

Transverse dynamics of quarks in the proton –
from Ji to Jaffe-Manohar orbital angular momentum

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Acknowledgments:

M. Burkardt, S. Liuti,
and members of the Lattice TMD Collaboration

Gauge ensembles provided by:

MILC Collaboration

Proton spin decompositions

$$\frac{1}{2} = \frac{1}{2} \sum_q \Delta q + \sum_q L_q + J_g \quad (\text{Ji})$$

$$\frac{1}{2} = \frac{1}{2} \sum_q \Delta q + \sum_q \mathcal{L}_q + \Delta g + \mathcal{L}_g \quad (\text{Jaffe-Manohar})$$

...and many more (in fact, we will see a continuous interpolation between the two ...)

There isn't one unique way of separating quark and gluon orbital angular momentum – the different decompositions have different, legitimate meanings.

Quark orbital angular momentum

Interpreting terms in the energy-momentum tensor:

$$L_q \sim -i\psi^\dagger(\vec{r} \times \vec{D})_z\psi$$

Can be obtained from $L_q = J_q - S_q$, where S_q and J_q can be related to GPDs (Ji sum rule) – this has been used in Lattice QCD.

$$\mathcal{L}_q \sim -i\psi^\dagger(\vec{r} \times \vec{\partial})_z\psi \quad \text{in light cone gauge}$$

Hitherto not accessed in Lattice QCD.

Direct evaluation of quark orbital angular momentum

$$L_3^{\mathcal{U}} = \int dx \int d^2 k_T \int d^2 r_T (r_T \times k_T)_3 \mathcal{W}^{\mathcal{U}}(x, k_T, r_T) \quad \text{Wigner distribution}$$

$$\frac{L_3^{\mathcal{U}}}{n} = \frac{\epsilon_{ij} \frac{\partial}{\partial z_{T,i}} \frac{\partial}{\partial \Delta_{T,j}} \langle p', S | \bar{\psi}(-z/2) \gamma^+ \mathcal{U}[-z/2, z/2] \psi(z/2) | p, S \rangle |_{z^+=z^-=0, \Delta_T=0, z_T \rightarrow 0}}{\langle p', S | \bar{\psi}(-z/2) \gamma^+ \mathcal{U}[-z/2, z/2] \psi(z/2) | p, S \rangle |_{z^+=z^-=0, \Delta_T=0, z_T \rightarrow 0}}$$

n : Number of valence quarks

$$p' = P + \Delta_T/2, \quad p = P - \Delta_T/2, \quad P, S \text{ in 3-direction, } P \rightarrow \infty$$

This is the same type of operator as used in TMD studies – generalization to off-forward matrix element adds transverse position information

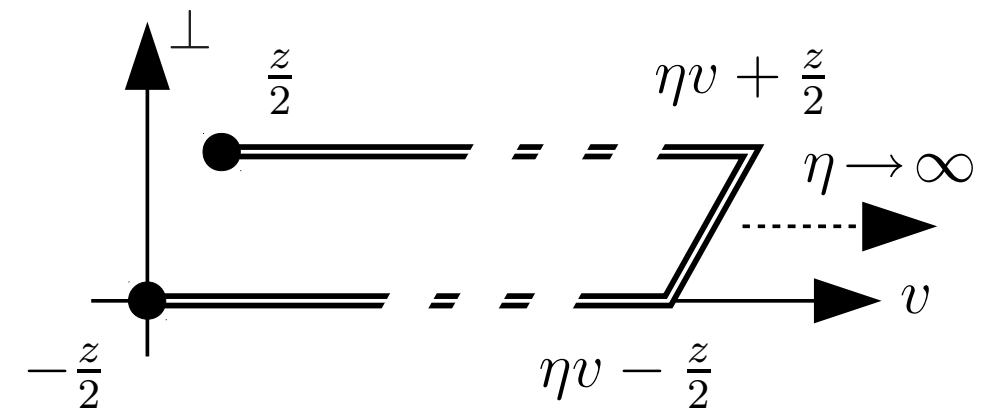
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Role of the gauge link \mathcal{U} :

Y. Hatta, M. Burkardt:

- Straight $\mathcal{U}[-z/2, z/2] \longrightarrow$ Ji OAM
- Staple-shaped $\mathcal{U}[-z/2, z/2] \longrightarrow$ Jaffe-Manohar OAM
- Difference is torque accumulated due to final state interaction



Direct evaluation of quark orbital angular momentum

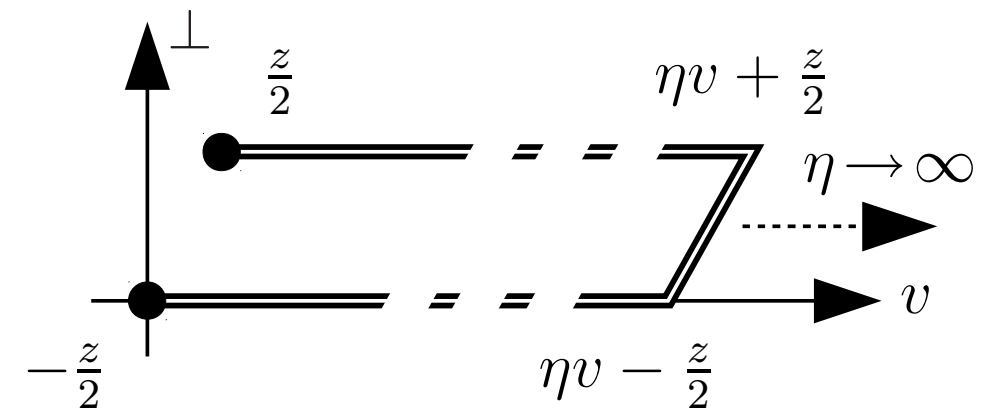
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Role of the gauge link \mathcal{U} :

Direction of staple taken off light cone (rapidity divergences)

Characterized by Collins-Soper parameter $\hat{\zeta} \equiv \frac{P \cdot v}{|P||v|}$

Are interested in $\hat{\zeta} \rightarrow \infty$; synonymous with $P \rightarrow \infty$ in the frame of the lattice calculation ($v = e_3$)



Ensemble details

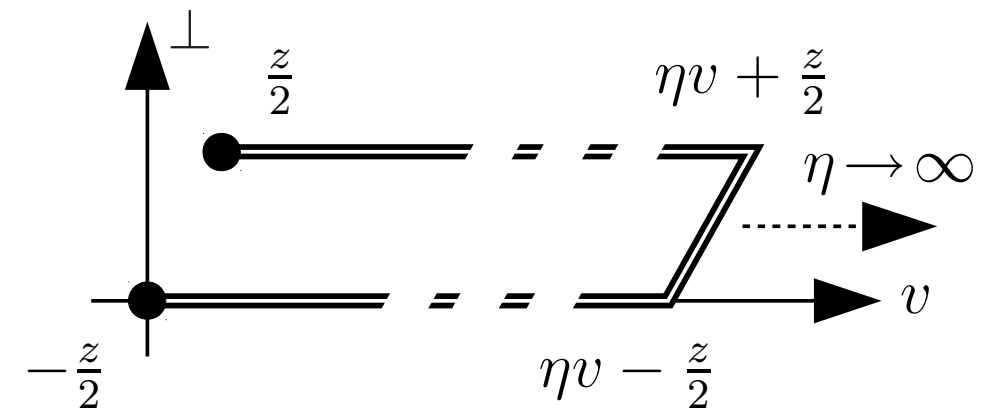
LHPC mixed action scheme: MILC asqtad configurations, domain wall valence quarks

$L^3 \times T$	$a(\text{fm})$	$am_{u,d}$	am_s	m_π^{DWF} (MeV)	m_N^{DWF} (GeV)	#conf.	#meas.
$20^3 \times 64$	0.11849(14)(99)	0.02	0.05	518.4(07)(49)	1.348(09)(13)	486	3888

Direct evaluation of quark orbital angular momentum

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Parameters to consider: $\Delta, \hat{\zeta}, z, \eta$



Direct evaluation of quark orbital angular momentum

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Dataset contains only one value of $|\Delta_T| = 4\pi/aL \approx 1$ GeV

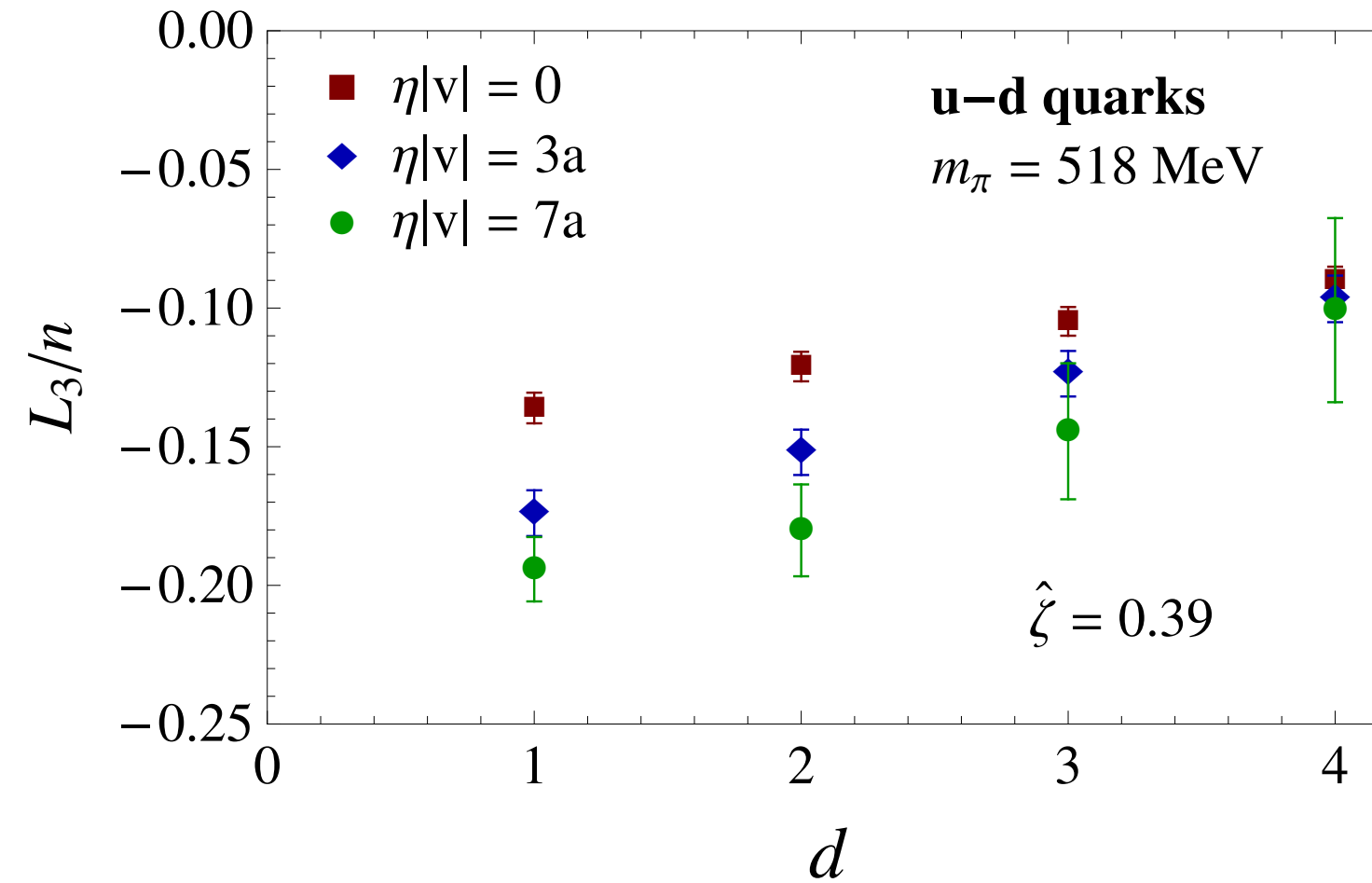
Substantial underestimate of $\partial f / \partial \Delta_T$ by using

$$\left. \frac{\partial f}{\partial \Delta_{T,j}} \right|_{\Delta_{T,j}=0} = \frac{1}{2\Delta_{T,j}} (f(\Delta_{T,j}) - f(-\Delta_{T,j}))$$

This issue is eliminated in current production run – direct derivative method
(G. M. de Divitiis, R. Petronzio, N. Tantalo)

- Take derivative a priori, analytically
- Direct Monte Carlo sampling of differentiated quantity

Direct evaluation of quark orbital angular momentum

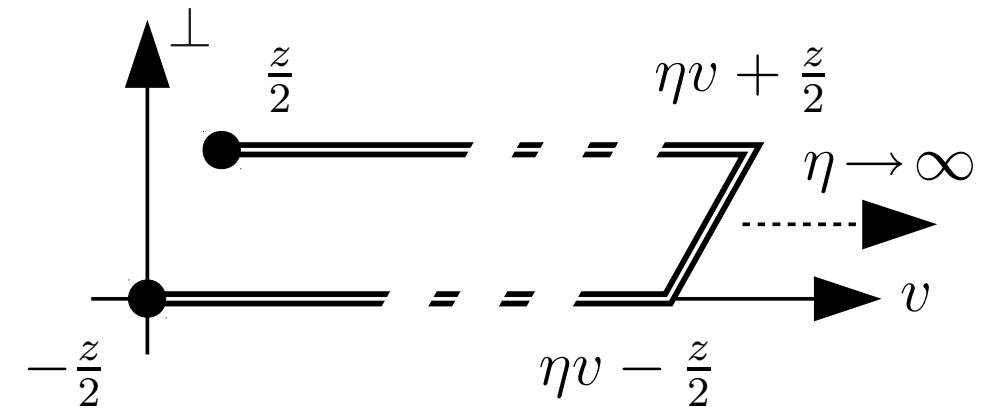


$$\left. \frac{\partial f}{\partial z_{T,i}} \right|_{z_{T,i}=0} = \frac{1}{2da} (f(dae_i) - f(-dae_i))$$

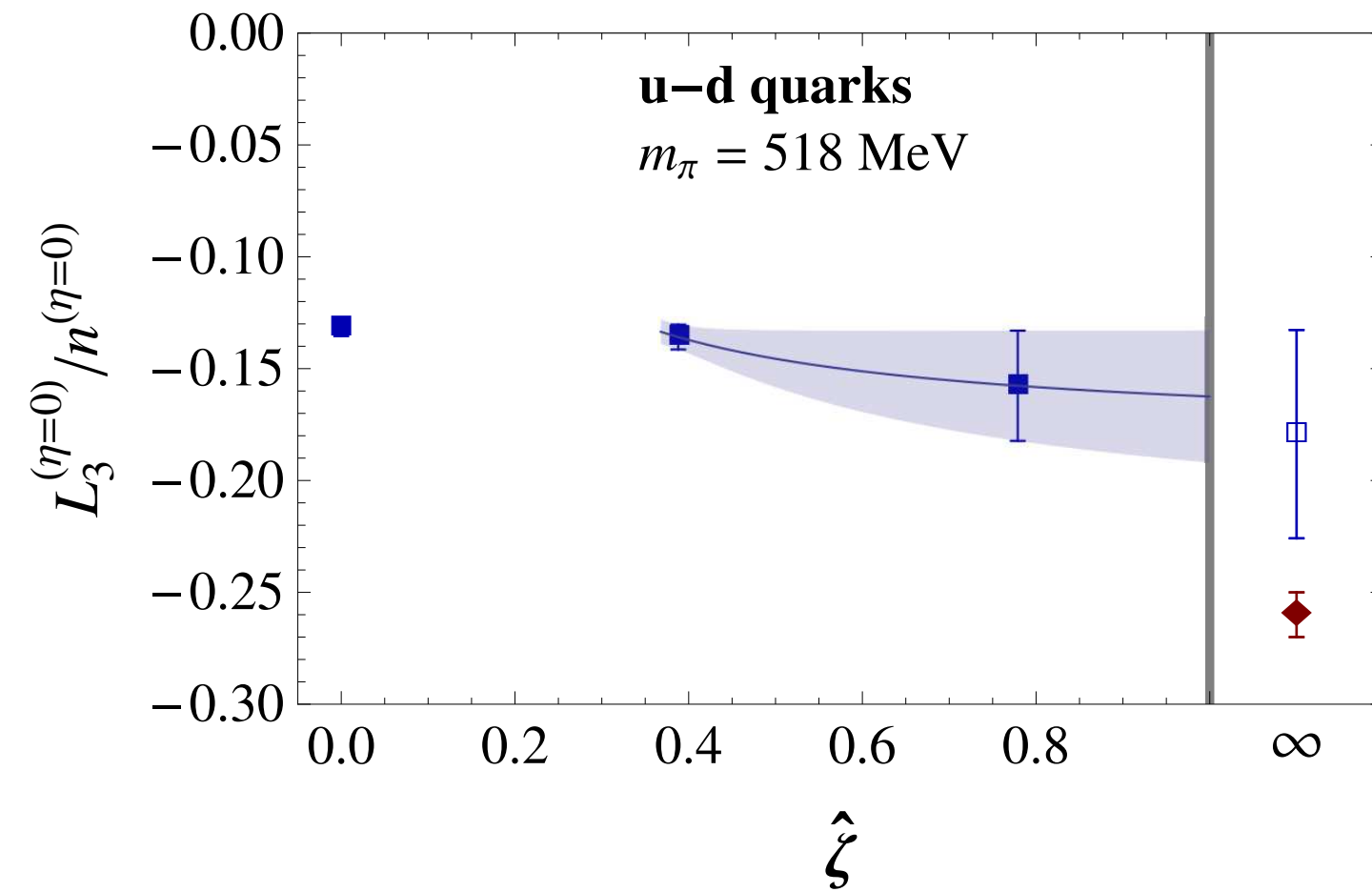
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Remaining parameters to consider: $\hat{\zeta}, \eta$

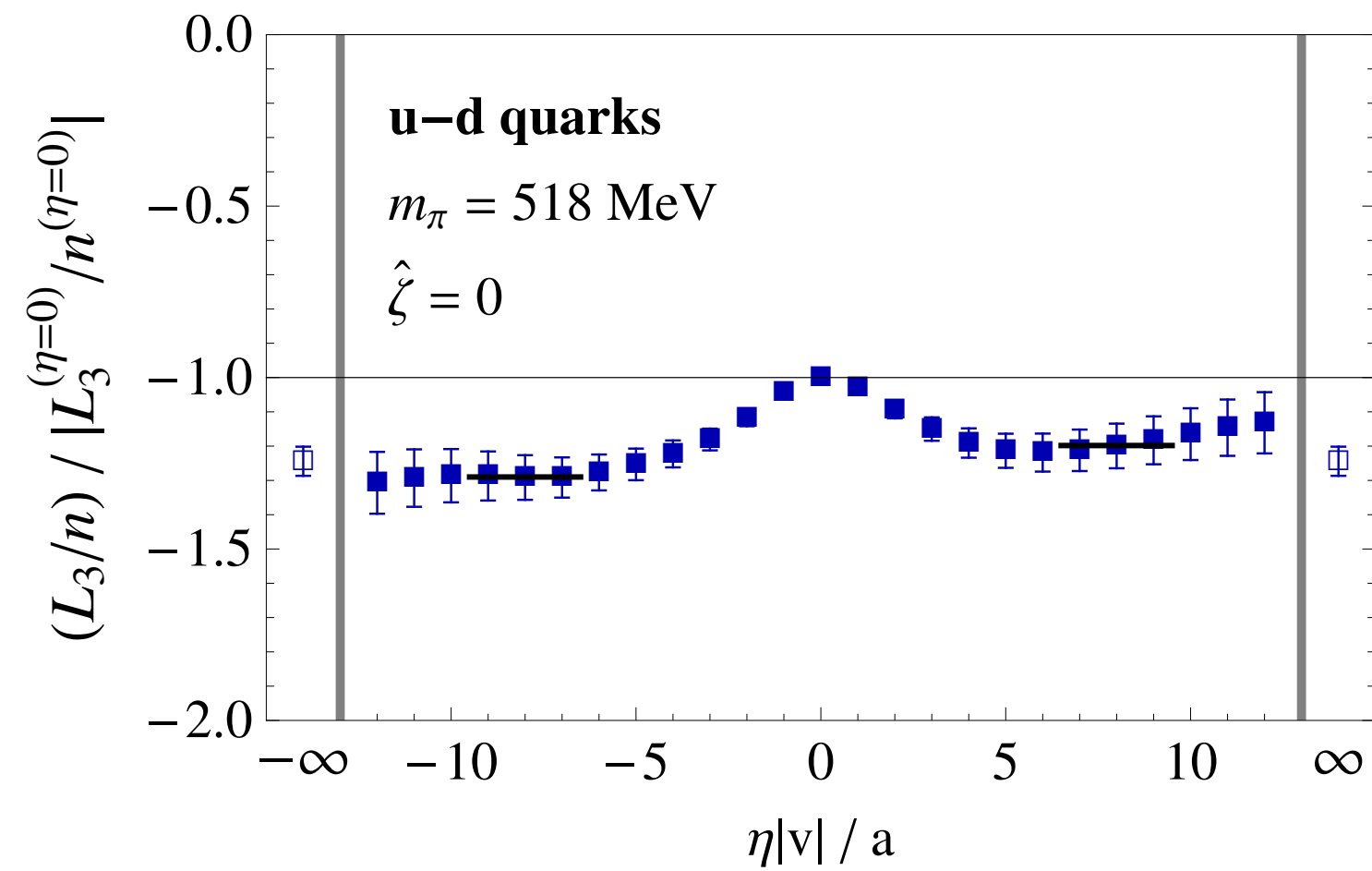


Ji quark orbital angular momentum: $\eta = 0$

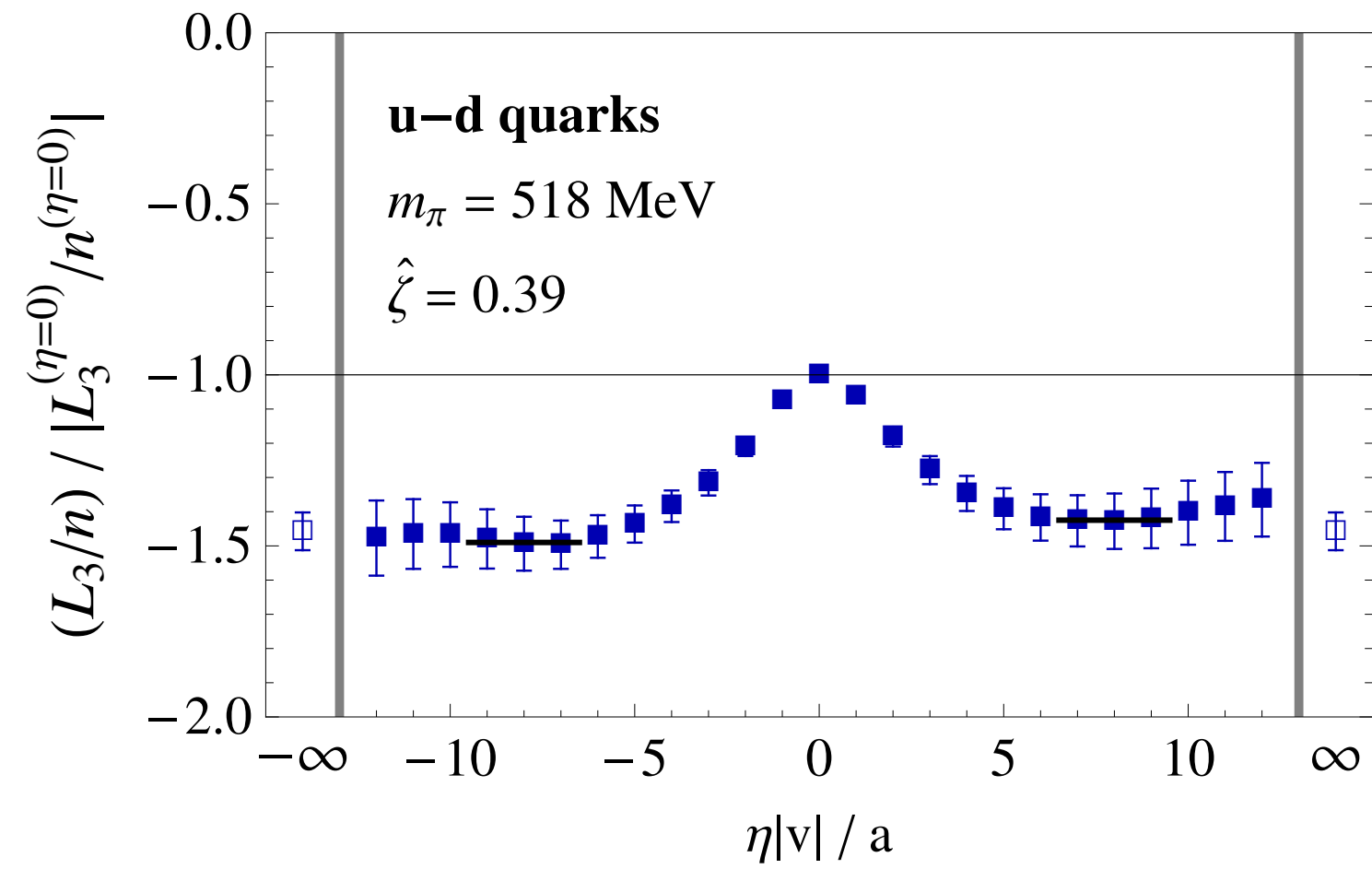


→ Signature of underestimate of $\partial f / \partial \Delta_T$

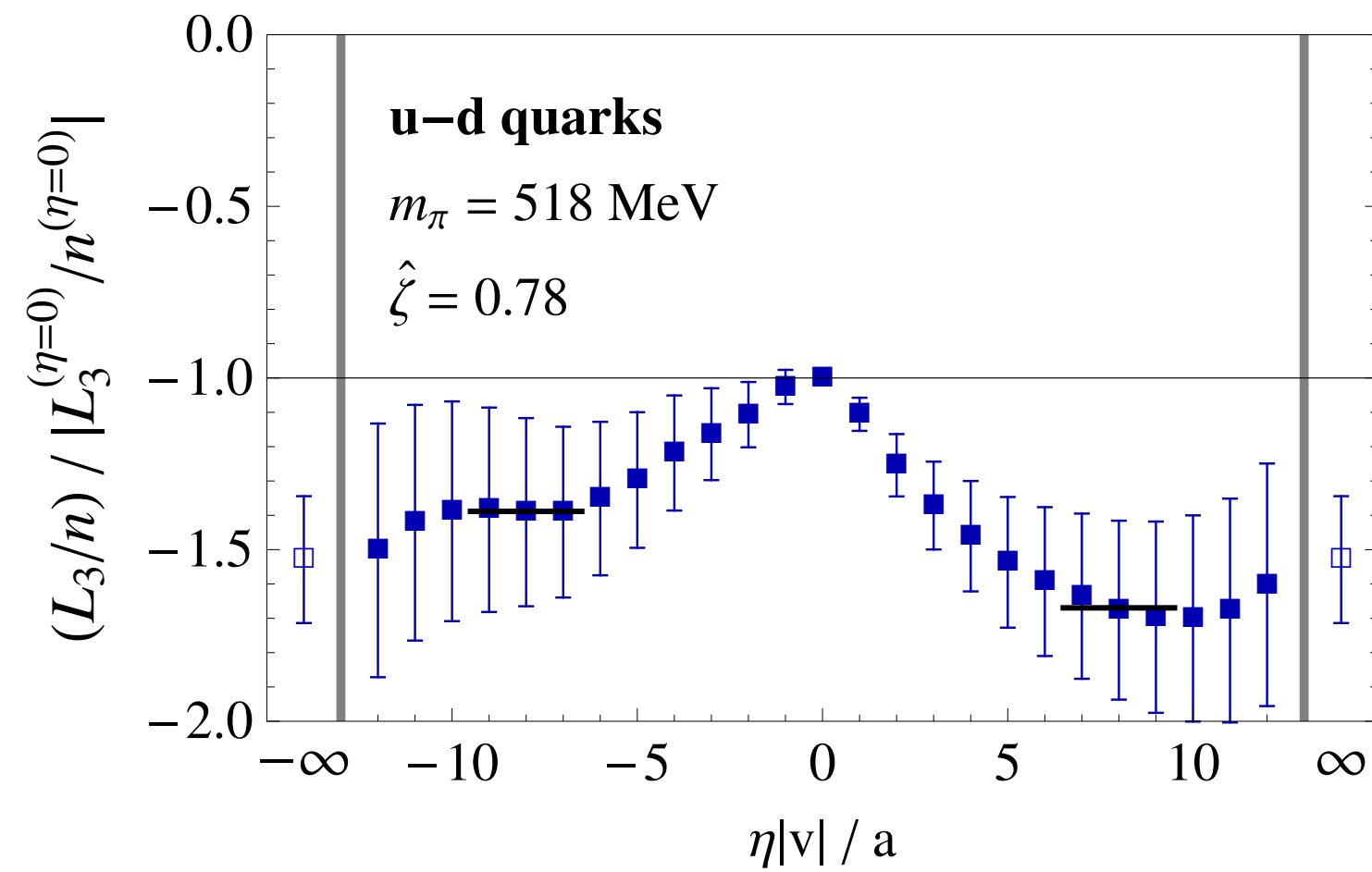
From Ji to Jaffe-Manohar quark orbital angular momentum



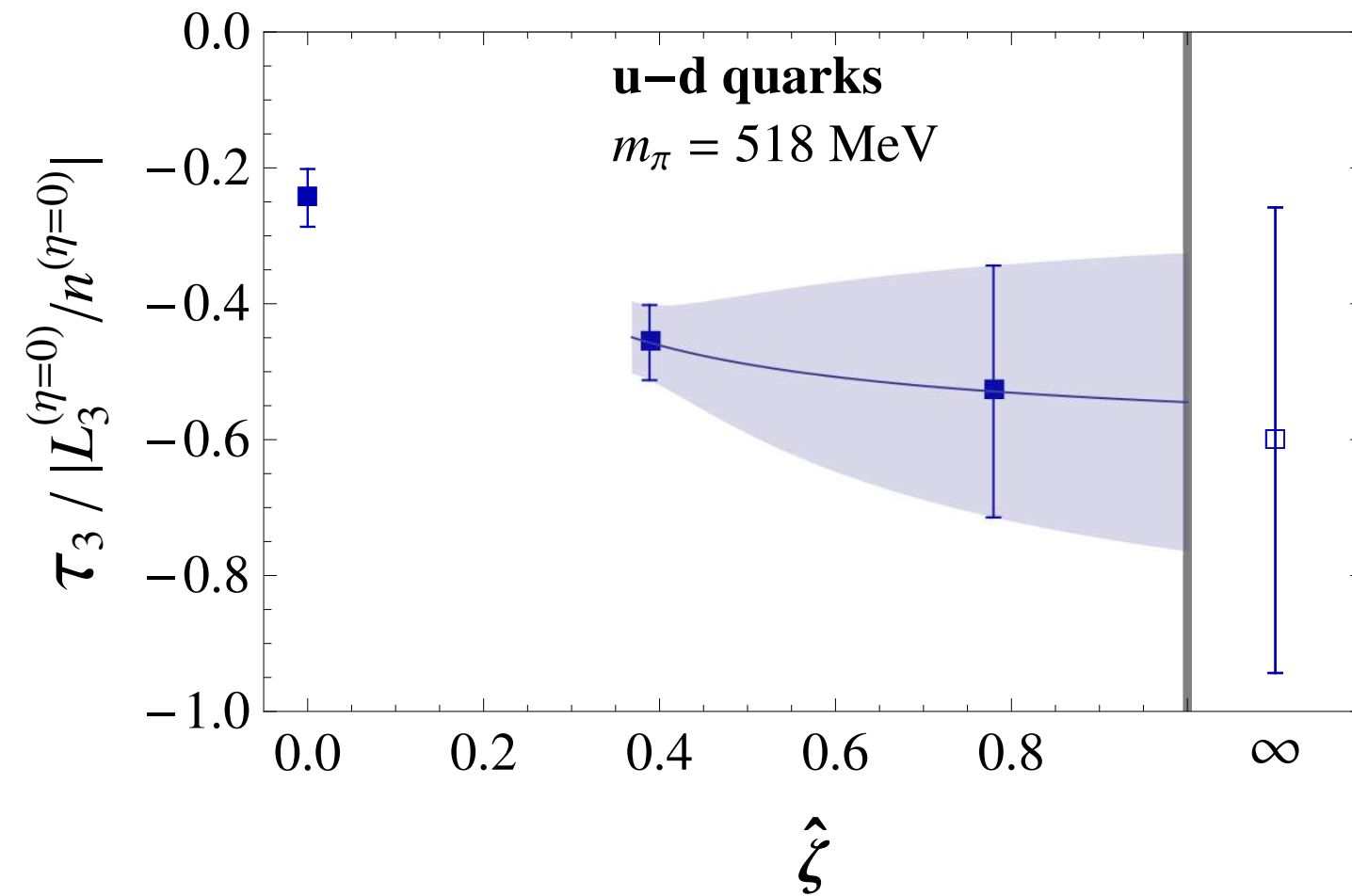
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From Ji to Jaffe-Manohar quark orbital angular momentum



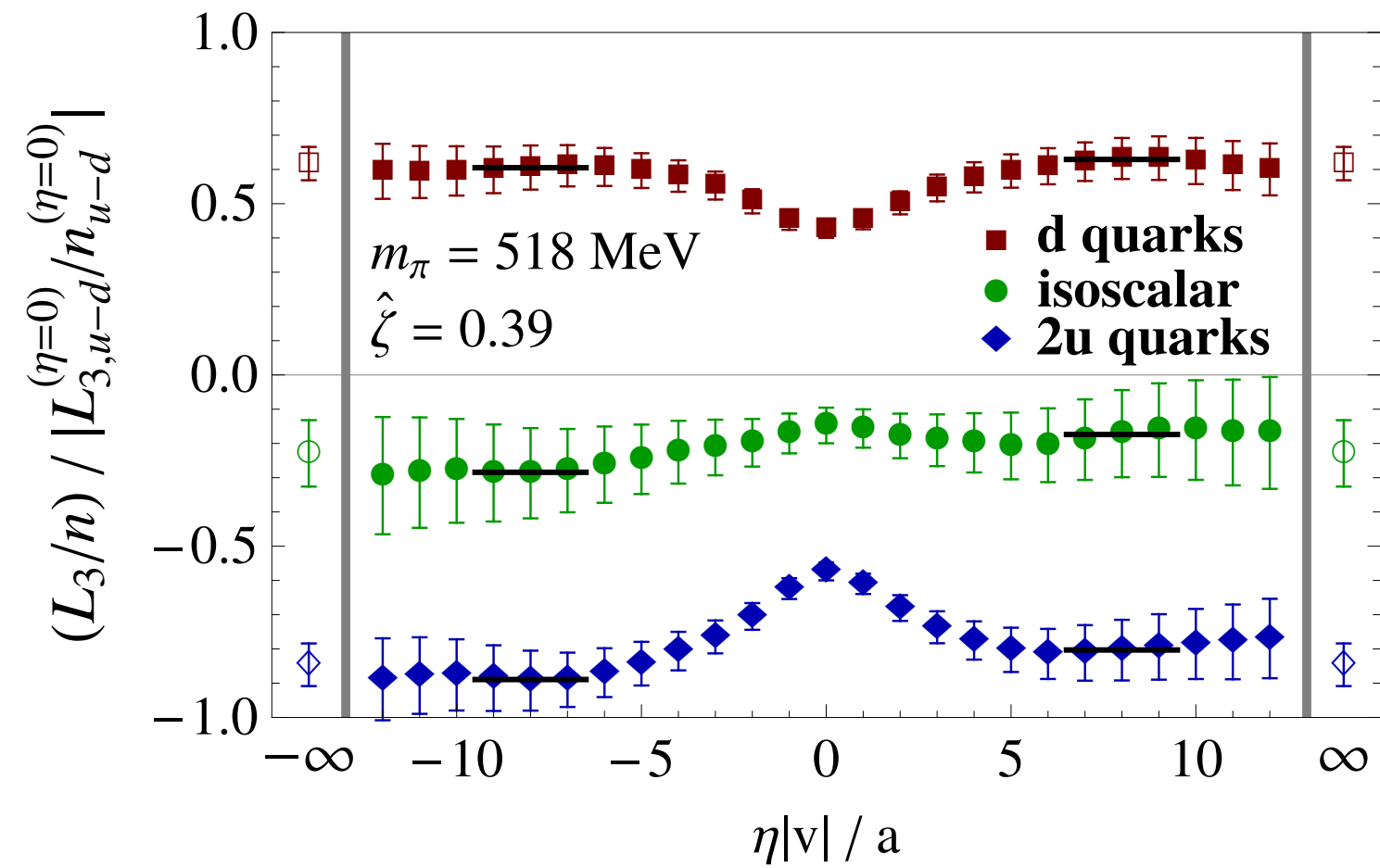
Burkardt's torque – extrapolation in $\hat{\zeta}$



$$\tau_3 = (L_3^{(\eta=\infty)} / n^{(\eta=\infty)}) - (L_3^{(\eta=0)} / n^{(\eta=0)})$$

Integrated torque accumulated by struck quark leaving proton

Flavor separation – from Ji to Jaffe-Manohar quark orbital angular momentum



Conclusions

- Quark orbital angular momentum can be accessed directly in Lattice QCD, continuously interpolating between the Ji and Jaffe-Manohar definitions.
- In the gathered exploratory dataset, the difference between the Ji and Jaffe-Manohar definitions, i.e., the torque accumulated by the struck quark leaving a proton, is clearly resolvable, sizeable ($\sim 50\%$ of the original Ji OAM), and leads to an enhancement of Jaffe-Manohar OAM relative to Ji OAM.
- Principal shortcoming of present dataset is too large momentum transfer. This practical issue is eliminated by using a direct derivative method in current data production.
- Plan to explore quark orbital angular momentum also via twist-3 GPD \tilde{E}_{2T} .