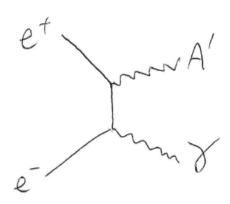
Opportunities with Phase 2: Dark Sector

E. Graziani

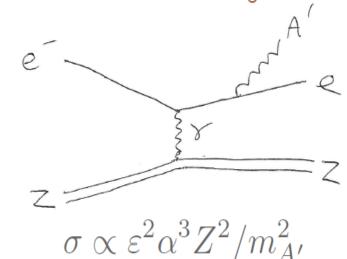
Dark photon production mechanisms

e+e- annihilation



$$\sigma \propto \varepsilon^2 \alpha^2 \left(1 - m_{A'}^2 / E_{CM}^2\right) / E_{CM}^2$$

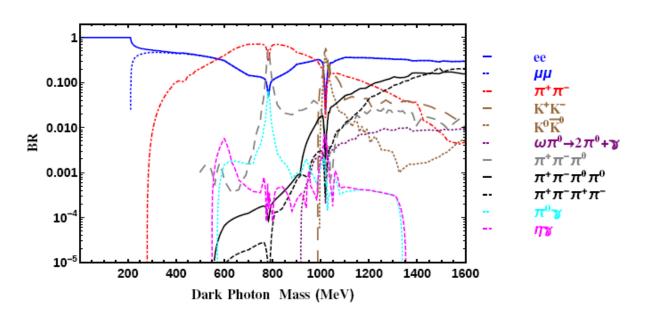
e- bremsstrahlung



meson decay

$$\mathcal{B} = 2\varepsilon^2 \left(1 - m_{A'}^2 / m_{\pi^0}^2 \right)^3$$

Dark photon visible decays



- The dark photon visible decay sector is quite crowded.
- Huge and impressive harvest of results during the last ~ 7 years, from colliders, fixed target, meson decay experiments.
- The possibility of explaining $(g-2)_{\mu}$ is now closed (mostly due to NA48/2 $\pi^0 \rightarrow \gamma e^+e^-$ search)
- Still vast and interesting regions to explore, but they require luminosities well beyond Phase 2.

From SLAC Dark Sectors Workshop executive summary

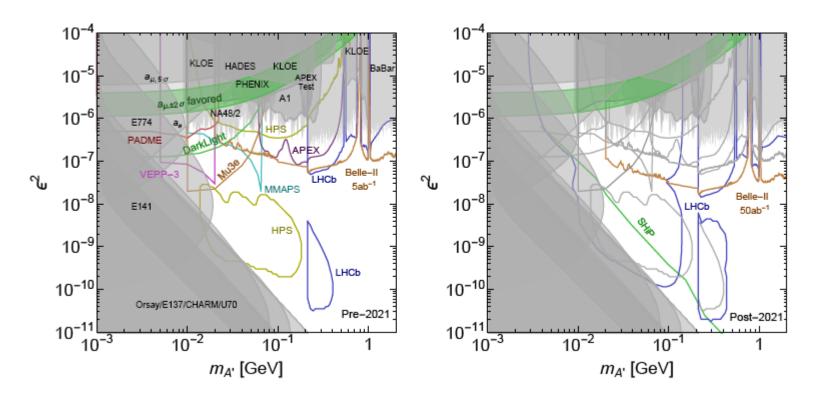


FIG. 4: Sensitivity to A' for exclusive experiments seeking visible decay modes $A' \to \ell^+\ell^-$. Left: Experiments capable of delivering results over the next 5 years to 2021. Shaded regions show existing bounds. Green band shows 2σ region in which an A' can explain the discrepancy between the calculated and measured value for the muon g-2. Right: Longer term prospects beyond 2021 for experimental sensitivity. All projections on left plot are repeated in gray here. Note that LHCb and Belle-II can probe to higher masses than 2 GeV and SHIP can probe to lower values of ϵ than indicated.

One real golden channel at Phase 2: search for the invisible decays of the dark photon

- Much more unconstrained wrt visible decay case
- May still probe (g-2)_μ anomaly
- May give access to light dark matter
- Requires specific triggers (not guaranteed at higher luminosities)
- Both signal and background come from the continuum →
- ullet ightarrow can be searched at any c.m. energy, provided machine conditions are good

From B2TIP report (draft)

14.6 Search for a Dark Photon decaying into Light Dark matter (C. Hearty, T. Ferber)

Dark sectors are an exciting topic in particle physics. These theories introduce new particles that interact gravitationally with standard model matter, but do not interact directly via the SM electroweak or strong forces. Such particles would be the dark matter that is observed astronomically.

14.6.1 Theory

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There are a variety of theories that involve a dark sector. One of the simplest includes a dark photon A' that mixes with strength ε to the standard model photon [59]. Annihilation of heavy dark matter fermions would produce an A', which would decay to standard model particles, if the A' is the lightest dark sector particle. This process could explain the positron excess observed by PAMELA, Fermi LAT, and AMS [60, 61, 62]. These observations are consistent with an A' mass $M_{A'}$ in the MeV/c² to GeV/c² range. With this mass, the A' could be radiatively produced in e^+e^- collisions, $e^+e^- \rightarrow \gamma A'$.

The cross section for this process is proportional to $\varepsilon^2\alpha^2/E_{CM}^2$ [63]. The decay branching fractions of the A' are the same as a virtual photon of mass $M_{A'}$ (i.e. $e^+e^- \rightarrow \gamma^* \rightarrow X$).

A significant number of experiments have recently published results of A' searches where the A' decays into charged lepton pairs. Several other dedicated experiments will proceed over the next several years. A recent search by BaBar for the radiative production of the A' in the e^+e^- and $\mu^+\mu^-$ final states used 514 fb⁻¹ of data [64]. The standard model rates for $e^+e^- \to \gamma e^+e^-$ and $e^+e^- \to \gamma \mu^+\mu^-$ are large, and the search for the A' consists of a search for a narrow peak in the dilepton mass spectrum on top of a large

background.

If the A' is not the lightest dark sector particle, 2120 it will dominantly decay into light dark matter 21211 via $A' \rightarrow \chi \bar{\chi}$. Since the interaction probability of 2122 dark matter with the detector is negligible, the 2122 experimental signature of such a decay will be 21.24 a mono-energetic ISR photon γ_{ISR} with energy nm $E_{\gamma} = (E_{CM}^2 - M_{A'}^2)/(2E_{CM})$. This search requires must a L1 trigger that is sensitive to single photons which 2127 was not available at Belle and only partially avail- nm able at BaBar. BaBar recorded about 57 fb⁻¹ of 2120 data with various single photon triggers in its final 21220 year of operations, including a scan above $\Upsilon(4S)$. 21221 28 fb⁻¹ of this data recorded at the $\Upsilon(3S)$ reso-max nance was used in a search for invisible decays of 21222 the light Higgs A^0 , $\Upsilon(3S) \rightarrow \gamma A^0$, $A^0 \rightarrow \text{invisible.} 2124$ The result was presented in a conference note [65], 21225 but has not yet been published.

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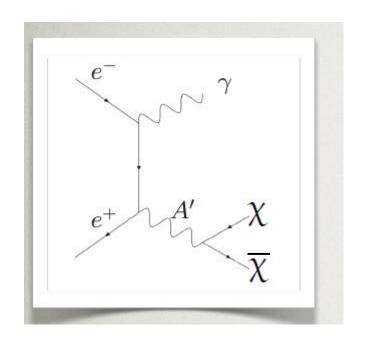
14.6.2 Experiment

Monte Carlo Simulation

Signal MC events $(e^+e^- \rightarrow \gamma A', A' \rightarrow \chi \bar{\chi})$ are mass generated using MadGraph [66] and a model 21200 based on [67] that includes a dark photon A' and nmfermionic dark matter y. Each signal sample 2122 is generated using a fixed dark photon mass 21222 $m_{A'} = [0.1, 0.5, 1.0, 2.0, 3.0, 4.0, 5.0, 6.0, 7.0, 8.0, 21224]$ 8.5, 8.75, 9.0, 9.25, 9.5, 9.75 GeV and contains 2123 50000 events. Events are generated for a maximal nms photon pseudo-rapidity of η_{π}^* < 1.681, which nwcorresponds to $|\cos(\theta_n^*)| = 0.933$. The beam 2123 energy is set to $E^* = 10.58 \,\text{GeV}$. We assume a dark nmmatter mass $m_{\nu} = 1 \,\text{MeV}$, and we set the coupling 21340 to $g_e = g_{\gamma}$. The decay width of the dark photon is m_{γ} set to the tree-level width which increases slowly 2132 with $m_{A'}$ and is of O(MeV). We assume that all 21342 decays of the A' are into $\chi \bar{\chi}$ and set the kinetic mass mixing parameter to $\epsilon = 1$. The resulting cross 2136 section, including vacuum polarization corrections 2136 (up to about 10 %), is shown in Fig. 14.3.

The background in this analysis is domi-nated nated by high cross section QED processes nzo $e^+e^- \rightarrow e^+e^-\gamma(\gamma)$ and $e^+e^- \rightarrow \gamma\gamma(\gamma)$ that nzo

The process

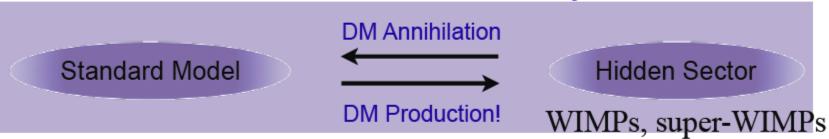


monochromatic photon
$$E_{\gamma} = (s-m_{A'}^2)/2\sqrt{s}$$

Kinematically equivalent description in terms of E_{γ} or $m_{A'}$ (missing mass)

If light dark matter $m_{A'} > 2m_{\chi}$ exists, then all BRs to SM particles $\sim \epsilon^2 \Rightarrow$ BR(A' $\to \chi \overline{\chi}$) $\approx 100\%$, with χ escaping detection

Possible connection to WIMP-y dark matter



Mediators (SM Z, h etc or dark force)

Heavy WIMP/heavy mediators: - "mainstream" literature

Light WIMPs/light mediators: Boehm et al; Fayet; MP, Ritz, Voloshin; Hooper,
Zurek; others

Heavy WIMPs/light mediators: Finkbeiner, Weiner; Pospelov, Ritz, Voloshin (secluded DM); Arkani-Hamed et al., many others

Light WIMPs/heavy mediators: does not work. (Except for super-WIMPs; or non-standard thermal history)

M. Pospelov

Light WIMPs due to light mediators direct production/detection

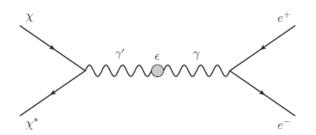
(Boehm, Fayet; MP, Riz, Voloshin ...) Light dark matter is not ruled out if one adds a light mediator.

WIMP paradigm:
$$\sigma_{\text{annih}}(v/c) \sim 1 \text{ pbn } \implies \Omega_{\text{DM}} \simeq 0.25,$$

Electroweak mediators lead to the so-called Lee-Weinberg window,

$$\sigma(v/c) \propto \left\{ \begin{array}{ll} G_F^2 m_\chi^2 & \text{for } m_\chi \ll m_W, \\ 1/m_\chi^2 & \text{for } m_\chi \gg m_W. \end{array} \right. \implies \text{few GeV} < m_\chi < \text{few TeV}$$

If instead the annihilation occurs via a force carrier with light mass, DM can be as light as ~ MeV (and not ruled out by the CMB if it is a scalar).



Dark matter and structure of galaxies







Core/cusp problem: Galaxies and clusters are less dense than cold dark matter (WIMPs) predictions

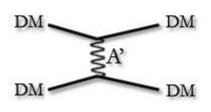
Moore (1994), Flores & Primack (1994)

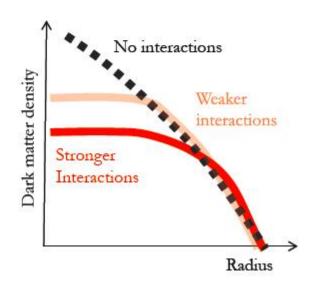
Self-interacting dark matter

Spergel & Steinhardt (2000)

MeV-GeV scale dark force

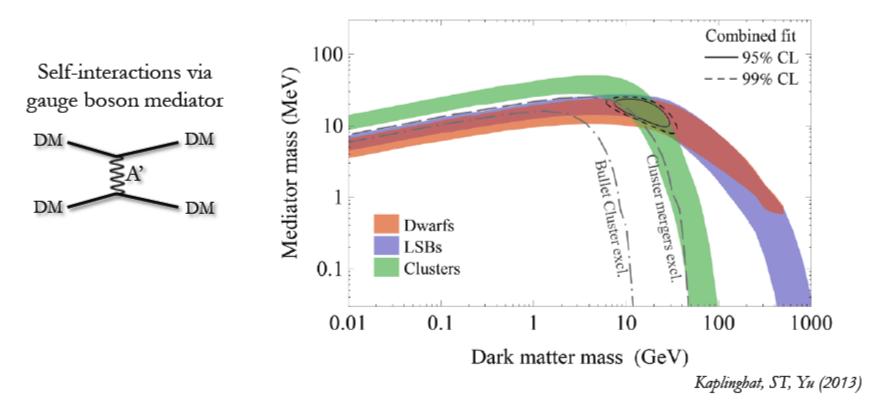
J—GeV scale dark ford ST, Yu, Zurek (2013)







Dark matter with dark photon

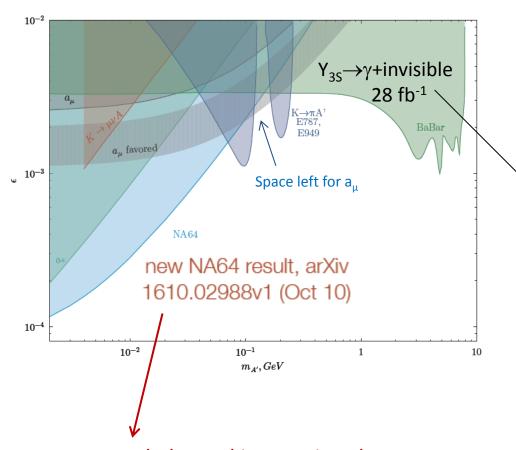


Model-dependent: Dark sector parameters can be fit from astrophysical data. Not fixed how dark photon couples to Standard Model (kinetic mixing unknown)

Model-independent: Dark sector particles below GeV scale to get large enough cross section



Dark photon to invisible: present situation



Two week data taking run in July (criticized result)

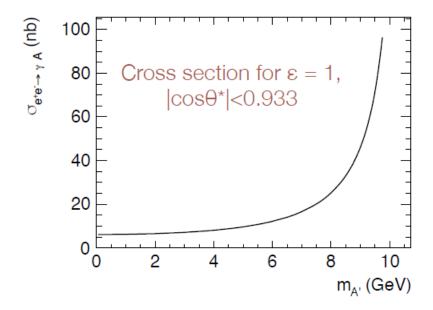
BaBar took a small portion (~54 fb⁻¹) of the full data set with a single photon trigger.

Belle never had one

Unpublished (conference paper) reinterpretation of a $Y_{3S} \rightarrow \gamma + h_0$ search. Angular distributions are quite different (vector vs scalar). Some refinement of the measurement expected. Former Belle II projections were based on this result.

MC samples

- All produced by Torben, at CLUMEQ (McGill), Phase 2 geometry and background levels.
 - Bhabhas and radiative Bhabhas scaled by 0.025; others by 0.1, using 12th BG campaign (no two-photon)
- Signal: MadGraph, various A' masses

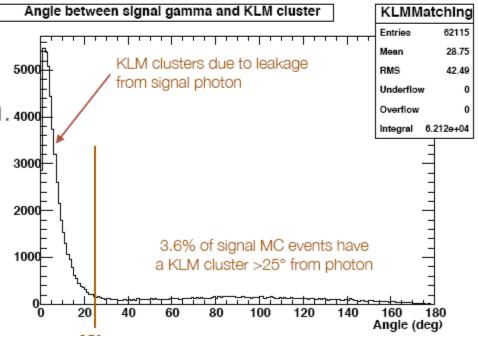


Backgrounds

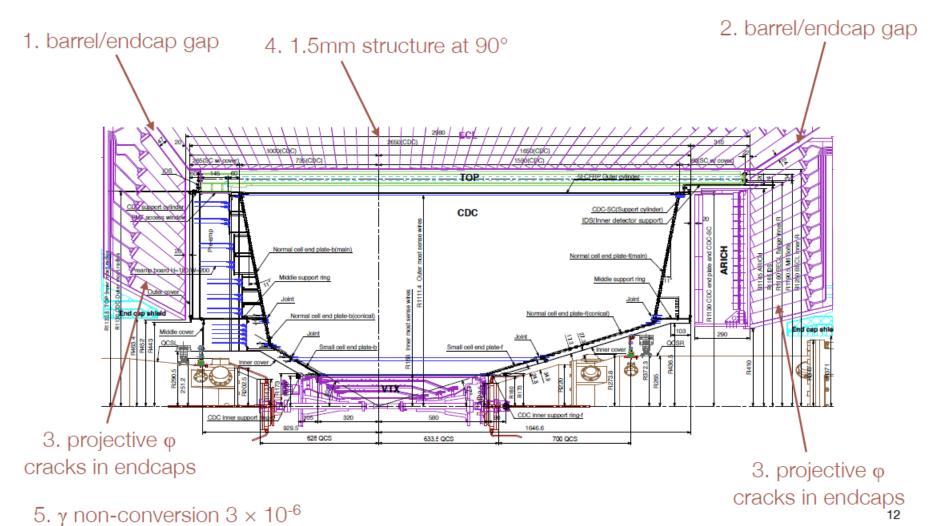
- Irreducible: final state has only one photon in [12°,157°]
 - 2M events with E*>1.8 GeV and 22° < θ < 139° in 20 fb⁻¹ (0.1 nb); E and θ are strongly correlated.
 - 85% due to radiative Bhabhas
 - 15% due to e⁺e⁻ $\rightarrow \gamma \gamma (\gamma)$ (minimum 3 photons in final state)
- Reducible: two photons in detector, one of which is missed in both ECL and KLM

Event selection (other than final cut on θ_{lab})

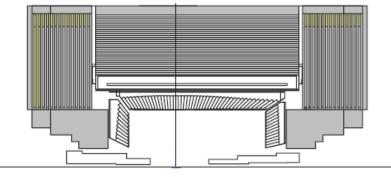
- E* >1.8 GeV with no other ECL cluster with E* > 0.1 GeV
- No tracks with p_t>0.2 GeV/c
- No KLM cluster >25°
 (3D angle in center of mass) from signal photon. 4000
- Not optimized.

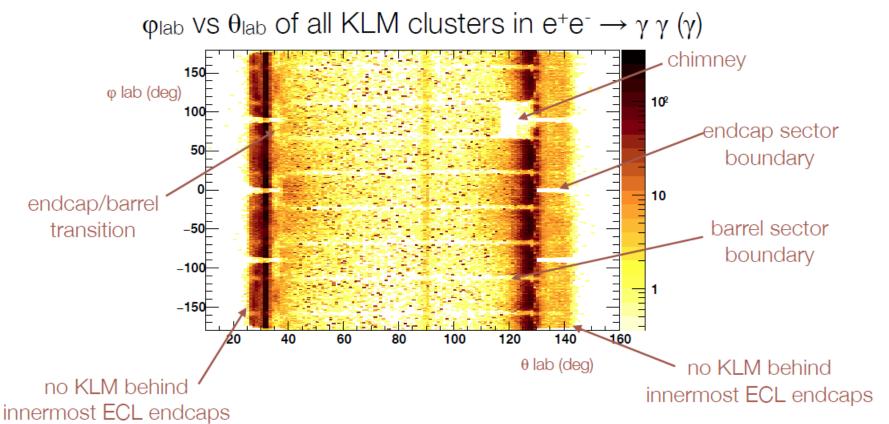


Sources of ECL inefficiency (in order of importance)



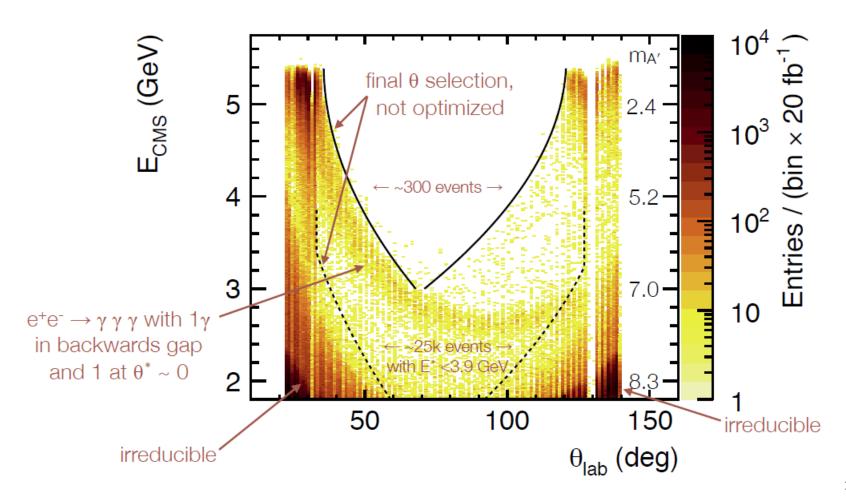
Sources of KLM inefficiency





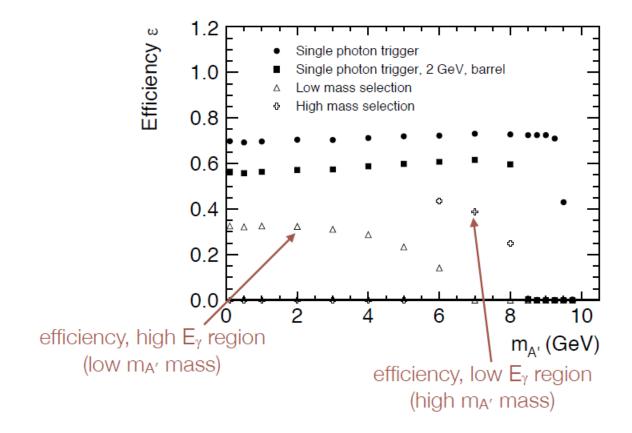
From C. Hearty presentation at October Belle II GM Predicted backgrounds in 20 fb⁻¹

• Final sample is almost entirely $e^+e^- \rightarrow \gamma \gamma (\gamma)$ with $\geq 3\gamma$



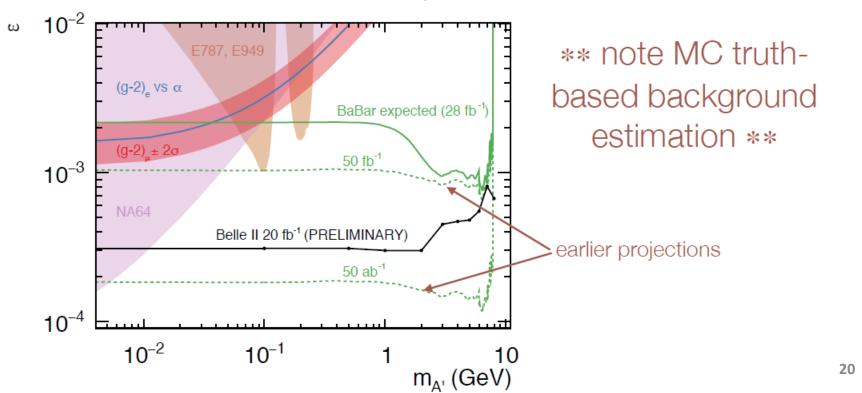
From C. Hearty presentation at October Belle II GM Efficiency for signal events

Signal inefficiency is dominated by acceptance.



Projected upper limits on ε as a function of $m_{A'}$ for 20 fb⁻¹ Phase 2 data set

- Limits are much better than earlier projections derived from BaBar because of better detector performance.
 - no projective cracks in φ; fully functioning muon system



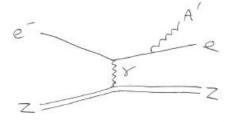
Big improvement wrt former projections based on BaBar: mostly due to the non-projective cracks of ECL

Probably optimistic, but anyway a solid indication of a very interesting potential reach. Main concerns are about:

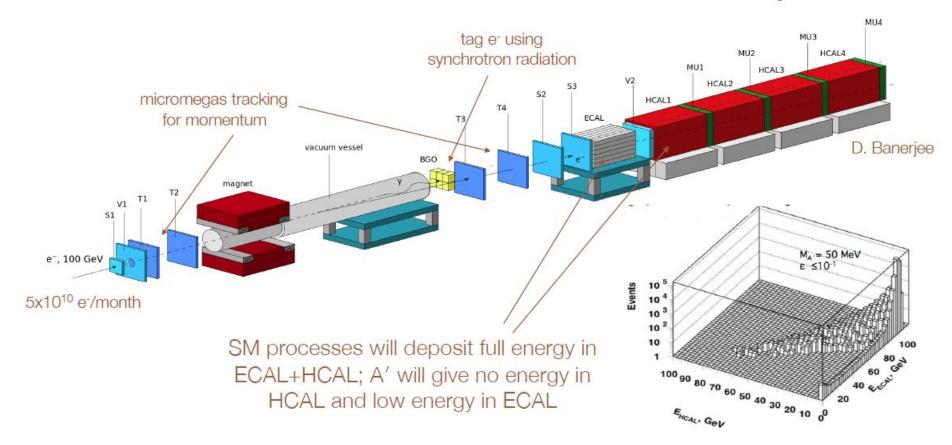
- Truth Monte Carlo vs real life
- Machine conditions
- Trigger performances (or, more generally, sustainability)

Anyway, let's take literally this estimate and locate it in time, somewhere in late 2018

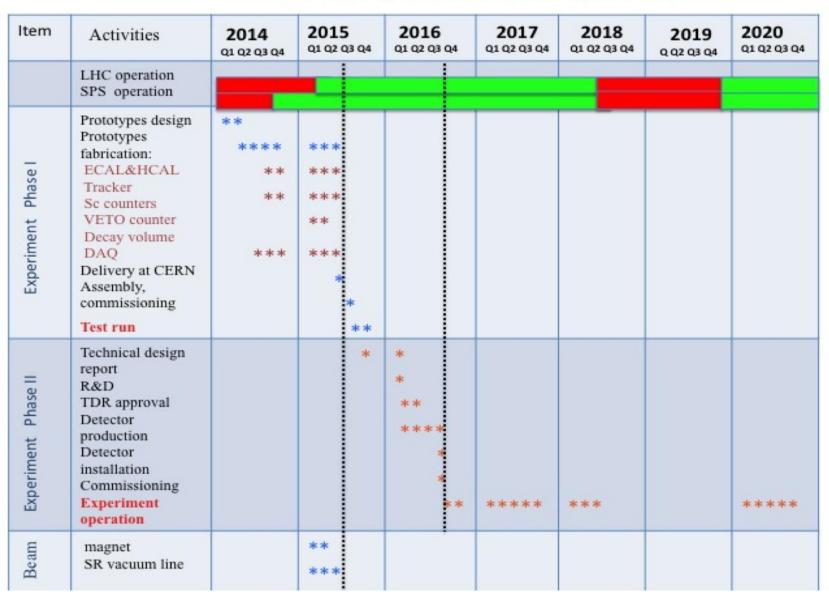
NA64 (CERN): search for invisible A' decays using an active beam dump



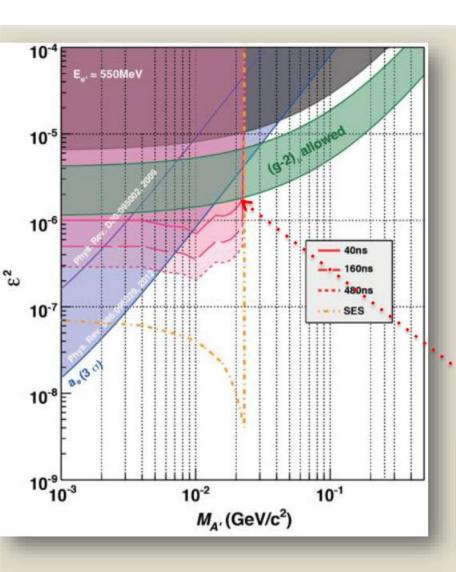
100 GeV e⁻ from H4 beam line. 2 week run in July



NA64: Preliminary schedule of the experiment



PADME (LNF) will start by ~ mid 2018 (after KLOE-2 roll out)



- Based on 2.5x10¹⁰ fully GEANT4 simulated
 550MeV e+ on target events
 - Number of BG events is extrapolated to 1x10¹³ electrons on target
- Using $N(A'\gamma) = \sigma(N_{BG})$
- δ enhancement factor $\delta(M_{A'}) = \sigma(A\Box\gamma)/\sigma(\gamma\gamma)$ with ϵ =1 due to A' mass

$$\frac{\Gamma(e^{+}e^{-} \to A'\gamma)}{\Gamma(e^{+}e^{-} \to \gamma\gamma)} = \frac{N(A'\gamma)}{N(\gamma)} \frac{Acc(\gamma\gamma)}{Acc(A'\gamma)} = \varepsilon \cdot \delta$$

 $10^{13} EOT = 6000 e^{+}/bunch \times 3.1 \cdot 10^{7} s \cdot 49 Hz$

PADME can explore in a *model-independent way the* favorite by $(g-2)_{\mu}$ band up to $M^2_{A'}=2m_eE_{e+}$

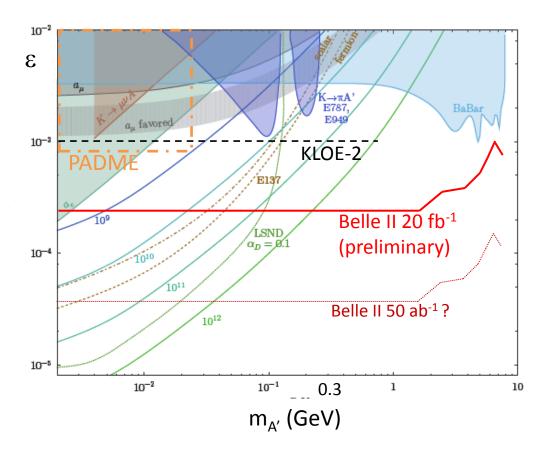
 E_{e+} =550 MeV: $M_{A'}$ < 23.7 MeV/ c^2

 E_{e+} =750 MeV: $M_{A'}$ < 27.7 MeV/ c^2

 E_{e+} =1 GeV: $M_{A'}$ < 32 MeV/ c^2

KLOE-2 implemented a single photon trigger (\sim 350 MeV threshold, barrel only) right now and will (hopefully) run with it up to the end of the data taking (approximately additional 2 fb⁻¹ by mid 2018). Rough potential reach \sim 10⁻³, but more studies are needed (it might be sustantially better or worse, depending on how well machine background is rejected).

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A very interesting picture for us, both along horizontal and vertical scales

Don't take these numbers to the last digit. Lines drawn 'by eye'

Venendo a noi ...

La misura è per ora saldamente nelle mani del gruppo di Vancouver/British Columbia (Chris & Torben)

Un possibile spazio di inserimento potrebbe riguardare l'utilizzo del KLM, nelle cui attività di commissioning saremo in ogni caso coinvolti.

KLM può essere usato sia per rivelare fotoni, sia per sopprimere i fondi derivanti da inefficienze di ECL in regioni specifiche.

Inoltre, campioni di $e^+e^- \to \mu^+ \, \mu^- \, e \, e^+e^- \to \mu^+ \, \mu^- \, \gamma$ possono consentire una «radiografia» di alta precisione del calorimetro, essendo questa ricerca molto sensibile ad inefficienze, cracks, ecc ...

Su questa base abbiamo avuto un breve approccio con Chris allo scorso GM, per il momento senza seguito.

Il gruppo di Roma3 conterà su un nuovo dottorando ad inizio di gennaio (G. De Pietro). Se decidesse di scegliere questo come argomento di tesi, naturalmente, l'attività potrebbe subire un boost.

BACKUP SLIDES

Single photon triggers

- Two level 1 triggers for single photon physics. Both exclude lowest angle ECL trigger towers (angles updated since June).
- 1 GeV: E* > 1 GeV, 2nd cluster must have E* < 0.2 GeV
 - cut on 2nd cluster is new since June B2GM
- 2 GeV: E* > 2 GeV, not a Bhabha, not γγ
 - "aggressive" Bhabha veto:
 p₁* > 3 GeV/c, p₂* > 1 GeV/c, angle >143°,
 ≥1 track associated with an ECL cluster E*>3 GeV
 - $\gamma\gamma$ veto: $E_1^*>2.5$ GeV, $E_2^*>2.5$ GeV, $>150^\circ,$ no long tracks

Single photon trigger background cross sections nb

	γγ	wide angle Bhabhas	TEEGG	Total nb
1 GeV	0.2	0.4	1.6	2.2
2 GeV	0.5	2.9	0.1	3.5

both e[±] have $\theta^* > 1^\circ$ 1 e[±] with $\theta^* < 1^\circ$

- 1 GeV trigger is dominated by true 1γ events; 2 GeV by radiative Bhabhas.
- 2 GeV trigger is also used for ISR physics