Laser preparation of intense beams <u>&</u> <u>targets</u> of vectorially polarized protons and tensorially polarized deuterons via molecular quantum beats. FORTH Institute of Electronic Structure and Laser



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- Hyperpolarization our approach
 - Physical processes
 - Technological aspects

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- Hyperpolarization our approach
 - Physical processes
 - Molecular Photodissociation
 - Hyperfine polarization quantum beats
 - Technological aspects
 - Detection of spin polarization
 - Storage in B fields
 - Status of high power (table-top) laser technology
 - Why time is ripe?

- Motivation:
 - Nuclear/particle physics
 - Reaction dynamics
 - NMR/MRI





- Motivation:
 - Nuclear/particle physics
 - Reaction dynamics
 - NMR/MRI
 - Polarized Fusion
 - Wednesday 12 September A11: Prof. RAKITZIS, T. Peter

• Some vocabulary: Polarization





- Laser polarization (orientation/alignment) of nuclei?
 - Via the electron
 - Two-step process
 - Polarize the electron laser dissociation
 - Polarization quantum beating
 - Transfer polarization to the nuclei

- Polarize the electron: Photodissociation
- R. J. van Brunt and R. N. Zare, J. Chem. Phys. 48, 4304 (1968)
 - Angular momentum projection \rightarrow conserved during the dissociation process
 - Circularly polarized light \rightarrow increase angular momentum projection (of Ω) by 1



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• In a <u>prompt</u> dissociation process this increase can be found in the fragments





- Polarize the nucleus: Polarization beating
- Symmetric double well
 - Supports two eigenstates: S & A
 - Electron R (L) : S ± A
 - If $E_S \neq E_A \rightarrow$ time dependence: Right \leftrightarrow Left

$$|S\rangle = \frac{1}{\sqrt{2}} (|S\rangle + |A\rangle) \xrightarrow{\text{Time}} |R\rangle = \frac{1}{\sqrt{2}} (e^{-iE_{S}t/\hbar} |S\rangle + e^{-iE_{A}t/\hbar} |A\rangle) \qquad S + I \rightarrow F$$

$$|L\rangle = \frac{1}{\sqrt{2}} (|S\rangle - |A\rangle) \xrightarrow{\text{Time}} |L\rangle = \frac{1}{\sqrt{2}} (e^{-iE_{S}t/\hbar} |S\rangle - e^{-iE_{A}t/\hbar} |A\rangle) \qquad \frac{|1-1\rangle}{|00\rangle} \xrightarrow{|1-1\rangle}$$

$$t = \pi\hbar/\Delta E \qquad 100$$

- Photodissociation does not resolve hyperfine levels
- Parallel/antiparallel = superposition F states

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 $\begin{aligned} \uparrow \downarrow &= \frac{1}{\sqrt{2}} \left(|10\rangle + |00\rangle \right) \\ \downarrow \uparrow &= \frac{1}{\sqrt{2}} \left(|10\rangle - |00\rangle \right) \end{aligned}$

Electronic Structure



- Very attractive combination timescales
- (prompt) Photodissociation ~ 100 fs
- Polarization beating (H, D) ~ ns
- State-of-the-art: SEOP and Stern Gerlach $\sim \mu s$
- Dissociation projected 10¹⁷ cm⁻³ with 10 ns (alkali-H depolarization cross sections)







- Very attractive combination timescales
- Needs to be proved!
 - Means detecting the polarization
 - Detect electron (nucleus) polarization in high density





- High density SPD and SPH
 - Milner et al., Phys. Rev. Lett. **118**, 243201 (2017)
 - Use an Optical centrifuge to polarize O₂ molecules
 - Detect the electron polarization via the EMF induced in a pick-up coil



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- High density SPD and SPH
 - Milner et al., Phys. Rev. Lett. **118**, 243201 (2017)
 - Detect the electron polarization via the EMF induced in a pick-up coil
 - Sudden polarization : create magnetization
 - Depolarization : exponential decay
 - EMF proportional to dP/dt





- High density SPD and SPH
 - If polarization exchange to nucleus
 - Magnetization vanishes (μ_B/μ_N ~1836)







- High density SPD and SPH
 - Sofikitis et al., Phys. Rev. Lett. **121**, 083001 (2018)
 - We project signals ~ 100 μ V with densities 10¹⁸ cm⁻³
 - However, UV noise ~ mV
 - PEM-Signal reversal
 - Dissociation : RCP/LCP shot-to-shot basis
 - Accumulate RCP and LCP separately



- High density SPD and SPH
 - Sofikitis et al., Phys. Rev. Lett. 121, 083001 (2018)
 - EMF signals:
 - Sharp rising edge ~ 150 ps photolysis
 - Exponential decay depolarization
 - ns polarization exchange
 - Similar for H from HBr
 - FFT \rightarrow Hyperfine frequencies



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- High density SPD and SPH
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 - Vary laser intensity multiphoton contributions limited
 - Deplete the signal densities up to 10¹⁹ cm⁻³
 - Simulate for 100 mJ
 - Laser fusion at NIF





PULSED TARGET – INTENSE PULSE BEAMS

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 - 2 MJ
 - 2.3 ns
 - @350 nm
 - Focus at 100 μm
 - For 10^{18} cm⁻³ \rightarrow 10^7 (10⁵) neutrons/pulse for DT (DD)
 - Sofikitis et al., PRL **118** 233401 (2017)
 - State-of-the-art (magnetic): ~0.01 neutrons/s
 - H. Paetz gen. Schieck Eur. Phys. J. A 44, 321–354 (2010)



- Obviously simple photodissociation cannot do it
 - Electron is $s = \frac{1}{2} \rightarrow no alignment$
- Can we find an 'alignment reservoir' linked with nuclei?
- Molecular rotation







- Molecular rotation
- Aligned in lab frame
 - IR (coherent) excitation
 - (IR pumping)
- HF beating
 - Transfer alignment to I
 - And back
 - microseconds

Τ ~ μs











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- Polarization transfer from molecular rotation to nuclear spin
- A generic experiment:
- (molecular beams)
- Four stages:





- Polarization transfer from molecular rotation to nuclear spin
- A generic experiment:
- (molecular beams)
- Four stages:



Storage?

- Polarization transfer from molecular rotation to nuclear spin
 - Storage in Hight density molecules
 - Coupling to J sensitive to collisions
 - Storage in B fields
 - Decoupling of hyperfine interaction
 - B(t, /) is not a step function
 - Make it as close as possible



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Polarization transfer from molecular rotation to nuclear spin

Storage?

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- Example H₂
- Timing IR B field can be critical



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Polarization transfer from molecular rotation to nuclear spin

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Storage?

- Polarization transfer from molecular rotation to nuclear spin
 - Storage in B fields
 - Decoupling of hyperfine interaction
 - B(t, /) is not a step function
 - Make it as close as possible
 - Example H₂
 - Timing IR B field can be critical
 - Experimentally adjust timing
 - v ~ 1000 m/s \rightarrow mm accuracy



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D: Numbers

- Polarization transfer from molecular rotation to nuclear spin
 - Tunable lasers IR (pumping) P up to 5 W ($\lambda \sim \mu m 10^{19}$ ph/s)
 - Fixed λ lasers (dissociation) P up to 200 W (UV)
 - Fiber laser technology





D: Numbers

- Polarization transfer from molecular rotation to nuclear spin
 - For an IR power ~ 5 W N_{ph} ~ 5x10¹⁹ photons/s
 - $D_{MB} > 10^{14} \text{ cm}^{-3} \rightarrow \text{limited by the number of photons}$
 - IR pumped molecules $> 10^{19}$ mol/s
 - $\sigma_{diss} \simeq 10^{-18} 10^{-17} \text{ cm}^2 \rightarrow P_{diss} \simeq 10^{-2}$



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Conclusions:

- Laser-molecule interactions can provide *oriented* and *aligned* H, D
- IR pumping and/or photodissociation
 - Demonstration of 10¹⁹ SPD/cm³ (orientation)
 - Alignment from IR pumped molecules
 - High power lasers high SPD flux (up to 10²⁰ SPD/s)
 - Orientation AND alignment
 - Storage in B fields

ectronic Structure

Stereodynamic effects in the dissociation of small molecules

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- Collaborators
 - T. P. Rakitzis IESL FORTH, UOC Physics
 - H. Kannis
 - G. Boulogiannis
 - Luis Rubio Lago









THANK YOU!!!

Points

- Pulsed Targets 10¹⁹ 10¹⁶ per shot- Pulsed intense beams
- Talk about particles per second
- Tune B field in cm

bitor, ability one 2 mil miller [10].

For the D-T or D-³He nuclear fusion reactions, the angular distribution of the neutron or proton products $D(\theta, \phi)$ about the quantization axis, as a function of the nuclear vector polarizations p_z of D, T, or ³He is well approximated by [25]:

$$D(\theta, \phi) = \frac{\sigma_0}{3} [(2+p) - (2p + p_{zz})P_2(\cos\theta)]/4\pi \quad (10)$$

where $p = p_z(D)\overline{p_z}(Y)$, Y = T or ³He, p_z is the nuclear vector polarization, p_{zz} is the tensor polarization for D nuclei, σ_0 is the fusion cross section through the intermediate ⁵He or ⁵Li $\frac{3}{2}^+$ state for the D-T and D-³He reactions, respectively, and $P_2(x)$ is the 2nd Legendre Polynomial. The first term in Eq. (10) is proportional to the integrated product signal, so that for maximal nuclear polarization, with p = 1, the product integrated intensity is increased by 50% compared to p = 0; also, in this case $p_{zz} = 1$, and hence $D(\theta, \phi) \sim 1 - P_2(\cos\theta) \sim \sin^2\theta$.

For the D-³He reaction performed with $p_z({}^{3}\text{He}) = 0.76 \ [26, 27], \ p_z(D, t = 1.5 \, ns) = 0.12$ (reported here), and $p_{zz}(D) = 0$, we predict a 14% variation in the angular distribution (between $\theta = 0^{\circ}$ and 90°) and a 4.5% increase in the integrated intensity, whereas if bond alignment is used prior to dissociation [28], $p_z(D, t = 1.5 \, ns) = 0.5$ can be produced [8], leading to a 70% variation in the product angular distribution and a 19% increase in integrated intensity. In contrast, the effect of

$$\langle m_{\mathcal{S}}(\mathbf{D}) \rangle = \frac{16}{27} \sin^2 \left(\frac{\Delta E}{2\hbar} t \right),$$

$$\langle m_I(\mathbf{D}) \rangle = \frac{1}{2} - \frac{16}{27} \sin^2 \left(\frac{\Delta E}{2\hbar} t \right),$$

Laser-fusion neutron-yield calculations were performed

using a modified [46] MEDUSA code [47], which include the following physical processes: initial laser pulse absorption, propagation of a heat wave, ionization, one-dimensional hydrodynamics of plasma by two-temperature approximation, nonlinear thermal conductivity and limitation, classical and resonant absorption of laser radiation, thermal radiation emission and its absorption by plasma, and thermonuclear reaction yield. The calculations predict that the irradiation of 10^{18} cm⁻³ deuterium and tritium atoms and 10^{19} cm⁻³ ³He with a 2 MJ, 2.3 ns pulse at 350 nm focused to $\sim 100 \ \mu m$ (National Ignition Facility at Lawrence Livermore National Laboratory) will heat the resulting ions to average collision (thermal) energies of ~ 10 keV and lead to the production of $\sim 10^7$ neutrons/pulse from the D-T reaction and $\sim 10^5$ neutrons/pulse each from the D-D and D-³He reactions. Irradiation with a 6 kJ (with $\lambda \sim 1 \mu m$), 1 ps laser pulse [48,49], focused to 10 μ m, will produce ~10⁴ neutrons/pulse from D-T reactions and ~100 neutrons/pulse from D-D or D-³He fusion reactions. These neutron yields compare well with the ones foreseen for ongoing polarized fusion experiments, of ~0.01 neutrons/s [50], and show that the study of polarized fusion with high signals is possible, using high-density SPD and laser-initiated fusion.





- High density SPD and SPH
 - Sofikitis et al., Phys. Rev. Lett. 121, 083001 (2018)
 - But 10¹⁹ cm⁻³???
 - 10¹⁷ cm⁻³ projection done using alkali-H cross sections
 - Depolarization no two body process
 - Simplest process fitting the data
 - Inert gas breaks the complex (SF₆)
 - Minimum depolarization from I

$$D^{\uparrow} + DI \xleftarrow{k_1}{k_{-1}} DI - D^{\uparrow}$$

Complex formation

$$D^{\uparrow} + DI \xrightarrow{k} D + DI$$

Direct depolarization DI

$$DI-D^{\uparrow} \xrightarrow{k_d} DI-D$$

intramolecular

$$DI-D^{\uparrow}+DI \xrightarrow{k_2} D^{\uparrow}+2DI$$

Collisional dissociation DI

$$K = \frac{k_1 k_d [DI]}{k_{-1} + k_d + k_2 [DI]}$$

a (s) 0.20 y y 0.15 0.10 0.05 0.00 200 400 600 800 DI Pressure (mbar)