

EICUG 2017



eSTARlight

An event generator for the EIC

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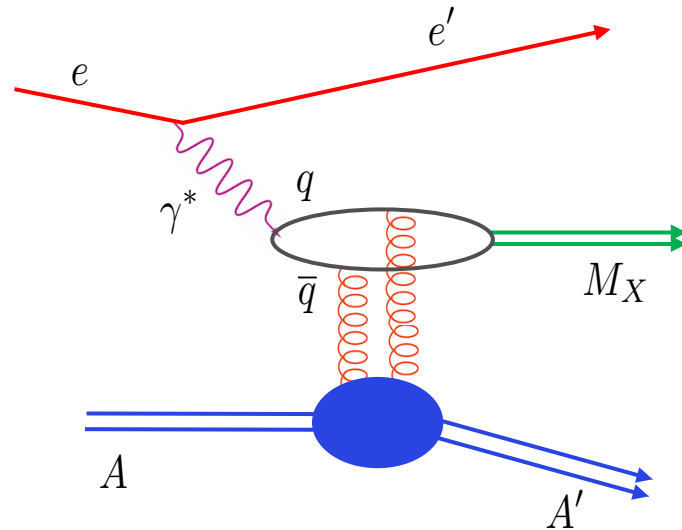
EIC UG 2017, Trieste Italy

Outline

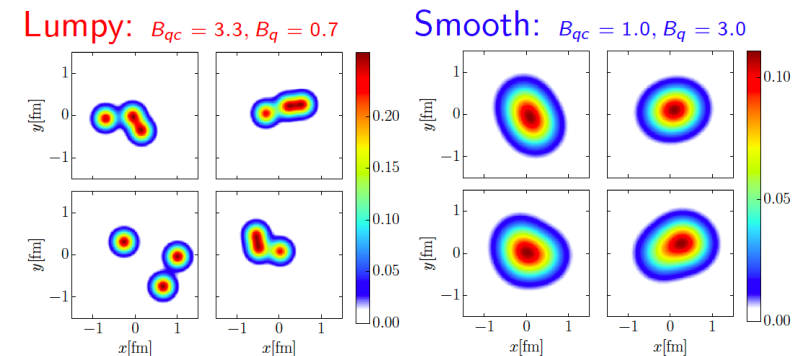
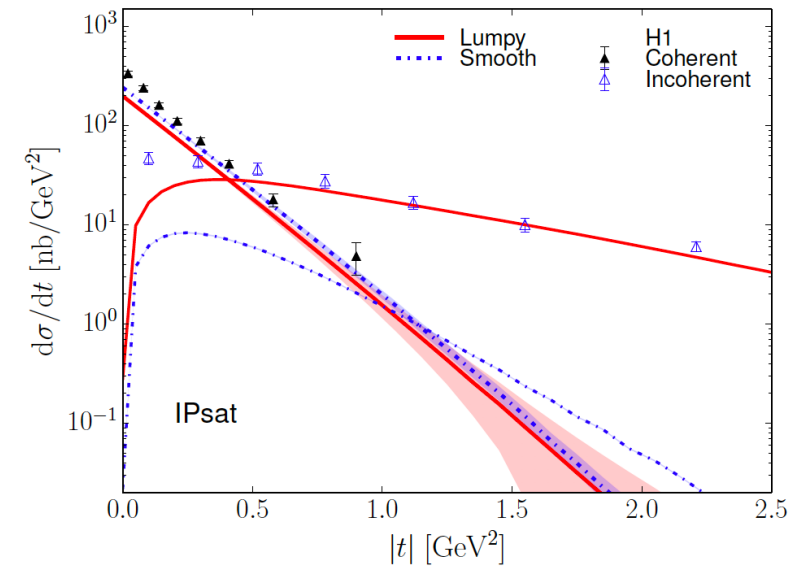
- Motivation
- STARlight: our predecessor
- Photo-nuclear interaction in eSTARlight:
 - $\gamma p \rightarrow V.M. + p$ vector meson production
 - Comparison to data
 - Extension to γA
- Vector meson production at an EIC
- Estimating rates
- Caveats and outlook
- Summary

Diffraction at an EIC

- How are the quarks and gluons distributed within the nucleon? What about nucleus?
- Initial state geometry is necessary input for hydro:
 - Initial state (IS) geometry \rightarrow final state collectivity
 - Collective phenomena has been observed in p-p and p-A collisions



- Diffractive processes (no color exchange) can probe:
 - Target remains intact: *Coherent* (small $|t|$) probes average gluon density and their spatial distribution
 - Target break up: *Incoherent* (large $|t|$) is sensitive to fluctuations in IS



H. Mäntysaari, QM17 parallel session 5.2

Electron-ion collisions

Some important kinematic quantities to consider in electron-ion collisions:

- Resolving power

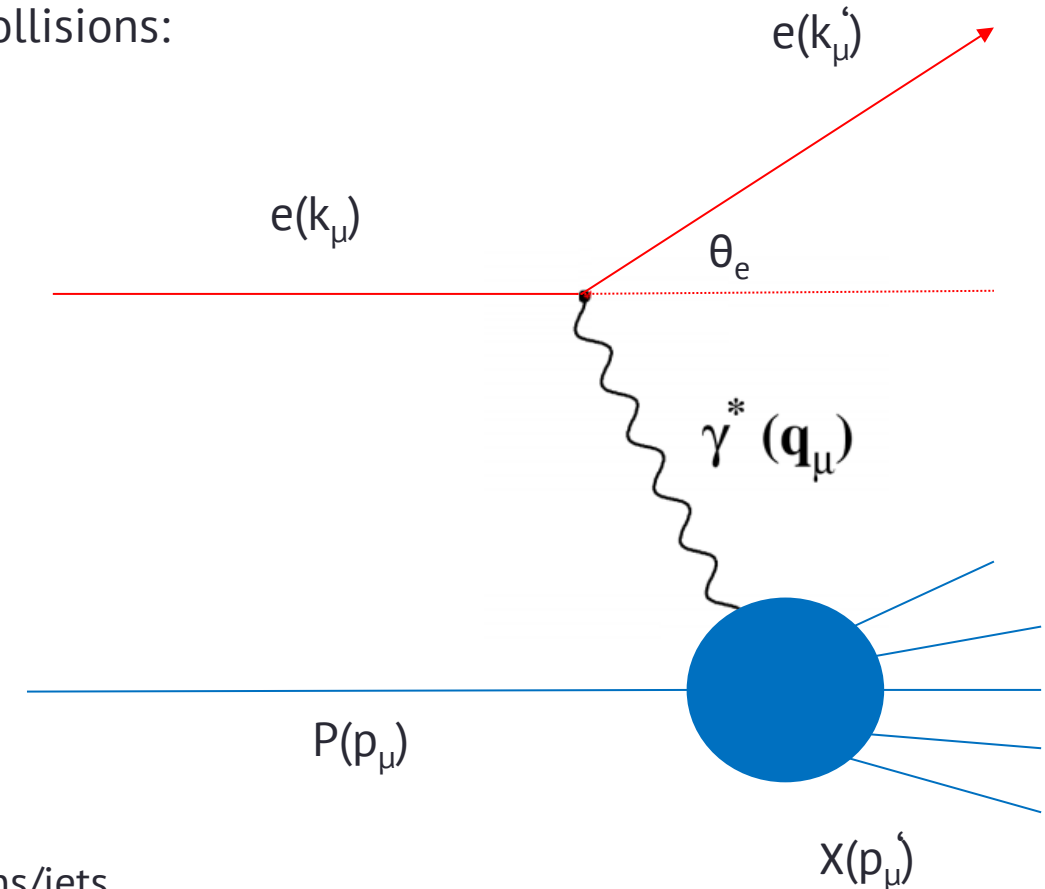
$$Q^2 = -q^2 = -(k_\mu - k'_\mu)^2$$

$$= 2E_e E'_e (1 - \cos \theta_e)$$
- Momentum of struck quark

$$x = \frac{Q^2}{2pq} = \frac{Q^2}{sy}$$
- Measure of inelasticity

$$y = \frac{pq}{pk} = 1 - \frac{E'_e}{E_e} \cos^2 \left(\frac{\theta_e}{2} \right)$$

$$s = 4E_p E_e$$



Inclusive: Detect scattered lepton.

$$e + p/A \rightarrow e' + X$$

Semi-inclusive: Detect scattered lepton in coincidence with identified hadrons/jets.

$$e + p/A \rightarrow e' + h + X$$

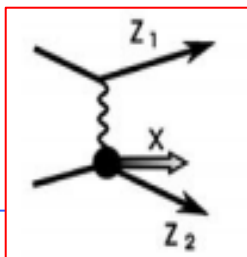
Exclusive: Detect scattered lepton, id'd hadrons/jets and target fragments.

Our predecessor - STARlight

Widely used M.C. in ultra-peripheral collisions (UPCs). Choice of channels mostly driven by experimental accessibility

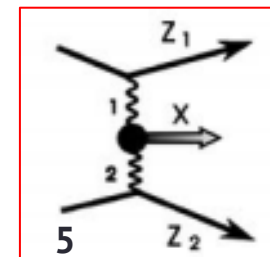
Photonuclear: Photon-Pomeron (γP) final states

- Light vector mesons: $\rho, \omega, \phi, \rho' \rightarrow \pi\pi\pi\pi$,
- Heavy vector mesons (VM): $J/\psi, \psi', \Upsilon(1S), \Upsilon(2S), \Upsilon(3S)$
- VM follows photon polarization (along beam)
 - Correct angular distributions
- p_T distributions include (optionally) 2-site interference
- General photonuclear interactions, via DPMJET interface
- Direct $\pi^+\pi^-$



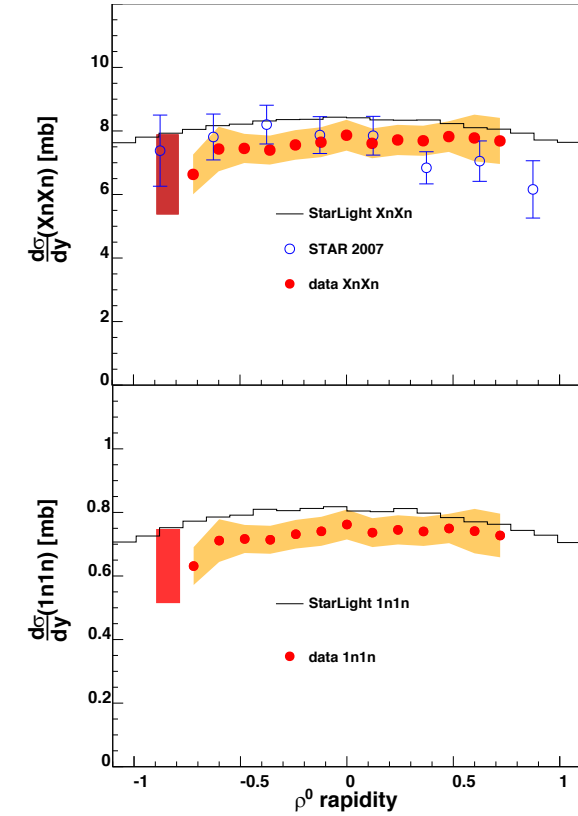
$\gamma\gamma$ final states

- Lepton pairs $e^+e^-, \mu^+\mu^-, \tau^+\tau^-$
- Single mesons: $\eta, \eta', f_0(980), f_2(1270), f_2'(1525), \eta_c$
 - “Simple” decays in STARlight ensure correct angular distribution
 - “Complex” final states decayed via PYTHIA 8
- $\rho^0\rho^0$
- Heavy axions

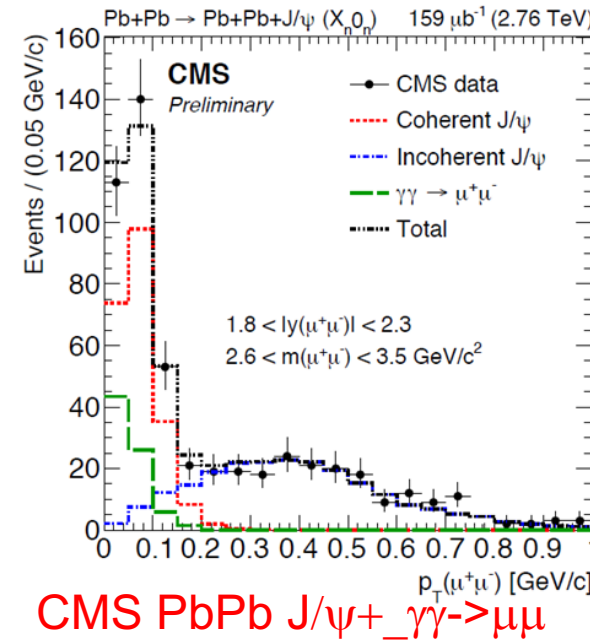
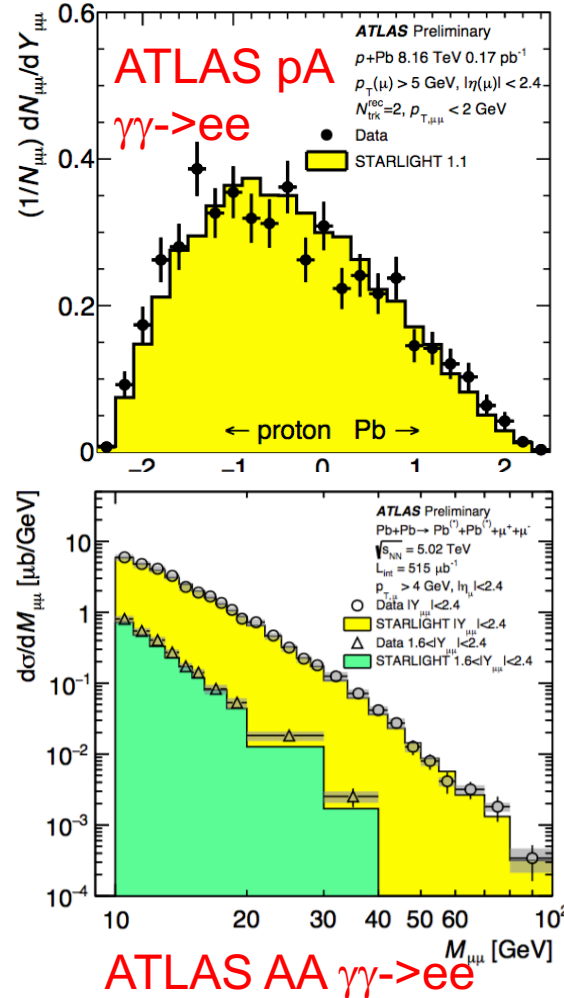


STARlight versus data

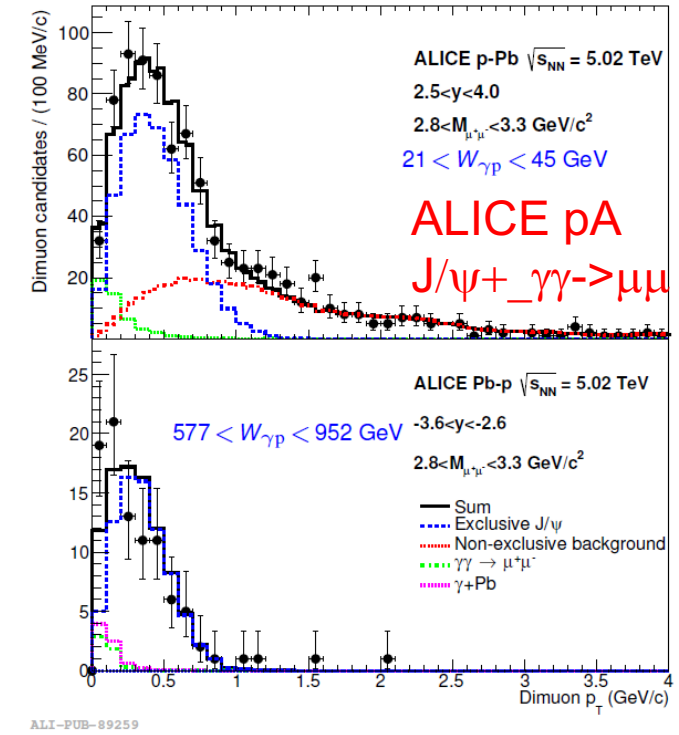
STAR
AuAu $\rightarrow \rho^0$



CMS - arXiv:1605.06966
ALICE - PRL113 232504 (2014)
ALICE - QM17, M. Dyndal
STAR - DIS16



CMS PbPb J/ψ+ $\gamma\gamma\rightarrow\mu\mu$



- Good track record proven by many experiments and collision energies in UPC collision
- Confidence in description of nucleus targets: helpful in eSTARlight given the lack of eA data

Electro-nuclear interactions

$$\sigma(e + X \rightarrow e + X + V.M.) = \int dQ^2 \int dE_\gamma \boxed{\frac{dN_\gamma(E_\gamma, Q^2)}{dE_\gamma dQ^2}} \boxed{\sigma_{\gamma X}(W, Q^2)}$$

Photon flux Photonuclear cross section

$$= \int dQ^2 \int dE_\gamma \boxed{\frac{dN_\gamma(E_\gamma, Q^2)}{dE_\gamma dQ^2}} \boxed{\left(\frac{1}{1 + Q^2/M_v^2} \right)^n \sigma_{\gamma X}(W)}$$

• Differential x-section sampled to produce events:

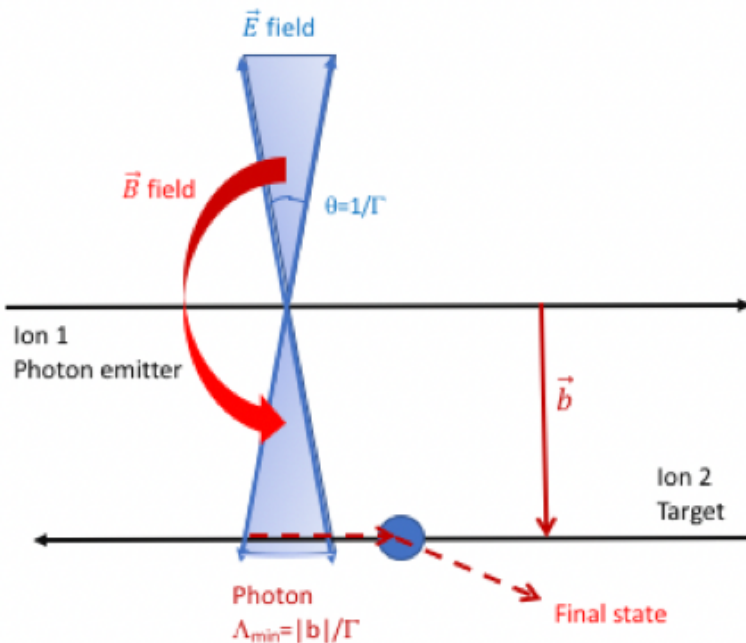
- Rejection sampling used to draw from distributions
- Two look up tables used to speed up event generation:
 1. Draw E_γ and Q^2 pair from $f(E_\gamma, Q^2)$ to construct e- γ vertex
 2. Sample $\sigma_{\gamma X}(W)$ with input photon and desired V.M. state

Effective Photon flux
 $f(E_\gamma, Q^2)$

Reduced photonuclear
cross section

Photon flux

- Using equivalent photon approach (EPA):
 - Electric field radially outwards from e and Lorentz contracted in target (p,A) frame => Looks like a pancake in transverse direction
 - Magnetic field circles around beam axis
- Perpendicular fields resemble EM radiation: treat each source as dressed by photons



- Include necessary corrections to account for the virtuality of the exchanged photon¹:

$$n = \frac{d^2 N(E_\gamma, Q^2)}{d(Q^2) dE_\gamma} = \frac{\alpha}{\pi} \frac{1}{E_\gamma |Q^2|} \left[1 - \frac{E_\gamma}{E_e} + \frac{1}{2} \left(\frac{E_\gamma}{E_e} \right)^2 - \left(1 - \frac{E_\gamma}{E_e} \right) \left| \frac{Q_{min}^2}{Q^2} \right| \right]$$

- And the available q^2 range:

$$Q_{min}^2 = \frac{m_e^2 E_\gamma^2}{E_e(E_e - E_\gamma)} \leq Q^2 \leq 4E_e(E_e - E_\gamma)$$

1: Phys.Rept. 15 181-281 (1975)

Photonuclear Cross Section

$$\sigma_{\gamma p}(W, Q^2) = \left(\frac{1}{1 + Q^2/M_v^2} \right)^n \sigma_{\gamma p}(W)$$

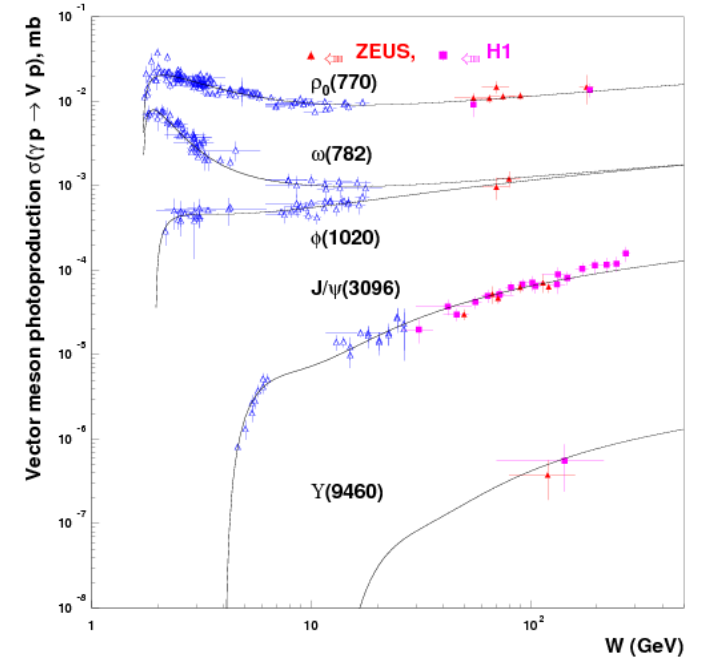
- Interactions are done, mostly, with parameterization from HERA¹ for $\gamma p \rightarrow V p$ in terms of the γp center of mass energy $W_{\gamma p}$.

$$\sigma_{\gamma p}(W) = \underbrace{\sigma_P \cdot W^\epsilon}_{\text{Pomeron exchange}} + \underbrace{\sigma_M \cdot W^\eta}_{\text{Meson exchange}}$$

- Only ρ and ω can be produced via meson exchange ($\sigma_M > 0$)
- For the J/ψ , ψ' and Y states the power law is supplemented with factor to account for the near-threshold decrease in cross section

$$\sigma_{\gamma p}(W) = \sigma_P \cdot \left[1 - \frac{(m_p + m_{VM})^2}{W^2} \right]^2 \cdot W^\epsilon$$

- The n is also obtained from fits to data²

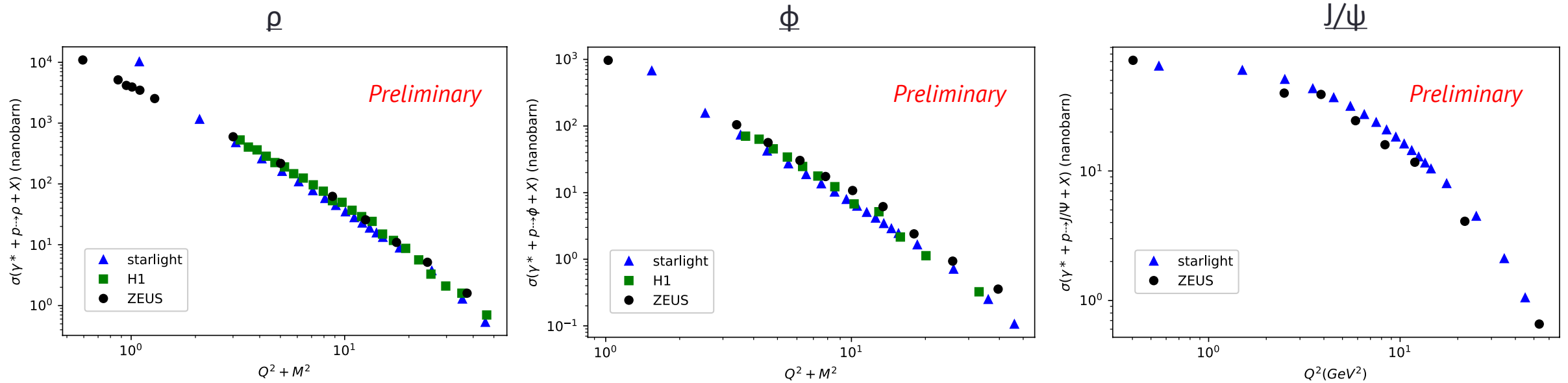


Vector Meson	σ_P	ϵ	σ_M	η
ρ^0 & ρ'	$5.0 \mu\text{b}$	0.22	$26.0 \mu\text{b}$	1.23
ω	$0.55 \mu\text{b}$	0.22	$18.0 \mu\text{b}$	1.92
ϕ	$0.34 \mu\text{b}$	0.22	—	—
J/ψ	4.06 nb	0.65	—	—
$\psi(2S)$	0.674 nb	0.65	—	—
$Y(1S)$	6.4 pb	0.74	—	—
$Y(2S)$	2.9 pb	0.74	—	—
$Y(2S)$	2.1 pb	0.74	—	—

¹: S.R.Klein, *et. al.*, Computer Physics Communications **212** 0010-4655 (2017)

²: Phys.Lett.B377:259-272(1996), JHEP 1005:032(2010), Nucl.Phys.B695:3-37(2004)

HERA comparison to data: $\sigma(\gamma^* + p \rightarrow \text{V.M.} + p)$



- Gamma-proton cross-sections obtained following same procedure in experiment:

$$\sigma_{\gamma p} = \frac{\int dE_\gamma \int dQ^2 \frac{d^2 N}{dE_\gamma d(Q^2)} \sigma_{\gamma p}(E_\gamma, Q^2)}{\int dE_\gamma \int dQ^2 \frac{d^2 N}{dE_\gamma d(Q^2)}}$$

- $\sigma_{\gamma p}$ measured at HERA is well described by eSTARlight over a broad Q^2 range

ϕ : Phys.Lett.B377:259-272(1996)

ρ : JHEP 1005:032(2010)

J/ψ : Nucl.Phys.B695:3-37(2004)

Coherent Photonuclear Cross Section $\sigma(\gamma A \rightarrow VA)$

- Extrapolate photonuclear cross section from γp to γA using Quantum Glauber calculation to take into account the nuclear form factor:

$$\sigma_{tot}(VA) = \int d^2b \left[2 \cdot \left(1 - e^{-\sigma_{tot}(Vp)T_{AA}(b)/2} \right) \right]$$

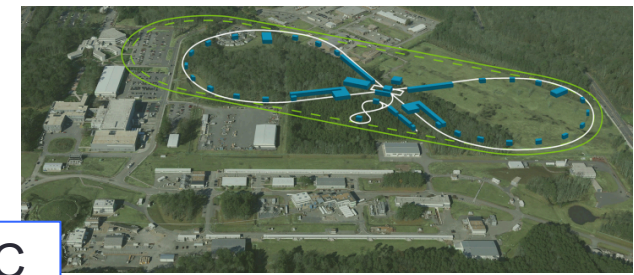
- Generalized vector dominance model and optical theorem used to obtain the photo-nuclear cross section

$$\sigma(\gamma A \rightarrow VA) = \frac{d\sigma(\gamma A \rightarrow VA)}{dt} \bigg|_{t=0} \int_{t_{min}}^{\infty} dt |F(t)|^2$$

- There are two modes in eSTARlight to calculate the cross sections:
 - Assuming a narrow resonance for the vector mesons, in which case: $t_{min} = (M_V^2/4k\gamma)^2$
 - Convoluting the spectrum with a Breit-Wigner shape. The difference can be substantial (~ 5% reduction in cross section ρ^0 in heavy-ion collisions)

EIC prediction: $\sigma(e+X \rightarrow e+V.M.+X)$

- eSTARlight can be used to estimate $e+X \rightarrow e+X+V.M.$ cross sections for:
 - Different center of mass energies (accelerator facilities)
 - Different V.M. species
 - Different collision systems ($X = p, Au$, etc.)
 - Arbitrary virtuality Q^2
- Heavy-ions provide considerable production (larger CM energy) at the cost of Q^2 due to kinematic reach
- Still have considerable production $Q^2 > 1 \text{ GeV} \rightarrow$ electrons sufficiently scattered away from beam for exclusive measurements



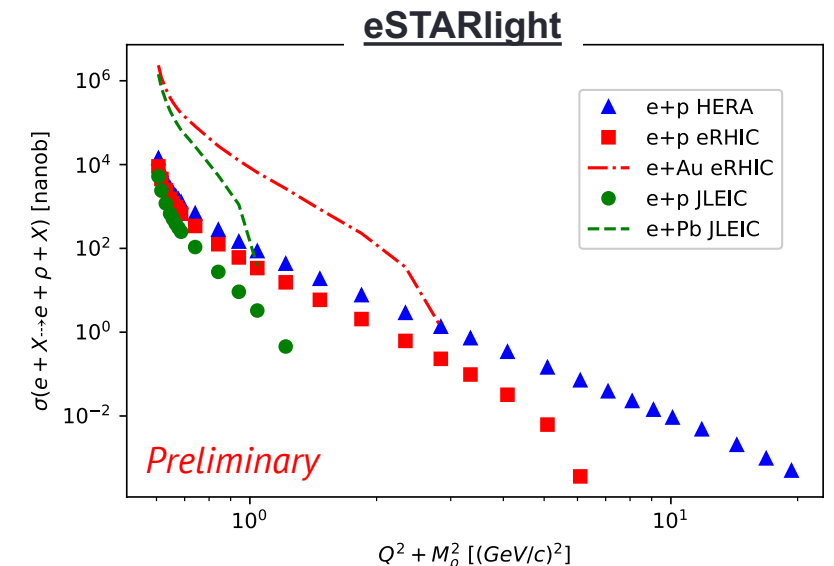
JLEIC



eRHIC

	JLEIC (JLAB)	eRHIC (BNL)	HERA (DESY)
electron e [GeV]	5	10	27.5
proton p [GeV]	60	250	920
Heavy-Ion A [GeV/n]	$^{208}\text{Pb}^{82+}$	$^{197}\text{Au}^{79+}$	n/a
	40	100	n/a

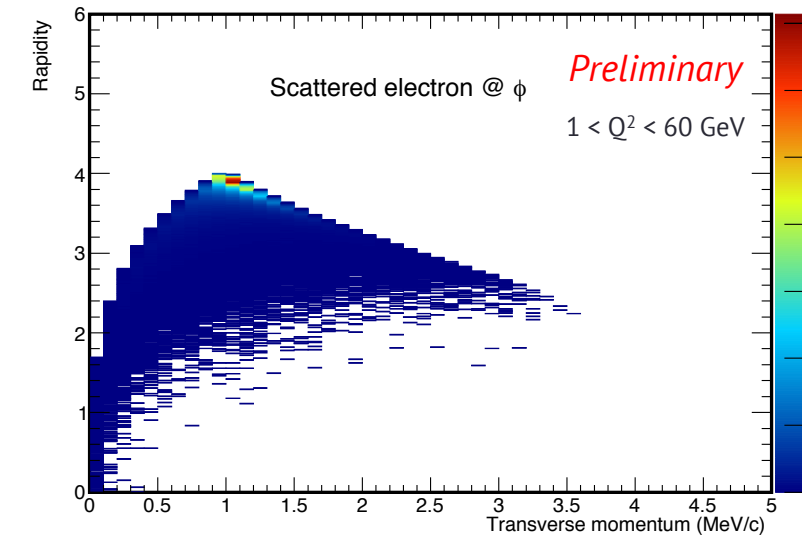
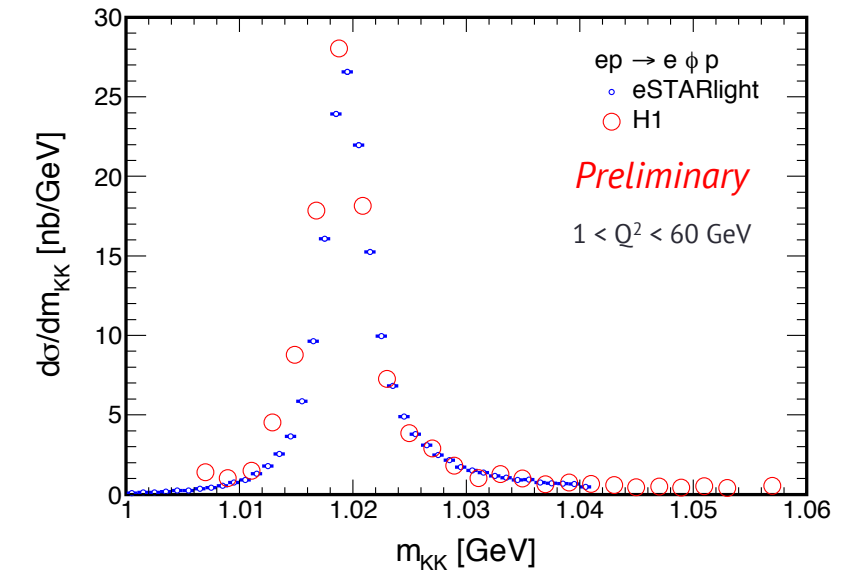
Table 1: EIC energies from whitepaper A. Accardi *et. al.*, *Eur. Phys. J. A*, 52 9(2016). HERA energies for comparison.



Event generation

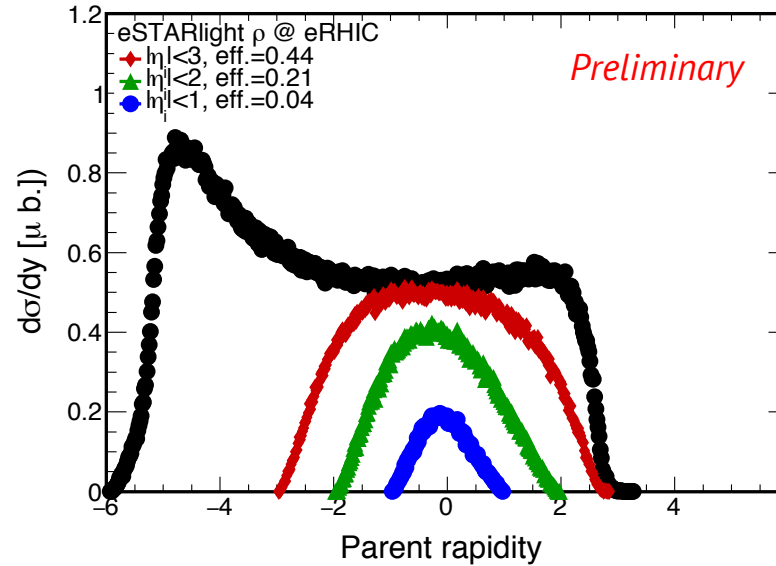
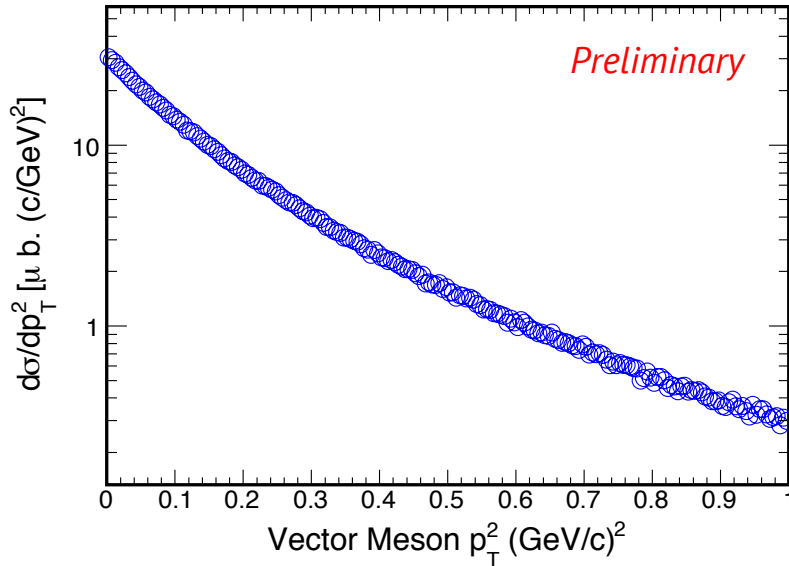
- Coherent or incoherent(**pending**) final states
 - Coherent: couples to nucleus as a whole
 - p_T distribution of Pomeron determined by nuclear form factor
 - Incoherent: couples to a single nucleon in the nucleus
 - p_T distribution of Pomeron determined by proton form factor (same as proton targets)
- Generates final states for input to GEANT to simulate detector response.
- Track outgoing electron and target(**pending**) for semi-inclusive and exclusive measurements
- Simple decays handled by eSTARlight. Can couple to PYTHIA for more complex decays
 - Can put requirements on decay product η , p_T

H1 data corrected for exp. resolution and background subtraction JHEP05(2010)032

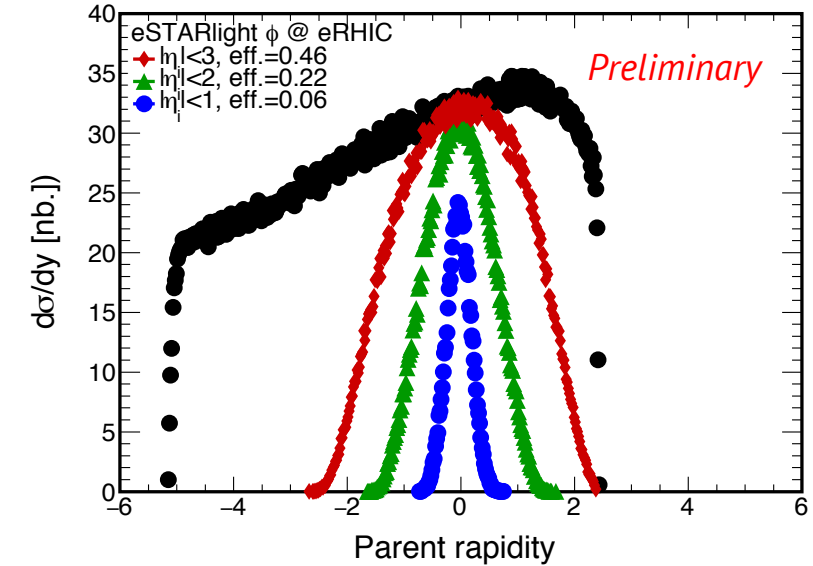


Event generation and distributions eRHIC

ρ



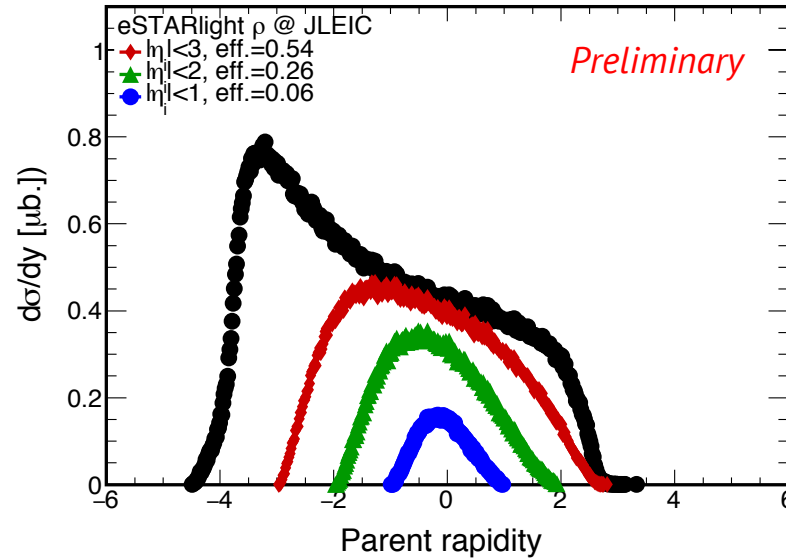
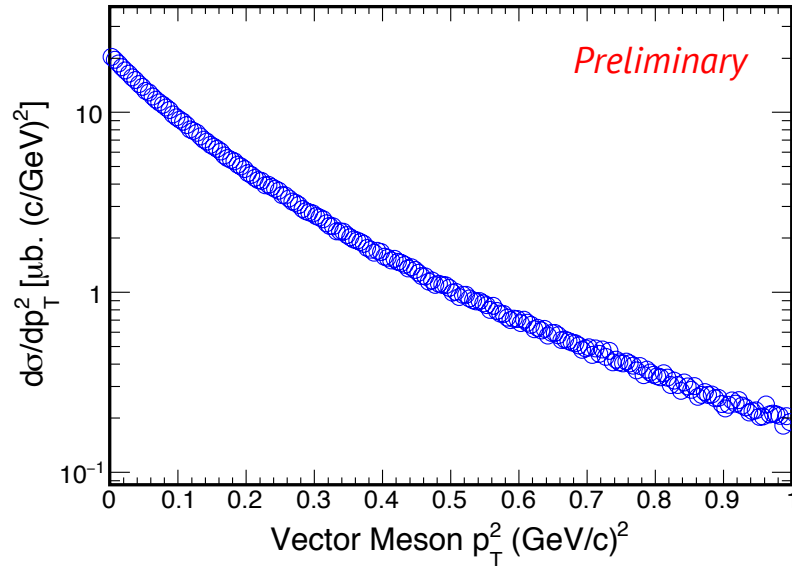
ϕ



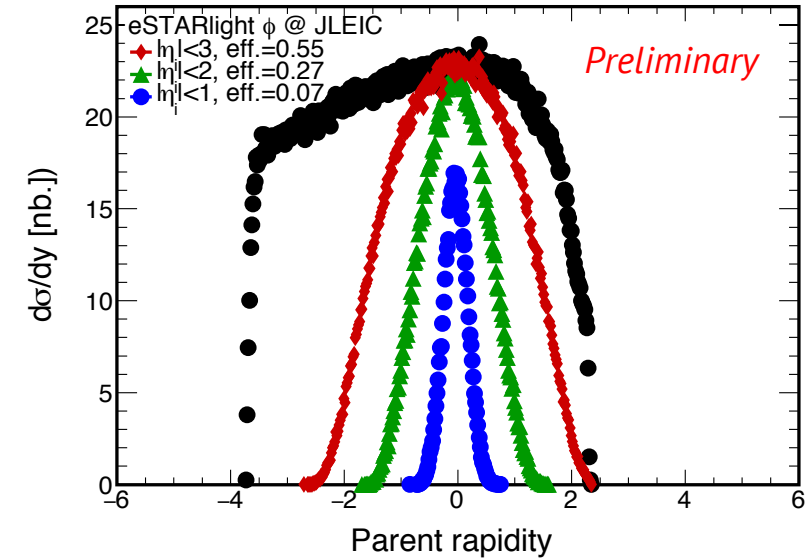
- Vector meson production occurs over a large rapidity window (roughly matches photon energy).
- ρ peak at $y = -4.5$ due to photon-meson exchange (mostly near threshold). Not present in ϕ (production only occurs through photon-pomeron exchange)
- Color curves illustrate V.M.'s with both daughters within a given pseudo-rapidity (η) range.
- Mid-rapidity detectors would sample between $\sim 50\%$ ($|\eta| < 5$) and $< 5\%$ ($|\eta| < 1$) of production:
 - eSTARlight kinematics can inform detector design

Event generation and distributions JLEIC

ρ



ϕ



- Generally lower production at JLEIC relative to eRHIC (different scales).
- Mid-rapidity detectors sample a larger fraction of produced vector-mesons (relative to eRHIC)
- Color curves illustrate V.M.'s with both daughters within a given pseudorapidity (η) range.
- Mid-rapidity detectors would sample between $\sim 55\%$ ($|\eta| < 5$) and $< 10\%$ ($|\eta| < 1$) of production:
 - eSTARlight kinematics can inform detector design

Estimating rates for an EIC

eRHIC¹:

electrons*: 10 GeV

protons: 250 GeV

¹⁹⁷Au⁷⁹⁺: 100 GeV/n

	$\sigma(e + X \rightarrow e + \text{V.M.} + X)$		Rates ($L_{\text{int}} = 10 \text{ fb}^{-1}/\text{A}$)	
V.M.	X = p	X = Au	X = p	X = Au
ρ	4.68 μb	1.4 mb	46.8×10^9	71×10^9
ϕ	211 nb	80 μb	2.1×10^9	4×10^9
J/ψ^*	7.8 nb	2.3 μb	78×10^6	120×10^6
$\Upsilon(1S)^*$	11.3 pb	3.5 nb	113×10^3	180×10^3

JLEIC¹:

electrons: 5 GeV

protons: 60 GeV

²⁰⁸Pb⁸²⁺: 40 GeV/n

	$\sigma(e + X \rightarrow e + \text{V.M.} + X)$		Rates ($L_{\text{int}} = 10 \text{ fb}^{-1}/\text{A}$)	
V.M.	X = p	X = Pb	X = p	X = Pb
ρ	3.1 μb	1.0 mb	31×10^9	48×10^9
ϕ	121 nb	55 μb	1.2×10^9	2.6×10^9
J/ψ^{**}	2.2 nb	930 nb	22×10^6	45×10^6
$\Upsilon(1S)^{**}$	2.0 pb	710 pb	20×10^3	34×10^3

Don't account for branching ratios

¹: EIC whitepaper: A. Accardi et. al., Eur. Phys. J. A, 52 9(2016)

*: Latest design uses 21.5 GeV electron beam @ eRHIC. Rates need to be revisited

**: Likely overestimated: Doesn't account for loss of longitudinal coherence

Estimating rates for an EIC ($Q^2 > 1 \text{ [GeV}^2/\text{c}^2]$)

eRHIC¹:

electrons*: 10 GeV

protons: 250 GeV

¹⁹⁷Au⁷⁹⁺: 100 GeV/n

	$\sigma(e + X \rightarrow e + \text{V.M.} + X)$		Rates ($L_{\text{int}} = 10 \text{ fb}^{-1}/A$)	
V.M.	X = p	X = Au	X = p	X = Au
ρ	16 nb	5.1 μb	160×10^6	260×10^6
ϕ	1.7 nb	660 nb	17×10^6	34×10^6
J/ψ^{**}	440 pb	150 nb	4×10^6	7.7×10^6
$\Upsilon(1S)^{**}$	1.7 pb	550 pb	17×10^3	28×10^3

JLEIC¹:

electrons: 5 GeV

protons: 60 GeV

²⁰⁸Pb⁸²⁺: 40 GeV/n

	$\sigma(e + X \rightarrow e + \text{V.M.} + X)$		Rates ($L_{\text{int}} = 10 \text{ fb}^{-1}/A$)	
V.M.	X = p	X = Pb	X = p	X = Pb
ρ	12 nb	4 μb	120×10^6	190×10^6
ϕ	1.1 nb	490 nb	11×10^6	24×10^6
J/ψ^{**}	150 pb	68 nb	1.5×10^6	3.3×10^6
$\Upsilon(1S)^{**}$	330 fb	130 pb	3×10^3	6.3×10^3

- Q^2 restriction affects V.M. species to different degree:
 ρ ($\times 10^{-3}$) vs. J/ψ ($\times 10^{-1} - 10^{-2}$)

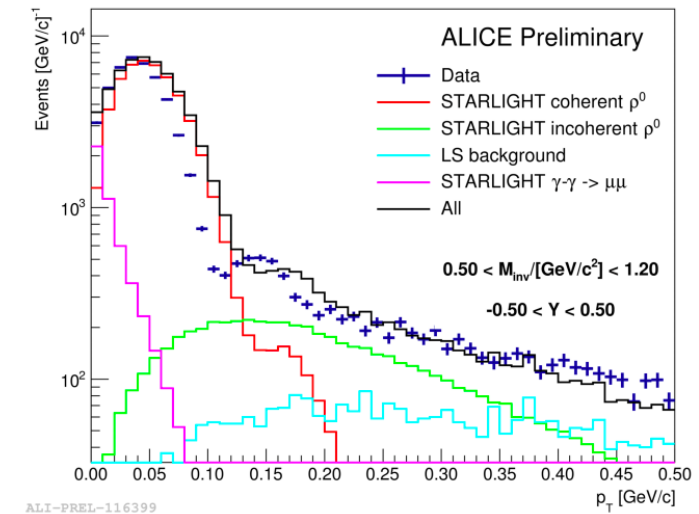
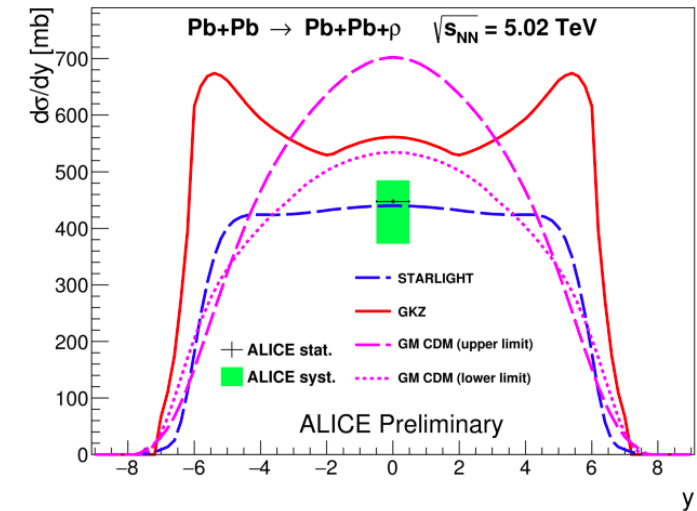
Don't account for branching ratios

Some caveats:

- Classical vs. quantum Glauber calculation:
 - Quantum calculations predict dips but over estimate cross sections (factor 2)¹ for UPC ρ^0 photoproduction at RHIC and LHC.
 - Option for classical Glauber available.
- Assumes complete longitudinal coherence
 - Under investigation, should only affect near threshold (small y)

Next steps:

- No neutron skin:
 - Form factor in eSTARlight from electron-nucleus elastic scattering measurements \rightarrow only probes the protons.
 - Shift in ρ coherent peak.
- Vet all angular distributions – (also parametrization to data)
- Implement γ - γ interactions
- Production of exotic final states
- Include charge exchange interactions, i.e. $\gamma+p \rightarrow Z_c(4430)+n$



1: J. High Energy. Phys. (2015) 2015: 95
2: arXiv:1611.05471 (2016)

Summary

- eSTARlight can simulate a wide variety of final states:
 - Evaluate feasibility (cross sections, rates, ...) of different physics topics to be studied
- Vector mesons are produced over a wide rapidity range, roughly corresponding to photon energies
 - Drive detector design by understanding final state kinematics
 - Detector covering +/- 2 units of η would sample ~20-25% of vector meson cross sections
 - Forward detector would sample high energy photons, while backward rapidity samples near threshold
- Generally high enough production rates for $Q^2 > 1$ GeV events
- High enough $Y(1S)$ production at an EIC to allow limited studies ($Q^2 > 1$ GeV rates are somewhat low)
- Are the rates are large enough to study rare events (exotica)?

Thank you

Back – ups (or might get removed)