

# Enhanced Betatron Radiation from a Laser Wakefield Accelerator in a Long Focal Length Geometry

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# Acknowledgements

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- Betatron Radiation from Laser Wakefield Accelerators.
- Greater than 10-fold increase in Betatron Brightness with long focal length driver.
- Explanation via PIC simulation.

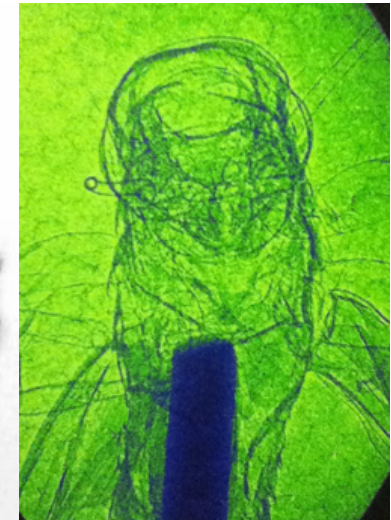
# Importance of X-ray Sources

Imaging: absorption and phase contrast.

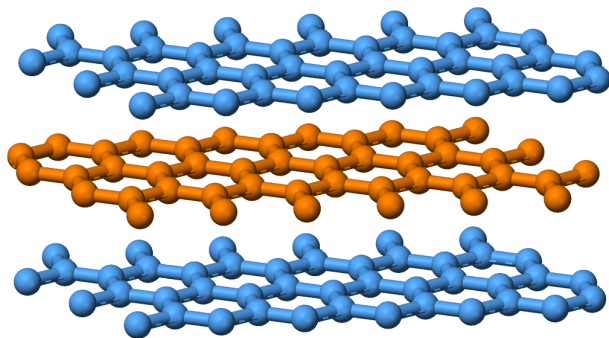
$$\eta = 1 - \delta + i\beta$$

$\delta \gg \beta$  at high E

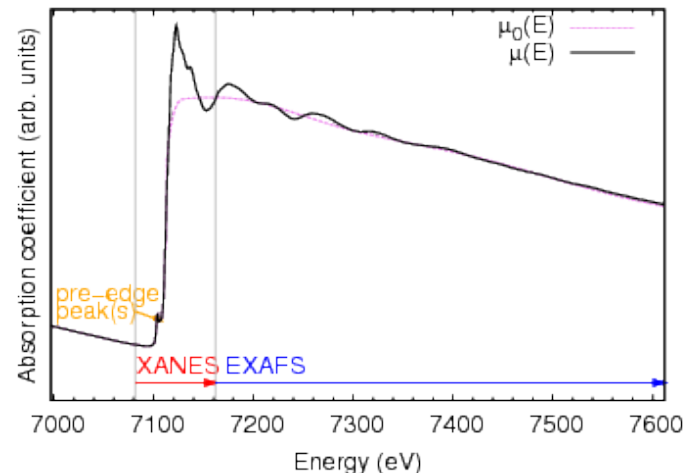
If the x-rays are coherent you can do phase-contrast imaging (xPCI). Pick up gradients in  $\delta$ , get edge enhancement.



Material structure at the atomic scale:



E.g. graphite: distance between planes  $d = 0.335$  nm (3.7 keV)  $\rightarrow$  hard x-ray diffraction.



Atomic structure from 'ringing' in absorption spectrum near K-edge.

# Importance of Bright X-ray Sources

$$B / \text{photons/s/mm}^2 / \text{mrad}^2 / 0.1\% \text{BW}$$

Look at fast processes

Small source size  
means high res.  
lens-less imaging in  
compact geometry.

High flux- high  
SNR or high  
magnification

Important for  
structural studies,  
less so for imaging.

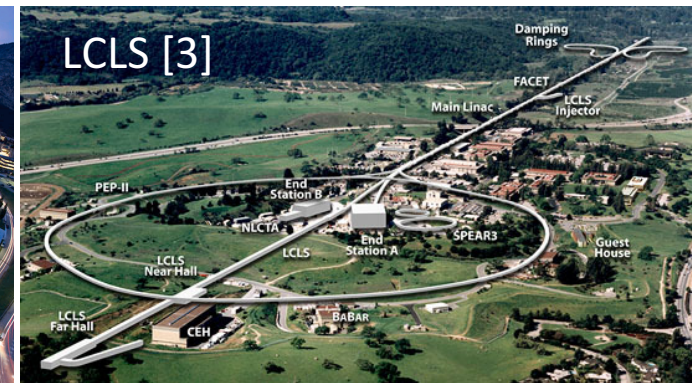
Achievable with sources based on  
relativistic electron beams.

ESRF:  $B = 10^{23}$ - $10^{24}$  @ 10-20 keV

LCLS:  $B = 10^{31}$ - $10^{34}$  @ 10 keV [1]



ESRF [2]



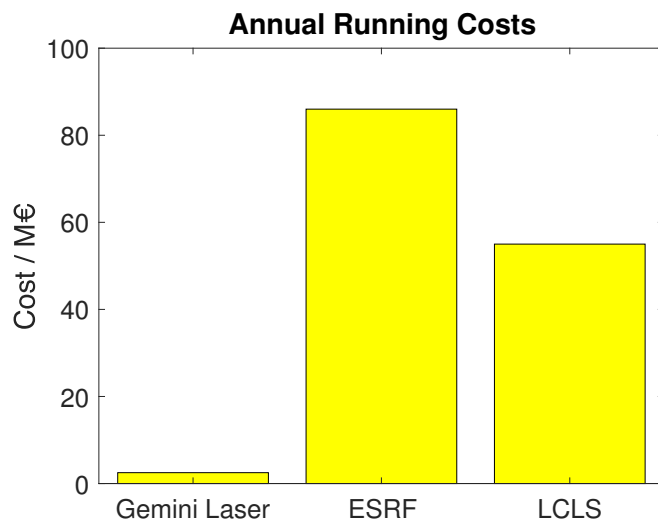
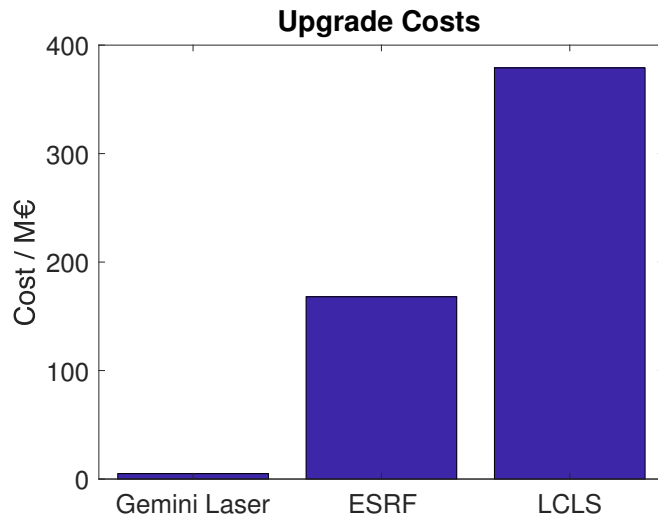
LCLS [3]

[1] [http://photon-science.desy.de/research/studentsteaching/sr\\_and\\_fel\\_basics/fel\\_basics/tdr\\_spectral\\_characteristics/index\\_eng.html](http://photon-science.desy.de/research/studentsteaching/sr_and_fel_basics/fel_basics/tdr_spectral_characteristics/index_eng.html)

[2] <https://www.engelskagymnasiet.se/sites/engelskagymnasiet.se/files/esrf-011.jpg>

[3] [http://web.stanford.edu/group/suncat/cgi-bin/suncat/sites/default/files/lcls\\_0.jpg](http://web.stanford.edu/group/suncat/cgi-bin/suncat/sites/default/files/lcls_0.jpg)

# Why a Laser Based Source?



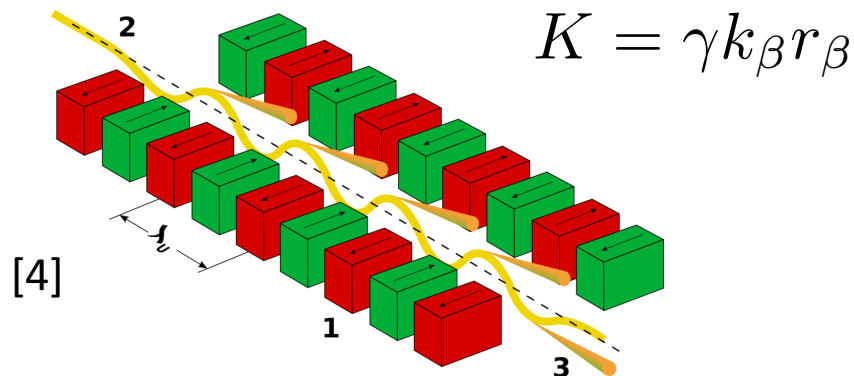
- High rep rate of conventional sources is not always required, especially for dynamic experiments e.g. HEDP.
- Laser pulse duration 10s fs (XFEL is 10s fs, synchrotron is 10s-100 ps).
- Large market: 47 synchrotron light sources worldwide, 3 current hard x-ray FEL's. 4 more being built.





# Betatron Radiation = Plasma Wiggler Radiation

Focussing forces of bubble means it acts as a plasma wiggler for electrons.

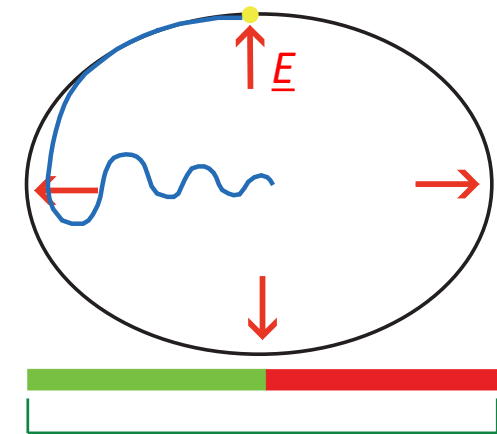


Well collimated. Emission cone angle  $\theta \approx K/\gamma$

$K > 1$  for betatron source: hard, broadband radiation.

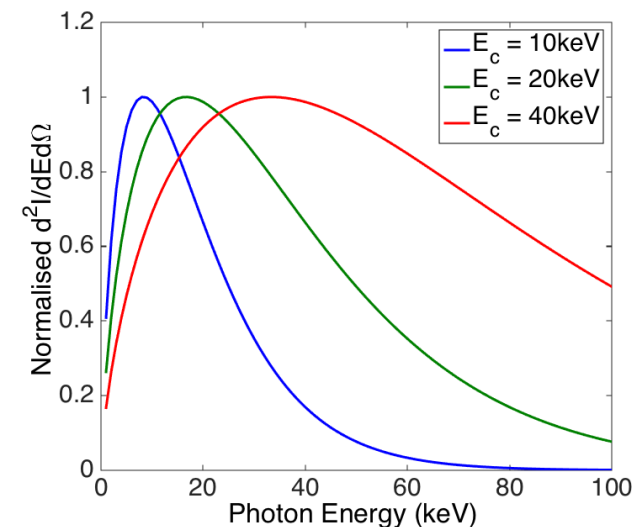
Well approximated by on axis synchrotron spectrum [5].

Spectrum defined by critical energy  $E_c = \frac{3\hbar}{4c} \gamma^2 \omega_p^2 r_{\beta}$



$$\lambda_b \approx \lambda_p \sqrt{a_0}$$

$$a_0 = \gamma_q \beta_q \approx 0.86 \sqrt{I_{18} \lambda^2}$$



[4] CC BY-SA 3.0, <https://commons.wikimedia.org/w/index.php?curid=537945>

[5] E. Esarey et. Al. Physical Review E, 65(5):056505, May 2002.

# High Brightness Betatron Experiments

Laser	$a_0$	$E_c$ / keV	No. photons	Brightness ph./s/mm <sup>2</sup> / mrad <sup>2</sup> /0.1% BW	Source size (um)
Hercules [6]	4.7	8-21	$5 \times 10^7$	$1 \times 10^{22}$	1
Gemini [7]	3.4	20-33	$1.3 \times 10^9$	$\sim 1 \times 10^{23}$	1.6
Texas Petawatt [8]	-	8-18	$10^8$ - $10^9$	Of order $10^{22}$ - $10^{23}$	-
Jupiter Laser Facility [9]	$\sim 3$	15	$5 \times 10^8$	$1 \times 10^{23}$	



Ideal way to improve brightness: increase number of photons/mrad<sup>2</sup>. Consequences:

Improve single-shot image quality, probe matter via XANES/ EXAFS (require  $>10^6$  photons/eV for good stats [10]), Laue diffraction?

[6] S. Kneip et. al. Nature Physics 6(12):980-983,2010.

[8] X. Wang et. al. Nature Communications 4:1988, 2013.

[10] F. Albert et al. Plasma Physics and Controlled Fusion 56(8):084015.

[7] J.M. Cole, J.C. Wood et. al. Scientific Reports 5:13244, 2015.

[9] Yan W. et al. PNAS 111(16):5825-5830, 2014.



# How to Increase the Number of Betatron Photons

- Total emitted energy:  $W_{tot} \propto N_e \gamma^{5/2} r_\beta^2$
- Maximum electron energy:  $\gamma \propto \frac{1}{n_e}$
- Matched spot size for good laser guiding:

$$k_p w_0 \approx 2\sqrt{a_0} \implies w_0 \propto 1/\sqrt{n_e}$$

- Long focal lengths (big focal spots) allow us to be matched at lower density, increasing maximum energy gain.

# How to Increase the Number of Betatron Photons

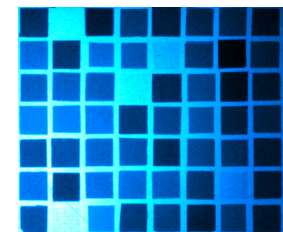
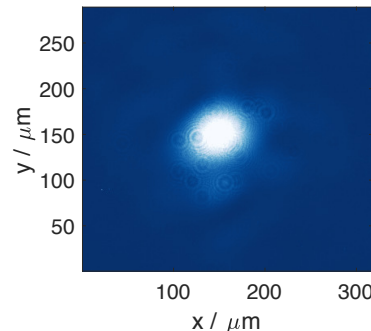
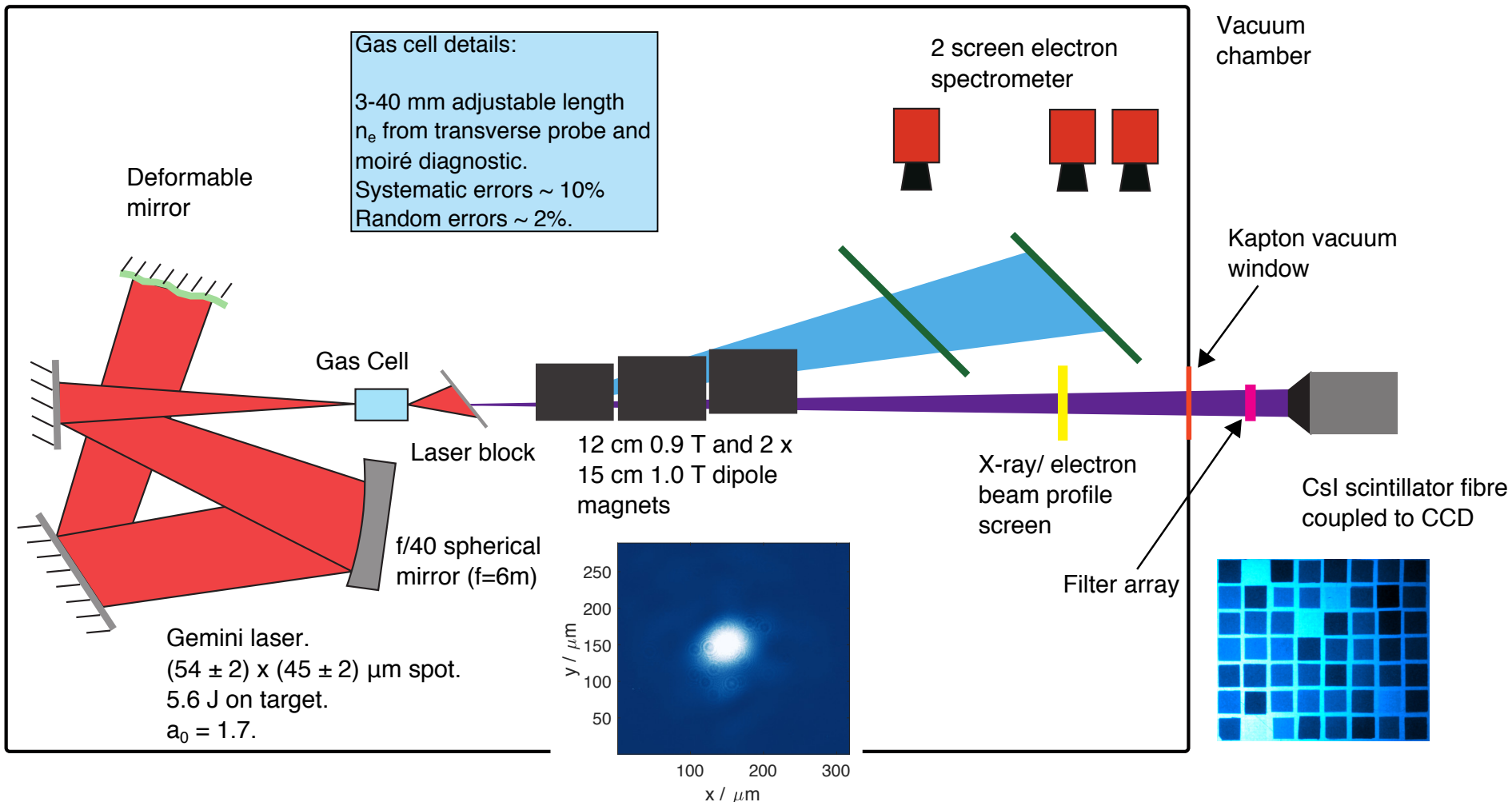
- Total emitted energy:  $W_{tot} \propto N_e \gamma^{5/2} r_\beta^2$
- Maximum electron energy:  $\gamma \propto \frac{1}{n_e}$
- Matched spot size for good laser guiding:

Typical electron energy at Gemini with f/20 focussing:  
 $0.7 \pm 0.1$  GeV.  
In optimal conditions:  
 $E \gtrsim 1$  GeV.

$$k_p w_0 \approx 2\sqrt{a_0} \implies w_0 \propto 1/\sqrt{n_e}$$

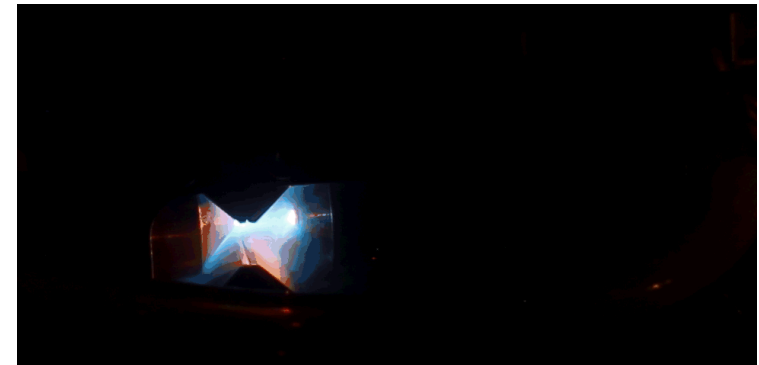
- Long focal lengths (big focal spots) allow us to be matched at lower density, increasing maximum energy gain.

# Experiment Set-up

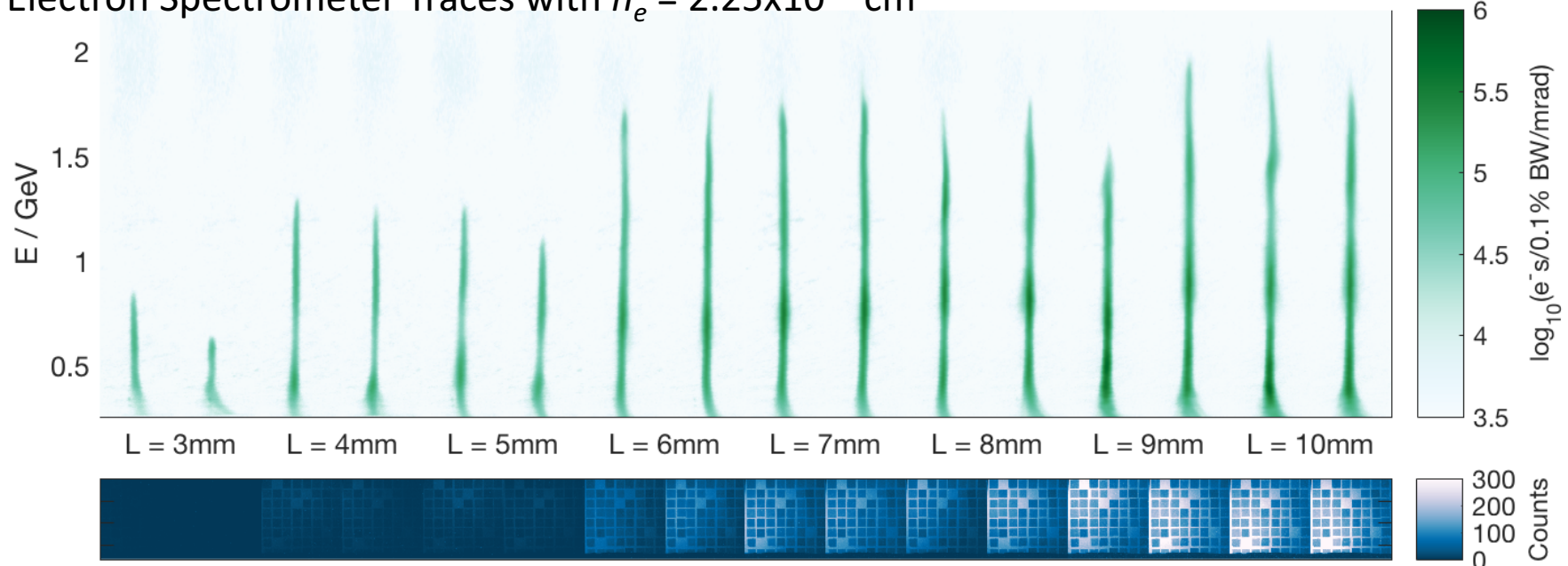


# Electron and X-ray Results I

Performed scans at constant plasma density changing the length of the He plasma in the range 4-40 mm.

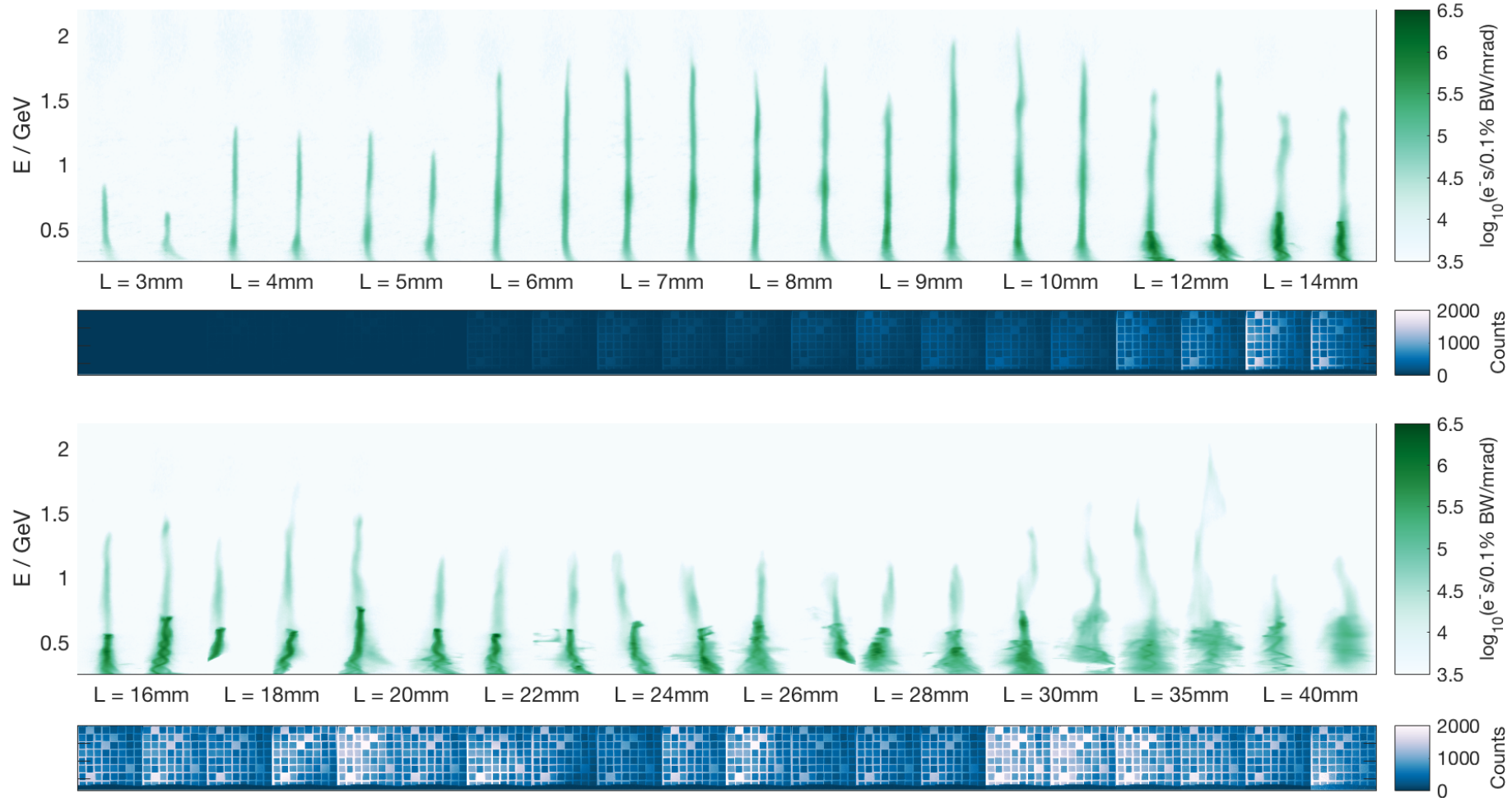


Electron Spectrometer Traces with  $n_e = 2.25 \times 10^{18} \text{ cm}^{-3}$



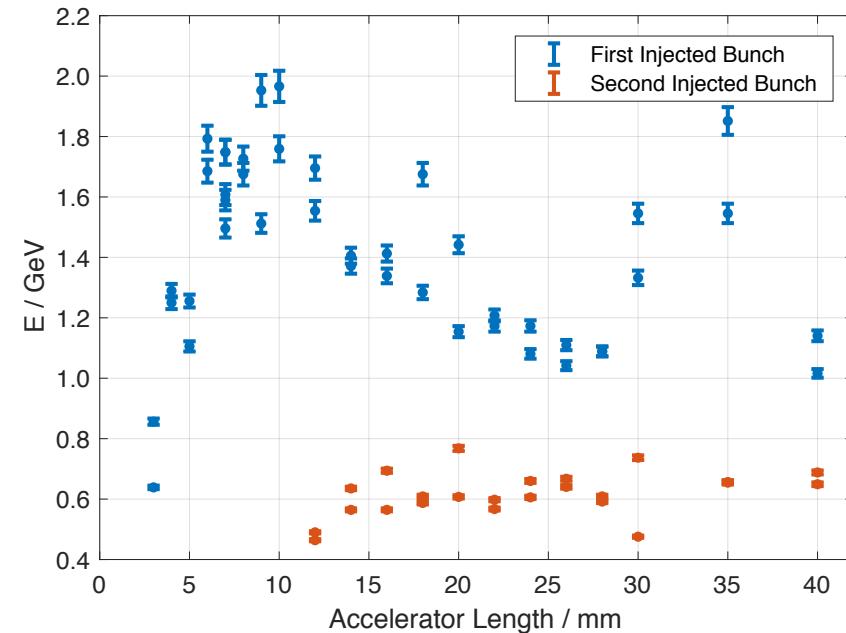
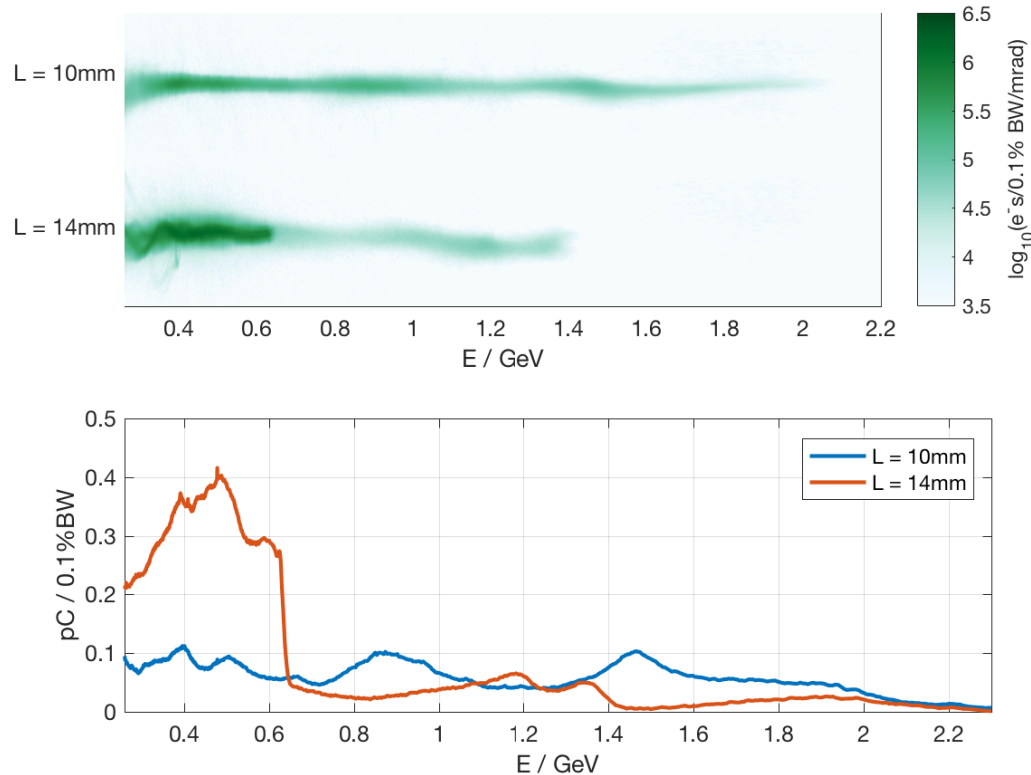
Betatron x-ray Images of Filter Array

# Electron and X-ray Results II



Injection of 2<sup>nd</sup> bunch with high charge/0.1%BW and large transverse momentum.  
Second injection correlates with a large increase in the number of x-rays being detected.

# Electron Beam Energy Spectrum



Above: Peak electron energies vs. accelerator length. Errors from uncertainty in screen position.

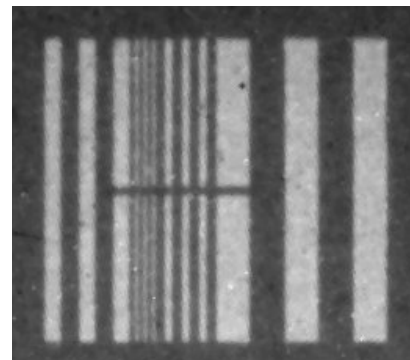
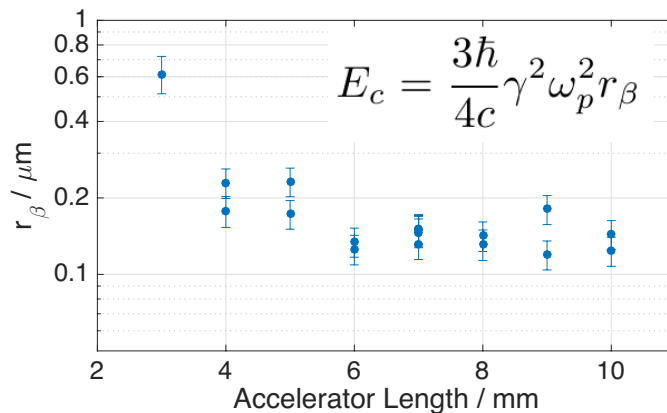
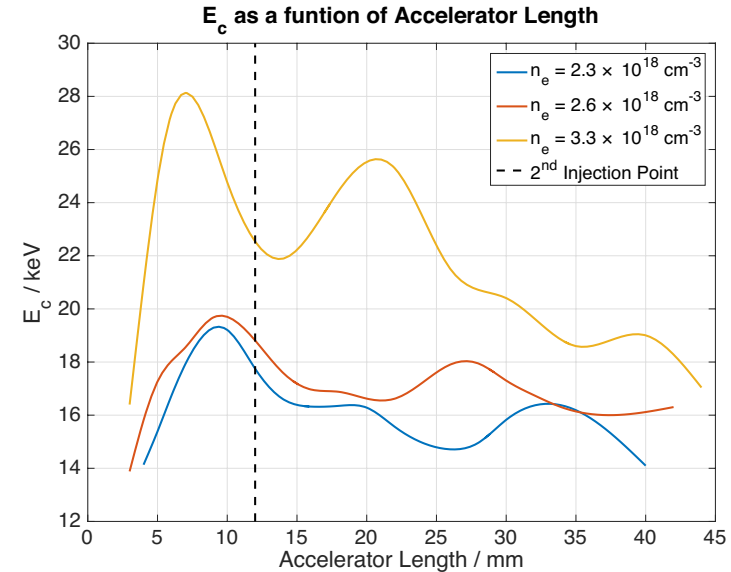
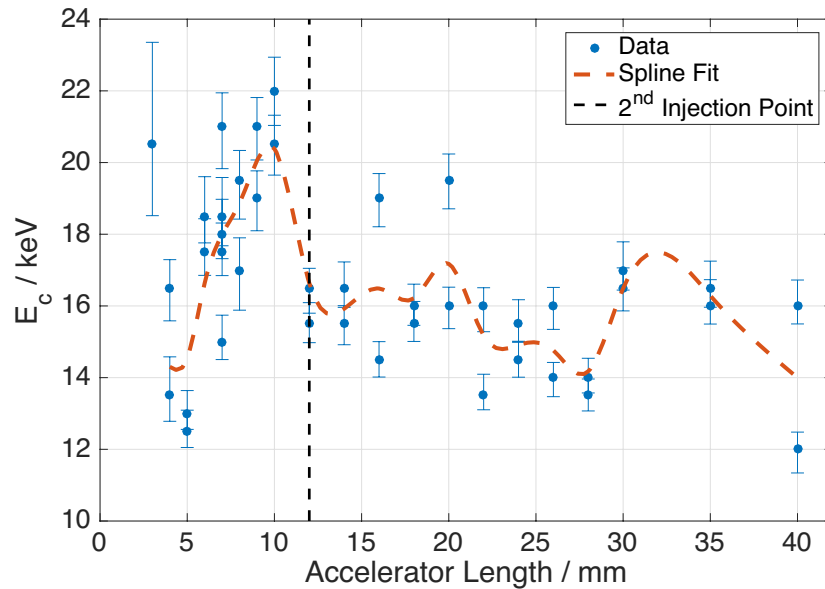
Second bunch has up to 4x the charge/ unit bandwidth.

Injection close to dephasing point of first bunch.

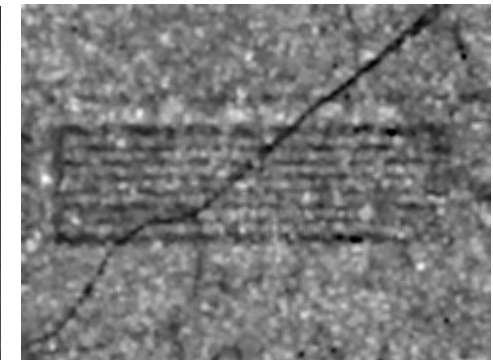
Limited acceleration of second bunch: never beyond 800 MeV.



# Betatron Spectral Characteristics

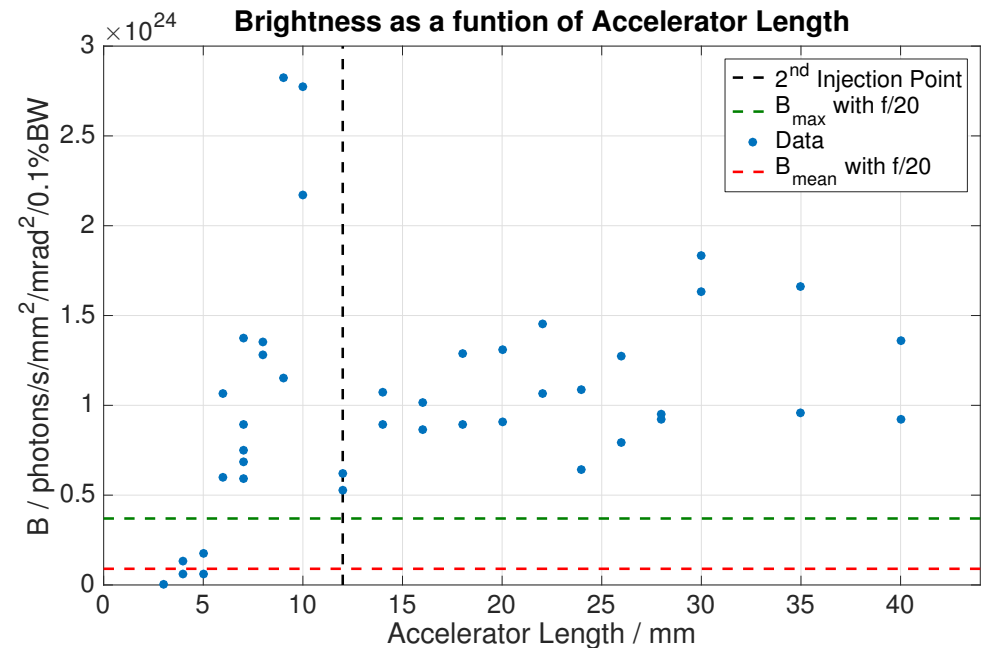
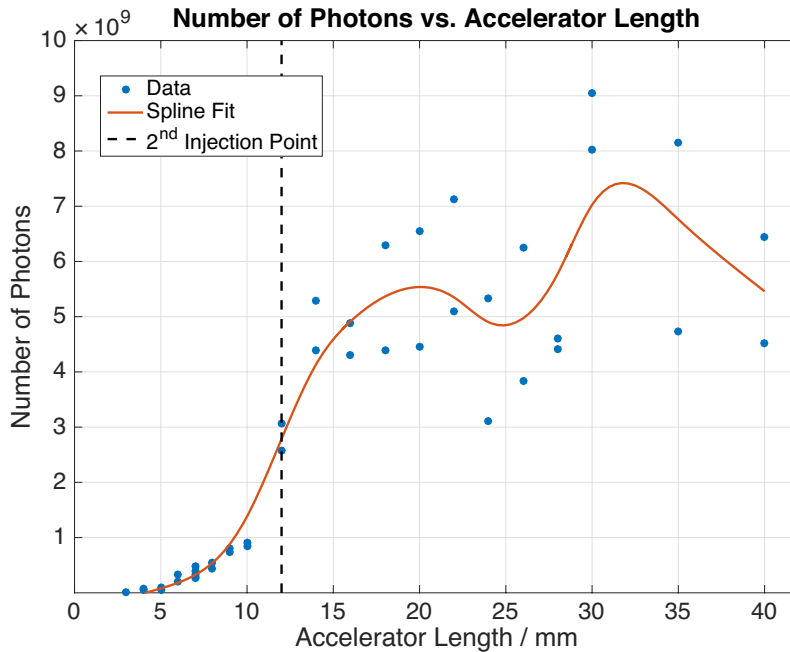


10  $\mu\text{m}$  bars, M=11



4  $\mu\text{m}$  bars, M=29

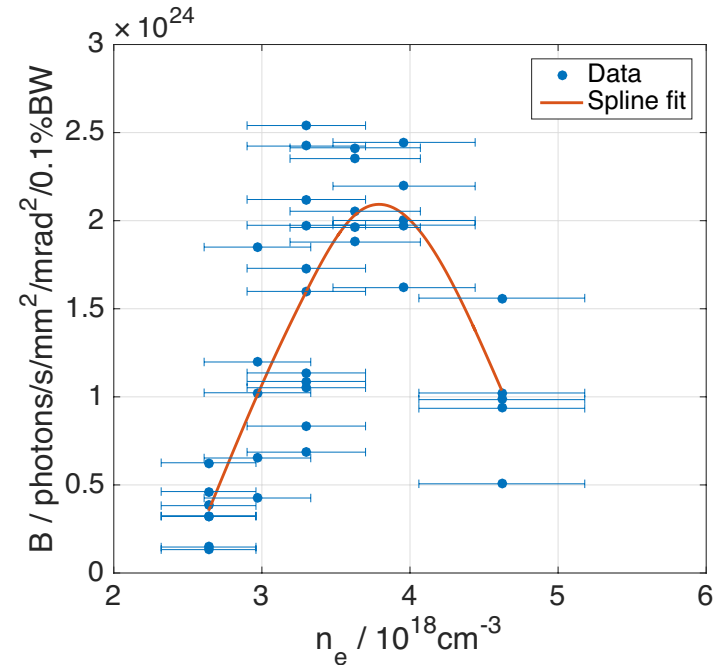
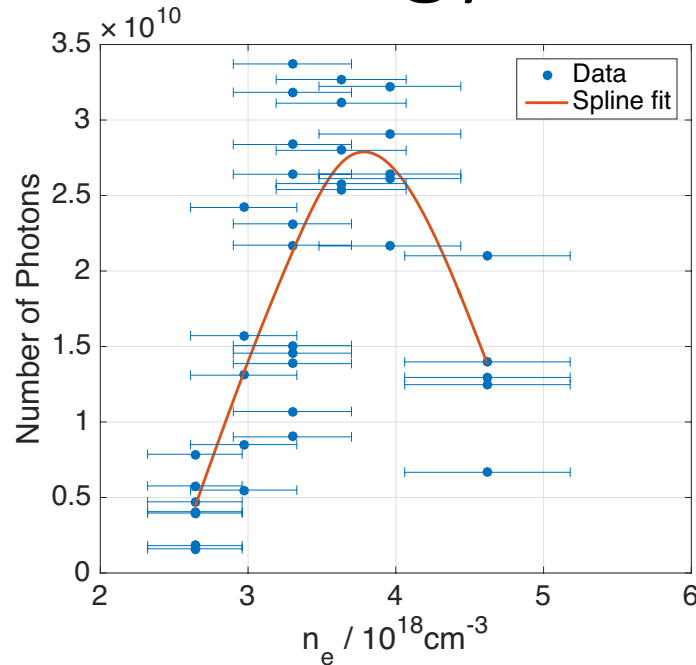
# Betatron Brightness



First injected bunch: high brightness driven by small source size.

With second bunch: relative reduction in brightness from rapid increase in source size.  
But we see a dramatic increase in the number of photons- good for applications!

# Further Brightness Increase with Laser Energy



Increase in laser energy: 5.6 J to 11.3 J.  $a_{0,max}$ : 1.7 to 2.4

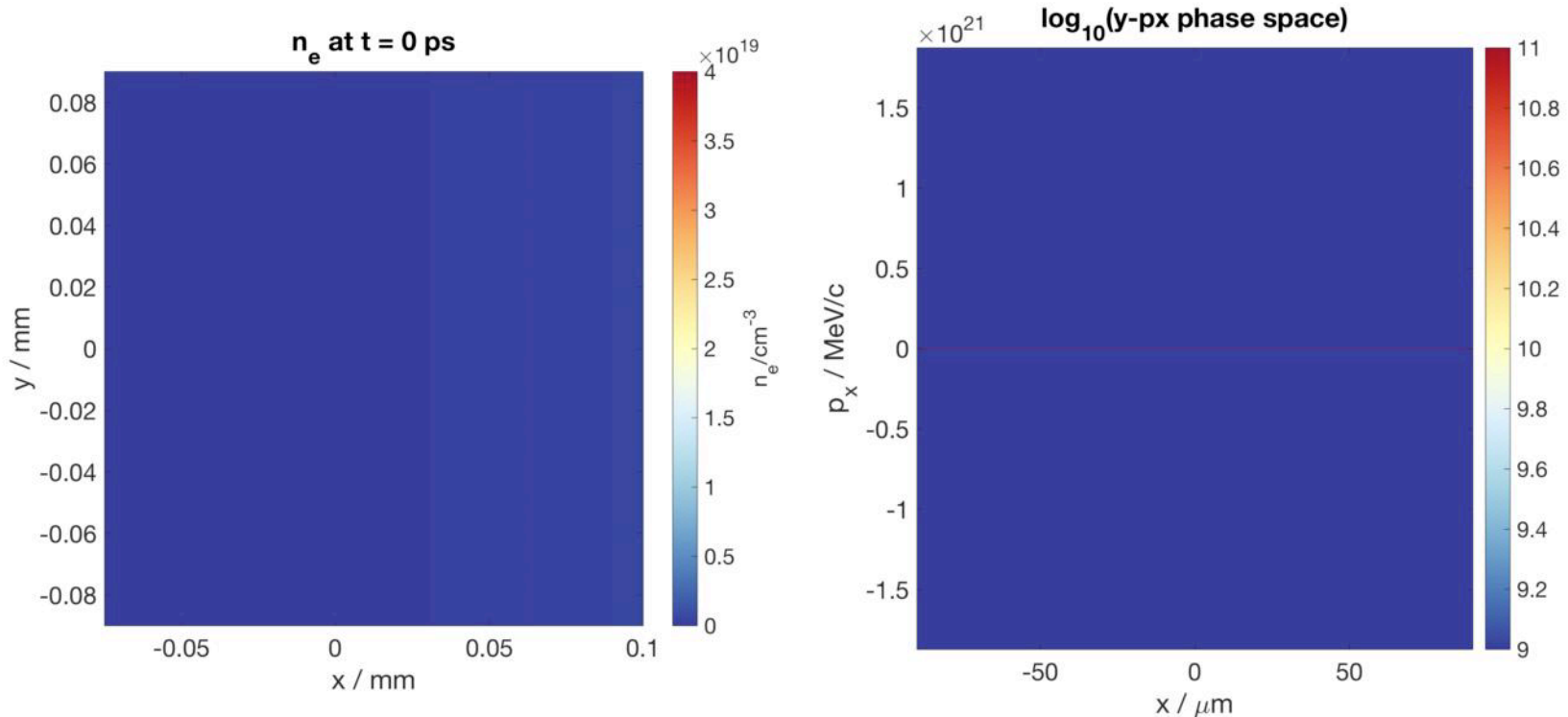
~70% increase in brightness in high photon regime. Approx. 6x more photons.

Good for applications.  $1.4 \times 10^6$  photons per eV @ 6-8 keV. B. Kettle talk, few shot XANES spectra.

J. Cole talk, high quality medically relevant imaging. Both in WG4 right now.

# PIC Simulations I

$$n_e = 2.5 \times 10^{18} \text{ cm}^{-3}, a_0 = 3.8$$



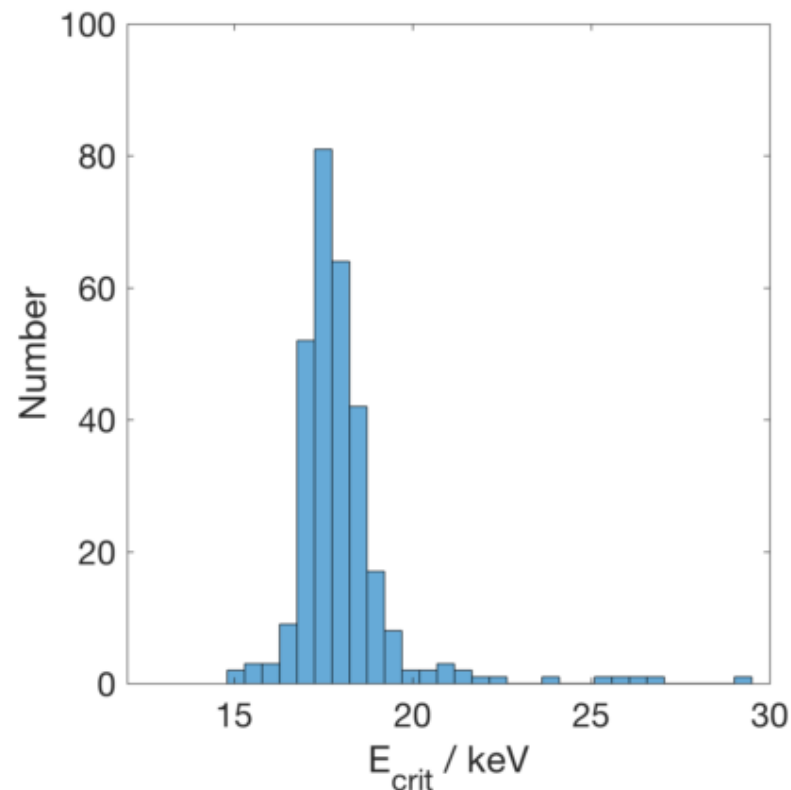
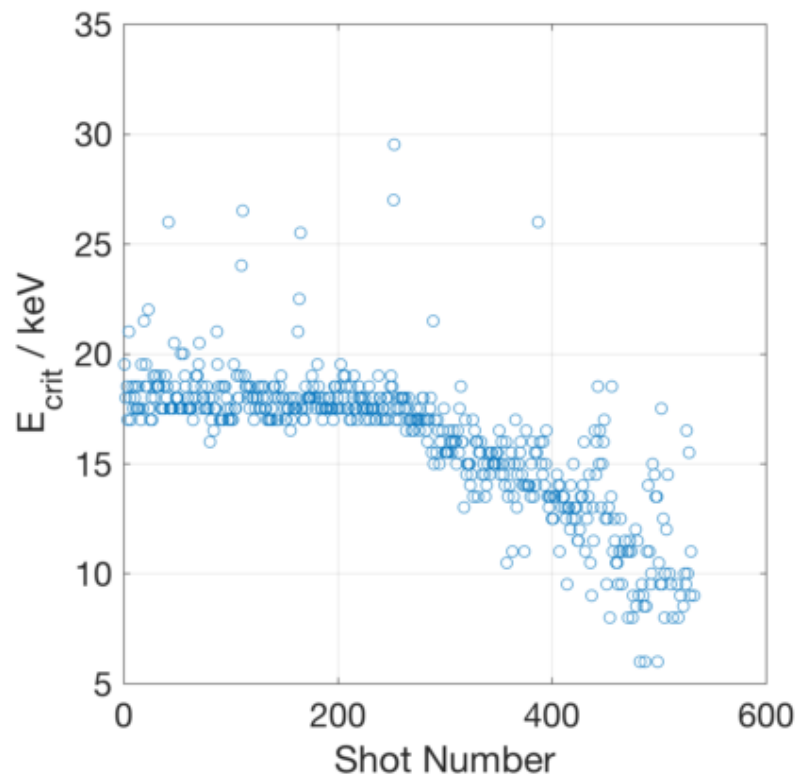
- First bunch injected very early- during wakefield formation. Second, high charge/ unit BW, bunch injected due to bubble expansion [11][12].
- Second bunch undergoes large oscillations ( $W_{\text{tot}} \propto r_{\beta}^2$ ).
- Oscillations seem linked to bubble oscillation- explains coherent oscillation structure of some beams from the experiment.

Aakash Sahai

[11] S. Kalmykov et al. Phys. Rev. Lett. 103(13):1-4, 2009.

[12] S. A. Yi et al., PPCF 19(3):014012, 2012.

# Reproducibility



- Dynamic injection and laser behaviour.
- Reproducible betatron x-ray spectrum.
- Excellent reproducibility for 300 shots.
- Mean: 18.0 keV. Standard Deviation: 1.6 keV.

# Conclusions

- Long focal length wakefield experiment produced high brightness betatron beams in 2 regimes.
- First injection: few hundred nm source size.
- Second injection: high brightness and photon number.  
>  $10^6$  photons/eV,  
>  $10^{24}$  photons/s/mm<sup>2</sup>/mrad<sup>2</sup>/0.1%BW
- Results backed up by PIC simulation.

