

TECHNISCHE HOCHSCHULE NÜRNBERG
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Seit 1456

Kinetic Modeling of Electrostatic Ion thrusters

Workshop on Ion Propulsion and Accelerator Industrial Applications

Julia Duras^{*,+}, R.Schneider⁺, D.Kahnfeld⁺, P.Mathias⁺, G.Bandelow⁺,
K.Lüskow⁺, K.Mayash⁺, S.Kemnitz^{§,+}, N.Koch^{*}

^{*}Nuremberg Institute of Technology

⁺Institute of Physics, University of Greifswald

[§]Institute of Computational Science, University of Rostock

THALES

Bayerisches Staatsministerium für
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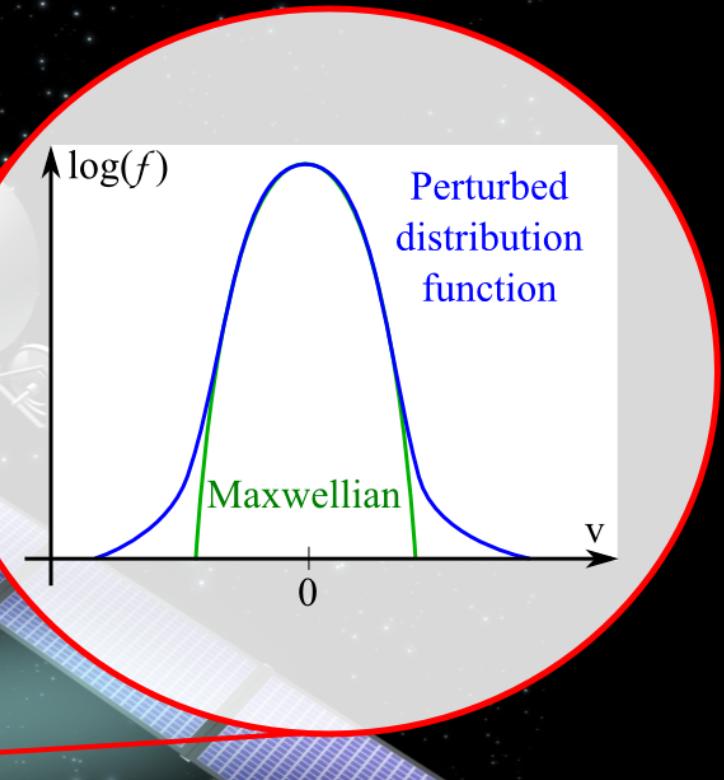
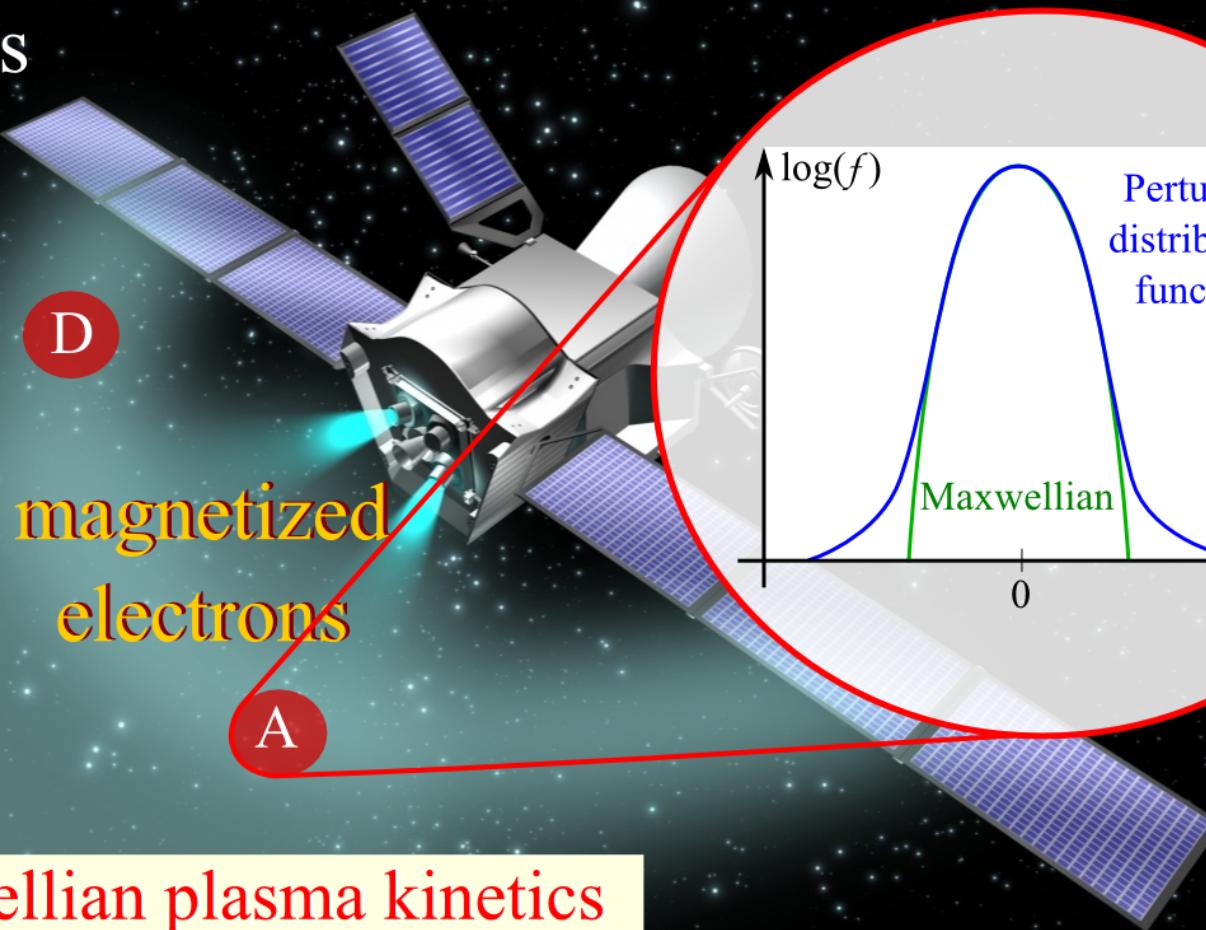


Ion thruster plume interaction with satellites

solar panels

plume
expansion

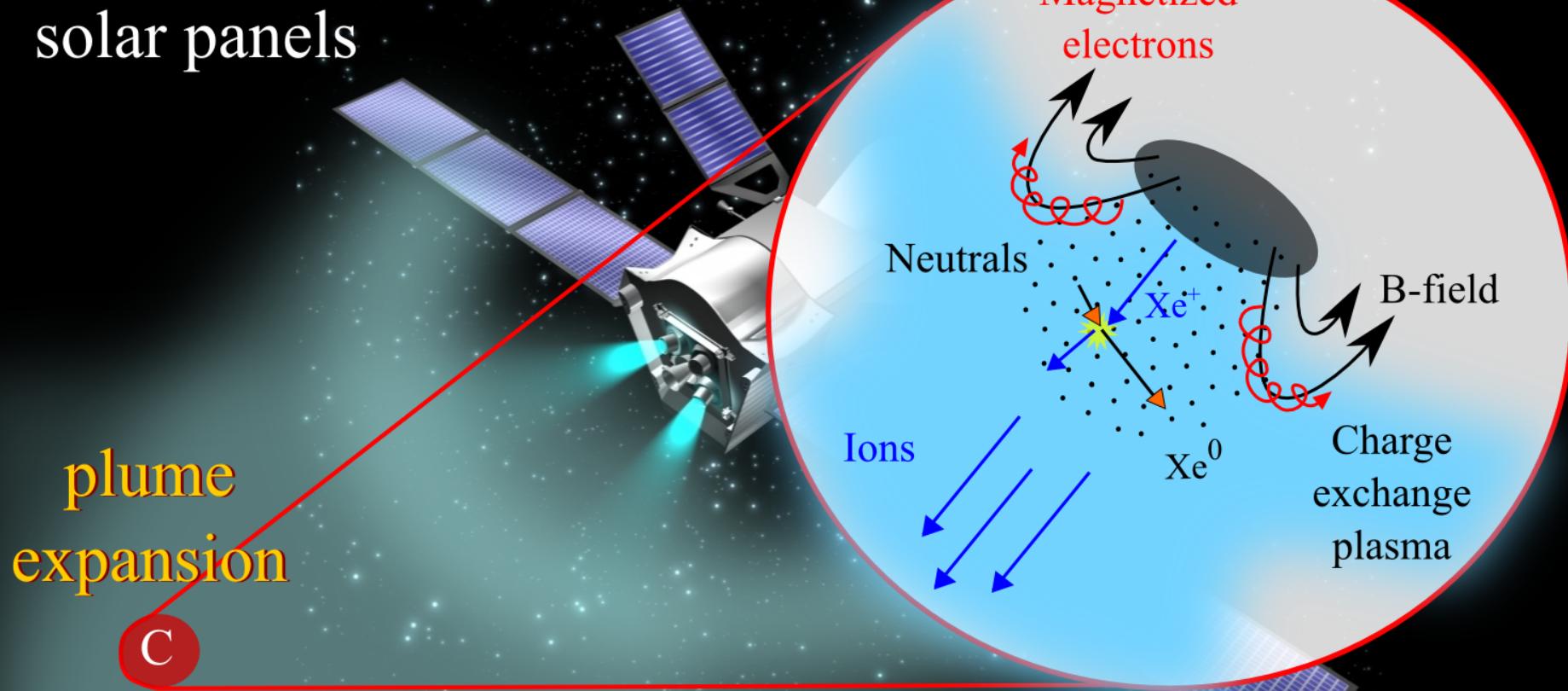
C



- (A) Non-Maxwellian plasma kinetics
- (B) Plasma-wall interaction
- (C) Non-magnetized expansion dynamics
- (D) Plasma & electromagnetic waves

Ion thruster plume interaction with satellites

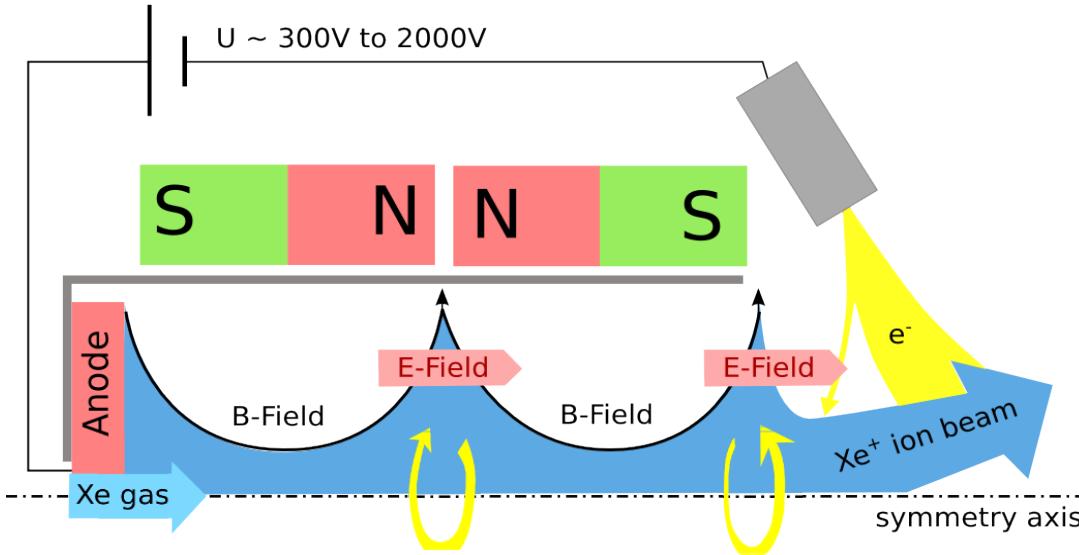
solar panels



- A Non-Maxwellian plasma kinetics
- B Plasma-wall interaction
- C Non-magnetized expansion dynamics
- D Plasma & electromagnetic waves

The HEMP thruster

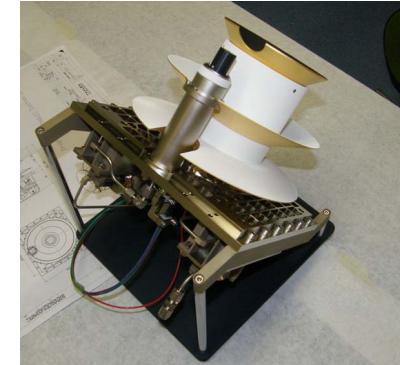
patented by THALES Electron Devices
with an initial patent filed in 1998



Self-consistent kinetic simulation: $\lambda_e^{mfp} \approx L_{system}$

Self consistent plasma + neutral dynamic
Anomalous transport
Secondary electron emission
Integrated model of channel and plume

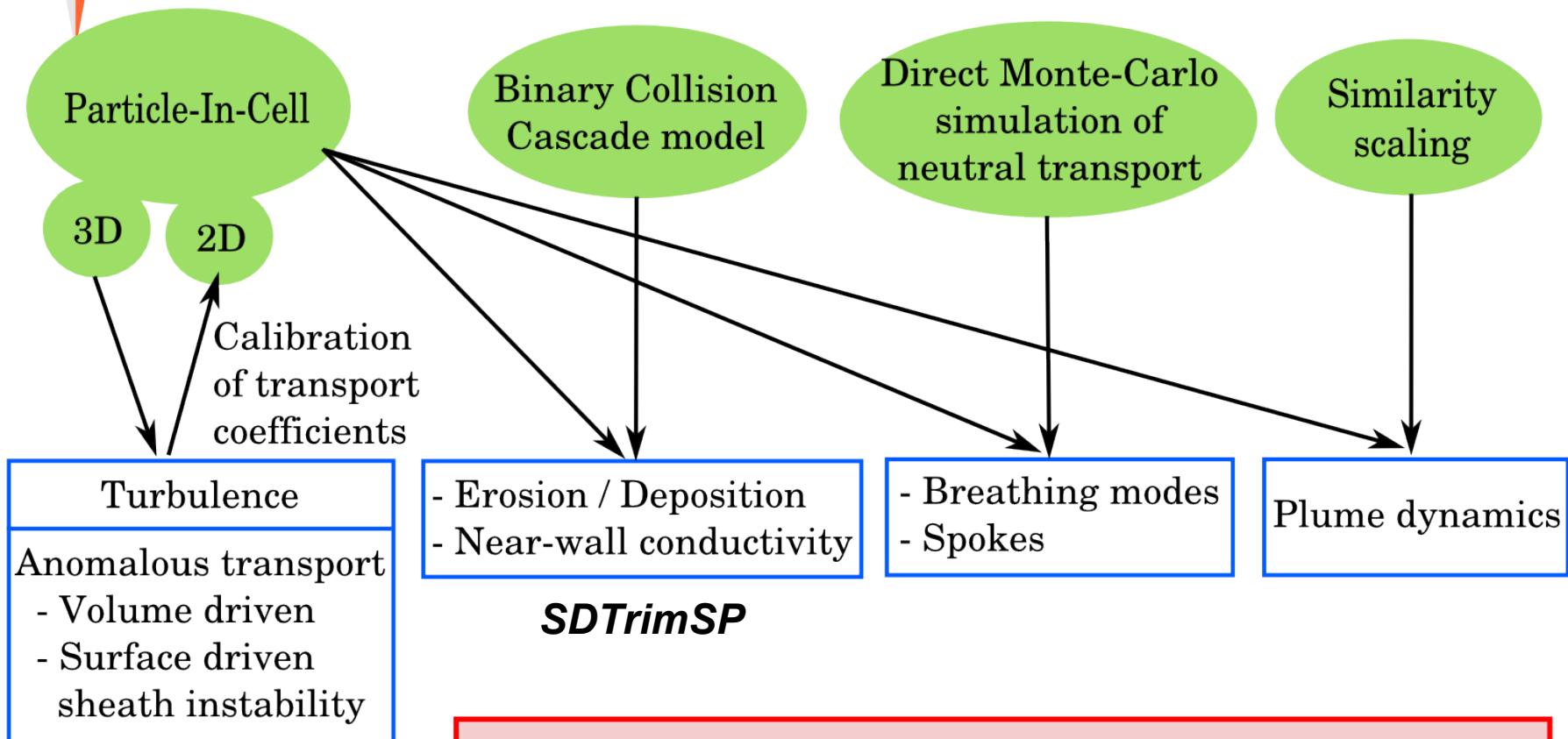
DLR Projects:
50 RS 0804, 50 RS1101



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**System without grid
and
reduced wall contact**

Simulation tools

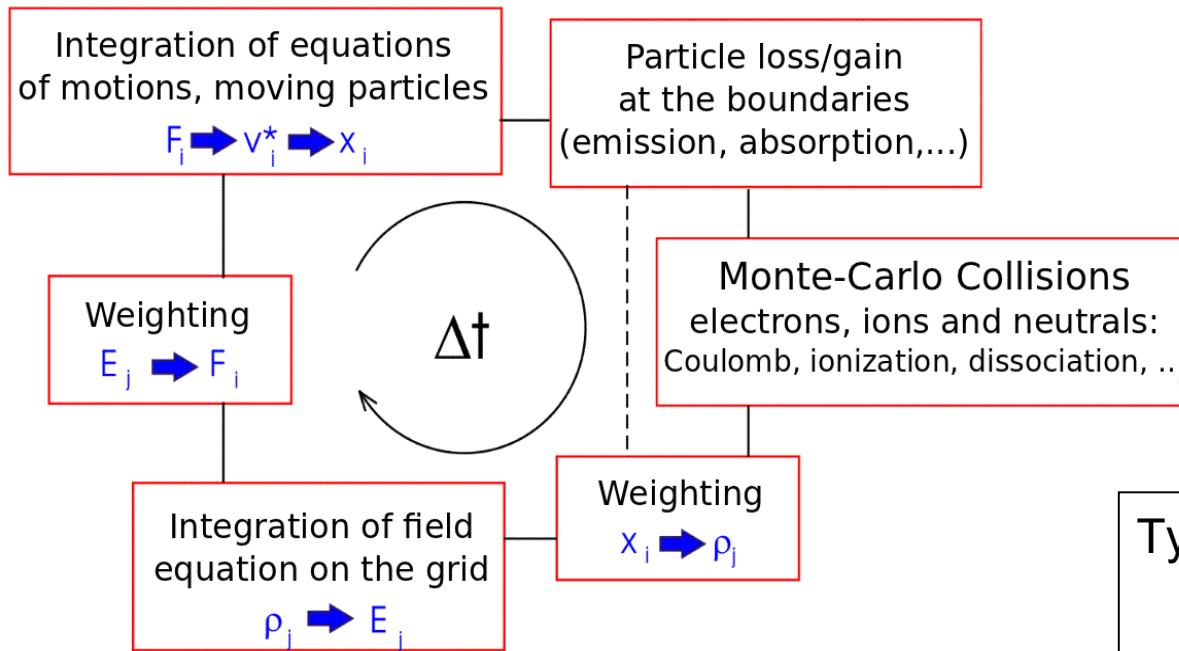


Own codes!

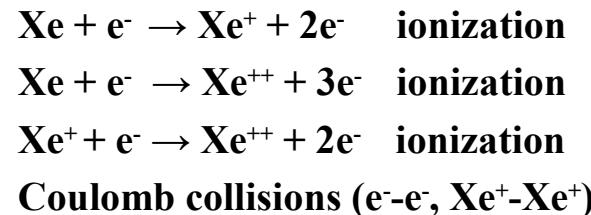
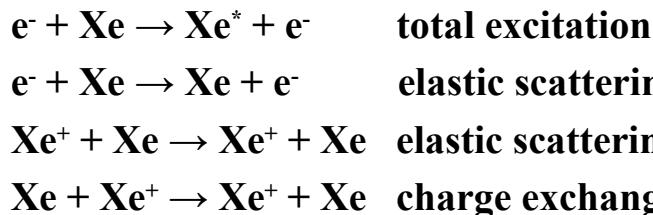
Particle-In-Cell: [1], [2]
 SDTrimSP: [3], [4], [5]
 Similarity scaling: [6]

- Self consistent coupling of all methods
- Reduction of empirical parameters
- Deduction from higher hierarchical models

MCC Particle-In-Cell simulation



Direct Monte-Carlo Collisions:



Requirements:

$$\Delta t = 0.2/\omega_{p,e} \quad \Delta x = 0.5\lambda_{D,e}$$

equidistant mesh (self force!)

Similarity scaling:

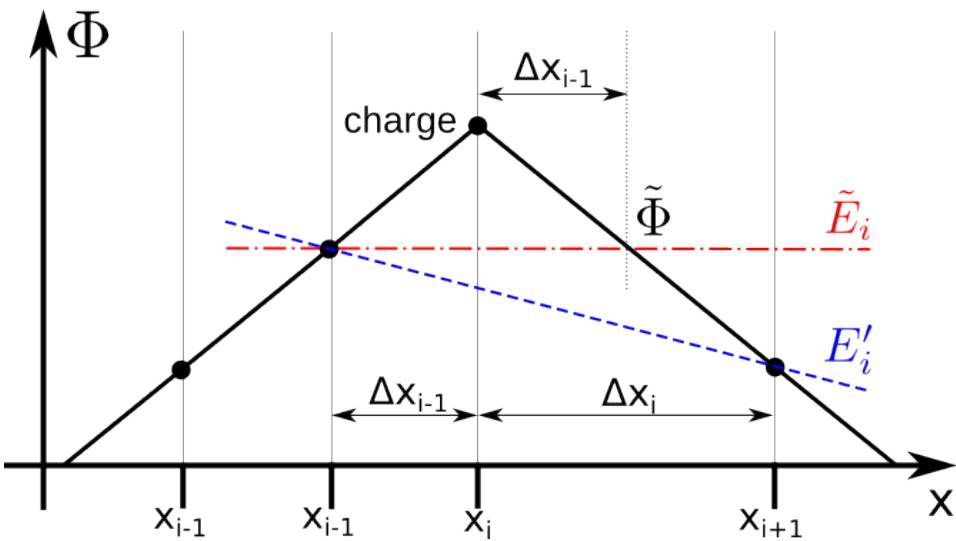
$$\frac{r_{L,e}}{L} = const. \quad \frac{\lambda_{e,N}^{mfp}}{L} = const.$$

surface/volume limit

Typical parameters for HEMP:
 10^5 cells, 10^8 time steps,
 10^8 pseudo-particles
→ computing time ≈ 2 weeks

Non-equidistant PIC

Test: 1D PIC, 1 electron



different E-field interpolation

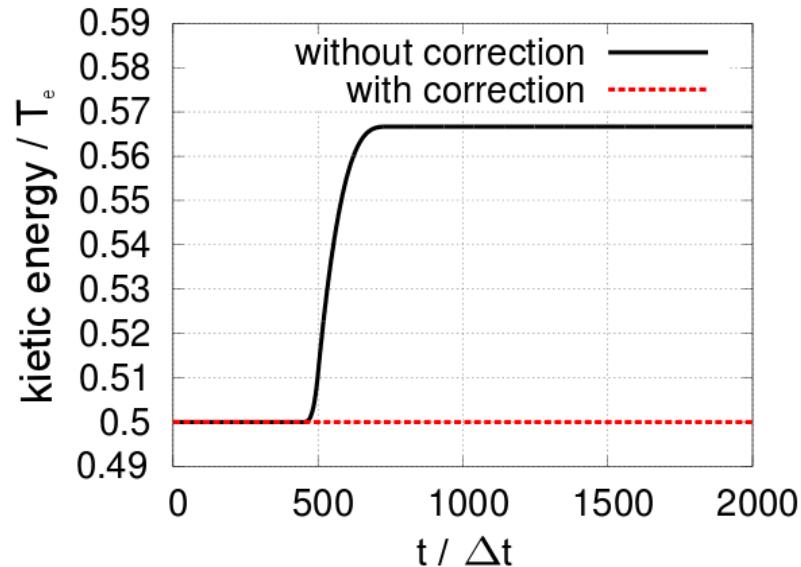


self force disappears

**non-symmetric situation
for non-equidistant grids**



artificial self force

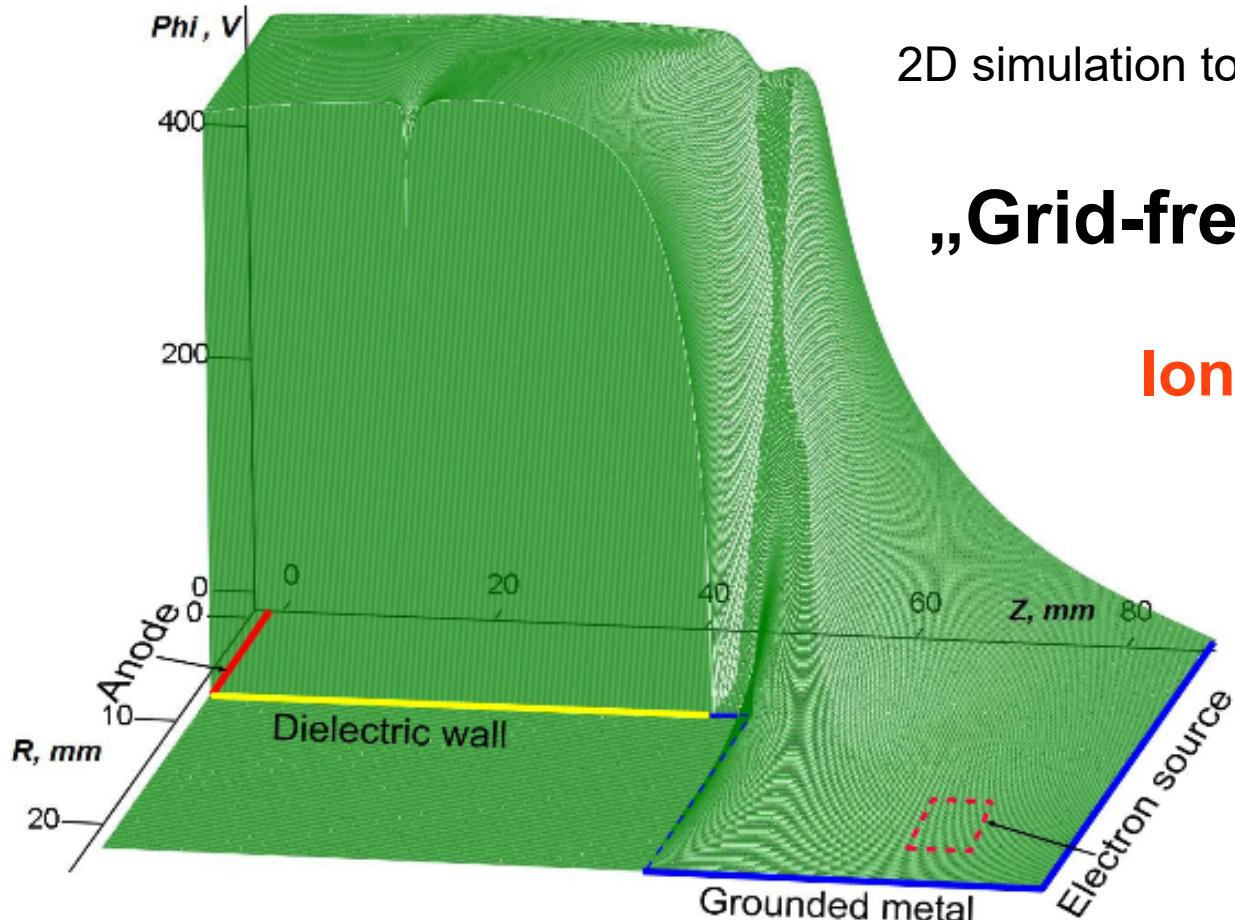


[7] Duras, Matyash; CPP (2014)

Simulation results

HEMP DM3a thruster

Potential profile:



2D simulation to reduce computational effort

„Grid-free grid thruster”

Ions accelerated
at the exit

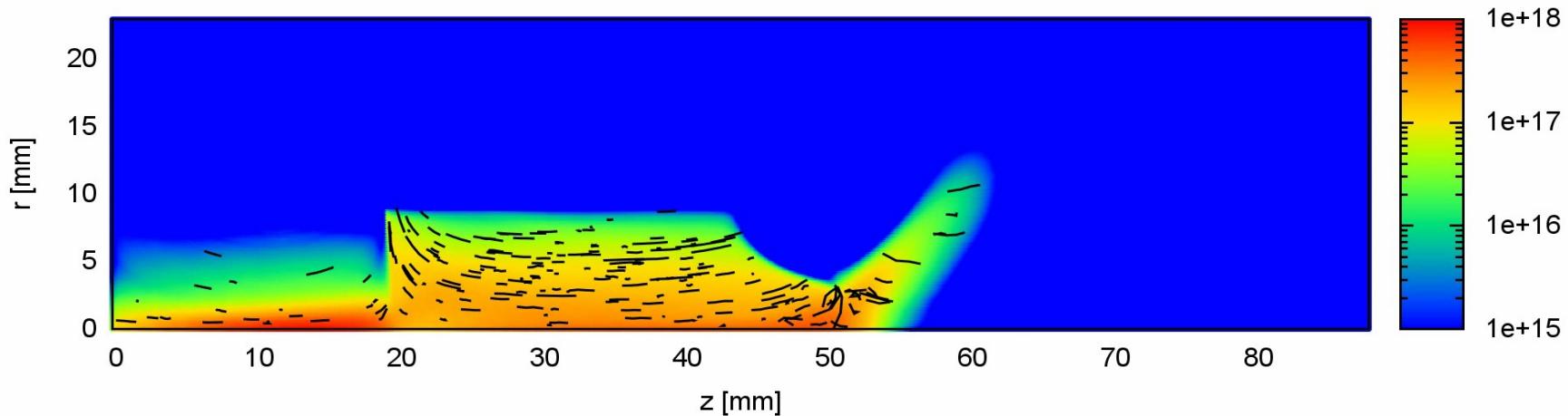
$$R_{\text{thruster}} = 9 \text{ mm}$$
$$L_{\text{thruster}} = 51.2 \text{ mm}$$

[8] Matyash, Schneider; IEEE (2010).
[9] Kalentev, Matyash; CPP (2014).

Simulation results

HEMP DM3a thruster

Electron density with example particles:

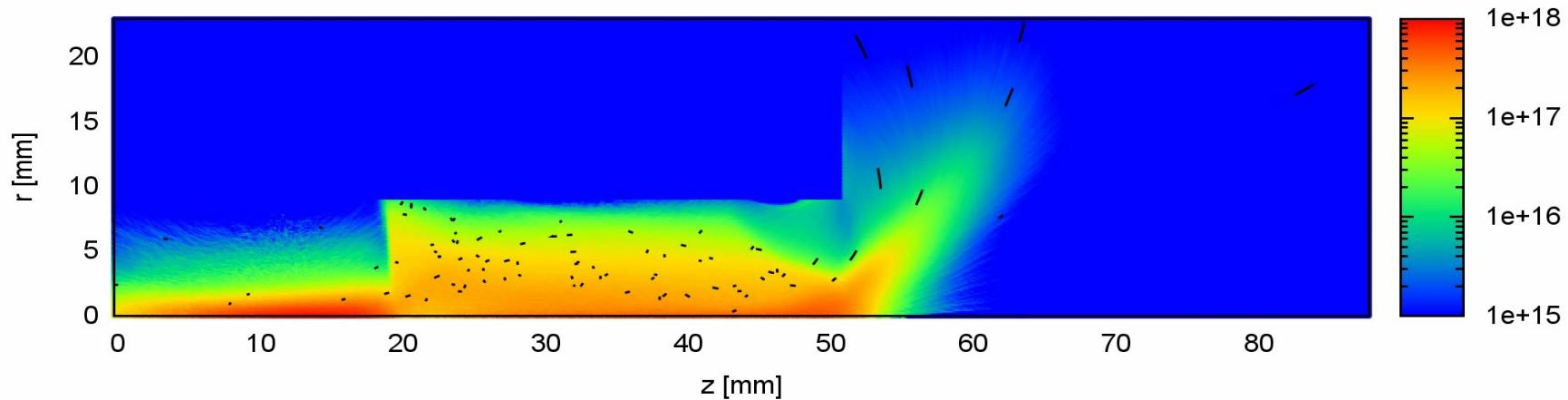


- Magnetised electrons $50 \cdot \Delta t_e$
- Non-Maxwellian velocity distribution $\lambda_e^{mfp} \approx L_{\text{system}}$
- Oscillating in front of the exit

Simulation results

HEMP DM3a thruster

Ion density with example particles:



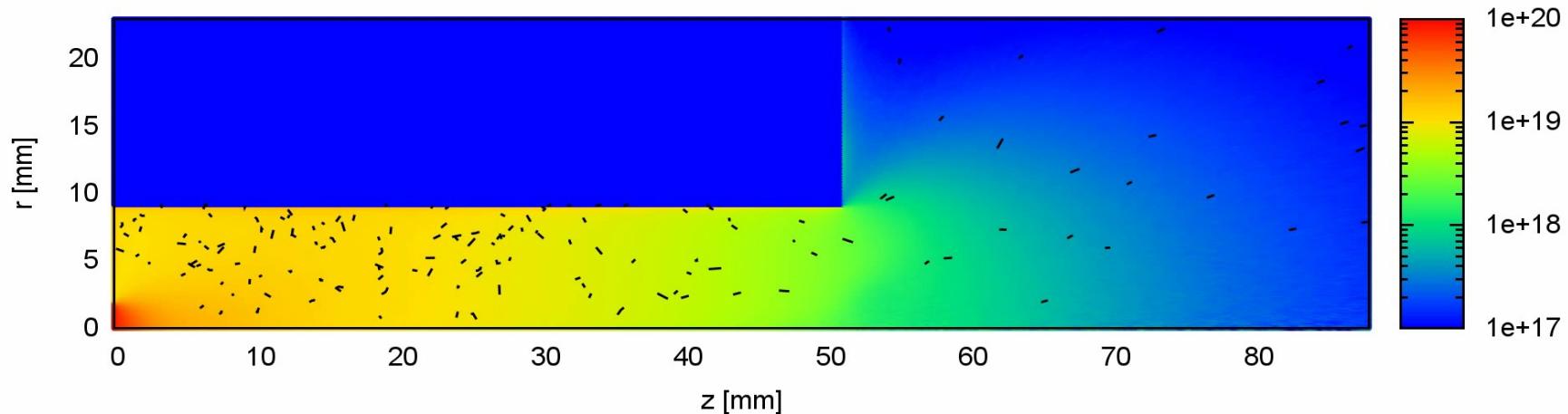
- Not magnetized
- Thermal velocity distribution in the channel
- Acceleration at the thruster exit

$$2.000 \cdot \Delta t_e$$

Simulation results

HEMP DM3a thruster

Neutral density with example particles:



$$10.000 \cdot \Delta t_e$$

- Thermal velocity in the channel
- Expansion at the channel exit

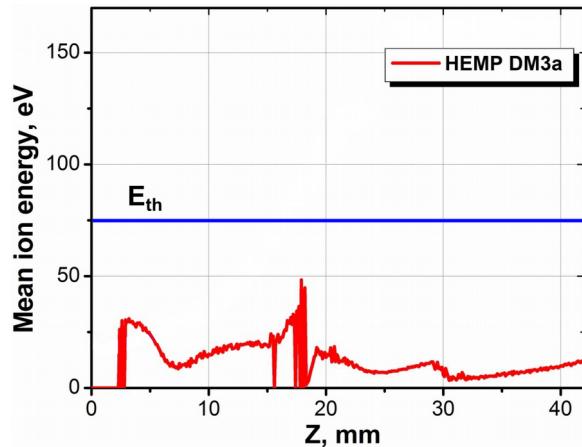
Diagnostics

HEMP DM3a thruster

Validation with experimental data

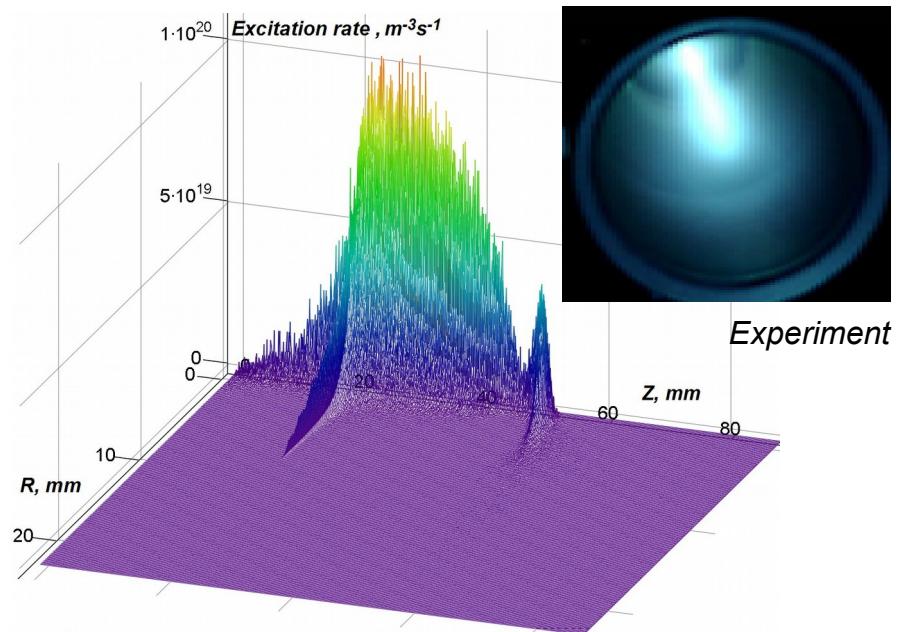
Erosion:

- Ion flux non-negligible in HEMP only at cusp positions
- Energy below the sputtering threshold ($E_{th} \sim 75$ eV)



„Grid-free grid thruster
with minimized erosion”

Total excitation:

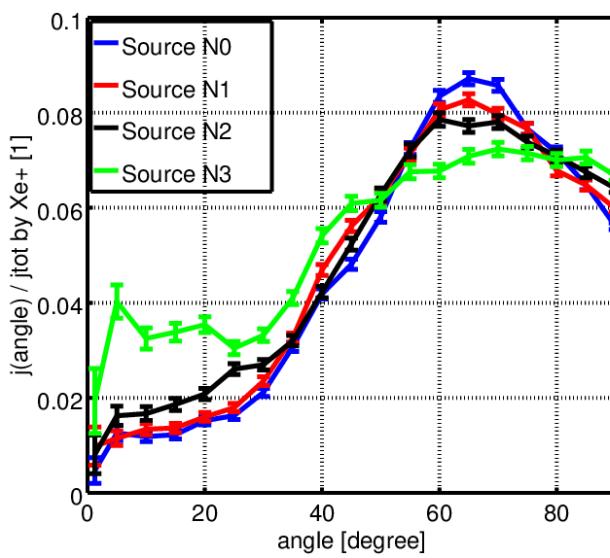
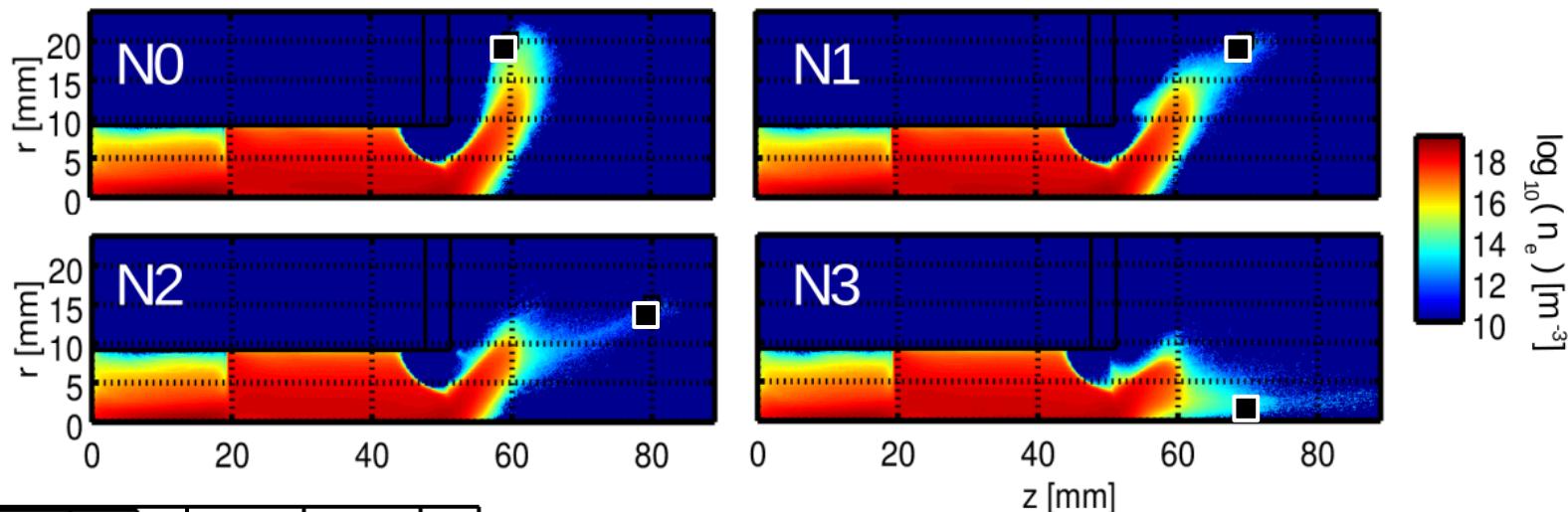


Light emission from axis
and cusps

[8] Matyash, Schneider; IEEE (2010).

Influence of electron source

Thermal source with $T_e = 2\text{eV}$ and $I_{src} = 0.3\text{mA}$



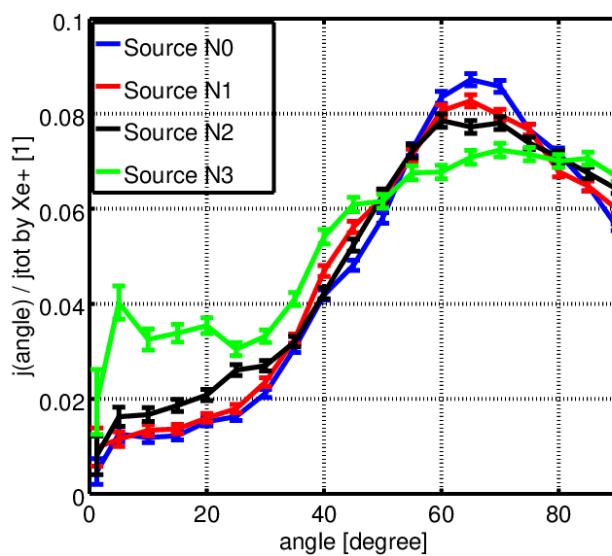
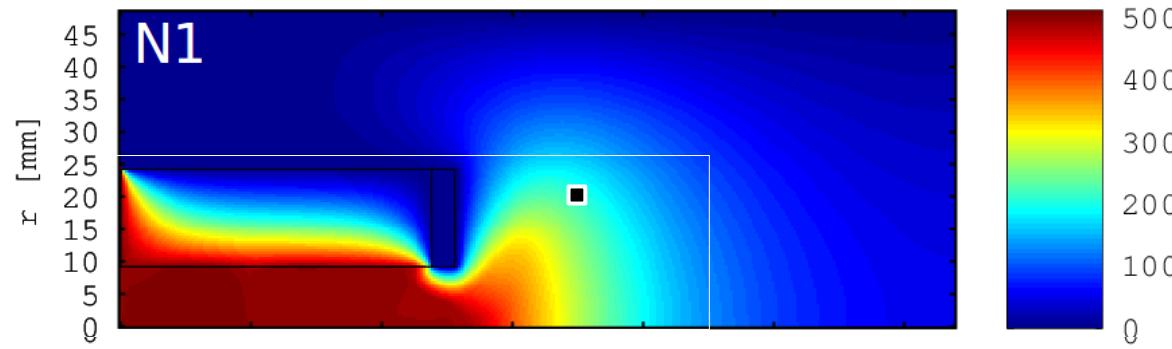
Due to B-field, only axial source changes electron density → angular ion distribution

Axial electrons enter channel directly
Non axial electrons = reservoir

[10] Duras, Schneider; PPT (2016).

Influence of electron source

Thermal source with $T_e = 2\text{eV}$ and $I_{src} = 0.3\text{mA}$



No plume plasma

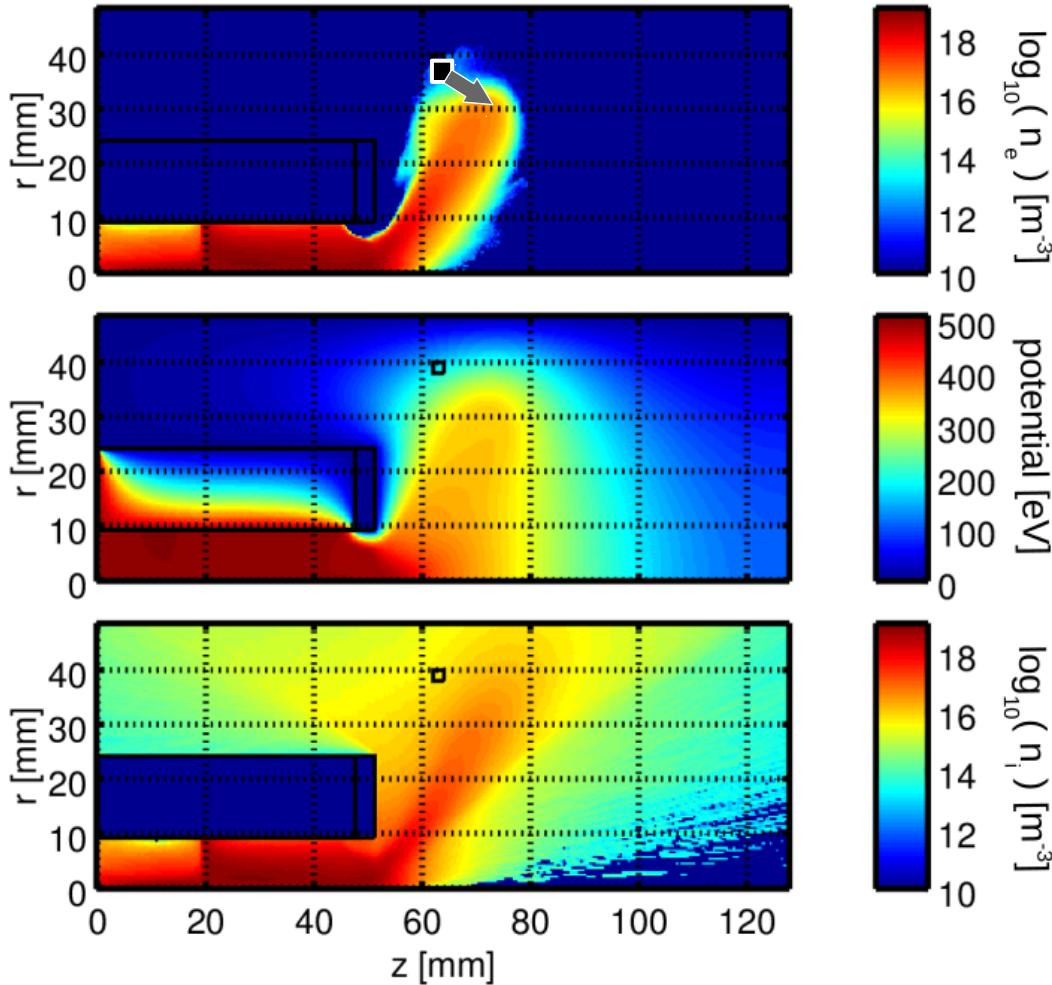


No screening of domain boundary potential
→ influences angular ion distribution

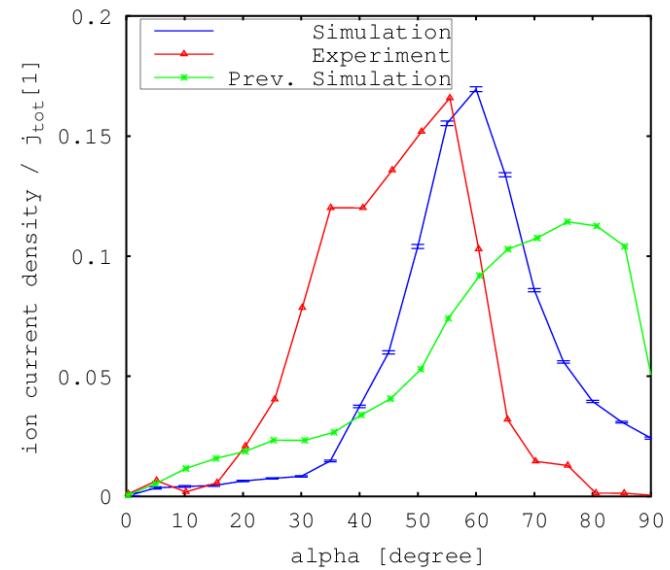
[10] Duras, Schneider; PPT (2016).

Influence of electron source

Directed source with drift velocity $v_{\text{drift,e}} = 20 \text{ eV}$, $T_e = 0.1 \text{ eV}$ and $I_{\text{src}} = 0.3 \text{ mA}$



**Plasma screens
top boundary potential
= more realistic**



[10] Duras, Schneider; PPT (2016).

Summary

- **Self-consistent kinetic simulation are needed**
→ computationally costly → deduction from higher hierarchies
- **Correction of E-field calculation for non-equidistant grids**
- **Validation with experiment: wall erosion & total emission**
- **Studies of electron sources:**
 - Ion beam divergence is mainly determined by B-field
 - Screening of domain potential is important for simulation

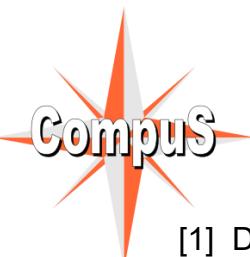
Acknowledgment:

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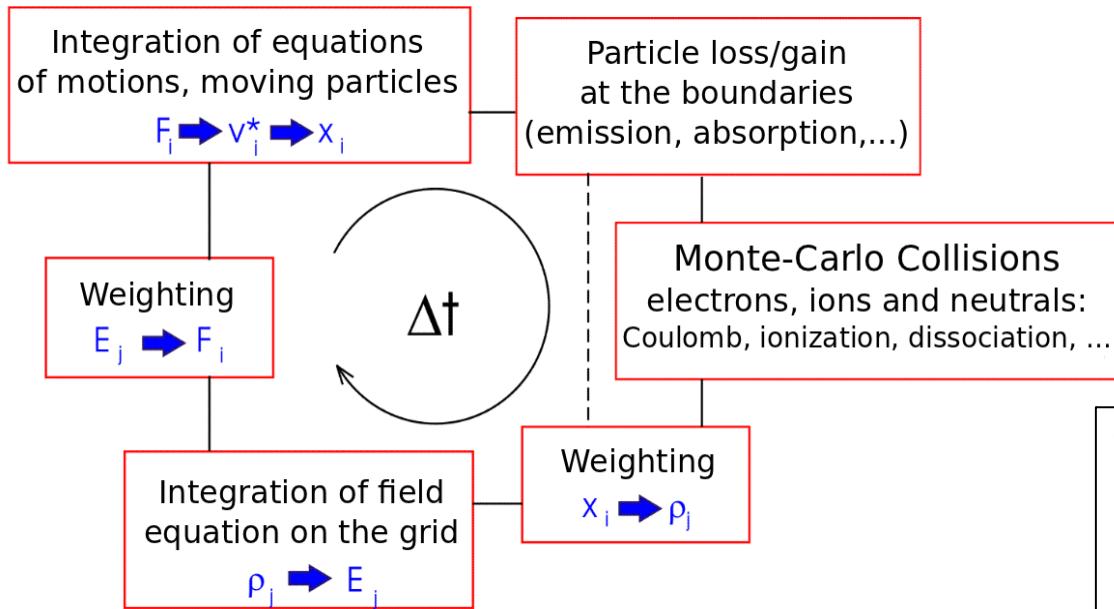




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Particle-In-Cell simulation of ion thrusters



Requirements:

$\Delta t = 0.2/\omega_{p,e}$ $\Delta x = 0.5\lambda_{D,e}$
 equidistant mesh for momentum conservation

Typical parameters for HEMP:
 10^5 cells, 10^8 time steps,
 10^8 pseudo-particles
 → computing time ≈ 2 weeks

Length	$L = f L^*$
Magnetic field	$B = f^1 B^*$
Cross section	$\sigma = f^1 \sigma^*$
Plasma density	$n_p = n_p^*$
Neutral density	$n_n = n_n^*$
Electric potential	$\Phi = \Phi^*$
Temperature	$T = T^*$

Similarity scaling:

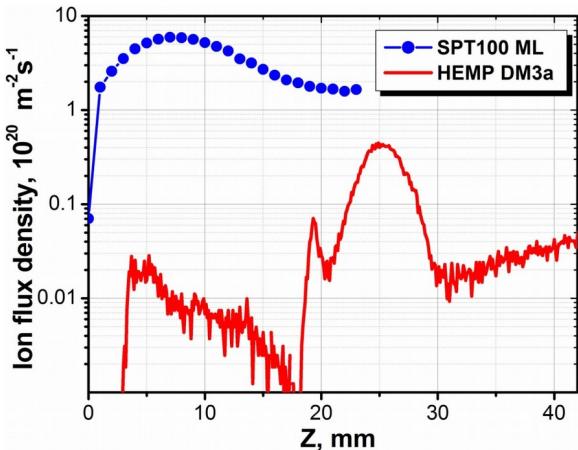
$$\frac{r_{L,e}}{L} = \text{const.} \quad \frac{\lambda_{e,N}^{mfp}}{L} = \text{const.} \quad \rightarrow \text{surface/volume limit}$$

Energy not scaled → atomic physics preserved
Debye length unchanged → reduction of grid size

Channel erosion

HEMP DM3a thruster

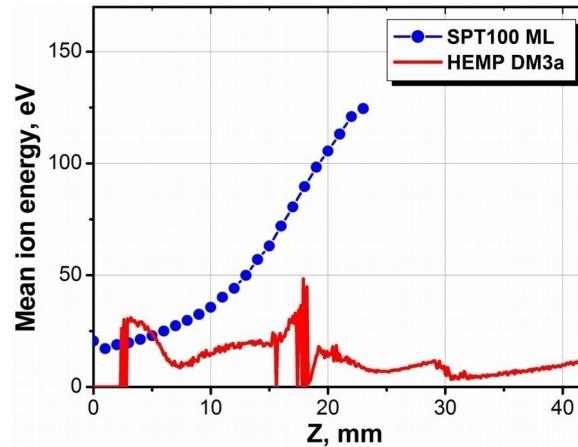
Ion fluxes to the channel wall



Ion flux non-negligible in HEMP only at cusp position Energy below the sputtering threshold ($E_{th} \sim 75$ eV)

[7] Matyash, Schneider; IEEE (2010).

Mean energy of the ions



Erosion rate

