



SSA in $ep^\uparrow \rightarrow hX$ and contribution of quasi-real photon

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based on: U. D'Alesio, CF, F. Murgia, PRD95 (2017)

Outline

1. Introduction

2. Quasi-real photon exchange

Weizsäcker-Williams approximation

Weizsäcker-Williams contribution to SSA

3. Phenomenological analysis

Comparison with HERMES data

Predictions

JLab

COMPASS

EIC

4. Conclusions

Introduction

- Inclusive $p^\uparrow p$ scattering: $A_N = \frac{d\sigma^\uparrow - d\sigma^\downarrow}{d\sigma^\uparrow + d\sigma^\downarrow} = \frac{d\Delta\sigma}{2d\sigma_{\text{unp}}}$
[see e.g. Aschenauer's, M. Boer's talks]

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- **two-scale** vs **single-scale** processes [see e.g. Prokudin's talk]

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test for TMD factorisation

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test for TMD factorisation

- two approaches:
 - **Generalize Parton Model (GPM)** [Anselmino, Boglione, D'Alesio, Melis, Murgia, Prokudin, PRD81(2010) & PRD89 (2014)]
 - **collinear twist-3** [Gamberg, Kang, Metz, Pitonyak, Prokudin, PRD90(2014)]

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⇒ extend previous GPM phenomenological studies

Previous studies

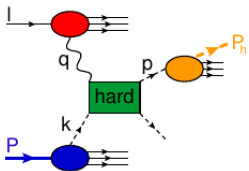
- Detailed kinematics and first predictions [*Anselmino et al. (2010)*]
- first comparison with a selection ($Q^2 \gtrsim 1 \text{ GeV}^2$) of HERMES data [*PLB 728 (2014)*]: TMD-LO study, $\ell q \rightarrow \ell q$ only partonic channel [*Anselmino et al. (2014)*];

$$\begin{aligned}
 d\sigma_{LO}^{\uparrow} - d\sigma_{LO}^{\downarrow} = & \sum_q \left\{ \Delta^N f_{q/p\uparrow} \otimes d\hat{\sigma} \otimes D_{h/q} \right. && \text{quark Sivers} \\
 & + h_1^{q/p} \otimes d\Delta\hat{\sigma} \otimes \Delta^N D_{h/q\uparrow} \cos\phi_C && \text{Collins I} \\
 & \left. + h_{1T}^{\perp q/p} \otimes d\Delta\hat{\sigma} \otimes \Delta^N D_{h/q\uparrow} \cos(\phi_C - 2\phi_q) \right\} && \text{Collins II}
 \end{aligned}$$

$$d\hat{\sigma} \simeq e_q^2 \frac{\hat{s}^2 + \hat{u}^2}{\hat{t}^2}, \quad d\Delta\hat{\sigma} \simeq -e_q^2 \frac{\hat{s}\hat{u}}{\hat{t}^2}$$

- anti-tagged ($Q^2 \approx 0$) events not included

Weizsäcker-Williams approximation

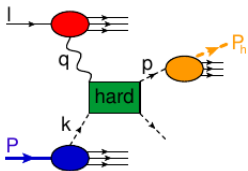


$l \rightarrow l\gamma$ with final lepton almost collinear $\Rightarrow Q^2 \sim 0$ (quasi-real γ)

lepton as a source of real photons

$$d\sigma(pl \rightarrow hX) = \int_0^1 dy f_{\gamma/e}(y) d\sigma(p\gamma \rightarrow hX)$$

Weizsäcker-Williams approximation



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$$d\sigma(pl \rightarrow hX) = \int_0^1 dy \, f_{\gamma/e}(y) \, d\sigma(p\gamma \rightarrow hX)$$

WW distribution [*Hinderer, Schlegel, Vogelsang, PRD92 (2015); 93 (2016)*]:

$$f_{\gamma/e}(y, \mu) = \frac{\alpha}{2\pi} \frac{1 + (1-y)^2}{y} \left[\ln\left(\frac{\mu^2}{y^2 m_\ell^2}\right) - 1 \right] + \mathcal{O}(\alpha^2)$$

New partonic channels: $q\gamma \rightarrow qg$, $g\gamma \rightarrow q\bar{q}$

WW contribution to A_N for $p^\uparrow \ell \rightarrow hX$

$$A_N = \frac{d\Delta\sigma^{\text{LO}} + d\Delta\sigma^{\text{WW}}}{2[d\sigma^{\text{LO}} + d\sigma^{\text{WW}}]}$$

$$\begin{aligned}
 d\Delta\sigma^{\text{WW}} &= d\sigma_{p\gamma}^\uparrow - d\sigma_{p\gamma}^\downarrow = \\
 &= \sum_q \left\{ \Delta^N f_{q/p}^\uparrow \cos\phi_q \otimes [d\hat{\sigma}^{q\gamma \rightarrow q} \otimes D_{h/q} + d\hat{\sigma}^{q\gamma \rightarrow g} \otimes D_{h/g}] \quad \text{quark Sivers} \right. \\
 &+ \Delta^N f_{g/p}^\uparrow \cos\phi_g \otimes [d\hat{\sigma}^{g\gamma \rightarrow q} \otimes D_{h/q} + d\hat{\sigma}^{g\gamma \rightarrow \bar{q}} \otimes D_{h/\bar{q}}] \quad \text{gluon Sivers} \\
 &+ h_1^{q/p} \otimes d\Delta\hat{\sigma}^{q\gamma \rightarrow q} \otimes \Delta^N D_{h/q}^\uparrow \cos\phi_C \quad \text{Collins I} \\
 &\left. + h_{1T}^{\perp q/p} \otimes d\Delta\hat{\sigma}^{q\gamma \rightarrow q} \otimes \Delta^N D_{h/q}^\uparrow \cos(\phi_C - 2\phi_q) \right\} \quad \text{Collins II}
 \end{aligned} \tag{1}$$

$$\begin{aligned}
 d\hat{\sigma}^{q\gamma \rightarrow q} &= -\frac{4}{3}e_q^2 \frac{\hat{s}^2 + \hat{u}^2}{\hat{s}\hat{u}}, \quad d\hat{\sigma}^{q\gamma \rightarrow g} = d\hat{\sigma}^{q\gamma \rightarrow q}(\hat{u} \leftrightarrow \hat{t}), \quad d\hat{\sigma}^{g\gamma \rightarrow q} = d\hat{\sigma}^{g\gamma \rightarrow \bar{q}} = e_q^2 \frac{\hat{u}^2 + \hat{t}^2}{\hat{t}\hat{u}} \\
 d\Delta\hat{\sigma}^{q\gamma \rightarrow q} &= \frac{8}{3}e_q^2
 \end{aligned}$$

WW contribution to A_N for $p^\uparrow \ell \rightarrow hX$

$$A_N = \frac{d\Delta\sigma^{\text{LO}} + d\Delta\sigma^{\text{WW}}}{2[d\sigma^{\text{LO}} + d\sigma^{\text{WW}}]}$$

$$d\Delta\sigma^{\text{WW}} = d\sigma_{p\gamma}^\uparrow - d\sigma_{p\gamma}^\downarrow =$$

$$= \sum_q \left\{ \Delta^N f_{q/p\uparrow}^N \cos\phi_q \otimes [d\hat{\sigma}^{q\gamma\rightarrow q} \otimes D_{h/q} + d\hat{\sigma}^{q\gamma\rightarrow g} \otimes D_{h/g}] \right. \quad \text{quark Sivers}$$

$$+ \Delta^N f_{g/p\uparrow}^N \cos\phi_g \otimes [d\hat{\sigma}^{g\gamma\rightarrow q} \otimes D_{h/q} + d\hat{\sigma}^{g\gamma\rightarrow \bar{q}} \otimes D_{h/\bar{q}}] \quad \text{gluon Sivers}$$

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$$d\Delta\hat{\sigma}^{q\gamma\rightarrow q} = \frac{8}{3}e_q^2$$

\hat{u} dependence absent at LO!

Phenomenological analysis

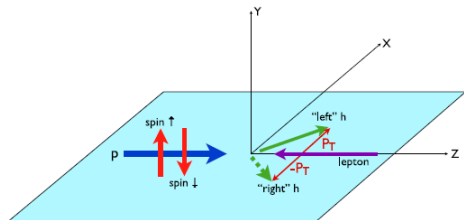
Parametrisation and data selection

- TMD parametrisation: usual factorised form

$$N_q(x) \times [\text{Coll. PDF(FF)}] \times [\text{Gaussian } \mathbf{k}_\perp(\mathbf{p}_\perp) \text{ factor}]$$

- tested two WW distribution and two different μ scale (P_T and $\frac{\sqrt{s}}{2}$): almost no differences $\Rightarrow \mu = P_T$
- **SIDIS1** and **SIDIS2** extractions for Sivers and Collins functions [*see Pisano's talk*]
- gluon Sivers function from fit to SSA for $p^\uparrow p \rightarrow \pi^0 X$ data at midrapidity [*D'Alesio, Murgia, Pisano, JHEP 09 (2015)*]
- HERMES A_N data [*PLB 728 (2014)*]:
 - one bin with $\langle P_T \rangle \simeq 1$ GeV vs. x_F
 - one bin at fixed x_F vs. $P_T \geq 1$ GeV (anti-tagged data)

Kinematics for HERMES



HERMES: lp^\uparrow c.m. frame

$$d\sigma = d\sigma_{UU} \left[1 + S_T \sin \psi A_{UT}^{\sin \psi} \right]$$

$$\sin \psi = \mathbf{S}_T \cdot (\hat{\mathbf{P}}_T \times \hat{\mathbf{k}}) \text{ and } \hat{\mathbf{k}} = -\hat{\mathbf{p}}$$

- \uparrow (\downarrow) pols. are along $+Y_{\text{cm}}$ ($-Y_{\text{cm}}$) in both $p^\uparrow l$ or lp^\uparrow c.m. frames; recall $x_F = 2P_L/\sqrt{s}$, so:

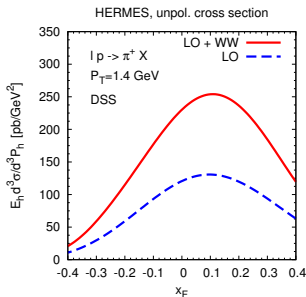
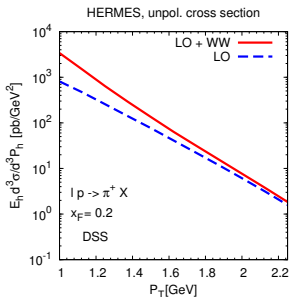
$$A_{UT}^{\sin \psi}(x_F, P_T) = A_N^{p^\uparrow l \rightarrow hX}(-x_F, P_T)$$

In $p^\uparrow l$ c.m. frame:

$$A_N = \frac{d\sigma^\uparrow - d\sigma^\downarrow}{d\sigma^\uparrow + d\sigma^\downarrow}$$

$$\begin{aligned} A(\phi_S, S_T) &= \mathbf{S}_T \cdot (\hat{\mathbf{p}} \times \hat{\mathbf{P}}_T) A_N \\ &= S_T \sin \phi_S A_N \end{aligned}$$

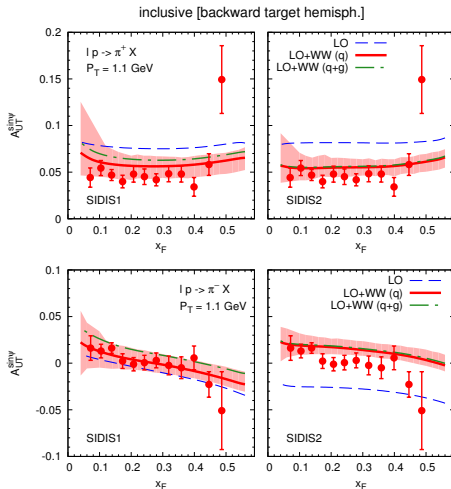
Unpolarized cross sections - HERMES



- $x_F > 0$ proton backward region
- WW contribution more important at $x_F > 0$: naively real photon ($Q^2 \approx 0$) for forward production but, **in backward region** $|\hat{u}| \ll |\hat{t}|$ and

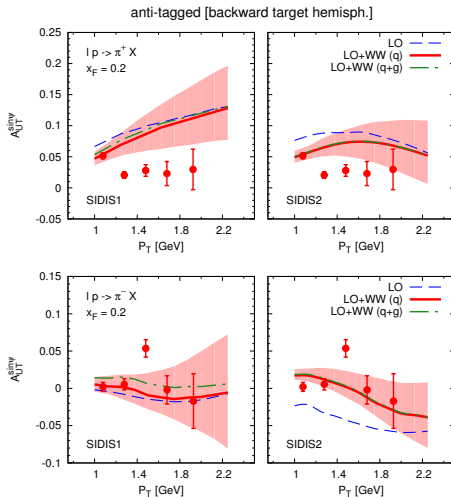
$$d\hat{\sigma}_{LO} \sim \frac{1}{\hat{t}^2}, \quad d\hat{\sigma}_{WW} \sim \frac{1}{\hat{s}\hat{u}}$$

A_N vs. x_F and comparison with HERMES data



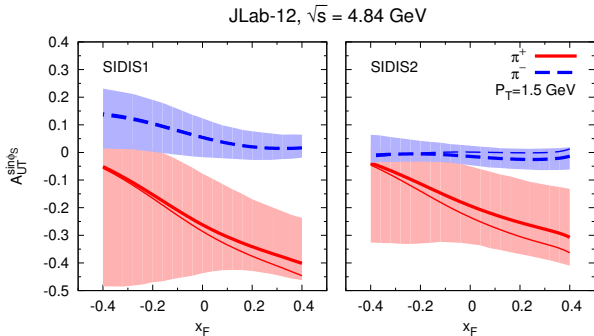
- Collins and gluon Sivers effects negligible
- bands are for LO+WW(q)

A_N vs. P_T and comparison with HERMES data



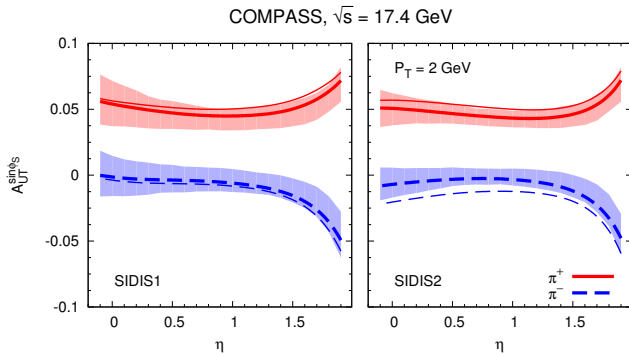
- discrepancies for π^+

A_N predictions - JLab



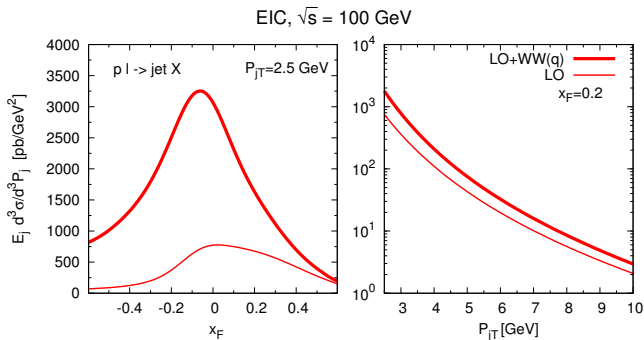
- role of u- and d-quark Sivers function exchanged while adopting SU(2) symmetry for neutron
- large $A_N(\pi^+)$ at large x_F due to the probed values of quark's light-cone momentum fraction (up to 0.1 at $x_F > 0$)
- wider uncertainty bands due to large-x region probed at moderate energies (extractions of Sivers function are unconstrained at large x)

A_N predictions - COMPASS



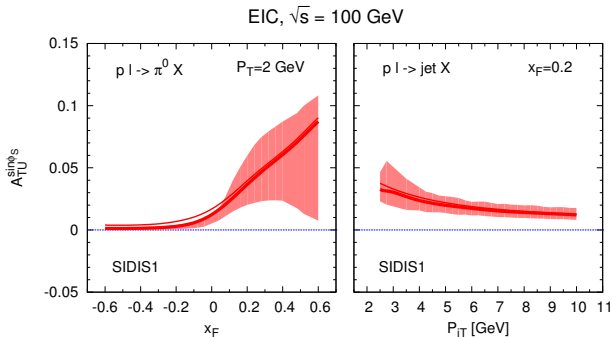
- sizeable SSA for π^+ with quite narrow error bands \Rightarrow need data to test the approach

$d\sigma_{\text{unp}}$ predictions - jet production at EIC



- here, $x_F > 0$ corresponds to the forward proton hemisphere
- pure Siverts effect - no fragmentation process! ($D_{h/q} \rightarrow 1$ and $\Delta^N D_{h/q\uparrow} \rightarrow 0$)

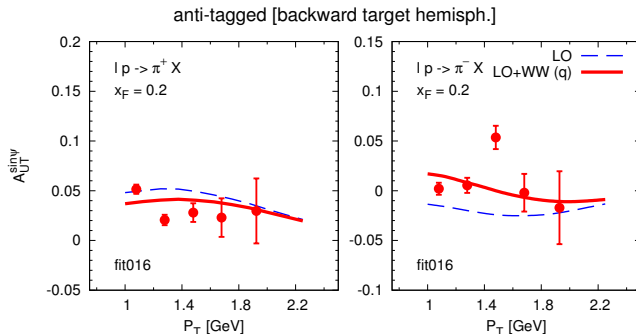
A_N predictions - EIC



- gluon Sivers effect completely negligible
- WW contribution does not change LO behaviour (both enters with the same structure in the SSA)
- interesting flat P_T behaviour - also measurable!

A_N vs. x_F and comparison with HERMES data - new fit

Using the most recent Sivvers fit [*Anselmino, Boglione, D'Alesio, Murgia, Prokudin, JHEP04 (2017), see Boglione's talk*]:



- different Gaussian widths and PDFs (CTEQ6L)

Conclusions

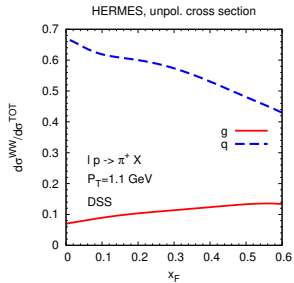
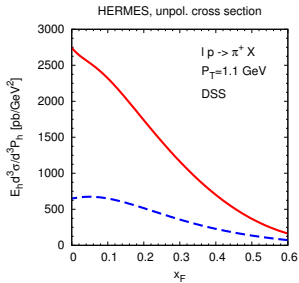
- computed (analytically and numerically), in the spirit of a **unified TMD approach**, WW contribution to A_N for $\ell p^\uparrow \rightarrow hX$
- role of quasi-real photon exchange: **huge** in unpolarised cross sections, **crucial** for SSAs description
- comparison with HERMES data **improves significantly** from previous LO analysis when quasi-real photon contribution is included
- found sizable predictions for JLab and COMPASS for π^+ production and same behaviour as $p^\uparrow p \rightarrow hX$ at EIC: fundamental to assess the **validity of TMD approach**
- further investigation (both experimental and theoretical) are required

Thanks a lot for your attention!

**Gratzias meda po' s'attenzione
bostra!**

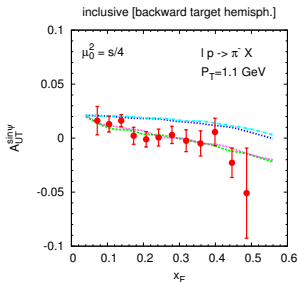
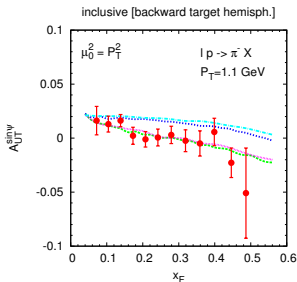
Backup

$d\sigma_{\text{unp}}$: LO and WW relative weights



- $\text{WW}/\text{TOT} \sim 70 \div 75\%$ at small x_F and still $\sim 60\%$ at large x_F
 \Rightarrow **dominance of WW contribution**
- gluon channel about 10%: γq and ℓq channels dominate

A_N : WW distribution and scale choices (I)

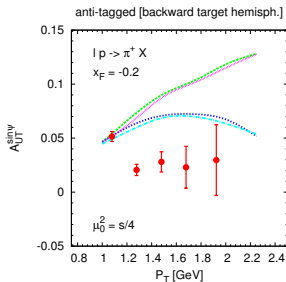
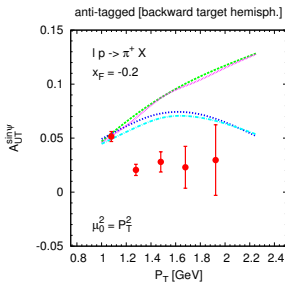


[Vogelsang et al., SIDIS1] [Vogelsang et al., SIDIS2]
 [Mukherjee et al., SIDIS1] [Mukherjee et al., SIDIS2]

[Brodsky, Kinoshita, Terazawa (1971), Kniehl (1990)]:

$$f_{\gamma/e}(y, E) = \frac{\alpha}{\pi} \left\{ \frac{1 + (1 - y)^2}{y} \left[\ln \frac{E}{m} - \frac{1}{2} \right] + \frac{y}{2} \left[\ln \left(\frac{2}{y} - 2 \right) + 1 \right] + \frac{(2 - y)^2}{2y} \ln \left(\frac{2 - 2y}{2 - y} \right) \right\}$$

A_N : WW distribution and scale choices (II)

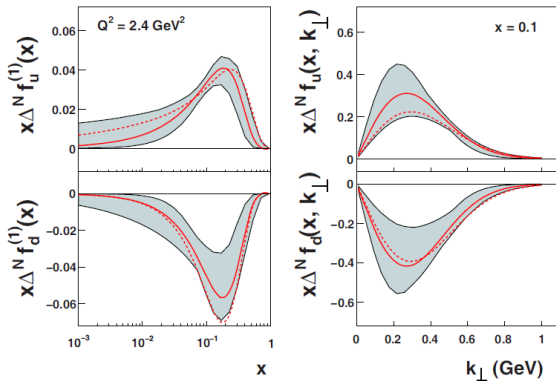


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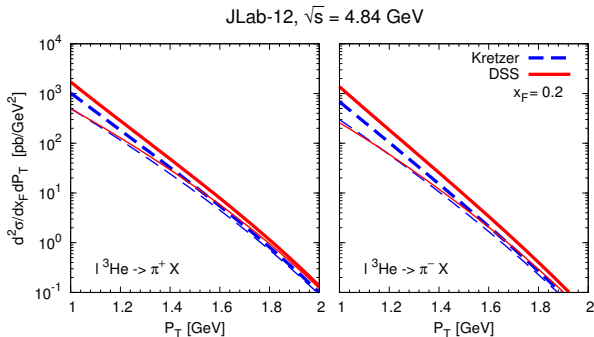
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Sivers function from SIDIS extractions



- Dashed line: SIDIS1 extraction (GRV98LO PDFs + Kretzer FFs)
- Solid line + band: SIDIS2 extraction (GRV98LO PDFs + DSS 2007 LO FFs)

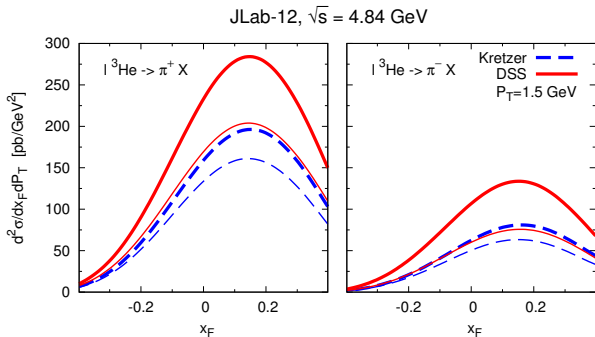
$d\sigma_{\text{unp}}$ predictions - JLAB (I)



$$\frac{d^2\sigma}{dx_F dP_T} = \frac{2\pi P_T}{\sqrt{x_F^2 + x_T^2}} E_\pi \frac{d^3\sigma}{d^3\mathbf{P}_\pi},$$

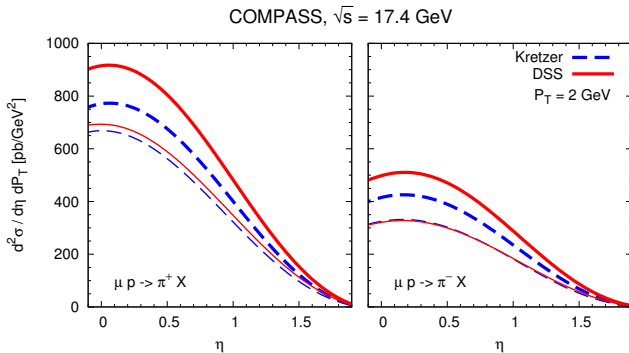
where $x_T = 2P_T/\sqrt{s}$.

$d\sigma_{\text{unp}}$ predictions - JLAB (II)



- LO different results due to large- z region explored more relevant for DSS FFs

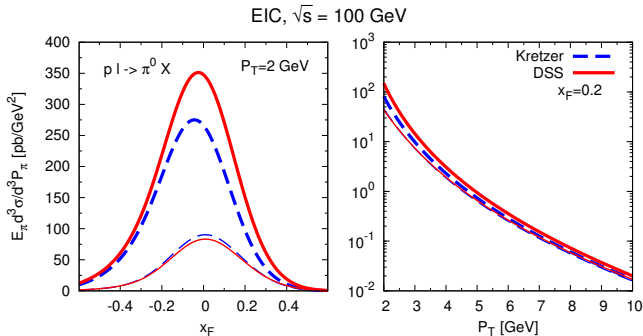
$d\sigma_{\text{unp}}$ predictions - COMPASS



$$\frac{d^2\sigma}{d\eta dP_T} = 2\pi P_T E_\pi \frac{d^3\sigma}{d^3\mathbf{P}_\pi}$$

- log term in WW distribution is smaller ($m_\mu \sim 200m_e$)

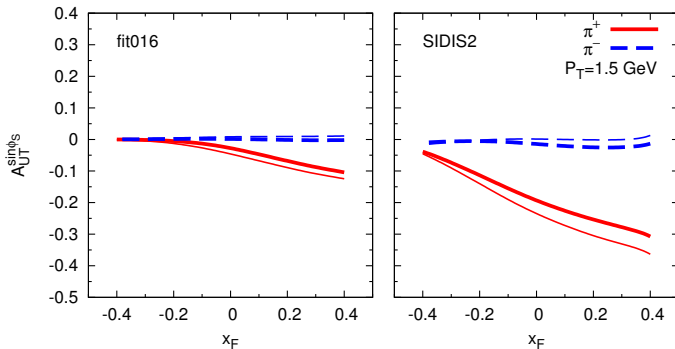
$d\sigma_{\text{unp}}$ predictions - π^0 at EIC



- here, $x_F > 0$ corresponds to the forward proton hemisphere
- pure Siverts effect - no fragmentation process!

A_N predictions - sign change fit at JLab

JLab-12, $\sqrt{s} = 4.84$ GeV



A_N predictions - sign change fit at COMPASS

COMPASS, $\sqrt{s} = 17.4$ GeV

