

Measuring quark polarizations at ATLAS and CMS

Yevgeny Kats



Based on:

JHEP 11 (2015) 067 [arXiv:1505.02771]

with Galanti, Giammanco, Grossman, Stamou, Zupan

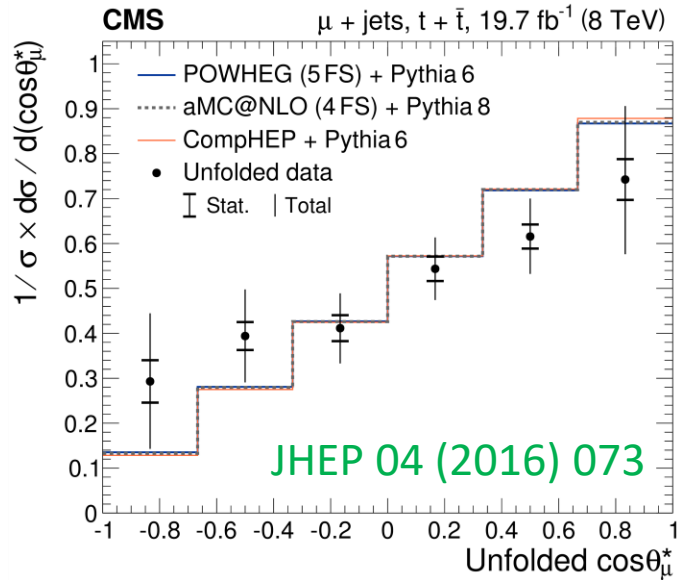
Phys. Rev. D 92, 071503 (2015) [arXiv:1505.06731]

JHEP 11 (2016) 011 [arXiv:1512.00438]

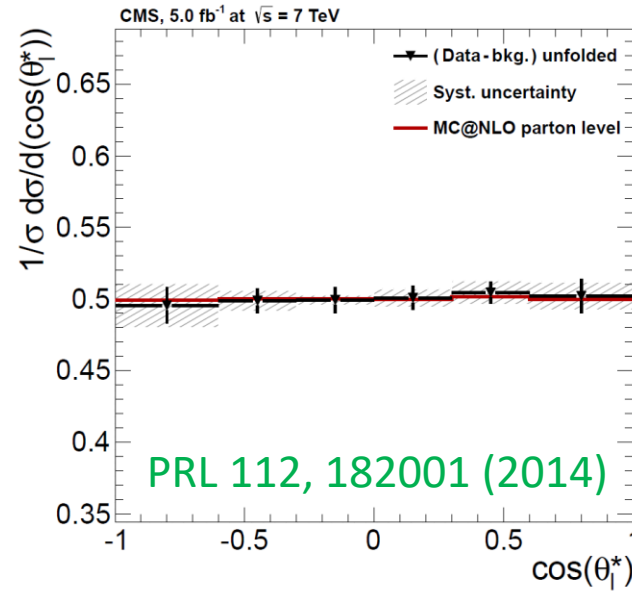
Motivation

ATLAS and CMS already measure **top quark polarization**.

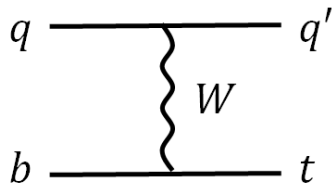
single top production



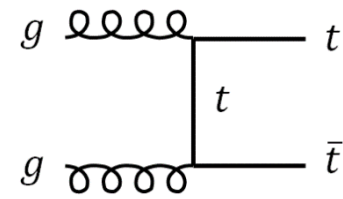
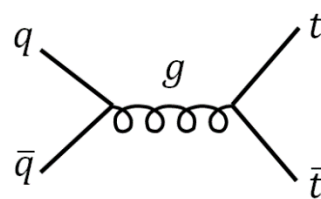
top pair production



EW process → polarized



QCD process → unpolarized

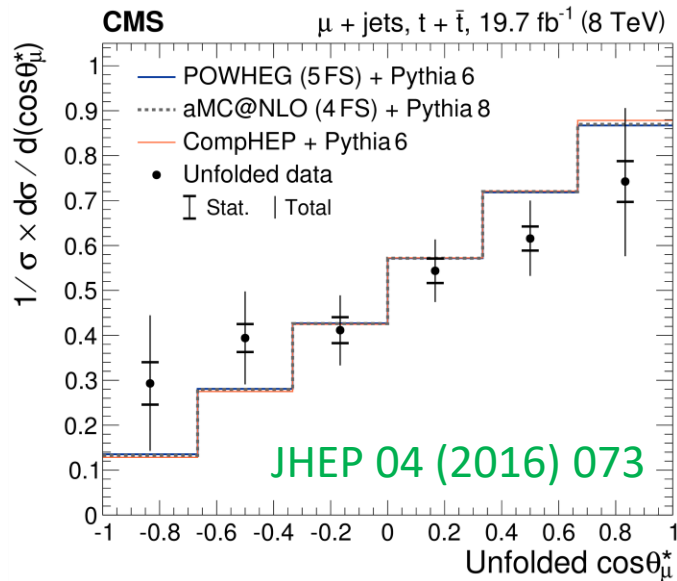


etc.

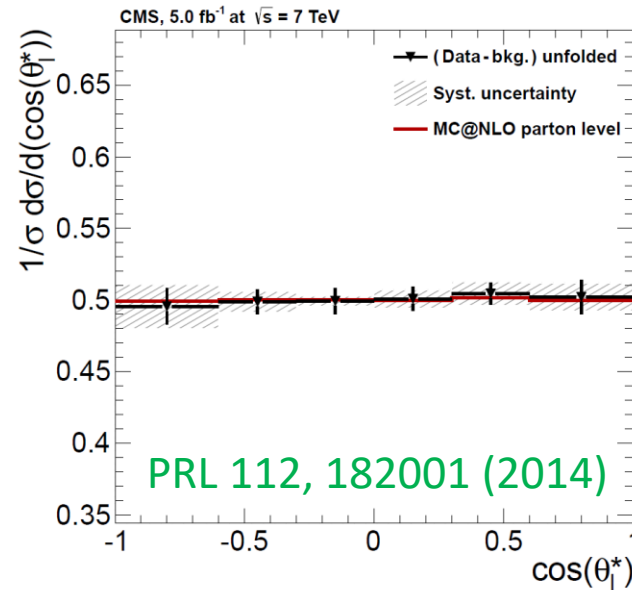
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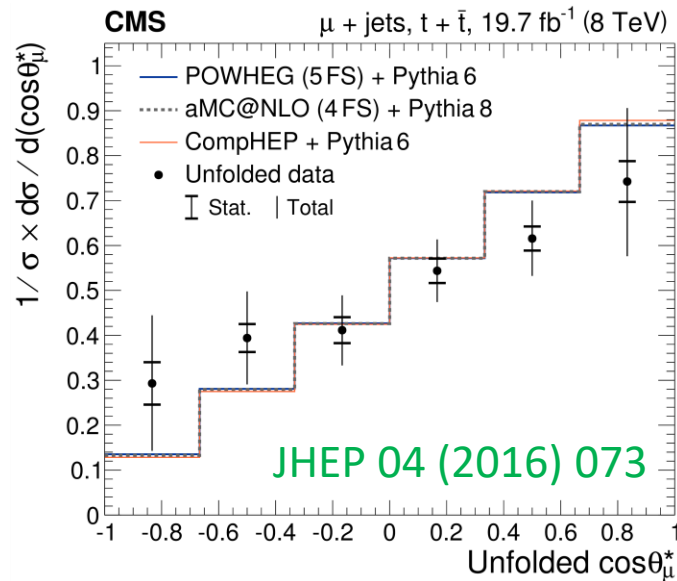
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Polarization of tops from **new physics** processes will teach us about their production mechanism.

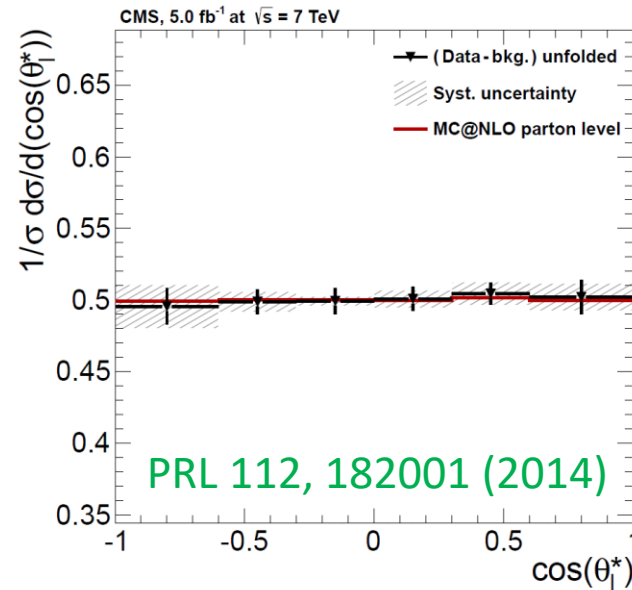
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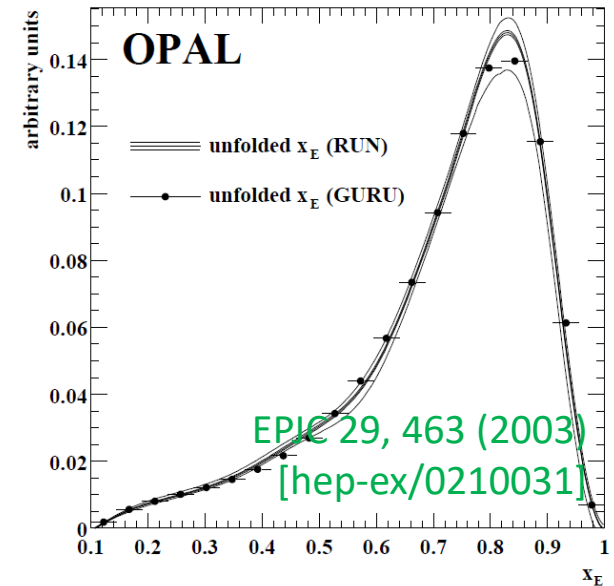
Polarization of tops from **new physics** processes will teach us about their production mechanism.

Can we do analogous measurements for the **other quarks**?

Heavy quarks (b, c)

For heavy quarks, $m_q \gg \Lambda_{\text{QCD}}$

- The quark is carried by a **very energetic** heavy-flavored hadron.



Heavy quarks (b, c)

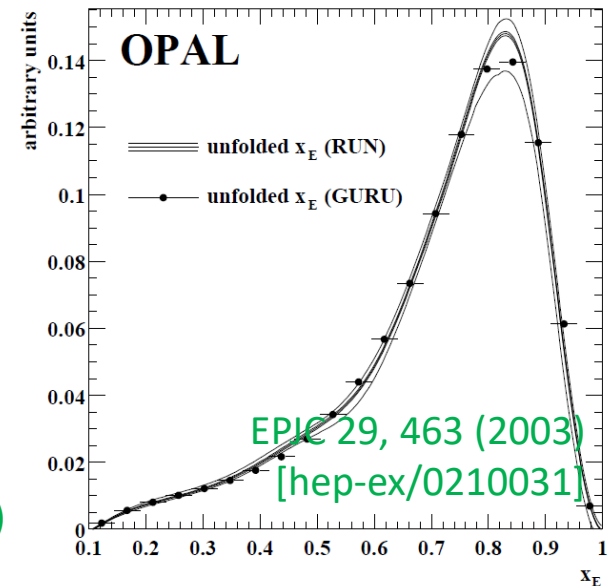
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Mannel and Schuler, PLB 279, 194 (1992)

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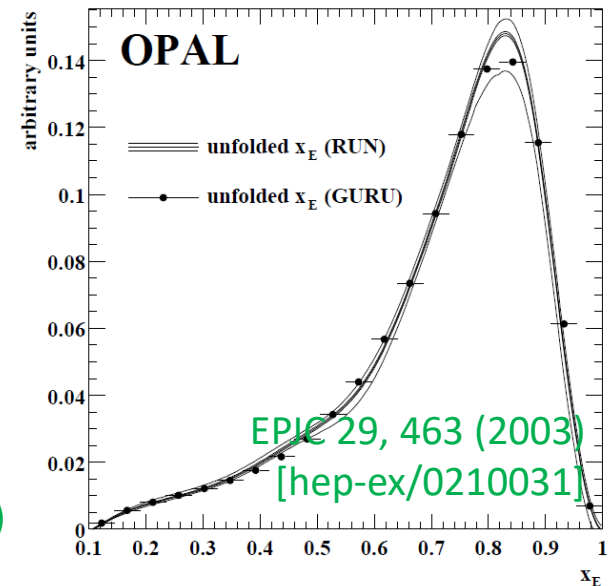
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Evidence observed at **LEP** via Λ_b ($\approx bud$) baryons in $Z \rightarrow b\bar{b}$.

$$\mathcal{P}(\Lambda_b) = -0.23^{+0.24}_{-0.20} \begin{matrix} +0.08 \\ -0.07 \end{matrix} \quad (\text{ALEPH}) \quad \text{PLB 365, 437 (1996)}$$

$$\mathcal{P}(\Lambda_b) = -0.49^{+0.32}_{-0.30} \pm 0.17 \quad (\text{DELPHI}) \quad \text{PLB 474, 205 (2000)}$$

$$\mathcal{P}(\Lambda_b) = -0.56^{+0.20}_{-0.13} \pm 0.09 \quad (\text{OPAL}) \quad \text{PLB 444, 539 (1998) [hep-ex/9808006]}$$

b-quark polarization retention

chromomagnetic
moment

$$\mu_b \propto \frac{1}{m_b}$$

$$m_b \gg \Lambda_{\text{QCD}}$$

b spin **preserved**
during hadronization

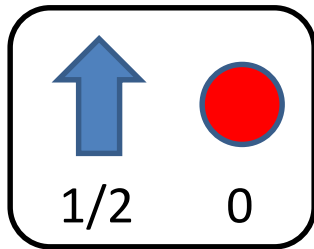
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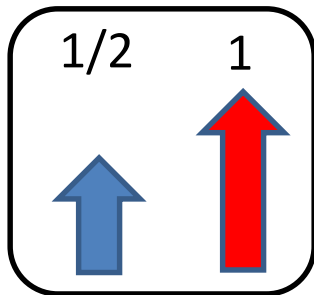
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Λ_b

b spin **preserved**
also during lifetime

b *qq*



Σ_b, Σ_b^*

b spin **oscillates**
during lifetime

Λ_b sample contaminated
by $\Sigma_b^{(*)} \rightarrow \Lambda_b \pi$

fragmentation fraction $f(b \rightarrow \text{baryons}) \approx 8\%$

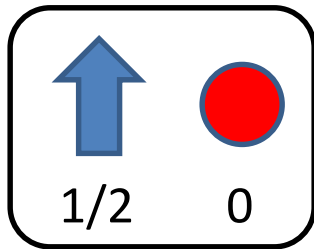
c-quark polarization retention

chromomagnetic
moment

$$\mu_c \propto \frac{1}{m_c}$$

$m_c \gg \Lambda_{\text{QCD}}$ as a rough
approximation

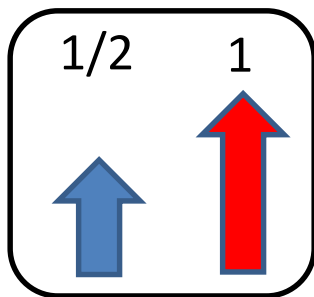
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during hadronization



Λ_c

c spin **preserved**
also during lifetime

c *qq*



$\Sigma_c(2455)$

Σ_c, Σ_c^*

$\Sigma_c(2520)$

c spin **oscillates**
during lifetime

Λ_c sample contaminated
by $\Sigma_c^{(*)} \rightarrow \Lambda_c \pi$

fragmentation fraction $f(c \rightarrow \text{baryons}) \approx 6\%$

b-quark polarization retention

Dominant polarization loss effect

$\Sigma_b^{(*)} \rightarrow \Lambda_b \pi$ decays

$$r \equiv \frac{\mathcal{P}(\Lambda_b)}{\mathcal{P}(b)} = ?$$

b -quark polarization retention

Dominant polarization loss effect

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$$r \equiv \frac{\mathcal{P}(\Lambda_b)}{\mathcal{P}(b)} = ?$$

$$\begin{aligned} |\Lambda_{b,+1/2}\rangle &= |b_{+1/2}\rangle |S_0\rangle \\ |\Sigma_{b,+1/2}\rangle &= -\sqrt{\frac{1}{3}} |b_{+1/2}\rangle |T_0\rangle + \sqrt{\frac{2}{3}} |b_{-1/2}\rangle |T_{+1}\rangle \\ |\Sigma_{b,+1/2}^*\rangle &= \sqrt{\frac{2}{3}} |b_{+1/2}\rangle |T_0\rangle + \sqrt{\frac{1}{3}} |b_{-1/2}\rangle |T_{+1}\rangle \\ |\Sigma_{b,+3/2}^*\rangle &= |b_{+1/2}\rangle |T_{+1}\rangle \end{aligned}$$

“diquarks”

S	T
spin-0	spin-1
isosinglet	isotriplet

Production as a b spin eigenstate.

Decay as a Σ_b or Σ_b^* mass eigenstate.

e.g. $|b_{+1/2}\rangle |T_0\rangle = -\sqrt{\frac{1}{3}} |\Sigma_{b,+1/2}\rangle + \sqrt{\frac{2}{3}} |\Sigma_{b,+1/2}^*\rangle$

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“diquarks”

S	T
spin-0	spin-1
isosinglet	isotriplet

$$A = \frac{\text{prob}(\Sigma_b^{(*)})}{\text{prob}(\Lambda_b)} = 9 \frac{\text{prob}(T)}{\text{prob}(S)}$$

$$w_1 = \frac{\text{prob}(T_{\pm 1})}{\text{prob}(T)} \quad \text{along axis of fragmentation}$$

Falk and Peskin
PRD 49, 3320 (1994)
[hep-ph/9308241]

b -quark polarization retention

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$$r \approx \frac{1 + (1 + 4w_1)A/9}{1 + A}$$

b -quark polarization retention

More precisely, need to account for $\Sigma_b^{(*)}$ widths (interference).

Can do it by considering $\Sigma_b^{(*)}$ propagation:

Parameter	(MeV)
Γ_{Σ_b}	7 ± 3
$\Gamma_{\Sigma_b^*}$	9 ± 2
$m_{\Sigma_b^*} - m_{\Sigma_b}$	21 ± 2

$$|E\rangle \propto \int d\cos\theta d\phi \sum_{J,M} \langle J, M | \frac{1}{2}, +\frac{1}{2}; 1, m \rangle \frac{p_\pi(E)}{E - m_J + i\Gamma(E)/2} \times$$

$$\times \sum_s \langle \frac{1}{2}, s; 1, M - s | J, M \rangle Y_1^{M-s}(\theta, \phi) |\theta, \phi\rangle |s\rangle$$

$$\rho(E) \propto \text{Tr}_{\theta, \phi} |E\rangle \langle E|$$

$$\rho \propto \int_{m_{\Lambda_b} + m_\pi}^{\infty} dE p_\pi(E) \exp(-E/T) \rho(E)$$

phase space

statistical hadronization model ($T \approx 165$ MeV)

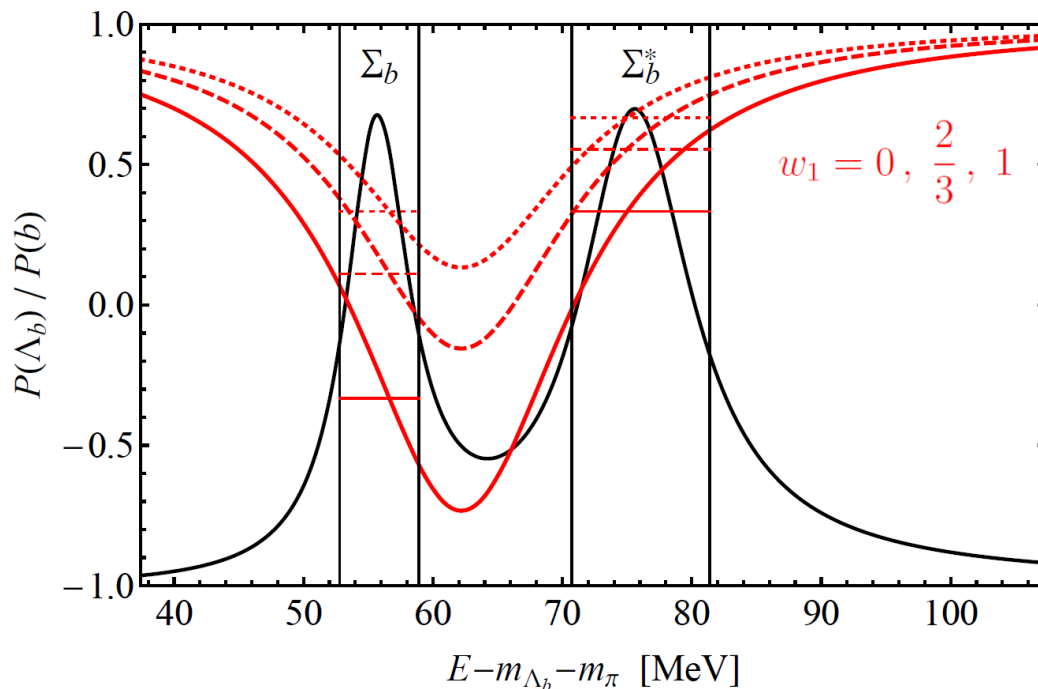
review: PLB 678, 350 (2009) [arXiv:0904.1368]

pion
momentum

Λ_b spin

b -quark polarization retention

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$$r \equiv \frac{\mathcal{P}(\Lambda_b)}{\mathcal{P}(b)} \approx \frac{1 + (0.23 + 0.38w_1)A}{1 + A}$$

b-quark polarization retention

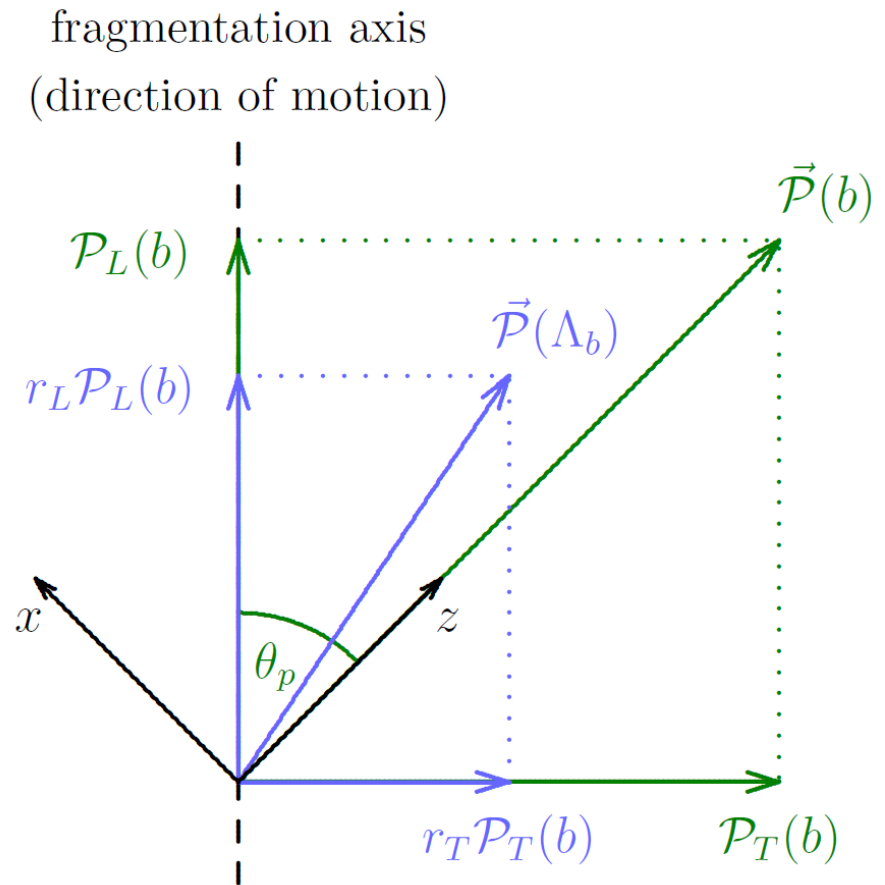
$$r_L \approx \frac{1 + (0.23 + 0.38w_1)A}{1 + A}$$

$$r_T \approx \frac{1 + (0.62 - 0.19w_1)A}{1 + A}$$

Directional dependence, since

$$w_1 = \frac{\text{prob}(T_{\pm 1})}{\text{prob}(T)}$$

holds along the fragmentation axis.



Heavy quark polarization retention

$$r_L \approx \frac{1 + (0.23 + 0.38w_1)A}{1 + A}$$

$$A = \frac{\text{prob}(\Sigma_b^{(*)})}{\text{prob}(\Lambda_b)} = 9 \frac{\text{prob}(T)}{\text{prob}(S)}$$

$$r_T \approx \frac{1 + (0.62 - 0.19w_1)A}{1 + A}$$

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What is known about A and w_1 (for both b and c quarks)?

Heavy quark polarization retention

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What is known about A and w_1 (for both b and c quarks)?

Pythia tunes $0.24 \lesssim A \lesssim 0.45$ (based on light hadron data)

DELPHI (LEP) $1 \lesssim A \lesssim 10$ (b) $w_1 = -0.36 \pm 0.30 \pm 0.30$ (b)
DELPHI-95-107

E791 $A \approx 1.1$ (c) **CLEO (CESR)** $w_1 = 0.71 \pm 0.13$ (c)
PLB 379, 292 (1996) PRL 78, 2304 (1997)

Statistical hadronization $A \approx 2.6$ (b and c)
review: PLB 678, 350 (2009)

Adamov & Goldstein $A \approx 6$ (b and c) $w_1 \approx 0.41$ (b), 0.39 (c)
PRD 64, 014021 (2001)

Heavy quark polarization retention

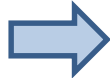
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What is known about A and w_1 (for both b and c quarks)?

Overall, $A \sim \mathcal{O}(1)$, $0 \leq w_1 \leq 1$  $r_L, r_T \sim \mathcal{O}(1)$

r_L consistent with Λ_b

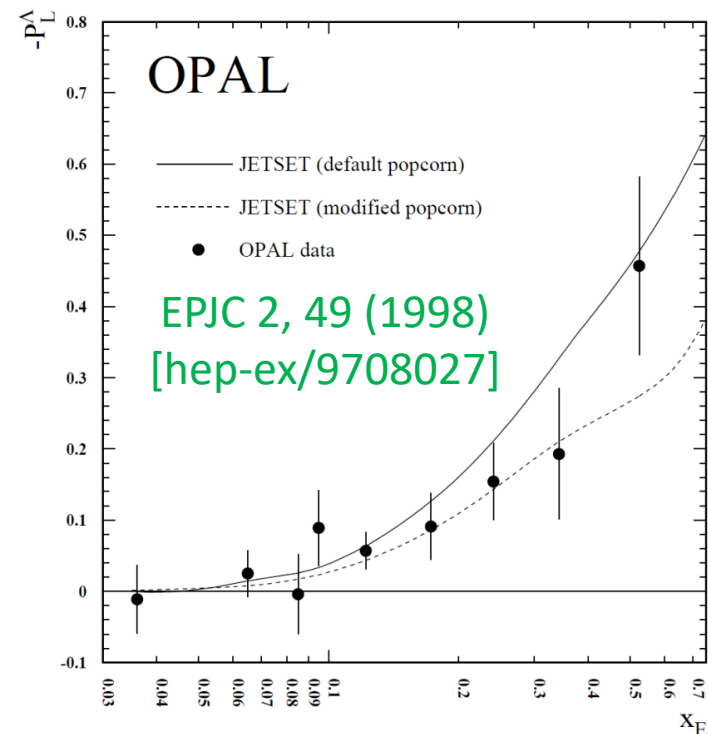
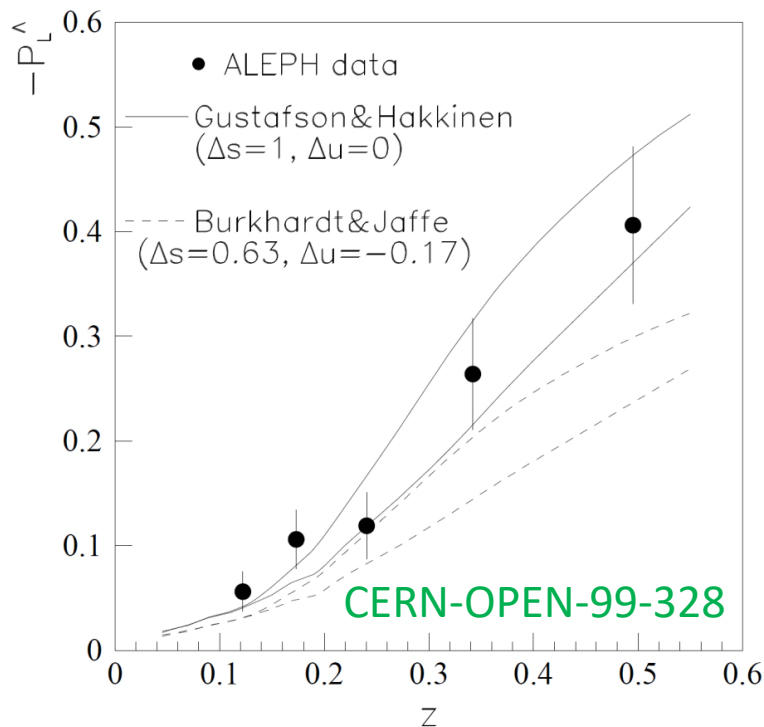
results from LEP

***s*-quark polarization retention?**

- Cannot argue for polarization retention using heavy-quark limit.
Cannot argue for polarization loss either!

s -quark polarization retention!

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- Λ polarization studies were done in Z decays at LEP.



s-quark polarization retention!

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For $z > 0.3$:

$$\mathcal{P}(\Lambda) = -0.31 \pm 0.05 \quad \text{ALEPH, CERN-OPEN-99-328}$$

$$\mathcal{P}(\Lambda) = -0.33 \pm 0.08 \quad \text{OPAL, EPJC 2, 49 (1998) [hep-ex/9708027]}$$

Contributions from all quark flavors are included.

For strange quarks only (non-negligible modeling uncertainty):

$$-0.65 \lesssim \mathcal{P}(\Lambda) \lesssim -0.49$$

Sizable polarization retention!

Nice sources of polarized quarks

Top pair production $pp \rightarrow t\bar{t}$

- $t \rightarrow W^+ b$ produces polarized ***b*** quarks.
 - ↪ $c\bar{s}, u\bar{d}$ produces polarized ***c, s, u, d*** quarks.
- Easy to select a clean $t\bar{t}$ sample (e.g., in lepton + jets).
- Kinematic reconstruction and charm tagging enable studying the different quark flavors separately.
- Statistics in Run 2 is as large as in Z decays at LEP.

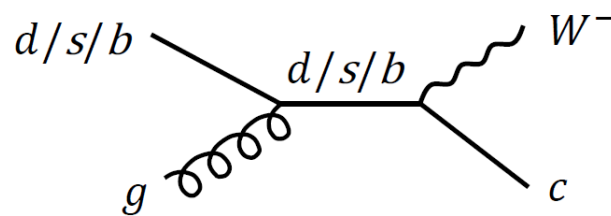
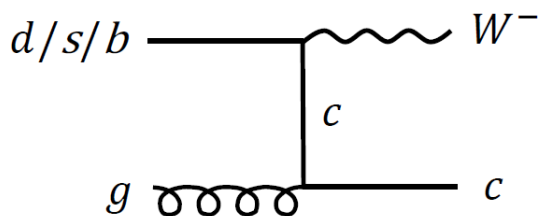
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$W+c$ production $pp \rightarrow W^- c$

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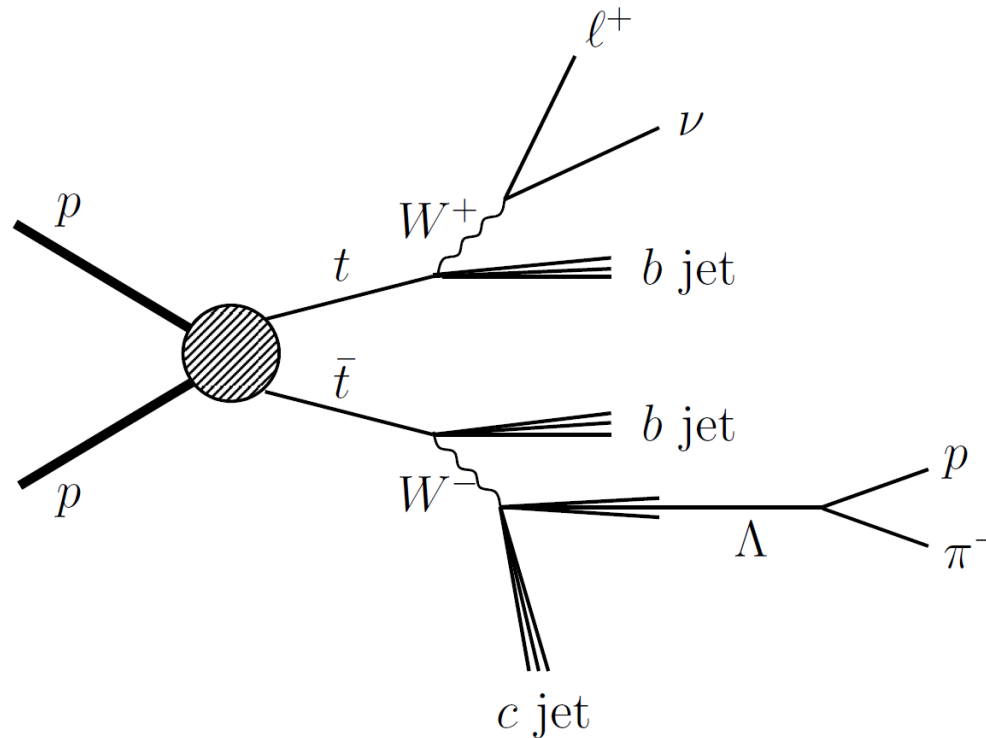
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$W+c$ production $pp \rightarrow W^- c$

- Polarized ***c*** quarks.
- Order-of-magnitude higher statistics than $t\bar{t}$, although backgrounds are higher too.

Measurement of s polarization in $t\bar{t}$

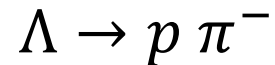


Main steps:

- Typical single-lepton $t\bar{t}$ selection
- Typical kinematic reconstruction and global event interpretation
- Charm tagging
- Λ reconstruction and polarization measurement

Λ polarization measurement

In the Λ rest frame, the decay

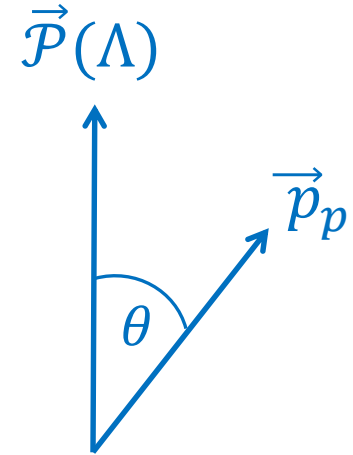


has the angular distribution

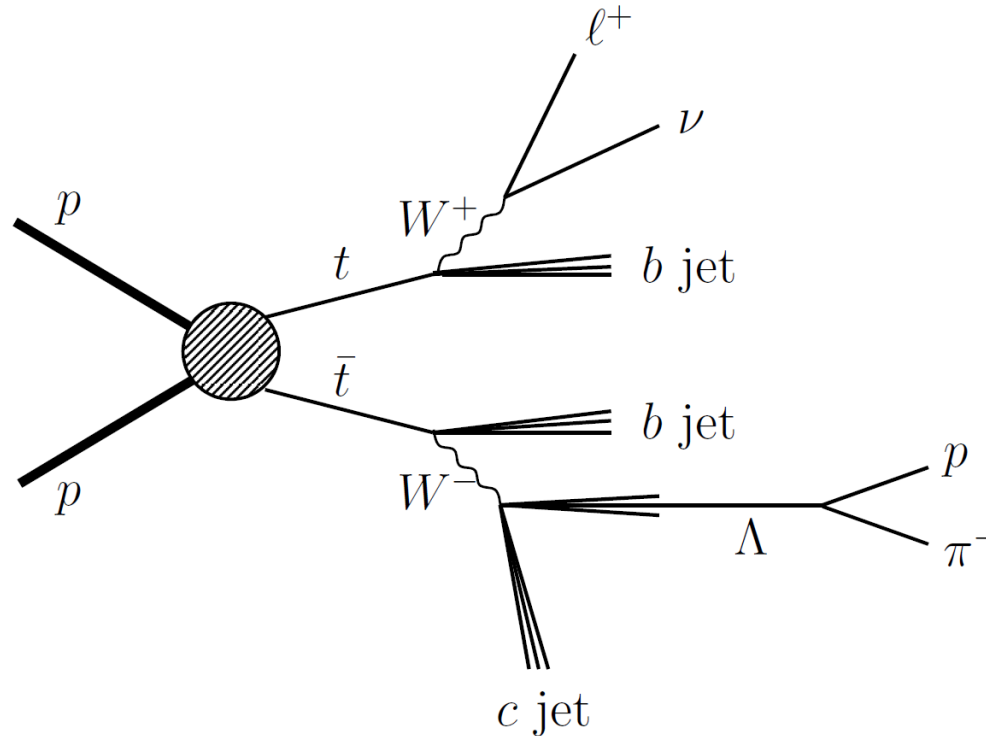
$$\frac{1}{\Gamma} \frac{d\Gamma}{d \cos \theta} = \frac{1}{2} (1 + \alpha \mathcal{P}(\Lambda) \cos \theta)$$

where

$$\alpha = 0.642 \pm 0.013$$

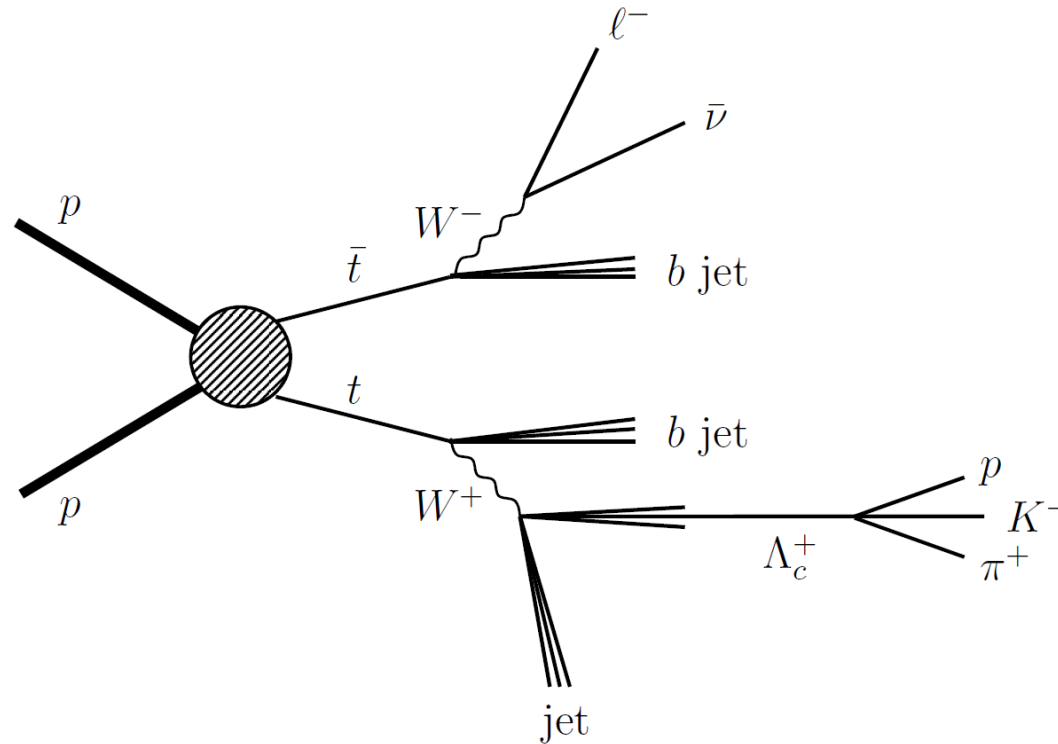


Measurement of s polarization in $t\bar{t}$



Statistical precision of roughly 16% possible
at ATLAS/CMS in Run 2 (with 100 fb^{-1} of data).

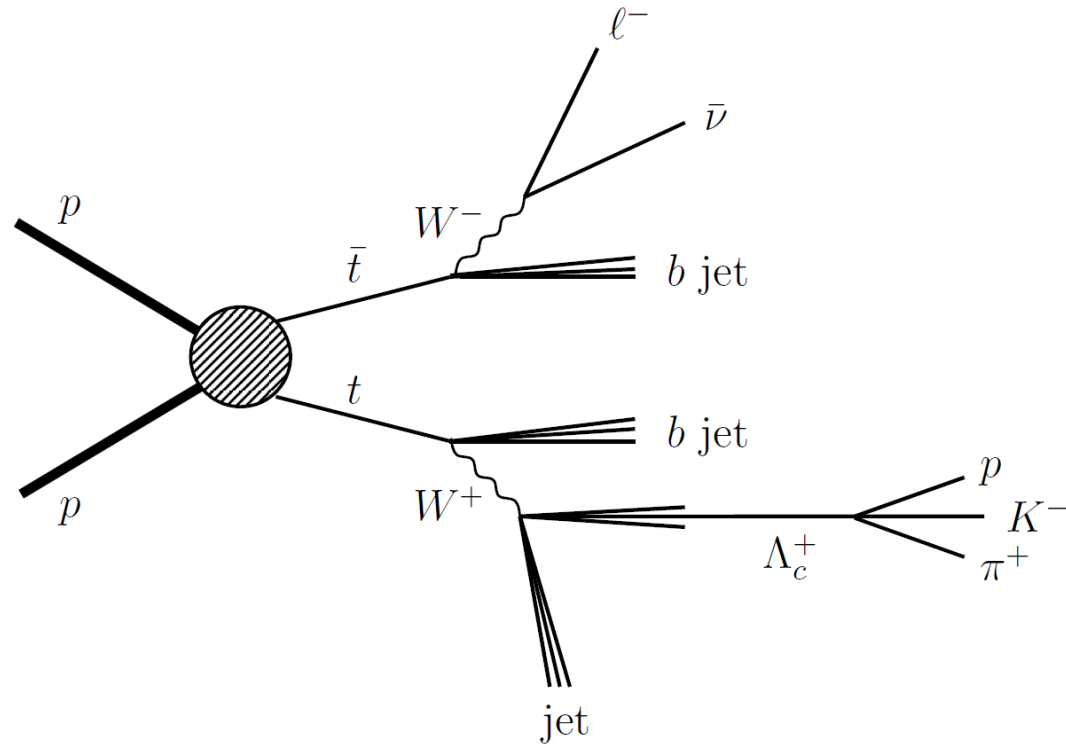
Measurement of c polarization in $t\bar{t}$



Main steps:

- Typical single-lepton $t\bar{t}$ selection
- Typical kinematic reconstruction and global event interpretation
- Λ_c reconstruction and polarization measurement

Measurement of c polarization in $t\bar{t}$

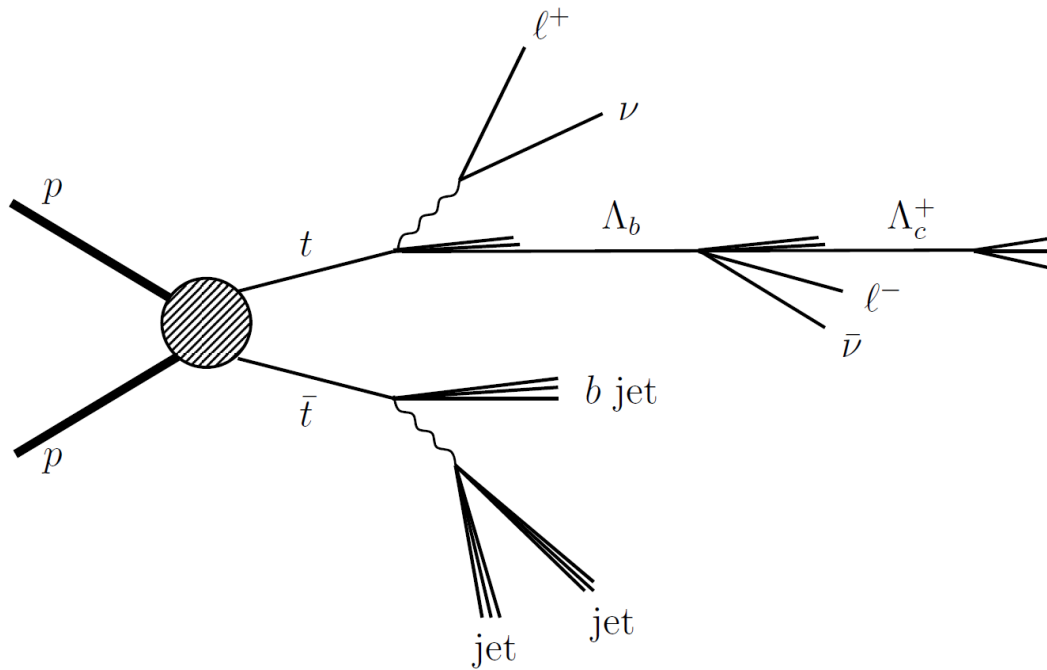


$$\alpha_i r_L = 0.6$$

Selection	Expected events	Purity (example)	$\Delta\mathcal{A}_{FB}/\mathcal{A}_{FB}$
Baseline	$1.7 \times 10^6 t\bar{t} + \mathcal{O}(10^5)$ bkg		
$\Lambda_c^+ \rightarrow pK^-\pi^+$	$810 \times (\epsilon_{\Lambda_c}/25\%)$	20%	26%
		100%	11%

Statistical precision of order 10% possible
at ATLAS/CMS in Run 2 (with 100 fb^{-1} of data).

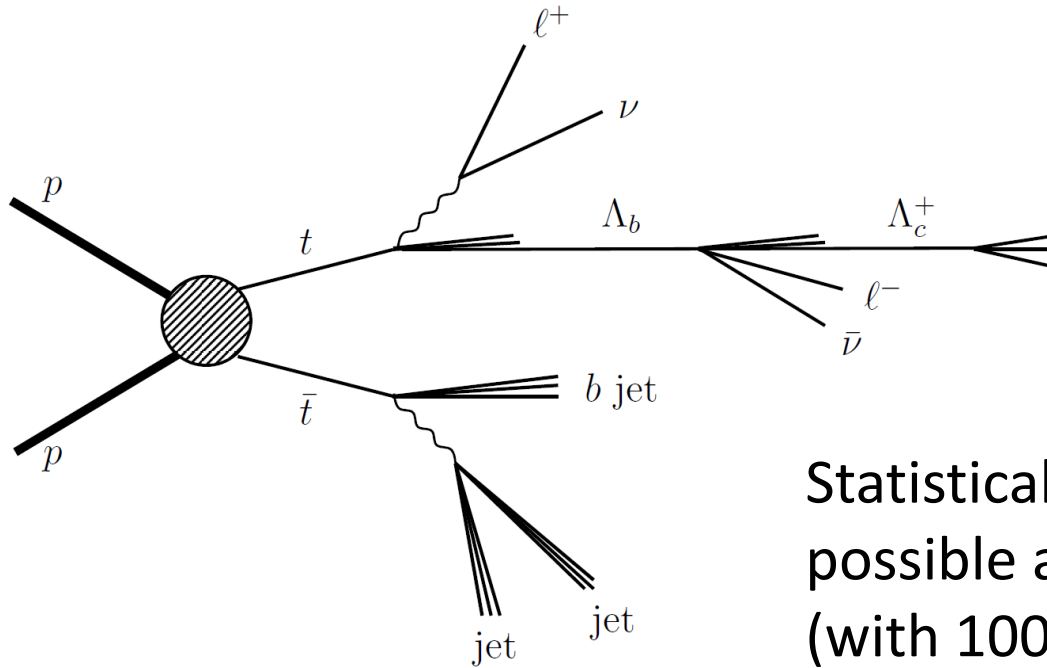
Measurement of b polarization in $t\bar{t}$



Main steps:

- Typical single-lepton $t\bar{t}$ selection (w/soft-muon b tag)
- Typical kinematic reconstruction and global event interpretation
- Λ_b reconstruction (using inclusive, semi-inclusive or exclusive approach) and polarization measurement

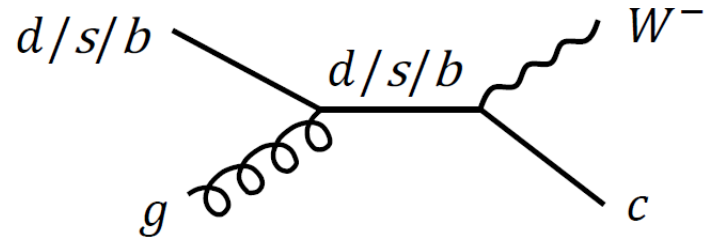
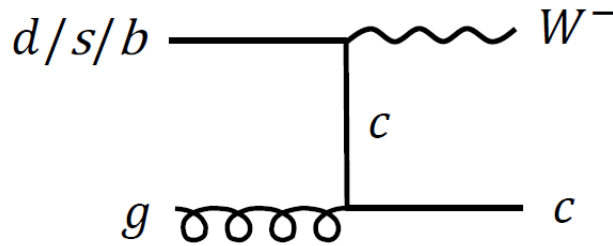
Measurement of b polarization in $t\bar{t}$



Statistical precision of about 10% possible at ATLAS/CMS in Run 2 (with 100 fb^{-1} of data)

Selection	Expected events		
Baseline	$3 \times 10^6 t\bar{t} + \mathcal{O}(10^6) \text{ bkg}$		
Soft-muon b tagging	$5 \times 10^5 t\bar{t} + \mathcal{O}(10^4) \text{ bkg}$		$r_L = 0.6$
Signal events ($t \rightarrow b \rightarrow \Lambda_b \rightarrow \mu\nu X_c$)	Purity (example)		$\Delta\mathcal{A}_{FB}/\mathcal{A}_{FB}$
Inclusive	34 400	$\mathcal{O}(f_{\text{baryon}})$ (e.g., 7%)	$\pm 7\%$
Semi-inclusive	$2300 \times (\epsilon_{\Lambda}/30\%)$	70%	$\pm 8\%$
Exclusive	$1040 \times (\epsilon_{\Lambda_c}/25\%)$	30%	$\pm 19\%$
		100%	$\pm 10\%$

Measurement of c polarization in $W+c$



ATLAS and CMS measured $W+c$ cross section at 7 TeV

ATLAS, JHEP 1405, 068 (2014) [arXiv:1402.6263]

CMS, JHEP 1402, 013 (2014) [arXiv:1310.1138]

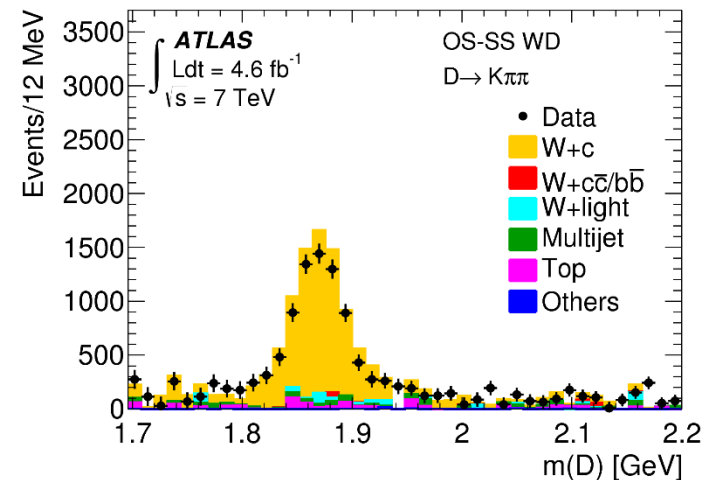
in particular by relying on the decays

$$D^+ \rightarrow K^- \pi^+ \pi^+$$

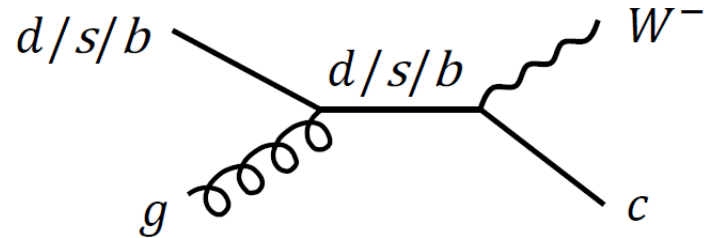
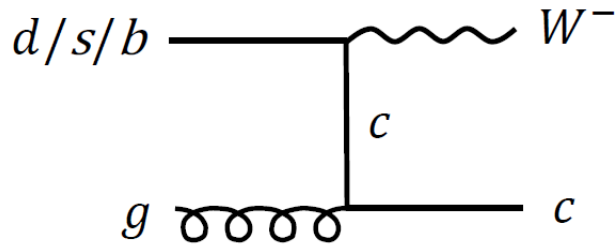
Similar to our decay of interest

$$\Lambda_c^+ \rightarrow p K^- \pi^+$$

(See backup slides for more details.)



Measurement of c polarization in $W+c$



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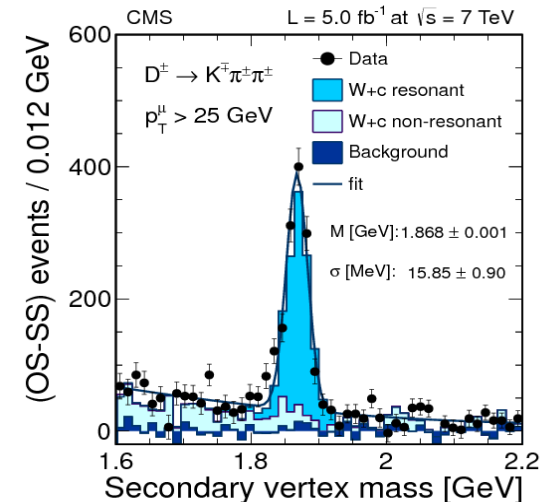
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Measurement of b polarization in QCD events

Inclusive QCD production: $pp \rightarrow b\bar{b} + X$

- Enormous cross section, but **unpolarized** at the leading order.

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- At NLO \rightarrow **transverse** polarization

(an opportunity to measure r_T)

\rightarrow strong kinematic dependence

\rightarrow suppressed at high momenta

$$\mathcal{P}(b) \sim \alpha_s \frac{m_b}{p_b}$$

Bernreuther, Brandenburg, Uwer, PLB 368, 153 (1996)

Dharmaratna and Goldstein, PRD 53, 1073 (1996)

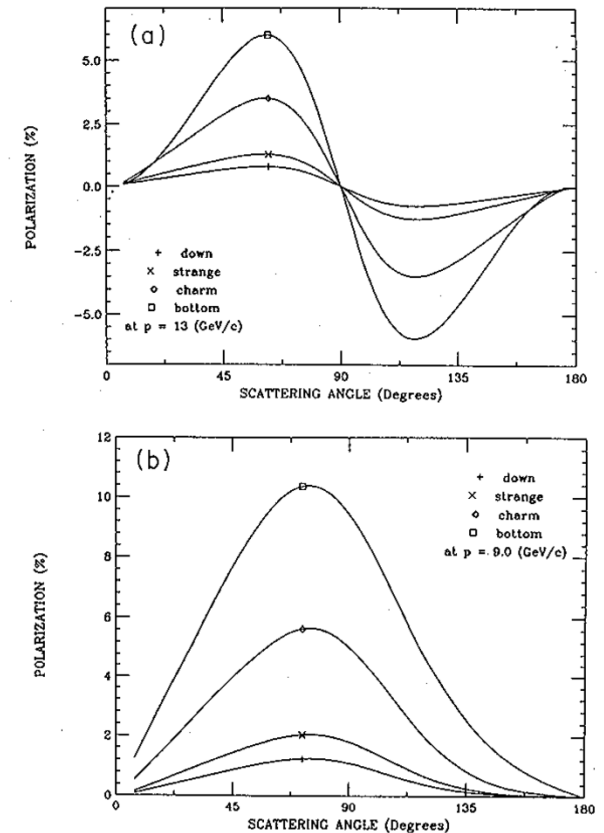


FIG. 7. Polarization of up, strange, charm, and bottom quarks at the subprocess CM momentum of (a) 13 GeV/c for gluon fusion and (b) 9 GeV/c for annihilation. Other parameters are identical to Fig. 5.

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Existing LHCb analysis:

Measurements of the $\Lambda_b^0 \rightarrow J/\psi \Lambda$ decay amplitudes and the Λ_b^0 polarisation in pp collisions at $\sqrt{s} = 7 \text{ TeV}$

PLB 724, 27 (2013)
[arXiv:1302.5578]

$$\mathcal{P}(\Lambda_b) = 0.06 \pm 0.07 \pm 0.02$$

Suboptimal due to inclusiveness over the kinematics.

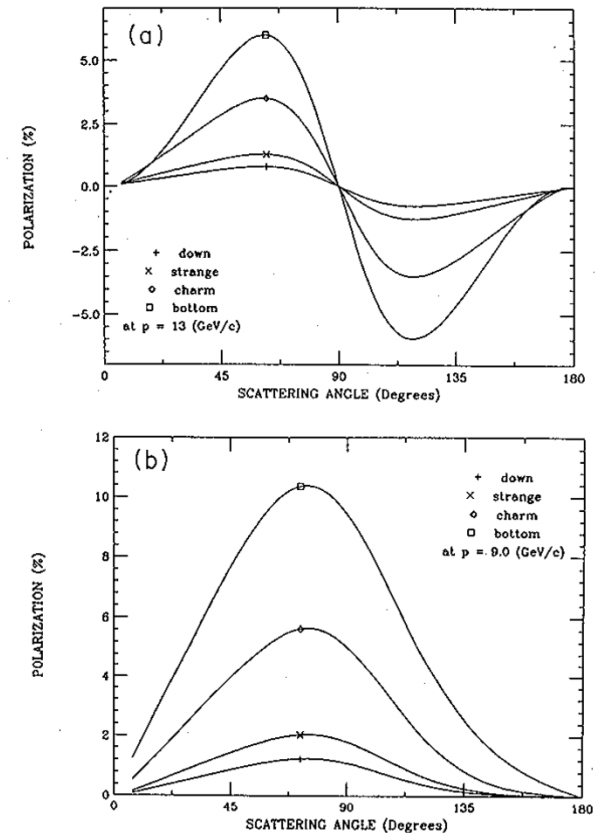


FIG. 7. Polarization of up, strange, charm, and bottom quarks at the subprocess CM momentum of (a) 13 GeV/c for gluon fusion and (b) 9 GeV/c for annihilation. Other parameters are identical to Fig. 5.

Measuring A directly

A is simply the ratio of the $\Sigma_b^{(*)}$ and direct Λ_b yields, independent of the b polarization:

$$A = \frac{\text{prob}(\Sigma_b^{(*)})}{\text{prob}(\Lambda_b)} = 9 \frac{\text{prob}(T)}{\text{prob}(S)}$$

Can be measured by any experiment that can reconstruct

$$\Sigma_b^{(*)\pm,0} \rightarrow \Lambda_b \pi^{\pm,0}$$

In particular, LHCb, ATLAS, CMS in inclusive QCD samples.

Could have been done even at the Tevatron.

CDF, PRL 99, 202001 (2007) [arXiv:0706.3868]

CDF, PRD 85, 092011 (2012) [arXiv:1112.2808]

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$$\Sigma_b^{(*)\pm,0} \rightarrow \Lambda_b \pi^{\pm,0}$$

In particular, LHCb, ATLAS, CMS in inclusive QCD samples.

Same holds for

$$\Sigma_c^{(*)++,+0} \rightarrow \Lambda_c^+ \pi^{\pm,0}$$

where B factories can also help.

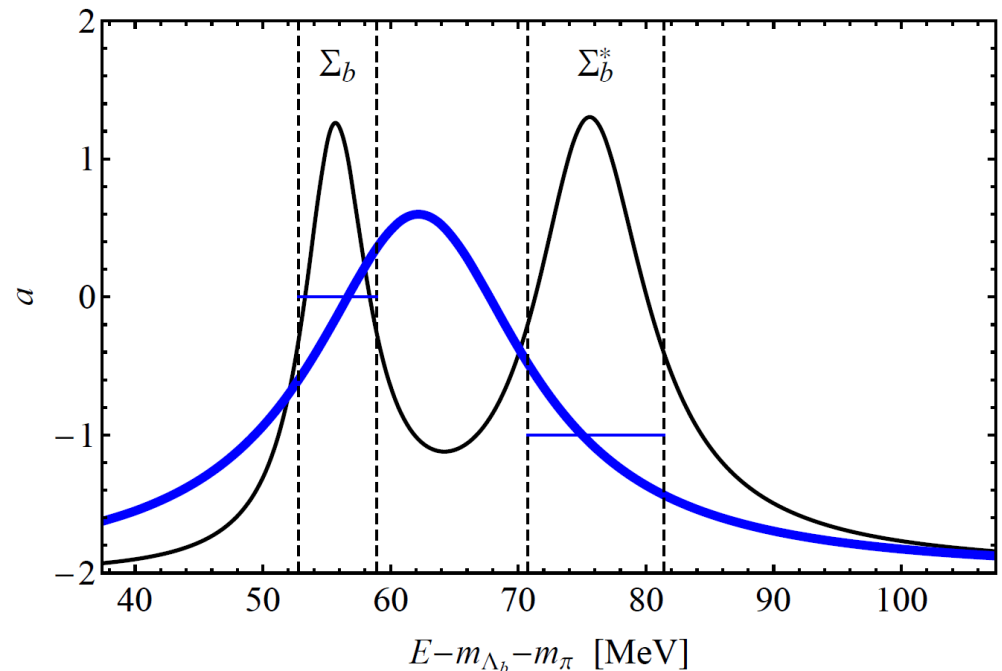
Measuring w_1 directly

The angular distribution of $\Sigma_b^{(*)} \rightarrow \Lambda_b \pi$ is sensitive to w_1 , independent of the b polarization:

$$\frac{1}{\Gamma} \frac{d\Gamma}{d \cos \theta} = \frac{1}{2} + \frac{9}{8} a \left(w_1 - \frac{2}{3} \right) \left(\cos^2 \theta - \frac{1}{3} \right)$$

where a is given in the plot.

Can be measured by any experiment that can reconstruct these decays (e.g., LHCb, ATLAS, CMS).



Same holds for $\Sigma_c^{(*)}$ and Λ_c .

Summary: motivated measurements

In $t\bar{t}$ production (ATLAS, CMS)

- Longitudinal Λ_b polarization in b jets $\rightarrow r_L$ for bottom
- Longitudinal Λ_c polarization in c jets $\rightarrow r_L$ for charm
- Longitudinal Λ polarization in s jets \rightarrow long. pol. FF for strange
- (far future) Longitudinal Λ polarization in u, d jets \rightarrow
long. pol. FFs for up, down

In $W+c$ production (ATLAS, CMS, maybe LHCb)

- Longitudinal Λ_c polarization $\rightarrow r_L$ for charm
(Esp. LHCb may also try separating out the $\Sigma_c^{(*)}$ contributions.)

In QCD production (LHCb, ATLAS, CMS)

- Transverse Λ_b (and Λ_c ?) polarization
as a function of the event kinematics $\rightarrow r_T$ for bottom (charm?)

Summary: motivated measurements

In QCD production (LHCb, ATLAS, CMS)

- $\Sigma_b^{(*)}$ yields (relative to direct Λ_b) $\rightarrow A$ for bottom
and pion angular distribution $\rightarrow w_1$
- $\Sigma_c^{(*)}$ yields (relative to direct Λ_c) $\rightarrow A$ for charm
and pion angular distribution $\rightarrow w_1$

In new-physics samples, once discovered (ATLAS, CMS)

- Measure quark polarizations \rightarrow learn about the new physics
(Statistics will likely be a severe limitation.)

In $t\bar{t}$ and $W+c$ production in the long term (ATLAS, CMS, LHCb)

- Measurements of polarized fragmentation functions.

Thank You!

Supplementary Slides

Mass splittings and widths

bottom system

$$m_{\Lambda_b} = 5619.5 \pm 0.4 \text{ MeV}$$

Parameter	(MeV)
$m_{\Sigma_b} - m_{\Lambda_b}$	194 ± 2
$m_{\Sigma_b^*} - m_{\Lambda_b}$	214 ± 2
$\Delta \equiv m_{\Sigma_b^*} - m_{\Sigma_b}$	21 ± 2
Γ_{Σ_b}	7 ± 3
$\Gamma_{\Sigma_b^*}$	9 ± 2

charm system

$$m_{\Lambda_c} = 2286.5 \pm 0.2 \text{ MeV}$$

Parameter	(MeV)
$m_{\Sigma_c} - m_{\Lambda_c}$	167.4 ± 0.1
$m_{\Sigma_c^*} - m_{\Lambda_c}$	231.9 ± 0.4
$\Delta \equiv m_{\Sigma_c^*} - m_{\Sigma_c}$	64.5 ± 0.5
Γ_{Σ_c}	2.2 ± 0.2
$\Gamma_{\Sigma_c^*}$	15 ± 1

Measurement of b polarization in Z decays

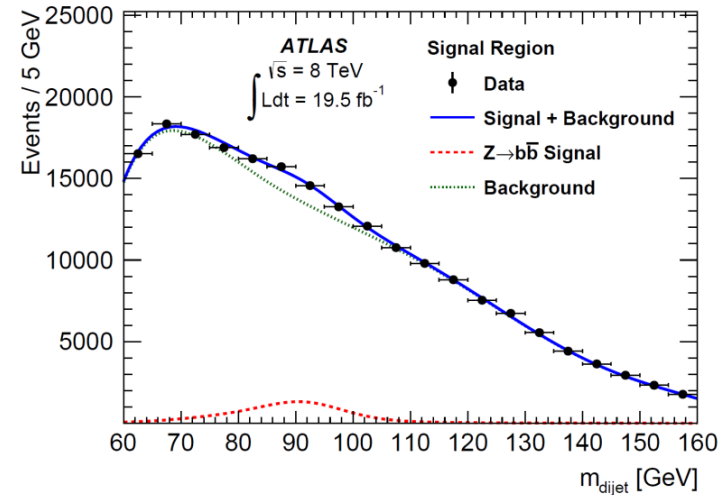
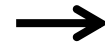
Z production: $pp \rightarrow Z \rightarrow b\bar{b}$

- Longitudinally polarized b quarks (similar to $t\bar{t}$)
- Large cross section

$$\frac{\sigma(pp \rightarrow Z \rightarrow b\bar{b})}{\sigma(pp \rightarrow t\bar{t} \rightarrow W^+W^-b\bar{b})} \sim 10$$

- Large QCD background (at 8 TeV, $S/B \approx 1/15$ even for $p_T^Z > 200$ GeV) dilutes the asymmetry.

Probably less effective than $t\bar{t}$.



PLB 738, 25 (2014)
[arXiv:1404.7042]

Λ_b polarization measurement

Which Λ_b decay to use?

We picked **semileptonic** mode **inclusive** in charm hadrons (large BR, no hadronic uncertainties).

	Mode	Fraction (Γ_i/Γ)
Γ_1	$J/\psi(1S)\Lambda \times B(b \rightarrow \Lambda_b^0)$	$(5.8 \pm 0.8) \times 10^{-5}$
Γ_2	$\rho D^0 \pi^-$	$(5.9^{+4.0}_{-3.2}) \times 10^{-4}$
Γ_3	$\rho D^0 K^-$	$(4.3^{+3.0}_{-2.4}) \times 10^{-5}$
Γ_4	$\Lambda_c^+ \pi^-$	$(5.7^{+4.0}_{-2.6}) \times 10^{-3}$
Γ_5	$\Lambda_c^+ K^-$	$(4.2^{+2.6}_{-1.9}) \times 10^{-4}$
Γ_6	$\Lambda_c^+ a_1(1260)^-$	seen
Γ_7	$\Lambda_c^+ \pi^+ \pi^- \pi^-$	$(8^{+5}_{-4}) \times 10^{-3}$
Γ_8	$\Lambda_c(2595)^+ \pi^-$, $\Lambda_c(2595)^+ \rightarrow \Lambda_c^+ \pi^+ \pi^-$	$(3.7^{+2.8}_{-2.3}) \times 10^{-4}$
Γ_9	$\Lambda_c(2625)^+ \pi^-$, $\Lambda_c(2625)^+ \rightarrow \Lambda_c^+ \pi^+ \pi^-$	$(3.6^{+2.7}_{-2.1}) \times 10^{-4}$
Γ_{10}	$\Sigma_c(2455)^0 \pi^+ \pi^-$, $\Sigma_c^0 \rightarrow \Lambda_c^+ \pi^-$	$(6^{+5}_{-4}) \times 10^{-4}$
Γ_{11}	$\Sigma_c(2455)^{++} \pi^- \pi^-$, $\Sigma_c^{++} \rightarrow \Lambda_c^+ \pi^+$	$(3.5^{+2.8}_{-2.3}) \times 10^{-4}$
Γ_{12}	$\Lambda K^0 2\pi^+ 2\pi^-$	
Γ_{13}	$\Lambda_c^+ \ell^- \bar{\nu}_\ell$ anything	[a] $(9.9 \pm 2.2) \%$
Γ_{14}	$\Lambda_c^+ \ell^- \bar{\nu}_\ell$	$(6.5^{+3.2}_{-2.5}) \%$
Γ_{15}	$\Lambda_c^+ \pi^+ \pi^- \ell^- \bar{\nu}_\ell$	$(5.6 \pm 3.1) \%$
Γ_{16}	$\Lambda_c(2595)^+ \ell^- \bar{\nu}_\ell$	$(8 \pm 5) \times 10^{-3}$
Γ_{17}	$\Lambda_c(2625)^+ \ell^- \bar{\nu}_\ell$	$(1.4^{+0.9}_{-0.7}) \%$
Γ_{18}	$\Sigma_c(2455)^0 \pi^+ \ell^- \bar{\nu}_\ell$	
Γ_{19}	$\Sigma_c(2455)^{++} \pi^- \ell^- \bar{\nu}_\ell$	
Γ_{20}	ρh^-	[b] $< 2.3 \times 10^{-5}$
Γ_{21}	$\rho \pi^-$	$(4.1 \pm 0.8) \times 10^{-6}$
Γ_{22}	ρK^-	$(4.9 \pm 0.9) \times 10^{-6}$
Γ_{23}	$\Lambda \mu^+ \mu^-$	$(1.08 \pm 0.28) \times 10^{-6}$
Γ_{24}	$\Lambda \gamma$	$< 1.3 \times 10^{-3}$

Λ_b polarization measurement

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Includes also:

$$\Lambda_b \rightarrow p D^0 \ell^- \bar{\nu}_\ell \quad \text{small contribution}$$

	Mode	Fraction (Γ_i/Γ)
Γ_1	$J/\psi(1S) \Lambda \times B(b \rightarrow \Lambda_b^0)$	$(5.8 \pm 0.8) \times 10^{-5}$
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Γ_{17}	$\Lambda_c(2625)^+ \ell^- \bar{\nu}_\ell$	$(1.4^{+0.9}_{-0.7}) \%$
Γ_{18}	$\Sigma_c(2455)^0 \pi^+ \ell^- \bar{\nu}_\ell$	
Γ_{19}	$\Sigma_c(2455)^{++} \pi^- \ell^- \bar{\nu}_\ell$	
Γ_{20}	$p h^-$	[b] $< 2.3 \times 10^{-5}$
Γ_{21}	$p \pi^-$	$(4.1 \pm 0.8) \times 10^{-6}$
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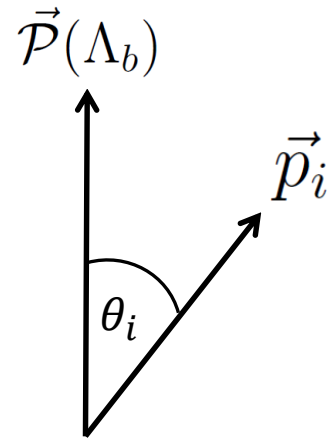
Λ_b polarization measurement

For the inclusive semileptonic decays

$$\Lambda_b \rightarrow X_c \ell^- \bar{\nu}$$

Λ_b polarization is encoded in the angular distributions

$$\frac{1}{\Gamma_{\Lambda_b}} \frac{d\Gamma_{\Lambda_b}}{d \cos \theta_i} = \frac{1}{2} (1 + \alpha_i \mathcal{P}(\Lambda_b) \cos \theta_i) \quad i = \ell \text{ or } \nu$$



where

$$\alpha_\ell = \frac{-\frac{1}{3} + 4x_c + 12x_c^2 - \frac{44}{3}x_c^3 - x_c^4 + 12x_c^2 \log x_c + 8x_c^3 \log x_c}{1 - 8x_c + 8x_c^3 - x_c^4 - 12x_c^2 \log x_c} \approx -0.26$$

$$\alpha_\nu = 1$$

$x_c = \frac{m_c^2}{m_b^2}$

$\mathcal{O}(\Lambda_{\text{QCD}}/m_b)$ corrections are absent, and α_s corrections are few %.

Manohar, Wise
PRD 49, 1310 (1994)
[hep-ph/9308246]

Czarnecki, Jezabek, Korner, Kuhn, PRL 73, 384 (1994)
Czarnecki, Jezabek, NPB 427, 3 (1994)

Λ_b polarization measurement

$$\Lambda_b \rightarrow X_c \ell^- \bar{\nu} \quad (\text{BR} \approx 10\% \text{ per flavor})$$

- Soft-muon b tagging e.g. CMS-PAS-BTV-09-001

- Neutrino reconstruction using...
 - Λ_b mass constraint Dambach, Langenegger, Starodumov
 - Λ_b flight direction NIMA 569, 824 (2006) [hep-ph/0607294]

- Neutrino A_{FB} measurement (in the Λ_b rest frame)

- Approaches regarding semileptonic B -meson background:
 - Inclusive** keep it
 - Semi-inclusive** demand $\Lambda \rightarrow p\pi^-$ among decay products
 - Exclusive** demand a fully-reconstructible Λ_c decay

See paper for many additional details...

Λ_c polarization measurement

$$\Lambda_c^+ \rightarrow pK^- \pi^+ \quad (\text{BR} \approx 6.7\%)$$

- Three tracks reconstructing the Λ_c mass.
- Backgrounds under the mass peak can be suppressed in various ways.
- Spin analyzing powers α_i seem to be large for K^- , small for p and π^+ .

NA32: Jeżabek, Rybicki, Ryłko, PLB 286, 175 (1992)

Precise values not essential for new physics samples if SM calibration samples are available.

Also, α_i can be determined (e.g., in LHCb) from a sample of

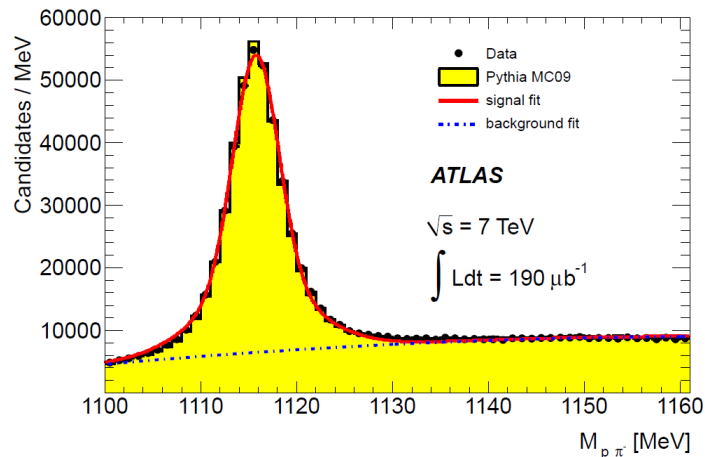
Λ_c 's produced from inclusive b -hadron decays

by calibrating on $\Lambda_c^+ \rightarrow \Lambda \pi^+$ (where $\alpha_\Lambda = -0.91 \pm 0.15$).

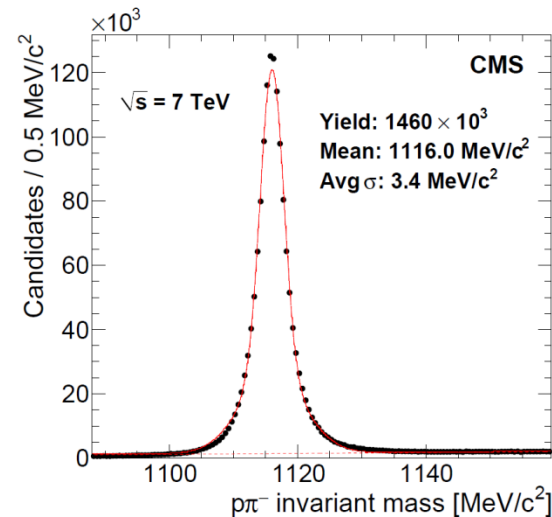
Λ polarization measurement

$$\Lambda \rightarrow p \pi^- \quad (\text{BR} \approx 64\%)$$

- Pair of tracks from a highly displaced vertex reconstructing the Λ mass.
- Spin analyzing power $\alpha \approx 0.64$
- ATLAS and CMS already have experience with Λ 's

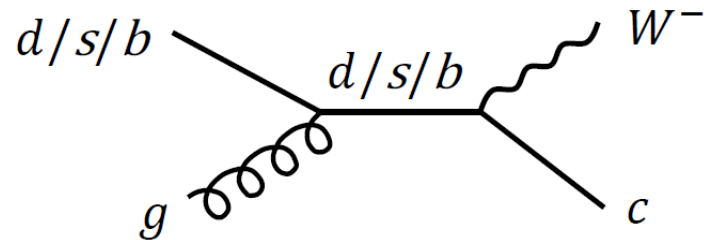
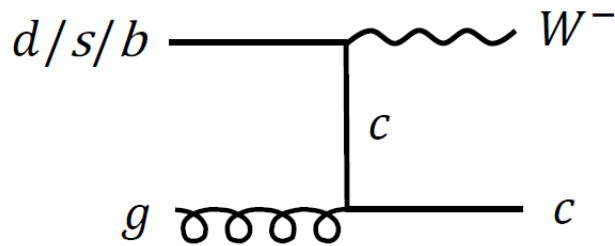


PRD 85, 012001 (2012)
[arXiv:1111.1297]



JHEP 05, 064 (2011) [arXiv:1102.4282]

Measurement of c polarization in $W+c$



ATLAS and CMS measured $W+c$ cross section at 7 TeV

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CMS, JHEP 1402, 013 (2014) [arXiv:1310.1138]

in particular by relying on the decays

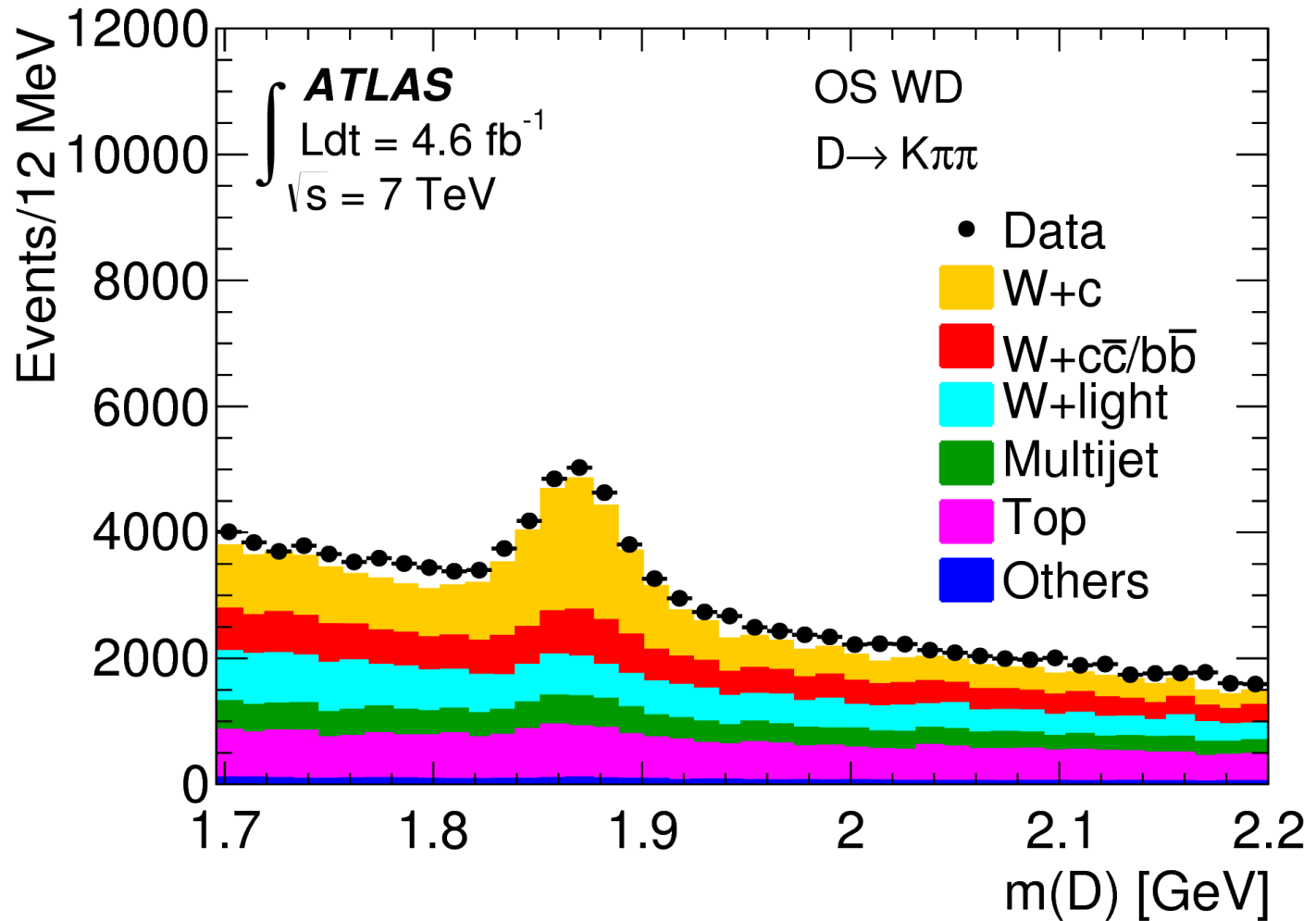
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Similar to our decay of interest

$$\Lambda_c^+ \rightarrow p K^- \pi^+$$

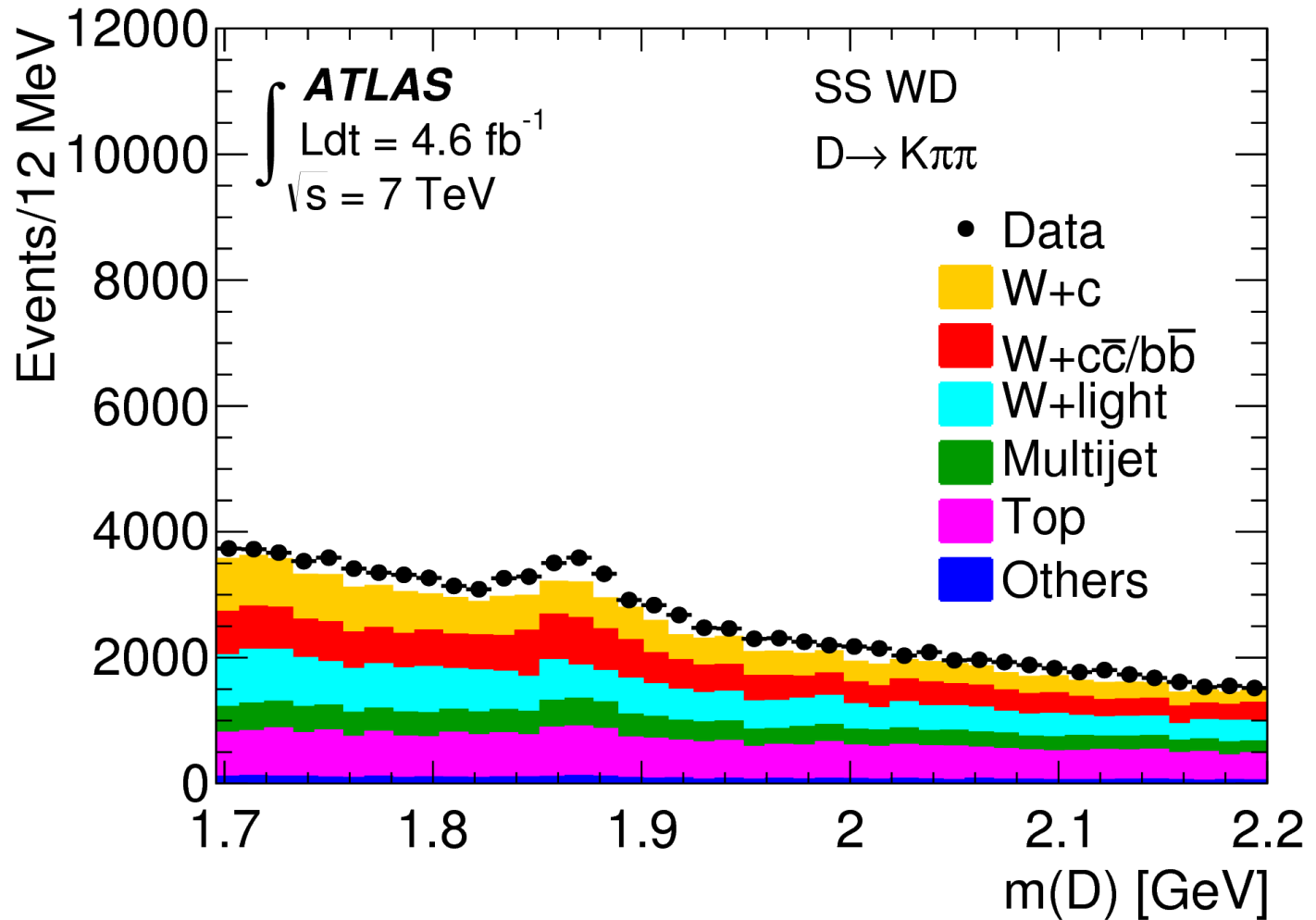
Example: $D^+ \rightarrow K^- \pi^+ \pi^+$ in ATLAS

ATLAS, JHEP 1405, 068 (2014) [arXiv:1402.6263]



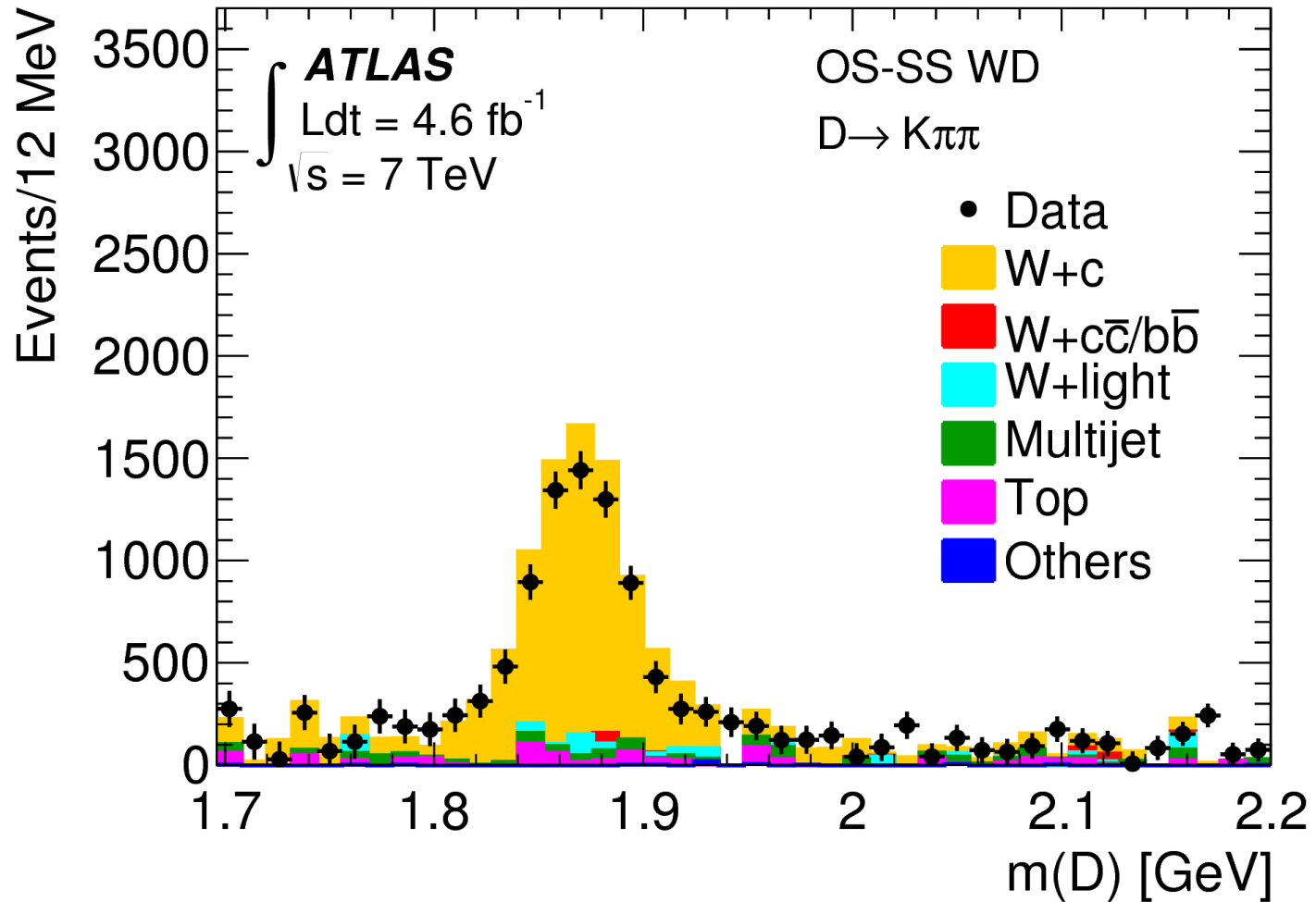
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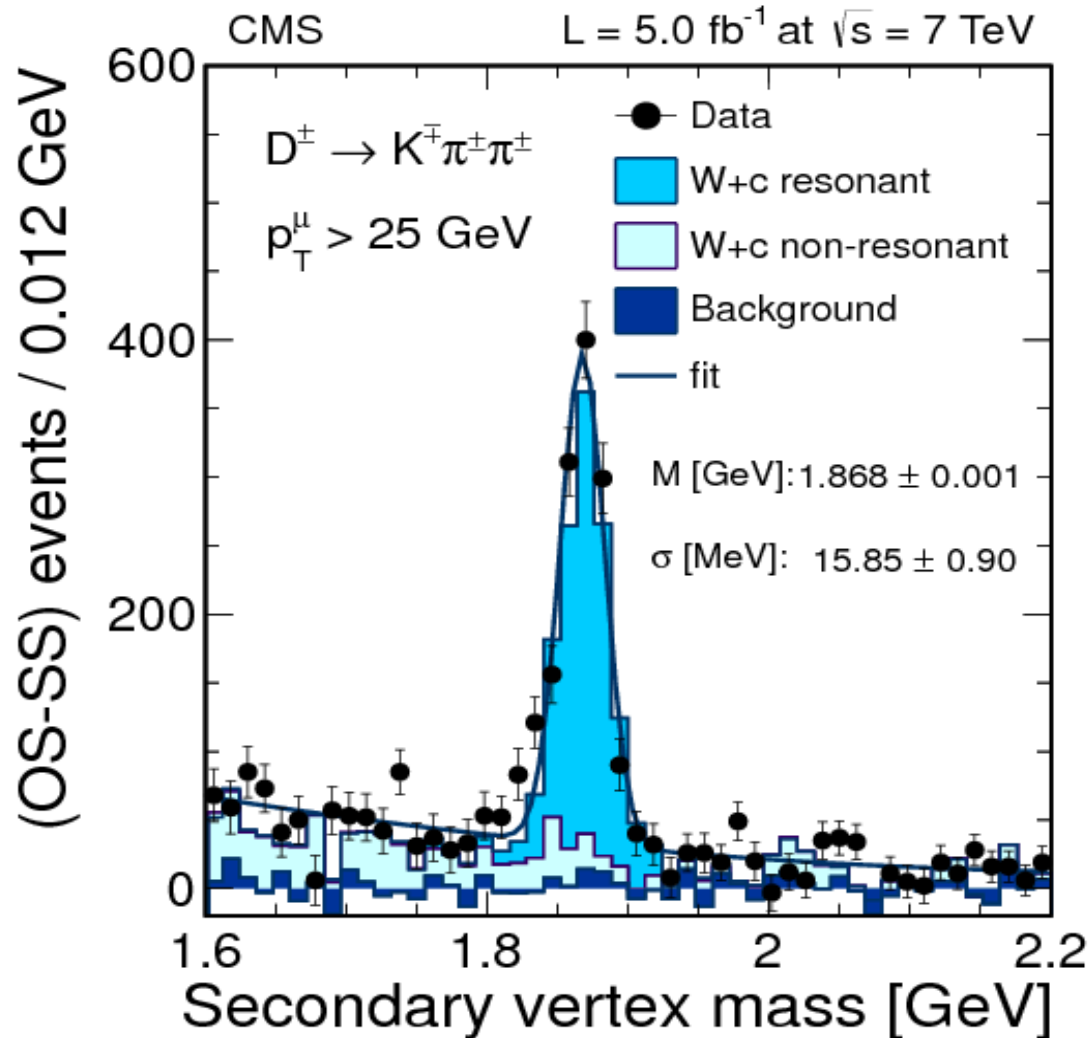
Example: $D^+ \rightarrow K^- \pi^+ \pi^+$ in ATLAS

ATLAS, JHEP 1405, 068 (2014) [arXiv:1402.6263]



Example: $D^+ \rightarrow K^- \pi^+ \pi^+$ in CMS

CMS, JHEP 1402, 013 (2014) [arXiv:1310.1138]



$$\Lambda_c^+ \rightarrow pK^- \pi^+ \text{ vs. } D^+ \rightarrow K^- \pi^+ \pi^+$$

Same signature (3-prong displaced vertex, mass peak), but:

➤ The $\Lambda_c^+ \rightarrow pK^- \pi^+$ signal peak is smaller:

$$\frac{f(c \rightarrow D^+) \mathcal{B}(D^+ \rightarrow K^- \pi^+ \pi^+)}{f(c \rightarrow \Lambda_c^+) \mathcal{B}(\Lambda_c^+ \rightarrow pK^- \pi^+)} \approx 5.3$$

while background is roughly the same.

Ambiguity resolution: in the lab frame, $|\vec{p}(p)| > |\vec{p}(\pi^+)|$.

➤ The Λ_c^+ vertex is less displaced:

$$\tau_{\Lambda_c^+} \approx \frac{\tau_{D^0}}{2} \approx \frac{\tau_{D_s^+}}{2.5} \approx \frac{\tau_{D^+}}{5}$$

e.g., in CMS analysis, < 20% of events had a good secondary vertex
(events contain about 61% D^0 , 24% D^+ , 8% D_s^+ , 6% Λ_c^+)

Improvements for $W+c$ in Run 2

- **Statistics x 60** (cross section x 3, luminosity x 20)

(S/B remains similar because cross sections increase by similar factors.)

- **Upgrades to ATLAS and CMS pixel detectors**

ATLAS: installed IBL

Innermost layer at: **3.3 cm (vs. 5.0 cm in Run 1)**

Smaller pixel size: **50 x 250 (vs. 50 x 400) μm^2**

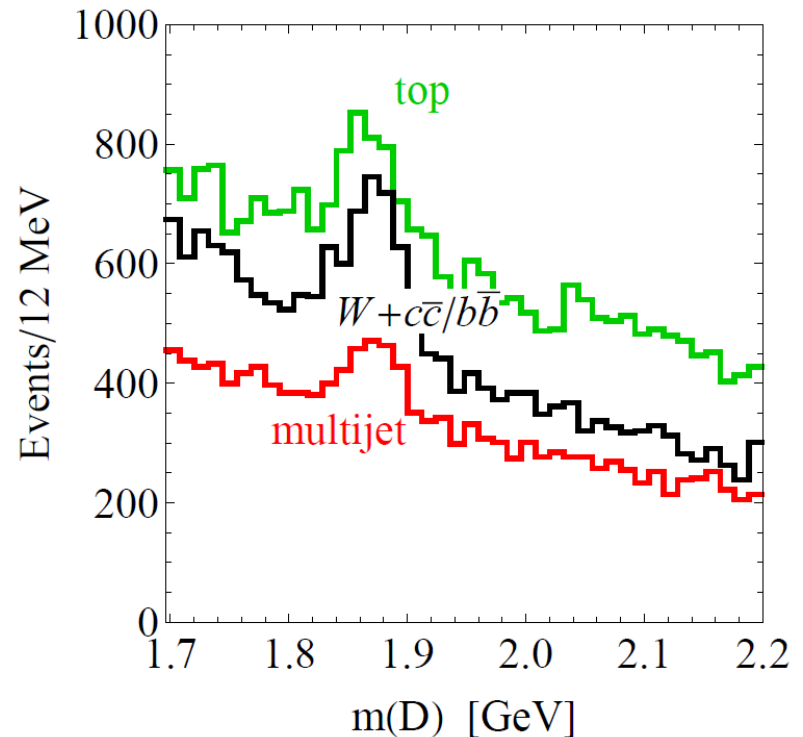
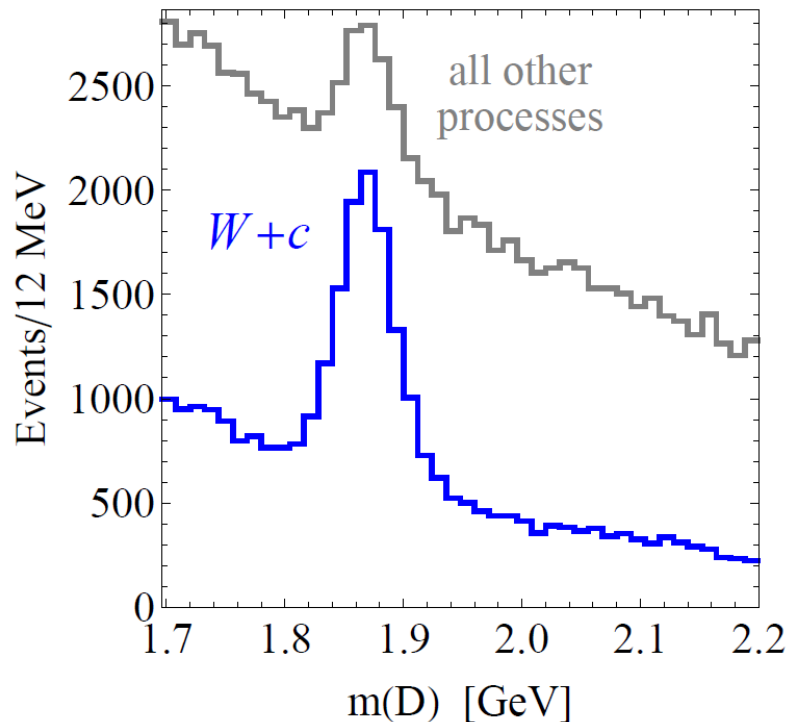
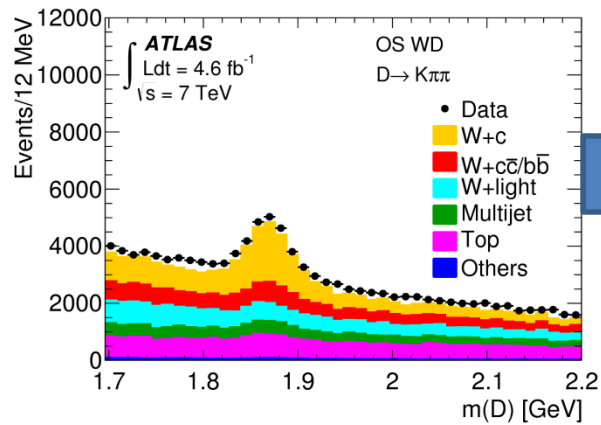
CMS: pixel detector upgrade in winter 2016-2017

Innermost layer at: **3.0 cm (vs. 4.4 cm now)**

Pixel size unchanged: **100 x 150 μm^2**

Backgrounds: ATLAS D^+ example

ATLAS, JHEP 1405, 068 (2014) [arXiv:1402.6263]



Λ_c polarization backgrounds in $W+c$

PEAKING COMPONENTS (REAL Λ_c)

- **c 's in multijet:** unpolarized
- **$W^+ \rightarrow c\bar{s}$ in top:** polarized like the signal
- **b 's in top, $W+b\bar{b}$, multijet:** polarization due to
electroweak $b \rightarrow \Lambda_c, \Sigma_c^{(*)}$
Control region with highly-displaced Λ_c 's ($\tau_b \approx 7\tau_{\Lambda_c}$).
- **$W+c\bar{c}$:** can be estimated from the wrong-sign sample

SMOOTH COMPONENTS (FAKE Λ_c)

At the very least, can be extrapolated from sidebands
(up to a certain systematic uncertainty).

Statistical precision for $W+c$ in Run 2

A variable sensitive to the polarization:

$$A_{\text{FB}} = \frac{N(\cos \theta_{K^-} > 0) - N(\cos \theta_{K^-} < 0)}{N}$$

Statistical uncertainty:

$$\sigma(A_{\text{FB}}) = \sqrt{\frac{1 - A_{\text{FB}}^2}{N}} \approx \frac{1}{\sqrt{N}}$$

The signal contribution:

$$A_{\text{FB},S} = \frac{\alpha_{K^-} \mathcal{P}(\Lambda_c) S}{2N}$$

Significance of observing non-zero $\mathcal{P}(c)$:

$$\frac{|A_{\text{FB},S}|}{\sigma(A_{\text{FB}})} = \frac{|\alpha_{K^-} \mathcal{P}(\Lambda_c)|}{2} \frac{S}{\sqrt{N}}$$

Statistical precision for $W+c$ in Run 2

A ballpark figure

- Start with the ATLAS D^+ peak.
- Account for the difference between the $D^+ \rightarrow K^- \pi^+ \pi^+$ and $\Lambda_c^+ \rightarrow p K^- \pi^+$ rates.
- Assume Run 2 statistics (100 fb^{-1})

Without the displacement issue, $S/\sqrt{N} \approx 47$.

For, e.g., $|\alpha_{K^-} \mathcal{P}(\Lambda_c)| = 0.4$, this gives 11% precision.

Suppose that relaxed displacement requirements increase N by a factor of 2 while still losing 1/2 of S .
Still, 3σ significance for observing non-zero $\mathcal{P}(c)$.

u, d polarizations

Cannot use decays of protons or neutrons, but can again consider the Λ ($\approx sud$).

Naïve quark model: all the Λ spin is on the s ☹️

Nucleon DIS + flavor SU(3): u and d carry about -20% each 😊

Burkardt and Jaffe, PRL 70, 2537 (1993) [hep-ph/9302232]

Jaffe, PRD 54, 6581 (1996) [hep-ph/9605456]

Further inputs possible in the future from:

- Polarized DIS and polarized pp collisions

e.g., COMPASS, EPJC 64, 171 (2009)

Deng (STAR), Phys.Part.Nucl. 45, 73 (2014)

- Lattice QCD

QCDSF, PLB 545, 112 (2002) [hep-lat/0208017]

CSSM and QCDSF/UKQCD, PRD 90, 014510 (2014) [arXiv:1405.3019]

Chambers et al., PRD 92, 114517 (2015) [arXiv:1508.06856]

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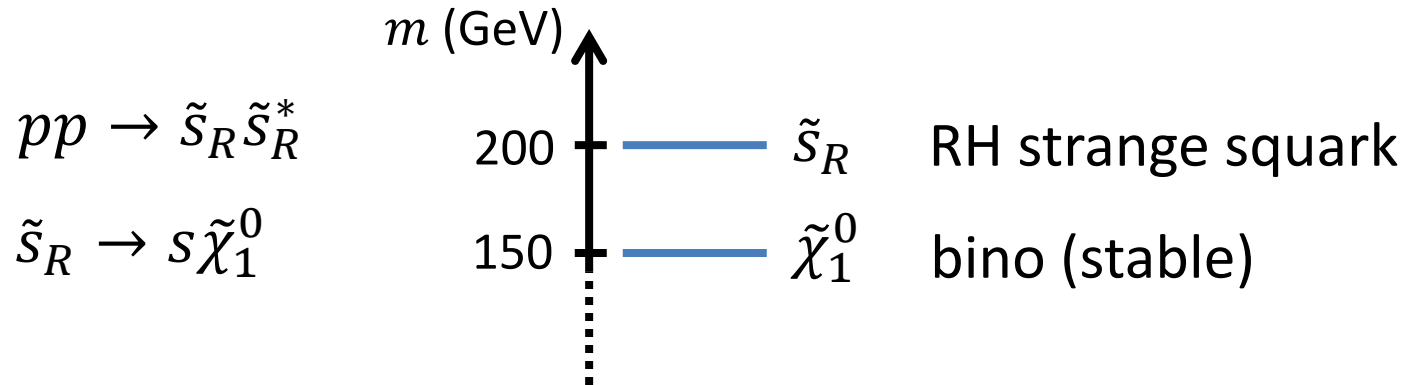
Jaffe, PRD 54, 6581 (1996) [hep-ph/9605456]

Studies of u, d jets in $t\bar{t}$ samples will require **much more statistics** than s , also because:

- No u or d tagging; c -tag veto only partially effective
(Can define separate u and d samples, contaminated by c and s respectively, using W_{leptonic} charge.)
- Fragmentation fractions of $u, d \rightarrow \Lambda$ smaller than $s \rightarrow \Lambda$

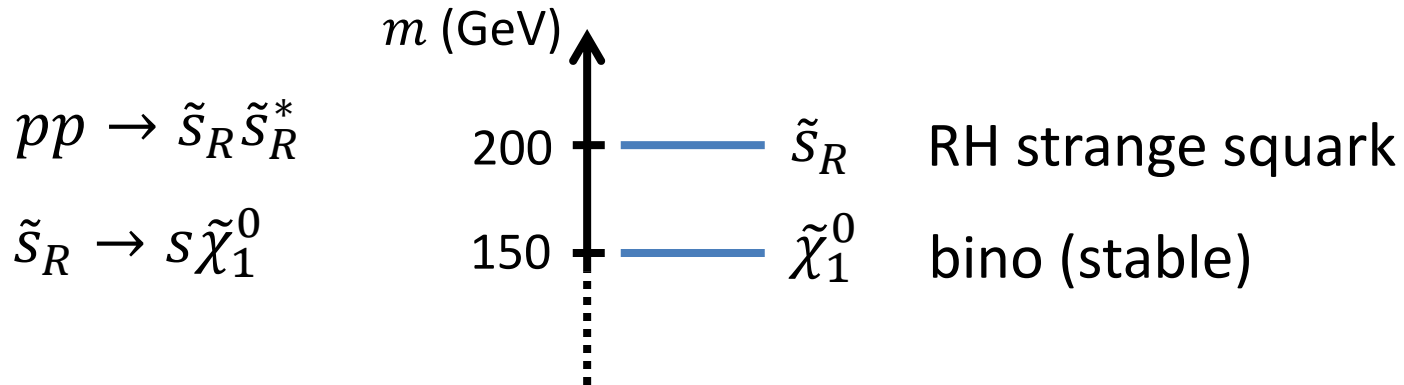
New physics example

Suppose a jets + MET excess is being attributed to:

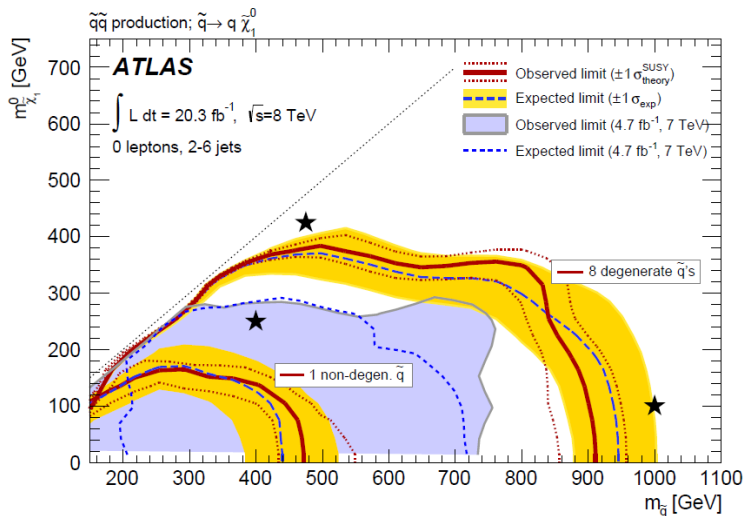


New physics example

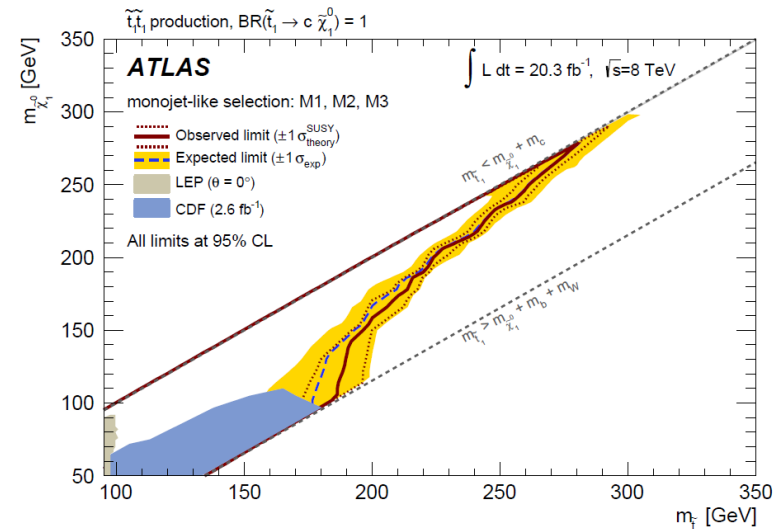
Suppose a jets + MET excess is being attributed to:



This scenario was barely beyond the reach of Run 1.



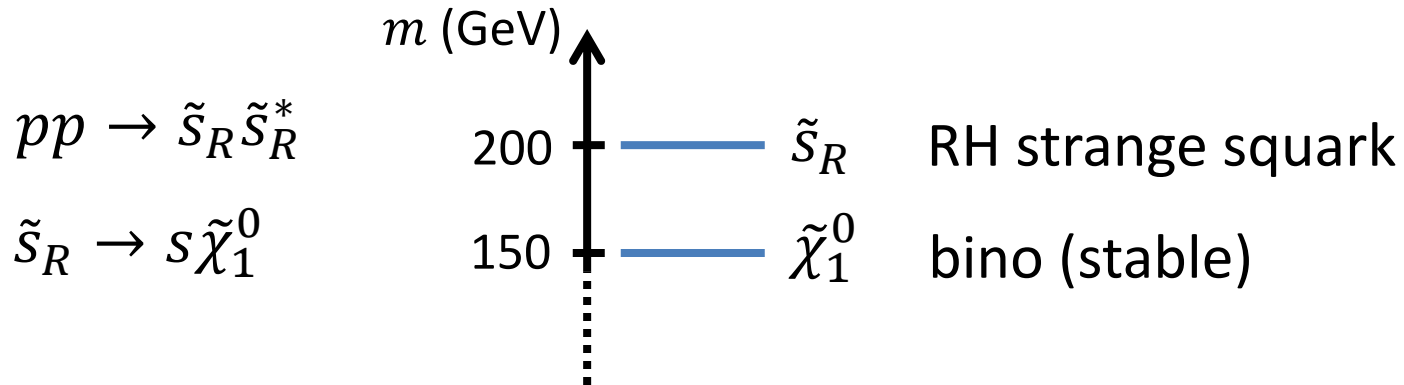
JHEP 09, 176 (2014) [arXiv:1405.7875]



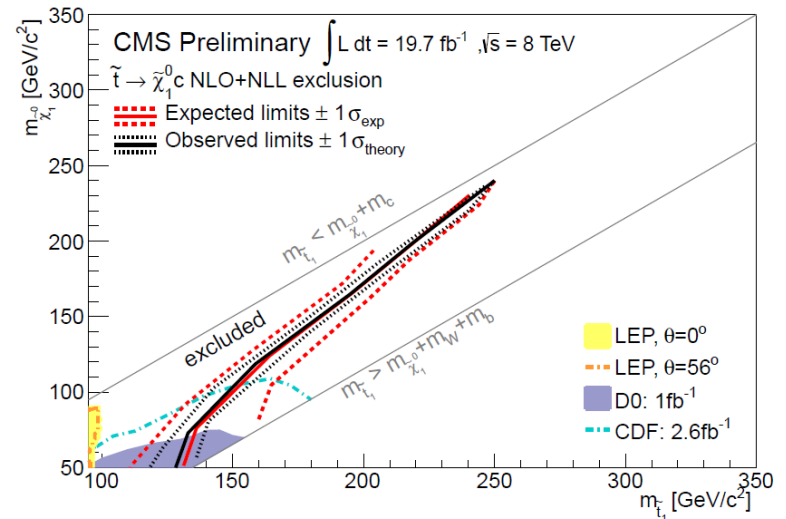
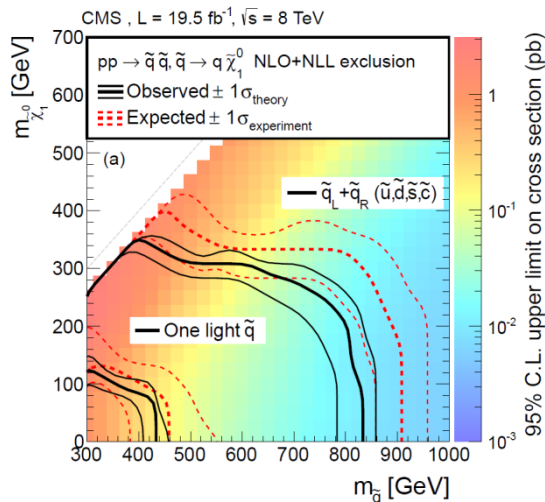
PRD 90, 052008 (2014) [arXiv:1407.0608]

New physics example

Suppose a jets + MET excess is being attributed to:



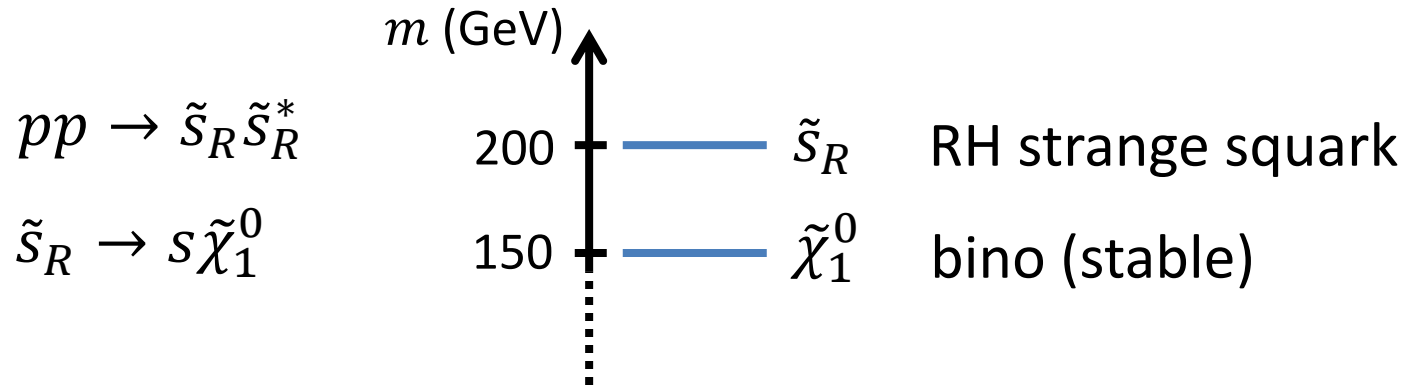
This scenario was barely beyond the reach of Run 1.



*The masses of interest are unfortunately not shown.

New physics example

Suppose a jets + MET excess is being attributed to:



Test this interpretation by measuring the s -quark polarization.

Rough estimate (see paper for details):

for 3 ab^{-1} of 14 TeV data: statistical precision of better than **30%**
(even without optimization of selection cuts, without accounting for the expected detector upgrades, and without combining ATLAS and CMS)