

Transversity in inclusive DIS

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Transversity 2017

Laboratori Nazionali di Frascati,

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Based on: *Accardi, Bacchetta, PLB 773 (2017) 632*
Accardi, Signori, work in progress

Overview

- ❑ **Inclusive DIS with jet correlators**
 - Quarks are not asymptotic states
 - Cross section & g_2 revisited
- ❑ **Twist-3 TMD sum rules**
 - New results, old ones revisited
- ❑ **Measuring the jet correlator**
- ❑ **Concluding thoughts**

Inclusive DIS with jet correlators

Accardi, Bacchetta, PLB 773 (2017) 632

TMDs in spin $\frac{1}{2}$ targets

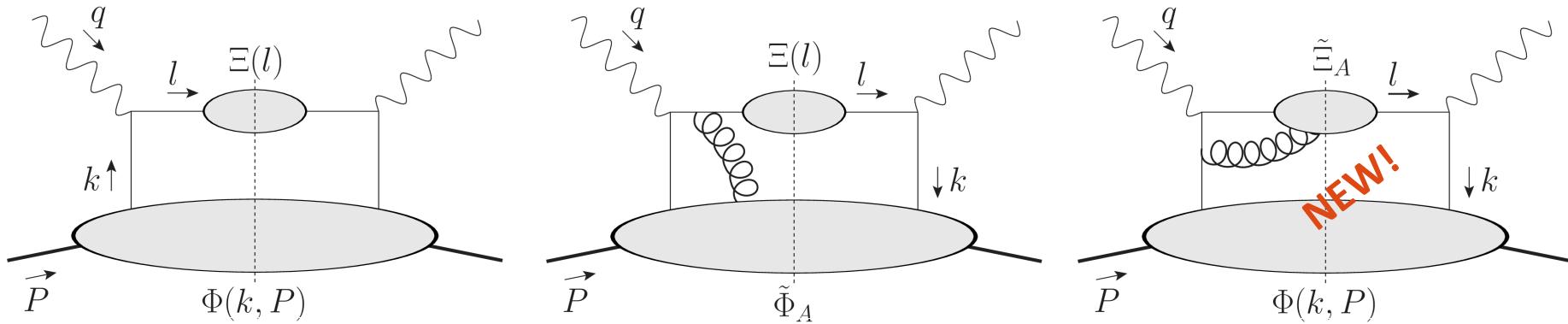
		PARTON SPIN		
QUARKS		γ^+	$\gamma^+ \gamma_5$	$\gamma^+ \gamma^\alpha \gamma_5$
TARGET SPIN	U	f_1		h_1^\perp
	L		g_1	h_{1L}^\perp
	T	f_{1T}^\perp	g_{1T}	$h_1^\perp h_{1T}^\perp$

→ P. Mulders, QCDev2017

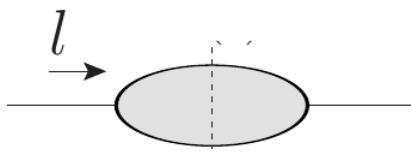
- Integrated (collinear) correlator: only circled ones survive
- Christ-Lee theorem (1970): h_1 not observable in inclusive DIS
- Not quite true:
 - Vacuum fluctuations can flip the spin of the struck quark
 - Large contribution $\sim h_1$ pops up in the $g_2 - g_2^{WW}$ structure function

Inclusive DIS with jet correlators

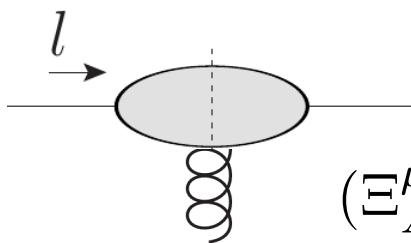
AA, Bacchetta, PLB 773 ('17) 632



Jet correlators: → non-asymptotic quark states



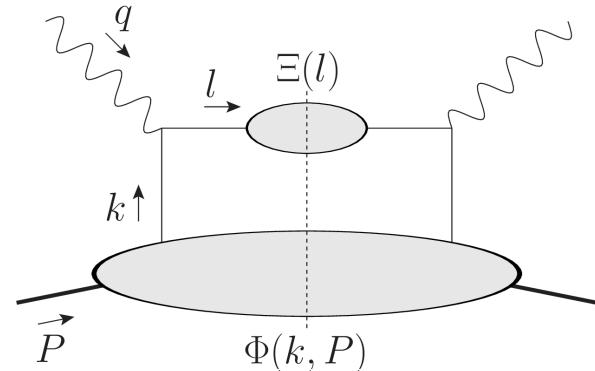
$$\Xi_{ij}(l, n_+) = F.T. \langle 0 | \mathcal{U}_{(+\infty, \eta)}^{n_+} \psi_i(\eta) \bar{\psi}_j(0) \mathcal{U}_{(0, +\infty)}^{n_+} | 0 \rangle$$



$$(\Xi_A^\mu)_{ij} = F.T. \langle 0 | \mathcal{U}_{(+\infty, \eta)}^{n_+} g A^\mu(\eta) \psi_i(\eta) \bar{\psi}_j(0) \mathcal{U}_{(0, +\infty)}^{n_+} | 0 \rangle$$

Factorization

- At order $1/Q$, neglect k^- compared to q^-
 - The cross section only depends on the **integrated jet correlator**



$$\Xi(l^-, l_T) \equiv \int \frac{dl^2}{2l^-} \Xi(l) = \frac{\Lambda}{2l^-} \xi_1 + \xi_2 \frac{\not{k}}{2} + \text{twist-4 terms}$$

- Coefficients can be interpreted in terms of quark spectral functions:

$$\xi_1 = \int d\mu^2 \frac{\mu}{\Lambda} J_1(\mu^2) \equiv \boxed{\frac{M_q}{\Lambda}} \quad \begin{matrix} \leftarrow \\ \text{Spin-flip average "jet" mass} \\ \rightarrow \text{can couple to transversity!} \end{matrix}$$

$$\xi_2 = \int d\mu^2 J_2(\mu^2) = 1 \quad \leftarrow \quad \text{Exactly}$$

- Positivity constraints imply

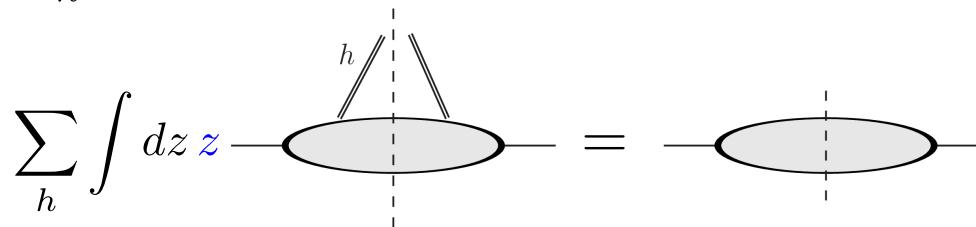
$$0 < M_q < \int d\mu^2 \mu J_2(\mu^2) \implies \boxed{M_q = O(100 \text{ MeV})}$$

Much larger than m_q !

Jet and TMD sum rules

- Utilize the following jet correlator sum rule:

$$\sum_h \int d^2 p_{hT} \frac{d\bar{p}_h}{2\bar{p}_h} \textcolor{blue}{p_h^-} \Delta^h(l, p_h) = \textcolor{blue}{l^-} \Xi(l)$$



- At TMD level, this implies:

$$\sum_h \int dz d^2 p_{hT} z D_1^h(z, p_{hT}) = \xi_2 = 1$$

$$\sum_h \int dz d^2 p_{hT} E^h(z, p_{hT}) = \xi_1 = \boxed{\frac{M_q}{\Lambda}}$$

$$\sum_h \int dz d^2 p_{hT} \tilde{E}^{q,h}(z, p_{hT}) = \boxed{\frac{M_q - m_q}{\Lambda}}$$

Novel TMD sum rules

$M_q^{\text{pert}} = m_q \Rightarrow$ Old ones

(see later for more...)

Finally, the DIS cross section

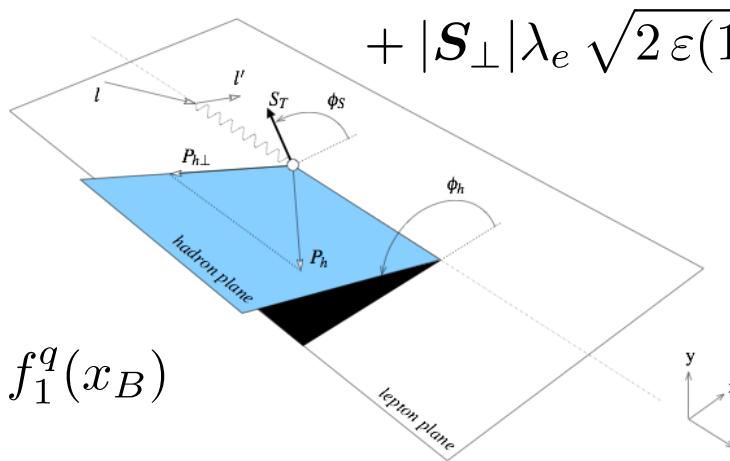
- ❑ Inclusive DIS

$$\frac{d\sigma}{dx_B dy d\phi_S} \propto \left\{ F_T + \varepsilon F_L + S_{||} \lambda_e \sqrt{1 - \varepsilon^2} F_{LL} + |S_{\perp}| \lambda_e \sqrt{2\varepsilon(1 - \varepsilon)} \cos \phi_S F_{LT}^{\cos \phi_S} \right\}$$

- ❑ Integrating the SIDIS structure functions:

$$F_T = x_B \sum_q e_q^2 f_1^q(x_B)$$

$$F_L = 0$$



$$F_{LL} = x_B \sum_q e_q^2 g_1^q(x_B)$$

$$F_{LT}^{\cos \phi_S} = -x_B \sum_q e_q^2 \frac{2M}{Q} \left(x_B g_T^q(x_B) + \boxed{\frac{M_q - m_q}{M} h_1^q(x_B)} \right)$$

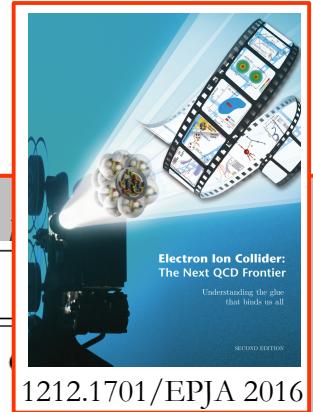
Transversity in inclusive DIS!

Finally, the DIS cross section

Inclusive DIS

$$\frac{d\sigma}{dQ^2} \propto \left\{ F_T + \varepsilon F_L + S_{\parallel} \lambda_e \sqrt{1 - \varepsilon^2} F_L \right\}$$

Deliverables	Observables	What we learn
Sivers & unpolarized TMD quarks and gluon	SIDIS with Transverse polarization; di-hadron (di-jet)	Quantum Interference & Spin-Orbital 3D Imaging of quark's motion: valence + sea 3D Imaging of gluon's motion QCD dynamics in a unprecedented Q^2 (P_{hT}) range
Chiral-odd functions: Transversity; Boer-Mulders	SIDIS with Transverse polarization	3 rd basic quark PDF: valence + sea, tensor charge Novel spin-dependent hadronization effect QCD dynamics in a chiral-odd sector with a wide Q^2 (P_{hT}) coverage



1212.1701/EPJA 2016

Table 2.2: Science Matrix for TMD: 3D structure in transverse momentum space: (upper) golden measurements; (lower) the silver measurements.

$$F_{LT}^{\cos \phi_S} = -x_B \sum_q e_q^2 \frac{2M}{Q} \left(x_B g_T^q(x_B) + \boxed{\frac{M_q - m_q}{M} h_1^q(x_B)} \right)$$

Transversity in inclusive DIS!



g2 structure function revisited

AA, Bacchetta, PLB 773 ('17) 632

- Using EOM, Lorentz Invariance Relations, can show that

$$g_2(x_B) - g_2^{WW}(x_B) \equiv g_2^{\text{quark}} = \frac{1}{2} \sum_a e_a^2 \left(g_2^{q,\text{tw3}}(x_B) + \frac{m_q}{M} \left(\frac{h_1^q}{x} \right)^* (x_B) + \boxed{\frac{M_q - m_q}{M} \frac{h_1^q(x_B)}{x_B}} \right)$$

Color force distribution Transversity in inclusive DIS!

- Consequences:

- h1 accessible in inclusive DIS! \leftrightarrow Potentially large signal
- new background to extraction of qGq effects
- Measuring the tensor charge

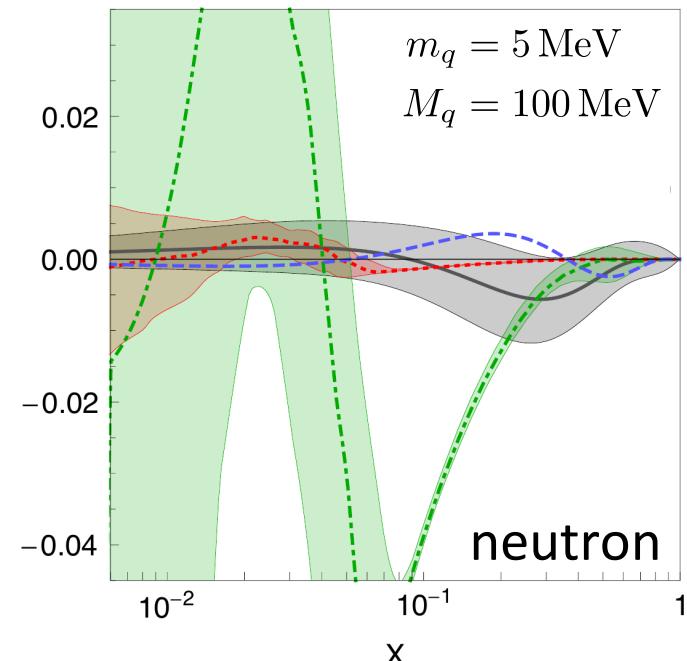
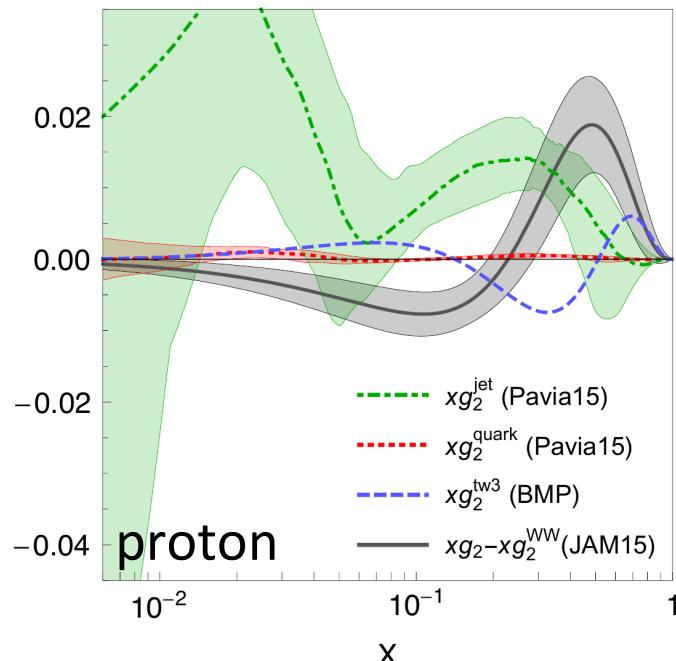
$$f^*(x) = -f(x) + \int_x^1 \frac{dy}{y} f(y)$$

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Novel non-perturbative sum rules

AA, Bacchetta, PLB 773 ('17) 632

- Taking moments of g_2 with $M_u \approx M_d \equiv M_{jet}$

Burkardt-Cottingham

$$\int_0^1 g_2(x) = M_{\text{"jet"}} \int_0^1 dx \frac{h_1(x)}{x}$$

- unlikely to still be zero!
- Is BC breaking finite?
Perturbatively, yes ($h \sim x$) but...

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→ Small- x asymptotics:

$$g_1^{NS} \sim x^{\epsilon_g} \quad \epsilon_g = -\sqrt{\alpha_s N_c / \pi} \approx -0.6$$

→ Kovchegov, Pitonyak, Sievert
PRD(2017)93

But h_1 is also non-singlet, expect

$$h_1 \sim x^{\epsilon_h} \quad \epsilon_h = \epsilon_g < 0!!$$

- Is BC badly broken? $1/N_c$ corrections non negligible? Or ...?

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How does spin propagate to small x ?

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Novel non-perturbative sum rules

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$$\int_0^1 g_2(x) = M^{“jet”} \int_0^1 dx \frac{h_1(x)}{x}$$

Efremov-Teryaev-Leader

$$\int_0^1 x g_2^{q-\bar{q}}(x) = 2 M^{“jet”} \underbrace{\int_0^1 dx h_1^{q-\bar{q}}(x)}_{\text{Tensor charge } \delta_T}$$

→ Novel way to measure the tensor charge!

Novel non-perturbative sum rules

AA, Bacchetta, PLB 773 ('17) 632

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Color polarizability

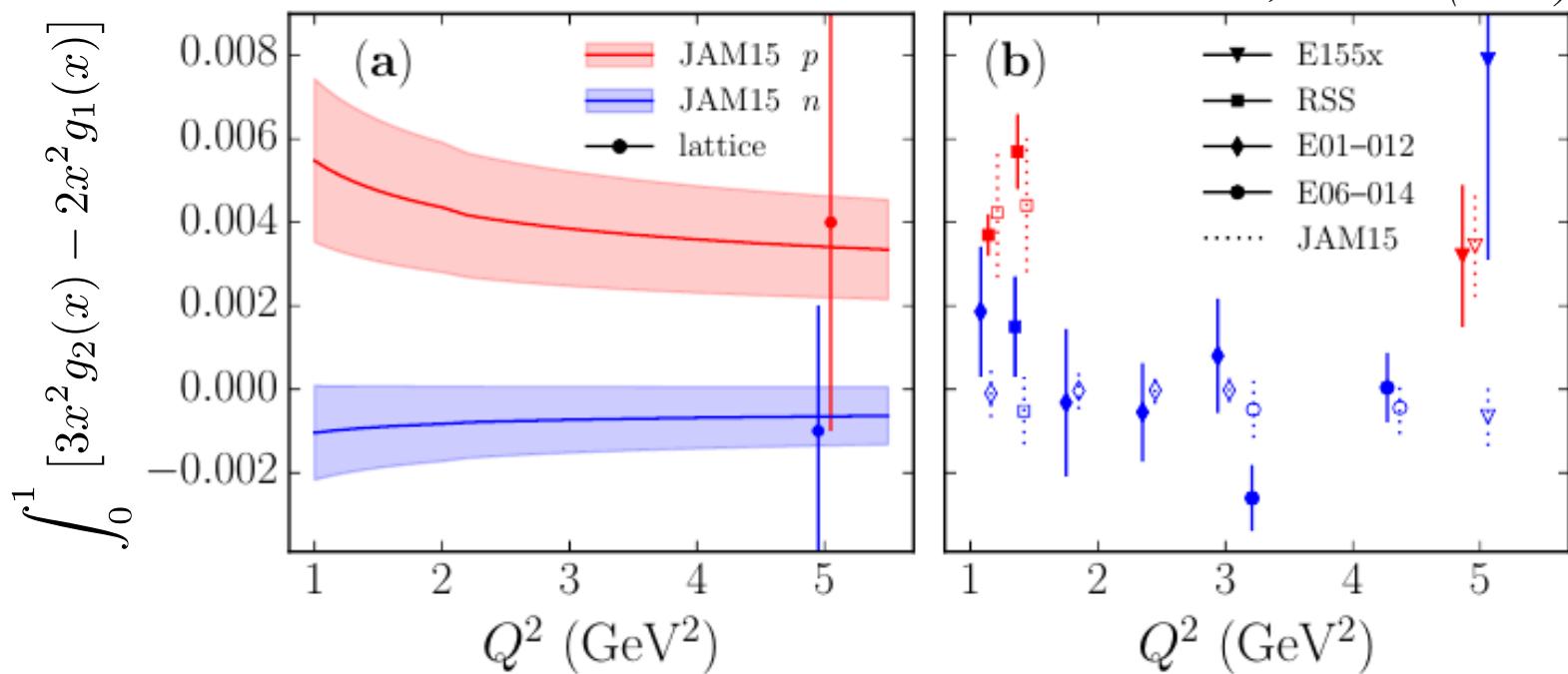
$$\int_0^1 [3x^2 g_2(x) - 2x^2 g_1(x)] = d_2 + 3 M^{“jet”} \underbrace{\int_0^1 dx x h_1(x)}_{\text{“background”}} + O(m_q)$$

Color force
– Burkardt $\sim \langle P | \bar{\psi} \gamma^+ F^{+\alpha} \psi | P \rangle$

“pure twist-3”

Higher twist and color polarizability

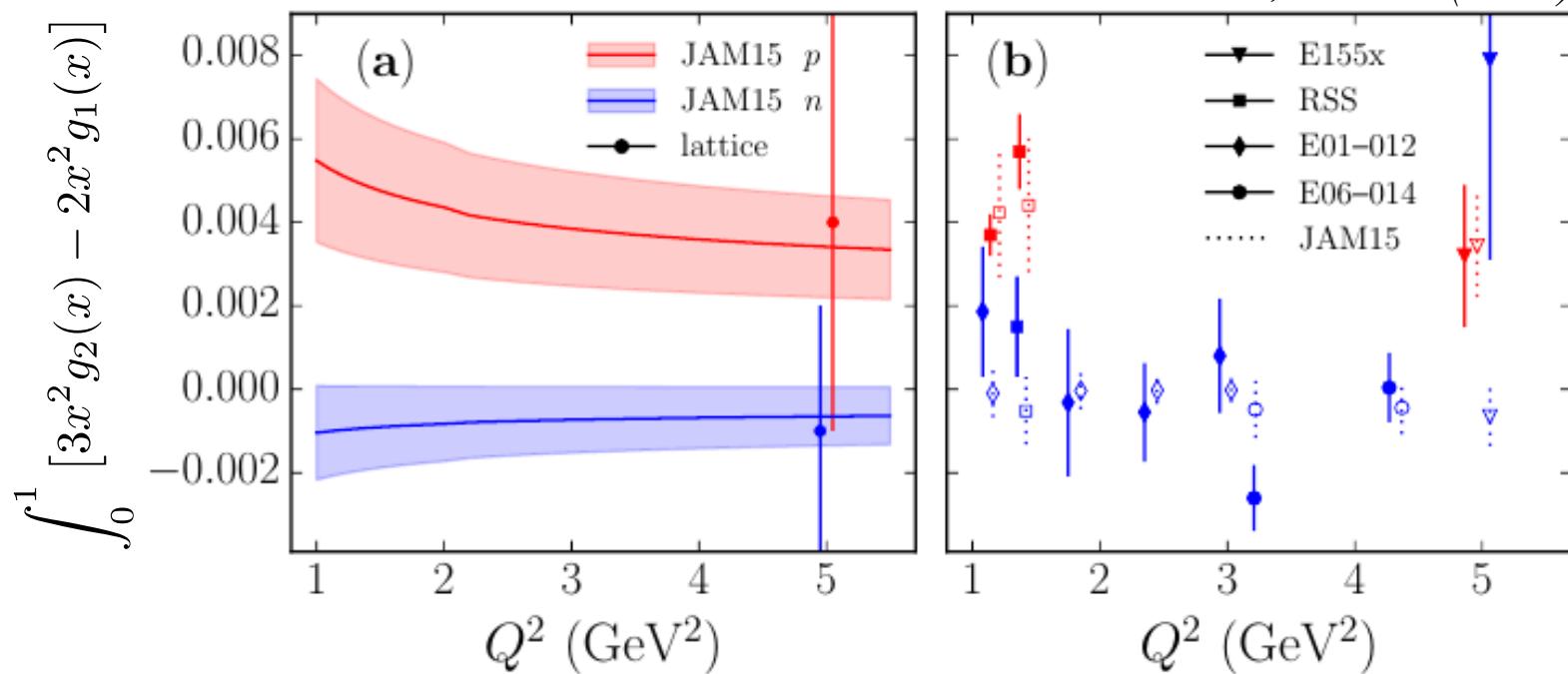
N.Sato *et al.*, PRD93 (2015) 074005



Lattice calcs from Goeckele et al. 2005 – time to revisit?

Higher twist and color polarizability

N.Sato *et al.*, PRD93 (2015) 074005



Lattice calcs from Goeckele et al. 2005 – time to revisit?

Can one also calculate Mjet on the lattice?

New TMD sum rules

Accardi, Signori, in preparation

Quark-quark TMD sum rules

- General jet correlator sum rule:

*AA, Bachetta '17
Meissner, Metz, Pitonyak '10*

$$\sum_h \int d^2 p_{hT} \frac{dp_h^-}{2p_h^-} p_h^\mu \Delta^h(l, p_h) = l^\mu \Xi(l) \quad \mu = \begin{cases} - & \text{longitudinal} \\ \alpha & \text{transverse} \end{cases}$$

- At TMD level, take suitable traces and find:

	Longitudinal	Transverse
Twist-2	$\sum_h \int dz z D_{1h}(z) = 1$ <i>Collins-Soper</i>	$\sum_h \int dz z H_{1h}^{\perp(1)}(z) = 0$ <i>Schaefer-Teryaev</i>
Twist-3	$\left\{ \begin{array}{l} \sum_h \int dz E_h(z, p_{hT}) = \frac{M_q}{\Lambda} \\ \sum_h \int dz H_h(z) = 0 \end{array} \right.$ NEW!	$\sum_h \int dz G_h^{\perp(1)}(z) = 0$ NEW! $\sum_h \int dz D_h^{\perp(1)}(z) =$ IN PROGRESS

Quark-gluon-quark TMD sum rules

- Combine q-q sum rules using Equation of Motion relations:

$$\sum_h \int dz \tilde{E}_h(z) = \frac{M_q - m_q}{\Lambda} \text{NEW!} \implies \text{Transversity in DIS!}$$

$$\sum_h \int dz \tilde{H}_h(z) = 0 \text{NEW!} \implies \int dz z F_{UT}^{\sin \phi_S}(x, z) = 0 \quad \textit{Diehl-Sapeta}$$

$$\sum_h \int dz \tilde{D}^\perp(z) = \text{IN PROGRESS}$$

NEW: proven at correlator level

Measuring the jet correlator

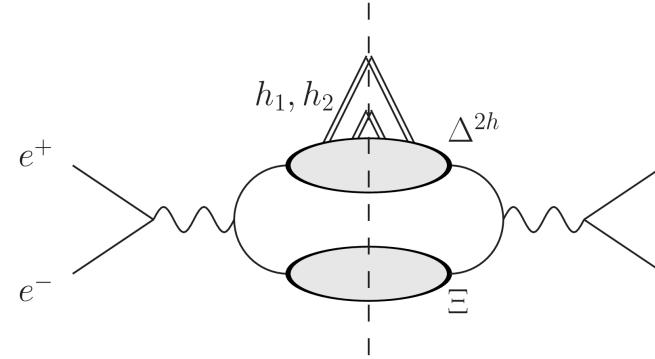
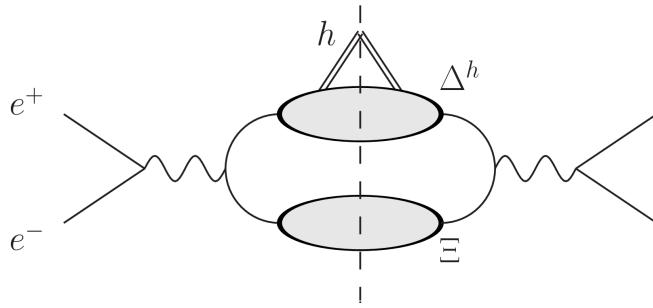
In collaboration with
A.Bacchetta, M.Radici, A.Signori

Measuring the jet correlator

Accardi, Bacchetta, Radici, *in prep.*

Related to confinement, mass generation [C.Roberts]

- ☐ Jet mass M_{jet} can be measured in polarized $e^+ + e^-$:



- Needs **LT asymmetry** in semi-inclusive **Lambda** production

$$\frac{d\sigma^L(e^+e^- \rightarrow \Lambda X)}{d\Omega dz} = \frac{3\alpha^2}{Q^2} \lambda_e \sum_a e_a^2 \left\{ \frac{C(y)}{2} \lambda_h G_1 + D(y) |\mathbf{S}_T| \cos(\phi_S) \frac{2M_h}{Q} \left(\frac{G_T}{z} + \frac{M_q - m_q}{M_h} H_1 \right) \right\}$$

- Similarly a **LU asymmetry** in unpolarized **dihadron** production

Need polarized e+e- colliders!

- Are existing facilities enough?

adapted from Particle Data Book

	BEPC	super KEKB	ILC	JLab/BNL
E beam [GeV]	1.9	4 (e ⁻) 7 (e ⁻)	250	?
\sqrt{s} [GeV]	3 – 5	10	500	?
polarization	?	maybe	80% e ⁻ 60% e ⁺	YES!

- Can we get a (polarized) e+ e- collider at JLab / BNL?
 - At JLab12 ? JLEIC ? eRHIC?
- What else is interesting to study?
 - Factorization tests for FFs (low s, unpol), ... Ideas?

Final thoughts

Final thoughts

□ Jet correlators open up new theory and phenomenology

- Transversity contributes to inclusive g_2
- Extended BC and ETL sum rules
 - New handle on proton tensor charge
- Open question: spin transport to small x !
- New twist-3 TMD sum rules

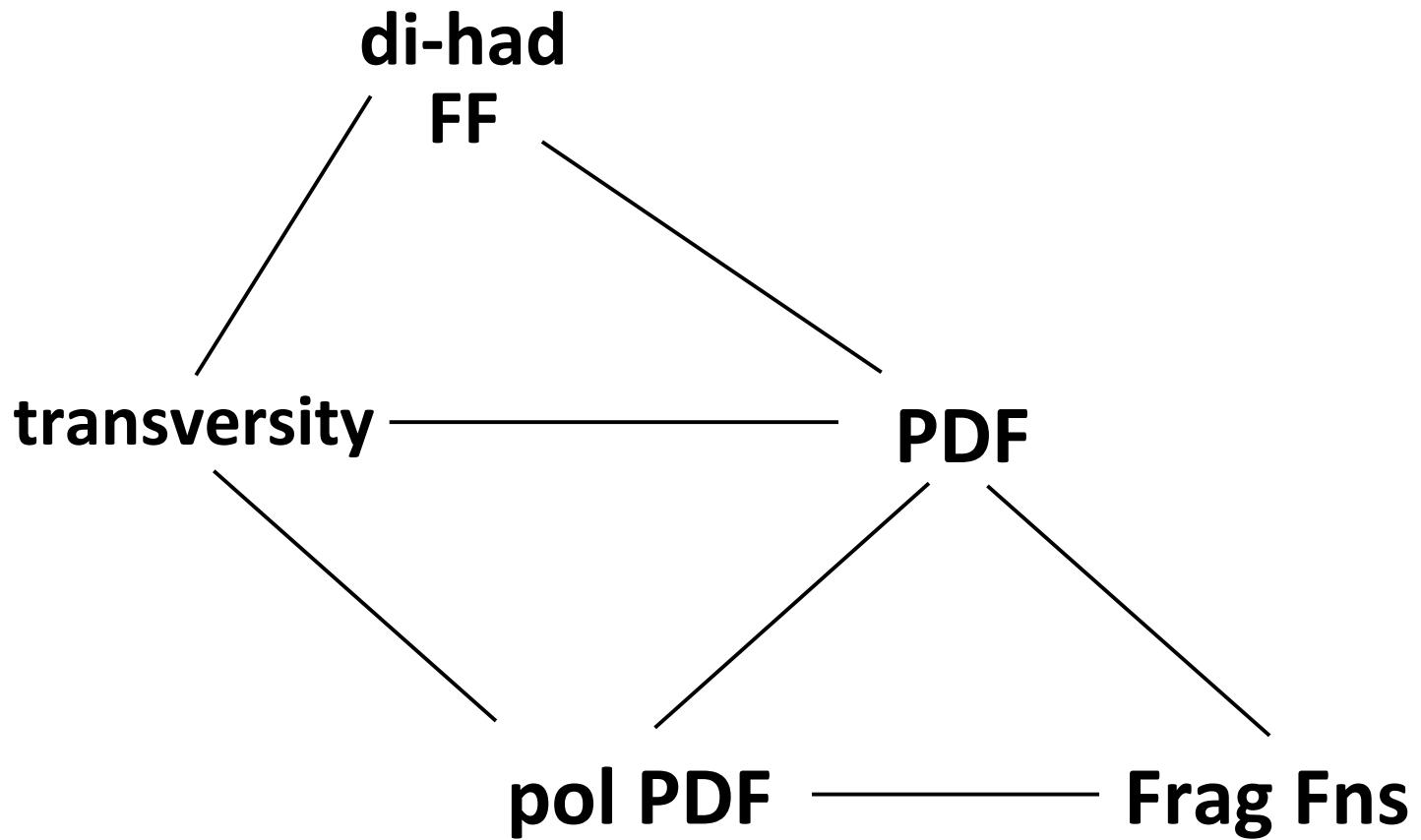
□ How to measure jet correlators?

- Polarized e+e- collider
- Observables in p+p ??

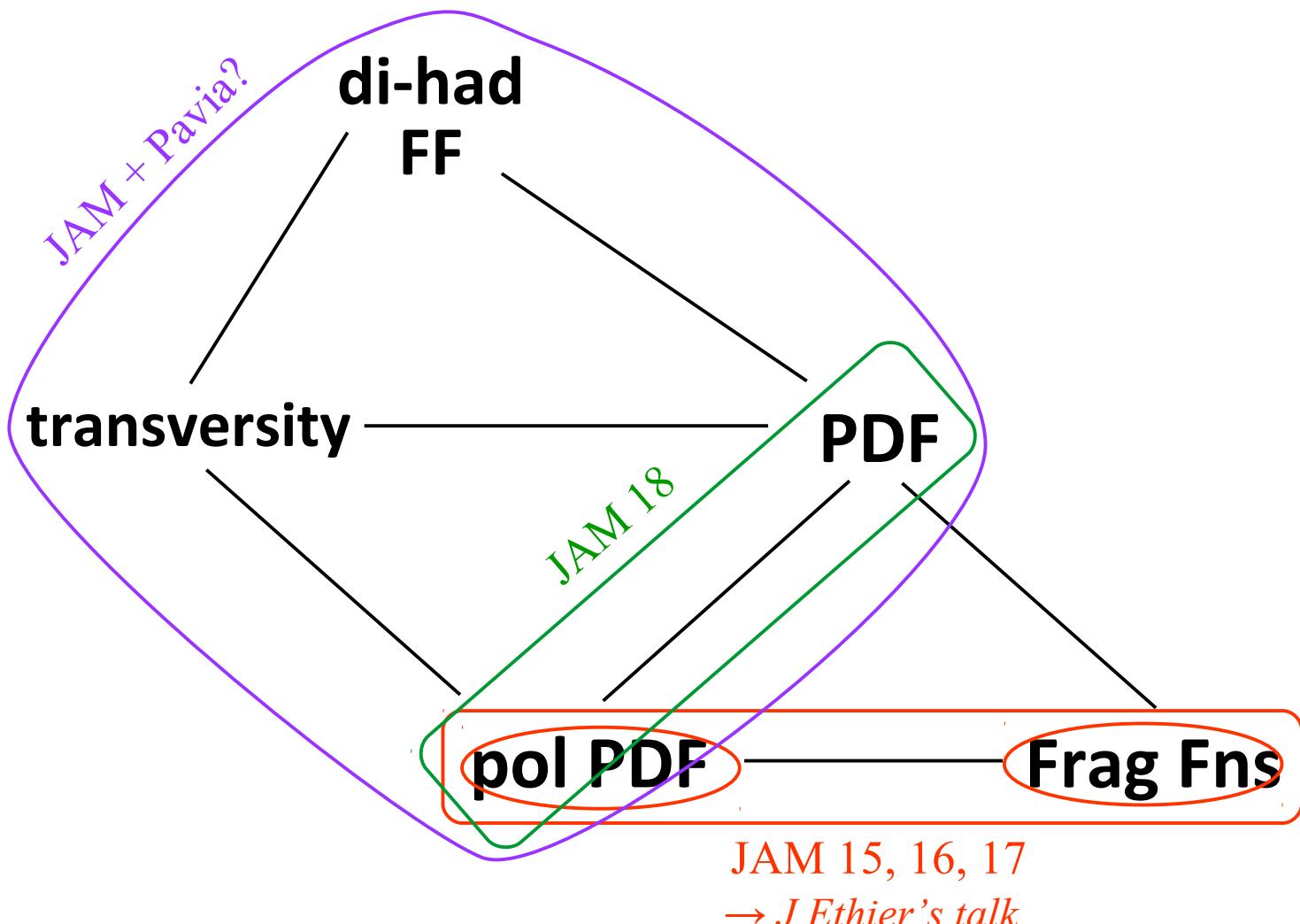
□ How to calculate jet correlators?

- In lattice QCD?
- Any relation to chiral condensate?
- What can we learn from OPE?

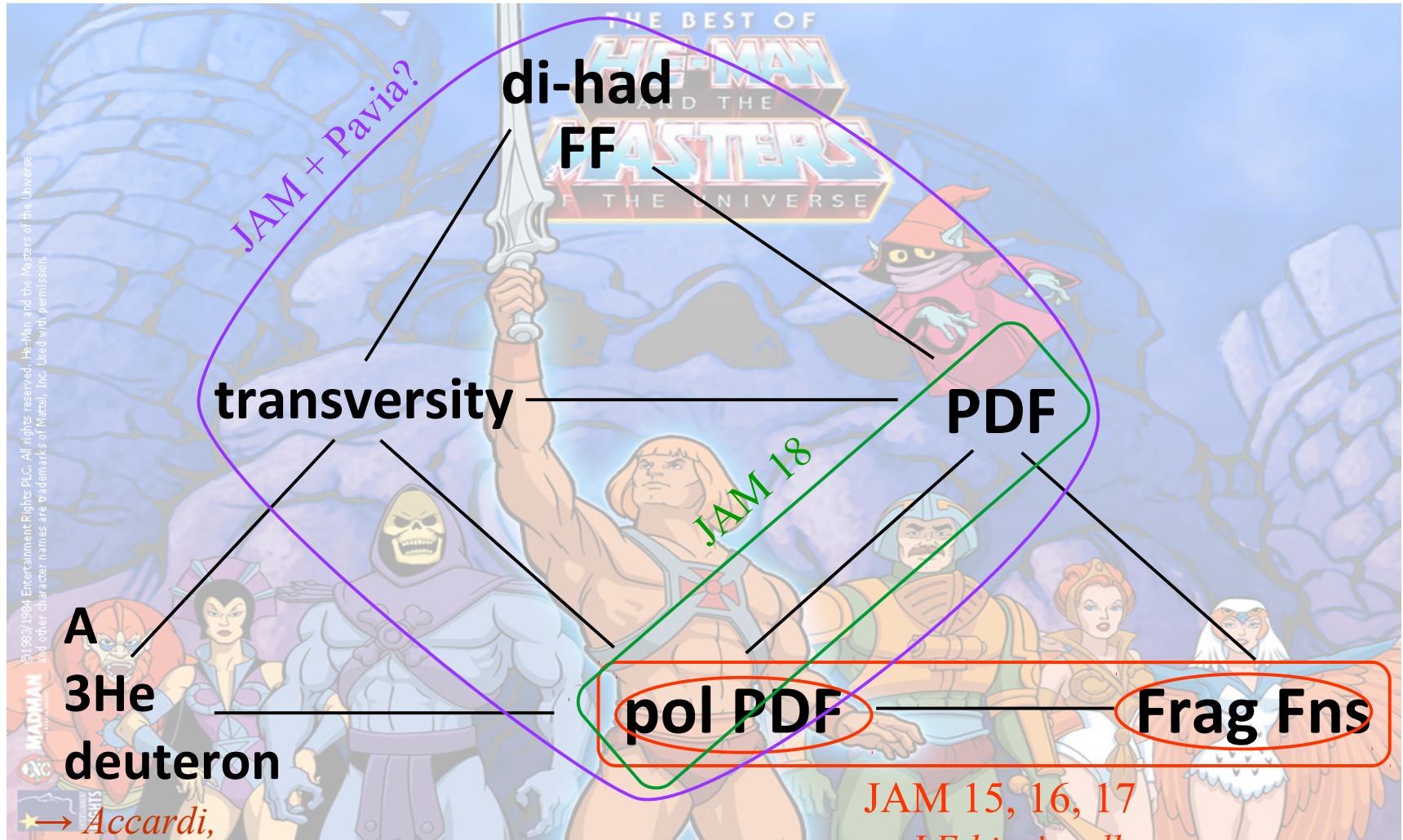
Towards “universal” fits



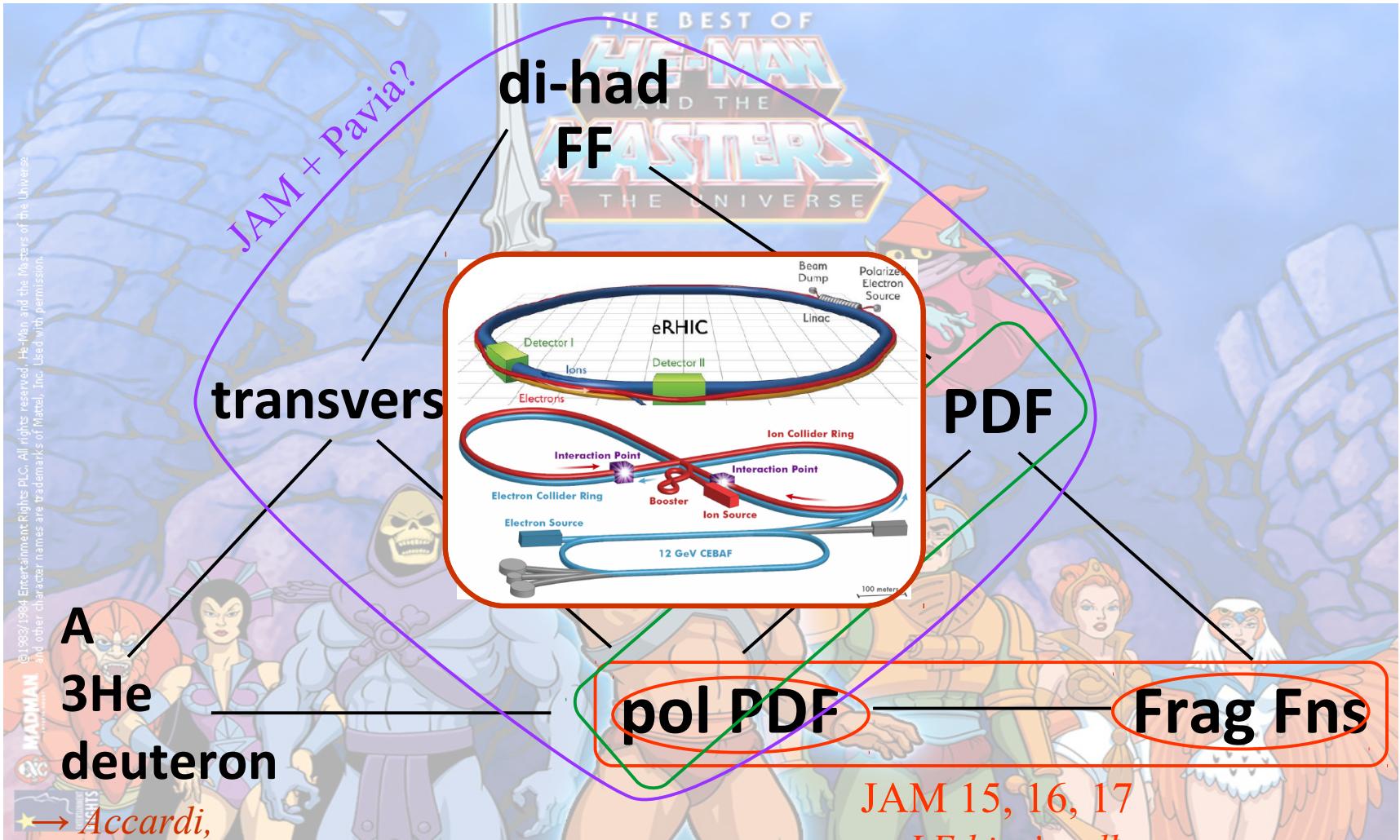
Towards “universal” fits



Masters of the Universe



Masters of the Universe



Thank you!