Measurement of Vector and Tensor Asymmetries in Quasielastic (e,e'p) Electron Scattering from Deuterium

- Deuteron structure theoretically calculable to high precision
- The BLAST Experiment D.K. Hasell *et al.*, Ann. Rev. Nucl. Part. Sci. **61**, 409 (2011)
- Quasielastic (e,e'p) Results A. DeGrush et al. Phys. Rev. Lett. 119, 182501 (2017)
- Conclusions



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AUGUST 1996

#### Femtometer toroidal structures in nuclei

J.L. Forest, V.R. Pandharipande, S.C. Pieper, R.B. Wiringa, R. Schiavilla, A. Arriaga

#### Nucleon-nucleon potential

 $v_{0,1} = v_{0,1}^{c}(r) + \underline{v}_{0,1}^{t}(r)S_{ij} + v_{0,1}^{ls}(r)\mathbf{L} \cdot \mathbf{S} + v_{0,1}^{l2}(r)L^{2}$  $+ v_{0,1}^{ls2}(r)(\mathbf{L} \cdot \mathbf{S})^{2},$ 

#### tensor force

- θ is the polar angle of **r** wrt the spin quantization axis
- Strong effect of the tensor force on NN potential depending on different M<sub>s</sub> substates
- Equidensity surfaces have very different structures depending on M<sub>s</sub>



# Equidensity Surfaces having $\rho_d^{\pm 1} = 0.24 \text{ fm}^{-3}$ (A) and $\rho_d^{0} = 0.24 \text{ fm}^{-3}$ (B)

In the absence of the tensor force, the equidensity surfaces are concentric spheres





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# **Predictions for Quasielastic (e,e'p)**



FIG. 12. The calculated values of  $\tilde{d}(e,e'p)n$  cross section for the kinematics described in the text. Hollow and full symbols indicate results of complete calculations without and with meson-exchange currents.

#### Zero around 300 MeV/c!

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#### Measurement of the Vector and Tensor Asymmetries at Large Missing Momentum in Quasielastic $(\vec{e}, e'p)$ Electron Scattering from Deuterium

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### **MIT-Bates Linear Accelerator Center**

DOE Nuclear Physics National User Facility 1974-2005

> DOE Nuclear Physics Research & Engineering Center of Excellence 2005 - present

SPIN 2018 Ferrara, Italy

500m



- Siberian snake
- Spin flipper
- Compton polarimeter



**South Hall Ring** 

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SPIN 2018 Ferrara, Italy

 $L \sim 10^{32}$  electron-atoms cm<sup>2</sup>s<sup>-1</sup>

### **Stored beam for BLAST**



- Accelerator complex and BLAST experiment fully automated
- Stored currents: routinely fill to 225 mA, lifetime of 35 minutes at 100mA
- Beam Polarization:
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- ~65% with possibility of rapid reversal (flipper) SPIN 2018 10

Ferrara, Italy

### **South Hall Ring Polarization**

#### Electron beam energy: 850 MeV



#### **Compton polarimeter data from Dec. 2003 – Dec. 2004 Mean polarization of 66% measured**

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### **BLAST Detector**



- Toroidal magnetic field
  - 3.8 kG max
- Drift Chambers
  - 3 chambers/sector
  - 2 superlayers/chamber (±5°)
  - 3 sense layers/superlayer
  - 18 tracking layers/sector
  - 954 sense wires
  - Cerenkov Detectors
    - 1 cm thick aerogel
    - Electron identification
    - Time of flight scintillators
      - 16 vertical bars, 2.5 cm thick
      - Trigger and relative timing
- Neutron detectors
  - 10 cm thick in left sector
  - 25-30 cm thick in right sector
- 2 level, 8 channel trigger system
  - Concurrent data acquisition

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NEUTRON

### **BLAST Toroid**





- 8 copper coils – 6730 A
  - 3700 G
- field mapped (3D)
  - coil position adjusted
  - ±1% of calculated
  - minimize target field
  - tracking



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# **BLAST Polarized <sup>2</sup>H Target**



D. Cheever et al., Nucl. Instr. Meth. A 556, 410 (2006)

- Atomic beam source embedded in ≈2 kG BLAST toroidal magnetic field
- Required custom shielding of many components
- Flux:  $5 \ge 10^{16}$  atoms/sec
- Drifilm coated, cryogenically cooled target cell
- Cycled through M<sub>s</sub> sub-states
- Target spin-states switched every 5 minutes
- $h = 0.656 \pm 0.007(\text{stat}) \pm 0.04$  (sys)
- $hPz = 0.580 \pm 0.0034$  (stat)  $\pm 0.0054$ (sys)
- $Pzz = 0.683 \pm 0.015(stat) \pm 0.013(sys)$

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#### H & D2 collected charge for BLAST, 2004-2005 = 3.25 MegaCoulomb



### **Elastic Electron Scattering Primer**



$$W^2 = (q + p)^2 = M^2 + 2Mv - Q^2$$

Elastic scattering:  $W^2 = M^2 \implies Q^2 = 2 Mv$ 

$$\frac{d\sigma}{d\Omega} = \left(\frac{d\sigma}{d\Omega}\right)_{Mott} \bullet \left[\frac{G_E^{\ \ p2} + \tau G_M^{\ \ p2}}{1 + \tau} + 2\tau G_M^{\ \ p2} \tan^2 \frac{\theta}{2}\right]$$

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### BLAST scientific motivation: nucleon and nuclear structure at low Q<sup>2</sup>

- Pion is essential to understanding both nucleon and nuclear structure
- In low energy elastic electron-nucleon scattering one would expect effects of mesons to occur at

 $r \sim 2 \text{ fermi} \implies Q^2 \sim 0.1 (GeV/c)^2$ 

- Search for effects of meson cloud on long distance structure of nucleon
- Seek precise determination of neutron electric form factor with low systematic uncertainties
- Spin structure of deuterium is a stringent test of our understanding of the nucleon-nucleon interaction in nuclei
- **Optimal experimental technique**: precision experiments possible using polarized gas target internal to electron storage ring

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## **Polarized Elastic Electron Scattering**

- In elastic scattering, polarizing both the electron and the target proton, allows determination of the ratio  $G_E/G_M$  with low systematic uncertainty.
- With a tensor polarized deuteron target, elastic scattering yields  $T_{20}$ .
- With a vector polarized deuteron target, and by detecting the neutron in coincidence, the ratio  $G_E/G_M$  can be determined for the neutron.
- Spin is used as a "knob" to access scattering from the neutron.
- By detecting quasielastic <sup>2</sup>H(e,e'n) scattering, G<sup>n</sup><sub>E</sub>(Q<sup>2</sup>) was determined at MIT-Bates at low Q<sup>2</sup> using the Bates Large Acceptance Spectrometer Toroid (BLAST).

### **Charge distribution of neutron**



### Interpretation of Neutron Charge Distribution



#### Figure 9

 $4\pi r^2 \rho_{\text{Breit}}^n(r)$  showing the relative contributions of the various vector mesons from the GKex model together with the perturbative quantum chromodynamics (pQCD) contribution.

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Precision measurement of deuteron tensor analyzing powers with BLAST

C. Zhang *et al.,* Phys. Rev. Lett. **107**, 252501 (2011)

Precision data validate effective field theory

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# Quasielastic <sup>2</sup>H(e,e'p) Reaction

**Cross section** 

$$\frac{d\sigma}{d\omega d\Omega_e d\Omega_{pn}^{\rm CM}} = S_0 [1 + P_z A_d^V + P_{zz} A_d^T + h(A_e + P_z A_{ed}^V + P_{zz} A_{ed}^T)]$$



(e)

Missing momentum

$$\boldsymbol{p}_m \equiv \boldsymbol{q} - \boldsymbol{p}_f$$

- Electron beam polarized
- <sup>2</sup>H both vector and tensor pol.
- **q** and **p**<sub>f</sub> and thus **p**<sub>m</sub> determined by BLAST spectrometer
- Large acceptance => p<sub>m</sub> measured up to 500 MeV/c
- Data taking simultaneous with G<sup>n</sup><sub>E</sub> and T<sub>20</sub> measurements

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(d)

# Quasielastic <sup>2</sup>H(e,e'p) Data Taking



FIG. 1. Histograms of the yields versus missing mass for target spin angle  $\approx 31^{\circ}$  without (red) and with (black) Čerenkov cuts for  $0.1 < Q^2 < 0.5 \ (\text{GeV}/c)^2$  for *opposing* (left) and *same* (right) sector kinematics.

- Data taken in two separate running periods
- Simultaneous with BLAST measurements of  $G^{n}_{E}$  and  $T_{20}$
- Average target spin angles were 31.3±0.43° and 47.4±0.45°
- Target spin angle was in horizontal plane pointing into the left sector
- Determined using T<sub>20</sub> data
  - Electrons scattered into the right (left) sector delivered momentum transfer predominantly parallel (perpendicular) to the target spin vector, the so-called *same sector* (*opposing sector*) kinematics

# Quasielastic <sup>2</sup>H(e,e'p) Vector Asymmetries



FIG. 2. Beam-vector asymmetries  $A_{ed}^V$  for  $0.1 < Q^2 < 0.5$  (GeV/c)<sup>2</sup> vs  $p_m$ . Panels (a) and (c) refer to same sector kinematics for target spin angles  $\approx 31^\circ$  and  $\approx 47^\circ$ . Panels (b) and (d) refer to opposing sector kinematics for the same target spin angles.

$$\mathbf{A^{V}_{ed}} \approx \mathbf{hP_{z}} \text{ at } \mathbf{p_{m}} = \mathbf{0}$$
  
$$P = \sqrt{\frac{2}{3}} P_{z} \left( P_{s} - \frac{1}{2} P_{D} \right)$$

#### Zero crossing at about 320 MeV/c

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# Quasielastic <sup>2</sup>H(e,e'p) Tensor Asymmetries



FIG. 3. Tensor asymmetries  $A_d^T$  for  $0.1 < Q^2 < 0.5$  (GeV/c)<sup>2</sup> vs  $p_m$ . Panels (a) and (c) refer to *same sector* kinematics for target spin angles  $\approx 32^\circ$  and  $\approx 47^\circ$ . Panels (b) and (d) refer to *opposing sector* kinematics for the same target spin angles.

- At low p<sub>m</sub>, S-state dominates so A<sup>T</sup><sub>d</sub> is small
- At high p<sub>m</sub>, D-state dominates
- However, at high p<sub>m</sub> there is a strong influence of FSI: tensor component of NN force

# **Summary**

- Newly published data are reported for  $A_{ed}^{V}$  and  $A_{d}^{T}$  spin asymmetries for  $0.1 < Q^2 < 0.5$  (GeV/c)<sup>2</sup>.
- Mapped out in quasielastic kinematics  ${}^{2}H(e,e'p)$  over  $0 < p_{m} < 500 \text{ MeV/c}.$
- Polarized electron beam incident on vector and tensor polarized <sup>2</sup>H target.
- Large acceptance detector allows large range in  $Q^2$  and  $p_m$ .
- D-state contribution is clearly evident as  $p_m$  increases.
- $A_{ed}^{V}$  has a zero crossing at  $p_m \approx 320$  MeV/c, as predicted.
- A<sup>T</sup><sub>d</sub> in same and opposing sector kinematics probe the proton-neutron interaction over a large spatial range.
- Theoretical understanding validated by experiment.