

The exciting adventure of condensed-matter and optical analogs of gravitational systems

The tale of Navier and Stokes meeting Heisenberg at Hawking's place

Iacopo Carusotto

INO-CNR BEC Center, Trento, Italy

Work done in collaboration with:

- Stefano Finazzi, Pjotrs Grisins, Pierre-Élie Larré, Alessio Recati, S. Butera (BEC Center, Trento, Italy)
- Roberto Balbinot, Serena Fagnocchi (Università di Bologna, Italy)
- Alessandro Fabbri (IFIC - Univ. de Valencia and CSIC, Spain)
- Nicolas Pavloff, Renaud Parentani (Université Paris-Sud, Orsay, France)
- Dario Gerace (Università di Pavia, Italy)
- Hai-Son Nguyen, Alberto Amo, Jacqueline Bloch (LPN-CNRS, Marcoussis, France)
- M. Tettamanti, A. Parola, S. Cacciatori (Univ. Insubria, Como, Italy)

Acoustic (and fishic) black hole horizon

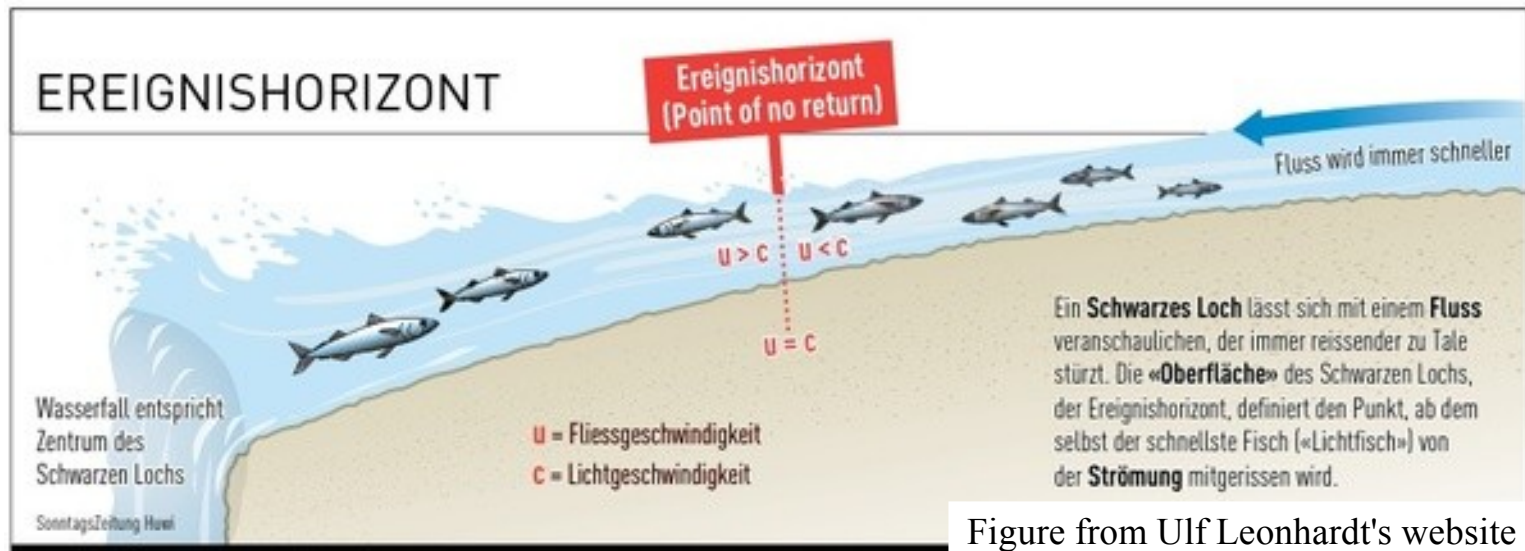


Figure from Ulf Leonhardt's website

Excitations (i.e. fish) propagate (i.e swim) at $v=c_s \pm v_{\text{flow}}$

- **Horizon region** separating **sub-sonic** (i.e. **sub-fishic**) flow (upstream) from **super-sonic** (i.e. **sub-fishic**) flow (downstream)
- **Excitations** (i.e. **fish**) in super-sonic (i.e. super-fishic) region **can not travel** (i.e. **swim**) **back** through **horizon**
- What happens with **quantum fishes** ? **Hawking radiation** ?

Behavior analogous to **astrophysical black hole horizon**

Mathematical framework

Sonic dispersion of phonons in superfluid → relativistic eq for BEC phase

$$\frac{1}{\sqrt{-G}} \partial_\mu \left[\sqrt{-G} G^{\mu\nu} \partial_\nu \right] \phi(x, t) = 0$$

mathematically equivalent to light propagation in curved space-time metric

$$ds^2 = G_{\mu\nu} dx^\mu dx^\nu = \frac{n(x)}{c_s(x)} \left[-c_s(x)^2 dt^2 + (d\vec{x} - \vec{v}(x) dt)(d\vec{x} - \vec{v}(x) dt) \right]$$

Once quantized → quantum field theory in a curved space time

Simplest analog black hole geometry:

- one-dimensional geometry, flow in the +x direction
- $v(x)/c(x)$ increases along +x direction; horizon where $v(x_H) / c(x_H) = 1$

As in astrophysical black holes → Hawking emission at $T_H = \frac{\hbar c^3}{8\pi k_B G M}$

What happens in analogs ?

Some examples of analog models

- Surface waves on flowing classical fluids, e.g. in water tanks
- Nonlinear optical systems:
moving refractive index perturbations created by strong optical pump
- Other systems:
 - Ion chains in ring traps
 - Arrays of circuit-QED cavities
- Quantum fluids:
 - BECs of ultra-cold atoms
 - quantum fluids of light, e.g. gas of (dressed) photons in microcavity
 - propagating light in bulk nonlinear media → see Victor's talk later on

1 - Surface waves on a classical fluid



Phase and group velocities of **surface waves** $c_s = \sqrt{gh}$

Flow speed v also varies with h

→ both modulated by **tank bottom profile**

Classical Hawking processes:

- positive to negative **mode conversion**
- conversion amplitude claimed to follow **thermal law**
- Serious concerns in L.-P. Euvé's PhD thesis (Poitiers 2017)
See also arXiv:1409.3830 (with Parentani's group)

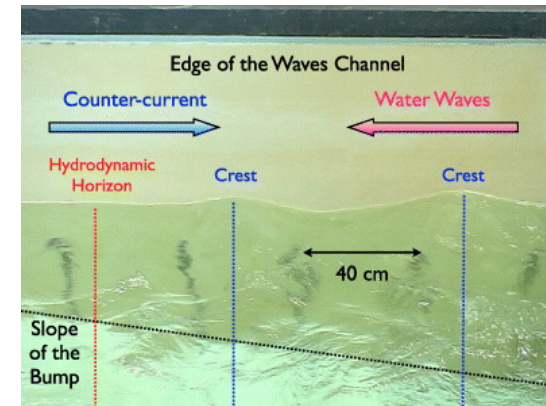
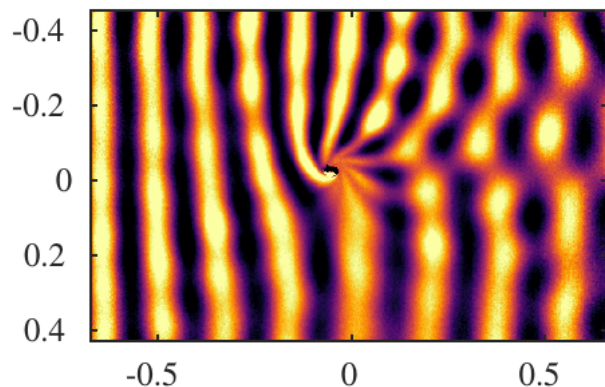


Figure: Rousseaux et al., NJP 2010

Limited hope of detecting quantum features



Expt. observ. of superradiant scattering from rotating BH

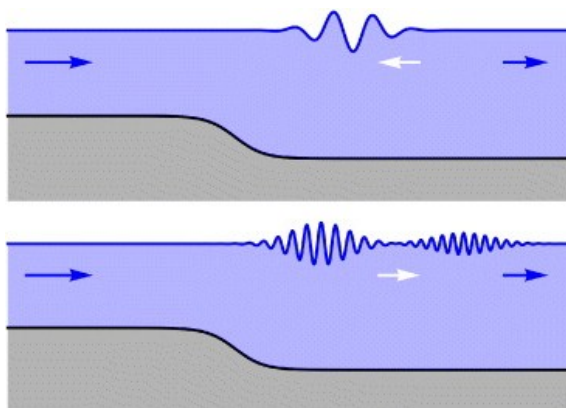
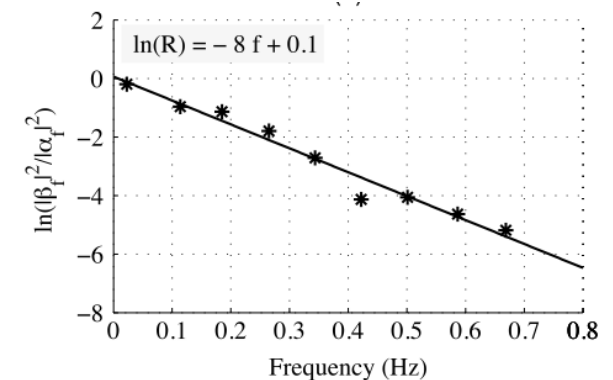


Figure: Leonhardt, Robertson, NJP 2012



Classical Hawking processes reported in Weinfurter *et al.* PRL 2011.

Correlation expts in arXiv:1511.08145

2 – Nonlinear optical systems

The idea

(Philbin et al., Science 2008)

Strong optical pulse in nonlinear crystal:

- propagates at v_g
- optical $\chi^{(3)}$ of medium modifies $n \rightarrow n + \delta n$

If $c/n > v_g > c / (n + \delta n)$

- other frequencies feel **horizon**
- **quantum fluctuations** \rightarrow **analog HR**

Experimental claim of HR from Como group
(Belgiorno et al. PRL 2010)

Not yet full consensus on interpretation:

- Dispersion $n(\omega)$ of **silica** very complicate
- **Other emission processes** present...
...with not too different **spectral features**

Is observed emission really Hawking?

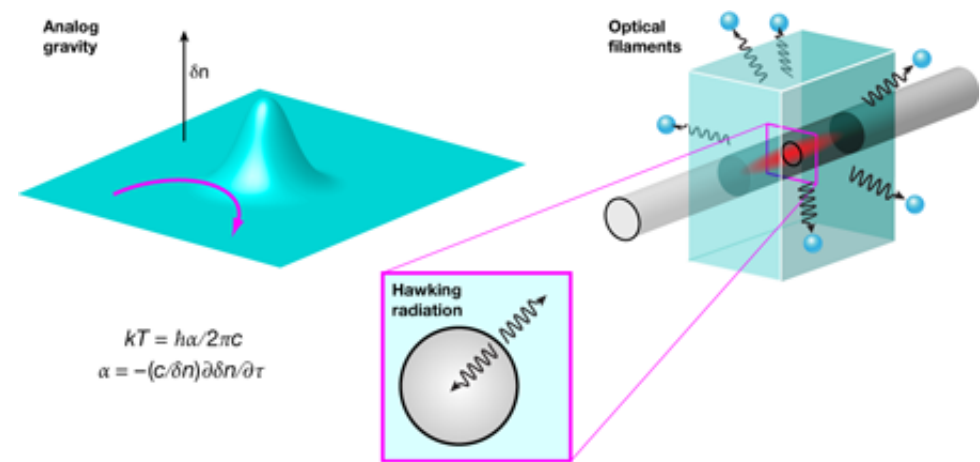
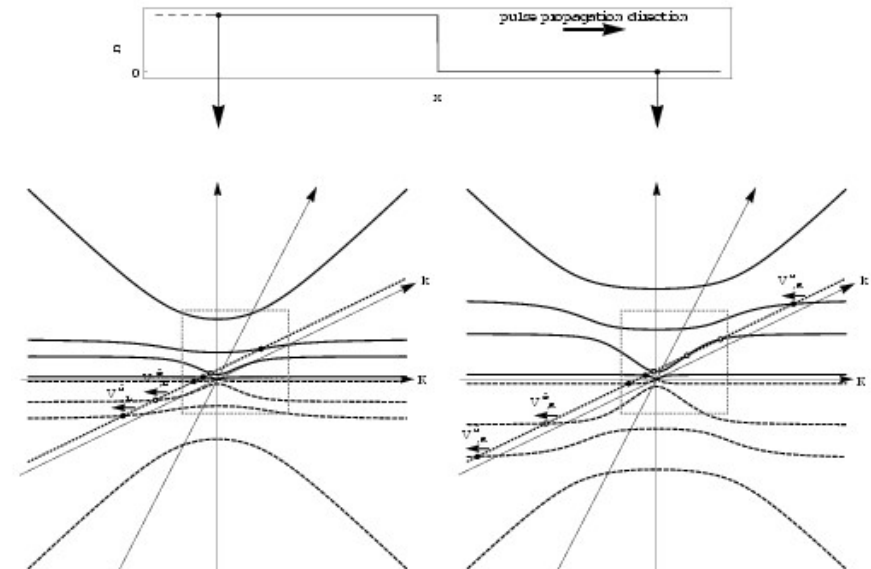


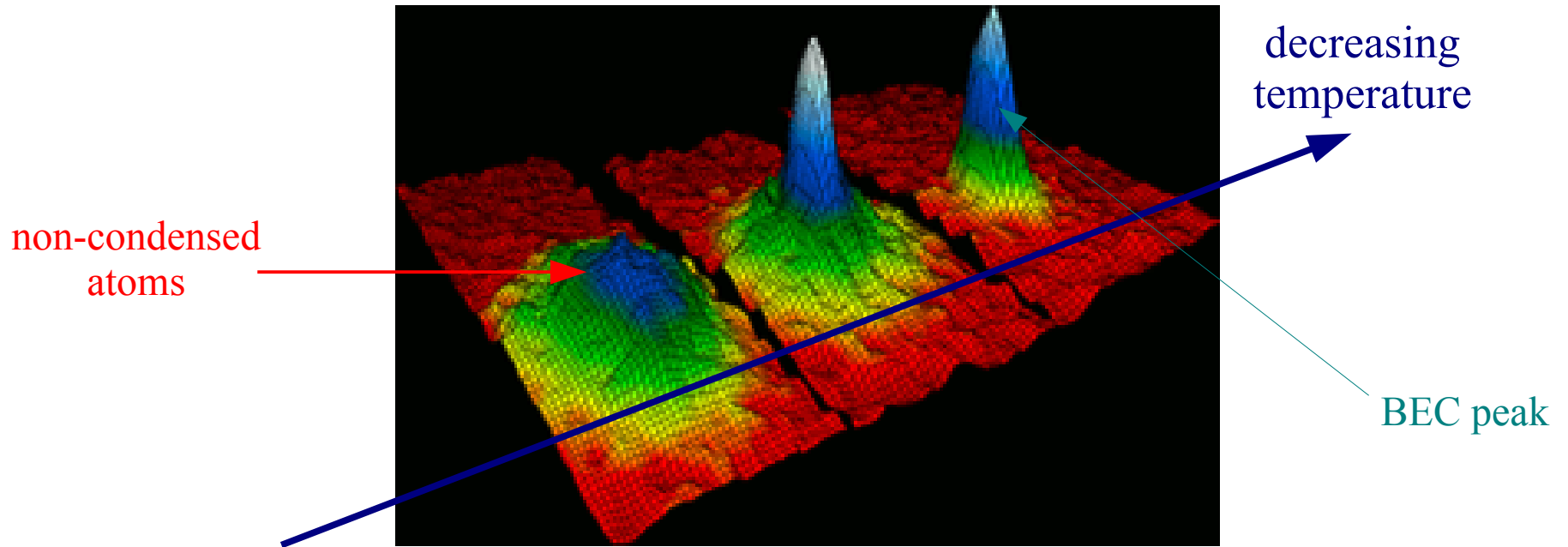
Figure adapted from Belgiorno et al. PRL 2010



Part 1:

Superfluid ultracold atomic gases

What are superfluid ultracold atomic gases?



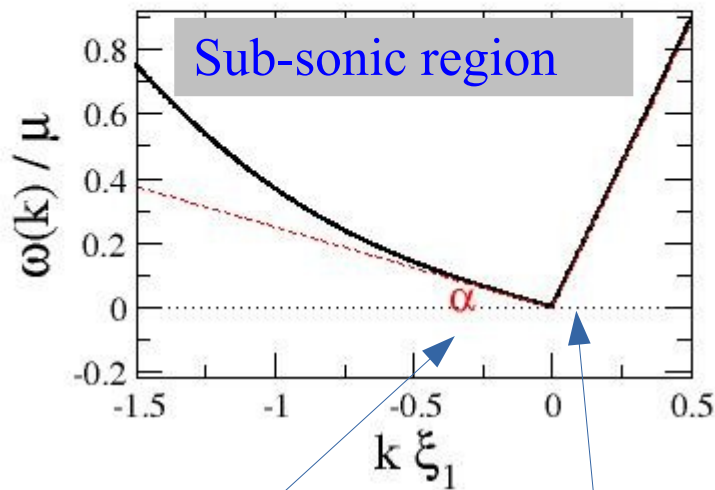
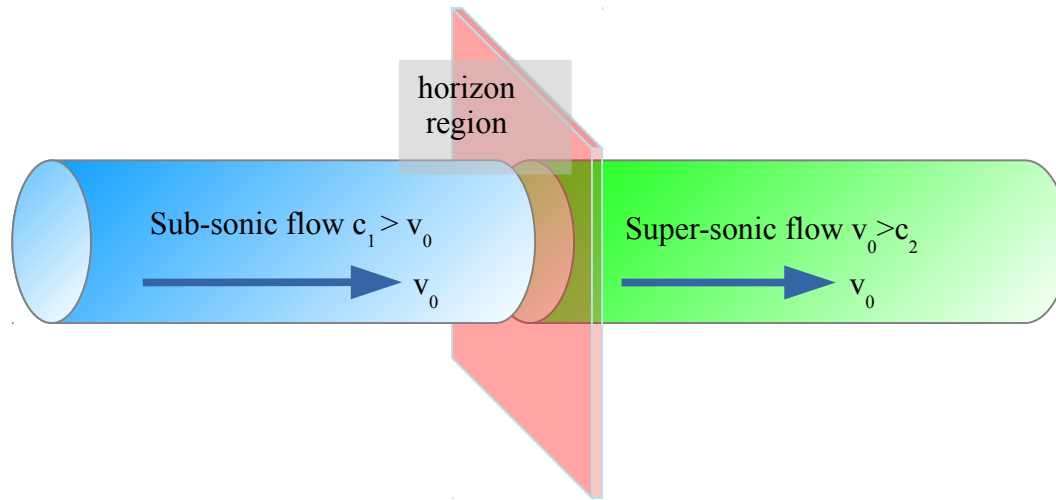
BEC in Rb atomic cloud via momentum distribution images
Figure from JILA group

At nano-K temperatures of laser + evaporation cooled atomic clouds:

- Bose-Einstein condensate with macroscopic occupation of $k=0$ orbital
- superfluid features with phonon mode with sonic dispersion
- Ultra-low temperature: long-lived quantum coherence of excitations
- Trans-sonic flow \rightarrow Hawking T_H in nK range for μm -size clouds

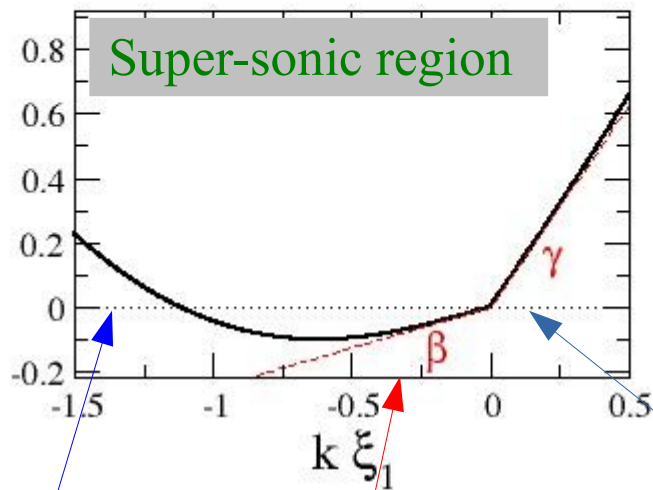
$$T_H = \frac{\hbar}{4\pi k_B c_s} \left[\frac{d}{dx} (c_s^2 - v^2) \right]_H$$

A significant detail...



$$v_g = v_0 - c_1 < 0$$

$$v_g = v_0 + c_1 > 0$$



$$v_g = v_0 - c_2 > 0$$

$$v_g = v_0 + c_2 > 0$$

Sound can not propagate back !

New feature of atomic BEC: single particle excitations can emerge from black hole !!

This raises interesting fundamental questions...

Standard derivations of Hawking radiation often assume:

- linear dispersion $\omega(k) = c |k|$ at all length scales
- infinite blue shift of modes at horizon
- relativity and QFT valid up to arbitrary energies

These assumptions violated in BEC-based analogs:

- is HR robust w/r to deviation from hydrodynamic dispersion?
- what is role of single particle nature of high-k excitations?
- at closer look: microscopic mechanism of HR completely different...

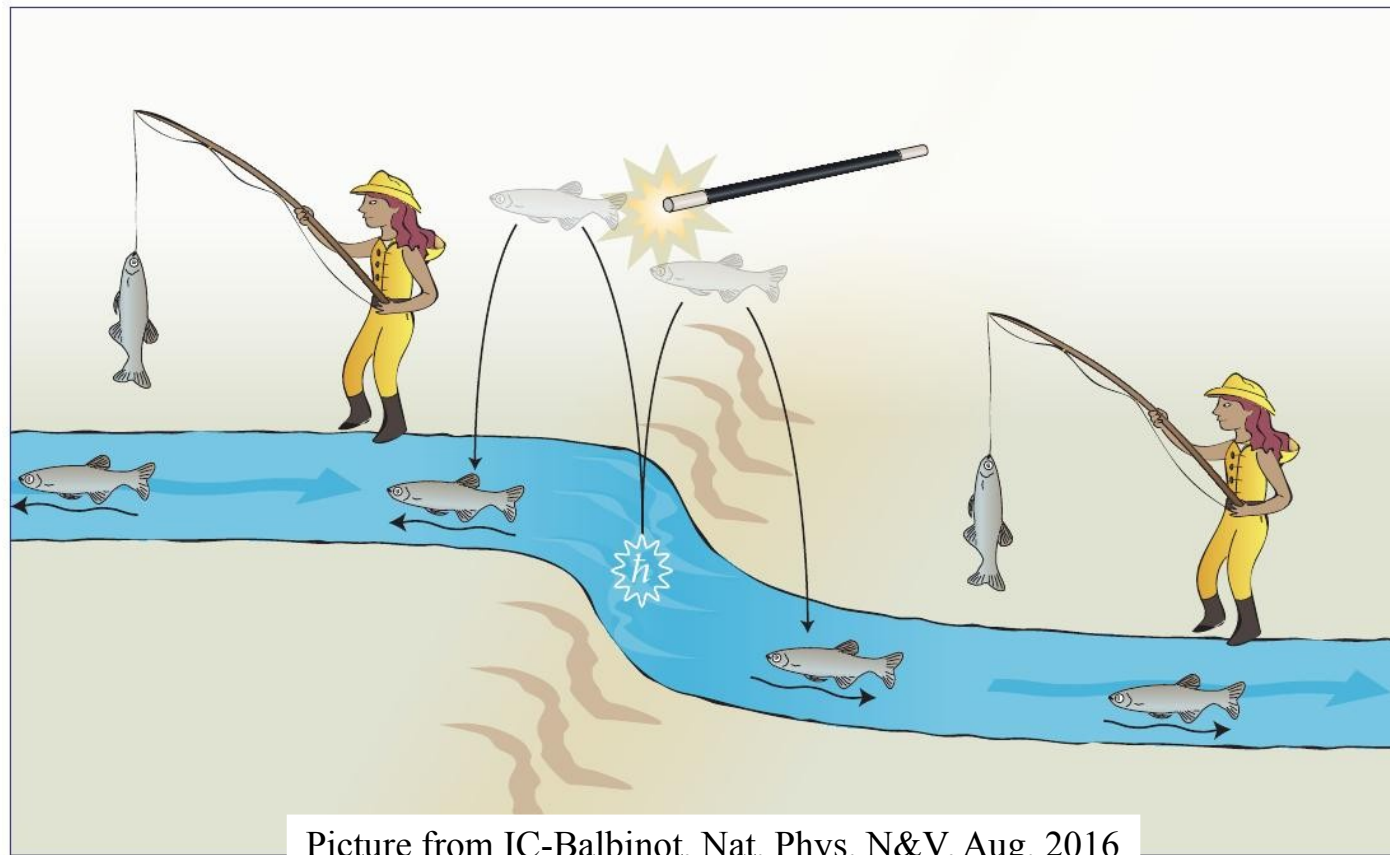
Some open questions:

- thermal HR spectrum modified by “Planck-scale” physics?
- does this provide new features in BH signal in LHC ?
(and possibly contribute to save the world)

W. G. Unruh, Phys. Rev. D 51, 2827 (1995); R. Brout et al., Phys. Rev. D 52, 4559 (1995)
T. A. Jacobson and R. Parentani, *An echo of Black Holes*, Scientific American, Dec. 2005.

Persone	Società	Ricerca
	Rössler, la Cassandra della Fisica	Aveva previsto tutti i pericoli ma nessuno l'ha ascoltato. Oggi l'amara rivincita. "Avrei preferito avere torto". IMMAGINI - VIDEO

How to detect Hawking radiation in atomic gases?



- Hawking radiation → correlated pairs generated simultaneously at the horizon
- HR isolated from background of thermal and noise phonons by measuring correlations on opposite sides of horizon
- In the picture: Hawking fish are caught simultaneously by the two fisherwomen!

This idea put into formulas...

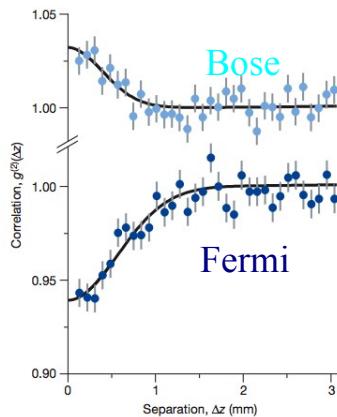
Density-density correlation function $G_2(x, x') = \frac{\langle : \delta n(x) \delta n(x') : \rangle}{\langle n(x) \rangle \langle n(x') \rangle}$

Prediction of **gravitational analogy** (Balbinot *et al.* PRA 2008):

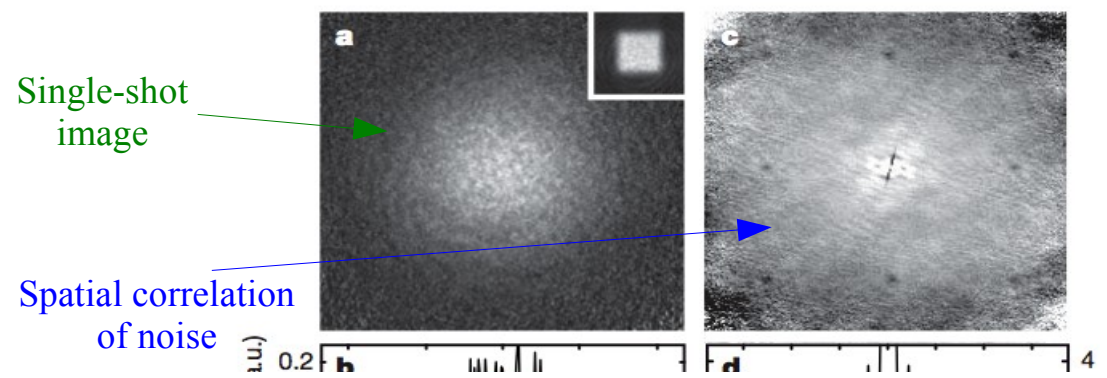
– **entanglement in Hawking pairs** gives **long-range in/out correlations**

$$G_2(x, x') = - \frac{\xi_1 \xi_2}{16 \pi c_1 c_2} \frac{k^2}{\sqrt{n^2 \xi_1 \xi_2}} \frac{c_1 c_2}{(c_1 - v)(v - c_2)} \cosh^{-2} \left[\frac{k}{2} \left(\frac{x}{c_1 - v} + \frac{x'}{v - c_2} \right) \right]$$

– allows to **isolate Hawking phonons** from background of **incoherent thermal phonons**



Jeltes *et al.*, Nature 445, 402 (2007)



Rom *et al.*, Nature 444, 733 (2006)

Measurement of density correlations **experimentally demonstrated**:

- **Atomic HB-T**: positive correlation due to **thermal Bose atoms** (negative for fermions)
- **Noise correlations** in TOF picture after **expansion from lattice**

Ab initio numerics: Wigner-Monte Carlo

At $t=0$, homogeneous system:

- **Condensate** wavefunction in plane-wave state
- **Quantum + thermal fluctuations** in plane wave Bogoliubov modes
- **Gaussian α_k** , variance $\langle |\alpha_k|^2 \rangle = [2 \tanh(E_k / 2k_B T)]^{-1} \rightarrow 1/2$ for $T \rightarrow 0$.

$$\psi(x, t=0) = e^{i k_0 x} \left[\sqrt{n_0} + \sum_k \left(u_k e^{i k x} \alpha_k + v_k e^{-i k x} \alpha_k^* \right) \right]$$

At later times: conservative (for atoms!) evolution under GPE

$$i \hbar \partial_t \psi(x) = -\frac{\hbar^2}{2m} \partial_x^2 \psi(x) + V(x) \psi(x) + g(x) |\psi(x)|^2 \psi(x)$$

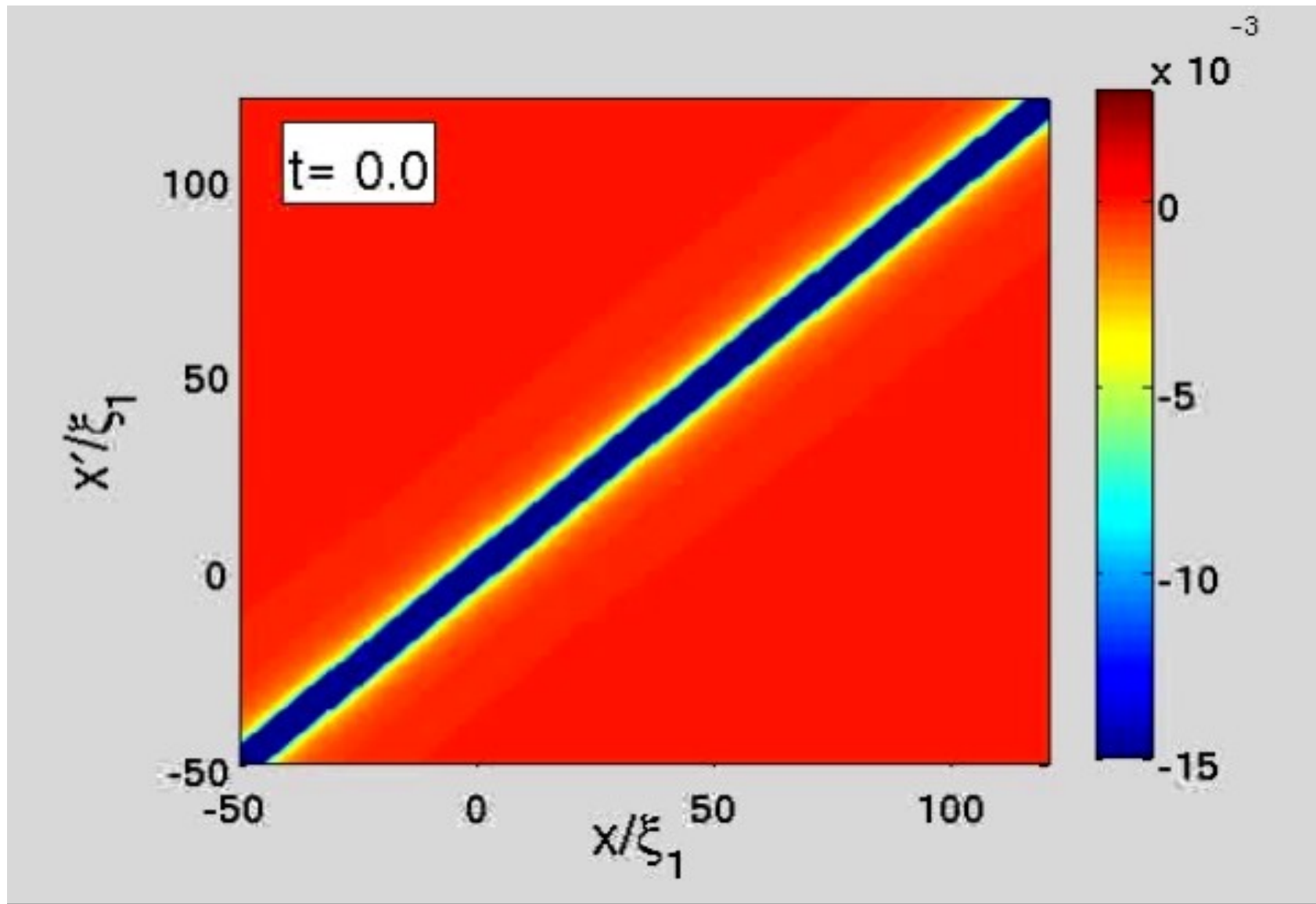
Expectation values of observables:

- Average over noise provides **symmetrically-ordered observables**

$$\langle \psi^*(x) \psi(x') \rangle_W = \frac{1}{2} \langle \hat{\psi}^\dagger(x) \hat{\psi}(x') + \hat{\psi}(x') \hat{\psi}^\dagger(x) \rangle_Q$$

Equivalent to Bogoliubov, but can explore longer-time dynamics

Density correlations: the movie

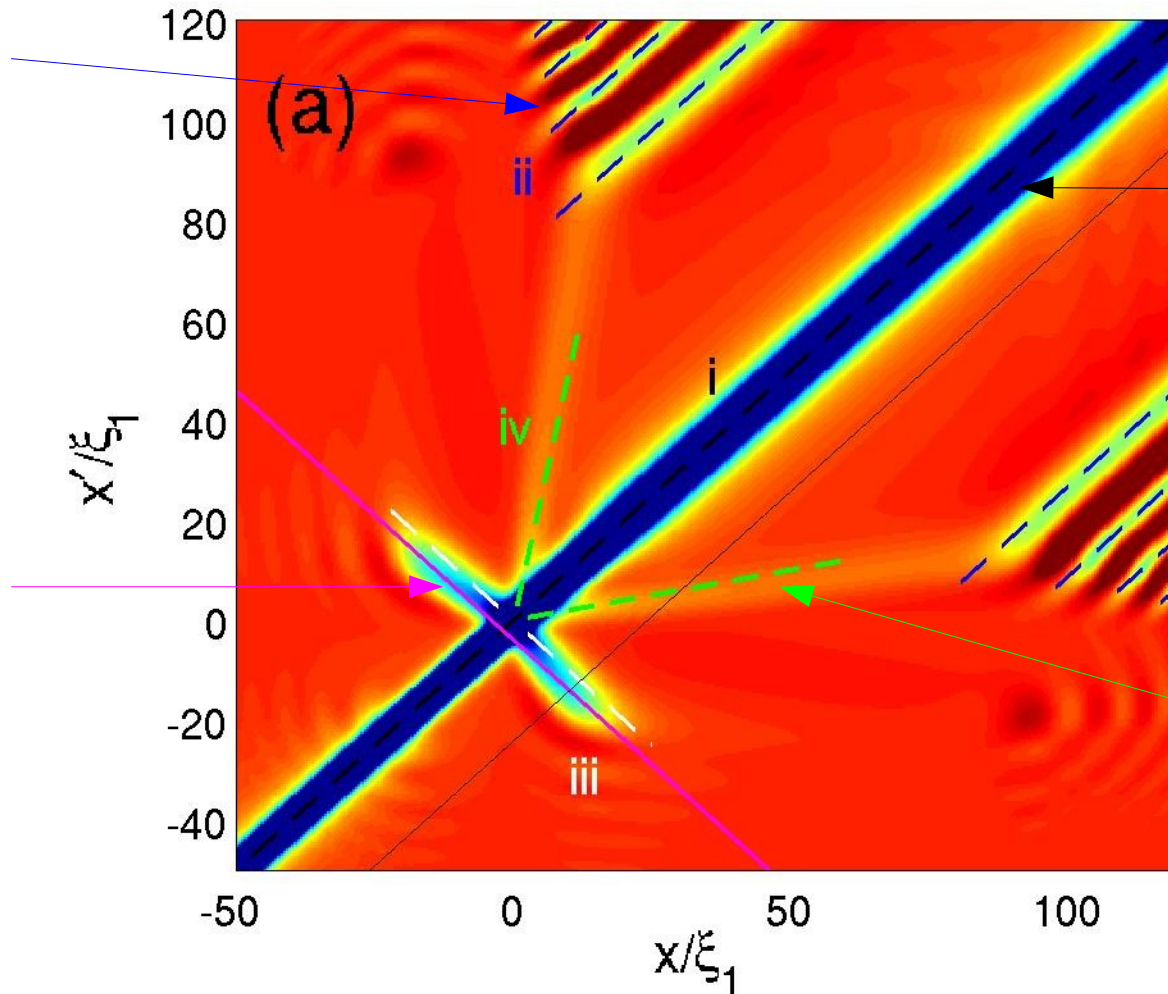


A snapshot of density correlations

Density plot of : $(n \xi) * [G^{(2)}(x,x') - 1]$

(ii)
Dynamical
Casimir
emission

(iii)
Hawking
in / out

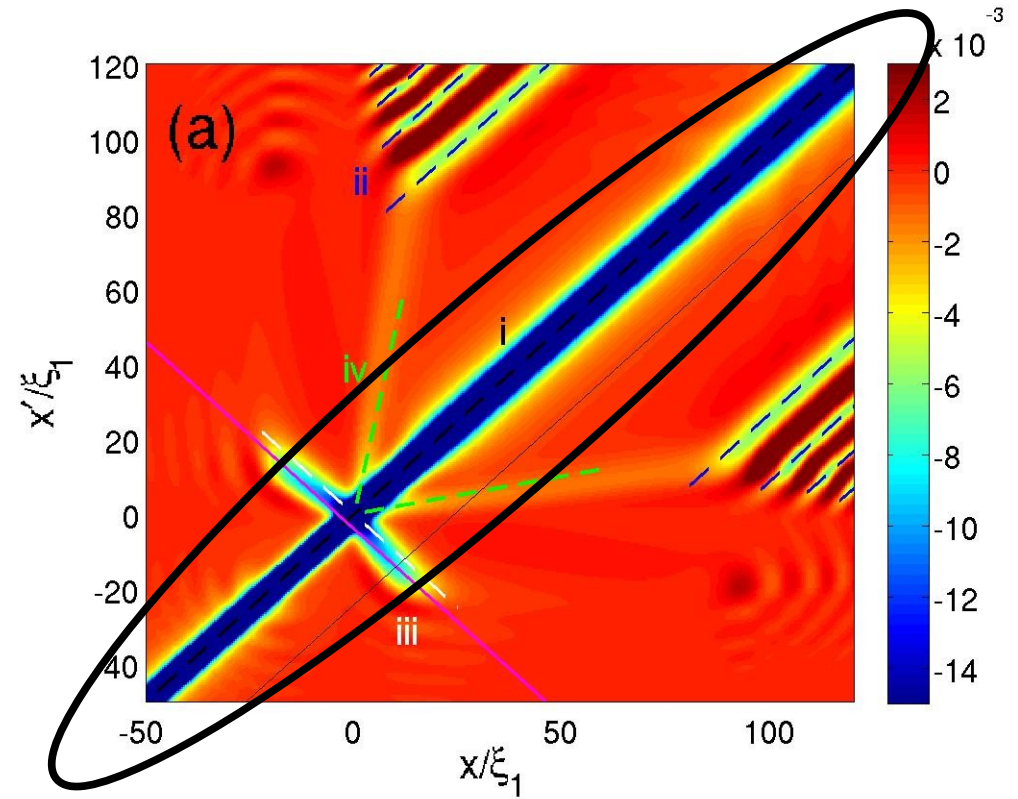


(i)
Many-body
antibunching

(iv)
Hawking
in / in

Feature (i) : Many-body antibunching

- present at all times
- due to **repulsive interactions**
- almost unaffected by flow



See e.g.: M. Naraschewski and R. J. Glauber, PRA 59, 4595 (1999)

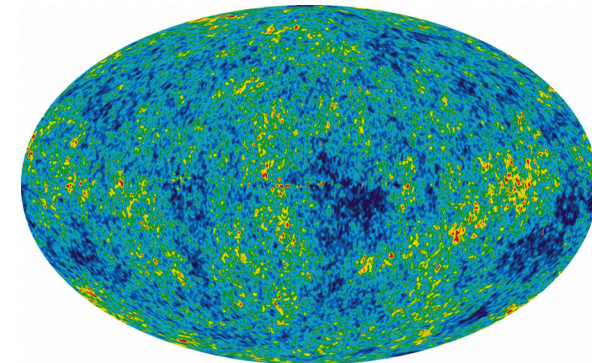
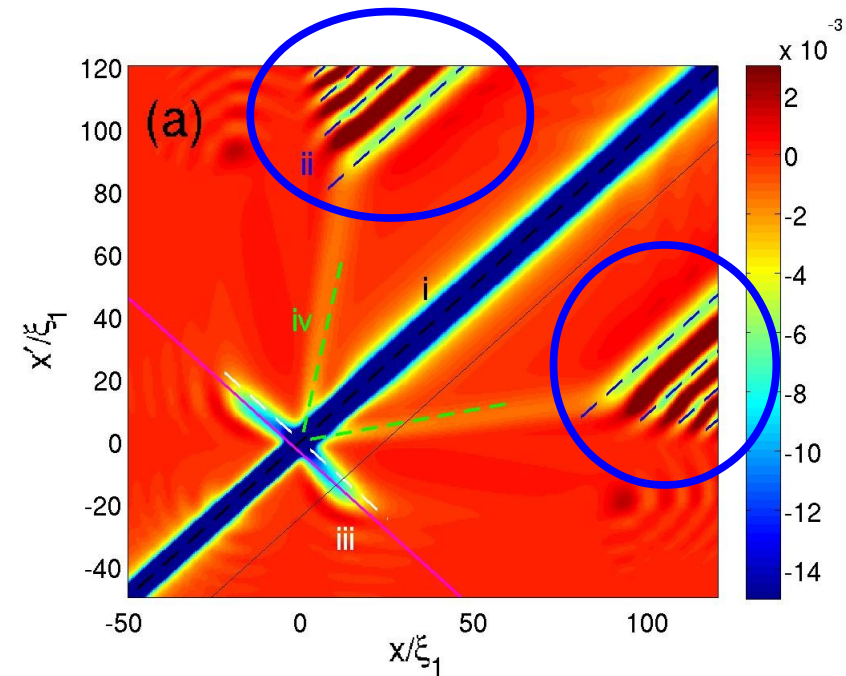
Feature (ii): Dynamical Casimir emission of phonons

Fringes parallel to main diagonal

- intensity depends on speed of switch-on
- only in $x > 0$ region, move away in time
- do not depend on flow pattern, also present in homogeneous system

Physical interpretation:

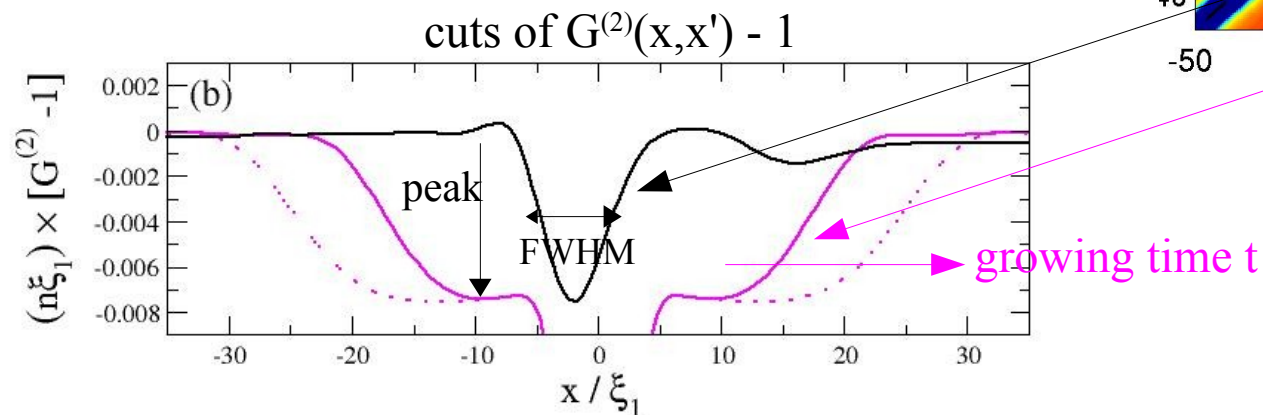
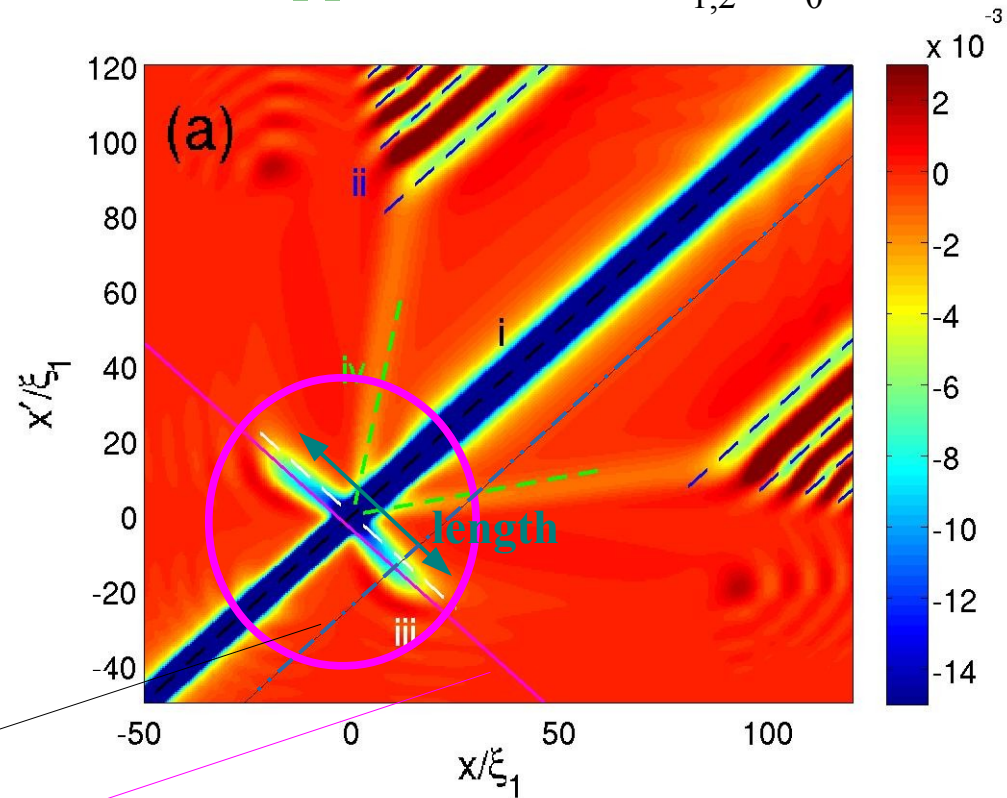
- in $x > 0$ region $g_1 \rightarrow g_2$ within short time σ_t : non-adiabatic modulation of Bogoliubov vacuum
- fringes depend on $|x-x'|$: counter-propagating correlated pairs emitted at $t=0$ at all points $x > 0$
- density correlations propagate away at speed $\geq 2c_s$.
- model of amplification of metric fluctuations during cosmological inflation period



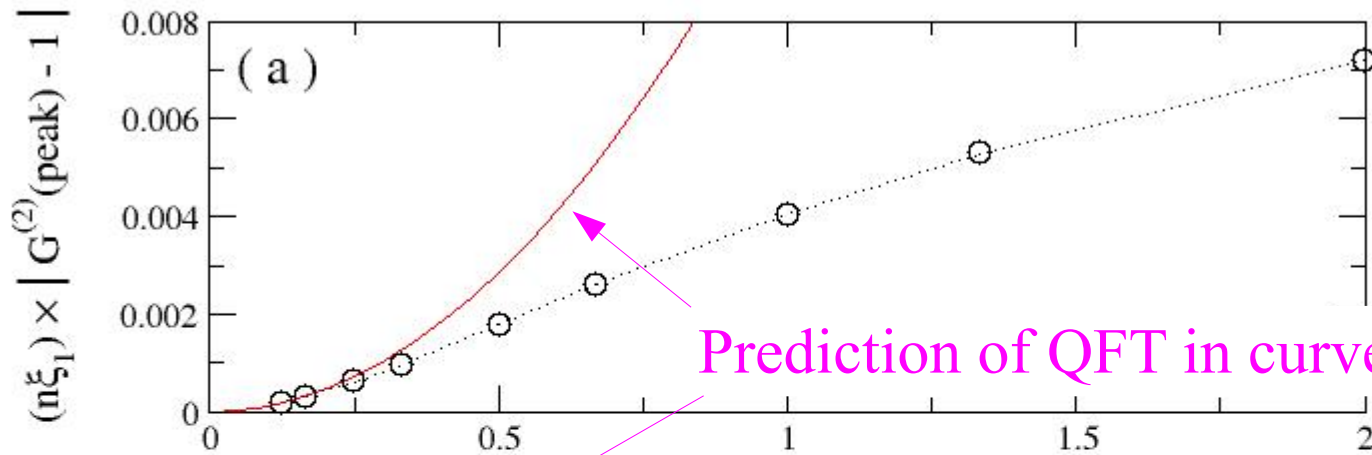
Feature (iii) : The Hawking signal

Negative correlation tongue extending from the horizon $x=x'=0$

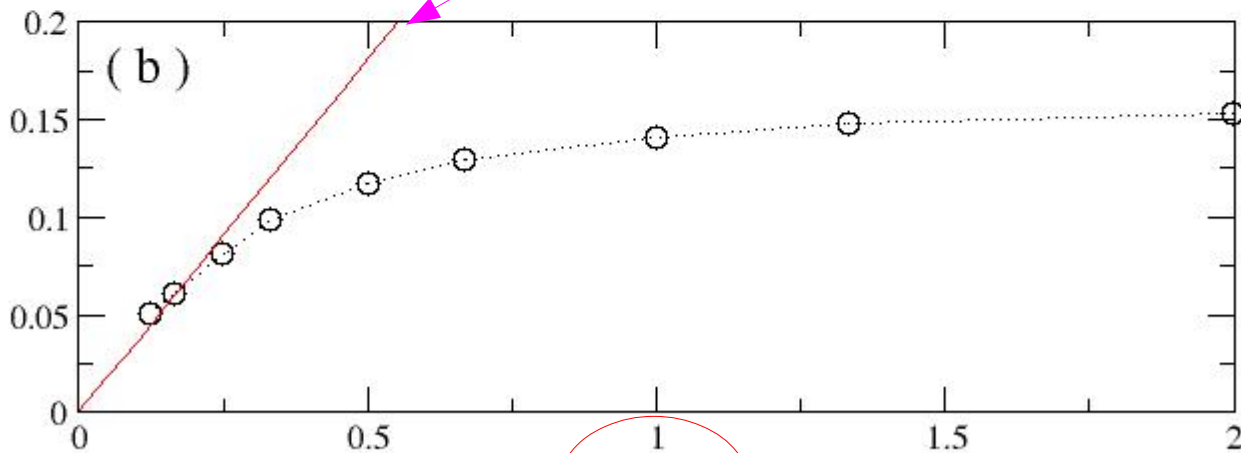
- long-range in/out density correlation which disappears if both $c_{1,2} < v_0$
- length grows linearly in t
- peak height, FWHM constant in t
- slope $\frac{v_0 - c_2}{v_0 - c_1}$ agrees with theory
 - pairs emitted at all t from horizon
 - propagate at sound speed



Quantitative analysis



$\xi_1 / \Delta x_{\text{FWHM}}$

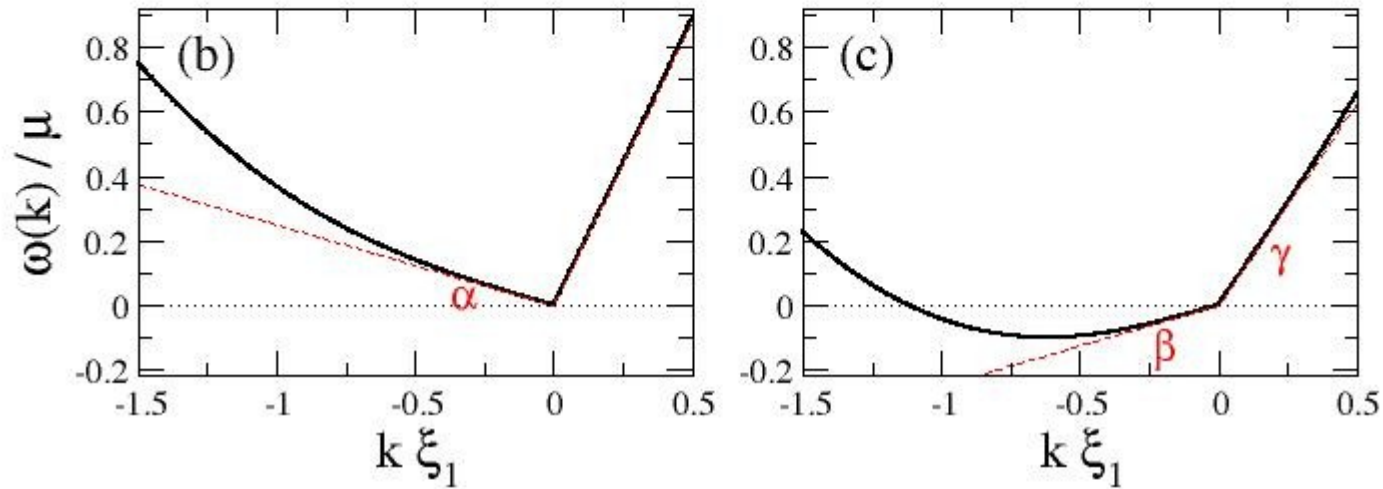


Proportional to emission temperature

Proportional to surface gravity

Analog model prediction quantitatively correct in hydrodynamic limit $\xi_1 / \sigma_x \ll 1$
 Significant discrepancies for strong surface gravity

Effect of UV dispersion



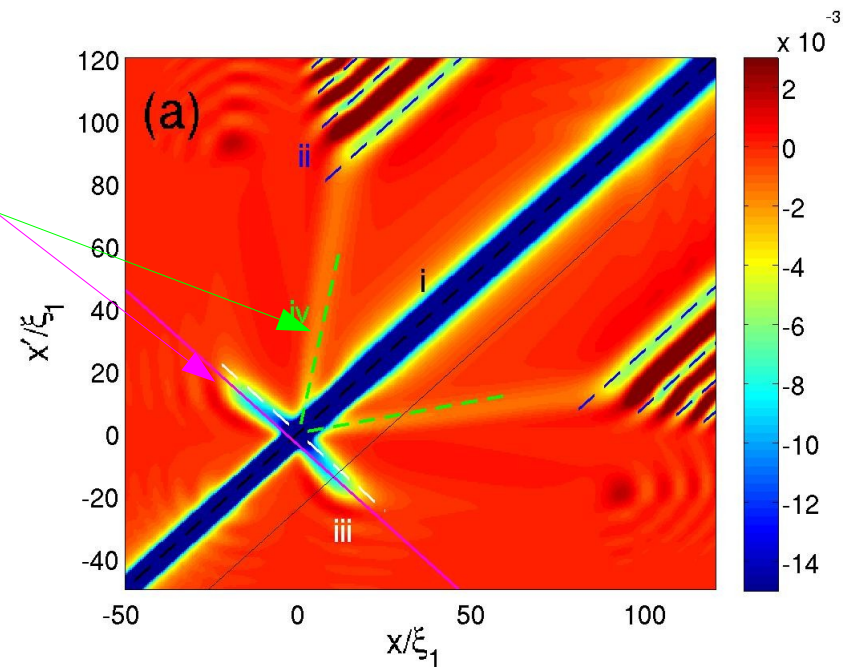
Two parametric “Hawking” processes:

- **in/out**: vacuum $\rightarrow \alpha + \beta$ (feature iii)
- **in/in**: vacuum $\rightarrow \beta + \gamma$ (feature iv)
- third tongue α - γ hardly visible here

Energy conserved only if **sub/super-sonic**

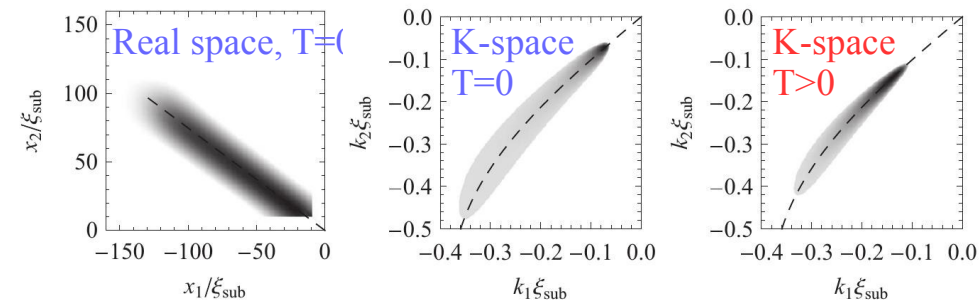
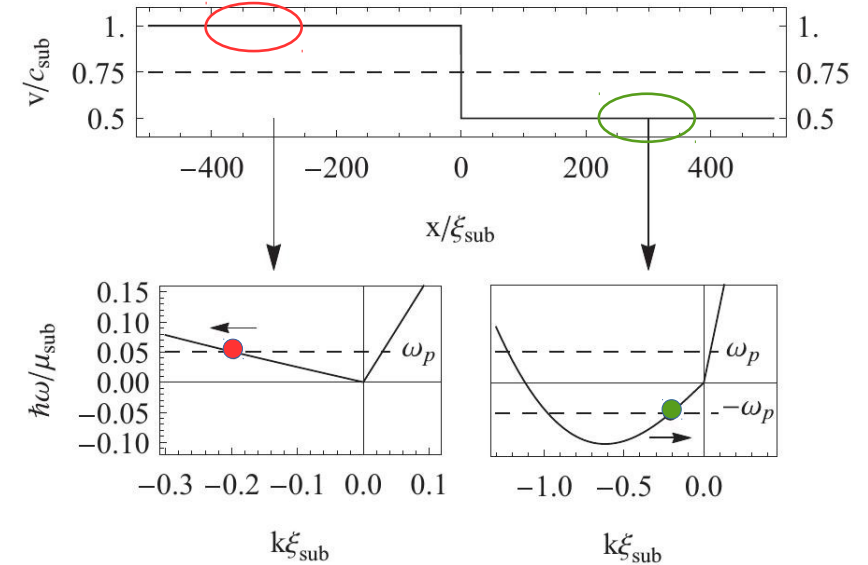
Momentum provided by horizon

Slope of tongues $\frac{v_0 - c_2}{v_0 - c_1} \simeq -1$, $\frac{v_0 - c_2}{v_0 + c_2} \simeq \frac{1}{5}$



How to assess quantum nature of HR in atomic gas ?

- Signal in density/intensity correlations reinforced at finite T by **stimulated Hawking emission**.
→ **Not a signature of quantum origin of emission**
- **Peres-Horodecki criterion** for entanglement in bipartite systems
→ **correlations of quadratures of phonon operators** on either side of the horizon
- Phonon wavepacket operators localized in **real-** and **momentum spaces**
- Need to measure **both phonon quadratures** with **spatial** and **spectral selectivity** (or make strong assumptions on the correlation function)



Part 1-b:

The recent experiments

1- Black hole lasing (2014)

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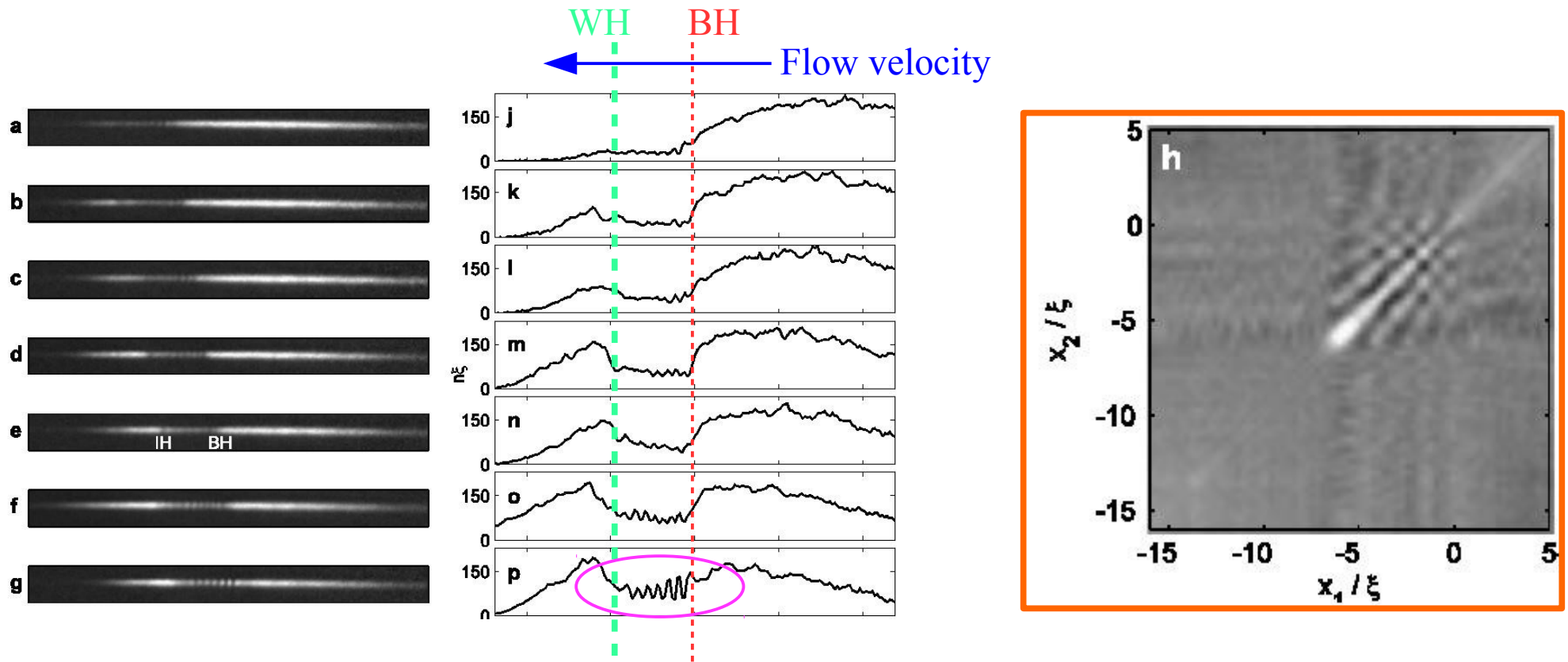
PUBLISHED ONLINE: 12 OCTOBER 2014 | DOI: 10.1038/NPHYS3104

nature
physics

Observation of self-amplifying Hawking radiation in an analogue black-hole laser

Jeff Steinhauer

By a combination of quantum field theory and general relativity, black holes have been predicted to emit Hawking radiation. Observation from an actual black hole is, however, probably extremely difficult, so attention has turned to analogue systems in the search for such radiation. Here, we create a narrow, low density, very low temperature atomic Bose-Einstein condensate, containing an analogue black-hole horizon and an inner horizon, as in a charged black hole. We report the observation of Hawking radiation emitted by this black-hole analogue, which is the output of the black-hole laser formed between the horizons. We also observe the exponential growth of a standing wave between the horizons, which results from interference between the negative-energy partners of the Hawking radiation and the negative-energy particles reflected from the inner horizon. We thus observe self-amplifying Hawking radiation.



Pair of **BH/WH horizons** in **flowing BEC** against a **waterfall+harmonic potential**:

- **BH horizon** at potential edge. **WH horizon** further downstream by growing harmonic potential.
- **Black hole lasing instability** → exponential growth of **density modulation**: “self-amplifying HR”
- Checkerboard/striped pattern in **correlation function of density fluctuations**
- **Open questions**: is initial seed **classical** or **quantum** ? What is precise mechanism for amplification?

Numerical analysis of experiment

GPE simulation with experimental parameters

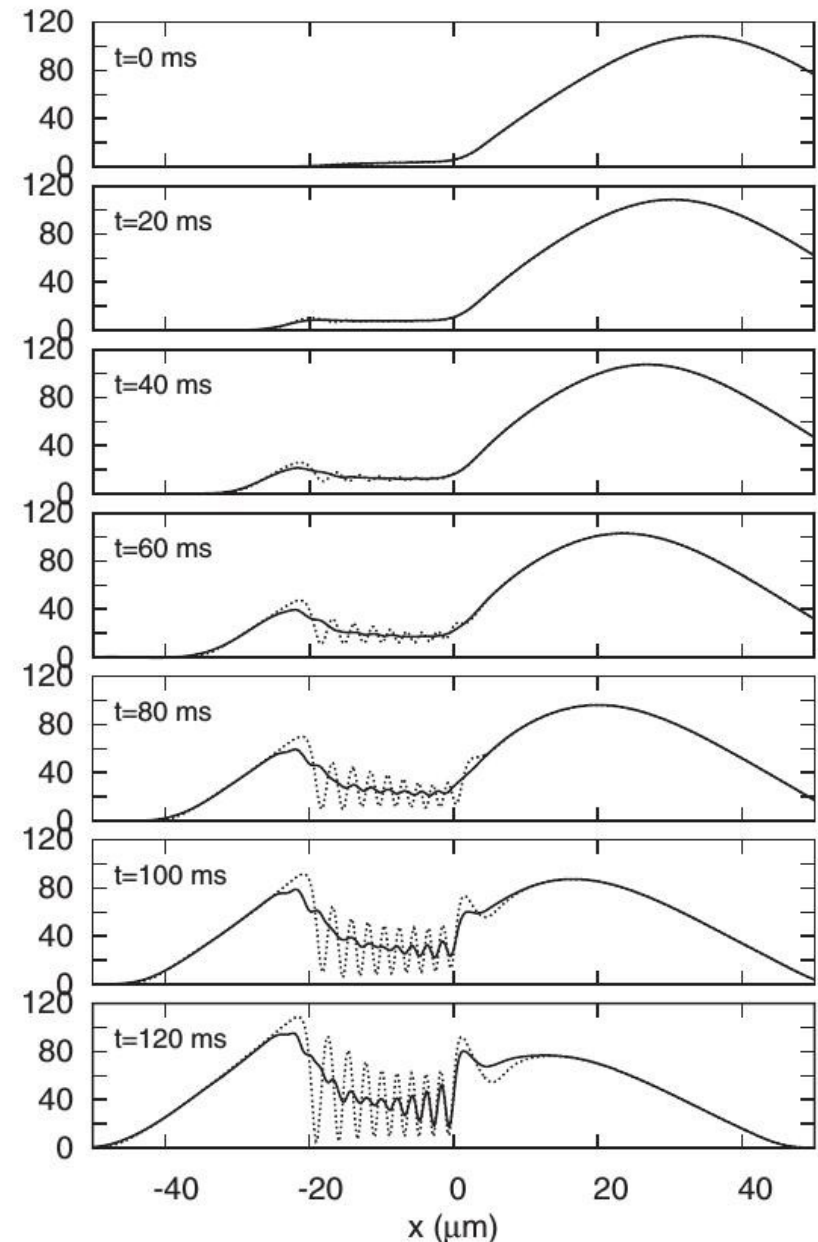
(Tettamanti et al., EPL 2016)

Reproduces quantitatively experimental features

- Observed effect appears to be classical
- Quantum fluctuations don't play crucial role

Further Maryland's concerns about BH lasing mechanism,
Jacobson-Clark's group, PRA 96, 023616 (2017):

- WH region emits Cherenkov-Bogoliubov waves
- Amplitude of CB grows in time following atomic density at WH
- Interpretation validated by our numerics



Tettamanti et al., EPL (2016)

2- Hawking radiation (2016)

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PUBLISHED ONLINE: 15 AUGUST 2016 | DOI: 10.1038/NPHYS3863

Observation of quantum Hawking radiation and its entanglement in an analogue black hole

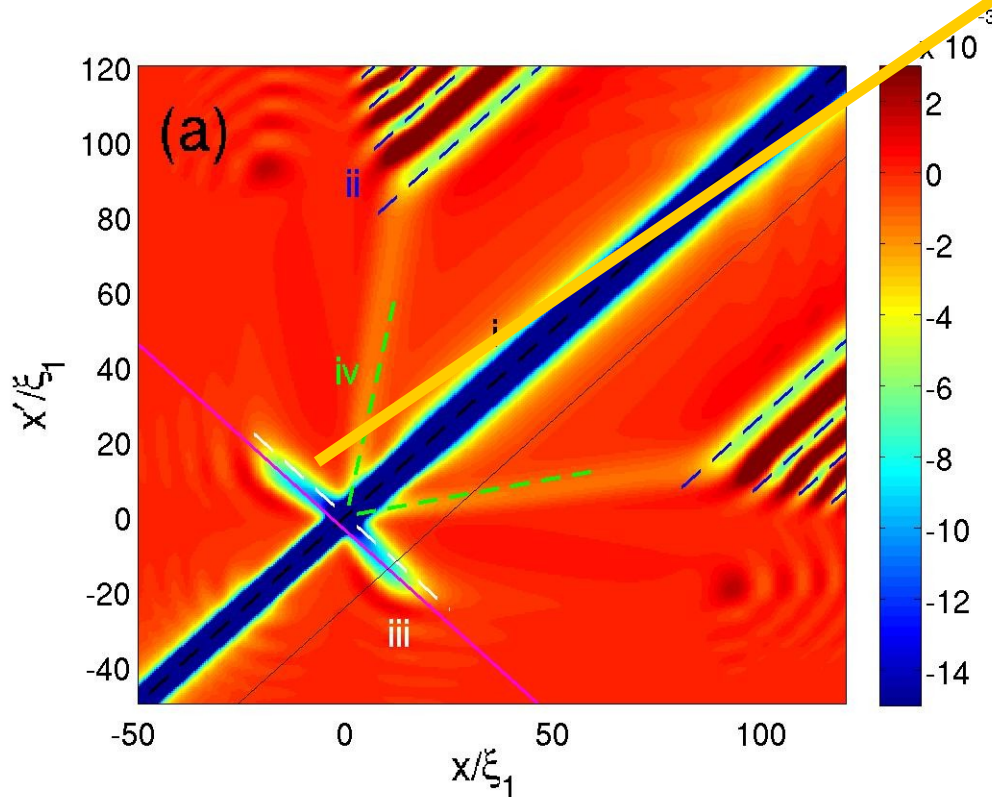
Jeff Steinhauer

We observe spontaneous Hawking radiation, stimulated by quantum vacuum fluctuations, emanating from an analogue black hole in an atomic Bose-Einstein condensate. Correlations are observed between the Hawking particles outside the black hole and the partner particles inside. These correlations indicate an approximately thermal distribution of Hawking radiation. We find that the high-energy pairs are entangled, while the low-energy pairs are not, within the reasonable assumption that excitations with different frequencies are not correlated. The entanglement verifies the quantum nature of the Hawking radiation. The results are consistent with a driven oscillation experiment and a numerical simulation.

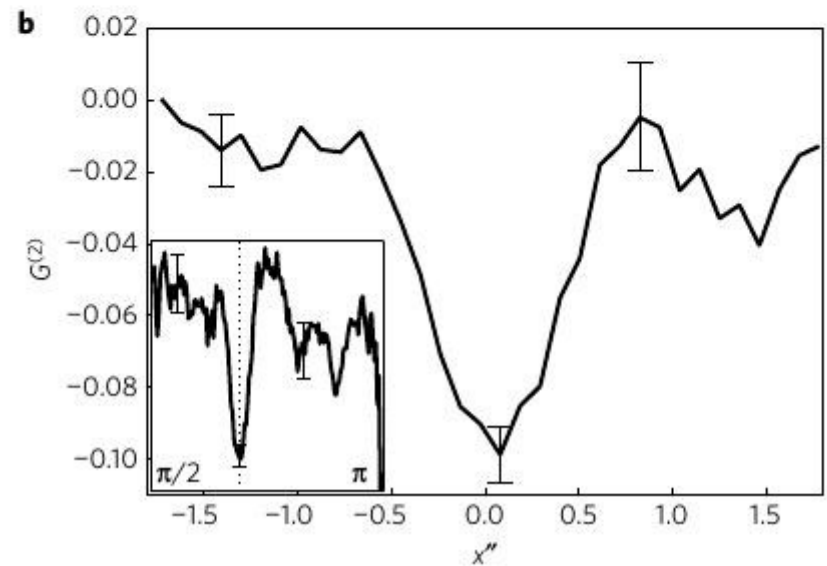
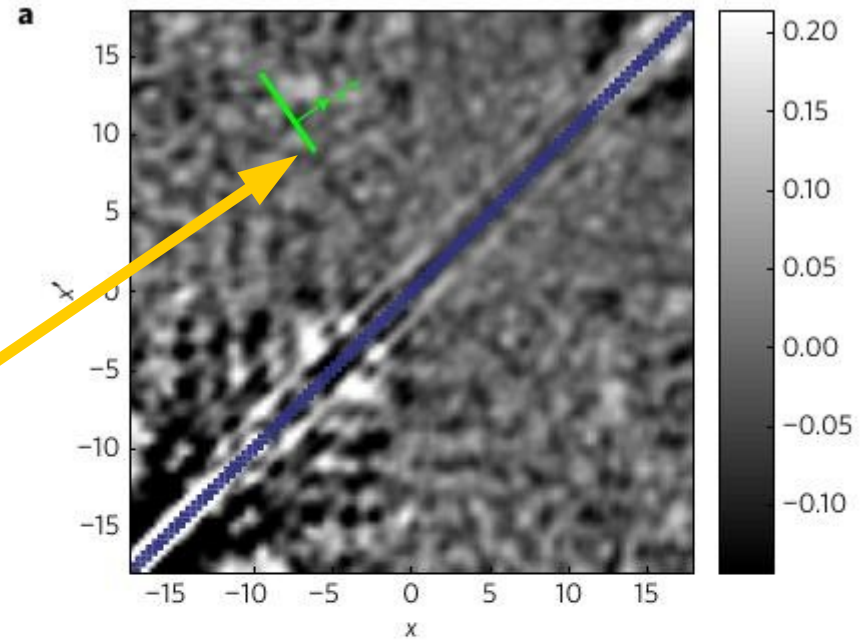
Analog Hawking radiation detected

Analog black hole configuration obtained by sending 1D atomic BEC against optical potential

Experimental evidence of HR based on **Balbinot-Fabrizi moustache** in correlation function of density fluctuations



Theory: IC et al., NJP 2008



Expt: Steinhauer Nat. Phys. '16

Experimental evidence of entanglement

Peres-Horodecki criterion for non-separability

$$\Delta \equiv \underbrace{\langle \hat{b}_{k_{HR}}^\dagger \hat{b}_{k_{HR}} \rangle \langle \hat{b}_{k_P}^\dagger \hat{b}_{k_P} \rangle}_{\text{Population}} - \underbrace{|\langle \hat{b}_{k_{HR}} \hat{b}_{k_P} \rangle|^2}_{\text{Anomalous correlation}} < 0$$

Assumption of **uncorrelated initial fluctuations**

→ simplified Finazzi-IC protocol to
extract anom. corr. from density fluct.

Entanglement visible in **intermediate k-range**

→ HR from zero-point fluctuations
→ produces entangled phonon pairs

Are data statistically significant?

No other possible explanation?

Long-term perspectives

- Quantum Hydrodynamics: Navier-Stokes eqs.
with hats on macroscopic hydrodynamic variables
- Entangled states of a macroscopic fluid

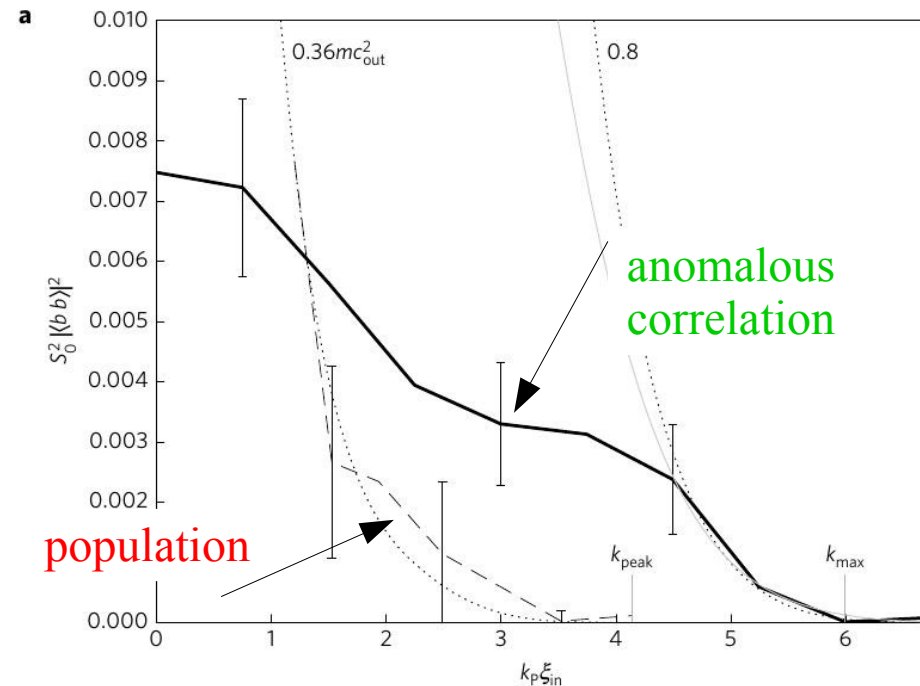


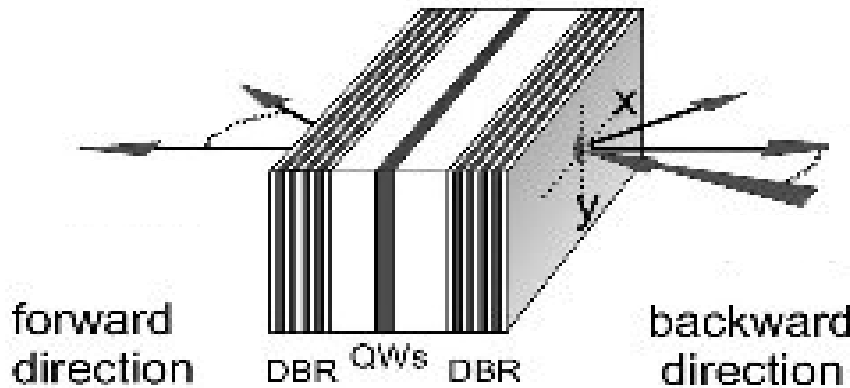
Figure from Steinhauer, Nat. Phys. '16

Much more theoretical work by de Nova,
Sols, Parentani, Bruschi, Fuentes, etc.

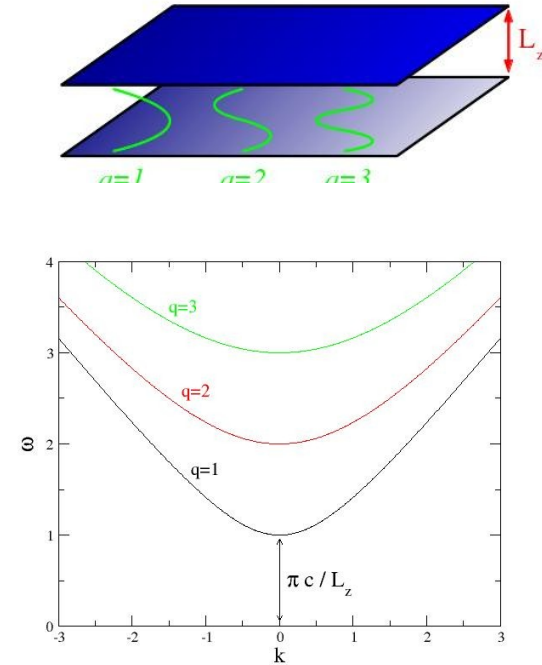
Part 2:

Quantum fluids of light

What are quantum fluids of light?



$$\sigma \sim \alpha^4 \frac{\hbar^2}{m^2 c^2} \left(\frac{\hbar \omega}{mc^2} \right)^6$$



- Photons confined to propagate along microcavity plane
- Spatial confinement along z provides photon mass
- $\chi^{(3)}$ optical nonlinearity (due to excitonic component) provides binary interactions
- Laser pump coherently injects photons:
 - radiative losses determine non equilibrium steady state
 - coherence of polariton fluid guaranteed by coherent pump
- All properties of in-cavity photon fluid transferred to secondary emitted light
- Alternative platform: propagating light in bulk nonlinear media → see Victor's talk later on

Experimental observation of superfluid behaviour

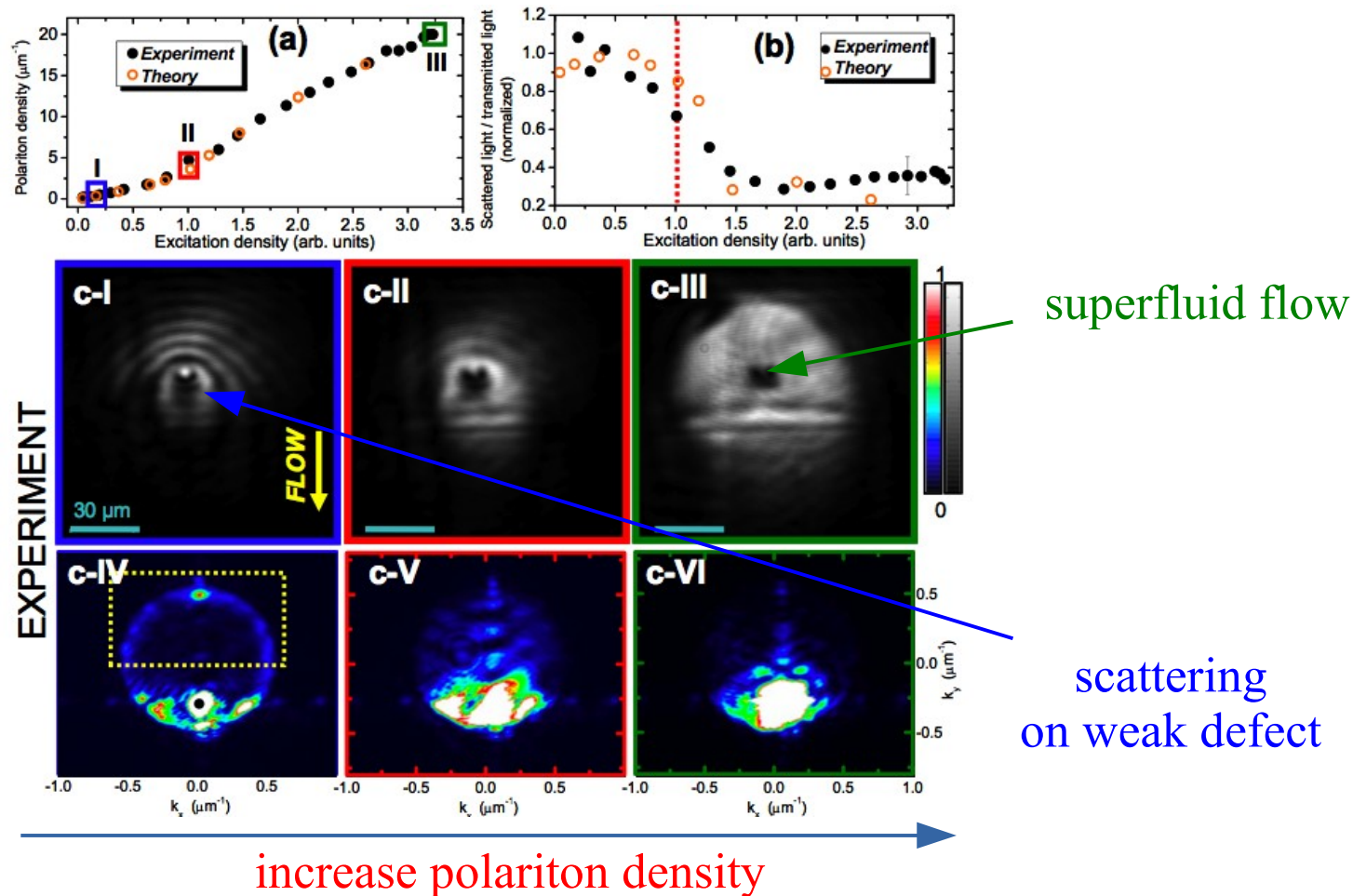


Figure from LKB-P6 group:

A.Amo, J.Lefrère, S.Pigeon, C.Adrados, C.Ciuti, IC, R. Houdré, E.Giacobino, A.Bramati, *Observation of Superfluidity of Polaritons in Semiconductor Microcavities*, Nature Phys. **5**, 805 (2009)

Theory: IC and C. Ciuti, PRL **93**, 166401 (2004)

Acoustic horizons in fluid of light

Polariton-polariton interactions

- Bogoliubov phonon dispersion on top of polariton condensate

Pump at an angle

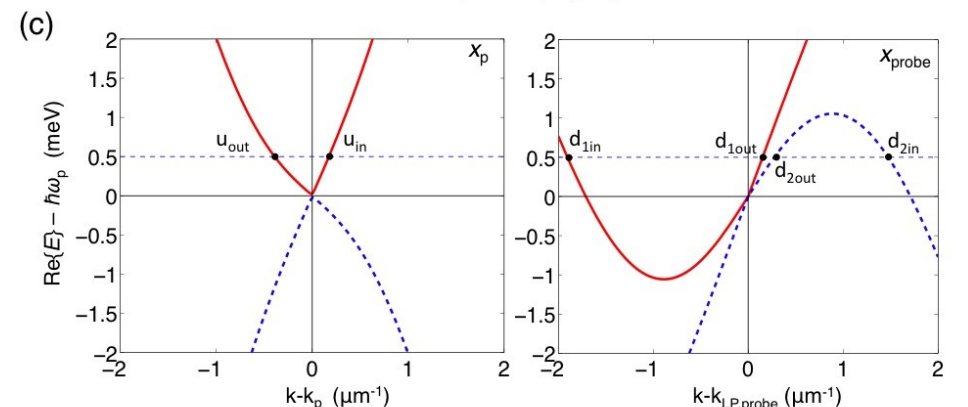
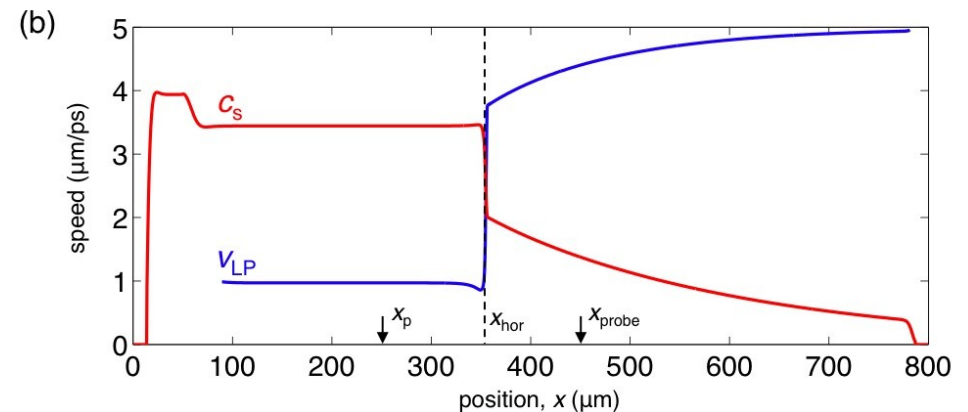
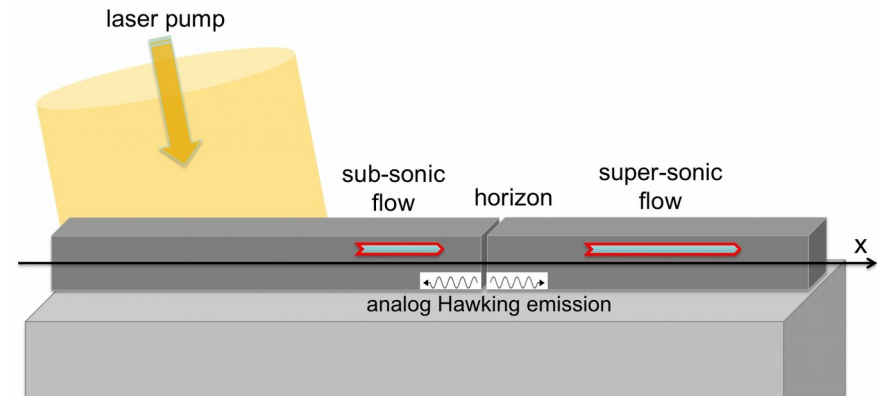
- finite in-plane wavevector, so condensate is flowing

Tailored pump spot + Defect

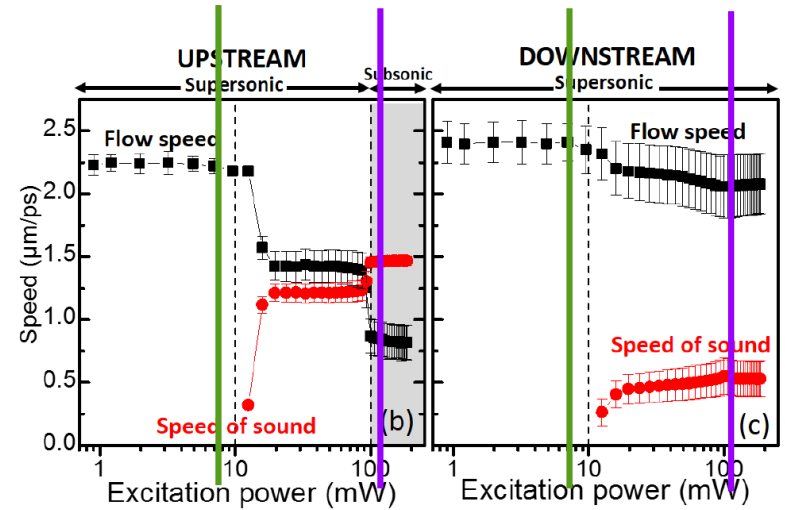
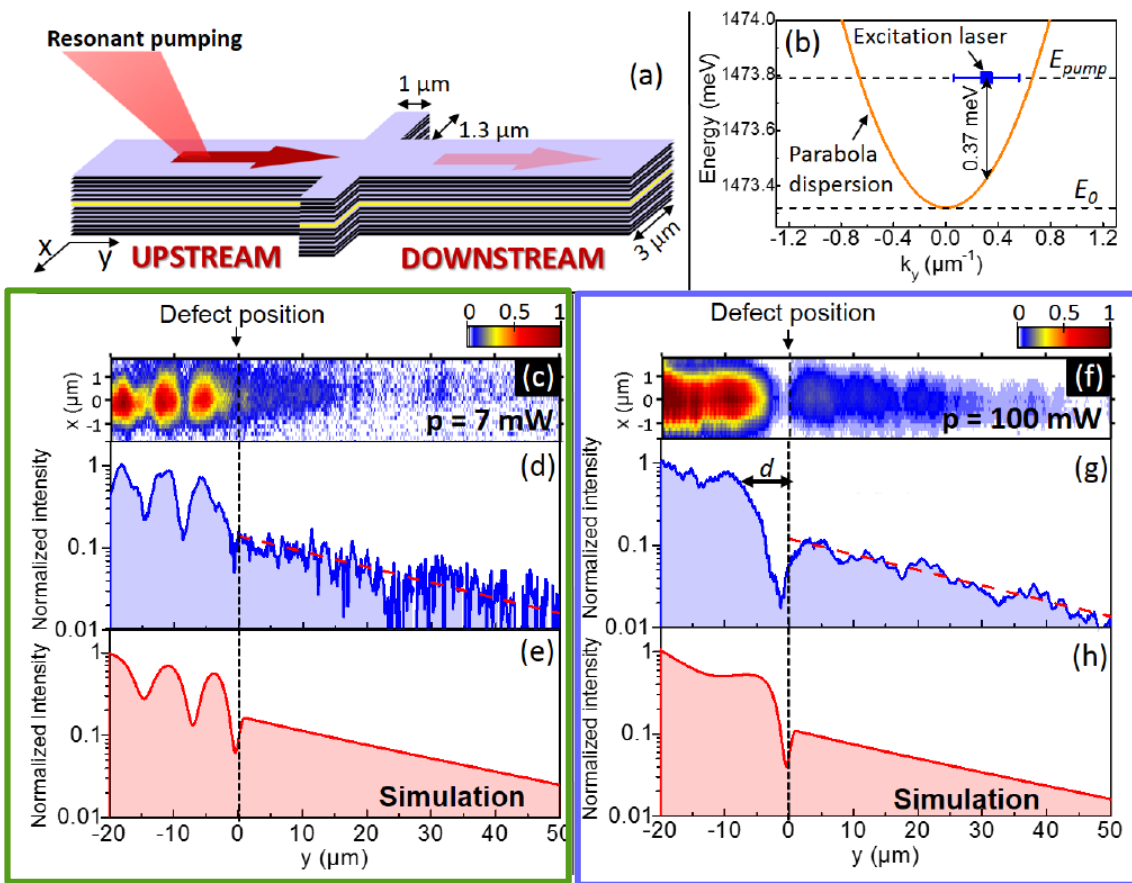
- Horizon with large surface gravity

Hawking emission

- phonons on photon fluid
- correlations of emitted light
- much higher T_H thanks to small photon mass first proposed by F. Marino, PRA **78**, 063804 (2008)



Experimental results @ LPN



BH created!

The hunt for
Hawking radiation
is now open!!

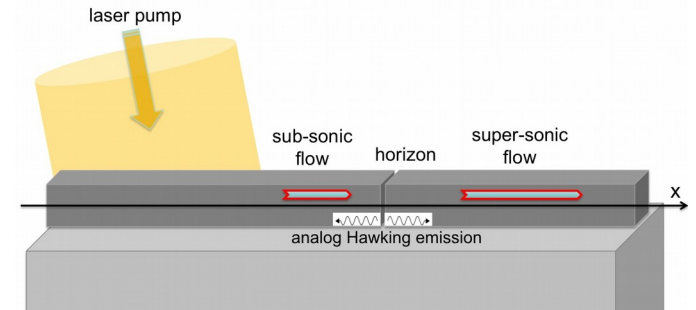
Hawking emission in dissipative photon fluids

- Wigner-MC simulation with driving/losses:

$$i dE = \left\{ \omega_o - \frac{\hbar \nabla^2}{2m} + V_{ext} + g |E|^2 - \frac{i}{2} \gamma \right\} E dt + F_{ext}(x, t) dt + dW$$

- Near-field emission pattern from wire :
Correlation function of intensity noise
at different positions (x, x')

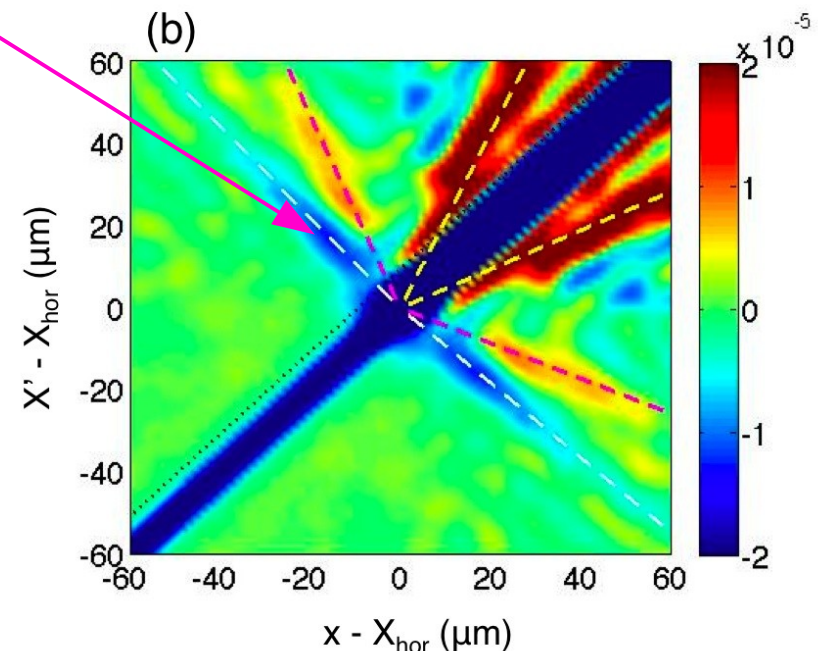
- Signature of Hawking radiation processes:
“Balbinot-Fabbri” correlation tongues
Conversion of zero-point fluctuations
into correlated pairs of Bogoliubov phonons
propagating away from horizon



- In optics language:
parametric emission of entangled photons
flow+horizon play role of pump
photons dressed by fluid into phonons

- Proposed experiment (in progress):

- steady state under cw pumping
- collect near-field emission
- measure intensity noise
- integrate over long time to extract signal out of shot noise



Conclusions and perspectives

Push forward experiments with **artificial black holes** in atoms and polaritons:

- Confirm presence of instability in atomic experiment (Steinhauer, Nat. Phys. 2014)
- Describe quantitatively and understand Hawking emission experiment (Steinhauer, Nat. Phys. 2016)
- HR in spinor condensates → promising to highlight back-reaction (Butera, Ohberg, IC, PRA 2017)
- (In progress) Pump/probe stimulated Hawking radiation in microcavity photons system (Nguyen et al., PRL 15, Grisins et al., PRB 16)
- (Longer run: theory+experiments) Assess quantum origin of spontaneous HR from zero-point fluctuations at the horizon
- Explore conceptually new strategy: **propagating fluids of light**
 - Theory: conservative dynamics under z/t mapping (Fouxon, **Fleurov** et al., EPL 2010; Larré, IC, PRA 2015)
 - Experiments @ Heriot-Watt (Faccio) and @ Tel Aviv (Bar Ad & **Fleurov**)
- Investigate more complex geometries
 - Black hole lasing in BH/WH configurations
 - Vortex configurations: Superradiance, ergoregions around rotating BHs (quantum features inaccessible to surface-wave experiments)
- quantum simulator for curved-space-time QFT, give insight on high-energy questions

The future:

Back-reaction effects

(towards BH evaporation)

The little I understand about back-reaction in astrophysics & quantum gravity

What is the long-term fate of a BH?

HR carries away energy, so horizon
must (very slowly) shrink to
conserve energy/mass

According to some theories, BH horizon
may eventually disappear

- What is left once BH has disappeared ?
- Is there any remnant of
what has fallen into the BH ?

Our approach:

- Analog models simulate well QFT on curved space-time...
- ...but Einstein eqs. (coupling of matter/energy to metric) not implemented

Still, any hint from higher order couplings of quantum fluctuations to macroscopic flow?
What can a quantum optician's point of view teach on this physics?

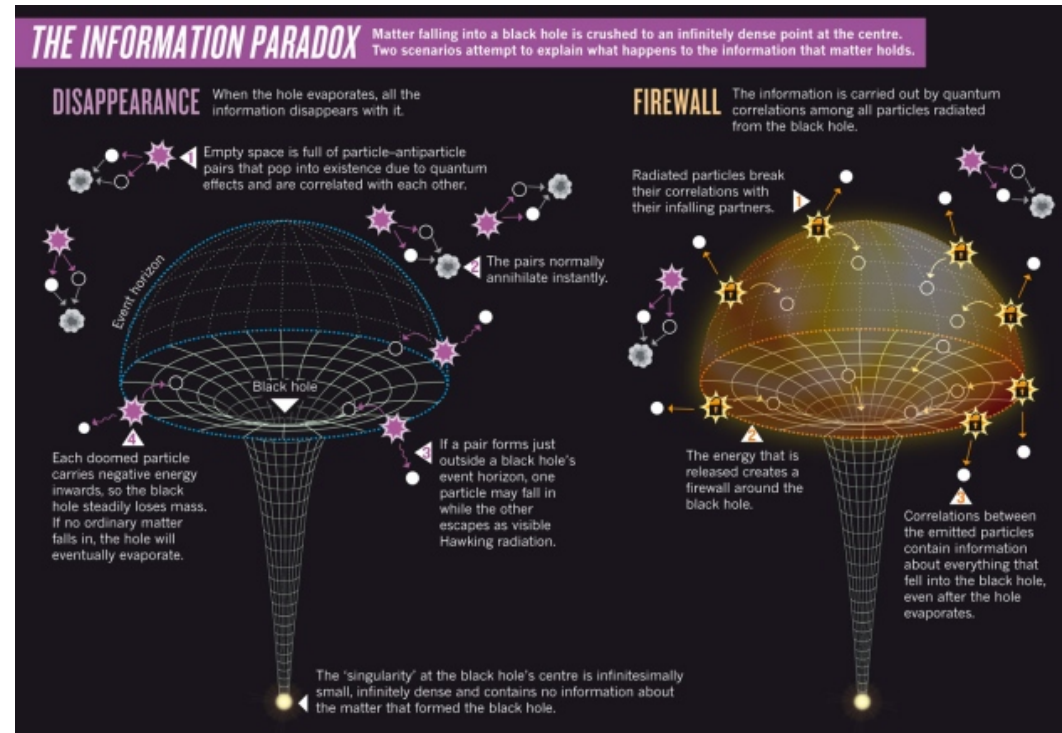
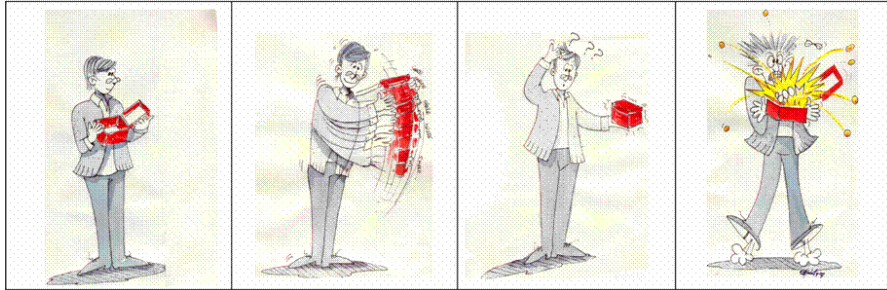


Figure from Nature, 496 (2013)

My favourite toy model

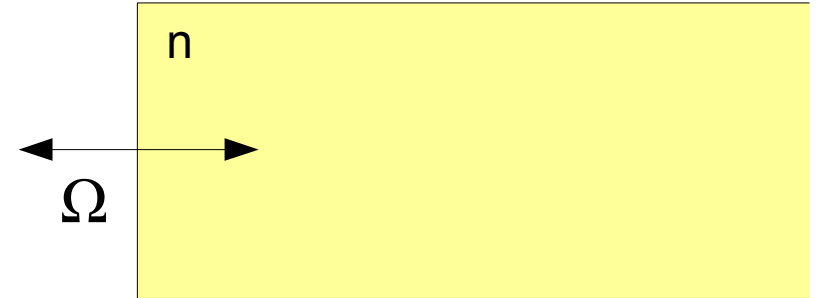
Back-reaction effect of dynamical Casimir emission



Take an optical cavity
in the e.m. vacuum state

Mechanically
shake it very fast

Beware when you open it again:
(a few) photons may burn you !!



Simplest configuration:

- Half-space slab of refractive index n and mass M
- Mechanically oscillating at frequency Ω
- Prediction for the **dissipated energy** within 1D scalar model:

$$Q^{-1} = \frac{\tau}{2\pi E_{osc}} \frac{dE_{diss}}{dt} = \frac{1}{6} \left(\frac{n-1}{n} \right)^2 \frac{\hbar \Omega}{Mc^2}$$

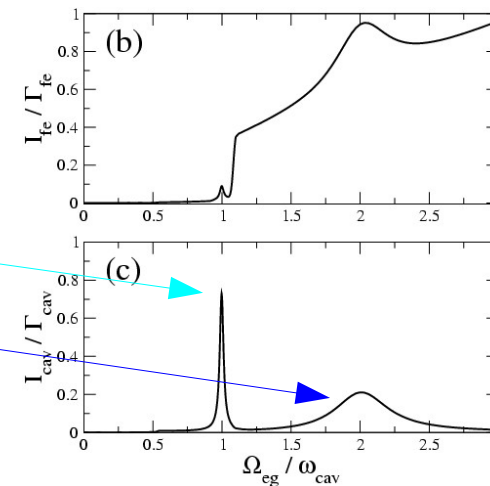
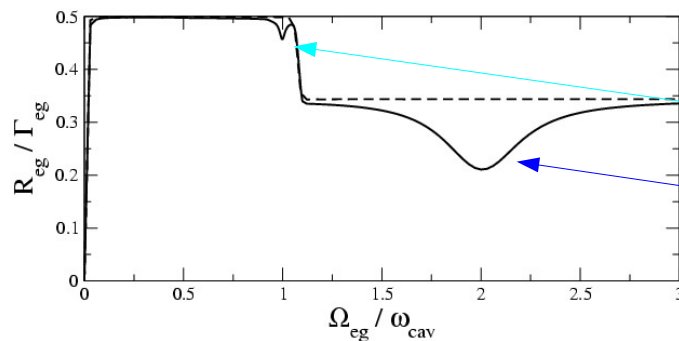
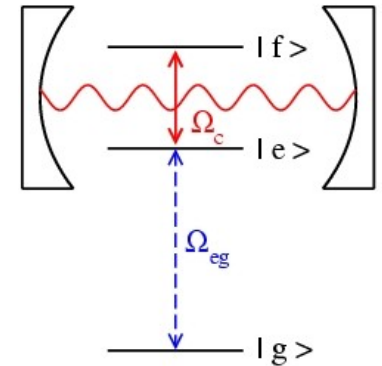
(from Barton and Eberlein, Ann. Phys. 227, 222 (1993))

- value is **ridiculously small**
- hopeless experimental observation by mechanical means, but...

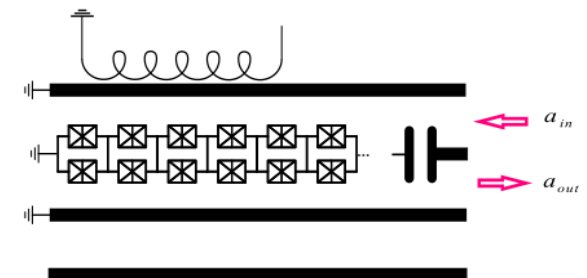
All-optical back-reaction effect

Coherently-driven three-level emitter embedded in optical cavity

- Drive laser on $g \leftrightarrow e$ transition experiences **absorption**
- Absorbed energy $R_{eg} = 2\Omega_{eg} \text{Im}\{\text{Tr}[\hat{c}_{eg}^\dagger \rho_{ss}]\}$.
- Peaks in DCE give **dip in absorption**:
stronger “friction” reduces absorption rate



- “Easily” observed with optical or μ -wave (circuit-QED) techniques
- Theoretical challenge: extend to analog BH’s!



If you wish to know more...

JSF Julian Schwinger
Foundation

BEC
BOSE EINSTEIN CONDENSATION

[Living Reviews in Relativity](#)

December 2011, 14:3 | [Cite as](#)

Analogue Gravity

Authors [Authors and affiliations](#)

Carlos Barceló , Stefano Liberati, Matt Visser

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First Online: 11 May 2011

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Nat. Phys., Aug.15h, 2016

news & views

QUANTUM HYDRODYNAMICS

Acoustic Hawking radiation

A milestone for quantum hydrodynamics may have been reached, with experiments on a black hole-like event horizon for sound waves providing strong evidence for a sonic analogue of Hawking radiation.

Iacopo Carusotto and Roberto Balbinot

REVIEWS OF MODERN PHYSICS, VOLUME 85, JANUARY–MARCH 2013

Quantum fluids of light

Iacopo Carusotto*

INO-CNR BEC Center and Dipartimento di Fisica, Università di Trento, I-38123 Povo, Italy

Cristiano Ciuti†

Laboratoire Matériaux et Phénomènes Quantiques, Université Paris Diderot-Paris 7 et CNRS, Bâtiment Condorcet, 10 rue Alice Domon et Léonie Duquet, 75205 Paris Cedex 13, France

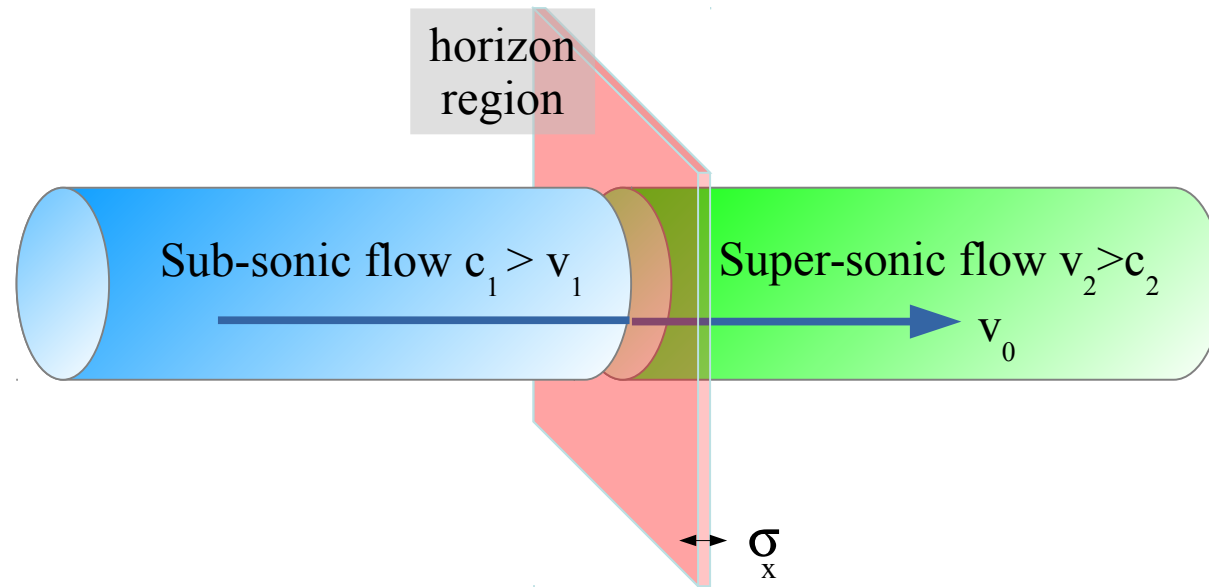
(published 21 February 2013)



Come and visit us in Trento!

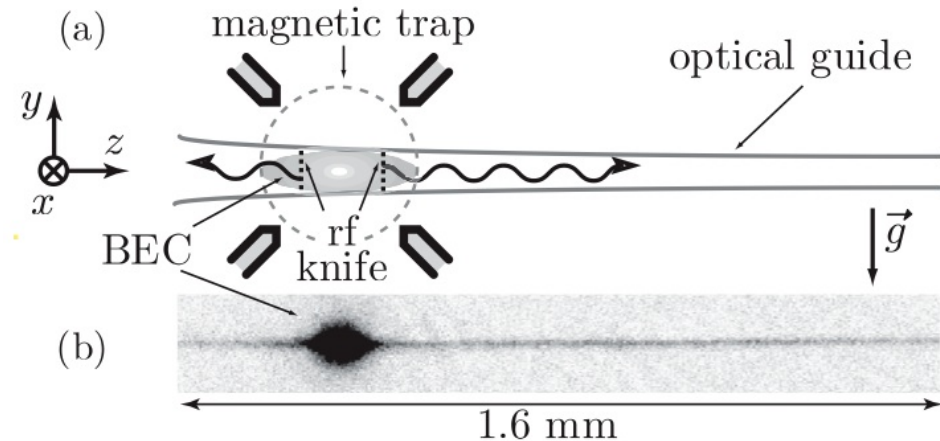
I. Carusotto and C. Ciuti, Reviews of Modern Physics **85**, 299 (2013)

How to generate and study an acoustic black hole ?

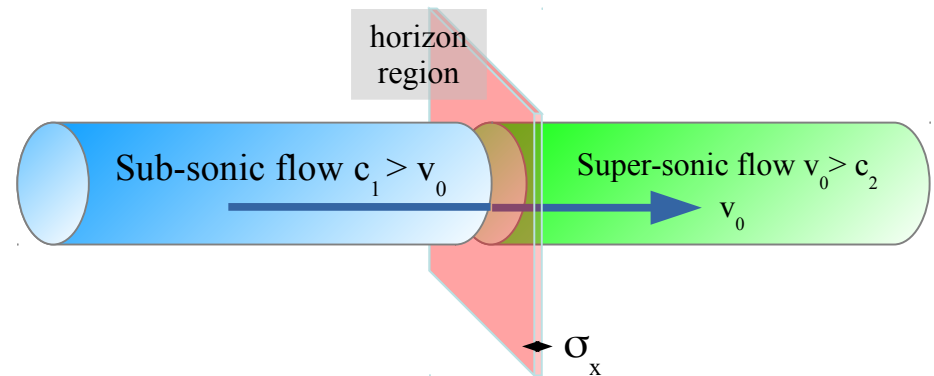


- start from some **uniform flow**
- switch on **horizon** at $t=0$ and go to black-hole regime $c_1 > v_1$, $v_2 > c_2$
 - minimize **deterministic disturbances**, e.g. **Landau processes** (in super-sonic region) and **soliton shedding** during and after switch-on
- concentrate on **quantum fluctuations**
 - **isolate** (thermal) **Hawking emission** from **background phonons** (also thermal)

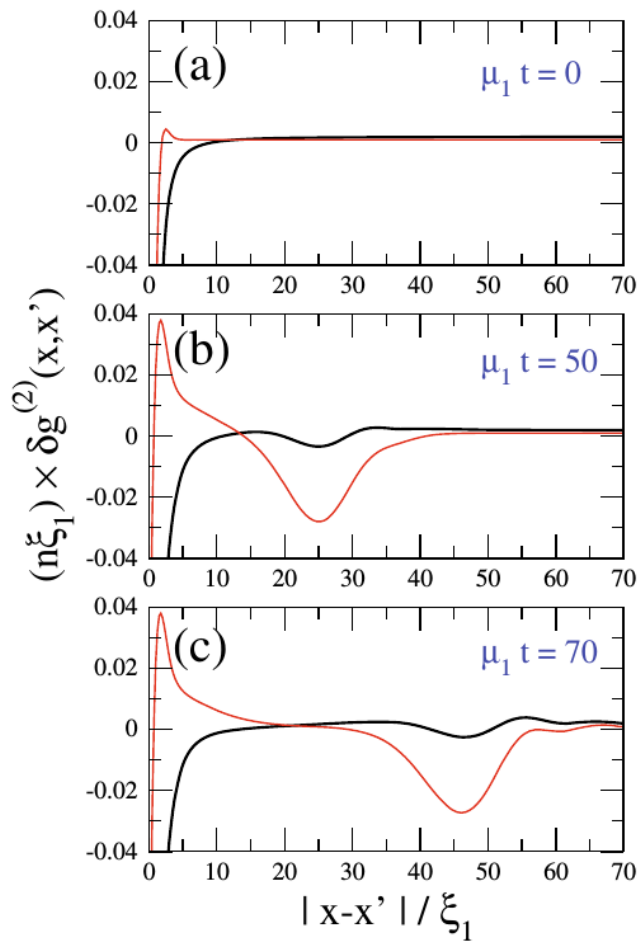
Space-dependent Feshbach resonance



From: W. Guerin *et al.*, PRL **97**, 200402 (2006)

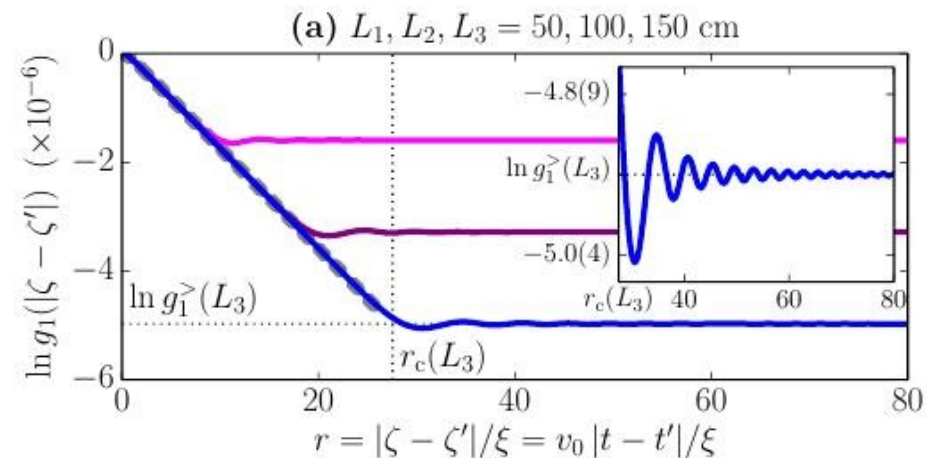
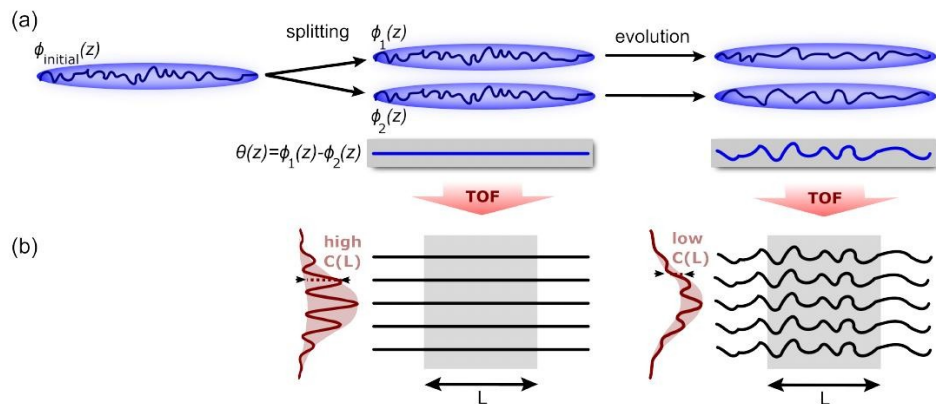


- **Out-coupled atom laser beam**: uniform density and velocity v_0
- **Atom-atom interaction** constant initially uniform and equal to g_1
- **Within σ_x around $t=0$** : modulation $g_1 \rightarrow g_2$ and $V_1 \rightarrow V_2$ in $x > 0$ region only
via: Feshbach resonance (g depends on applied B) or modify transverse confinement
- Step in nonlinear coupling constant $g \rightarrow$ step in sound speed c .
- Black-hole formed if $c_1 > v_0 > c_2$. Arbitrarily large surface gravity via thickness σ_x of crossover region
- **Chemical potential jump** to be compensated by external potential $V_1 + ng_1 = V_2 + ng_2$
allows to avoid Cerenkov-Landau phonon emission, soliton shedding
- Experiments in Steinhauer, Nat. Phys. 10, 864 (2014) \rightarrow slightly different geometry w/o Feshbach.
Upper bound in surface gravity?

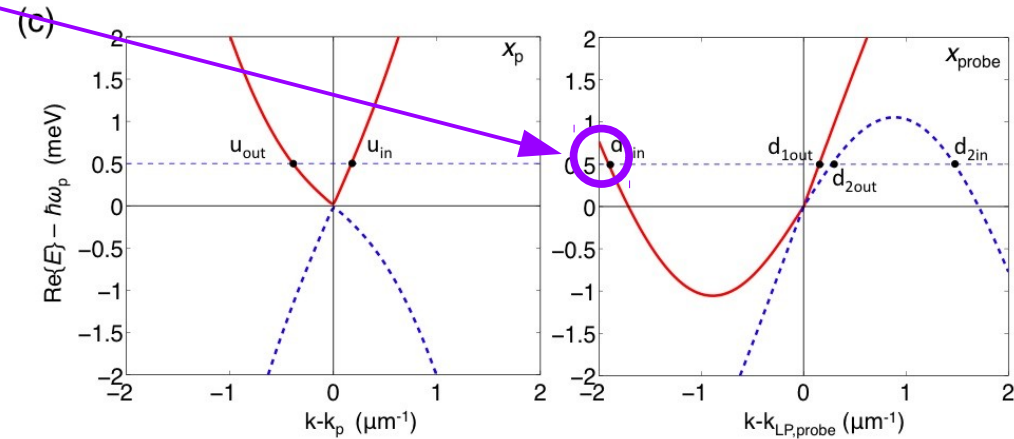
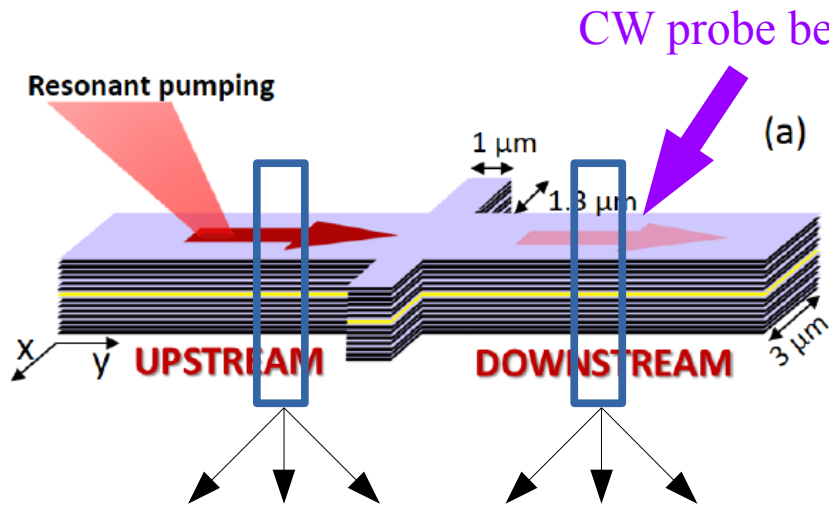


Dynamical Casimir signal on $g^{(2)}(x-x')$:

- density counterpart of **phase decoherence** in split atomic BEC's
- continuity eq. $\partial_x v + \partial_t n = \partial_x (i\partial_x \theta) + \partial_t n = 0$ links dip in $g^{(2)}(x-x')$ to sharp corner in $g^{(1)}(x)$ (within hydrodynamic limit)

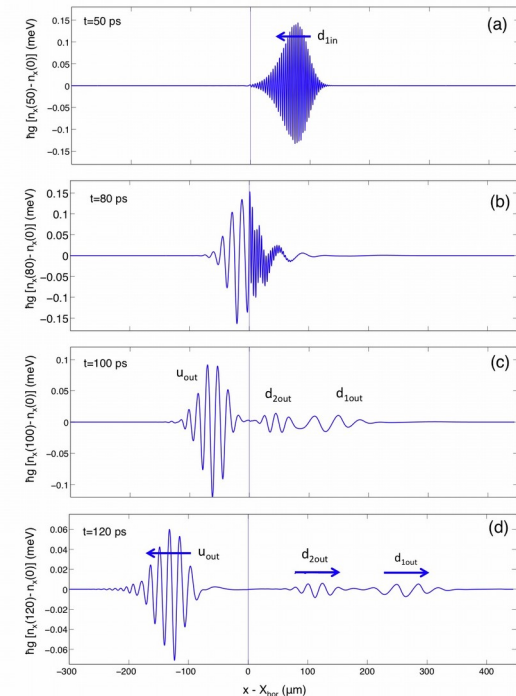


Pump-probe detection of (classical) HR

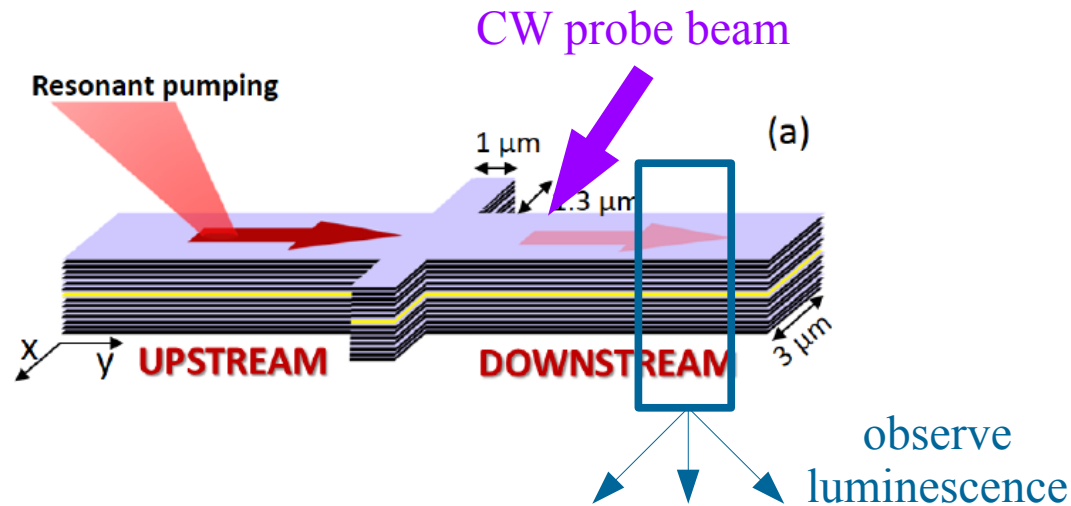


- CW probe at ω_{probe} , frequency resolved detect at ω_{probe} and FWM signal @ $2\omega_{pump} - \omega_{probe}$
- Stimulated Hawking on mode d_{2out} \rightarrow peak in angular distribution
- Scattering matrix $S(\omega)$ \rightarrow signature of thermal Hawking emiss.
- In contrast to pulse expt, no need for temporal resolution

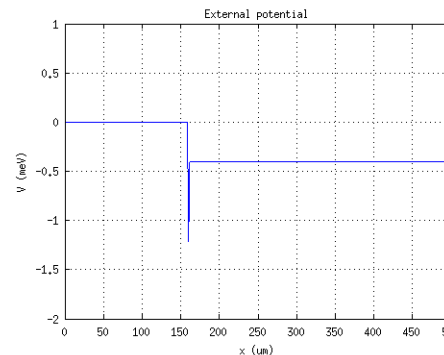
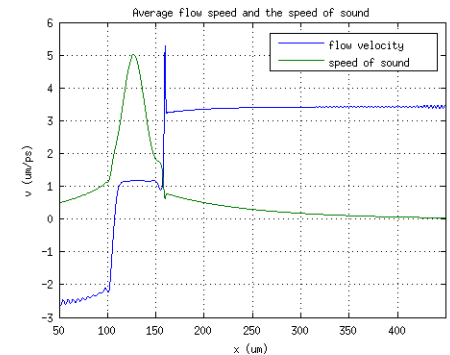
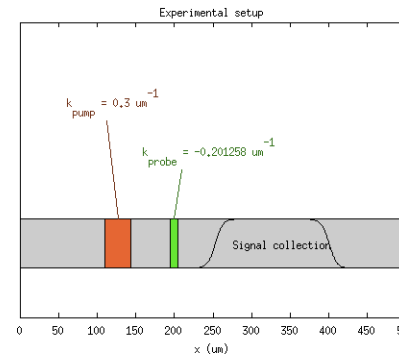
Expt with surface waves on water (Weinfurter, Unruh, PRL 2010) appears not conclusive as no horizon present, new expt in progress (Rousseaux)



Numerical simulations



- Parameters extracted from LPN experiment
- Add “waterfall” potential downstream of defect to facilitate HR observation
- Sample under fabrication @ LPN



Parameters	
LP lifetime $2\pi/\gamma_{LP}$	= 88,232 ps
Pump frequency ω_P	= 2233.26 ps ⁻¹ = 1473.95 meV
Pump intensity FF_P	= 166,667 ps ⁻¹ = 110 meV
Probe frequency ω_{probe}	= 2231.97 ps ⁻¹ = 1473.1 meV
Probe intensity F_{probe}	= 0,00151515 ps ⁻¹ = 0,001 meV
$\omega_{probe} - \omega_P$	= -1,28805 ps ⁻¹ = -0,850115 meV
Pump wavenumber k_P	= 0,3 μm ⁻¹
Probe wavenumber k_{probe}	= -0,201258 μm ⁻¹
$2 \frac{u}{v^2}$	= 70597,7

Classical Hawking signal

Collected frequency-selected momentum distribution:

- Probe frequency ω_{pr} (Blue)
- FWM frequency $\omega_{FWM} = 2\omega_{pump} - \omega_{pr}$ (Green)

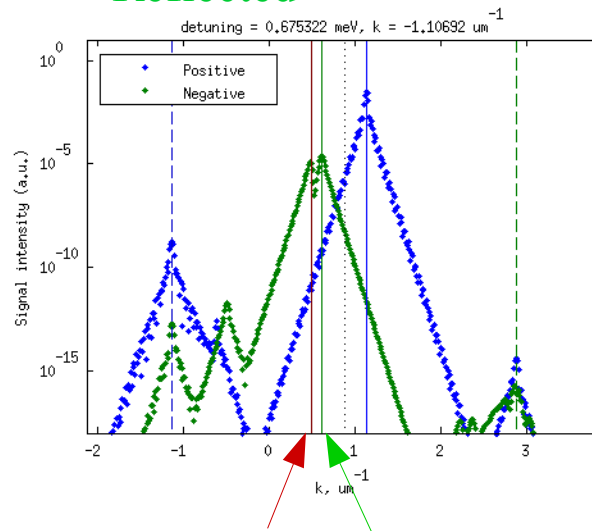
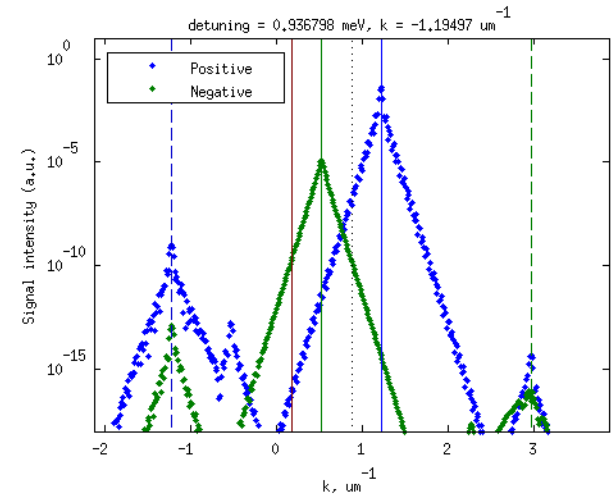
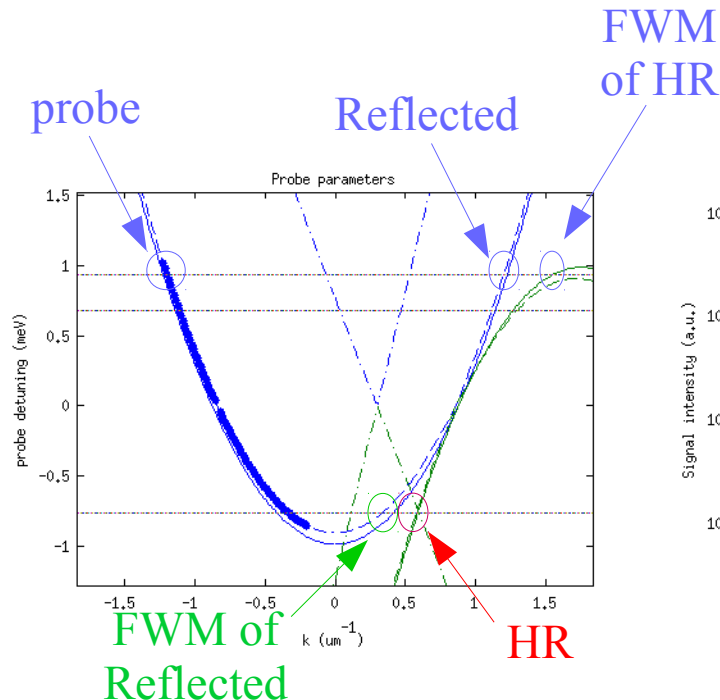
Hawking processes:

→ additional scattering channel

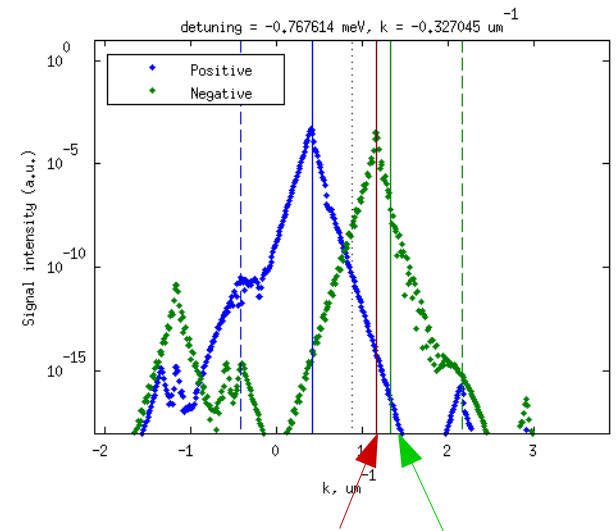
Reflected signal dominates @ ω_{pr}

HR comparable to FWM of reflection @ ω_{FWM}

HR dominates for $\omega_{pr} < \omega_{pump}$



HR Reflection

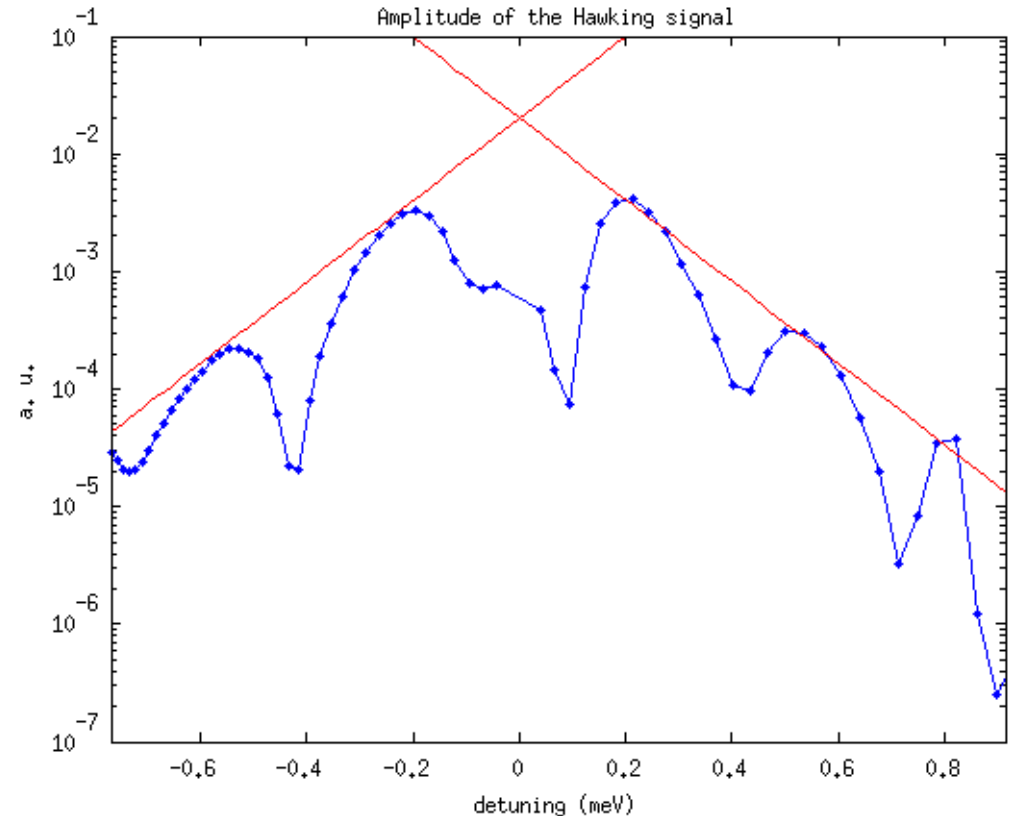


HR Reflection

Estimated Hawking temperature

ω_{pr} dependence of Hawking peak
→ thermal tail @ T_{H}

- Numerical → $T_{\text{H}} \sim 1.4 \text{ K}$
(after correcting for propagating losses)
- Not far from theoretical prediction $T_{\text{H}} \sim 1 \text{ K}$
from gravitational analogy

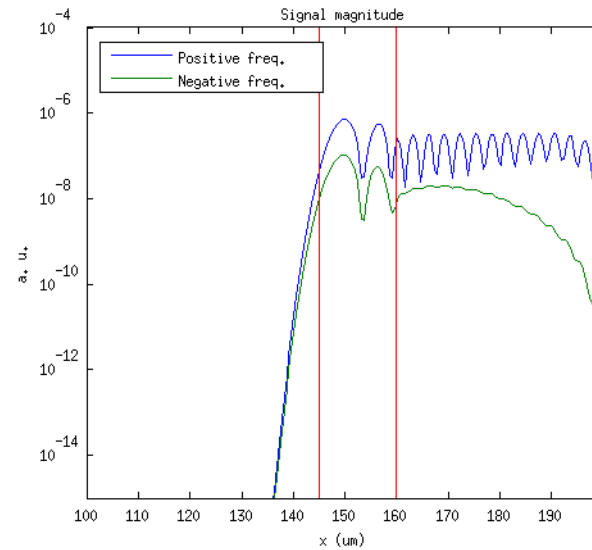
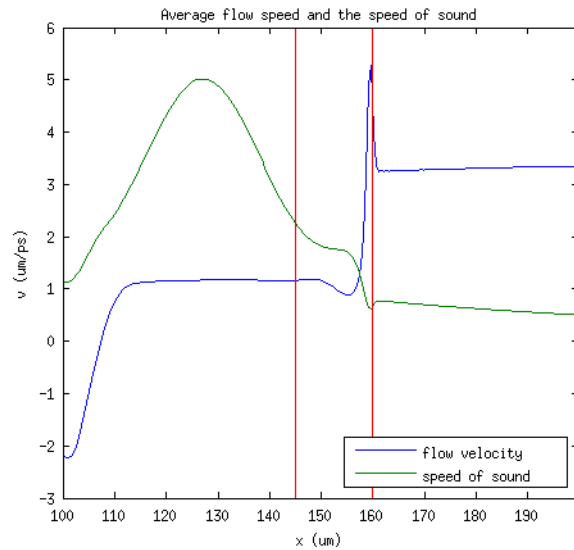
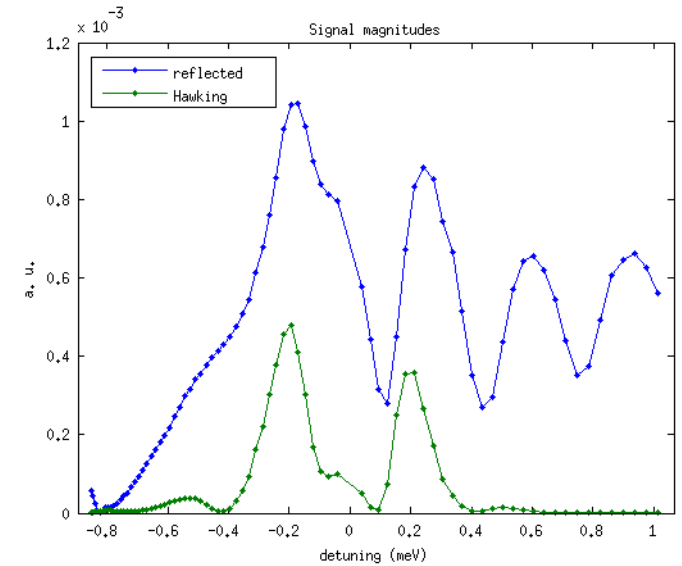


Why oscillations?

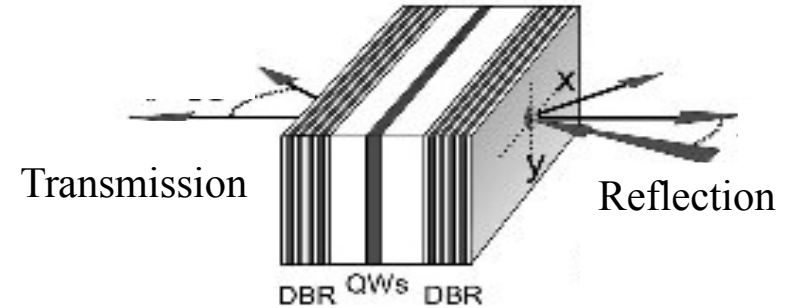
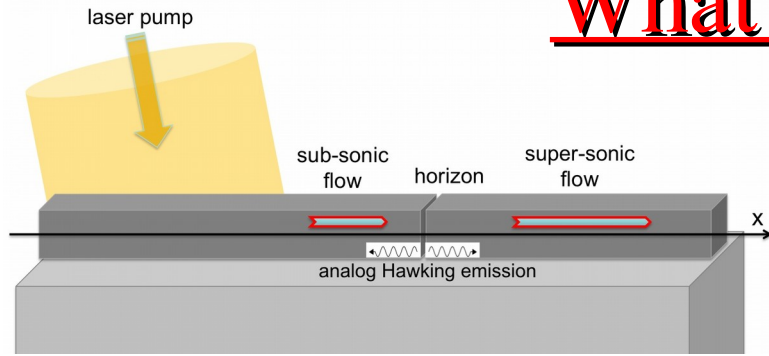
BH horizon and strong pump form cavity for phonons

Reflected and Hawking signals oscillate with ω_{pr}

Optical analogs of “eternal BHs” in thermal equilibrium with environment, e.g. with mirror right above horizon



What about dissipation?



Light fluid in cavity coupled to dissipation baths:

- radiative emission \rightarrow photon decay at rate γ ; used to observe field, compensated by pump
- description in terms of master equation for density matrix ρ

$$\frac{d\rho}{dt} = -\frac{i}{\hbar} [H, \rho] + \frac{\gamma}{2} \int dr \left[2\psi(r) \rho \psi^\dagger(r) - \psi^\dagger(r) \psi(r) \rho - \rho \psi^\dagger(r) \psi(r) \right]$$

$$H = \int dr \left[\psi^\dagger(r) \left(\omega_o - \frac{\hbar \nabla^2}{2m} \right) \psi(r) + \frac{g}{2} \psi^\dagger(r) \psi^\dagger(r) \psi(r) \psi(r) + F(r, t) \psi^\dagger(r) + F^*(r, t) \psi(r) \right]$$

Fluctuation-dissipation theorem:

- quantum noise acts back onto quantum field
- accurate description in terms of classical field eqs. + stochastic noise

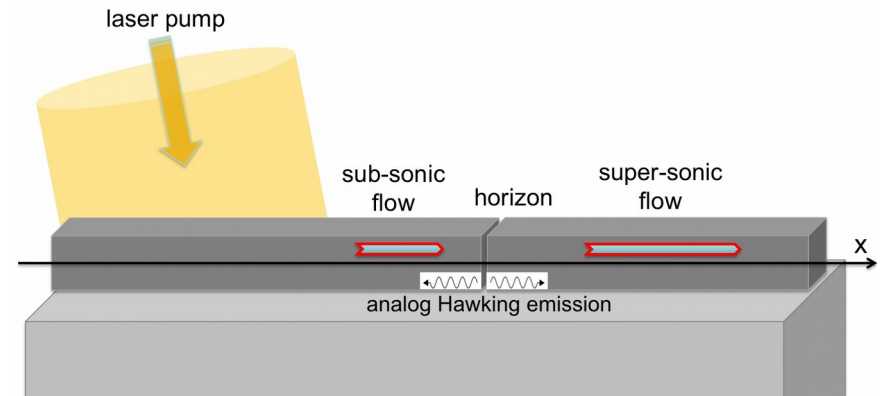
Properties of photon fluid affected by non-equilibrium nature:

- richer dispersion of Bogoliubov excitations; condensation at $k \neq 0$; ...
- steady-state population of Bogoliubov modes

What about acoustic horizons in fluids of light?

Polariton-polariton interactions

- Bogoliubov phonon dispersion on top of polariton condensate

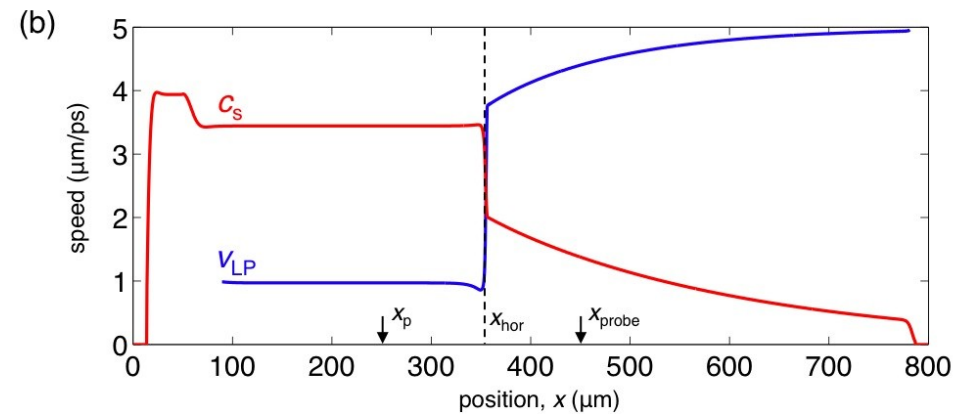


Pump at an angle

- finite in-plane wavevector, so condensate is flowing

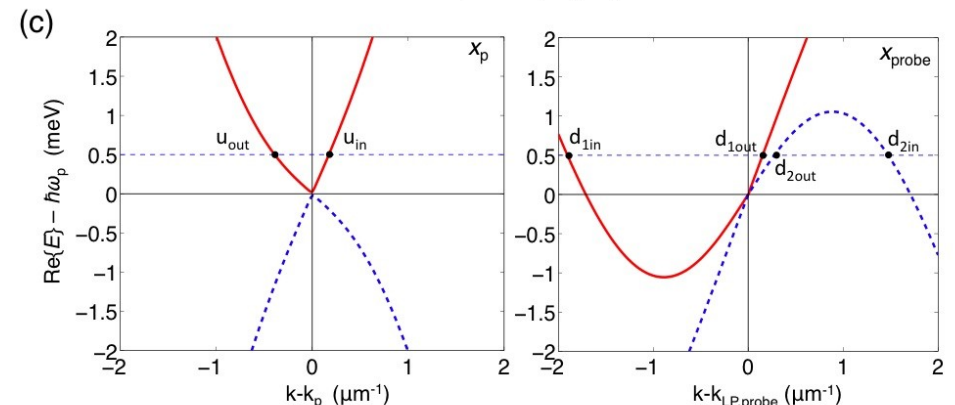
Tailored pump spot + Defect

- Horizon with large surface gravity



Hawking emission

- phonons on photon fluid
- correlations of emitted light



Is HR in microcavities a quantum process?

Unavoidable losses of microcavity device generate **excitations in the fluid**:
lost photon \rightarrow creates hole \rightarrow Bogoliubov excitation

Spurious excitations up to $\omega \sim gn$, comparable to T_H

X.Busch, R.Parentani, IC, *Spectrum and entanglement of phonons in quantum fluids of light*, PRA 2014

Might be tamed by **taylorizing ω -dependent reflectivity of cavity mirror**
(as it is likely occurring in atom outcoupling from trap)

Stimulate Hawking processes giving rise to **“thermal” Hawking signal**:

- **Density correlation signal** reinforced
- **Quantum entanglement** still present ?
 - How to detect it ?
 - Hong-Ou-Mandel on emitted light spatially + wavevector-selected ?

