

Superconducting Qubits

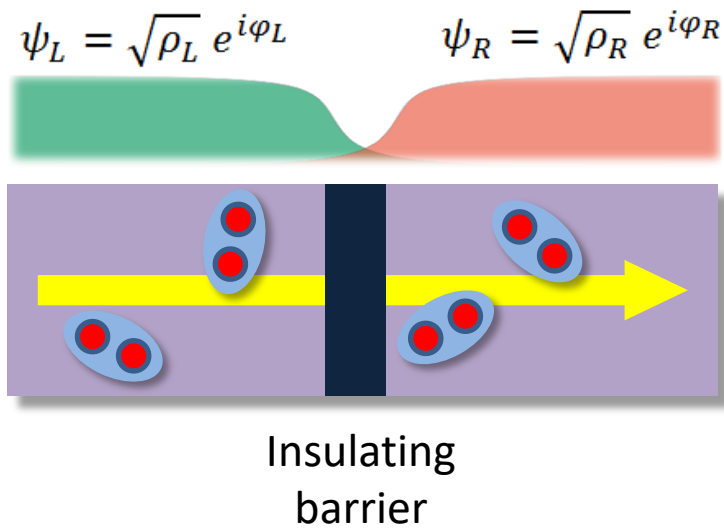
Fabio Chiarello

Institute for Photonics and Nanotechnologies

IFN – CNR Rome

The Josephson's Lego bricks box





Phase difference

$$\varphi = \varphi_R - \varphi_L$$

$$\phi_0 = \frac{h}{2e}$$

$$\phi_b = \frac{\hbar}{2e}$$

Josephson equations

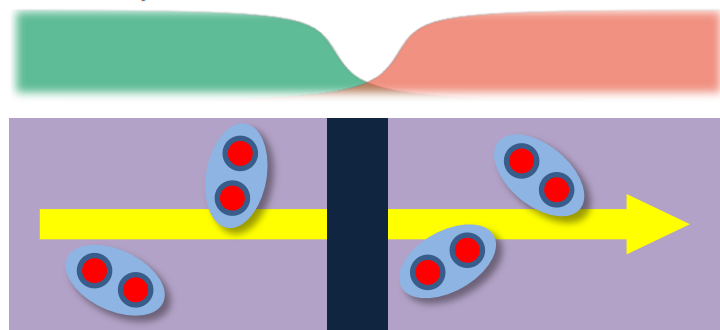
$$I = I_0 \sin \varphi$$

$$V = \frac{\hbar}{2e} \frac{d\varphi}{dt}$$

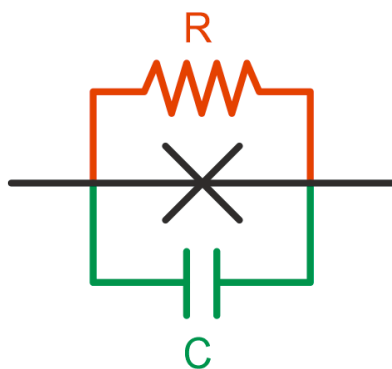
Symbol



$$\psi_L = \sqrt{\rho_L} e^{i\varphi_L} \quad \psi_R = \sqrt{\rho_R} e^{i\varphi_R}$$



Insulating
barrier



Phase difference

$$\varphi = \varphi_R - \varphi_L$$

$$\phi_0 = \frac{h}{2e}$$

$$\phi_b = \frac{\hbar}{2e}$$

Josephson equations

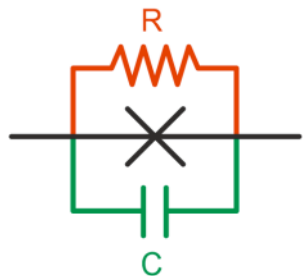
$$I = I_0 \sin \varphi + C \frac{dV}{dt} + \frac{V}{R}$$

$$V = \frac{\hbar}{2e} \frac{d\varphi}{dt}$$



$$C \phi_b \frac{d^2 \varphi}{dt^2} + \frac{\phi_b}{R} \frac{d\varphi}{dt} + I_0 \sin \varphi = I$$

Mechanical equivalent



$$C\phi_b \frac{d^2\varphi}{dt^2} + \frac{\phi_b}{R} \frac{d\varphi}{dt} + I_0 \sin\varphi = I$$

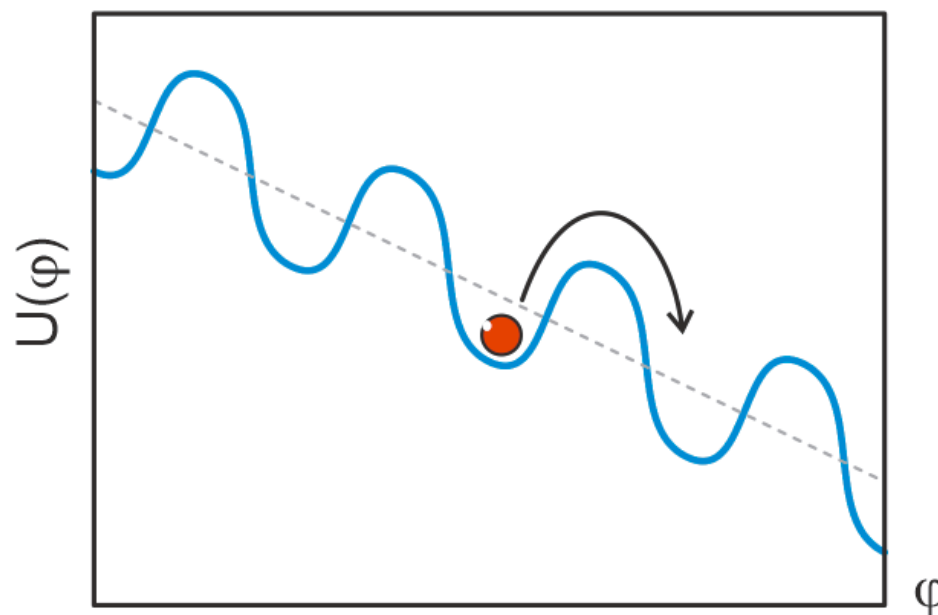
Motion equations

$$M \frac{d^2\varphi}{dt^2} + M\gamma \frac{d\varphi}{dt} = -\frac{dU}{d\varphi}$$

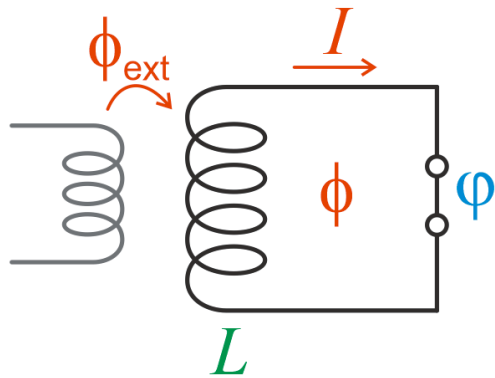
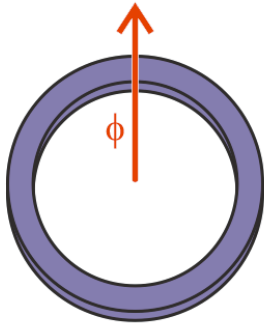
Effective potential

$$U = -E_j(\cos\varphi + I/I_0)$$

$$\begin{aligned} M &= C\phi_b^2 \\ E_j &= I_0\phi_b^2 \\ \gamma &= 1/RC \end{aligned}$$



$$V = \frac{\hbar}{2e} \frac{d\varphi}{dt}$$



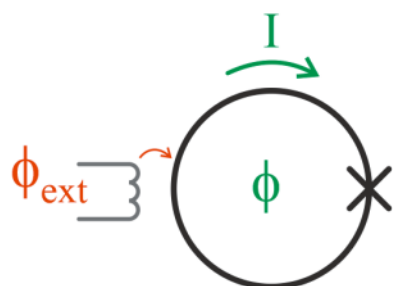
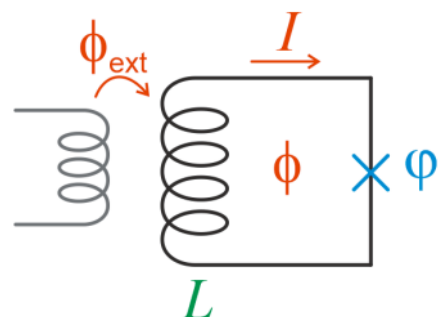
$$\phi = \phi_{ext} - LI$$

$$\phi = \frac{\hbar}{2e} \varphi$$

$$\varphi = 2\pi n$$



$$\phi = \phi_{ext} - LI = n\phi_0$$



$$C\phi_b \frac{d^2\varphi}{dt^2} + \frac{\phi_b}{R} \frac{d\varphi}{dt} + I_0 \sin\varphi = I$$

$$\phi = \phi_{ext} - LI$$

$$\phi = \frac{\hbar}{2e} \varphi$$



Motion equations

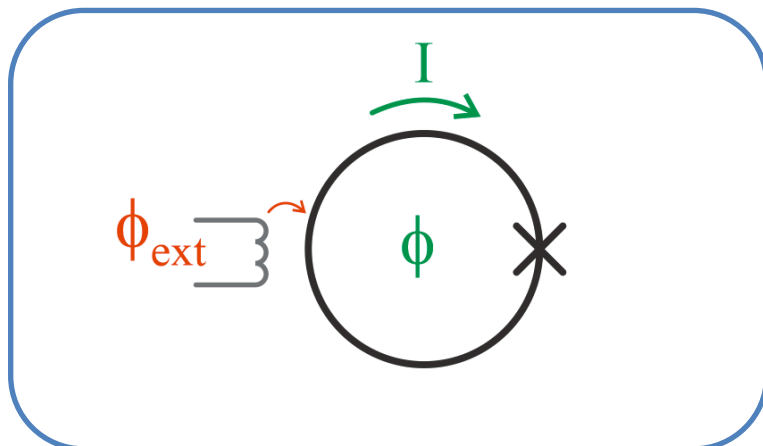
$$M \frac{d^2\varphi}{dt^2} + M\gamma \frac{d\varphi}{dt} = - \frac{dU}{d\varphi}$$

Effective potential

$$U = \frac{1}{2} E_L (\varphi - \phi_{ext}/\phi_0)^2 - E_j \cos\varphi$$

$$M = C\phi_b^2 \quad \gamma = 1/RC$$

$$E_j = I_0\phi_b^2 \quad E_L = \phi_b^2/L$$

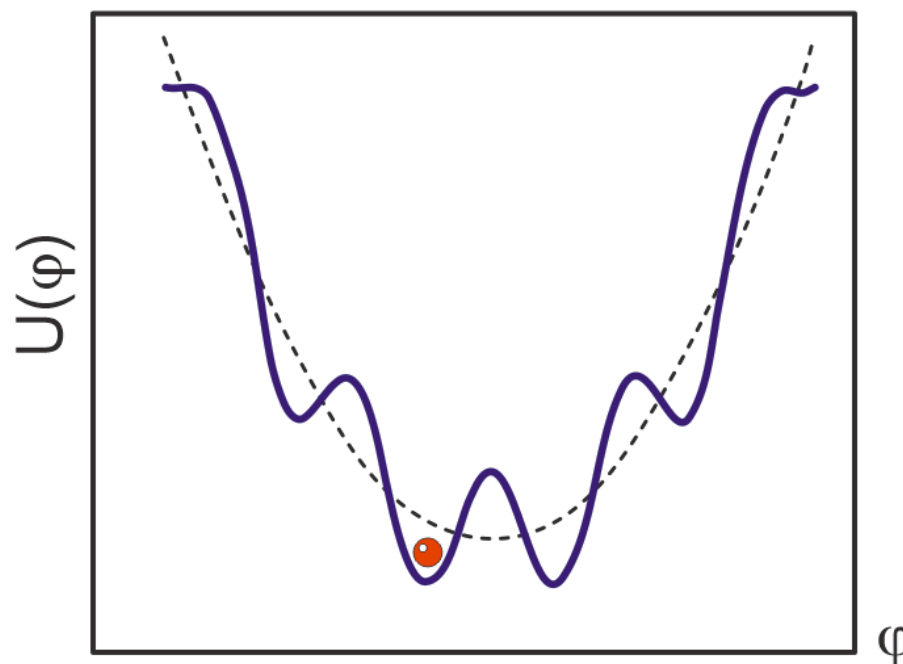


$$M \frac{d^2 \varphi}{dt^2} + M \gamma \frac{d\varphi}{dt} = - \frac{dU}{d\varphi}$$

$$U = \frac{1}{2} E_L (\varphi - \Phi_{ext} / \Phi_0)^2 - E_j \cos \varphi$$

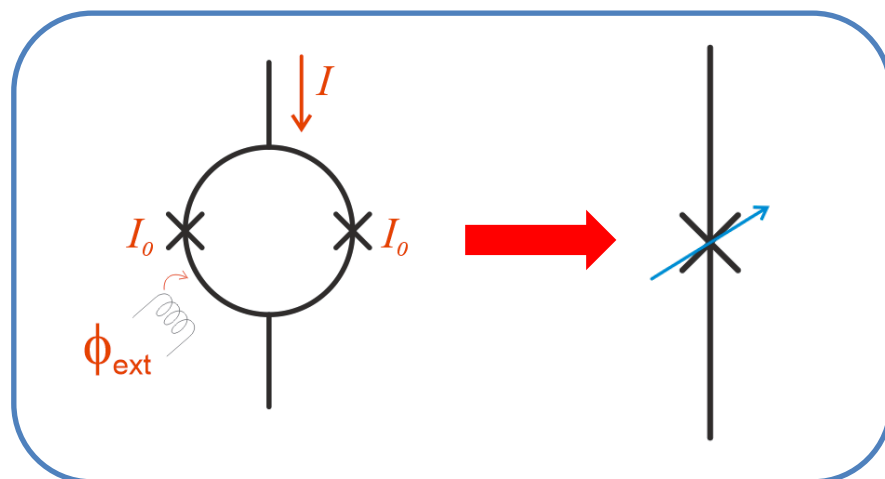
$$M = C \Phi_b^2 \quad \gamma = 1/RC$$

$$E_j = I_0 \Phi_b^2 \quad E_L = \Phi_b^2 / L$$



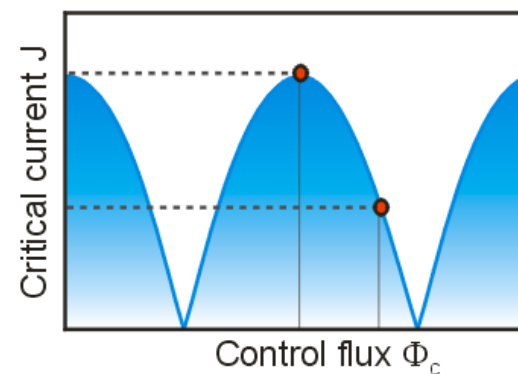
Torrioli's talk
on SQUIDs at 16:00

Tunable Josephson element



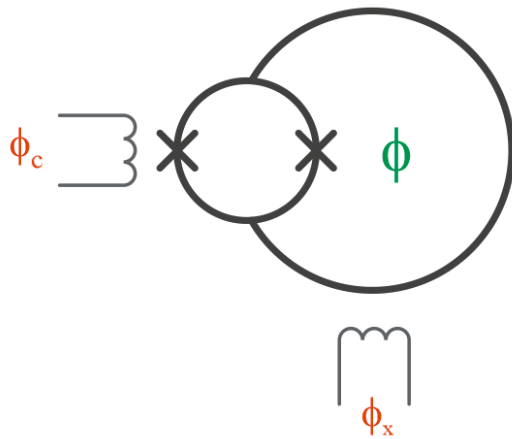
$$I_0(\phi_{ext}) \cong 2I_0 \cos(\pi\phi_{ext}/\phi_0)$$

For $LI_0 \ll \phi_b$

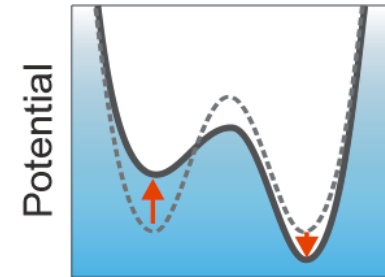
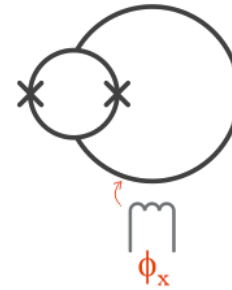


Torrioli's talk
on SQUIDs at 16:00

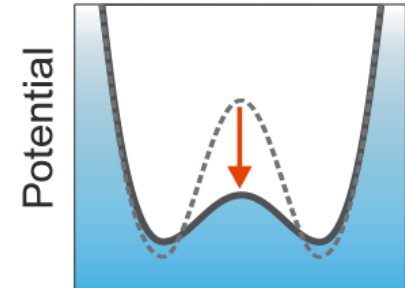
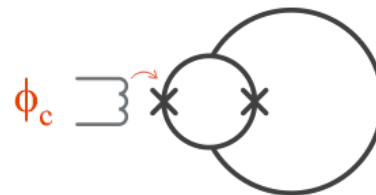
Tunable rf SQUID - two controls



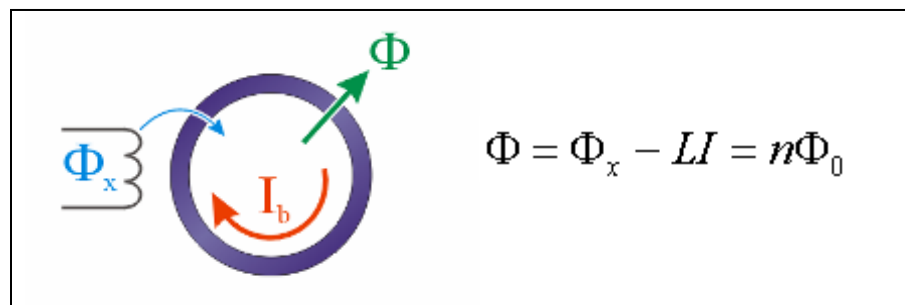
Large loop: symmetry



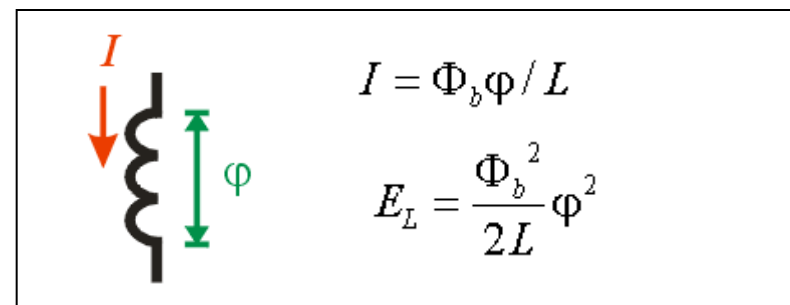
Small loop: barrier



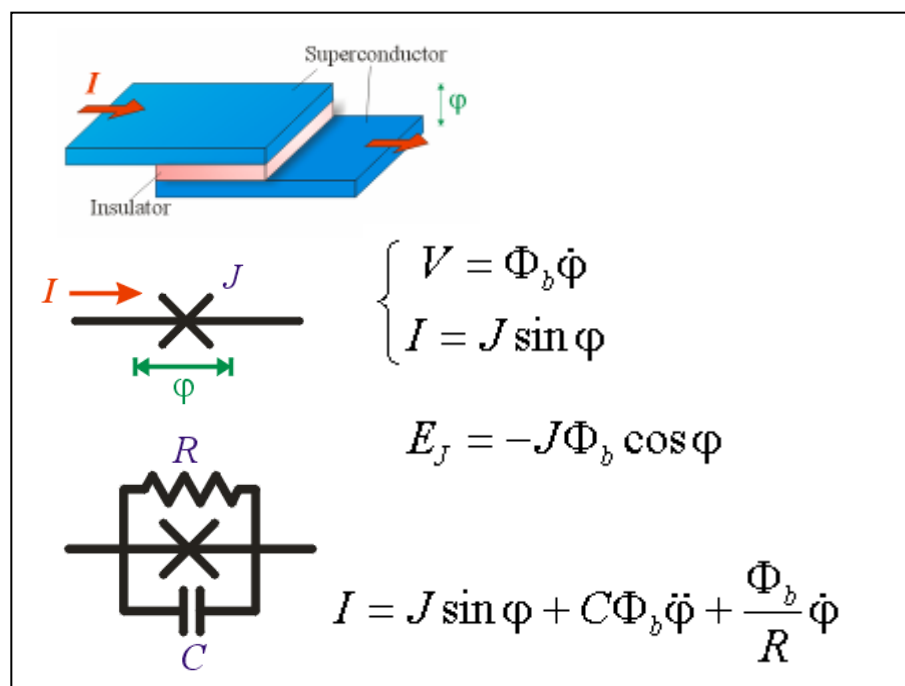
Flux quantization



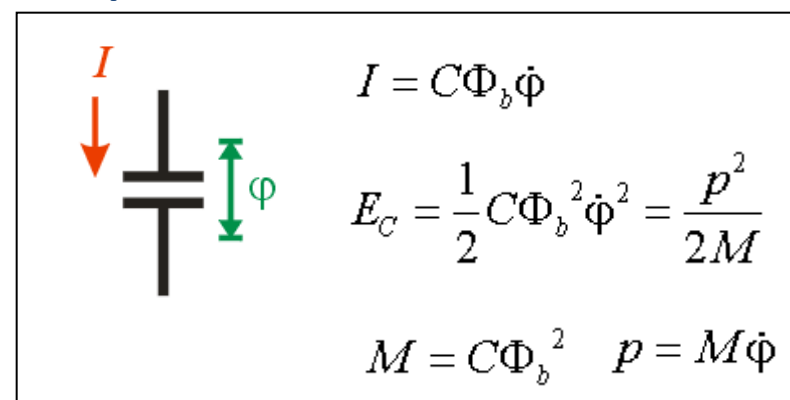
Inductance



Josephson junction

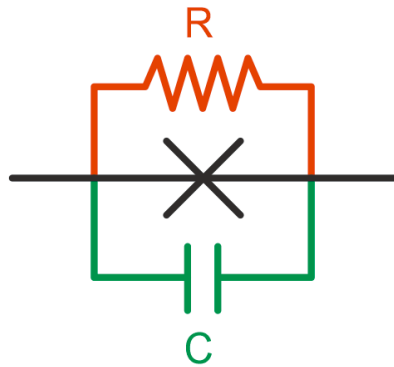


Capacitance



$$\Phi_0 = h/2e \cong 2.068 \times 10^{-15} \text{ Wb}$$

$$\Phi_b = \Phi_0 / 2\pi = h/2e \cong 3.291 \times 10^{-16} \text{ Wb}$$



$$M \frac{d^2 \varphi}{dt^2} + M \gamma \frac{d\varphi}{dt} = - \frac{dU}{d\varphi}$$

$$U = -E_j (\cos \varphi + I/I_0)$$

$$M = C \phi_b^2$$

$$E_j = I_0 \phi_b^2$$

$$\gamma = 1/RC$$

$$H = E_c n^2 + U$$

$$Q = CV = C \phi_b d\varphi/dt$$

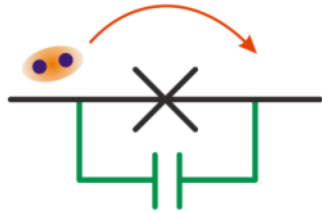
$$n = Q/(2e) = \text{Cooper pairs on } C$$

$$E_c = 2e^2/C$$

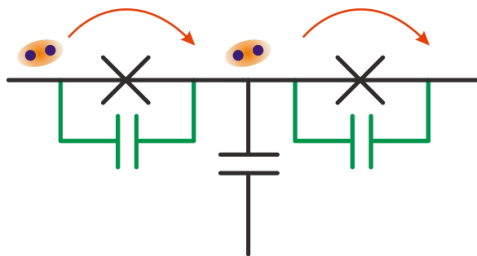
For $E_c \gg E_j$ control on single Cooper pair crossing (small junctions)

For $E_c \gg E_j$ control on single Cooper pair crossing

$$H = E_c n^2 + U$$

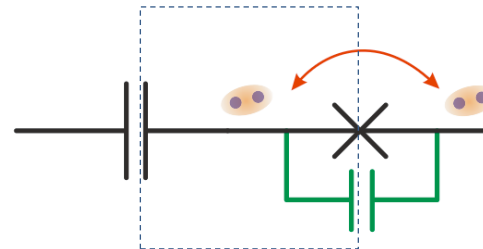


SET Transistor



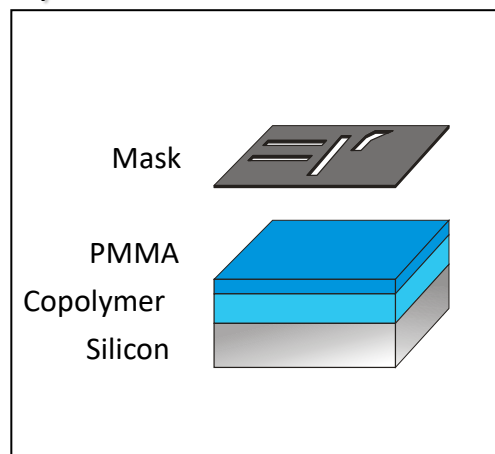
Flux of Cooper pairs
controlled one by one

Cooper Pair Box

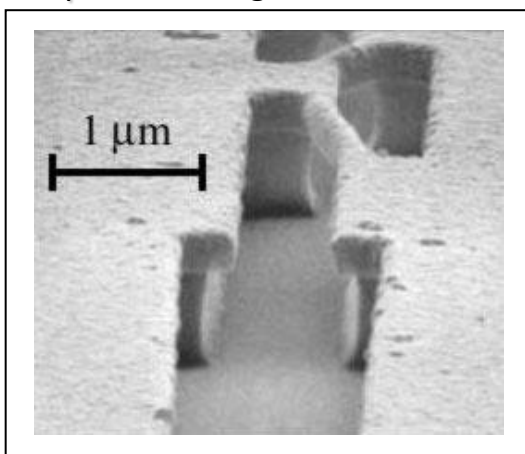


Can store a single
Cooper pair

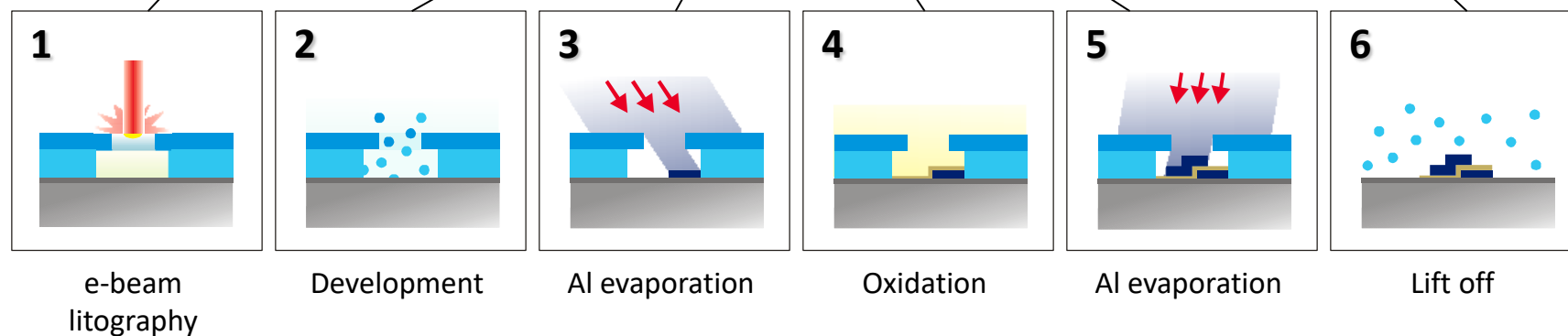
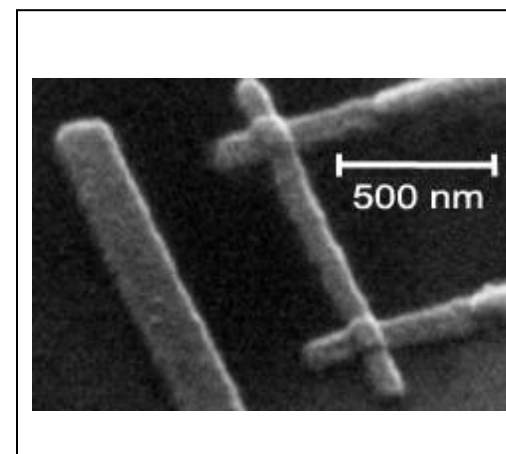
Layers



Suspended bridges



Devices

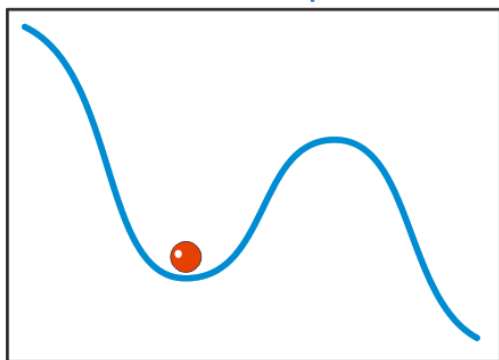


Quantum behavior



Is it possible a quantum description of the equivalent mechanical model?

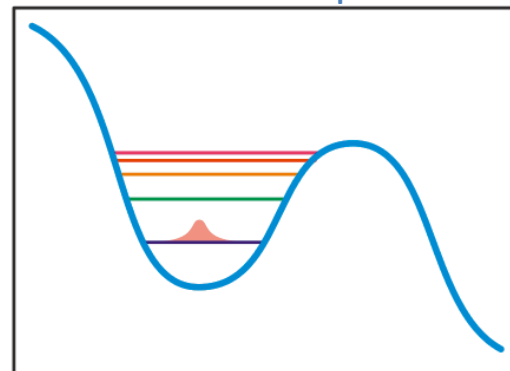
Classical description



$$M \frac{d^2 \varphi}{dt^2} = - \frac{dU}{d\varphi}$$



Quantum description



$$i\hbar \frac{\partial}{\partial t} \psi = H\psi$$

$$H = E_c n^2 + U$$

$$[n, \varphi] = -i$$

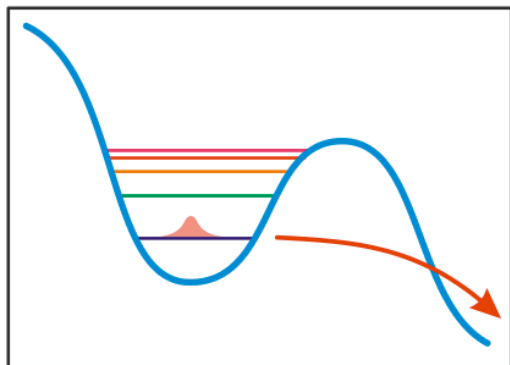
n = Cooper pairs on C

$$E_c = 2e^2 / C$$

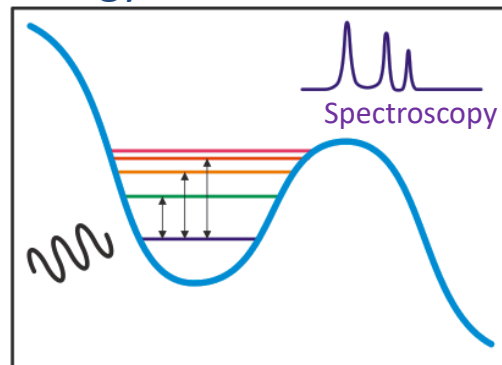
P.W. Anderson, in: E.R. Caianiello (Ed.), Lectures on the Many Body Problem, Vol. 2, Academic Press, New York, 1964.

Observed quantum effects in Josephson systems

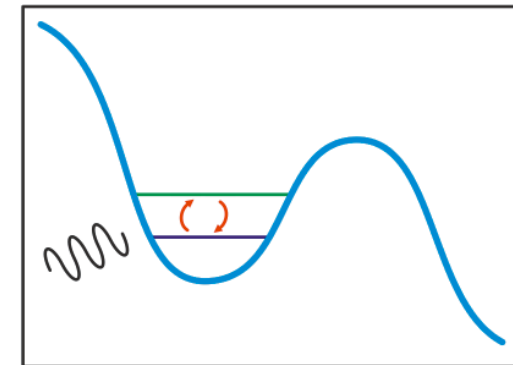
Tunnel Effect



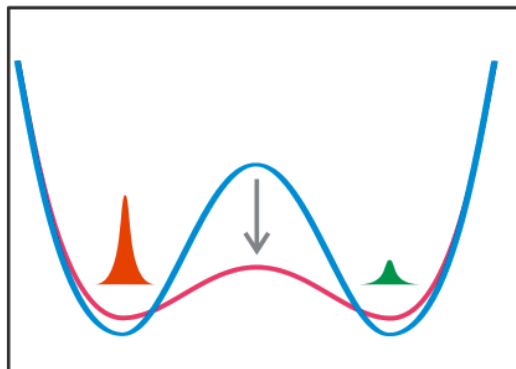
Energy Level Quantization



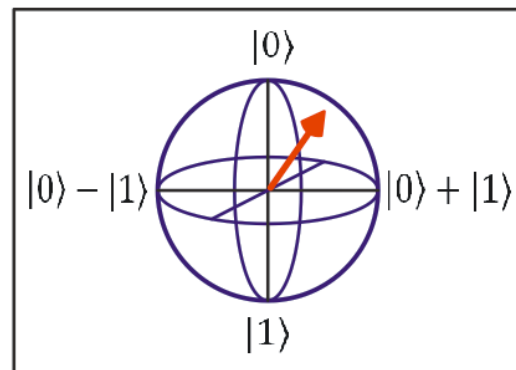
Rabi Oscillations



Nonadiabatic manipulation

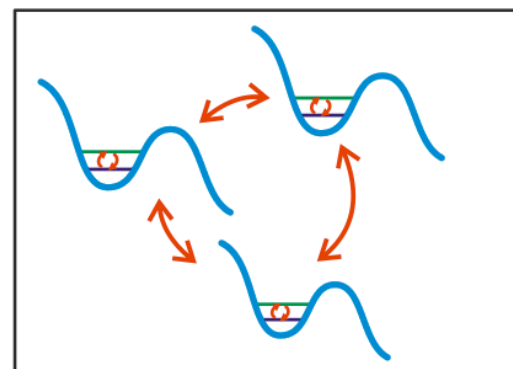


Block Sphere manipulation



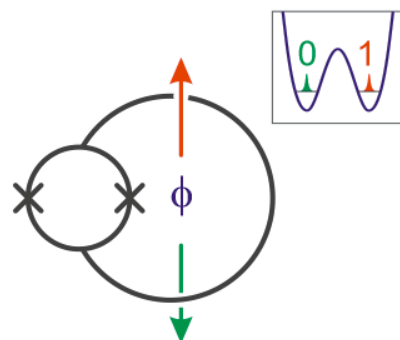
(Ramsey fringes,
Spin Echo, ...)

Entangled Systems



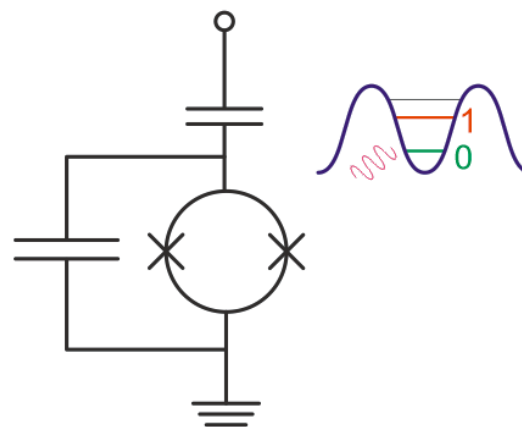
Qubit: quantum two state systems
which can be manipulated and coupled

Flux Qubit



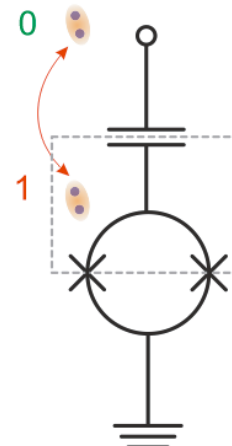
$$E_j > E_L \gg E_c$$

Transmon Qubit



$$E_j > E_c$$

Cooper Pair Box Qubit

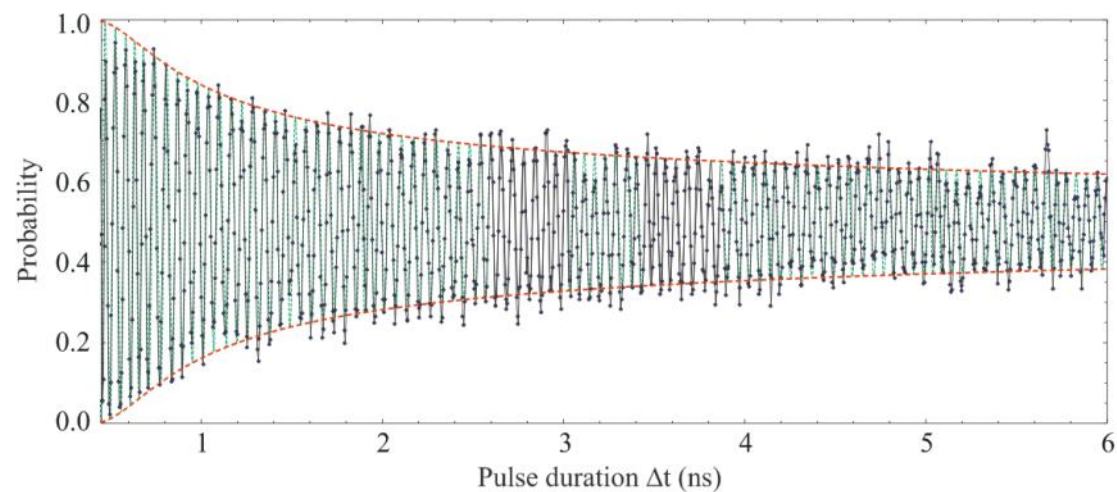
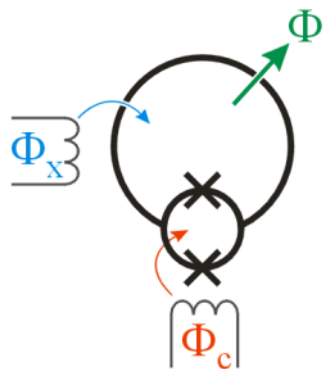


$$E_j < E_c$$

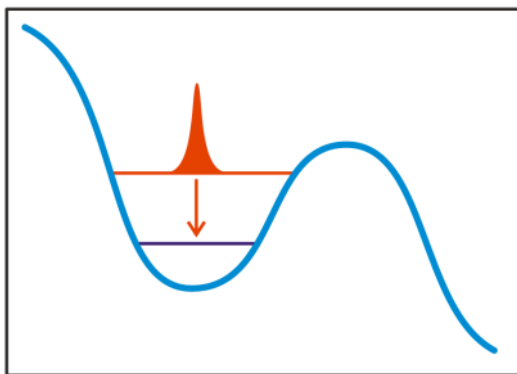
Single junction
3 junction flux qubit
Quantronium

...

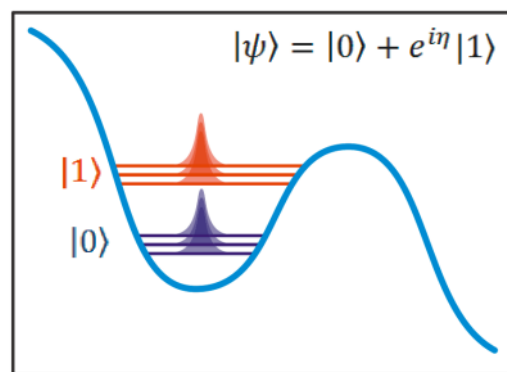
Effect of noise from different sources



Relaxation $T_1 = 1/\gamma_1$



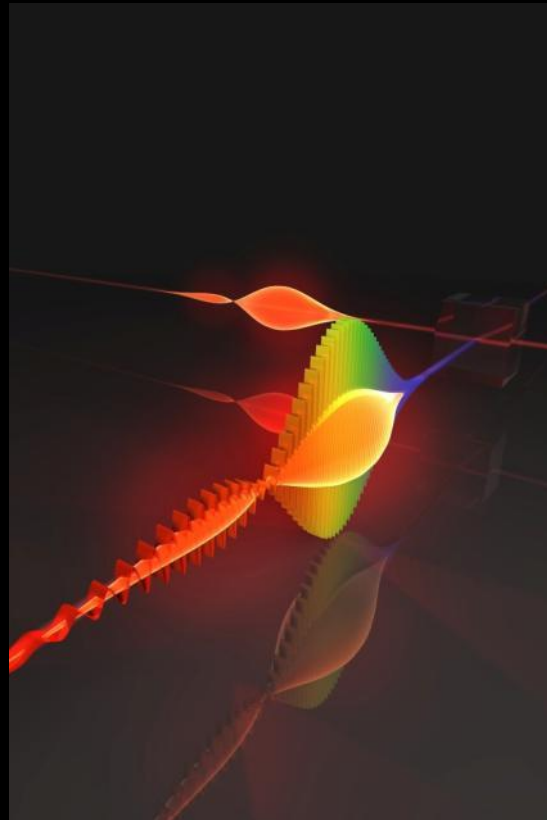
Dephasing $T_2 = 1/\gamma_2$



Ohmic noise

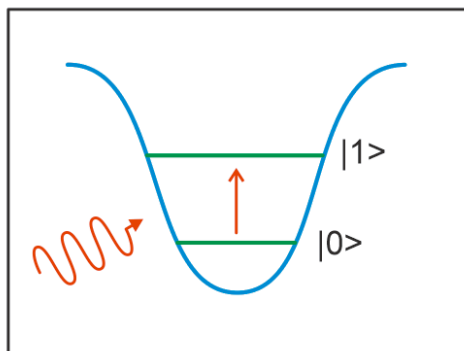
$$\gamma = \frac{\gamma_1}{2} + \gamma_2$$

Single photon detection

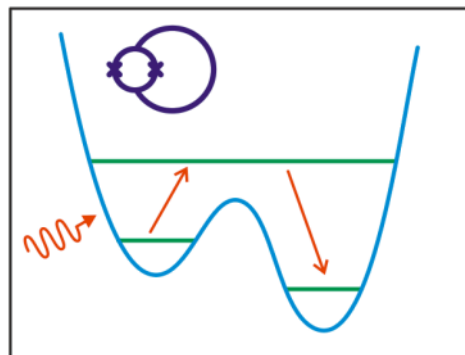


Absorption of a single (GHz) photon
Detection of the qubit state

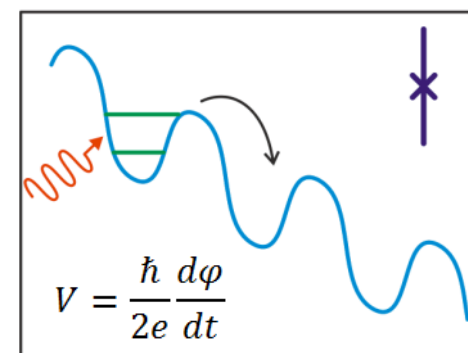
Artificial atom



Lambda transition



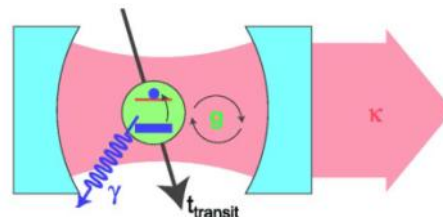
Activation in voltage state



$$\Delta f \sim \frac{\sqrt{E_j E_c}}{h} \sim 10 \text{ GHz}$$

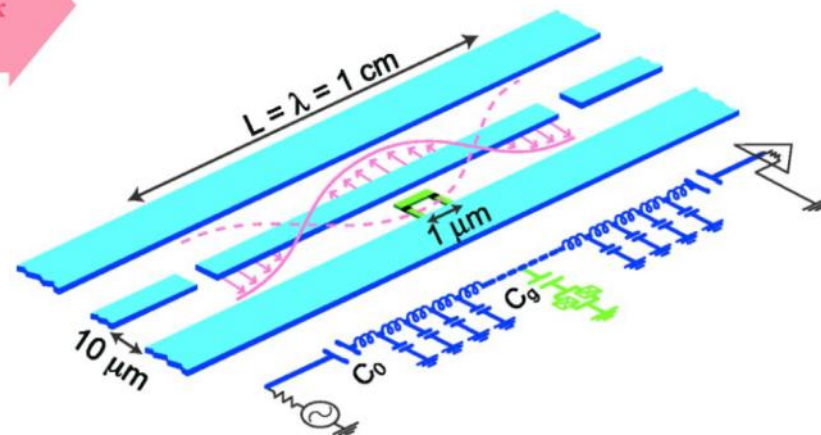
Cavity: coplanar waveguide

Atom: Cooper pair box qubit

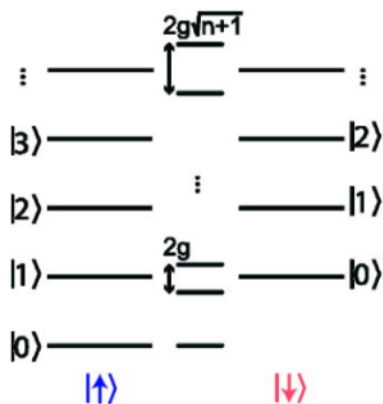


Jaynes-Cummings Hamiltonian

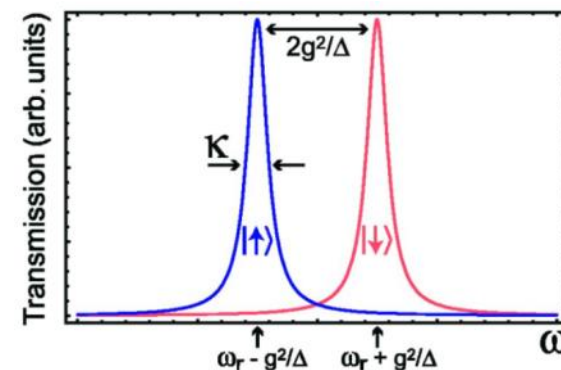
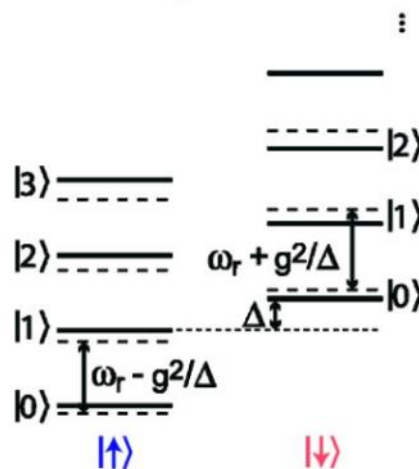
$$H = \hbar \omega_r \left(a^\dagger a + \frac{1}{2} \right) + \frac{\hbar \Omega}{2} \sigma^z + \hbar g (a^\dagger \sigma^- + \sigma^+ a)$$



Tuned $\omega_r = \Omega$

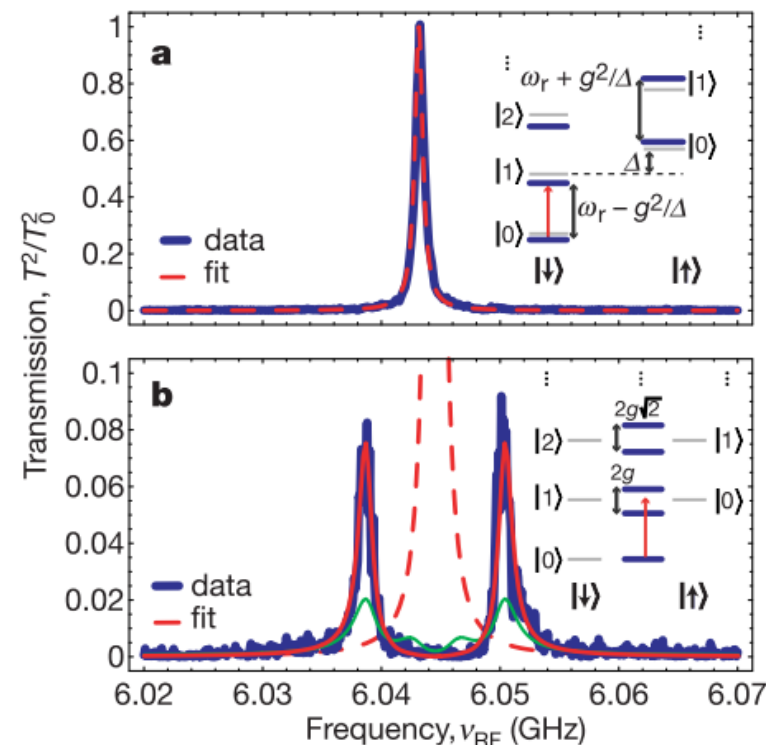
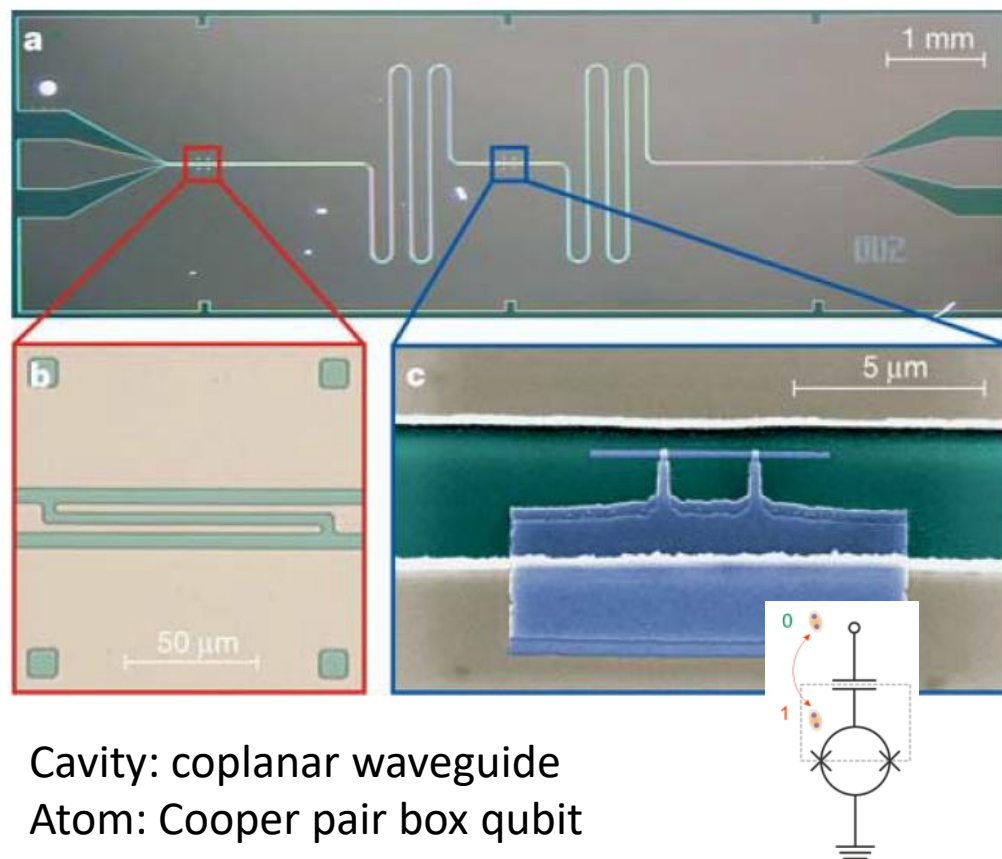


Detuned $\omega_r = \Omega - \Delta$



*A. Blais, et al., Phys. Rev. A **69**, 62320 (2004).*

Coupling between an “artificial atom” and a “cavity”

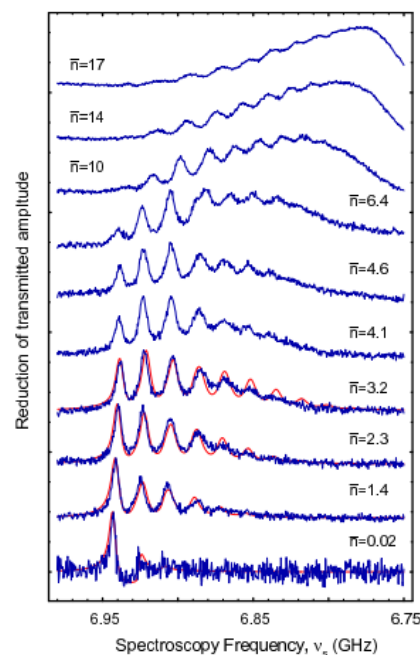
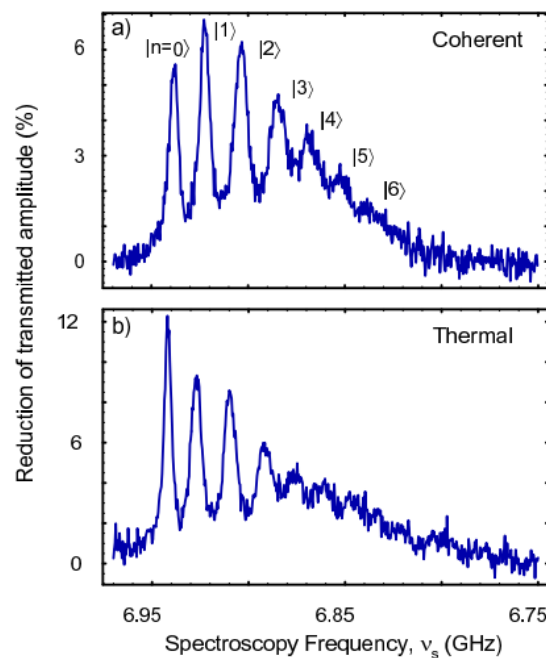
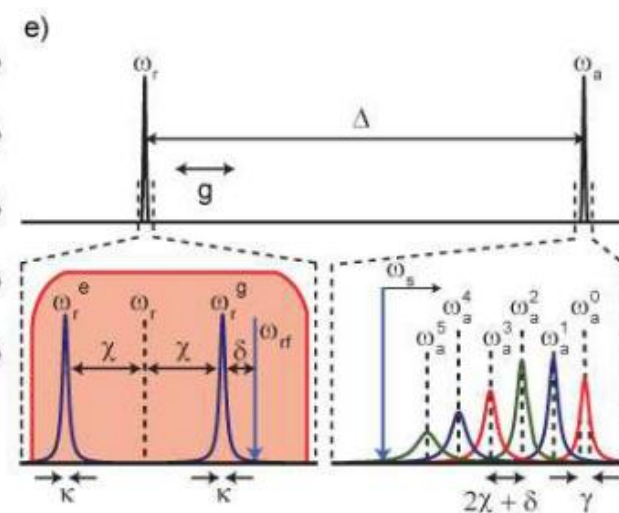
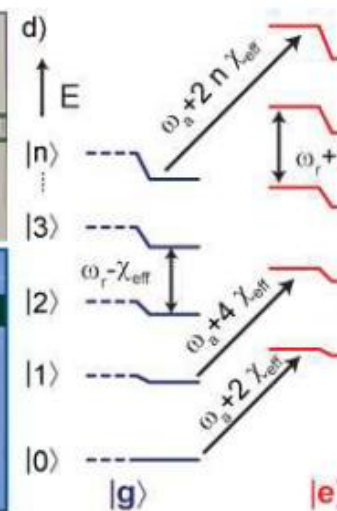
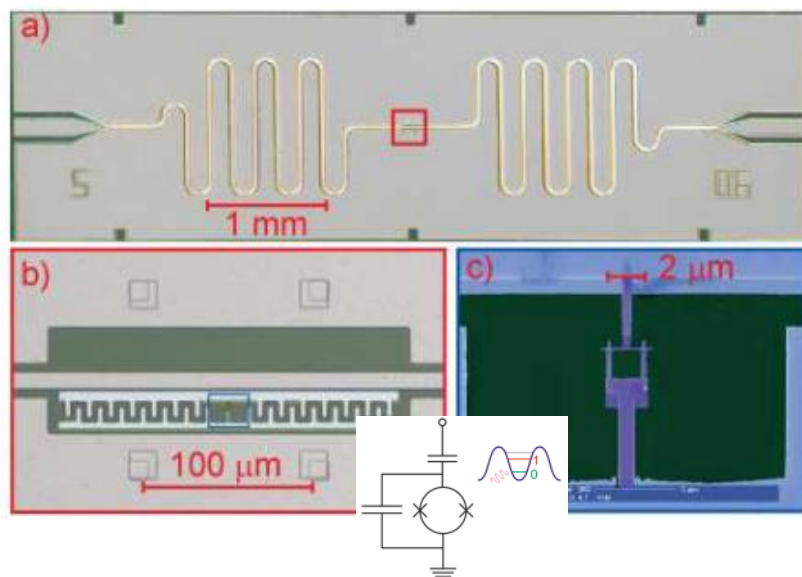


$$Q \approx 10\,000$$

$$T_k \approx 200\text{ ns}$$

$$T_\gamma \approx 230\text{ ns}$$

A. Wallraff et al., *Nature* **431**, 162–167 (2004).

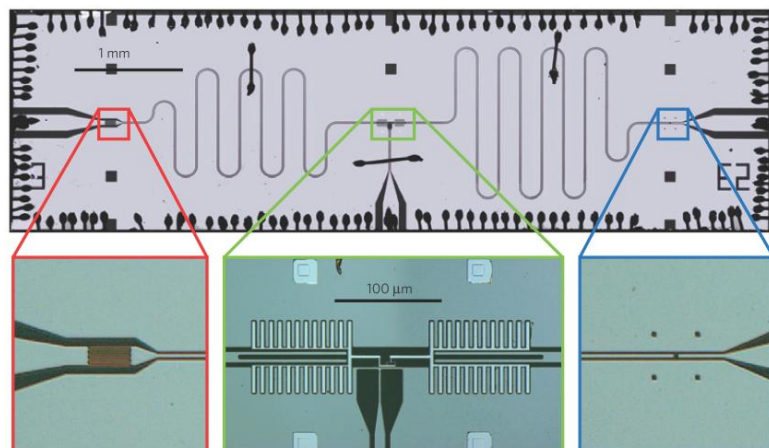
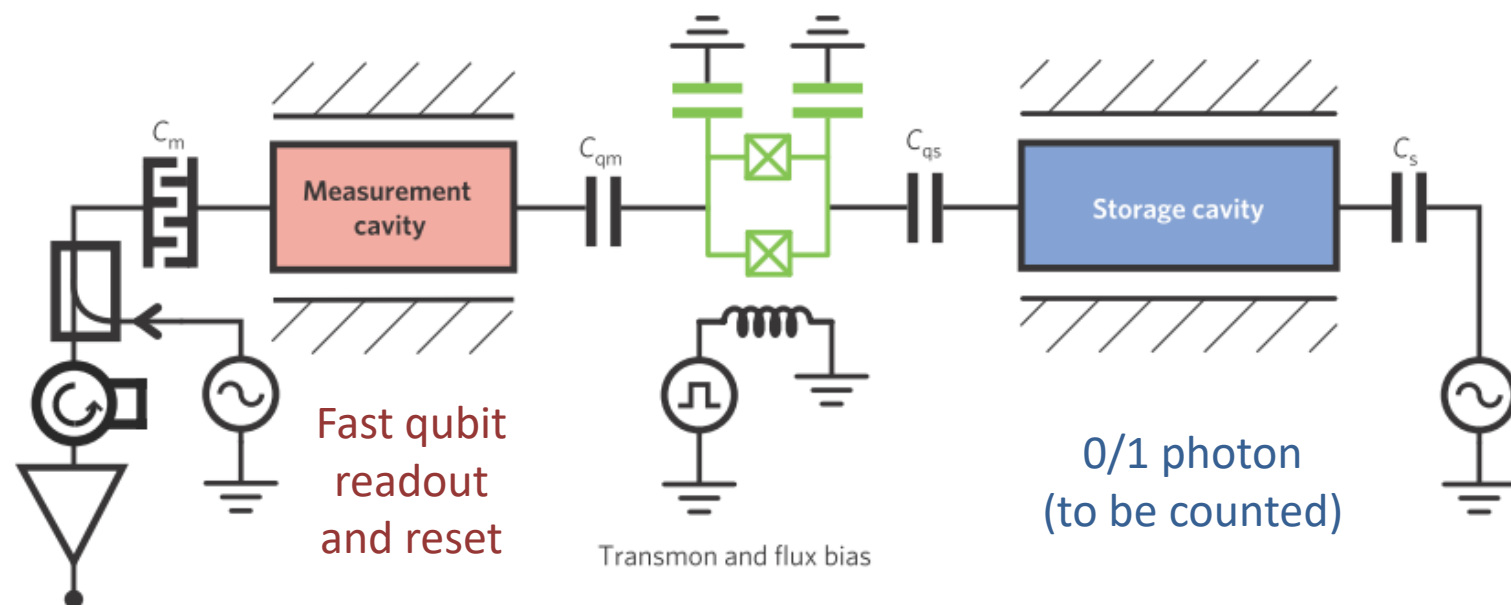


$$Q \approx 30\,000$$

$$T_k \approx 640\text{ ns}$$

$$T_\gamma \approx 84\text{ ns}$$

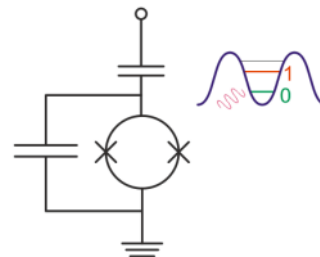
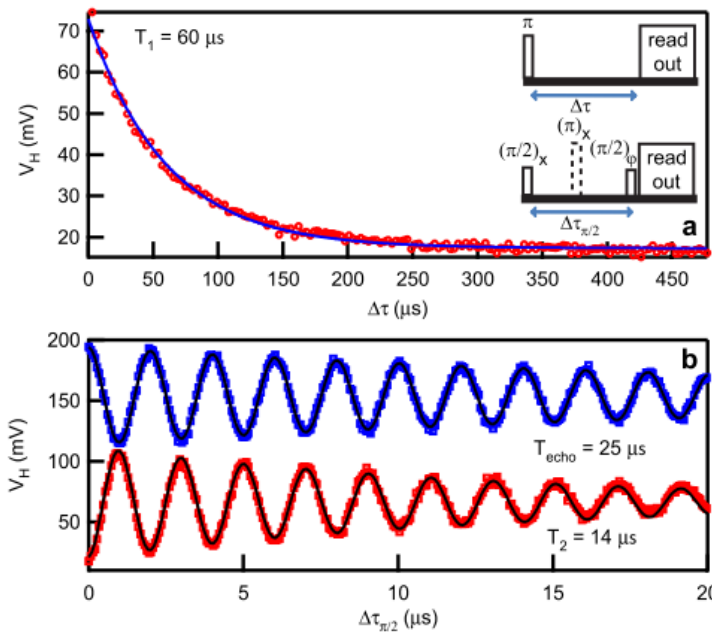
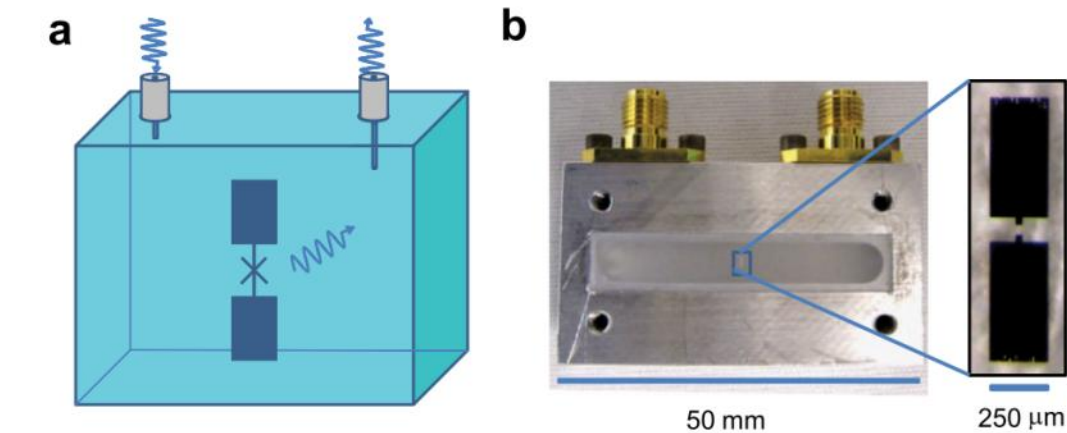
D. I. Schuster et al., Nature **445**, 515–518 (2007).



90% QND

B. R. Johnson et al., Quantum non-demolition detection of single microwave photons in a circuit, Nat Phys 6, 663–667 (2010).

Qubit in a 3D cavity



$$Q \approx 1\,000\,000$$

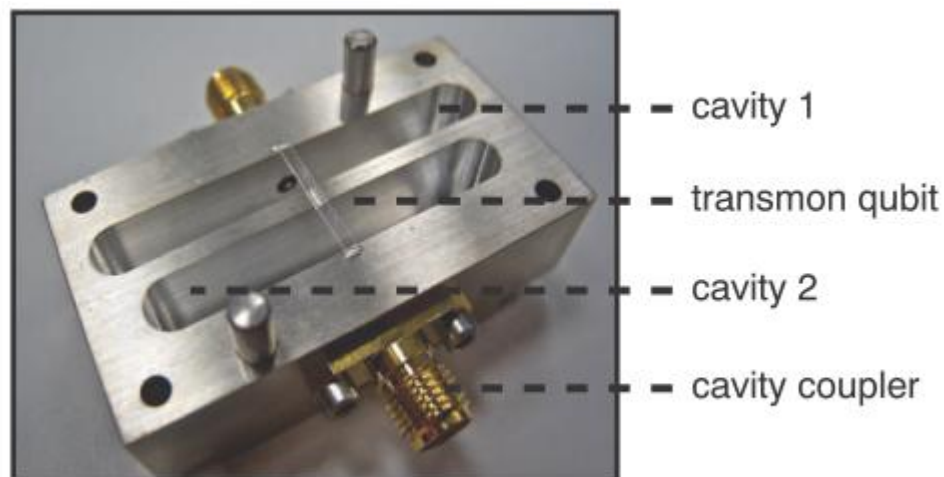
$$T_k \approx 20 \mu\text{s}$$

$$T_\gamma \approx 20 \mu\text{s}$$

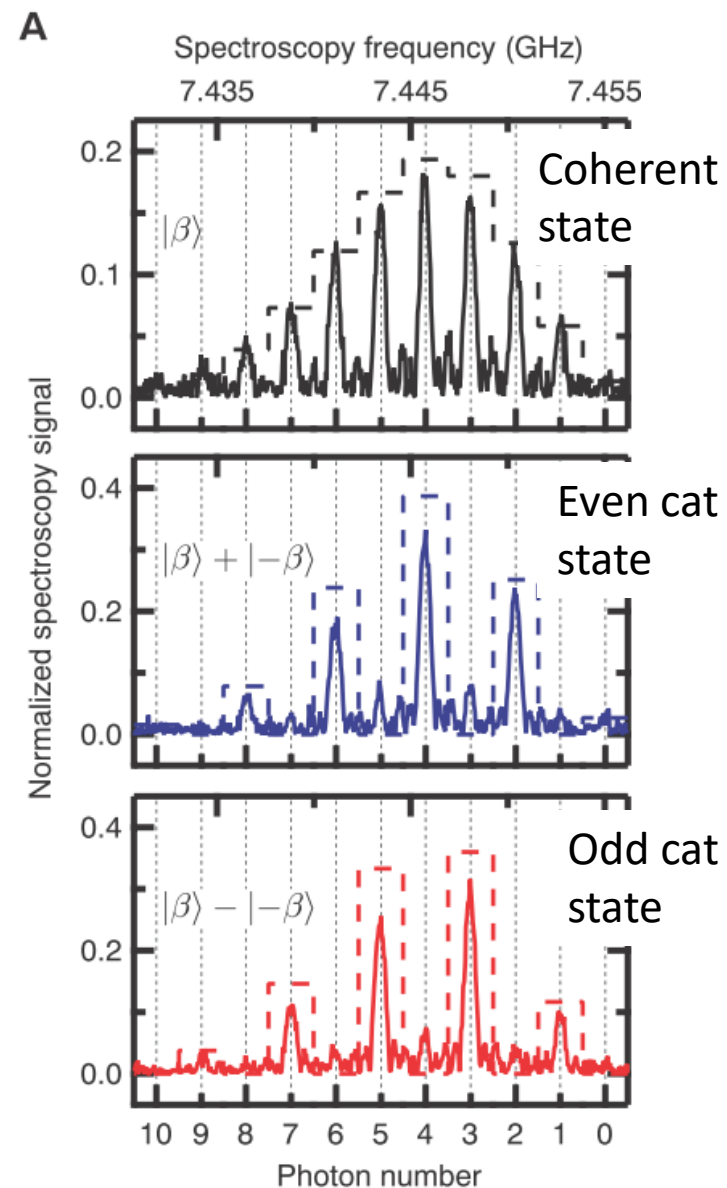
H. Paik, Phys. Rev. Lett. 107 (2011).

First efforts to use
a similar system
for axions

*Akash Dixit
Aaron Chou
Dave Schuster
University of
Chicago*

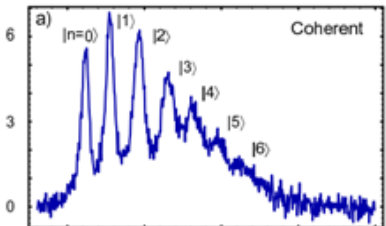
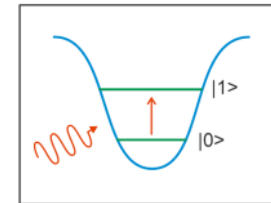
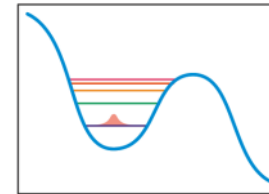
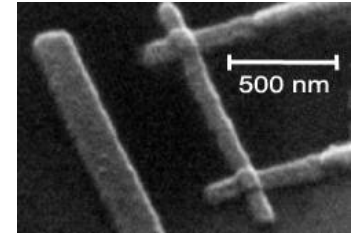
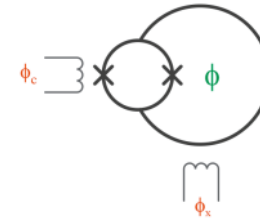


*B. Vlastakis et al., Deterministically Encoding Quantum Information Using 100-Photon Schrödinger Cat States, Science **342**, 607–610 (2013).*



Josephson devices:

- Realization of flexible systems
- Quantum behavior
- Detection of single photons at ~ 10 GHz
- Coupling with cavities at ~ 10 GHz



Thank you!

