Quark-gluon correlations in the twist-3 TMD using light-front wave functions

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European Research Council





* Twist-3: a particular TMD



***** Twist-3: a particular TMD

Decomposition at twist-3 and singular terms

OUTLINE

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- * A particular model calculation

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Lorcé, Pasquini, Schweitzer JHEP01(2015)103

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 $\overline{\psi}(0)\mathcal{W}(0;\xi)\psi(\xi) = \overline{\psi}_{+}(0)\mathcal{W}(0;\xi)\psi_{-}(\xi) + \overline{\psi}_{-}(0)\mathcal{W}(0;\xi)\psi_{+}(\xi)$

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$$\mathcal{W}(0;\xi)\,\psi_{-}(\xi))|_{\xi^{+}=0} = \mathcal{W}_{1}(0^{-},\mathbf{0}_{\perp};0^{-},\boldsymbol{\xi}_{\perp})\psi_{-}(0^{+},0^{-},\boldsymbol{\xi}_{\perp})$$

$$-i\int_{0^{-}}^{\xi^{-}} d\zeta^{-} \mathcal{W}_{1}(0^{-}, \mathbf{0}_{\perp}; \zeta^{-}, \boldsymbol{\xi}_{\perp}) \frac{\gamma^{+}}{2} \left(i\boldsymbol{\gamma}_{\perp} \cdot \boldsymbol{D}_{\perp} + m\right) \psi_{+}(0^{+}, \zeta^{-}, \boldsymbol{\xi}_{\perp})$$

$$e = e_{\text{sing.}} + e_{\text{mass}} + e_{\text{tw3}}$$

For the PDF: Efremov and Schweitzer JHEP 0308(2003)006

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$$\begin{split} e_{tw3} &= -\frac{P^+}{M} \frac{g_s}{2} \int \frac{d\xi^- d\boldsymbol{\xi}_\perp}{2(2\pi)^3} e^{ik^+ \xi^- - \boldsymbol{k}_\perp \cdot \boldsymbol{\xi}_\perp} \\ &\times \left(\int_{0^-}^{\xi^-} d\zeta^- \int_{\infty^-}^{\zeta^-} d\eta^- \langle P | \overline{\psi}(0) \mathcal{W}_1(0^-, \mathbf{0}_\perp; \eta^-, \boldsymbol{\xi}_\perp) G^+_j(0^+, \eta^-, \boldsymbol{\xi}_\perp) \right) \\ &\times \sigma^{j+} \mathcal{W}_1(\eta^-, \boldsymbol{\xi}_\perp; \zeta^-, \boldsymbol{\xi}_\perp) \psi(0^+, \zeta^-, \boldsymbol{\xi}_\perp) | P \rangle \\ &+ \int_{0^-}^{\xi^-} d\zeta^- \int_{0^-}^{\infty^-} d\eta^- \langle P | \overline{\psi}(0) \mathcal{W}_1(0^-, \mathbf{0}_\perp; \eta^-, \mathbf{0}_\perp) G^+_j(0^+, \eta^-, \mathbf{0}_\perp) \\ &\times \sigma^{j+} \mathcal{W}_1(\eta^-, \mathbf{0}_\perp; \zeta^-, \boldsymbol{\xi}_\perp) \psi(0^+, \zeta^-, \boldsymbol{\xi}_\perp) | P \rangle \end{split}$$

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In light-cone gauge

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$$\begin{aligned} e(x, \boldsymbol{k}_{\perp}) &= \frac{\delta(x)}{2M} \int \frac{d\boldsymbol{\xi}_{\perp}}{(2\pi)^2} e^{-i\boldsymbol{k}_{\perp} \cdot \boldsymbol{\xi}_{\perp}} \langle P | \overline{\psi}(0) \mathcal{W}(0; 0^+, 0^-, \boldsymbol{\xi}_{\perp}) \psi(0^+, 0^-, \boldsymbol{\xi}_{\perp}) | P \rangle \\ &+ \tilde{e}(x, \boldsymbol{k}_{\perp}) + \frac{m}{xM} f_1(x, \boldsymbol{k}_{\perp}) \\ &- \delta(x) \int_{-1}^1 dy \left(\tilde{e}(y, \boldsymbol{k}_{\perp}) + \frac{m}{yM} f_1(y, \boldsymbol{k}_{\perp}) \right) \end{aligned}$$

In light-cone gauge







$$\int dx d\mathbf{k}_{\perp} \ e(x, \mathbf{k}_{\perp}) = \frac{1}{2M} \langle P | \overline{\psi}(0) \psi(0) | P \rangle$$

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Brodsky, Pauli Phys.Rept.301(1998)299
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LFWFs depends on all the partons' quantum numbers.

$$\psi^{1} , \ \psi^{2} \ \ L_{z} = 0 \qquad \psi^{5} \ \ L_{z} = -1$$

 $\psi^{3} , \ \psi^{4} \ \ L_{z} = 1 \qquad \psi^{6} \ \ L_{z} = 2$

Ji, Ma, Yuan NPB652(2003)383

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3 independent (& leading) LFWAs for $L_z=0$ of 3q+g

Braun et al. PRD83(2011)094023

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$$\begin{split} \tilde{e}^{u}(x,k_{\perp}) &= \frac{4g_{s}}{Mx\sqrt{3}} \int \frac{[dx]_{1234}[dk]_{1234}}{\sqrt{x_{4}}} \Re \Biggl\{ -\delta(x-x_{4}-x_{1})\delta^{2}(\mathbf{k}_{\perp}-\mathbf{k}_{\perp4}-\mathbf{k}_{\perp1}) \\ &\times \left[\left(4\Psi^{(0)*}(-2-3,3,2) + 2\Psi^{(0)*}(2,3,-2-3) \right) \Psi^{1\uparrow}(1,2,3,4) \right. \\ &+ \left(2\Psi^{(0)*}(-2-3,2,3) + \Psi^{(0)*}(3,2,-2-3) \right) \Psi^{2\uparrow}(1,2,3,4) \\ &- 2\Psi^{(0)*}(2,-2-3,3)\Psi^{\downarrow}(1,2,3,4) \Biggr] + \delta(x-x_{4}-x_{2})\delta^{2}(\mathbf{k}_{\perp}-\mathbf{k}_{\perp4}-\mathbf{k}_{\perp2}) \\ &\times \left[2\Psi^{(0)*}(-1-3,1,3)\Psi^{2\uparrow}(1,2,3,4) - \Psi^{(0)*}(1,-1-3,3)\Psi^{\downarrow}(1,2,3,4) \Biggr] \\ &+ \delta(x-x_{4}-x_{3})\delta^{2}(\mathbf{k}_{\perp}-\mathbf{k}_{\perp4}-\mathbf{k}_{\perp3})\Psi^{(0)*}(-1-2,1,2)\Psi^{1\uparrow}(1,2,3,4) \Biggr\} \end{split}$$

Fock state expansion up to the 3q+g state (neglecting the anti-quarks)

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For each Fock state only leading function of L=0 are considered

An explicit form for the LFWFs is required

$$\Psi_N^{L=0} = \phi(x_1, \dots, x_N) \Omega_N(\{x_i\}, \{k_{\perp,i}\}, \{a_i\})$$

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From connection between LFWF and distribution amplitudes and by conformal expansion

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$$\Omega_N\left(\{x_i\}, \{k_{\perp,i}\}, \{a_i\}\right) = \mathcal{N}e^{-\sum_{i=1}^N a_i \frac{k_{\perp,i}^2 + m_i^2}{x_i}}$$

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 $\Omega_N(\{x_i\}, \{k_{\perp,i}\}, \{a_i\}) = \mathcal{N}e^{-\sum_{i=1}^N a_i \frac{k_{\perp,i}^2 + m_i^2}{x_i}}$ $a_i = a_N \quad \forall i = 1, \dots, N$





 $Q^2 = 1 \text{ GeV}^2$















Equation of motion to decompose the Twist-3 distributions



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Fock state expansion and model for LFWFs



- Equation of motion to decompose the Twist-3 distributions
- Fock state expansion and model for LFWFs
- The pure Twist-3 contribution is not-negligile, at least in the model

BACKBUP
Model

$$\begin{split} \phi &= 120 f_N x_1 x_2 x_3 \left(1 + A(x_1 - x_3) + B(x_1 + x_3 - 2x_2) \right) \\ \phi^{\downarrow} &= -\frac{M}{96g_s} \frac{8!}{2} \lambda_1^g x_1 x_2 x_3 x_4^2 \\ \psi &= -\frac{M}{96g_s} \frac{8!}{4} \left(\lambda_2^g + \lambda_3^g \right) x_1 x_2 x_3 x_4^2 \\ \theta &= -\frac{M}{96g_s} \frac{8!}{4} \left(\lambda_2^g - \lambda_3^g \right) x_1 x_2 x_3 x_4^2 \end{split}$$
PARAMETERS

$$f_N, \quad A, \quad B, \quad a_3,$$

 $\lambda_1, \quad \lambda_2, \quad \lambda_2, \quad \frac{a_4}{a_3}$

 $m_q, m_g,$

$$F_{\mathrm{UU}}^{\cos\phi_{h}} = \frac{2M}{Q} \mathcal{C} \left[\dots - \frac{P_{h} \cdot \boldsymbol{k}_{\perp}}{|\boldsymbol{P}_{h}|M} \left(xf^{\perp}D_{1} + \frac{M_{h}}{M}h_{1}^{\perp}\frac{\tilde{H}}{z} \right) \right]$$
$$F_{\mathrm{LU}}^{\sin\phi_{h}} = \frac{2M}{Q} \mathcal{C} \left[\dots - \frac{P_{h} \cdot \boldsymbol{p}_{\perp}}{|\boldsymbol{P}_{h}|M} \left(xeH_{1}^{\perp} + \frac{M_{h}}{M}f_{1}\frac{\tilde{G}^{\perp}}{z} \right) \right]$$

SIDIS

