

Recent results on NNPDF parton distribution and fragmentation functions

Electron Ion Collider User Group Meeting 2017

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Foreword

- ① Guiding principle: leading-twist collinear factorisation

$$\mathcal{O}_I = \sum_{f=q,\bar{q},g} C_{If}(y, \alpha_s(\mu^2)) \otimes f(y, \mu^2) + \text{p.s. corrections} \quad f \otimes g = \int_x^1 \frac{dy}{y} f\left(\frac{x}{y}\right) g(y)$$

- ② Parametrisation: general, smooth, flexible at an initial scale Q_0^2

$$x f_i(x, Q_0^2) = A_{f_i} x^{a_{f_i}} (1-x)^{b_{f_i}} \mathcal{F}(x, \{c_{f_i}\}) \quad \mathcal{F}(x, \{c_{f_i}\}) \text{ is a neural network}$$

- ③ A prescription to determine/compute expectation values and uncertainties

$$E[\mathcal{O}] = \int \mathcal{D}\Delta f \mathcal{P}(\Delta f | data) \mathcal{O}(\Delta f) \quad V[\mathcal{O}] = \int \mathcal{D}\Delta f \mathcal{P}(\Delta f | data) [\mathcal{O}(\Delta f) - E[\mathcal{O}]]^2$$

Monte Carlo: $\mathcal{P}(\Delta f | data) \longrightarrow \{\Delta f_k\}$

$$E[\mathcal{O}] \approx \frac{1}{N} \sum_k \mathcal{O}(\Delta f_k) \quad V[\mathcal{O}] \approx \frac{1}{N} \sum_k [\mathcal{O}(\Delta f_k) - E[\mathcal{O}]]^2$$

Unpolarised PDFs

NNPDF3.1

[arXiv:1706.00428]

Helicity PDFs

NNPDFPOL1.1

[NP B887 276]

Fragmentation functions

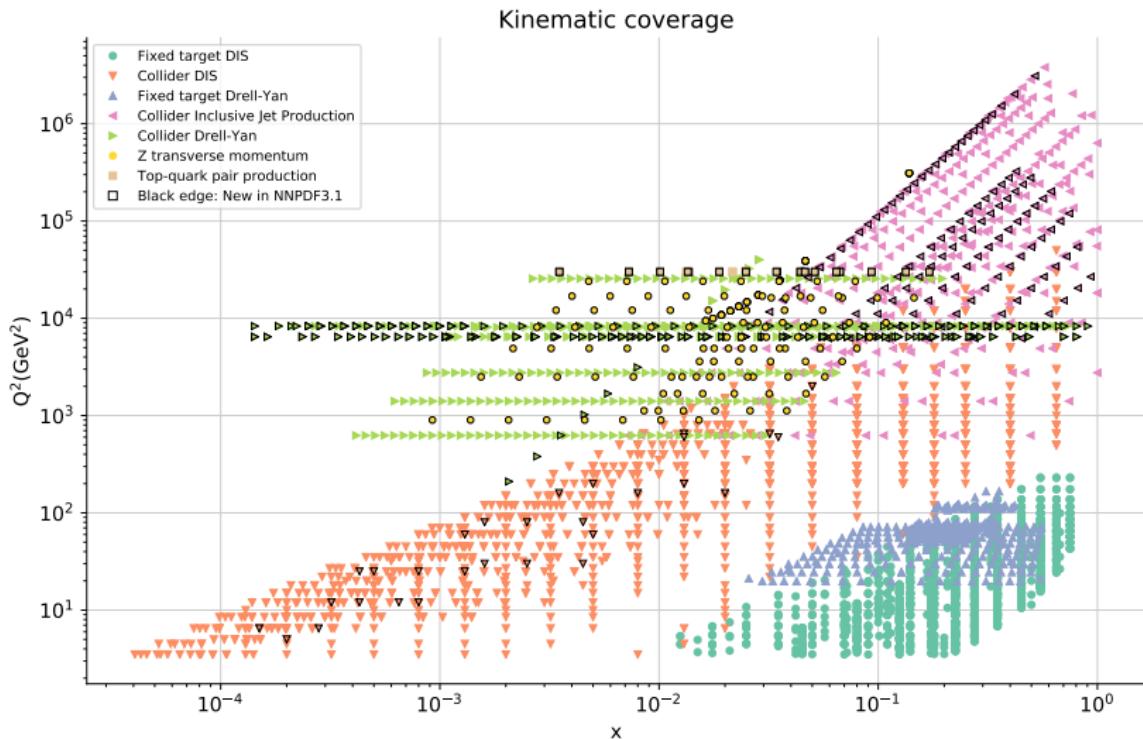
NNFF1.0 (π, K, p)

[arXiv:1706.07049]

ALL PDF/FF sets are available through LHAPDF: <http://lhapdf.hepforge.org/>

1. NNPDF3.1 [arXiv:1706:00428, submitted to EPJC]

The dataset

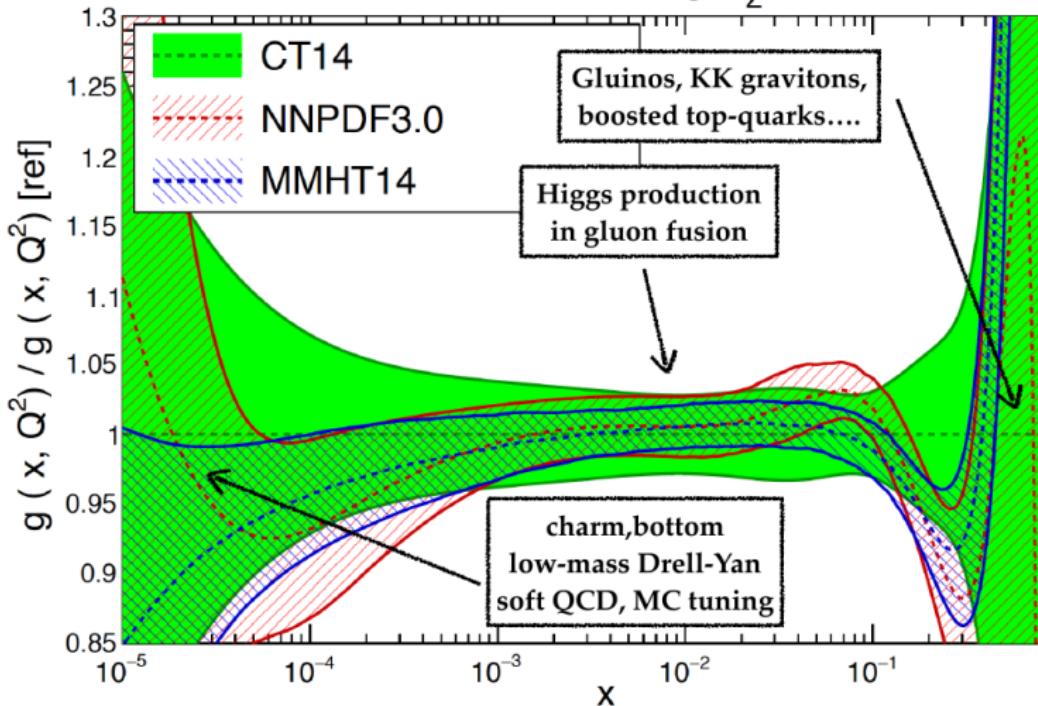


$$N_{\text{dat}} = 4295$$

$$\chi^2/N_{\text{dat}} = 1.168/1.148 \text{ (NLO/NNLO)}$$

The gluon PDF

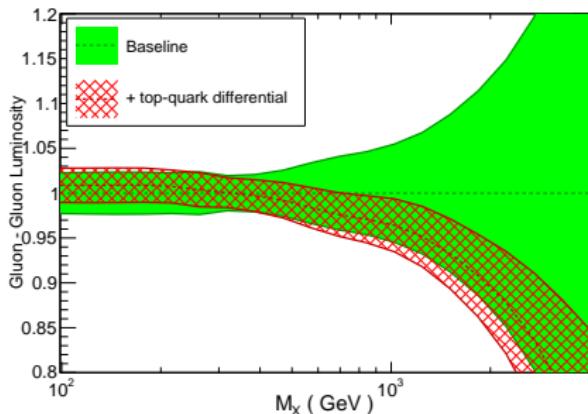
NNLO, $Q^2=100 \text{ GeV}^2$, $\alpha_S(M_Z)=0.118$



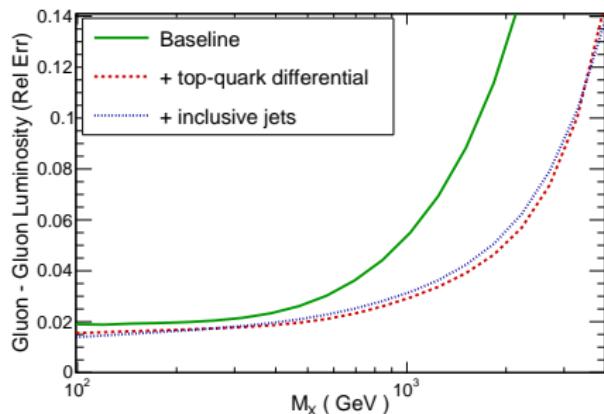
A precise knowledge of the gluon PDF is required over all the range in x to exploit the full potential of the LHC

The gluon PDF at large x : $t\bar{t}$ differential distributions

NNLO, global fits, LHC 13 TeV



NNLO, global fits, LHC 13 TeV



ATLAS and CMS rapidity distributions at $\sqrt{s} = 8$ TeV

Significant reduction of gg luminosity uncertainties at $M_X \geq \mathcal{O}(1)$ TeV
e.g., at $M_X \sim 2$ TeV, uncertainties decrease from 13% to 5%

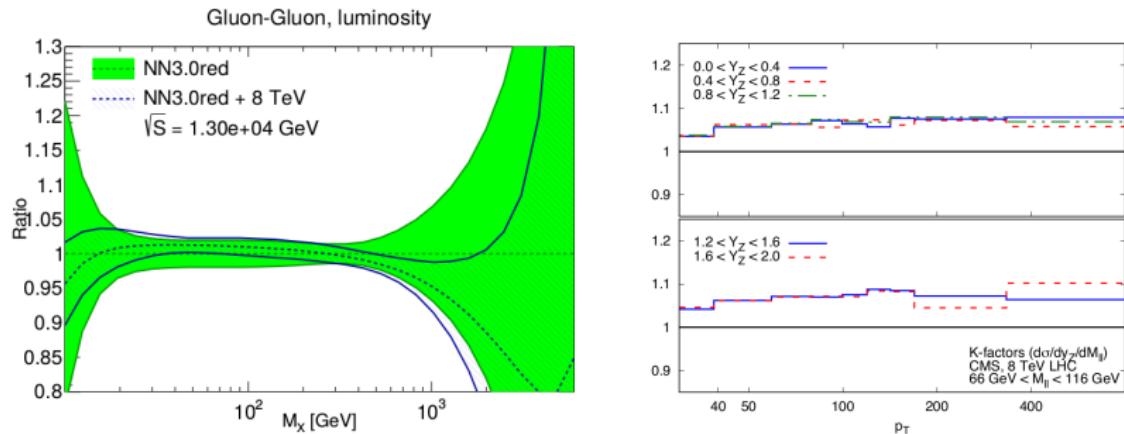
Impact of $t\bar{t}$ differential data similar to that of jet data

though jet data analysed neglecting NNLO QCD corrections in the matrix element

A precision determination of the gluon PDF at large x is now possible at NNLO
the situation should only improve thanks to the recent NNLO jet calculation
 $t\bar{t}$ differential distributions are included in the NNPDF3.1 PDF release

[see JHEP 1704 (2017) 044 and 1706.00428 for details]

The gluon PDF at medium x : the Z -boson p_T distribution



ATLAS and CMS p_T distributions at $\sqrt{s} = 8$ TeV
in various rapidity bins in the Z -peak region

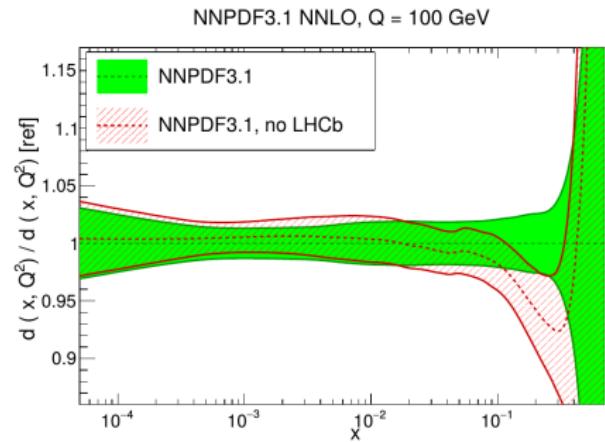
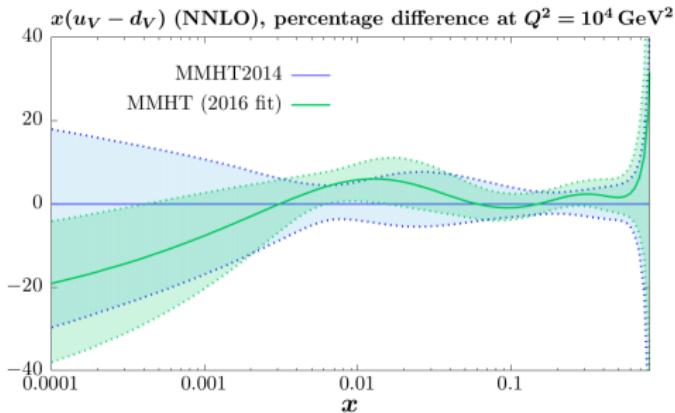
NNLO/NLO K -factors 5%-10% depending on the rapidity/invariant mass region
challenge: measurements have sub-percent experimental errors

Complementary information on the gluon PDF
e.g., at $M_X \sim 2$ TeV, uncertainties decrease from 13% to 8%

Z p_T distributions are included in the NNPDF3.1 PDF release

[see 1705.00343 and 1706.00428 for details]

Quark flavour separation from LHC data



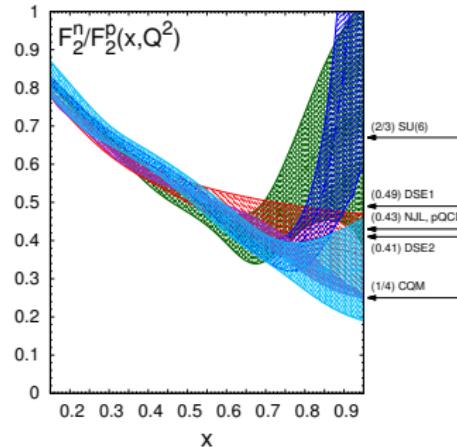
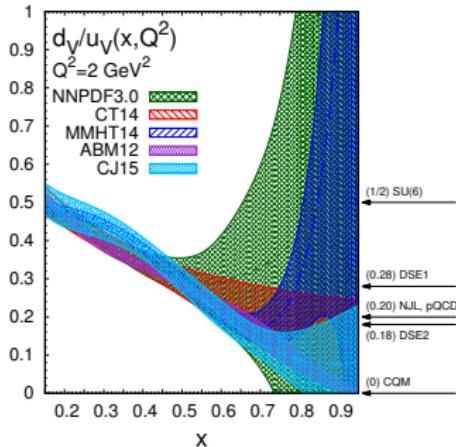
High-precision W and Z production data from ATLAS, CMS and LHCb handle on quark/antiquark flavour separation

Largest impact on light quarks at large x provided by LHCb data
error reduction by a factor 2 in NNPDF3.1 at $x \sim 0.1$

Combined effect of (LHC) CMS, LHCb and (Tevatron) D0 W, Z data
improved determination of $x(u_V - d_V)$

[see R. Thorne's talk at DIS2017 and 1706.00428 for details]

The PDF ratio d_V/u_V at large x [EPJC 76 (2016) 383]



$$\frac{d_V}{u_V} \xrightarrow{x \rightarrow 1} (1-x)^{b_{d_V}} e^{-b_{u_V}} \xrightarrow[\text{c. r.}]{b_{d_V} = b_{u_V}} k$$

$$\frac{F_2^n}{F_2^p} \xrightarrow{x \rightarrow 1} \frac{4(1-x)^{b_{u_V}} + (1-x)^{b_{d_V}}}{(1-x)^{b_{u_V}} + 4(1-x)^{b_{d_V}}} \xrightarrow[\text{c. r.}]{b_{d_V} = b_{u_V}} 1$$

$$\text{case } b_{u_V} \gg b_{d_V} : \frac{d_V}{u_V} \xrightarrow{x \rightarrow 1} \infty; \quad \frac{F_2^n}{F_2^p} \xrightarrow{x \rightarrow 1} 4 \quad \text{case } b_{u_V} \ll b_{d_V} : \frac{d_V}{u_V} \xrightarrow{x \rightarrow 1} 0; \quad \frac{F_2^n}{F_2^p} \xrightarrow{x \rightarrow 1} \frac{1}{4}$$

No predictive power from current PDF determinations, no discrimination among models unless $\frac{d_V}{u_V} \xrightarrow{x \rightarrow 1} k$ is built in the parametrisation (CT14, CJ16, ABM12)

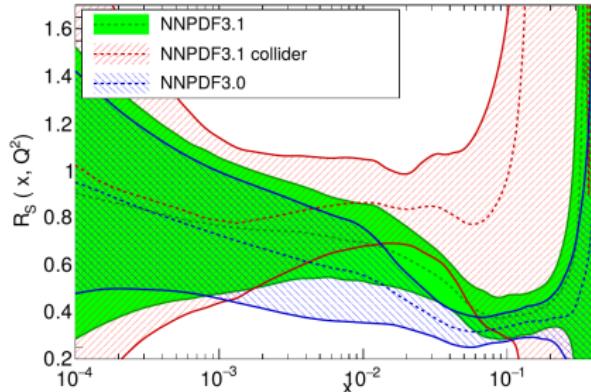
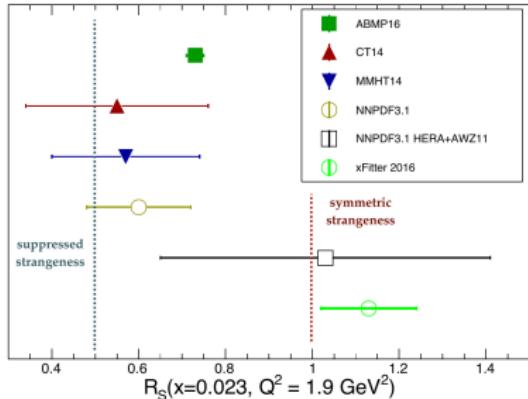
The EIC may measure the ratio F_2^n/F_2^p with high accuracy, provided neutron beams expected to be less prone to nuclear and/or higher twist corrections than fixed-target DIS

Complementary measurements from the LHC (DY) and (particularly) the LHeC (DIS)

The strange PDF from collider data

$$R_s = (s + \bar{s}) / (\bar{u} + \bar{d})$$

NNLO, $Q=1.38 \text{ GeV}$



In most PDF fits the strange PDF is suppressed w.r.t up and down sea quark PDFs effect mostly driven by neutrino dimuon data

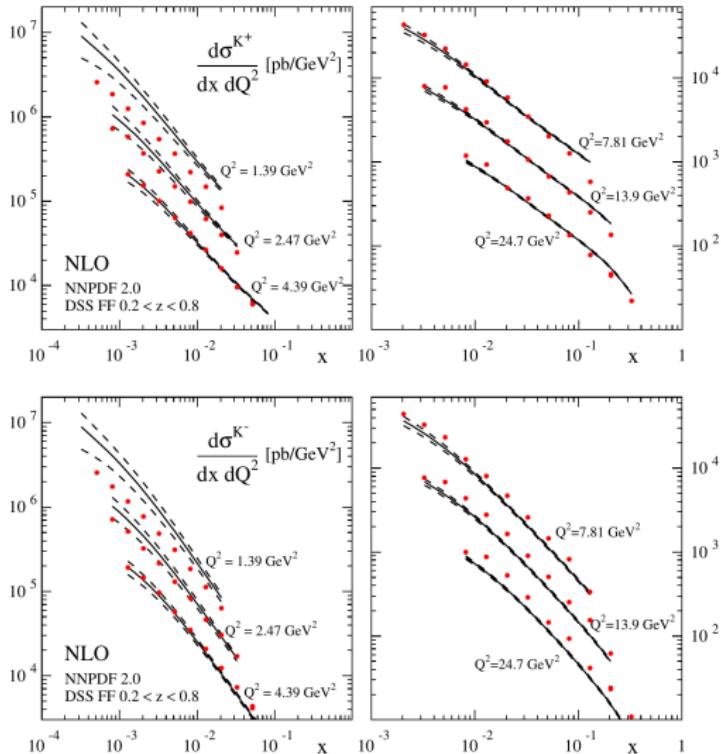
A symmetric strange sea PDF is preferred by recent collider data
in particular by ATLAS W, Z rapidity distributions (2011)

$$R_s(x, Q^2) = \frac{s(x, Q^2) + \bar{s}(x, Q^2)}{\bar{u}(x, Q^2) + \bar{d}(x, Q^2)} \left\{ \begin{array}{l} \sim 0.5 \text{ from neutrino and CMS } W + c \text{ data} \\ \sim 1.0 \text{ from ATLAS } W, Z \end{array} \right.$$

The new ATLAS data can be accommodated easily in the global fit
increased strangeness, though not as much as in a collider-only fit
some tension remains between collider and neutrino data

[see R. Thorne's talk at DIS2017 and 1706.00428 for details]

The strange PDF: K^\pm production in SIDIS at an EIC



[figure taken from arXiv:1108:1713]

red points: pseudodata at an EIC
(based on PYTHIA + JETSET)

black curves: theory predictions
(NNPDF2.0 + DSS07, NLO)

$0.01 \leq y \leq 0.95$, $\sqrt{s} = 70.7$ GeV
 z integrated in the range $[0.2, 0.8]$

small x : $d\sigma^{K^+} \approx d\sigma^{K^-}$

large x : $d\sigma^{K^+} \gg d\sigma^{K^-}$

may constrain s^+ and s^-

drawback: K^\pm fragmentation

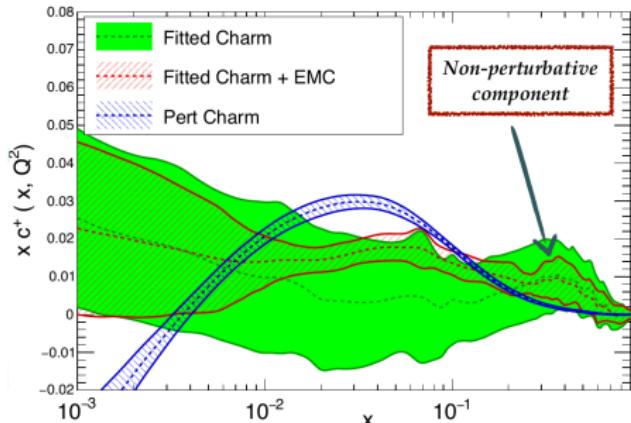
- study FFs separately
- analyse PDFs and FFs simultaneously

LHeC: direct sensitivity to s
charm tagging in CC DIS ($W + s \rightarrow c$)

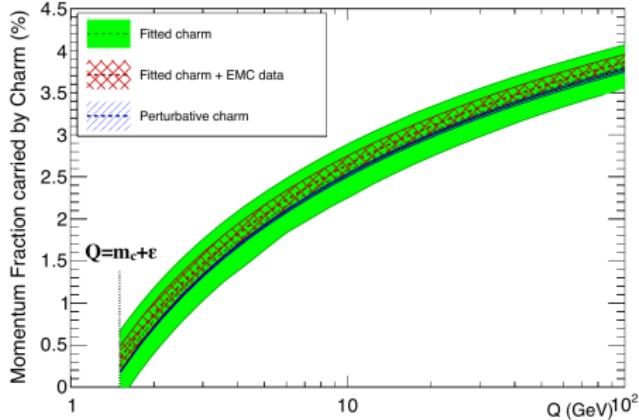
π^\pm production in SIDIS at an EIC
allow for a determination of $\bar{u} - \bar{d}$

The charm PDF: perturbative vs fitted

NNPDF3.1 NNLO, $Q = 1.7 \text{ GeV}$



NNPDF3.1 NNLO

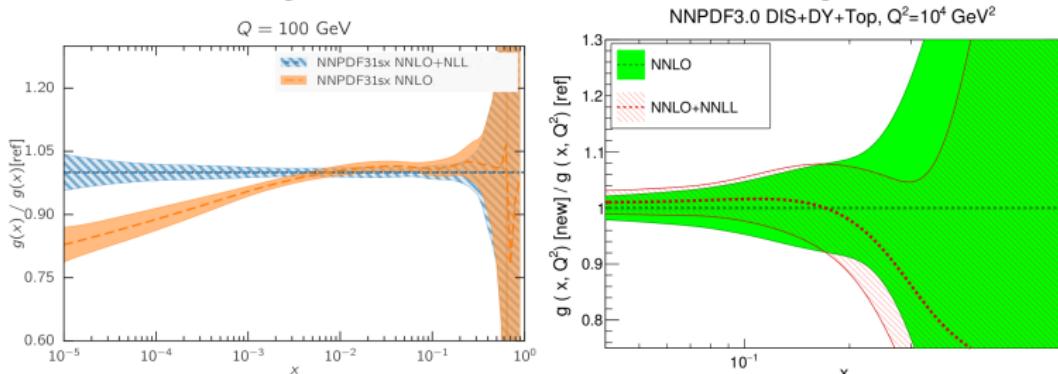


Parametrise the $c^+(x, Q_0^2)$, quark and gluon PDFs on the same footing
take into account massive charm-initiated contribution to the DIS structure functions
stabilise the dependence of LHC processes upon variations of m_c
quantify the nonperturbative charm component in the proton (BHPS? sea-like?)

Fitted charm found to differ from perturbative charm at scales $Q \sim m_c$ in NNPDF3.1
preference for a BHPS-like shape
shape driven by LHCb W, Z data + EMC data

At $Q = 1.65 \text{ GeV}$ charm carry $0.26 \pm 0.42 \%$ of the proton momentum
but it is affected by large uncertainties, especially if no EMC data are included

Beyond fixed-order accuracy



$$\text{small } x: \frac{1}{x} \ln^k x$$

high-energy gluon emission: single logs

$$\text{large } x: \left(\frac{\ln^k(1-x)}{(1-x)} \right)_+$$

soft gluon emission: double logs

Large logs $\alpha_s \ln \sim 1$ spoil the convergence of the perturbative series

PDFs with threshold resummation [JHEP 1509 (2015) 191] (only DIS, DY Z/γ , total $t\bar{t}$)

suppression in PDFs partially or totally compensates enhancements in partonic cross-sections
accuracy of the resummed fit competitive with the fixed-order fit, except for the large- x gluon

large uncertainties for MSSM particle resummed cross-sections [EPJC 76 (2016) 53]

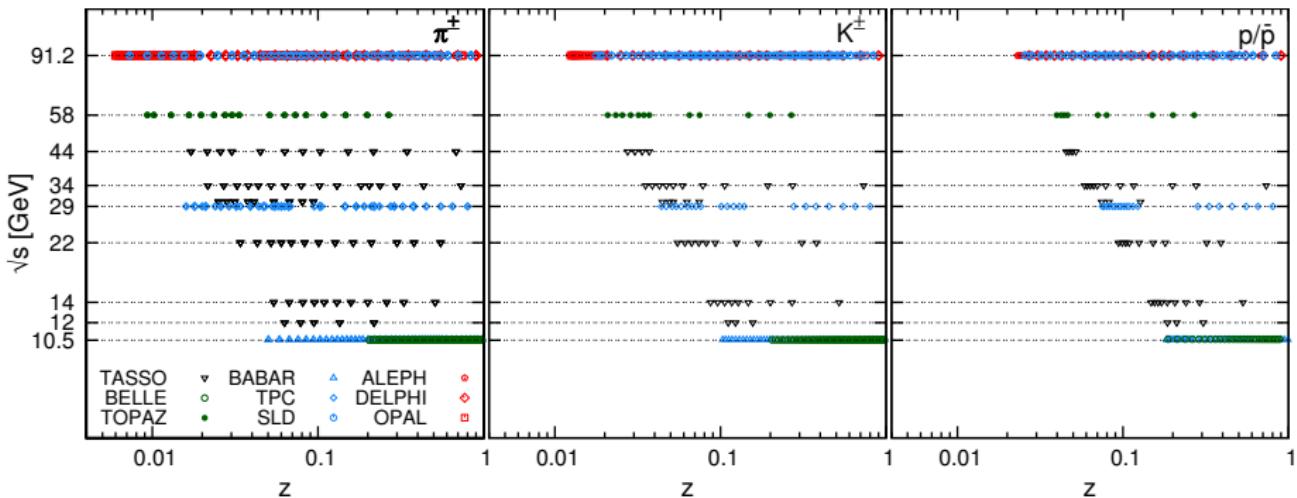
Towards PDFs with high-energy resummation [NNPDF, in progress] (DIS only)

N_{dat}	NLO+NLL	NLO	NNLO+NLL	NNLO
3102	1.111	1.113	1.108	1.126

Resummed PDFs enhanced at small x , uncertainties reduced

2. NNFF1.0 [arXiv:1706.07049, accepted for publication in EPJC]

The dataset



CERN-LEP: ALEPH [ZP C66 (1995) 353] DELPHI [EPJ C18 (2000) 203] OPAL [ZP C63 (1994) 181]

KEK: BELLE ($n_f = 4$) [PRL 111 (2013) 062002] TOPAZ [PL B345 (1995) 335]

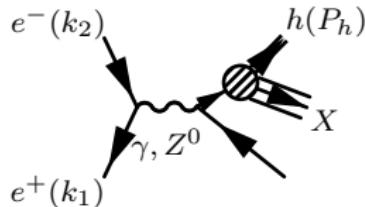
DESY-PETRA: TASSO [PL B94 (1980) 444, ZP C17 (1983) 5, ZP C42 (1989) 189]

SLAC: BABAR ($n_f = 4$) [PR D88 (2013) 032011] SLD [PR D58 (1999) 052001] TPC [PRL 61 (1988) 1263]

$$\frac{d\sigma^h}{dz} = \frac{4\pi\alpha^2(Q^2)}{Q^2} F_2^h(z, Q^2) \quad h = \pi^+ + \pi^-, K^+ + K^-, p + \bar{p} \quad \text{possibly normalised to } \sigma_{\text{tot}}$$

$$N_{\text{dat}} = 428 \text{ (pions)} \quad N_{\text{dat}} = 385 \text{ (kaons)} \quad N_{\text{dat}} = 360 \text{ (protons)}$$

From observables to fragmentation functions



$$e^+(k_1) + e^-(k_2) \xrightarrow{\gamma, Z^0} h(P_h) + X$$

$$q = k_1 + k_2 \quad q^2 = Q^2 > 0 \quad z = \frac{2P_h \cdot q}{Q^2}$$

$$F_2^h = \langle e^2 \rangle \left\{ C_{2,q}^S \otimes D_\Sigma^h + n_f C_{2,g}^S \otimes D_g^h + C_{2,q}^{NS} \otimes D_{NS}^h \right\}$$

$$\langle e^2 \rangle = \frac{1}{n_f} \sum_{q=1}^{n_f} \hat{e}_q^2 \quad D_\Sigma^h = \sum_{q=1}^{n_f} D_{q+}^h \quad D_{NS}^h = \sum_{q=1}^{n_f} \left(\frac{\hat{e}_q^2}{\langle e^2 \rangle} - 1 \right) D_{q+}^h \quad D_{q+}^h = D_q^h + D_{\bar{q}}^h$$

Coefficient functions and splitting functions known up to NNLO

[NPB 751 (2006) 18; NPB 749 (2006) 1; PLB 638 (2006) 61; NPB 845 (2012) 133]

No sensitivity to individual quark and antiquark FFs

Limited sensitivity to flavour separation via the variation of \hat{e}_q with Q^2
 $\hat{e}_u^2/\hat{e}_d^2(Q^2 = 10 \text{ GeV}) \sim 4 \Rightarrow D_{u+}^h, D_{d+}^h + D_{s+}^h$; $\hat{e}_u^2/\hat{e}_d^2(Q^2 = M_Z) \sim 0.8 \Rightarrow D_\Sigma^h$

Flavor separation between uds and c, b quarks achieved thanks to tagged data

Direct sensitivity to D_g^h only beyond LO, as $C_{2,g}^S$ is $\mathcal{O}(\alpha_s^2)$

Indirect sensitivity to D_g^h via scale violations in the time-like DGLAP evolution

Fit settings

Physical parameters: consistent with the NNPDF3.1 PDF set

$$\alpha_s(M_Z) = 0.118, \alpha(M_Z) = 1/128, m_c = 1.51 \text{ GeV}, m_b = 4.92 \text{ GeV}$$

Solution of DGLAP equations: numerical solution in z -space as implemented in APFEL
extensive benchmark performed up to NNLO [JHEP 1503 (2015) 046]

Parametrisation: each FF is parametrised with a feed-forward neural network (2-5-3-1)

$$D_i^h(Q_0, z) = \text{NN}(x) - \text{NN}(1), \quad Q_0 = 5 \text{ GeV}$$

$$h = \pi^+ + \pi^-$$

$$h = K^+ + K^-, h = p + \bar{p}$$

$$i = u^+, s^+, c^+, b^+, g$$

$$i = u^+, d^+, s^+, c^+, b^+, g$$

$$D_{u^+}^{\pi^\pm} = D_{d^+}^{\pi^\pm} \text{ (isospin symmetry)}$$

no further theoretical assumptions

we assume charge conjugation, from which $D_{q^+}^{\pi^+} = D_{q^+}^{\pi^-}$

we enforce positivity by construction, assuming quadratic NNs

initial scale above m_b , but below the lowest c.m. energy of the data, avoid threshold crossing

Heavy flavours: heavy-quark FFs are parametrised independently at the initial scale Q_0
a matched GM-VFNS (like FONLL) may be required if $Q_0 < m_c$ [PRD 94 (2016) 034037]

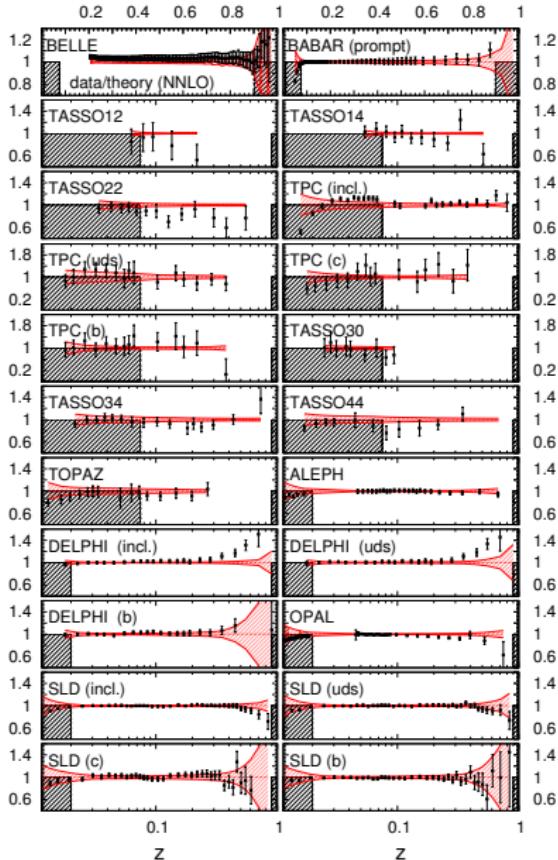
Kinematic cuts: $z \rightarrow 0$: contributions $\propto \ln z$; $z \rightarrow 1$: contributions $\propto \ln(1-z)$

$$z_{\min} = 0.075 (\sqrt{s} = M_Z); z_{\min} = 0.02 (\sqrt{s} = M_Z); z_{\max} = 0.9$$

kinematic corrections $\propto M_h/(sz^2)$ included exactly in the cross sections [PRD 73 (2006) 054020]

Minimisation: CMA-ES, verified with closure tests

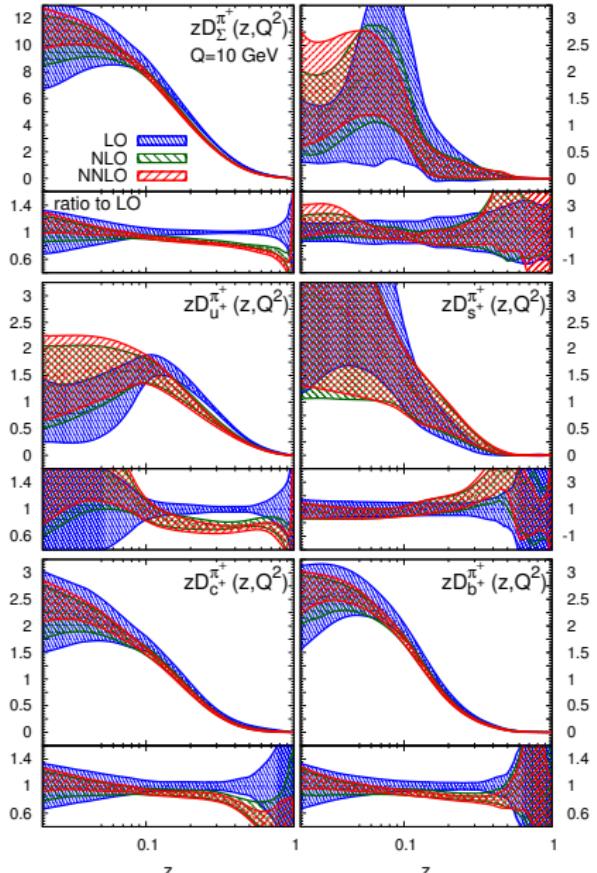
Fit quality: π^+



Exp.	N_{dat}	NNLO theory	
		χ^2/N_{dat}	remarks
BELLE	70	0.09	lack of correlations
BABAR	40	0.78	
TASSO12	4	0.87	
TASSO14	9	1.70	
TASSO22	8	1.91	} data fluctuations
TPC	13	0.85	
TPC-UDS	6	0.49	
TPC-C	6	0.52	
TPC-B	6	1.43	
TASSO30	—	—	not fitted
TASSO34	9	1.00	
TASSO44	6	2.34	data fluctuations
TOPAZ	5	0.80	
ALEPH	23	0.78	
DELPHI	21	1.86	tension with OPAL
DELPHI-UDS	21	1.54	
DELPHI-B	21	0.95	
OPAL	24	1.84	tension with DELPHI
SLD	34	0.83	
SLD-UDS	34	0.52	
SLD-C	34	1.06	
SLD-B	34	0.36	
TOTAL	428	0.87	

Overall good description of the dataset
 Signs of tension OPAL vs DELPHI (inclusive)
 Anomalously small χ^2/N_{dat} for BELLE

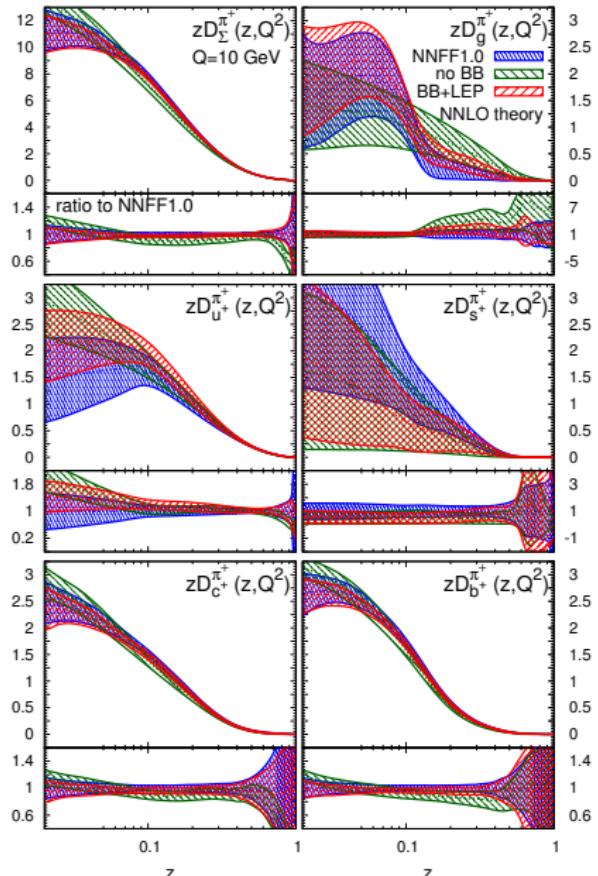
Dependence upon perturbative order: π^+



Exp.	N_{dat}	LO χ^2/N_{dat}	NLO χ^2/N_{dat}	NNLO χ^2/N_{dat}
BELLE	70	0.60	0.11	0.09
BABAR	40	1.91	1.77	0.78
TASSO12	4	0.70	0.85	0.87
TASSO14	9	1.55	1.67	1.70
TASSO22	8	1.64	1.91	1.91
TPC	13	0.46	0.65	0.85
TPC-UDS	6	0.78	0.55	0.49
TPC-C	6	0.55	0.53	0.52
TPC-B	6	1.44	1.43	1.43
TASSO30	-	-	-	-
TASSO34	9	1.16	0.98	1.00
TASSO44	6	2.01	2.24	2.34
TOPAZ	5	1.04	0.82	0.80
ALEPH	23	1.68	0.90	0.78
DELPHI	21	1.44	1.79	1.86
DELPHI-UDS	21	1.30	1.48	1.54
DELPHI-B	21	1.21	0.99	0.95
OPAL	24	2.29	1.88	1.84
SLD	34	2.33	1.14	0.83
SLD-UDS	34	0.95	0.65	0.52
SLD-C	34	3.33	1.33	1.06
SLD-B	34	0.45	0.38	0.36
TOTAL	428	1.44	1.02	0.87

Excellent perturbative convergence
FFs almost stable from NLO to NNLO
LO FF uncertainties larger than HO

Dependence upon the dataset: π^+

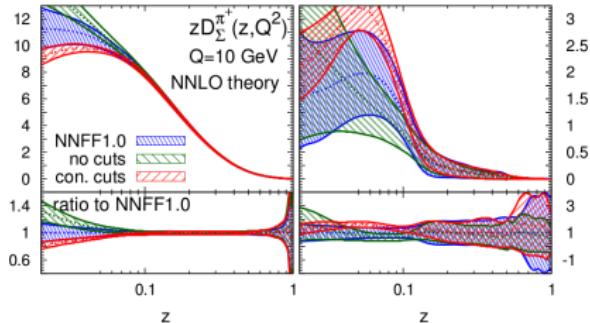


Exp.	NNLO theory	N _{dat}	NNFF1.0 χ^2/N_{dat}	no BB χ^2/N_{dat}	BB+LEP χ^2/N_{dat}
BELLE		70	0.09	(4.92)	0.09
BABAR		40	0.78	(144)	0.88
TASSO12		4	0.87	0.52	(0.87)
TASSO14		9	1.70	1.38	(1.71)
TASSO22		8	1.91	1.29	(2.15)
TPC		13	0.85	2.12	(0.81)
TPC-UDS		6	0.49	0.54	(0.77)
TPC-C		6	0.52	0.74	(0.58)
TPC-B		6	1.43	1.60	(1.48)
TASSO30		—	—	—	—
TASSO34		9	1.00	1.17	(1.38)
TASSO44		6	2.34	2.52	(2.97)
TOPAZ		5	0.80	0.92	(1.72)
ALEPH		23	0.78	0.57	0.74
DELPHI		21	1.86	1.97	1.82
DELPHI-UDS		21	1.54	1.56	1.42
DELPHI-B		21	0.95	1.01	0.95
OPAL		24	1.84	1.75	1.92
SLD		34	0.83	0.87	0.95
SLD-UDS		34	0.52	0.53	0.63
SLD-C		34	1.06	0.69	0.96
SLD-B		34	0.36	0.49	0.37
TOTAL		428	0.87	1.06	0.82

no BB: larger uncertainties; different gluon shape and different light flavour separation

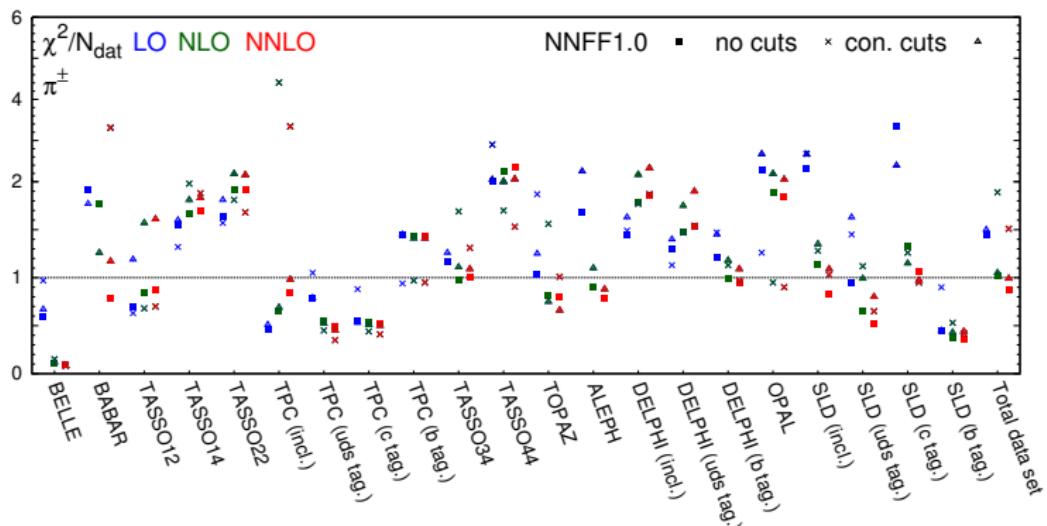
BB+LEP: comparable uncertainties; slightly different size of gluon and light quarks

Dependence upon kinematic cuts: π^+

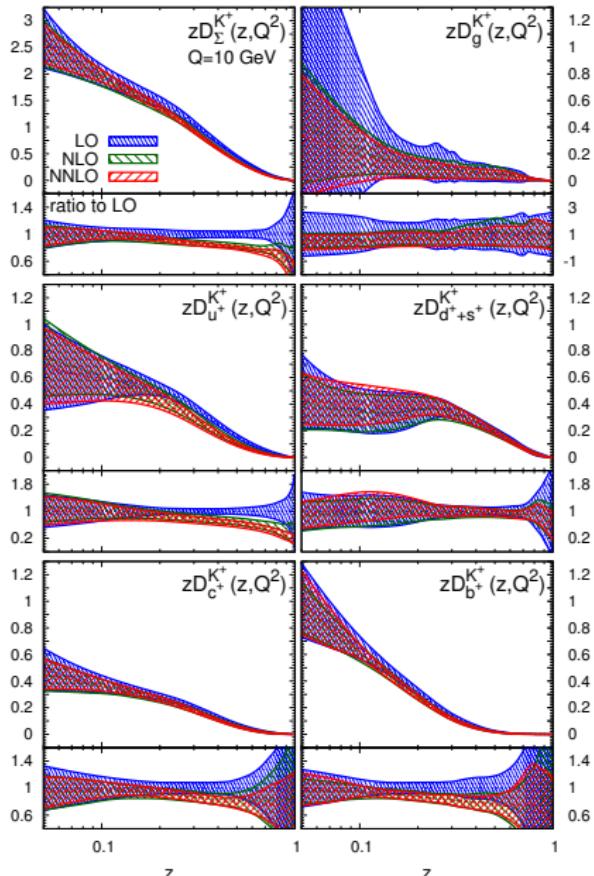


Excellent consistency in the overlapping region
Significantly varied FF shapes at low z
Stability upon inclusion of NNLO corrections

NNFF1.0 (M_Z)	$z_{\min}^{(M_Z)}$	no cuts (M_Z)	$z_{\min}^{(M_Z)}$	con. cuts (M_Z)	$z_{\min}^{(M_Z)}$
0.02	0.075	0.00	0.00	0.05	0.10



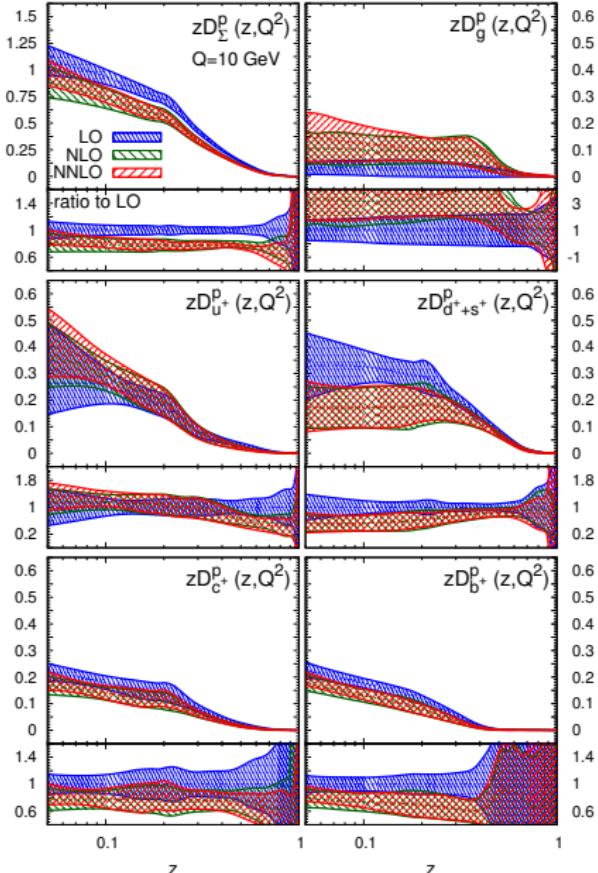
Fit quality vs perturbative order: K^+



Exp.	N_{dat}	LO χ^2/N_{dat}	NLO χ^2/N_{dat}	NNLO χ^2/N_{dat}
BELLE	70	0.21	0.32	0.33
BABAR	43	2.86	1.11	0.95
TASSO12	3	1.10	1.03	1.02
TASSO14	9	2.17	2.13	2.07
TASSO22	6	2.14	2.77	2.62
TPC	13	0.94	1.09	1.01
TPC-UDS	—	—	—	—
TPC-C	—	—	—	—
TPC-B	—	—	—	—
TASSO30	—	—	—	—
TASSO34	5	0.27	0.44	0.36
TASSO44	—	—	—	—
TOPAZ	3	0.61	1.19	0.99
ALEPH	18	0.47	0.55	0.56
DELPHI	22	0.28	0.33	0.34
DELPHI-UDS	22	1.38	1.49	1.32
DELPHI-B	22	0.58	0.49	0.52
OPAL	10	1.67	1.57	1.66
SLD	35	0.86	0.62	0.57
SLD-UDS	35	1.31	1.02	0.93
SLD-C	34	0.92	0.47	0.38
SLD-B	35	0.59	0.67	0.62
TOTAL	395	1.02	0.78	0.73

Excellent perturbative convergence
FFs almost stable from NLO to NNLO
LO FF uncertainties larger than HO

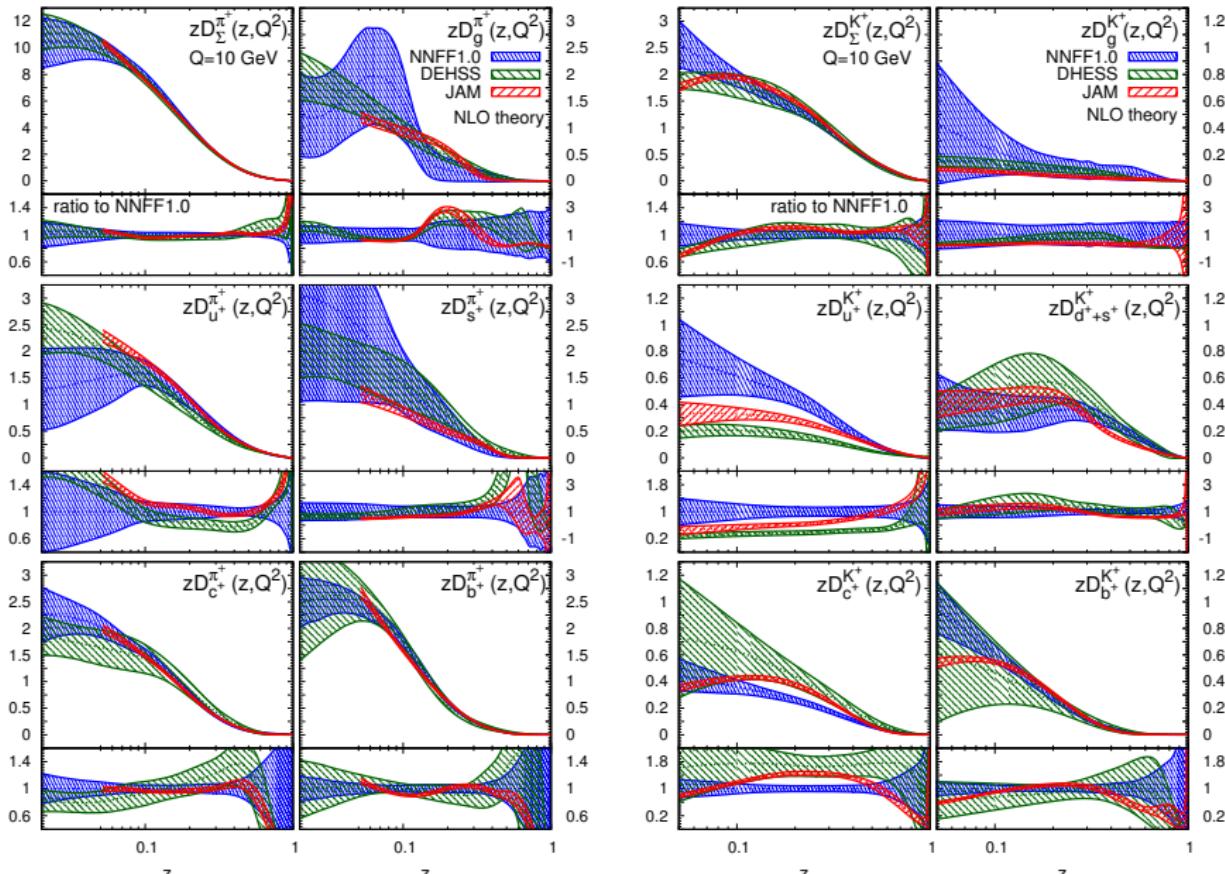
Fit quality vs perturbative order: p



Exp.	N_{dat}	LO χ^2/N_{dat}	NLO χ^2/N_{dat}	NNLO χ^2/N_{dat}
BELLE	29	0.10	0.31	0.50
BABAR	43	4.74	3.75	3.25
TASSO12	3	0.69	0.70	0.72
TASSO14	9	1.32	1.25	1.22
TASSO22	9	0.98	0.92	0.93
TPC	20	1.04	1.10	1.08
TPC-UDS	—	—	—	—
TPC-C	—	—	—	—
TPC-B	—	—	—	—
TASSO30	2	0.25	0.19	0.18
TASSO34	6	0.82	0.81	0.78
TASSO44	—	—	—	—
TOPAZ	4	0.79	1.21	1.19
ALEPH	26	1.36	1.43	1.28
DELPHI	22	0.48	0.49	0.49
DELPHI-UDS	22	0.47	0.46	0.45
DELPHI-B	22	0.89	0.89	0.91
OPAL	—	—	—	—
SLD	36	0.66	0.65	0.64
SLD-UDS	36	0.77	0.76	0.78
SLD-C	36	1.22	1.22	1.21
SLD-B	35	1.12	1.29	1.33
TOTAL	360	1.31	1.23	1.17

Excellent perturbative convergence
FFs almost stable from NLO to NNLO
LO FF uncertainties larger than HO

Comparison with other recent FF sets: π^+ and K^+



4. Conclusions

Summary

① Unpolarised PDFs: full exploitation of the LHC harvest

- ▶ increased precision over an extended kinematic region
- ▶ theoretical improvements (NNLO, resummations, EW corrections ...)
- ▶ complementarity with the EIC program

② Polarised PDFs: only a polarised EIC could provide a significant advancement

- ▶ unprecedented kinematic reach
- ▶ check modified evolution at small x
- ▶ meanwhile, make the most from the RHIC spin physics program

③ Fragmentation Functions: significant phenomenological effort

- ▶ improve the methodological sophistication of current analyses
- ▶ (ongoing) make the most from LHC data
- ▶ (ongoing) cross-talk between FFs and PDFs requires simultaneous analyses of both

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Thank you