



# Longitudinal and transverse ion cloud dynamics in an electron ring in presence of electromagnetic fields and gaps



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# Overview

## What

- We want to locate ion clouds, generated by ionisation of the residual gas, in an electron storage ring

## Why

- To minimize the beam dynamics degradation because of the ion cloud, a beam stability of 40  $\mu\text{m}$  is needed in the interaction point (IP)

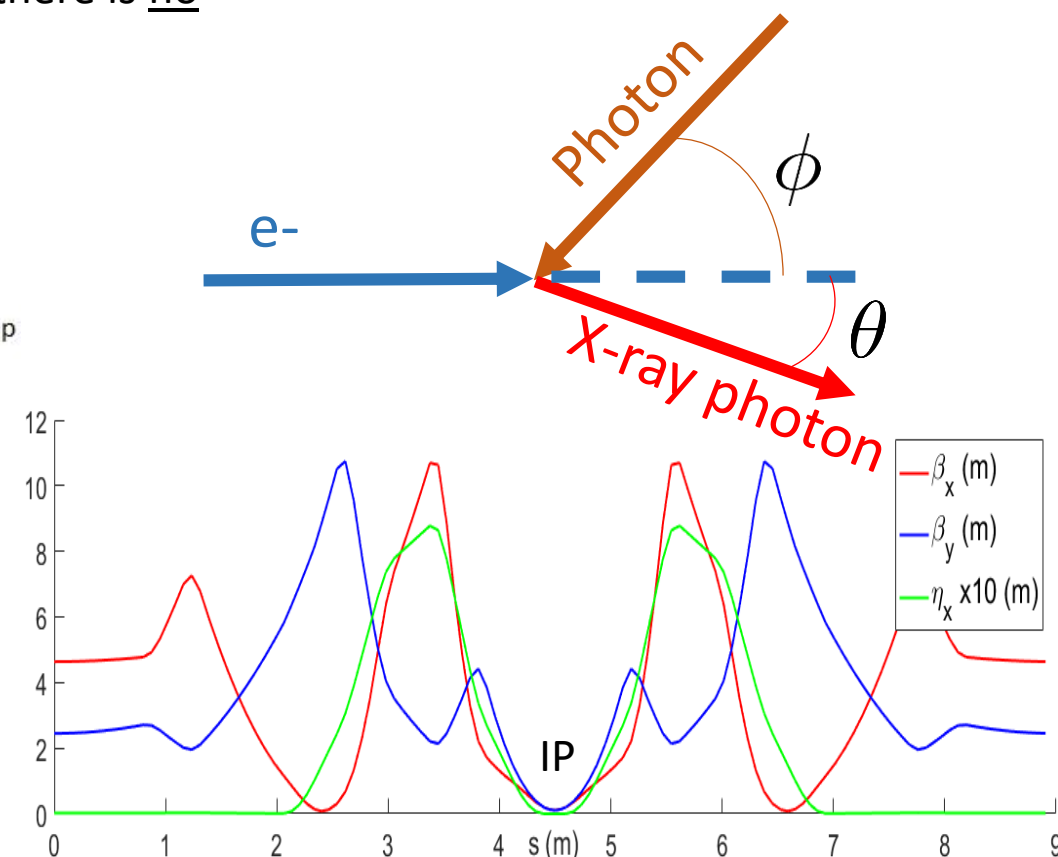
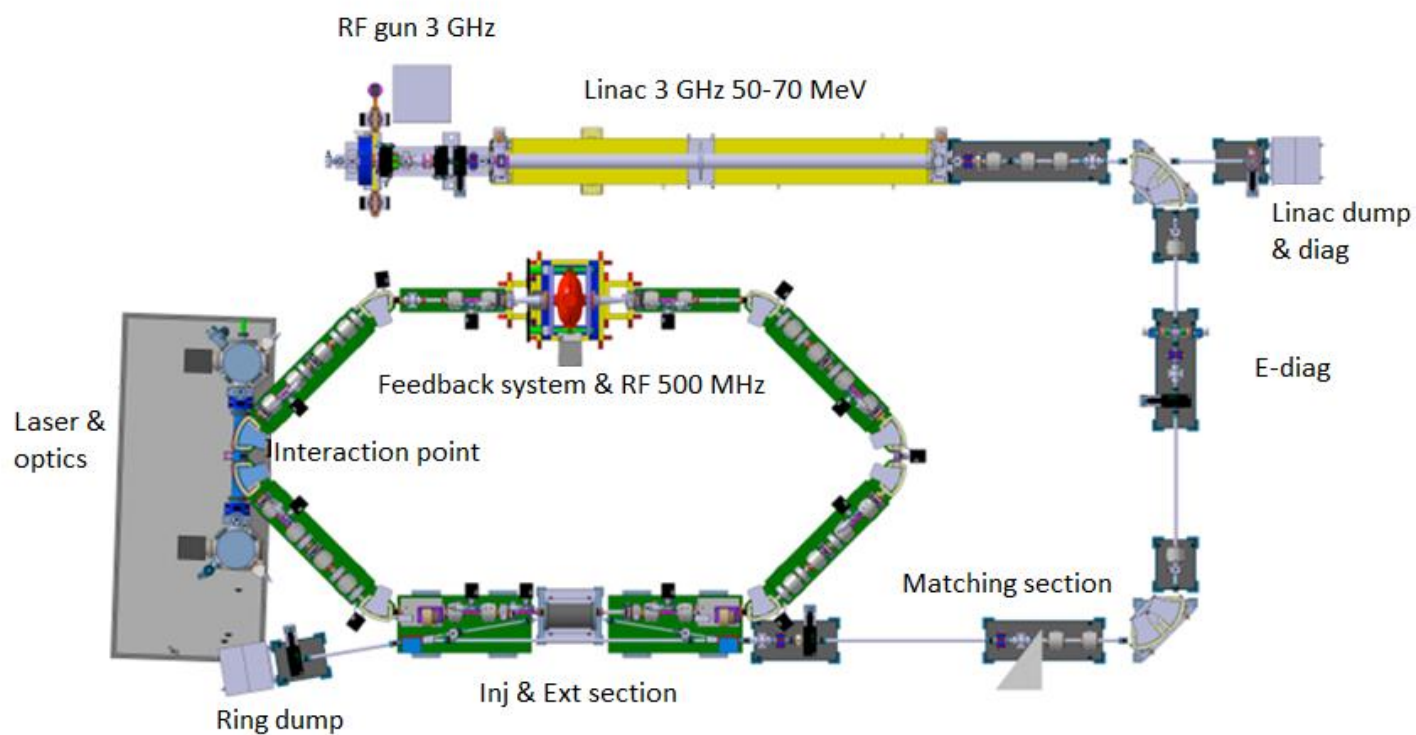
## How

- By understanding the physics of the beam-ion interaction, simulating the ion dynamics and using mitigations strategies

# ThomX, a storage ring without damping

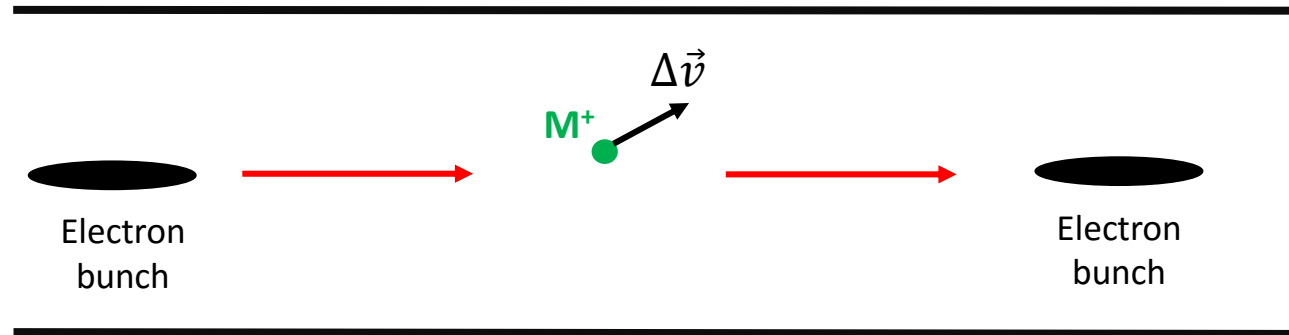
- ThomX is a Compton Backscattering Source (CBS) of X rays under construction at LAL.
- The e- bunch is stored for 20 ms then it is dumped while a new one is injected.
- Because of the short storage time and the low electron energy there is no synchrotron radiation damping.

Flux =  $10^{11-13}$  X/s  
 $E_X^{max}$  = 46-90 keV  
 Divergence = 10 mrad



# A model to describe the beam-ion interaction

The model gives the kick that an ion will feel when the electron beam is going through the beam pipe:



The model:

- Bassetti-Erskine formula<sup>1</sup> for transverse dynamics

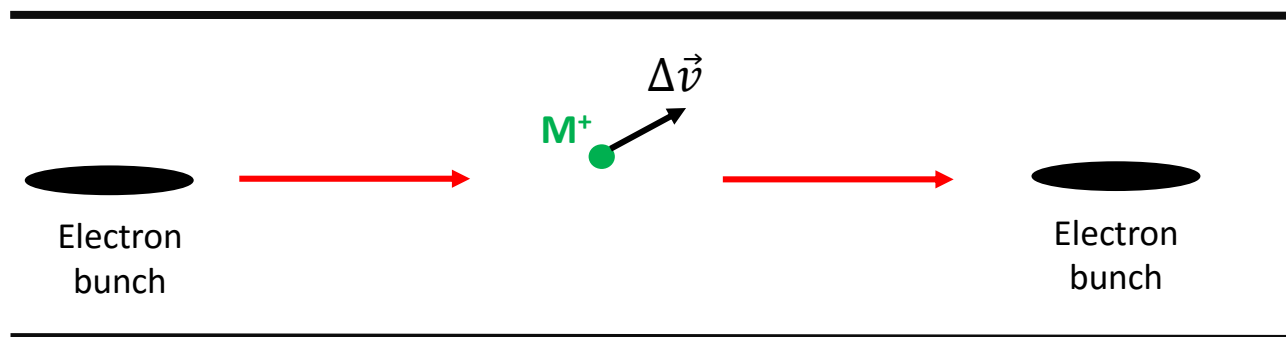
Assumptions:

- “Strong-weak” model of the beam-beam interaction
- Electron bunch is supposed Gaussian

$$i\Delta v_x + \Delta v_y = \frac{\overset{\text{Number of e- in a bunch}}{-NK\sqrt{\pi}}}{\sqrt{2(\sigma_x^2 - \sigma_y^2)}} \left( w \left( \frac{\overset{\text{Ion position}}{x + iy}}{\sqrt{2(\sigma_x^2 - \sigma_y^2)}} \right) - e^{-\left(\frac{x^2}{2\sigma_x^2} + \frac{y^2}{2\sigma_y^2}\right)} w \left( \frac{\overset{\text{Beam size at the ion position}}{x \frac{\sigma_y}{\sigma_x} + iy \frac{\sigma_x}{\sigma_y}}}{\sqrt{2(\sigma_x^2 - \sigma_y^2)}} \right) \right)$$

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The model gives the kick that an ion will feel when the electron beam is going through the beam pipe:



The model:

- Bassetti-Erskine formula<sup>1</sup> for transverse dynamics
- Sagan formula<sup>2</sup> for longitudinal dynamics

Assumptions:

- “Strong-weak” model of the beam-beam interaction
- Electron bunch is supposed Gaussian
- The beam trajectory is quasi-parallel to the longitudinal axis

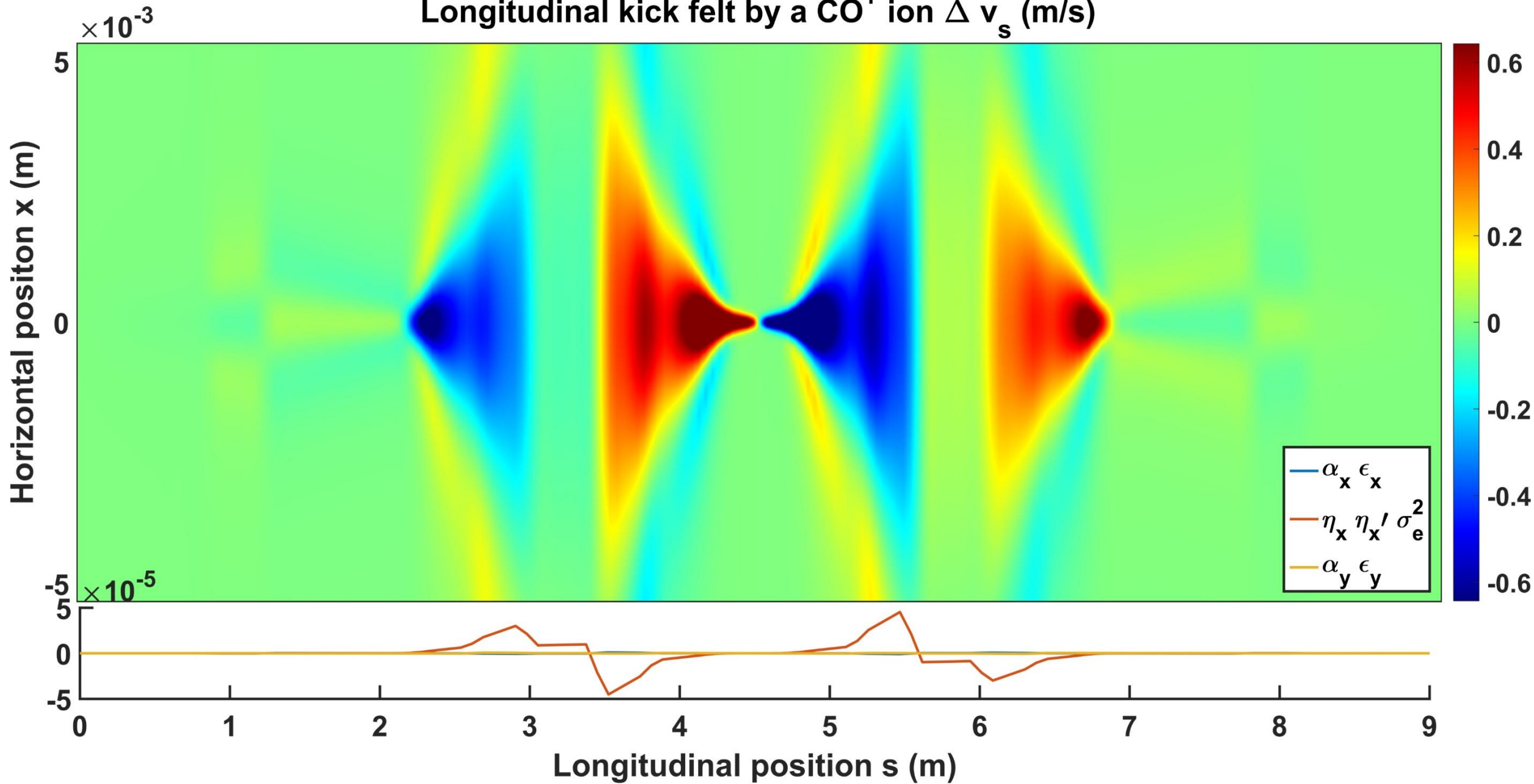
Energy spread

Dispersion function and derivative

$$\Delta v_s = (-\alpha_x \epsilon_x + \eta \eta' \sigma_\epsilon^2) \frac{\partial \Delta_x}{\partial x} - \alpha_y \epsilon_y \frac{\partial \Delta_y}{\partial y}$$

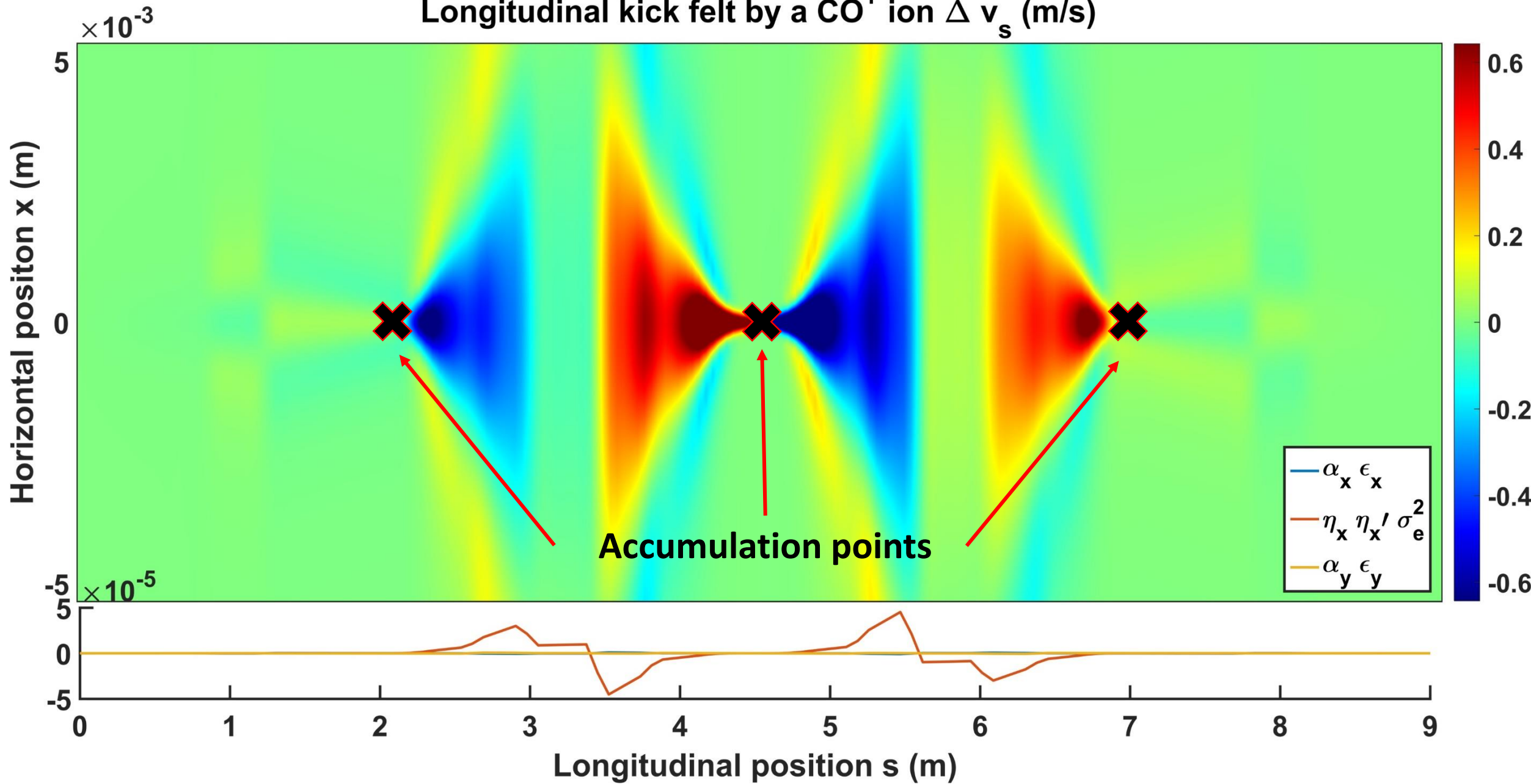
**The ion dynamics are determined by the optics and the lattice design**

# Longitudinal kick felt by a $\text{CO}^+$ ion $\Delta v_s$ (m/s)

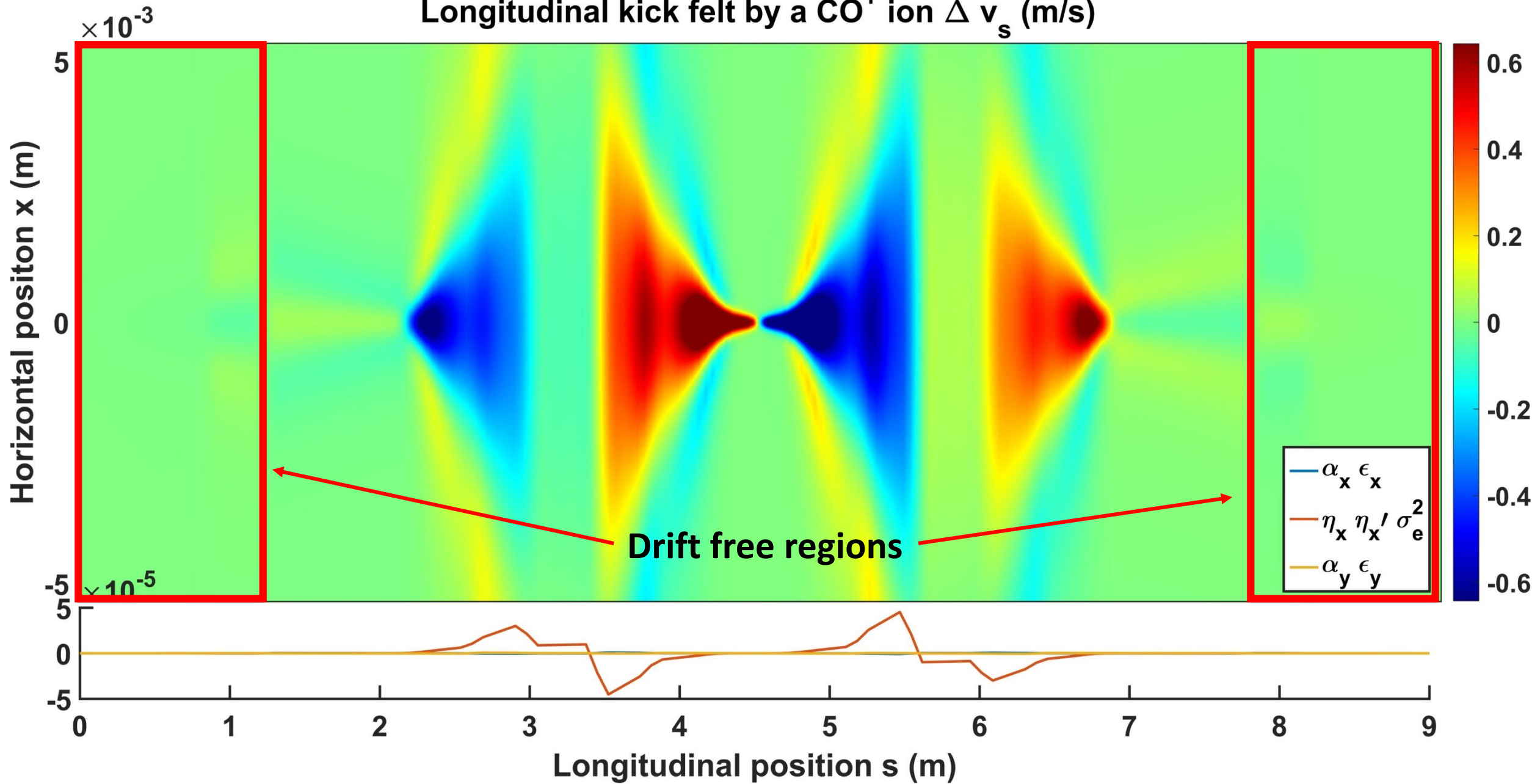




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# NUAGE, ion cloud tracking

The following effects impact the ion cloud dynamics:

- Beam-ion interaction
  - Ion trapping in magnetic field
  - Clearing electrodes
  - Gaps in bunch train
  - Ion cloud collective effects (self space charge, ...)
- Included in
- NUAGE a data parallel Matlab code for ion cloud tracking developed at LAL
- Not included in NUAGE
- 

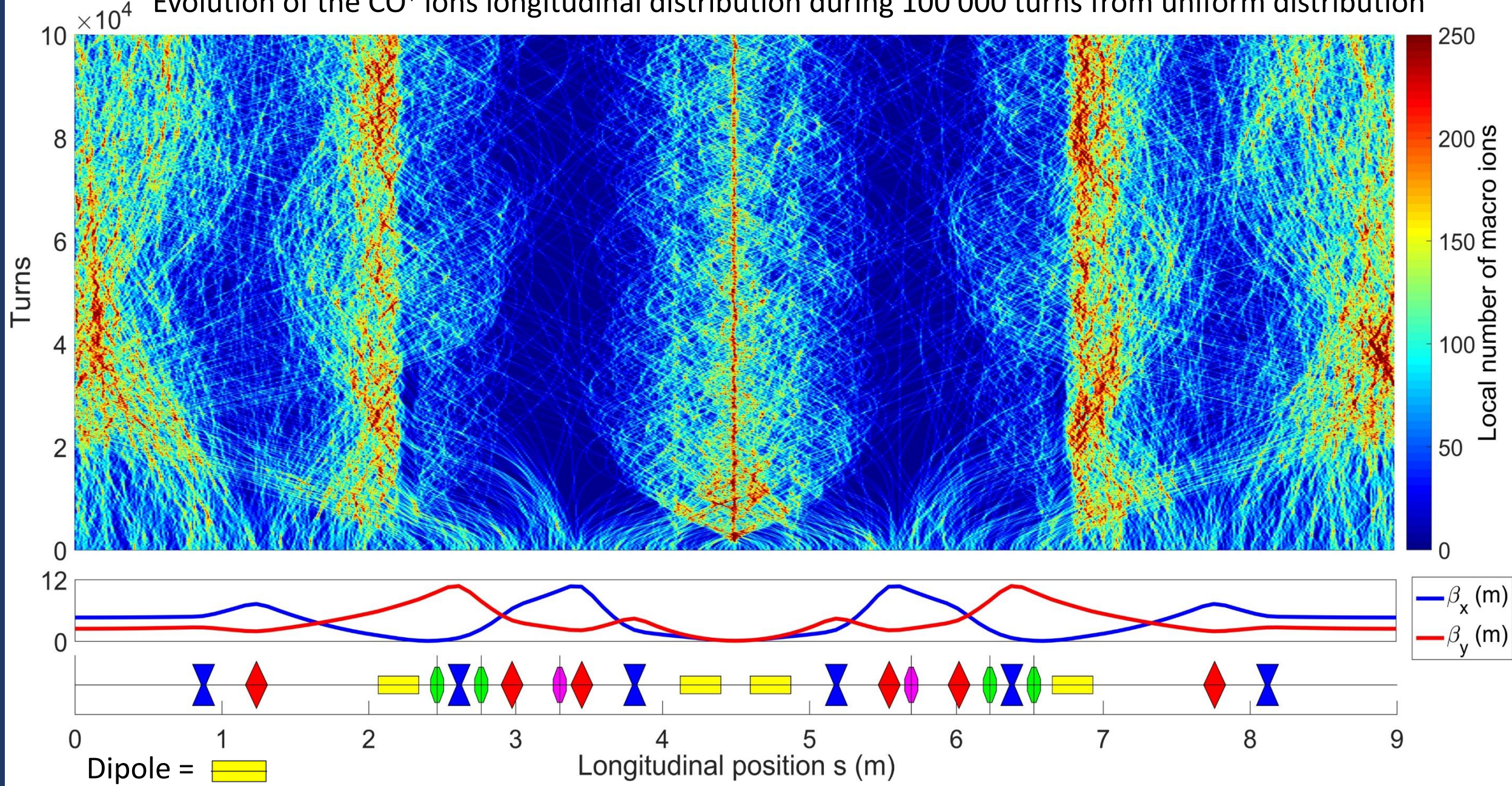
Tracking in magnetic elements (dipoles, quadrupoles) by solving ODE

Interpolation of the 3D field maps of the clearing electrodes

In NUAGE, the ion cloud is defined at the start of the simulation (composition, number of ions, ...) and tracked during a fixed time length. The simulation does not include the generation of new ions during the tracking.

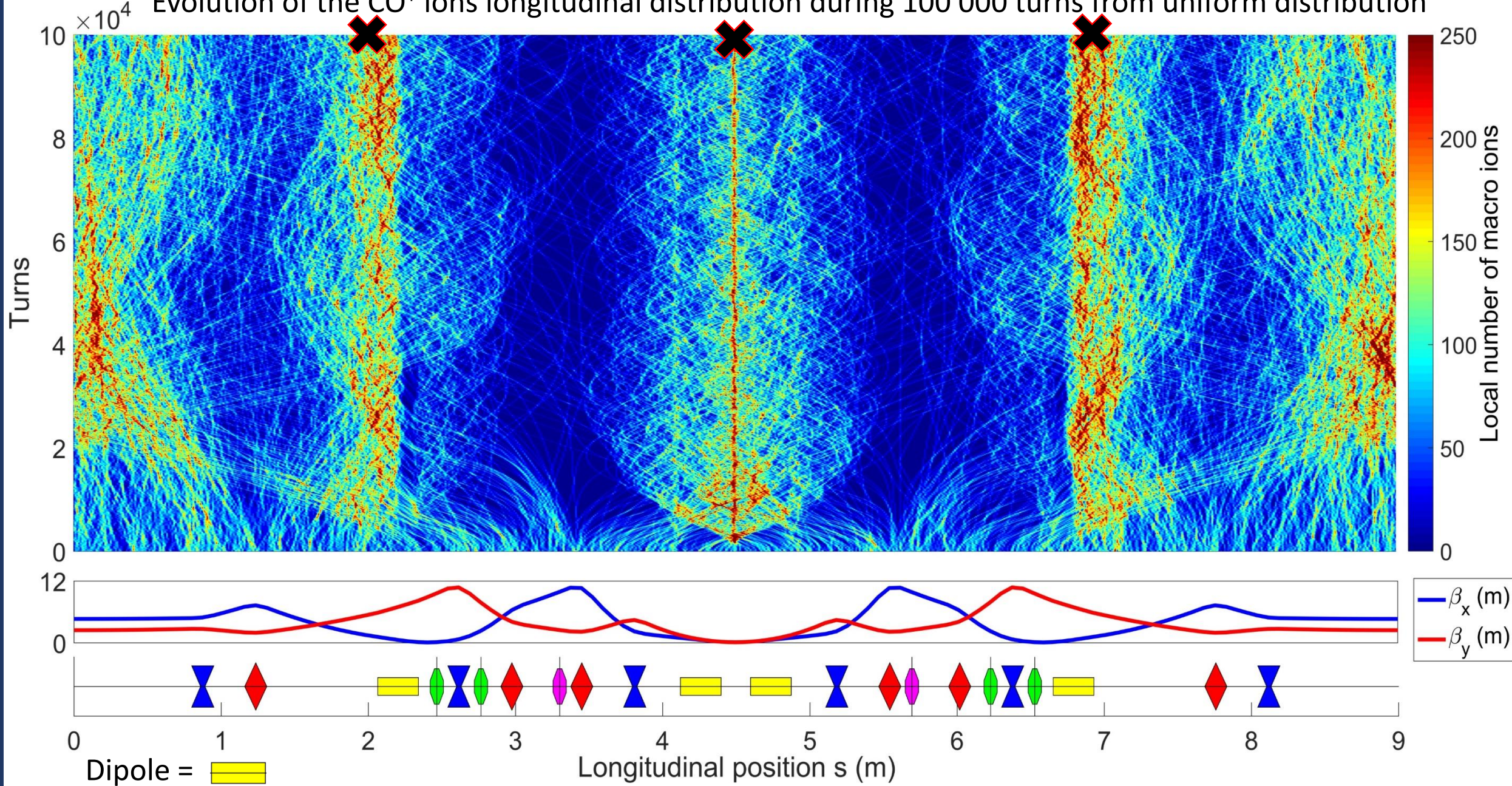


Evolution of the CO<sup>+</sup> ions longitudinal distribution during 100 000 turns from uniform distribution



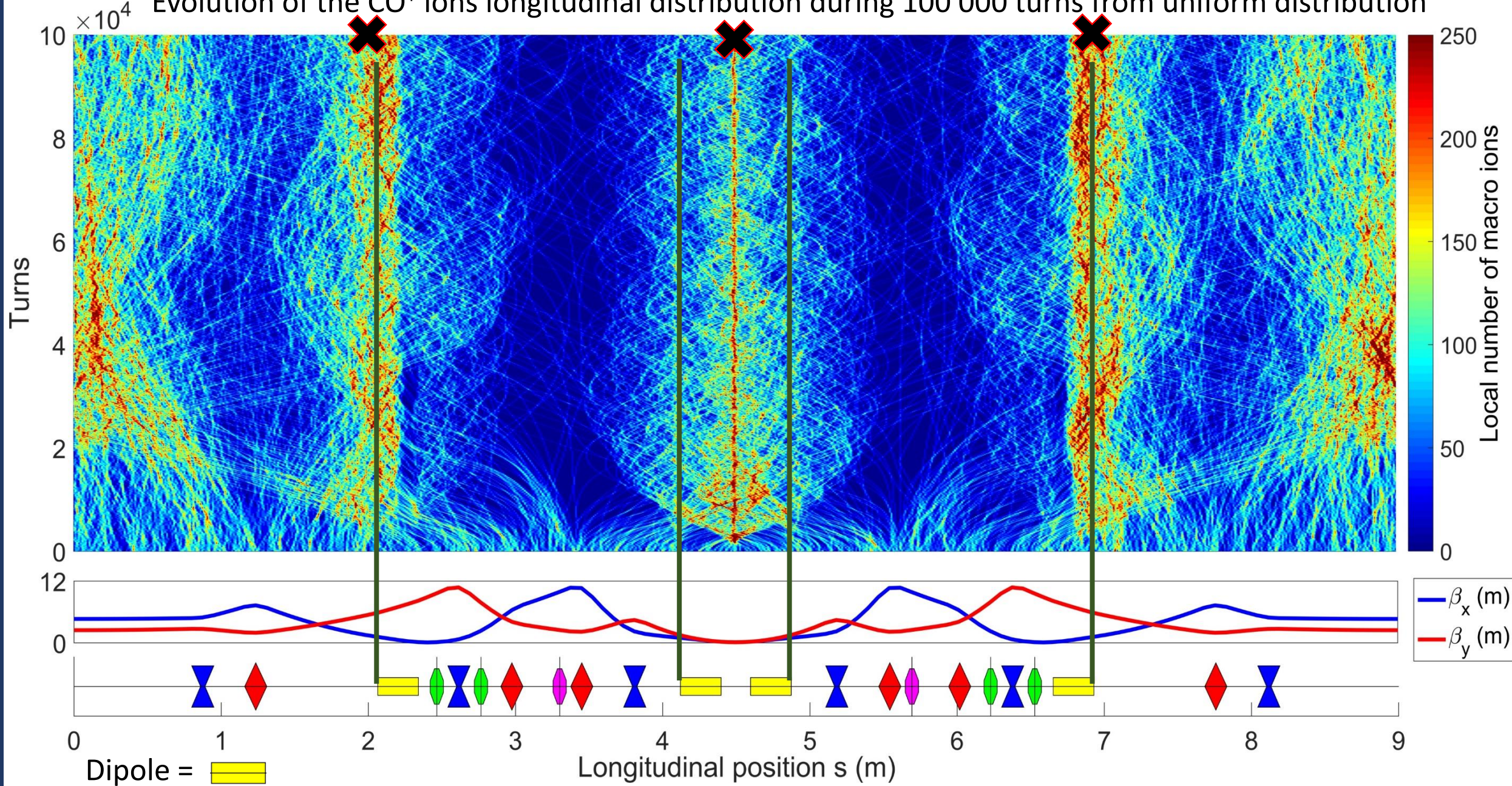


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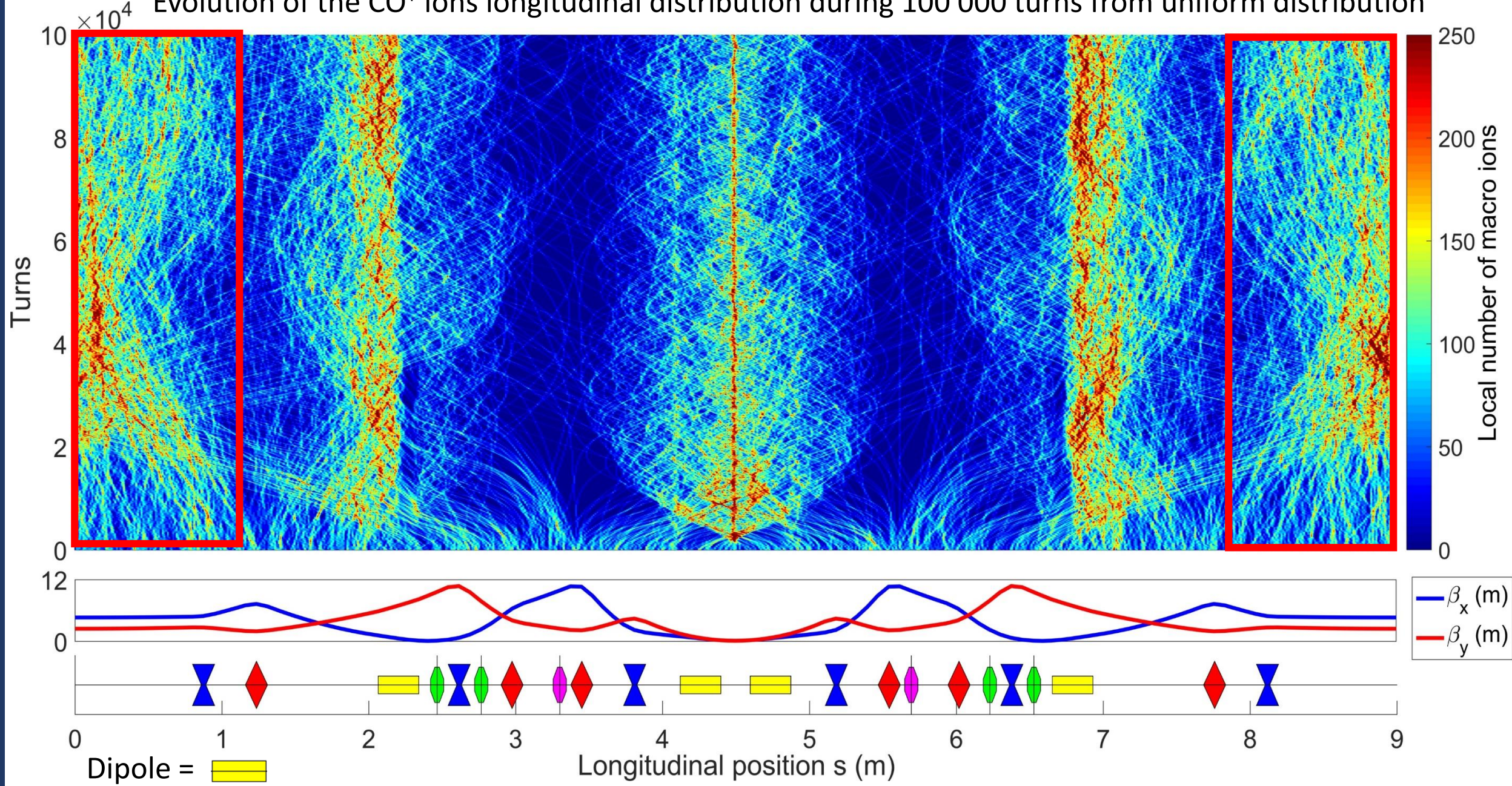


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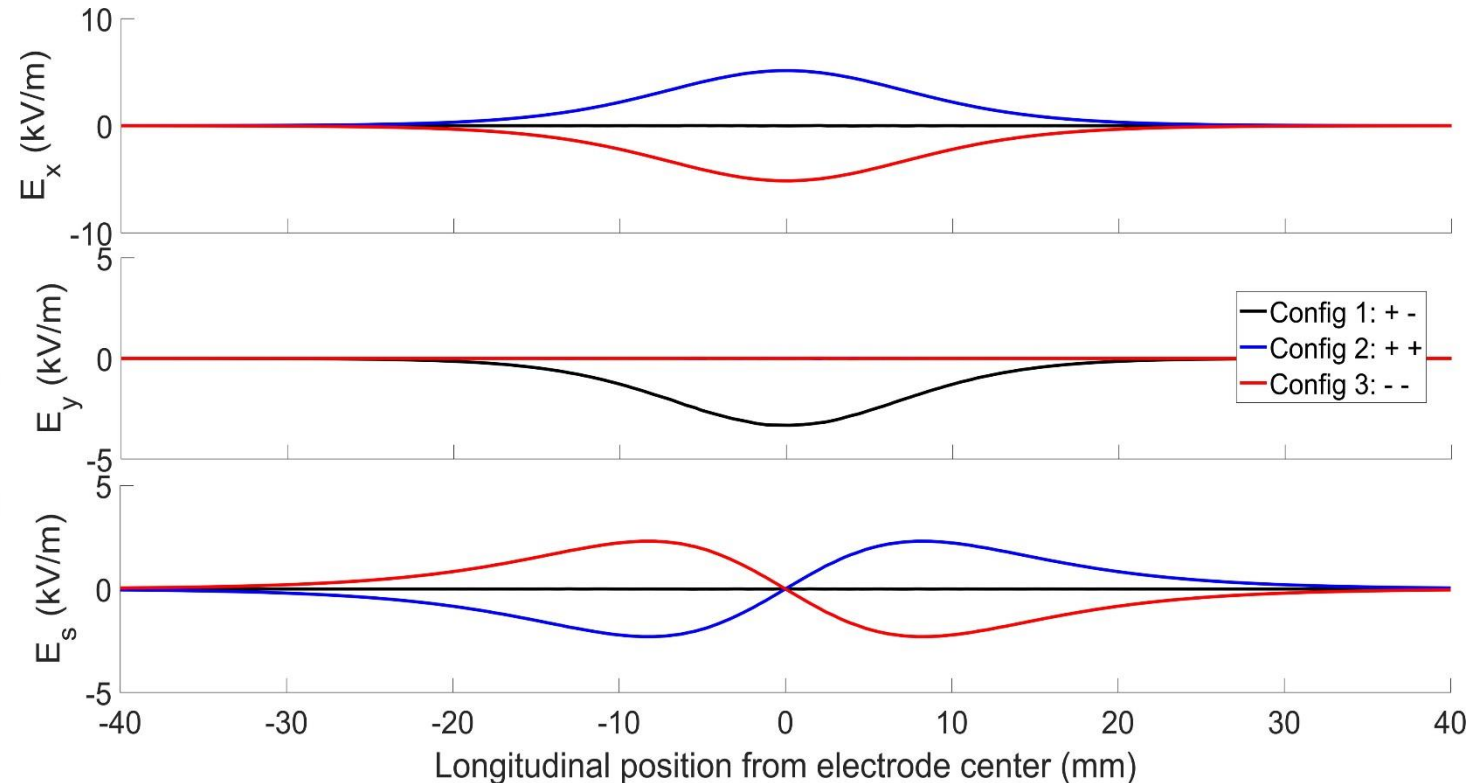
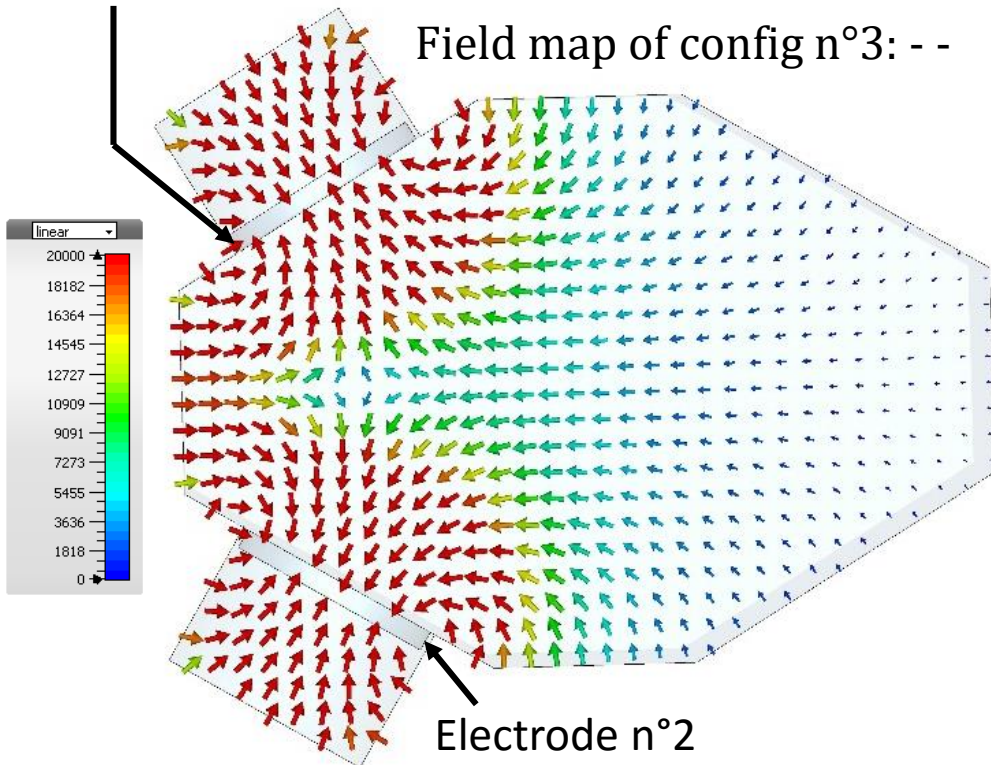




# Clearing electrode with longitudinal field

Config n°	Electrode polarity		Effect on axis	Colour
	1	2		
1	+	-	Transverse kick $E_y < 0$	Black
2	+	+	Transverse kick $E_x > 0$ and longitudinal expulsion point	Blue
3	-	-	Transverse kick $E_x < 0$ and longitudinal accumulation point	Red

Electrode n°1

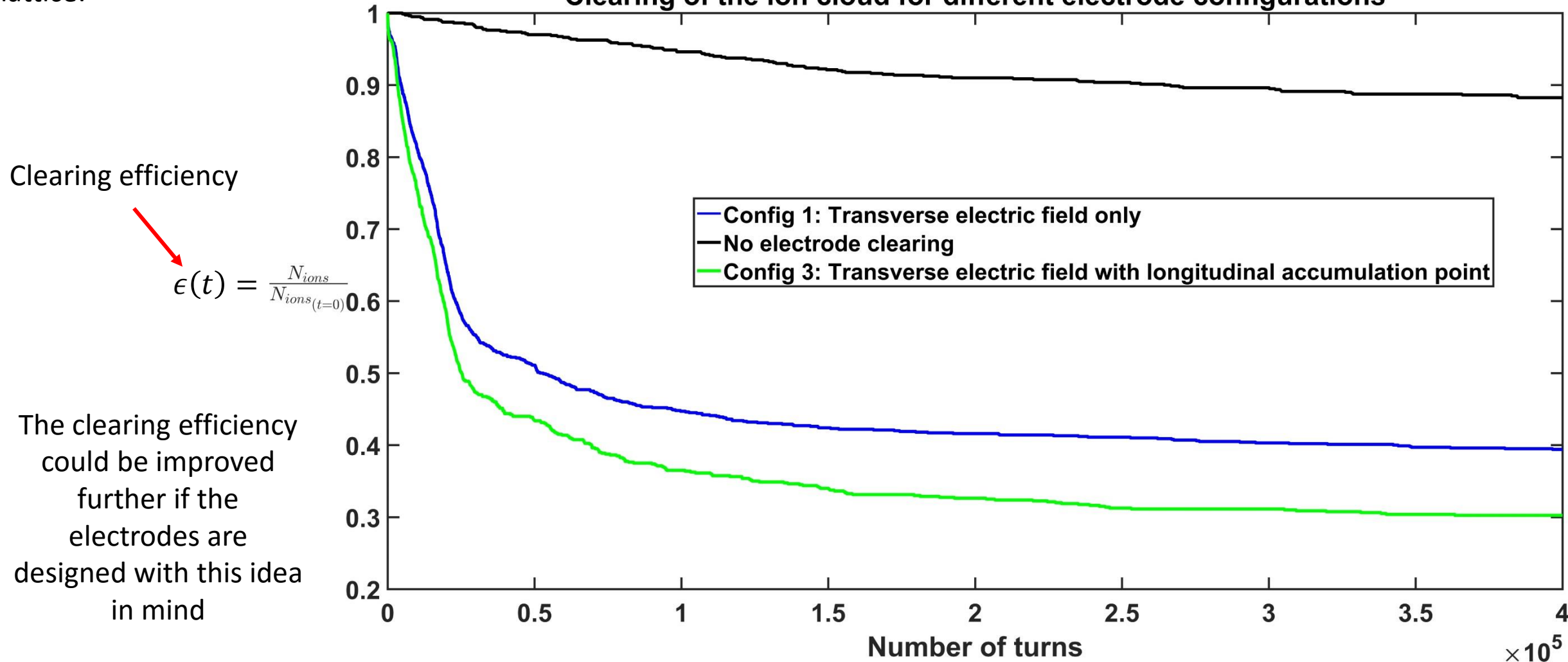




# Clearing electrode with longitudinal field

The ion cloud clearing can be optimised by carefully choosing the type of clearing field and its positioning in the lattice:

Clearing of the ion cloud for different electrode configurations

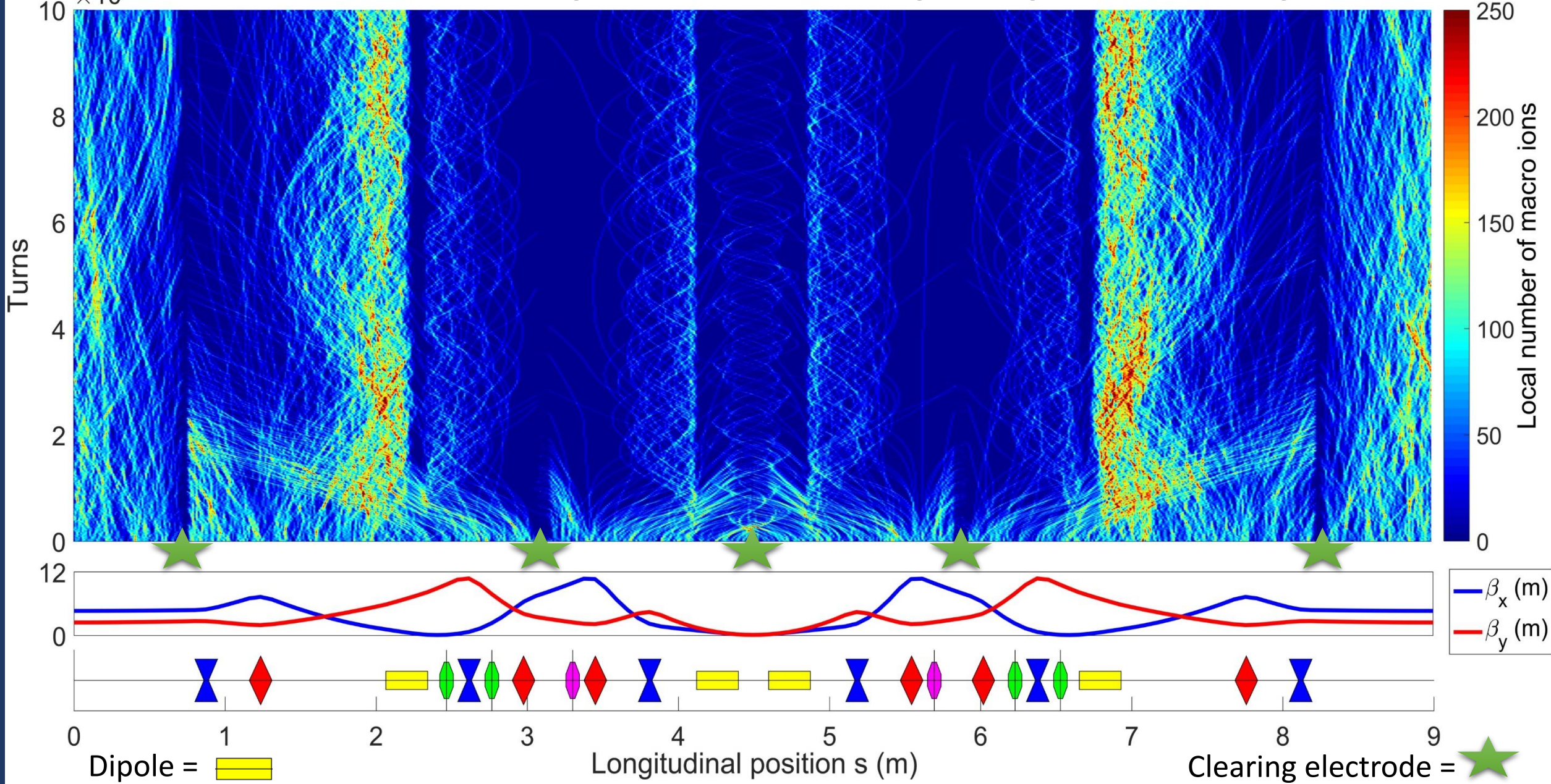


Clearing efficiency

$$\epsilon(t) = \frac{N_{ions}}{N_{ions}(t=0)}$$

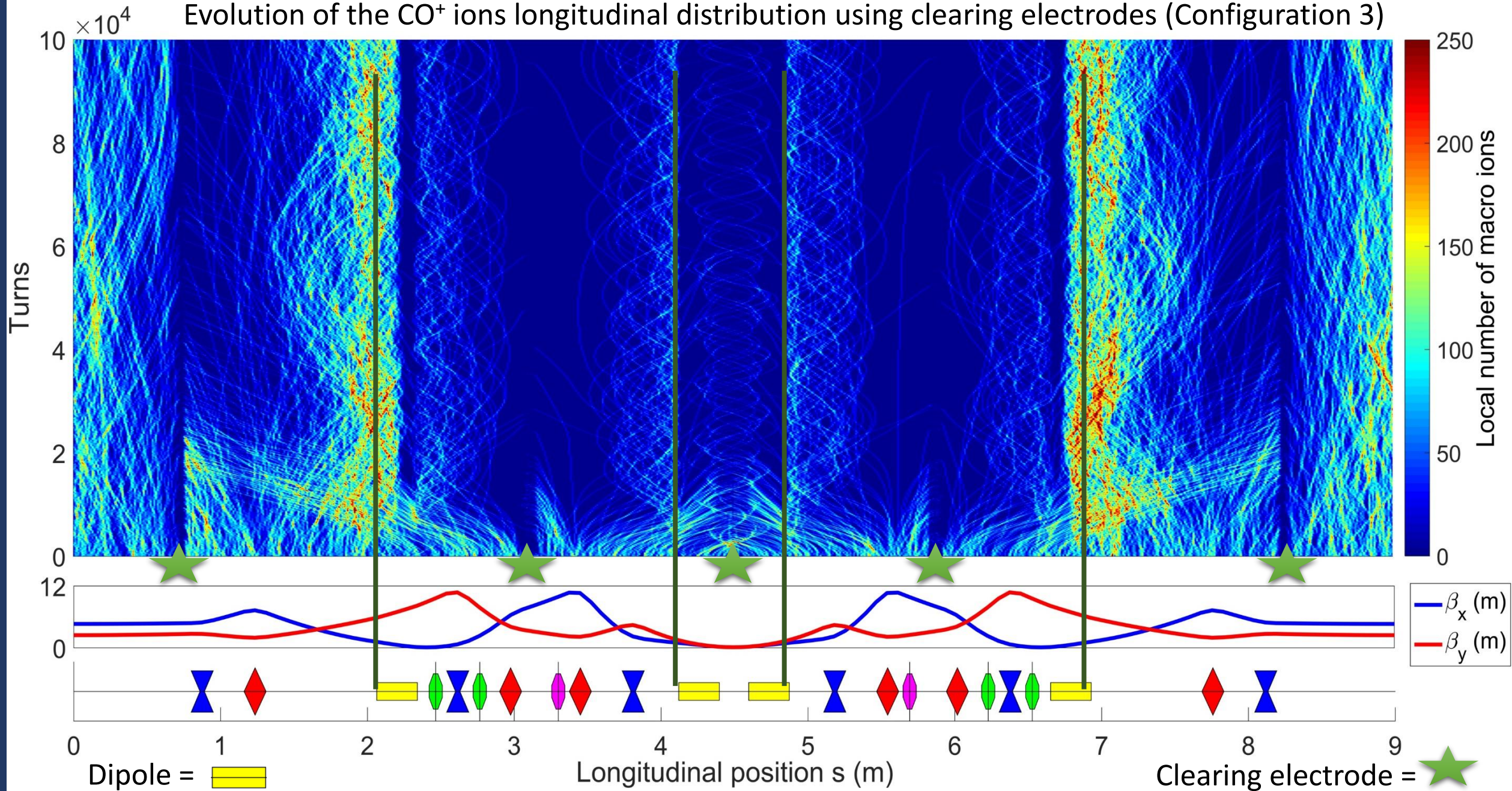
The clearing efficiency could be improved further if the electrodes are designed with this idea in mind

Evolution of the CO<sup>+</sup> ions longitudinal distribution using clearing electrodes (Configuration 3)



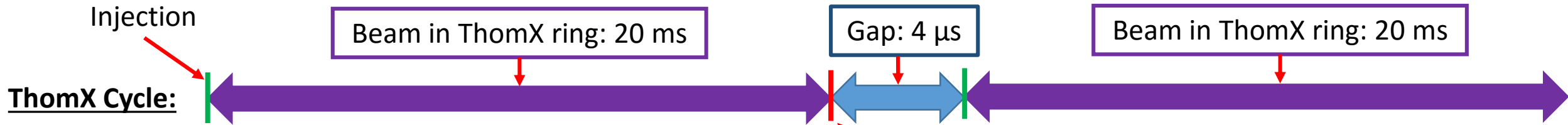


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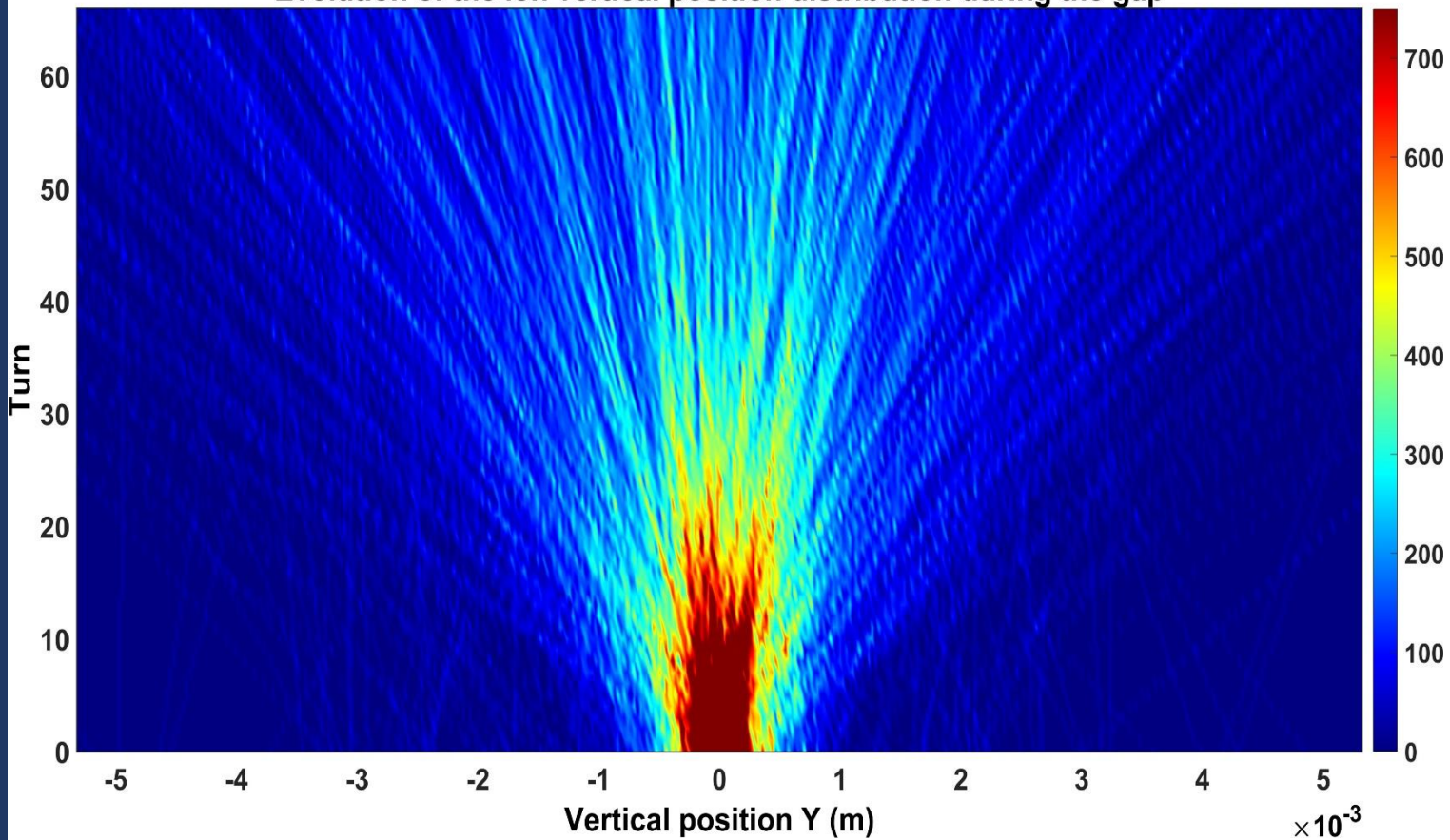




# Clearing gap, mixing ion positions

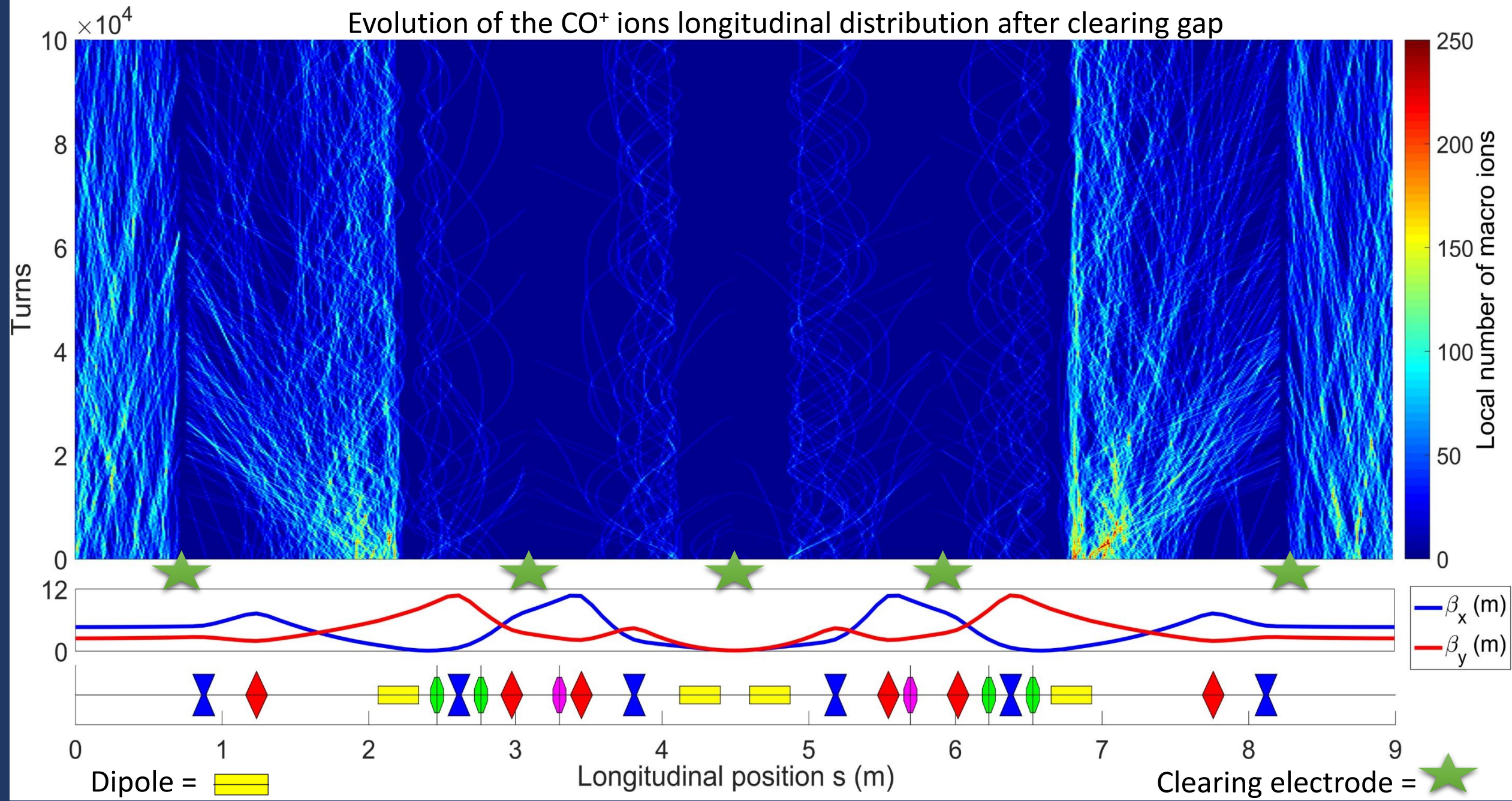


Evolution of the ion vertical position distribution during the gap



When there is no beam, ions are freed from beam potential well. The gap cleans a small part of the ions and « mix up » ion positions.





# Neutralisation factor

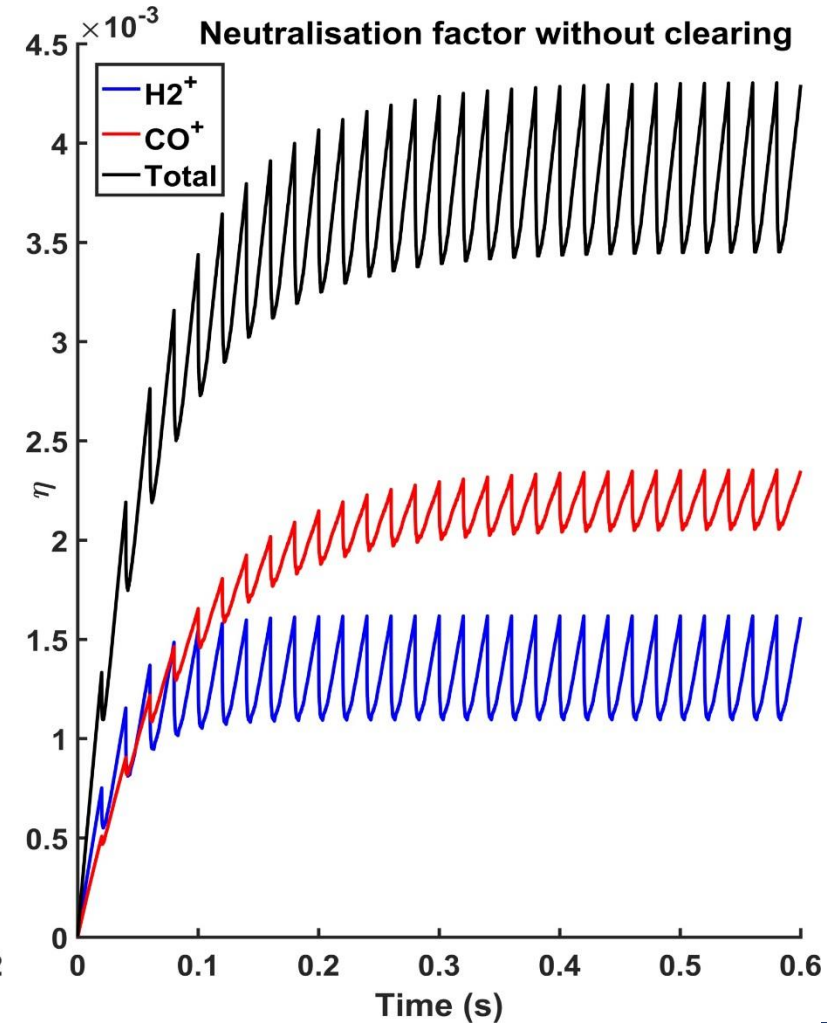
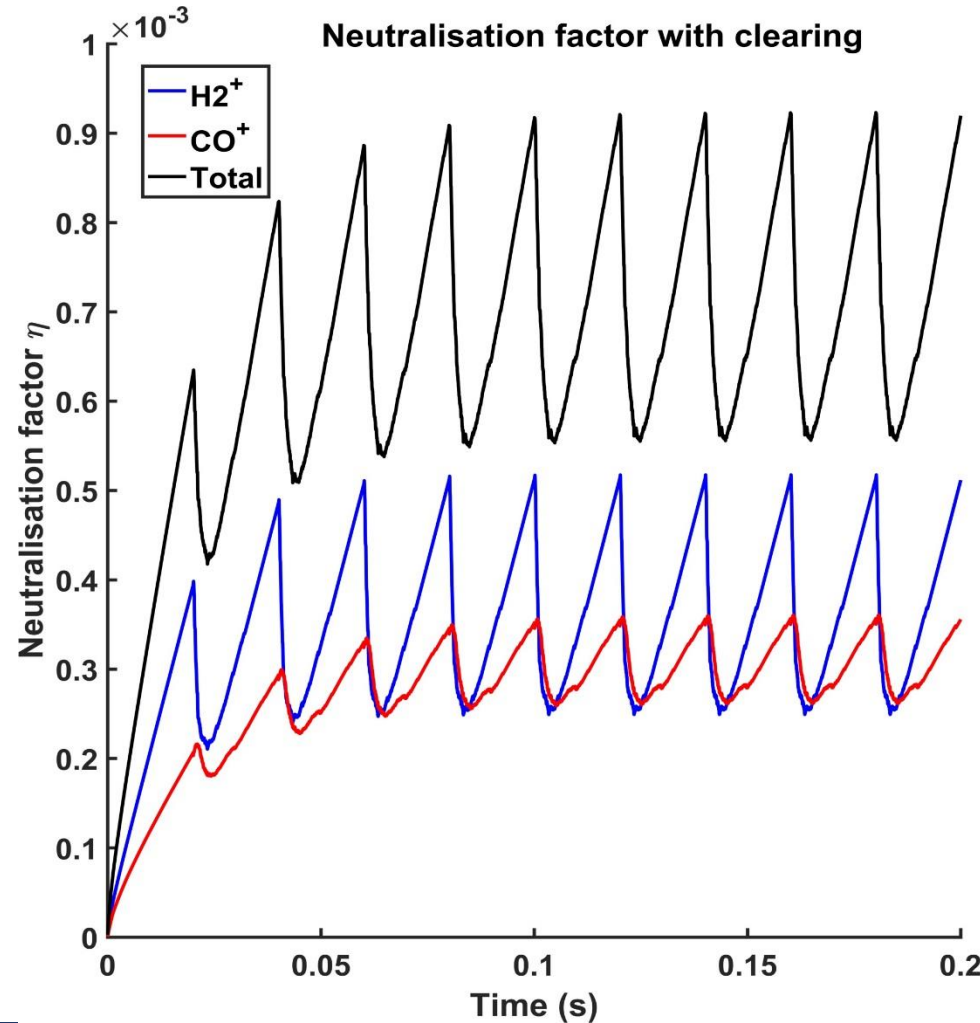
Using the clearing efficiencies  $\epsilon(t) = \frac{N_{ions}(t)}{N_{ions}(t=0)}$  computed by NUAGE and the known ionisation cross-section it is possible to get the neutralisation factor  $\eta = \frac{n_{ions}}{n_{electrons}}$  which is the relevant quantity for the beam dynamics:

Initial residual vacuum:

$$P_{tot} = 3 \cdot 10^{-10} \text{ mbar}$$

90 % H<sub>2</sub>  
10 % CO

$$\frac{\eta_{no\ clearing}}{\eta_{clearing}} \approx 4,5$$



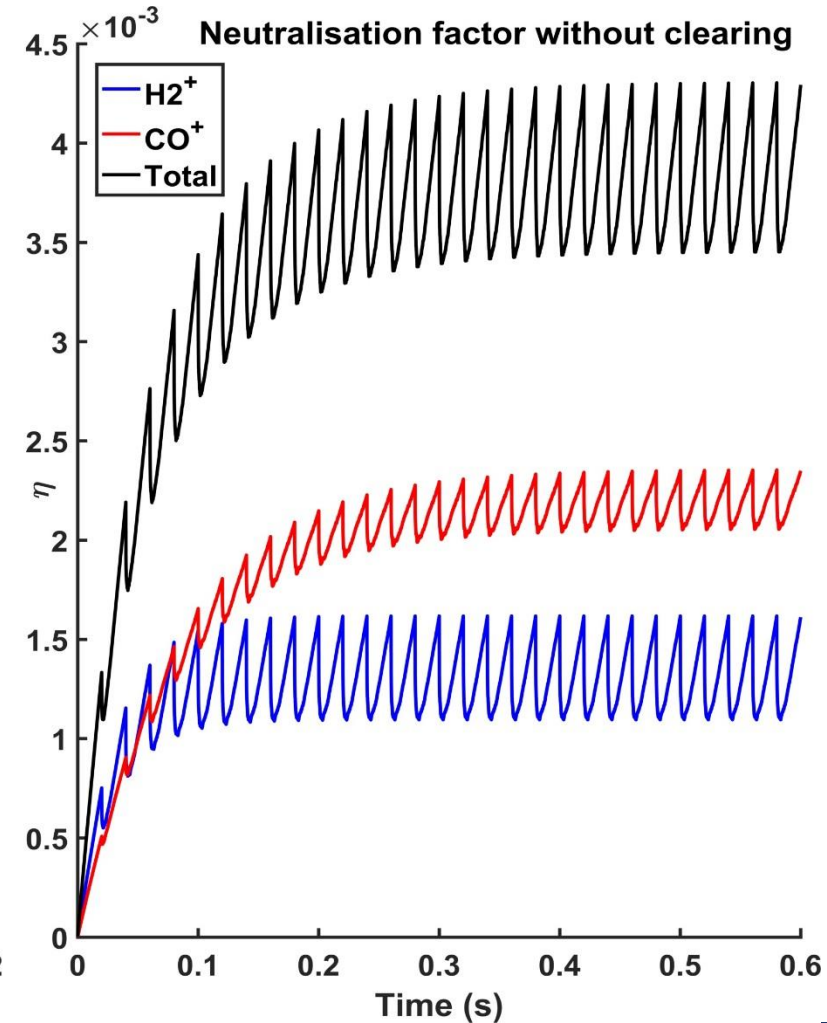
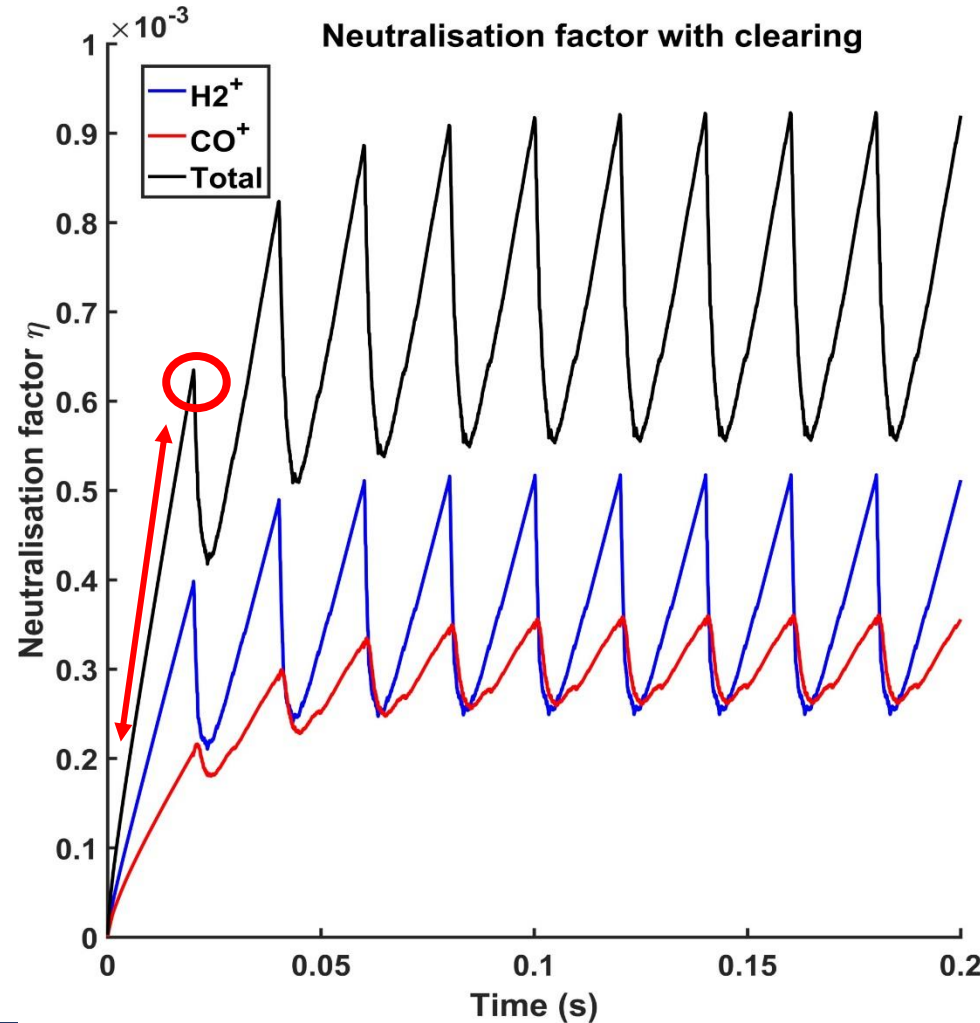


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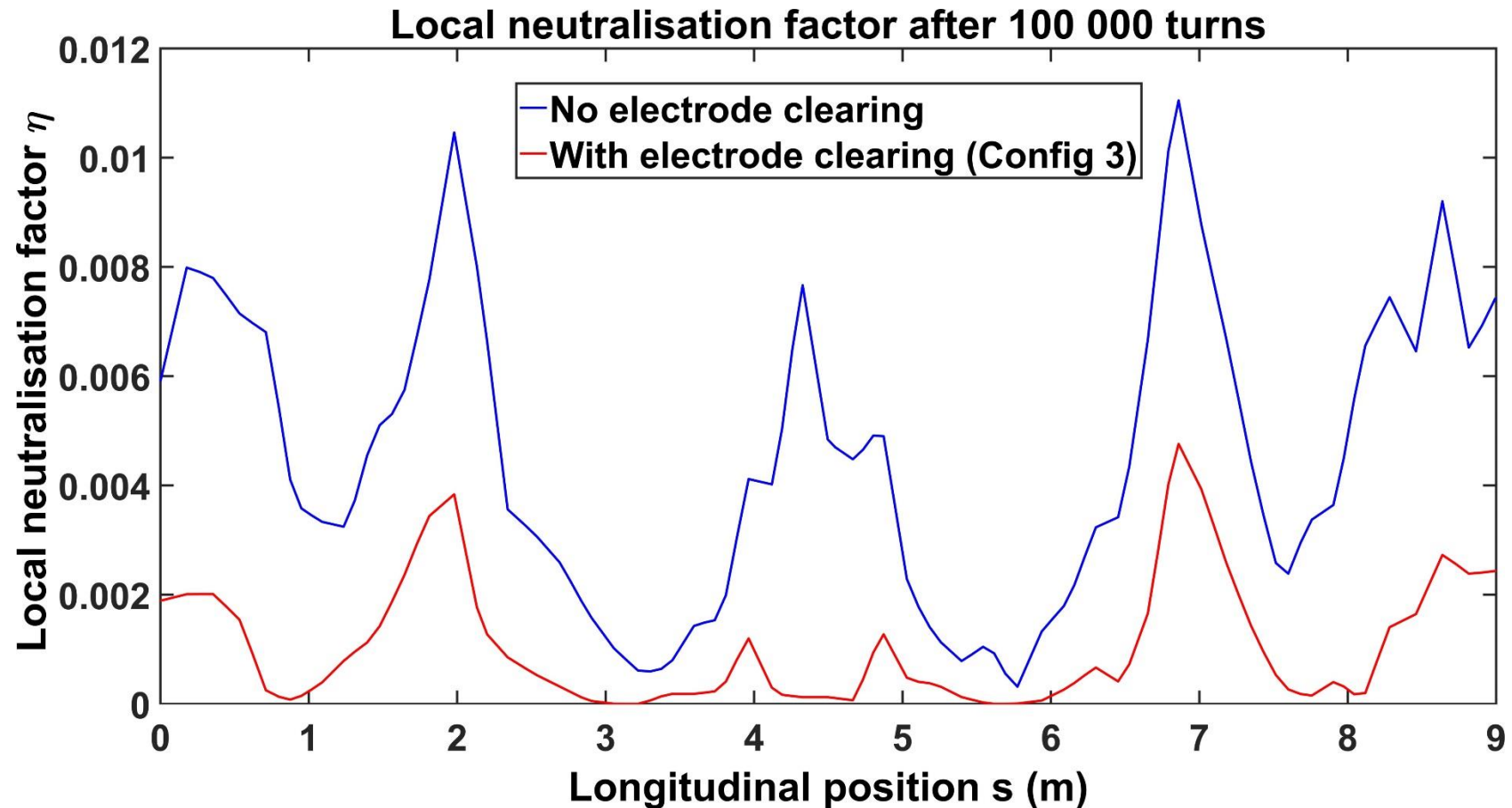
# Local neutralisation factor

It is possible to define a local neutralisation factor<sup>3</sup> which take into account the fact that the neutralisation is not homogeneous along the ring longitudinal position:

$$\eta(s) = \frac{L}{n_e} \frac{dn_i}{ds}$$

The mean value of the local neutralisation factor gives back the usual neutralisation factor:

$$\overline{\eta(s)} = \eta$$



# Induced tune shift

The tune shift induced by the ions can be computed by considering the ion force on the beam as an equivalent quadrupole. Assuming that the ion transverse distribution can be approximated<sup>4</sup> by a Gaussian distribution of  $\sigma_i = \frac{\sigma_e}{\sqrt{2}}$ , it gives:

$$\Delta Q_x = \frac{r_e n_e}{2\pi\gamma L} \int_0^L \frac{\beta_x \eta}{\sigma_{xe}(\sigma_{xe} + \sigma_{ye})} ds$$

Usually no information about the local neutralization factor is known so mean values are used:

$$\overline{\Delta Q_x} = \frac{r_e n_e}{2\pi\gamma} \frac{\overline{\beta_x \eta}}{\overline{\sigma_{xe}}(\overline{\sigma_{xe}} + \overline{\sigma_{ye}})}$$

Configuration	$\Delta Q_x$	$\overline{\Delta Q_x}$	Ratio	$\Delta Q_y$	$\overline{\Delta Q_y}$	Ratio
No clearing	9,94 E-4	1,52 E-4	6,5	1,14 E-3	5,84 E-4	2,0
Clearing	2,24 E-4	3,23 E-5	6,9	2,78 E-4	1,24 E-4	2,2

The same type of approach is possible for other effects induced by ions like emittance growth, tune spread, pressure increase, ...

# General conclusions from this study

- The longitudinal displacement of ion clouds is usually neglected but it can have an important impact. Understanding it helps to design effective mitigation strategies.
- Dipole magnets play an important role in the ion cloud dynamics.
- Possibility to use clearing electrode with longitudinal electric field for more effectiveness, to be demonstrated in ThomX.
- Use local values for the computation of the ion cloud induced effects in order not to minimize the effects, to be verified experimentally.

**Thank you !**

# Backups



# Horizontal kick felt by an $\text{CO}^+$ ion $\Delta v_x$ (m/s)

