



# Spin 2018

23<sup>RD</sup> INTERNATIONAL SPIN SYMPOSIUM  
FERRARA - ITALY

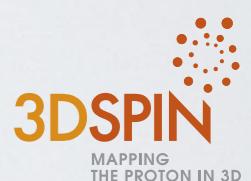
**FIRST extraction of  
TRANSVERSITY  
from data on  
lepton-hadron SCATTERING  
+ hadronic COLLISIONS**

**Marco Radici**  
INFN - Pavia

in collaboration with  
A. Bacchetta (Univ. Pavia)

based on  
**P.R.L. 120** (2018) 192001  
arXiv:1802.05212

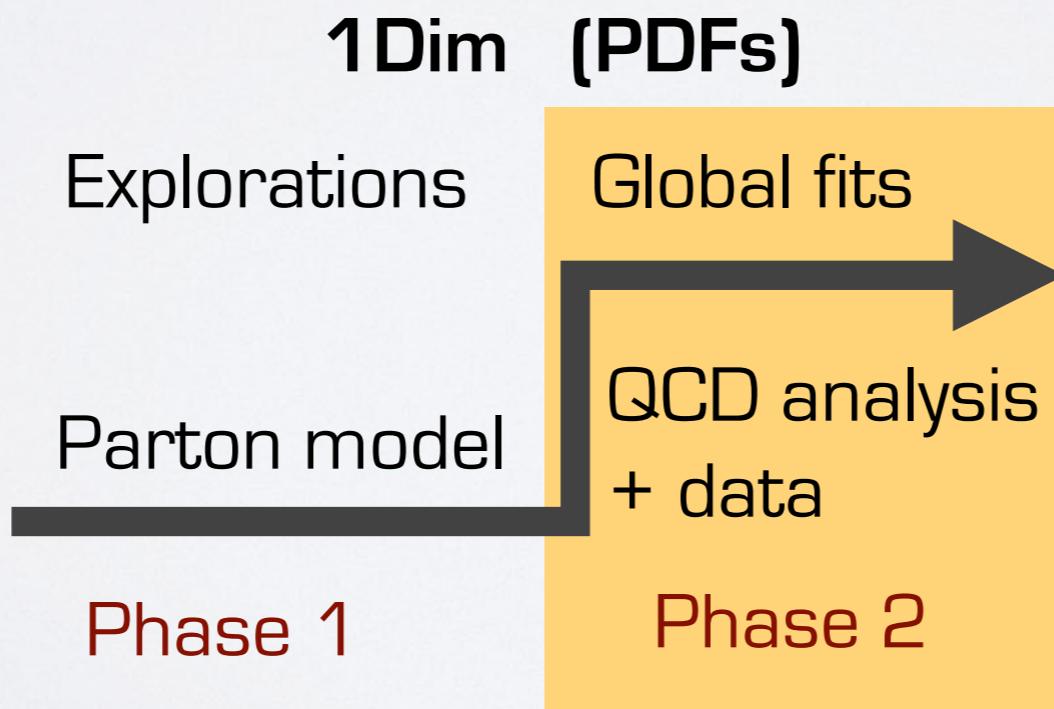
**plus updates**



# a phase transition

		quark polarization		
		U	L	T
nucleon polarization	U	$f_1$		$h_{1^\perp}$
	L		$g_{1L}$	$h_{1L^\perp}$
	T	$f_{1T^\perp}$	$g_{1T}$	$h_1$ $h_{1T^\perp}$

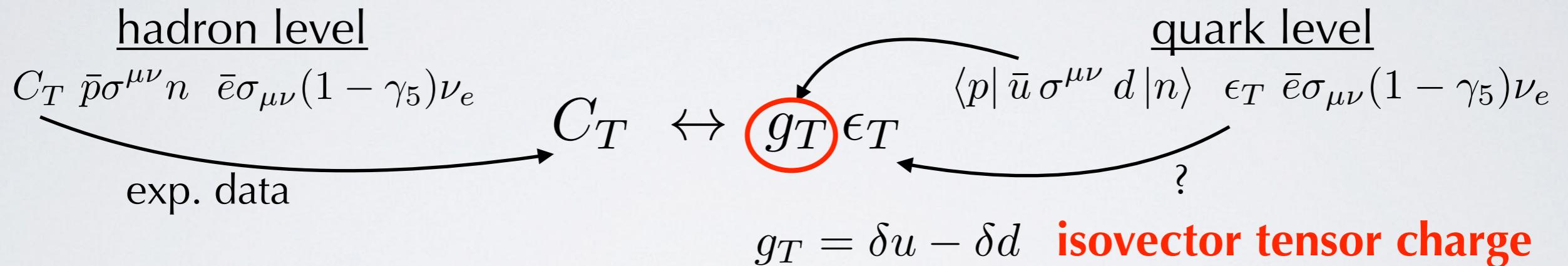
first global fit  
(= lepton-hadron scatt.  
and hadron collisions)  
of PDF  $h_1$



# Motivation

## searches for BSM New Physics

- **nuclear  $\beta$ -decay:** effective field theory including operators not in SM Lagrangian; for example, **tensor operator**



- **neutron EDM:** estimate CPV induced by quark chromo-EDM  $d_q$

$$d_n = \delta u d_u + \delta d d_d + \delta s d_s$$

**tensor charge**

**lattice**
**pheno**

$$\langle P, S | \bar{q} \sigma^{\mu\nu} q | P, S \rangle$$

$$= (P^\mu S^\nu - P^\nu S^\mu) \delta q$$

$$\delta q(Q^2) = \int_0^1 dx h_1^{q-\bar{q}}(x, Q^2)$$

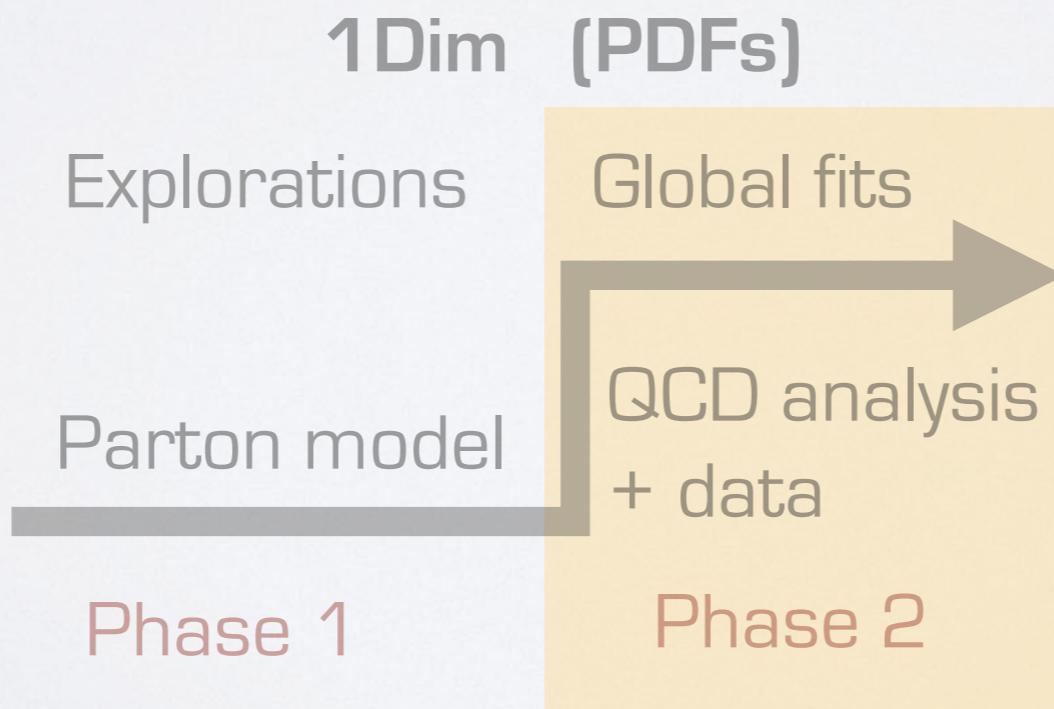
**transversity**

# a phase transition

		quark polarization		
		U	L	T
nucleon polarization	U	$f_1$		$h_{1^\perp}$
	L		$g_{1L}$	$h_{1L^\perp}$
	T	$f_{1T^\perp}$	$g_{1T}$	$h_1$ $h_{1T^\perp}$

chiral-odd  $\rightarrow$  SIDIS

first global fit  
(= lepton-hadron scatt.  
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of PDF  $h_1$

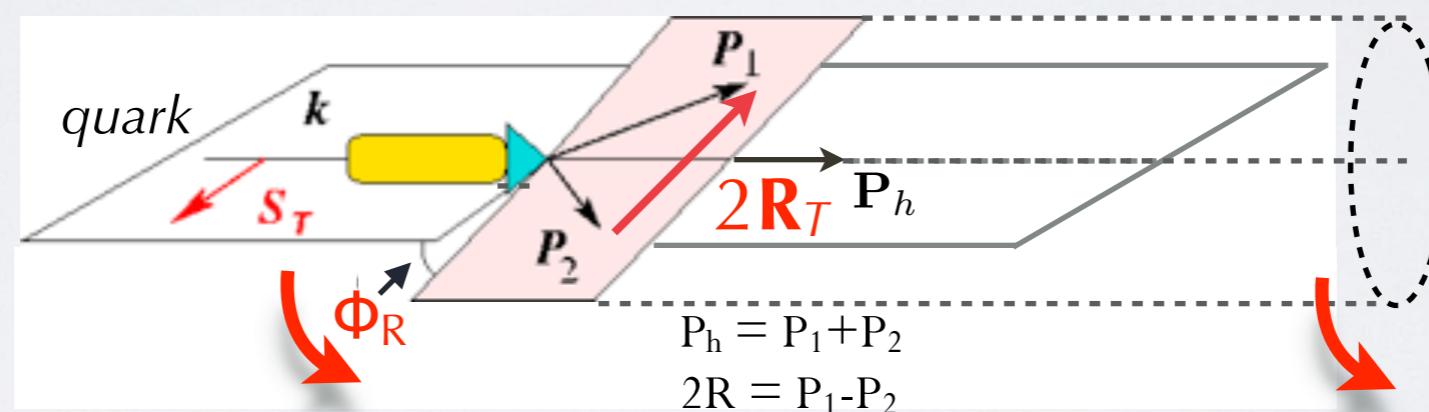


# 2-hadron-inclusive production

Collins, Heppelman, Ladinsky,  
N.P. **B420** (94)

$$R_T \ll Q \quad H_1^{\triangleleft}$$

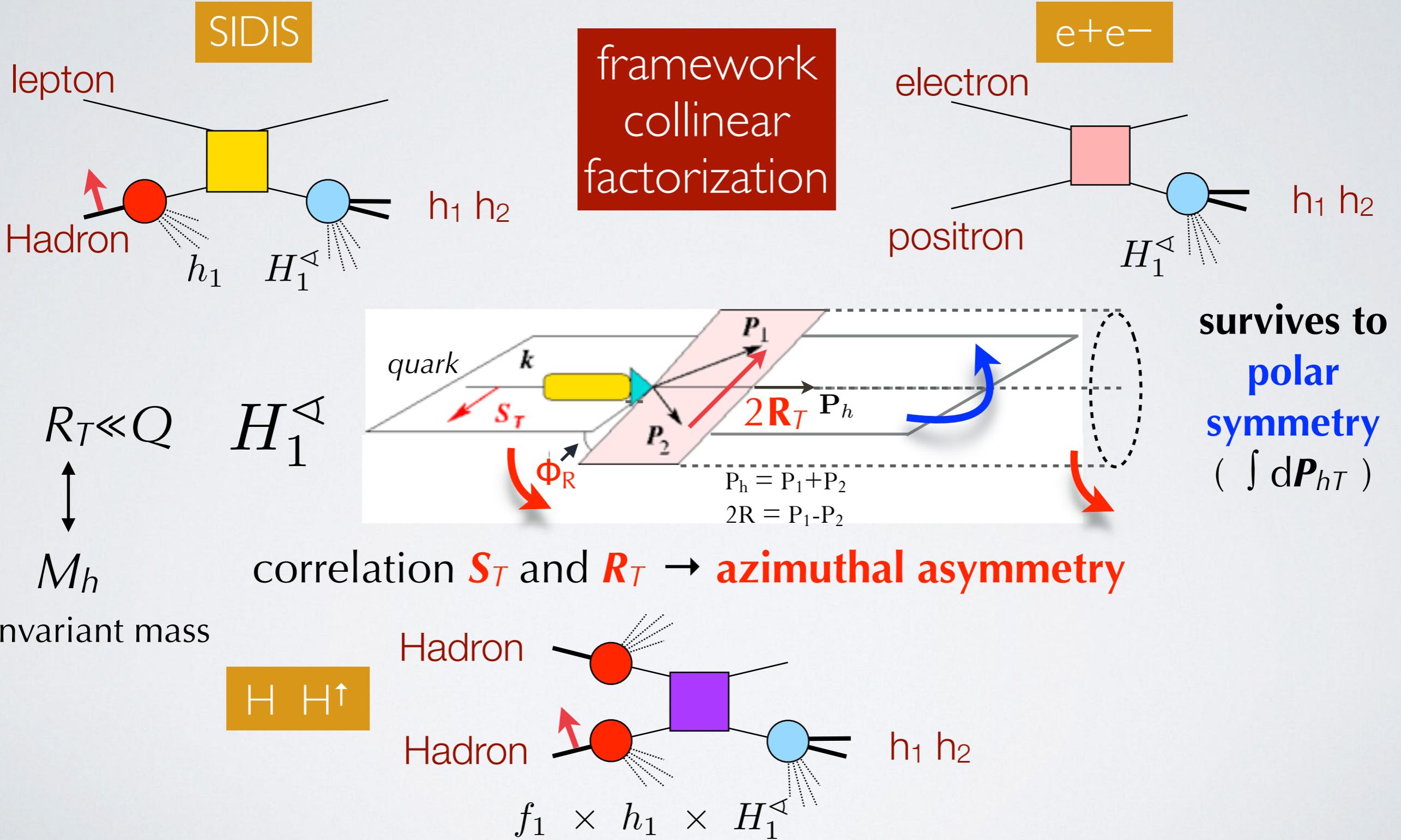
$\updownarrow$



correlation  $S_T$  and  $R_T \rightarrow$  **azimuthal asymmetry**

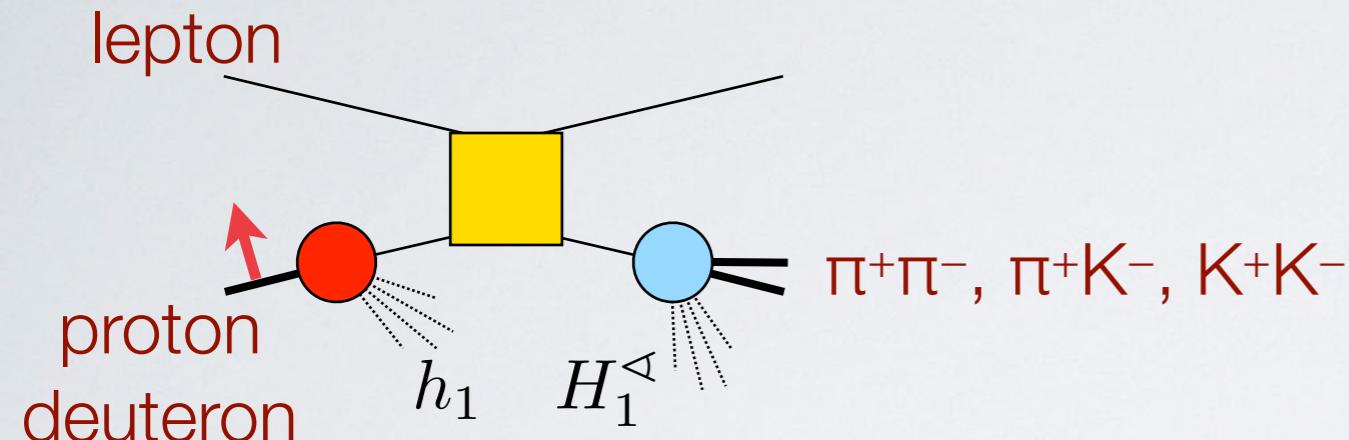
invariant mass

# 2-hadron-inclusive production



# exp. data for 2-hadron-inclusive production

SIDIS  $\ell^- H^\uparrow \rightarrow \ell^+ (h_1 h_2) X$

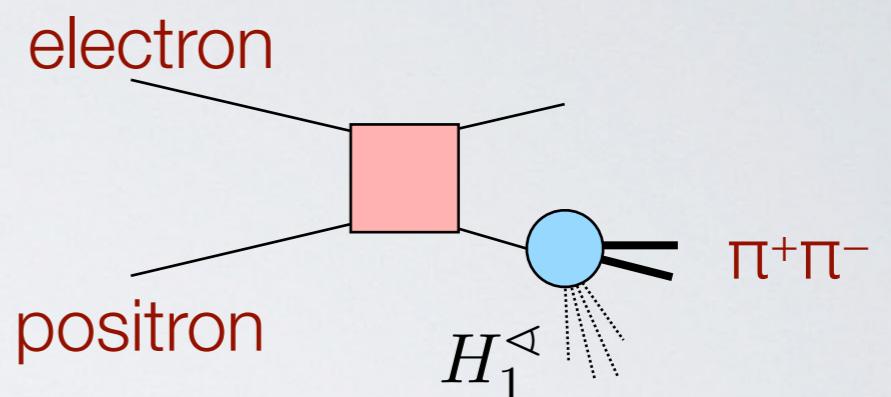


Airapetian et al.,  
*JHEP* **0806** (08) 017



Adolph et al., *P.L.* **B713** (12)  
Braun et al., *E.P.J. Web Conf.* **85** (15) 02018

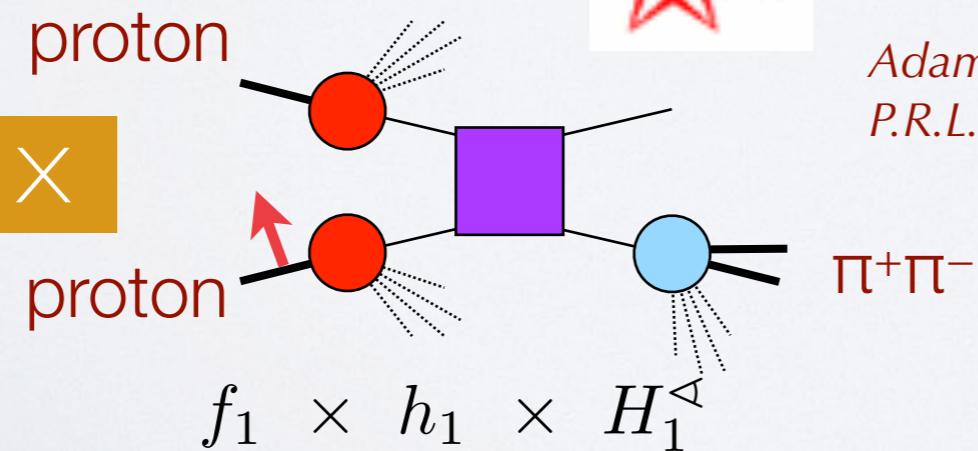
$e^+e^- \rightarrow (h_1 h_2) X$



Vossen et al., *P.R.L.* **107** (11) 072004

$D_1$  Seidl et al., *P.R.* **D96** (17) 032005

$H^- H^\uparrow \rightarrow (h_1 h_2) X$



run 2006 ( $s=200$ )

Adamczyk et al. (STAR),  
*P.R.L.* **115** (2015) 242501

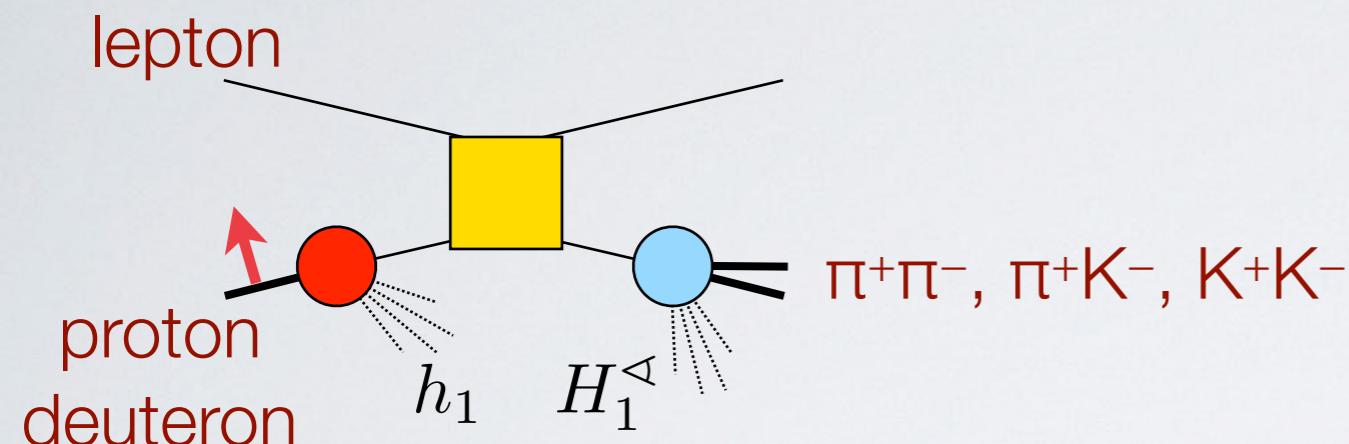
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Adamczyk et al. (STAR),  
*P.L.* **B780** (18) 332

$A_{UT}(\eta, M_h, P_T)$

# exp. data for 2-hadron-inclusive production

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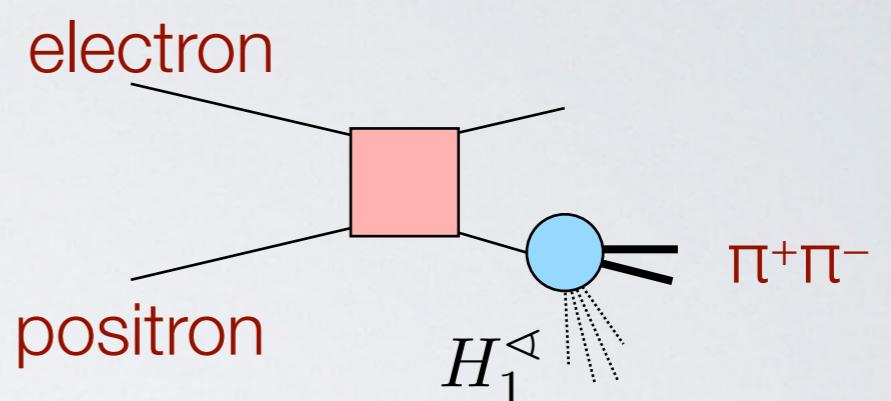


Airapetian et al.,  
*JHEP* **0806** (08) 017



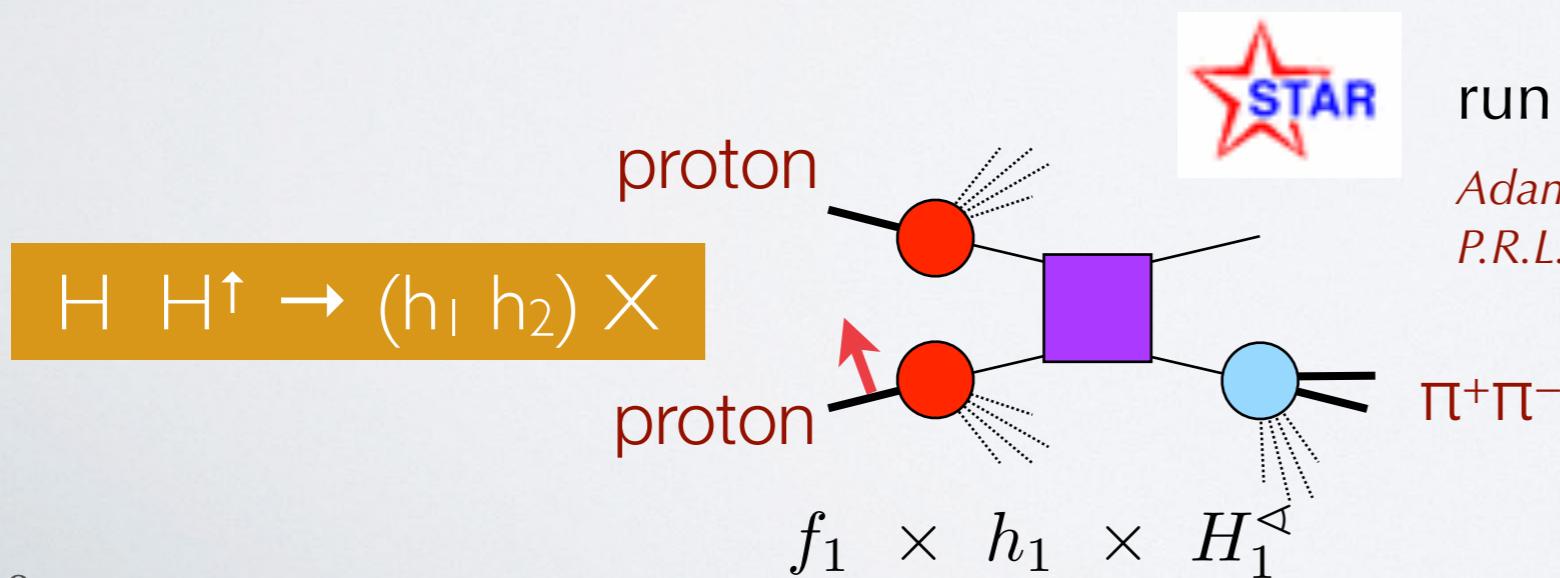
Adolph et al., *P.L.* **B713** (12)  
Braun et al., *E.P.J. Web Conf.* **85** (15) 02018

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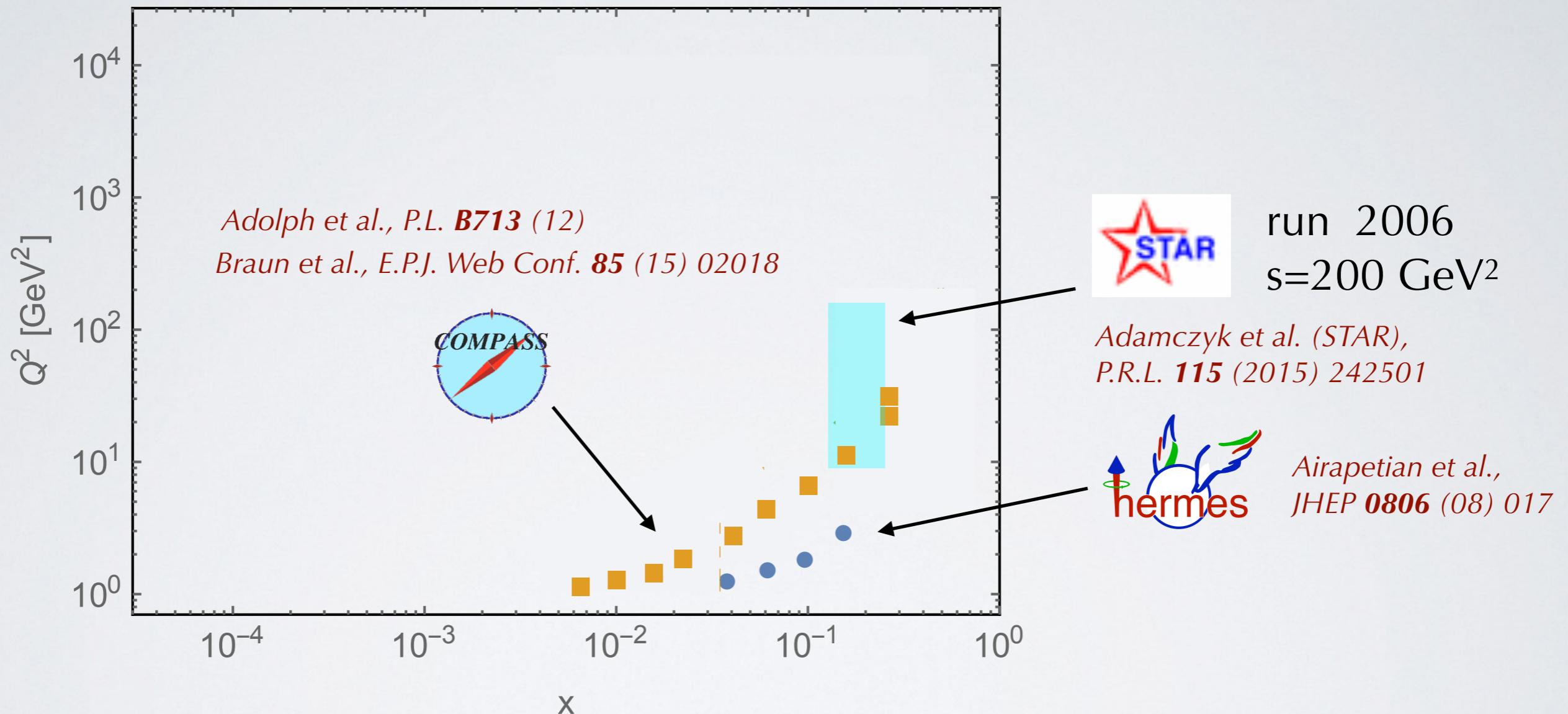


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Adamczyk et al. (STAR),  
*P.L.* **B780** (18) 332

$A_{UT}(\eta, M_h, P_T)$

# the kinematics



# choice of functional form

functional form whose Mellin transform can be computed analytically and complying with Soffer Bound at any x and scale  $Q^2$

$$h_1^{q_v}(x; Q_0^2) = F^{q_v}(x) \left[ \text{SB}^q(x) + \overline{\text{SB}}^{\bar{q}}(x) \right]$$

↓  
**Soffer Bound**

$$2|h_1^q(x, Q^2)| \leq 2 \text{ SB}^q(x, Q^2) = |f_1^q(x, Q^2) + g_1^q(x, Q^2)|$$

MSTW08      DSSV

↙

$$F^{q_v}(x) = \frac{N_{q_v}}{\max_x [|F^{q_v}(x)|]} x^{A_{q_v}} [1 + B_{q_v} \text{Ceb}_1(x) + C_{q_v} \text{Ceb}_2(x) + D_{q_v} \text{Ceb}_3(x)]$$

Ceb<sub>n</sub>(x) Cebyshev polynomial  
10 fitting parameters

constrain parameters

$$|N_{q_v}| \leq 1 \Rightarrow |F^{q_v}(x)| \leq 1 \quad \text{Soffer Bound ok at any } Q^2$$

# choice of functional form

$$h_1^{q_v}(x; Q_0^2) = F^{q_v}(x) \left[ \text{SB}^q(x) + \overline{\text{SB}}^{\bar{q}}(x) \right]$$

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constrain parameters

tensor charge  $\delta q(Q^2) = \int_{x_{\min}}^1 dx h_1^{q-\bar{q}}(x, Q^2)$

low-x behavior is important

if  $\lim_{x \rightarrow 0} x \text{SB}^q(x) \propto x^{a_q}$  then  $h_1^q(x) \underset{x \rightarrow 0}{\approx} x^{A_q + a_q - 1}$

1)  $\delta q$  finite  $\Rightarrow A_q + a_q > 0$

# choice of functional form

$$h_1^{q_v}(x; Q_0^2) = F^{q_v}(x) \left[ \text{SB}^q(x) + \overline{\text{SB}}^{\bar{q}}(x) \right]$$

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1)  $\delta q$  finite  $\Rightarrow A_q + a_q > 0$

2) small-x dipole picture  $\Rightarrow h_1^{q_v}(x) \underset{x \rightarrow 0}{\approx} x^{1-2\sqrt{\frac{\alpha_s(Q^2) N_c}{2\pi}}}$  at  $Q_0$   $A_q + a_q \sim 1$

Kovchegov & Sievert, arXiv:1808.10354

consistent with lattice

Kirschner et al.,  
Z.Phys. **C74** (97) 501

consistent with  
“massive” jet in DIS

Accardi and Bacchetta,  
P.L. **B773** (17) 632

$$\int_0^1 dx g_2(x) \propto \int_0^1 dx \frac{h_1(x)}{x} \longrightarrow A_q + a_q > 1$$

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Z.Phys. **C74** (97) 501

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Accardi and Bacchetta,  
P.L. **B773** (17) 632

our choice  $A_q + a_q > \frac{1}{3}$

grants also error  $O(1\%)$  for  
MSTW08  $x_{\min}=10^{-6}$

$$\int_0^1 dx g_2(x) \propto \int_0^1 dx \frac{h_1(x)}{x} \longrightarrow A_q + a_q > 1$$

# theoretical uncertainties

## unpolarized Di-hadron Fragmentation Function $D_1$

- **quark**  $D_{1q}$  is **well** constrained by  $e^+e^- \rightarrow (\pi^+\pi^-) X$  (Montecarlo)
- **gluon**  $D_{1g}$  is **not** constrained by  $e^+e^- \rightarrow (\pi^+\pi^-) X$  (currently, LO analysis)
- **no data** available yet for  $p p \rightarrow (\pi^+\pi^-) X$

we don't know anything about the gluon  $D_{1g}$

our choice: set  $D_{1g}(Q_0) = \begin{cases} 0 \\ D_{1u}(Q_0) / 4 \\ D_{1u}(Q_0) \end{cases}$

deteriorates our  $e^+e^-$  fit as  $\chi^2/\text{dof} = \begin{cases} 1.69 & 1.28 \\ 1.81 & 1.37 \\ 2.96 & 2.01 \end{cases}$

background     $\rho$     channels

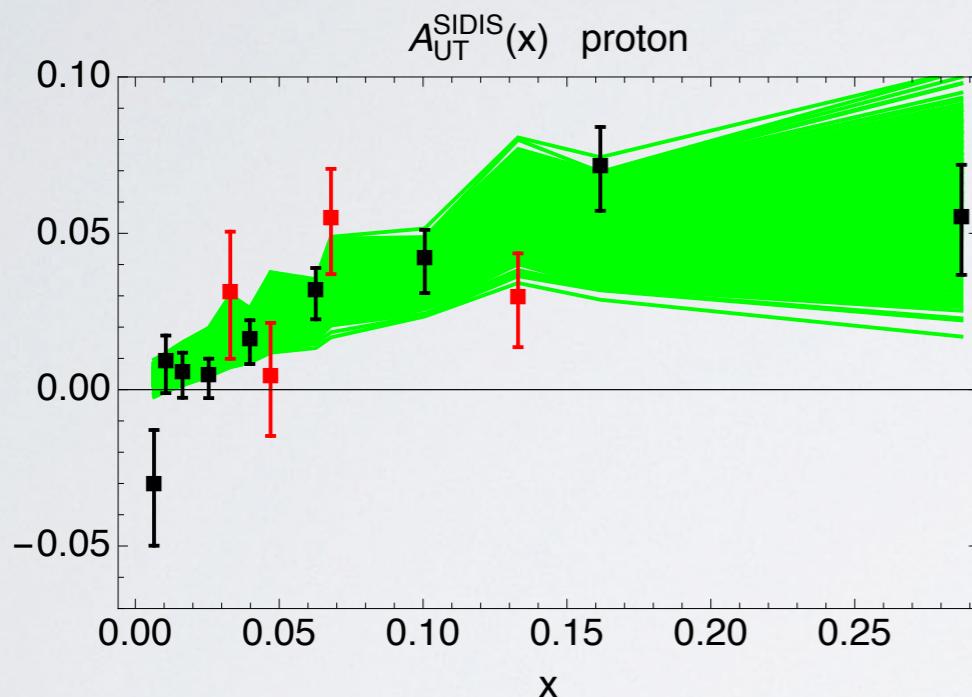
# statistical uncertainty



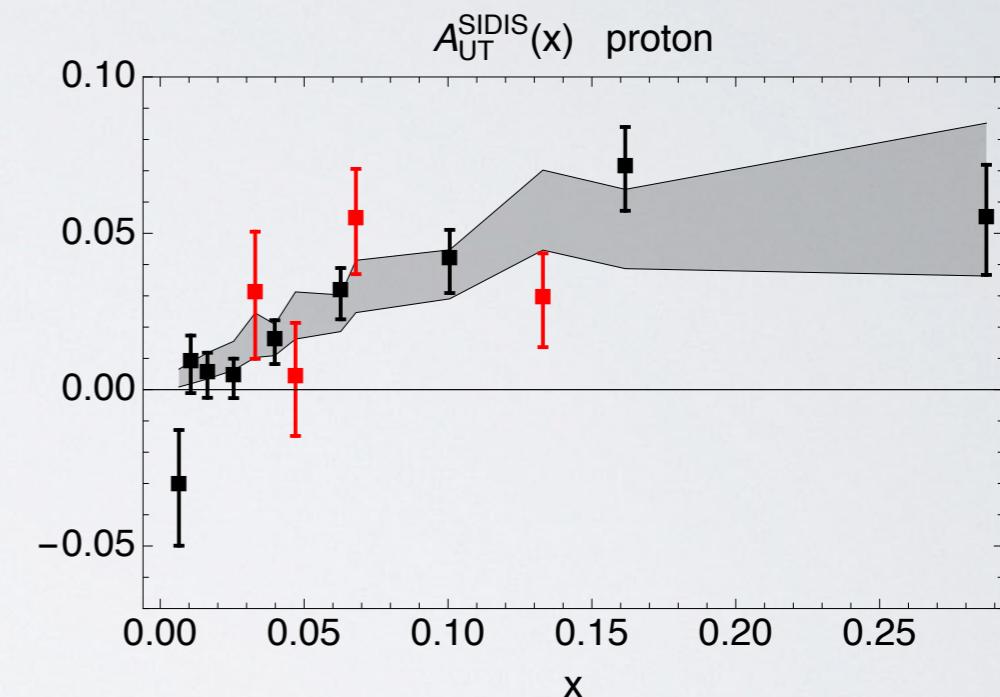
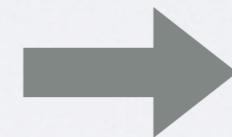
Braun et al., E.P.J. Web Conf. **85** (15) 02018



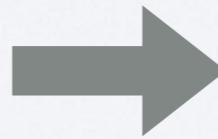
Airapetian et al., JHEP **0806** (08) 017



all 600 replicas



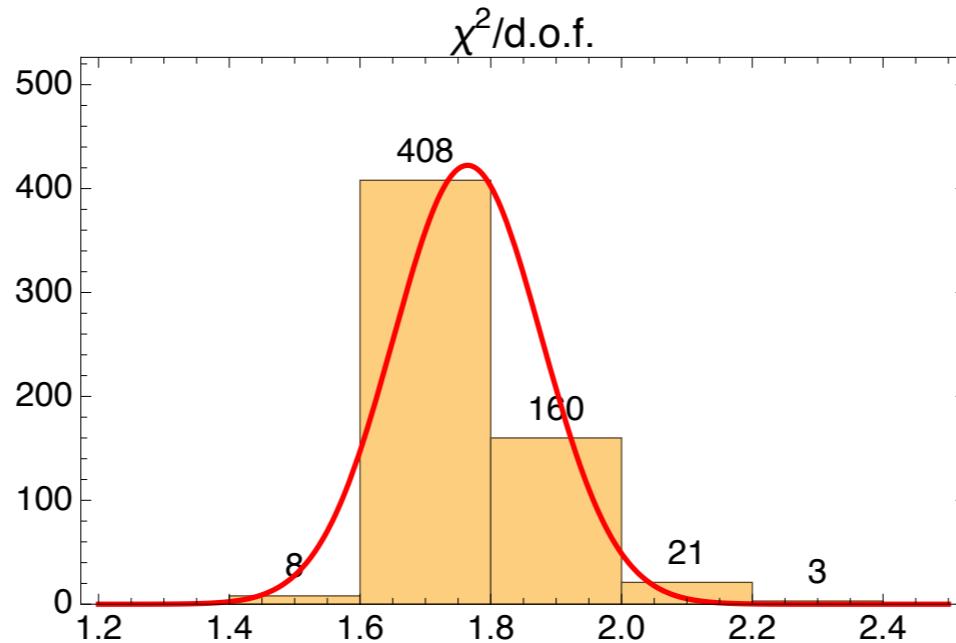
90% of replicas



the bootstrap method

# $\chi^2$ of the fit

$$\chi^2/\text{dof} = 1.76 \pm 0.11$$



**proton SIDIS**

13 data points = 4  + 9

**deuteron SIDIS**

9 data points =  + 9



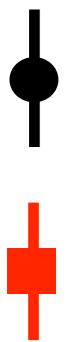
24 data points  $(4 \eta) \times \frac{4}{24} + (10 M_h) \times \frac{10}{24} + (10 p_T) \times \frac{10}{24}$

global fit

10 parameters

# results

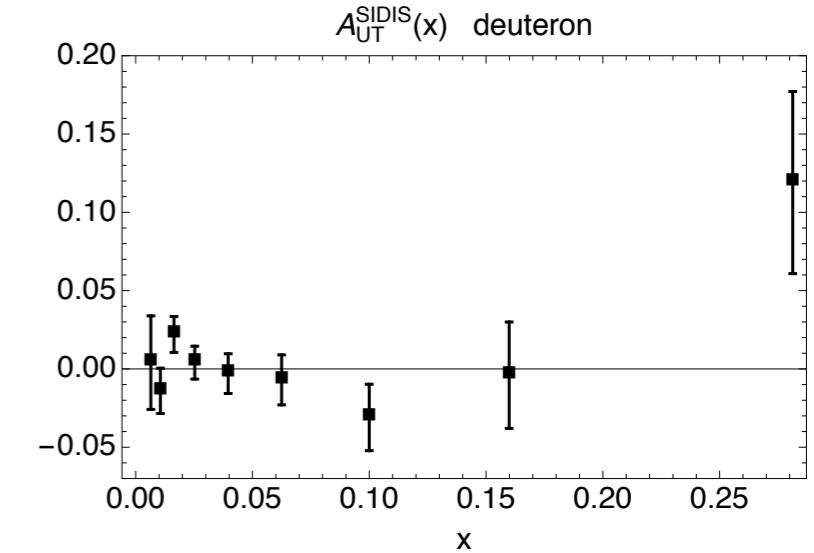
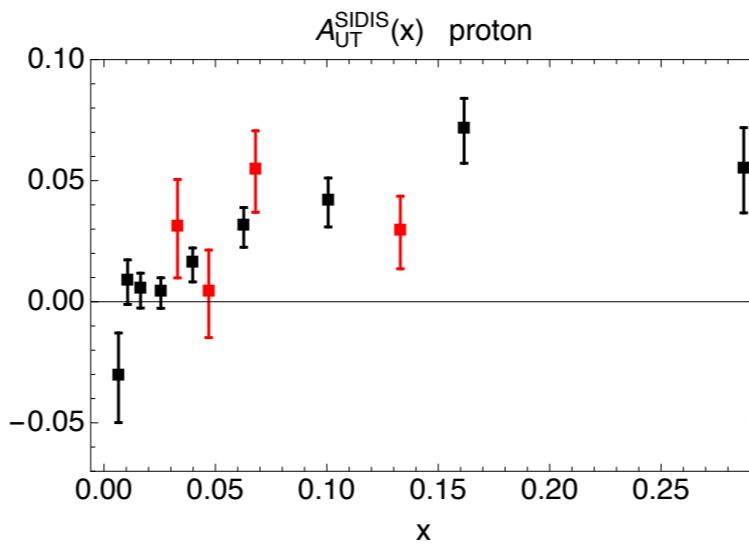
## SIDIS



*Adolph et al., P.L. B713 (12)*



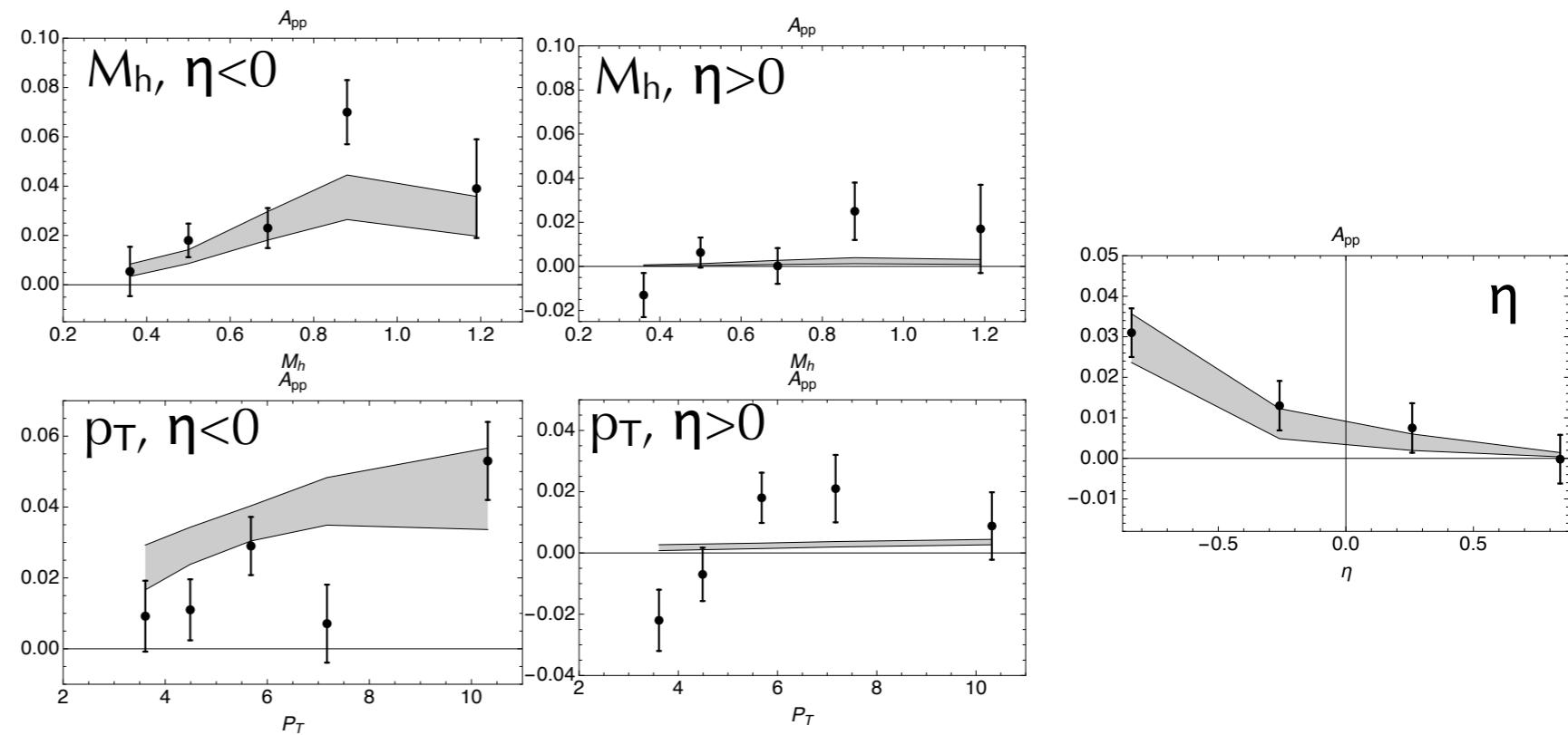
*Airapetian et al.,  
JHEP 0806 (08) 017*



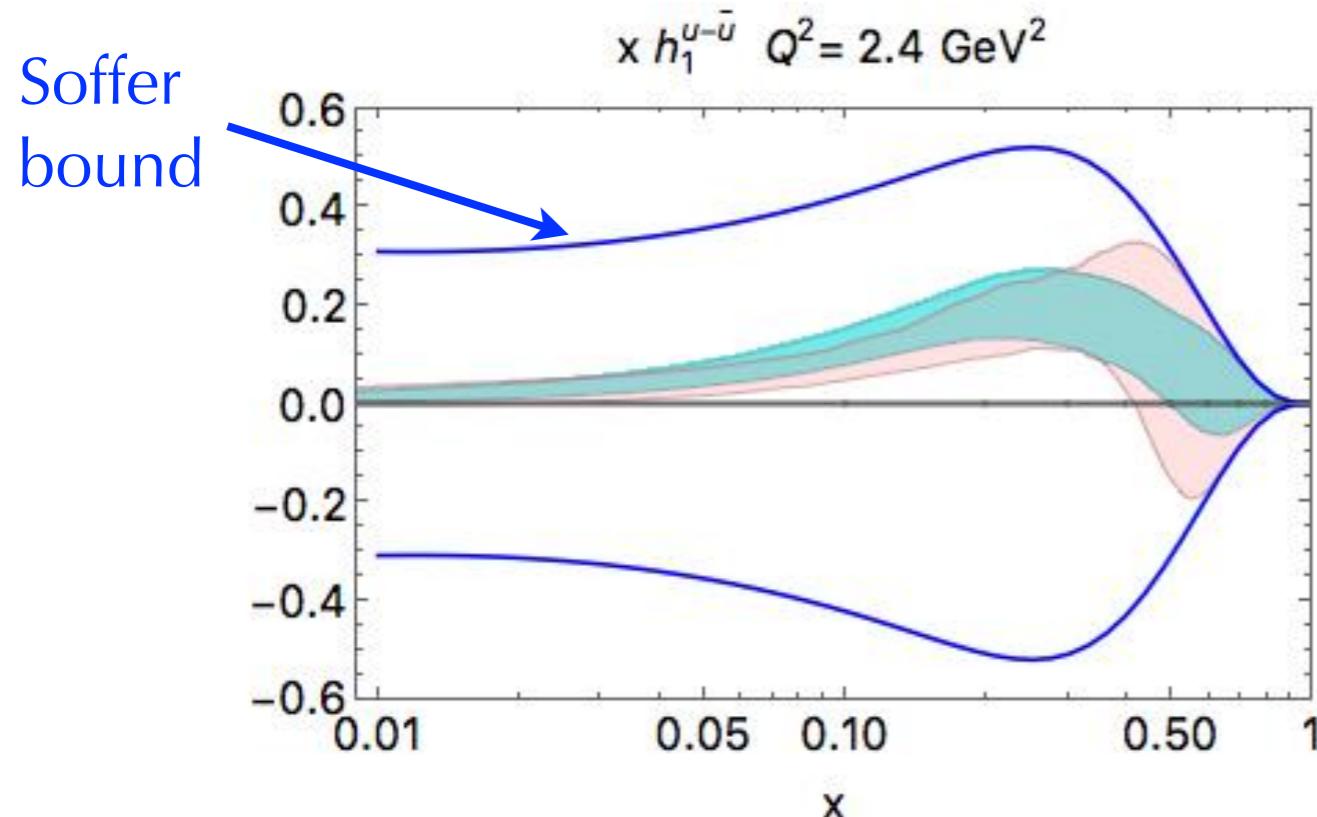
## pp collisions



*Adamczyk et al.,  
P.R.L. 115 (2015) 242501*



# comparison with previous fit



*Radici & Bacchetta,  
P.R.L. **120** (18) 192001*

global fit

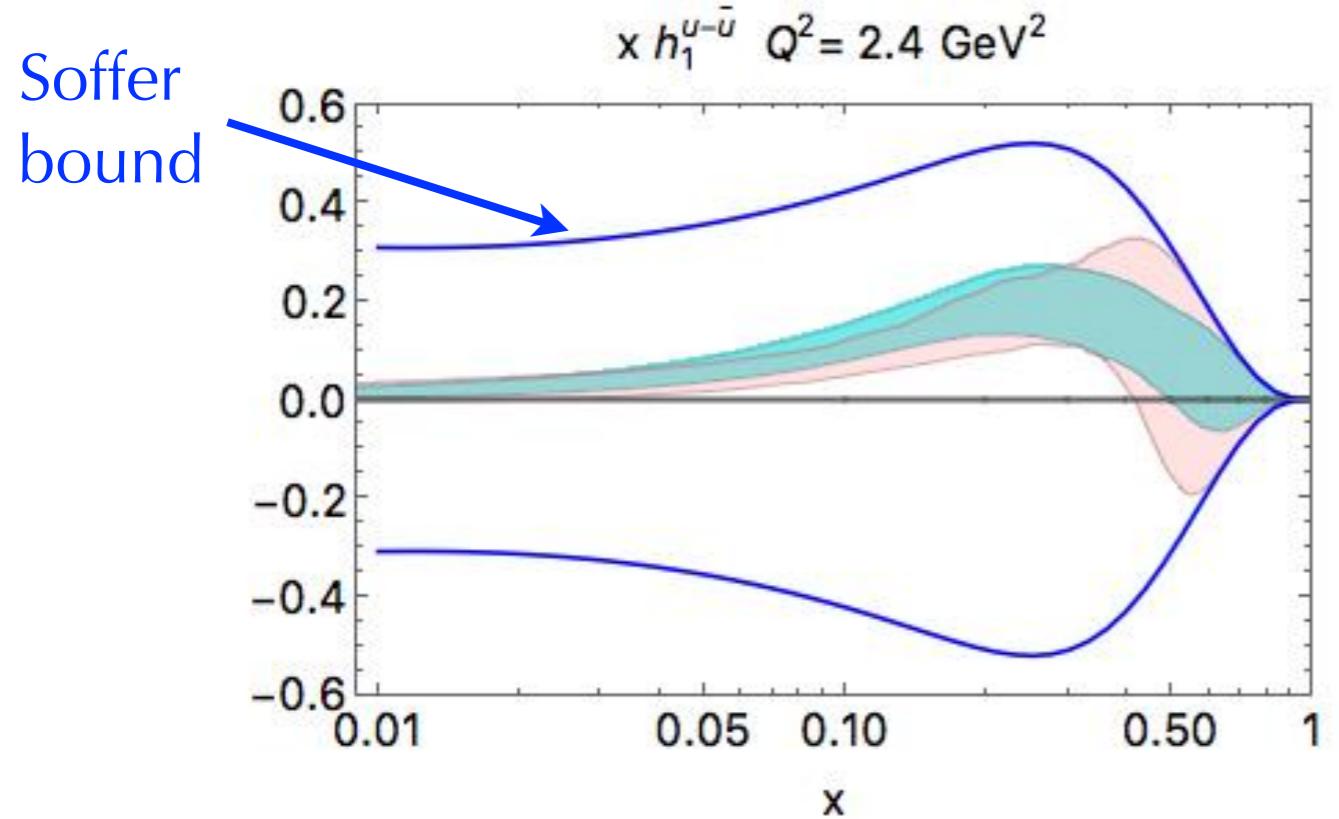
up

higher  
precision

old fit (only SIDIS data)

*Radici et al.,  
JHEP **1505** (15) 123*

# comparison with previous fit



*Radici & Bacchetta,  
P.R.L. **120** (18) 192001*

global fit

up

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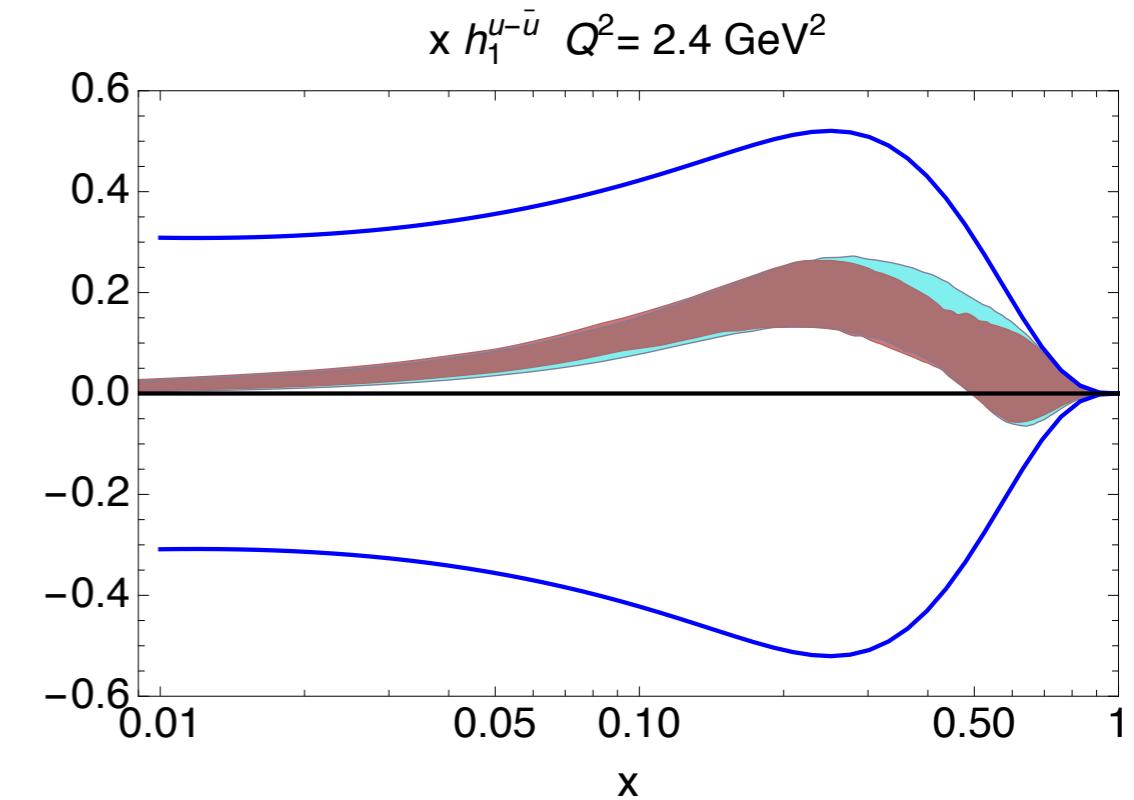
old fit (only SIDIS data)

*Radici et al.,  
JHEP **1505** (15) 123*

up

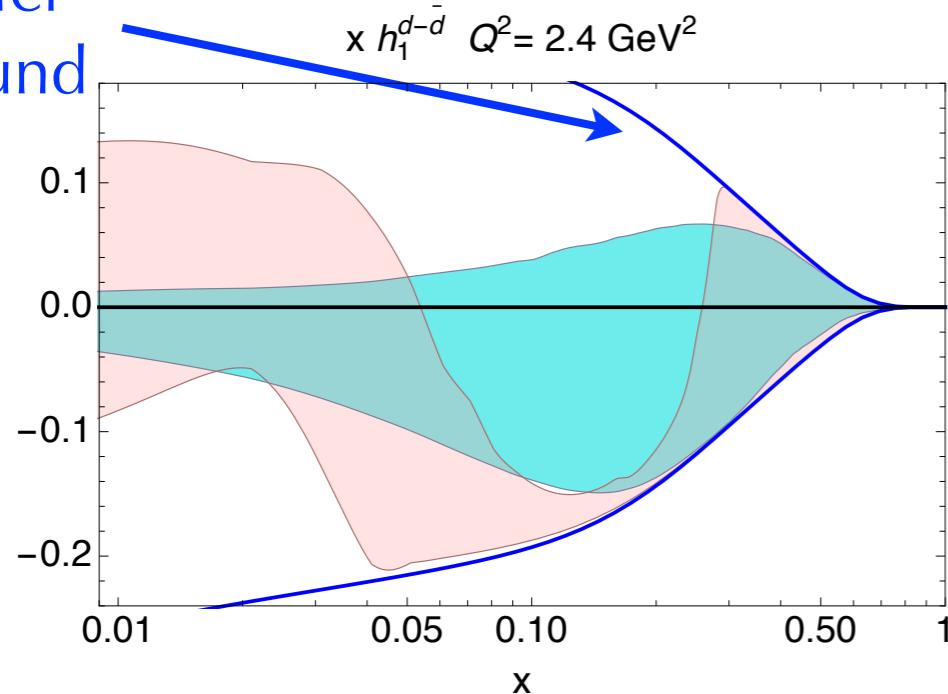
insensitive to  
uncertainty on  
gluon  $D_1$

$$D_{1g}(Q_0) = 0$$
$$D_{1g}(Q_0) = \begin{cases} 0 \\ D_1^u / 4 \\ D_1^u \end{cases}$$



# comparison with previous fit

Soffer  
bound



Radici & Bacchetta,  
P.R.L. 120 (18) 192001

global fit

down

old fit

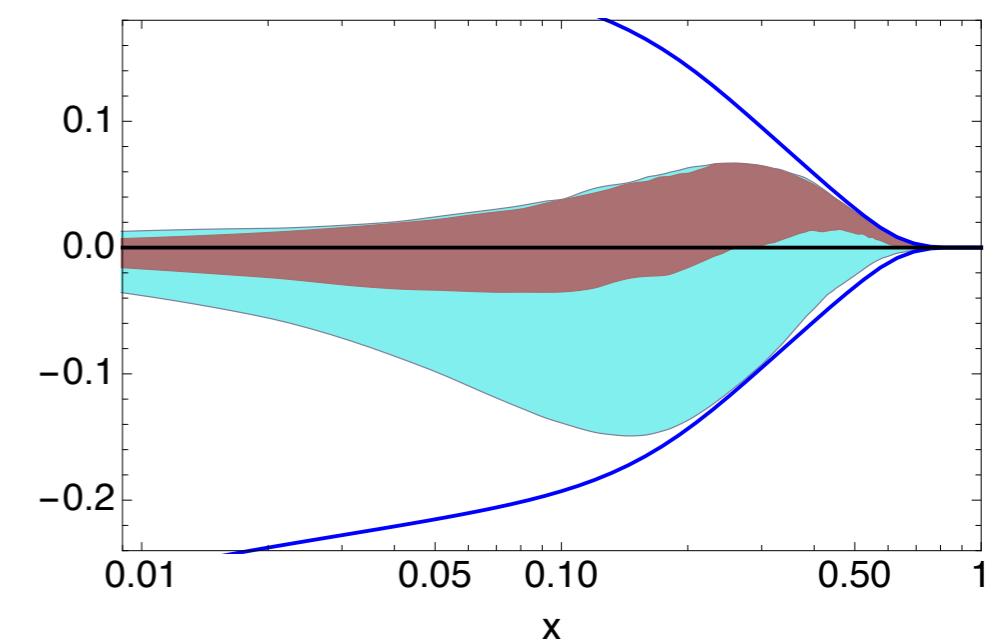
Radici et al.,  
JHEP 1505 (15) 123

down

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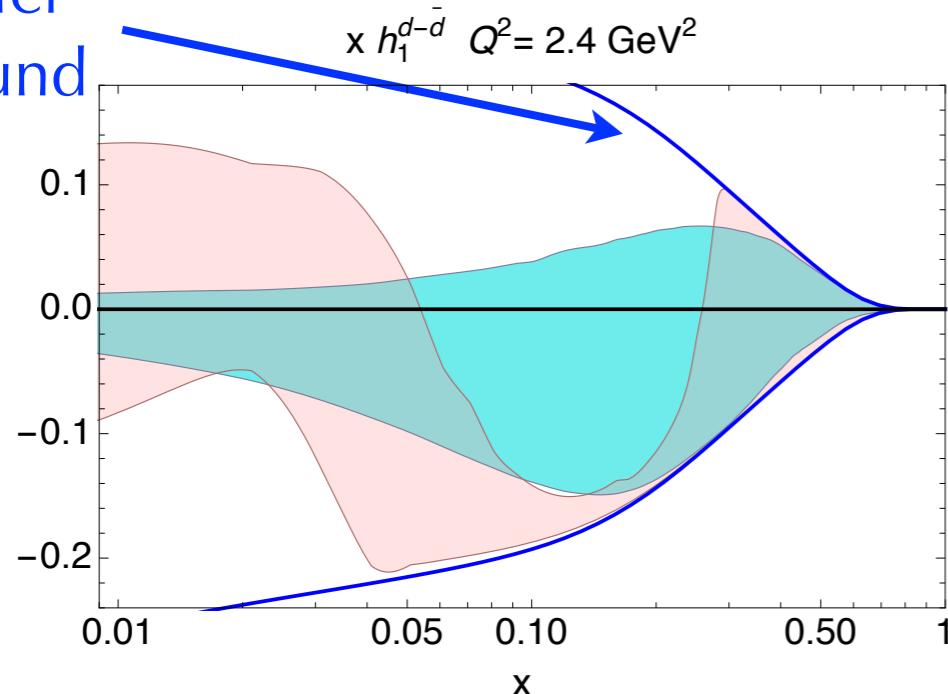
$$D_{1g}(Q_0) = 0$$
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$x h_1^{d-\bar{d}} \quad Q^2 = 2.4 \text{ GeV}^2$



# comparison with previous fit

Soffer  
bound



*Radici & Bacchetta,  
P.R.L. **120** (18) 192001*

global fit

old fit

*Radici et al.,  
JHEP **1505** (15) 123*

down

need better control on  
 $g \rightarrow \pi^+ \pi^-$

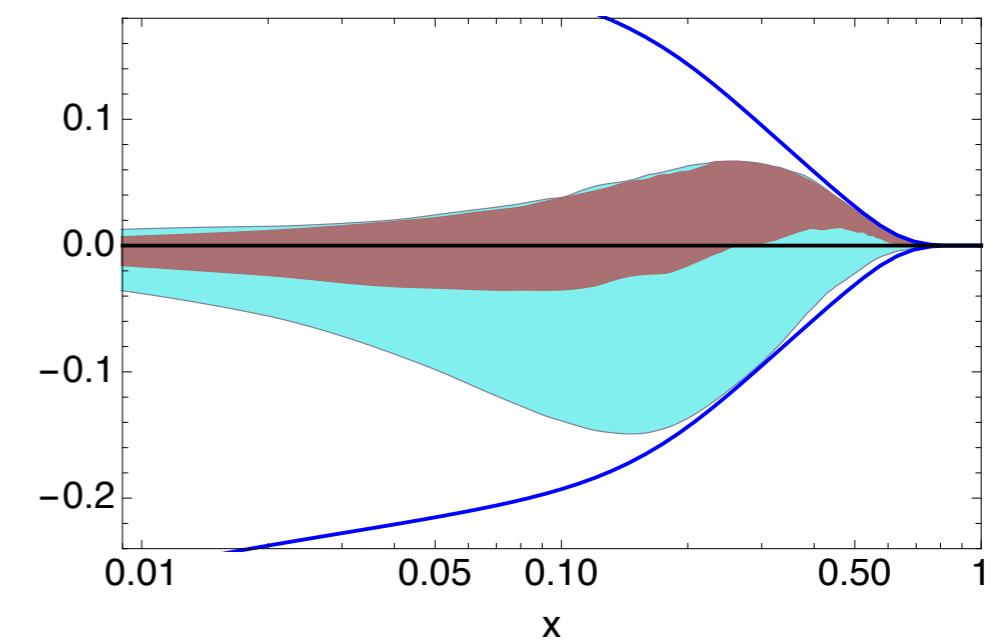
down

sensitive to  
uncertainty on  
gluon  $D_1$

$$D_1 g(Q_0) = 0$$

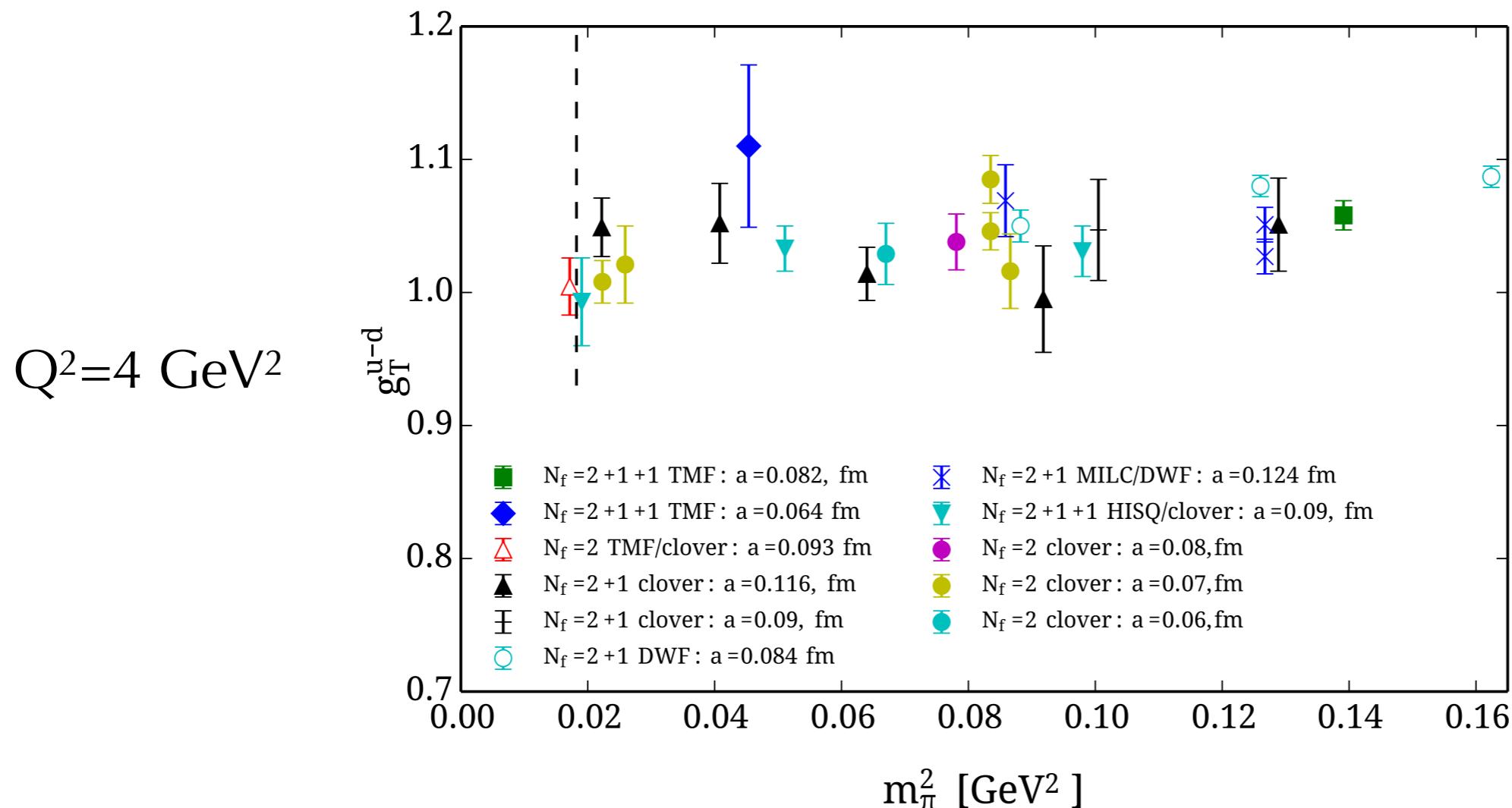
$$D_1 g(Q_0) = \begin{cases} 0 \\ D_1^u / 4 \\ D_1^u \end{cases}$$

$$x h_1^{d-\bar{d}} Q^2 = 2.4 \text{ GeV}^2$$



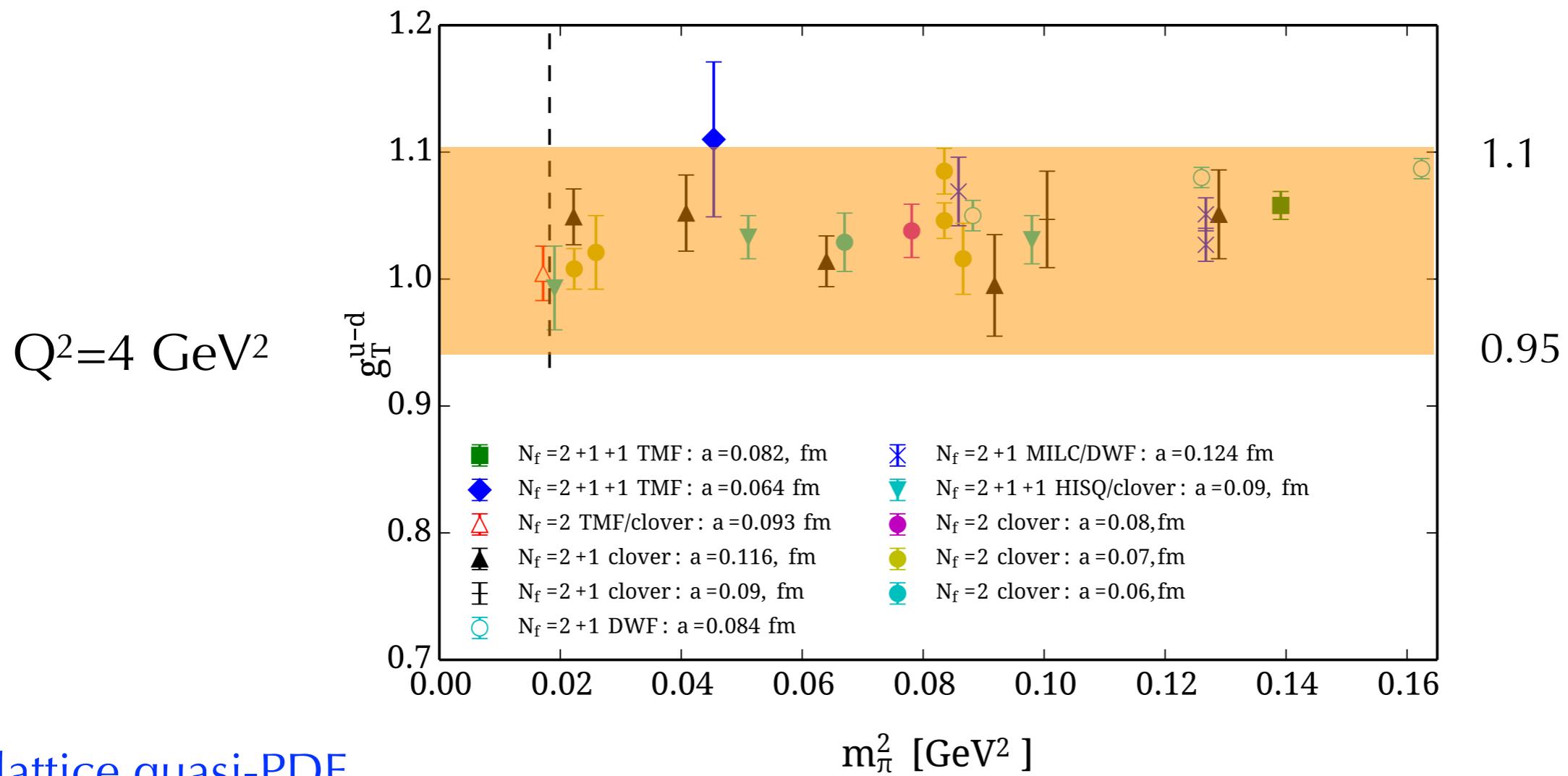
# isovector tensor charge $g_T = \delta u - \delta d$

lattice results  
with different  
discretization schemes, lattice spacings, volumes



# isovector tensor charge $g_T = \delta u - \delta d$

lattice results  
with different  
discretization schemes, lattice spacings, volumes



lattice quasi-PDF

see also talk by F. Steffens (ETMC)

see also arXiv:1803.04393 (LP<sup>3</sup>)

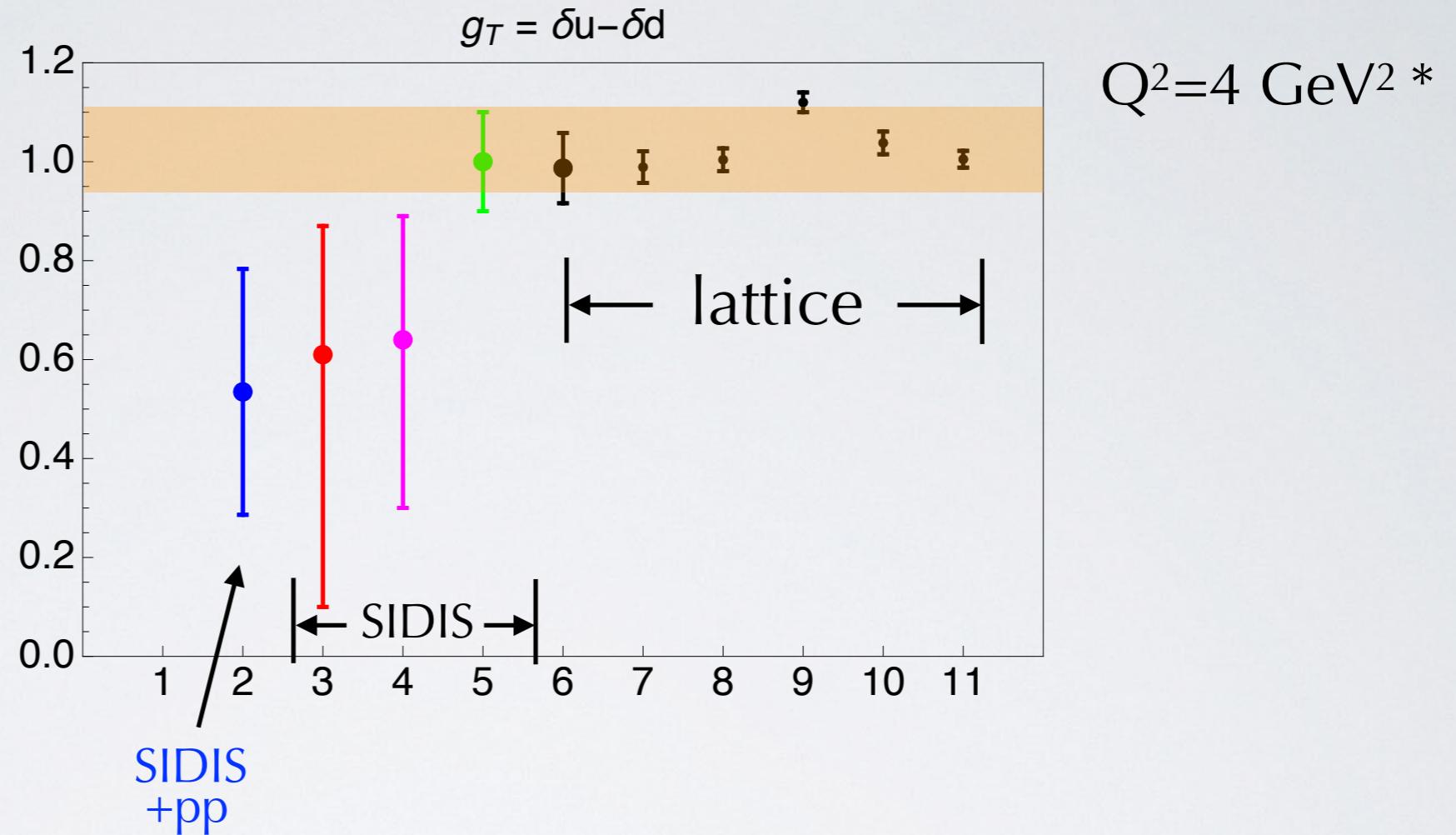
Alexandrou, arXiv:1612.04644

# isovector tensor charge $g_T = \delta u - \delta d$

?



JAM includes  
“lattice data”



*Radici & Bacchetta,  
P.R.L. 120 (18) 192001*

2) **global fit '17**

*Kang et al., P.R. D93 (16) 014009*

3) **“TMD fit” \*  $Q^2=10$**

*Anselmino et al., P.R. D87 (13) 094019*

4) **Torino fit \*  $Q^2=1$**

*Lin et al., P.R.L. 120 (18) 152502*

5) **JAM fit '17 \*  $Q_0^2=2$**

from GPD

see also talk by S. Liuti

6) PNDME '16

*Bhattacharya et al., P.R. D94 (16) 054508*

7) PNDME '18

*Gupta et al., P.R. D98 (18) 034503*

8) ETMC '17

*Alexandrou et al., P.R. D95 (17) 114514;  
E P.R. D96 (17) 099906*

9) NPLQCD '18

*Chang et al., P.R.L. 120 (18) 152002*

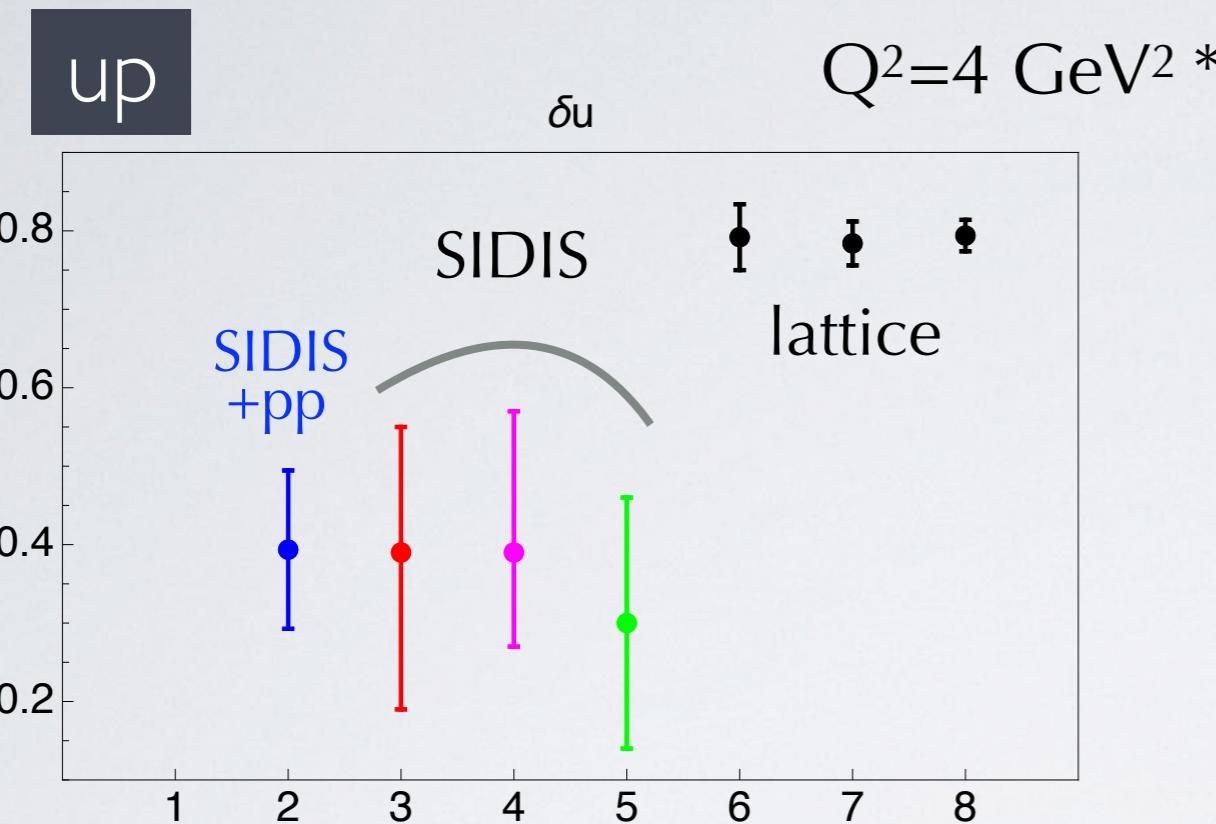
10) LHPC '12

*Green et al., P.R. D86 (12)*

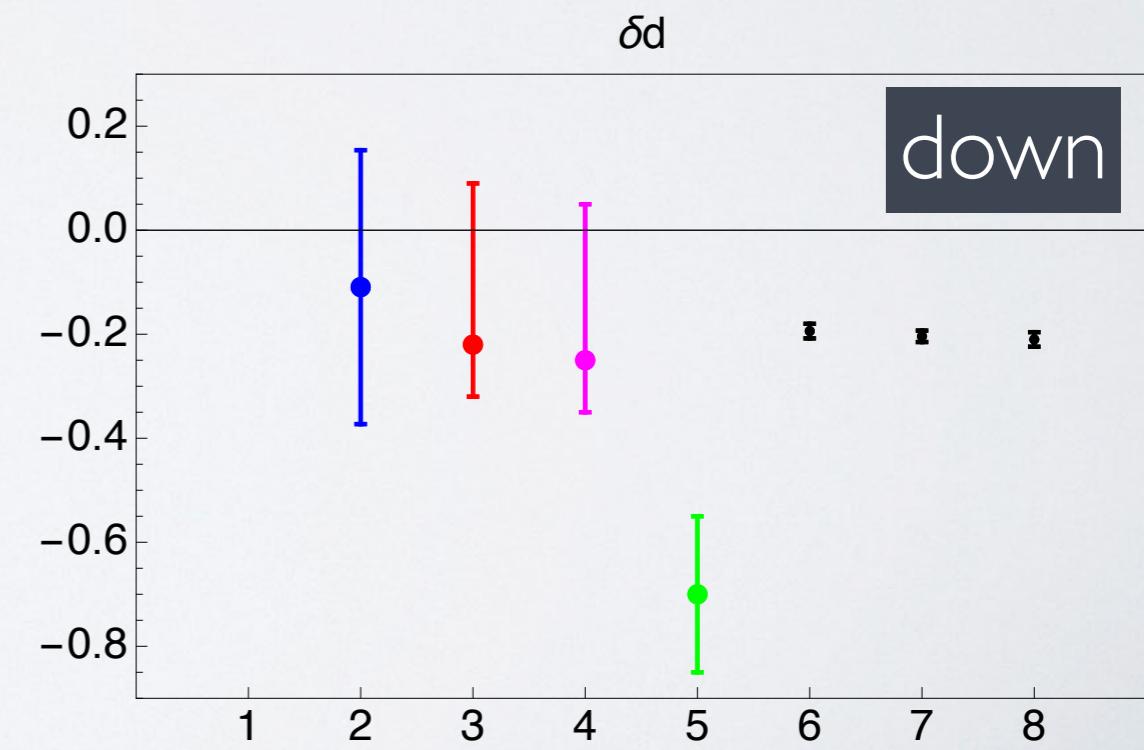
11) RQCD '14

*Bali et al., P.R. D91 (15)*

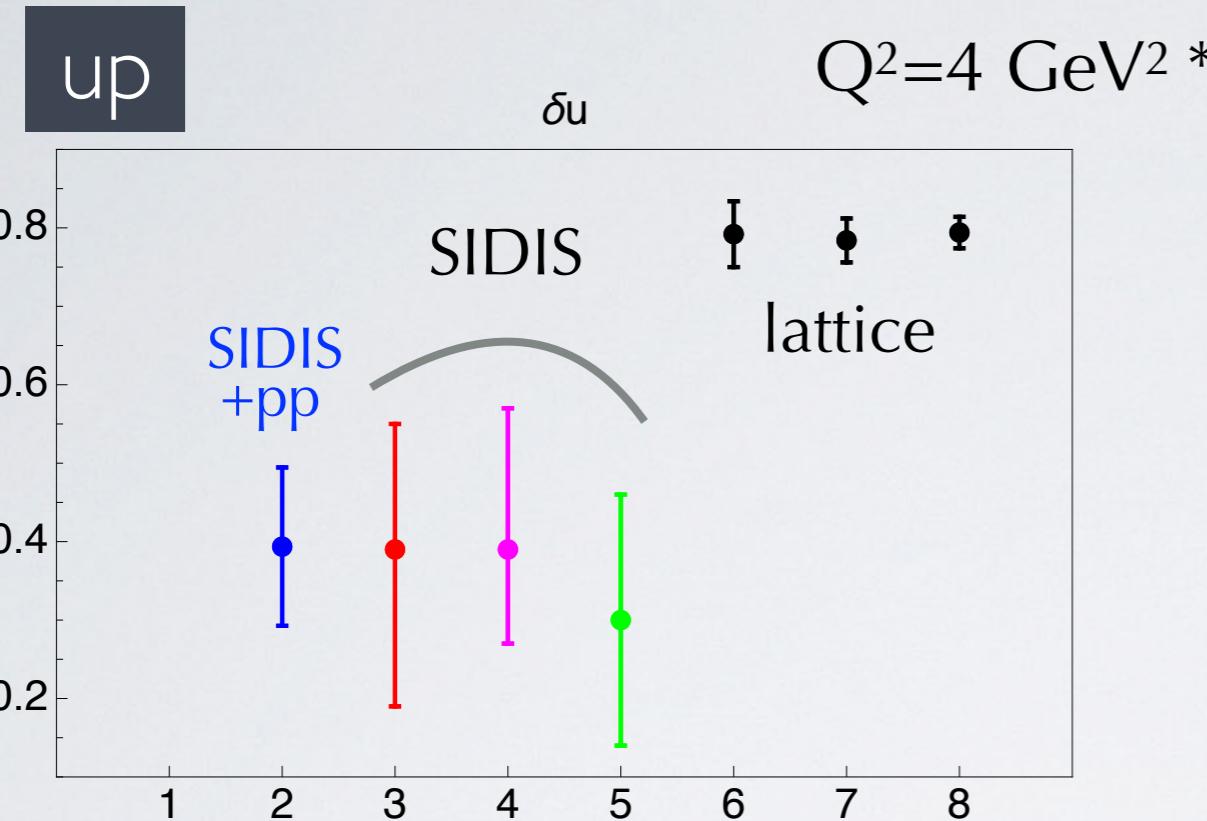
# tensor charge : separate flavors



- 2- global fit** *Radici & Bacchetta, P.R.L.120 (18) 192001*
- 3- TMD fit** *Kang et al., P.R.D93 (16) 014009* \*  $Q^2=10$
- 4- Torino** *Anselmino et al., P.R.D87 (13) 094019* \*  $Q^2=1$
- 5- JAM fit** *Lin et al., P.R.L.120 (18) 152502* \*  $Q_0^2=2$
- 6- PNDME16** *Bhattacharya et al., P.R.D94 (16) 054508*
- 7- PNDME18** *Gupta et al., arXiv:1808.07597*
- 8- ETMC17** *Alexandrou et al., P.R.D95 (17) 114514;*  
*E P.R.D96 (17) 099906*

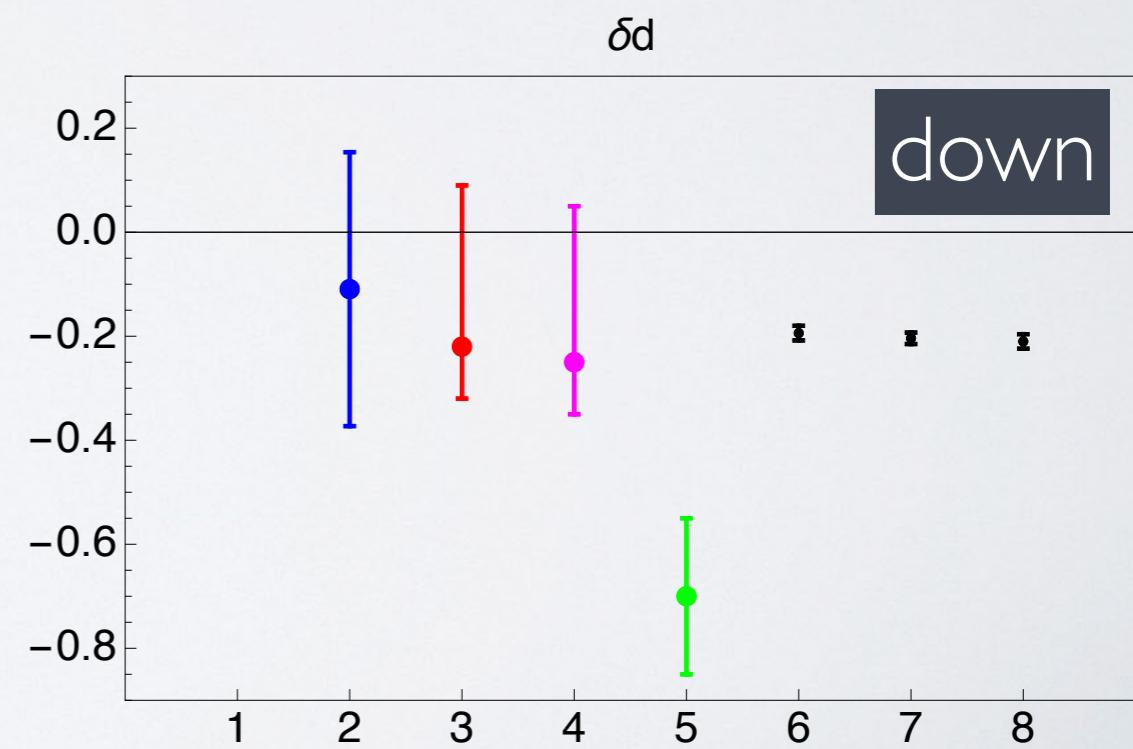


# tensor charge : separate flavors



incompatibility for up  
compatibility for down  
but within large errors  
(except JAM)

- 2- global fit** *Radici & Bacchetta, P.R.L.120 (18) 192001*
- 3- TMD fit** *Kang et al., P.R.D93 (16) 014009* \*  $Q^2=10$
- 4- Torino** *Anselmino et al., P.R.D87 (13) 094019* \*  $Q^2=1$
- 5- JAM fit** *Lin et al., P.R.L.120 (18) 152502* \*  $Q_0^2=2$
- 6- PNDME16** *Bhattacharya et al., P.R.D94 (16) 054508*
- 7- PNDME18** *Gupta et al., arXiv:1808.07597*
- 8- ETMC17** *Alexandrou et al., P.R.D95 (17) 114514;*  
*E P.R.D96 (17) 099906*

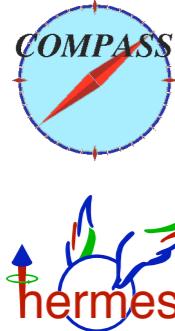
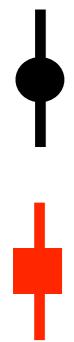


# “transverse-spin puzzle” ?

there seems to be no simultaneous compatibility  
about  $\delta u$ ,  $\delta d$ ,  $g_T = \delta u - \delta d$   
between lattice and  
phenomenological extractions  
of transversity

# results

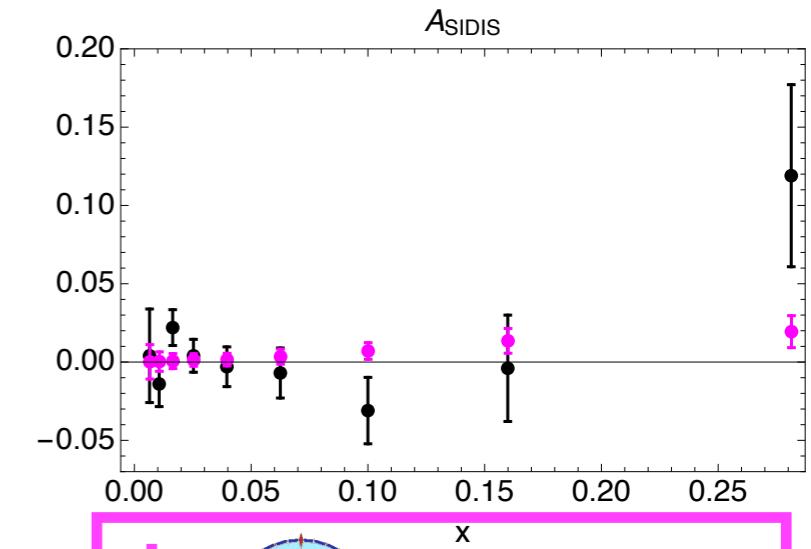
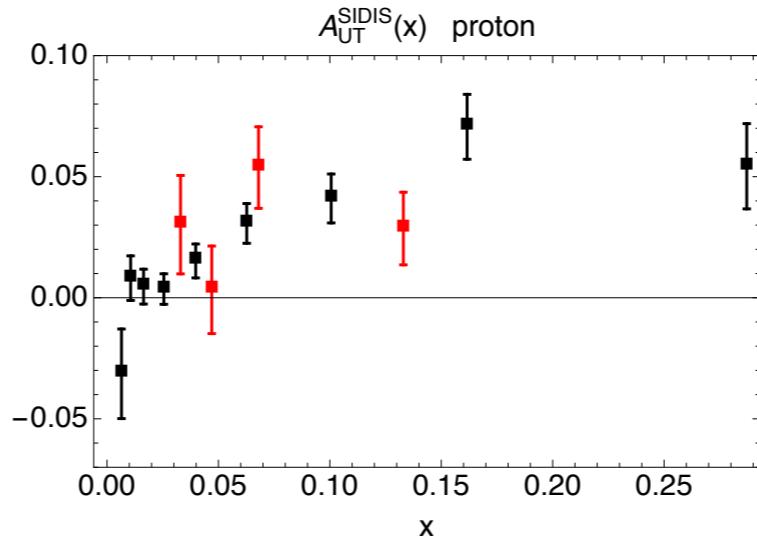
## SIDIS



*Adolph et al., P.L. B713 (12)*



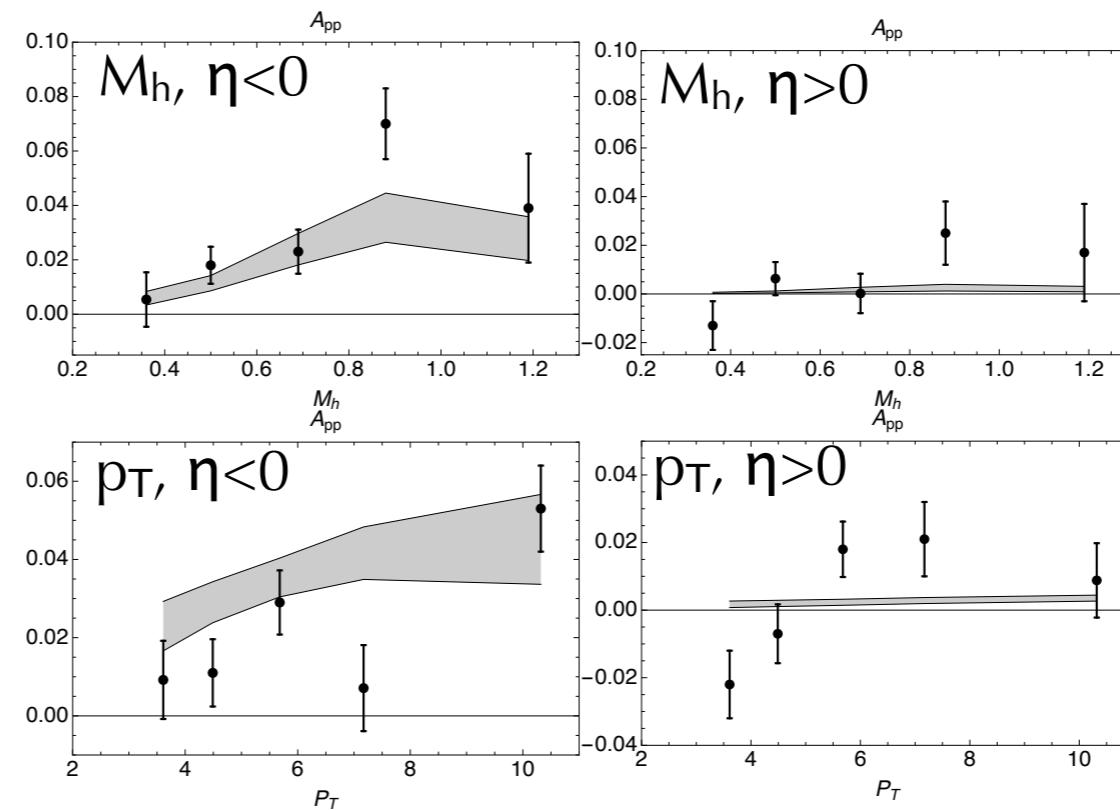
*Airapetian et al.,  
JHEP 0806 (08) 017*



## pp collisions



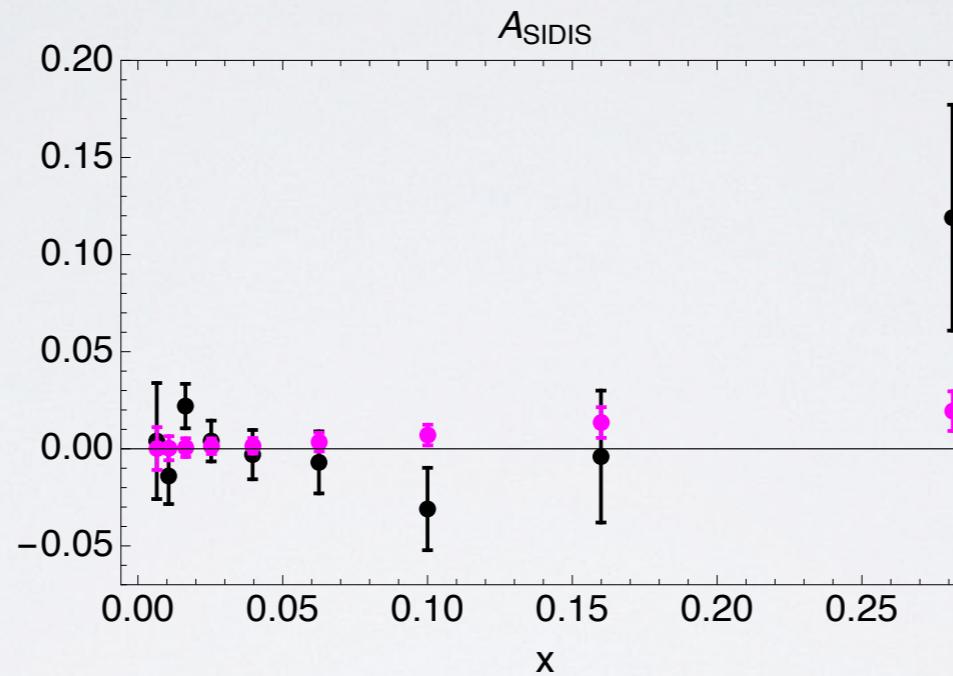
*Adamczyk et al.,  
P.R.L. 115 (2015) 242501*



# Compass pseudodata



Adolph et al., P.L. **B713** (12)



pseudodata

see next talk by F. Bradamante

statistical error  $\sim 0.6 \times$  [ error in 2010 proton run ]

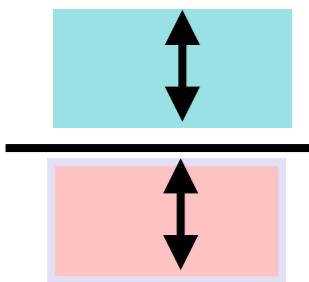
$\langle A \rangle$  = average value of replicas in previous global fit

# impact of pseudodata

global fit + pseudodata

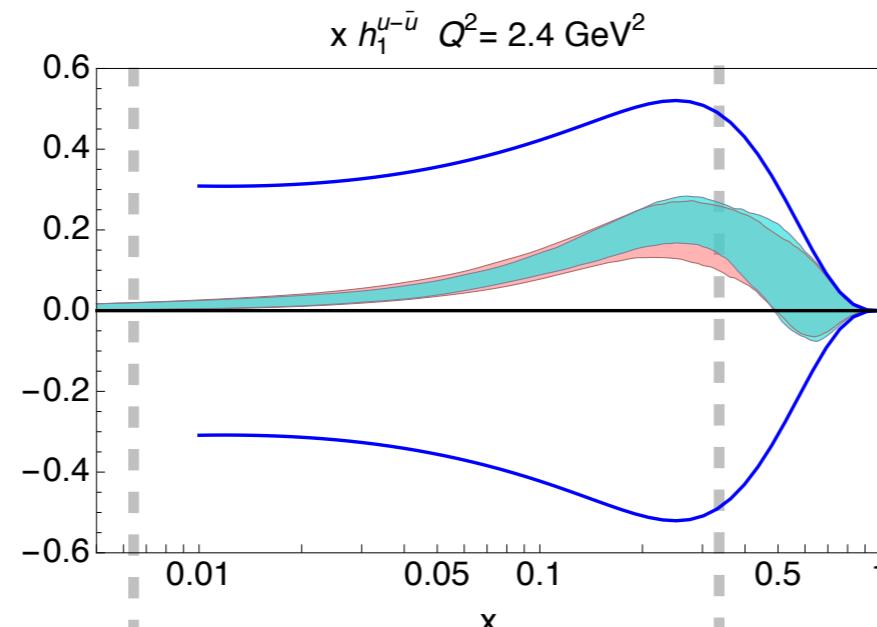
global fit

$$D_{1g}(Q_0) = \begin{cases} 0 \\ D_{1^u}/4 \\ D_{1^u} \end{cases}$$

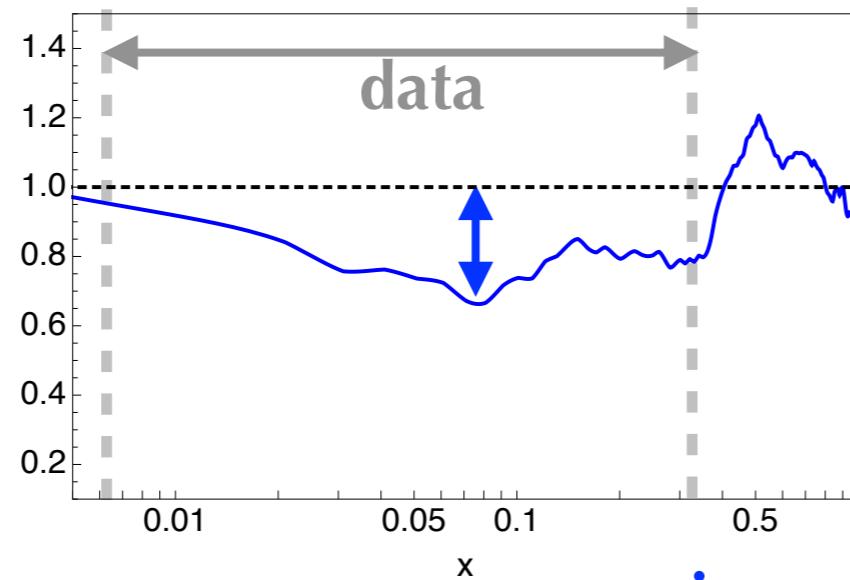
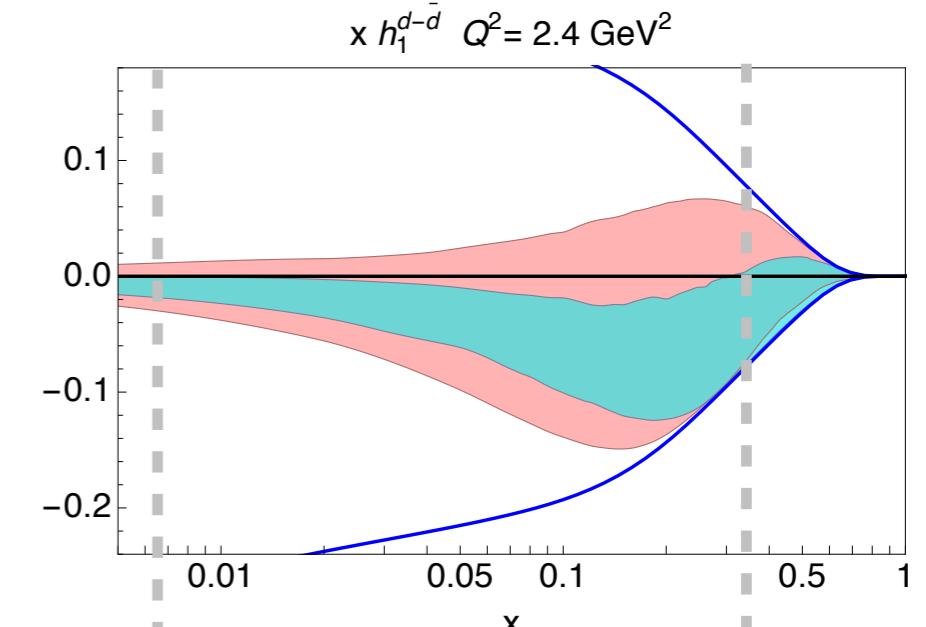


ratio of  
widths

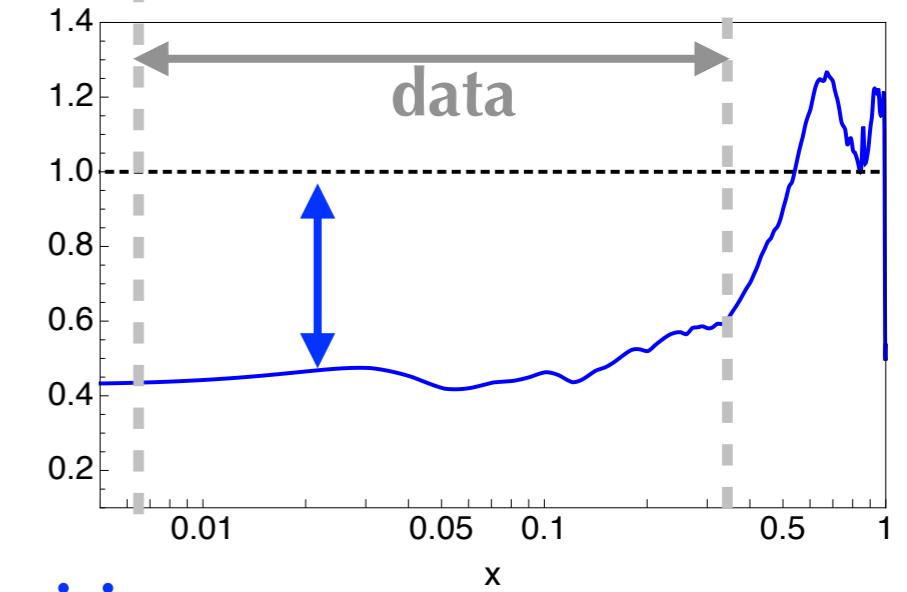
up



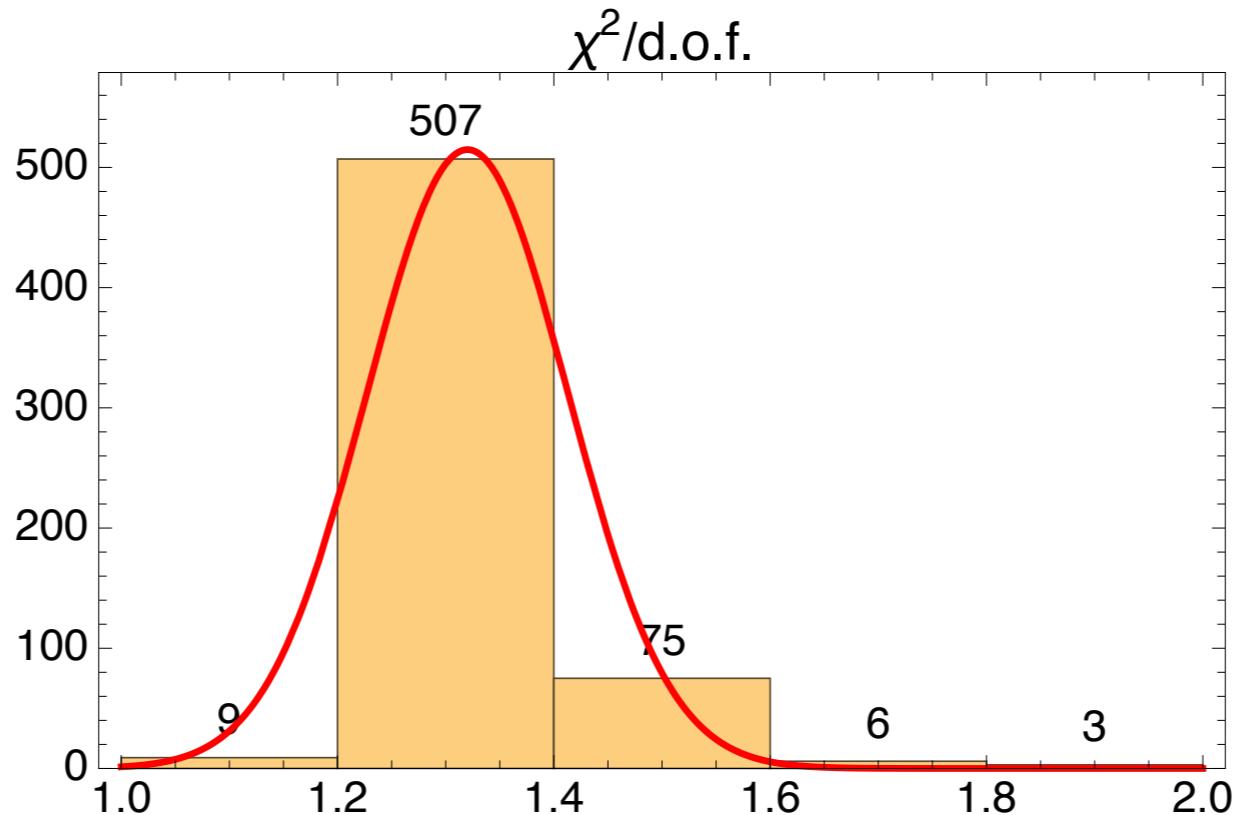
down



increase precision

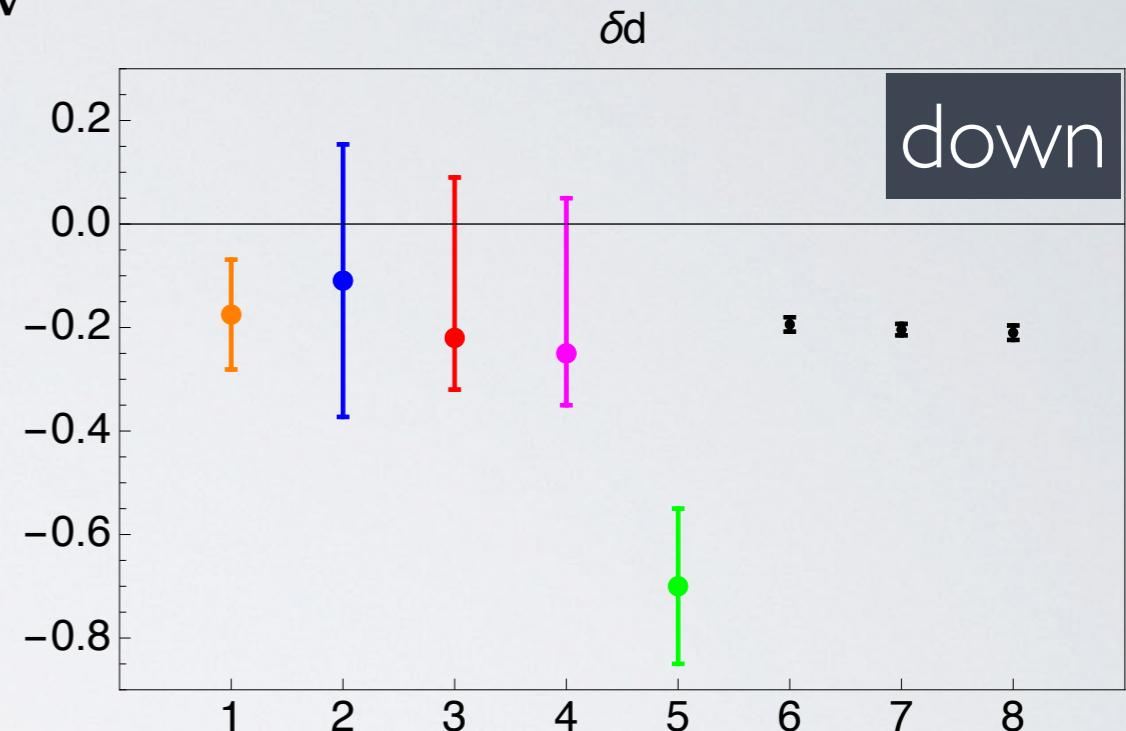
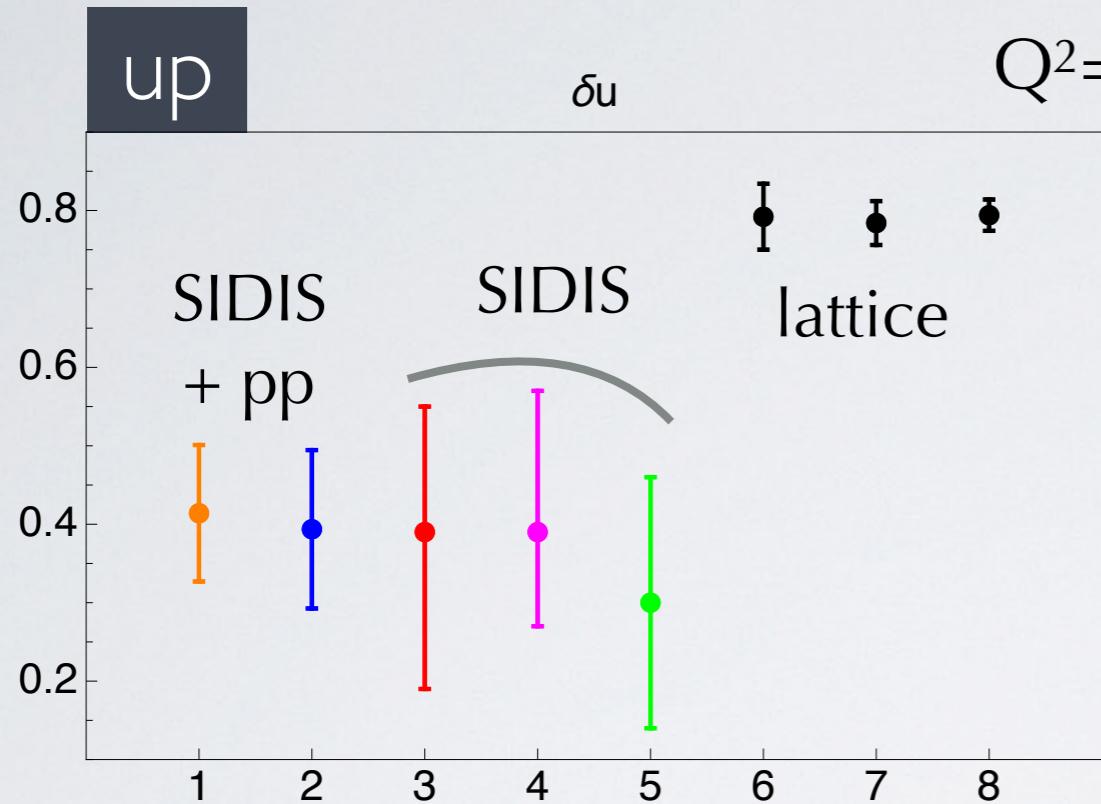


# better $\chi^2$



$$\chi^2/\text{dof} = 1.32 \pm 0.09$$

# impact of pseudodata



1- global fit + pseudodata

2- global fit *Radici & Bacchetta, P.R.L. 120 (18) 192001*

3- TMD fit *Kang et al., P.R. D93 (16) 014009* \*  $Q^2=10$

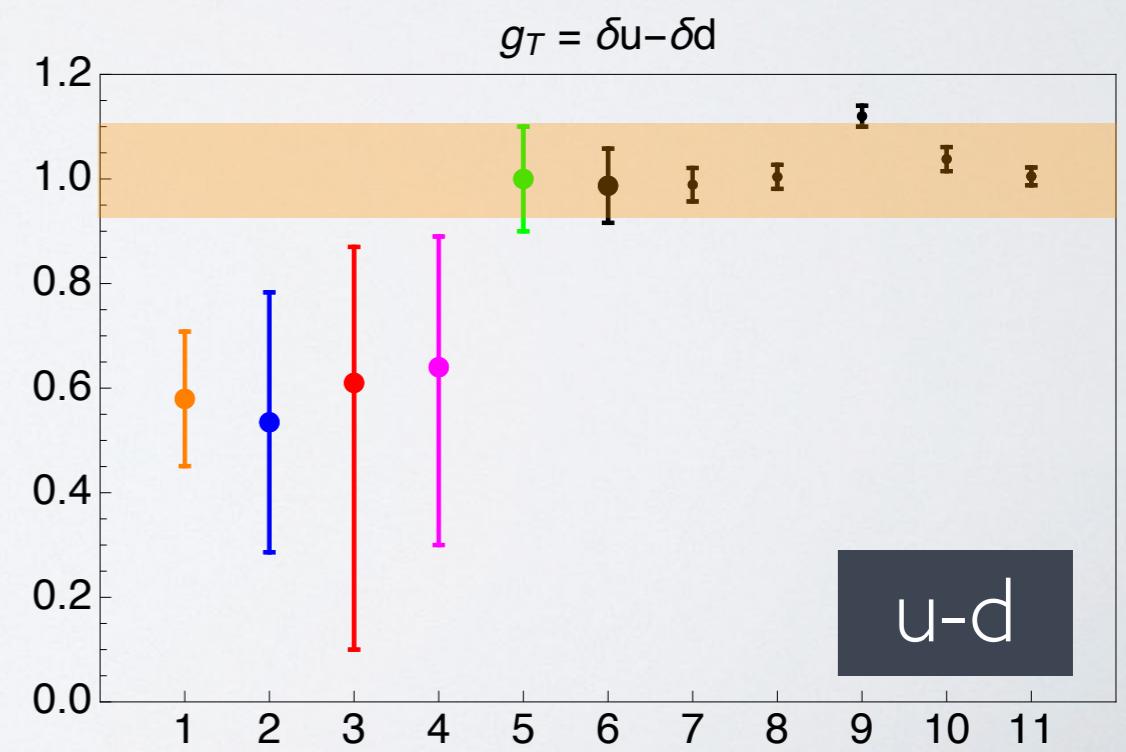
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5- JAM fit *Lin et al., P.R.L. 120 (18) 152502* \*  $Q_0^2=2$

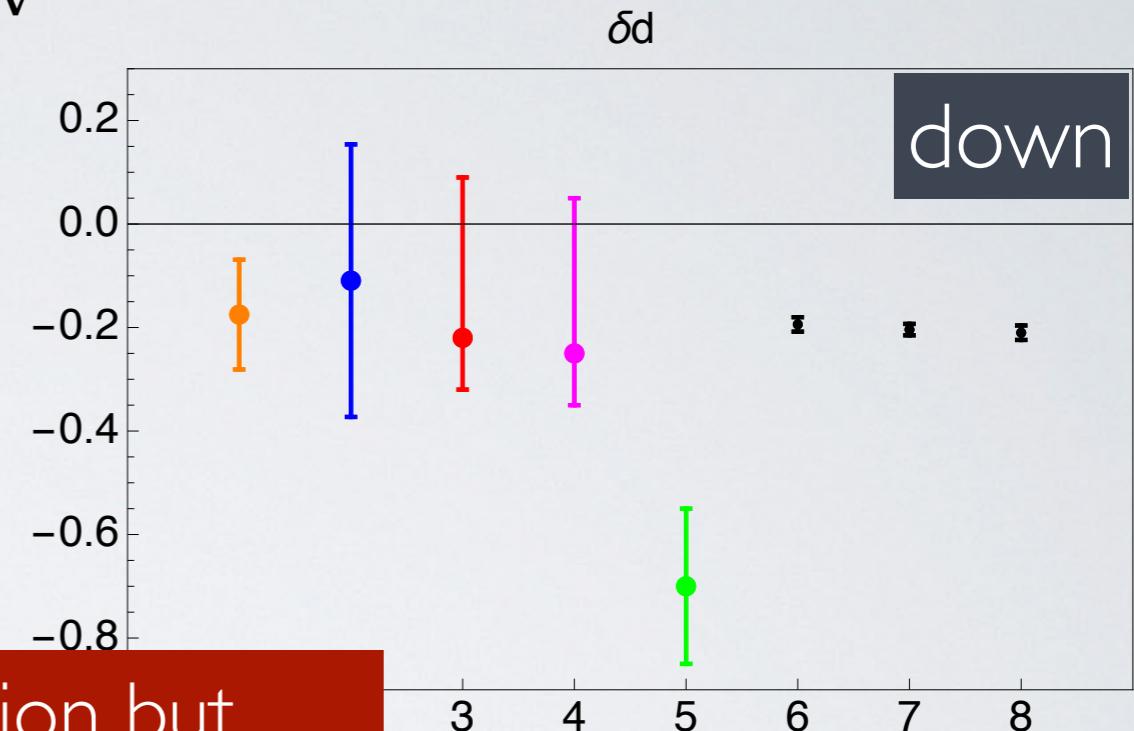
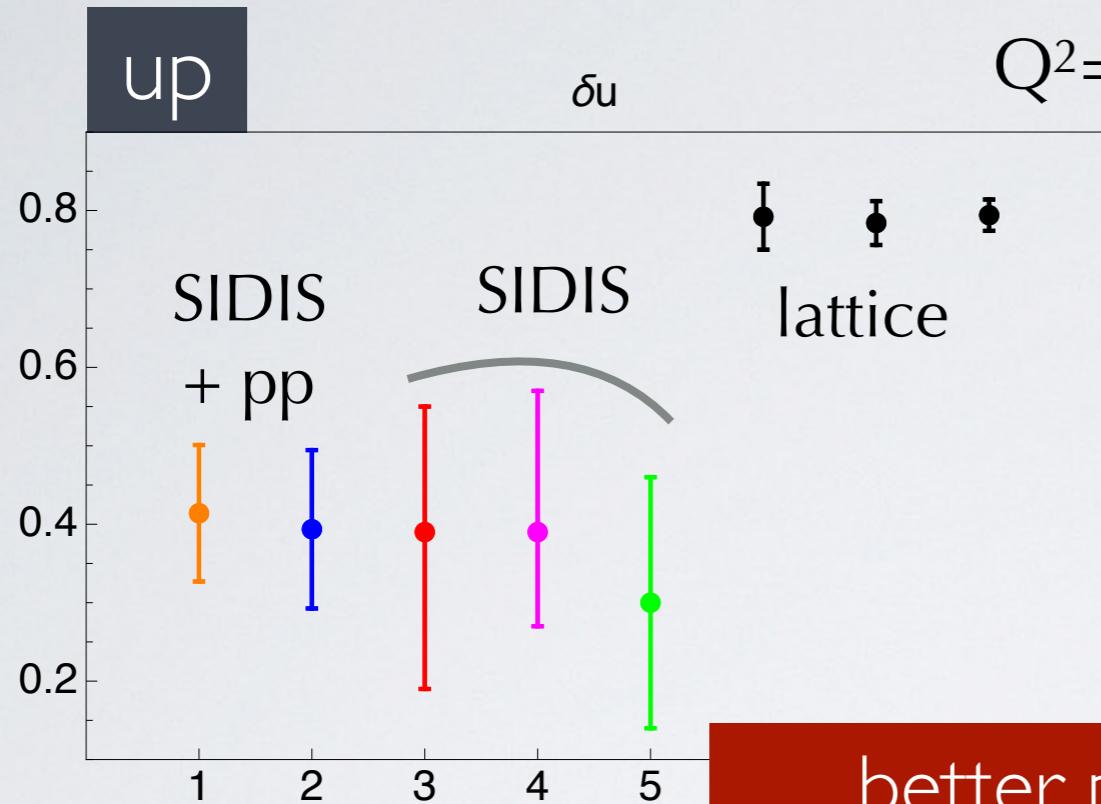
6- PNDME16 *Bhattacharya et al., P.R. D94 (16) 054508*

7- PNDME18 *Gupta et al., arXiv:1808.07597*

8- ETMC17 *Alexandrou et al., P.R. D95 (17) 114514;*  
*E P.R. D96 (17) 099906*



# impact of pseudodata



better precision but  
confirm general trend

1- global fit + pseudodata

2- global fit *Radici & Bacchetta, P.R.L. 120 (18) 192001*

3- TMD fit *Kang et al., P.R. D93 (16) 014009* \*  $Q^2=10$

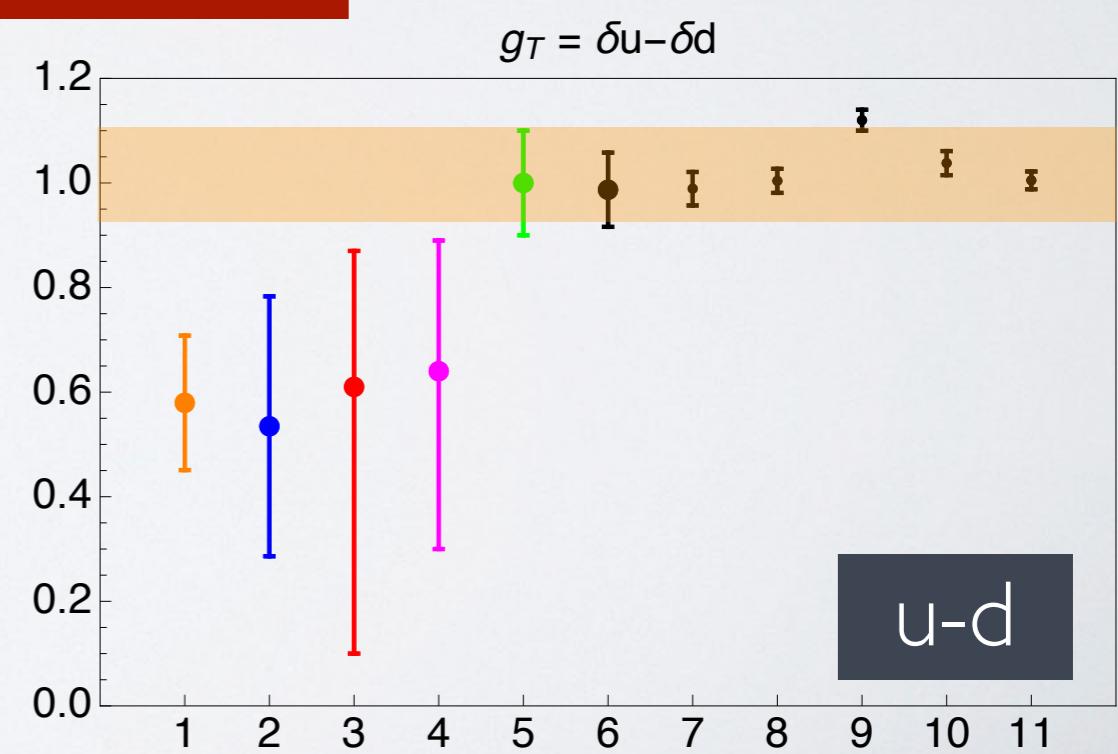
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6- PNDME16 *Bhattacharya et al., P.R. D94 (16) 054508*

7- PNDME18 *Gupta et al., arXiv:1808.07597*

8- ETMC17 *Alexandrou et al., P.R. D95 (17) 114514;  
E P.R. D96 (17) 099906*



# Conclusions

- first global fit of di-hadron inclusive data leading to extraction of transversity as a PDF in collinear framework
- inclusion of STAR p-p<sup>↑</sup> data increases precision of up channel; large uncertainty on down due to unconstrained gluon unpolarized di-hadron fragmentation function
- no apparent simultaneous compatibility with lattice for tensor charge in up, down, and isovector channels
- adding Compass pseudodata for deuteron increases precision, particularly for down, but seems to confirm this scenario
- are lattice moments really compatible with di-hadron inclusive small/large-x data ? What's going on at very small x ?

**THANK YOU**

# Back-up

# 2-hadron-inclusive production

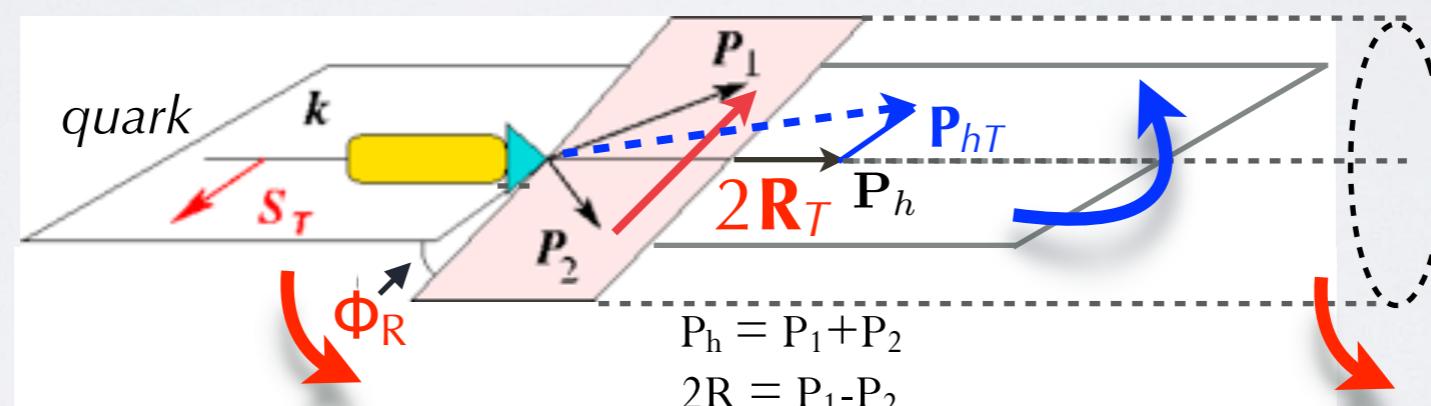
framework  
collinear  
factorization

Collins, Heppelman, Ladinsky,  
N.P. **B420** (94)

$$R_T \ll Q \quad H_1^{\triangleleft}$$

$\updownarrow$

invariant mass

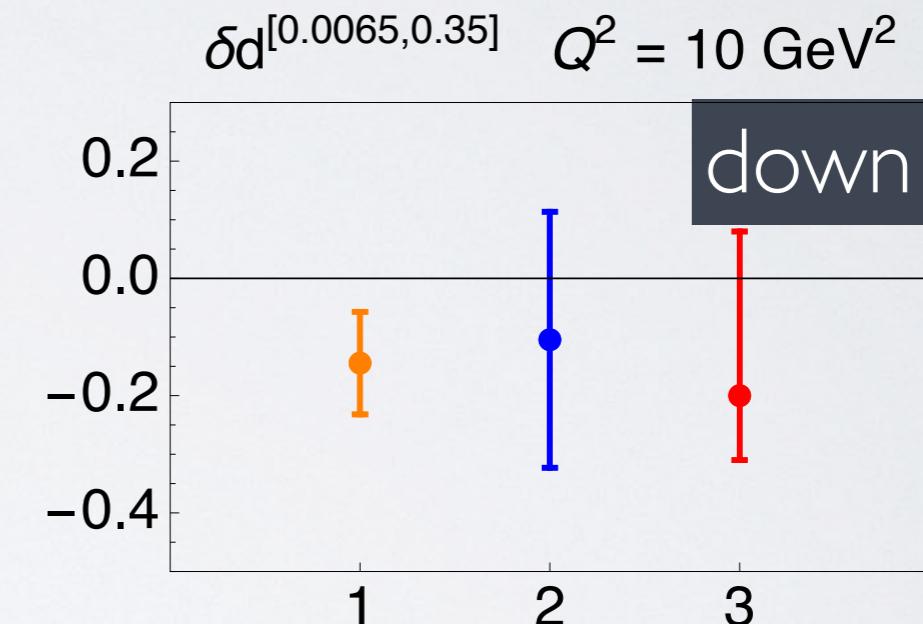
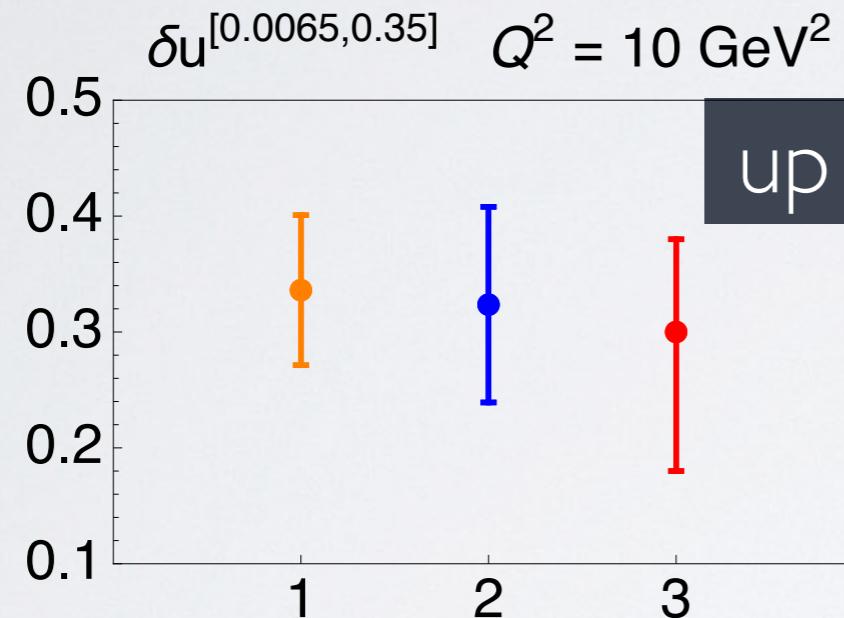


survives to  
polar  
symmetry  
(  $\int dP_{hT}$  )

correlation  $S_T$  and  $R_T \rightarrow$  azimuthal asymmetry

tensor charge  $\delta q(Q^2) = \int dx h_1 q\bar{q} (x, Q^2)$

truncated  
 $\delta q^{[0.0065, 0.35]} \quad Q^2 = 10$



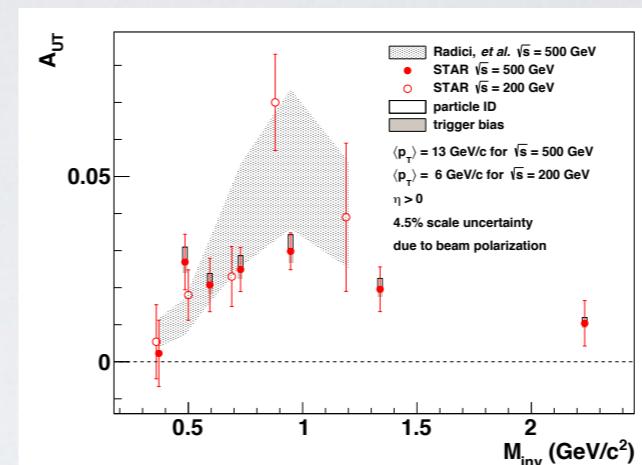
+  
pseudodata

global fit  
*Radici & Bacchetta,  
P.R.L. **120** (18) 192001*

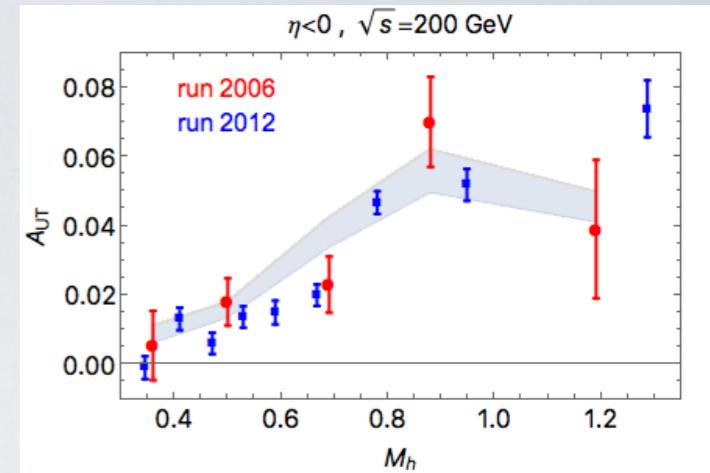
TMD fit  
*Kang et al.,  
P.R. D93 (16) 014009*

# To do list

- use also other (multi-dimensional) data from STAR run 2011 ( $s=500$ ) and (later) run 2012 ( $s=200$ )



Adamczyk et al. (STAR), P.L. **B780** (18) 332



Radici et al., P.R. **D94** (16) 034012

- need data on  $p+p \rightarrow (\pi\pi) X$  constrains gluon  $D_{1g}$
- refit di-hadron fragmentation functions using new data:  
 $e^+e^- \rightarrow (\pi\pi) X$  constrains  $D_{1q}$   
 (currently only by Montecarlo)
- use COMPASS data on  $\pi K$  and  $KK$  channels, and from  $\Lambda^\uparrow$  fragmentation:  
 constrain strange contribution ?
- explore other channels, like inclusive DIS via Jet fragm. funct.'s



Seidl et al.,  
P.R. **D96** (17) 032005

# more constraints on extrapolation

