



Gravitational tests using simultaneous atom interferometers

Gabriele Rosi

*Quantum gases, fundamental interactions
and cosmology conference*

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Outline

- Introduction to atom interferometry
- The history: MAGIA experiment (apparatus and G measurement)
- Test of the weak equivalence principle in its quantum formulation
- Geometry free determination of the gravity gradient

Atom Interferometry for **gravity** measurement

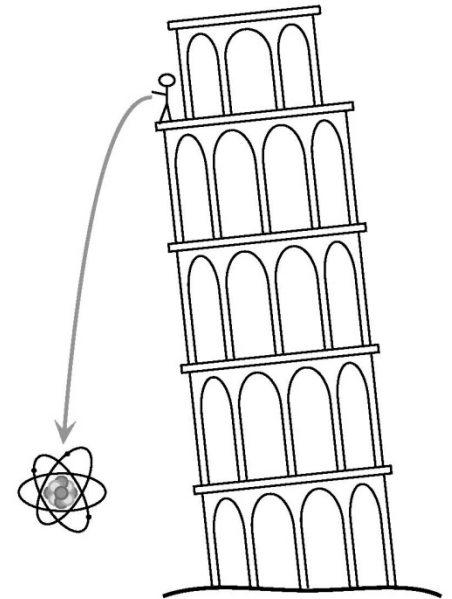
Atom Interferometry can measure accelerations

We use Cold Atoms as free falling microscopic masses

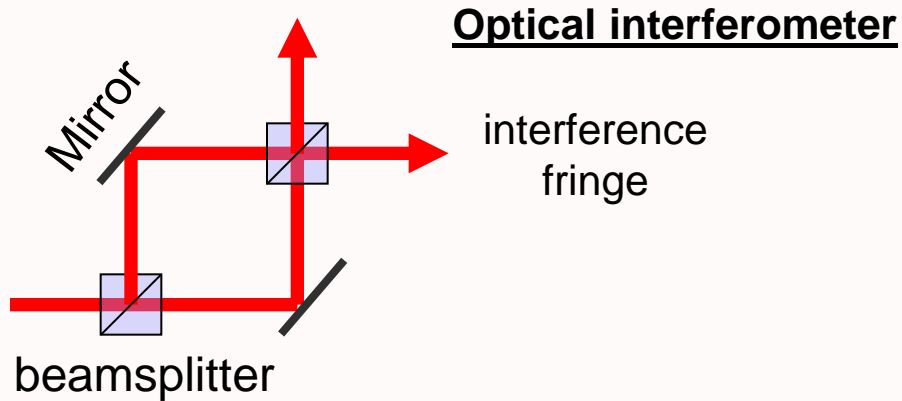
Quantum features of matter allow to improve the sensitivity (not just a time-of-flight measurement in the “classical way”)

Ingredients:

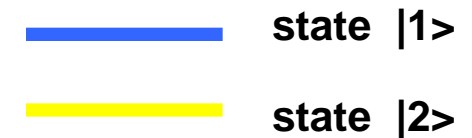
- A source of **Cold Atoms** ($\sim \mu$ K or less)
(the sample must be **slowly expanding** and **weakly interacting**)
- A laser system to cool the sample and to manipulate the wavepacket



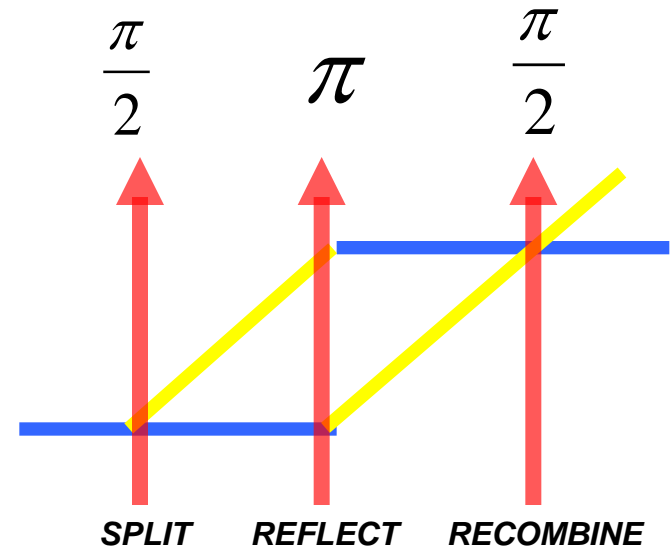
Atom / light Interferometry: the analogy



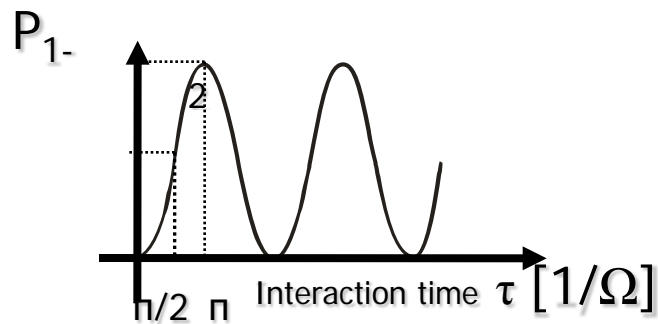
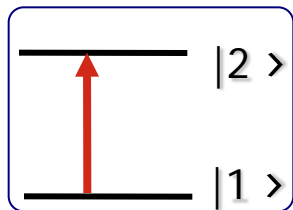
- $\pi/2$ pulse works like a beamsplitter
momentum transfer
- π pulse works like mirror;
momentum transfer



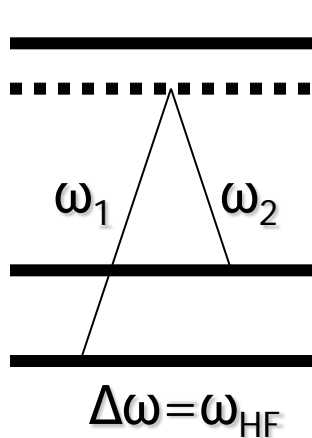
Atom interferometer



Once you have an **atomic two level system**
you can make the analogy with light looking
at the Rabi's population oscillations scheme



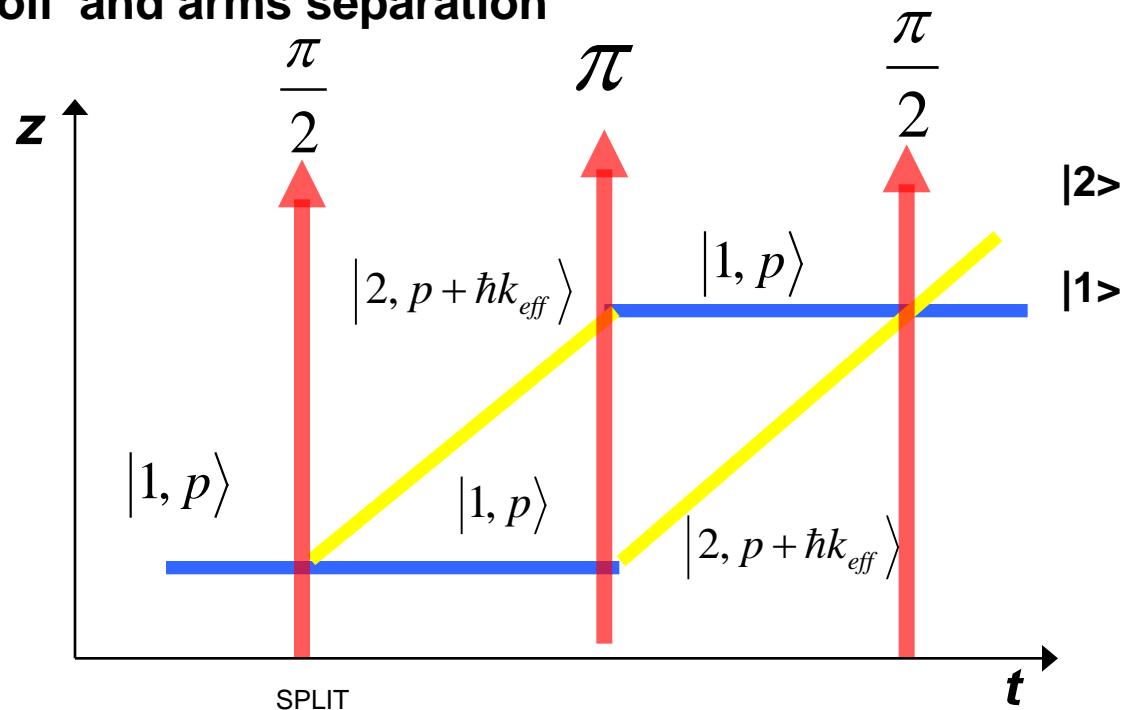
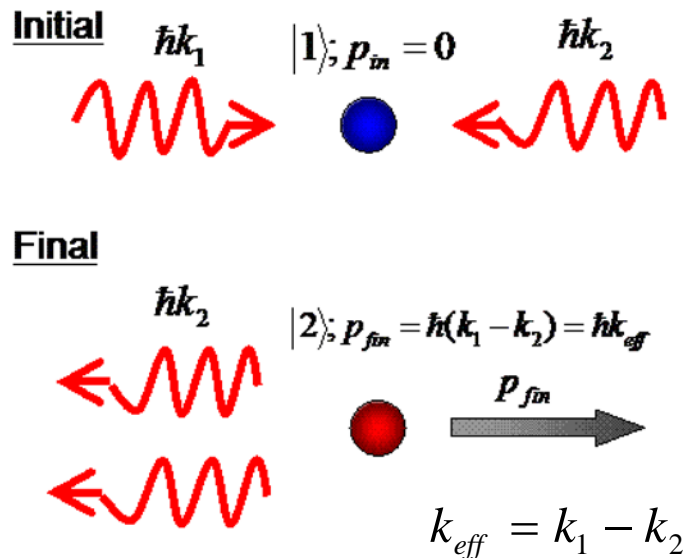
Atom Interferometry: theory



Two **hyperfine states** are coupled by **two photon RAMAN Transition** using two **couterpropagating beams**
frequency difference must be equal to hyperfine states separation

- We need to couple two **long-lived states**
- Why RAMAN: we need **large momentum recoil** (arms separation)

Momentum recoil and arms separation

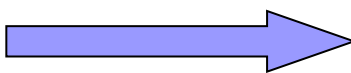


Atom Interferometry: theory

Wave function PHASE evolution

If $g = 0$ 

$$\Delta\Phi_{tot} = 0$$

$g \neq 0$ 

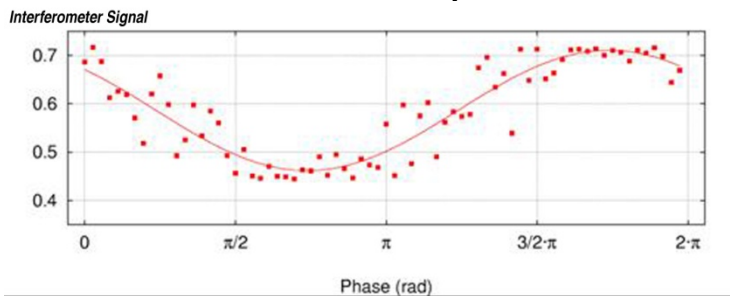
$$\Delta\Phi_{tot} = k_{eff} g T^2 + \Delta\Phi'$$

totally symmetric evolution
GRAVITY BREAKS THE SYMMETRY

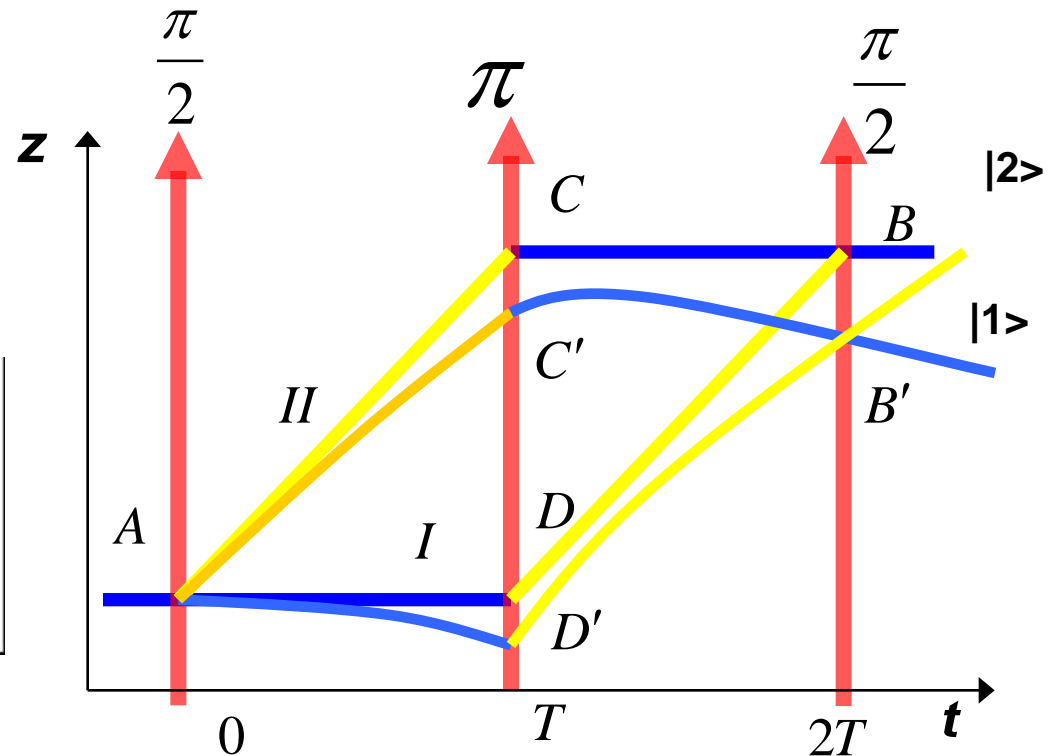
Population on final State depends
On the **interferometer phase**

$$P_1 = \frac{1}{2}(1 + \cos \Delta\Phi_{tot}) \quad P_2 = \frac{1}{2}(1 - \cos \Delta\Phi_{tot})$$

fringes



varying RAMAN laser's phase

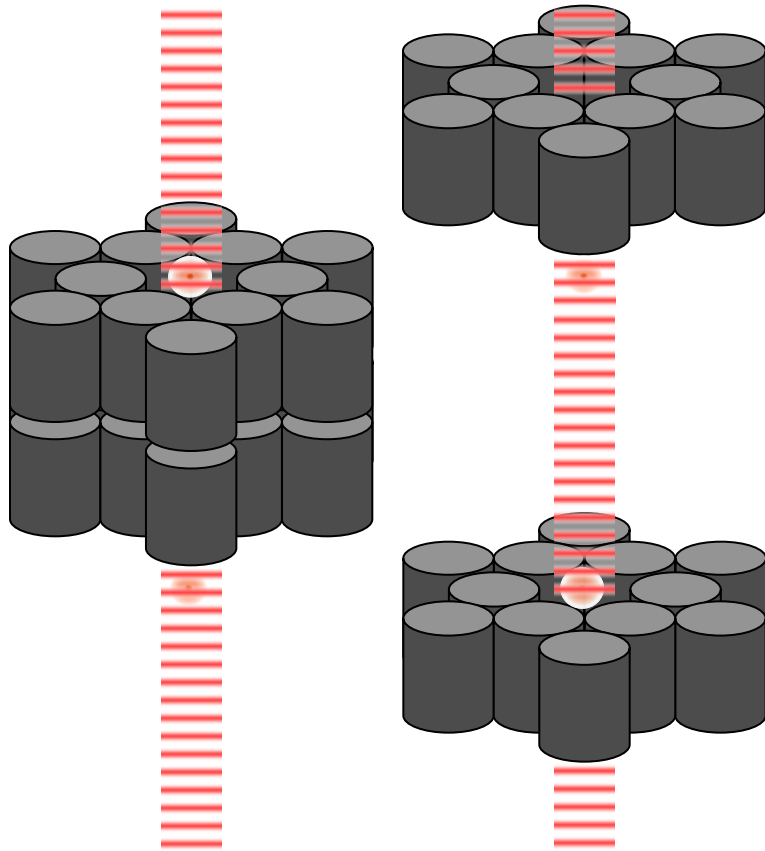




MAGIA

Misura Accurata di G mediante Interferometria Atomica

MAGIA - the procedure



SOURCE MASSES

Well-characterized tungsten cylinders

PROBE MASSES

Cold, freely falling ^{87}Rb atoms

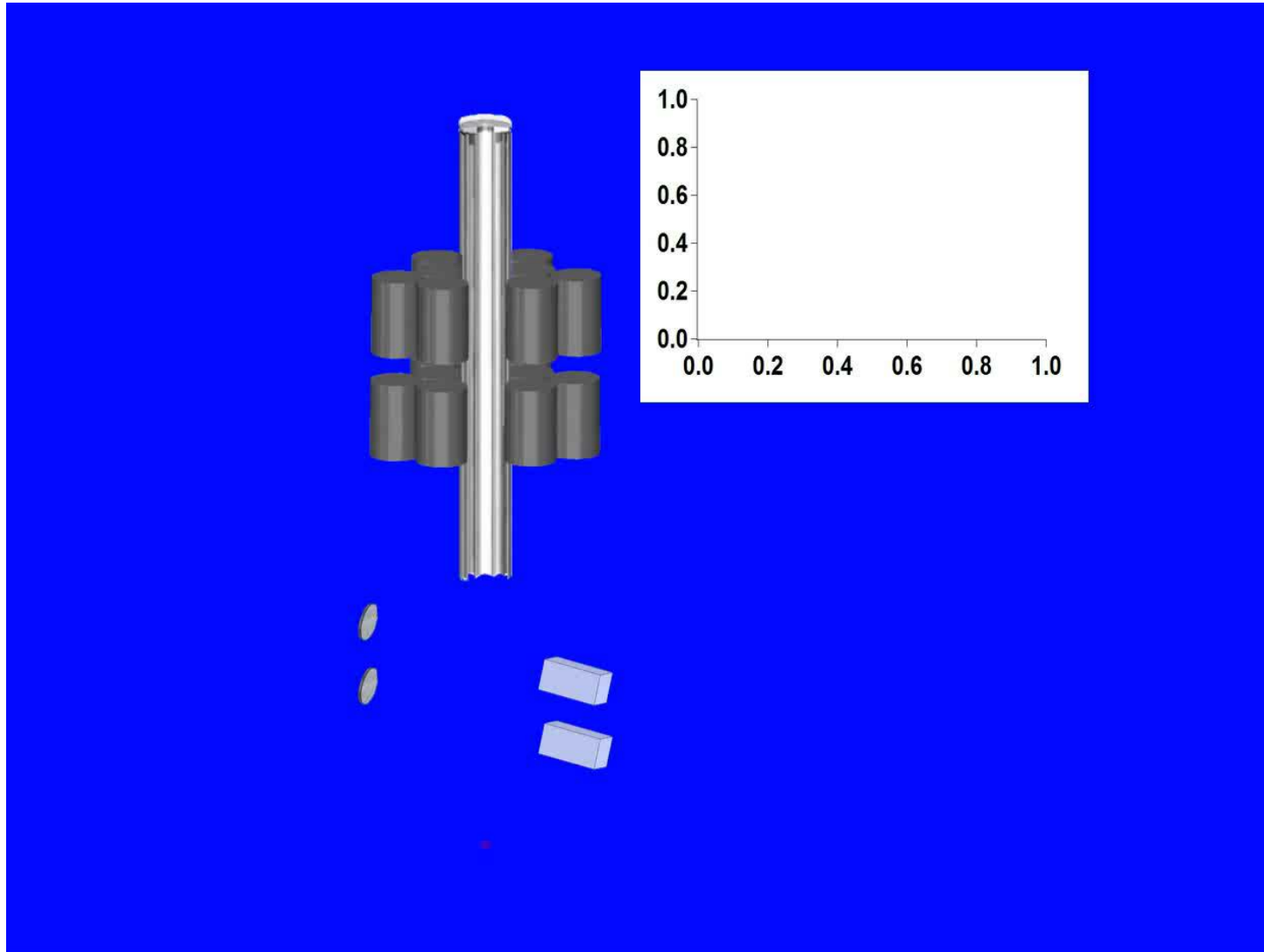
MEASUREMENT METHOD

Raman atom interferometry (local acc.)
Spatial & temporal differential scheme

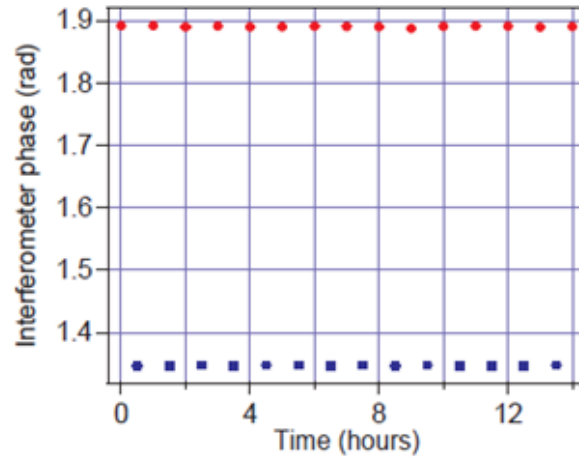
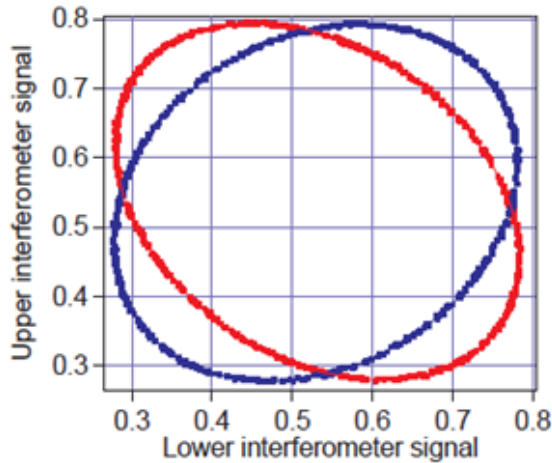
CALCULATION of gravitational attraction



MAGIA - the experimental sequence

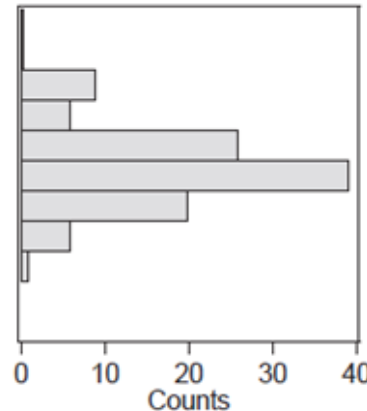
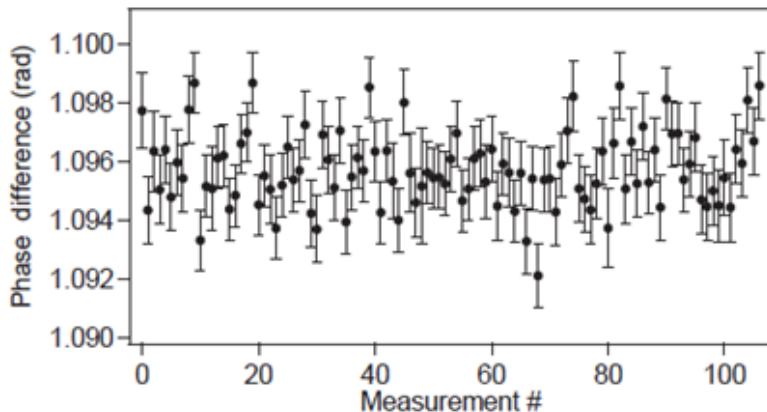


The G measurement



Features:

- Source masses modulation time: 30 mins
- Integration time: more than 100 hours over 2 weeks (July 2013)
- Sensitivity: $3 \times 10^{-9} \text{ g/Hz}^{1/2}$
- Final sensitivity: $\sim 10^{-11} \text{ g}$



$$G = 6.67191(77)_{\text{stat}}(62)_{\text{sys}} \times 10^{-11} \text{ m}^3 \text{ Kg}^{-1} \text{ s}^{-2}$$



MAGIA ADV

- Quantum test of the Weak Equivalence Principle
- Cancelling gravity gradients for future precision G measurements

Quantum formulation of the EEP

The Einstein Equivalence Principle plays a crucial role in our understanding of gravity. It can be organized into three conditions:

- Equivalence between the system's inertia and weight (WEP)
- Independence of local non-gravitational experiments from the velocity of the free falling reference frame (LLI)
- Independence of local non-gravitational experiments of their location (LPI)

How to implement EEP in a non-relativistic quantum theory?*

$$\hat{H}_{nr} = m_r c^2 + \frac{\hat{P}^2}{2m_i} + m_g \phi(\hat{Q}) \quad \leftarrow \text{Non-relativistic Hamiltonian with classical potential}$$

$$\hat{M}_\alpha := m_\alpha \hat{I}_{int} + \frac{\hat{H}_{int,\alpha}}{c^2} \quad \alpha = r, i, g,$$

Developing to the first order:

$$\hat{H}_{test}^Q = m_r c^2 + \hat{H}_{int,r} + \frac{\hat{P}^2}{2m_i} + m_g \phi(\hat{Q}) - \hat{H}_{int,i} \frac{\hat{P}^2}{2m_i^2 c^2} - \hat{H}_{int,g} \frac{\phi(\hat{Q})}{c^2}$$

Relativistic time dilation term
Gravitational time dilation term

*Zych et al. "Quantum formulation of the Einstein Equivalence Principle", arXiv:1502.00971 (2015)

Quantum formulation of the WEP

		EEP			
		WEP	LLI	LPI	# param.
Newtonian	classical & quantum	$m_i = m_g$	—	—	1
Newtonian +	classical	$m_i c^2 + E_i = m_g c^2 + E_g$	$E_r = E_i$	$E_r = E_g$	$2n - 1$
mass-energy equiv.	quantum	$m_i c^2 \hat{I} + \hat{H}_i = m_g c^2 \hat{I} + \hat{H}_g$	$\hat{H}_r = \hat{H}_i$	$\hat{H}_r = \hat{H}_g$	$2n^2 - 1$

The acceleration operator in the Heisenberg picture is:

$$\hat{a}_{\hat{H}_{test}^Q} := d^2 \hat{Q} / dt^2 = -\frac{1}{\hbar^2} [[\hat{Q}, \hat{H}_{test}^Q], \hat{H}_{test}^Q] = -\hat{M}_g \hat{M}_i^{-1} \nabla \phi(\hat{Q}) + \frac{i}{\hbar} [\hat{H}_{int,i}, \hat{H}_{int,r}] \frac{\hat{P}}{m_i c^2} + \mathcal{O}(1/c^4)$$

where

$$\hat{M}_g \hat{M}_i^{-1} = \hat{I}_{int} - \hat{\eta} \quad \text{and} \quad \hat{\eta} \approx m_g / m_i (\hat{I} + \hat{H}_{int,g} / m_g c^2 - \hat{H}_{int,i} / m_i c^2)$$

If $[\hat{H}_{int,i}, \hat{H}_{int,g}] \neq 0$ internal and external degrees of freedom can be entangled!

A quantum WEP test

Quantum formulation of WEP requires $\widehat{M}_g = \widehat{M}_i$

In QM a state of internal energy can involve superpositions of states \rightarrow for the validity of the quantum WEP we need equivalence between the off-diagonal elements of the operators.

- Let us consider a two level systems (in our case F=1 and F=2 hyperfine ground state of ^{87}Rb):

$$\widehat{M}_g \widehat{M}_i^{-1} \approx \begin{pmatrix} r_1 & r \\ r^* & r_2 \end{pmatrix}$$

r_1 and r_2 are real numbers

r is a complex number $r=|r|e^{i\varphi}$

Classical WEP is valid if $r_1 = r_2 = 1$

Quantum WEP holds if $r = 0$

A quantum WEP test

Instrument is sensitive to*:

$$a_1 = g \langle 1 | \widehat{M}_g \widehat{M}_i^{-1} | 1 \rangle = gr_1$$

$$a_2 = g \langle 2 | \widehat{M}_g \widehat{M}_i^{-1} | 2 \rangle = gr_2$$

$$a_s = g \langle s | \widehat{M}_g \widehat{M}_i^{-1} | s \rangle = g((r_1 + r_2)/2 + |r| \cos(\varphi + \gamma))$$
$$(|s\rangle = (|1\rangle + e^{i\varphi} |2\rangle)/\sqrt{2})$$

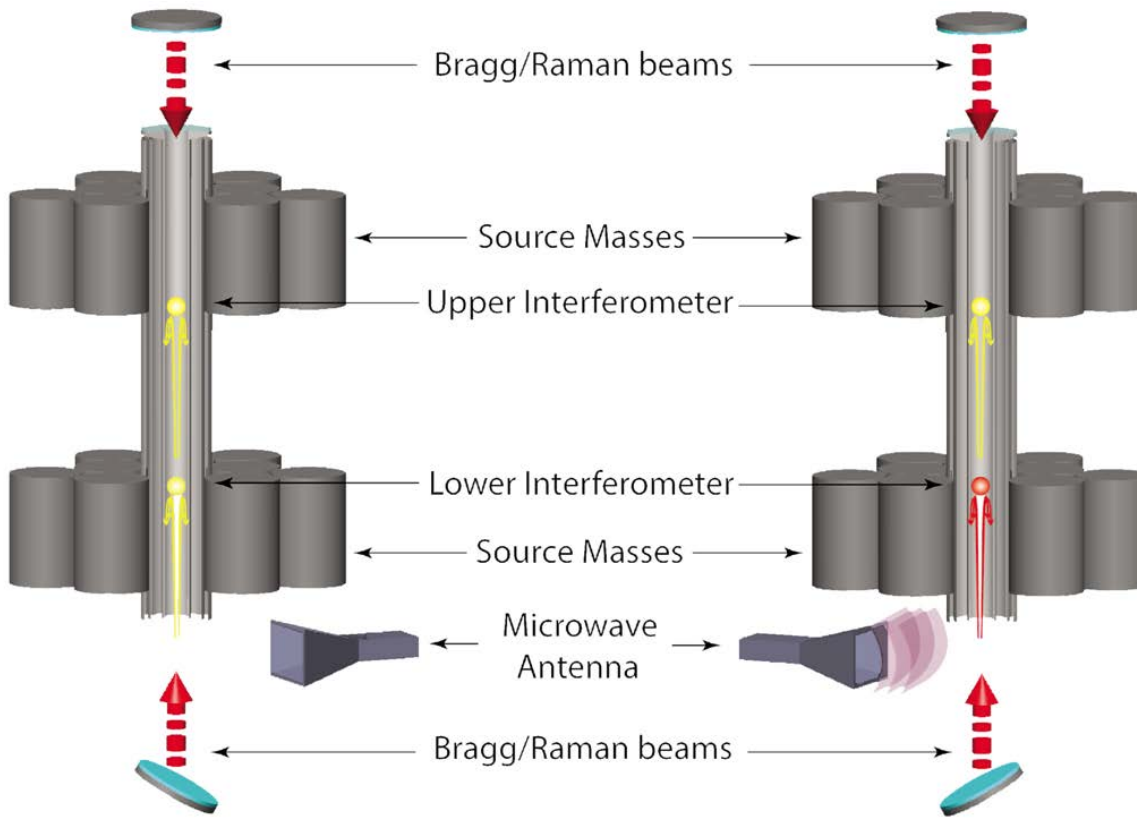
A **classical WEP violation** (introduced by diagonal elements $r_{1,2}$) emerges as a differential acceleration proportional to $r_1 - r_2$.

A **quantum WEP violation** would produce an excess phase noise on the acceleration measurements due to γ (random phase $\gg 2\pi$).

*G. Rosi et al. Nature Communications 8, 15529 (2017)

A quantum WEP test

With the **Bragg gradiometer** we compare the free fall accelerations for atoms prepared in pure hyperfine states ($F = 1$, $F = 2$) and atoms prepared in a **coherent superposition** of two different hyperfine states.



Superposition state is prepared with RF pulse
 $s = (|1\rangle + |2\rangle e^{i\gamma})/\sqrt{2}$
 γ : random phase introduced with RF pulse

G. Rosi et al. Nature Communications 8, 15529 (2017)

A quantum WEP test

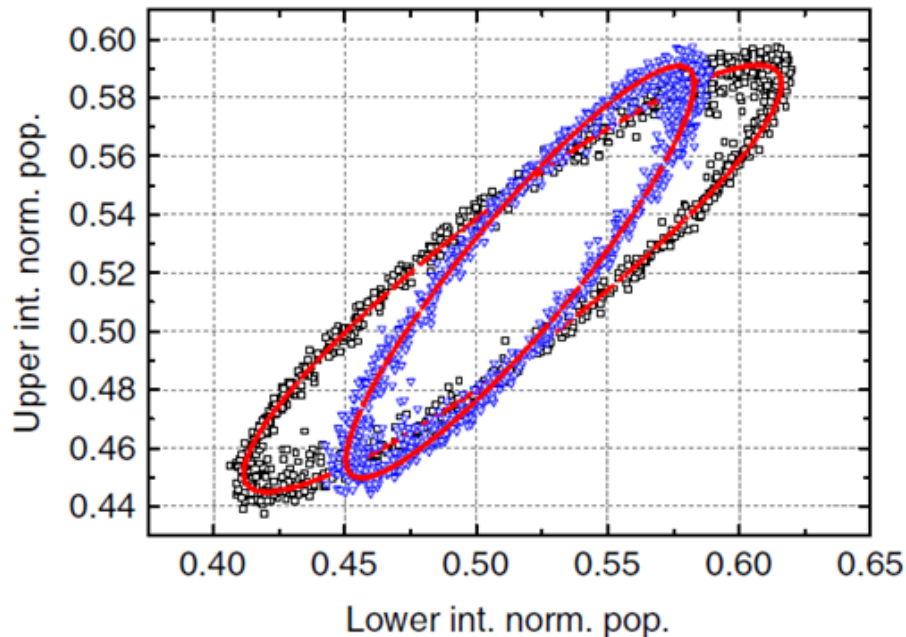
We realize three possible gradiometric configurations $\rightarrow \Phi_{1-1}, \Phi_{1-2}, \Phi_{1-s}$

Classical WEP test $\rightarrow \delta g_{1-2} \sim (\Phi_{1-1} - \Phi_{1-2}) \rightarrow \eta_{1-2} = (1,4 \pm 2,8) \times 10^{-9}$

Quantum WEP test \rightarrow Attributing all observed phase noise on 1-s ellipse to a WEP violation we estimate an upper limit for $|r| \rightarrow r \leq 5 \cdot 10^{-8}$.

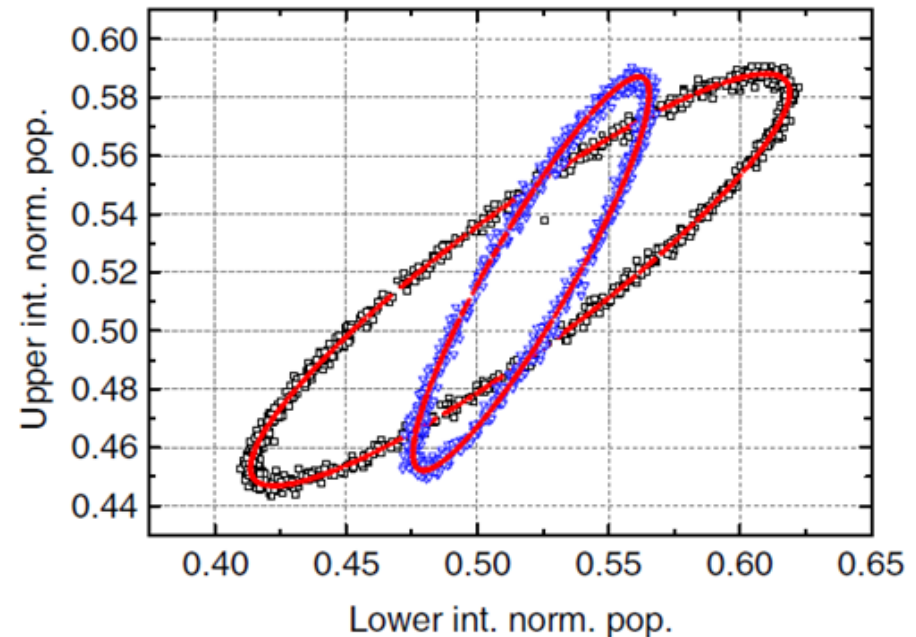
Black ellipse: 1 – 1 gradiometer

Blue ellipse: 1 – s gradiometer



Black ellipse: 1 – 1 gradiometer

Blue ellipse: 1 – 2 gradiometer



A quantum WEP test: prospects

- Energy difference between hyperfine state very tiny! (28 μeV)
- Likely, the commutator $[\hat{H}_{int,i}, \hat{H}_{int,g}]$ is proportional to the typical magnitude of H and therefore large ΔE yields to larger effects!
- Next step: **Sr interferometer on clock transition!* $\Delta E = 1.8 \text{ eV}$!**
- Next next step: entangled states between different isotopes?

*L.Hu et al., Submitted to PRL (2017)



MAGIA ADV

- Quantum test of the Weak Equivalence Principle
- Cancelling gravity gradients for future precision G measurements

Cancelling gravity gradients

■ The tidal forces on the atoms in a uniform gravity field and gradient modify the wavepacket trajectories. The gravimetric phase shift is

$$\varphi = k_{\text{eff}} g T^2 + k_{\text{eff}} \Gamma_{zz} (z_0 + v_0 T) T^2$$

z_0, v_0 initial atomic position and velocity.

■ The error on z_0 and v_0 is one of the **major sources of noise and systematics**:

• For WEP tests at 10^{-15} level is required a control on z_0 and v_0 of 1 nm and 0,3 nm/s.

• In the AI determination of $G^{[1]}$ one of the major sources of systematic error arises from the limited control on the thermal cloud degrees of freedom.

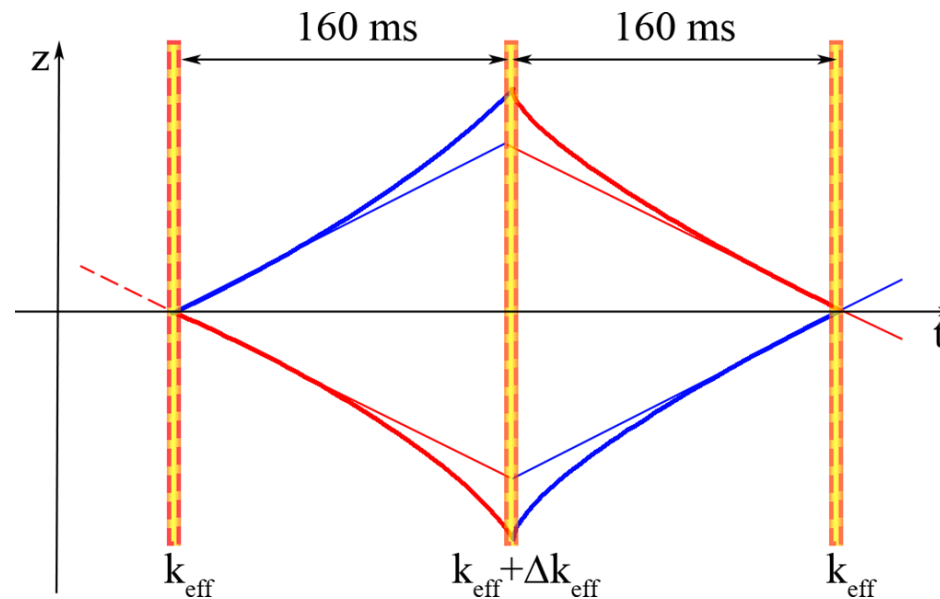
[1] G. Rosi, F. Sorrentino, L. Cacciapuoti, M. Prevedelli & G. M. Tino, "Precision measurement of the Newtonian gravitational constant using cold atoms", Nature **510**, 518-521 (2014).

Cancelling gravity gradients

Readapting the effective wave vector k_{eff} of the π pulse it is possible to **compensate the effect of $\Gamma_{zz}^{[1]}$**

$$\Delta k_{\text{eff}} = (\Gamma_{zz} T^2/2) k_{\text{eff}}$$

φ no longer depends on z_0 and v_0 .

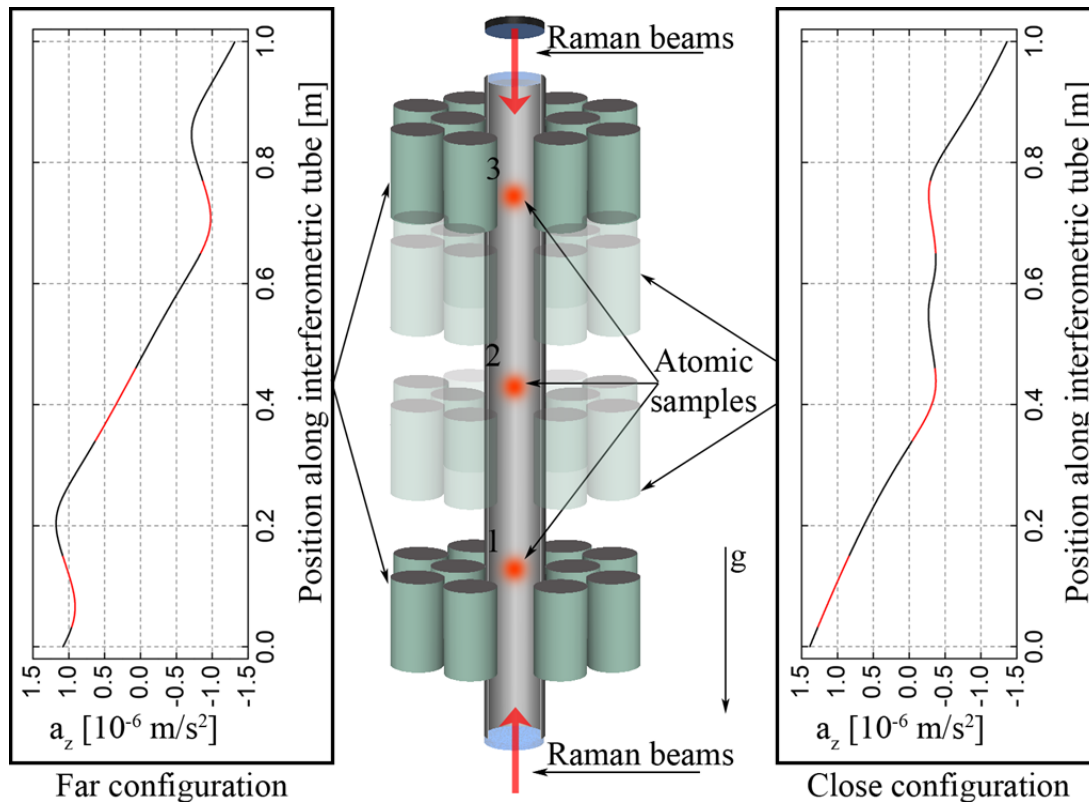


This procedure simulates the effect of a gravity gradient on the atomic trajectories. We implement it to measure **gravity gradients, gravity curvature**

[1] A. Roura, Phys. Rev. Lett. **118**, 160401 (2017).

Cancelling gravity gradients

We simultaneously interrogate **three clouds** with the Raman interferometer for **two source masses configurations**. During the π pulse the **frequency of the Raman lasers is changed** by $\Delta\nu^*$.

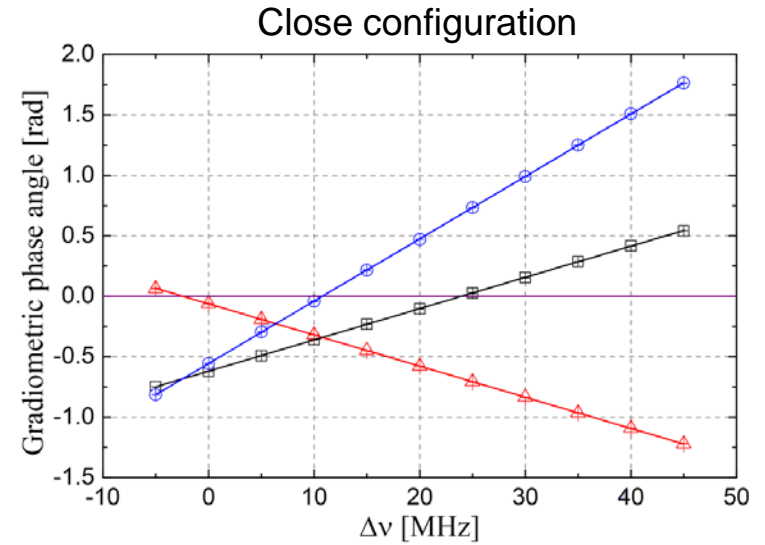
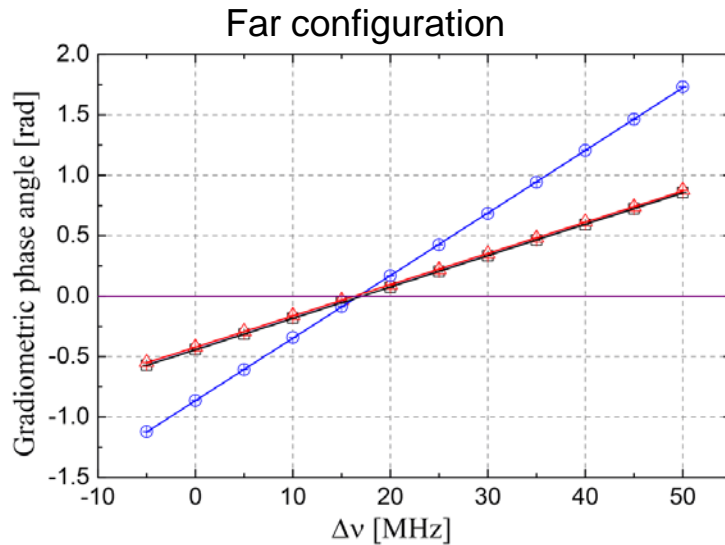


For the three gravity gradiometers (1-2, 2-3, 1-3) we measure the linear dependence $\Phi(\Delta\nu)$ vs $\Delta\nu$ (Φ gradiometric phase)

Gravity gradient is translated into a frequency!*

Cancelling gravity gradients

- Final results*:



	$\Gamma_{23} [10^{-6} \text{ s}^{-2}]$	$\Gamma_{12} [10^{-6} \text{ s}^{-2}]$	$\Gamma_{13} [10^{-6} \text{ s}^{-2}]$
Far	$-3,32 \pm 0,02$	$-3,48 \pm 0,01$	$-3,40 \pm 0,01$
Close	$0,497 \pm 0,006$	$-4,87 \pm 0,01$	$-2,193 \pm 0,006$

- Measurements also provide gravity gradient sign.
- As expected $\Gamma_{13} = (\Gamma_{12} + \Gamma_{23})/2$.
- Measured cloud distances and gravity curvature in close configuration:

$$d_{23} = (307,2 \pm 0,3)\text{mm} \quad d_{12} = (308,6 \pm 0,4)\text{mm}$$

$$\xi_{\text{close}} = (\Gamma_{23} - \Gamma_{12})/d = (1,743 \pm 0,004) \times 10^{-5} \text{ m}^{-1}\text{s}^{-2}$$

*D'amico et al. (paper submitted)

Conclusions and prospects

- In presence of a **linear** gravity gradient the zero crossing frequency is **independent** by clouds positions and velocities
- The main systematic effect in G determination with cold atoms arises from the **required knowledge** of atomic distribution
- Can we fabricate a source mass in order to produce an almost linear acceleration profile? **Yes**

➔ Determination of G at 10 ppm possible with thermal clouds!*

Recipe: double differential zero-crossing frequency determination

*G. Rosi, Metrologia (2017)



Thank you for the attention!

<http://www.coldatoms.lens.unifi.it>