ECT\*, Trento, 09/14/2017

LFC17

# LATTICE FIELD THEORY RESULTS ON NEW STRONG DYNAMICS

Enrico Rinaldi





## NEW STRONG DYNAMICS

### Composite Higgs

### Composite Dark Matter

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New SU(N<sub>c</sub>) gauge sector with N<sub>f</sub> fermions in the N<sub>r</sub> representation of the gauge group

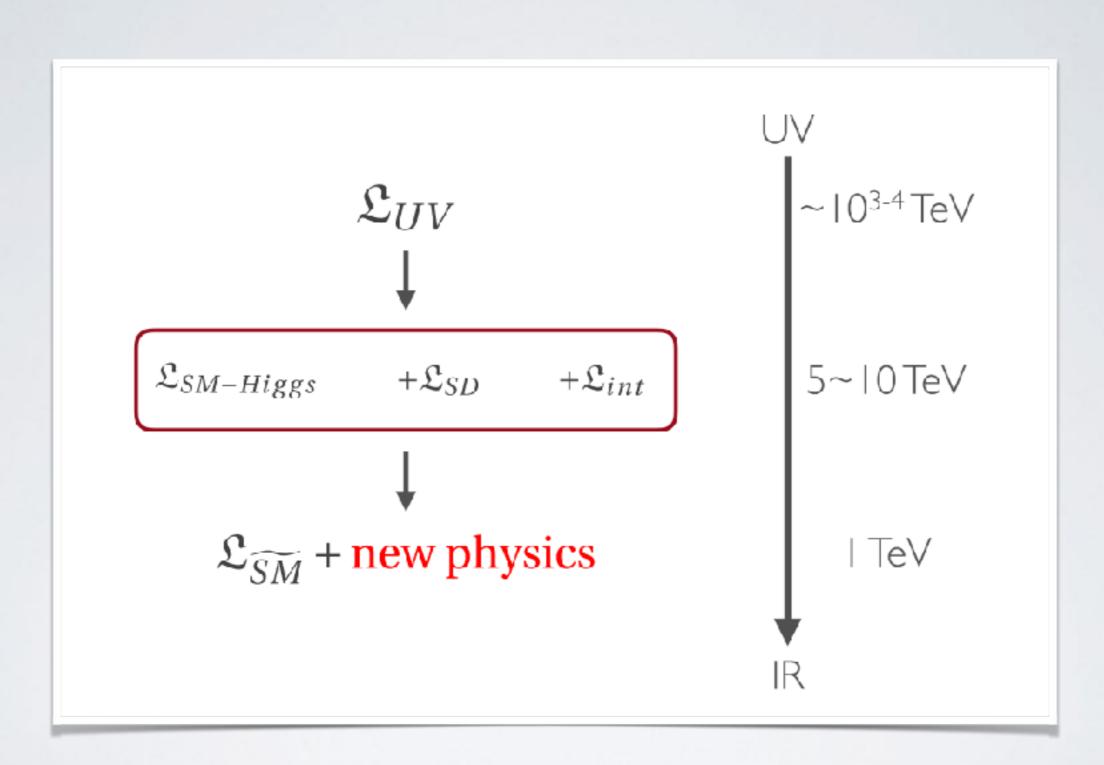
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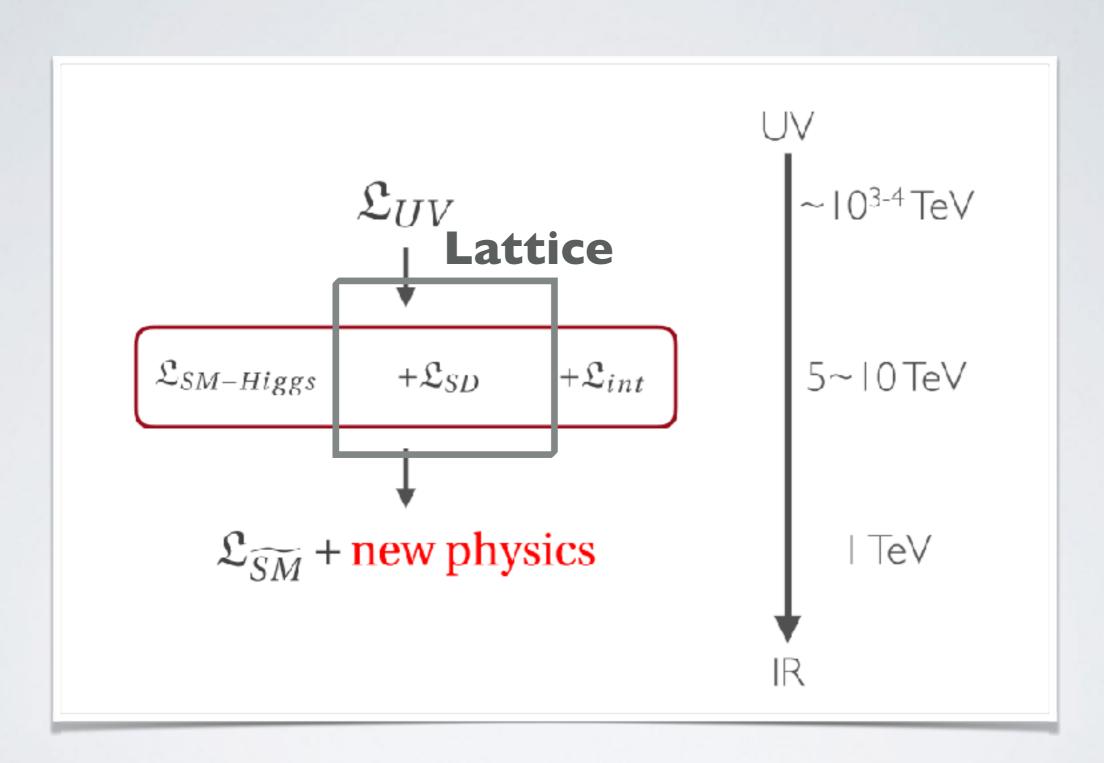
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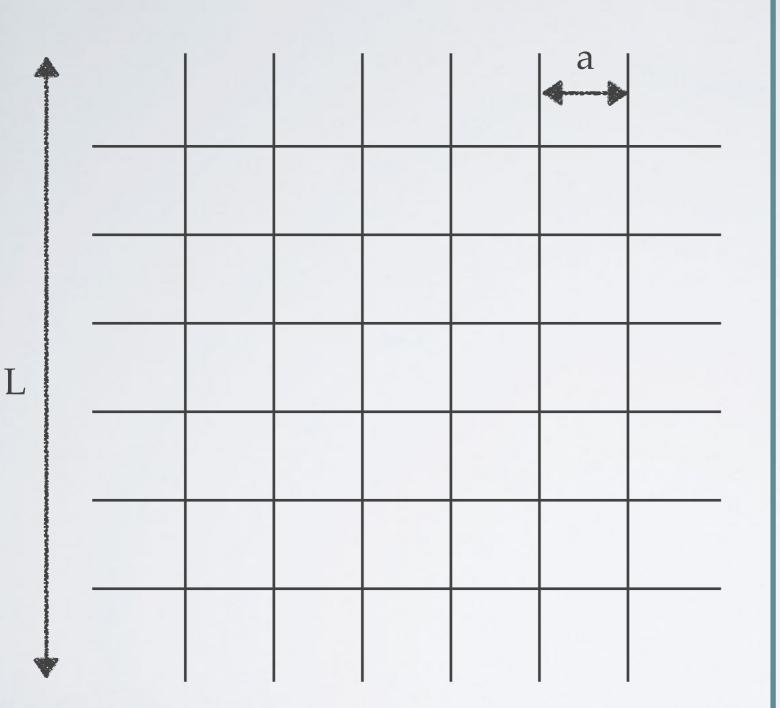
Most of the theory work is done using EFTs and there are only a handful of UV complete models



from Pica 1701.07782



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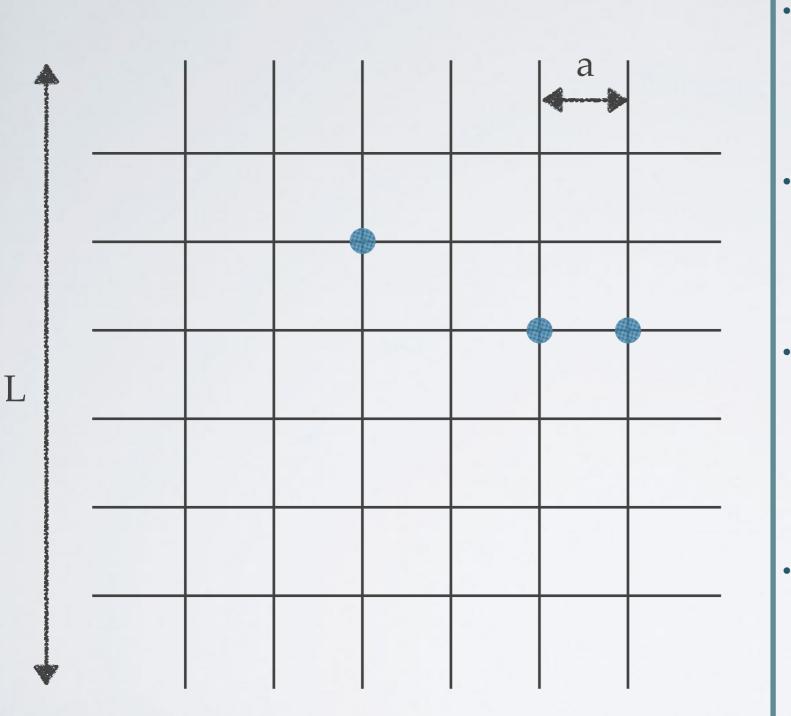


#### Discretize space and time

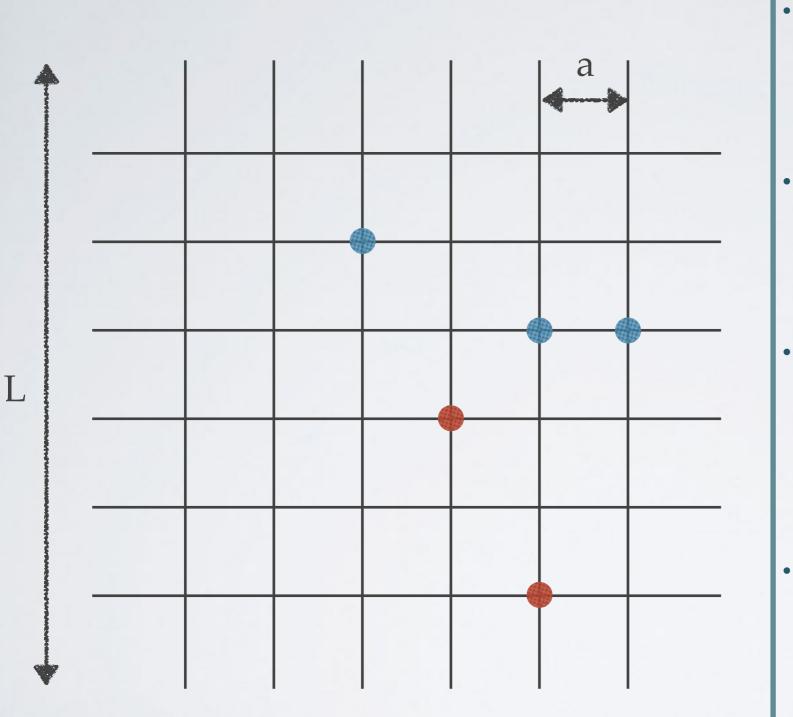
- lattice spacing "a"
- lattice size "L"

•

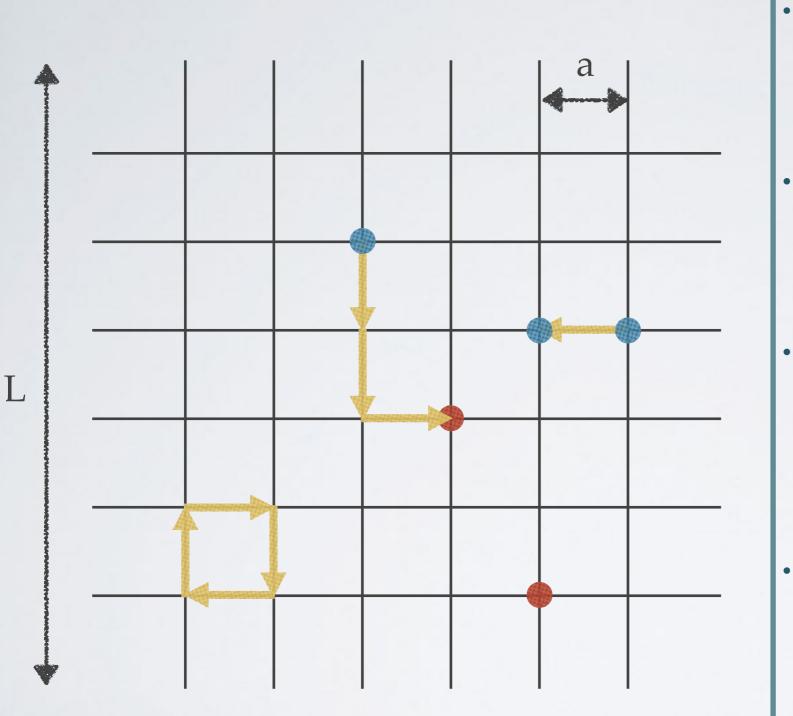
- Keep all d.o.f. of the theory
  - not a model!
  - no simplifications
- Amenable to numerical methods
  - Monte Carlo sampling
  - use supercomputers
- Precisely quantifiable and improvable errors
  - Systematic
  - Statistical



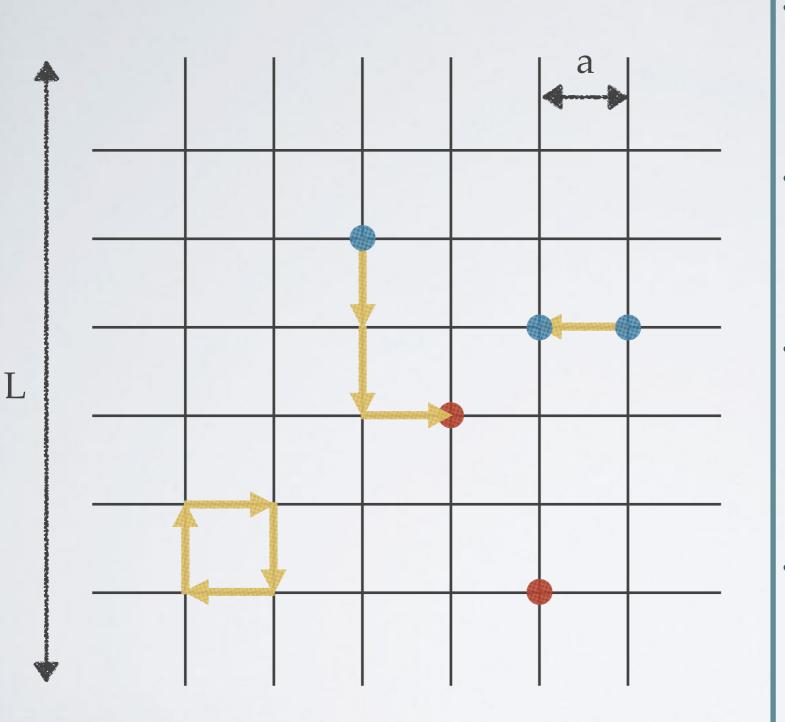
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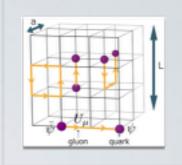


N<sub>c</sub> N<sub>f</sub> N<sub>r</sub> parameters that can be easily changed

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[KEK-Japan]

### Importance of lattice simulations



★Lattice simulations are needed to numerically <u>solve strong</u> <u>dynamics</u>

★ Controllable systematic errors and room for improvement

★Naive dimensional analysis and EFT approaches can miss important <u>non-perturbative</u> contributions

★EFTs inspired by QCD <u>might not work</u> when the dynamics is different

 $\star$ Lattice studies can reliably point out similarities or differences as the parameter space (N<sub>c</sub>,N<sub>f</sub>,N<sub>r</sub>) is scanned

# MOTIVATIONS

- Strongly interacting quantum field theory with different  $N_{\rm c},N_{\rm f}$  and  $N_{\rm r}$ 
  - is the dynamics different from QCD?
  - what is the hierarchy in the spectrum?
  - is there a light scalar singlet?
- Phenomenology of physics beyond the Standard Model
  - light Higgs from composite dynamics (pNGB or dilatonic nature)
  - large anomalous dimensions
  - expected (near-)conformal dynamics for consistency with experiments

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in QCD there is a broad resonance  $f_0(500)$ 

- This talk: SU(3) gauge theory with 8 degenerate fundamental fermions
  - focus on hierarchy of masses towards the chiral limit
  - focus on flavor-singlet states
    - scalar 0<sup>++</sup> (also called "the sigma")
    - pseudoscalar 0<sup>-+</sup> (also called "the eta prime")
  - focus on vector resonance (also called the  $\rho$  meson)

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Lattice Strong Dynamics collaboration



Argonne: Jin, Osborn

Bern: Schaich

Boston: Brower, Rebbi, Weinberg

Colorado: Hasenfratz, Neil

Edinburgh: Witzel

LLNL:Vranas

UC Davis: Kiskis

Yale: Appelquist, Fleming, Gasbarro KEK: Aoki, Kurachi, Shibata

Kyoto: Aoyama

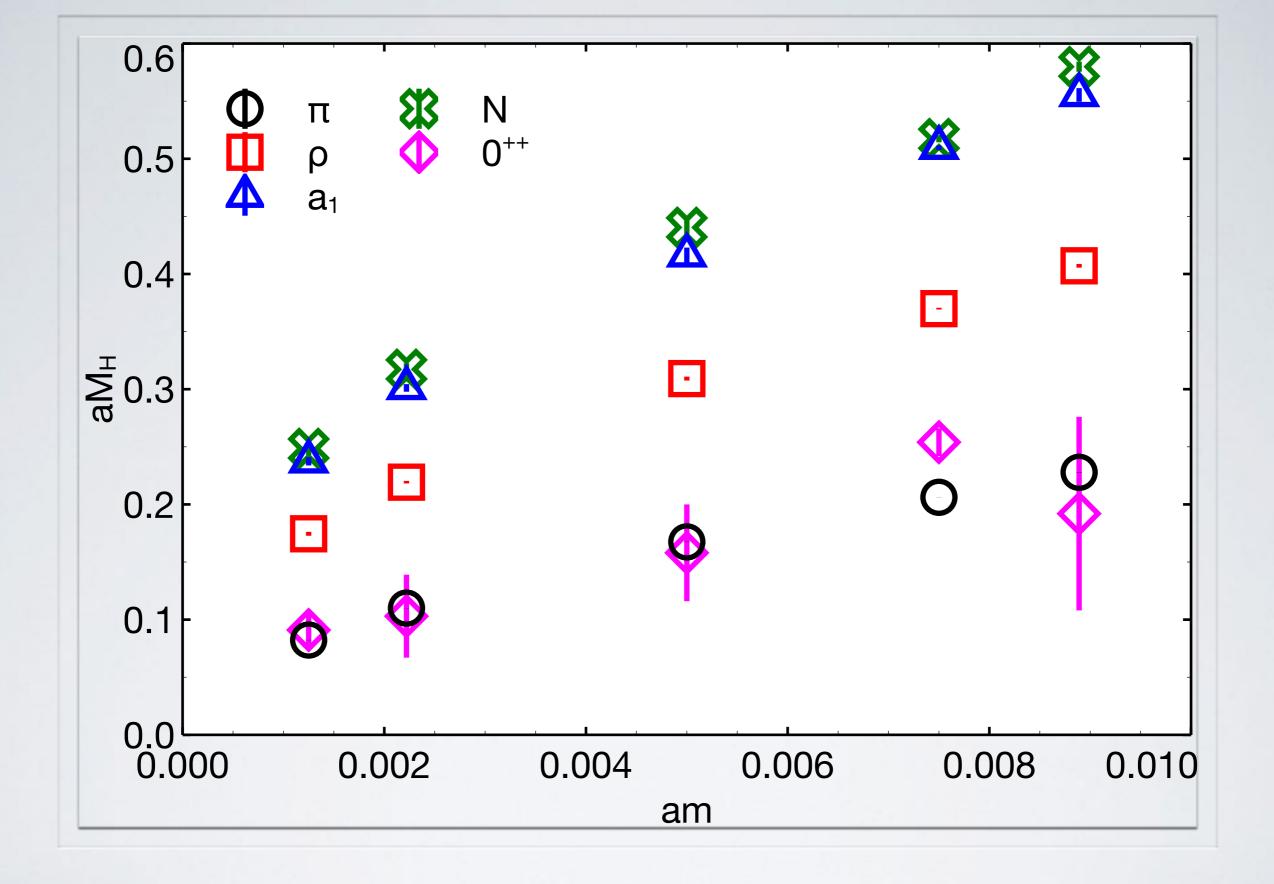
Nagoya: Maskawa, Yamawaki, Nagai

Nara: Ohki

Marseille: Miura

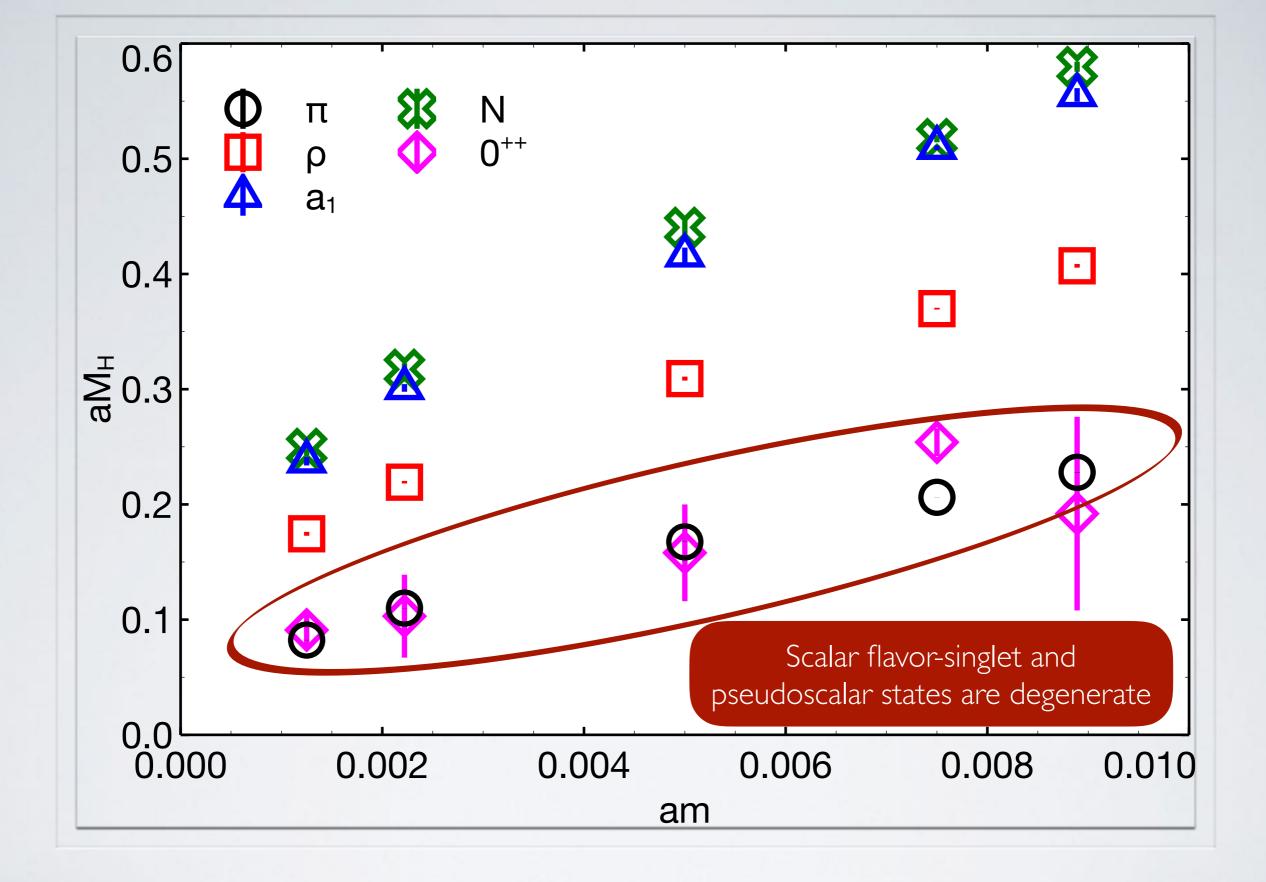
Swansea: Bennett

Tsukuba: Yamazaki



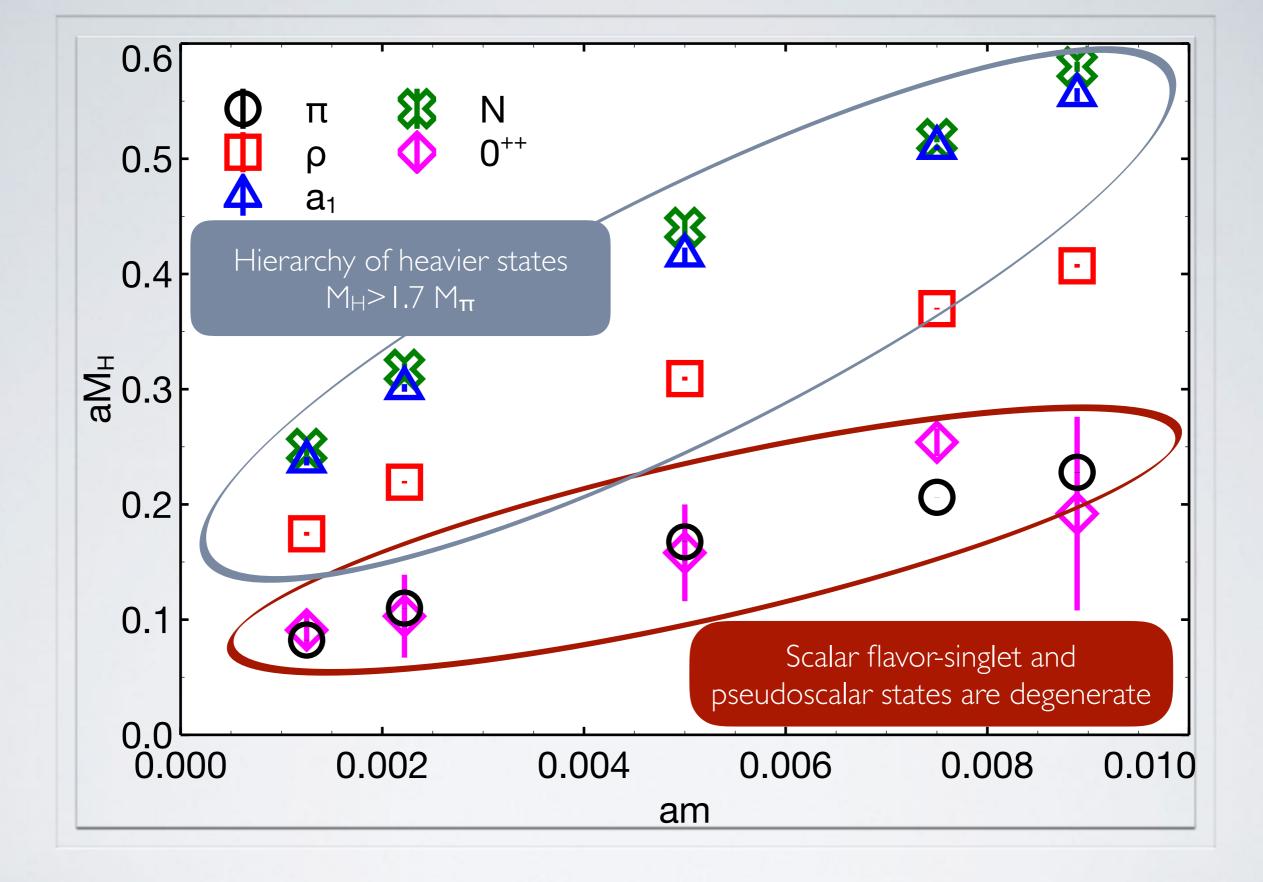
Lattice details: F+A Wilson plaquette action + nHYP smeared naive staggered quarks

LSD arxiv: 1601.04027 [scalar update, preliminary]



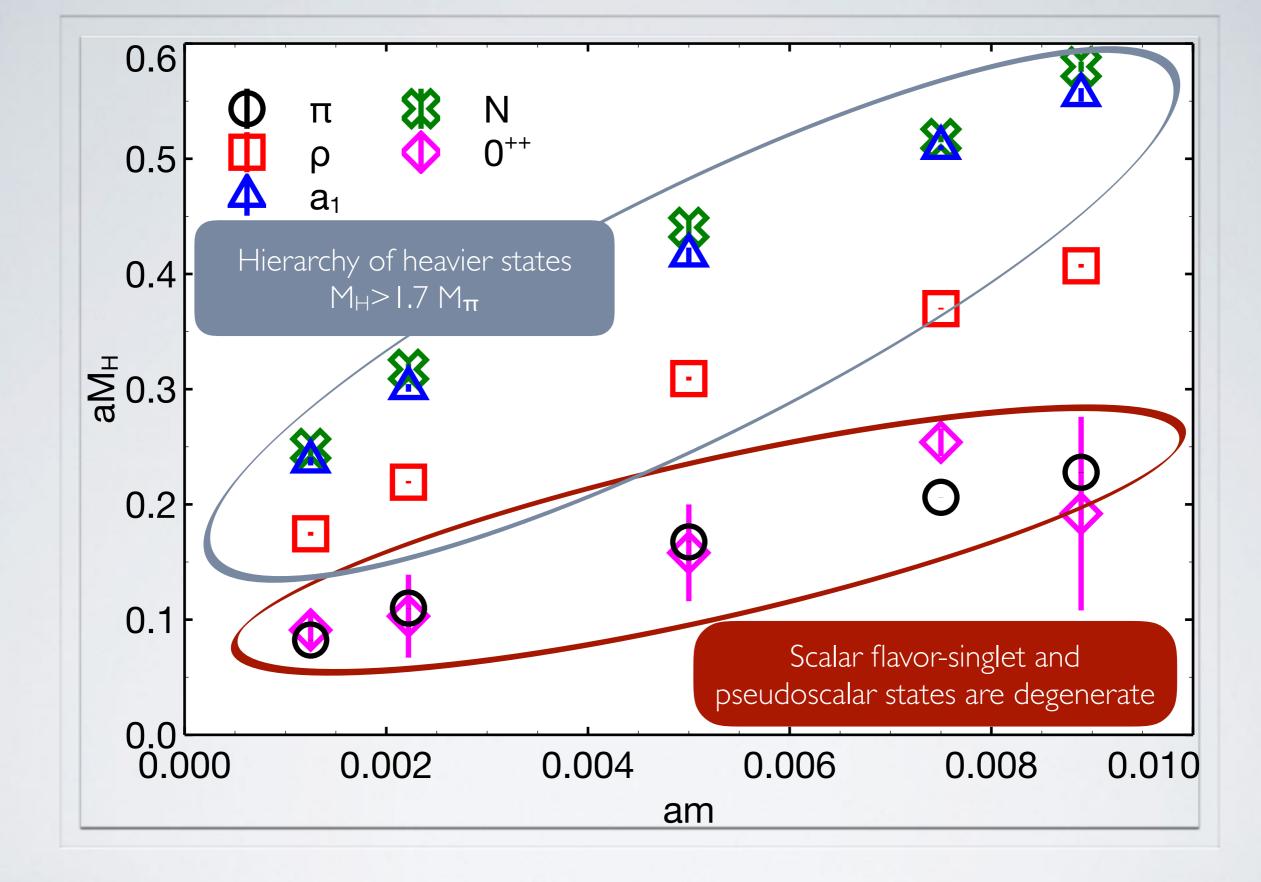
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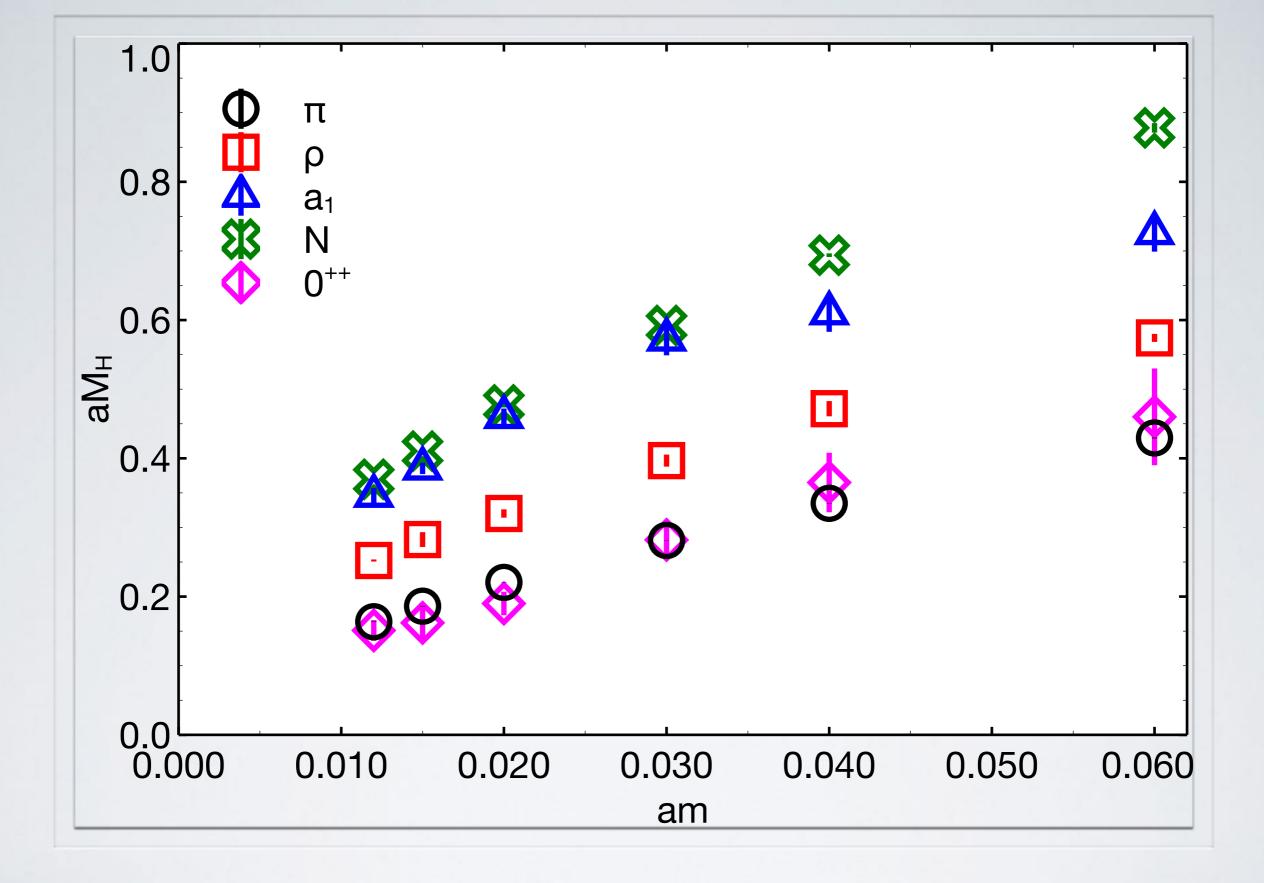


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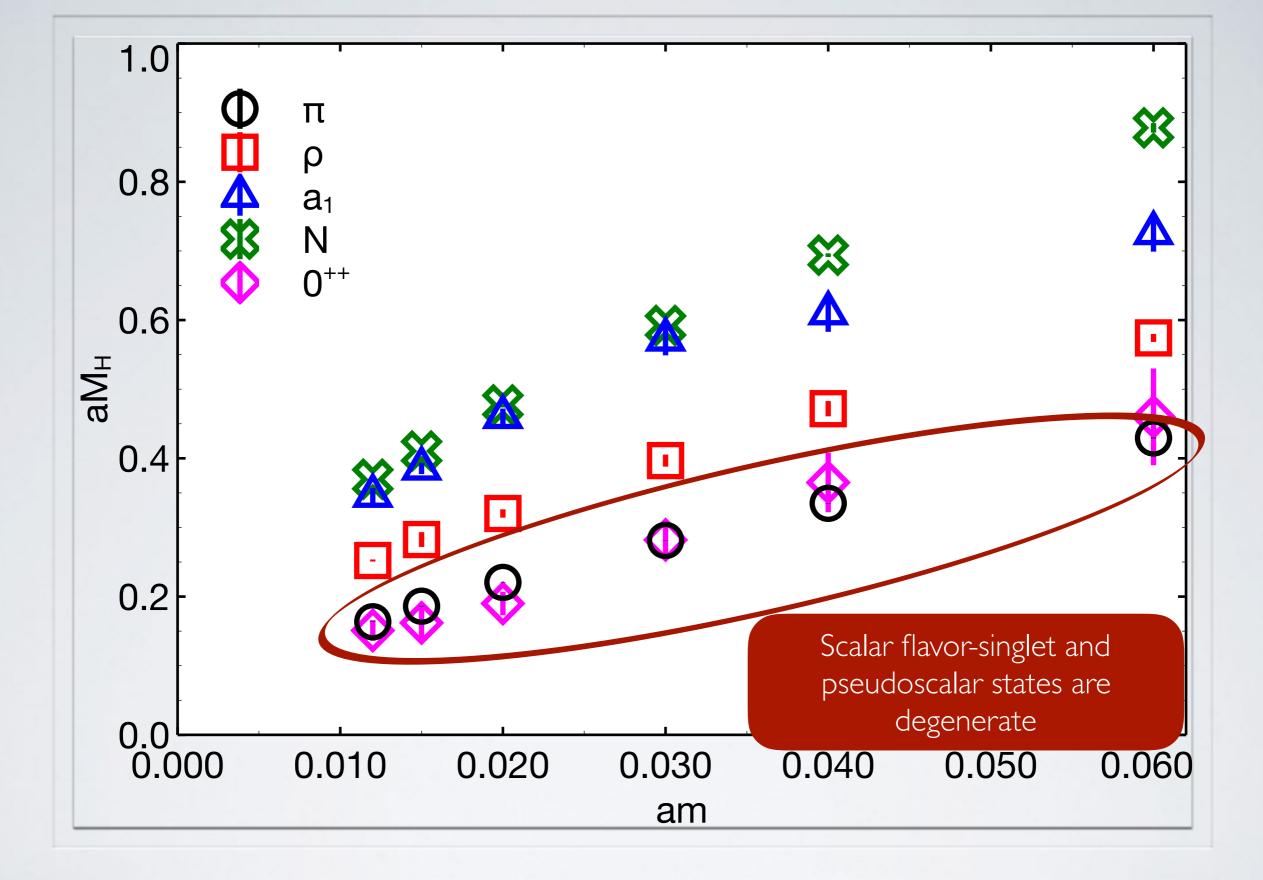
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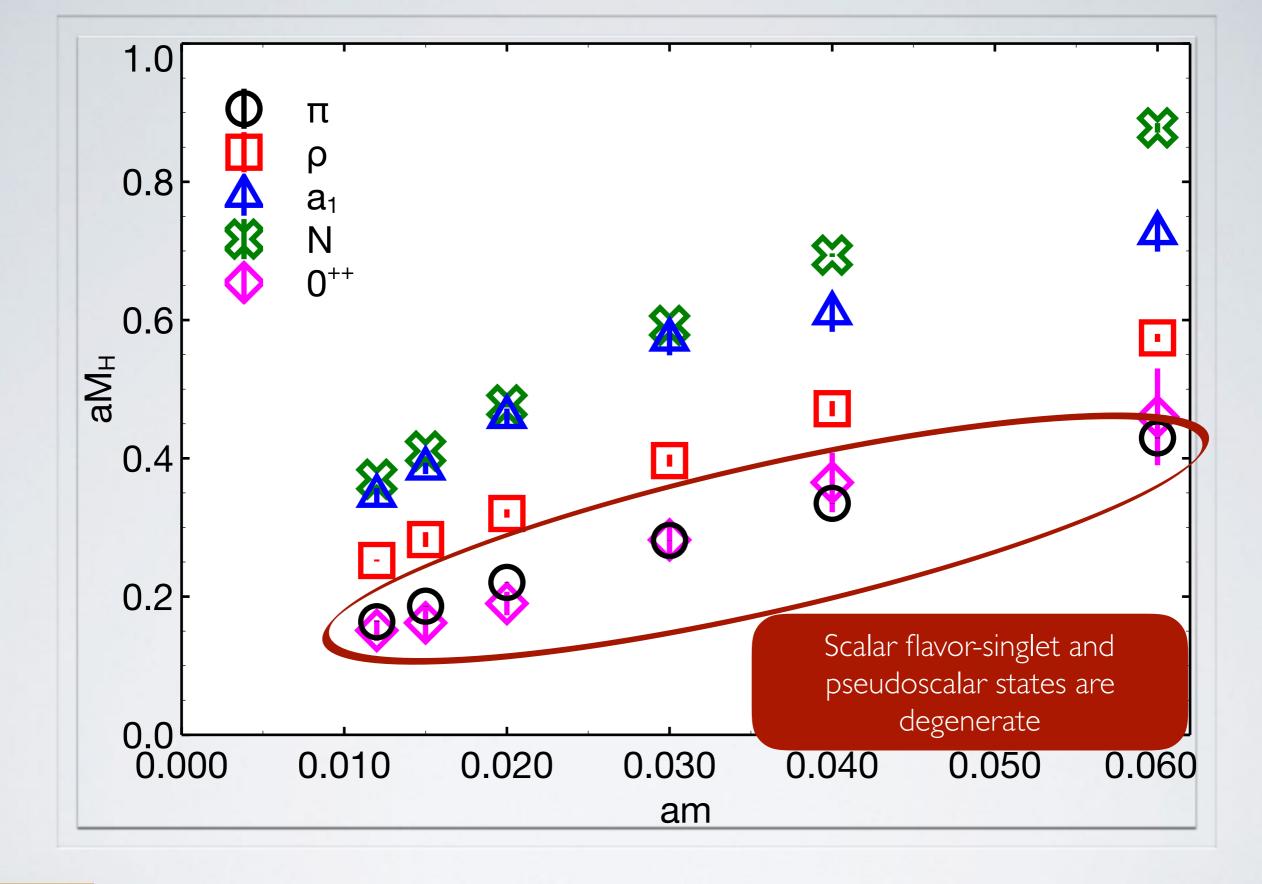
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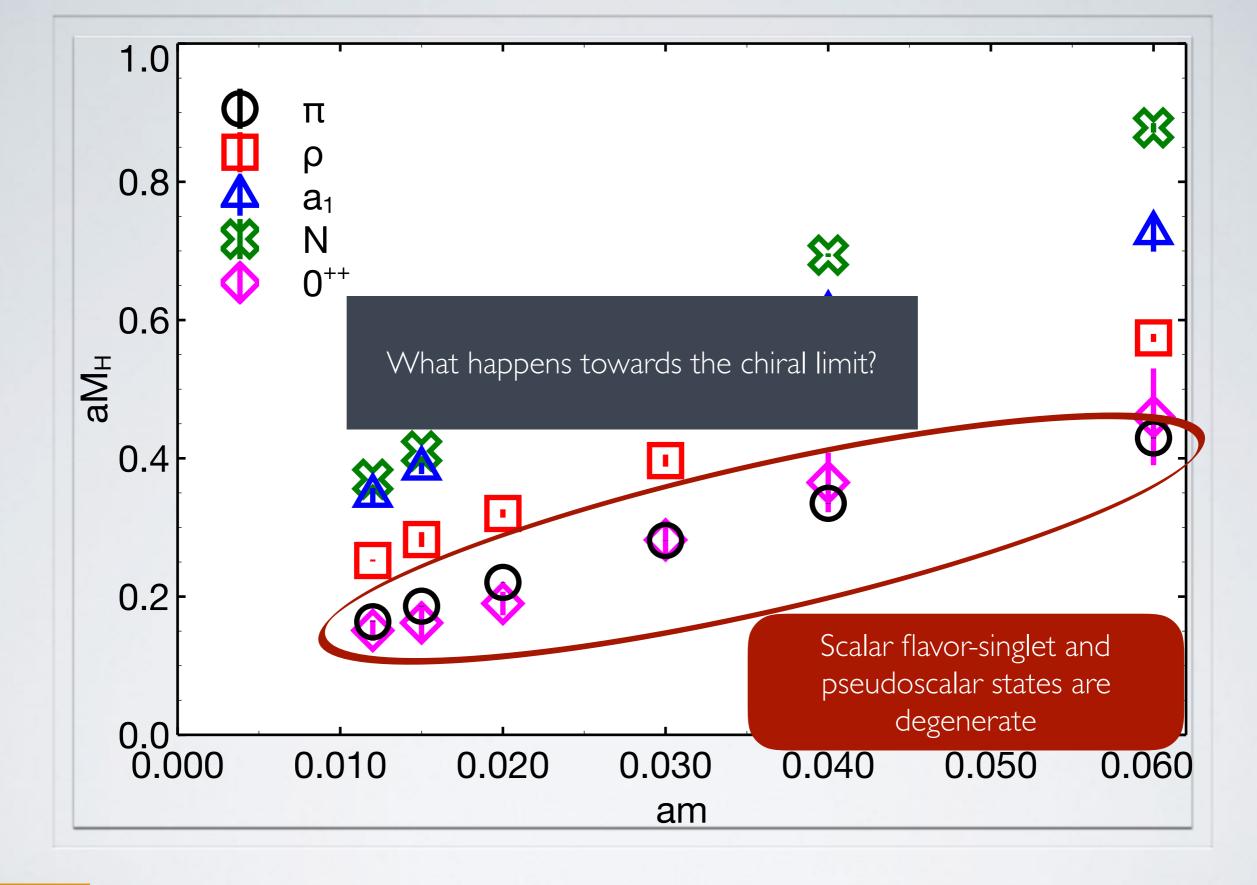


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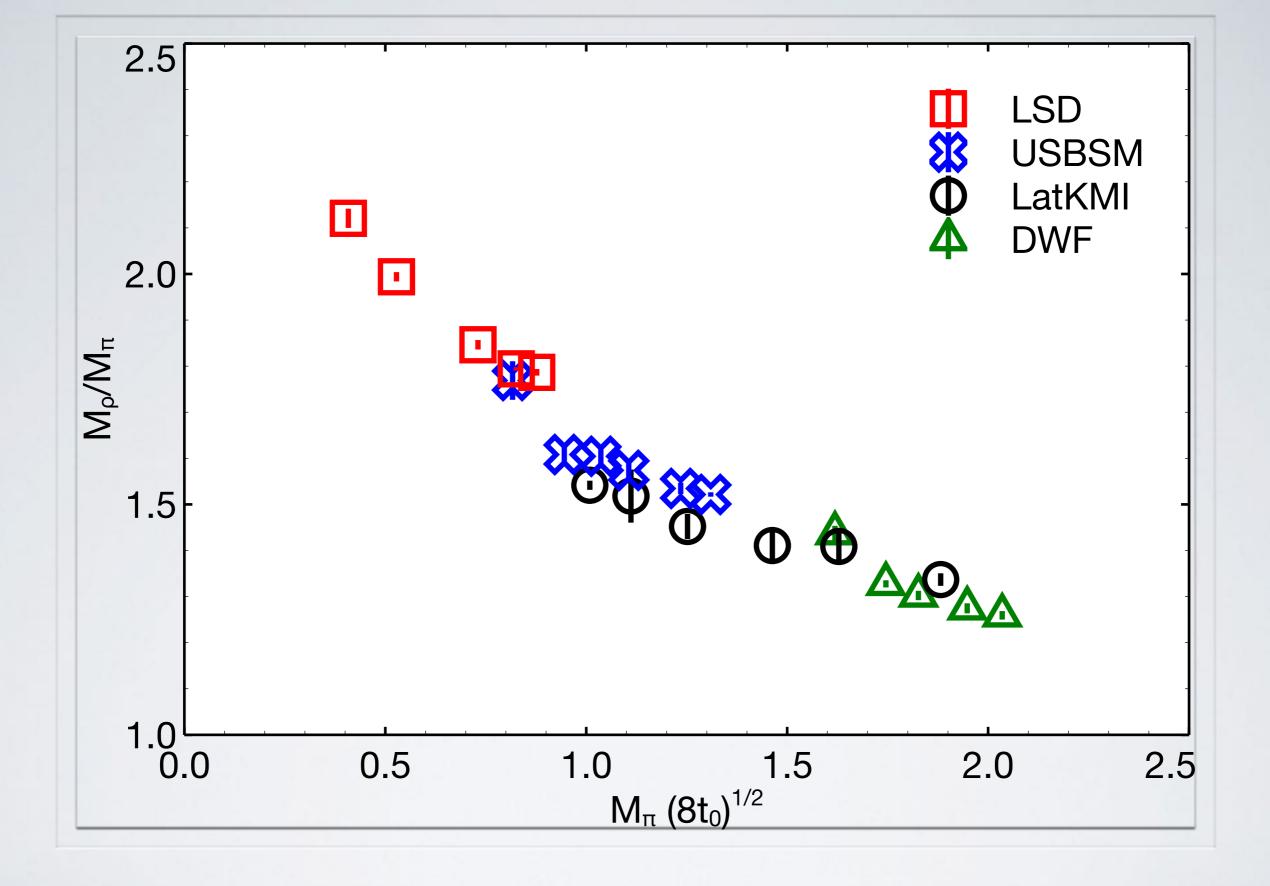
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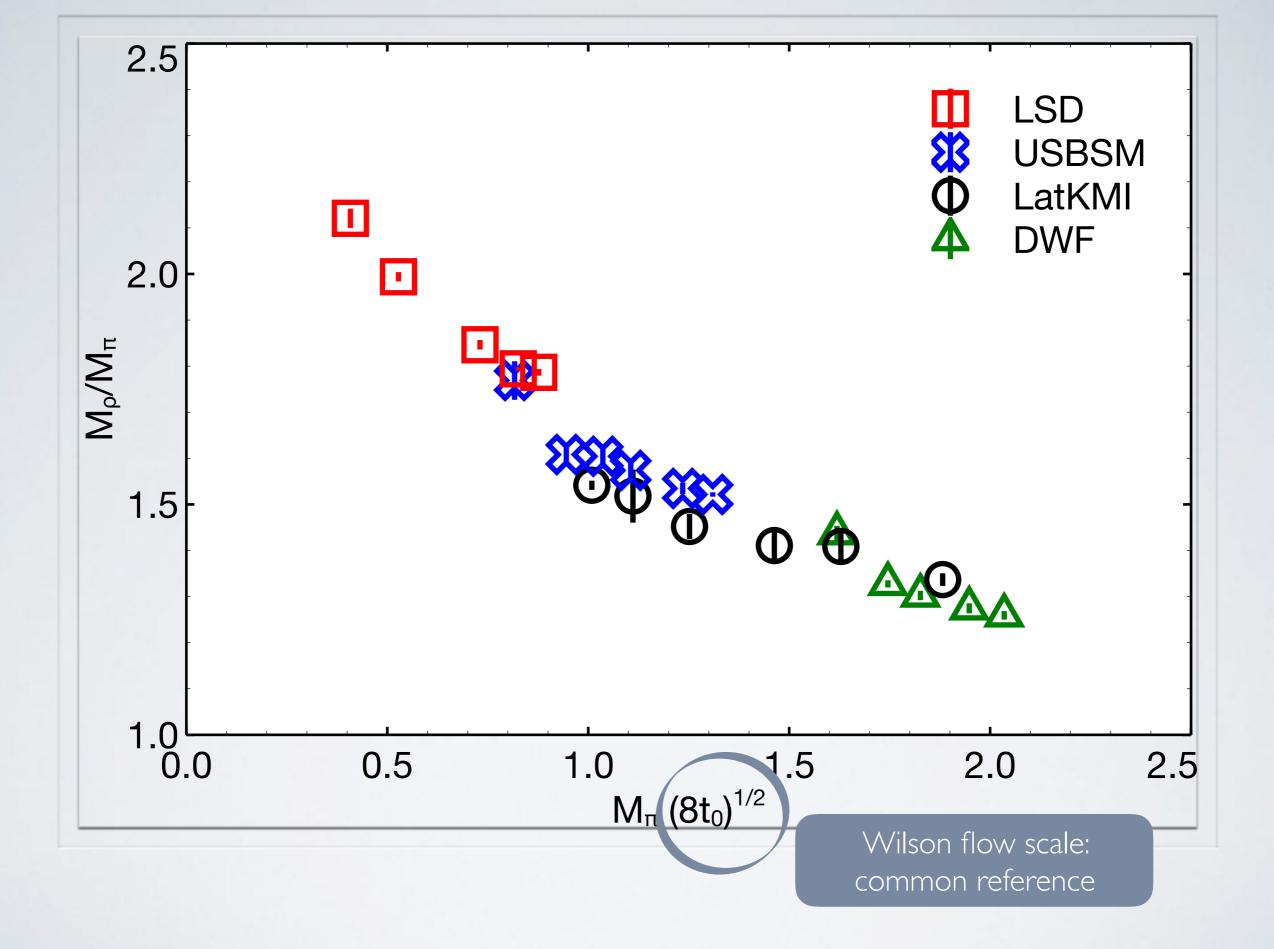
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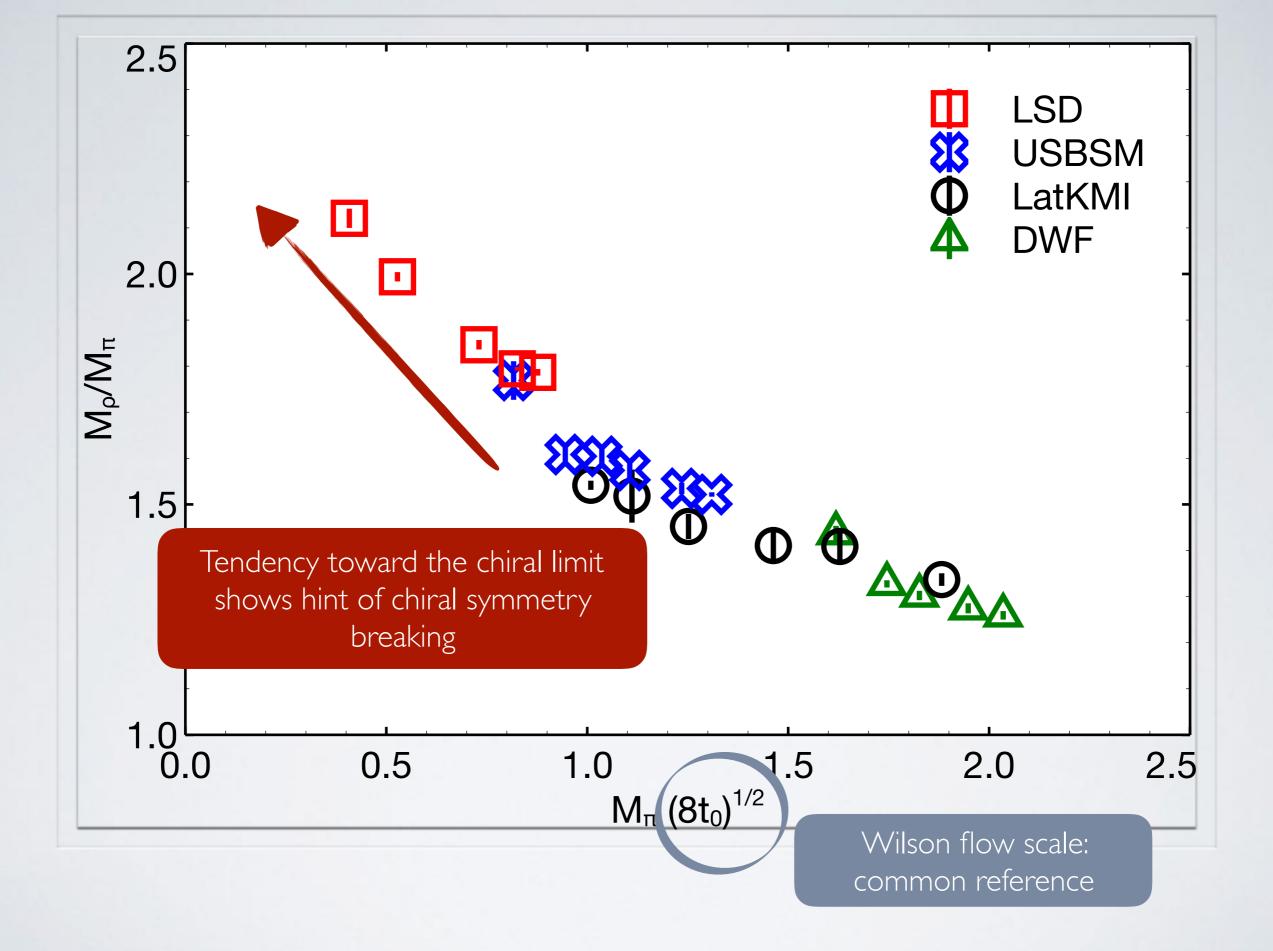


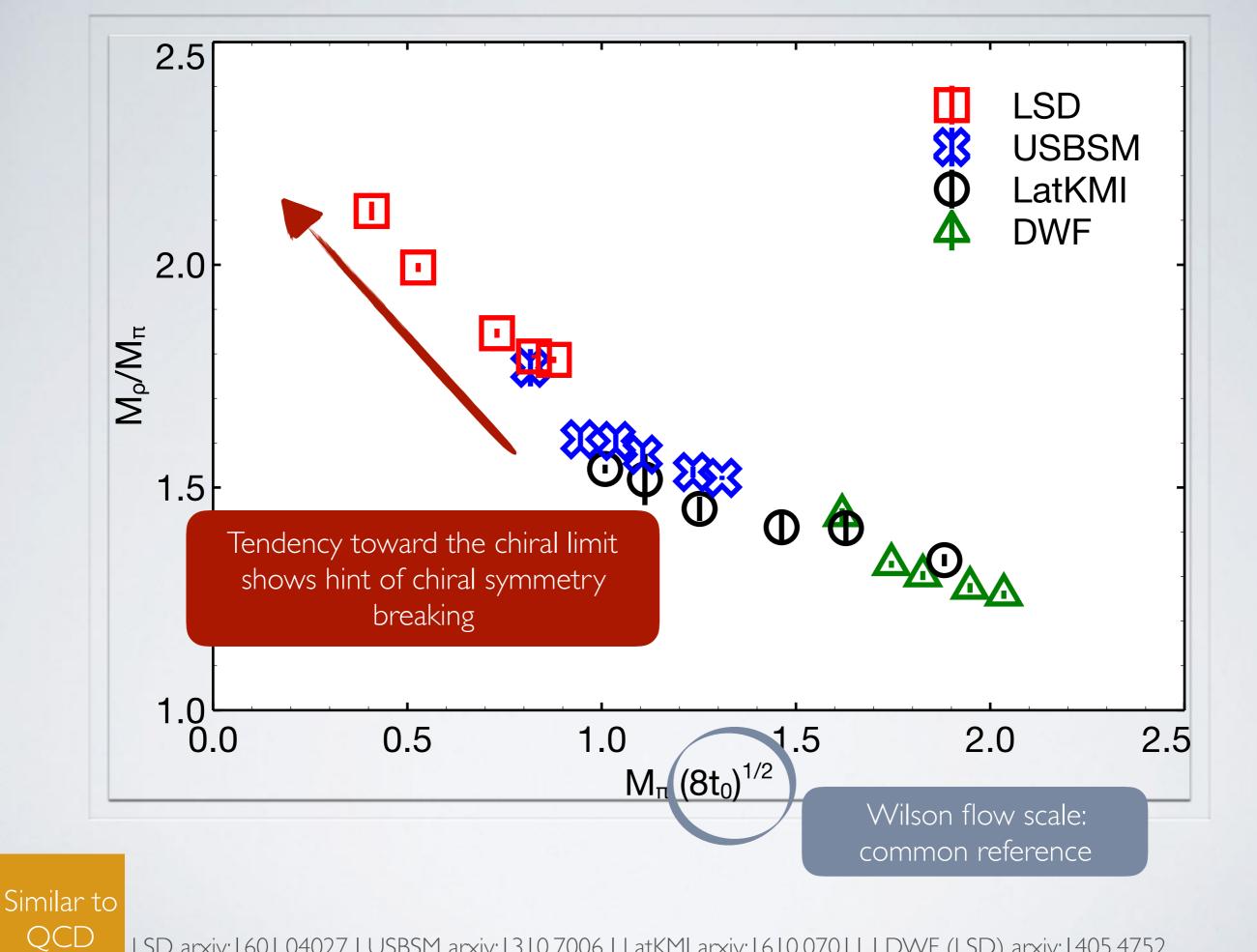
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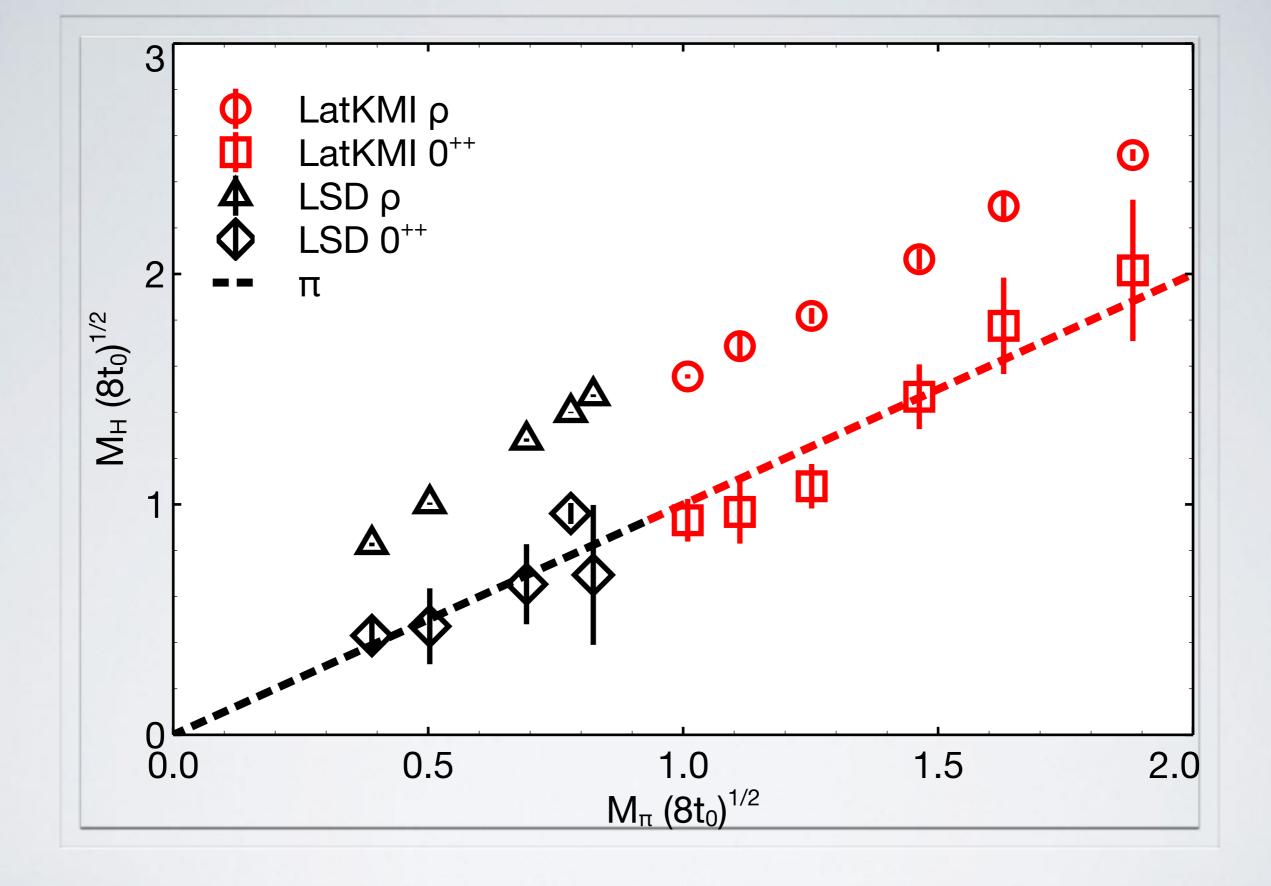
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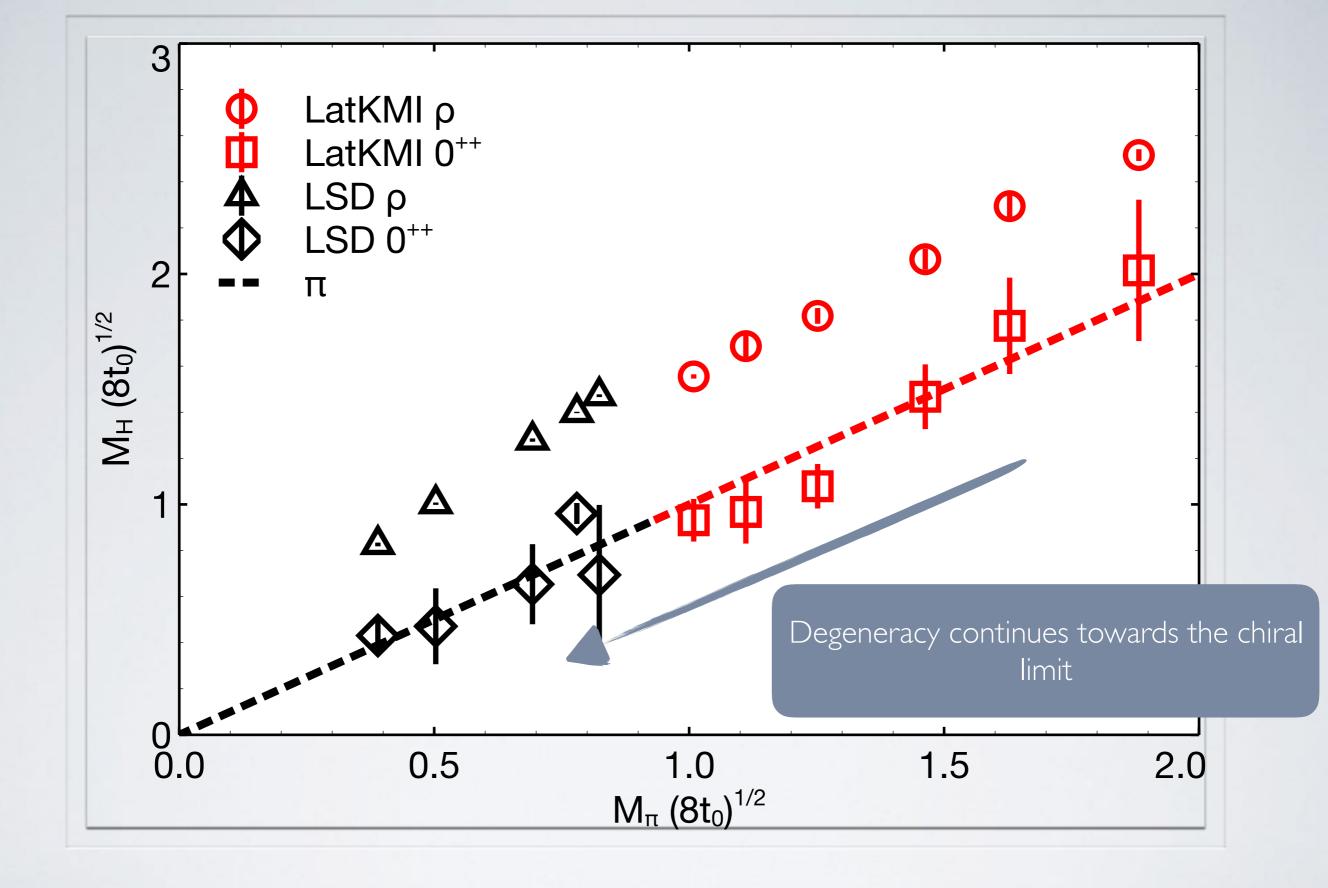




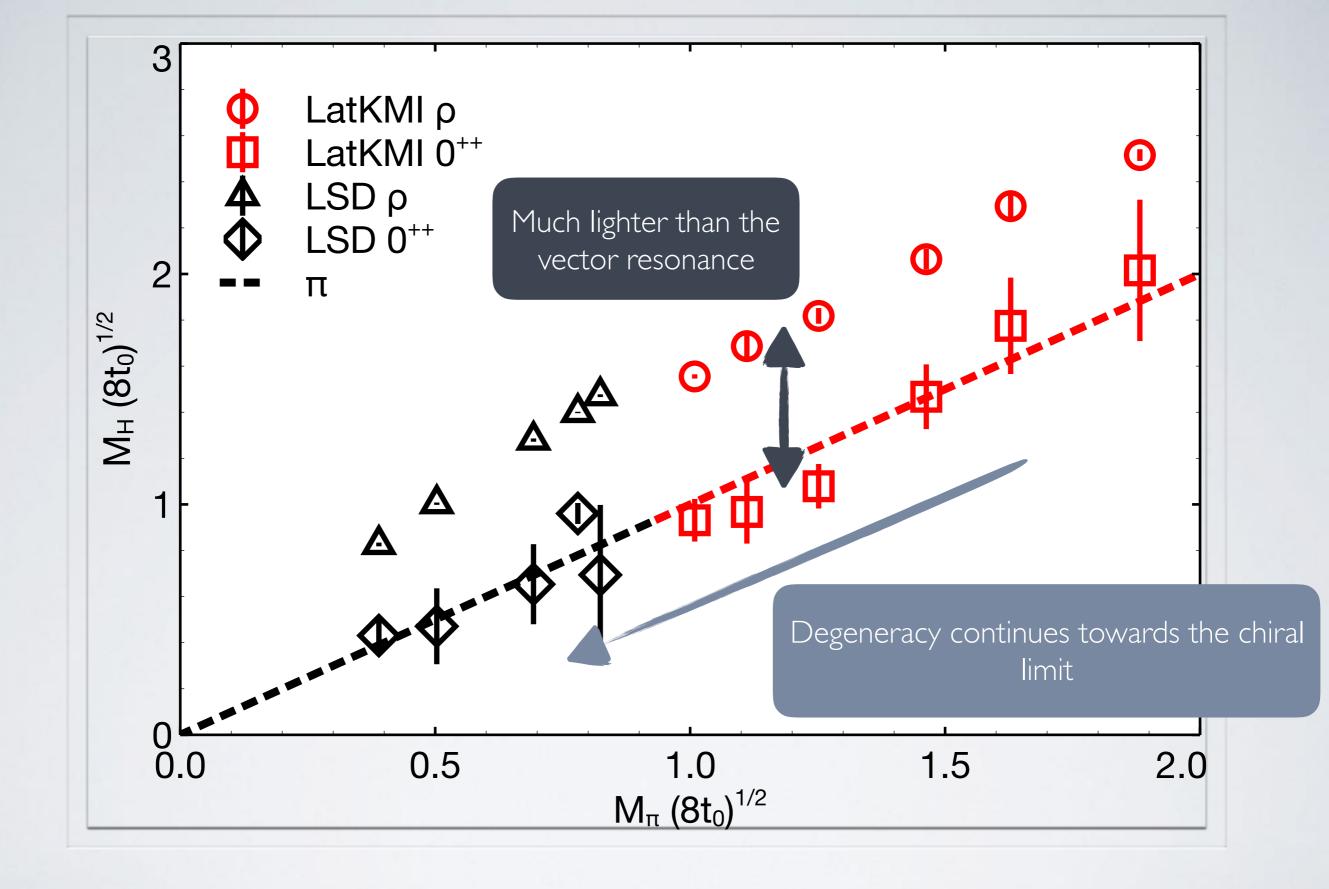




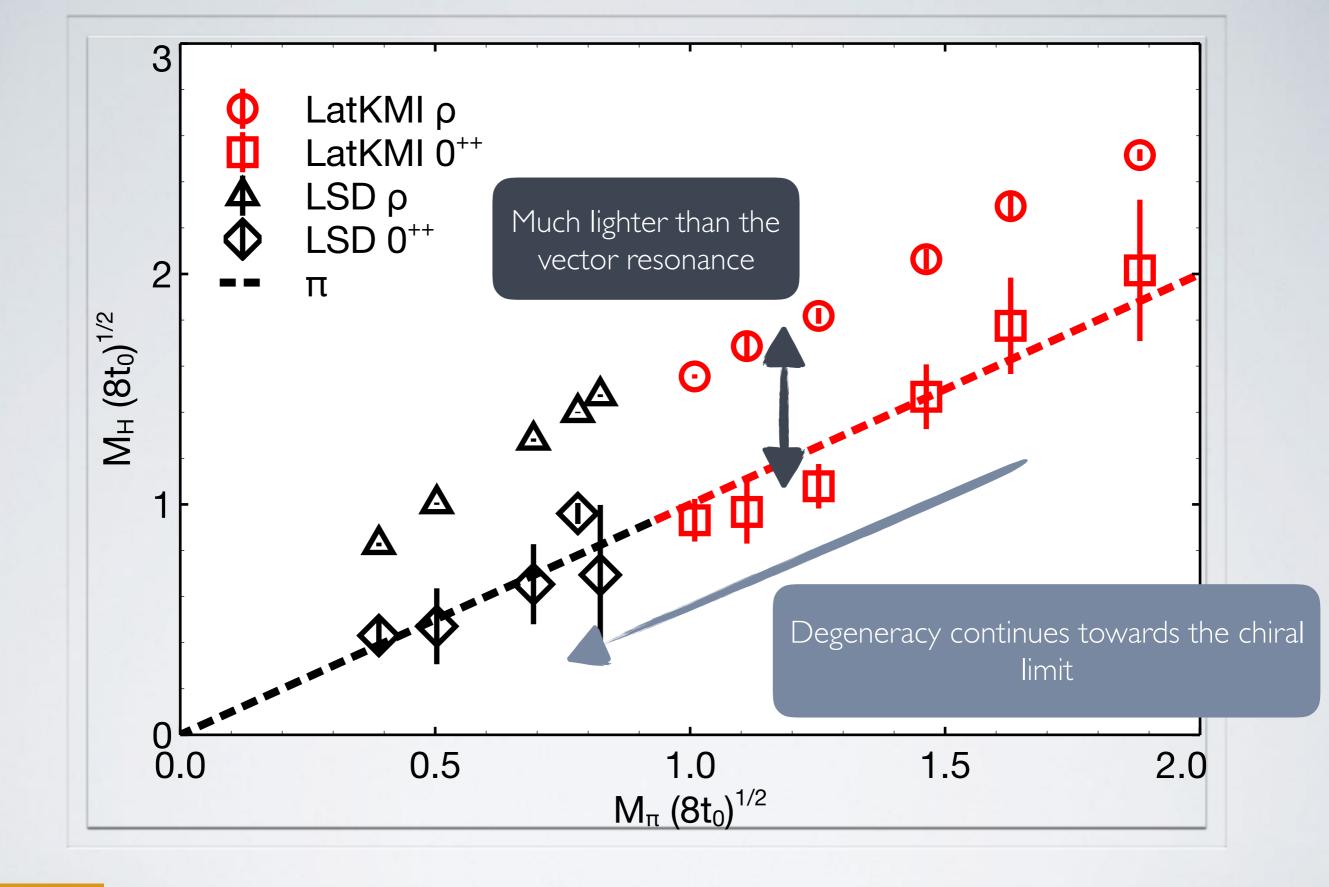
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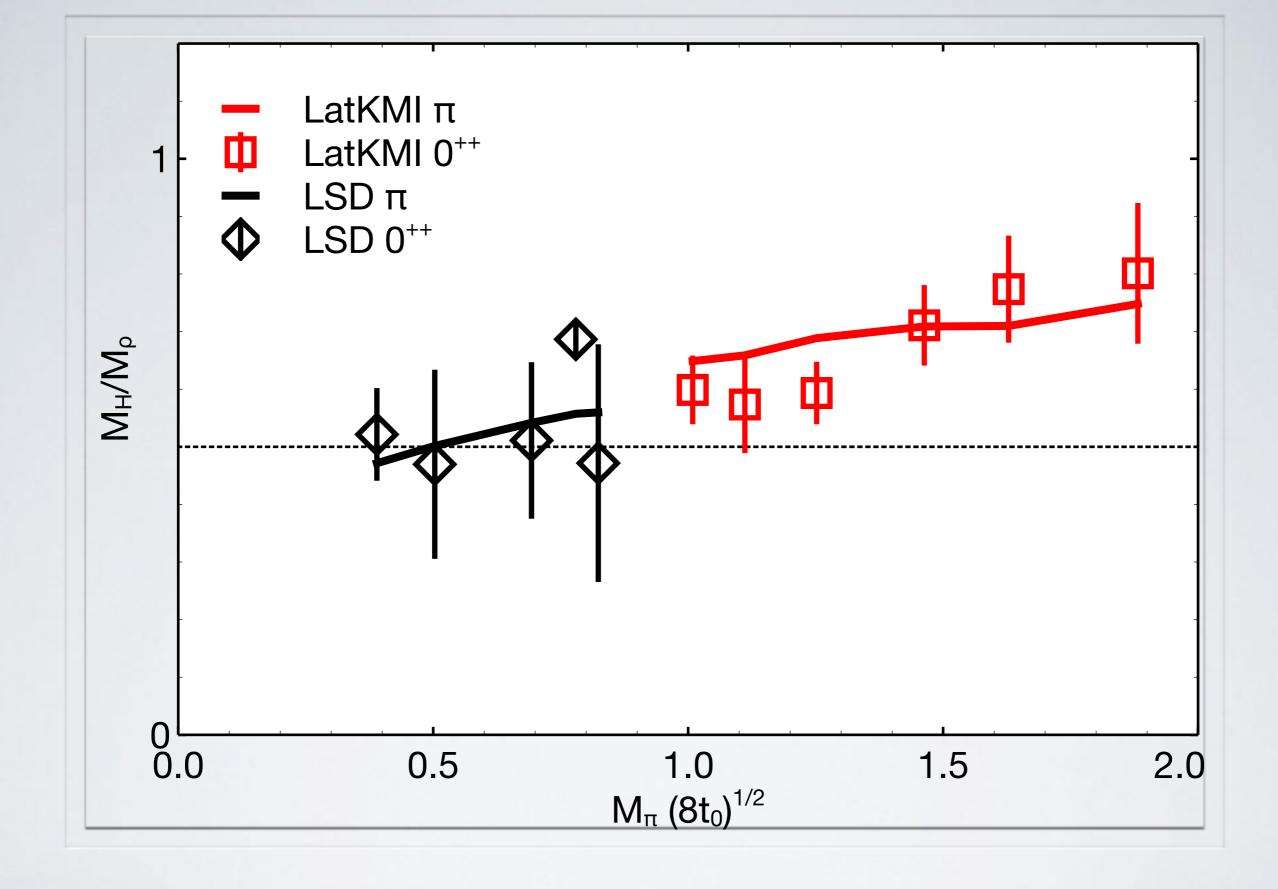


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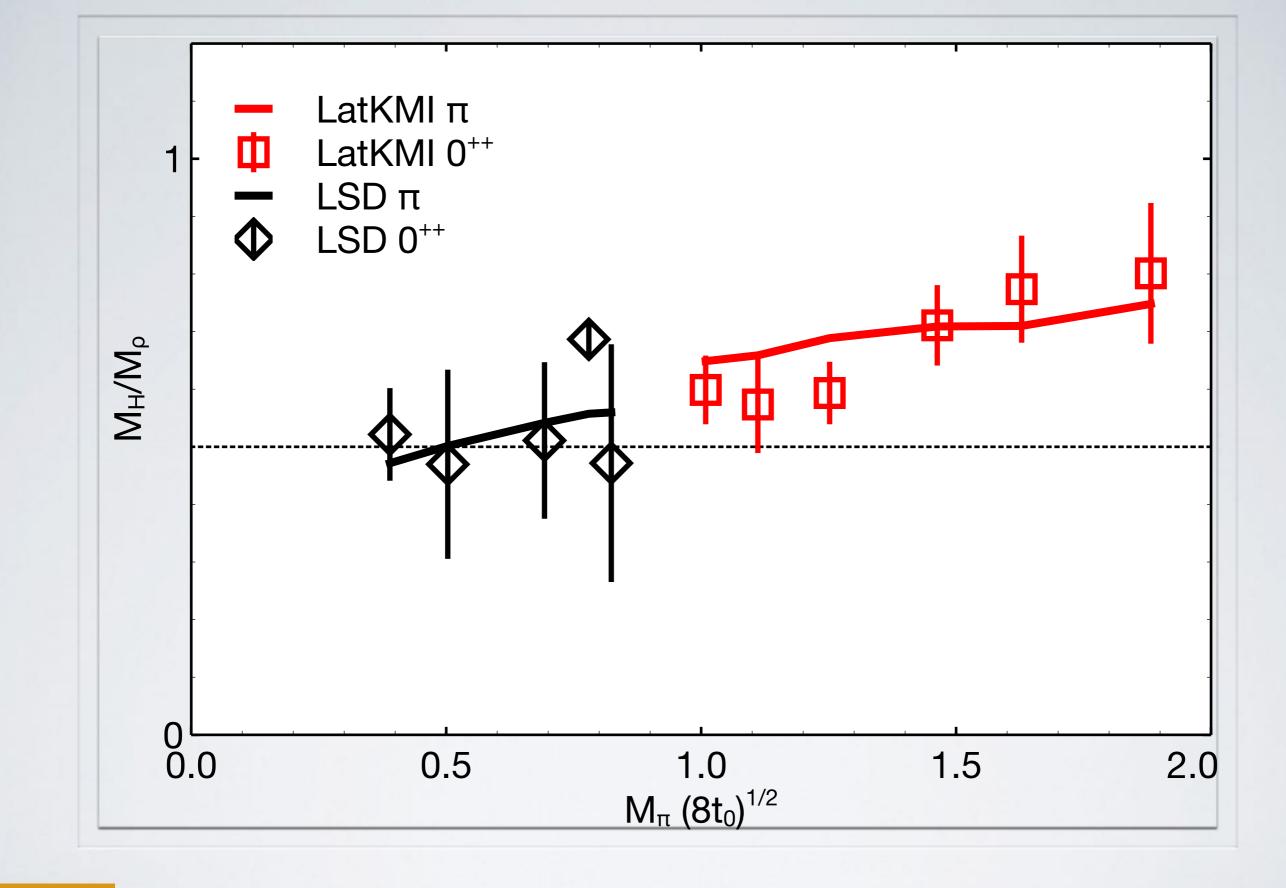


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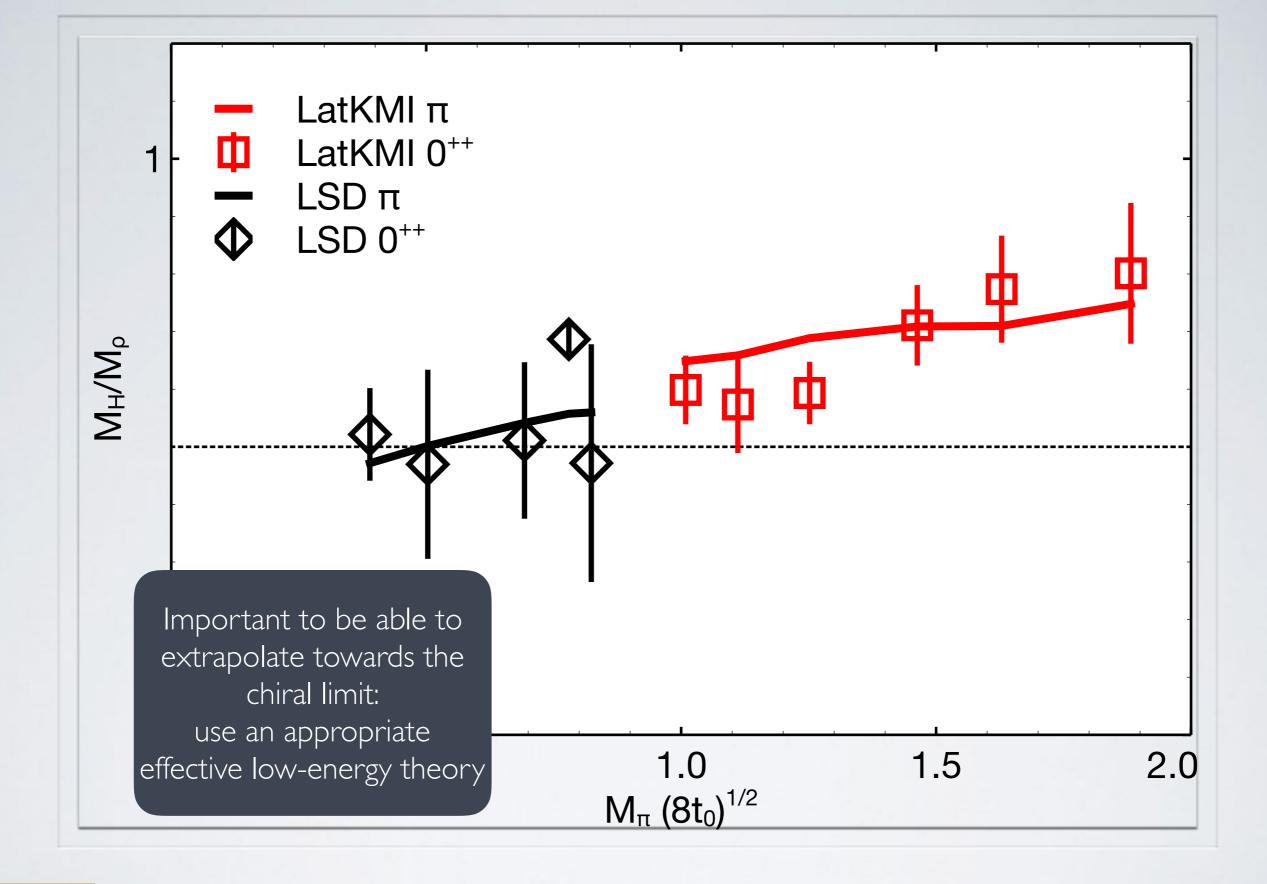


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#### • ChPT + "scalar" is a <u>new EFT</u> where the new light scalar d.o.f. is

treated as a "dilaton" [Goldberger, Grinstein & Skiba: 0708.1463, Matsuzaki & Yamawaki: 1311.3784, Golterman & Shamir: 1603.04575, Kasai, Okumura & Suzuki: 1609.02264, Hansen. Langae ble & Sannino: 1610.02904, Appelquist, Ingold by & Piai: 1702.04410]

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• Different forms of the Effective potential for the scalar and pseudoscalar can be directly tested against Lattice data [Appelquist,Ingoldby&Piai:1702.04410]

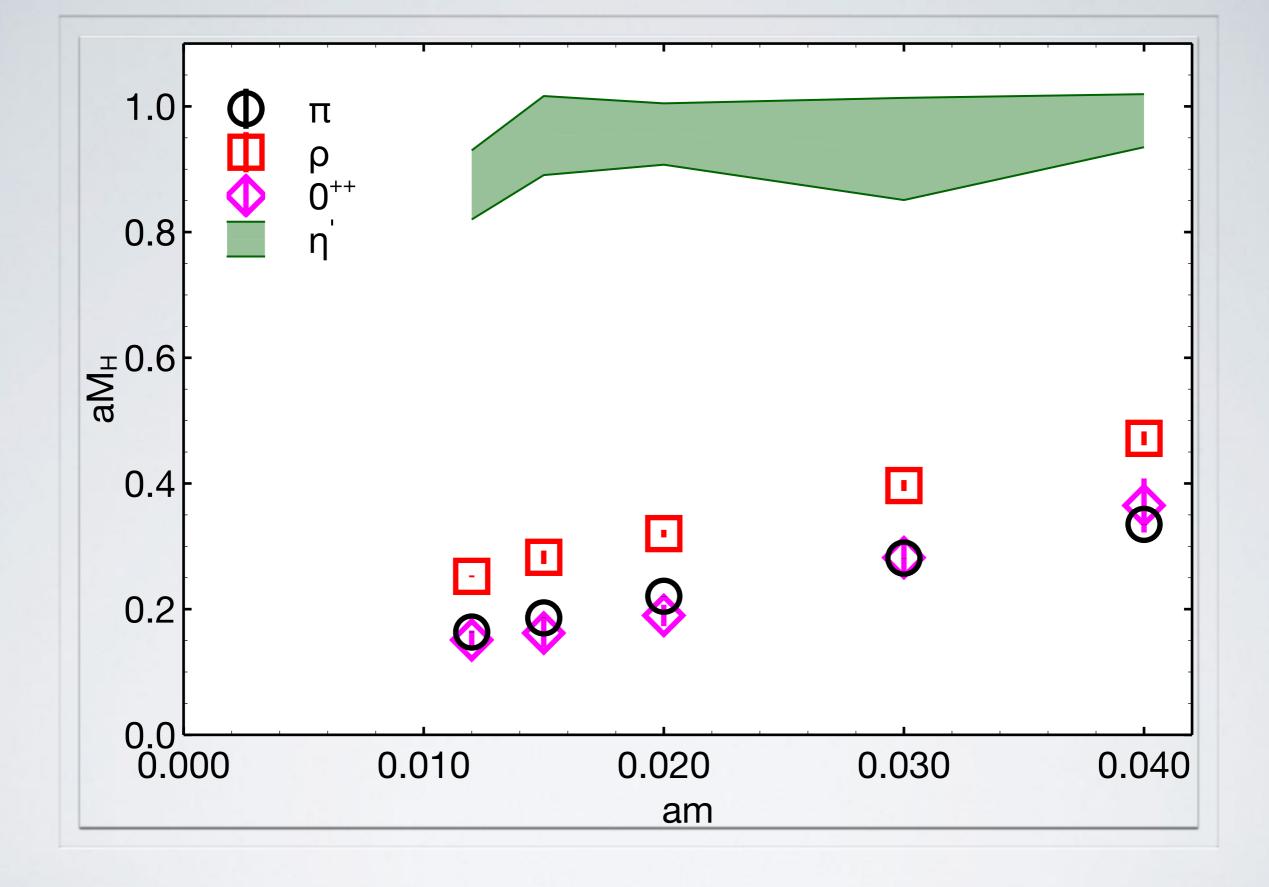


SU(3)

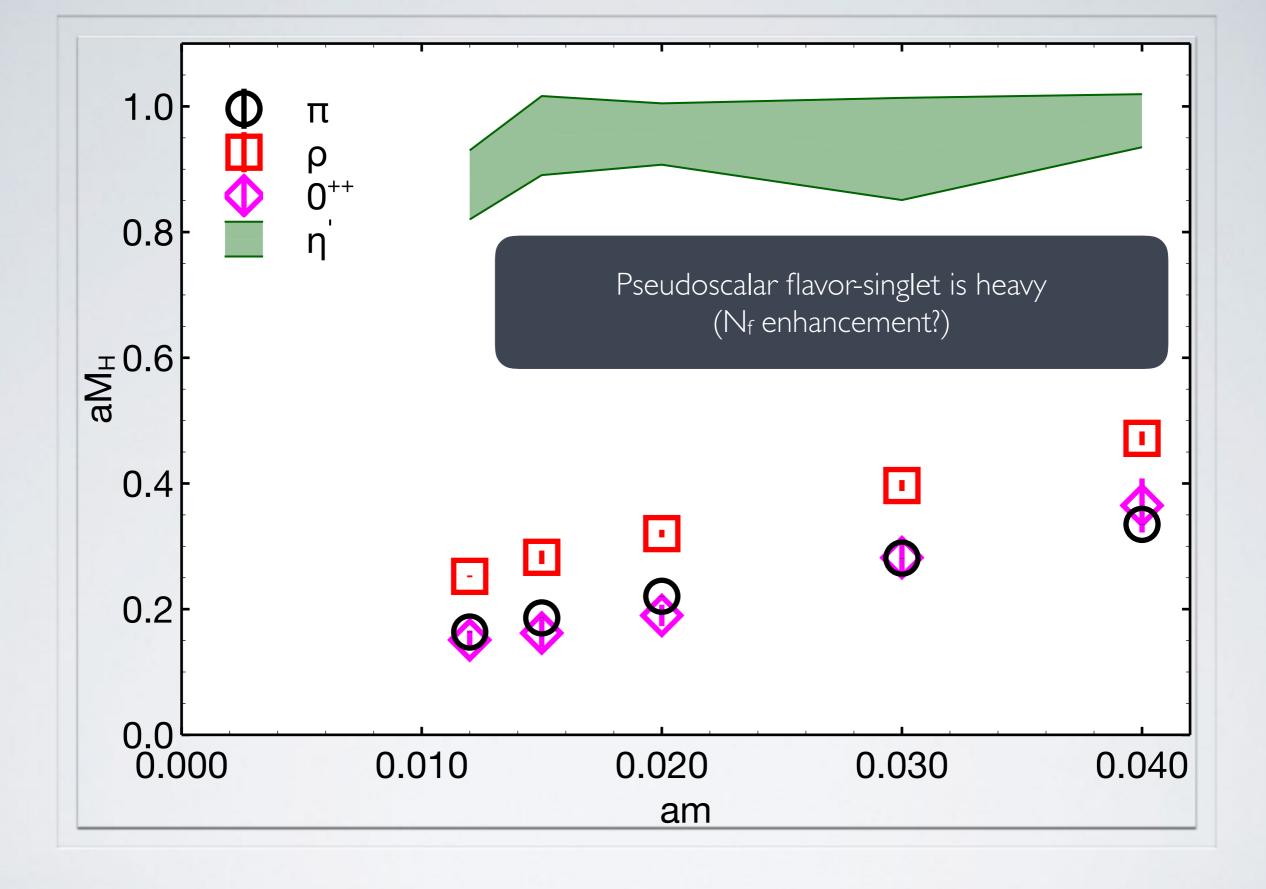
Nf=2 (S)

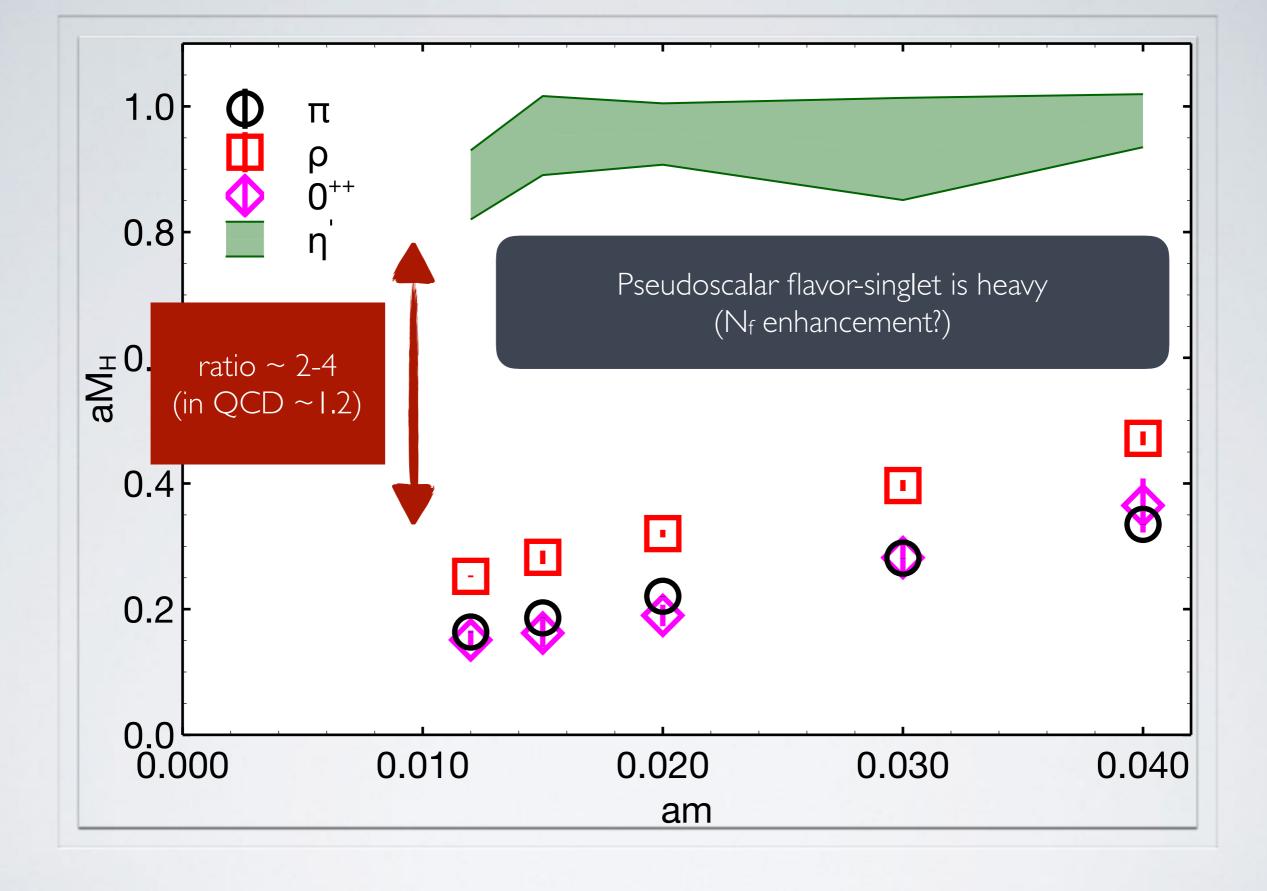
# SUMMARY RESULTS

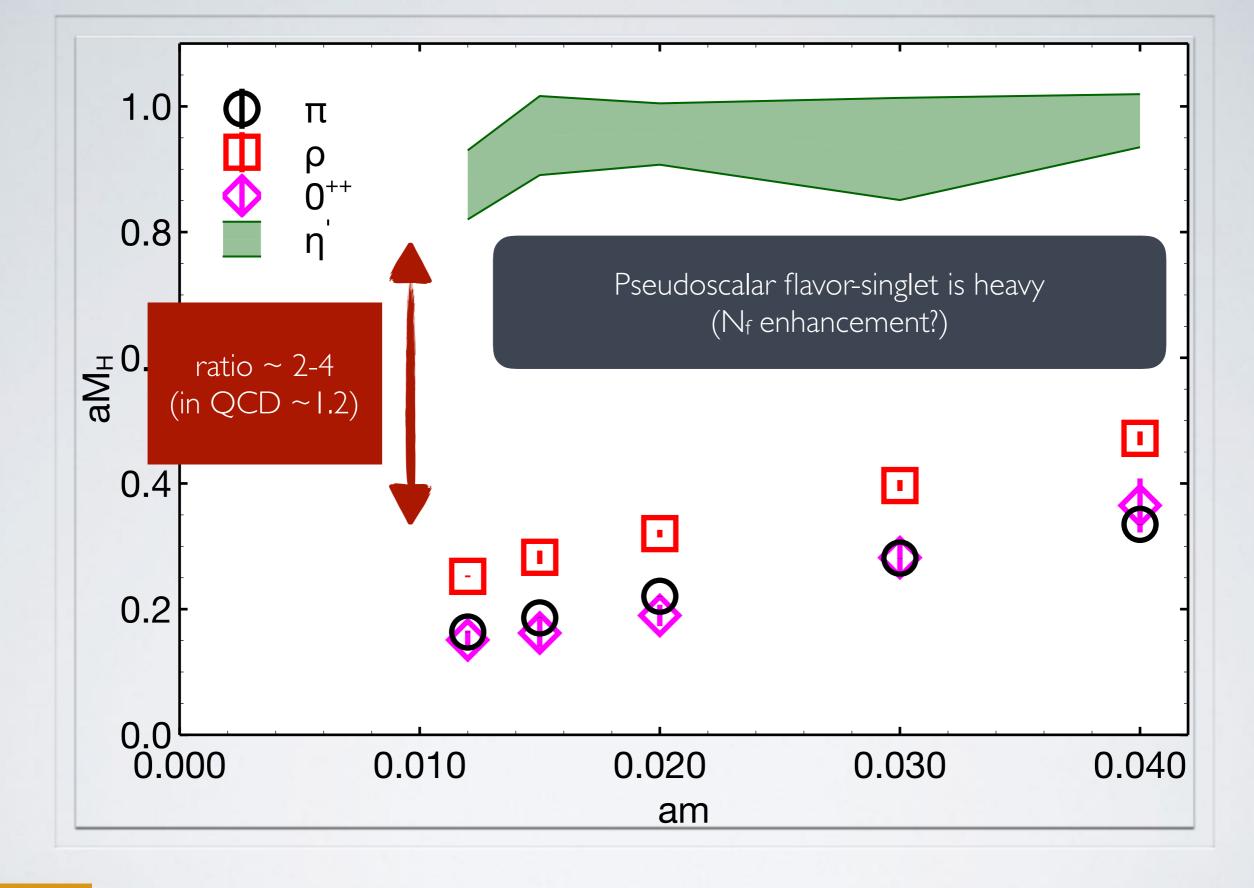
Template Models	Scalar	$\beta=3.20$ 14 Decreasing M <sub><math>\pi</math></sub> $n$ $a_1$	$\beta = 3.25$ $\bullet \pi$ $f_0$ Decreasing $M_{\pi}$ 14	
SU(2) N <sub>f</sub> =2 (F)		$12 \begin{bmatrix} 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 1 \\ 1 & 1 &$		
SU(2) N <sub>f</sub> =2 (A)	$\star$	$\begin{array}{c ccccccccccccccccccccccccccccccccccc$	$\begin{bmatrix} \bullet & \bullet & \bullet & \bullet \\ \bullet & \bullet & \bullet & \bullet & \bullet & \bullet \\ \bullet & \bullet &$	
SU(2) N <sub>f</sub> =1 (A)	$\star$	$\begin{array}{c} 4 \\ \bullet \\ 2 \\ \bullet \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0$		
SU(3) N <sub>f</sub> =12 (F)	$\star$	$0.05  0.04  0.03  0.02  0.01  0  0  0  M_{\pi}^2 \qquad M / T$	$M_{\pi}^{2}$	
SU(3) N <sub>f</sub> =8 (F)	$\bigstar$		→ 0 <sup>-</sup> vector meson → 0 <sup>+</sup> glueball SU Nf=	J(2) I (A)
SU(3) N <sub>f</sub> =4 (F)		mState /		
SU(3) N <sub>f</sub> =2 (S)			.08 0.1 0.12 0.14 0.16 PCAC	



LatKMI arxiv:1610.07011 [+ in preparation]

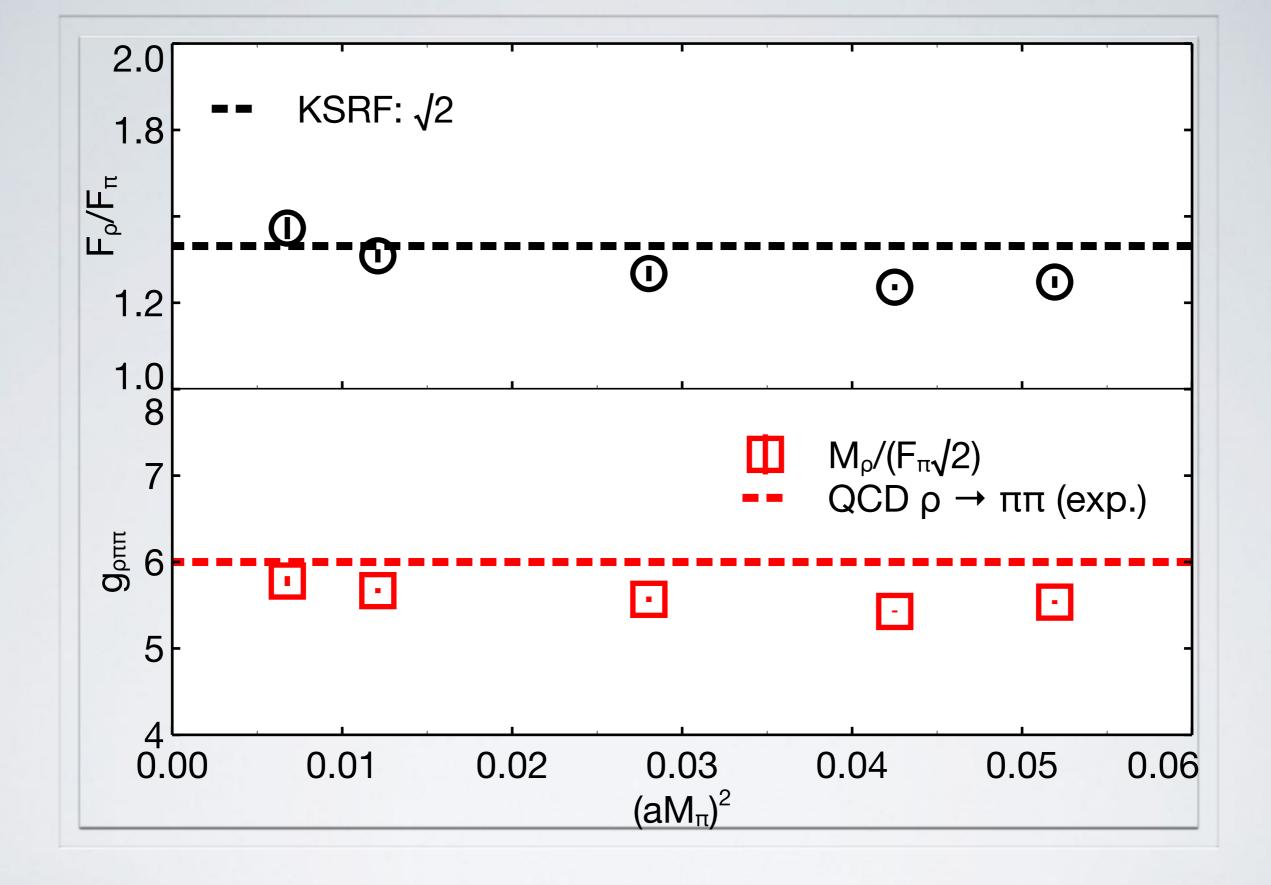


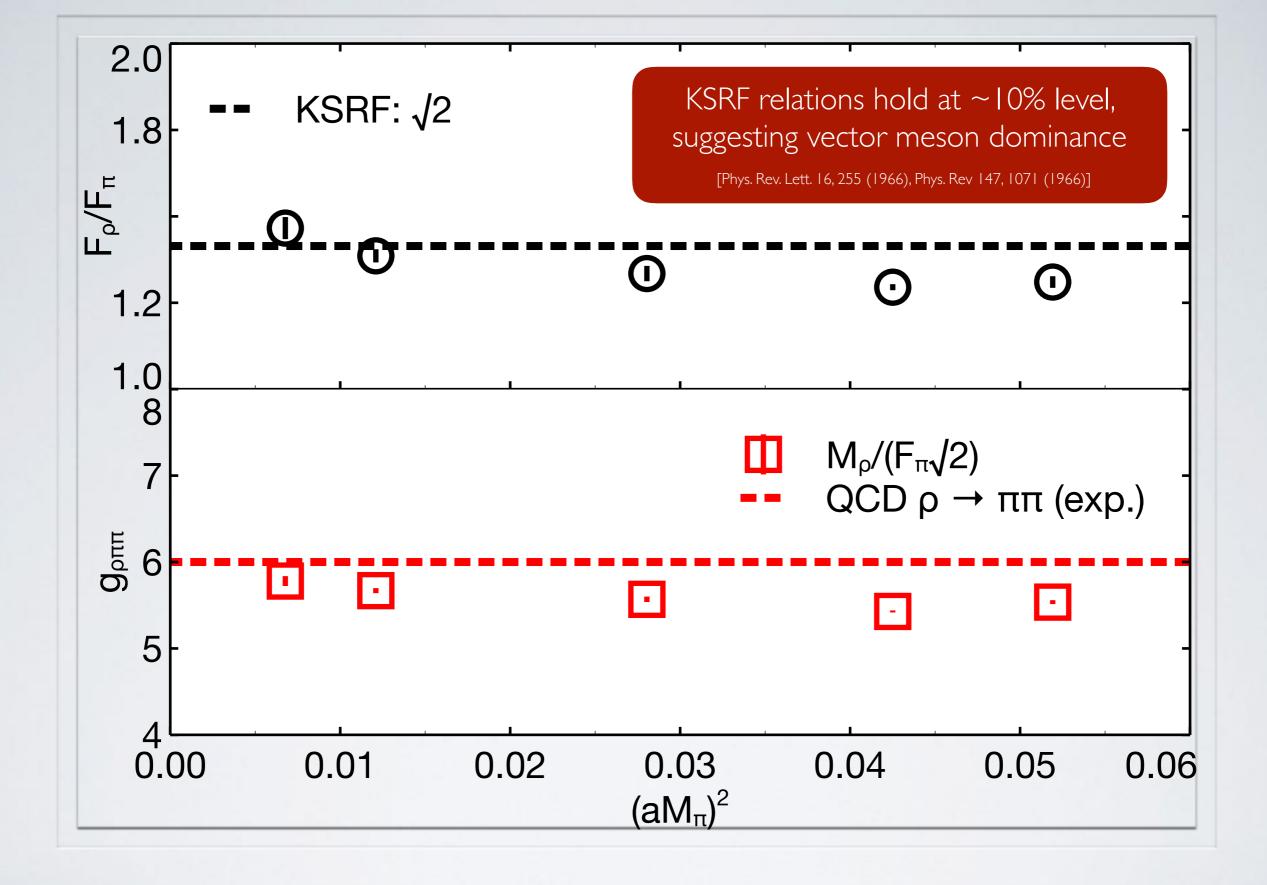


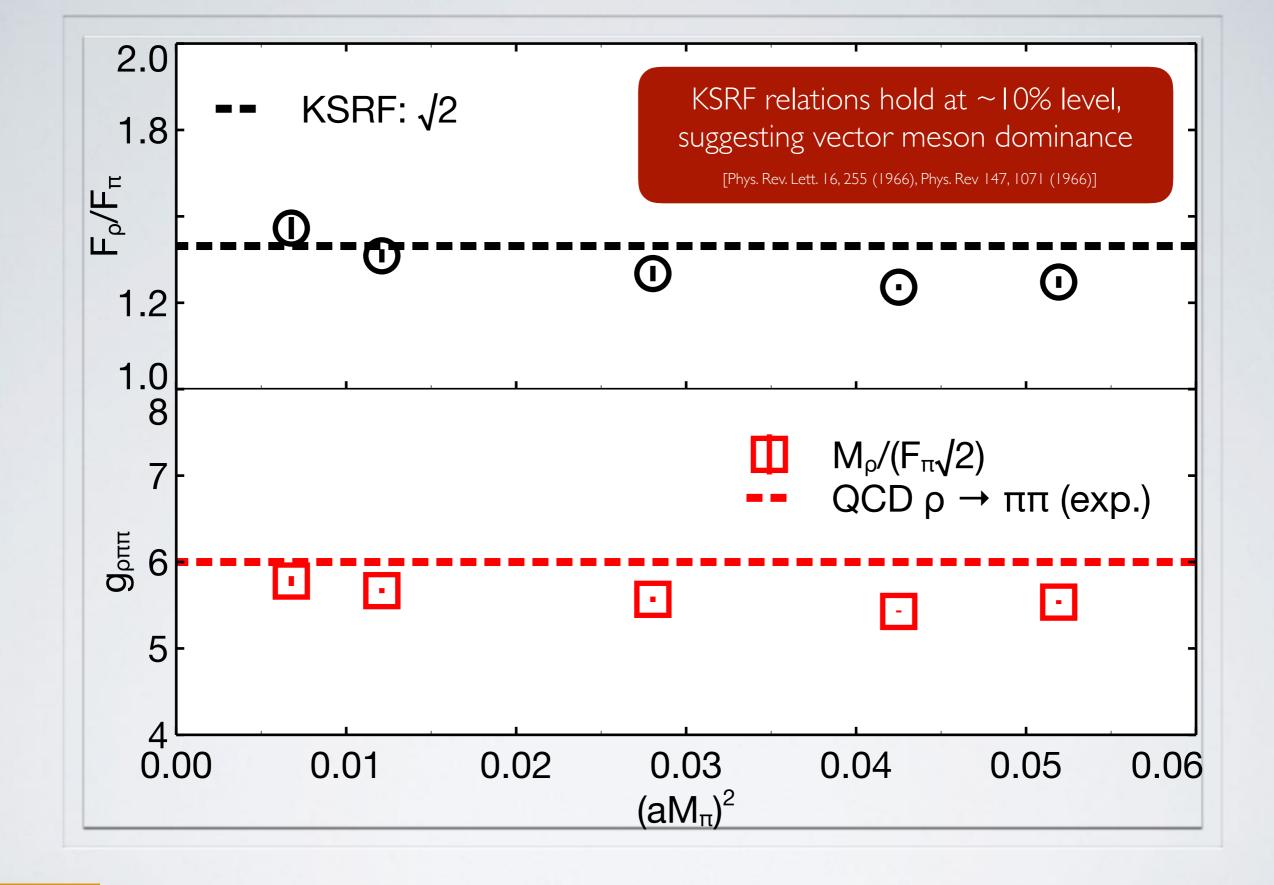


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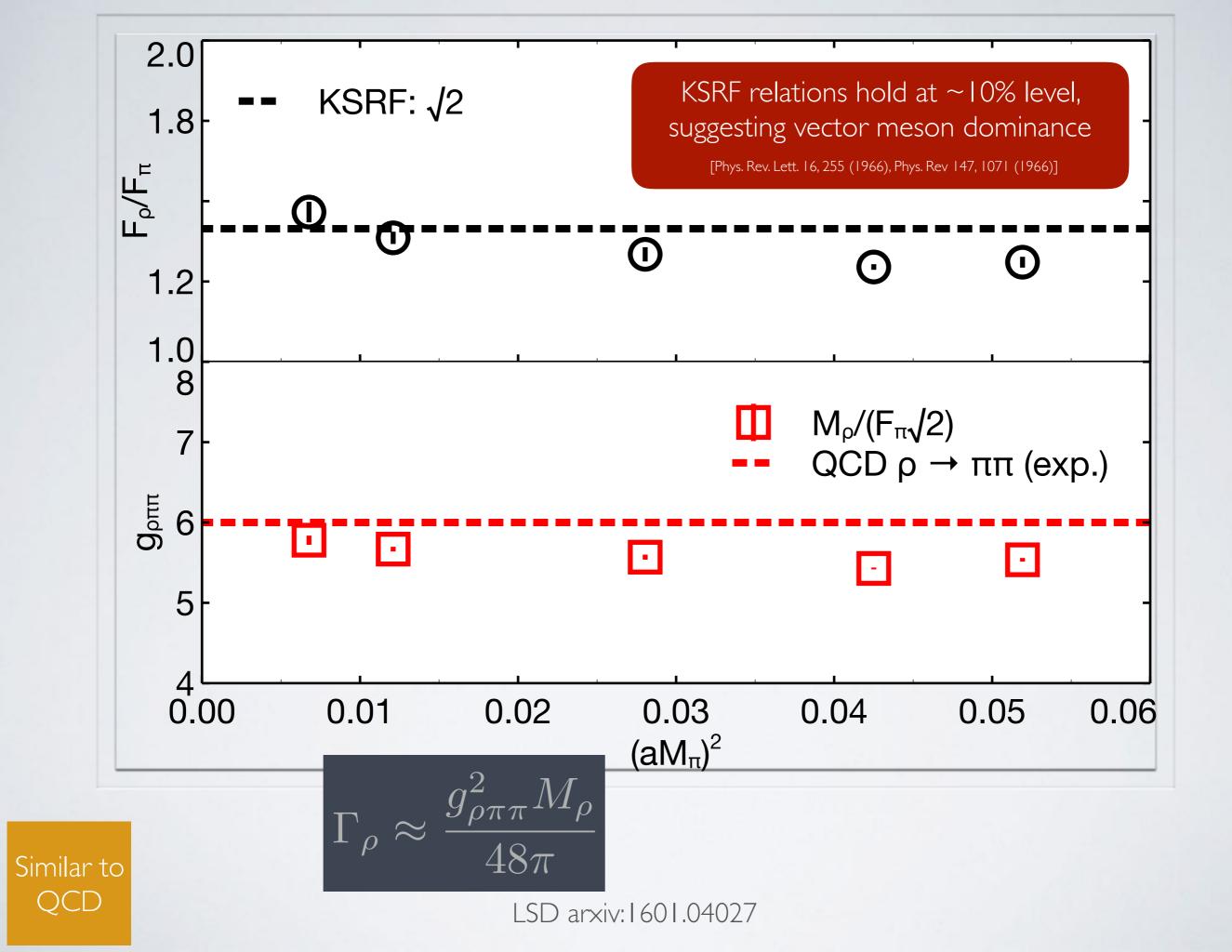


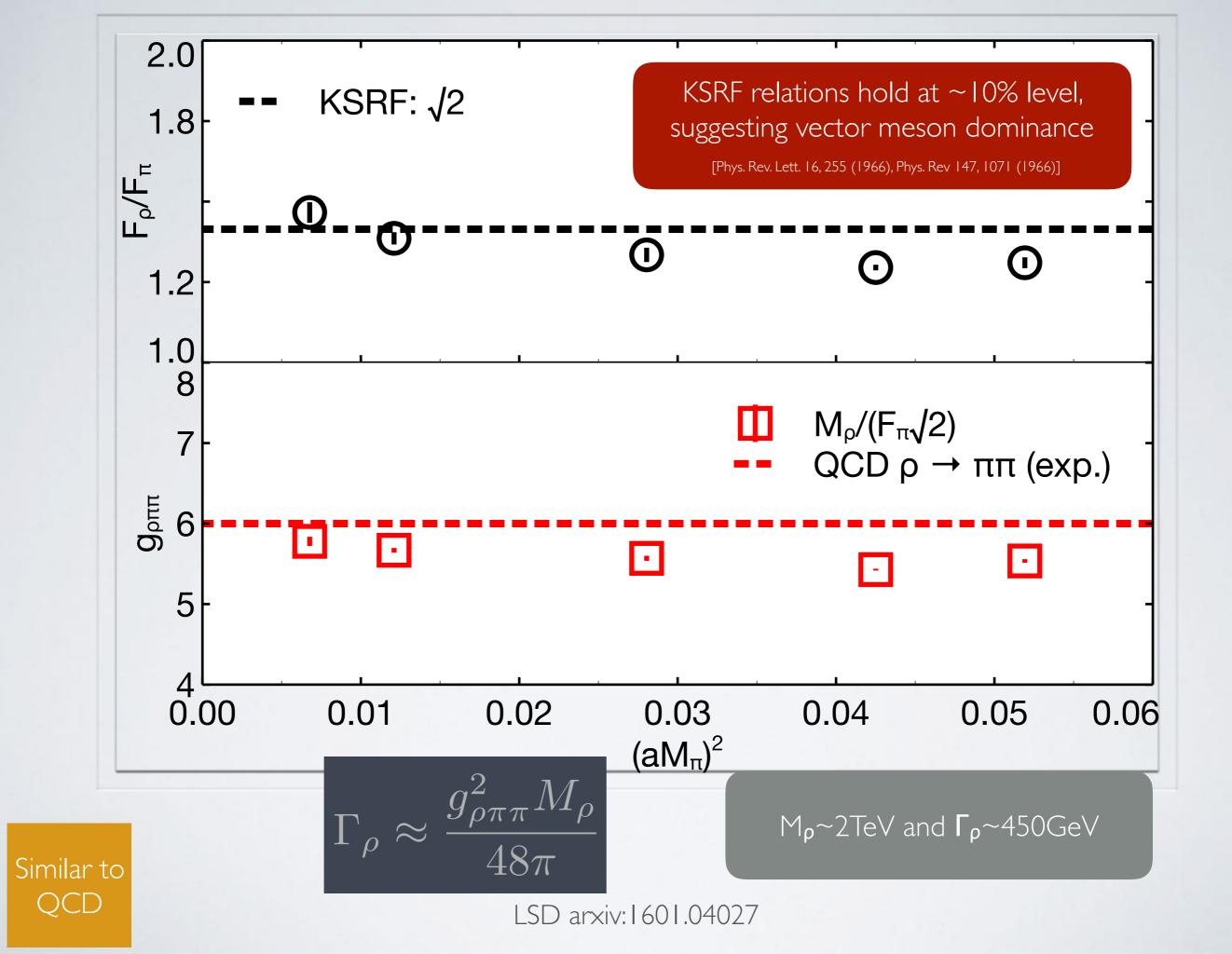




Similar to QCD

LSD arxiv:1601.04027





# CONCLUSIONS

- Lattice results for the SU(3) N<sub>f</sub>=8 theory using different discretizations are painting a consistent picture: <u>there is a light scalar flavor-singlet state</u>, the dynamics is different form QCD, but other heavier resonances behave similarly to QCD
- <u>A light scalar</u> has been observed in other systems that are very close or inside the conformal window
- Lattice studies can give information about other hadronic quantities:
  - S-parameter [LSD arxiv:1405.4752, LatKMI arxiv:1602.00796]
  - couplings between flavor-singlet scalar and pseudoscalars, e.g.  $\pi\pi$  scattering in the scalar channel [LSD arXiv:1702.00480 + in prep.]
  - Higgs coupling to SM fermions, e.g. dilaton decay constant [LatKMI arxiv:1610.07011]
  - anomalous dimensions [LSD arxiv:1405.4752, LatKMI arxiv:1610.07011]

Caveat: infinite volume limit, continuum limit and chiral limit need to be worked on!

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extra slides

★ Pion-like (dark quark-antiquark)

- ♦ pNGB DM [Hietanen et al., 1308.4130]
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- "Technibaryons" [LSD, 1301.1693]
- ✦ Stealth DM [LSD, 1503.04203-1503.04205]
- One-family TC [LatKMI, 1510.07373]

✦ Sextet CH [LatHC, 1601.03302]

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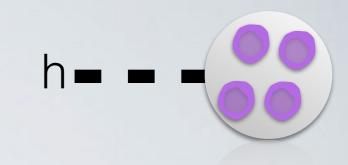
Template Models	Spectrum	Higgs	Mag. Dip.	Charge r.	Polariz.
$SU(2) N_{f}=1$	$\star$				
SU(2) N <sub>f</sub> =2	$\star$	$\star$		$\star$	$\star$
SU(3) N <sub>f</sub> =2,6	$\star$		$\star$	$\star$	
SU(3) N <sub>f</sub> =8	$\star$	$\star$			
SU(3) N <sub>f</sub> =2 (S)	$\star$				
$SU(4) N_f=4$	$\star$	$\mathbf{\star}$			$\star$
SO(4) N <sub>f</sub> =2 (V)	$\star$				
SU(N) Nf=0	$\star$				



# LATTICE RESULTS

Template Models	Spectrum	Higgs	Mag. Dip.	Charge r.	Polariz.
SU(2) N <sub>f</sub> =1	$\star$				
SU(2) N <sub>f</sub> =2	$\star$	$\star$	forbidden in pNGB DM	$\star$	$\star$
SU(3) N <sub>f</sub> =2,6	$\star$			$\star$	
SU(3) N <sub>f</sub> =8	$\star$	*			
SU(3) N <sub>f</sub> =2 (S)	$\star$				
SU(4) N <sub>f</sub> =4	$\bigstar$	$\star$	forbiddon in	Steelth DM	$\star$
SO(4) N <sub>f</sub> =2 (V)	$\star$		Torbidden in	Stealth DM	
SU(N) Nf=0			forbidden i	n SUNonia	

### Bounds from Higgs exchange

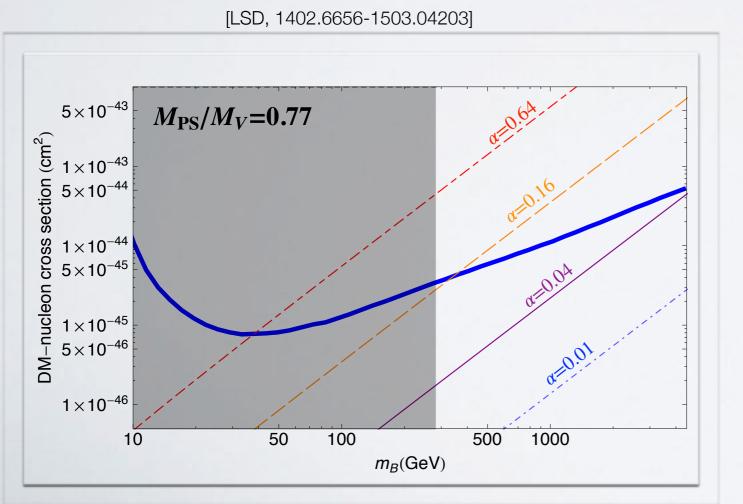


#### SU(3) N<sub>f</sub>=8 "technibaryon"

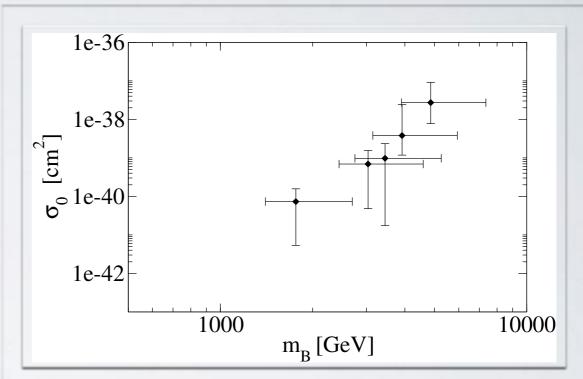
Lattice results for the cross-section are compared to experimental bounds

 Coupling space in specific models can be vastly constrained

SU(4) Nf=4 Stealth DM

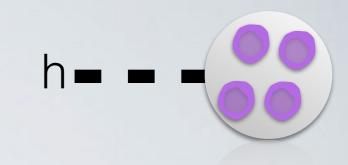


[LatKMI, 1510.07373]



- Some candidates can be excluded as dominant sources of dark matter
- There is lattice evidence for universality of dark scalar form factors [DeGrand et al., 1501.05665]

### Bounds from Higgs exchange



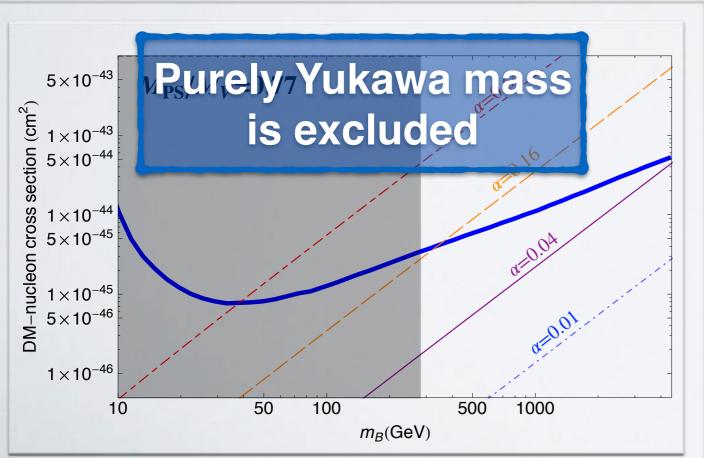
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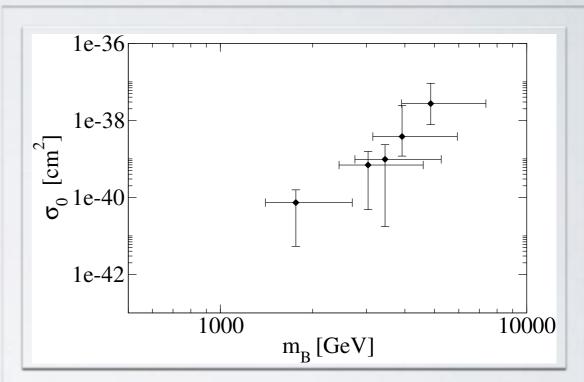
 Coupling space in specific models can be vastly constrained

SU(4) Nf=4 Stealth DM

[LSD, 1402.6656-1503.04203]

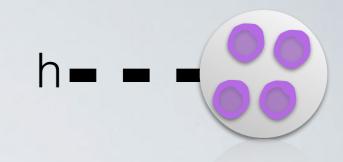


[LatKMI, 1510.07373]



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### Bounds from Higgs exchange



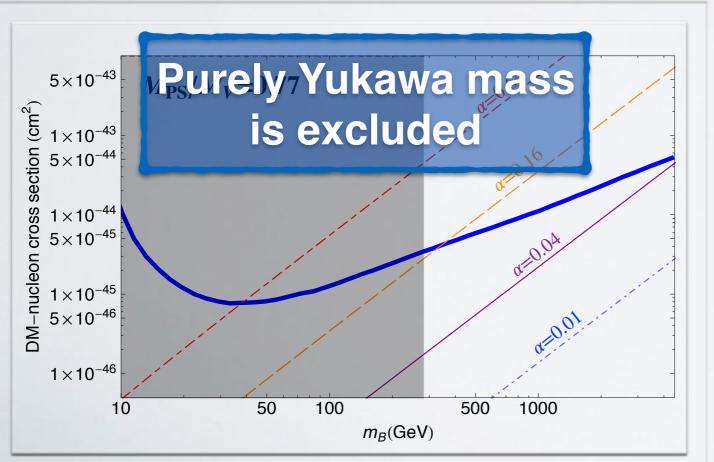
10000

#### SU(3) N<sub>f</sub>=8 "technibaryon"

- Lattice results for the cross-section are compared to experimental bounds
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SU(4) Nf=4 Stealth DM

[LSD, 1402.6656-1503.04203]



[LatKMI, 1510.07373]

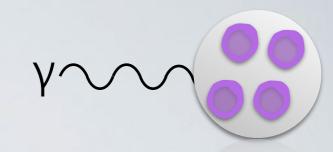
m<sub>R</sub> [GeV]

 Some candidates can be excluded as dominant sources of dark matter

1000

There is lattice evidence for universality of dark scalar form factors [DeGrand et al., 1501.05665]

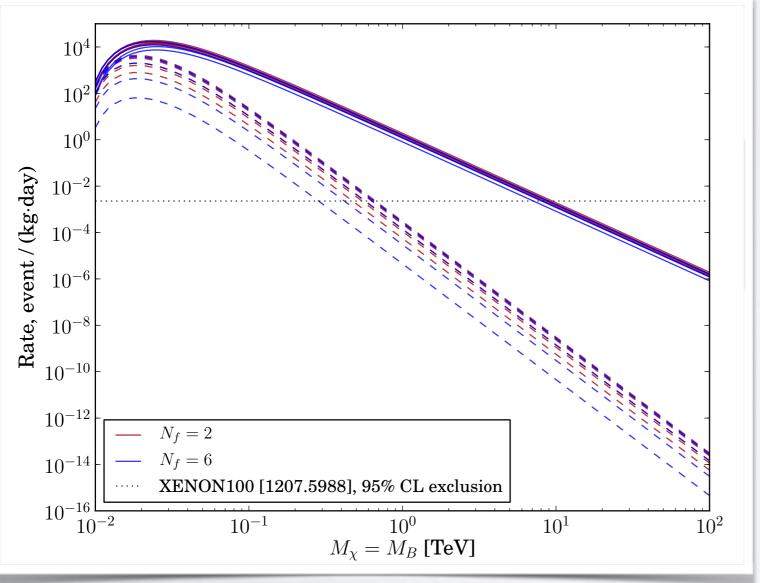
### Bounds from EM moments



Mesonic and Baryonic EM form factors directly from lattice simulations

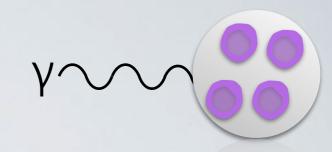
#### SU(3) N<sub>f</sub>=2,6 dark fermionic baryon





★ baryon similar to QCD neutron
★ dark quarks with Q=Y
★ calculate connected 3pt
★ scale set by DM mass
★ magnetic moment dominates
★ results independent of N<sub>f</sub>

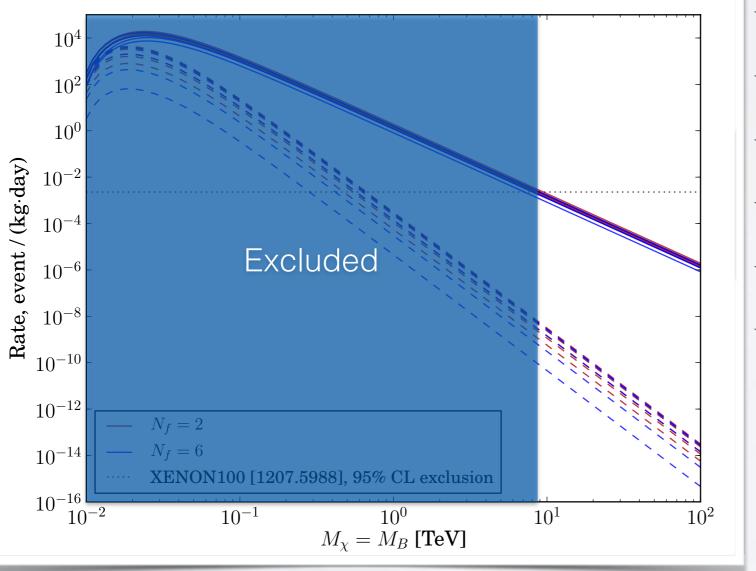
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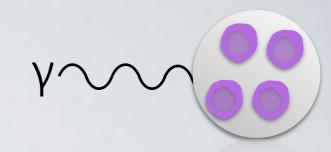




★ baryon similar to QCD neutron
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★ magnetic moment dominates
★ results independent of N<sub>f</sub>

 $M_B > \sim 10 \text{ TeV}$ 

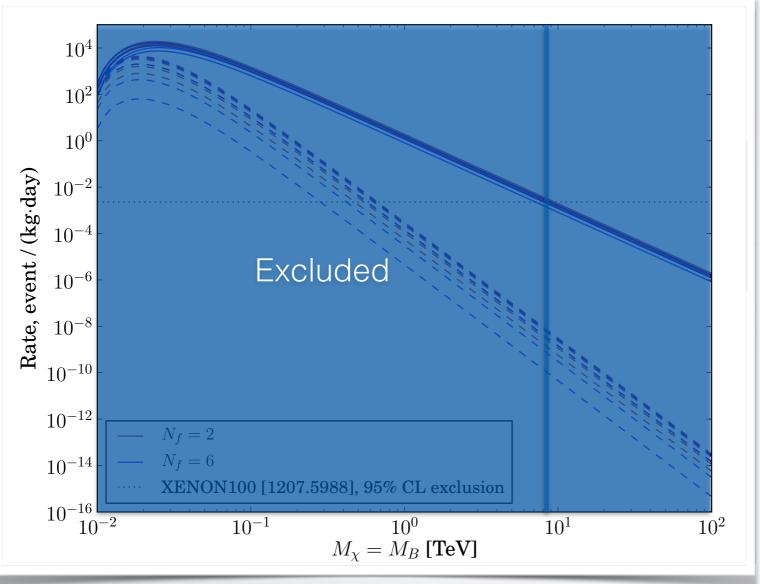
### Bounds from EM moments



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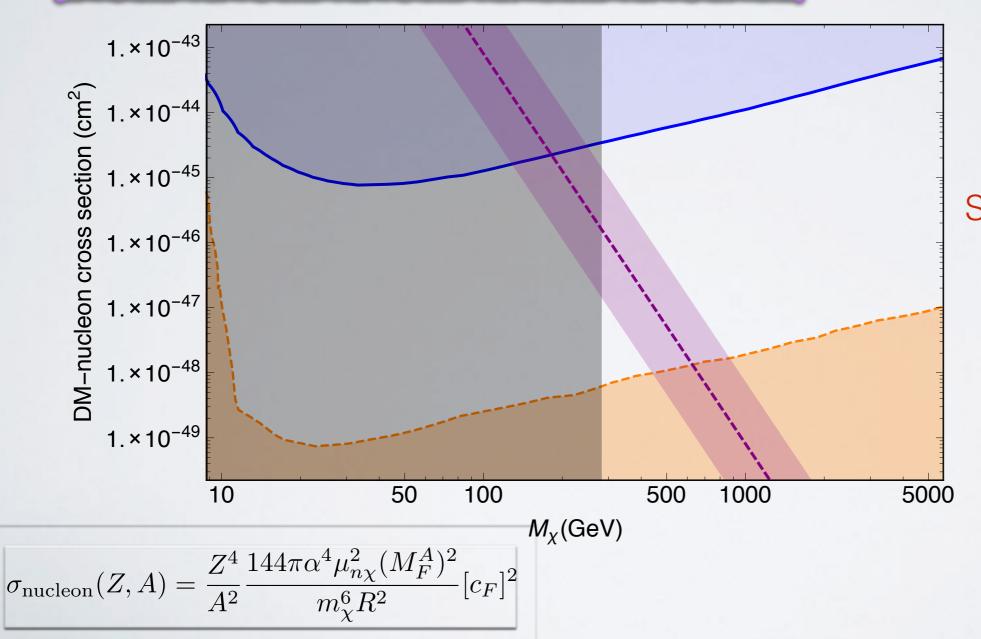
M<sub>B</sub> >~ 10 TeV pushed to ~100 TeV with new LUX

## Lowest bound from EM polarizability

[see also Drach et al.,1511.04370 for SU(2)  $N_f=2$ ]



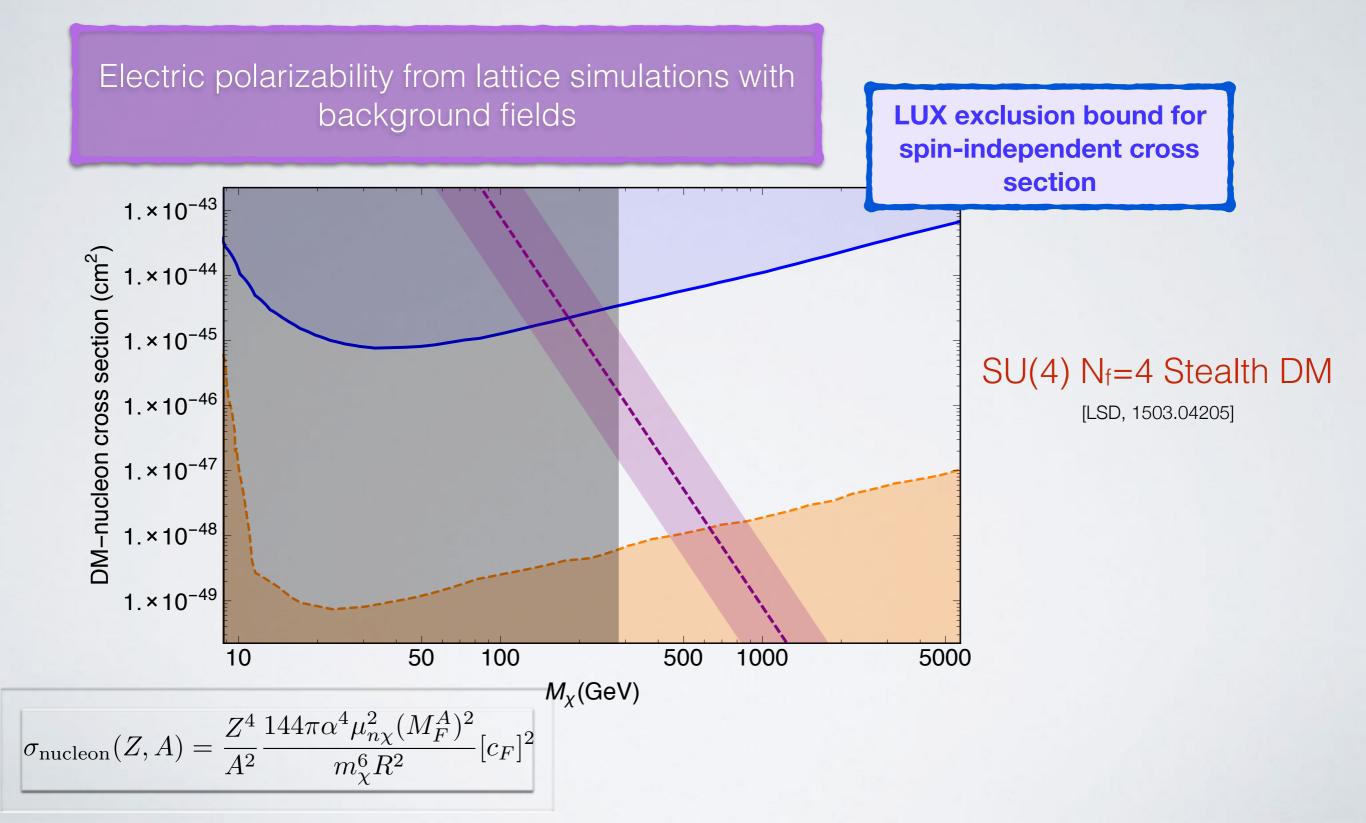
Electric polarizability from lattice simulations with background fields



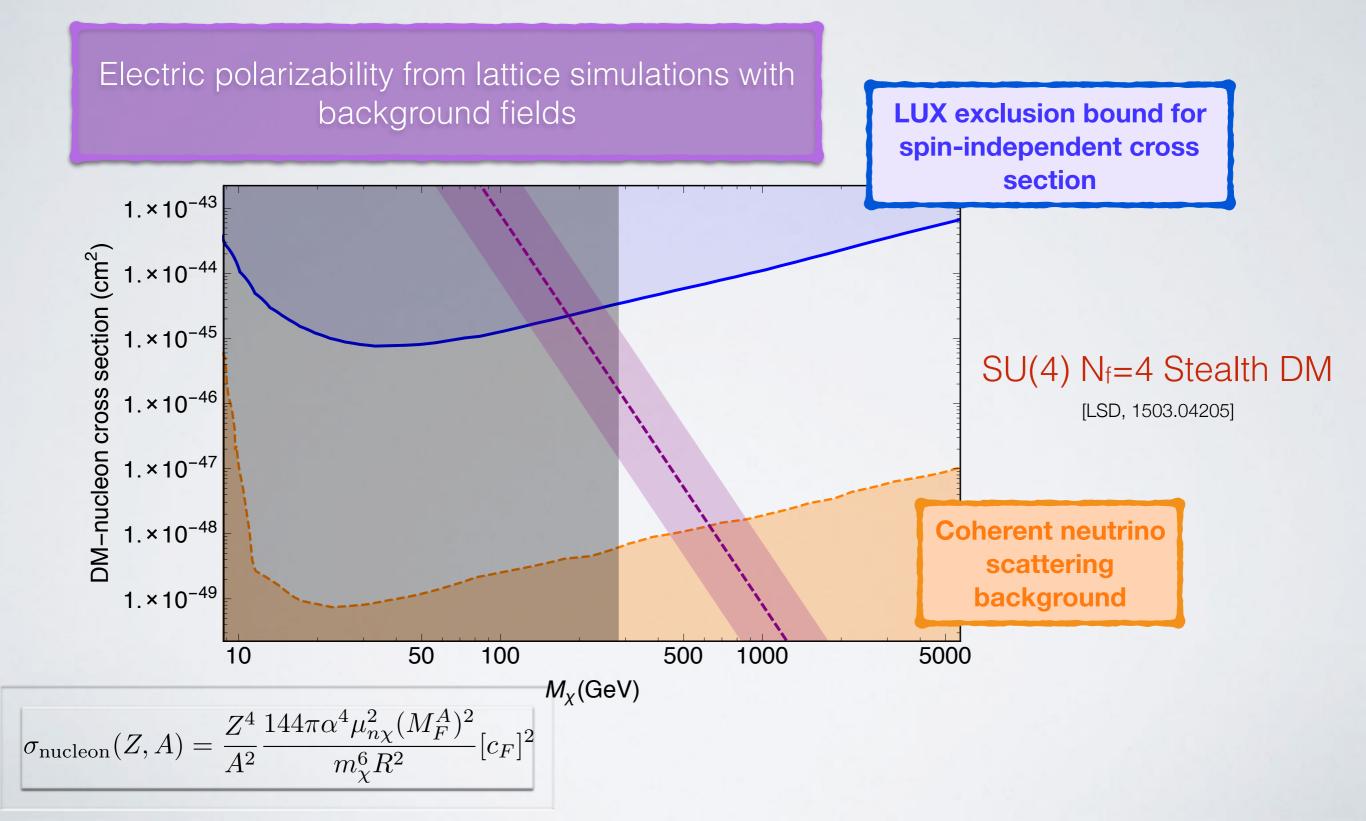
SU(4) N<sub>f</sub>=4 Stealth DM

<sup>[</sup>LSD, 1503.04205]

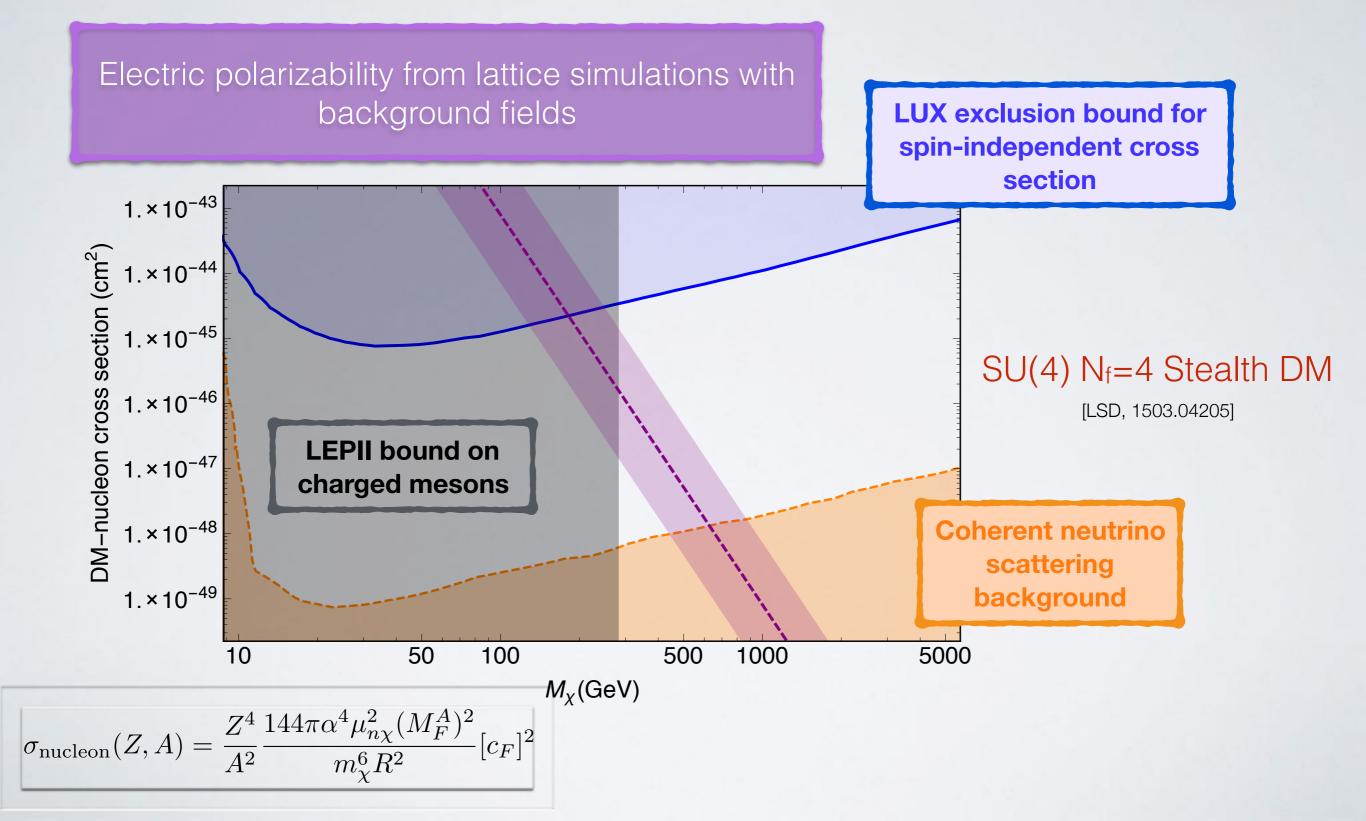




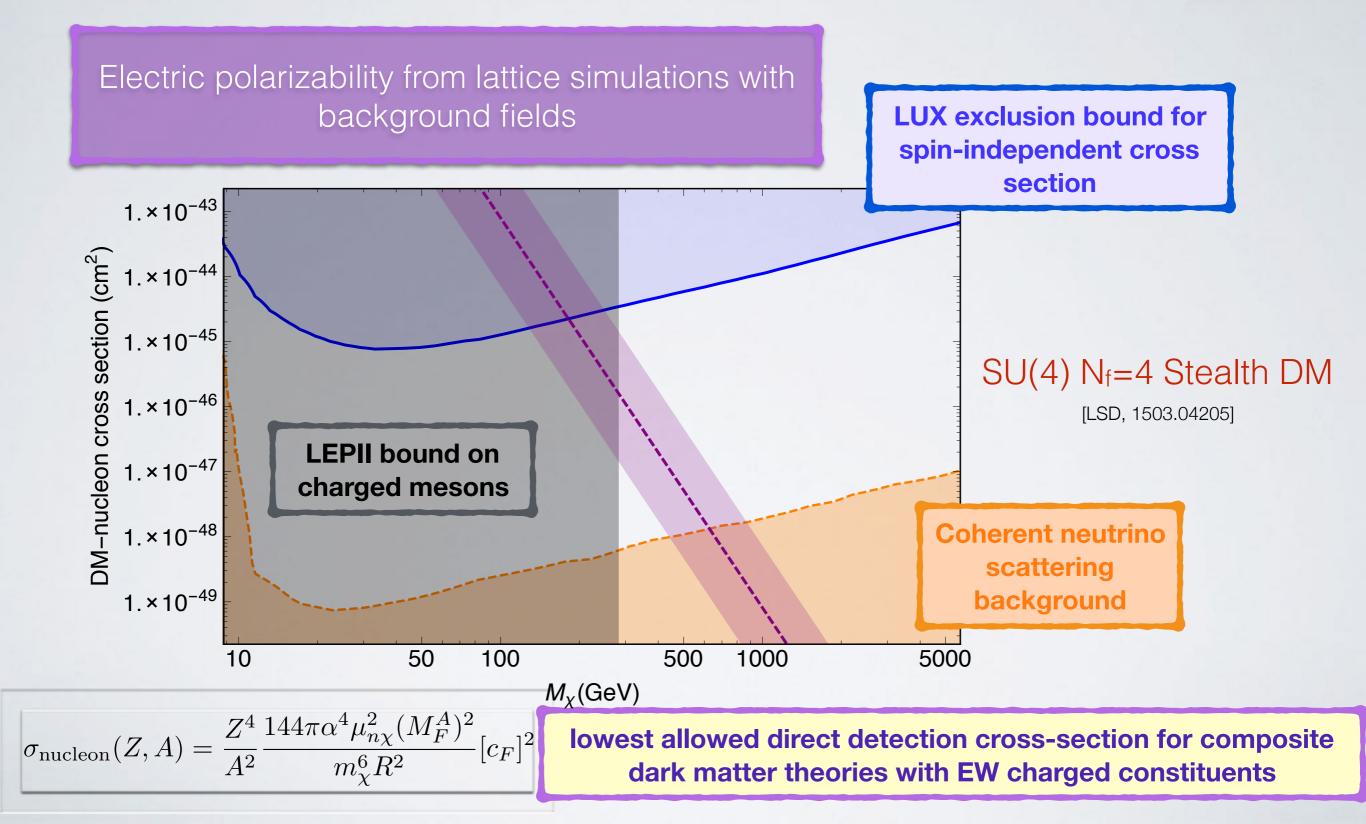












## "Stealth Dark Matter" model

[LSD collab., Phys. Rev. D92 (2015) 075030]

• The field content of the model  $\operatorname{SU}(N)_D \left| (\operatorname{SU}(2)_L, Y) \right|$ Field Qconsists in 8 Weyl fermions +1/2 $F_1 =$ (2, 0)Ν  $F_1^d$ -1/2 $\begin{bmatrix} F_2^u \\ F_2^d \end{bmatrix}$ +1/2 $\overline{\mathbf{N}}$  $({\bf 2}, 0)$ Dark fermions interact with the SM Higgs and obtain current/chiral  $F_3^u$ Ν (1, +1/2)+1/2 $F_3^d$ (1, -1/2)masses N -1/2 $\overline{\mathbf{N}}$  $F_4^u$ (1, +1/2)+1/2 $F_4^d$  $\overline{\mathbf{N}}$ (1, -1/2)-1/2Introduce vector-like masses for dark fermions that do not break EW symmetry  $\mathcal{L} \supset -y_{14}^{u} \epsilon_{ij} F_{1}^{i} H^{j} F_{4}^{d} + y_{14}^{d} F_{1} \cdot H^{\dagger} F_{4}^{u}$  $-y_{23}^{d}\epsilon_{ij}F_{2}^{i}H^{j}F_{3}^{d}-y_{23}^{u}F_{2}\cdot H^{\dagger}F_{3}^{u}+h.c.$ Diagonalizing in the mass eigenbasis gives 4 Dirac fermions  $\mathcal{L} \supset M_{12} I_{ij} F_1^i F_2^j - M_{34}^u F_3^u F_4^d + M_{34}^d F_3^d F_4^u + h.c.$ • Assume custodial SU(2) symmetry  $y_{14}^{u} = y_{14}^{d}$   $y_{23}^{u} = y_{23}^{d}$   $M_{34}^{u} = M_{34}^{d}$ arising when  $\boldsymbol{u} \leftrightarrow \boldsymbol{d}$ 

### "Stealth Dark Matter" model [LSD collab., Phys. Rev. D92 (2015) 075030] EW interactions

The field content of the model consists in *8 Weyl fermions*

 Dark fermions interact with the SM Higgs and obtain current/chiral masses

 Introduce vector-like masses for dark fermions that do not break EW symmetry

 Diagonalizing in the mass eigenbasis gives 4 Dirac fermions

 Assume custodial SU(2) symmetry arising when *u* ↔ *d*

Field	$ $ SU $(N)_D$	$(\mathrm{SU}(2)_L, Y)$	Q
$F_1 = \begin{pmatrix} F_1^u \\ F_1^d \end{pmatrix}$	N	( <b>2</b> ,0)	$\left(\begin{array}{c} +1/2\\ -1/2 \end{array}\right)$
$F_2 = \begin{pmatrix} F_2^u \\ F_2^d \end{pmatrix}$	$\overline{\mathbf{N}}$	( <b>2</b> ,0)	$\begin{pmatrix} +1/2 \\ -1/2 \end{pmatrix}$
$F_3^u$	N	(1, +1/2)	+1/2
$F_3^d$	N	(1, -1/2)	-1/2
$F_4^u$	$\overline{\mathbf{N}}$	(1, +1/2)	+1/2
$F_4^d$	$\overline{\mathbf{N}}$	(1, -1/2)	-1/2

 $\mathcal{L} \supset - y_{14}^{u} i_{j} F_{1}^{i} H^{j} F_{4}^{d} + y_{14}^{d} F_{1} \cdot H^{\dagger} F_{4}^{u}$  $- y_{23}^{d} \epsilon_{ij} F_{2}^{i} H^{j} F_{3}^{d} - y_{23}^{u} F_{2} \cdot H^{\dagger} F_{3}^{u} + h.c.$ 

 $\mathcal{L} \supset M_{12} F_{ij}^{i} F_{1}^{j} F_{2}^{j} - M_{34}^{u} F_{3}^{u} F_{4}^{d} + M_{34}^{d} F_{3}^{d} F_{4}^{u} + h.c.$ 

 $y_{14}^{u} = y_{14}^{d}$   $y_{23}^{u} = y_{23}^{d}$   $M_{34}^{u} = M_{34}^{d}$ 



#### Materia oscura "stealth"

Quark oscuri tenuti insieme da un'interazione forte a sua volta oscura. Ecco come la dark matter riuscirebbe a eludere a ogni tentativo d'incastraria. Enrico Rinaldi (LLNL): «Esiste la possibilità che questo "mondo oscuro", con le sue nuove particelle, possa essere rivelato dagli esperimenti in corso al Large Hadron Collider al CERN di Ginevra»

#### di Marco Malaspina 🛛 🚽 Segui Omolomiae

venerdì 25 settembre 2015 @ 16:15

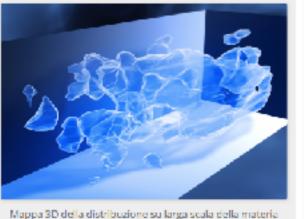
isiness

News

About

Careers

Stealth come furtiva. Stealth come imprendibile. Stealth come quei minacciosi aerei da guerra dal profilo sagomato così da essere invisibili al radar. Da guanto emerge dai calcoli dei fisici dell'LLNL, il Lawrence Livermore National Laboratory californiano, e dai modelli dati in pasto a Vulcan (un supercomputer per il calcolo parallelo in grado masticare numeri al ritmo dei petaflop), sarebbe questa la natura della materia oscura: steolthy, appunto. Per forza non c'è ancora esperimento che sia riuscito a incastrarla.



oscura ricostruita da misure di lente gravitazionale debole utilizzando il telescopio spaziale Hubble.

This 3D map illustrates the large-scale distribution of dark matter, reconstructed from measurements of weak gravitational lensing by using the Hubble Space Telescope. (Download Image)

### New 'stealth dark matter' theory may explain mystery of the universe's missing mass

Lawrence Livermore National Laboratory (LLNL) scientists have come up with a new theory that may identify why dark matter has evaded direct. detection in Earth-based experiments.

Anne M Stark stark8@lini.gov 🖾 925-422-9799

https://www.llnl.gov/news/new-stealth-dark-matter-

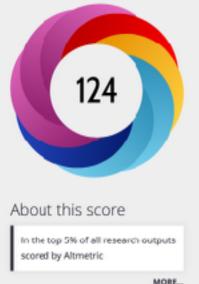
Community theory-may-explain-mystery-universes-missing-mass

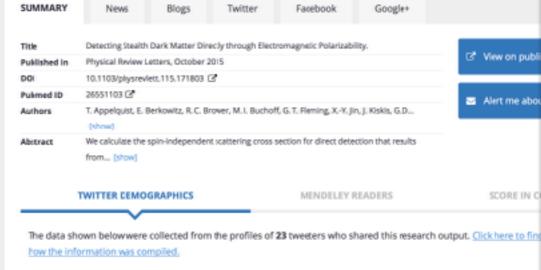


Un nuovo modello per la materia oscura

### Detecting Stealth Dark Matter Directly through Electromagnetic Polarizability.

Overview of attention for article published in Physical Review Letters, October 2015

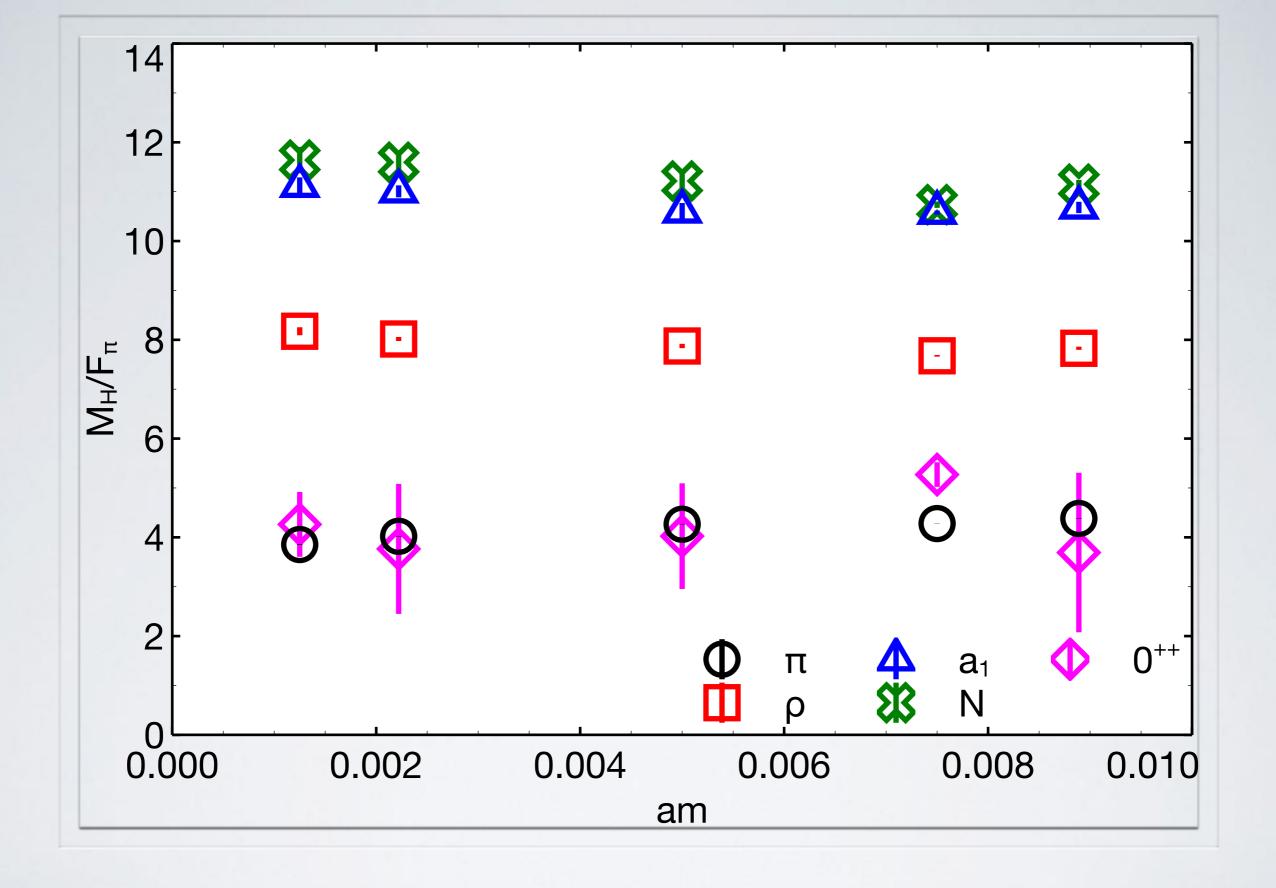




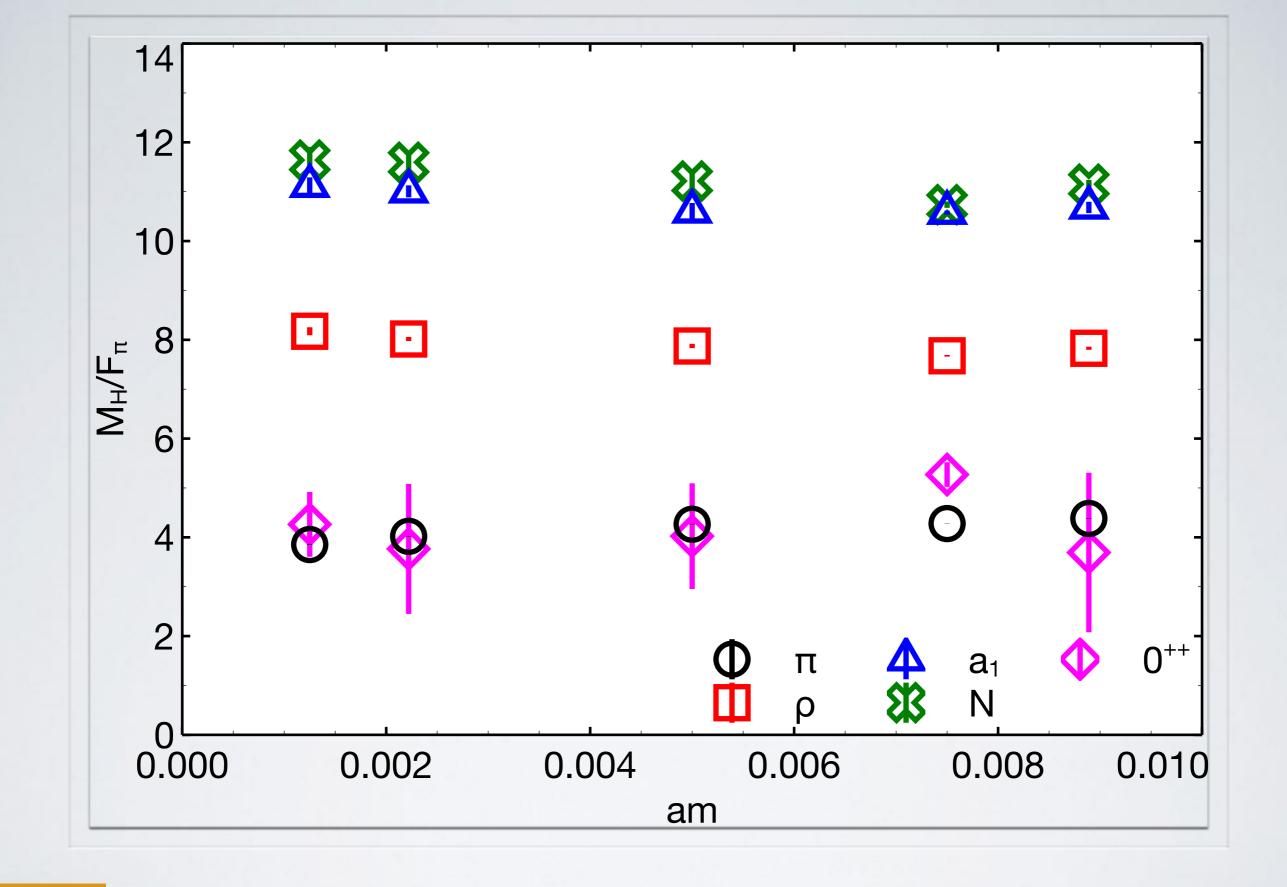
Questa forma misteriosa di materia potrebbe avere una struttura composita come la materia ordinaria, con "quarkoscuri" aggregati e tenuti insieme da un analogo della forza che permette ai normali nuclei di rimanere stabili. I componenti di guesto tipo di materia oscura, definita stealthmatter, potrebbero essere studiati in modo indiretto dal collisore Large Hadron Collider del CERN di Ginevra (red)

Cortesia Lawrence Livermore National Laboratory

28 settembre 2015

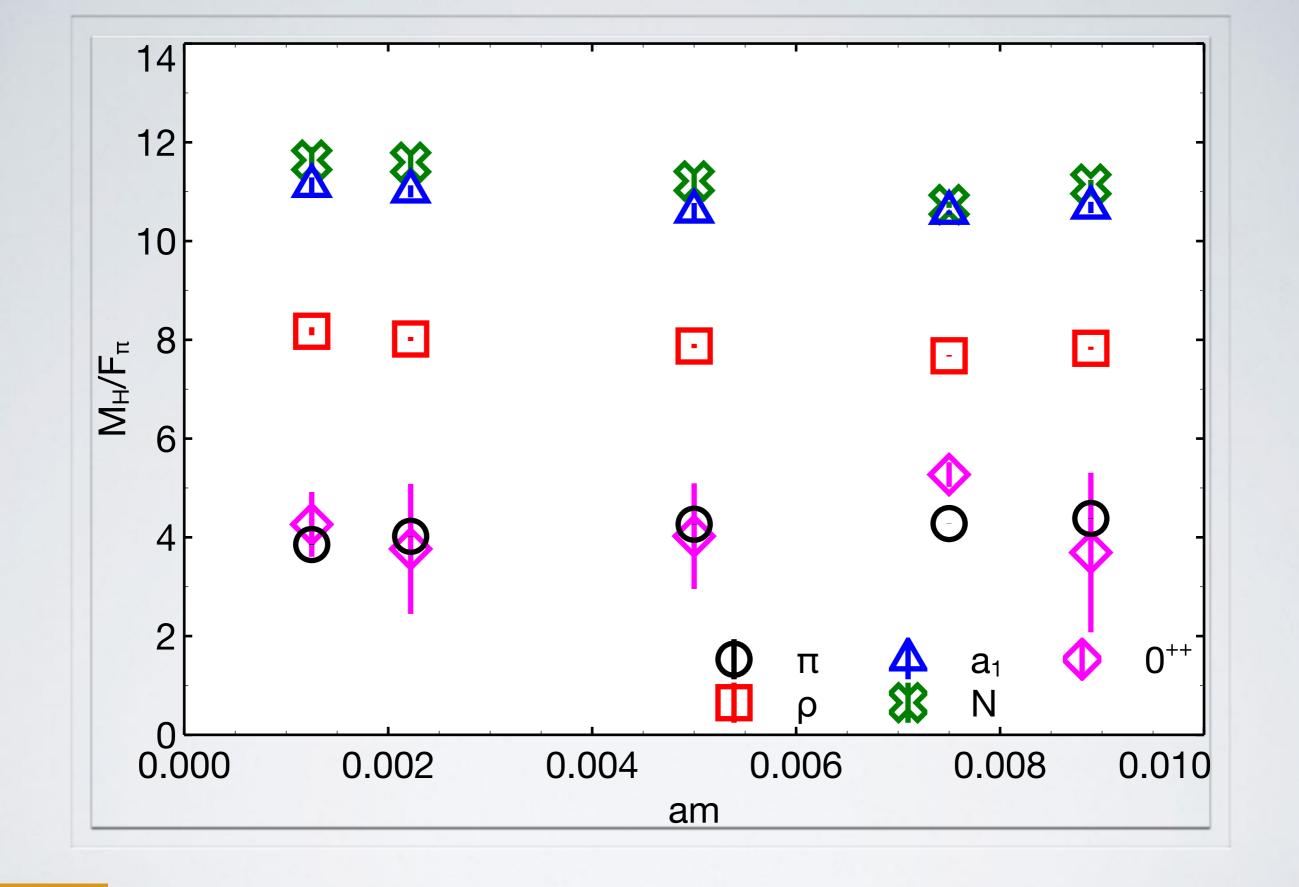


LSD arxiv: 1601.04027 [scalar update, preliminary]



Similar to QCD?

LSD arxiv: 1601.04027 [scalar update, preliminary]

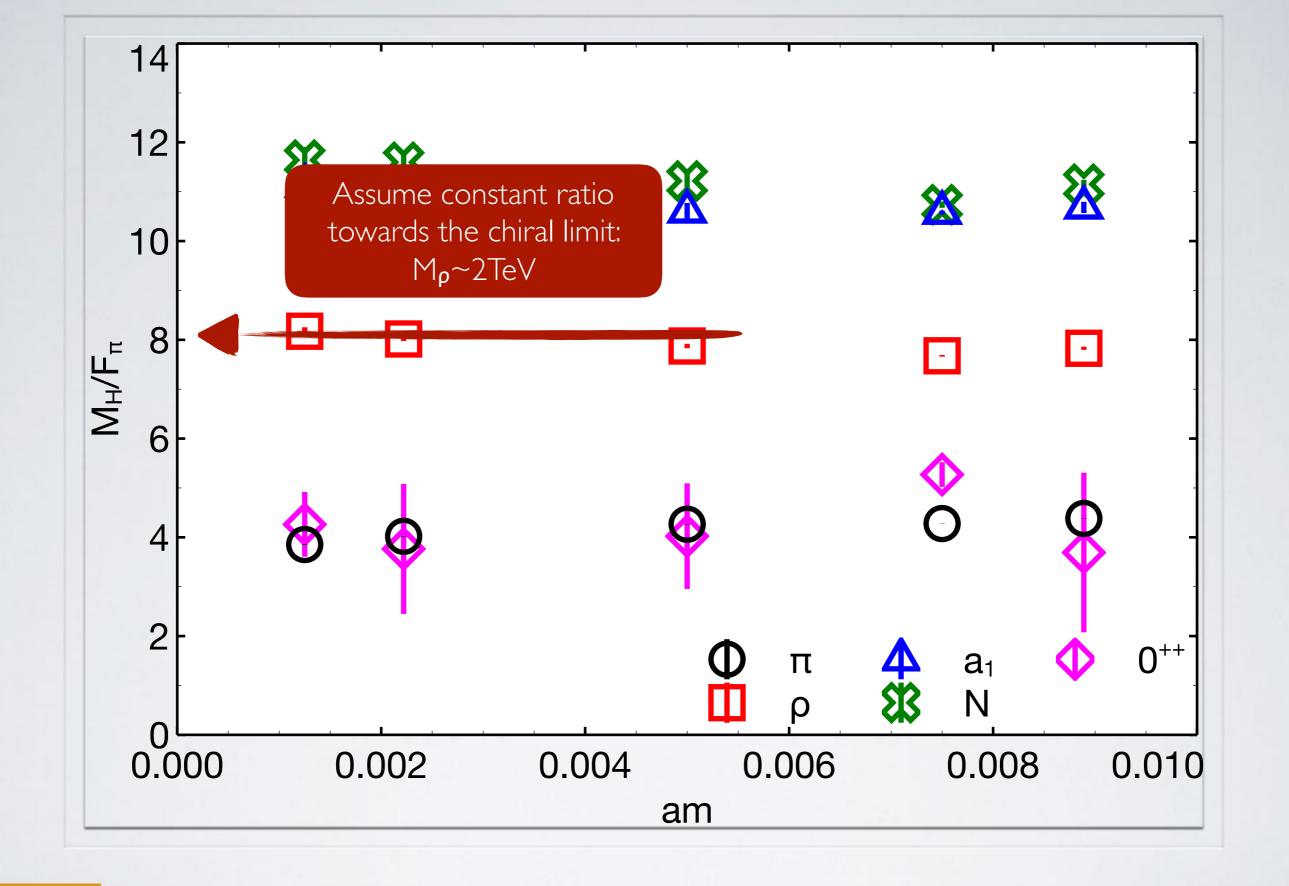


Assume Fπ~246GeV

LSD arxiv: 1601.04027 [scalar update, preliminary]

Similar to

QCD?

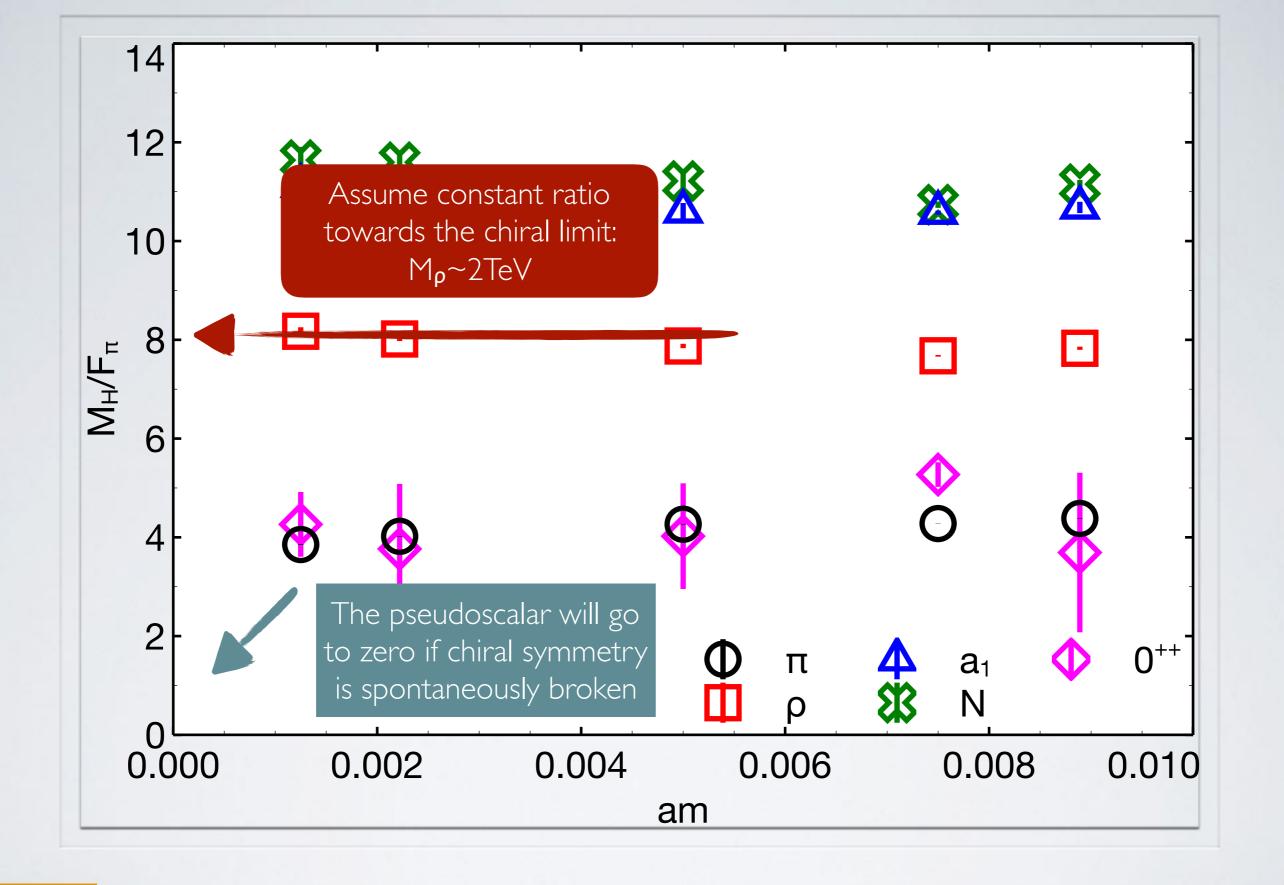


Assume Fπ∼246GeV

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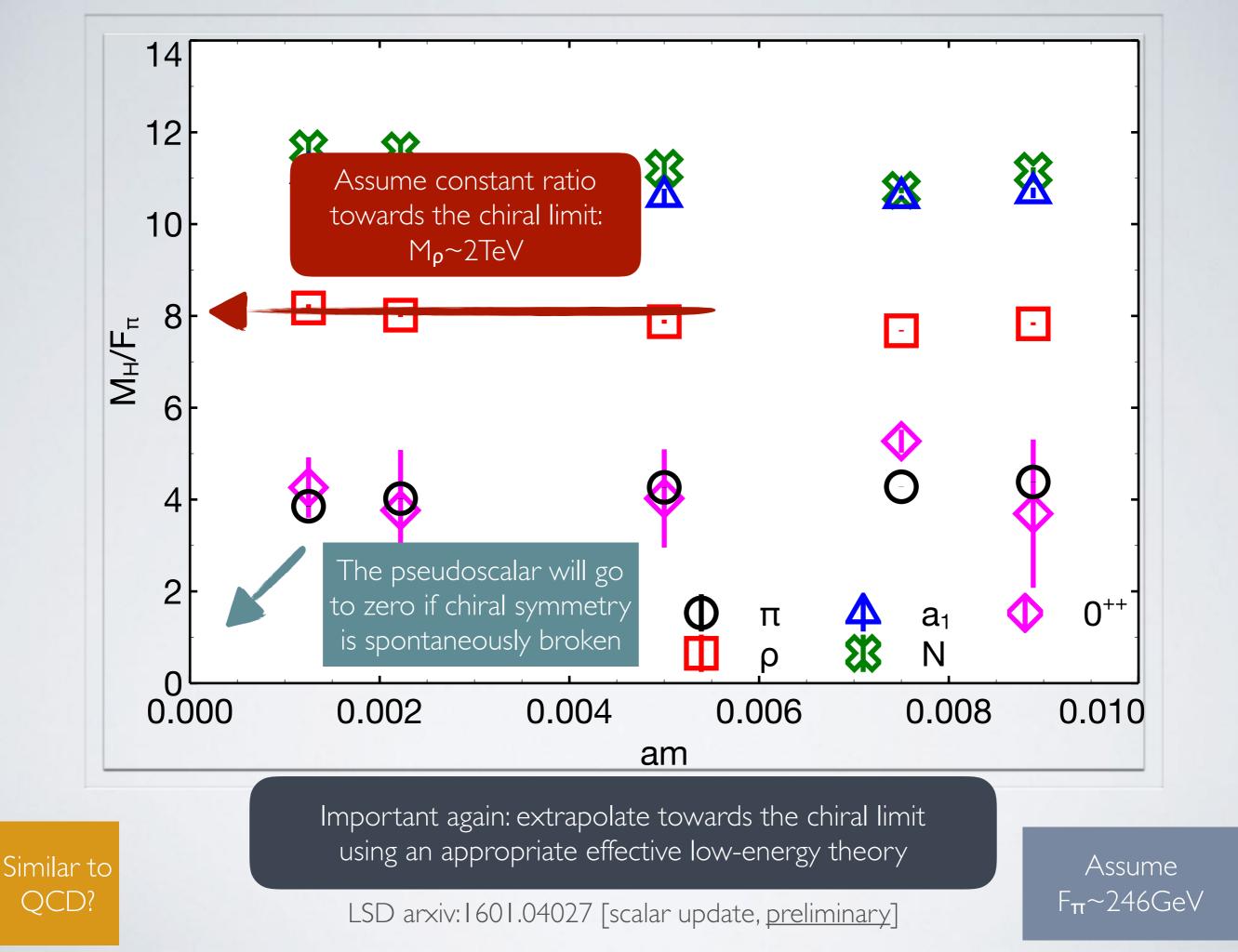


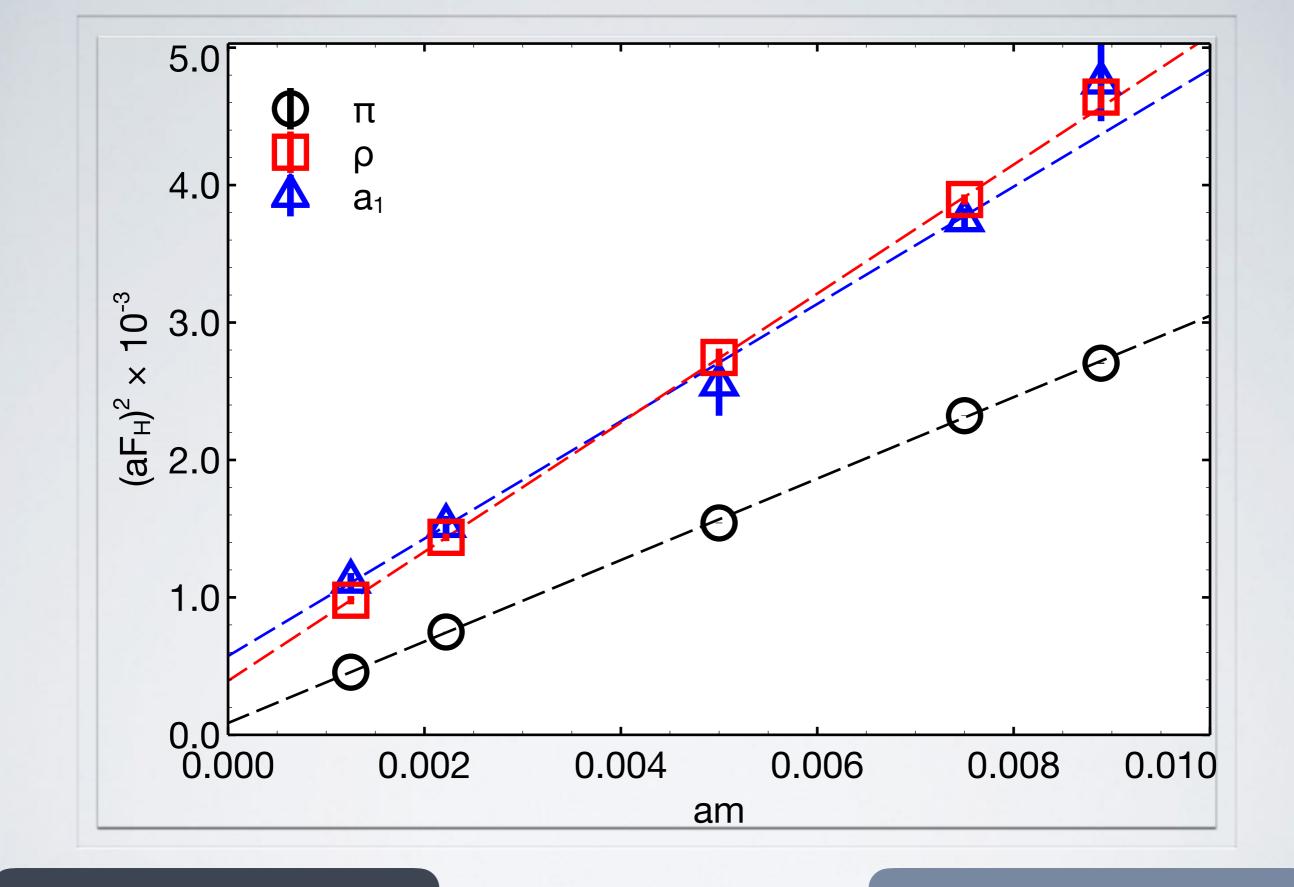
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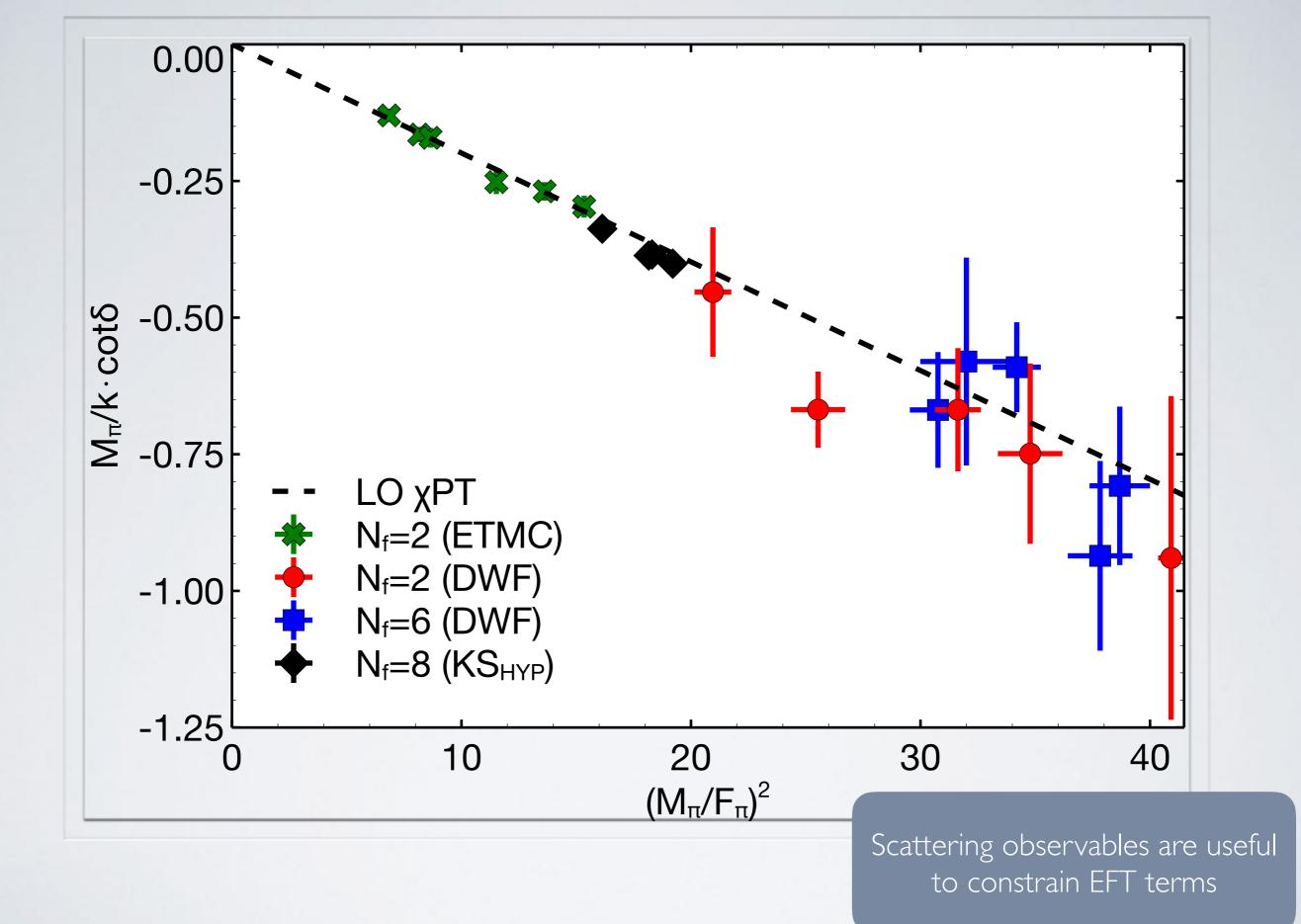




Useful to study vector meson width in the VMD picture

LSD arxiv:1601.04027

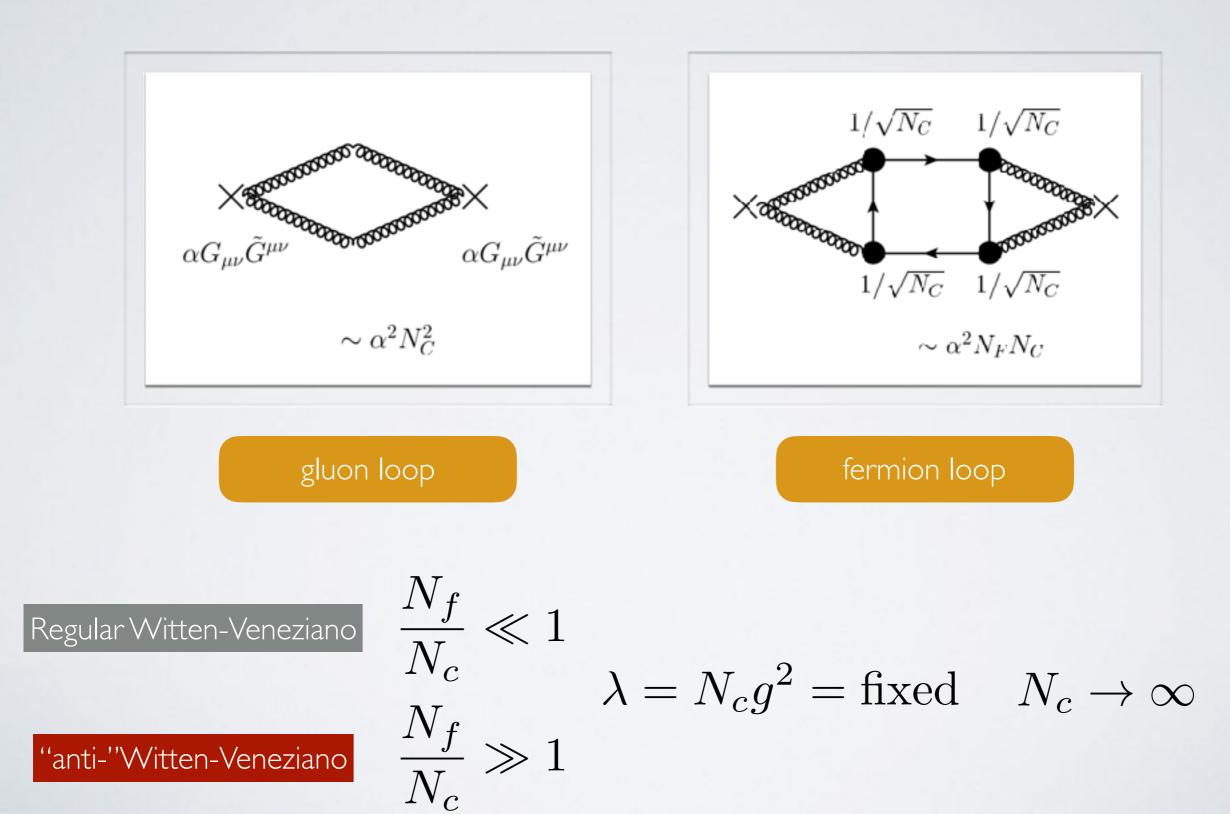
Useful to study EFT predictions [Appelquist et al. 1702.04410]



Gasbarro and Fleming (LSD collaboration) 1702.00480

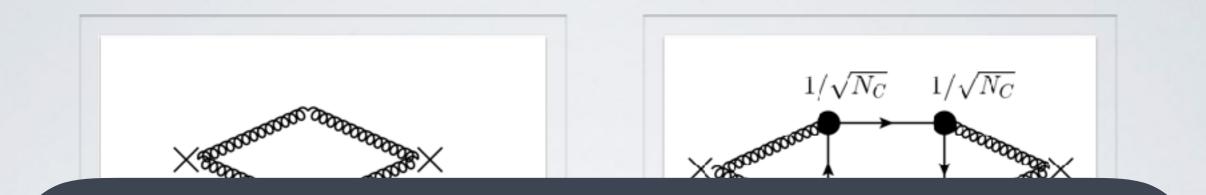
[based on ladder SD and WTI: Matsuzaki, Yamawaki arxiv: 1508.07688]

# WITTEN-VENEZIANO



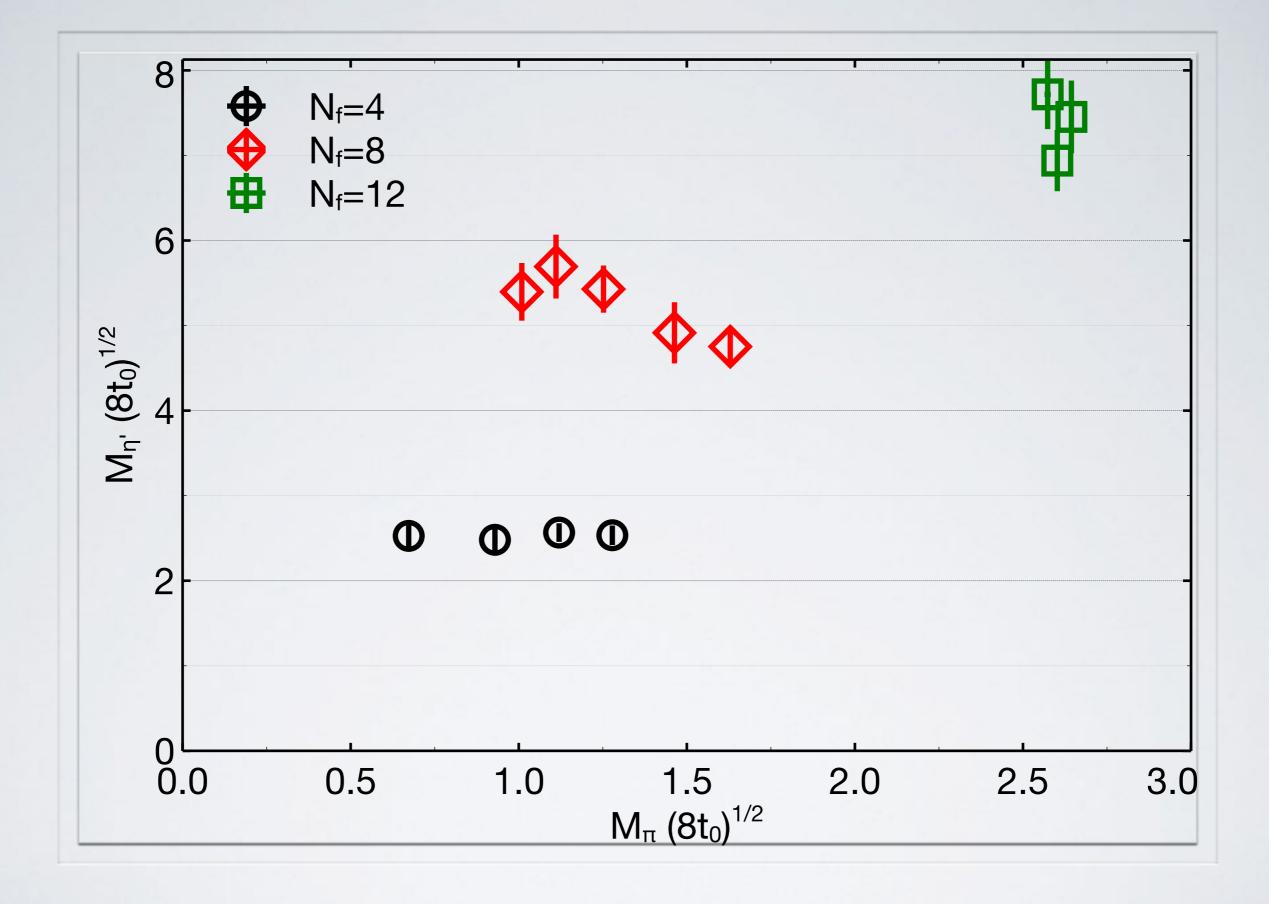
[based on ladder SD and WTI: Matsuzaki, Yamawaki arxiv: 1508.07688]

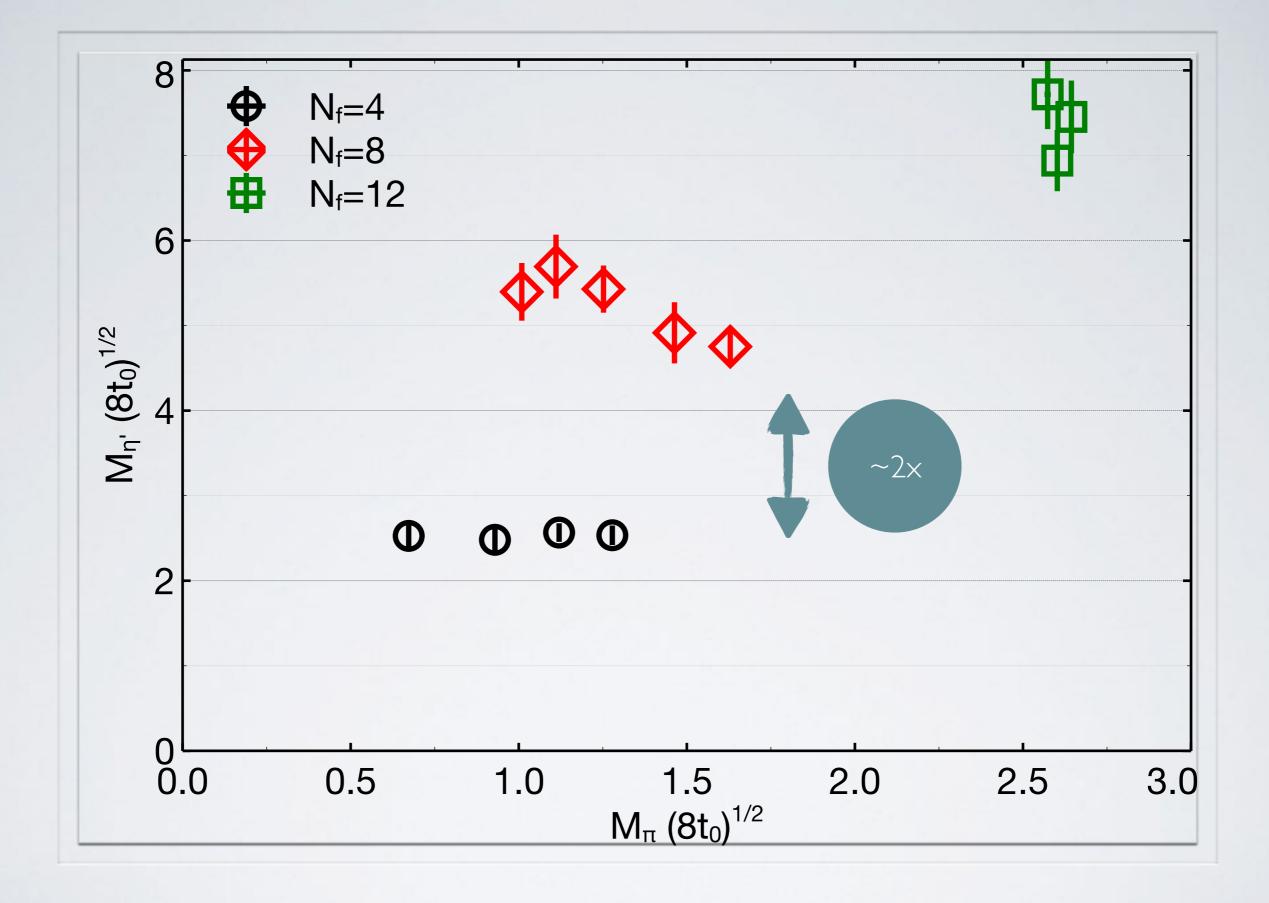
# WITTEN-VENEZIANO

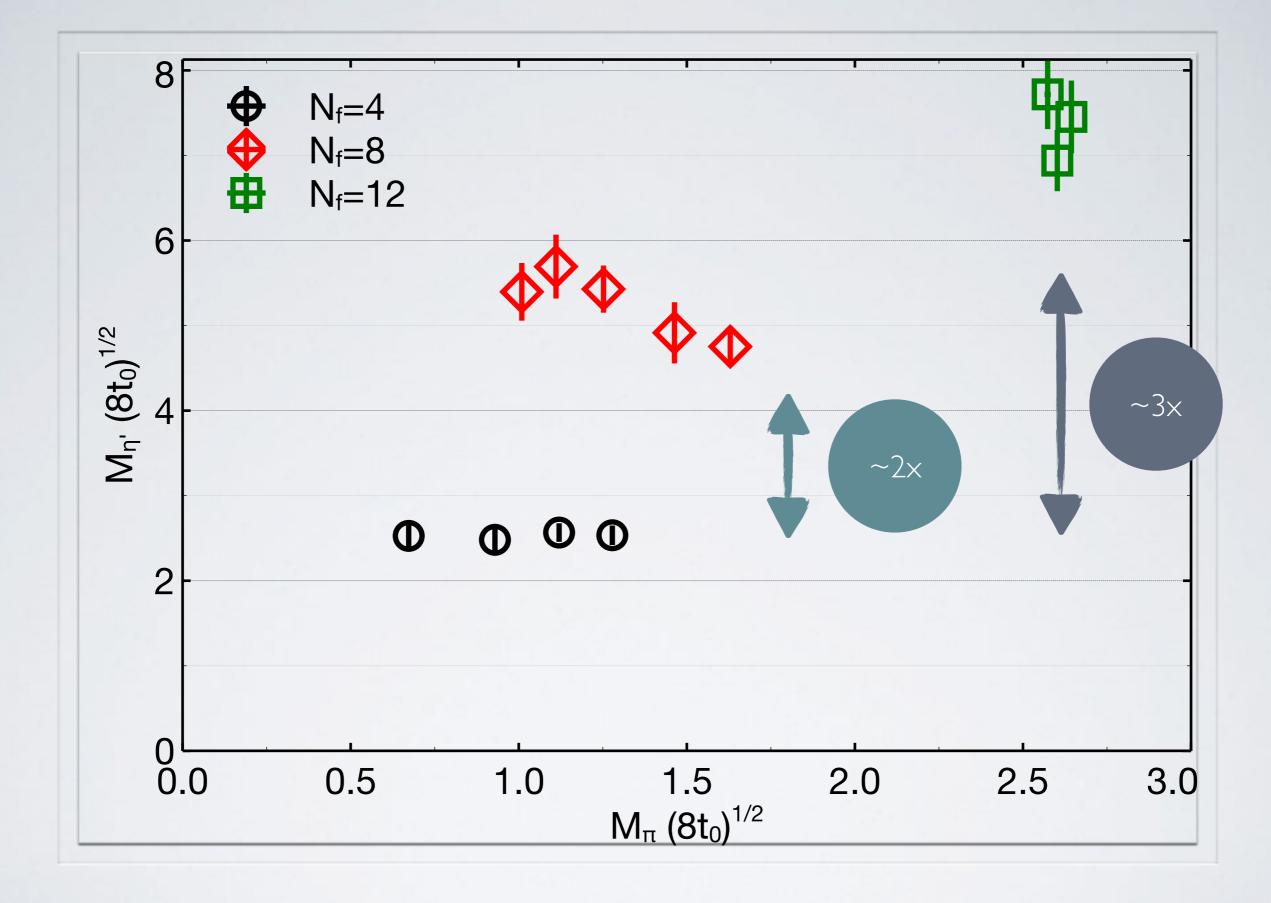


lattice results seem to align with expectations from ladder-SD analyses [arxiv:1508.07688]: a flavor-singlet scalar in a near conformal theory is light similarly to a flavor-singlet pseudoscalar in the Witten-Veneziano limit, but a flavor-singlet pseudoscalar is heavier in the ''anti''-Witten-Veneziano limit (large N<sub>f</sub>/N<sub>c</sub>)

Regular Witten-Veneziano $\frac{N_f}{N_c} \ll 1$ <br/> $\lambda = N_c g^2 = \text{fixed}$  $N_c \to \infty$ "anti-"Witten-Veneziano $\frac{N_f}{N_c} \gg 1$ 

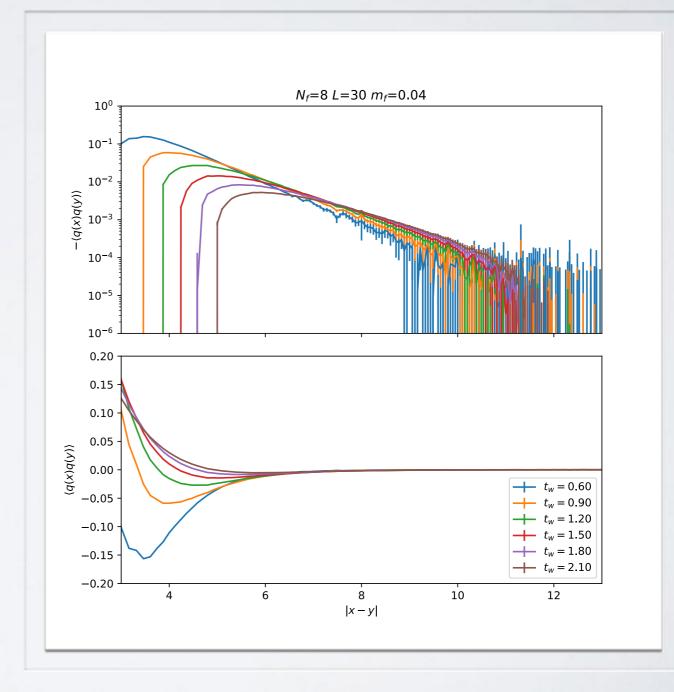






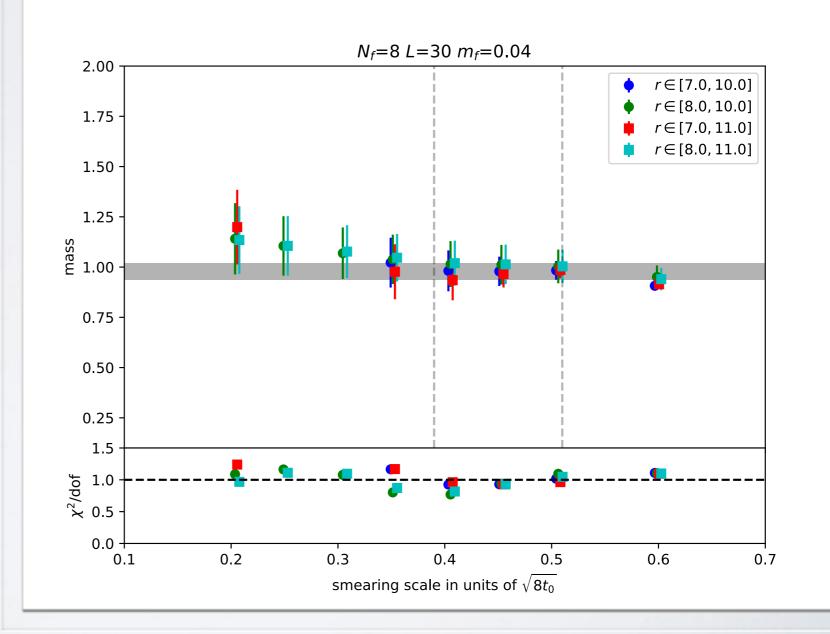
# METHODOLOGY FOR 0-+

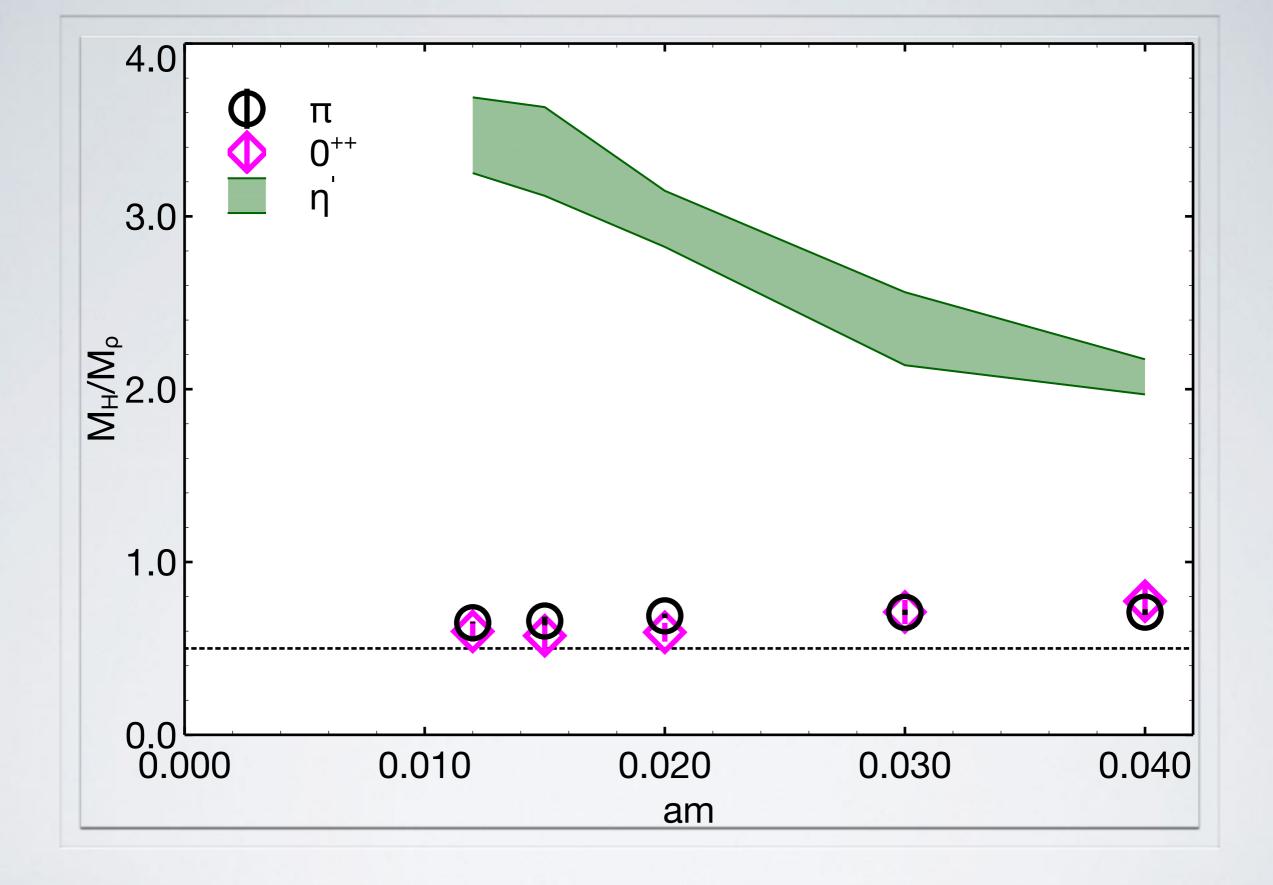
- Use a gluonic operator with 0 quantum numbers: topological charge density
- Use Wilson flow smearing as a technique to ameliorate the signal-to-noise problem
- Entirely similar to previous studies in SU(3) YM [arxiv: 1409.6459] and QCD [arxiv: 1509.00944]
- Main difficulty is estimating the systematics due to smearing and excited states

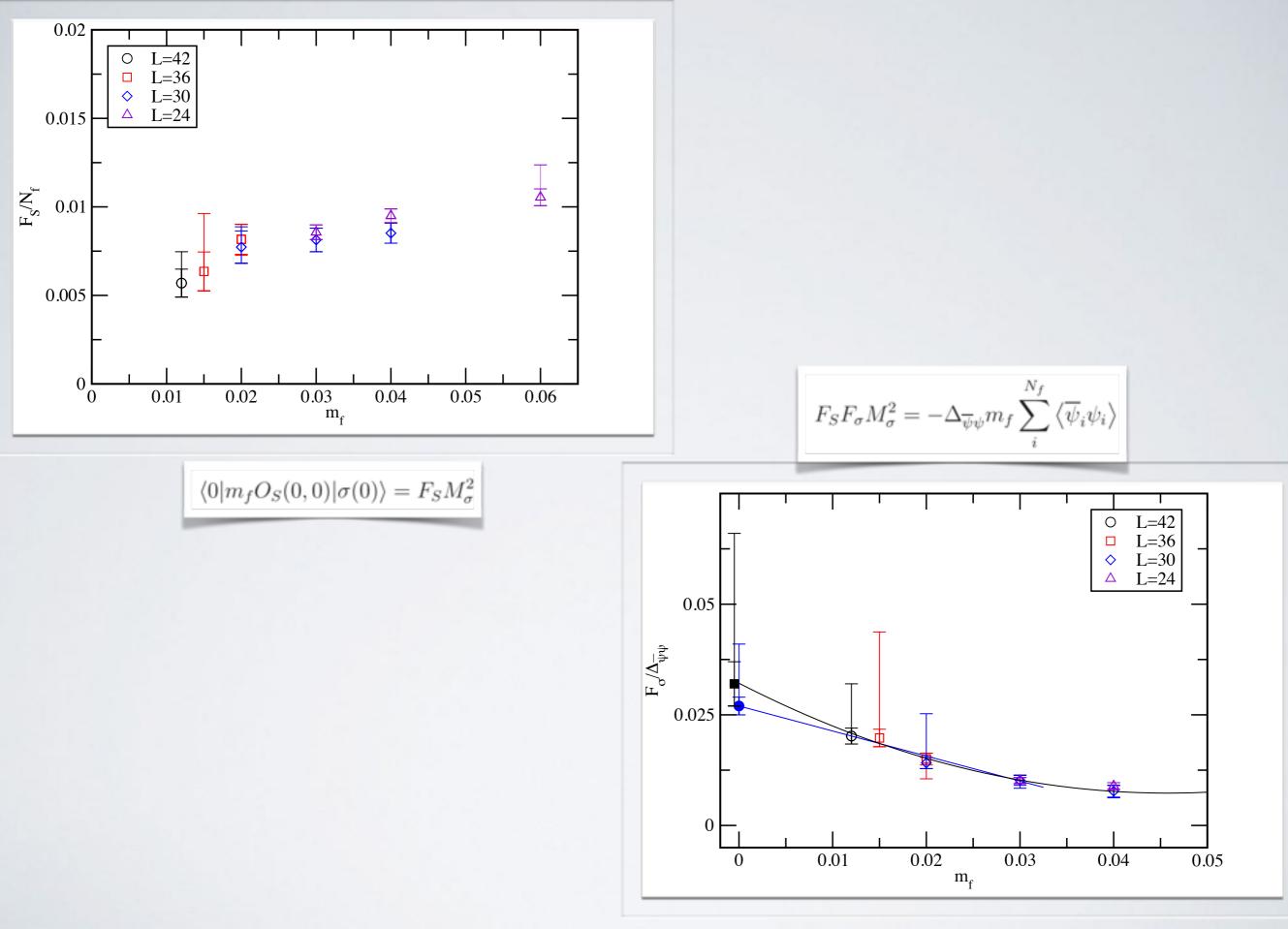


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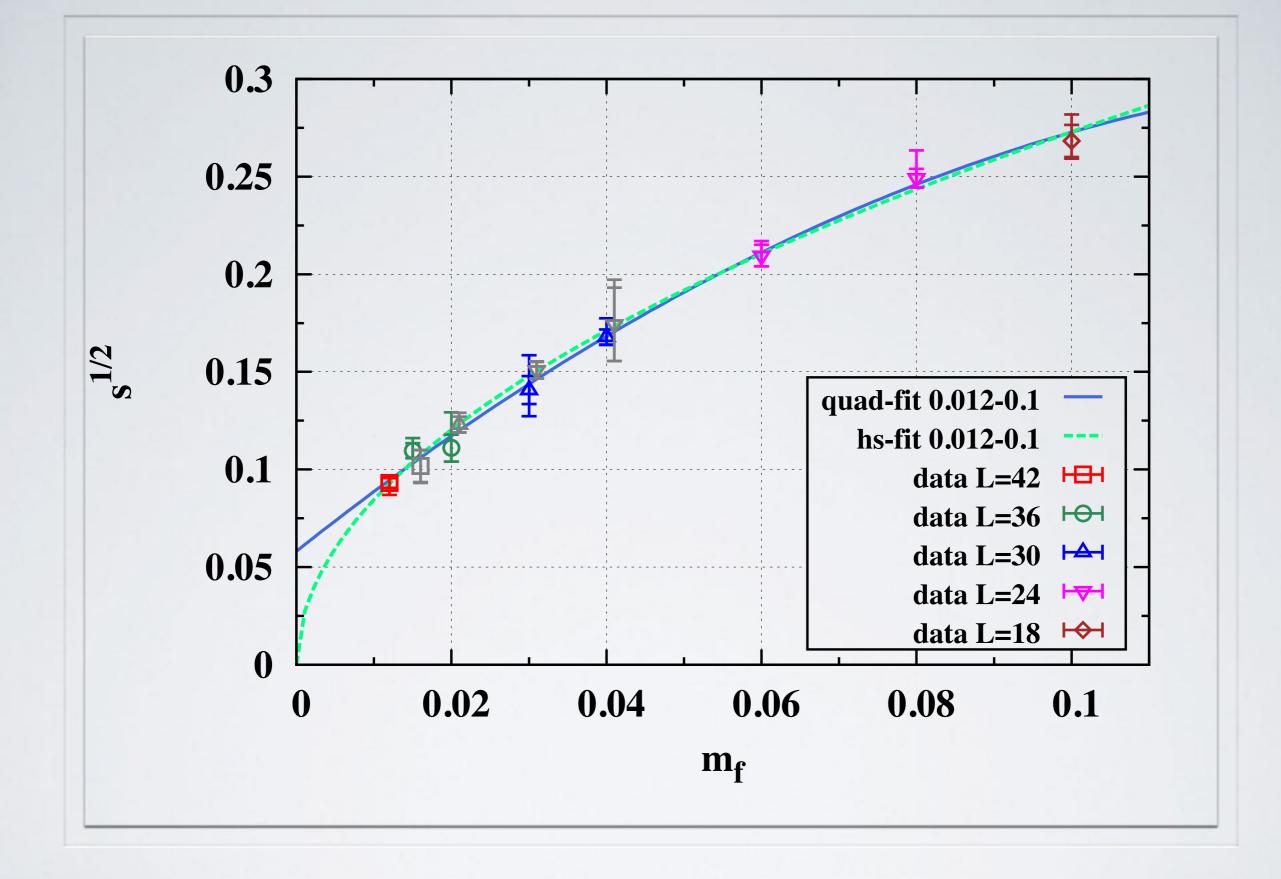
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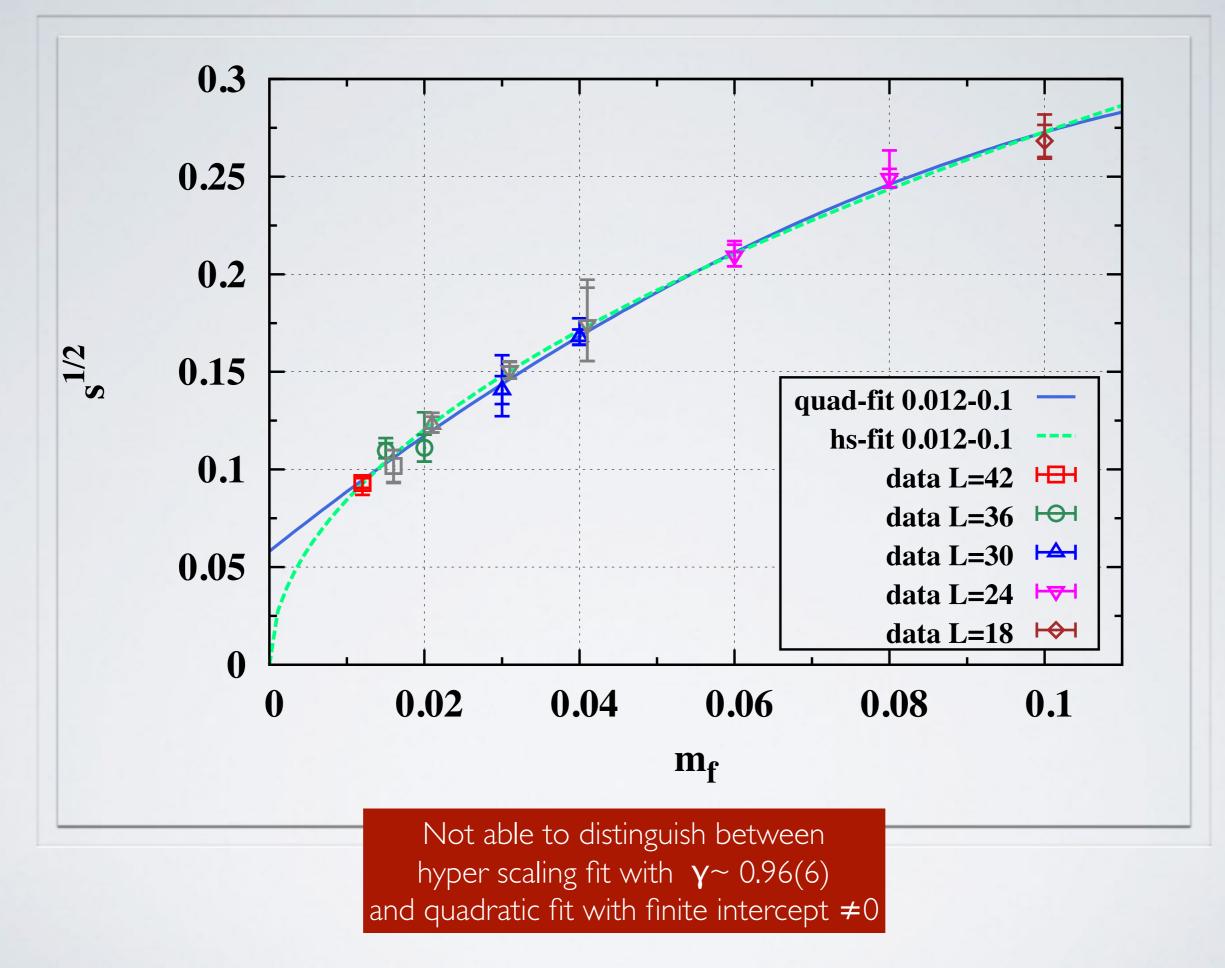




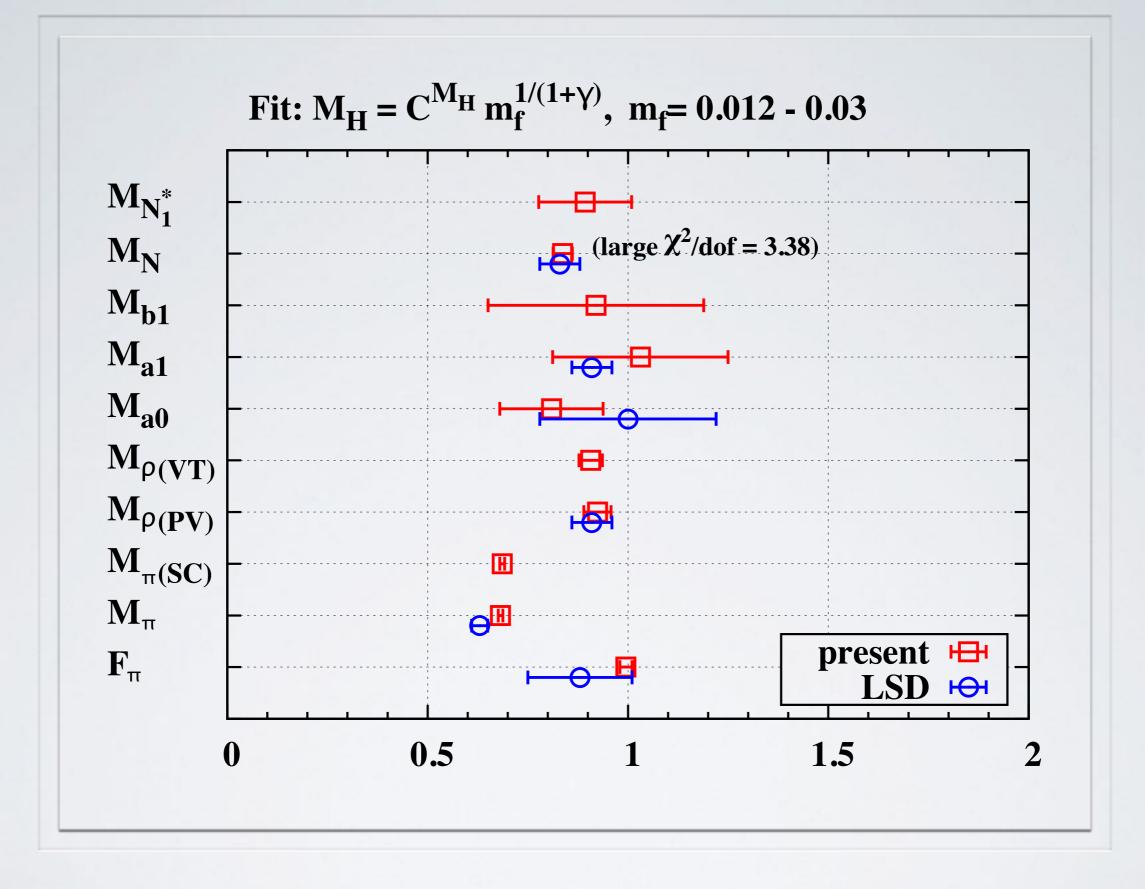


LatKMI arxiv:1610.07011

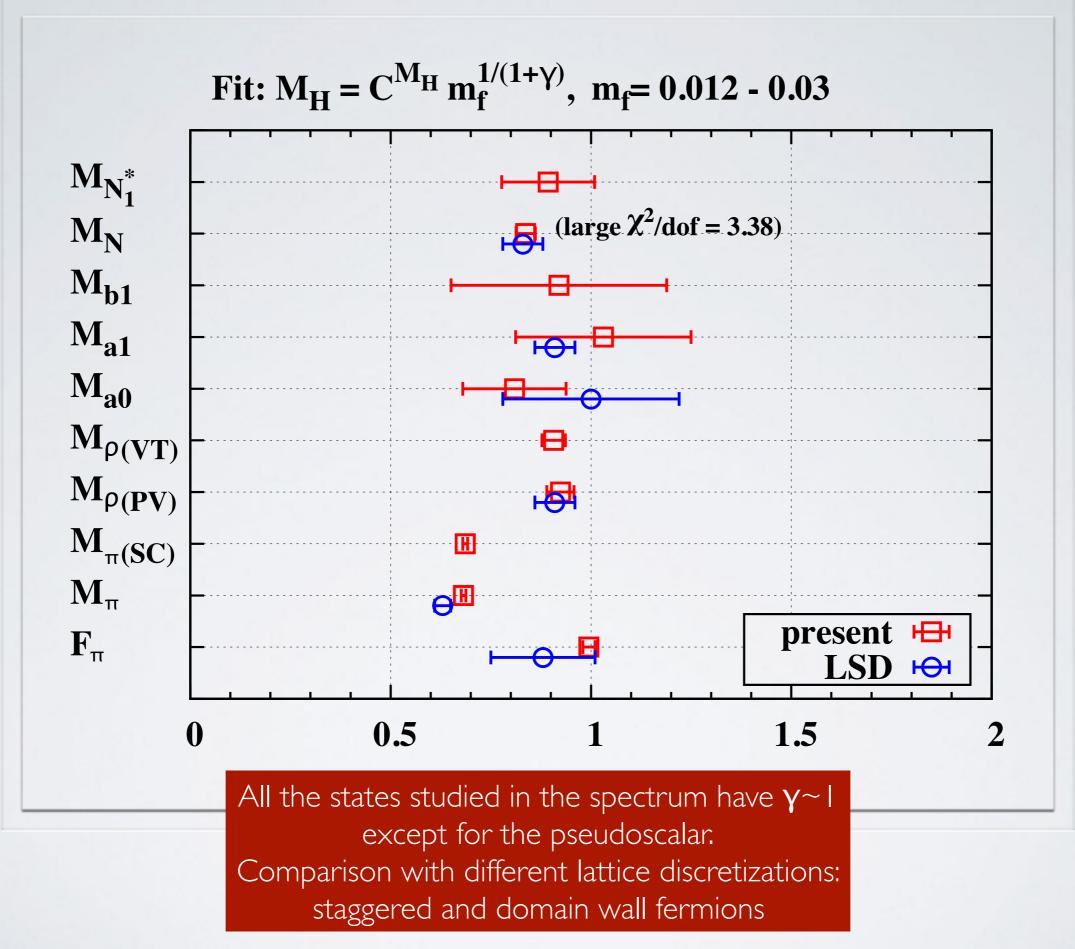




LatKMI arxiv:1610.07011



LatKMI arxiv:1610.07011 | LSD arxiv:1405.4752



LatKMI arxiv:1610.07011 | LSD arxiv:1405.4752