



Transversity and Λ polarization at COMPASS

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on behalf of the COMPASS Collaboration

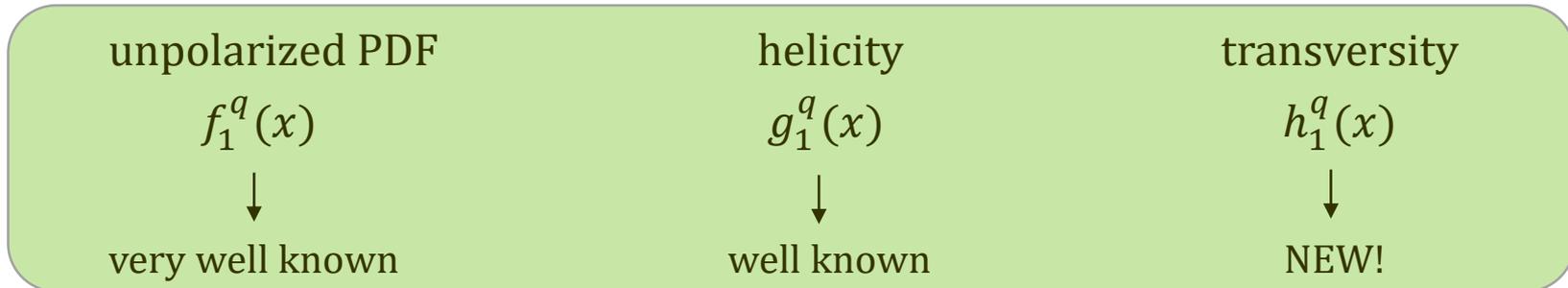


Content of this talk

- The physics case
- Data analysis
- Results
- Interpretation
- Conclusions and perspectives

The physics case

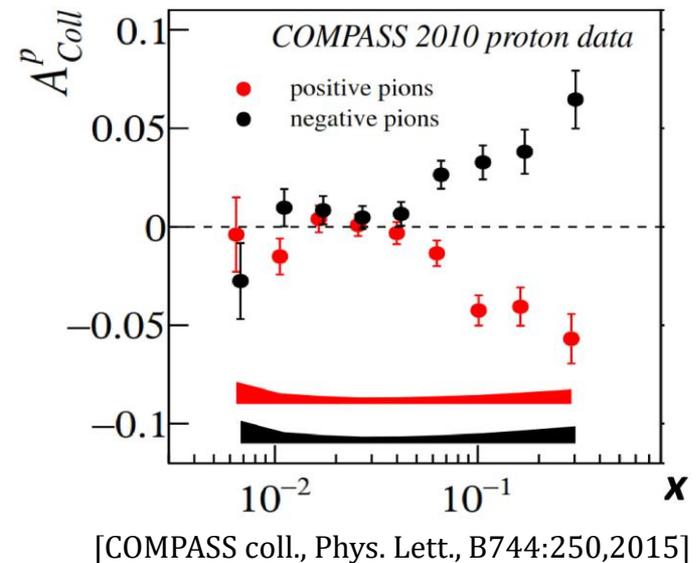
At leading order in collinear QCD, nucleon structure is described by three PDFs:

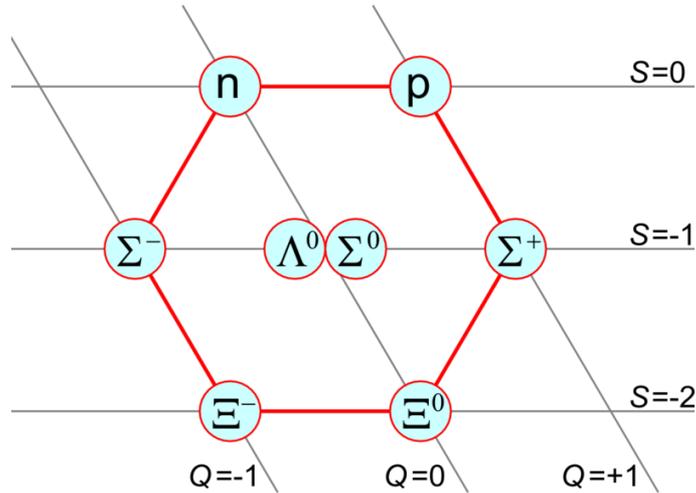


Transversity, introduced by Ralston and Soper in 1979 and rediscovered in early Nineties, accessible in SIDIS looking at:

- **Collins and dihadron asymmetry**
results from HERMES and COMPASS
- **Λ polarimetry**
so far only preliminary results from COMPASS

[Baldracchini et al., 1981] [Artru and Mekhfi, 1990] [Jaffe and Ji, 1992]





Λ MAIN PROPERTIES

- Mass $M_\Lambda = 1115.7 \text{ MeV}/c$
- Spin-parity $J^P = \frac{1}{2}^+$
- Isospin $I = 0$
- Valence quark content uds
- $\Lambda \rightarrow p \pi^-$ (BR 63.9%)
- $\tau = (2.632 \pm 0.020) 10^{-10} \text{ s}$

Λ s reveal their polarization P_Λ through an angular asymmetry in the emission of decay protons (self-analyzing decay)

$$\frac{dN}{d \cos \theta} \propto 1 + \alpha P_\Lambda \cos \theta$$

$\alpha = 0.642 \pm 0.013$ weak decay asymmetry parameter

θ angle between Λ spin and proton momentum in Λ rest frame.

In the **SIDIS** process $\ell p^\uparrow \rightarrow \ell' \Lambda X$,

- with target **nucleon transversely polarized** and
- knowing that **transversity is different from zero**

the quark polarization can be transmitted to the Λ according to the expression

$$P_\Lambda^{raw}(x, z) = f P_T D_{NN} \frac{\sum e_q^2 h_1^{q(\bar{q})} H_1^{\Lambda, q(\bar{q})}(z)}{\sum e_q^2 f_1^{q(\bar{q})} D_1^{\Lambda, q(\bar{q})}(z)}$$

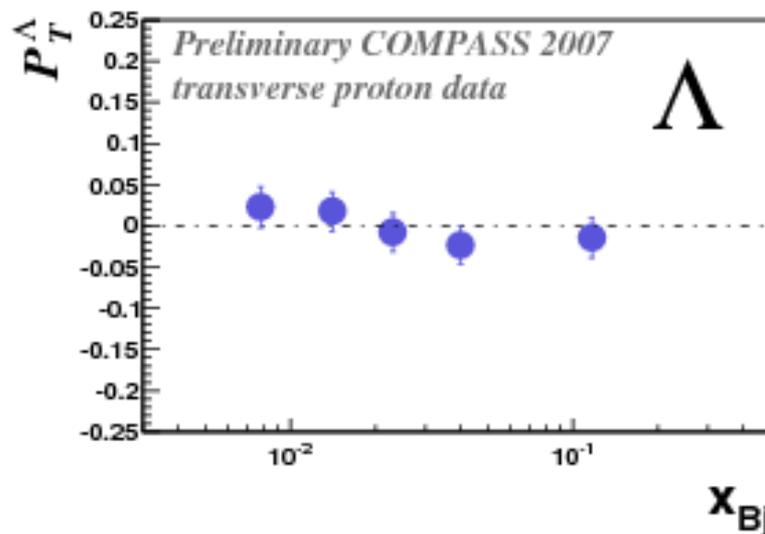
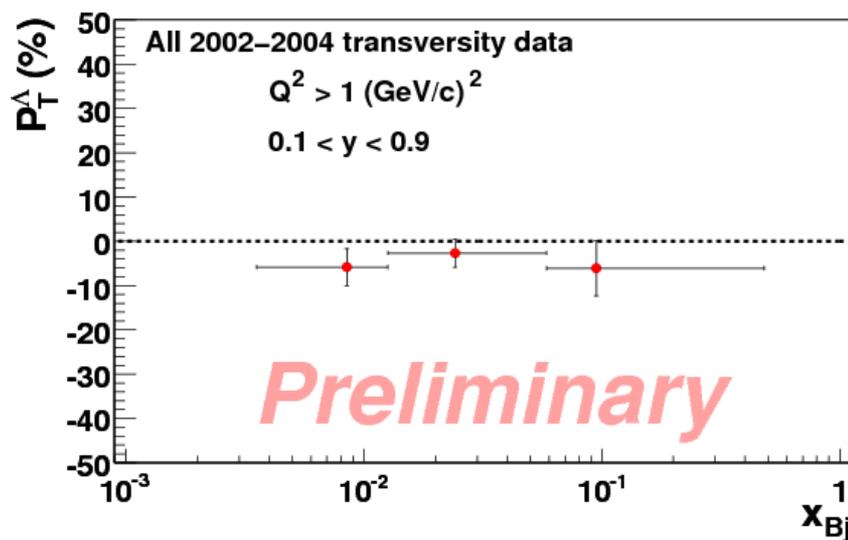
being f the dilution factor, P_T the target polarization and D_{NN} the depolarization factor.

Two remarks on this expressions. It holds true:

- assuming collinear kinematics ($\Lambda \parallel \gamma$, low p_t)
 - in the current fragmentation region
→ our choice $z > 0.2, x_F > 0$

It's a statistically limited measurement, but still interesting,

So far, only preliminary results from COMPASS
(on polarized deuteron and proton target – 2007 only)

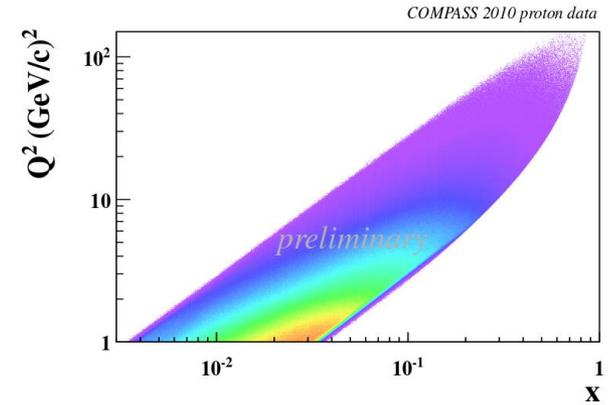
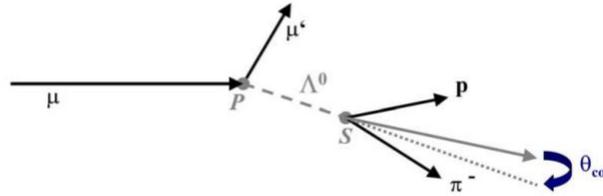


In this talk: results from the complete COMPASS transversely polarized proton data set.

Data analysis

DIS events:

- $Q^2 > 1 \text{ (GeV/c)}^2$
- $W > 5 \text{ (GeV/c)}^2$
- $0.1 < y < 0.9$
- $x > 0.003$



Final state candidates: two charged particles from the decay vertex (V^0 s) with

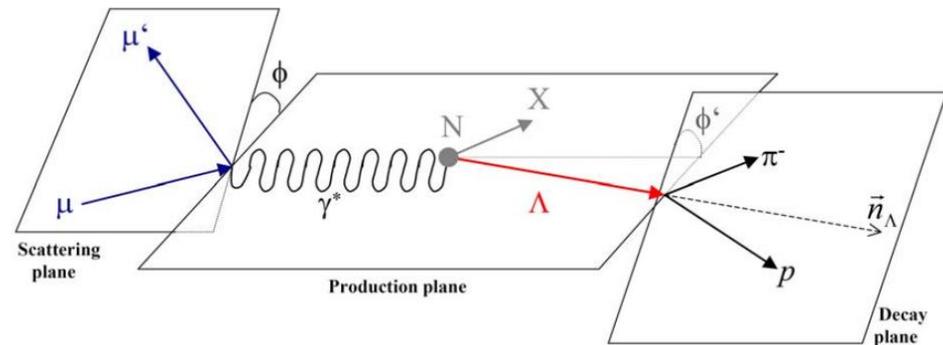
- opposite charge
- momenta $> 1 \text{ GeV/c}$
- $p_T > 23 \text{ MeV/c}$

to reject $e^+ e^-$ from γ conversion

- $\theta_{coll} = \arccos \frac{\mathbf{p}_\Lambda \cdot \mathbf{P}S}{|\mathbf{p}_\Lambda \cdot \mathbf{P}S|} < 7 \text{ mrad}$

- PID with RICH detector

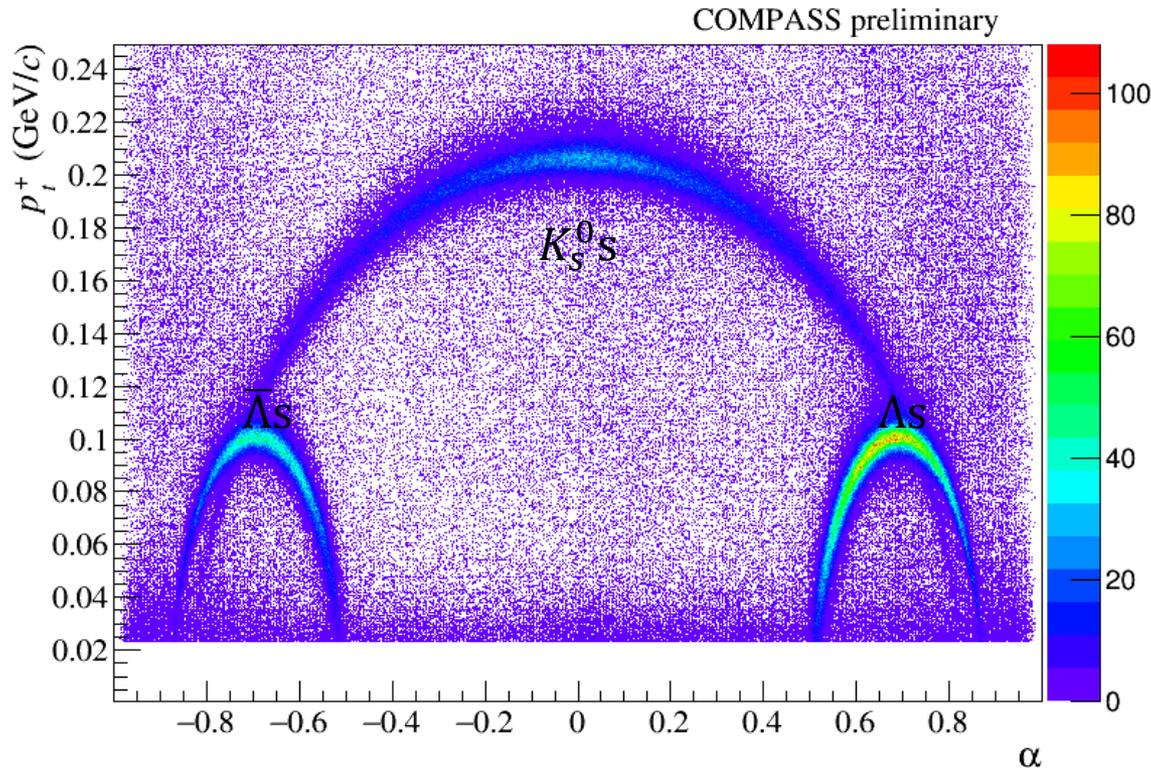
no direct identification of proton but veto on other particles



Longitudinal momentum asymmetry $\alpha = \frac{p_L^+ - p_L^-}{p_L^+ + p_L^-}$ VS transverse momentum p_T of one the decay particles in the V^0 rest frame.

Λ s on the rightmost arc, $\bar{\Lambda}$ s on the leftmost.

The leftover K_S^0 s appear on the large symmetric arc (anyway removed by the forthcoming mass cut).

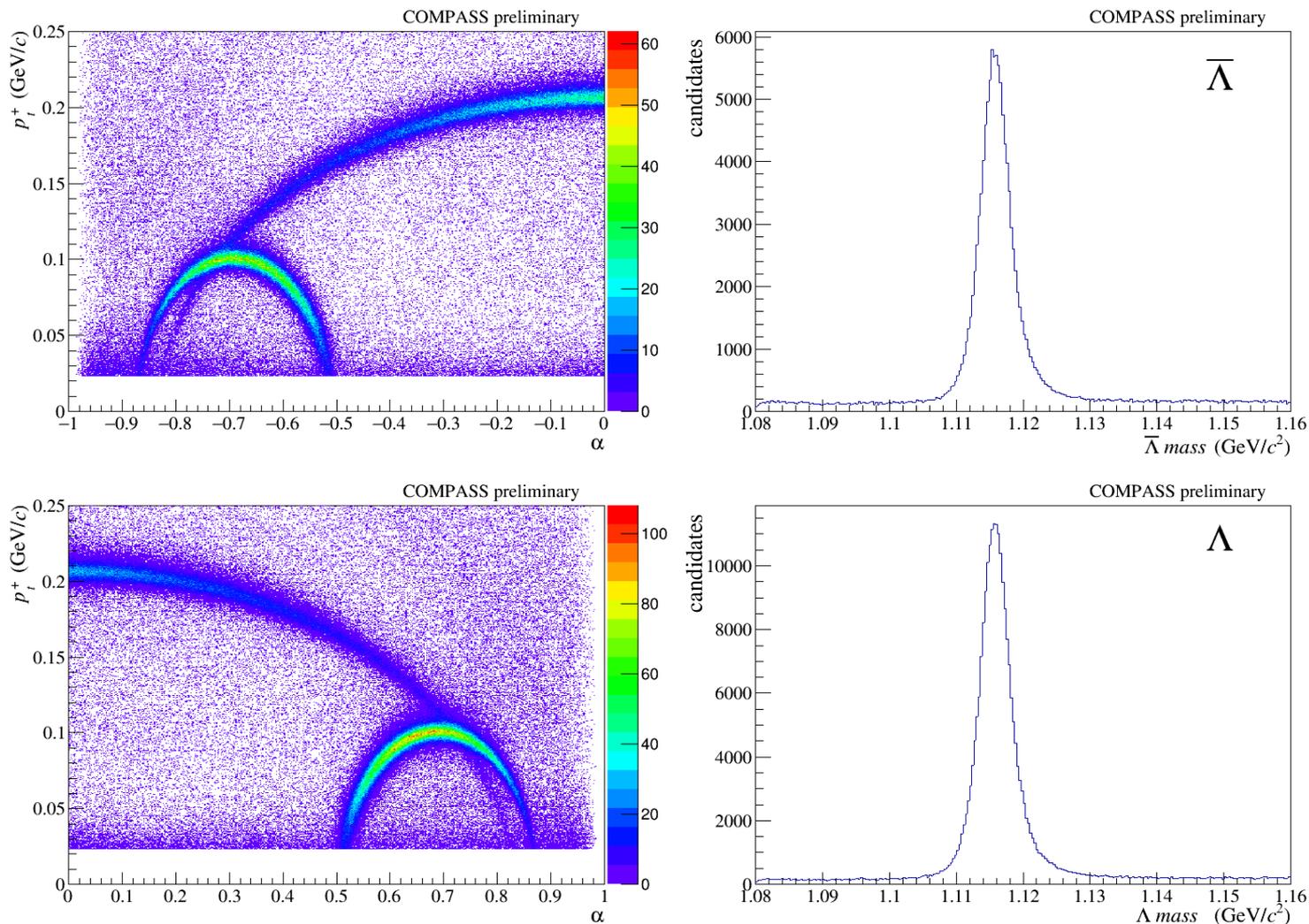


Final Λ – $\bar{\Lambda}$ candidates

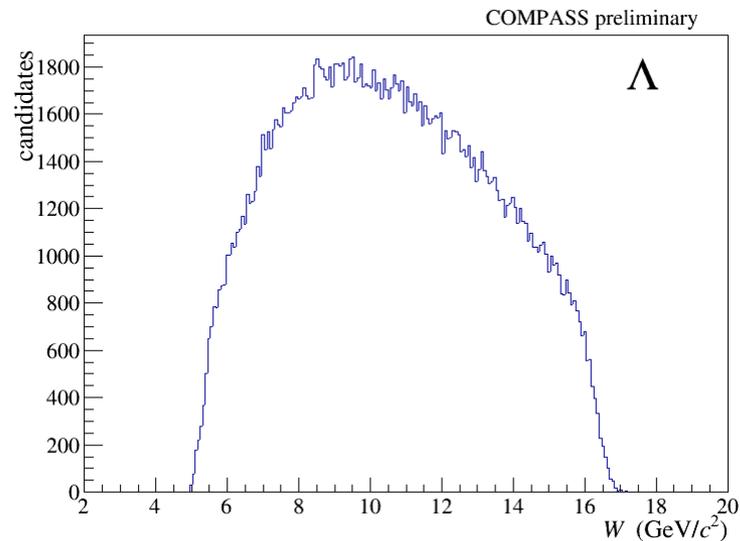
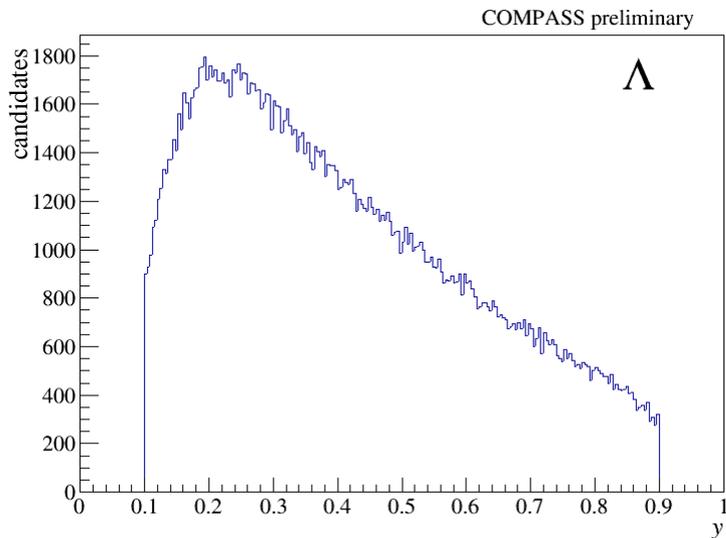
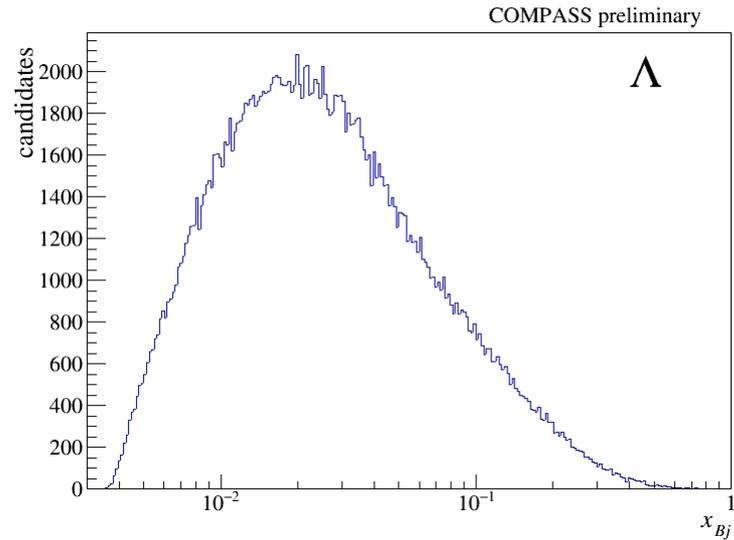
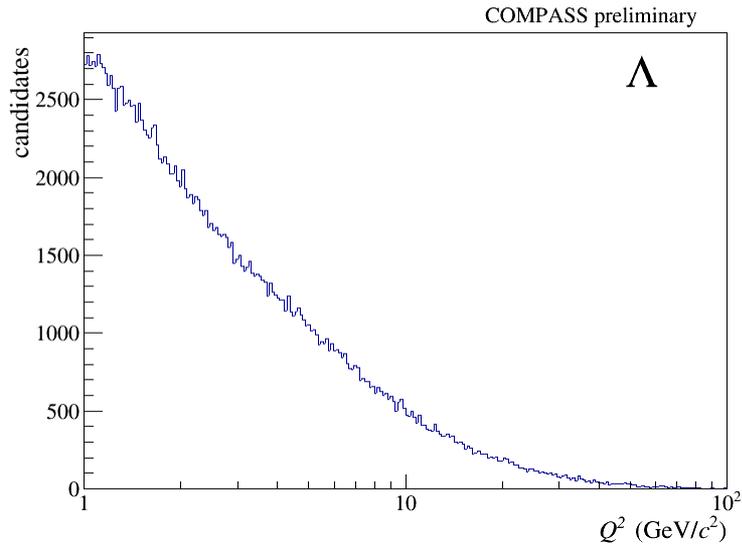


In the mass peak $\sim 305\,000$ Λ s, $154\,000$ $\bar{\Lambda}$ s.

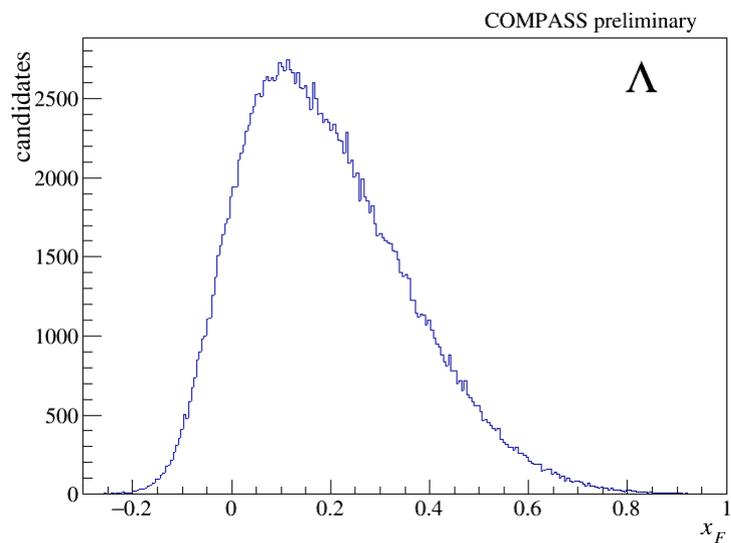
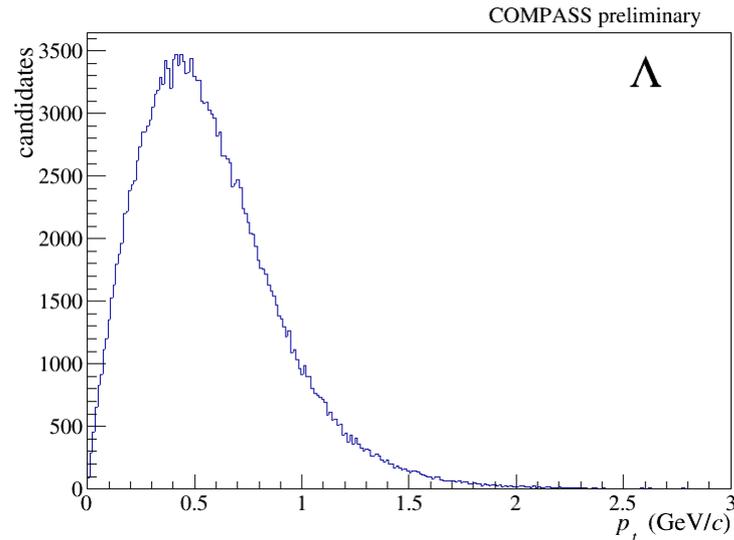
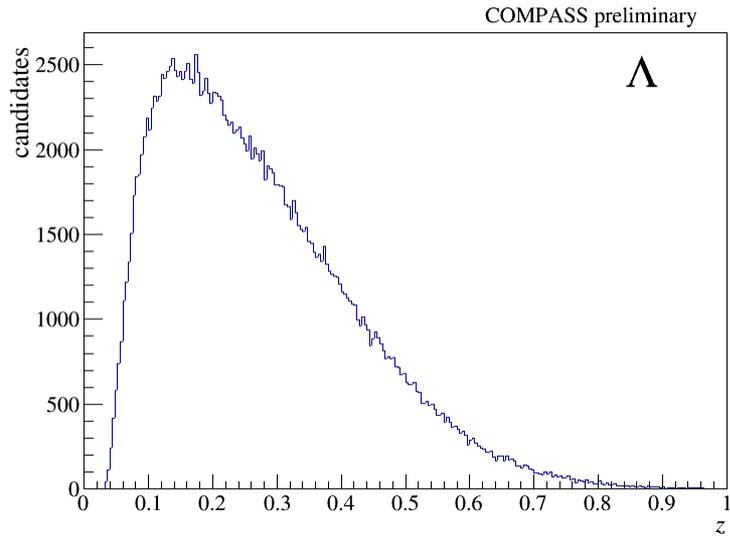
Very clean signal with low background (anyway taken into account)



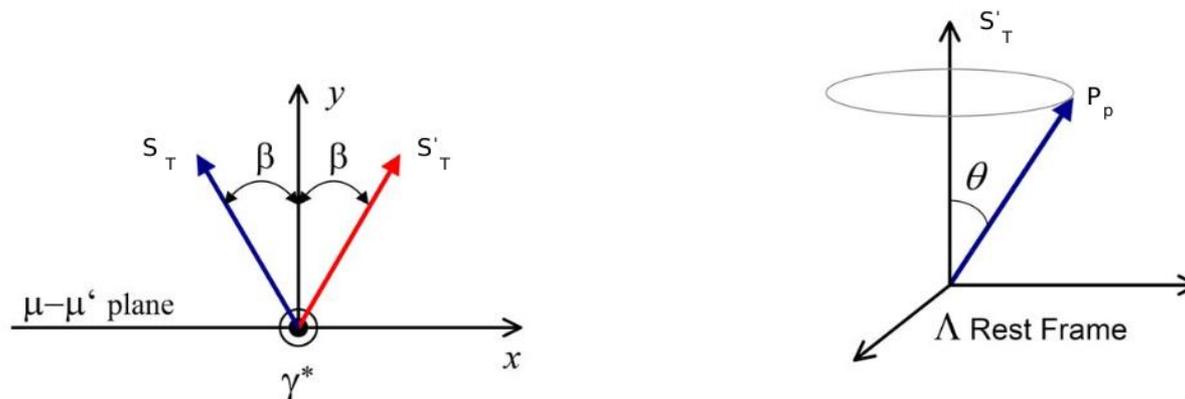
Kinematic distributions: Q^2 , x , y , W (Λ s)



Kinematic distributions: z , p_t , x_F (Λ s)



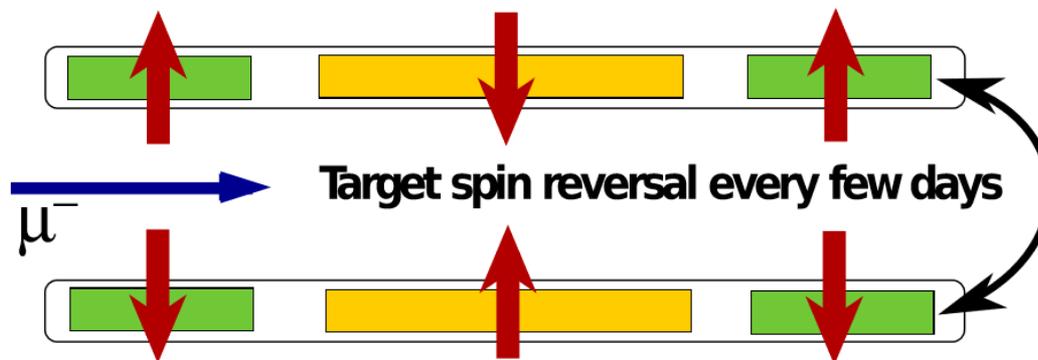
P_Λ is to be measured in the Λ rest frame as an angular asymmetry in the distribution of the proton wrt the outgoing quark spin direction [Mulders – Tangerman, 1996]



- Initial quark spin S_T parallel to the target polarization vector (transverse in the lab)
- Final quark spin S'_T reflection of S_T wrt normal to the scattering plane
- Event by event procedure

Polarization extracted using standard COMPASS methods that take advantage of:

- Polarized target geometry and
- Polarization reversal during data taking - to get rid of the spectrometer acceptance
- Standard studies on systematic effects give $\sigma_{syst} < 0.8 \sigma_{stat}$.



Results

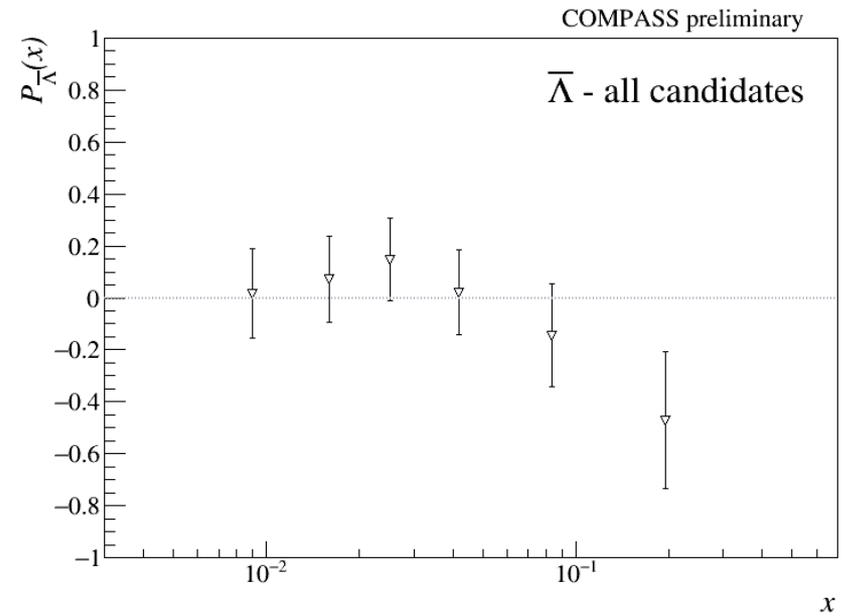
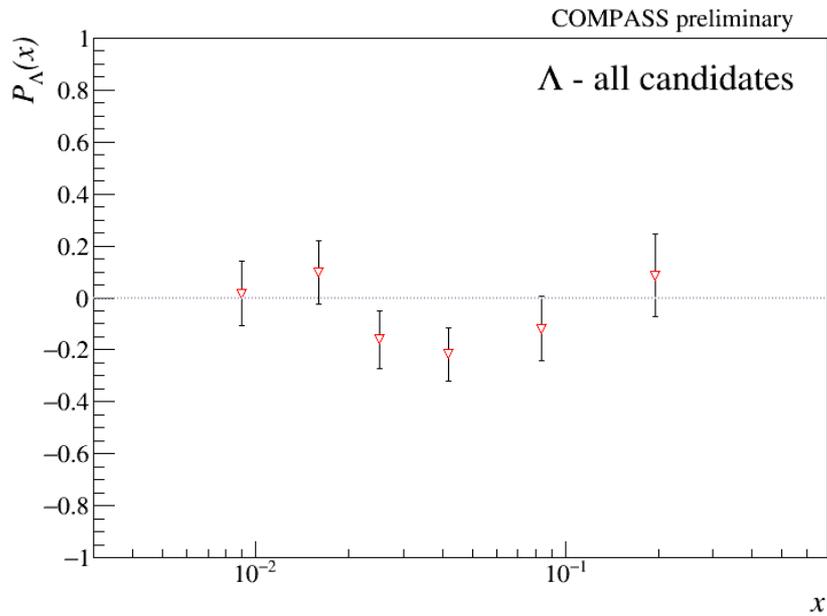
Polarization has been measured as a function of x , z and p_t for both Λ s and $\bar{\Lambda}$ s :

$$P_{\Lambda}^{raw}(x, z) = f P_T D_{NN} \frac{\sum e_q^2 h_1^{q(\bar{q})} H_1^{\Lambda, q(\bar{q})}(z)}{\sum e_q^2 f_1^{q(\bar{q})} D_1^{\Lambda, q(\bar{q})}(z)}$$

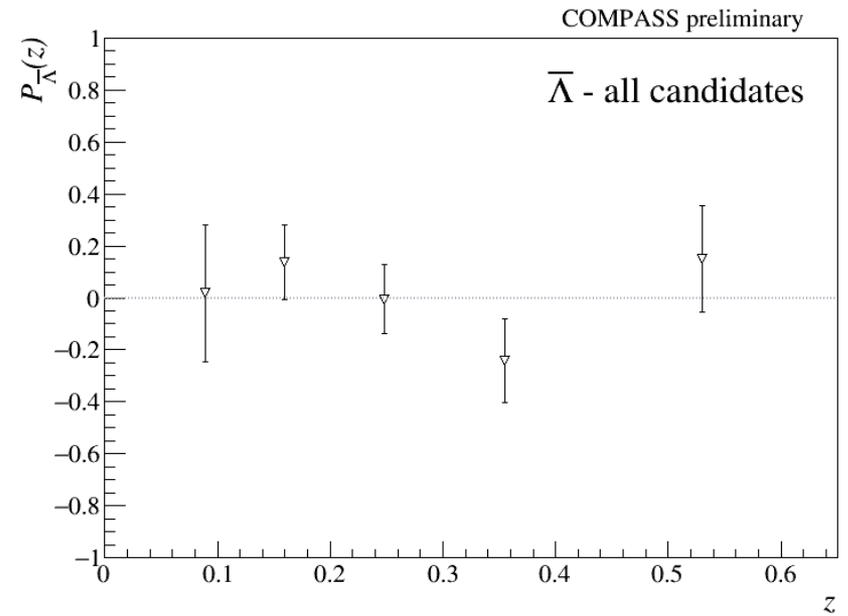
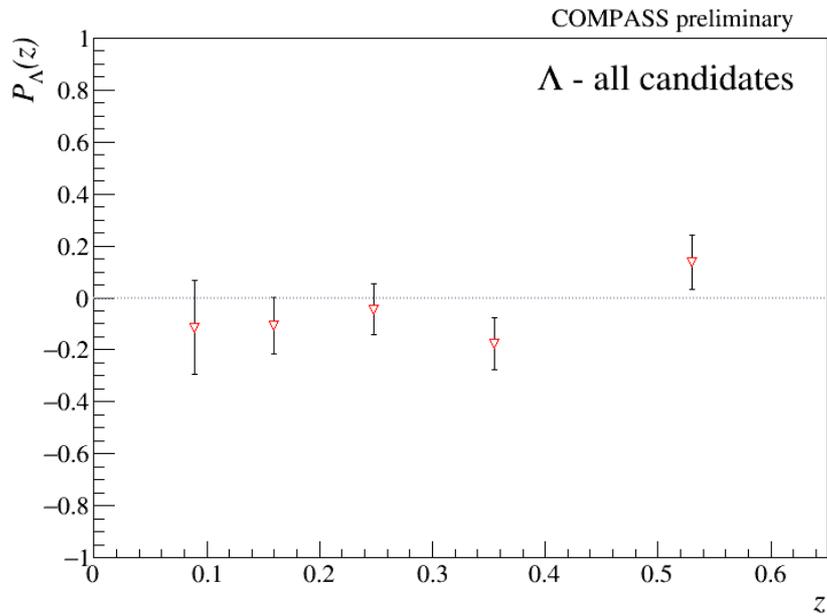
Note: Polarization plots are given here divided by f , P_T and D_{NN} (**spin transfer**).

$$P_{\Lambda}(x, z) = \frac{\sum e_q^2 h_1^{q(\bar{q})} H_1^{\Lambda, q(\bar{q})}(z)}{\sum e_q^2 f_1^{q(\bar{q})} D_1^{\Lambda, q(\bar{q})}(z)}$$

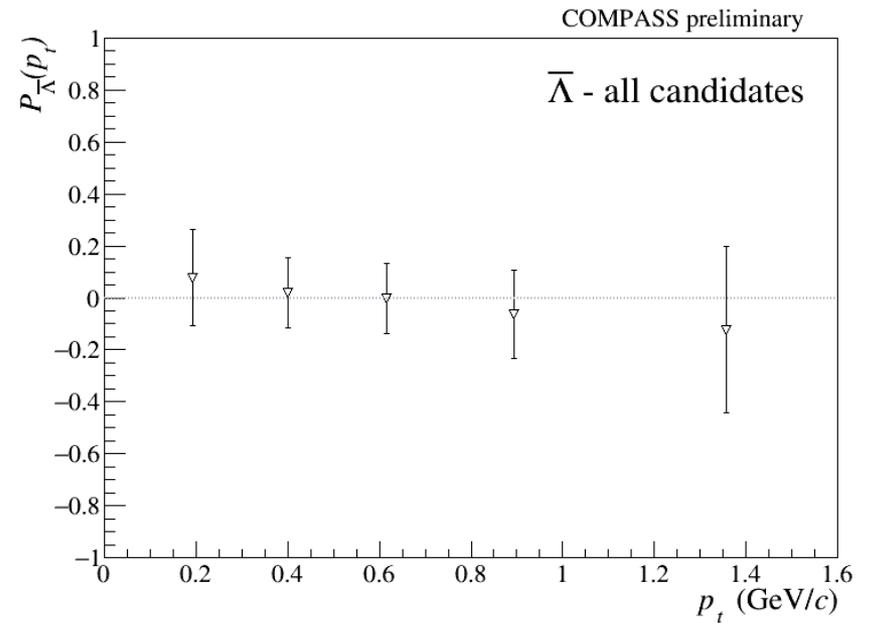
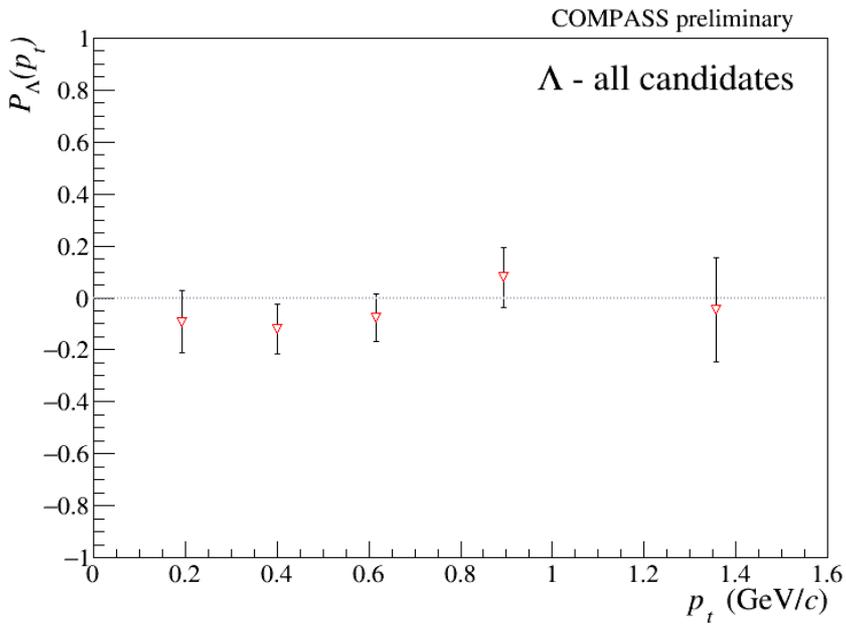
$$P_{\Lambda(\bar{\Lambda})}(x)$$



$$P_{\Lambda(\bar{\Lambda})}(z)$$



$$P_{\Lambda(\bar{\Lambda})}(p_t)$$



Polarization has also been measured in six other kinematic regions:

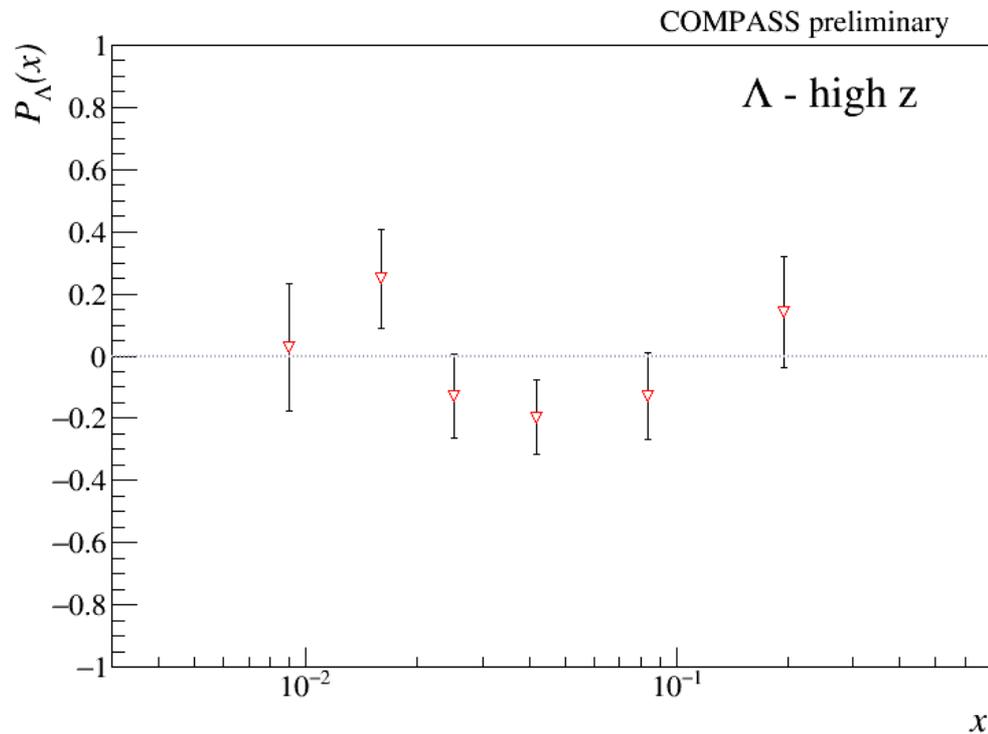
- High z : $z > 0.2$ and $x_F > 0$ (“current” fragmentation region)
- Low z : $z < 0.2$ or $x_F < 0$ (“target” fragmentation region)
- High x : $x > 0.032$ ($\mathbf{h}_1^{u,d}(x)$ different from zero)
- Low x : $x < 0.032$
- High p_t : $p_t > 1 \text{ GeV}/c$
- Low p_t : $p_t < 1 \text{ GeV}/c$

In general, as in the case of all Λ s and $\bar{\Lambda}$ s,
polarizations are found compatible with zero.

Interpretation

Interpretation for directly produced Λ : we neglect the spin transfer to Λ from heavier hyperons

$$P_\Lambda(x, z) = \frac{\sum e_q^2 h_1^{q(\bar{q})} H_1^{\Lambda, q(\bar{q})}(z)}{\sum e_q^2 f_1^{q(\bar{q})} D_1^{\Lambda, q(\bar{q})}(z)}$$



We **KNOW** that:

- $h_1^u(x)$ and $h_1^d(x)$ are different from zero at large x ;
- $h_1^{\bar{u}}(x)$ and $h_1^{\bar{d}}(x)$ are compatible with zero.

We can **ASSUME** that:

- $h_1^{\bar{s}}(x) \sim 0$;
- negligible contribution from \bar{q} in unpolarized fragmentation process;
- isospin symmetry at work: $D_1^{\Lambda,u}(z) = D_1^{\Lambda,d}(z)$ and $H_1^{\Lambda,u}(z) = H_1^{\Lambda,d}(z)$;
- $D_1^{\Lambda,s}(z) = c_1 D_1^{\Lambda,u}(z)$ with constant c_1 ;
- Analogously, if $H_1^{\Lambda,u}(z) \neq 0$, $H_1^{\Lambda,s}(z) = c_2 H_1^{\Lambda,u}(z)$.

The quantity $1/c_1$ is usually referred to as strangeness suppression factor. In [J.-J. Yang, Phys Rev D65 2002] it is put at 0.44.

With these ingredients we can write a simplified expression for P_Λ :

$$P_\Lambda(x, z) = \frac{[4h_1^u(x) + h_1^d(x)]H_1^{\Lambda,u}(z) + h_1^s(x)H_1^{\Lambda,s}(z)}{[4f_1^u(x) + f_1^d(x) + c_1f_1^s(x)]D_1^{\Lambda,u}(z)}$$

Now, we can interpret the data according to **three different hypotheses**:

1. Transversity is a valence object
2. Polarization is entirely due to the s quark (SU(6) approach)
3. Quark-diquark model [J.-J. Yang, Nucl Phys A699:562-578, 2002]

Note: simplified expression looks interesting with a deuteron target:

$$P_\Lambda(x, z) = \frac{5(h_1^u(x) + h_1^d(x))H_1^{\Lambda,u}(z) + 2h_1^s(x)H_1^{\Lambda,s}(z)}{5(f_1^u(x) + f_1^d(x))D_1^{\Lambda,u}(z) + 2f_1^s(x)D_1^{\Lambda,s}(z)} \approx \frac{2h_1^s(x)H_1^{\Lambda,s}(z)}{5(f_1^u(x) + f_1^d(x) + 2c_1f_1^s(x))D_1^{\Lambda,u}(z)}$$

If transversity is a valence object, then $h_1^s(x) \approx 0$ and

$$P_\Lambda(x) = \frac{[4h_1^u(x) + h_1^d(x)]}{[4f_1^u(x) + f_1^d(x) + c_1 f_1^s(x)]} \frac{\int dz H_1^{\Lambda,u}(z)}{\int dz D_1^{\Lambda,u}(z)}$$

$$\rightarrow R(x) = \frac{\int dz H_1^{\Lambda,u}(z)}{\int dz D_1^{\Lambda,u}(z)} = \frac{[4f_1^u(x) + f_1^d(x) + c_1 f_1^s(x)]}{[4h_1^u(x) + h_1^d(x)]} P_\Lambda(x)$$

c_1	$\langle R \rangle$
2	-0.39 ± 0.73
3	-0.38 ± 0.75
4	-0.37 ± 0.76

First extraction of $R(x)$, largely compatible with zero, weak dependence on c_1 .

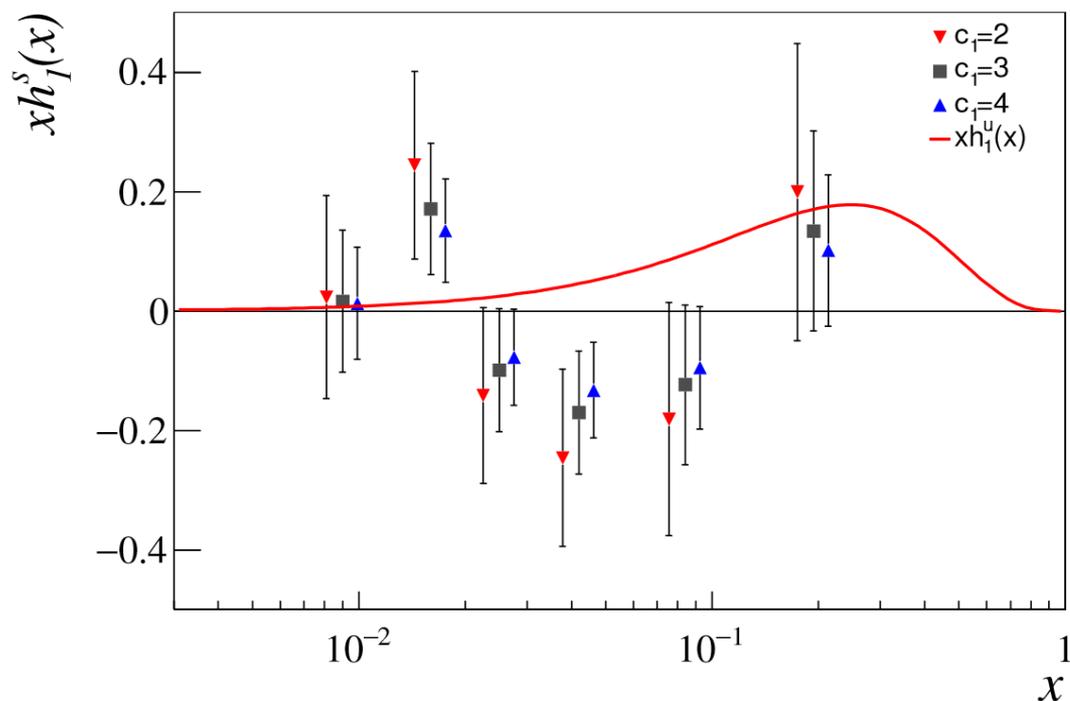
Hypothesis 2: polarization due to s quark only



If P_Λ is due to s quark only, then $H_1^{\Lambda,u}(z) = H_1^{\Lambda,d}(z) = 0$. Assuming $H_1^{\Lambda,s}(z) = D_1^{\Lambda,s}(z)$,

$$P_\Lambda(x) = \frac{c_1 h_1^s(x) D_1^{\Lambda,u}(z)}{[4f_1^u(x) + f_1^d(x) + c_1 f_1^s(x)] D_1^{\Lambda,u}(z)}$$

$$\rightarrow h_1^s(x) = \left(\frac{4f_1^u(x) + f_1^d(x)}{c_1} + f_1^s(x) \right) P_\Lambda(x)$$

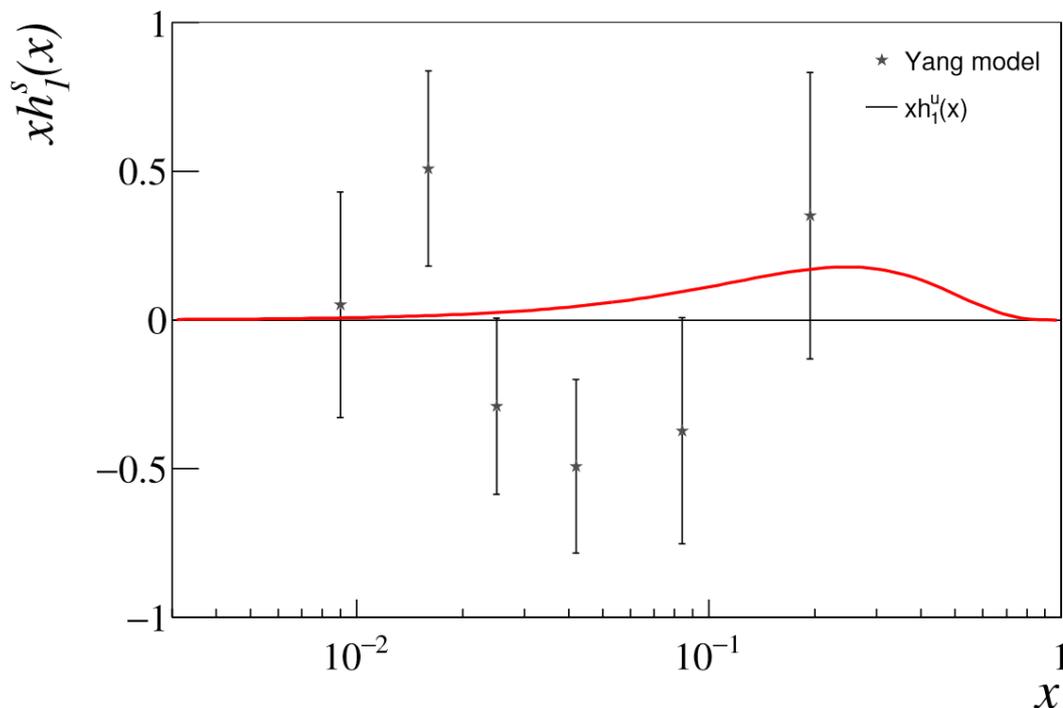


Hypothesis 3: Quark-diquark Yang model



P_Λ is written in terms of flavor (F) and spin structure functions (\widehat{W}). $h_1^s(x)$ can be accessed.

$$P_\Lambda(x, z) = \frac{\left(4h_1^u(x) + h_1^d(x)\right) \frac{1}{4} \left[\widehat{W}_s^{(u)}(z)F_s(z) - \widehat{W}_V^{(u)}(z)F_M(z)\right] + \mathbf{h_1^s(x)} \widehat{W}_s^{(s)}(z)}{\left(4f_1^u(x) + f_1^d(x)\right) \frac{1}{4} [F_s(z) + 3F_M(z)] + f_1^s(x)}$$



Conclusions and perspectives

- Transversity-induced polarization of Λ hyperons in SIDIS measured using the whole COMPASS transversely polarized proton data set.
- Λ and $\bar{\Lambda}$ polarizations evaluated in their rest frame along the outgoing quark spin; measured in seven kinematic regions, generally compatible with zero.
- Three main hypothesis to interpret the results:
 1. the first (transversity a valence object) gives the integrated ratio of the fragmentation functions $H_1^{\Lambda,u}(z)$ and $D_1^{\Lambda,u}(z)$, compatible with zero;
 2. the second (only s quark counts) allows for an extraction of $xh_1^s(x)$ dependent on the parameter $c_1 = D_1^{\Lambda,s}(z)/D_1^{\Lambda,u}(z)$;
 3. the third (quark-diquark model) again gives $xh_1^s(x)$ without assumptions on the fragmentation functions.
- Even if definite conclusions cannot be drawn, mainly due to the statistical uncertainty, this is a contribution to a longstanding issue
- Ratios of fragmentation functions are extracted here for the first time
- Much more could be studied with new deuteron data.

thank you!