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Outline



- Helicity amplitude ratios for hard exclusive ρ^0 production
- Spin-asymmetries in semi-inclusive DIS on transversely polarized hydrogen target
- Beam-helicity asymmetry in semi-inclusive DIS on unpolarized hydrogen and deuterium targets

Helicity amplitude ratios for exclusive ρ^0



.e

γ

t

N(p)

ρ0

N(p')



Exclusive meson production



Exclusive meson production

- probe various types of GPDs with different sensitivity and different flavour combinations
- complementary to DVCS

Target polarization state

- unpolarized target: nucleon-helicity-non-flip GPDs H, \tilde{H} and $\overline{E}_T=2H_T+\widetilde{E}_T$.
- transversely polarized target: nucleon-helicity-flip GPDs E, \tilde{E} and H_T.

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 $\gamma^*(\lambda_{\gamma}) + N(\lambda_N) \to V(\lambda_V) + N(\lambda'_N)$

• Helicity amplitude $F_{\lambda_V \lambda'_N \lambda_\gamma \lambda_N}$



 $\gamma^*(\lambda_{\gamma}) + N(\lambda_N) \to V(\lambda_V) + N(\lambda'_N)$







$$\gamma^*(\lambda_{\gamma}) + N(\lambda_N) \to V(\lambda_V) + N(\lambda'_N)$$

• Helicity amplitude
$$F_{\lambda_V\lambda_N'\lambda_\gamma\lambda_N}$$



$$F_{\lambda_V \lambda'_N \lambda_\gamma \lambda_N} = T_{\lambda_V \lambda'_N \lambda_\gamma \lambda_N} + U_{\lambda_V \lambda'_N \lambda_\gamma \lambda_N}$$

$$\gamma^*(\lambda_{\gamma}) + N(\lambda_N) \to V(\lambda_V) + N(\lambda'_N)$$



$$\gamma^*(\lambda_{\gamma}) + N(\lambda_N) \to V(\lambda_V) + N(\lambda'_N)$$



 $\propto F_{\lambda_V \lambda'_N \lambda_\gamma \lambda_N} \Sigma^{\alpha}_{\lambda_\gamma \lambda'_\gamma} F^*_{\lambda'_V \lambda'_N \lambda'_\gamma \lambda_N}$

Exclusive ρ^0 production: angular distribution

$$e + N \rightarrow e + N + \rho^0$$

 $\rho^0 \rightarrow \pi^+ + \pi^-$

Exclusive ρ^0 production: angular distribution



Fit angular distribution of decay pions $\mathcal{W}(\Phi,\phi,\Theta,\Psi)$ and extract either Spin Density Matrix Elements (SDMEs) or helicity amplitude ratios

Exclusive ρ^0 production: angular distribution

- transversely polarized H target
- longitudinally polarized $e^{\scriptscriptstyle\pm}$
- 8741 hard exclusive ρ^0 events

 $3.0 \text{ GeV} \leq W \leq 6.3 \text{ GeV}$ $1.0 \text{ GeV}^2 \leq Q^2 \leq 7.0 \text{ GeV}^2$ $0.0 \text{ GeV}^2 \leq -t' \leq 0.4 \text{ GeV}^2$

Fit angular distribution of decay pions $\mathcal{W}(\Phi,\phi,\Theta,\Psi)$ and extract either Spin Density Matrix Elements (SDMEs) or helicity amplitude ratios



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2% uncertainty beam polarization

already obtained in EPJ C71 (2011) 1609

extracted for first time

• 5 classes of helicity amplitude ratios



^{2%} uncertainty beam polarization



2% uncertainty beam polarization



2% uncertainty beam polarization



model for protons - S. Goloskokov and P. Kroll, Eur. Phys. J. C 50 (2007) 829; 53 (2008) 367, Eur. Phys. J. A 50 (2014) 146

$$F_{\lambda_{V}} \underbrace{\frac{1}{2}}_{\gamma=\lambda_{V}\frac{1}{2}} \propto \sum_{q,g} \mathcal{I}\left[\mathcal{A} \times \left(H^{a}, \frac{\xi^{2}}{1-\xi^{2}}E^{a}\right) + \mathcal{A}' \times \left(\tilde{H}^{a}, \frac{\xi^{2}}{1-\xi^{2}}\tilde{E}^{a}\right)\right]$$
$$F_{\lambda_{V}} \underbrace{\frac{1}{2}}_{\gamma=\lambda_{V}\frac{1}{2}} \propto \sum_{q,g} \mathcal{I}\left[\mathcal{A} \times E^{a} + \mathcal{A}' \times \xi\tilde{E}^{a}\right]$$

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$$F_{\lambda_{V}\left(\frac{1}{2}\lambda\right)} = \lambda_{V}\frac{1}{2} \propto \sum_{q,g} \mathcal{I}\left[\mathcal{A} \times \left(H^{a}, \frac{\xi^{2}}{1-\xi^{2}}E^{a}\right) + \mathcal{A}' \times \left(\tilde{H}^{a}, \frac{\xi^{2}}{1-\xi^{2}}\tilde{E}^{a}\right)\right]$$

$$F_{\lambda_{V}\left(\frac{1}{2}\right)_{\gamma} = \lambda_{V}\frac{1}{2}} \propto \sum_{q,g} \mathcal{I}\left[\mathcal{A} \times E^{a} + \mathcal{A}' \times \xi\tilde{E}^{a}\right] \qquad \text{natural parity}$$

$$natural parity$$

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$$I = \lambda_{V}\frac{1}{2} \propto \sum_{q,g} \mathcal{I}\left[\mathcal{A} \times E^{a} + \mathcal{A}' \times \xi\tilde{E}^{a}\right]$$

Factorization only proven for $\gamma_L^* \to V_L$.

Assumed for other transitions.

IR singularities regularised by modified perturbative approach.

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$$\begin{split} F_{\lambda_{V}\frac{1}{2}\lambda} = &\lambda_{V}\frac{1}{2} \propto \sum_{q,g} \mathcal{I} \left[\mathcal{A} \times \left(H^{a}, \frac{\xi^{2}}{1 - \xi^{2}} E^{a} \right) + \mathcal{A}' \times \left(\tilde{H}^{a}, \frac{\xi^{2}}{1 - \xi^{2}} \tilde{E}^{a} \right) \right] \\ F_{\lambda_{V}\frac{1}{2}\gamma} = &\lambda_{V}\frac{1}{2} \propto \sum_{q,g} \mathcal{I} \left[\mathcal{A} \times E^{a} + \mathcal{A}' \times \xi \tilde{E}^{a} \right] \\ \end{split}$$

Factorization only proven for $\gamma_L^* \to V_L$.

Assumed for other transitions.

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$\pi\omega$ transition form factor extracted from ω SDMEs

$$u_1 = 1 - r_{00}^{04} + 2r_{1-1}^{04} - 2r_{11}^1 - 2r_{1-1}^1$$



with pion-pole contribution with pion-pole contribution pion-pole contribution seems to account completely for unnatural-parity exchange

$\pi\omega$ transition form factor extracted from ω SDMEs

$$u_1 = 1 - r_{00}^{04} + 2r_{1-1}^{04} - 2r_{11}^1 - 2r_{1-1}^1$$



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Eur. Phys. J. C 77 (2017) 378





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GPD H.



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• GPD \tilde{H} + pion pole.



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• Only pion pole. Positive form factor.



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• GPD H_T.

Re t 8 lm t Reu Im u lm Re u A: $\gamma_{T} \rightarrow \rho_{T}$ Im u **⊢●_**<u>|</u>__ Re t₀₀ B: $\gamma_{I} \rightarrow \rho_{L}$ lm Re t₀₁ ₽. $\operatorname{Im} t_{01}^{(\cdot)}$ Re t₀₁ **HERMES** data ٠ GK model, positive sign of $\pi\rho$ FF Δ Im t₀₁ \Box GK model, negative sign of $\pi \rho$ FF Re u₀₁ $C: \gamma_{\tau} \rightarrow \rho_{L}$ Im u₀₁ Re t (1) lm ť Re t₁₀ lm Re u D: $\gamma \rightarrow \rho_{\tau}$ Im u Re 1 Im Re lm (2) Re u[`]1-1 (2) $E: \gamma_T \rightarrow \rho_{-T}$ lm u'⊢′ 0.8 Amplitude ratios 0.2 0.4 0 -0.2 0.6

Eur. Phys. J. C 77 (2017) 378

• Only pion pole. Positive form factor.

Semi-inclusive DIS single-hadron production


Presented amplitudes

$$\begin{aligned} \sigma^{h}(\phi,\phi_{S}) &= \sigma_{UU}^{h} \left\{ 1 + 2\langle\cos(\phi)\rangle_{UU}^{h} \cos(\phi) + 2\langle\cos(2\phi)\rangle_{UU}^{h} \cos(2\phi) \\ &+ \lambda_{l} 2\langle\sin(\phi)\rangle_{LU}^{h} \sin(\phi) \\ &+ S_{L} \left[2\langle\sin(\phi)\rangle_{UL}^{h} \sin(\phi) + 2\langle\sin(2\phi)\rangle_{UL}^{h} \sin(2\phi) \\ &+ \lambda_{l} \left(2\langle\cos(0\phi)\rangle_{LL}^{h} \cos(0\phi) + 2\langle\cos(\phi)\rangle_{LL}^{h} \cos(\phi) \right) \right] \\ &+ S_{T} \left[2\langle\sin(\phi - \phi_{S})\rangle_{UT}^{h} \sin(\phi - \phi_{S}) + 2\langle\sin(\phi + \phi_{S})\rangle_{UT}^{h} \sin(\phi + \phi_{S}) \\ &+ 2\langle\sin(3\phi - \phi_{S})\rangle_{UT}^{h} \sin(3\phi - \phi_{S}) + 2\langle\sin(\phi_{S})\rangle_{UT}^{h} \sin(\phi_{S}) \\ &+ 2\langle\sin(2\phi - \phi_{S})\rangle_{UT}^{h} \sin(2\phi - \phi_{S}) \\ &+ \lambda_{l} \left(2\langle\cos(\phi - \phi_{S})\rangle_{LT}^{h} \cos(\phi - \phi_{S}) \\ &+ 2\langle\cos(\phi_{S})\rangle_{LT}^{h} \cos(\phi_{S}) + 2\langle\cos(2\phi - \phi_{S})\rangle_{LT}^{h} \cos(2\phi - \phi_{S}) \right) \right] \end{aligned}$$

Presented here

$$Q^2 > 1 \text{ GeV}^2$$

 $W^2 > 10 \text{ GeV}^2$
 $0.023 < x < 0.6$

Presented amplitudes

• Longitudinally polarized e⁺/e⁻ beam

- Transversely polarized H target: fit all amplitudes simultaneously
 - ✤ Results for charged pions, kaons, (anti-)protons, neutral pions

$$+ S_{T} \left[2 \langle \sin(\phi - \phi_{S}) \rangle_{UT}^{h} \sin(\phi - \phi_{S}) + 2 \langle \sin(\phi + \phi_{S}) \rangle_{UT}^{h} \sin(\phi + \phi_{S}) \right]$$

$$+ 2 \langle \sin(3\phi - \phi_{S}) \rangle_{UT}^{h} \sin(3\phi - \phi_{S}) + 2 \langle \sin(\phi_{S}) \rangle_{UT}^{h} \sin(\phi_{S})$$

$$+ 2 \langle \sin(2\phi - \phi_{S}) \rangle_{UT}^{h} \sin(2\phi - \phi_{S})$$

$$+ \lambda_{l} \left(2 \langle \cos(\phi - \phi_{S}) \rangle_{LT}^{h} \cos(\phi - \phi_{S}) \right)$$

$$+ 2 \langle \cos(\phi_{S}) \rangle_{LT}^{h} \cos(\phi_{S}) + 2 \langle \cos(2\phi - \phi_{S}) \rangle_{LT}^{h} \cos(2\phi - \phi_{S}) \right]$$

 $Q^2 > 1 \text{ GeV}^2$ $W^2 > 10 \text{ GeV}^2$ 0.023 < x < 0.6





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 $\mathcal{C}[h_{1T}^q \times H_1^{\perp,q}]$





data from Belle, Babar, COMPASS, HERMES, Jefferson Lab Hall A

Collins amplitudes

 $\mathcal{C}[h_{1T}^q \times H_1^{\perp,q}]$





Collins amplitudes

 $\mathcal{C}[h_{1T}^q \times H_1^{\perp,q}]$





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 $\mathcal{C}[h_{1T}^q \times H_1^{\perp,q}]$





Pretzelosity amplitudes $C[h_{1T}^{\perp,q} \times H_1^{\perp,q}]$

- Pretzelosity
 - requires non-zero orbital angular momentum (model)
 - models:

$$h_{1T}^{\perp(1),q}(x) = g_{1L}^q(x) - h_{1T}^q(x)$$

- → measure for relativistic effects
- suppressed as $1/P_{h\perp}^2 \, {\rm compared}$ to Collins amplitude



Pretzelosity amplitudes



 $\mathcal{C}[h_{1T}^{\perp,q} \times H_1^{\perp,q}]$

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15

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 - requires non-zero orbital angular momentum (model)

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Pretzelosity amplitudes





Sivers amplitudes

$\mathcal{C}[f_{1T}^{\perp,q} \times D_1^q]$





Sivers amplitudes

 $\mathcal{C}[f_{1T}^{\perp,q} \times D_1^q]$



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Sivers amplitudes

0.023 < x < 0.072

0.072 < x < 0.098

0.098 < x < 0.138

0.138 < x < 0.600





Sivers amplitudes





Largest at large x and z, where purest u-quark probe



Sivers amplitudes





Other hadrons, no clear kinematic dependencies in 3DNo 3D for antiprotons

Twist-3: $\langle \sin(\phi_S) \rangle_{UT}$

$$\begin{split} &\langle \sin(\phi_S) \rangle_{UT} \\ &\propto \mathcal{C}[f_T^q \times D_1^q, h_{1T}^q \times \tilde{H}^q, h_T^q \times H_1^{\perp,q}, g_{1T}^{\perp,q} \times \tilde{G}^{\perp,q}, h_T^{\perp,q} \times H_1^{\perp,q}, f_{1T}^{\perp,q} \times \tilde{D}^{\perp,q}] \\ & \text{twist-3} \end{split}$$

Twist-3:
$$\langle \sin(\phi_S) \rangle_{UT}$$

 $\langle \sin(\phi_S) \rangle_{UT} \\ \propto \mathcal{C}[f_T^q \times D_1^q, h_{1T}^q \times \tilde{H}^q, h_T^q \times H_1^{\perp,q}, g_{1T}^{\perp,q} \times \tilde{G}^{\perp,q}, h_T^{\perp,q} \times H_1^{\perp,q}, f_{1T}^{\perp,q} \times \tilde{D}^{\perp,q}]$ twist-3 integrate over hadron transverse momentum $P_{h\perp}$

$$\langle \sin(\phi_S) \rangle_{UT} = -x \frac{2M_h}{Q} \sum_q e_q^2 h_{1T}^q \frac{H^q}{z}$$

no convolution

Twist-3: $\langle \sin(\phi_S) \rangle_{UT}$



• Significant non-zero signal for π^- , increasing with x, z

Twist-3: $\langle \sin(\phi_S) \rangle_{UT}$



Presented amplitudes

$$\sigma^{h}(\phi,\phi_{S}) = \sigma^{h}_{UU} \left\{ 1 + 2\langle \cos(\phi) \rangle^{h}_{UU} \cos(\phi) + 2\langle \cos(2\phi) \rangle^{h}_{UU} \cos(2\phi) + \lambda_{l} 2\langle \sin(\phi) \rangle^{h}_{LU} \sin(\phi) \right\}$$

- Longitudinally polarized $e^{\scriptscriptstyle +}/e^{\scriptscriptstyle -}$ beam
- \bullet Unpolarized H and D target

→ Results for charged pions, kaons, (anti-)protons



Twist-3: $\langle \sin(\phi) \rangle_{LU}^{h}$

 $\langle \sin(\phi) \rangle_{LU}^h \propto \mathcal{C} \left[h_1^\perp \times \tilde{E}, x \, e \times H_1^\perp, x \, g^\perp \times D_1, f_1 \times \tilde{G}^\perp \right]$

$$\begin{aligned} & \mathsf{Twist-3:} \left\langle \sin(\phi) \right\rangle_{LU}^{h} \\ & \langle \sin(\phi) \rangle_{LU}^{h} \propto \mathcal{C} \left[h_{1}^{\perp} \times \tilde{E}, x \, e \times H_{1}^{\perp}, x \, g^{\perp} \times D_{1}, f_{1} \times \tilde{G}^{\perp} \right] \\ & \mathsf{Boer-Mulders PDF} \end{aligned}$$

Boer-Mulders PDF



$$e(x) = e^{WW}(x) + \bar{e}(x)$$

$$e(x) = e^{WW}(x) + \bar{e}(x)$$
Boer-Mulders PDF
$$e_2 \equiv \int_0^1 dx \, x^2 \bar{e}(x)$$
force on struck quark at t=0
M. Burkardt, arXiv:0810.3589
FSI: t=0 \rightarrow \infty

$$\begin{aligned} & \mathsf{Twist-3:} \left\langle \sin(\phi) \right\rangle_{LU}^h \\ & \langle \sin(\phi) \rangle_{LU}^h \propto \mathcal{C} \left[h_1^\perp \times \tilde{E}, x \, e \times H_1^\perp, x \, g^\perp \times D_1, f_1 \times \tilde{G}^\perp \right] \\ & \checkmark \end{aligned}$$

$$\begin{aligned} & \mathsf{Chiral-odd \ T-even} \\ & \mathsf{twist-3 \ PDF} \end{aligned}$$

$$\begin{aligned} & \mathsf{Collins \ FF} \end{aligned}$$





Twist-3: $\langle \sin(\phi) \rangle_{LU}^{h}$ $\langle \sin(\phi) \rangle_{LU}^h \propto \mathcal{C} \left[h_1^\perp \times \tilde{E}, x \, e \times H_1^\perp, x \, g^\perp \times D_1, f_1 \times \tilde{G}^\perp \right]$ spin-independent chiral-even, T-odd PDF twist-3 FF

Twist-3: $\langle \sin(\phi) \rangle_{LU}^h$



- Agreement H and D data
- Positive results for pions



- No clear kinematic dependencies en 3D
- No 3D for anti-protons

 \rangle_{LU}^{h} Twist-3: $\langle \sin(\phi) \rangle$



• Both measurements give compatible results



- Opposite behaviour for $\pi^{\scriptscriptstyle -}$ z projection due to different x range probed

• CLAS probes higher x region: more sensitive to $e \times H_1^{\perp}$? $\langle \sin(\phi) \rangle_{LU}^h \propto \mathcal{C} \left[h_1^{\perp} \times \tilde{E}, x \, e \times H_1^{\perp}, x \, g^{\perp} \times D_1, f_1 \times \tilde{G}^{\perp} \right]$

Summary





Summary




Summary





Summary









