

TRANSVERSITY 2017

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INFN - FRASCATI NATIONAL LABORATORIES

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Partonic dynamics and the 3D structure of the nucleon: a global view

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From 3D atomic structure to quantum world, ...



Discovery:
Tiny nucleus - less than 1 trillionth in volume of an atom
Quantum probability - the Quantum World!

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Quantum probability - the Quantum World!

□ Localized mass and charge centers – vast "open" space:



From 3D hadron structure to QCD, ...

A modern "Rutherford" experiment (about 50 years ago):



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From 3D hadron structure to QCD, ...

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Outline of the rest of my talk

□ Next frontier of QCD and hadron physics? **Overarching questions?** □ How to "see" the hadron structure? See the "unseen"? □ How to quantify the hadron structure in QCD? Probabilities to catch the parton and its interaction? □ What do we hope to learn from future facilities? From JLab 12 GeV to a future Electron-Ion Collider From Lattice QCD – Complementary to experiments Summary and outlook

Next frontier of QCD & hadron physics, ...

How did hadrons, the building blocks of visible world, emerge from quarks and gluons?

Necessary knowledge for understanding where and how did we come from following the "Big Bang?

□ What is the internal structure of hadrons, and the dynamics behind the structure?

Necessary knowledge for understanding what are we made of, and what hold us together, as well as how do we improve and move forward – femtotechnology?

□ What is the key for understanding color confinement?

Necessary knowledge for understanding what is the mother nature of the nonlinear, strongly interacting dynamics of the color force?

□ Facts:

Gluons are dark!

No modern detector has been able to see quarks and gluons in isolation!

□ The intellectual challenge:

How to probe the quark-gluon dynamics, quantify the hadron structure, study the emergence of hadrons, ..., if we cannot see quarks and gluons?

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□ Answer to the challenge:

Theory advances:

QCD factorization – matching the quarks/gluons to hadrons with *controllable approximations*!

Experimental breakthroughs:

Jets – Footprints of energetic quarks and gluons

Quarks – Need an EM probe to "see" their existence, ...

Gluons – Varying the probe's resolution to "see" their effect, ...

Technology improvements:

Energy, luminosity, unprecedented resolution, ...

□ Hard probes to "catch" the quantum fluctuation:

Lorentz invariant cross sections are frame independent But, the physical picture of what happened is frame dependent

In any frame in which the proton is moving very fast,



Hard probe (t ~ 1/Q < fm) Probability to "catch" the parton!

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- \diamond Longitudinal momentum fraction **x**: $xP \sim Q$ the hard scale!
- \diamond Transverse momentum confined motion: $1/R \sim \Lambda_{
 m QCD} \ll Q k_T / k_T$
- ♦ Momentum transfer in diffractive t :

Spatial imaging: Mith the collision size: $\sim 1/Q$

 $t \ll Q \Rightarrow 1/b_T$

хp

□ What do we need to know for the structure?

 \Rightarrow In theory: $\langle P, S | \mathcal{O}(\overline{\psi}, \psi, A^{\mu}) | P, S \rangle$ – Hadronic matrix elements

of all possible operators: $\mathcal{O}(\overline{\psi},\psi,A^{\mu})$



Correlations between any number of fields in QCD

 k_T

♦ BUT:

None of these matrix elements is a direct physical observable – color confinement!

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- BUT: None of these matrix elements is a direct physical observable color confinement!
- In practice: Accessible hadron structure
 = hadron matrix elements of quarks and gluons, satisfying
 - 1) can be related to physical cross sections of hadrons and leptons with controllable approximation; and/or

2) can be calculated/extracted from lattice QCD

♦ Hard scale: Wave vs. particle nature of quarks and gluons?



Need probes/observables with at least ONE large-scale, & additional controllable small scale(s) to "see" the structure!

❑ Matching the observed hadron to the caught partons:

Any cross section with identified hadron(s) is NON-Perturbative!

Cross section with ONE identified hadron & ONE large momentum transfer:



❑ Matching the observed hadron to the caught partons:

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Cross section with ONE identified hadron & ONE large momentum transfer:



if there are more than one identified hadron, ...



We have learned a lot for 1D structure, ...

General expansion of quark distribution:



must have general expansion in terms of P, i, i etc. $\stackrel{P}{=}$

$$\phi(x) = \frac{1}{2} \left[q(x)\gamma \cdot P + s_{\parallel} \Delta q(x)\gamma_5\gamma \cdot P + \delta q(x)\gamma \cdot P\gamma_5\gamma \cdot S_{\perp} \right]$$

□ 3-leading power quark parton distribution:

$$q(x) = \frac{1}{4\pi} \int dz^{-} e^{iz^{-}xP^{+}} \langle P, S | \bar{\psi}(0) \gamma^{+} \psi \left(0, z^{-}, \mathbf{0}_{\perp}\right) | P, S \rangle$$

$$\Delta q(x) = \frac{1}{4\pi} \int dz^{-} e^{iz^{-}xP^{+}} \langle P, S | \bar{\psi}(0) \gamma^{+} \gamma_{5} \psi \left(0, z^{-}, \mathbf{0}_{\perp}\right) | P, S \rangle$$

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"unpolarized" – "longitudinally polarized" – "transversity"

Probed with observables with a single-hard scale!
Similar for gluon density and helicity distributions

PDFs of a spin-averaged proton

□ Modern sets of PDFs @NNLO with uncertainties:



Consistently fit almost all data with Q > 2GeV

Role of LHC data

□ Inclusive jet production at 7 TeV:



- Cross sections span 12 orders of magnitude
- Almost negligible statistical error

Partonic luminosities – Our discovery potential

q - qbar

g - **g**



Uncertainties are mainly in large- & small-x regimes – need both JLab12 & EIC

Future large-x experiments – JLab12

□ NSAC milestone HP14 (2018):



Plus many more JLab experiments:

E12-06-110 (Hall C on ³He), E12-06-122 (Hall A on ³He), E12-06-109 (CLAS on NH₃, ND₃), ... Future EIC help fix small-x PDFs! and Fermilab E906, ... Can lattice QCD help large-x?

Generation (spin-averaged):

See talks by Constantinou and Orginos, ...

$$\tilde{q}(x,\mu^2,P_z) \equiv \int \frac{d\xi_z}{4\pi} e^{-ixP_z\xi_z} \langle P|\overline{\psi}(\xi_z)\gamma_z \exp\left\{-ig\int_0^{\xi_z} d\eta_z A_z(\eta_z)\right\} \psi(0)|P\rangle + \text{UVCT}(\mu^2)$$

Proposed matching:

$$\tilde{q}(x,\mu^2,P_z) = \int_x^1 \frac{dy}{y} Z\left(\frac{x}{y},\frac{\mu}{P_z}\right) q(y,\mu^2) + \mathcal{O}\left(\frac{\Lambda^2}{P_z^2},\frac{M^2}{P_z^2}\right)$$

Ji, arXiv:1305.1539

ξ_z

Ζ

- Power divergence renormalization?
- Mixing with lower dimension operators cannot be treated perturbatively, ...
- Exploratory effort:

Lin et al., arXiv:1402.1462

0



□ Facts:

Ma and Qiu, 2014, 2017

- PDFs are time-independent, so as the factorized cross sections!
- The operators, defining PDFs, located on the light-cone is a consequence of the approximation defining the twist-2 factorization More precisely, the collinear approximation

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- The operators, defining PDFs, located on the light-cone is a consequence of the approximation defining the twist-2 factorization More precisely, the collinear approximation
- Beyond the quasi- and pseudo-PDFs:
 - NOT try to calculate PDFs directly from lattice QCD calculations
 - Calculate a set of time-independent (fixed or integrated over time) and good single hadron matrix elements – "lattice cross sections"

$$\begin{split} \sigma_n(\xi^2, \omega, P^2) &= \langle P | T\{\mathcal{O}_n(\xi)\} | P \rangle \qquad \omega = P \cdot \xi \\ \text{with} \qquad \mathcal{O}_{j_1 j_2}(\xi) &\equiv \xi^{d_{j_1} + d_{j_2} - 2} Z_{j_1}^{-1} Z_{j_2}^{-1} j_1(\xi) j_2(0) \\ j_S(\xi) &= \xi^2 Z_S^{-1} [\overline{\psi}_q \psi_q](\xi), \ j_G(\xi) &= \xi^3 Z_G^{-1} [-\frac{1}{4} F_{\mu\nu}^c F_{\mu\nu}^c](\xi) \end{split}$$

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 $\begin{array}{l} \diamondsuit \text{ Derive PDFs from Global Analysis of "data" on lattice cross sections} \\ \overline{\sigma}_{\mathrm{E}}^{\mathrm{Lat}}(\xi_{z}, 1/a, P_{z}) \xleftarrow{\mathcal{Z}} \sigma_{\mathrm{E}}(\xi_{z}, \tilde{\mu}^{2}, P_{z}) & - \text{Renormalization/continuum limit} \\ & & & & & & & \\ & & & & & & & \\ & & & & & & & \\ & & & & & & & \\ & & & & & & & \\ & & & & & & \\ & & & & & & \\ & & & & & & \\ & & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & & \\ & & & \\ & & & \\ & & & & \\ & & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & & \\ & & & & \\ & & &$

Paradigm shift: 3D imaging of hadrons

□ Cross sections with two-momentum scales observed: $Q_1 \gg Q_2 \sim 1/R \sim \Lambda_{\rm QCD}$

 \diamond "Soft" scale: $Q_2 \;$ could be more sensitive to hadron structure, e.g., confined motion



Paradigm shift: 3D imaging of hadrons

 $xp_{\star}k_{\rm T}$

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Two-scale observables with the hadron broken:



 $\diamond\,$ Natural observables with TWO very different scales

TMD factorization: partons' confined motion is encoded into TMDs

Paradigm shift: 3D imaging of hadrons

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 \diamond "Soft" scale: Q_2 could be more sensitive to hadron structure, e.g., confined motion

Two-scale observables with the hadron unbroken:



♦ Natural observables with TWO very different scales

♦ GPDs: Fourier Transform of t-dependence gives spatial b_T-dependence

Unified description of nucleon structure



See talks by Pasquini, ...

Coordinate Space

GPDs

Spatial distribution

Unified description of nucleon structure



Theory is solid – *TMDs & SIDIS as an example*:

 \diamond Low P_{hT} (P_{hT} << Q) – TMD factorization:

 $\sigma_{\text{SIDIS}}(Q, P_{h\perp}, x_B, z_h) = \hat{H}(Q) \otimes \Phi_f(x, k_\perp) \otimes \mathcal{D}_{f \to h}(z, p_\perp) \otimes \mathcal{S}(k_{s\perp}) + \mathcal{O}\left[\frac{P_{h\perp}}{Q}\right]$

 \Rightarrow High P_{hT} (P_{hT} ~ Q) – Collinear factorization:

$$\sigma_{\text{SIDIS}}(Q, P_{h\perp}, x_B, z_h) = \hat{H}(Q, P_{h\perp}, \alpha_s) \otimes \phi_f \otimes D_{f \to h} + \mathcal{O}\left(\frac{1}{P_{h\perp}}, \frac{1}{Q_{h\perp}}, \frac$$

$\Rightarrow \mathbf{P_{hT} Integrated} - \mathbf{Collinear factorization:} \\ \sigma_{\mathrm{SIDIS}}(Q, x_B, z_h) = \tilde{H}(Q, \alpha_s) \otimes \phi_f \otimes D_{f \to h} + \mathcal{O}\left(\frac{1}{O}\right)$

 $\diamond \text{ Very high } P_{hT} >> Q - Collinear factorization: } \\ \sigma_{\text{SIDIS}}(Q, P_{h\perp}, x_B, z_h) = \sum_{abc} \hat{H}_{ab \to c} \otimes \phi_{\gamma \to a} \otimes \phi_b \otimes D_{c \to h} + \mathcal{O}\left(\frac{1}{Q}, \frac{Q}{P_{h\perp}}\right)$

See talks by Rogers, Gamberg, Prokudin, ...

Hadron structure and Q²-dependence

\Box Variation of Q^2 at the hard collision:

Gluon shower – liberation of confined states



Hadron structure and Q²-dependence

□ Variation of Q² at the hard collision: *Gluon shower* – *liberation of confined states* – *get contributions from all powers* P P P (*xP, k_T*) + $\frac{P_h}{z}$, k'_T Confined motion

□ Single hard scale – collinear factorization:

 \Rightarrow DGLAP evolution (LT-approx.) use data at Q > Q₀ to help pin down PDFs at Q₀

 \diamond PDFs at Q₀ carry all non-perturbative structure information

Hadron structure and Q²-dependence

□ Variation of Q² at the hard collision: *Gluon shower* – *liberation of confined states* – *get contributions from all powers* P P P *xP, k_T* $\stackrel{P_h}{\xrightarrow{xP, k_T}}$ *Emergence of a hadron* hadronization

□ Single hard scale – collinear factorization:

- \Rightarrow DGLAP evolution (LT-approx.) use data at Q > Q₀ to help pin down PDFs at Q₀
- \diamond PDFs at Q₀ carry all non-perturbative structure information
- **Two scales TMD factorization:**
 - \diamond TMD factorization valid when P_T << Q, CS evolution in conjugate b_T-space
- ♦ When TMDs dominated by small b_T, shower dominated by perturbative logs, CS evolution help use all Q > Q₀ data to fix TMDs at Q₀ (Q₀ large, x small)
- ♦ When TMDs dominated by large b_T region, rich N.P. structure information,
 CS evolution does not help extract TMDs models, new challenge!

Evolution of the Sivers effect

□ Sivers Effect:

- Quantum correlation between the spin direction of colliding hadron and the preference of motion direction of its confined partons
- QCD Prediction: Sign change of Sivers function from SIDIS and DY

□ "Prediction" and large uncertainty of QCD evolution:



Evolution of the Sivers effect

STAR Collab. Phys. Rev. Lett. 116, 132301 (2016)

See talk by Aschenauer, ...



Theory curves (KQ) = No TMD evolution!also see M. Anselmino et al.
JHEP 1704 (2017) 046Data from STAR collaboration on A_N for W-production are consistent
with a sign change between SIDIS and DY

Hint of the sign change from lattice QCD

□ Gauge link for lattice calculation:

Engelhardt@TMD Collaboration meeting

Staple-shaped gauge link $\mathcal{U}[0, \eta v, \eta v + b, b]$



\Box Normalized moment of Sivers function – at given b_T :



Hunting for GPDs: Exclusive DIS



\Box Much more complicated – (x, ξ , t) variables:

♦ Challenge to derive GPDs from data

Hunting for GPDs: Exclusive DIS



 \Box Much more complicated – (x, ξ , t) variables:

- Challenge to derive GPDs from data
- GPDs could tell us:
 - Orbital contribution to proton's spin
 - Proton radius of quark & gluon density
 - ♦ Hints for color confining radius/mechanism
 - ♦ Origin of nuclear force, ...
 ♦ ...



GPDs: just the beginning



Hard probes from high energy collisions

□ Lepton-lepton collisions:





Hadrons

♦ No hadron in the initial-state

- ♦ Hadrons are emerged from energy
- Not for studying hadron structure

Hard probes from high energy collisions

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Hadron-hadron collisions:





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Hadrons

- ♦ Initial hadrons broken collision effect, ... \diamond Hadron structure – motion of quarks, ...
- \diamond Emergence of hadrons, ...

Hard probes from high energy collisions

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Lepton-hadron collisions:

 e^+ $\gamma */Z^0$ f Har

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Hadrons

♦ Initial hadrons broken – collision effect, ..
♦ Hadron structure – motion of quarks, ...
♦ Emergence of hadrons, ...

Hard collision without breaking the initial-state hadron – spatial imaging, ...

Many complementary probes at one facility

□ The future "Rutherford" experiment:



 $Q^2 \rightarrow$ Measure of resolution

- $y \rightarrow$ Measure of inelasticity
- $X \rightarrow$ Measure of momentum fraction

of the struck quark in a proton $Q^2 = S \times V$

Inclusive events: $e+p/A \rightarrow e'+X$ Detect only the scattered lepton in the detector

(Modern Rutherford experiment!)

<u>Semi-Inclusive events</u>: $e+p/A \rightarrow e'+h(p,K,p,jet)+X$

Detect the scattered lepton in coincidence with identified hadrons/jets (Initial hadron is broken – confined motion! – cleaner than h-h collisions) <u>Exclusive events:</u> e+p/A → e'+ p'/A'+ h(p,K,p,jet) Detect every things including scattered proton/nucleus (or its fragments) (Initial hadron is NOT broken – tomography! – almost impossible for h-h collisions)

From JLab12 to Electron-Ion Collider (EIC)

□ A giant "Microscope" – "see" quarks and gluons by breaking the hadron



From JLab12 to Electron-Ion Collider (EIC)

□ A giant "Microscope" – "see" quarks and gluons by breaking the hadron

p

To see the dynamics of the glue!

A sharpest "CT" – "imagine" quark/gluon structure without breaking the hadron

- "cat-scan" the nucleon and nuclei with a better than 1/10 fm resolution
- "see" proton "radius" of quark/gluon density as a function of "x" – another tunable knob!

To discover color confining radius, hints on confining mechanism!



1/Q

< 1/10 fm

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JLab12 – valence quarks, EIC – sea quarks and gluons

Impact of future SoLID at JLab12

□ Transversity & Tensor charge:

Z. Ye et al. Phys Lett. B767, 91 (2017)



Bayesian statistics is used to estimate the improvement from new data Current knowledge corresponds to a fit with TMD evolution Kang et al., P.R. D93 (16) 014009

Order of magnitude improvement in determining the tensor charge!

"full" is contribution from $0 \le x \le 1$ region "truncated" is contribution from $0.05 \le x \le 0.6$

Why US-EIC can do what HERA can't do?

Quantum imaging:

See talks by R. Ent,

- ♦ HERA discovered: 15% of e-p events is diffractive Proton not broken!
- ♦ US-EIC: 100-1000 times luminosity Critical for 3D tomography!

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- Quantum interference & entanglement:
 - VS-EIC: Highly polarized beams Origin of hadron property: Spin, ... Direct access to chromo-quantum interference!



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Quantum imaging:

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Quantum interference & entanglement:

VS-EIC: Highly polarized beams – Origin of hadron property: Spin, ... Direct access to chromo-quantum interference!



 US-EIC: Light-to-heavy nuclear beams – Origin of nuclear force, ... Help Catch the transition from chromo-quantum fluctuation to chromo-condensate of gluons,
 Emergence of hadrons (femtometer size detector!),
 – "a new controllable knob" – Atomic weight of nuclei

Summary

- QCD has been extremely successful in interpreting and predicting high energy experimental data!
- But, we still do not know much about hadron structure – work just started!



Round table discussion on this Wednesday!

- Cross sections with large momentum transfer(s) and identified hadron(s) are the source of structure information
- QCD factorization is necessary for any controllable "probe" for hadron's quark-gluon structure!
- EIC is a ultimate QCD machine, will provide answers to many of our questions on hadron structure, in particular, the confined transverse motions (TMDs), spatial distributions (GPDs), and multi-parton correlations, ...

Thank you!

Backup slides

What could we learn about hadron structure?

No elastic color current form factor!



□ Single-parton structure "seen" by a short-distance probe:

♦ 5D structure – encoded into the following density distribusions:

 $\begin{array}{cccc} & & & \\ & & & & \\ & & & & \\ & & & \\ & & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & &$

What could we learn about hadron structure?

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Multi-parton, quark-gluon corretions:

Single-spin asymmetry: $\propto [\sigma(Q, \vec{s}) - \sigma(Q, -\vec{s})]$

