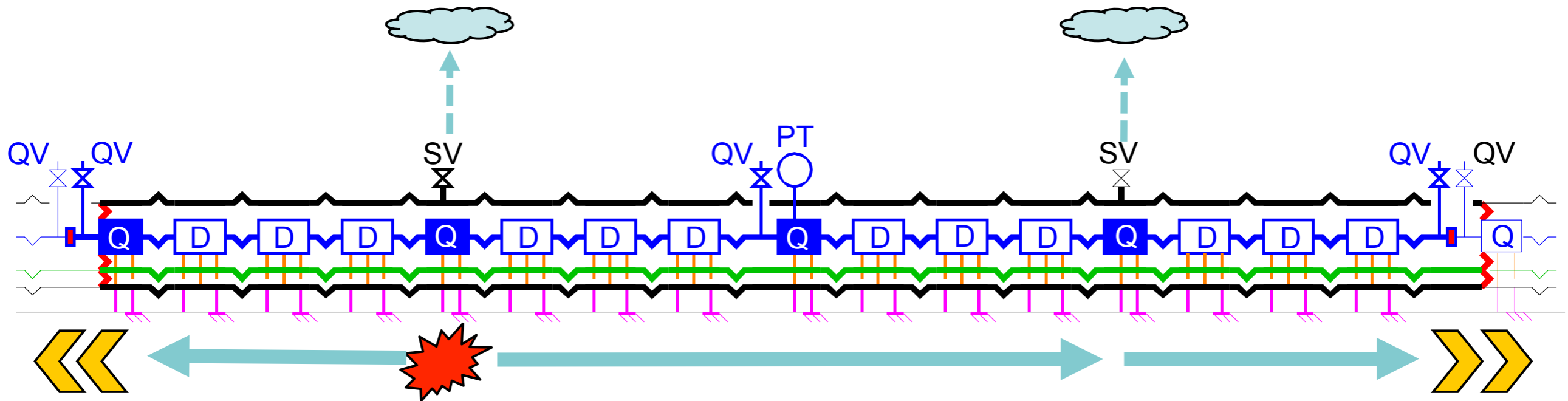
Around 400 MJ from a total of 600 MJ stored in the circuit were dissipated in the cold-mass and in electrical arcs.

- Large amounts of Helium were released into the insulating vacuum.

In total 6 tons of He were released.

This incident involved magnet powering, but no beam!

Helium pressure wave

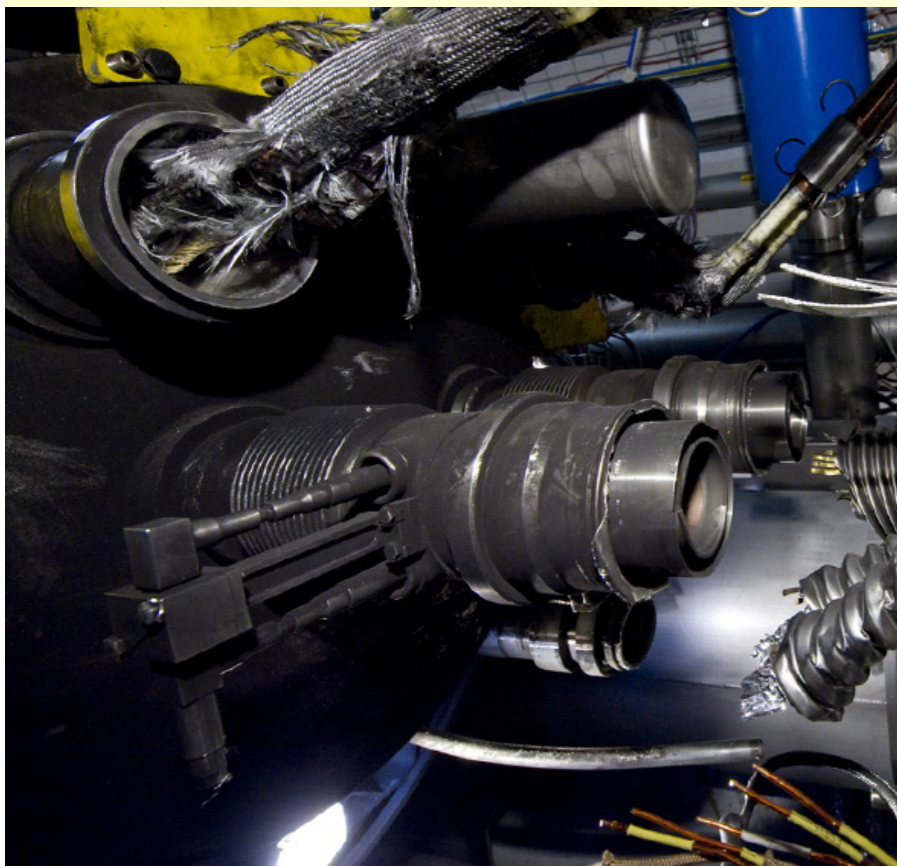


- Cold-mass
- Vacuum vessel
- Line E
- | Cold support post
- | Warm Jack
- ~ Compensator/Bellows
- ⚡ Vacuum barrier

- Pressure wave propagates along the magnets inside the insulating vacuum enclosure.
- Rapid pressure rise :
 - Self actuating relief valves could not handle the pressure.
designed for 2 kg He/s, incident ~ 20 kg/s.
 - Large forces exerted on the vacuum barriers (every 2 cells).
designed for a pressure of 1.5 bar, incident ~ 8 bar.
 - Several quadrupoles displaced by up to ~50 cm.
 - Connections to the cryogenic line damaged in some places.
 - Beam vacuum to atmospheric pressure.

Damage from 600MJ at the LHC

Arcing in the interconnection



53 magnets had to be repaired



Over-pressure

Magnet displacement

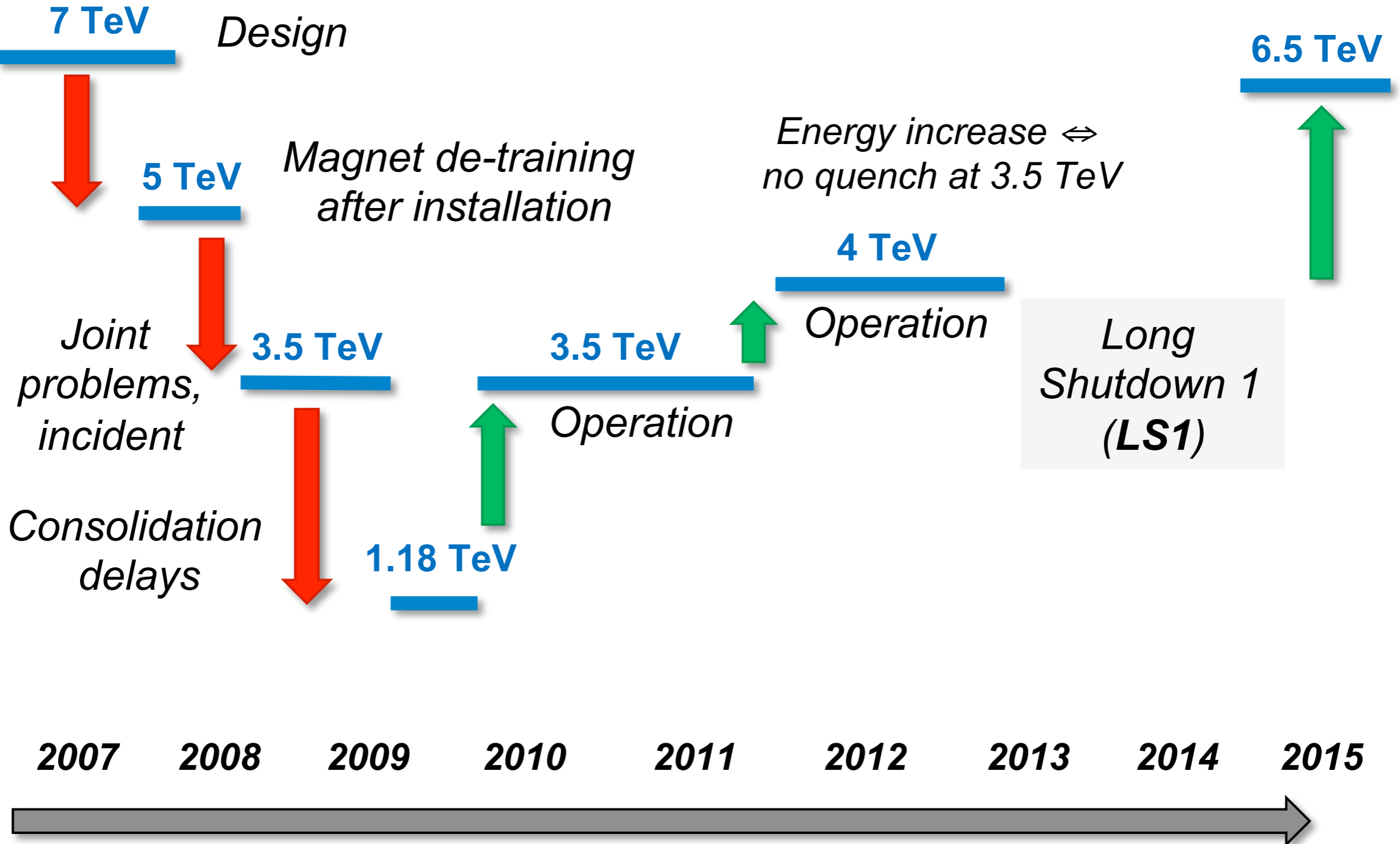


A major event for the LHC that caused (1) Nearly 1 year of delay in the startup with beam; (2) Severe limitations to the operating beam energy in Run I (2010-12); (3) Massive works in 2013-2014 to repair all the 10'000 interconnections! (4) Associated financial implications for delays and repair...

Consequences — energy evolution

Energy (TeV)

Consolidation of all interconnections



Beam damage — Relevant parameters

□ Momentum of the particle

□ Particle type

Activation is mainly an issue for hadron accelerators.

□ Energy stored in the beam

1 MJ can heat and melt 1.5 kg of copper.

1 MJ = energy stored in 0.25 kg of TNT.

□ Beam power

□ Beam size

□ Time structure of beam

One LHC beam = 360 MJ = ?



90 kg of TNT



8 litres of gasoline



15 kg of chocolate



Key factor :
how easily and how fast
the energy is released !!



Three P's for machine protection (MP³)



□ Protect the machine

- Highest priority is to avoid damage of the accelerator.

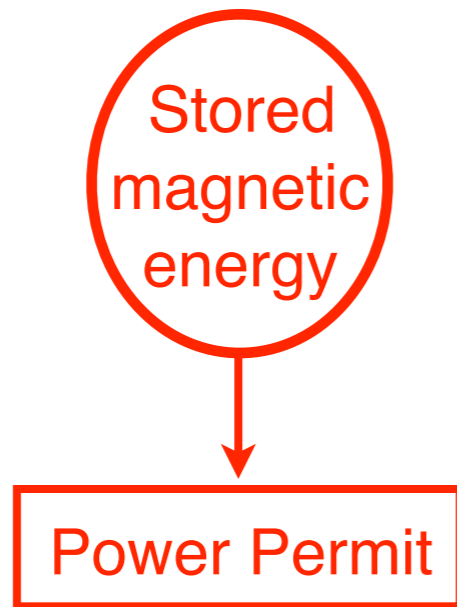
□ Protect the beam

- Complex protection systems reduce the availability of the accelerator, the number of “false” interlocks stopping operation must be minimized.
- Trade-off between protection and operation.

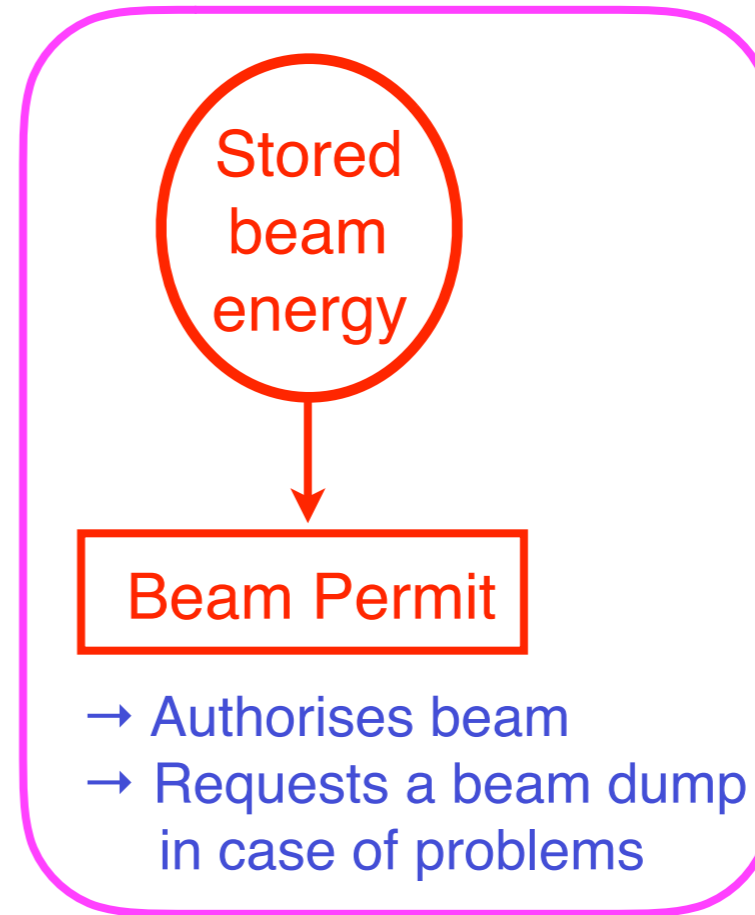
□ Provide the evidence

- Clear (post-mortem) diagnostics must be provided when:
 - the protection systems stop operation,
 - something goes wrong (failure, damage, but also ‘near miss’).

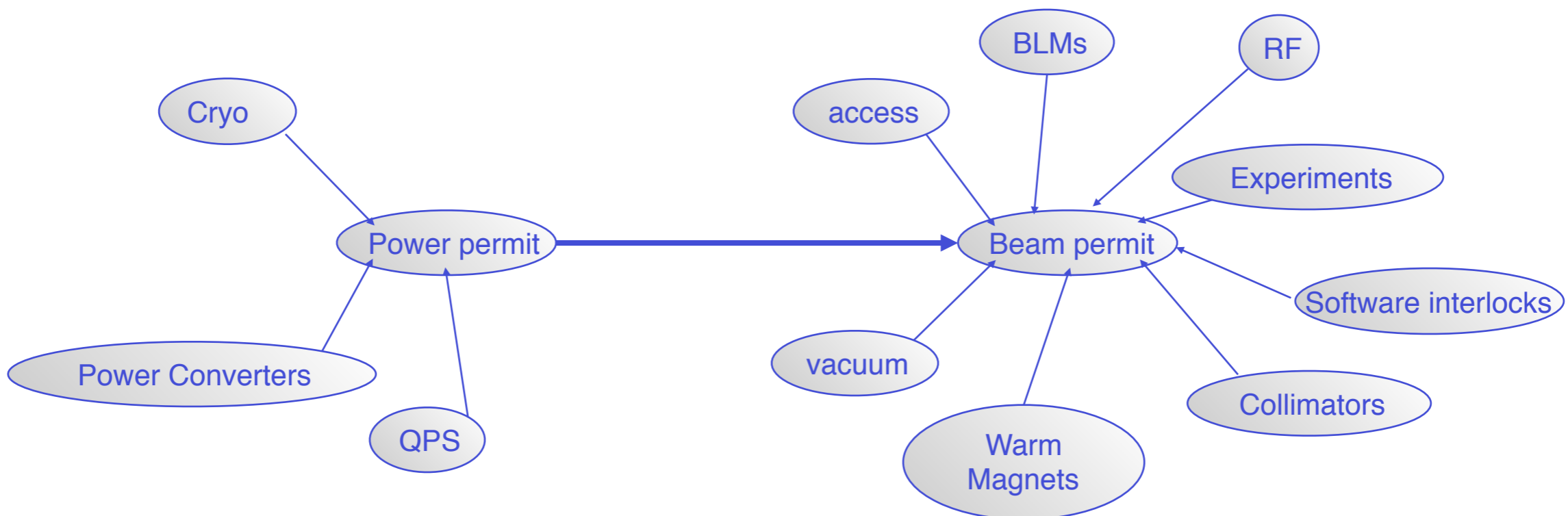
Machine protection philosophy



- Authorises power on
- Knocks power off in case of fault



- Authorises beam
- Requests a beam dump in case of problems



Active protection

- Equipment surveillance.
- Beam observation.
- Extraction (dump) kickers.

Detection of a failure directly on the equipment or by its effects on the beam.

Passive protection

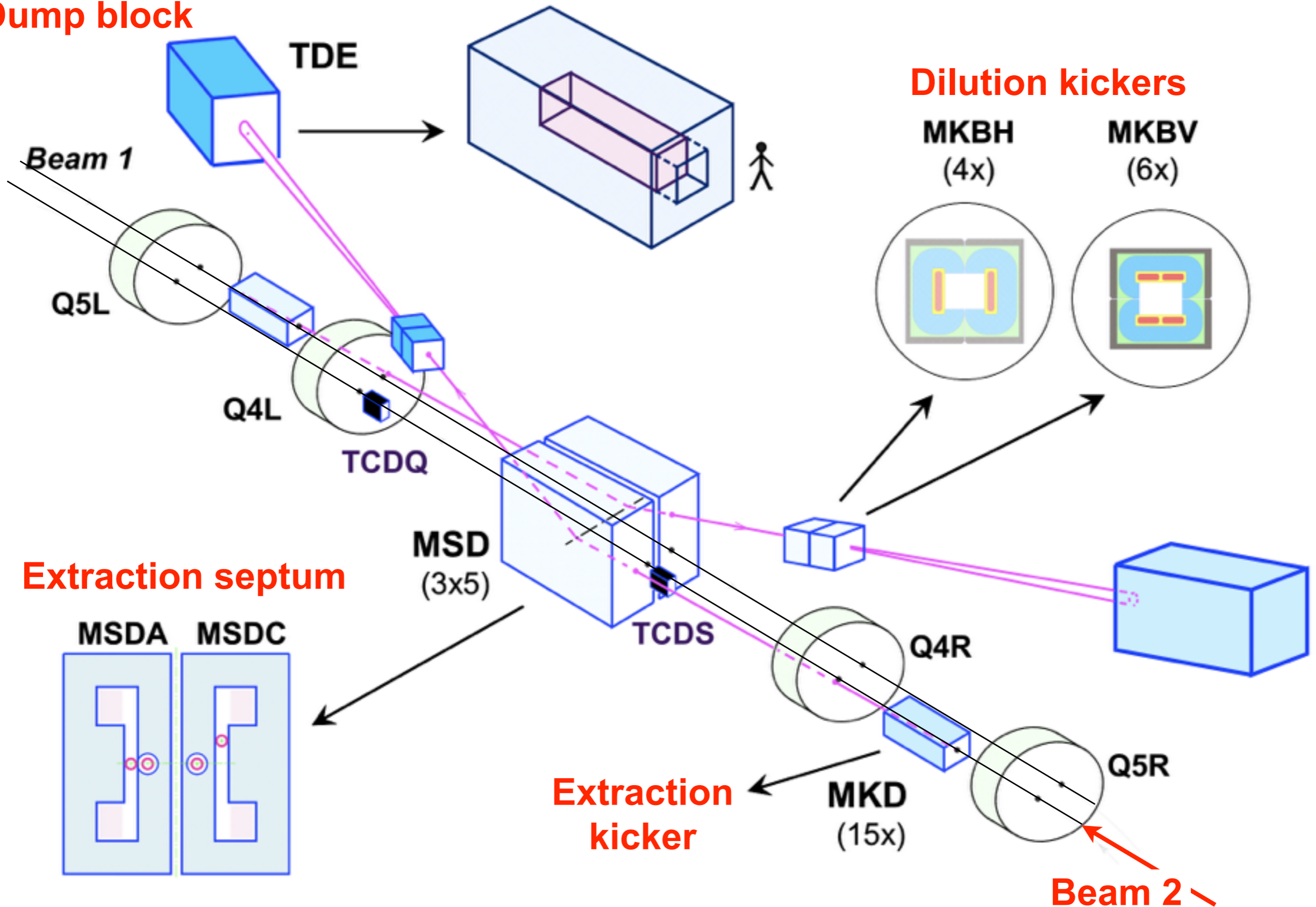
- Collimators.
- Masks.
- Absorbers.
- Dumps.

Obstacles to absorb/dilute the energy energy to mitigate risks of damage

Modern MP systems require both passive and active protection to cover all failure cases. The LHC system provided an unprecedented performance needed to meet the specific challenges of 362MJ beams!

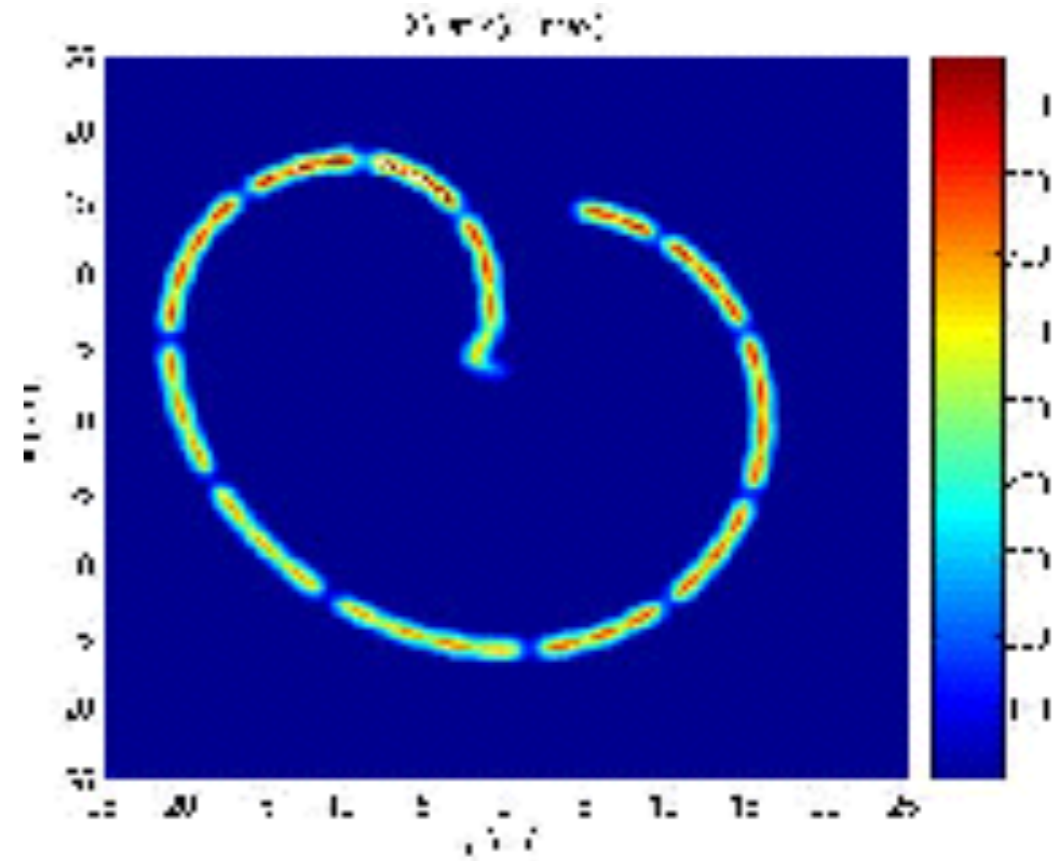
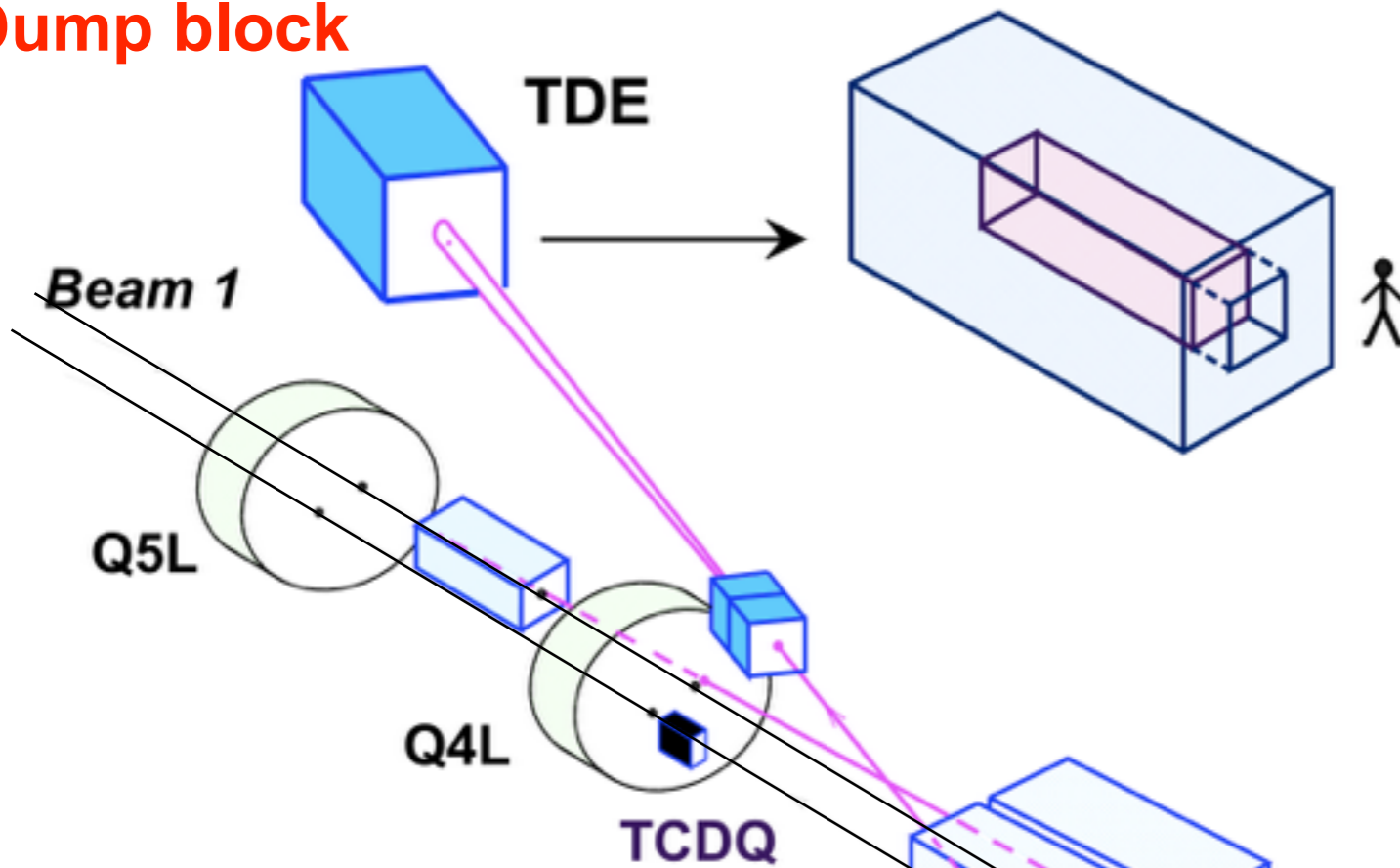
Recap.: LHC beam dump system

Dump block

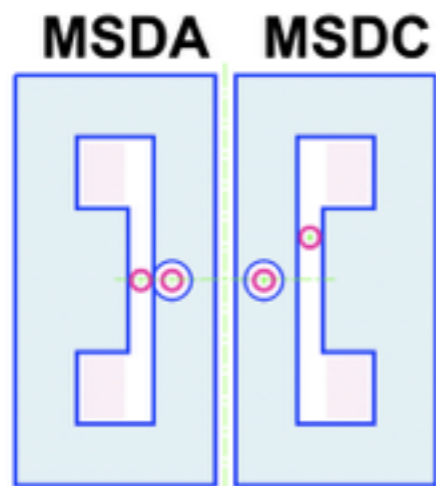


Recap.: LHC beam dump system

Dump block

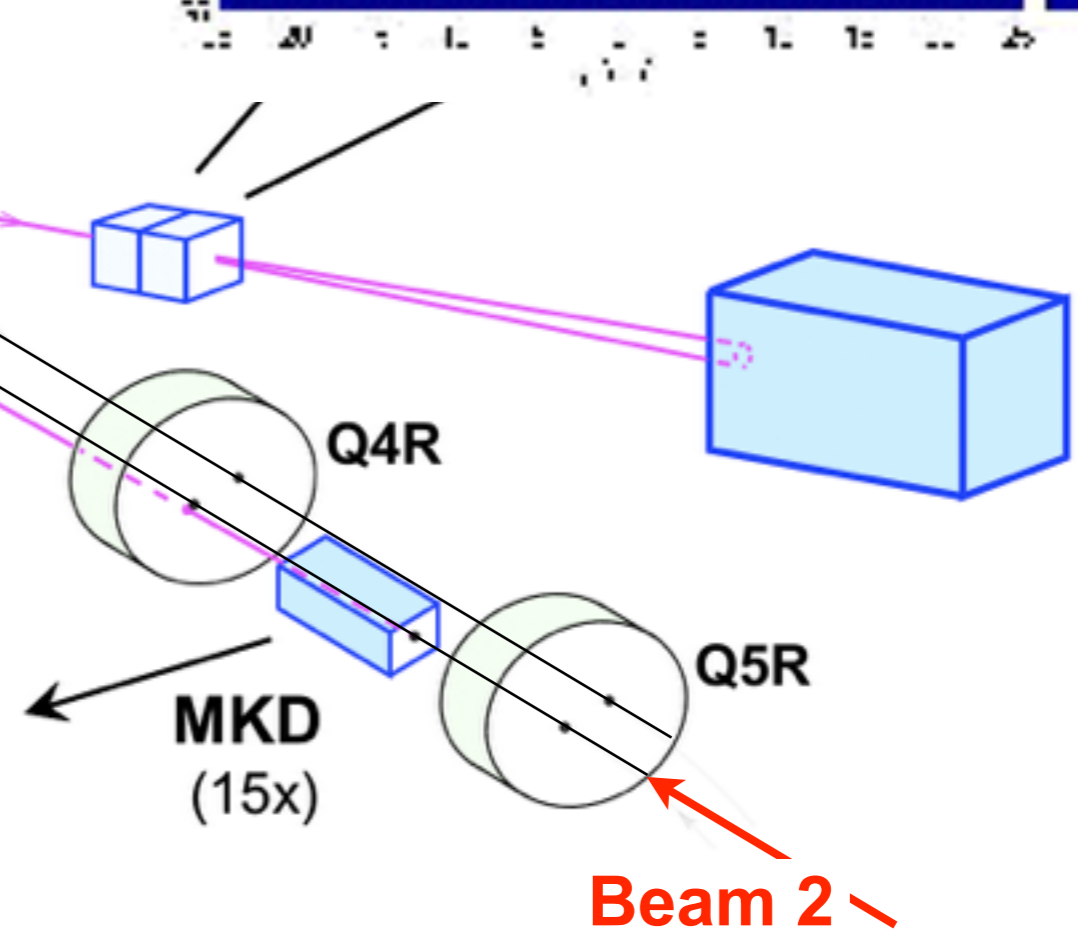


Extraction septum

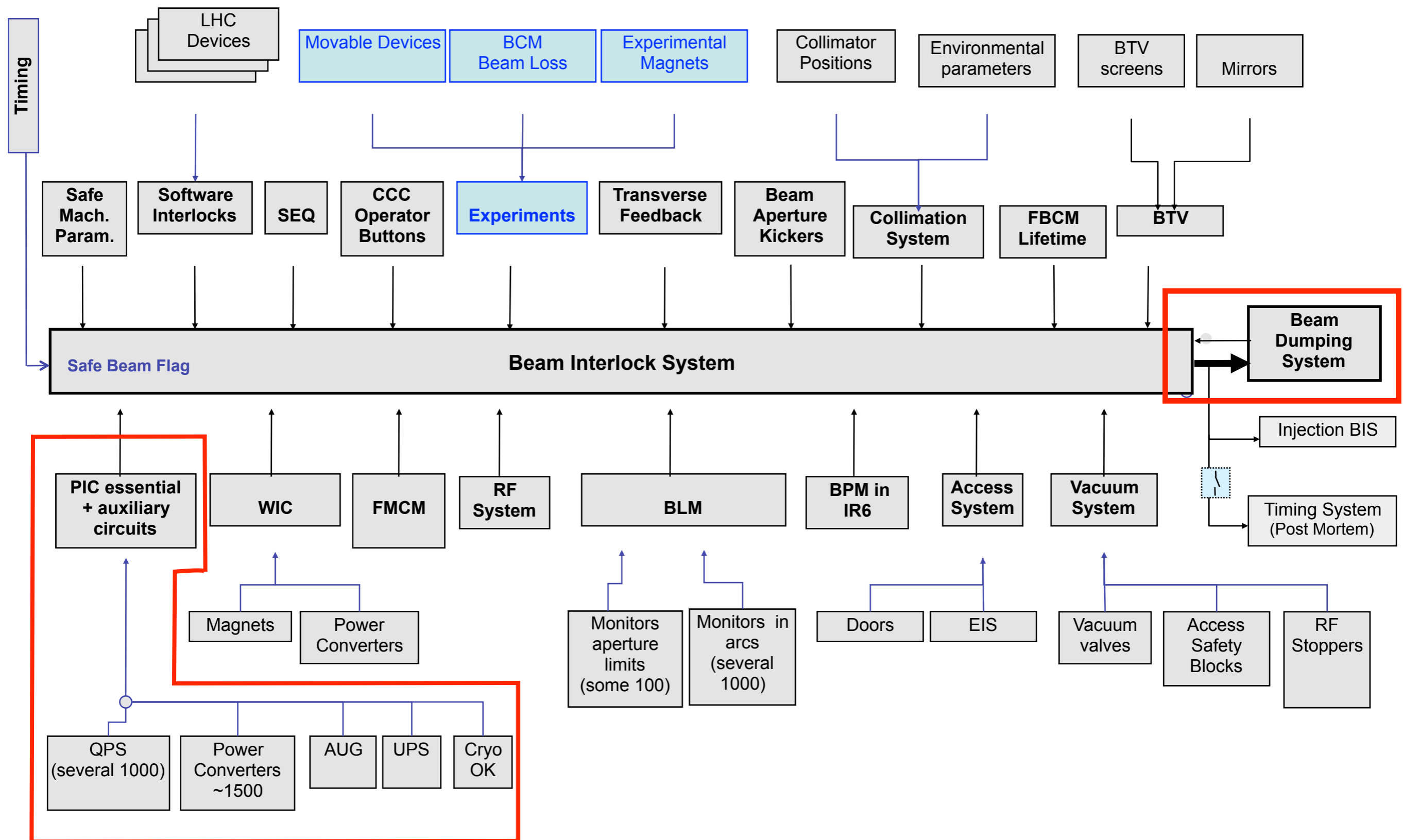


MSD
(3x5)

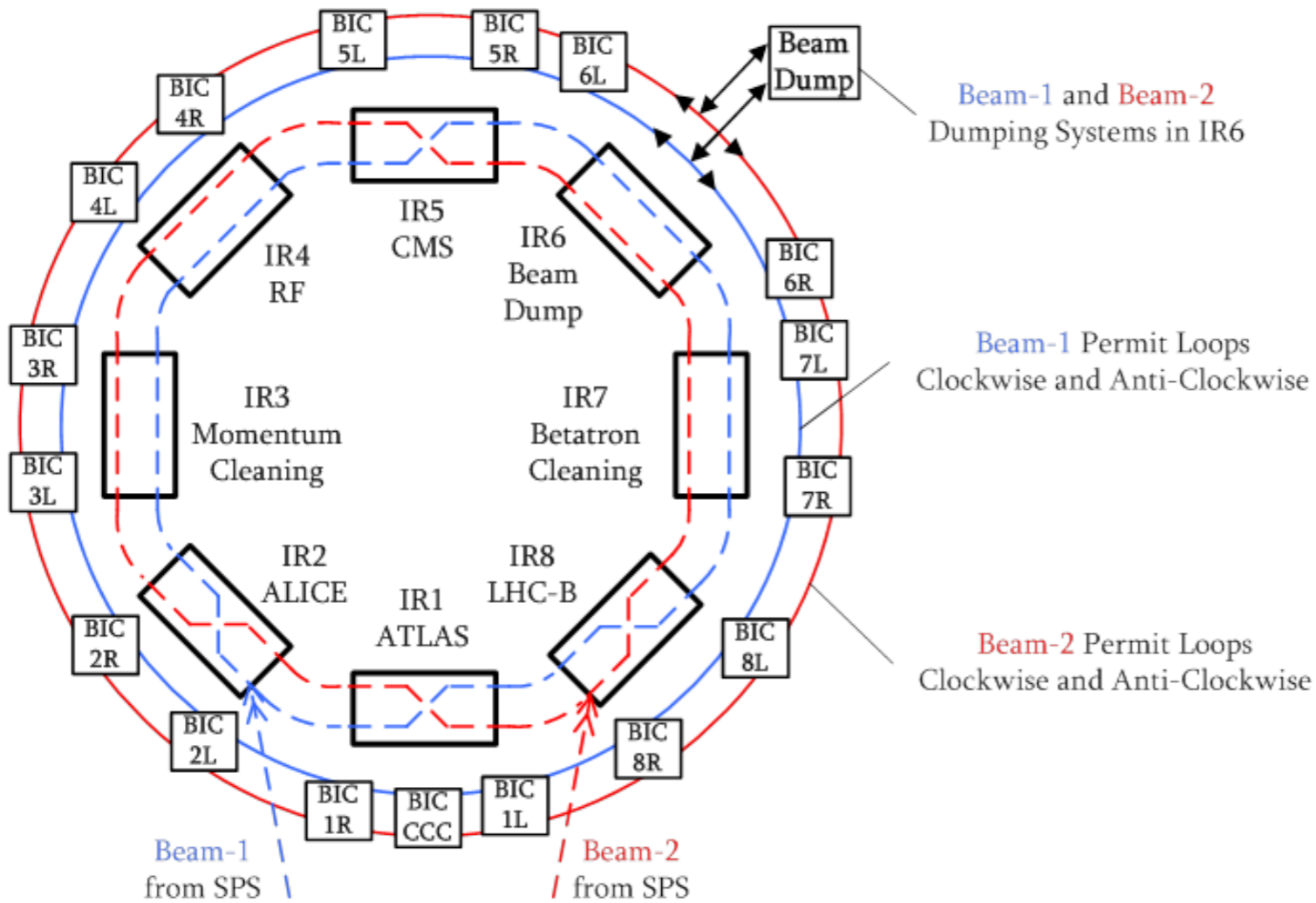
Extraction kicker



In practice....



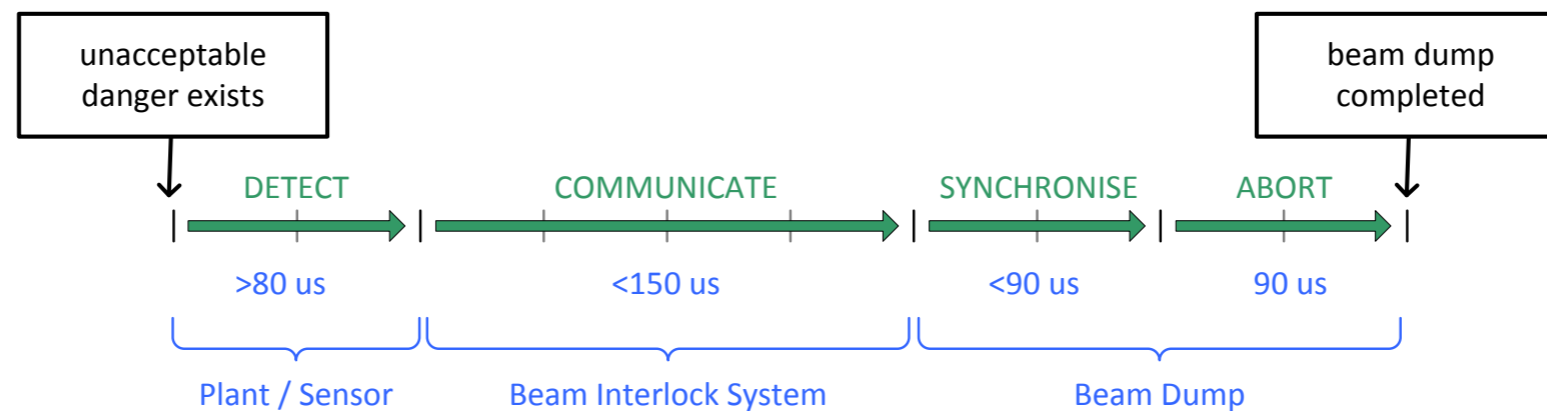
Beam interlock implementation



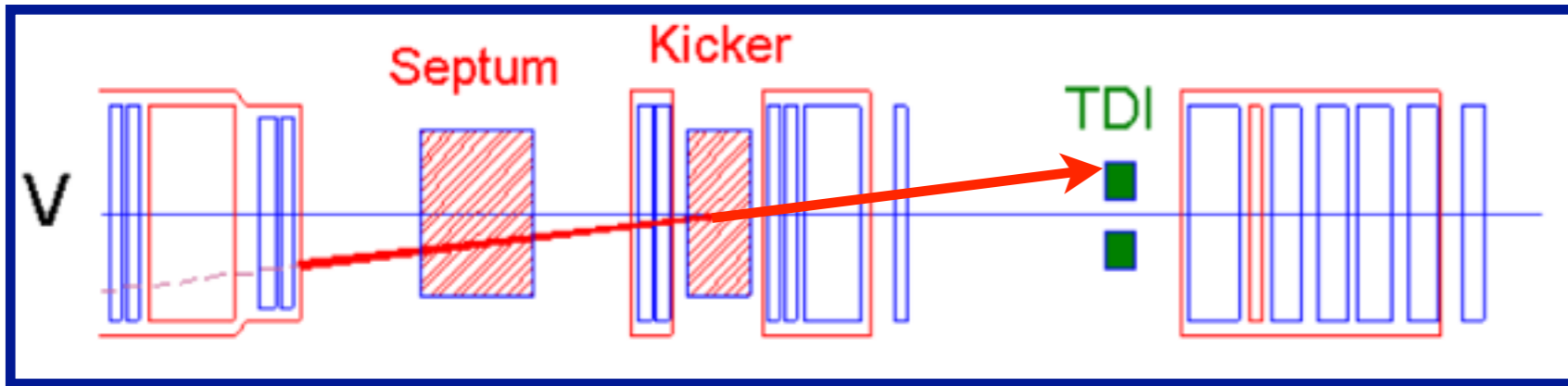
- The BPLs are connected to the LHC beam dumping system: a dump is triggered as soon as the signal of a single BPL is stopped.
- The BPLs are also connected to the LHC injection and SPS extraction interlock systems (same hardware design).

At the LHC the dump delay can reach ~3 turns ~300 μ s

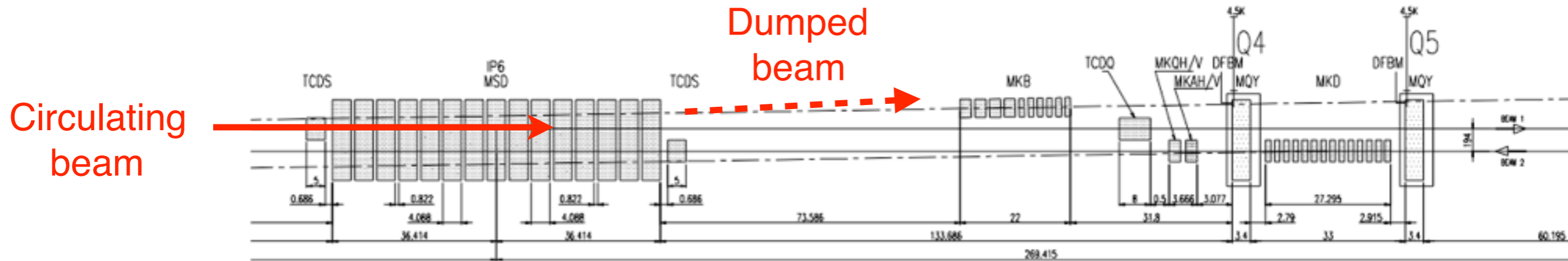
Revolution time
= 89 μ s



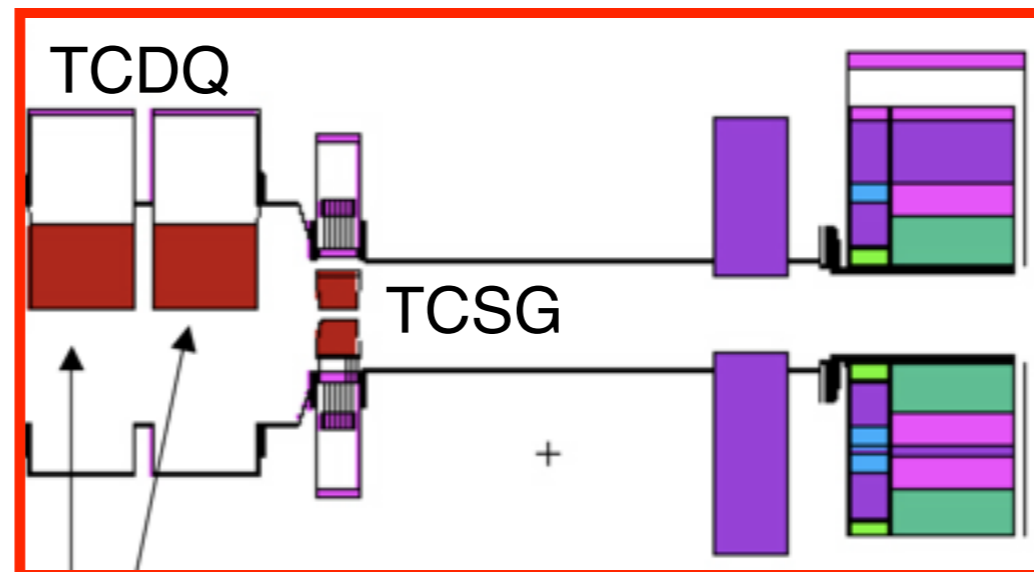
“Passive” protection



Injection protection
(introduced yesterday)



Dump protection
elements: movable
“TCDQ” and “TCSG”
collimators + fixed masks



TCDQ extraction protection



TCDQ = 9-m long collimator based on carbon that intercepts fractions of the beam that might be mis-kicked in case of problems with the beam dump!

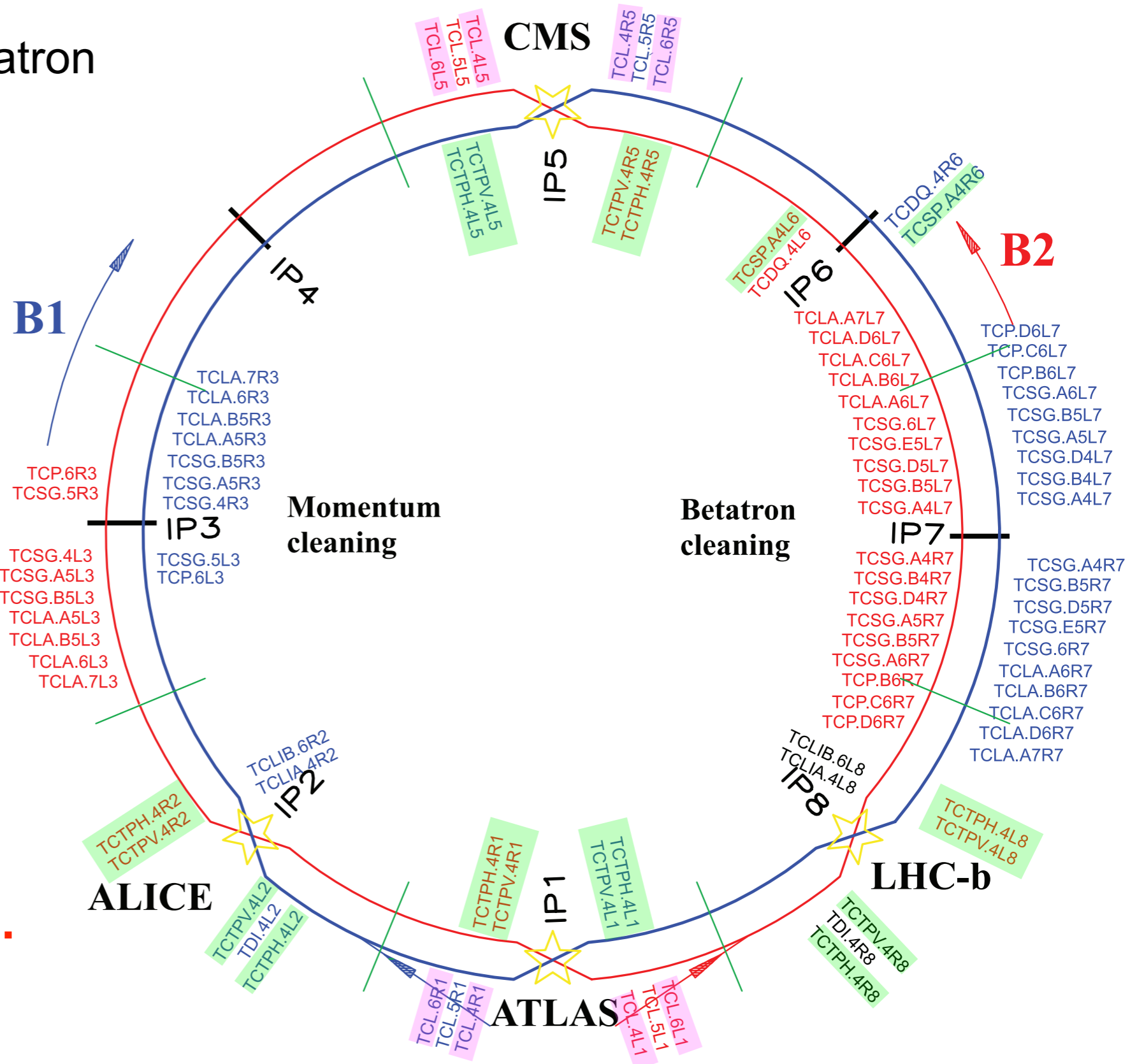
LHC collimation system

Dedicated insertions for betatron (IR7) and momentum (IR3) cleaning systems.

Cleaning of incoming beam in all experiments.

Physics debris collimation in the high-lumi IR1/5.

Total of 118 [was 108 in Run I] collimators (108 [was 100] movable).



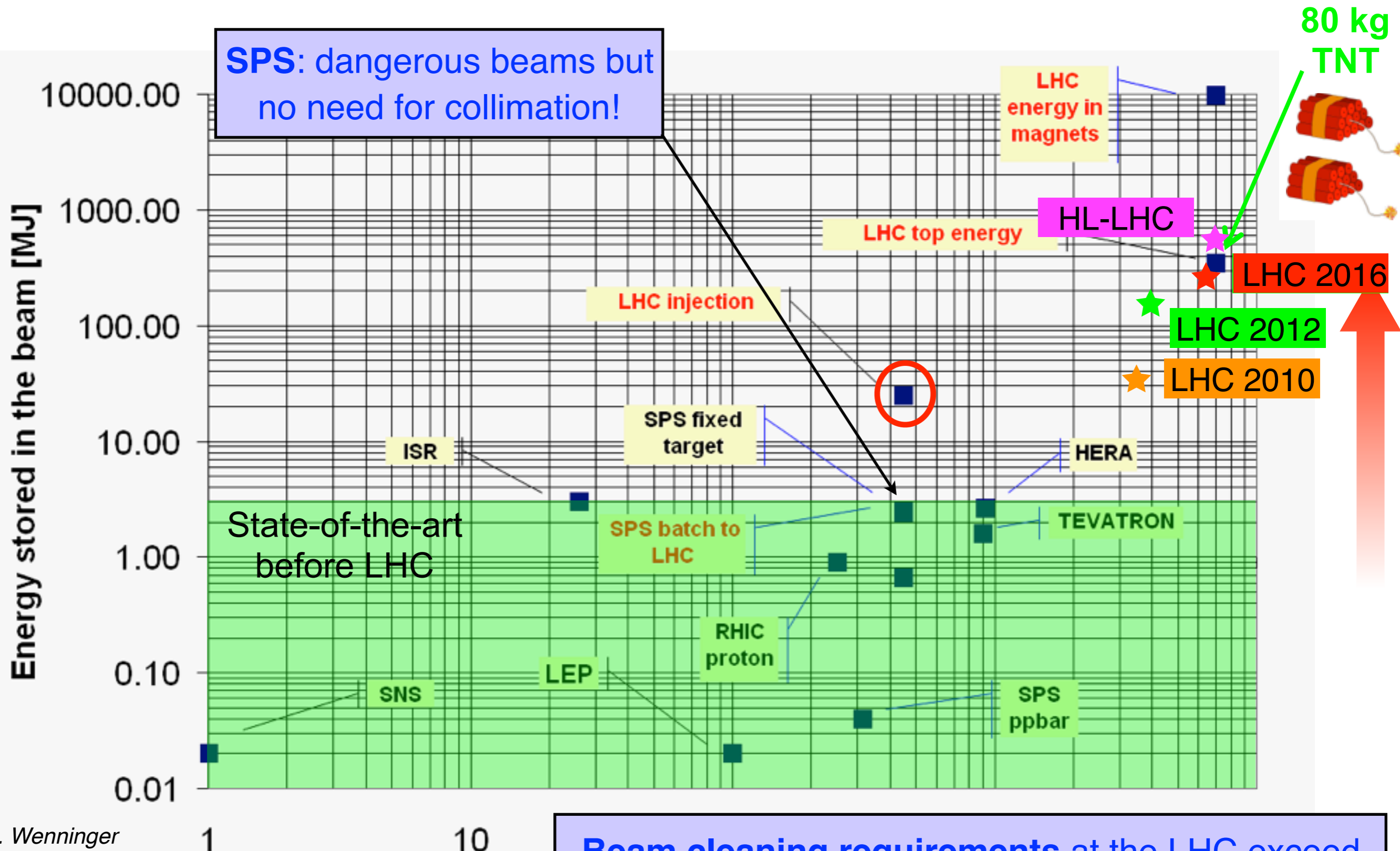
Where we are

- ☑ We have seen how the LHC beam requirements are met by the CERN accelerator complex
- ☑ We have introduced the main LHC accelerator systems
- ☑ We have introduced the key parameters for the LHC magnet system and for the LHC beam and seen how they determine the machine protection constraints. Driven by the quest for pushing luminosity of high-energy beams!
- ☑ We have presented the basic machine protection philosophy and some key implementations.
- ☑ We have introduced the collimation system as part of the passive protection.

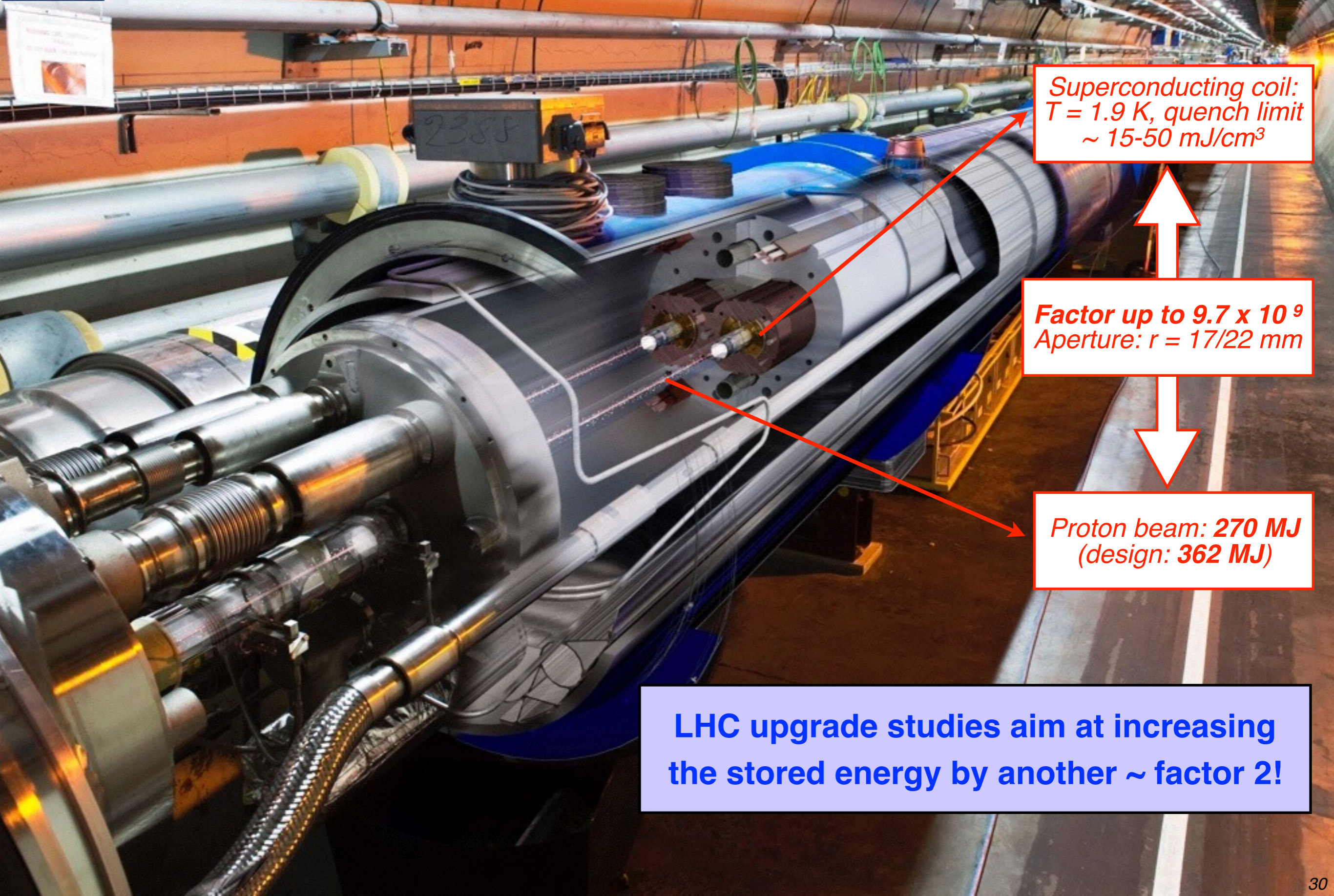
We will now see in detail the LHC collimation system!

- Main points from 1st lecture
- Machine protection and collimation
 - Concepts and LHC implementation
 - Case study: 2008 event
- **Beam losses and collimation**
 - **Roles of beam collimation systems**
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The stored energy challenge



1232 NbTi superconducting dipole magnets – each 15 m long
Magnetic field of 8.3 T (current of 11.8 kA) @ 1.9 K (super-fluid Helium)



*Superconducting coil:
T = 1.9 K, quench limit
~ 15-50 mJ/cm³*

*Factor up to 9.7×10^9
Aperture: r = 17/22 mm*

*Proton beam: 270 MJ
(design: 362 MJ)*

LHC upgrade studies aim at increasing the stored energy by another ~ factor 2!



The LHC collimator

Left jaw

Right jaw

1.0m+0.2m tapering



BEAM



The LHC collimator

Left jaw

Right jaw

1.0m+0.2m tapering

What is beam collimation and why we need it?
How do we design a collimation system?
How many collimators are needed?
Where are they located in the machine?
How are they built, with which materials?
How to measure and simulate cleaning?

BEAM

Beam halo collimation

*Controlled and safe disposal of **beam halo particles** produced by unavoidable beam losses.*

Achieved by reducing the transverse cross section of the beam.

Betatron (and off-momentum) **halo particles**

Particles with large betatron amplitudes (or energy deviations) with respect to the beam's reference particle.

Gaussian beams: typically, particles above 3 RMS beam sizes.

collimate /'kɒlɪ,meɪt/

VB (transitive)

1. to adjust the line of sight of (an optical instrument)
2. to use a collimator on (a beam of radiation or particles)
3. to make parallel or bring into line

Etymology: 17th Century: from New Latin *collimāre*, erroneously for Latin *collīnēare* to aim, from *com-* (intensive) + *līnēare*, from *līnea* line

collimator /'kɒlɪ,meɪtə/

N

1. a small telescope attached to a larger optical instrument as an aid in fixing its line of sight
2. an optical system of lenses and slits producing a nondivergent beam of light, usually for use in spectrosopes
3. any device for limiting the size and angle of spread of a beam of radiation or particles

Beam halo collimation

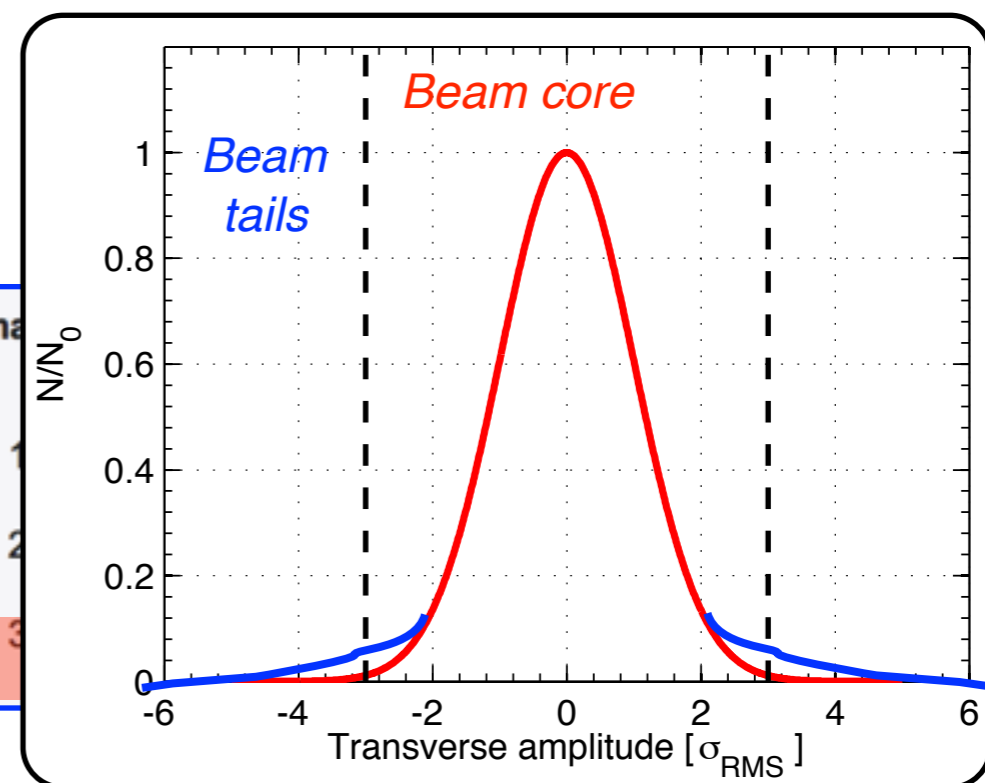
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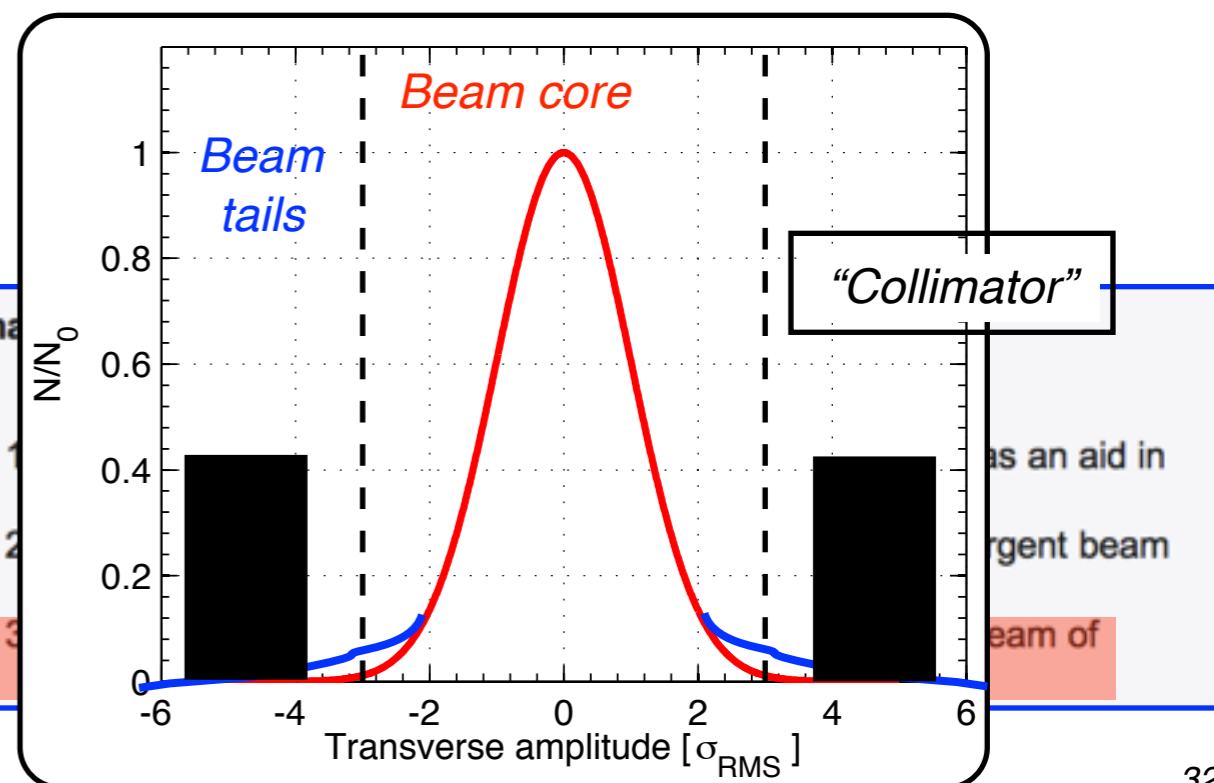
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Betatron (and off-momentum) **halo particles**

Particles with large betatron amplitudes (or energy deviations) with respect to the beam's reference particle.

Gaussian beams: typically, particles above 3 RMS beam sizes.

There are different goals of **collimation systems** depending on the machine.



collimate /'kɒlɪ, meɪt/

VB (transitive)

1. to adjust the line of sight of (an optical instrument)
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- **Halo cleaning** versus quench limits (super-conducting machines)
- Passive **machine protection**
First line of defence in case of accidental failures.
- **Concentration of losses/activation** in controlled areas
Ease maintenance by avoiding many distributed high-radiation areas.
- **Reduction total doses** on accelerator equipment
Provide local protection to equipment exposed to high doses (like the warm magnets in cleaning insertions)
- **Cleaning of physics debris** (physics products, in colliders)
Avoid magnet quenches close to the high-luminosity experiments
- Optimize **background** in the experiments
Minimize the impact of halo losses on quality of experimental data
- Beam tail/halo **scraping, halo**
Control and probe the transverse

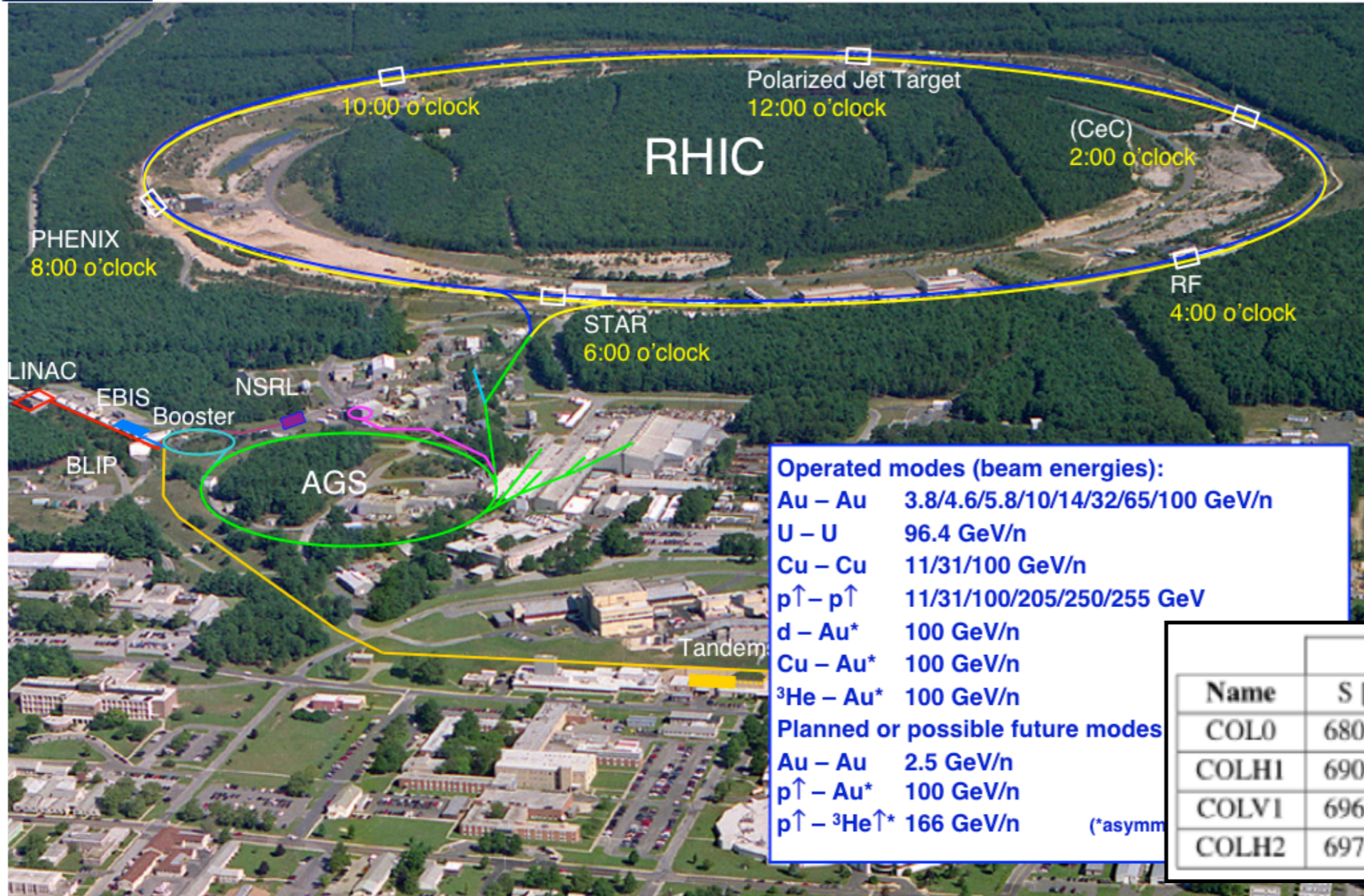
→ Main role of collimation in previous hadron colliders (SppS, Tevatron, ...)

This lecture: focus **collimation cleaning** functionality. LHC examples as a case study because all these roles are addressed !



***Why is the LHC
so special for
collimation
matters?***

RHIC collimation system



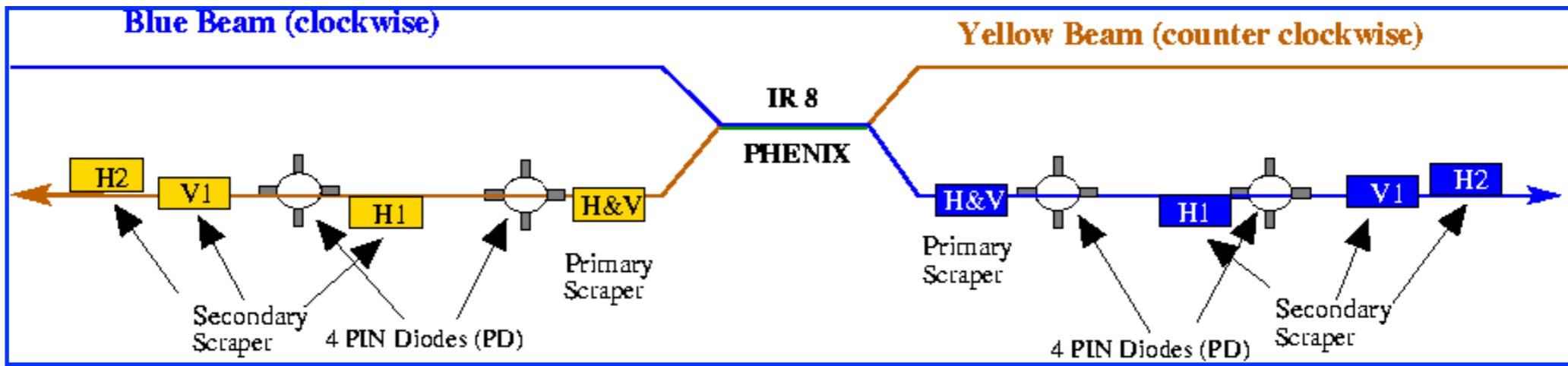
Operated modes (beam energies):
 Au – Au 3.8/4.6/5.8/10/14/32/65/100 GeV/n
 U – U 96.4 GeV/n
 Cu – Cu 11/31/100 GeV/n
 p[↑] – p[↑] 11/31/100/205/250/255 GeV
 d – Au* 100 GeV/n
 Cu – Au* 100 GeV/n
³He – Au* 100 GeV/n

Planned or possible future modes
 Au – Au 2.5 GeV/n
 p[↑] – Au* 100 GeV/n
 p[↑] – ³He[↑]* 166 GeV/n (*asymm)

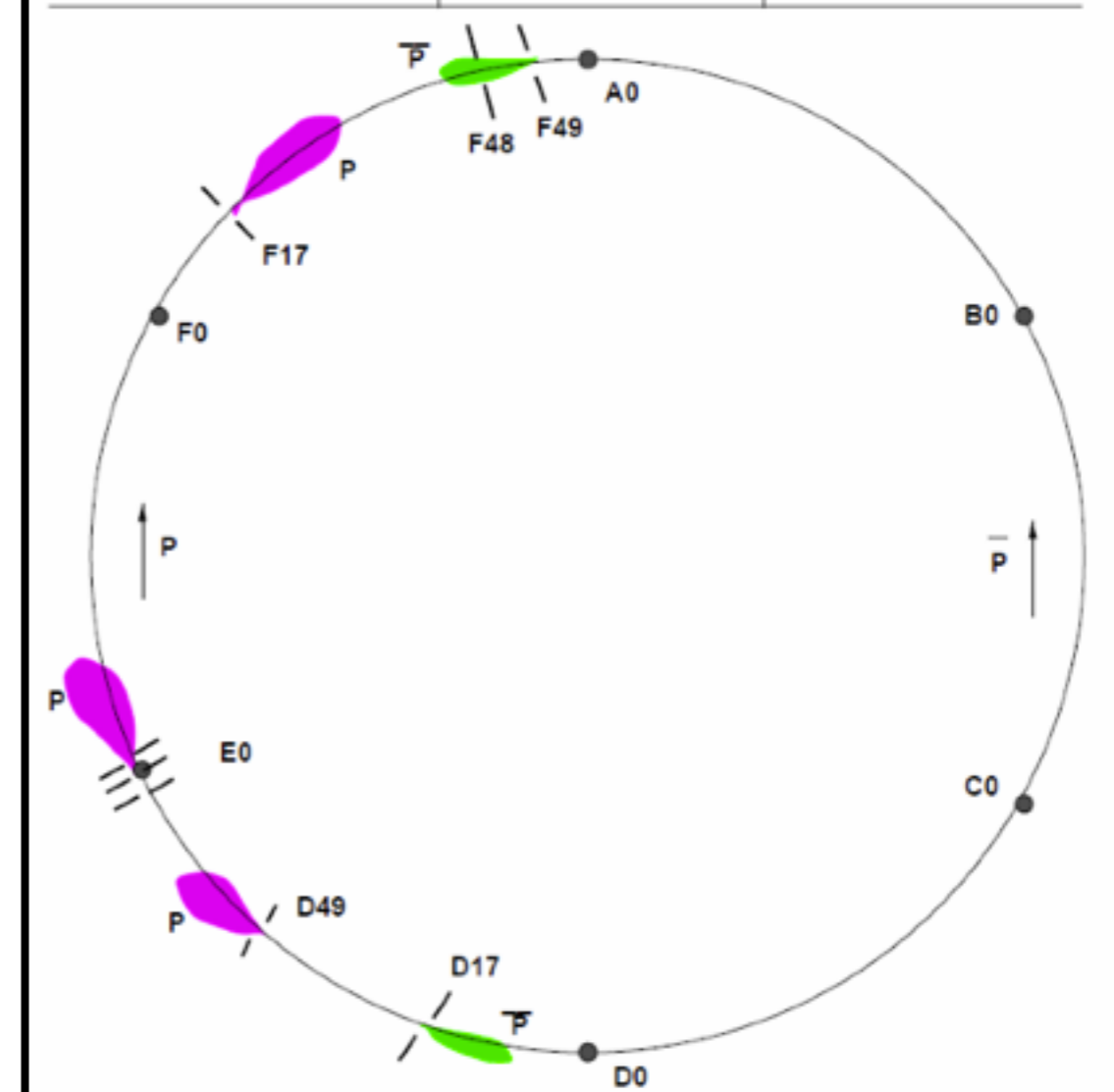
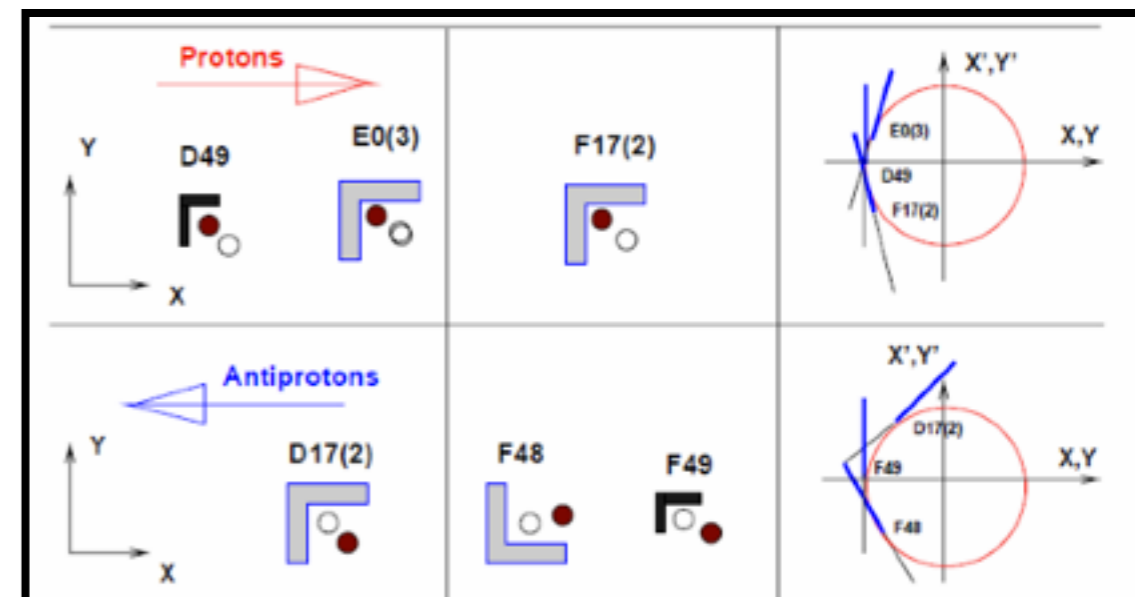
RHIC beam parameters [p]:
 $E_b = 250 \text{ GeV}$
 $N_{tot} = 110 \times 10^{11} p$
 $E_{stored} = \sim 440 \text{ kJ}$

Collimation system:
 8 collimators
 Some with L shape

Name	Blue		Yellow	
	S [m]	Plane	S [m]	Plane
COL0	680.752	Hor. + Vert.	3236.649	Hor. + Vert.
COLH1	690.533	Horizontal	3246.430	Horizontal
COLV1	696.706	Vertical	3252.603	Vertical
COLH2	697.728	Horizontal	3253.625	Horizontal



Tevatron Run II collimation system



Tevatron Run II parameters:

$$E_b = 1 \text{ TeV}$$

$$E_{\text{stored}} = \sim 2 \text{ MJ}$$

Collimation system:

13 collimators, L shape

26 positional degrees of freedom

Collimation of LEP collider

LEP parameters - e^+e^- collider:

$$E_b = 45-105 \text{ GeV}$$

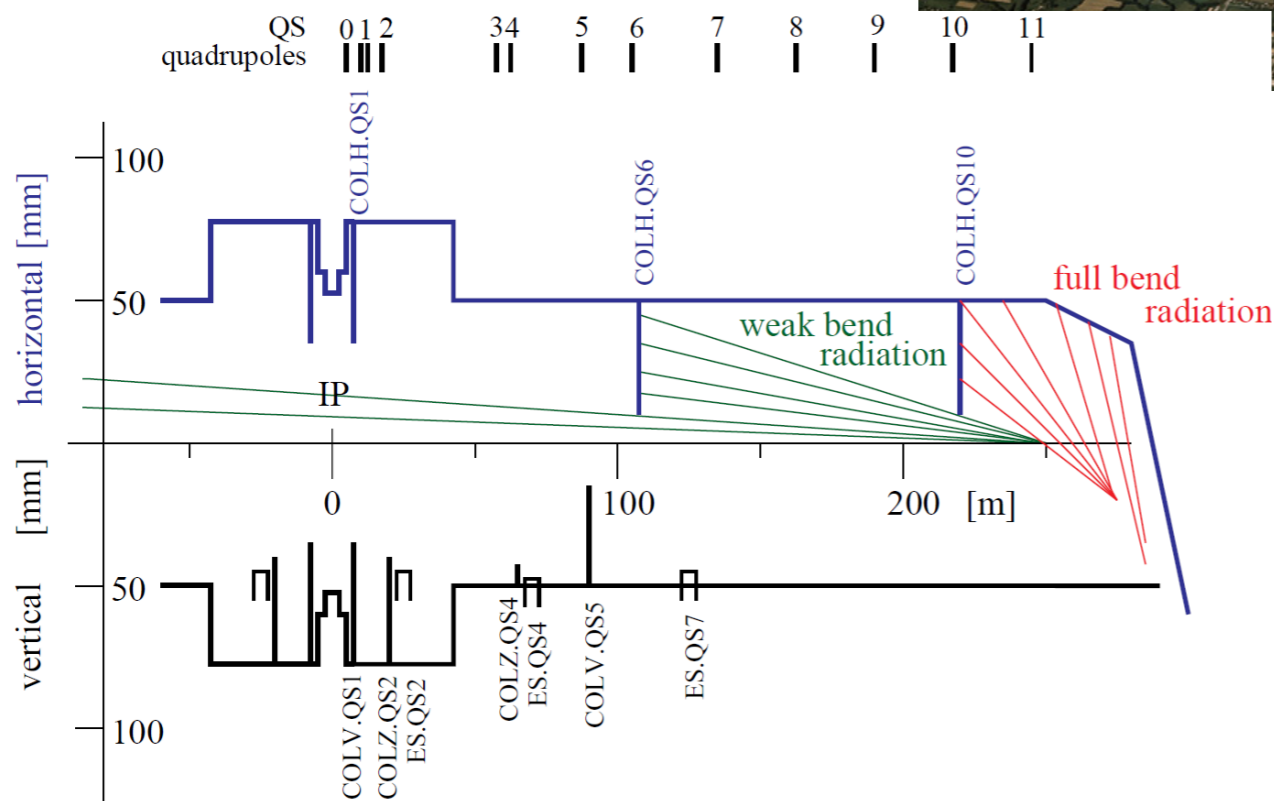
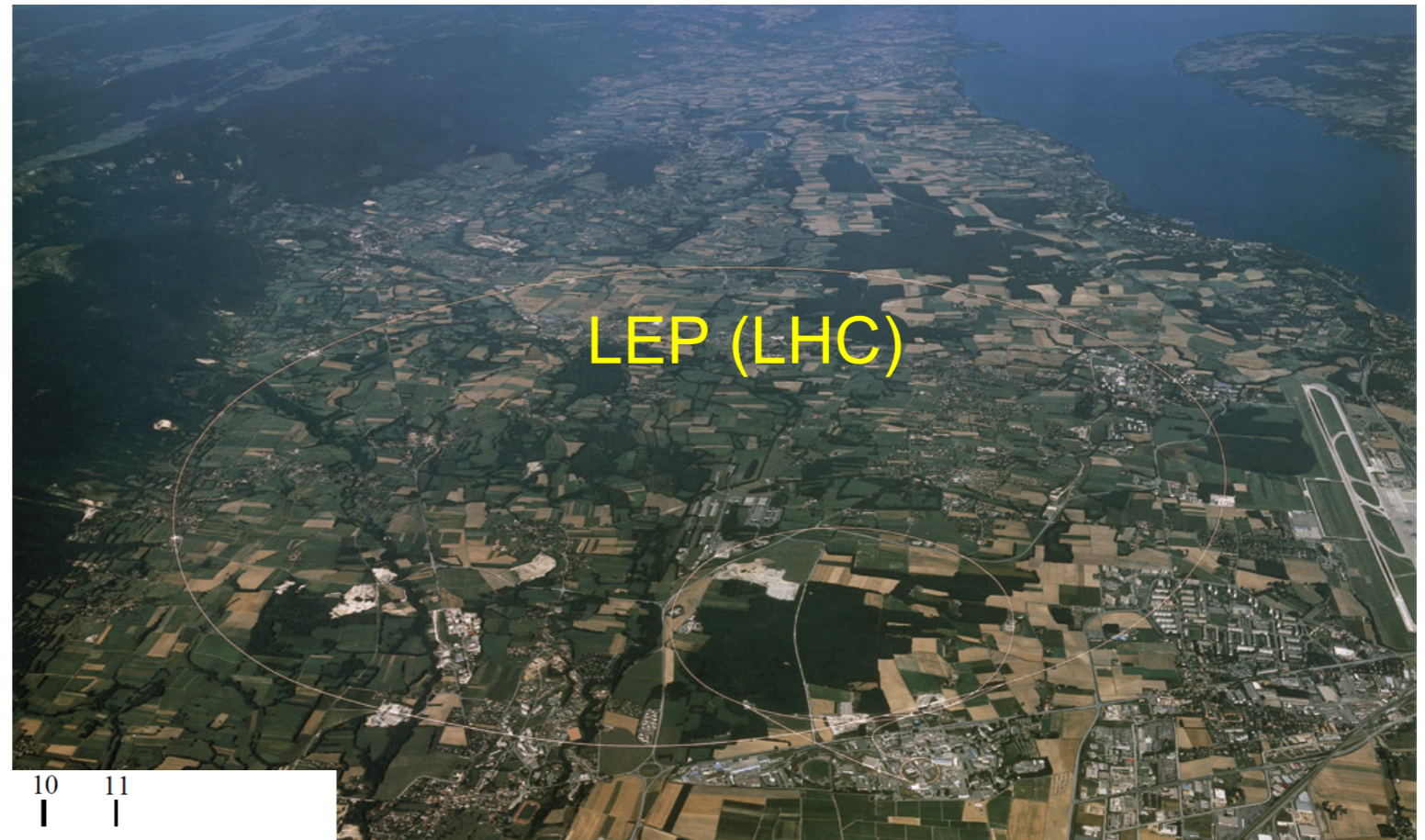
$$I_{\text{bunch}} = 4 \times 10^{11} \text{ e}^+/\text{e}^-$$

$$I_{\text{tot}} = 1.6 \times 10^{12} \text{ e}^+/\text{e}^-$$

$$E_{\text{stored}} = \sim 25 \text{ kJ}$$

$$\text{Bunch spacing} = 11 \mu\text{s}$$

$$\text{Synchrotron radiation power} \\ \sim 10 \text{ MW / beam}$$



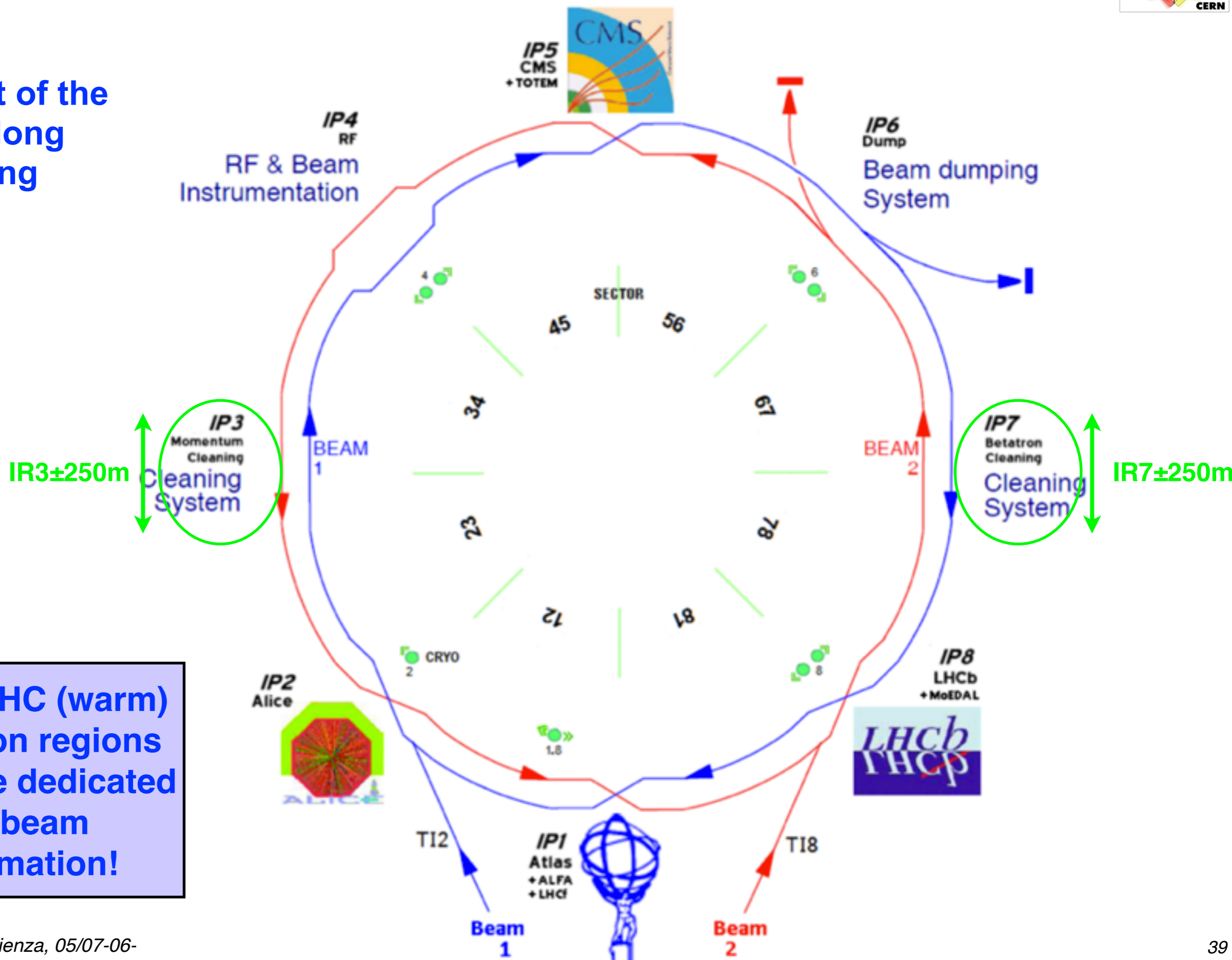
LEP collimation system:

96 collimators (mostly 2 jaw),
Betatron and off-energy,
Local masks at the experiments

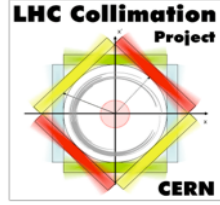
G. von Holtey et al, CERN-SL 97-40

LHC ring layout

Layout of the 27km-long LHC ring



2 of 8 LHC (warm) insertion regions (IRs) are dedicated to beam collimation!



LHC collimation layout

Collimation designed for nominal LHC design parameters:

$$E_b = 7 \text{ TeV}$$

$$I_{\text{bunch}} = 1.15 \times 10^{11} \text{ p}$$

$$I_{\text{tot}} = 3.2 \times 10^{14} \text{ p}$$

$$E_{\text{stored}} = 362 \text{ MJ}$$

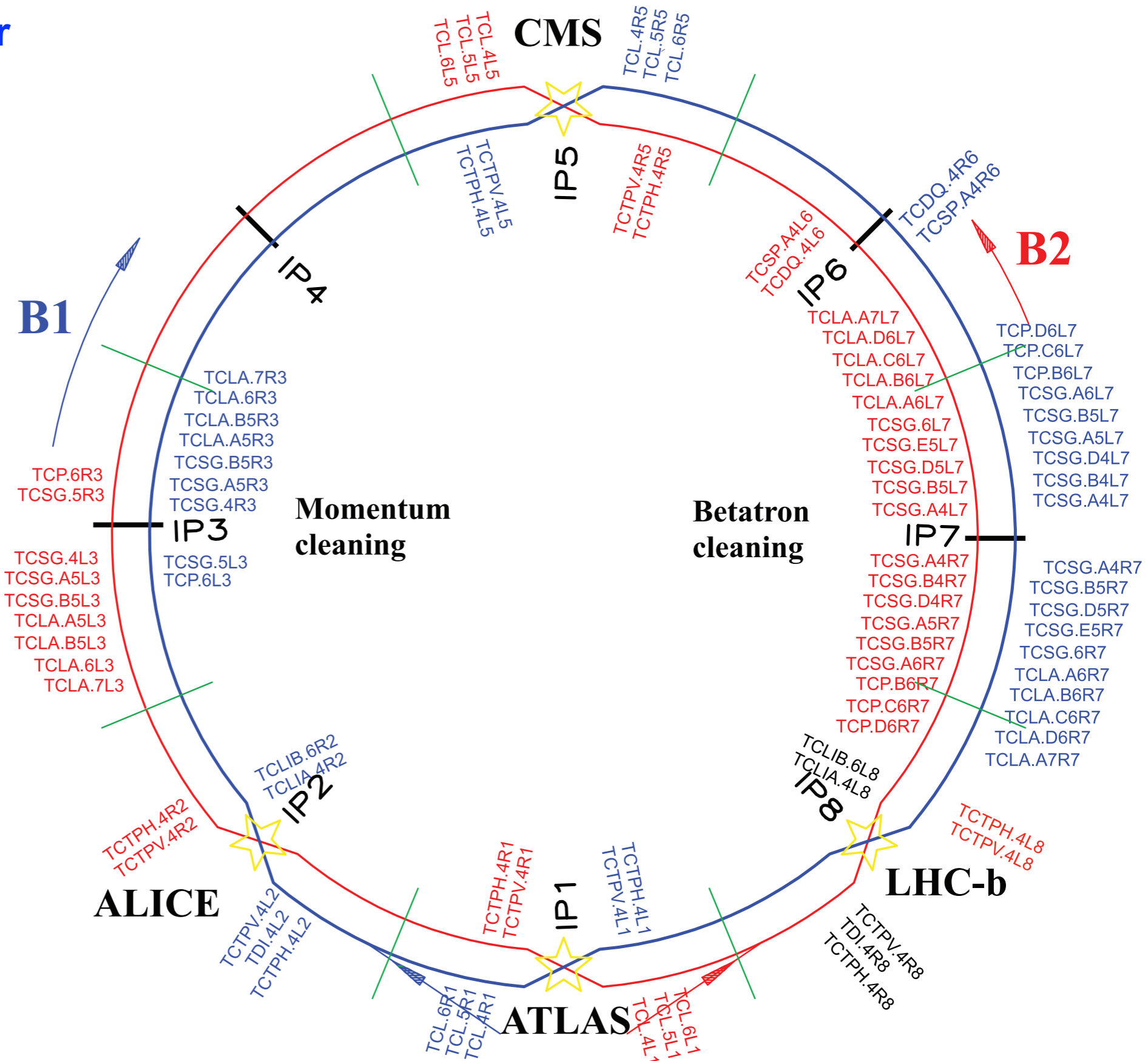
$$\text{Bunch spacing} = 25 \text{ ns}$$

Achieved:

$$E_b = 6.5 \text{ TeV}$$

$$E_{\text{stored}} = 270 \text{ MJ}$$

Total of 118 two-sided collimators (108 are movable, 4 motors each).



Why so many collimators?

It is **difficult to “stop”** high-energy hadrons and the energy that they carry!

You have seen that in previous lectures...

There are **many different loss mechanisms** that impose the deployment of **different solutions** for beam collimation, machine protection, optics scenarios etc.

Betatron losses in horizontal, vertical and diagonal planes require full “phase-space” coverage.

Momentum losses occur in different locations than betatron’s.

Different types of failures, slow and fast regimes, etc...

Collimators closest to the beams are made of **low-Z materials** (higher robustness at the expenses of absorption power).

Several collimators (respecting a well-defined hierarchy) are installed in ~500 m long warm insertions (LHC case).

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Ideal world (perfect machine): no beam losses throughout the operational cycle

Injection, energy ramp, betatron squeeze, collisions, beam dump.

No need for a collimation system!

In **real machines**, several effects cause **beam losses**:

- **Collisions** in the interaction points (beam burn up)
- Interaction with **residual gas** and **intra-beam scattering**
- **Beam instabilities** (single-bunch, collective, beam-beam)
- Dynamics changes during OP cycle (orbit drifts, optics changes, energy ramp, ...): “**operational losses**”
- Transverse **resonances**.
- Capture losses at beginning of the ramp.
- RF noise and out-of-bucket losses.
- Injection and dump losses.

We do not need to study all that in detail to understand beam collimation!

These effects can increase the **beam halo population** and ultimately cause beam losses!

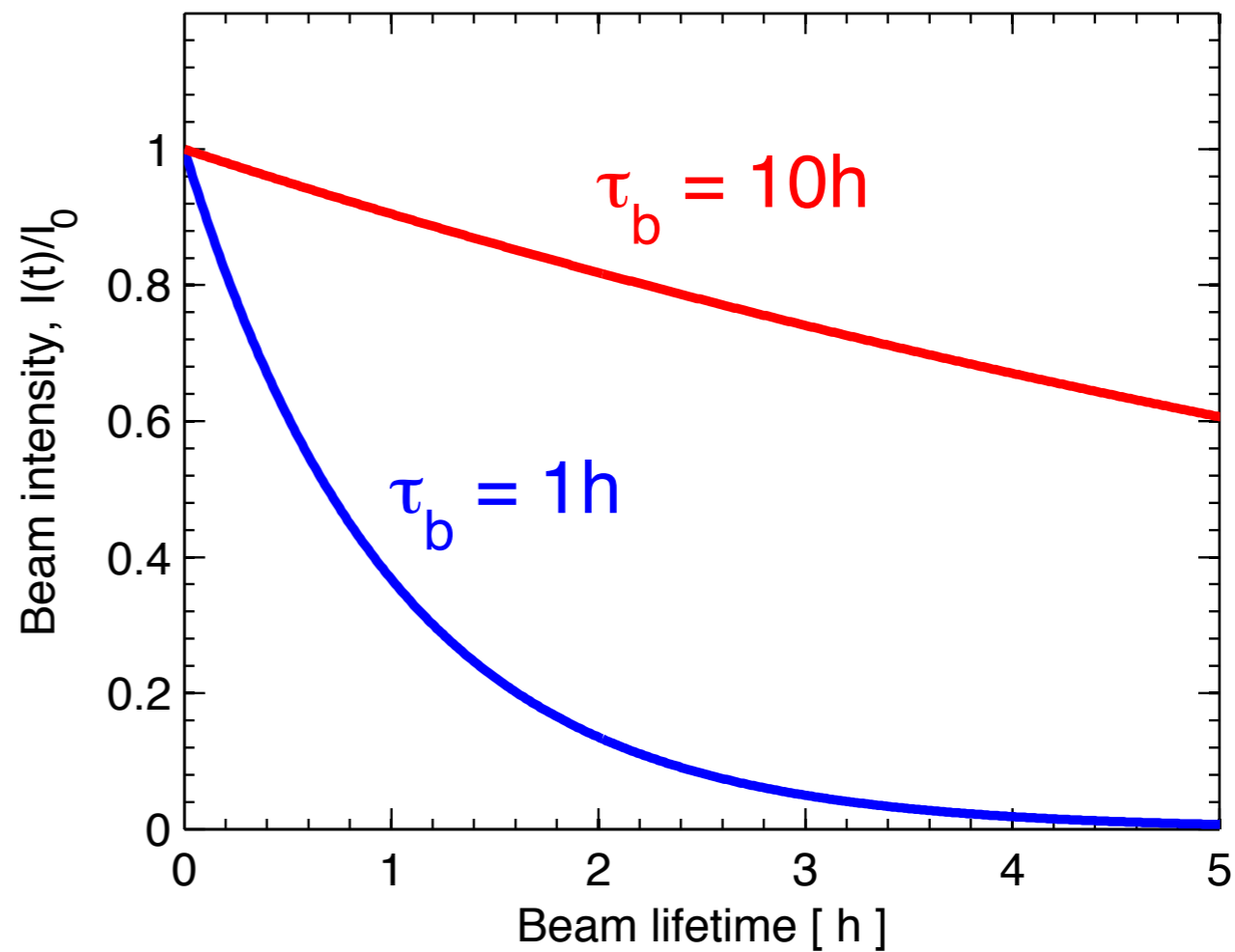
Beam loss mechanisms are modelled by assuming a non-infinite **beam lifetime**, τ_b

$$I(t) = I_0 \cdot e^{-\frac{t}{\tau_b}}$$

: Beam intensity versus time

$$-\frac{1}{I_0} \frac{dI}{dt} = \frac{1}{\tau_b}$$

: Proton loss rate

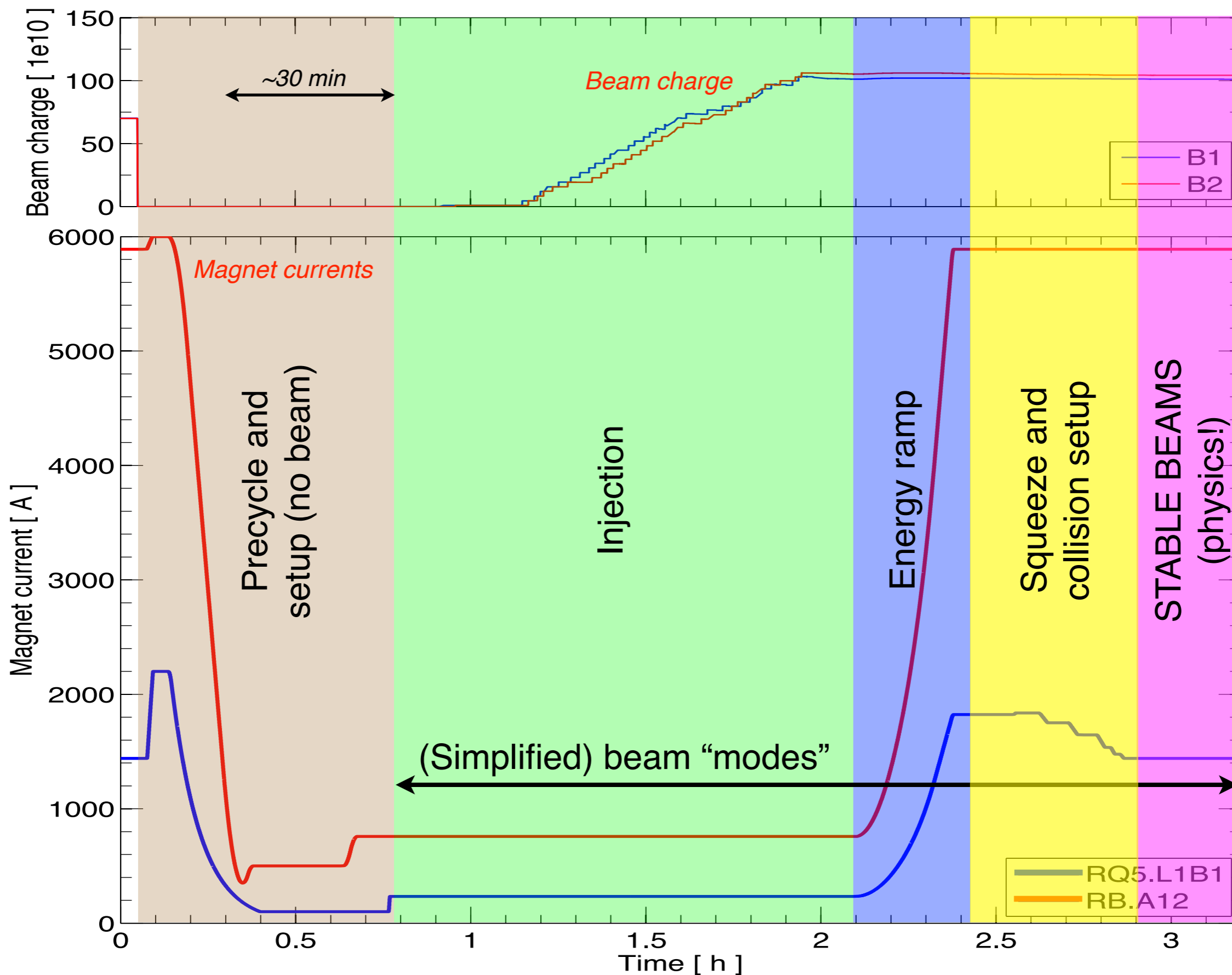


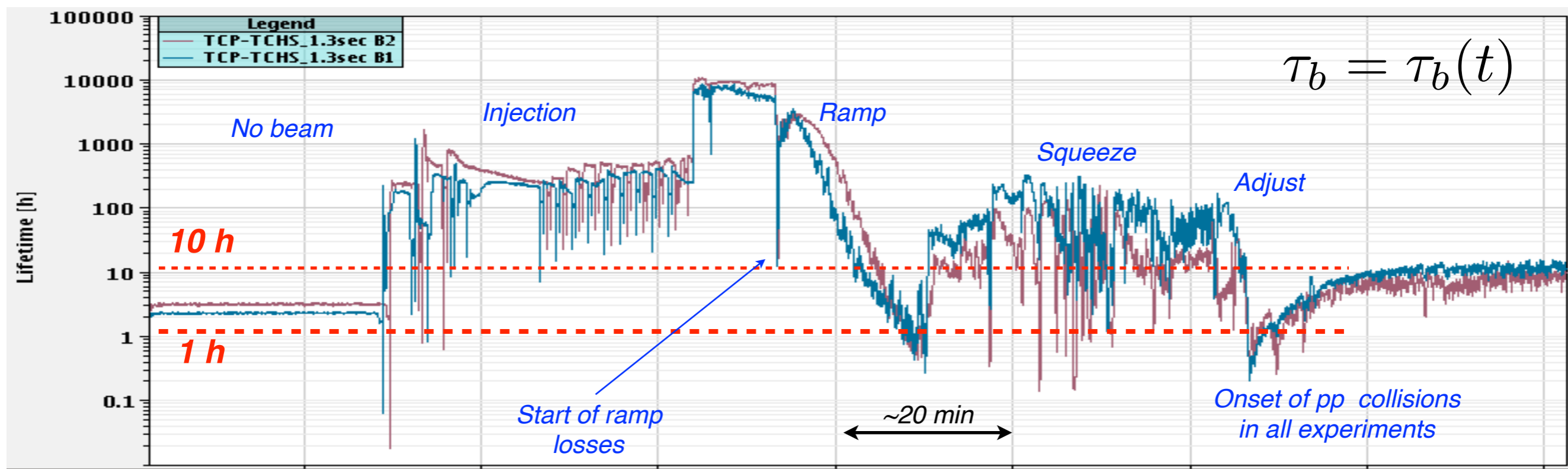
Beam losses mechanisms are characterized by a time-dependent **beam lifetime** during the machine cycle. This measures the **total beam losses** that a collimation system must handle.

*Example at 7 TeV: **1h lifetime** at the full intensity of 3.2×10^{14} protons (320 hundred trillion protons!) corresponds to a loss rate of about 90 billion proton per second, i.e. **0.1MJ/s = 100 kW!***



Operational cycle of a collider





Example of a typical physics fill in 2012.

What matters is the minimum lifetime → see peaks below 1 h!

At 7 TeV, this corresponds to peak losses **larger than 100 kW** that would be lost in the cold aperture. They **must be caught** before!!

Goal of a collimation system: catch these losses and ensure that a controlled fraction of them reaches sensitive equipment.

Collimation “**inefficiency**” → measures the fraction of beam losses that goes into sensitive equipment out of the total lost from the beam.



Key collimation design parameters



In *real* machines affected by beam losses, we need a **collimation system** that intercepts the **primary beam losses** (“primary halo”) and absorbs the energy that they carries.

Collimation designed to handle losses that otherwise would occur in an uncontrolled way around the machine.

Design loss rates are calculated from the **total beam intensity** and **beam energy** assuming a “**minimum allowed beam lifetime**” that can occur during operation.

A **collimation cleaning inefficiency** is defined to express the fraction of the total losses that goes into sensitive equipment.

Cold magnets, warm magnets, experiments (background), ...

Example: losses versus quench limits

N_{tot} : total beam populations [p]

$\frac{N_{\text{tot}}}{\tau_b}$: proton loss rate [p/s]

R_q : quench limit [p/m/s]

Condition to operate the machine: losses in the magnets remain below their quench limit

$$\frac{N_{\text{tot}}}{\tau_b} \times \tilde{\eta}_c < R_q$$

$\tilde{\eta}_c$: local cleaning inefficiency [1/m] → fraction of proton losses that is lost at a certain location.

$\tilde{\eta}_c = \tilde{\eta}_c(s)$: this is a function on the longitudinal coordinate (as seen later).

For the 1h lifetime case shown before, we get a loss rate at the LHC of 90×10^9 p/s. Assuming a quench limit of $R_q \sim 3.2 \times 10^7$ p/m/s at 7 TeV, one can calculate a **required inefficiency of a few 10⁻⁴!!**

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$\tilde{\eta}_c$: local cleaning inefficiency

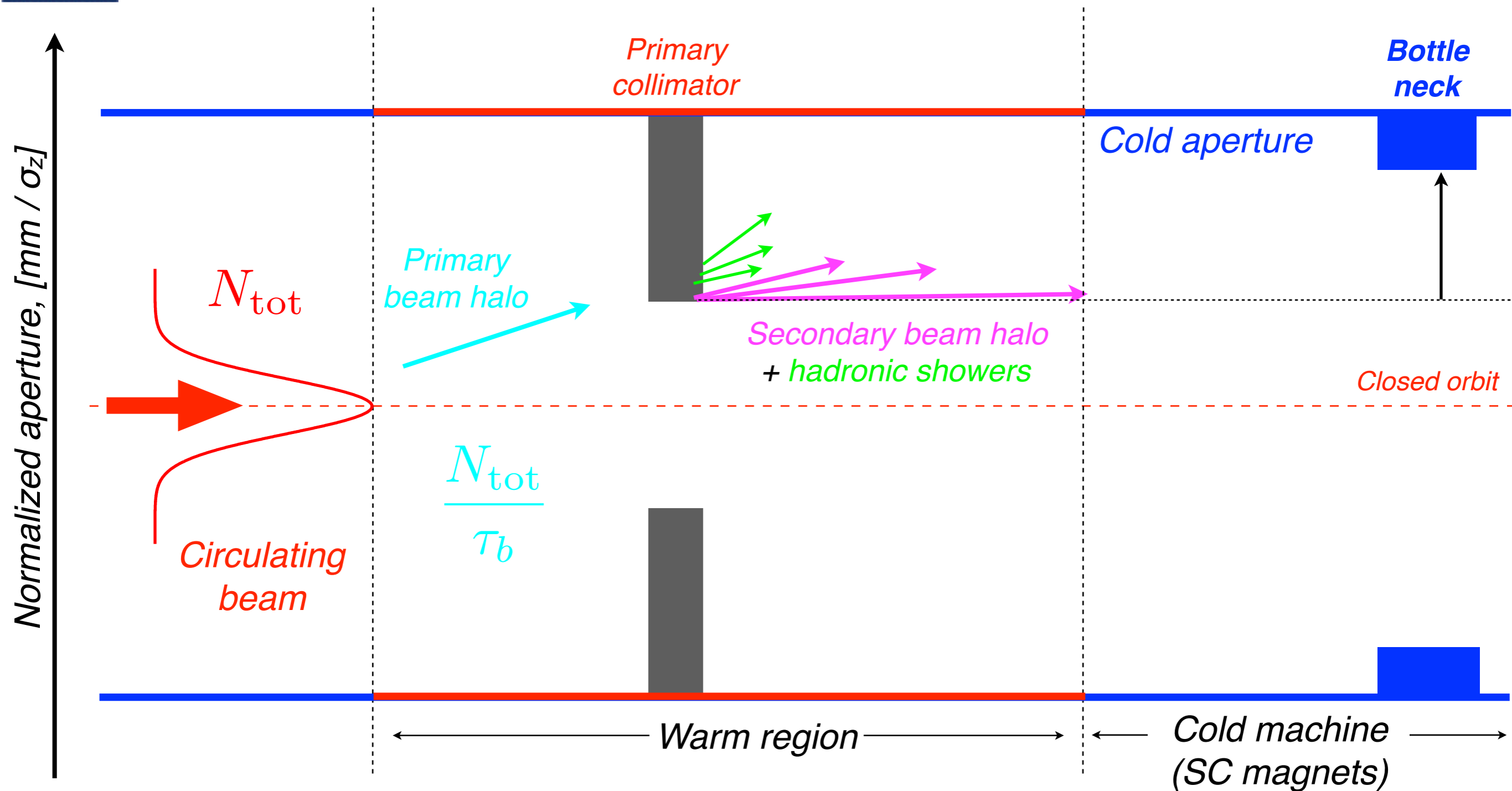
This is our **first specification** for the design of the collimation system. It can only be as good as the accuracy of “input” and “observable”...

$\tilde{\eta}_c = \tilde{\eta}_c(s)$: this is a function on the longitudinal coordinate (as seen later).

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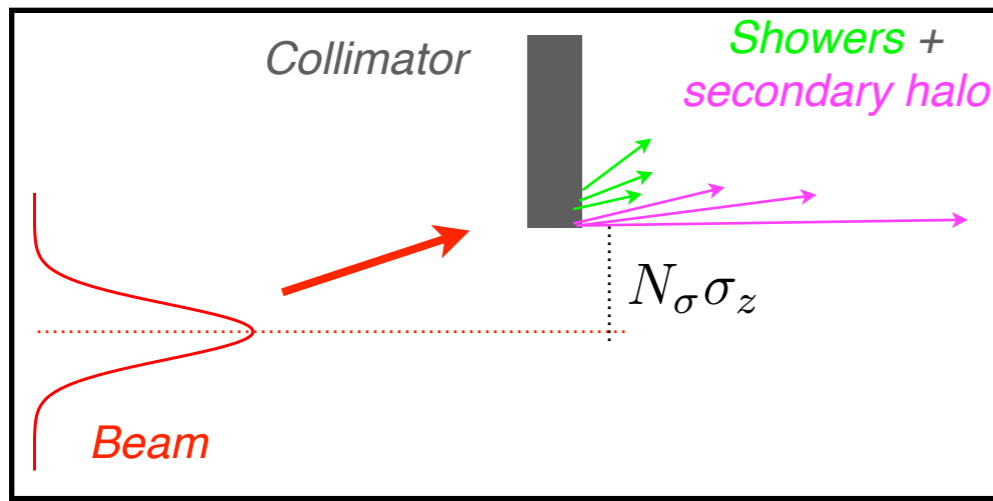
Aperture and single-stage cleaning



The particles lost from the beam core drift transversally and populate beam tails. Ultimately, they reach the machine *aperture bottleneck*.

Can we stop them with a single collimator that shields the cold aperture?

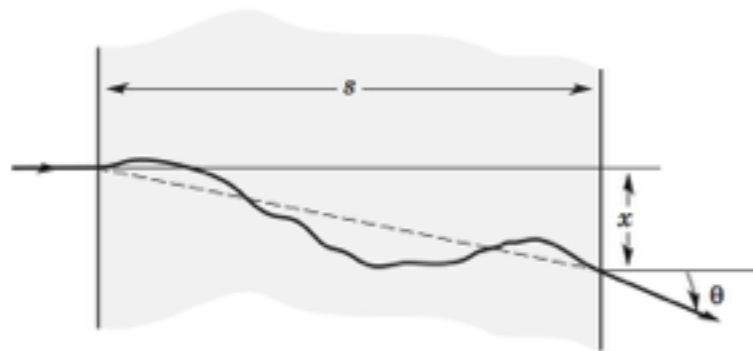
Particle interaction with collimators



If the “primary” collimator were a black absorber, it would be sufficient to shield the aperture by choosing a gap $N_\sigma \sigma_z$ smaller than the aperture bottleneck !

In reality, part of the beam energy and a fraction of the incident protons escape from the collimator!

For “cleaning” what matters is the energy leakage.

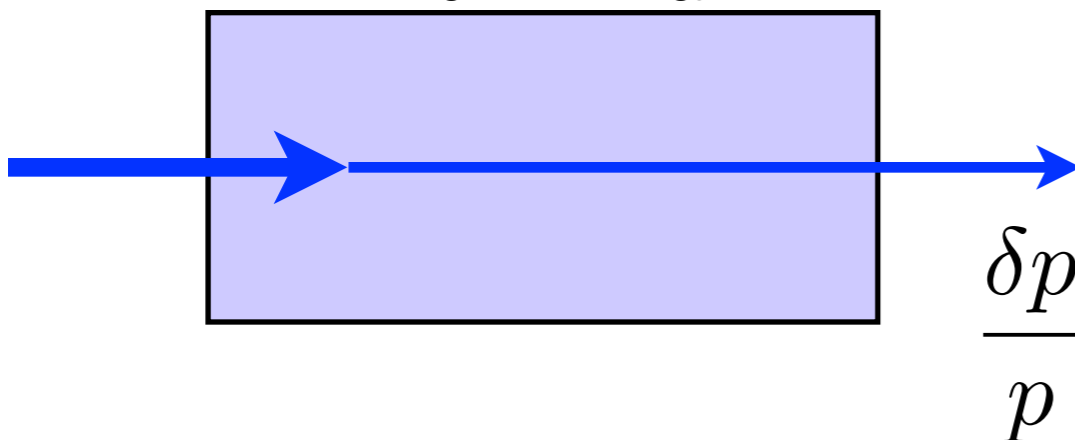


$$\sqrt{\langle \theta_p^2 \rangle} = \frac{13.6}{cp[\text{MeV}]} \sqrt{\frac{s}{\chi_0}} \left(1 + 0.038 \cdot \left(\frac{s}{\chi_0} \right) \right)$$

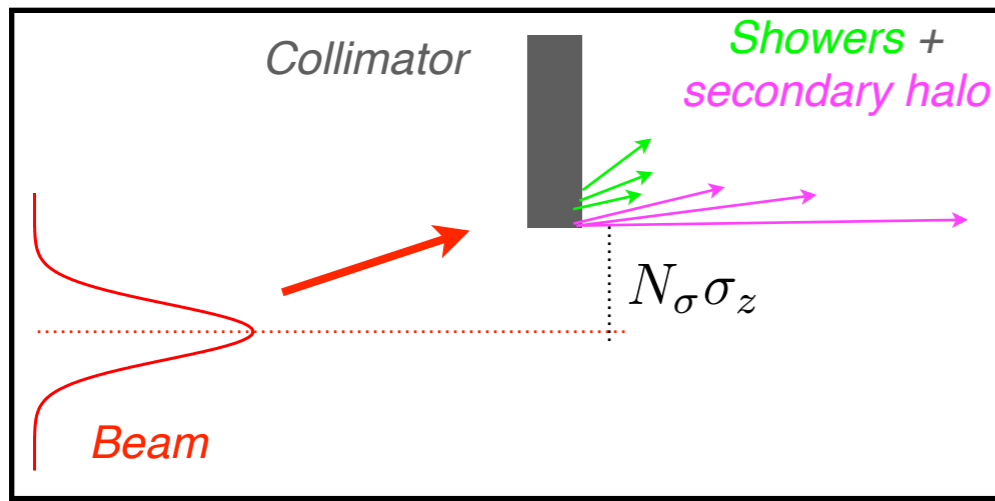
χ_0 : radiation length

Molière’s multiple-scattering theory: scattered particles gain a **transverse RMS kick**.

Single-diffractive interactions change the energy!



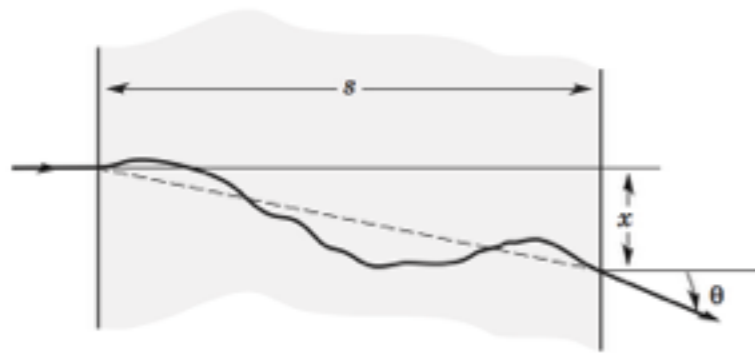
Some protons escape from the collimator with a reduced “rigidity” after losing energy through inelastic interactions.



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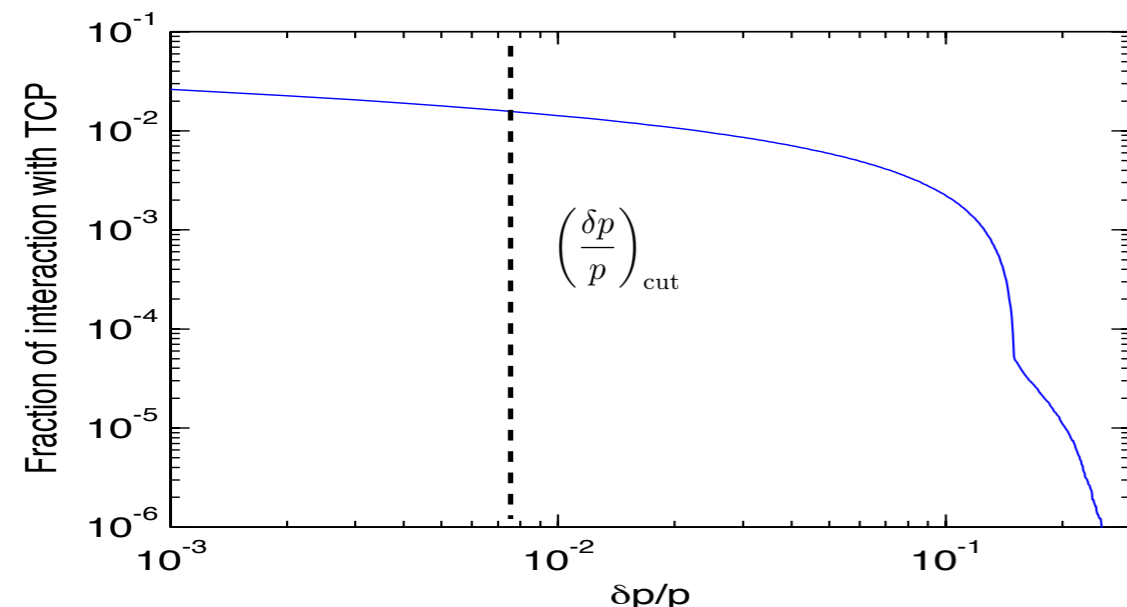


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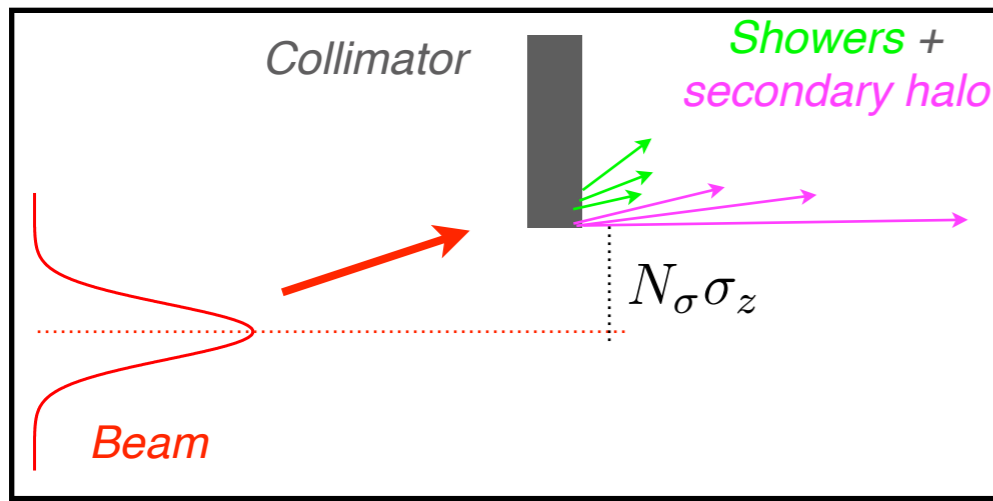
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Distribution of energy lost after multi-turn interaction with 60cm TCP



The interaction with collimator materials is itself a source of betatron and off-momentum halo (secondary halo).

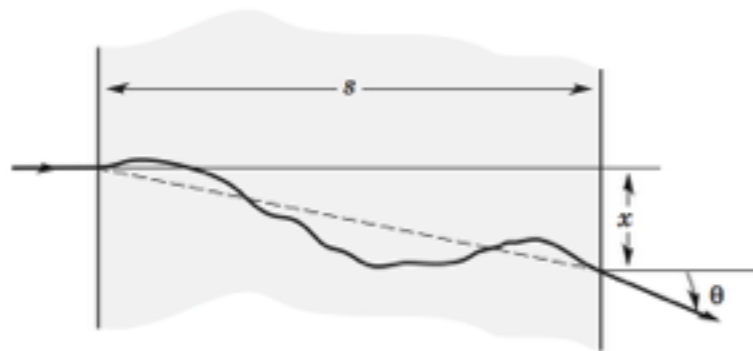
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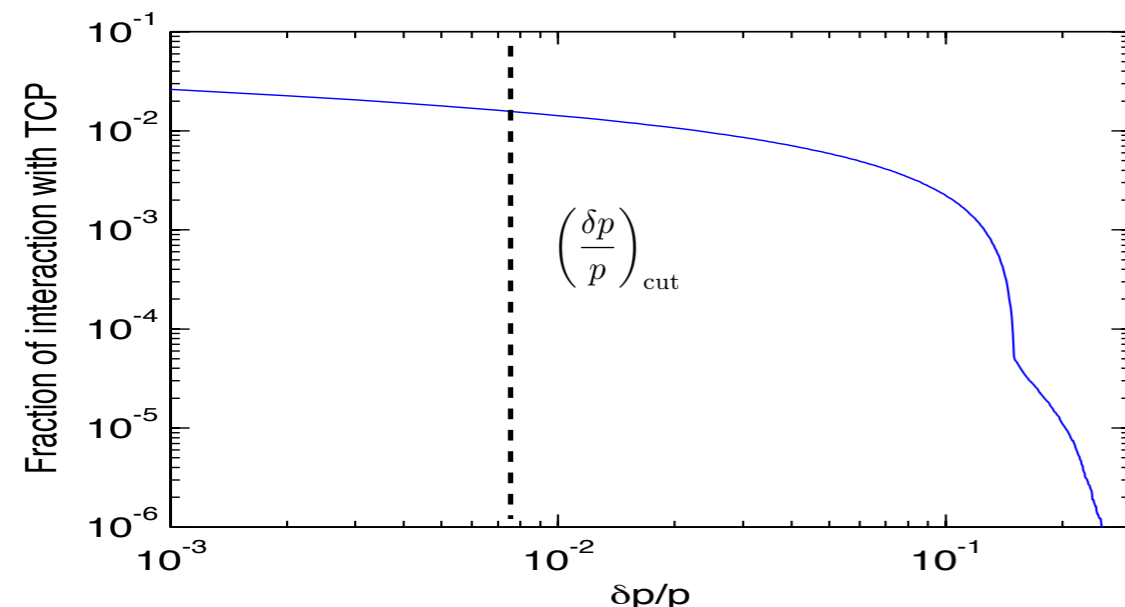


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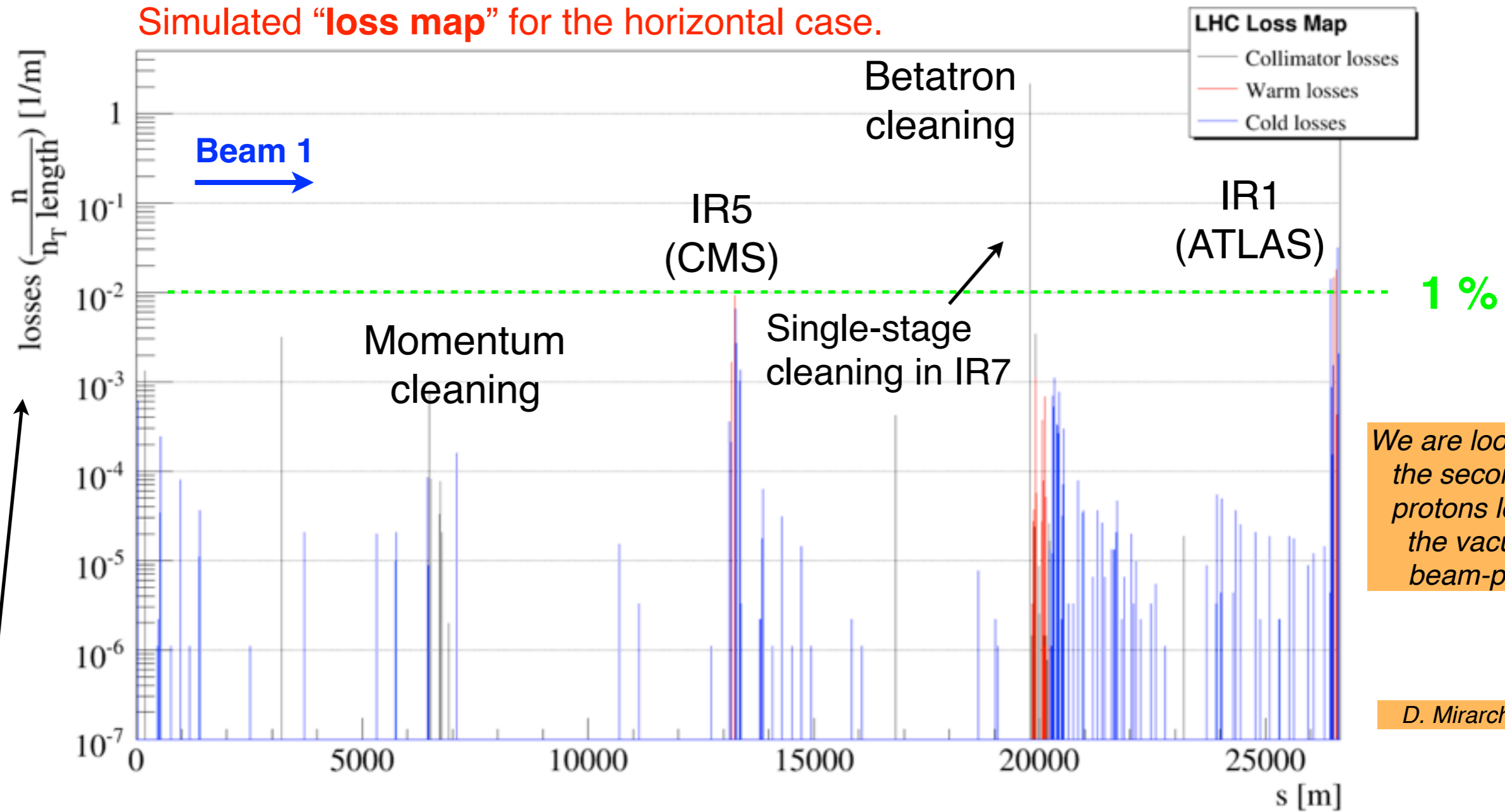
The interaction with collimator materials is itself a source of betatron and off-momentum halo (secondary halo).

Electro-magnetic and hadronic showers developed by the interaction carry an important fraction of the impacting beam energy that “escapes” from the collimator.

Note: multi-turn interactions occur with sub-micron impact parameters → this has an important effect on the absorption efficiency.

Single-stage cleaning - LHC at 7 TeV

Simulated “loss map” for the horizontal case.



We are looking at the secondary protons lost in the vacuum beam-pipe.

D. Mirarchi

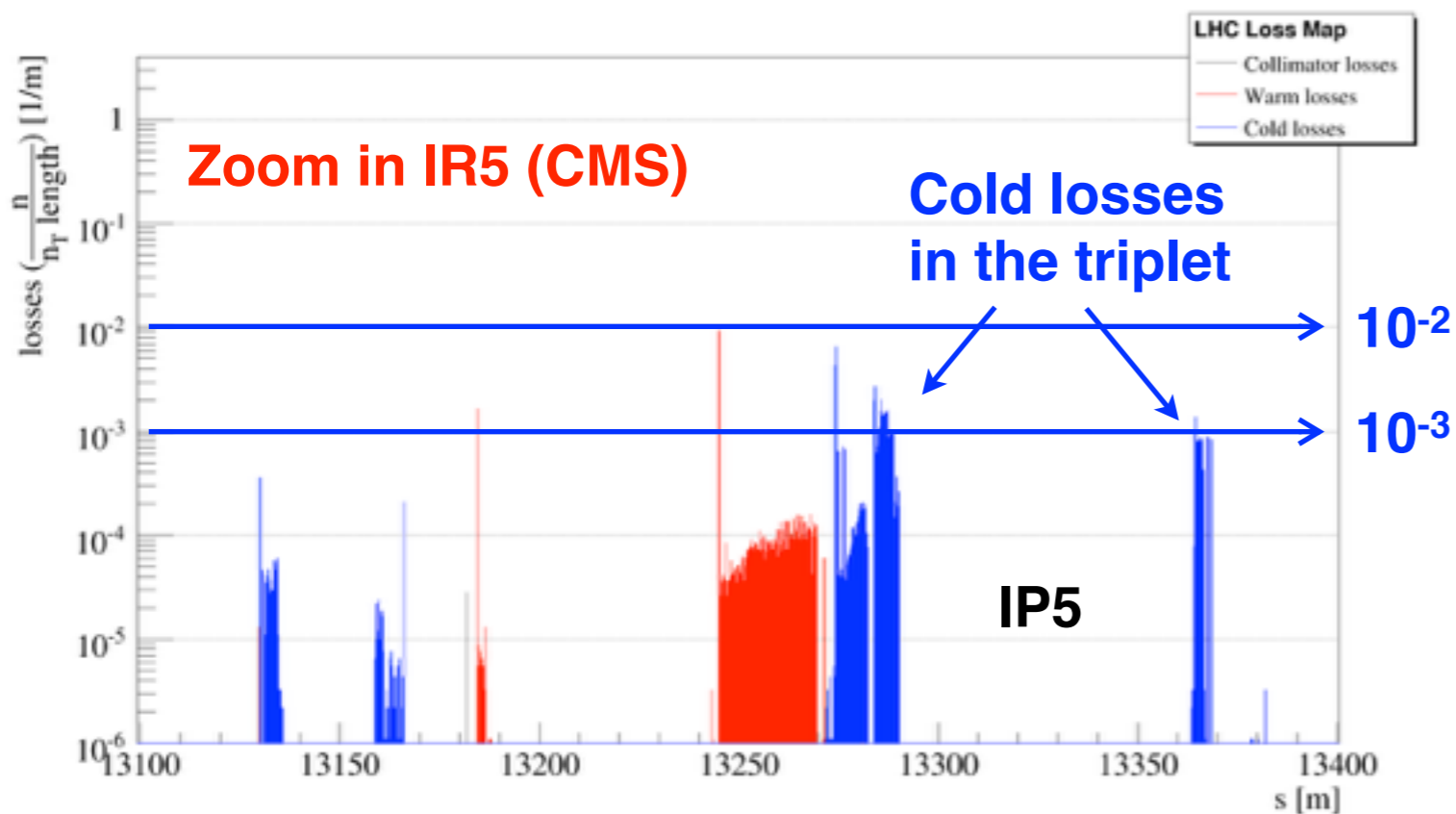
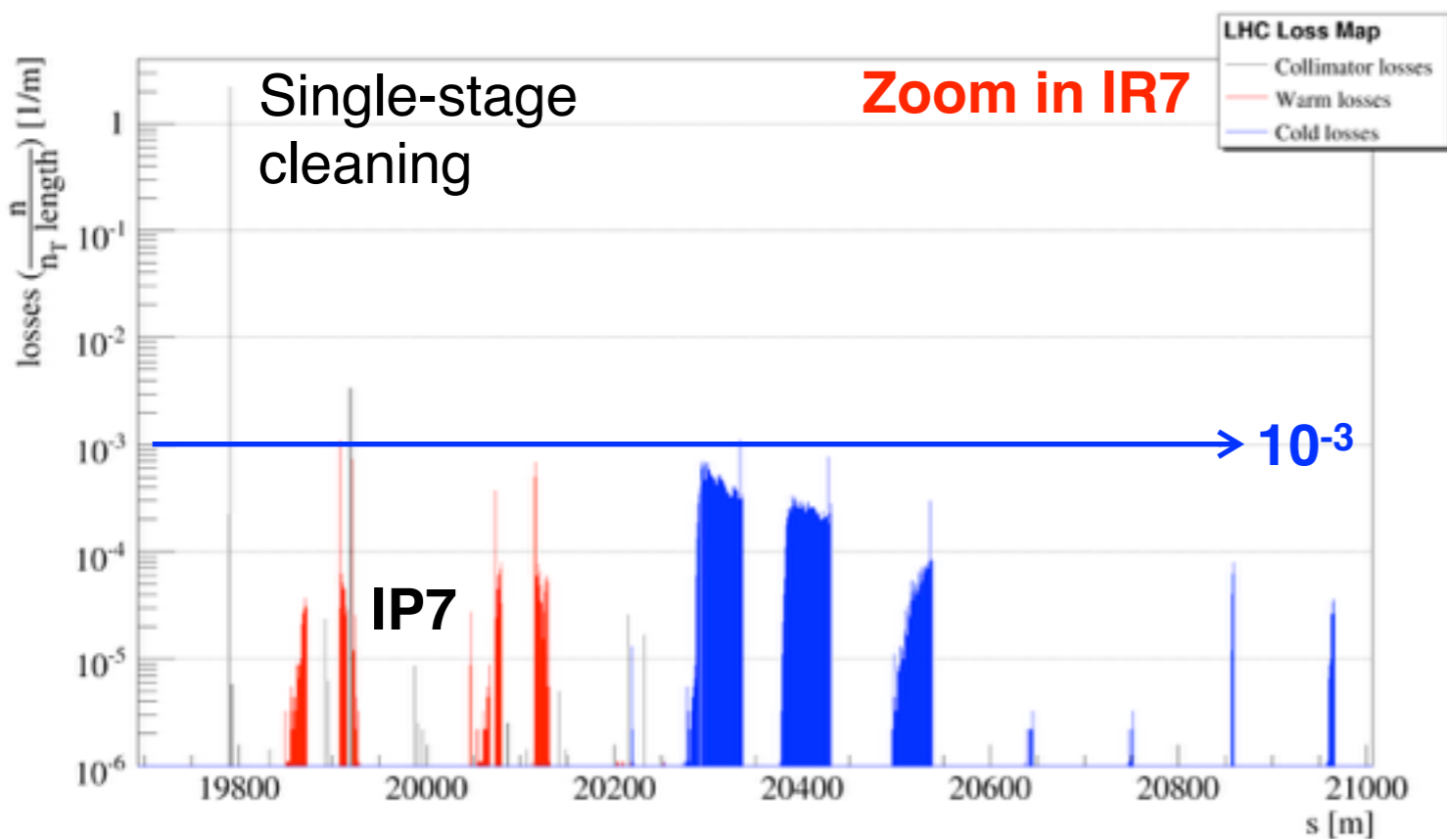
Local cleaning inefficiency

$$\tilde{\eta}_c(s) = \frac{1}{\Delta s} \frac{N_{\text{loss}}(s \rightarrow s + \Delta s)}{N_{\text{abs}}}$$

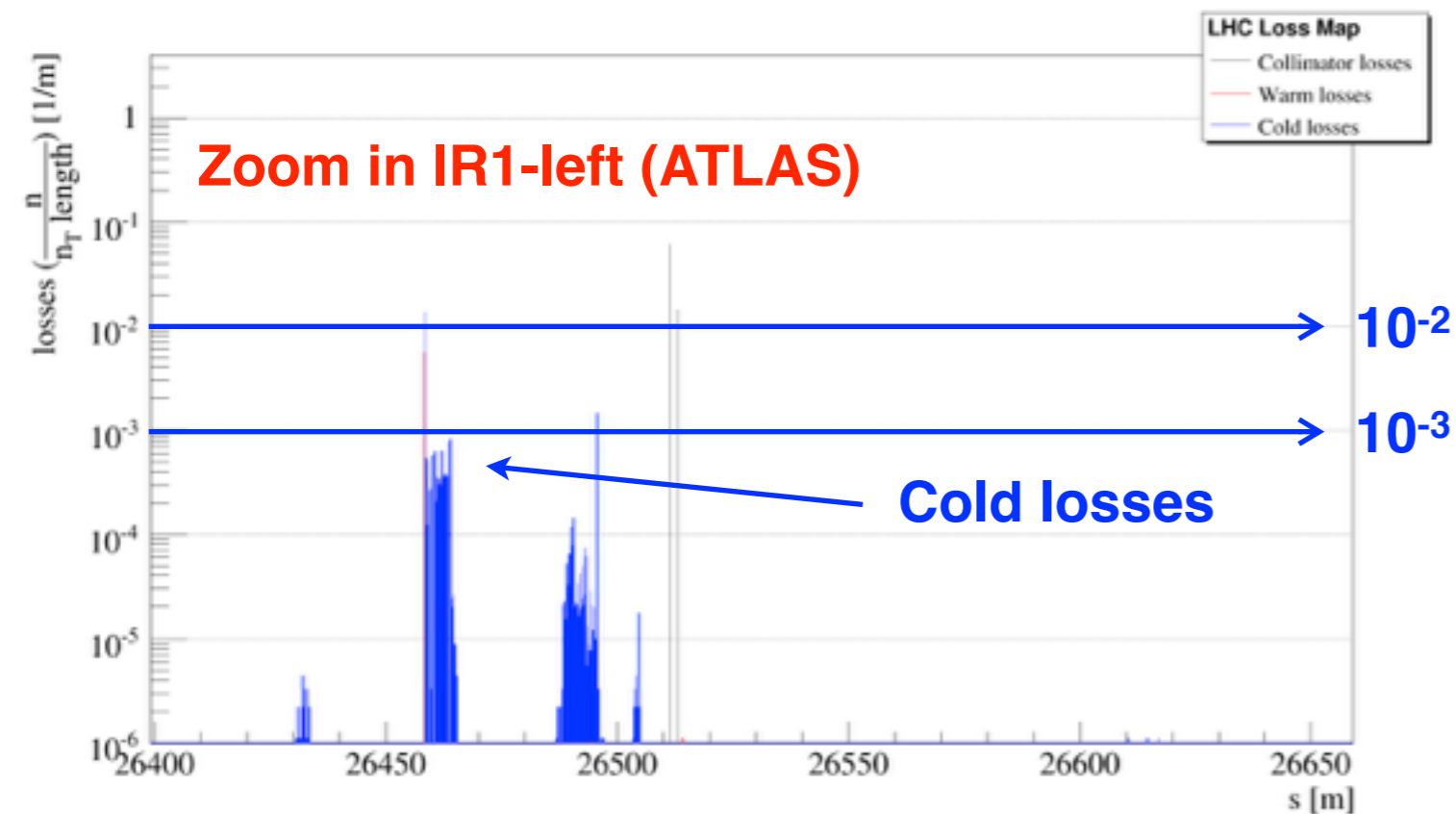
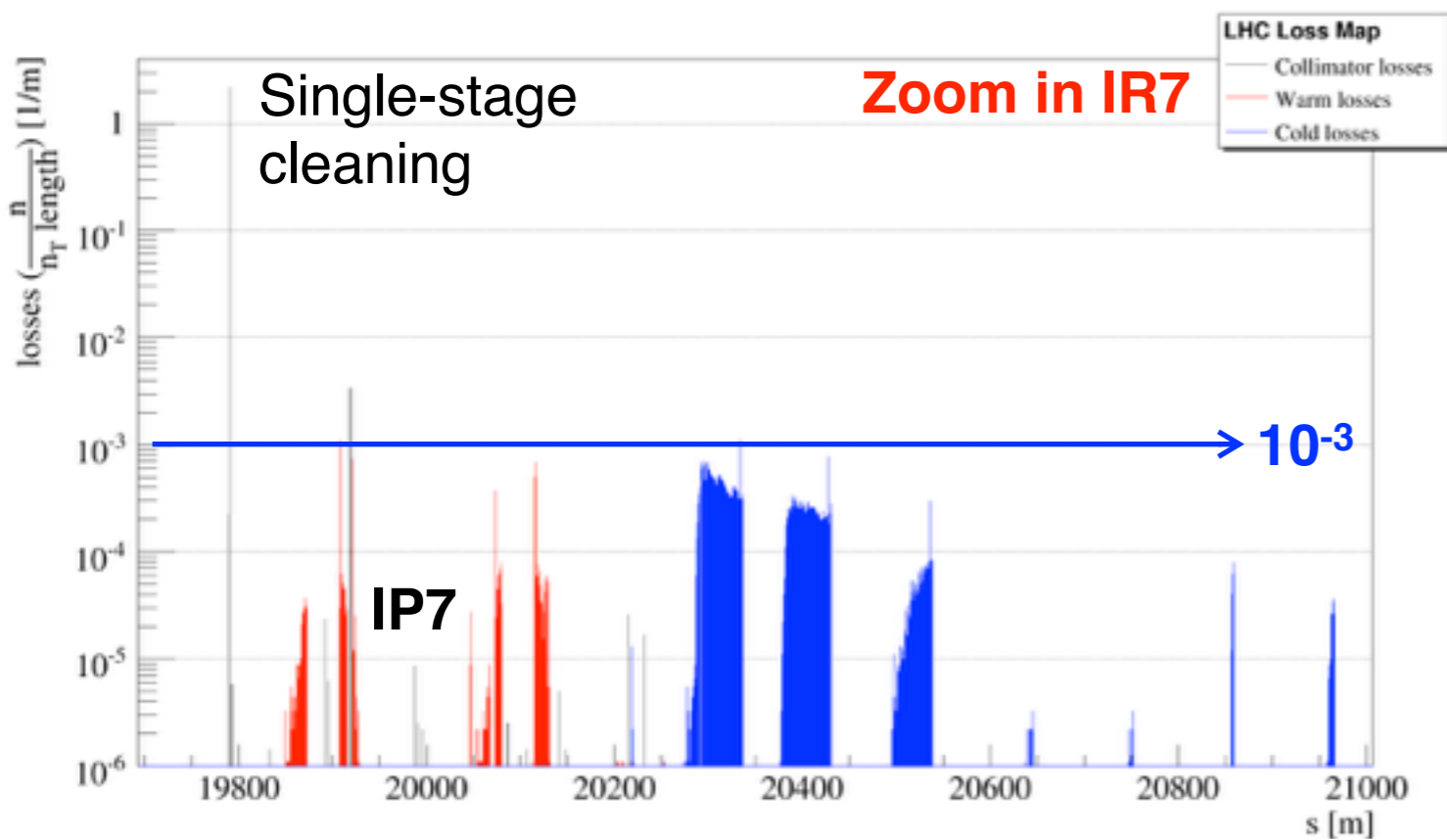
Fraction of proton lost per unit length.

Single-stage cleaning with one primary (H) collimator made 60 cm of Carbon: highest leakage in cold elements (blue spikes): **1-3 %**.

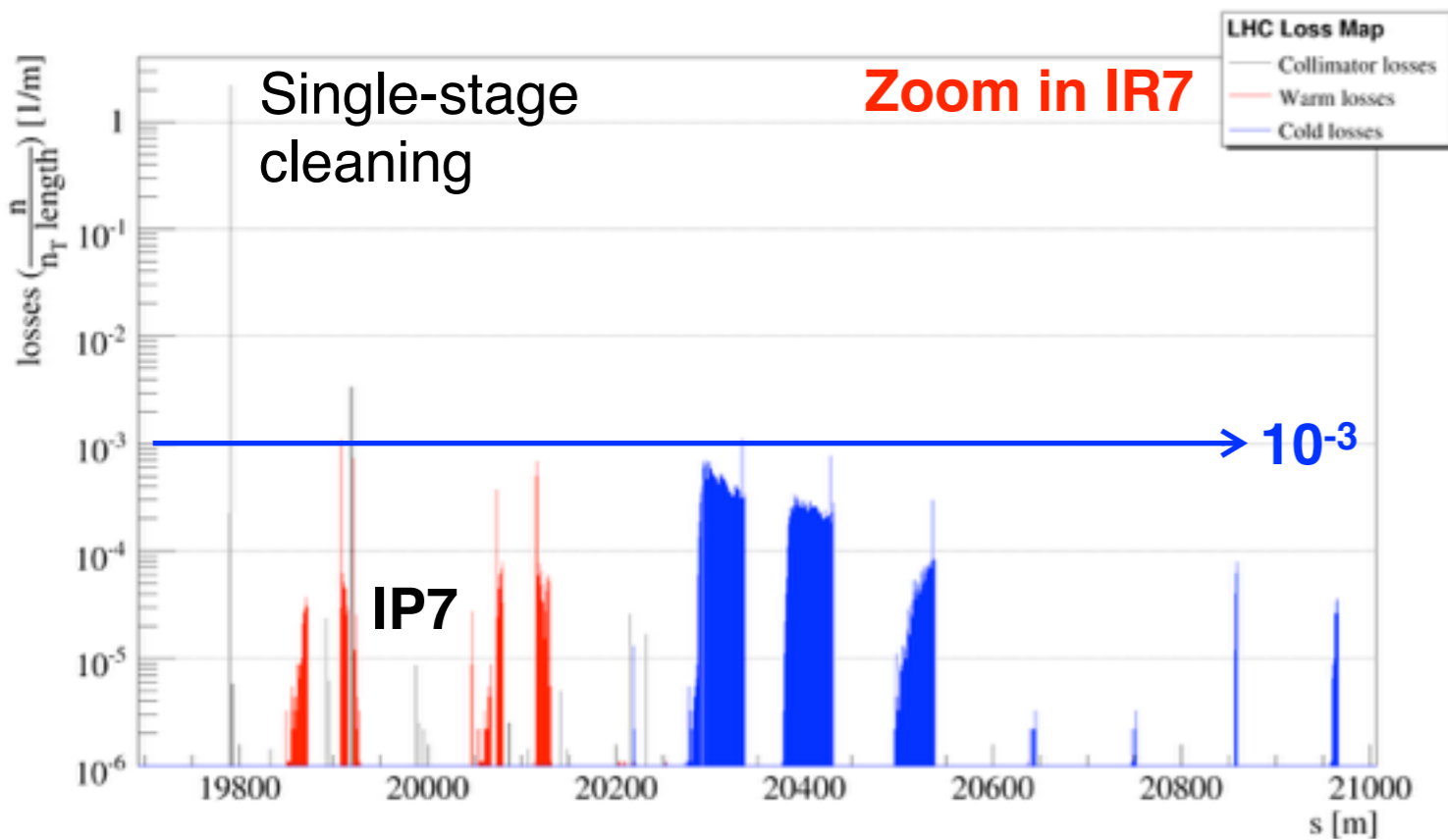
Comparison to quench limits



Comparison to quench limits



Comparison to quench limits

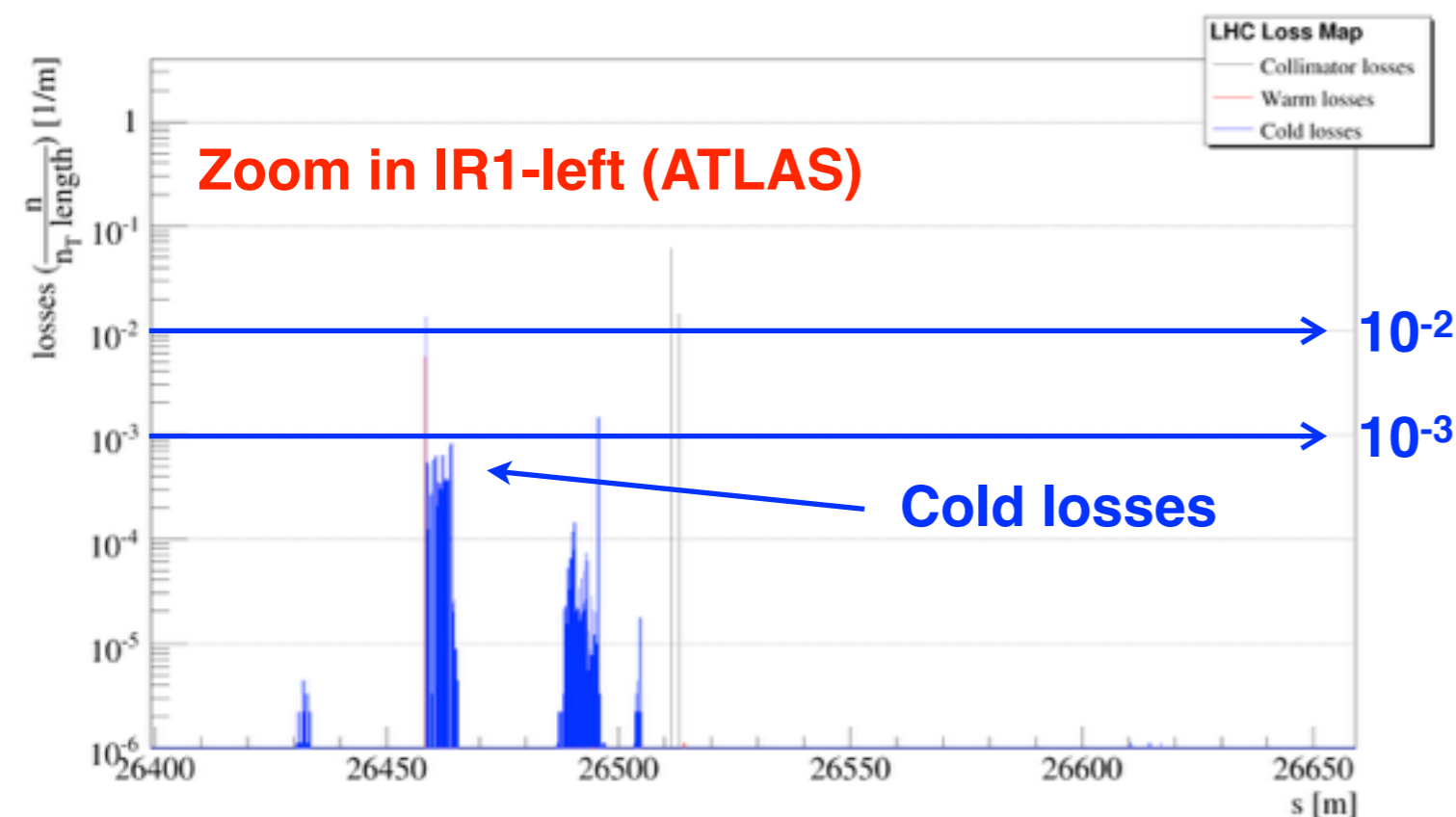


Typical assumed **quench limit** at 7 TeV for steady losses of \sim second timescales:

$$R_q (7 \text{ TeV}) = 3.2 \times 10^7 \text{ p/m/s}$$

With the single-stage cleaning predicted by this model, losses are up to:

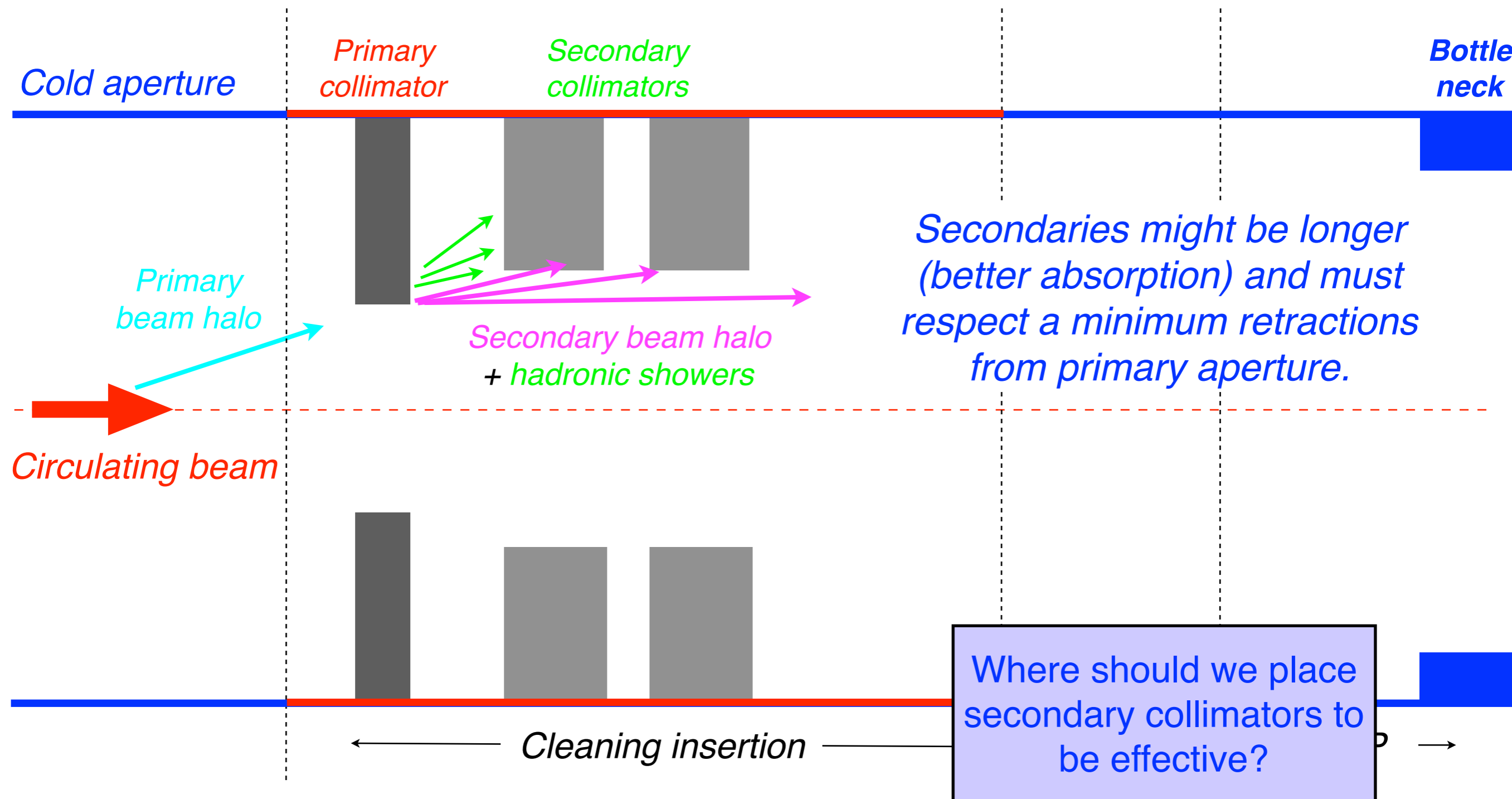
$$\begin{aligned} \tau_b = 1\text{h} &\rightarrow 90 \times 10^7 \text{ p/m/s} \text{ (30 x } R_q) \\ \tau_b = 0.2\text{h} &\rightarrow 450 \times 10^7 \text{ p/m/s} \text{ (150 x } R_q) \end{aligned}$$



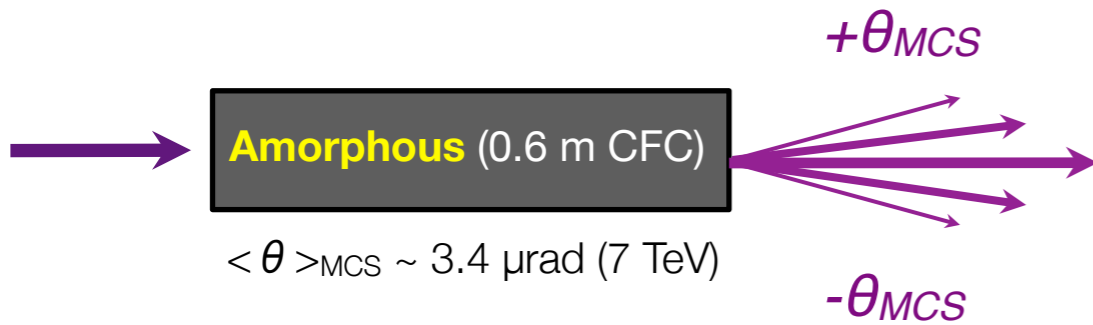
Single-stage cleaning is apparently not adequate for the LHC needs!

*Note: These are **approximated figures!** Detailed performance reach is estimated with more complex simulations including effects of showers!*

Two-stage collimation



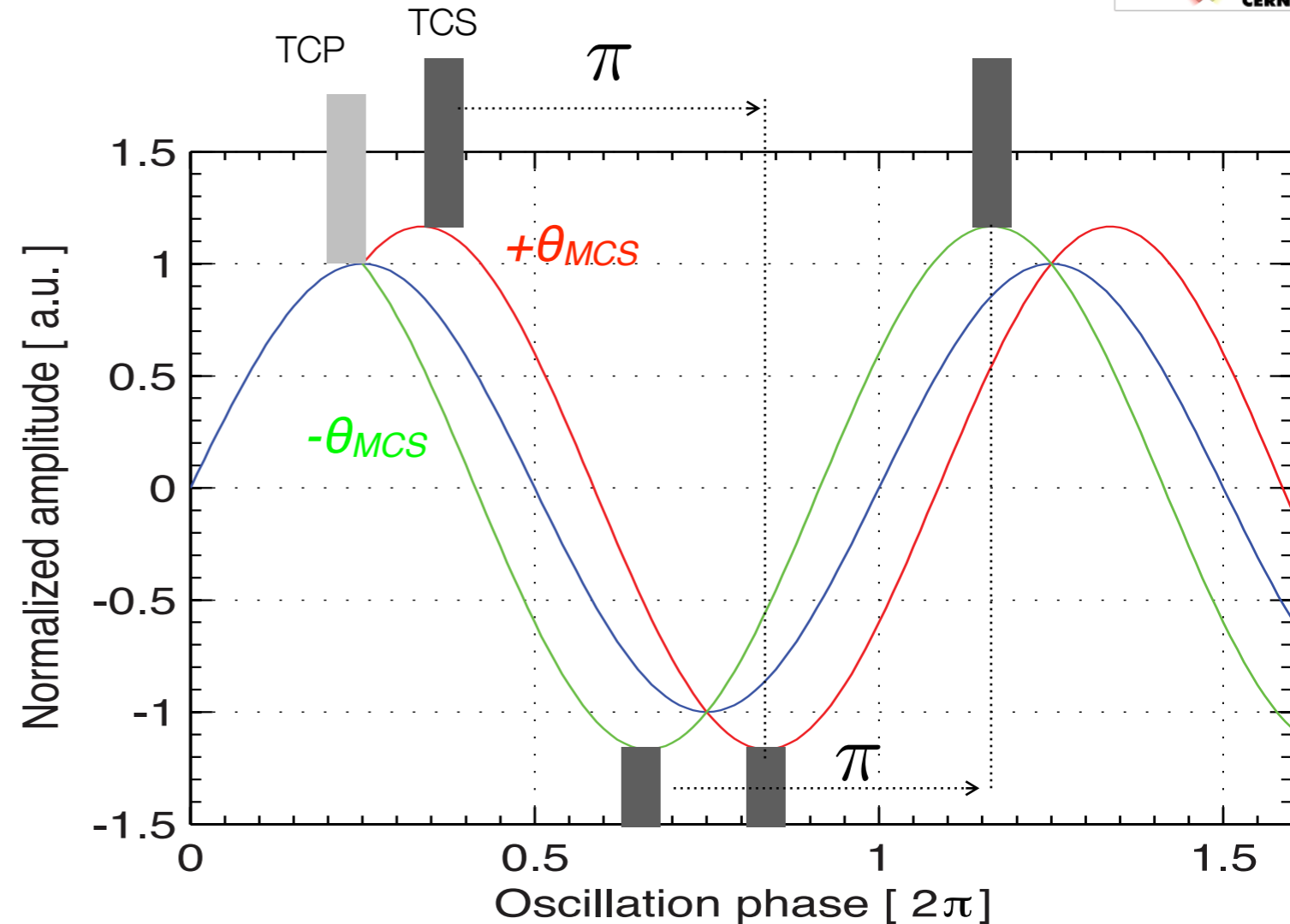
“Secondary” collimators (TCSs) can be added to intercept the secondary halo and the showers that leak out of the primary collimator.



There are two optimum phase locations to catch the debris from the primary collimators (TCPs).

Minimum: set of 2 secondary collimators (TCSs) covering $+\theta_{MCS}$ and $-\theta_{MCS}$.

Optimum: 4 TCSs (per plane) providing redundant coverage.



Betatron motion in $z \equiv (x, y)$

$$z_i(s) = \sqrt{\beta(s)\epsilon_i} \sin(\phi(s) + \phi_0)$$

$\beta(s)$: betatron function versus s

Secondary collimators must be placed at **optimum phase** locations where kicks from the TCP scattering translates into the largest offset.

Reality is a bit more complicated...

Optimum phases depend on TCP/TCS retraction

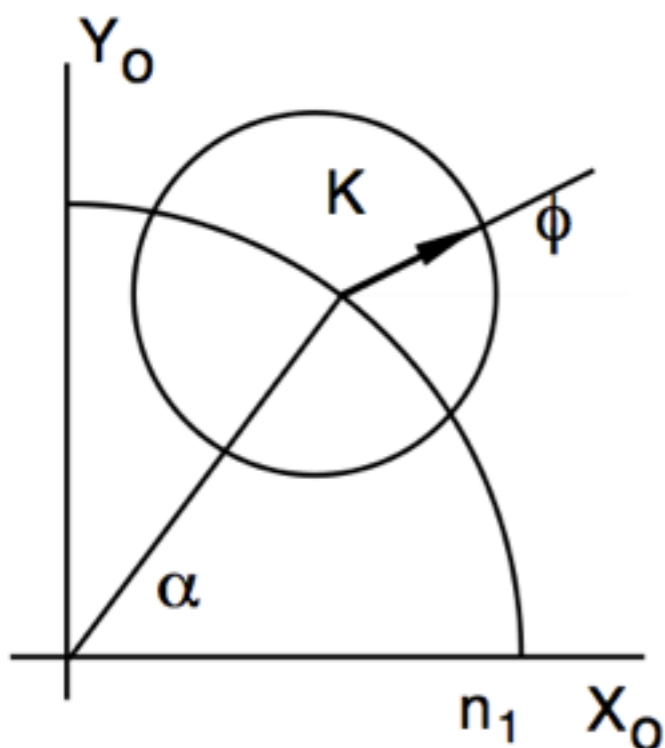
$$\tan \mu_x = \frac{\sqrt{n_{\text{TCP}}^2 - n_{\text{TCS}}^2}}{n_{\text{TCP}}^2} \frac{\cos \phi}{\cos \alpha}$$

$n_{\text{TCP}}, n_{\text{TCS}}$: TCP and TCS half-gap

α, ϕ : collimator plane and scattering angle

$$\cos \mu_0 = n_{\text{TCP}} / n_{\text{TCS}}$$

Phys.Rev.ST Accel.Beams 1:081001,1998



Optics of a two-stage collimation system

J. B. Jeanneret

CERN, CH-1211 Geneva, Switzerland

(Received 13 October 1998; published 21 December 1998)

Phase locations (μ_x, μ_y) and jaw orientation (α_J) to catch different scattering angle (ϕ) for horizontal ($\alpha=0$), vertical ($\alpha=\pi/2$) and skew ($\alpha=\pi/2$) scattering source locations.

α	ϕ	μ_x	μ_y	α_J
0	0	μ_0	—	0
0	π	$\pi - \mu_0$	—	0
0	$\pi/2$	π	$3\pi/2$	μ_0
0	$-\pi/2$	π	$3\pi/2$	$-\mu_0$
$\pi/4$	$\pi/4$	μ_0	μ_0	$\pi/4$
$\pi/4$	$5\pi/4$	$\pi - \mu_0$	$\pi - \mu_0$	$\pi/4$
$\pi/4$	$3\pi/4$	$\pi - \mu_0$	$\pi + \mu_0$	$\pi/4$
$\pi/4$	$-\pi/4$	$\pi + \mu_0$	$\pi - \mu_0$	$\pi/4$
$\pi/2$	$\pi/2$	—	μ_0	$\pi/2$
$\pi/2$	$-\pi/2$	—	$\pi - \mu_0$	$\pi/2$
$\pi/2$	π	$\pi/2$	π	$\pi/2 - \mu_0$
$\pi/2$	0	$\pi/2$	π	$\pi/2 + \mu_0$

Reality is a bit more complicated...

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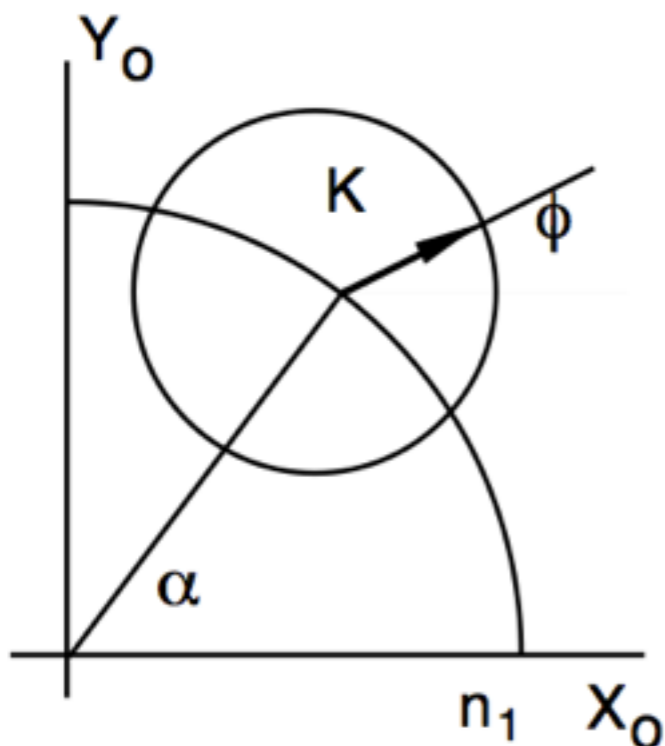
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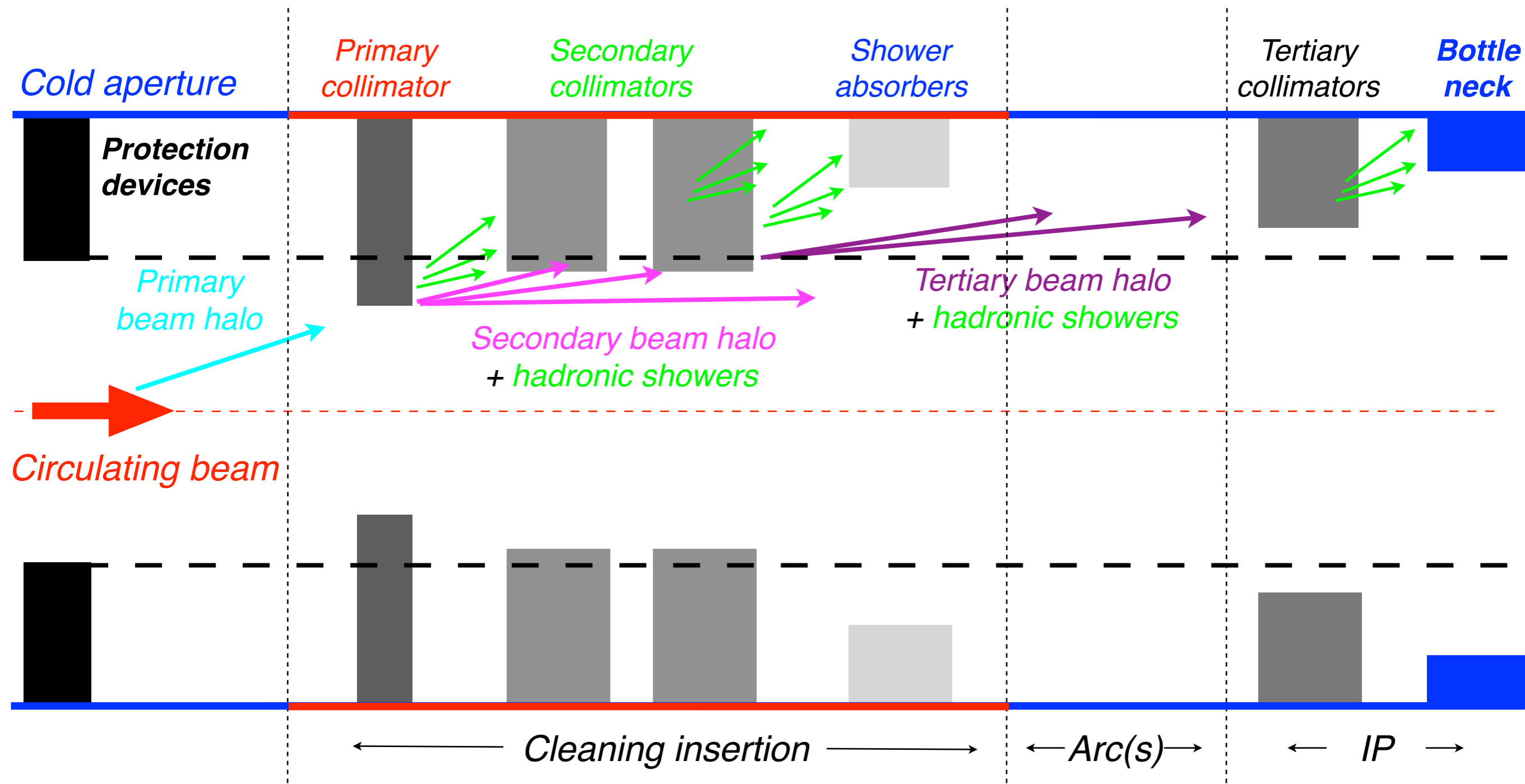
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0	0	μ_0	—	0
0	π	$\pi - \mu_0$	—	0
0	$\pi/2$	π	$3\pi/2$	μ_0
0	$-\pi/2$	π	$3\pi/2$	$-\mu_0$
$\pi/4$	$\pi/4$	μ_0	μ_0	$\pi/4$
$\pi/4$	$5\pi/4$	$\pi - \mu_0$	$\pi - \mu_0$	$\pi/4$
$\pi/4$				
$\pi/4$				
$\pi/2$				
$\pi/2$				
$\pi/2$				
$\pi/2$	0	$\pi/2$	π	$\pi/2 + \mu_0$

A finite number of secondary collimators can be used to catch efficiently the halo with three primary collimator orientation.

Multi-stage collimation at the LHC

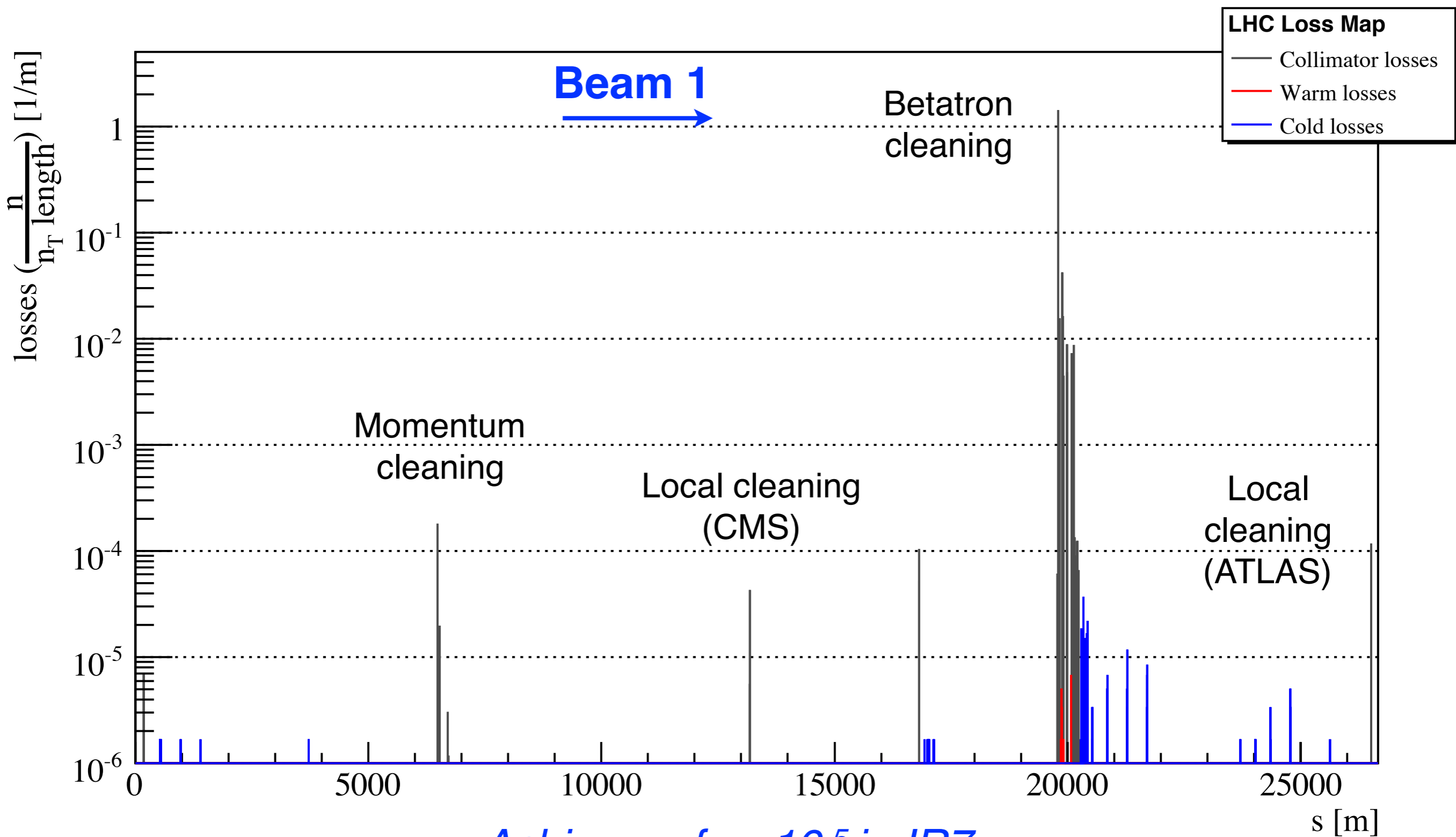


Including protection devices, a **5-stage cleaning** is required!

The system performance relies on achieving the well-defined **hierarchy** between different **collimator families** and **machine aperture**.



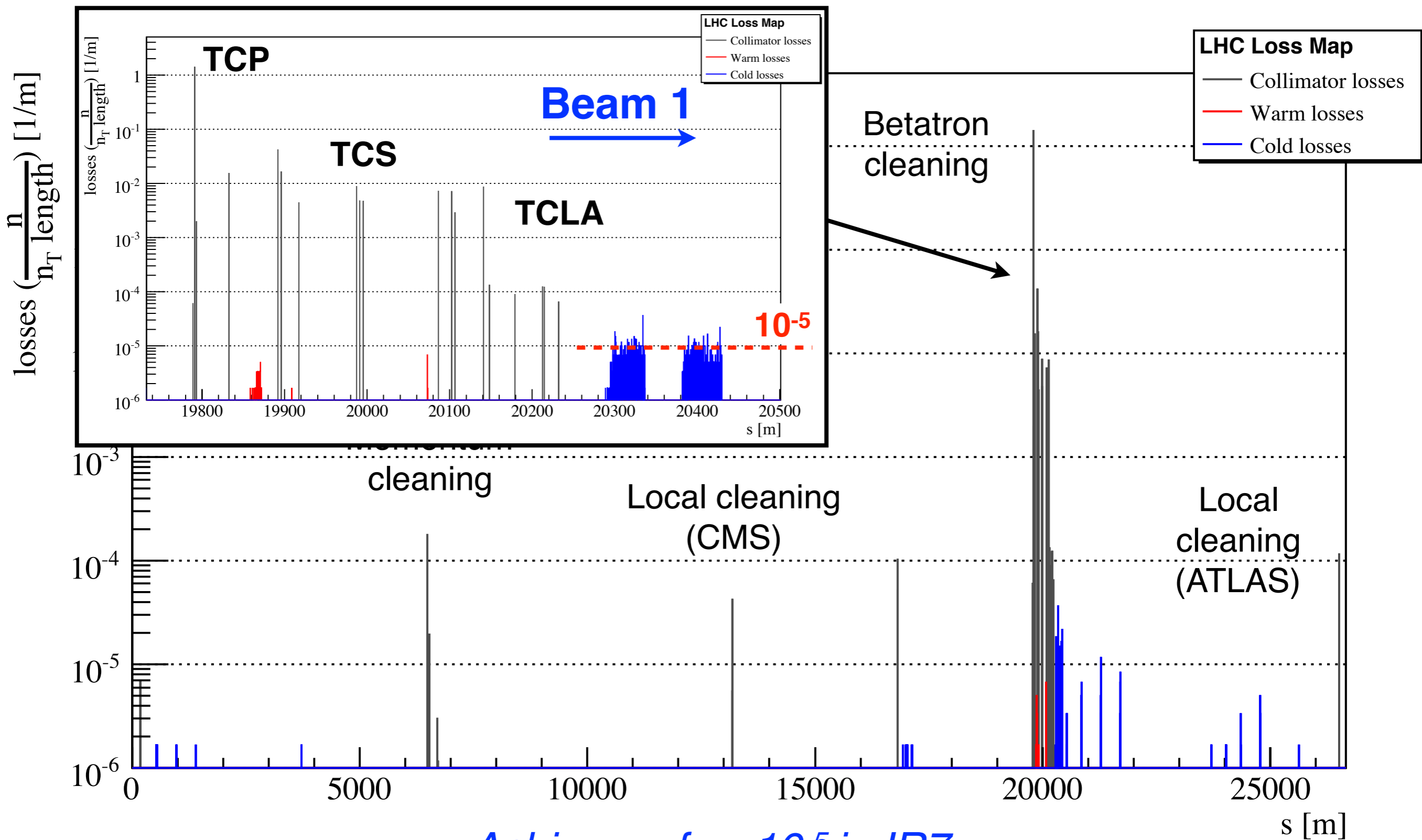
Simulated 7 TeV performance



Achieve a few 10^{-5} in IR7.

Cold losses in experiments removed by local protection.

Simulated 7 TeV performance

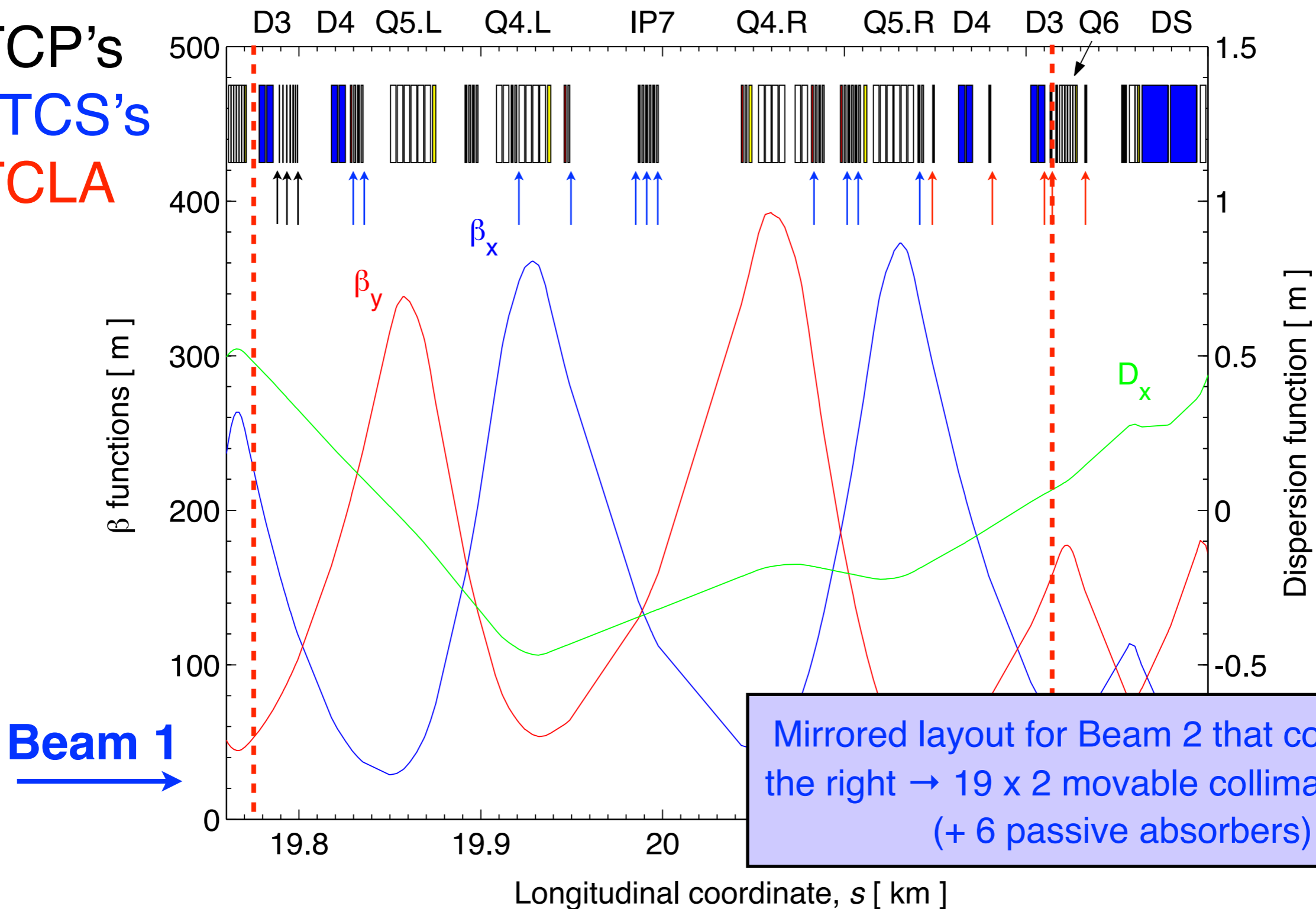


Achieve a few 10^{-5} in IR7.

Cold losses in experiments removed by local protection.

Betatron cleaning insertion

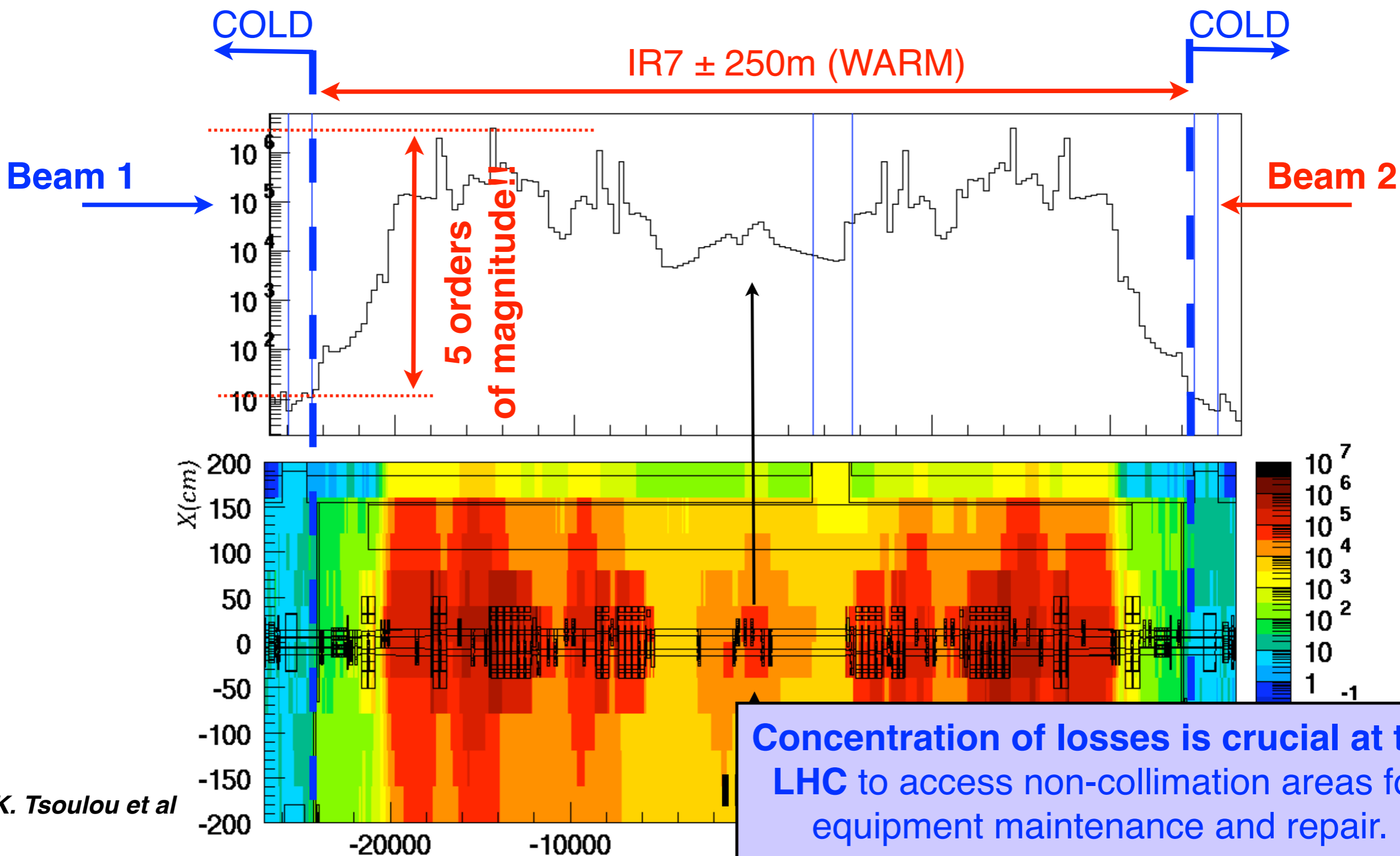
3 TCP's
11 TCS's
5 TCLAs



$$z_i(s) = \sqrt{\beta(s)\epsilon_i} \sin(\phi(s) + \phi_0) + \left(\frac{\delta p}{p}\right)_i D_z(s)$$

One full oscillation of the betatron motion to meet in the warm part the optimum phase conditions.

Radiation doses in collimation region



Activation from halo losses is basically confined within the warm insertions!

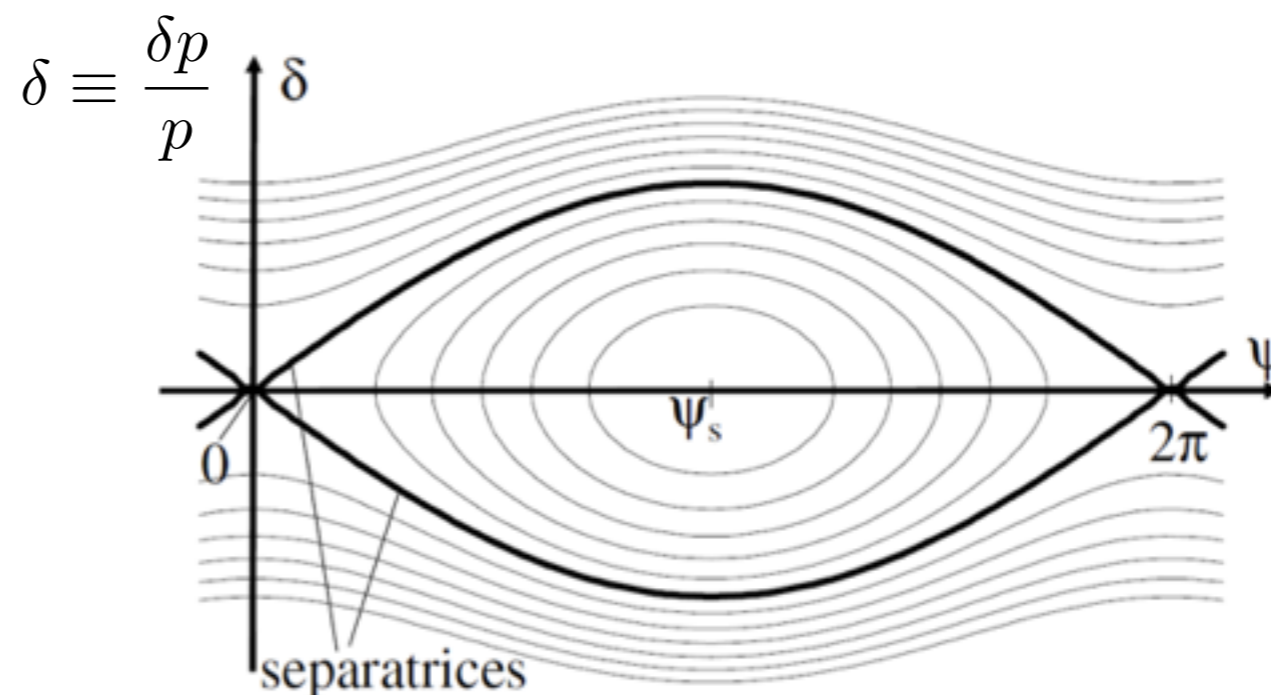
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“**Off-momentum losses**” = losses occurring when beam particles lose the energy matching compared to the reference particle.

$$z_i(s) = \sqrt{\beta(s)\epsilon_i} \sin(\phi(s) + \phi_0) + \left(\frac{\delta p}{p} \right)_i D_z(s)$$

Examples: trips or setting errors of RF system, capture losses at the start of ramp, synchrotron radiation losses of particle outside RF buckets, collision with other beams or with collimator materials.

How do we collimate these particles?



Catching off-momentum particles

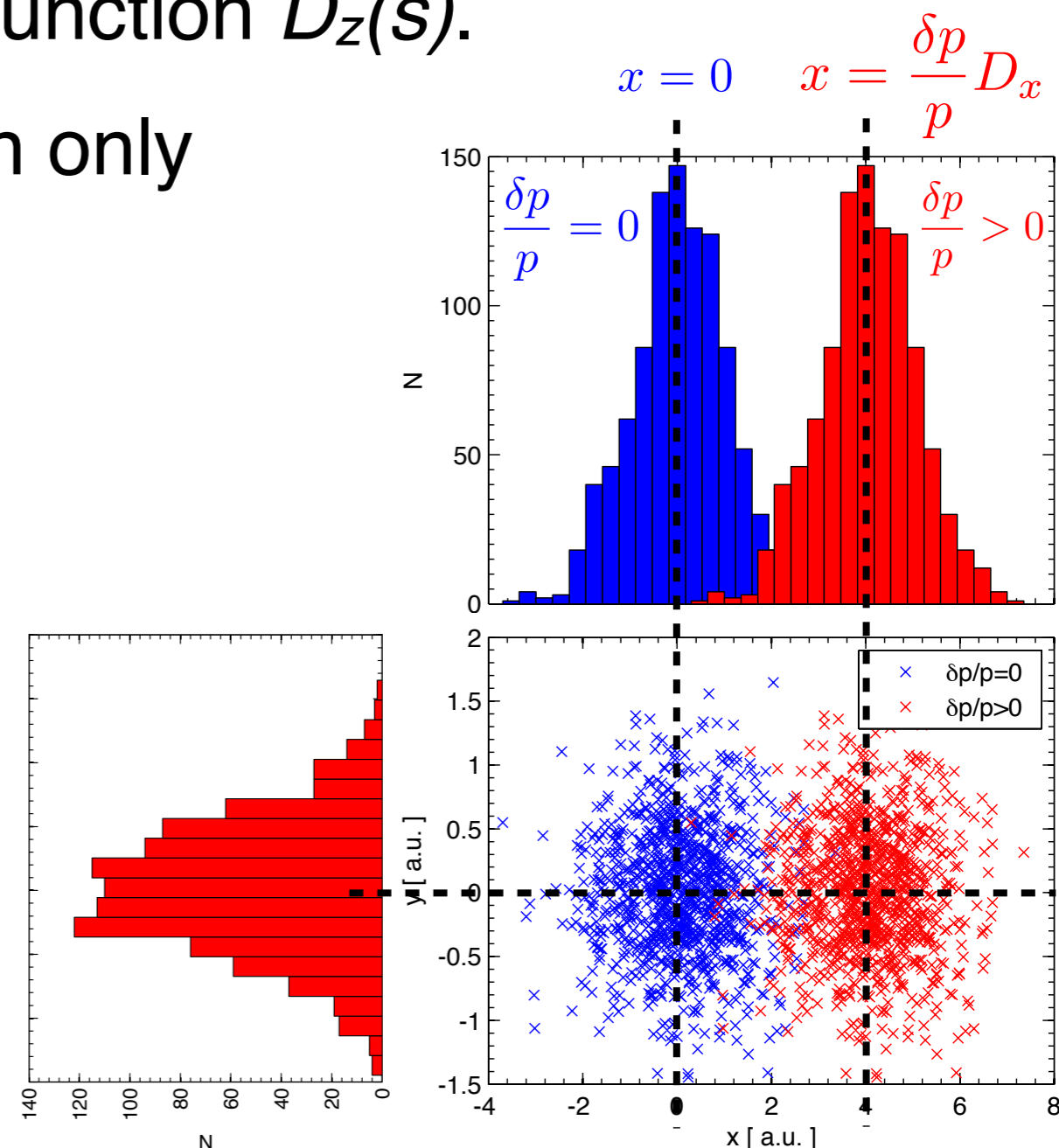
For all off-momentum loss cases, individual halo particles or the entire beam **maintain** their initial **betatron amplitude**.

The **mismatch in energy** translates into a **shift of position** that follows the periodic dispersion function $D_z(s)$.

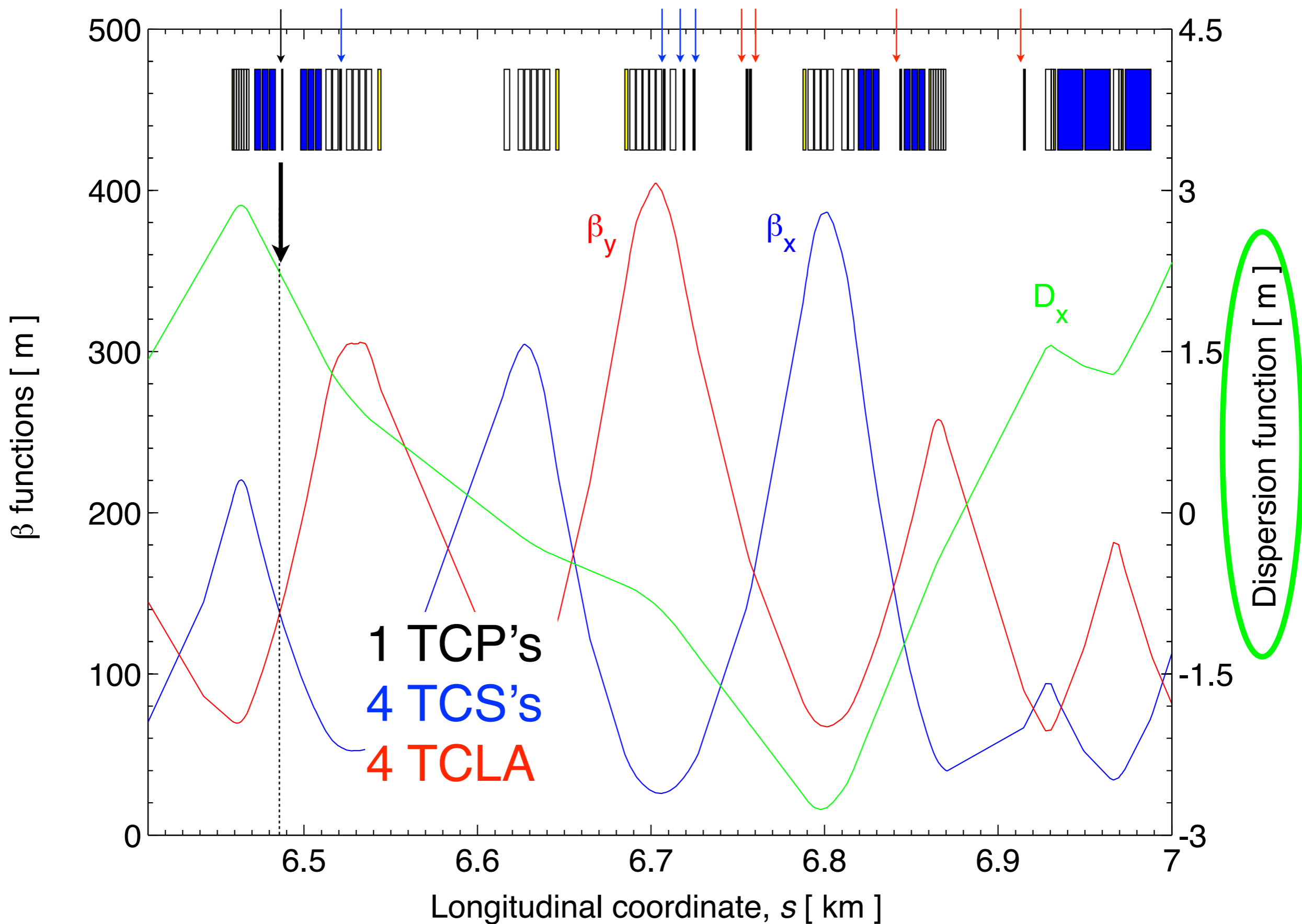
Circular accelerators have by design only horizontal dispersion

⇒ **only H momentum collimation!**

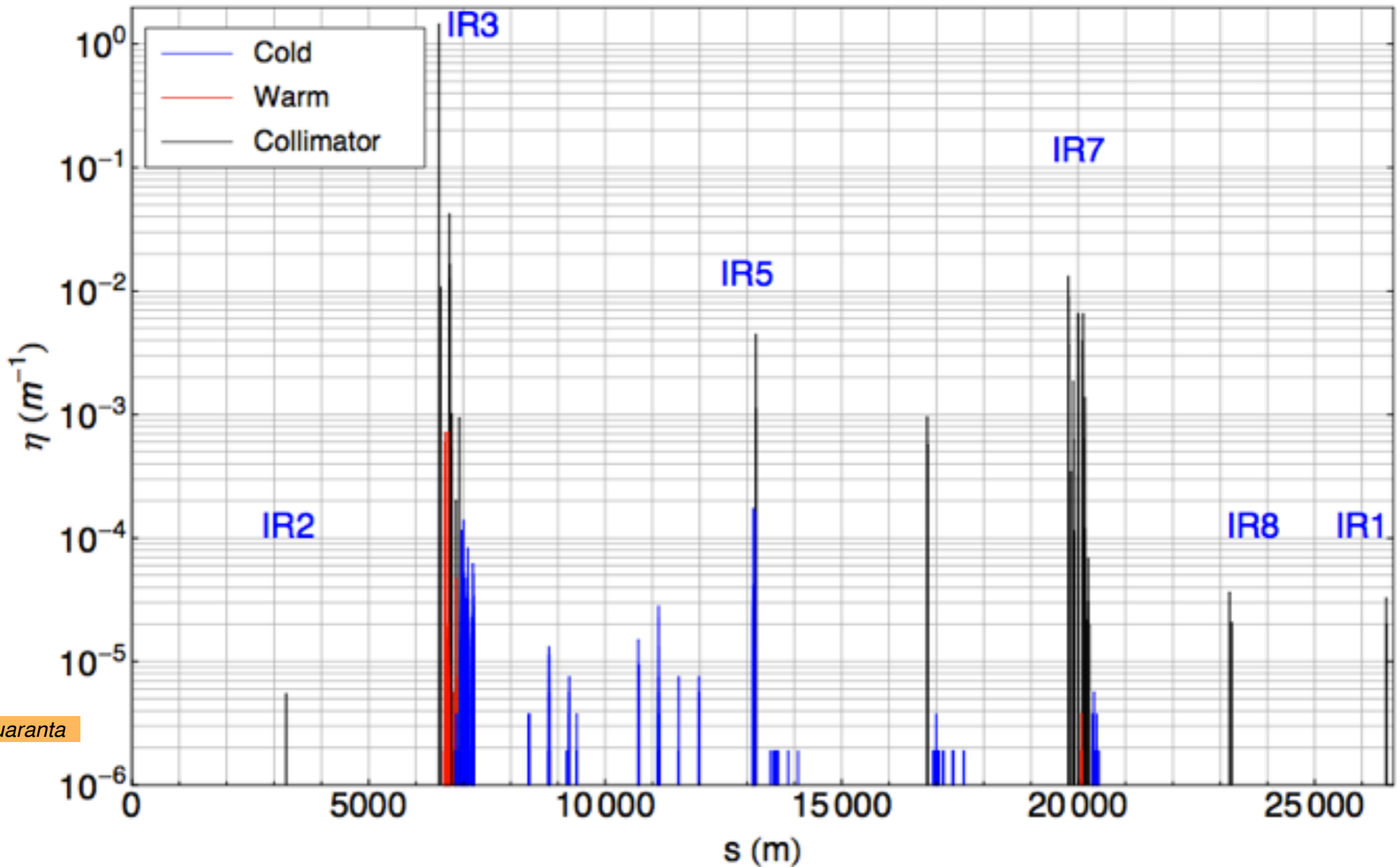
***Special optics conditions** in the momentum cleaning insertions ensure that the primary collimators are the “off-momentum bottleneck”. Otherwise, a **similar multi-stage approach** is used for cleaning.*



Momentum cleaning optics



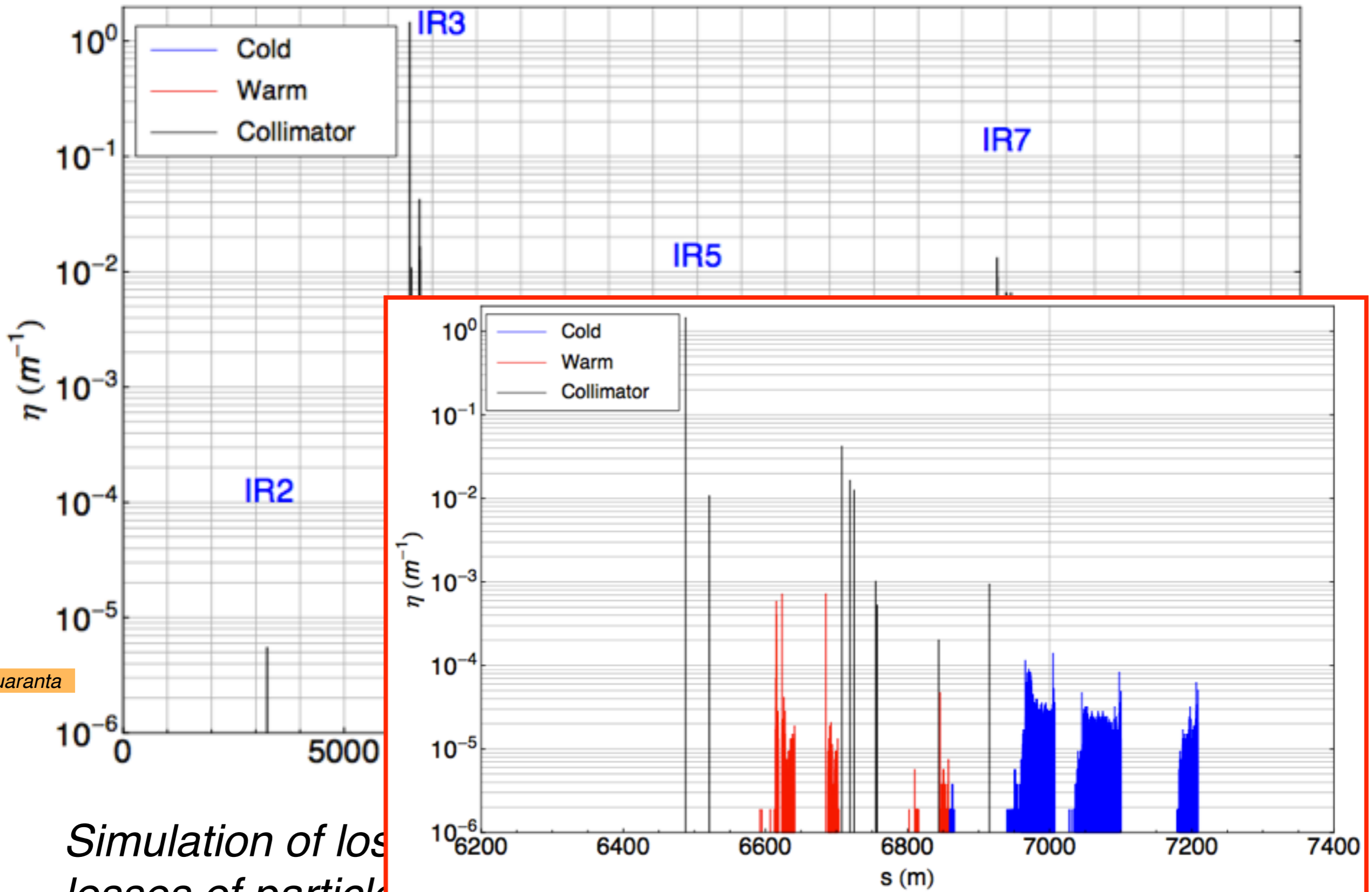
IR3 loss maps: synch. radiation losses



E. Quaranta

Simulation of losses in IR3 caused by synchrotron radiation losses of particles outside the RF buckets at the 7 TeV LHC.

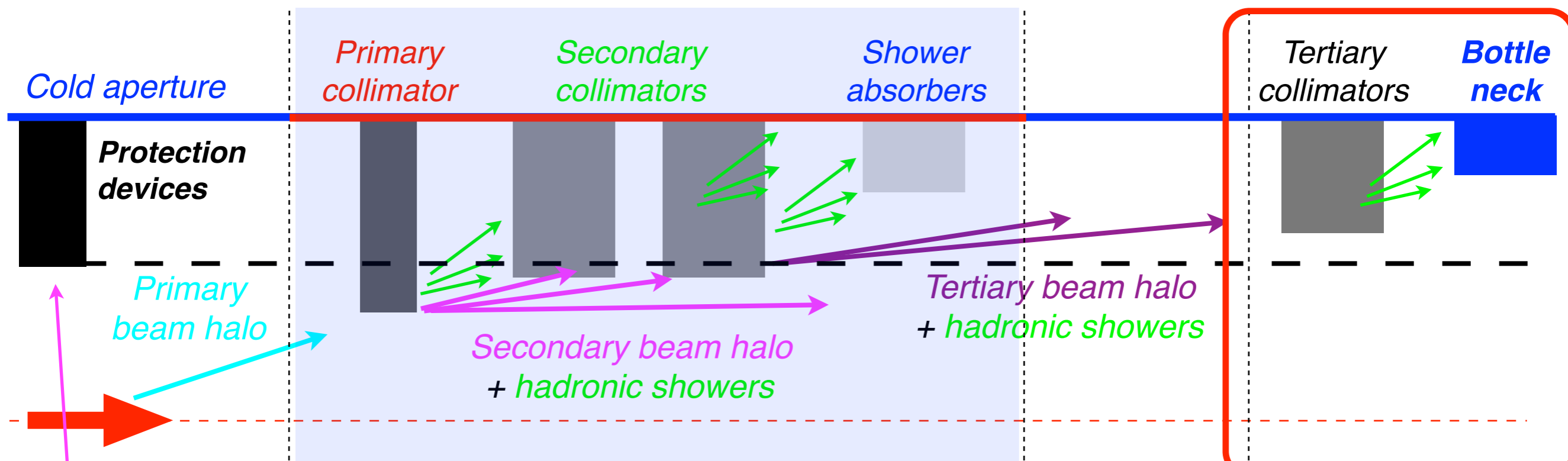
IR3 loss maps: synch. radiation losses



E. Quaranta

Simulation of losses of particles outside the 111 buckets at the 7 TeV LHC.

Local cleaning and protection



Protection devices covered in another lecture.

Note: all modern colliders had concerns with losses in the “low- β^* insertions”.

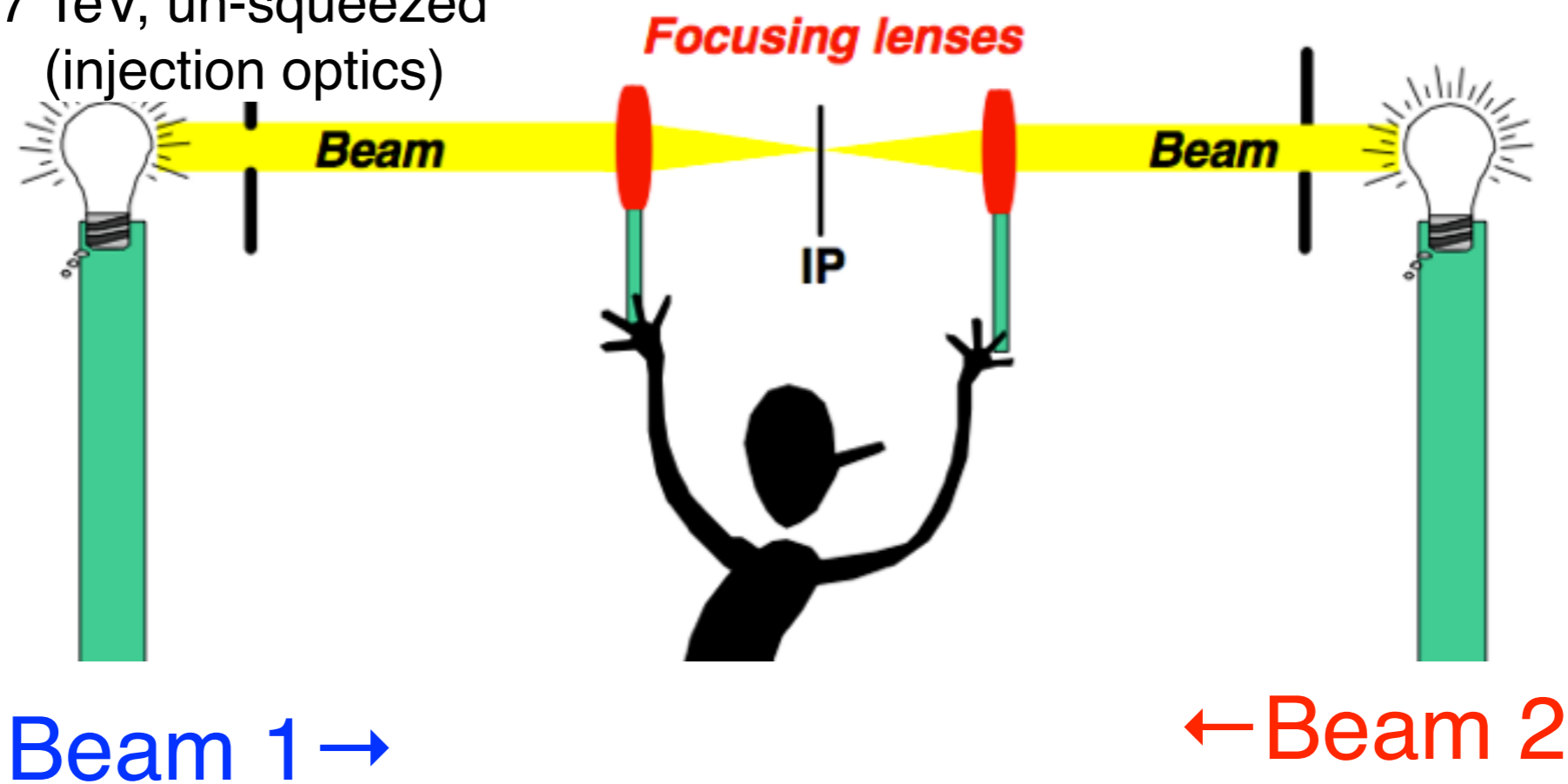
When do we need local protection?

How is the collimator position chosen in these cases?

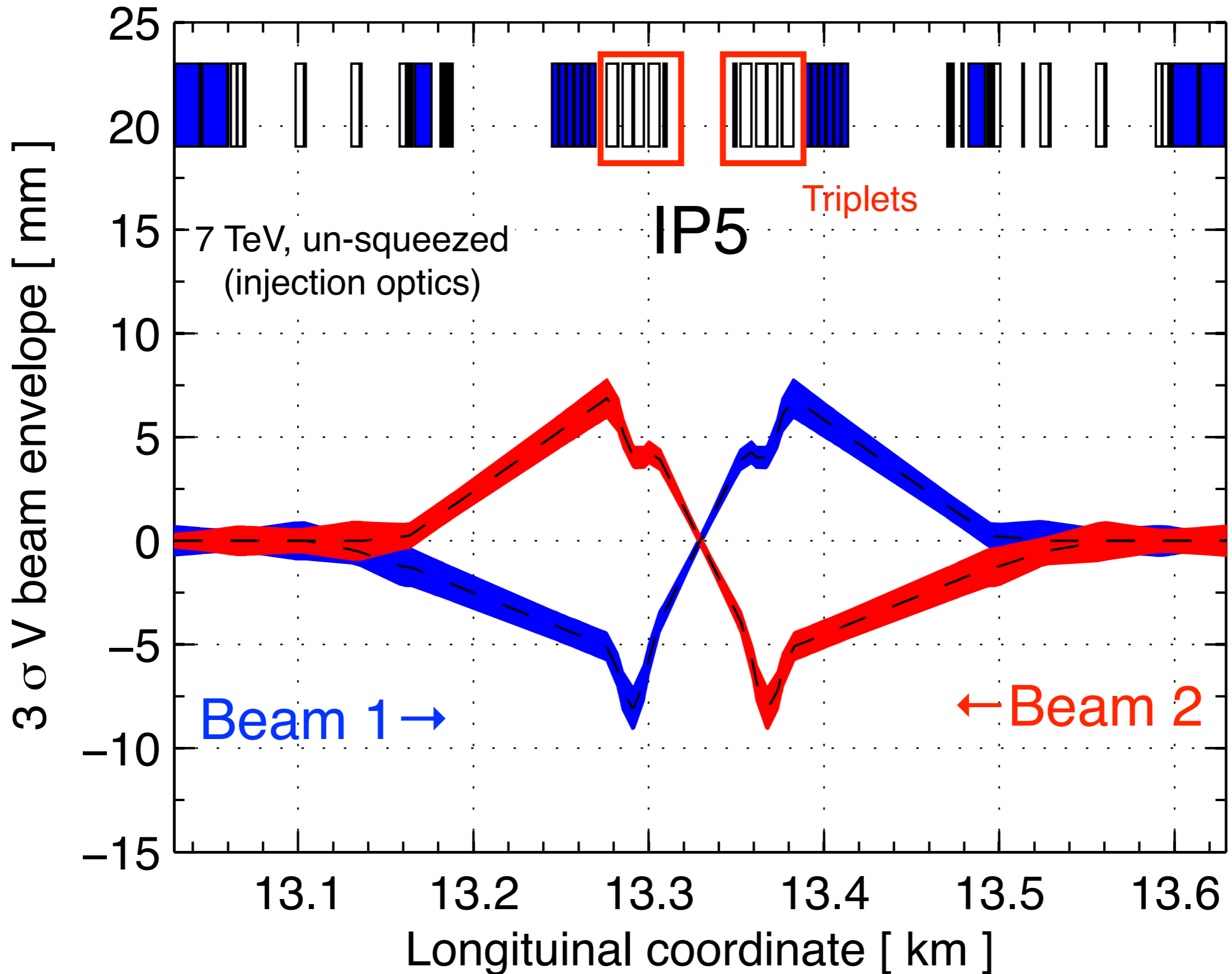
→ Briefly look at the **tertiary collimators** that protect the **inner triplet** in all experimental regions.

Optics in high-luminosity points

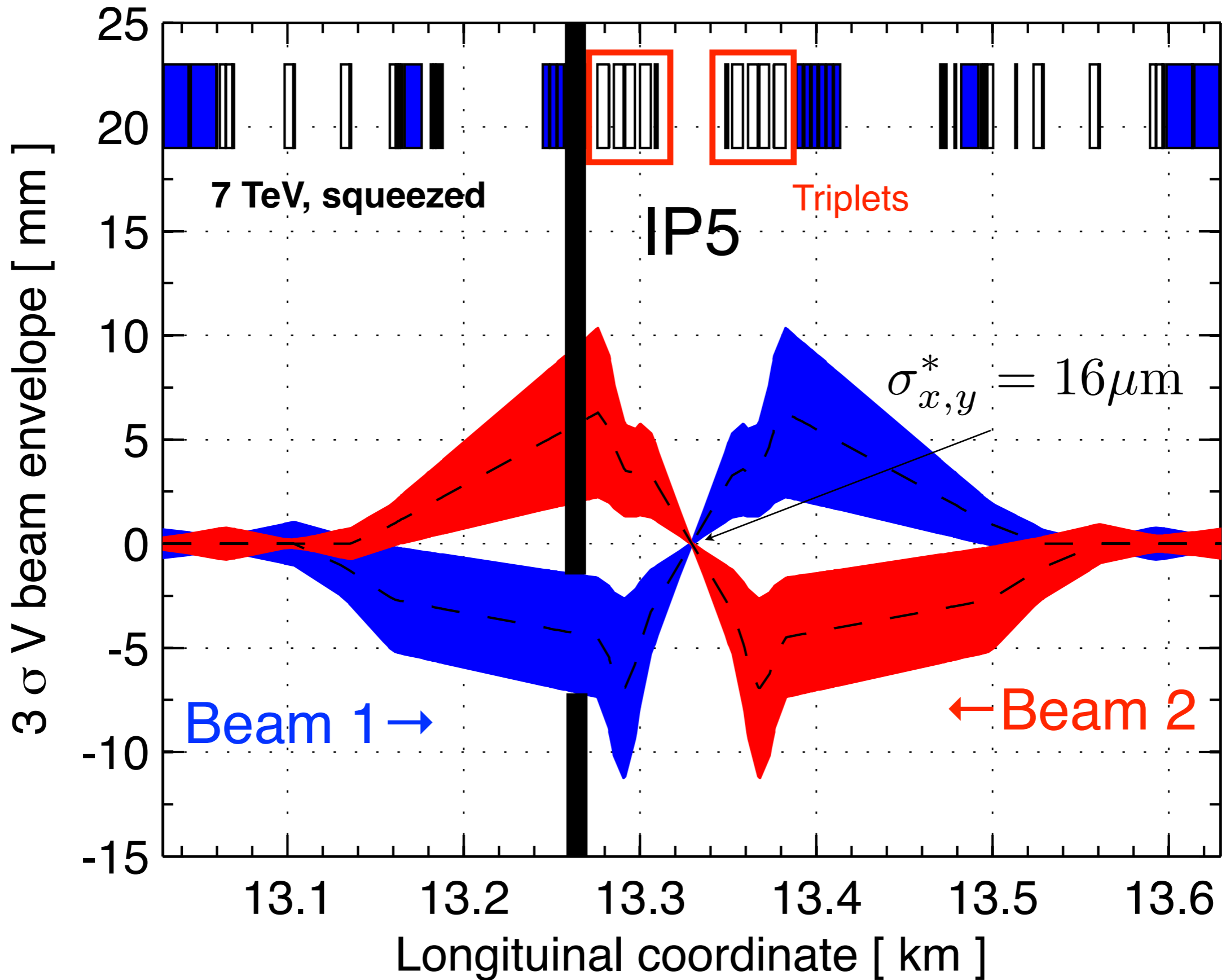
7 TeV, un-squeezed
(injection optics)



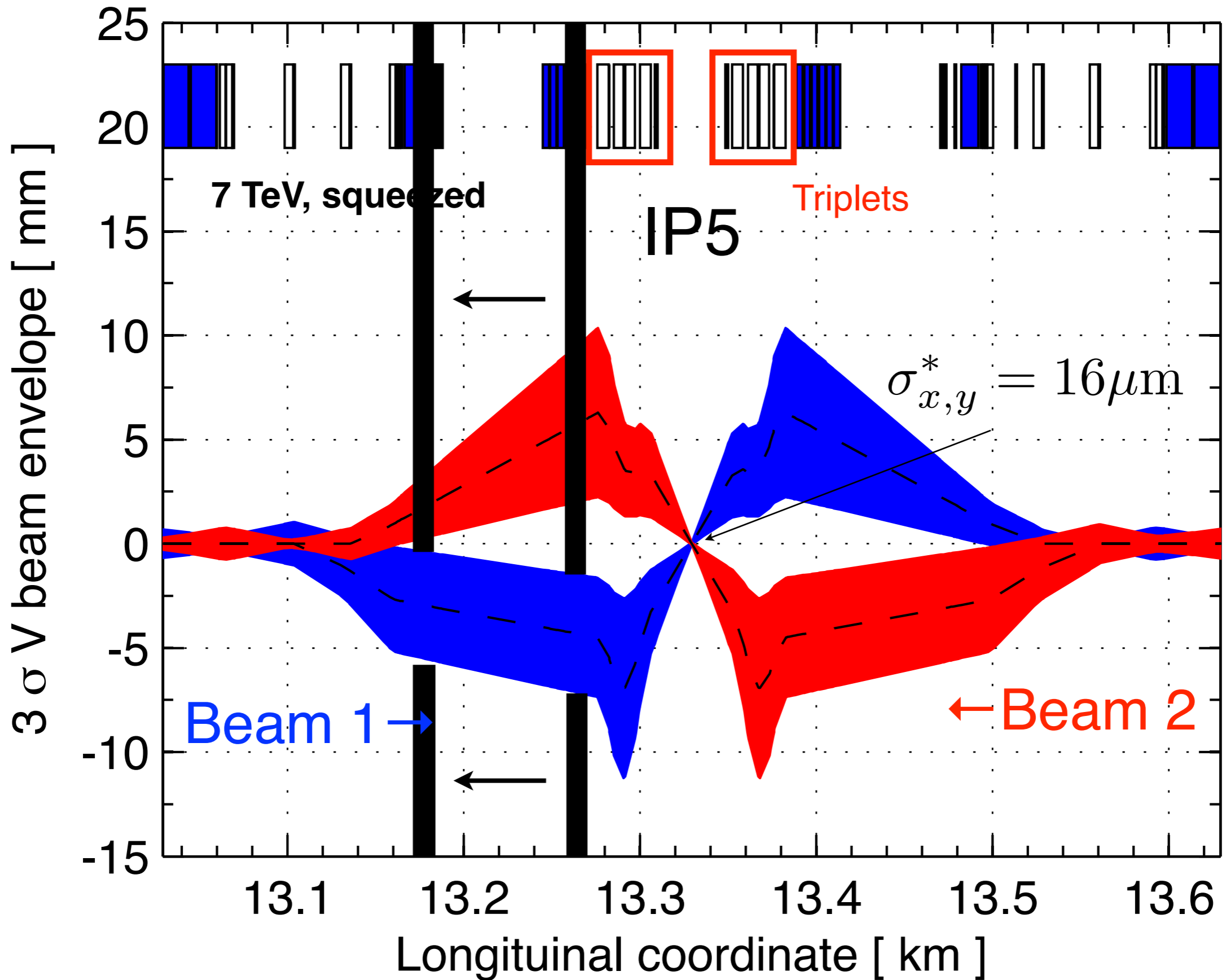
Optics in high-luminosity points



Optics in high-luminosity points



Optics in high-luminosity points



Tertiary collimators (TCT's) are part of the betatron collimation hierarchy and are used to protect the inner triplets of the low- β^* experiments

Clean the tertiary halo that leaks out of the cleaning insertions.

Protect the magnets in case of abnormal losses.

Tertiary collimators might be used to tune experiment backgrounds.

Triplet protection with “squeezed” beams is maximized by

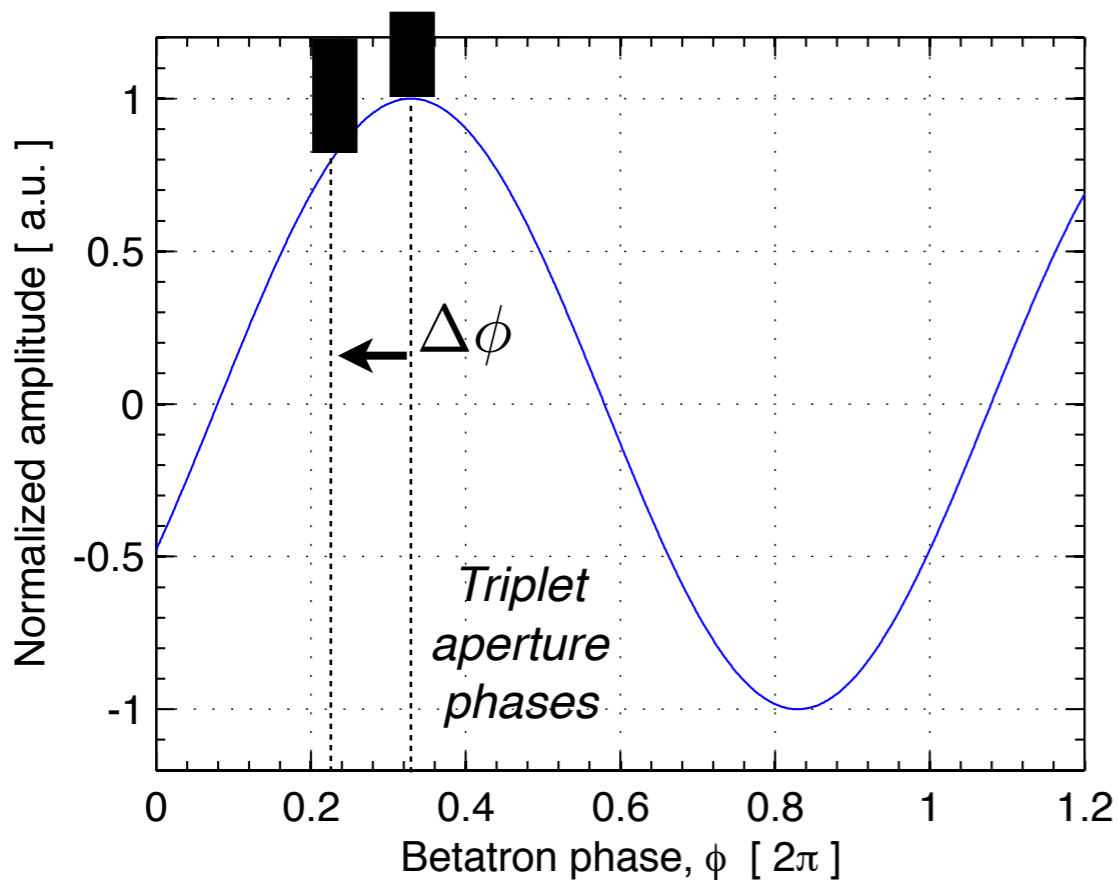
Minimizing the “betatron phase difference” to the TCT

Use high-Z material to maximize absorption → in case of catastrophic failures, better destroy the collimator than a magnet!

TCT's are located typically in **cold regions** → settings must guarantee that they are not exposed to large beam loads.

What if we cannot place TCT's at same phase of the triplet?

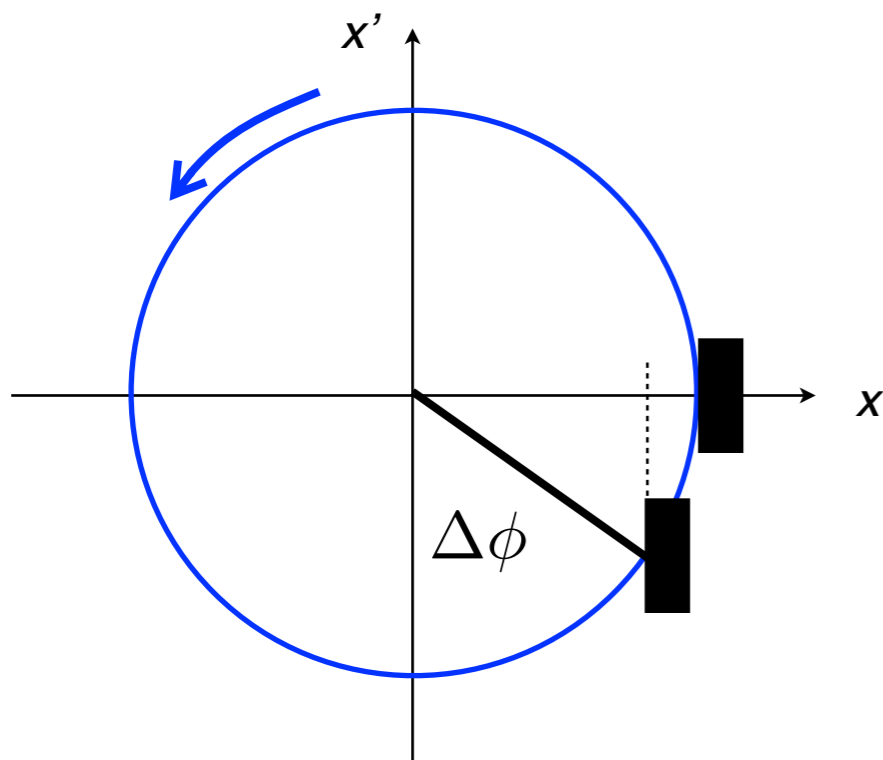
TCT settings versus aperture



*If one cannot install the TCT at the same phase at the aperture bottleneck, equivalent protection levels can only be achieved **closing the collimator to smaller gaps.***

Exercise: calculate the required TCT settings changes versus the phase difference.

*Who is more familiar with the beam dynamics, can also see the solution in the **normalized phase-space diagram.***



Change is small: with squeezed optics, $\Delta\phi \cong 0$ at the TCT location available!

- Main points from 1st lecture
- Machine protection and collimation
 - Concepts and LHC implementation
 - Case study: 2008 event
- Beam losses and collimation
 - Roles of beam collimation systems
 - Beam losses mechanisms
- Design of a multi-stage collimation system
 - Betatron collimation design
 - Advanced: off-momentum, local protection
- **The LHC beam collimation system**
 - **Detailed layouts**



LHC collimation system layout



**Two warm cleaning insertions,
3 collimation planes**

IR3: Momentum cleaning

- 1 primary (H)
- 4 secondary (H)
- 4 shower abs. (H,V)

IR7: Betatron cleaning

- 3 primary (H,V,S)
- 11 secondary (H,V,S)
- 5 shower abs. (H,V)

Local cleaning at triplets

8 tertiary (2 per IP)

Passive absorbers for warm magnets

Physics debris absorbers

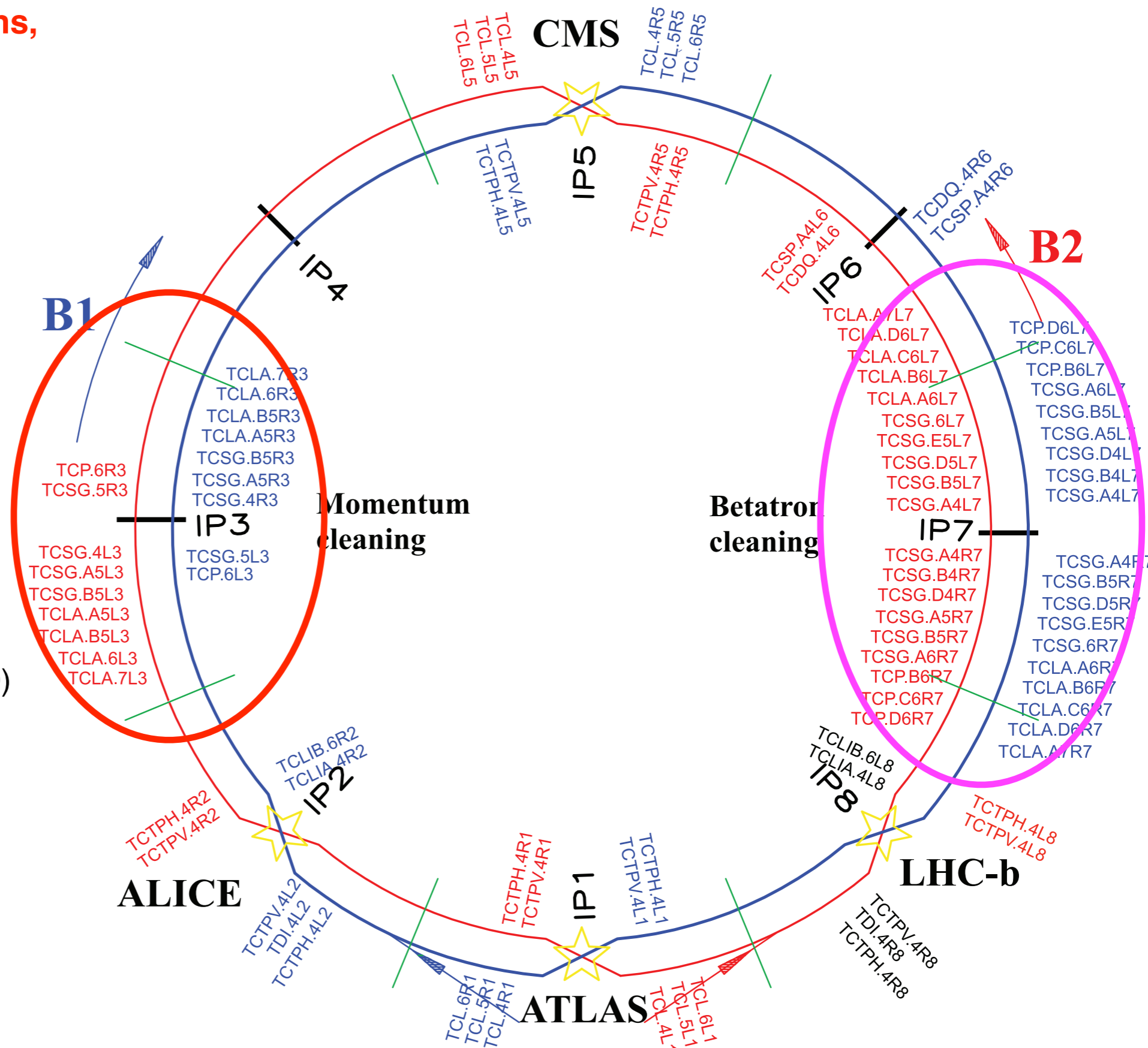
Transfer lines (13 collimators)

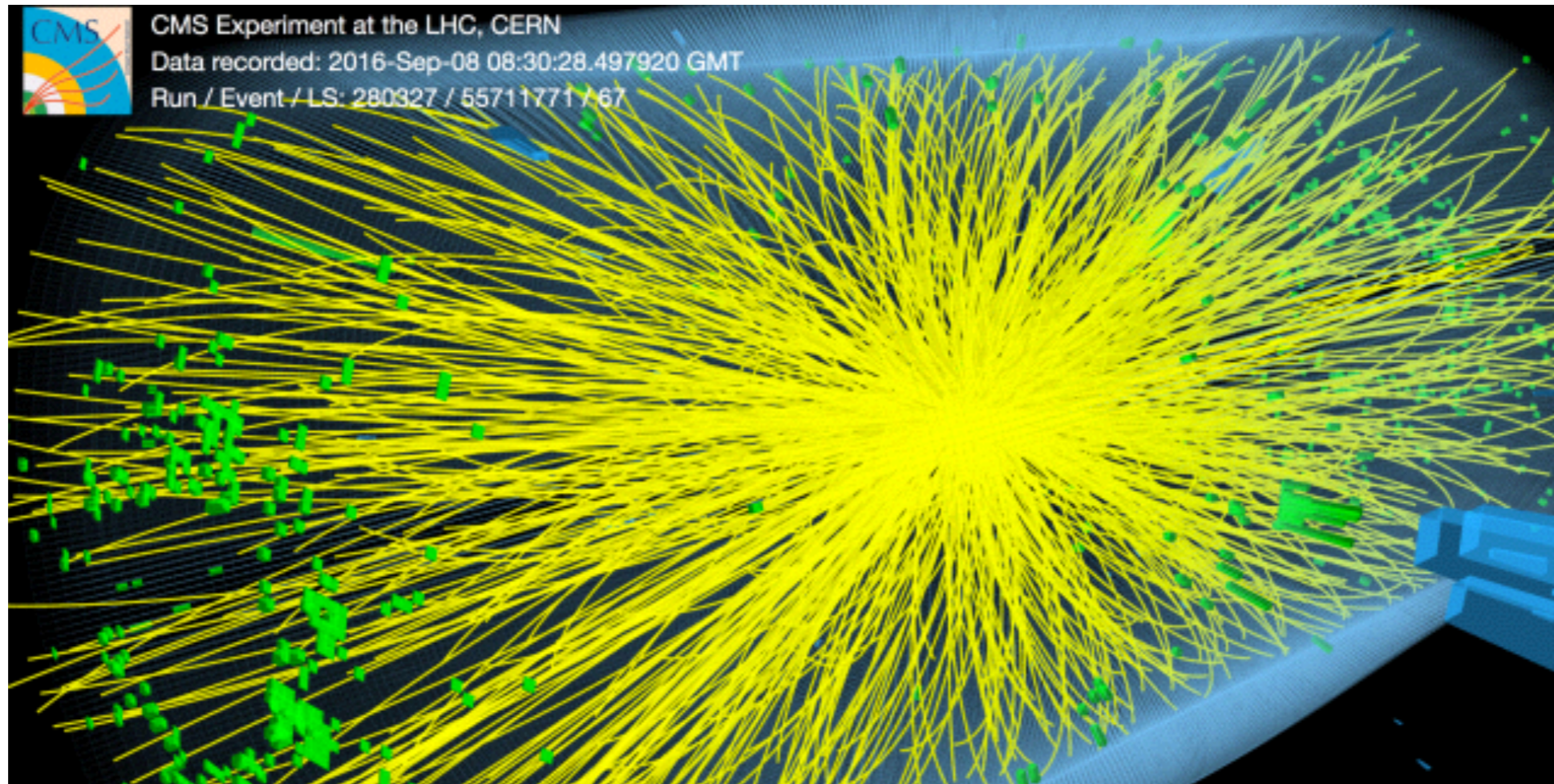
Injection and dump protection (10)

Total of 118 collimators

(108 movable).

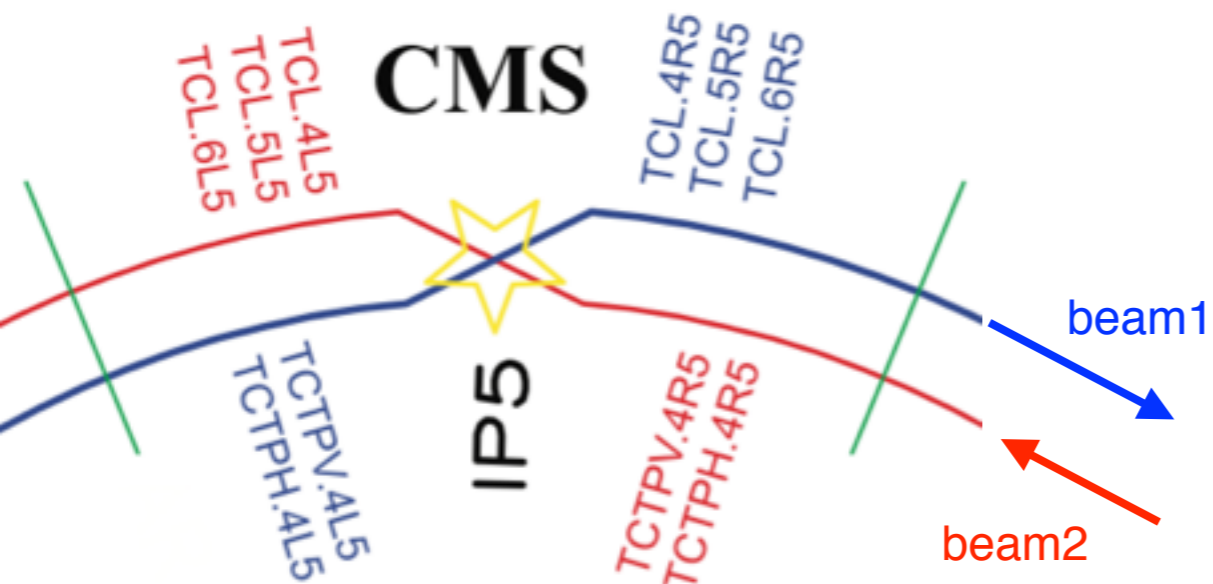
Two jaws (4 motors) per collimator!



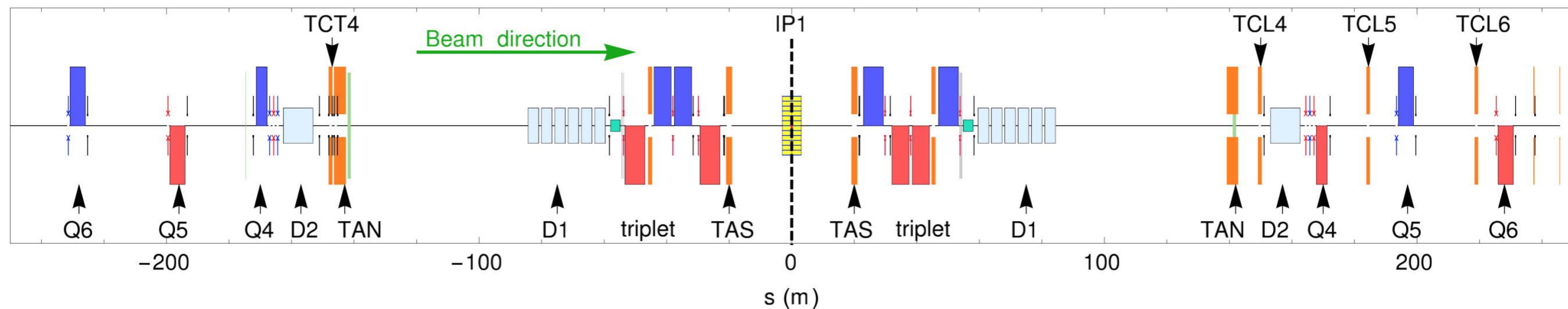


High-energy protons emerge from the collision points with perturbed trajectories — transverse kicks and energy deviations caused by elastic and inelastic collision with the opposite beam.

At the large luminosities, they risk to quench the cold magnets around ATLAS and CMS!

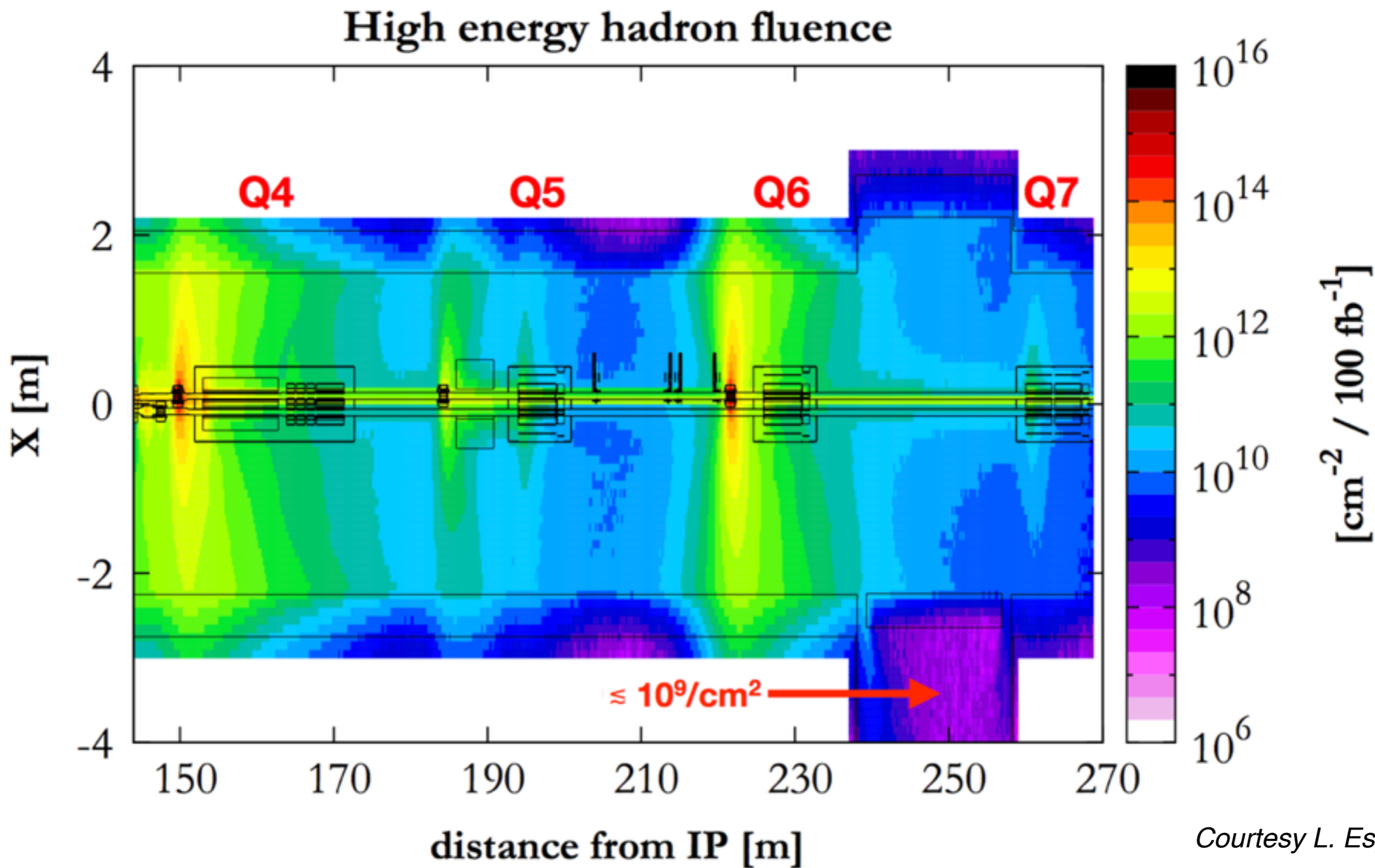


Protection of the “matching section” magnets and (partially) of the first magnets of the arcs around the experiments is done with dedicated collimators (3 per beam).



Courtesy R. Bruce

Physics debris collimation — iii



Courtesy L. Esposito