

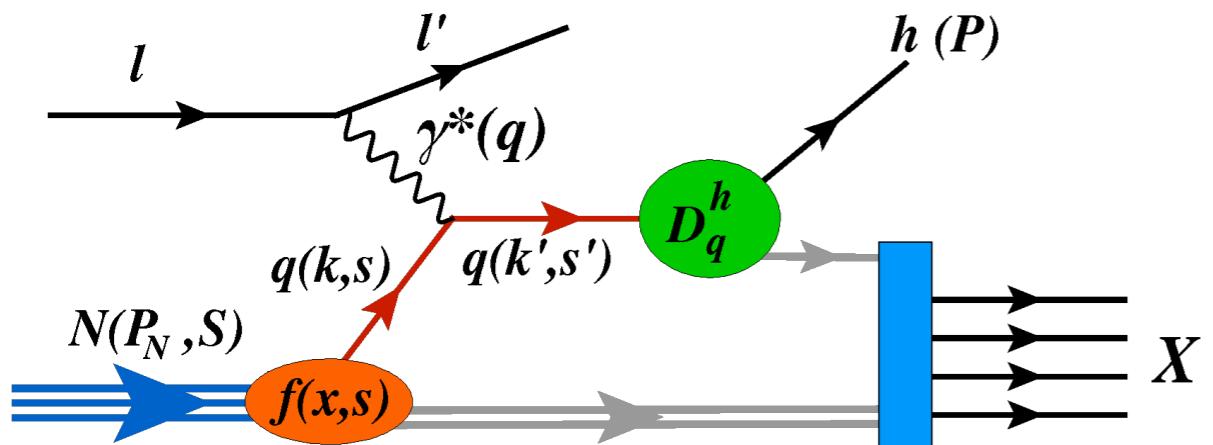
Summary of Parallel Sessions 2 and 5:

*Nucleon and nuclear structure and hadronization
Collective effects in nucleons and nuclei*

Convenors:

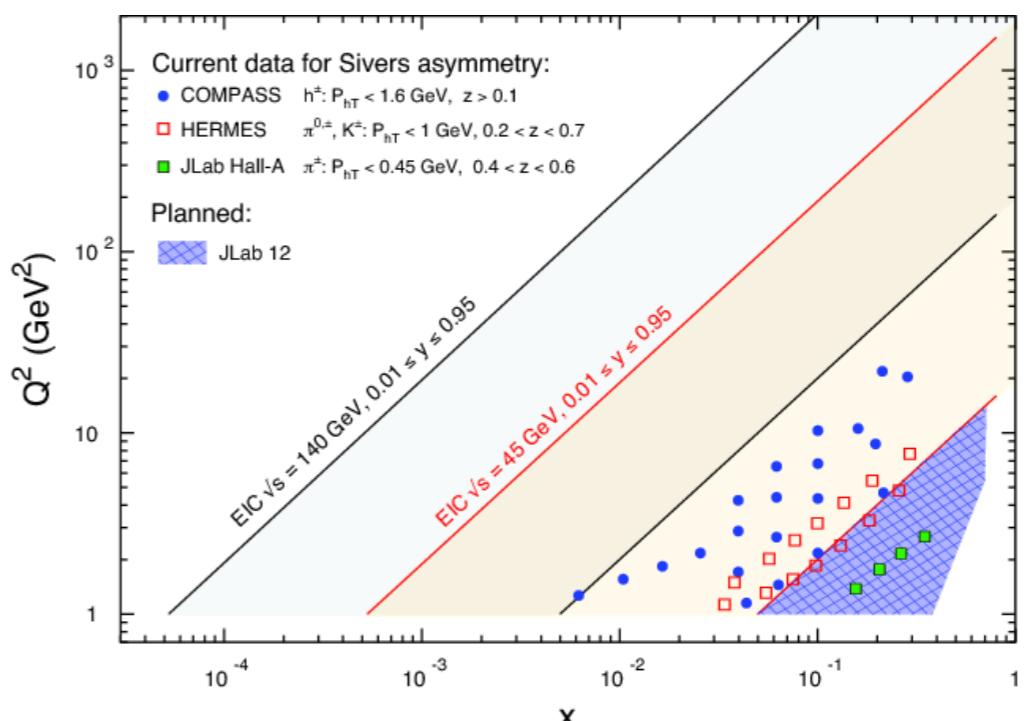
**Gunar Schnell, Hrayr Matevosyan,
Marta Ruspa, Yoshitaka Hatta**

SIDIS program in EIC



- Flavour separation in collinear PDFs.
- Polarised PDFs.
- TMD (polarized) PDFs.

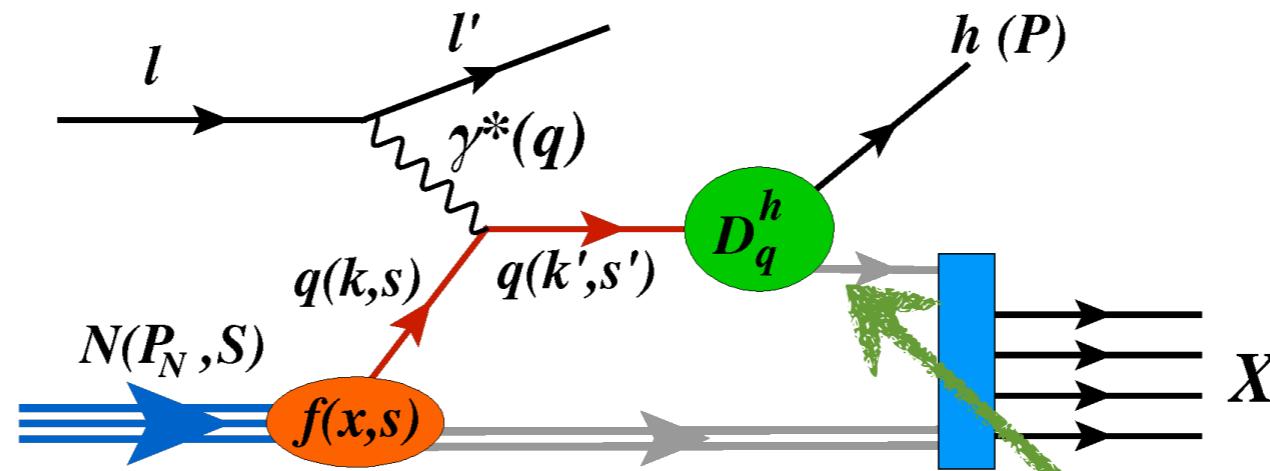
EIC extended
kinematic coverage



Accessing TMDs in SIDIS

- **Measurement of the transverse momentum of the produced hadron in SIDIS provides access to TMD PDFs/FFs.**

- *SIDIS Process with TM of hadron measured.*



- **TMD PDFs**

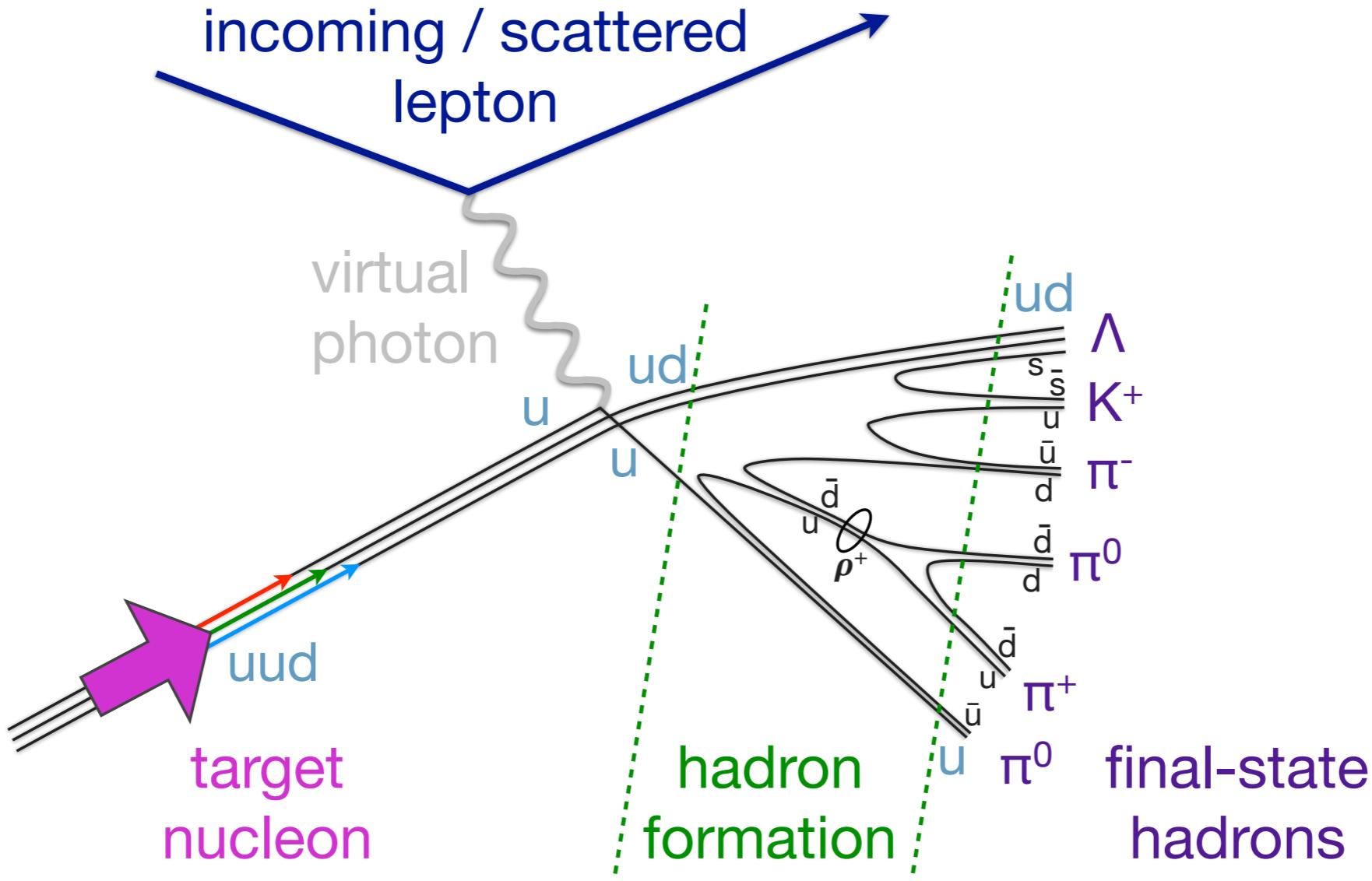
N/q	U	L	T
U	f_1		h_1^\perp
L		g_{1L}	h_{1L}^\perp
T	f_{1T}^\perp	g_{1T}^\perp	$h_1 h_{1T}^\perp$

- **TMD FFs**

q/h	U
U	D_1
L	
T	H_1^\perp

*unpol/spinless h !

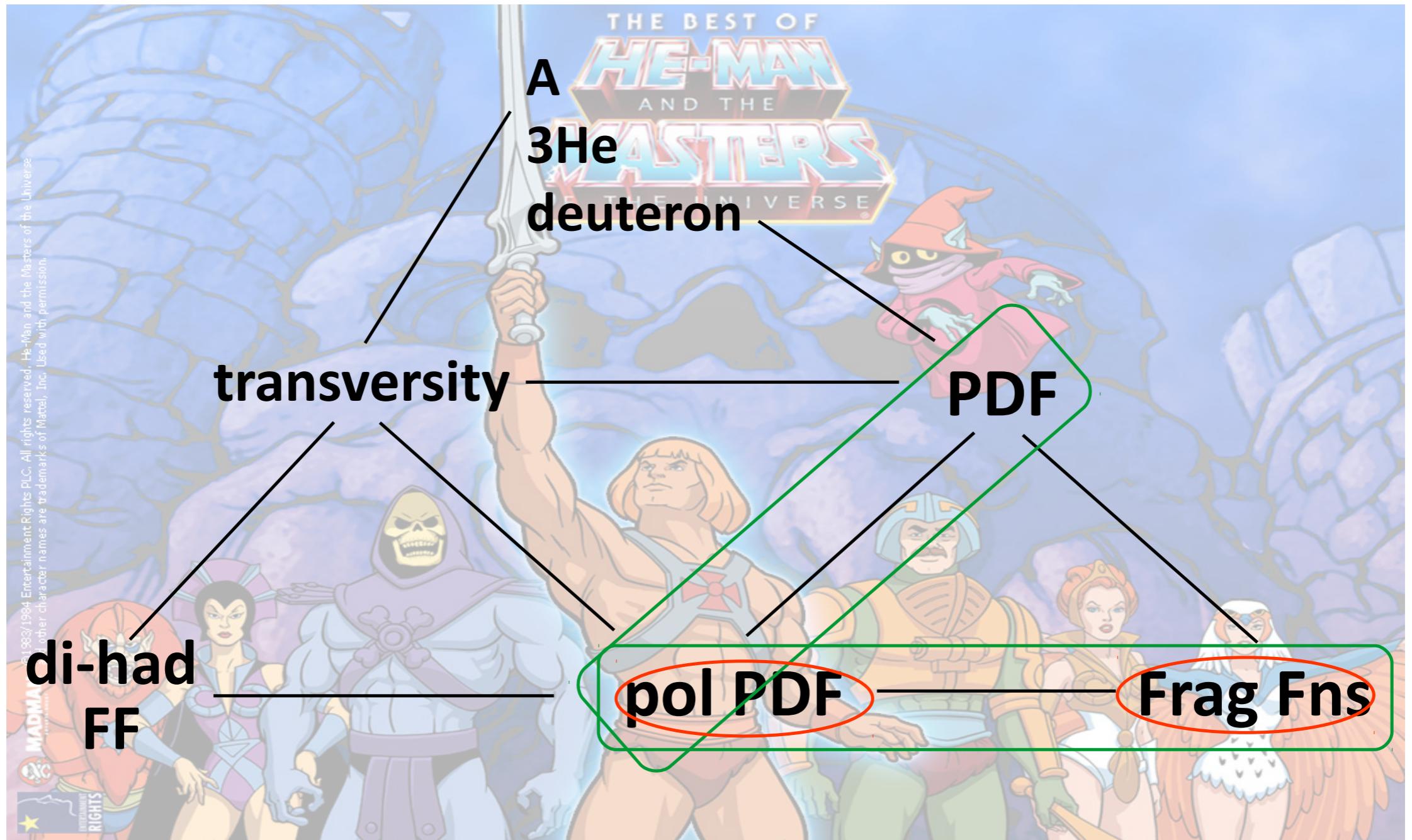
Accessing Parton Properties at EIC



- ♦ Need the knowledge on PDFs and (in particular) FFs.
- ♦ Their extractions not necessarily independent, but *naturally interlinked*.

Masters of the Universe

A. Accardi



accardi@jlab.org

EICUG meeting, Trieste – 21 July 2017

27

towards “universal” (combined) fits

5+ years: new fitting methods

A. Accardi

- More computing power, efficient implementations
 - New fitting, analysis methods
- In traditional fits:
 - Detailed χ^2 scans, refined statistical analysis
- Monte Carlo fitting methods:
 - NNPDF: bootstrap + neural network fit → *Nocera's talk*
 - JAM: bootstrap + Iterative Monte Carlo (IMC) approach
→ *Sato, Ethier et al (since 2015)*
- Self organizing maps → *Liuti et al.*



First simultaneous extraction of spin-dependent parton distributions and fragmentation functions from a global QCD analysis

J. J. Ethier,^{1,2} N. Sato,³ and W. Melnitchouk²

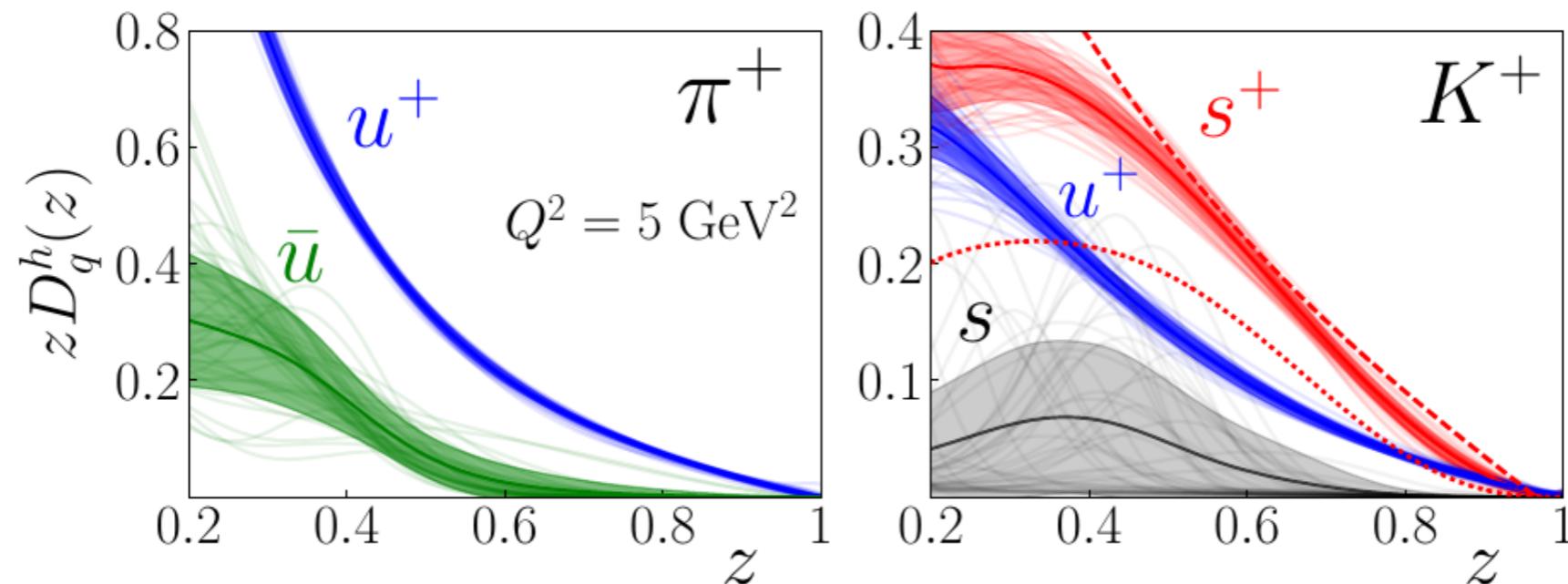
¹*College of William and Mary, Williamsburg, Virginia 23187, USA*

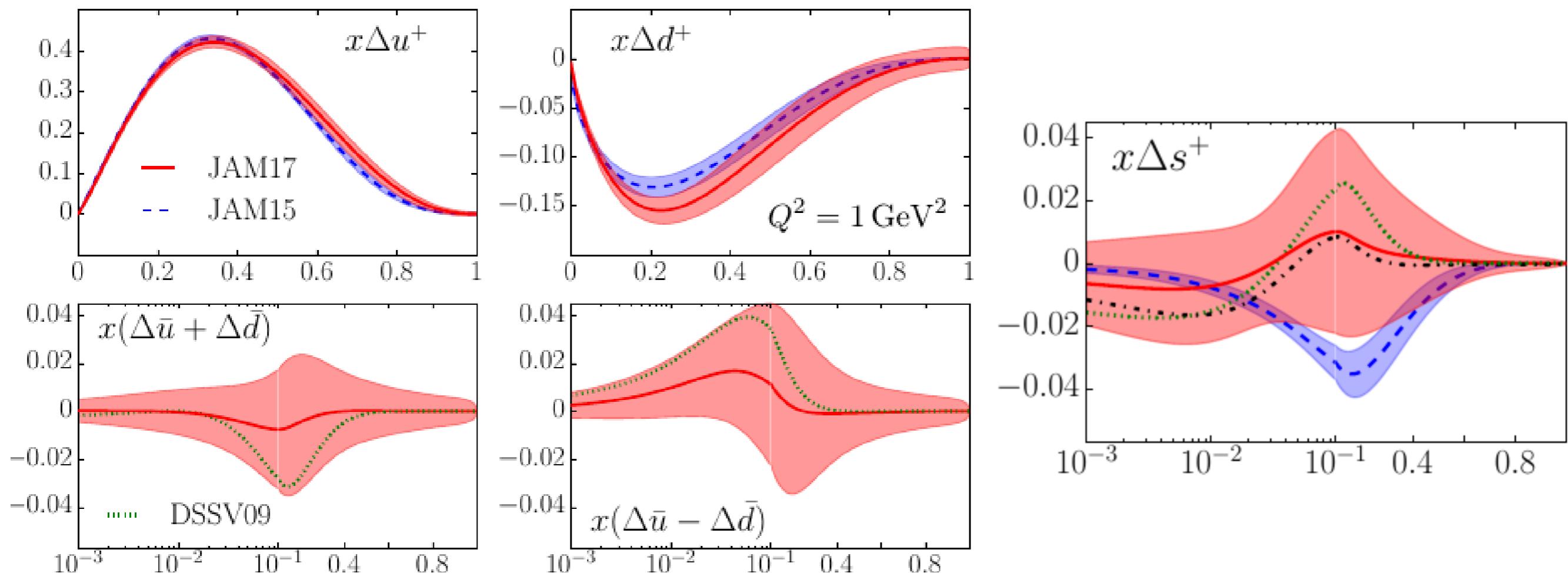
²*Jefferson Lab, Newport News, Virginia 23606, USA*

³*University of Connecticut, Storrs, Connecticut 06269, USA*

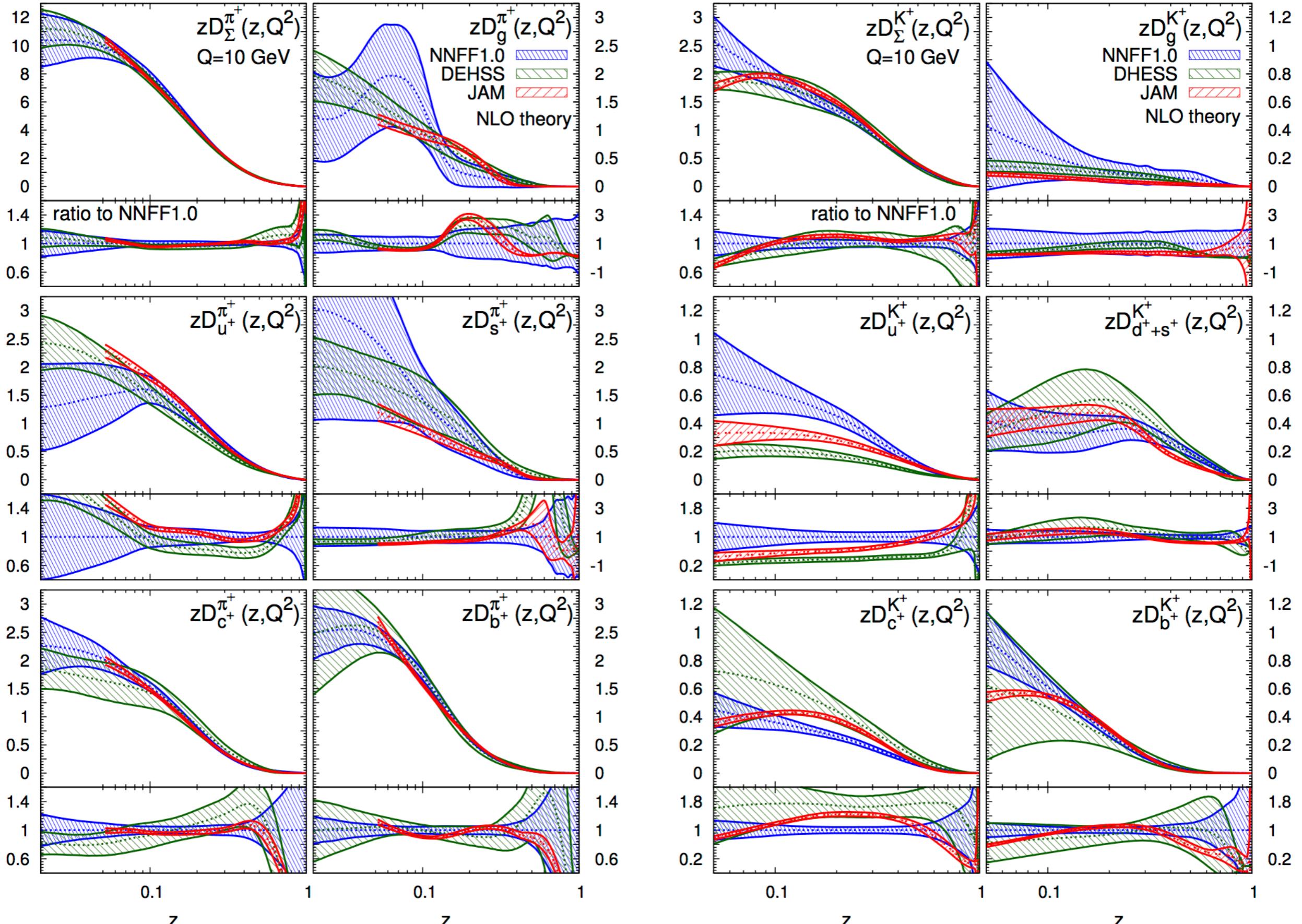
Jefferson Lab Angular Momentum (JAM) Collaboration

(Dated: May 18, 2017)

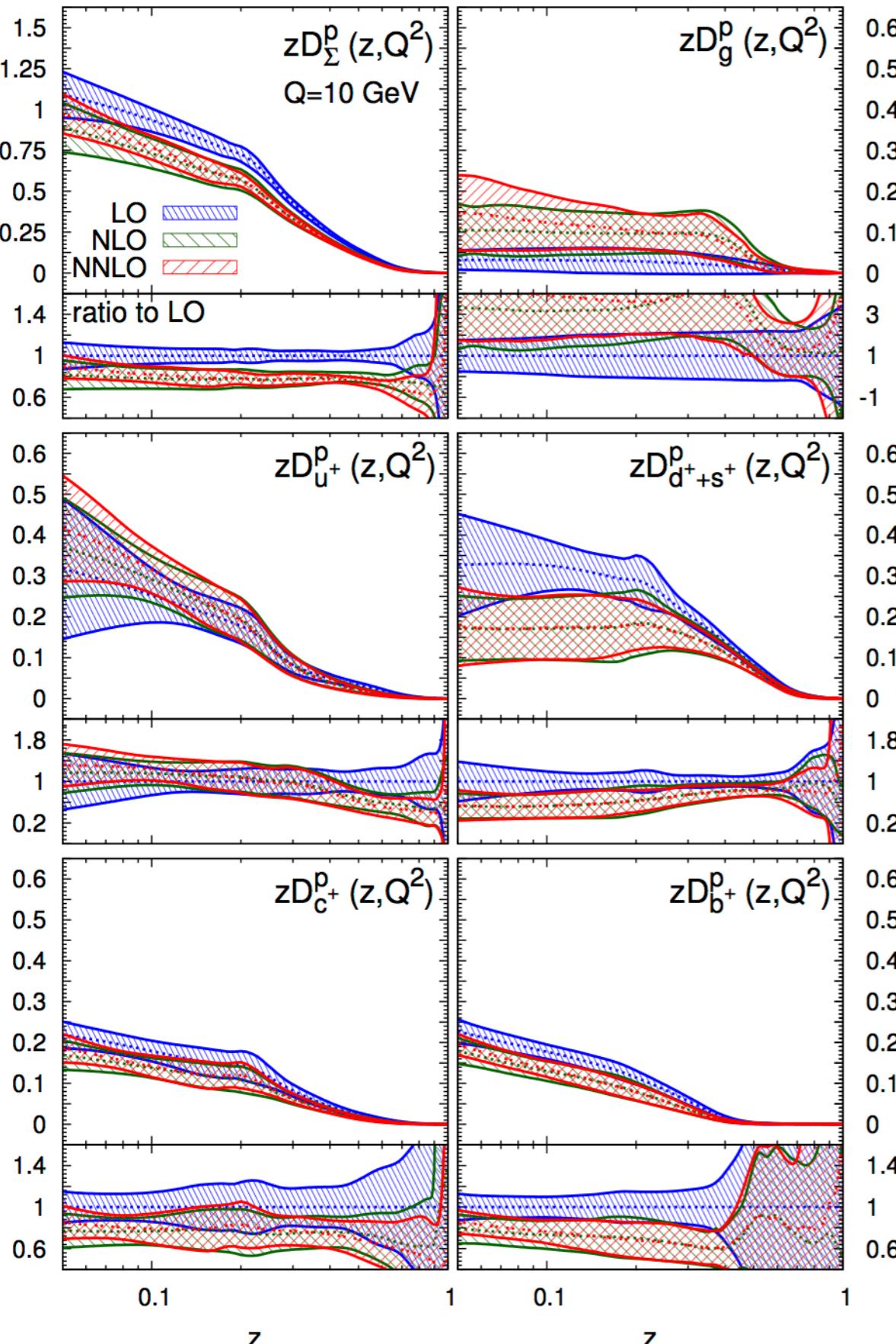




Comparison with other recent FF sets: π^+ and K^+



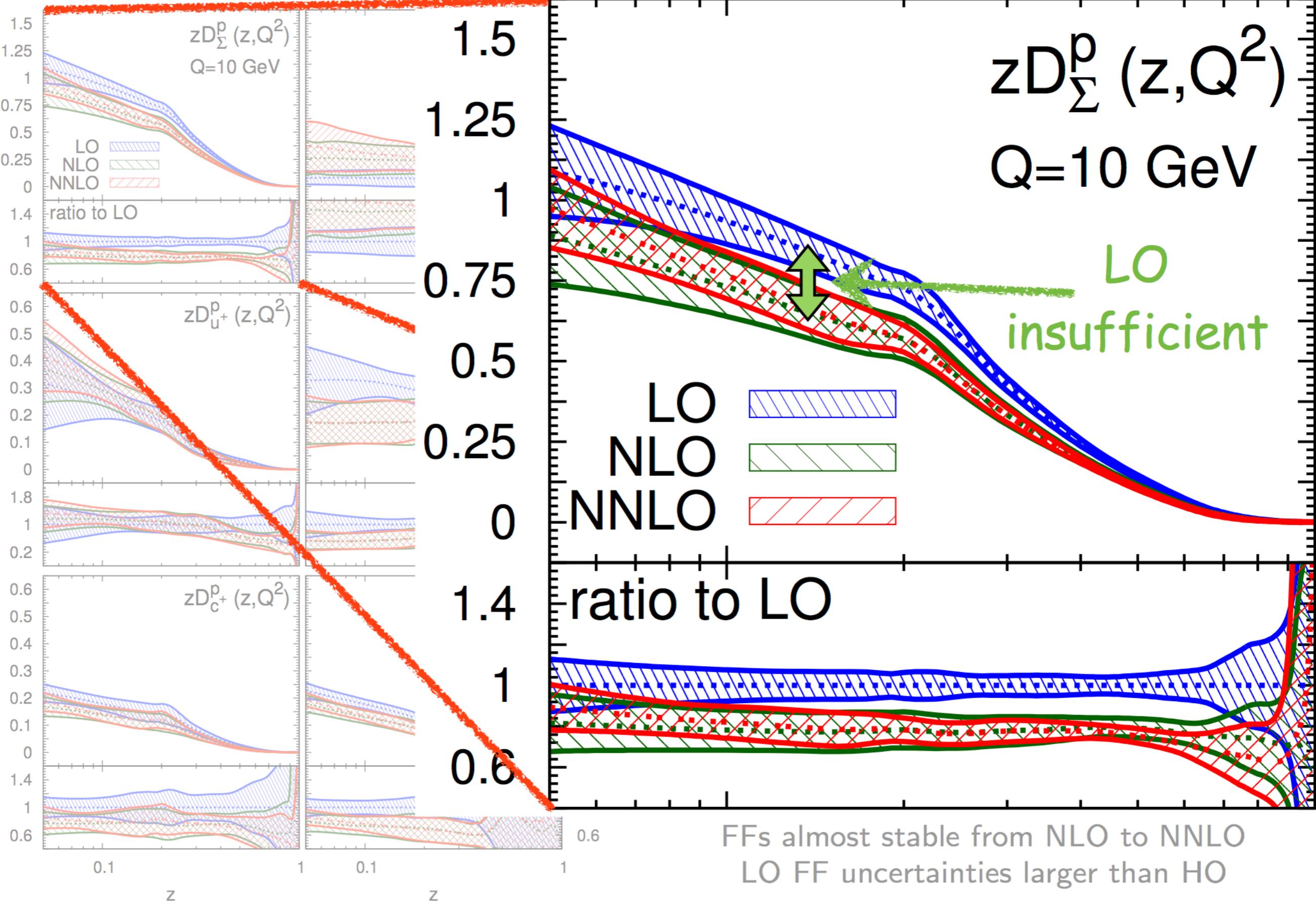
Fit quality vs perturbative order: p



Exp.	N_{dat}	LO χ^2/N_{dat}	NLO χ^2/N_{dat}	NNLO χ^2/N_{dat}
BELLE	29	0.10	0.31	0.50
BABAR	43	4.74	3.75	3.25
TASSO12	3	0.69	0.70	0.72
TASSO14	9	1.32	1.25	1.22
TASSO22	9	0.98	0.92	0.93
TPC	20	1.04	1.10	1.08
TPC-UDS	—	—	—	—
TPC-C	—	—	—	—
TPC-B	—	—	—	—
TASSO30	2	0.25	0.19	0.18
TASSO34	6	0.82	0.81	0.78
TASSO44	—	—	—	—
TOPAZ	4	0.79	1.21	1.19
ALEPH	26	1.36	1.43	1.28
DELPHI	22	0.48	0.49	0.49
DELPHI-UDS	22	0.47	0.46	0.45
DELPHI-B	22	0.89	0.89	0.91
OPAL	—	—	—	—
SLD	36	0.66	0.65	0.64
SLD-UDS	36	0.77	0.76	0.78
SLD-C	36	1.22	1.22	1.21
SLD-B	35	1.12	1.29	1.33
TOTAL	360	1.31	1.23	1.17

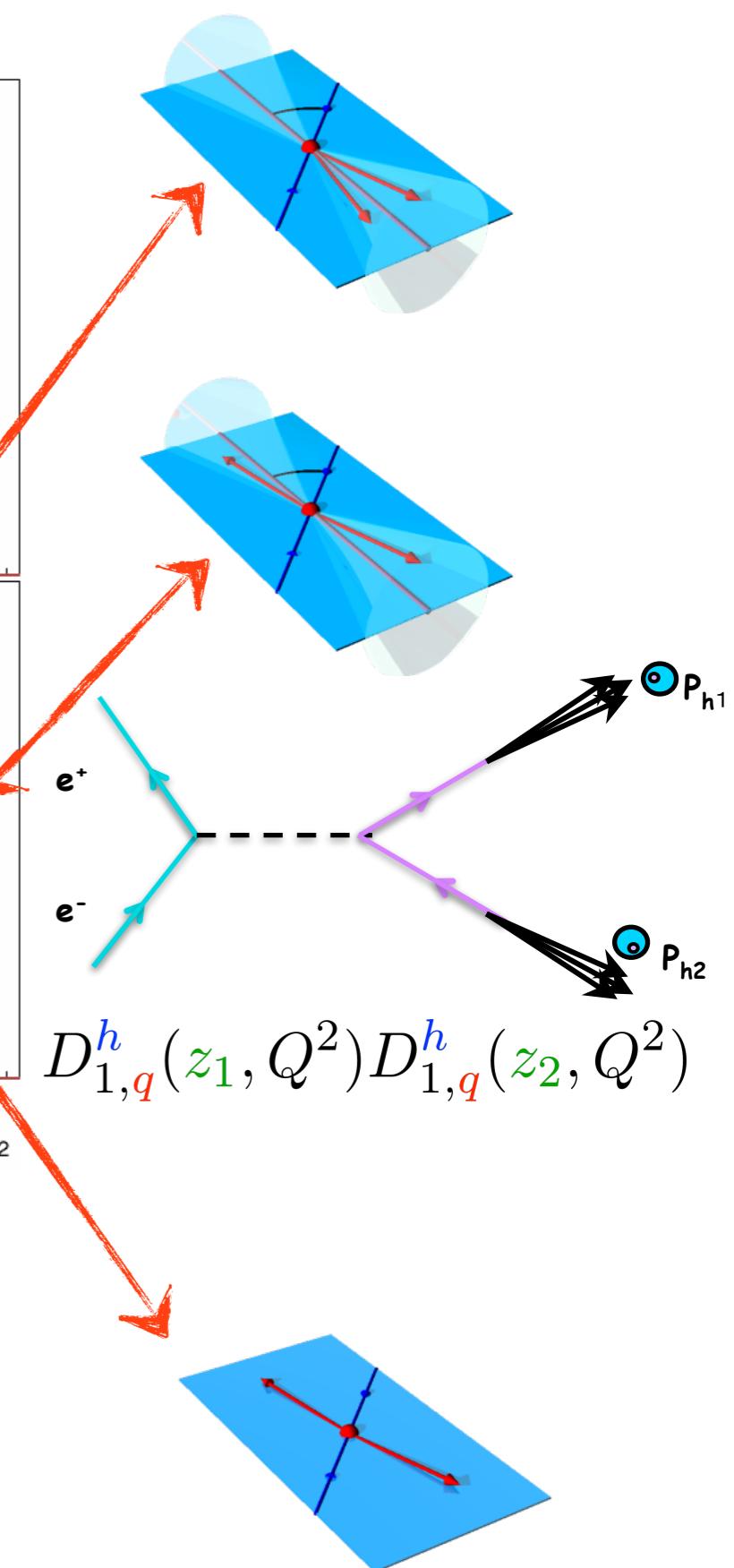
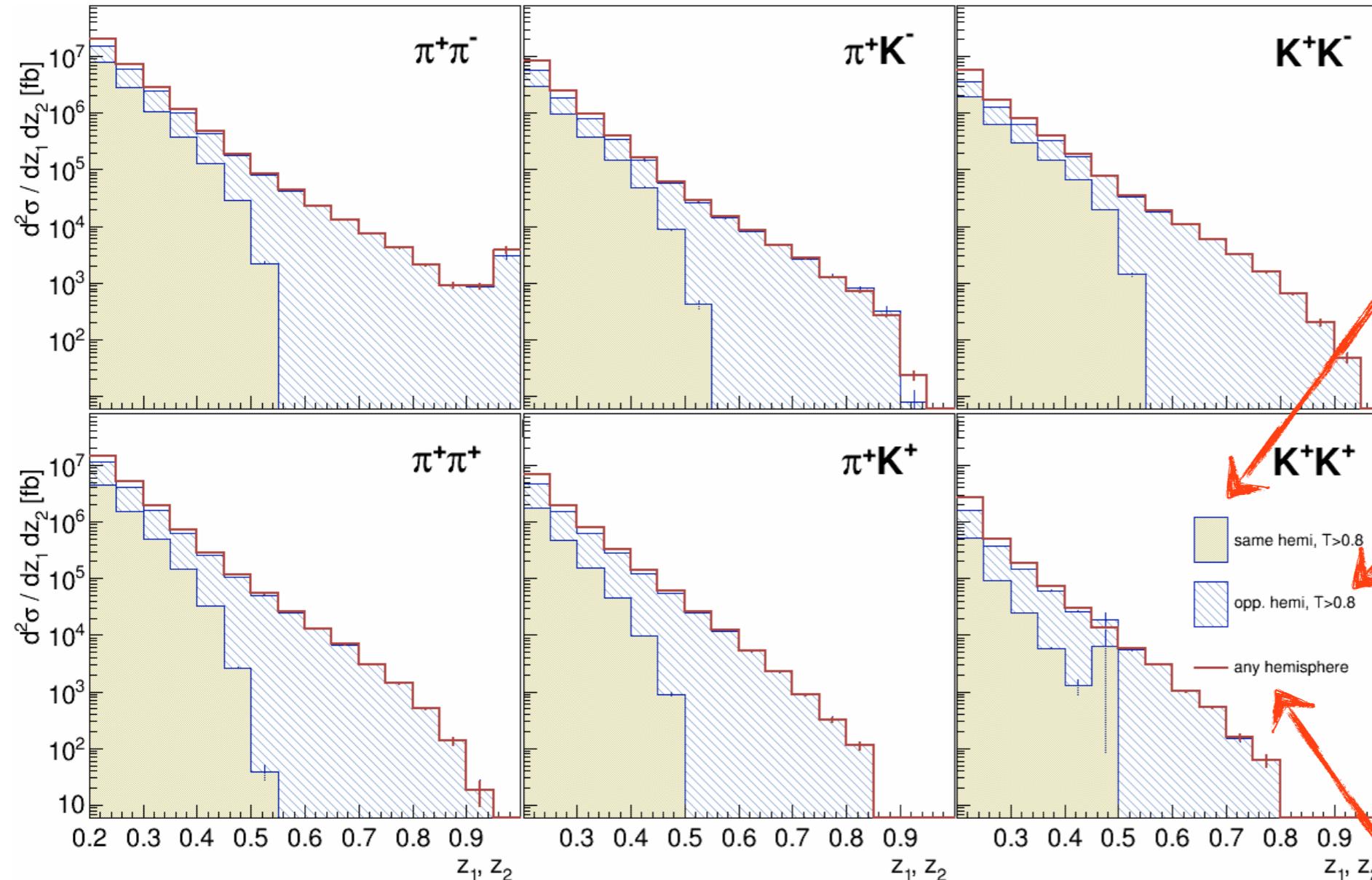
Excellent perturbative convergence
FFs almost stable from NLO to NNLO
LO FF uncertainties larger than HO

Fit quality vs perturbative order: p



Additional FF data from e+e-

R. Seidl

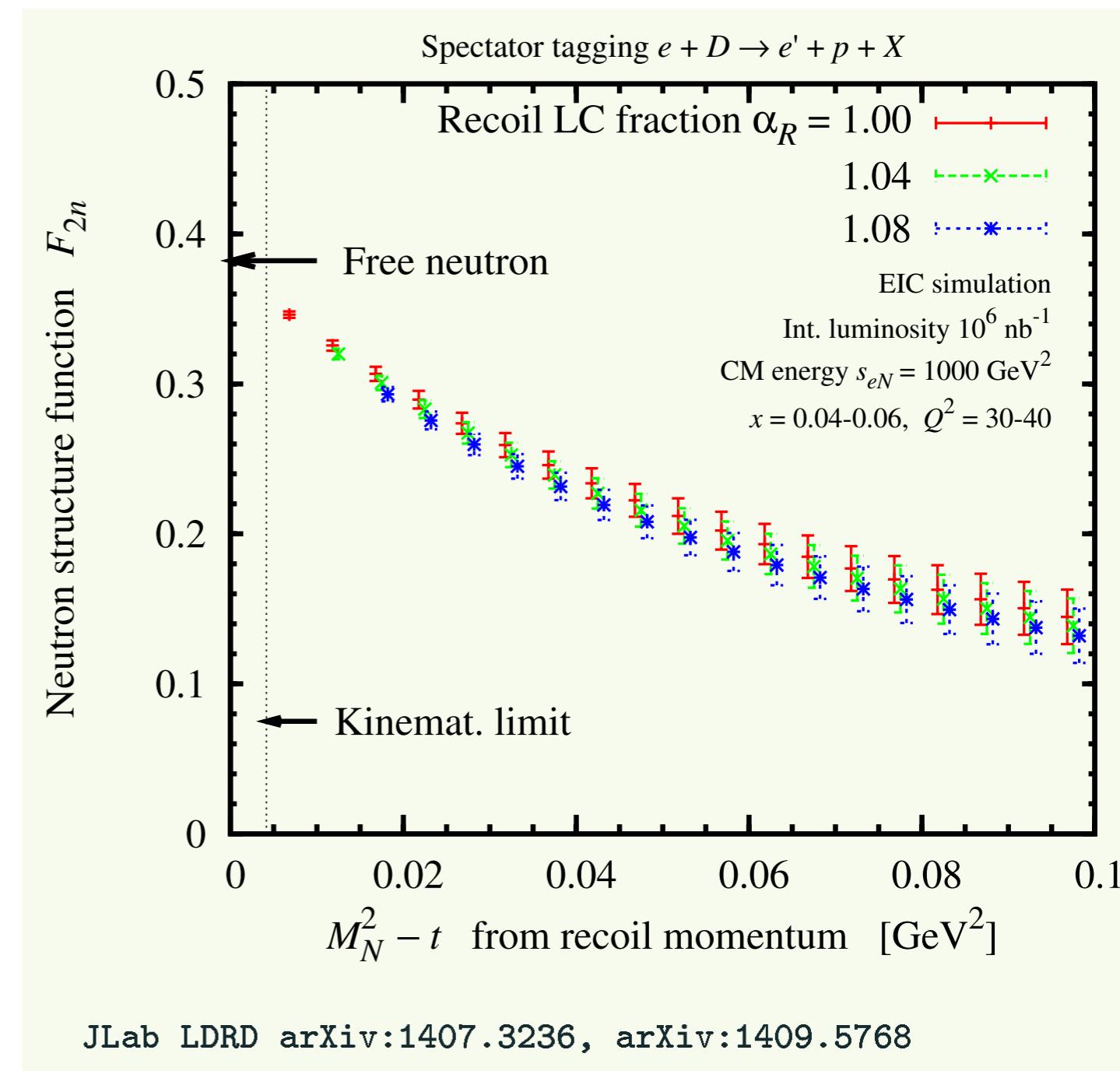
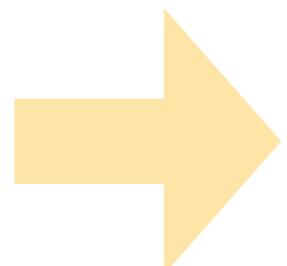
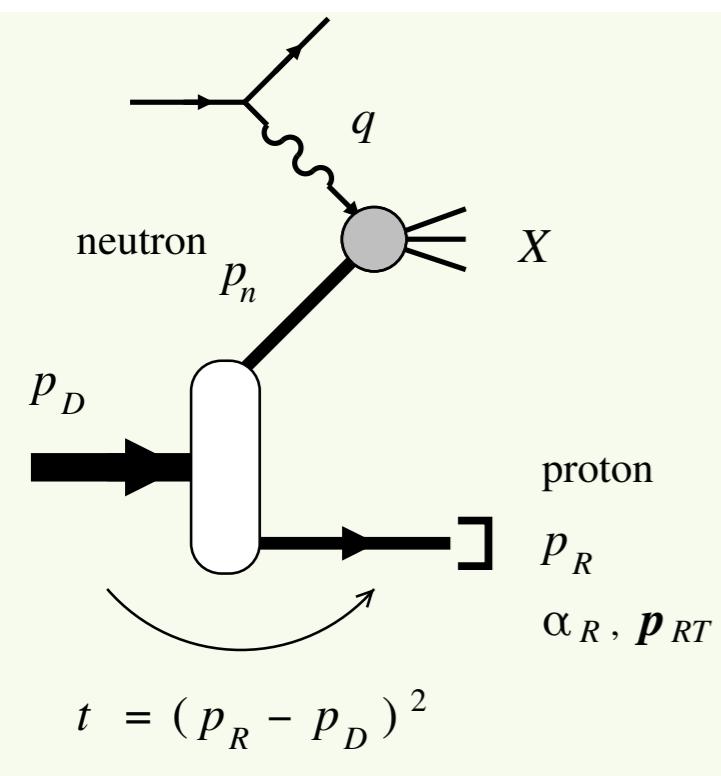


gain additional flavor sensitivity
by looking at pairs of hadrons

Tagged Structure Functions

W. Cosyn

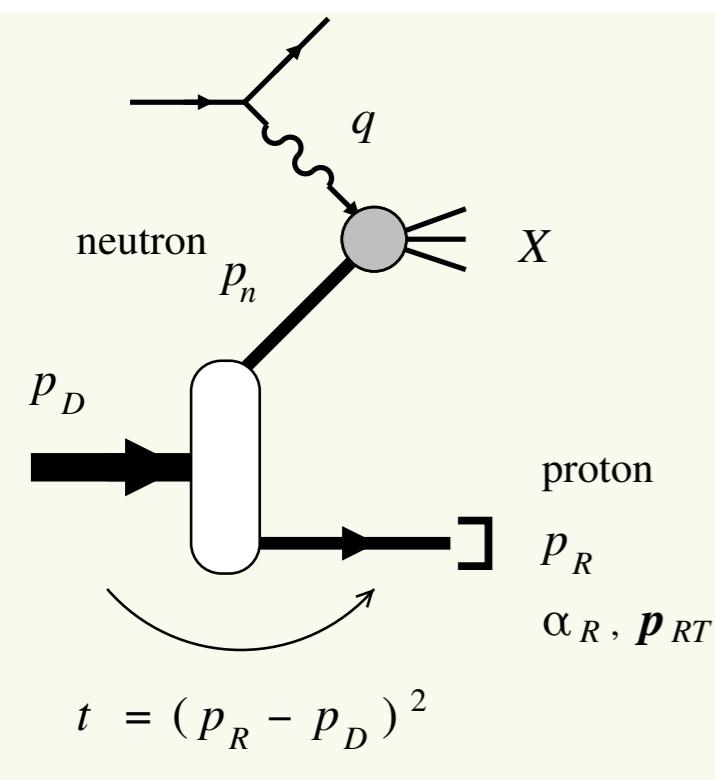
detect recoil proton from deuteron target to study neutron structure



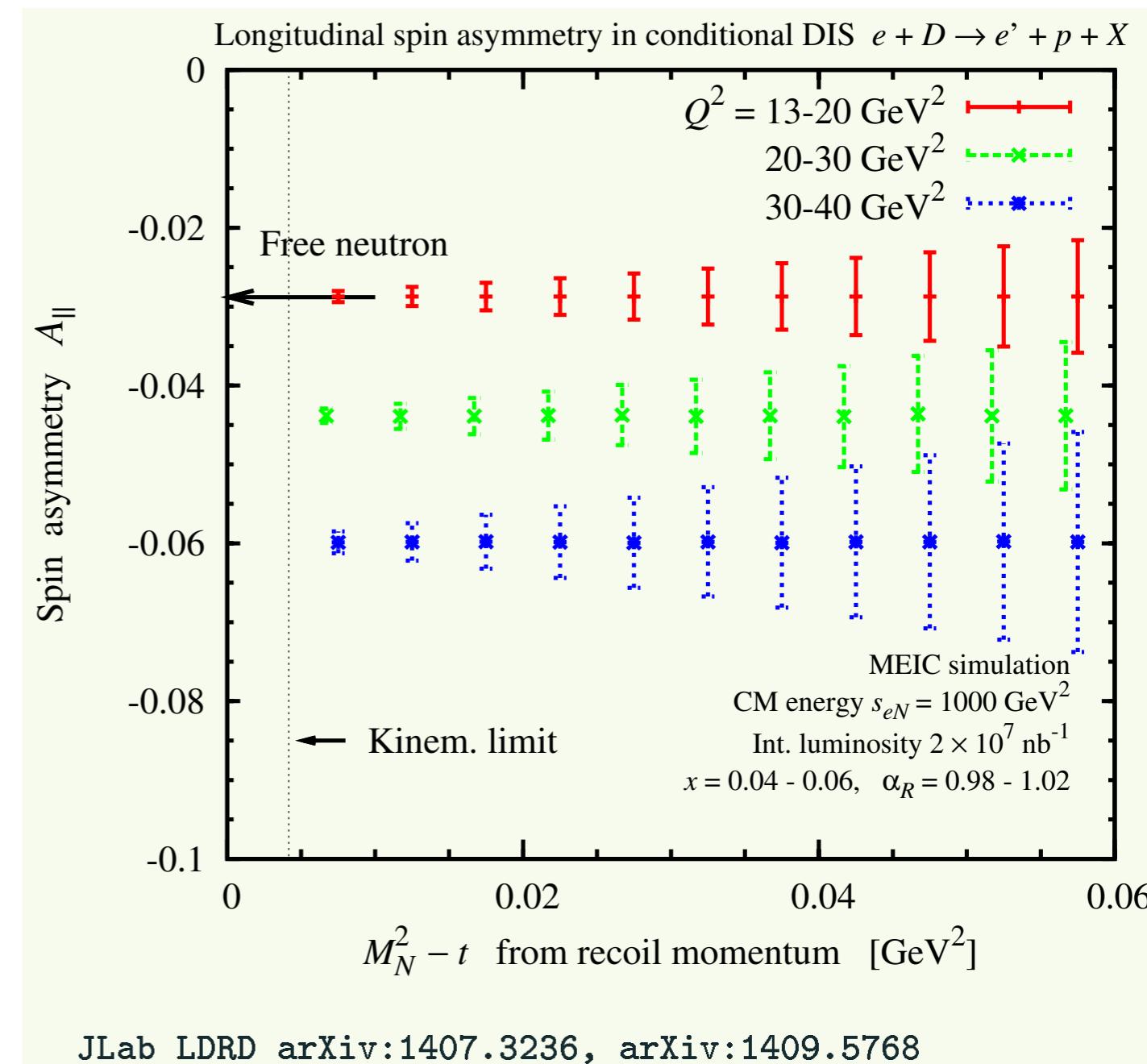
Tagged Structure Functions

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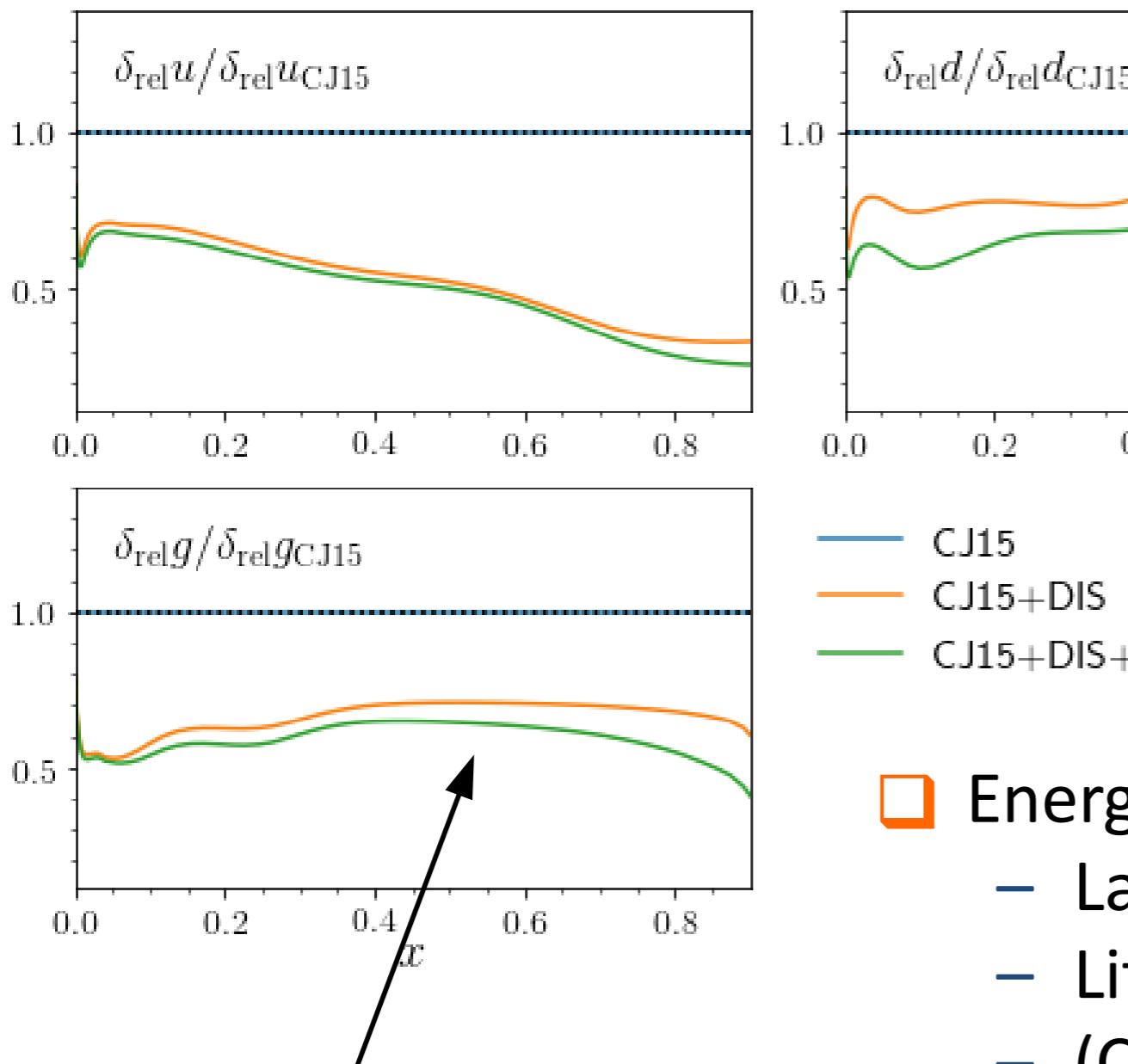


can also look at helicity distributions of neutron



Impact - summary

A. Accardi



40-50% better gluons

40-50% better
d-quarks

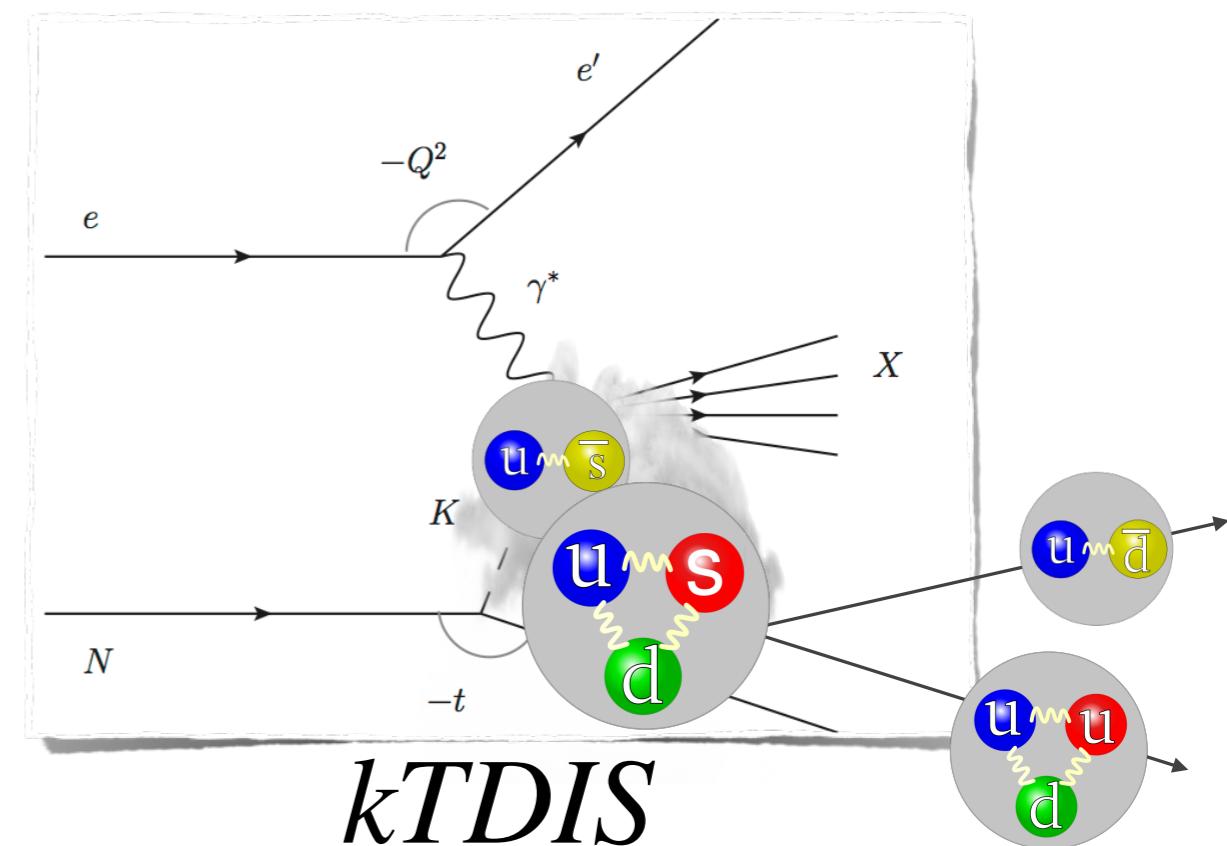
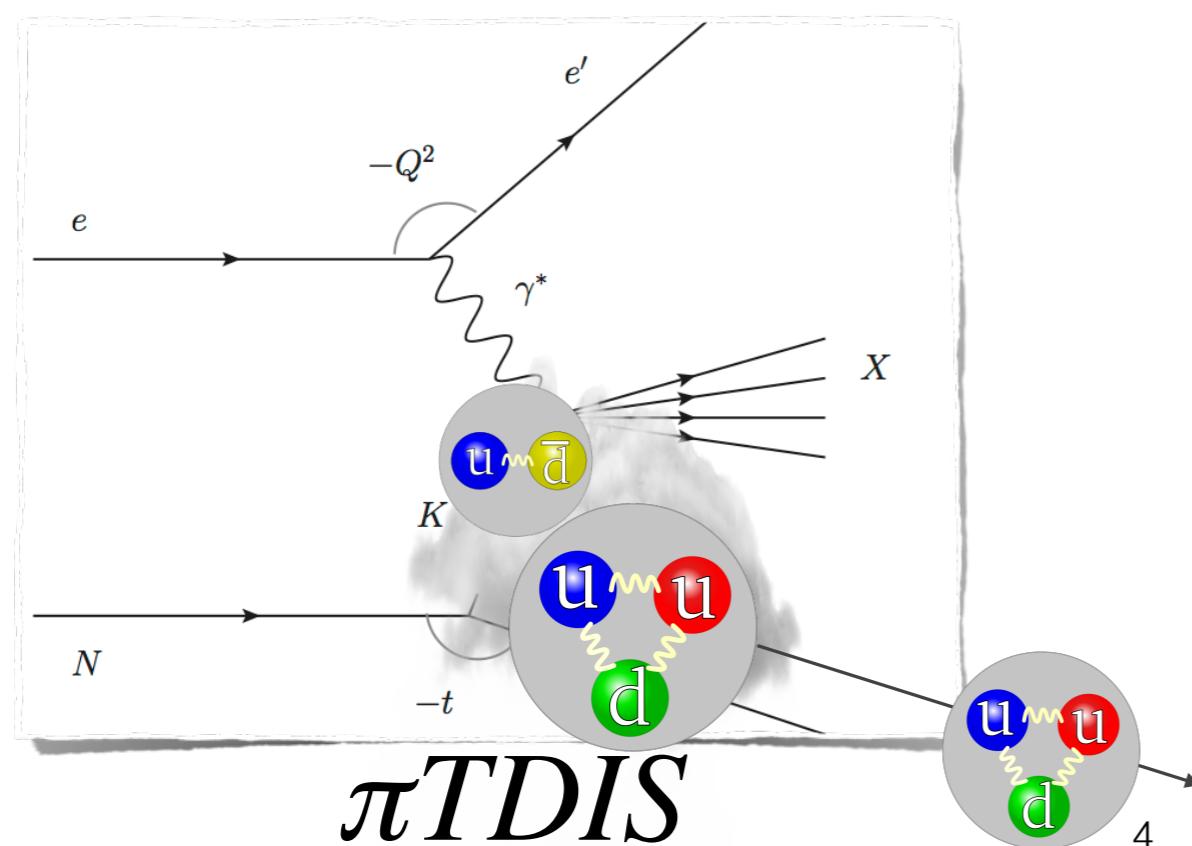
- Energy scan \leftrightarrow L/T separation
 - Large improvement in gluon uncertainty
 - Little effect on d -quarks
 - (CC have very little effects on u,d,g)
- Tagged neutrons
 - Noticeable improvement for d -quarks u
 - Some effects on gluons

Tagging method to study meson structure

K. Park

Tagged Deep Inelastic Scattering (TDIS)

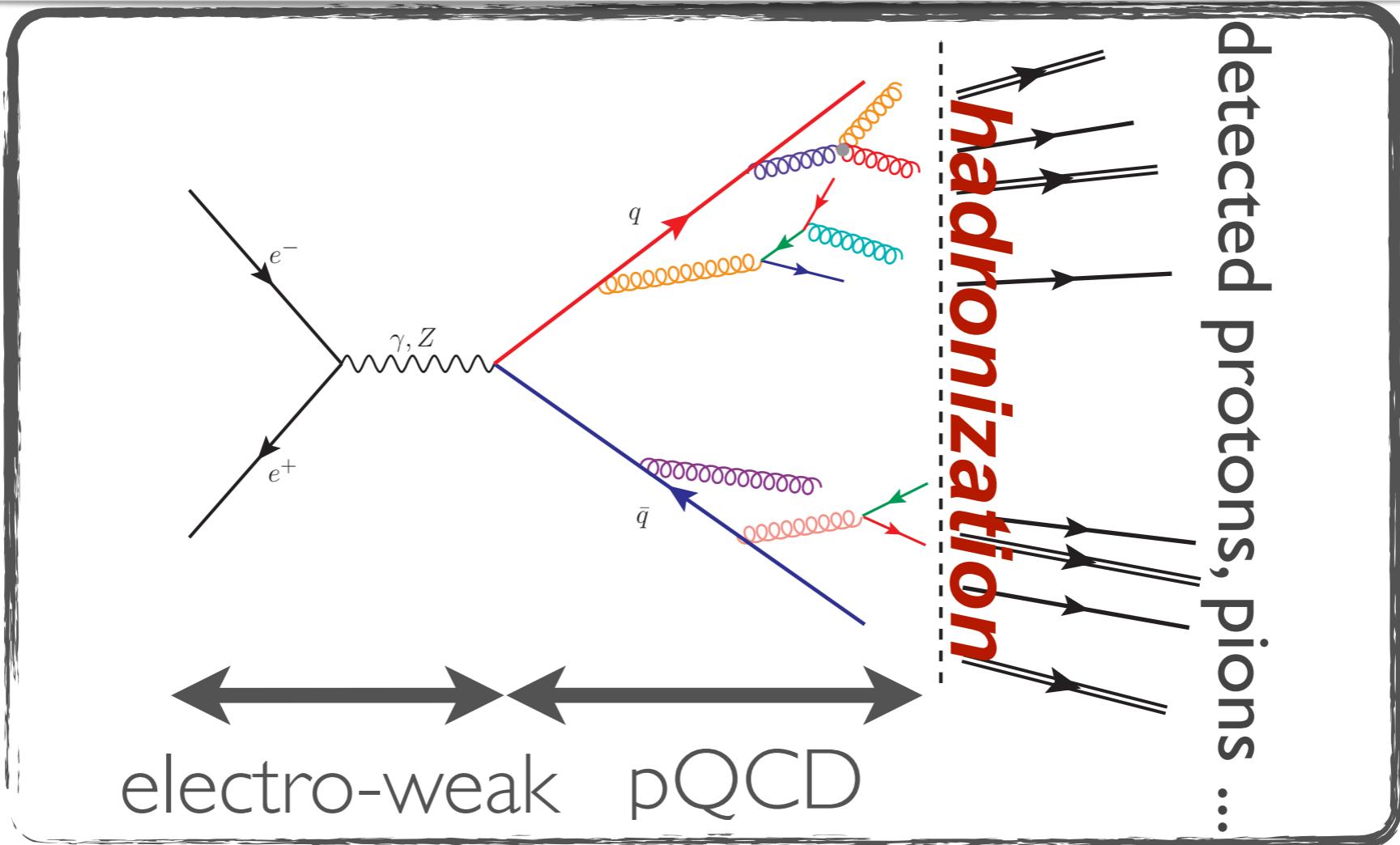
- Sullivan Process
 - .. provides reliable access to a meson target as t becomes space-like (the meson pole dominance of the process)
- Direct measure the mesonic-nucleon content



Hadronization

- The conjecture of Confinement:

♦ NO free quarks or gluons have been directly observed:
only HADRONS.



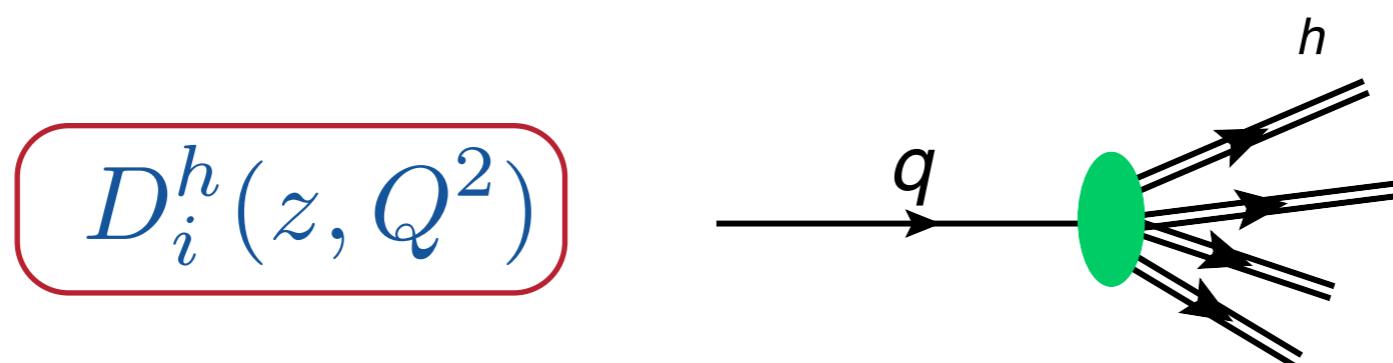
- ♦ Hadronization: describes the process where colored quarks and gluons form colourless hadrons (in deep inelastic scattering).

Fragmentation Functions

- The non-perturbative, universal functions encoding parton hadronization are the: ***Fragmentation Functions (FF)***.

$$\frac{1}{\sigma} \frac{d}{dz} \sigma(e^- e^+ \rightarrow hX) = \sum_i C_i(z, Q^2) \otimes D_i^h(z, Q^2)$$

- Unpolarized FF is the **number density** for parton *i* to produce hadron *h* with LC momentum fraction *z*.



$$D_i^h(z, Q^2)$$

- z is the light-cone mom. fraction of the parton carried by the hadron

$$z = \frac{p^-}{k^-} \approx z_h = \frac{2E_h}{Q}$$

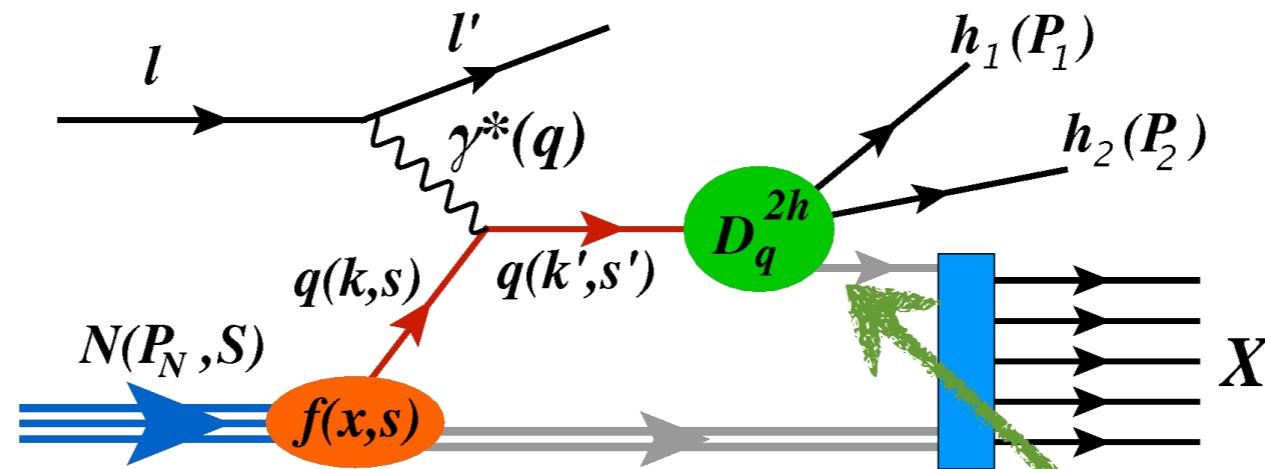
$$a^\pm = \frac{1}{\sqrt{2}}(a^0 \pm a^3)$$

Modelling Hadronization

- Need to understand the non-perturbative mechanism of hadron formation.
- To include the *TM dependence and polarization*.
- Connection between the *one hadron and dihadron FFs*.
- Implementation in MC generators (PYTHIA, etc).

TMD PDFs with Two-Hadron FFs

- Measuring two-hadron semi-inclusive DIS: an additional method for accessing TMD PDFs.
 - SIDIS Process with TM of hadrons measured.



- TMD PDFs

N/q	U	L	T
U	f_1		h_1^\perp
L		g_{1L}	h_{1L}^\perp
T	f_{1T}^\perp	g_{1T}^\perp	$h_1 h_{1T}^\perp$

- TMD DiFFs

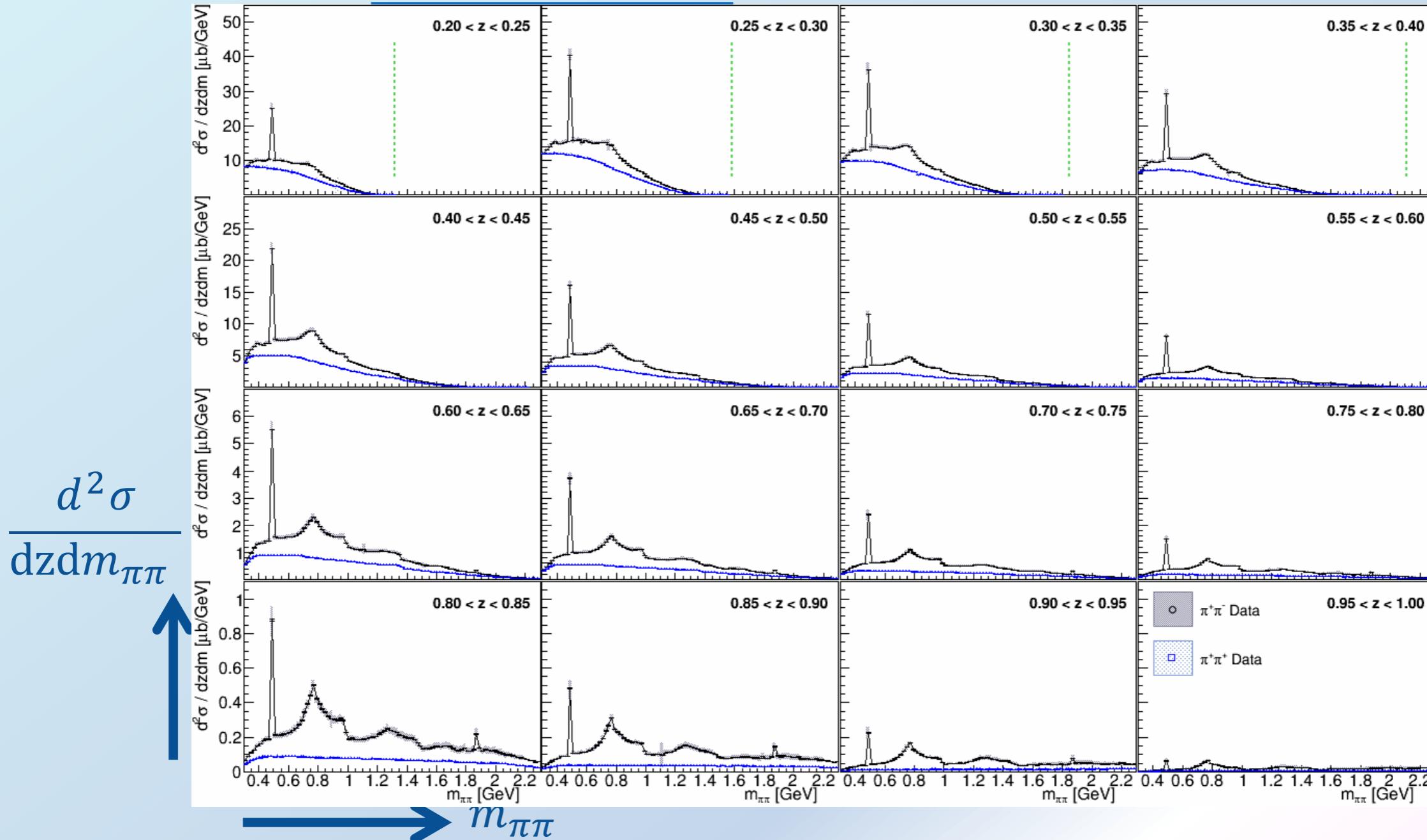
$q/h_1 h_2$	U
U	D_1
L	G_1^\perp
T	H_1^\perp H_1^\triangleleft

*unpol/spinless h !

Di-hadron mass dependence

Similar analysis in same hemisphere and mass – combined z binning. Important input for IFF based transversity global analysis

Belle: RS et.al. [arXiv:1706.08348](https://arxiv.org/abs/1706.08348)



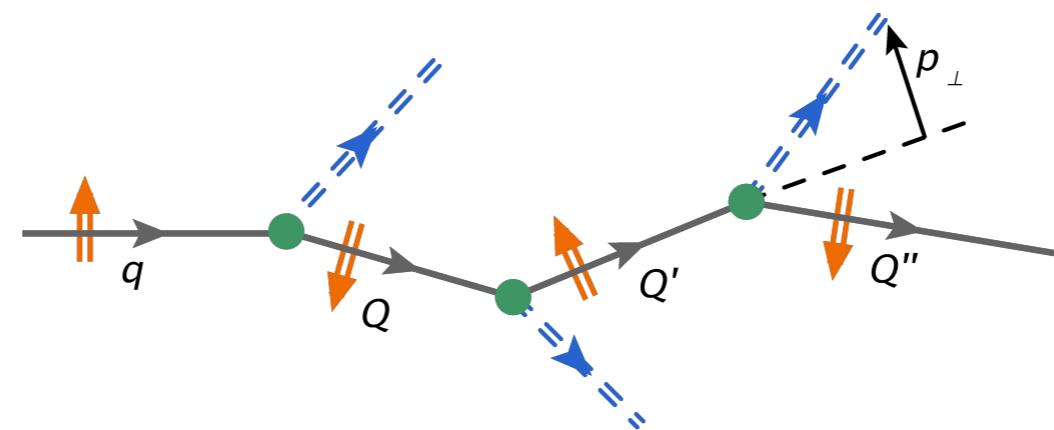
July 21, 2017

R.Seidl: Fragmentation in e+e- and EIC

Quark-jet Model

A. Kotzinian

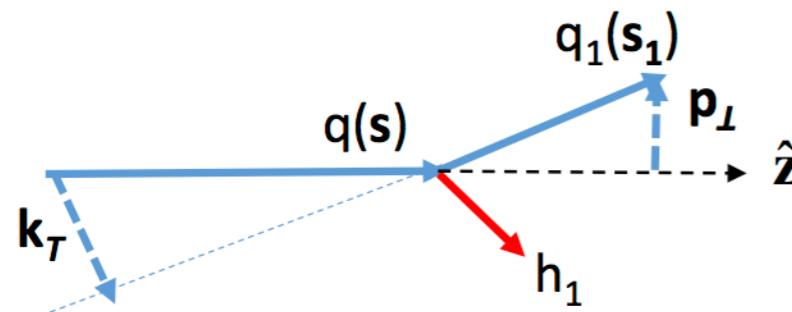
- ▶ Recursive framework for quark hadronization based on Field-Feynman model.



- ▶ Description of arbitrary quark polarization via spin density matrix formalism.
- ▶ Encode spin transfer via 8 TMD quark-to-quark elementary TMD FFs.
- ▶ Extraction of the *complete set of one- and two-hadron polarized FFs* for pions.

Quark-jet: Spin Transfer

A. Kotzinian



$$\langle \mathbf{s}_1 \rangle = \frac{\beta(z, \mathbf{p}_\perp; \mathbf{s})}{\alpha(z, \mathbf{p}_\perp; \mathbf{s})}$$

		Final quark polarization		
		U	L	T
Initial quark polarization	U	$D(z, \mathbf{p}_\perp^2)$		$-\frac{\mathbf{k}_T \times \hat{\mathbf{z}}}{M} D_T^\perp(z, \mathbf{p}_\perp^2)$
	L		$s_L G_L(z, \mathbf{p}_\perp^2)$	$s_L \frac{\mathbf{k}_T}{M} G_T(z, \mathbf{p}_\perp^2)$
	T	$-\frac{(\mathbf{k}_T \times \mathbf{s}_T) \cdot \hat{\mathbf{z}}}{M} H_T^\perp(z, \mathbf{p}_\perp^2)$	$\frac{\mathbf{k}_T \cdot \mathbf{s}_T}{M} H_L^\perp(z, \mathbf{p}_\perp^2)$	$\frac{\mathbf{s}_T H_T(z, \mathbf{p}_\perp^2)}{M} + \frac{\mathbf{k}_T (\mathbf{k}_T \cdot \mathbf{s}_T)}{M} H_T^\perp(z, \mathbf{p}_\perp^2)$

$$\begin{aligned} \alpha(z, \mathbf{p}_\perp; \mathbf{s}) &= D(z, \mathbf{p}_\perp^2) - \frac{1}{M} (\mathbf{k}_T \times \mathbf{s}_T) \cdot \hat{\mathbf{z}} H^\perp(z, \mathbf{p}_\perp^2) \\ \beta_L(z, \mathbf{p}_\perp; \mathbf{s}) &= s_L G_L(z, \mathbf{p}_\perp^2) - \frac{1}{M} (\mathbf{k}_T \cdot \mathbf{s}_T) H_L^\perp(z, \mathbf{p}_\perp^2) \\ \beta_\perp(z, \mathbf{p}_\perp; \mathbf{s}) &= -\frac{\mathbf{k}_T'}{M} D_T^\perp(z, \mathbf{p}_\perp^2) + s_L \frac{\mathbf{k}_T}{M} G_T(z, \mathbf{p}_\perp^2) \\ &\quad + \mathbf{s}_T H_T(z, \mathbf{p}_\perp^2) + \frac{\mathbf{k}_T}{M^2} (\mathbf{s}_T \cdot \mathbf{k}_T) H_T^\perp(z, \mathbf{p}_\perp^2) \end{aligned}$$

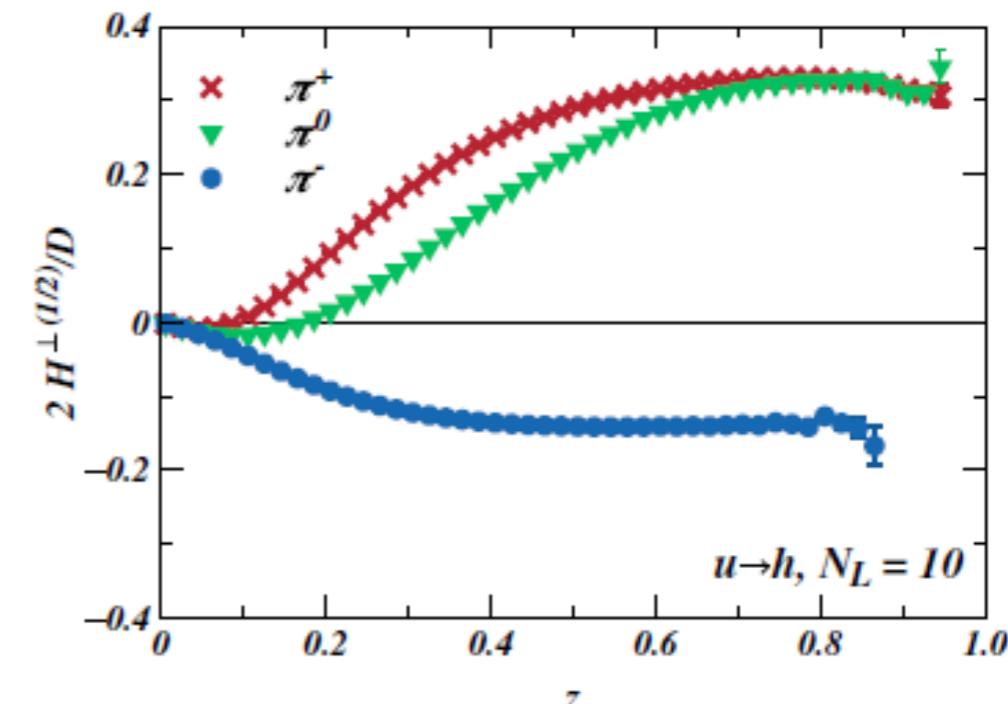
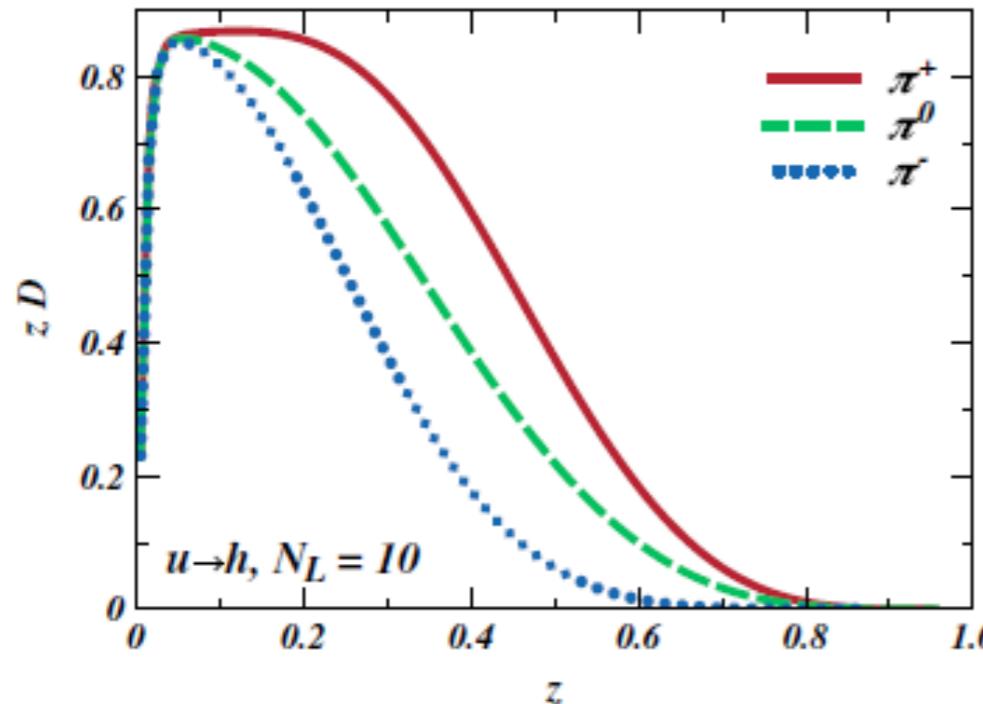
Polarized quark to unpolarized hadron SF

$$F^{q \rightarrow h_1}(z, \mathbf{p}_\perp; \mathbf{s}) = F^{q \rightarrow q_1}(1-z, -\mathbf{p}_\perp; \mathbf{s}_1 = 0, \mathbf{s}) = D(1-z, \mathbf{p}_\perp^2) + \frac{1}{M} (\mathbf{k}_T \times \mathbf{s}_T) \cdot \hat{\mathbf{z}} H^\perp(1-z, \mathbf{p}_\perp^2)$$

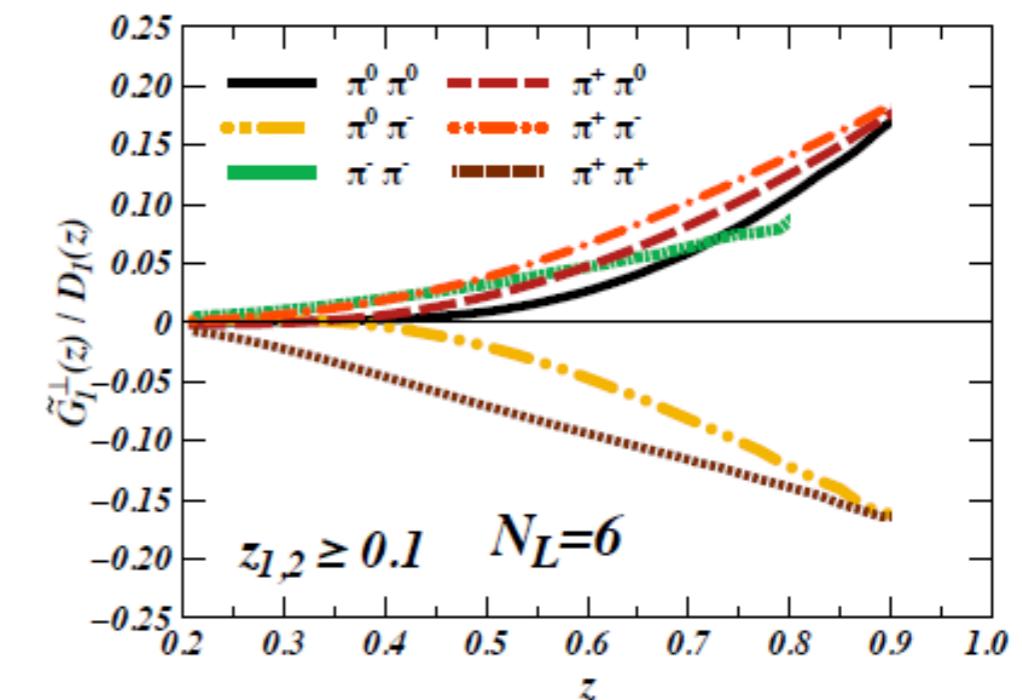
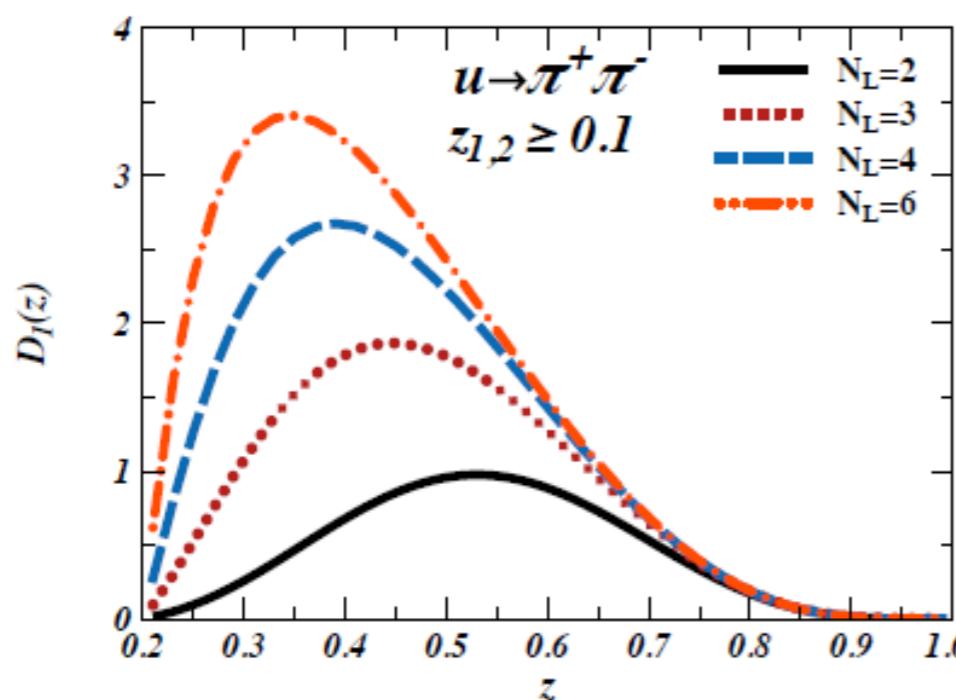
Quark-jet: Results

A. Kotzinian

One hadron: Collins effect

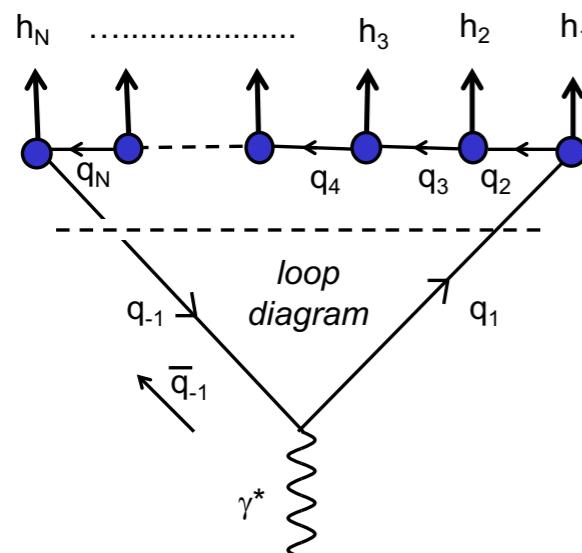


Two hadron: Helicity-dependent DiFF



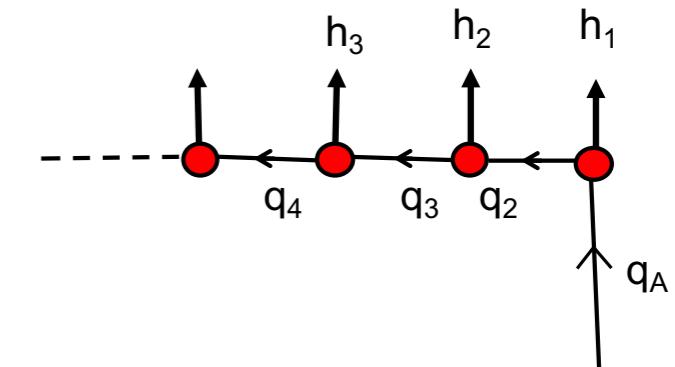
Recursive quark fragmentation with spin

- Complete description of hadronization in hard process: **Quark Multiperipheral model satisfying LR symmetry.**

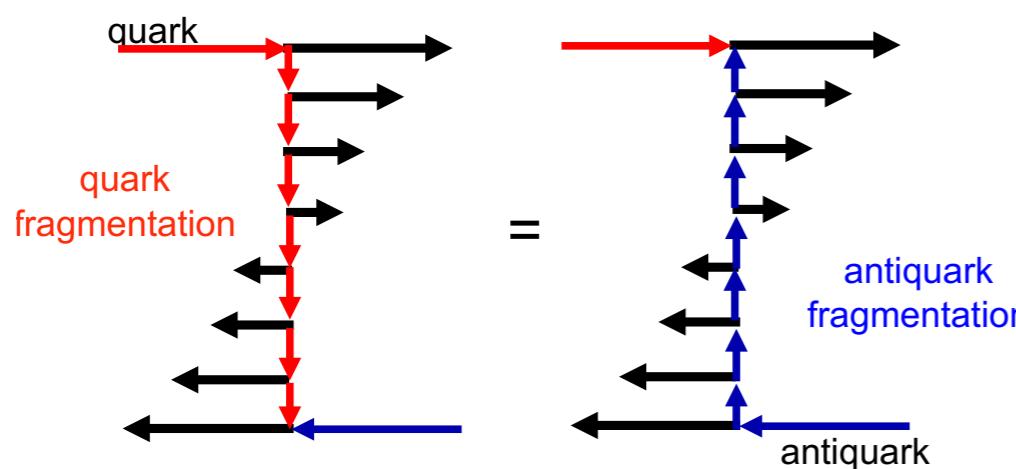


Recursive model

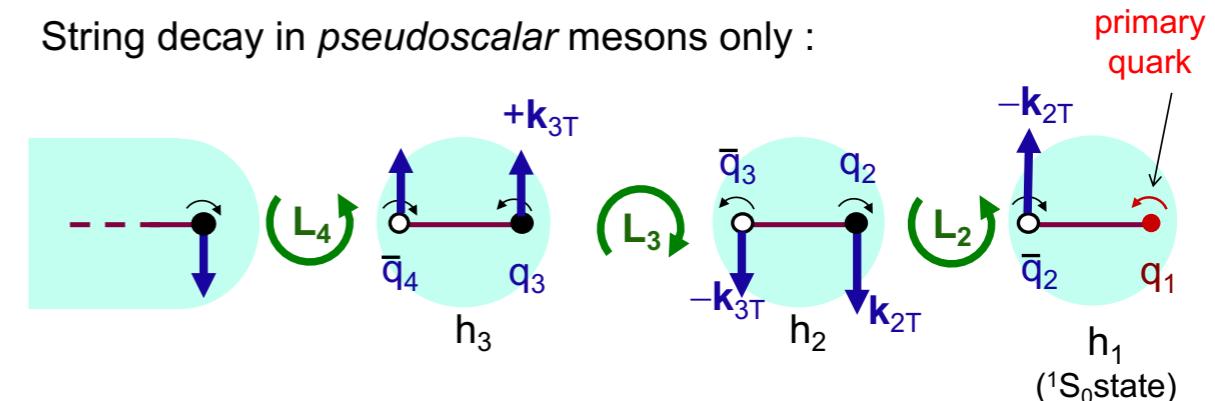
$$\begin{aligned} q_A(k_A) &\rightarrow h_1(p_1) + q_2(k_2), \\ q_2(k_2) &\rightarrow h_2(p_2) + q_3(k_3), \\ \dots \text{ etc.} \end{aligned}$$



Symmetry under *quark Line Reversal* (or “Left-Right” symmetry)



Iteration of the *string + 3P_0* mechanism



Synthesis of string and 3P0 inputs (pseudoscalar mesons only)

$$\begin{aligned}
 T_{q'hq}(k', k) &= C_{q'hq} && \xleftarrow{\text{Flavor}} \\
 \times \exp(-\frac{1}{2} b_T k'_T)^2 & && \xleftarrow{\text{transverse momentum cutoff}} \\
 \times (1-Z)^{a/2} \exp(-b_L \varepsilon^2 / 2Z) & && \xleftarrow{\text{string model}} \\
 \times (\mu + \sigma_z \boldsymbol{\sigma} \cdot \mathbf{k}'_T) & && \xleftarrow{^3P_0} \\
 \times \sigma_z & && \xleftarrow{\text{Pseudoscalar}} \\
 \times \hat{U}_q^{-1/2}(k_T) & && \xleftarrow{\text{reduced single quark density matrix}}
 \end{aligned}$$

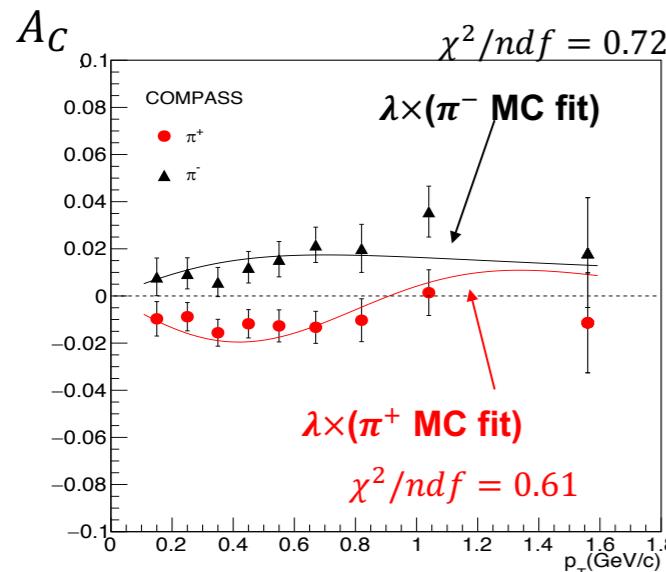
$$\hat{U}_q(k_T) = \sum_h \int dZ/Z \, d^2 p_T \, t^\dagger \, t$$

$$= [\hat{U}_0(k_T^2) + \hat{U}_1(k_T^2) \boldsymbol{\sigma} \cdot (\mathbf{z} \times \mathbf{k}_T) / |\mathbf{k}_T|]$$

($t = T$ without the last line)

Comparison with COMPASS data

Single hadron: Collins asymmetry



- λ is a scale parameter estimated minimizing

$$\chi^2 = \sum_i \frac{(A_{C,i} - \lambda f_{\pi^-}(p_{T,i}))^2}{\sigma_i^2}$$

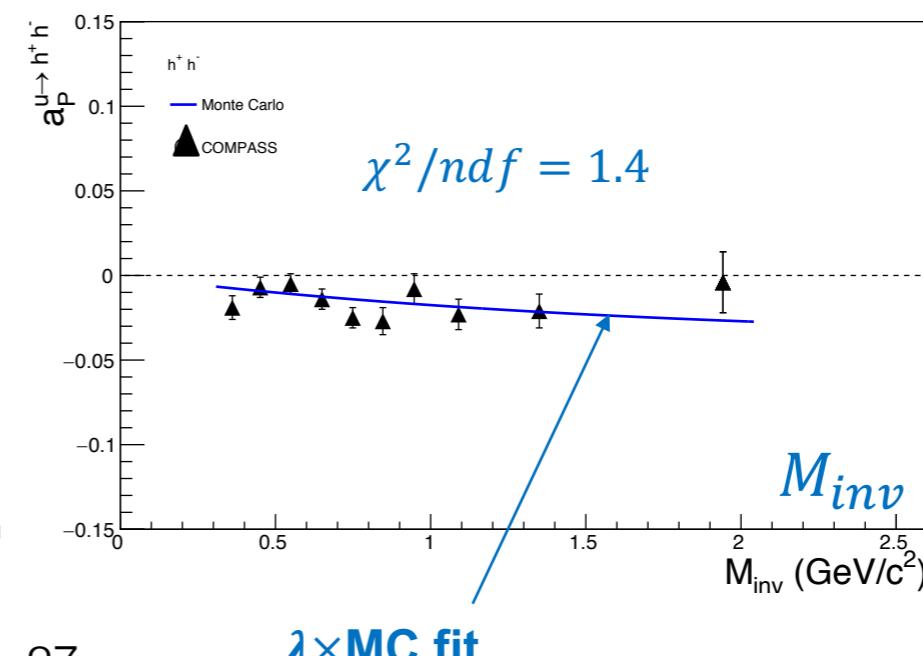
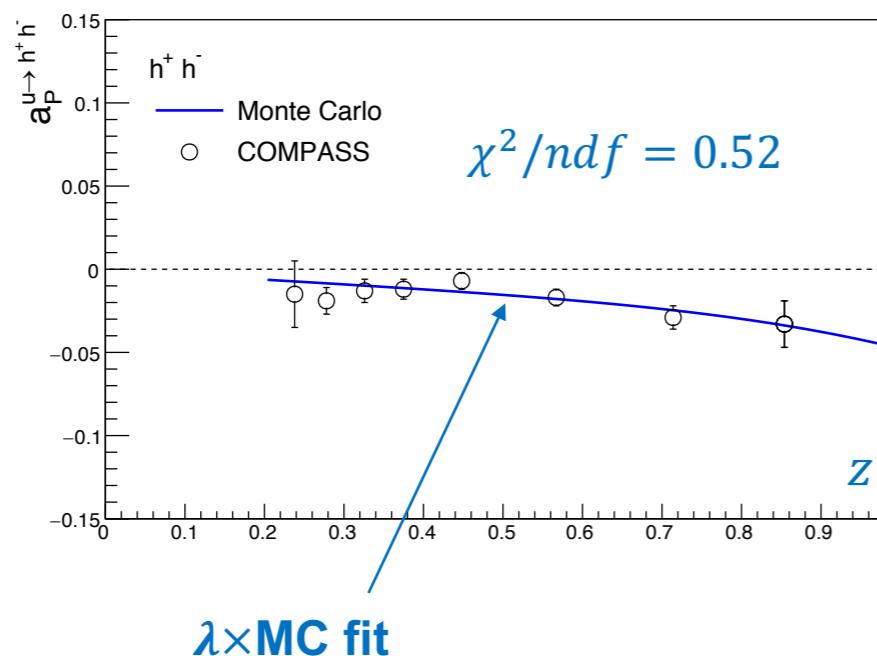
- $\lambda \sim \langle h_1^u/f_1^u \rangle$

- $f_{\pi^-}(p_T)$: preceding MC fit for π^-
- A_C : experimental value of the π^- asymmetry
- σ : statistical error

$$\lambda = 0.055 \pm 0.010$$

Dihadron: IFF

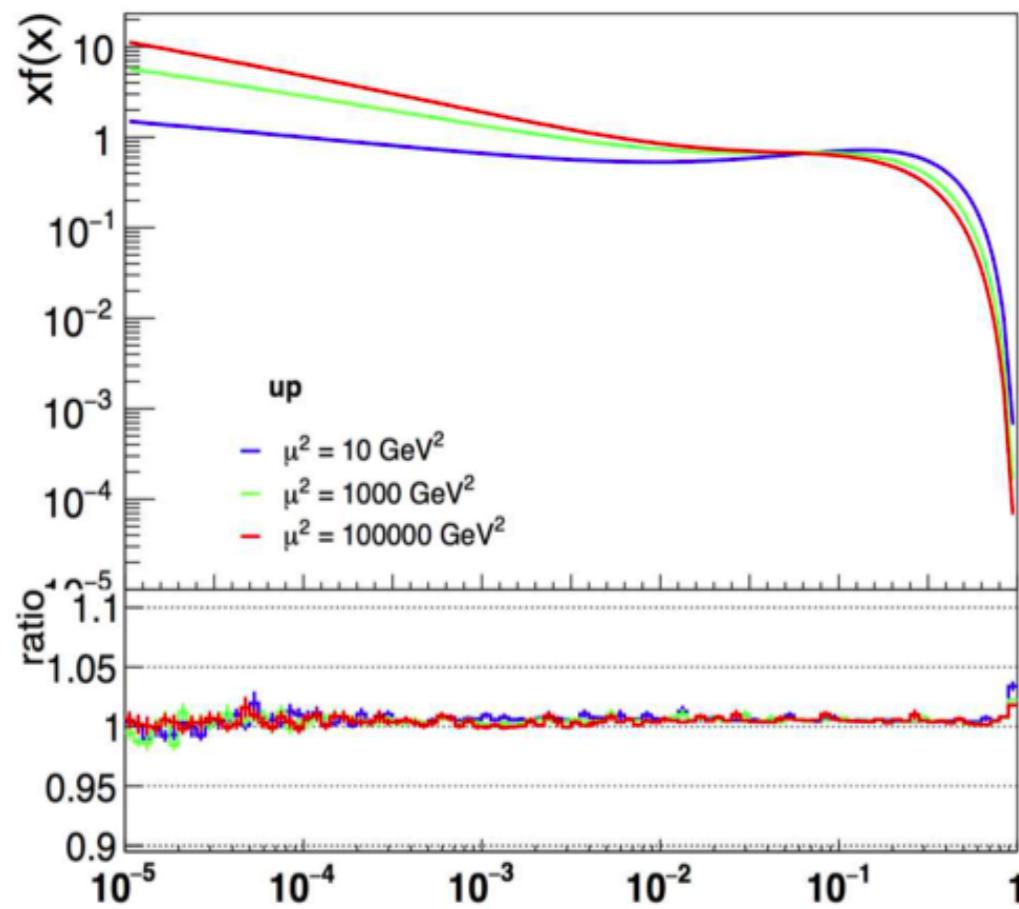
- Same λ also for 2h asymmetries



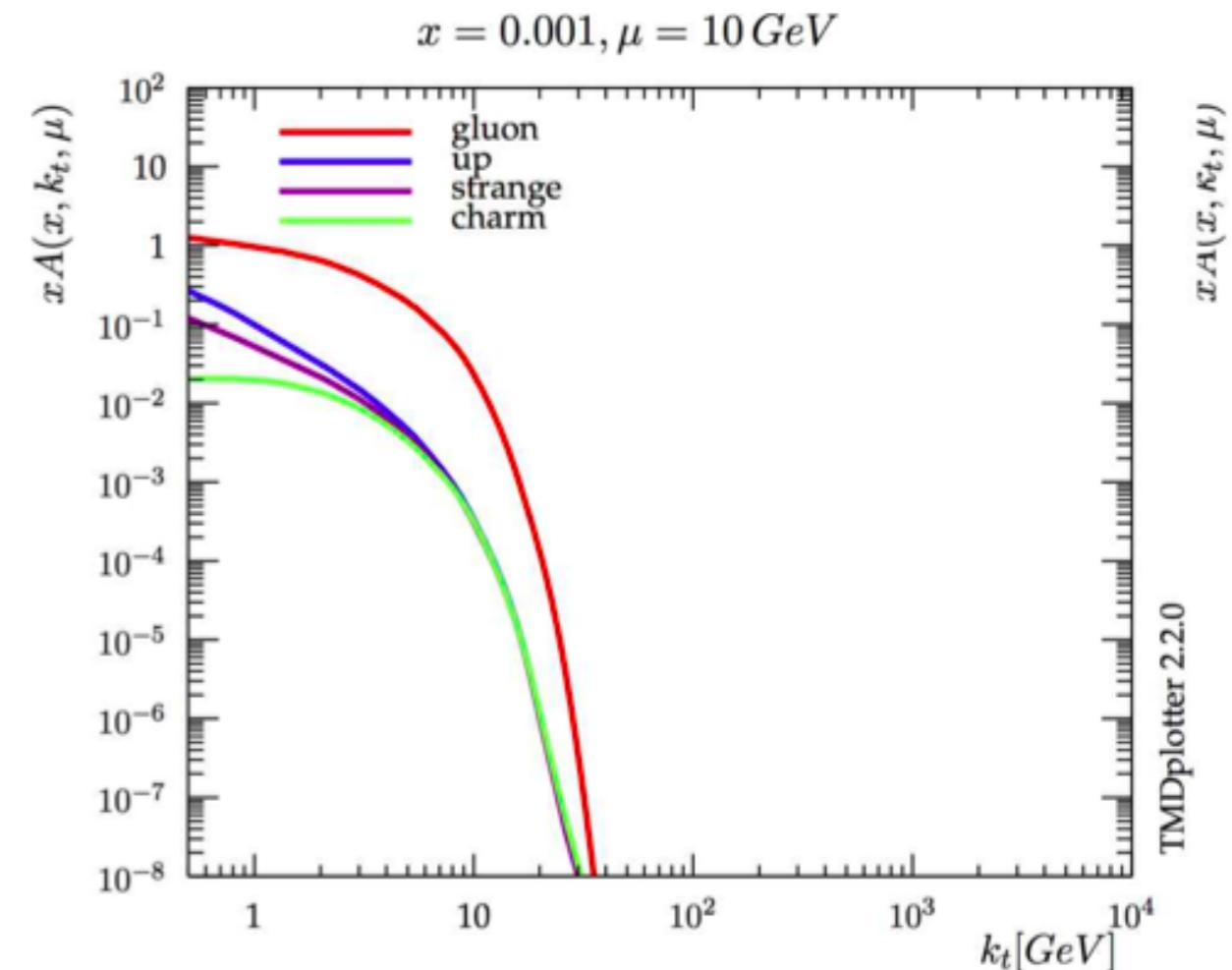
Parton Branching Solution for QCD Evolution Equations

- ❖ Consistent with “orthodox” solutions up to $NNLO$ at high precision.
- ❖ Provides the complete final partonic state with k_T .
- ❖ Determination of TMD for all flavours up to NLO: NO free parameters.
- ❖ TMD evolution implemented in xFitter - applicable for DIS processes.

Validation: vs DGLAP at NLO

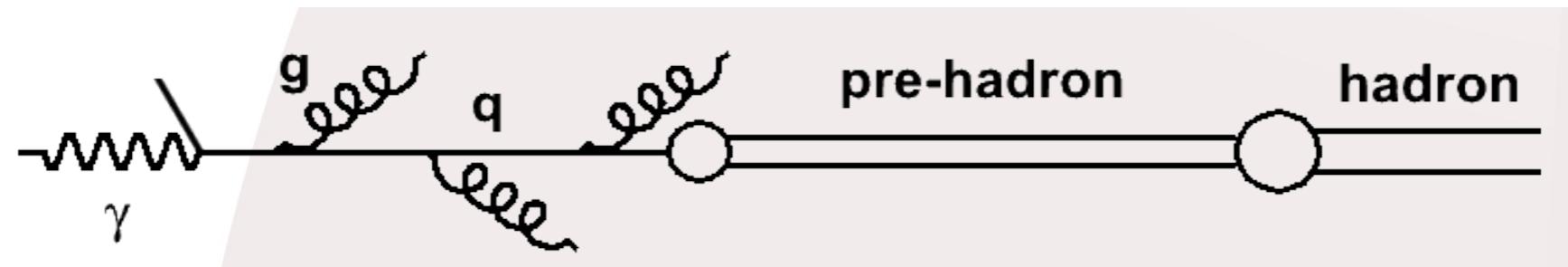


All flavours of TMDs

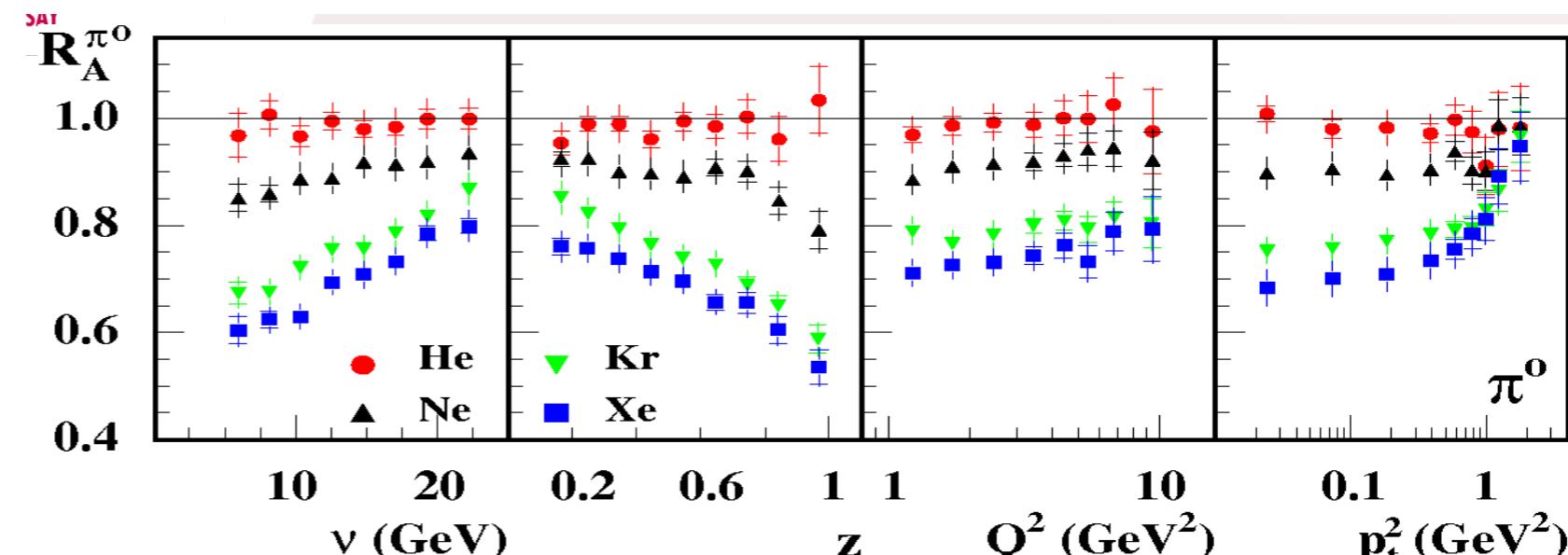


Hadronization in Nuclear Medium

R. Dupre

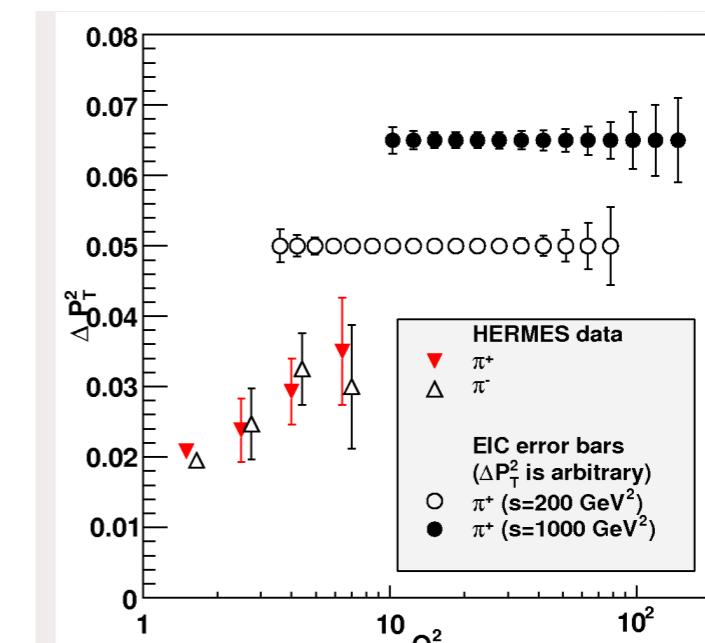
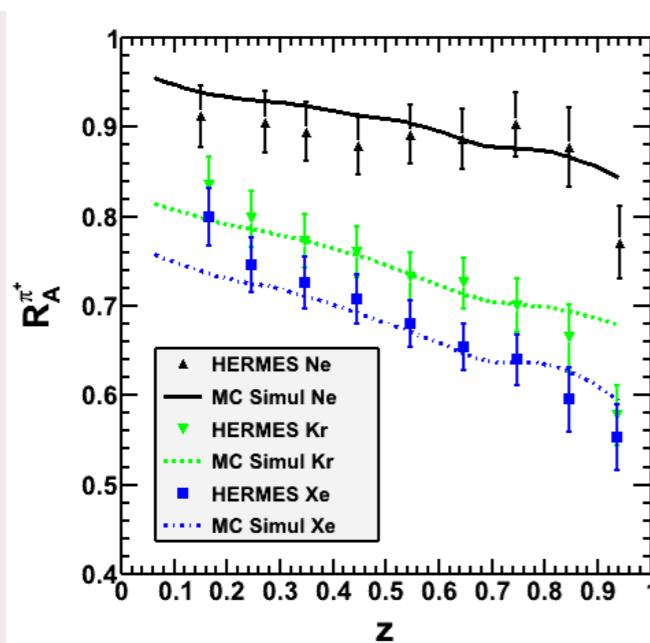


HERMES data on multiplicity ratios to D



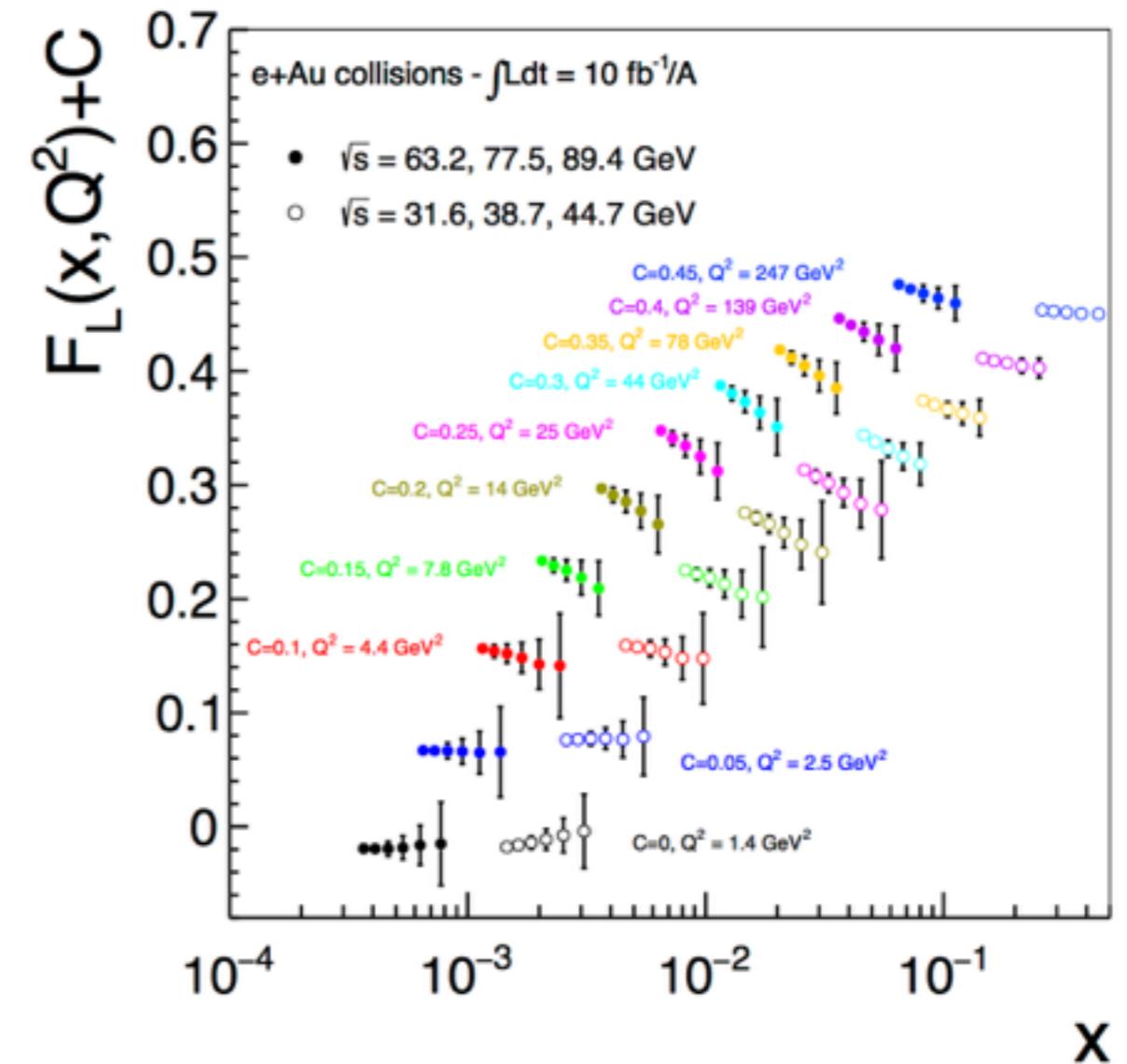
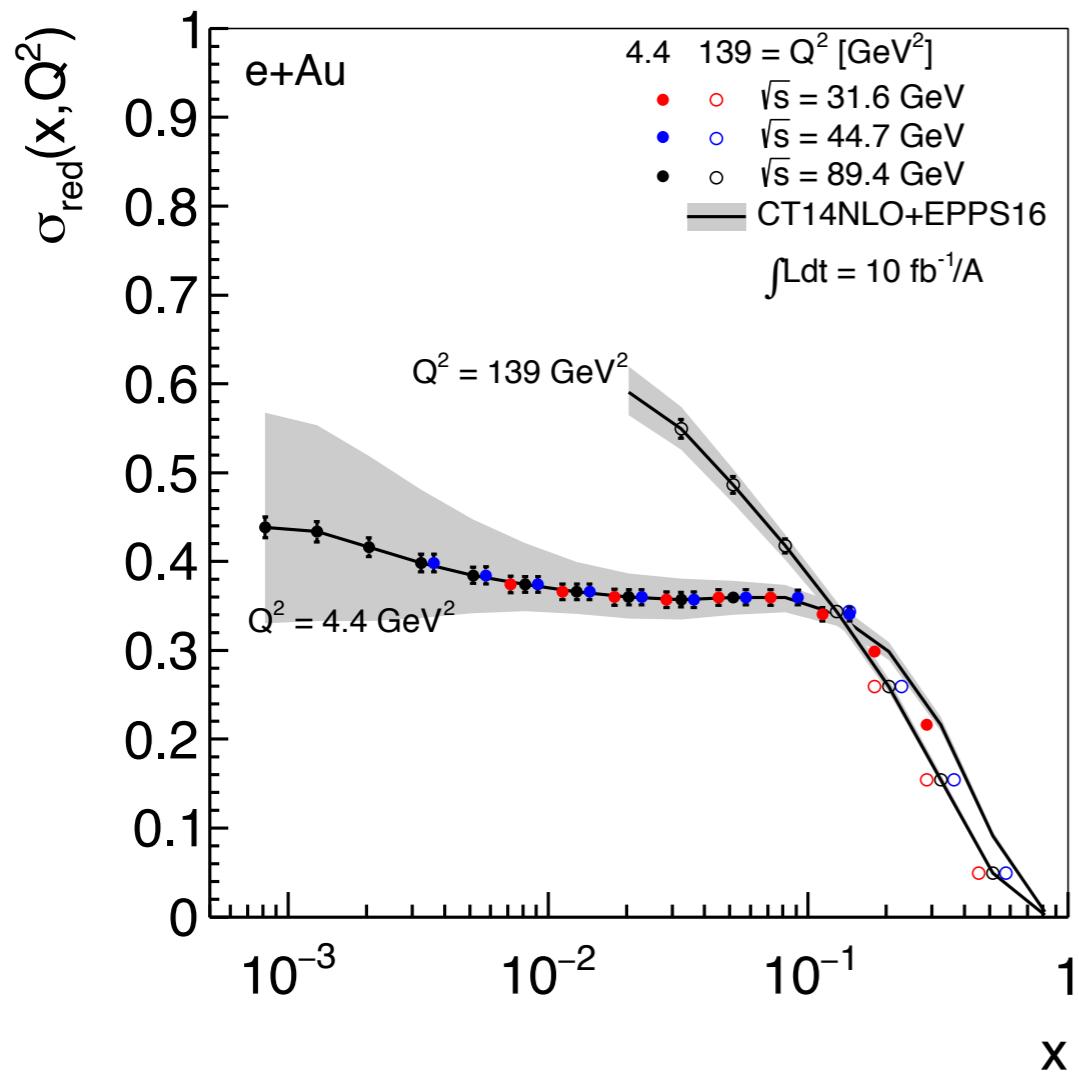
MC with parton energy loss in medium

Q^2 evolution with EIC projections



Impact of EIC on nuclear PDFs

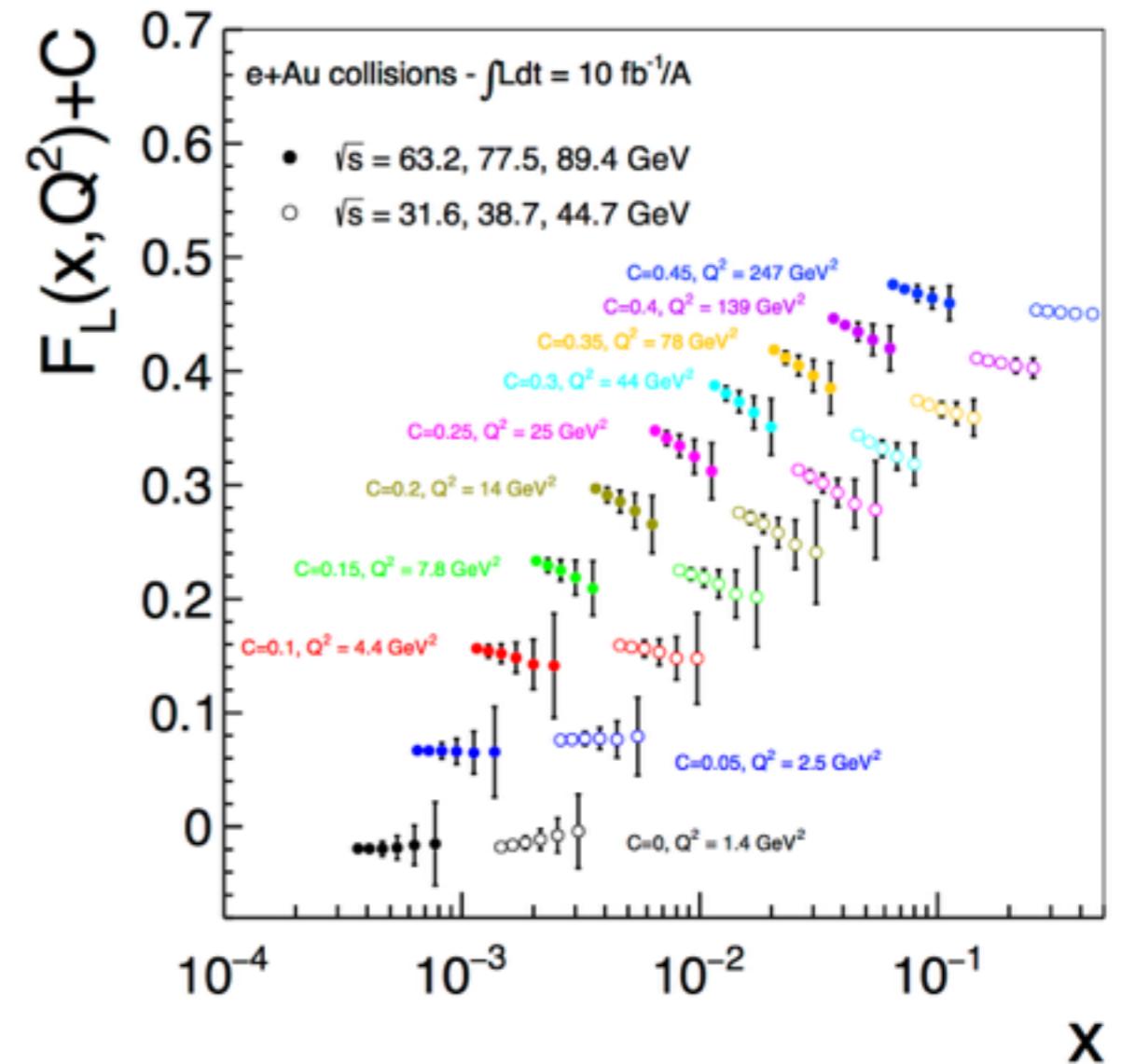
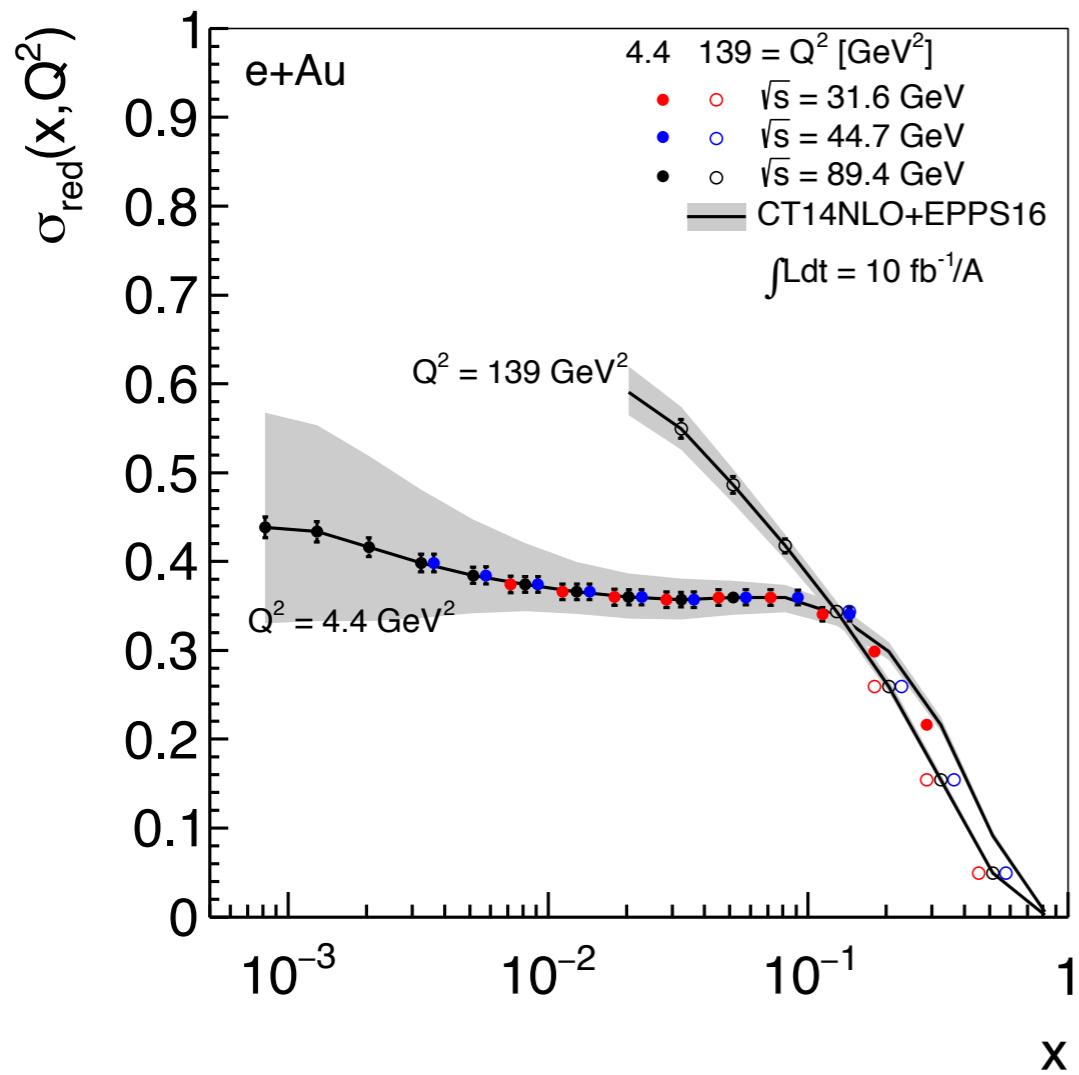
S. Fazio



An EIC at its highest energy provides a factor 10 larger reach in Q^2 and low- x compared to available data

Impact of EIC on nuclear PDFs

S. Fazio

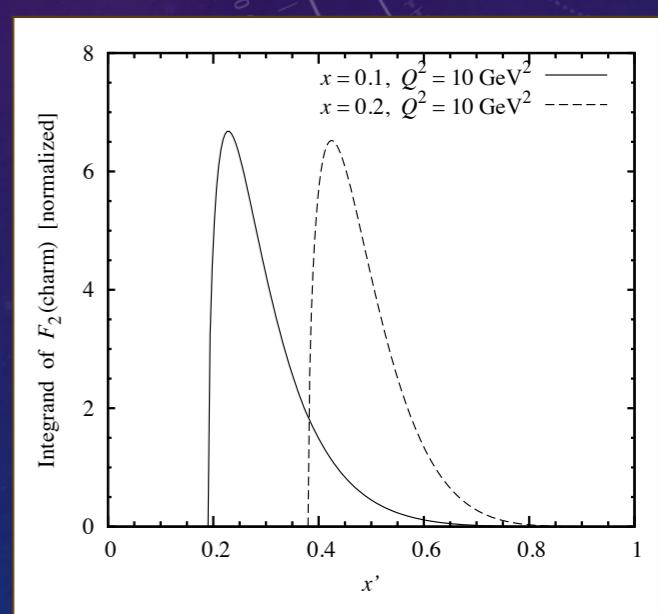
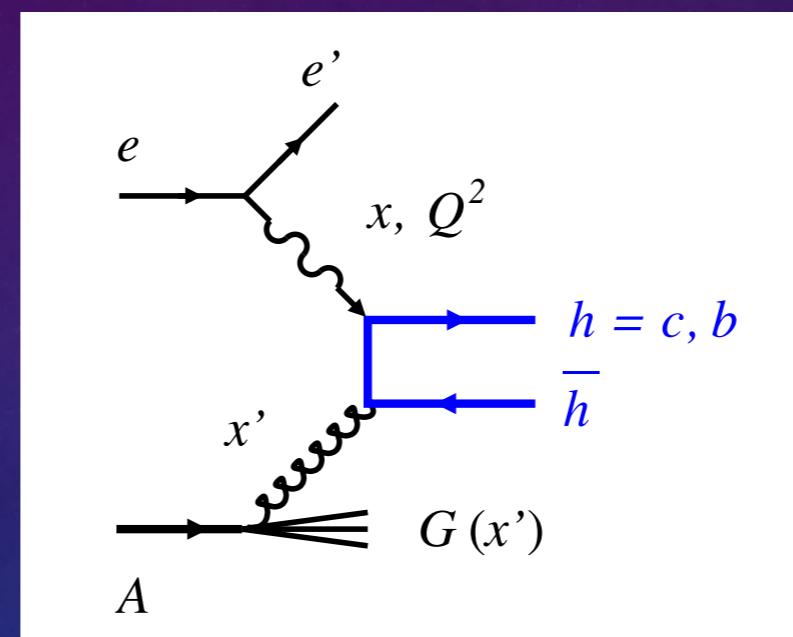


An EIC at its highest energy provides a factor 10 larger reach in Q^2 and low- x compared to available data
additional charm tagging will further help constraining gluon especially at high (gluon) x

Nuclear gluon densities via open charm

OPEN CHARM: A DIRECT PROBE OF GLUONS

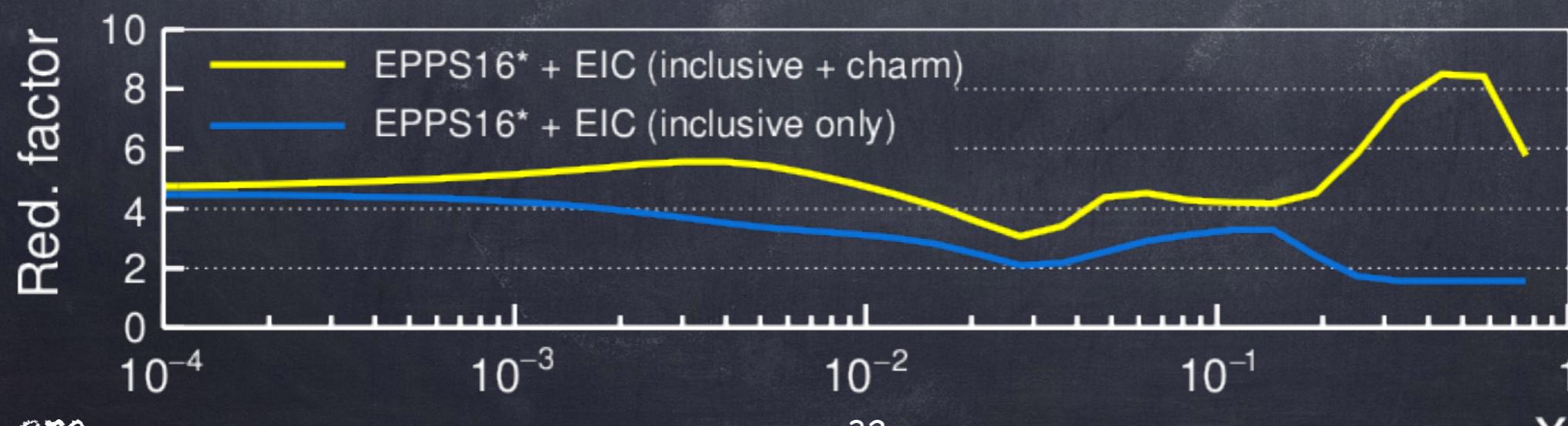
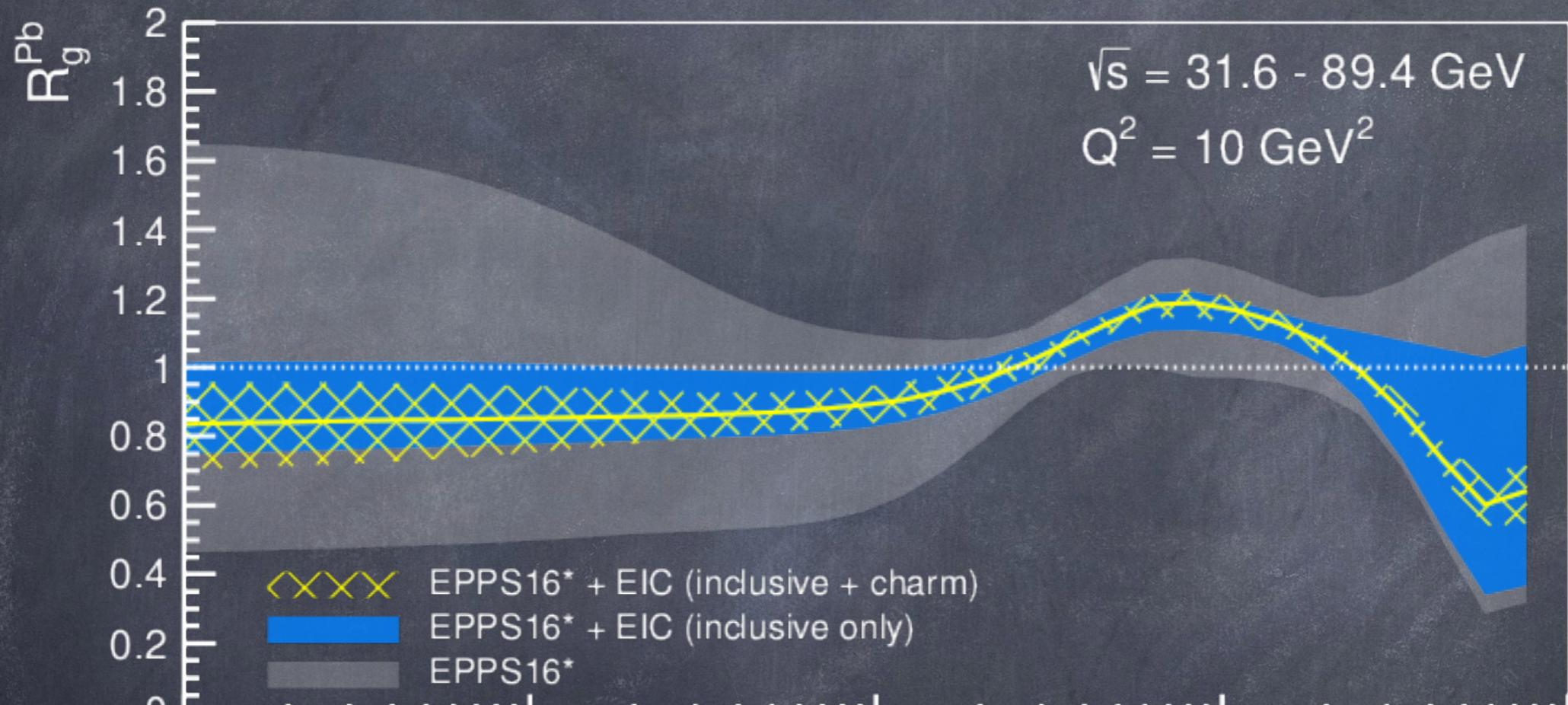
- EIC:
 - Nuclei
 - CM Energy
 - Luminosity
 - Count rate reach to “high”- x
 - Small branching fractions
 - Particle ID, vertexing



$$x_{\text{gluon}} = x' \geq \frac{4m_h^2 + Q^2}{W^2} \quad W^2 \gg Q^2$$

with a higher energy realization of an EIC

P. Zurita

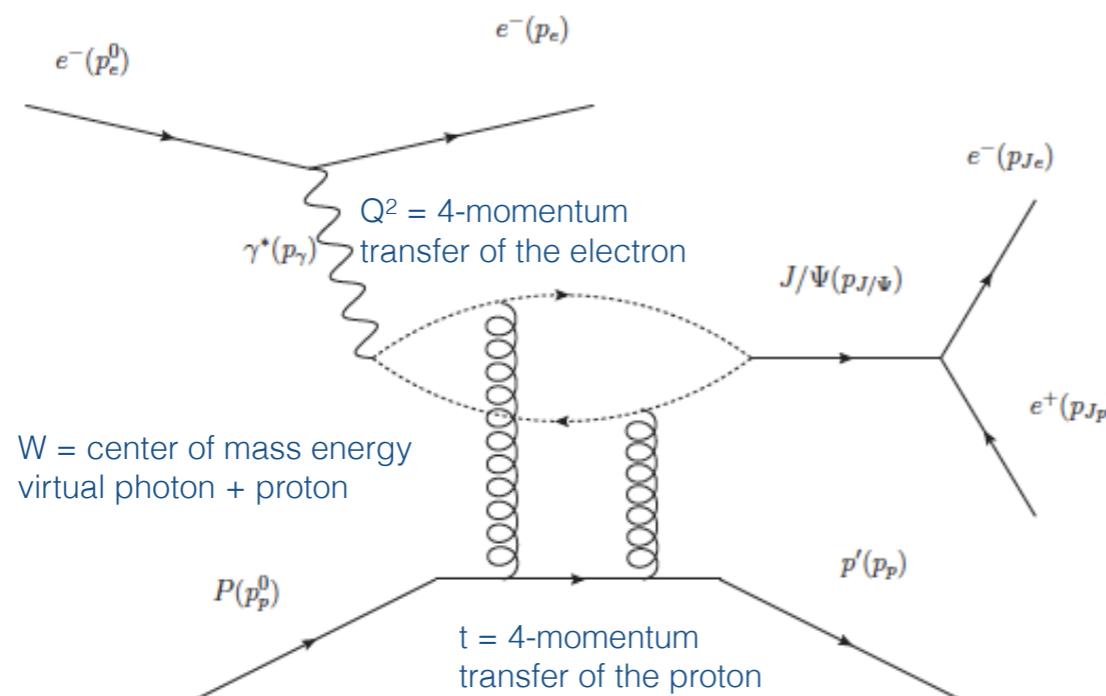


MODELS RELATE TRACE ANOMALY TO J/ Ψ PRODUCTION NEAR THRESHOLD

e.g. D. Kharzeev, EPJC9 459 (1999)

N. Fegege

Photo-production of J/ Ψ



J/ Ψ and Y production near threshold at EIC

