Azimuthal distributions in the Drell-Yan process

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Outline:

- Drell-Yan angular coefficients
- Fixed-order pQCD & phenomenology
- Resummation & phenomenology

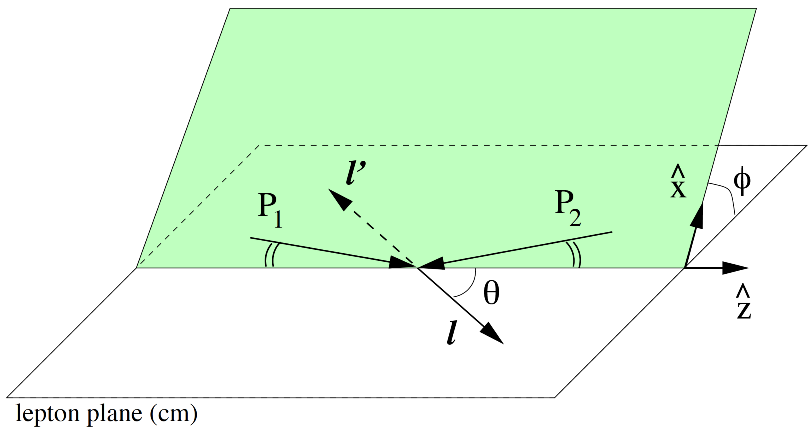
Will focus on “collinear pQCD perspective”
Drell-Yan extremely well explored

Collab. / discussions with M. Lambertsen, J. Steiglechner,
A. Bacchetta, G. Bozzi, F. Piacenza

Drell-Yan angular coefficients

Lepton angular distribution in Drell-Yan (photon exch.):

$$\frac{d\sigma}{d^4q\,d\Omega} = \frac{\alpha^2}{2\pi N_c Q^2 s^2}\Big(W_T\,(1+\cos^2\theta) + W_L\,(1-\cos^2\theta)$$
$$+W_\Delta \sin 2\theta \cos\phi + W_{\Delta\Delta}\sin^2\theta \cos 2\phi\Big)$$



(Collins-Soper frame)

$$\frac{d\sigma}{d^4q\, d\Omega} = \frac{\alpha^2}{2\pi N_c Q^2 s^2}\Big(W_T\,(1+\cos^2\theta) + W_L\,(1-\cos^2\theta)$$
$$+ W_\Delta\,\sin 2\theta\cos\phi + W_{\Delta\Delta}\sin^2\theta\cos 2\phi\Big)$$

$$= \frac{3\,\sigma_0}{4\pi}\,\frac{1}{\lambda+3}\left[1+\lambda\,\cos^2\theta + \mu\,\sin 2\theta\cos\phi + \frac{\nu}{2}\,\sin^2\theta\cos 2\phi\right]$$

$$= \frac{3\,\sigma_0}{16\pi}\left[1+\cos^2\theta + \frac{A_0}{2}\,(1-3\cos^2\theta) + A_1\,\sin 2\theta\cos\phi + \frac{A_2}{2}\,\sin^2\theta\cos 2\phi\right]$$

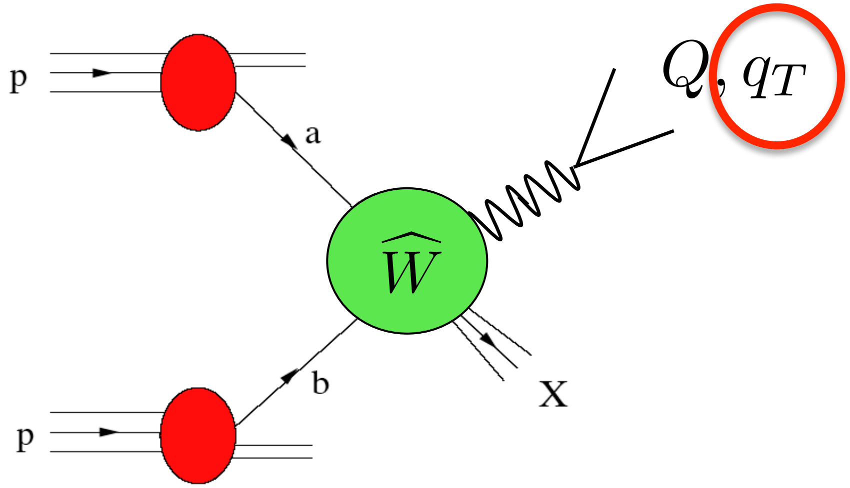
where:

$$\lambda = \frac{W_T - W_L}{W_T + W_L}\,, \quad \mu = \frac{W_\Delta}{W_T + W_L}\,, \quad \nu = \frac{2W_{\Delta\Delta}}{W_T + W_L}$$

$$A_0 = \frac{2W_L}{2W_T + W_L}\,, \quad A_1 = \frac{2W_\Delta}{2W_T + W_L}\,, \quad A_2 = \frac{4W_{\Delta\Delta}}{2W_T + W_L}$$

Fixed-order pQCD

Collinear factorization:



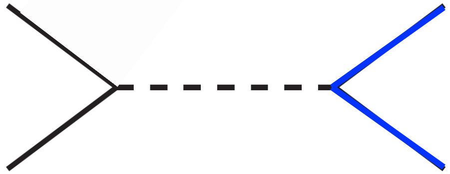
$$W_P = \sum_{a,b} \int dx_a dx_b \, f_a(x_a,\mu) \, f_b(x_b,\mu) \, \widehat{W}_{P,ab}(x_a P_a, x_b P_b, q, \alpha_s(\mu), \mu)$$

$$(P = T, L, \Delta, \Delta\Delta)$$

- $\widehat{W}_P$ partonic structure fcts.: perturbative

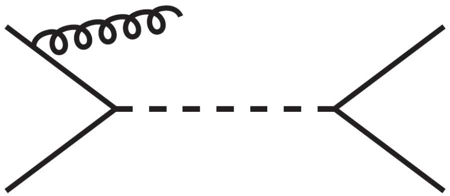
$$\widehat{W}_P = \frac{\alpha_s}{\pi} \widehat{W}_P^{(1)} + \left(\frac{\alpha_s}{\pi}\right)^2 \widehat{W}_P^{(2)} + \ldots \qquad \text{(fixed-order)}$$

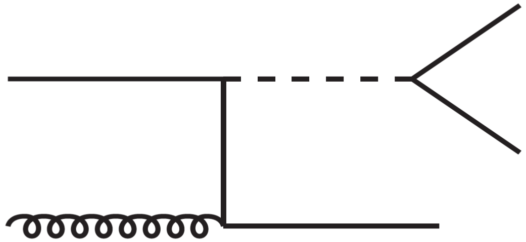
- zeroth order $q\bar{q} \to V \to \ell^+\ell^-$:



$q_T = 0$

- $q_T \neq 0$: first non-trivial order (= LO) $\mathcal{O}(\alpha_s)$



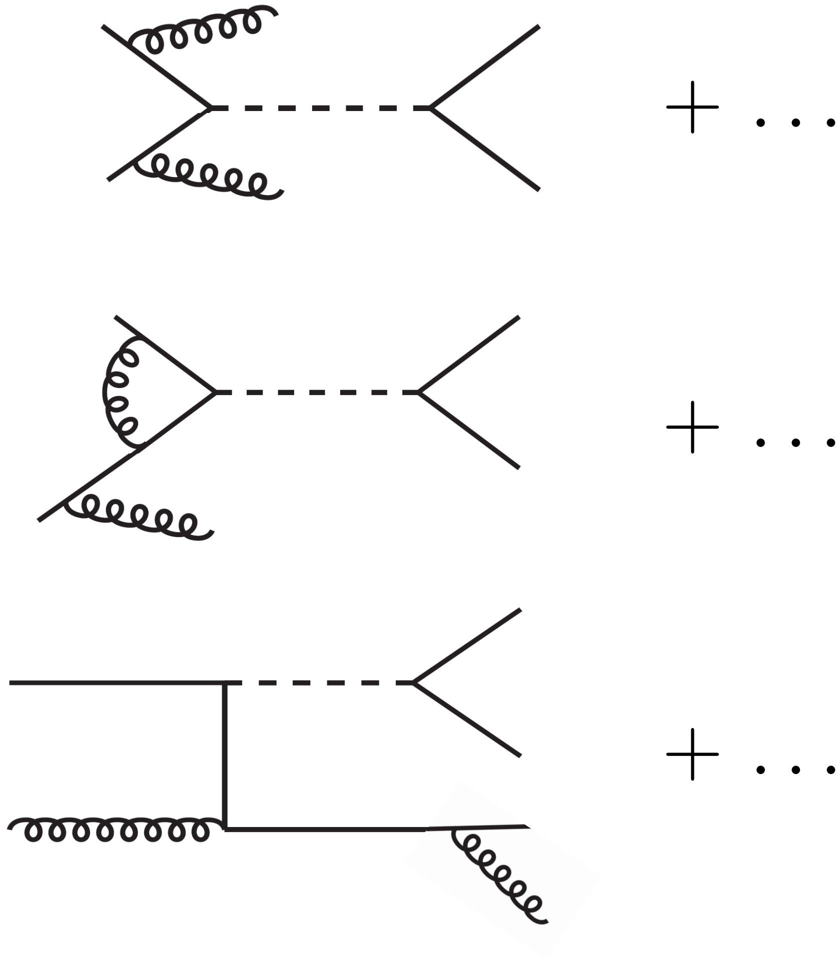


$$\lambda \neq 1, \ \mu \neq 0, \ \nu \neq 0$$

- but: $1 - \lambda - 2\nu \ = \ 0$ (Lam-Tung relation)

$$A_0 \ = \ A_2$$

NLO:



$\mathcal{O}(\alpha_s^2)$

first computed by Mirkes '92; Mirkes, Ohnemus '95

- a lot of work in recent 2 decades on $\mathcal{O}(\alpha_s^2)$ corrections for Drell-Yan process

  Hamberg, van Neerven, Matsuura; Harlander, Kilgore;
  Anastasiou, Dixon, Melnikov, Petriello; Melnikov, Petriello;
  Li, Petriello, Quackenbusch; Catani, Cieri, Ferrera, de Florian, Grazzini;
  Karlberg, Re, Zanderighi; …

- especially: $\mathcal{O}(\alpha_s^2)$ Monte-Carlo codes

  FEWZ: Melnikov, Petriello; Melnikov, Petriello;
  Li, Petriello, Quackenbusch

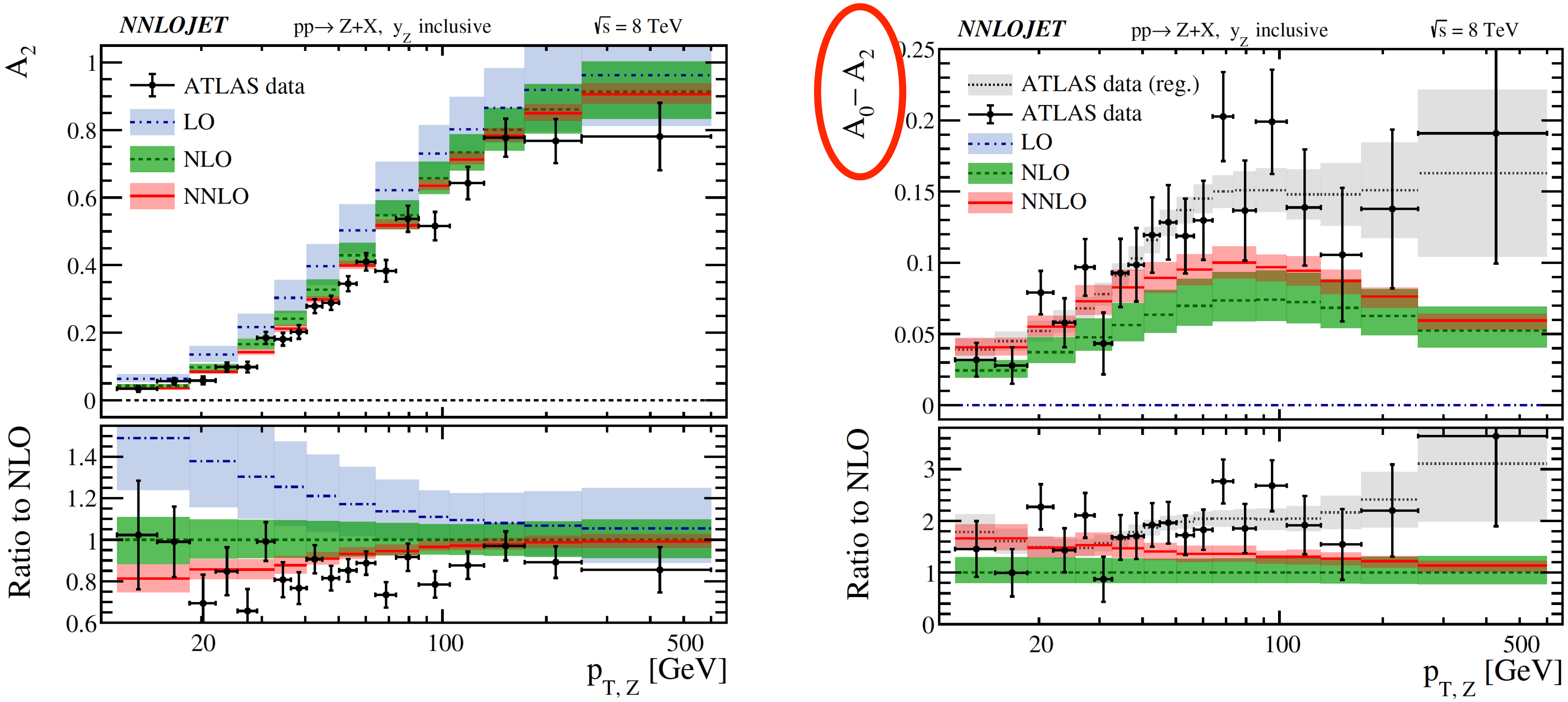
  DYNNLO: Catani, Cieri, Ferrera, de Florian, Grazzini

- most recently: $\mathcal{O}(\alpha_s^3)$

  Li, von Manteuffel, Schabinger, Zhu;
  Anastasiou, Duhr, Dulat, Herzog, Mistlberger;
  Gauld, Gehrmann-De Ridder, Gehrmann, Glover, Huss, Morgan

Fixed-order phenomenology

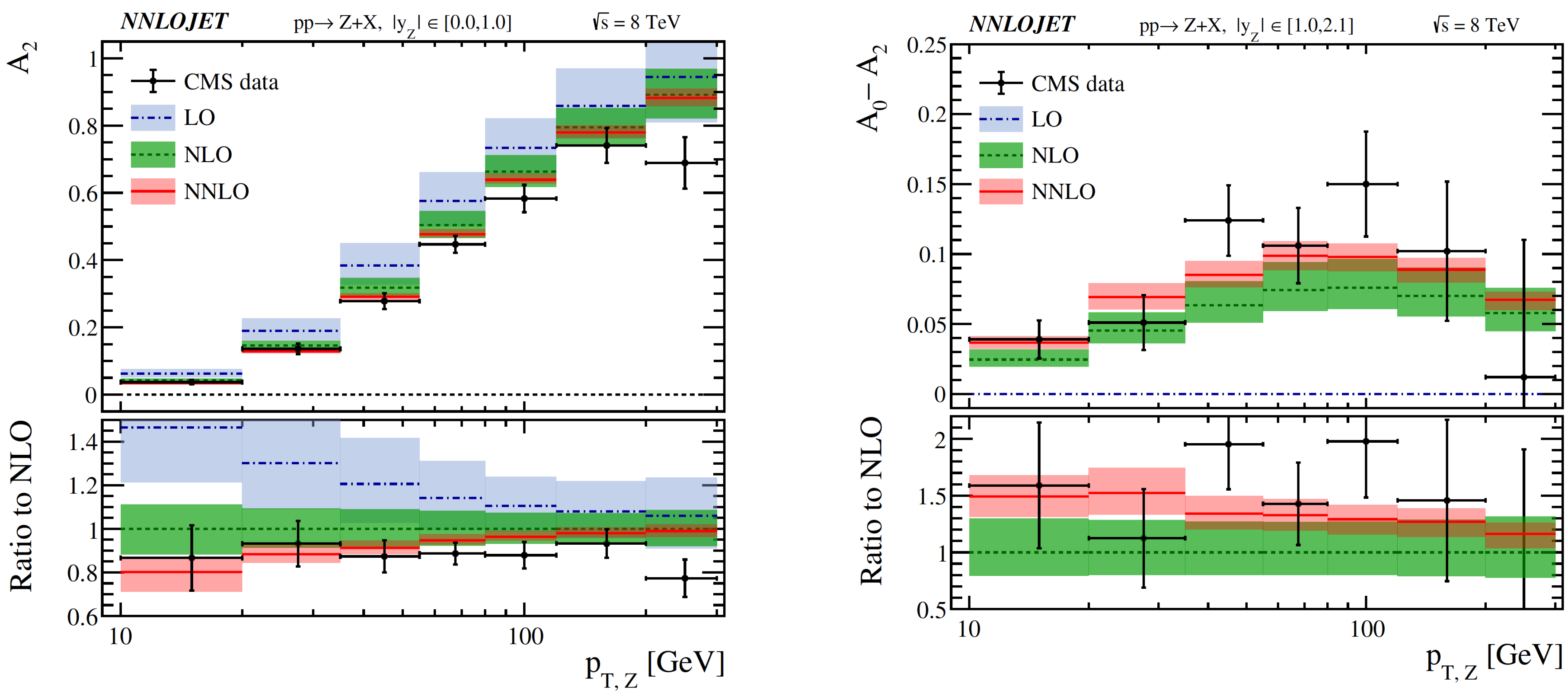
Gauld, Gehrmann-De Ridder, Gehrmann, Glover, Huss ‘17



$\mathcal{O}(\alpha_s^2)$ NLO (ATLAS): $\chi^2/\mathrm{N}_{\mathrm{data}} = 185.8/38 = 4.89$

$\mathcal{O}(\alpha_s^3)$ NNLO (ATLAS): $\chi^2/\mathrm{N}_{\mathrm{data}} = 68.3/38 = 1.80$

Gauld, Gehrmann-De Ridder, Gehrmann, Glover, Huss '17



$$\text{NLO (CMS):} \qquad \chi^2/\mathrm{N}_{\mathrm{data}} = 24.5/14 = 1.75$$

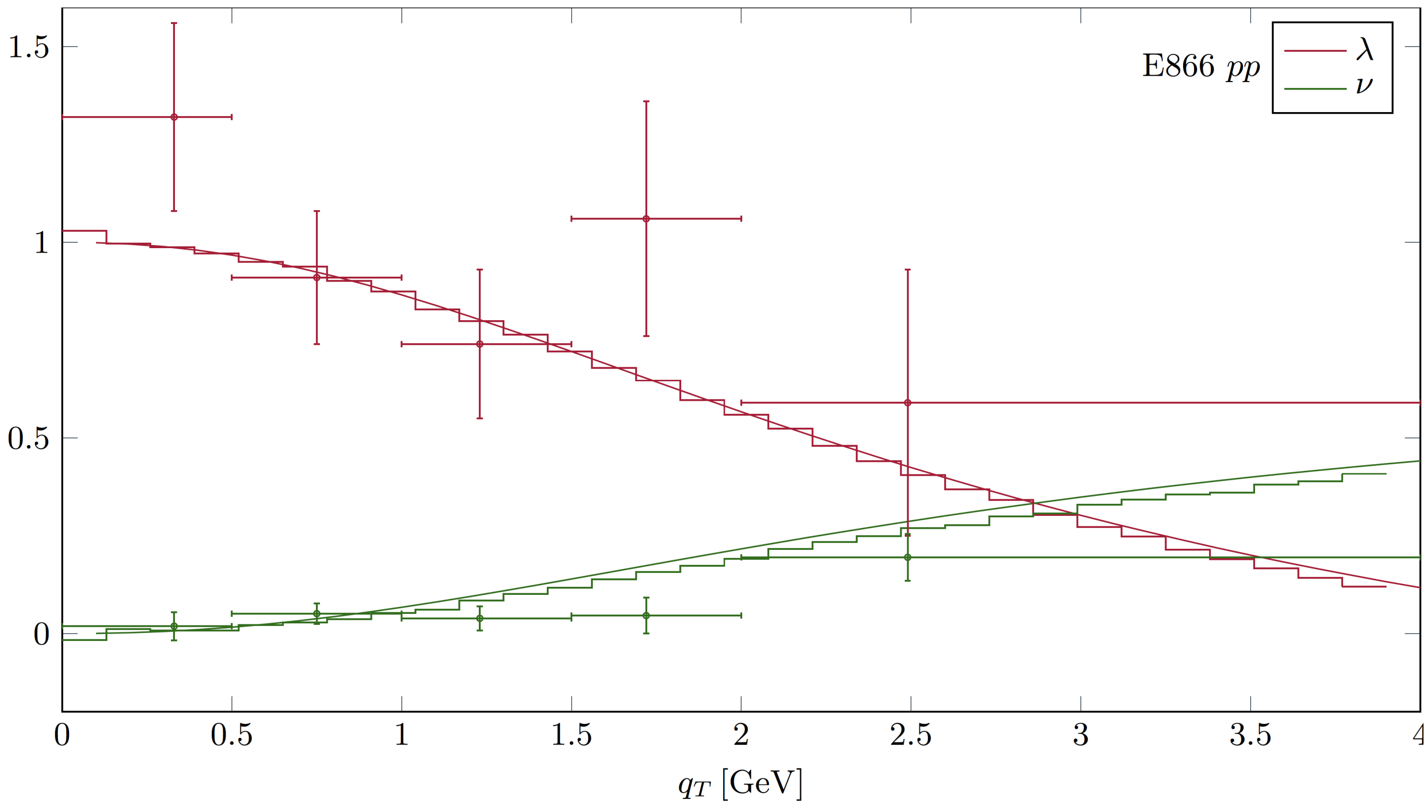
$$\text{NNLO (CMS):} \qquad \chi^2/\mathrm{N}_{\mathrm{data}} = 14.2/14 = 1.01$$

see also: Lambertsen, WV '16
Peng, Chang, McClellan, Teryaev '15

Lambertsen, WV ‘16

$pp, \ E = 800 \ \mathrm{GeV}$

E866

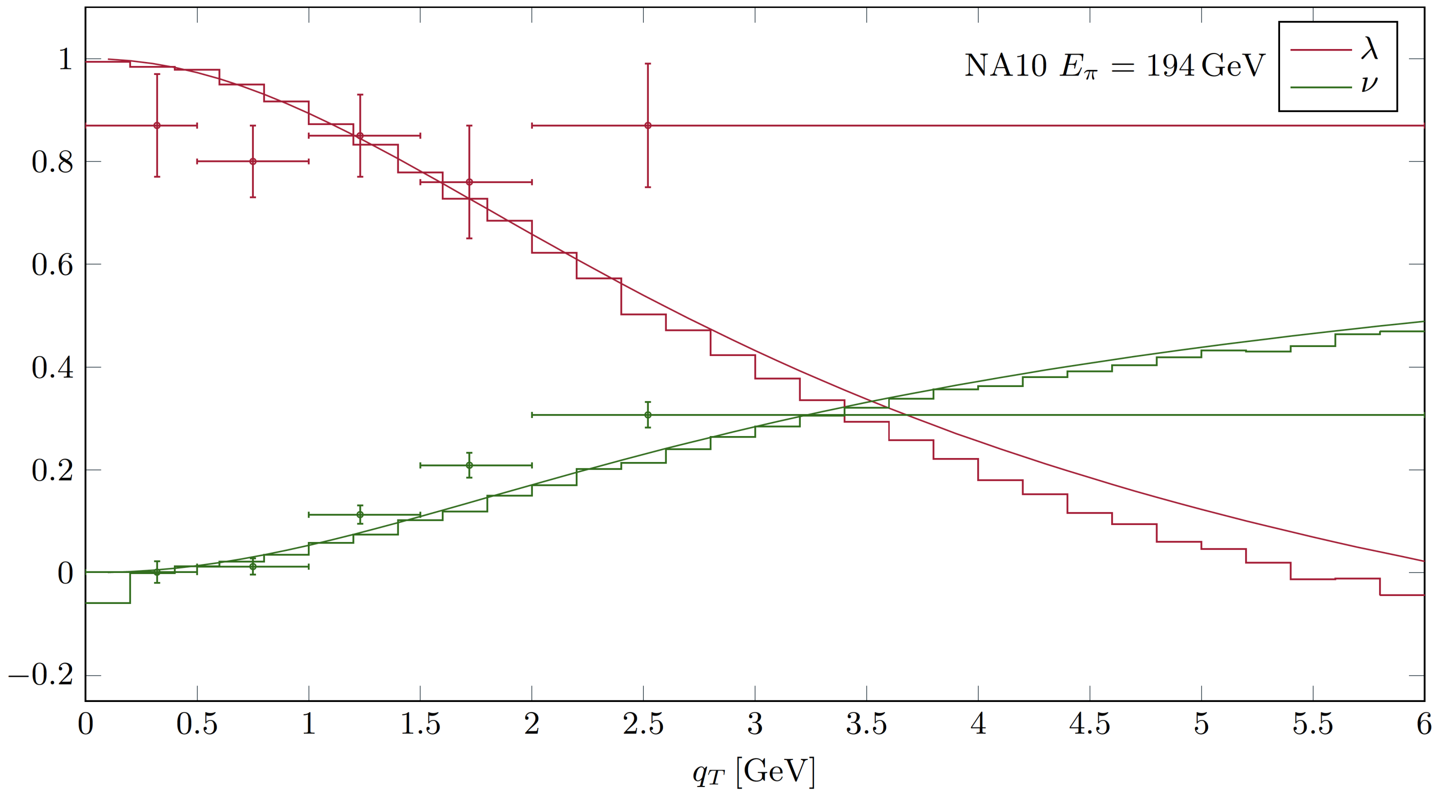


lines: LO $\mathcal{O}(\alpha_s)$

histograms: NLO $\mathcal{O}(\alpha_s^2)$

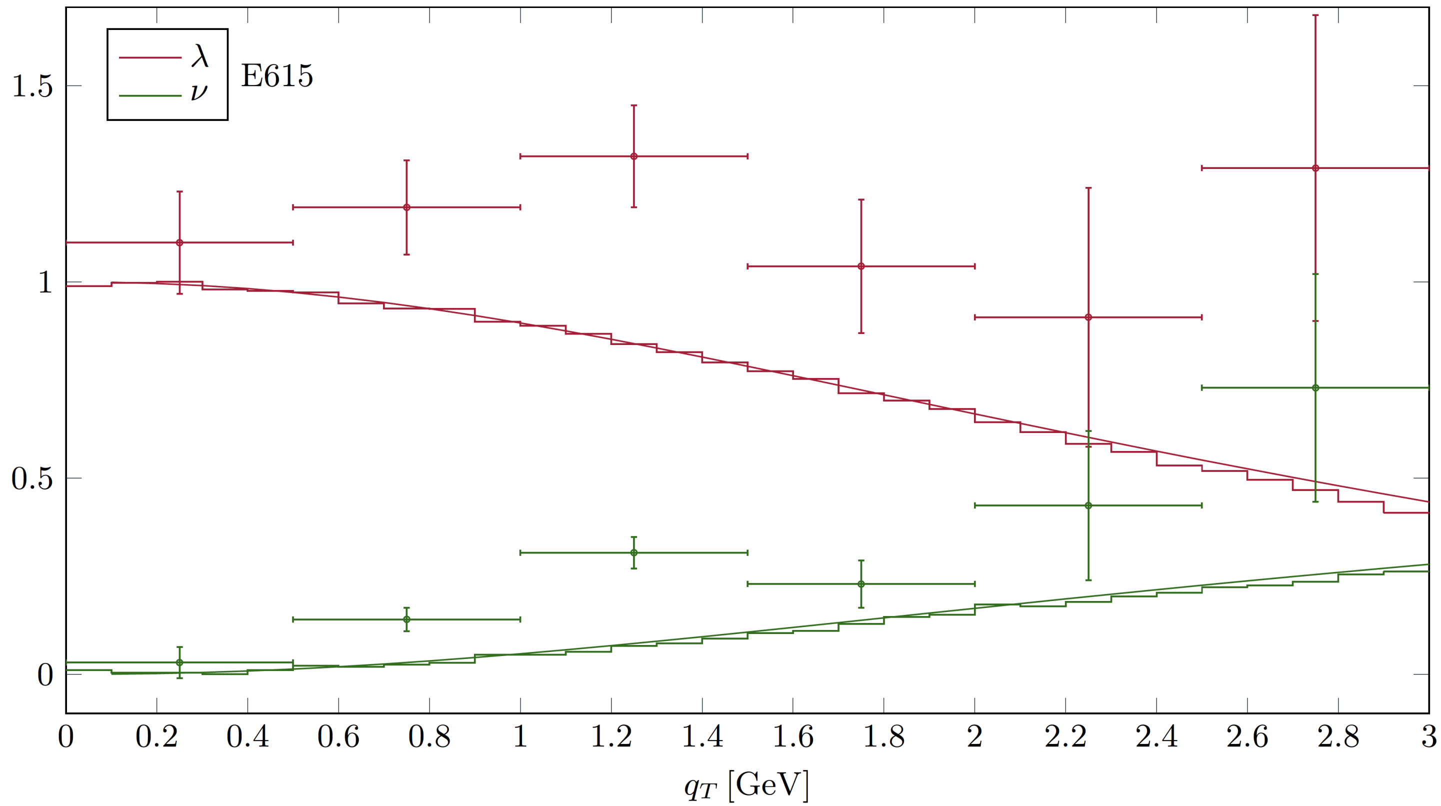
$\pi W,\ E_\pi = 194\,\mathrm{GeV}$

NA10



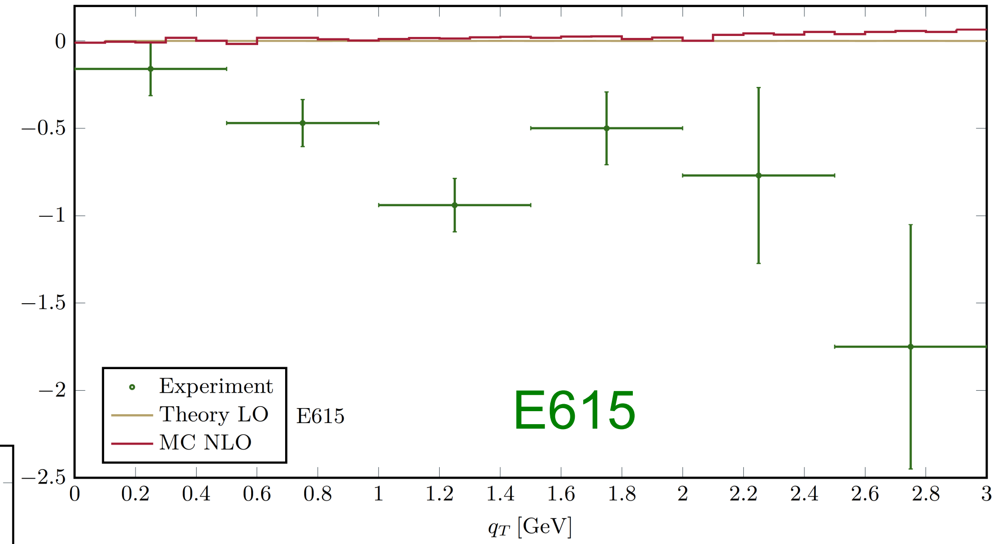
- “dispel myth” that pQCD cannot describe data: overall reasonable description

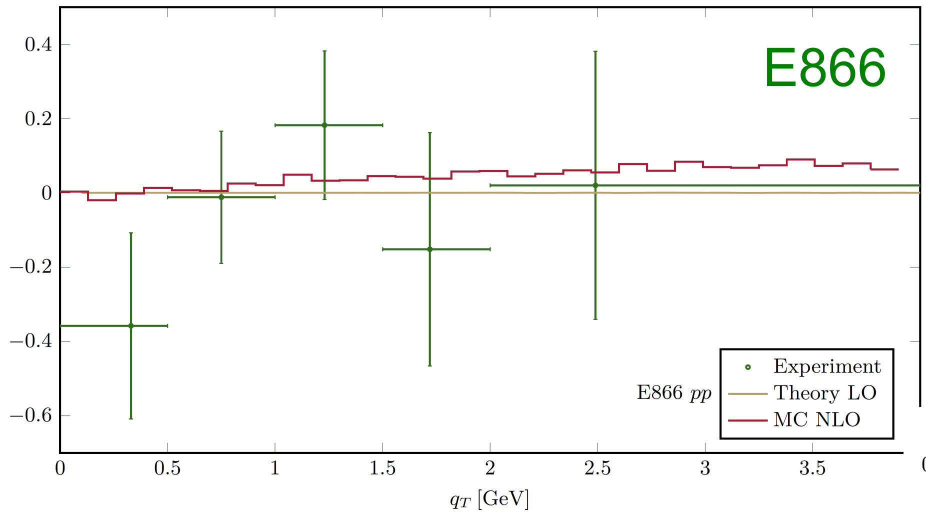
$\pi W,\ E_\pi = 252\,\mathrm{GeV}$ E615

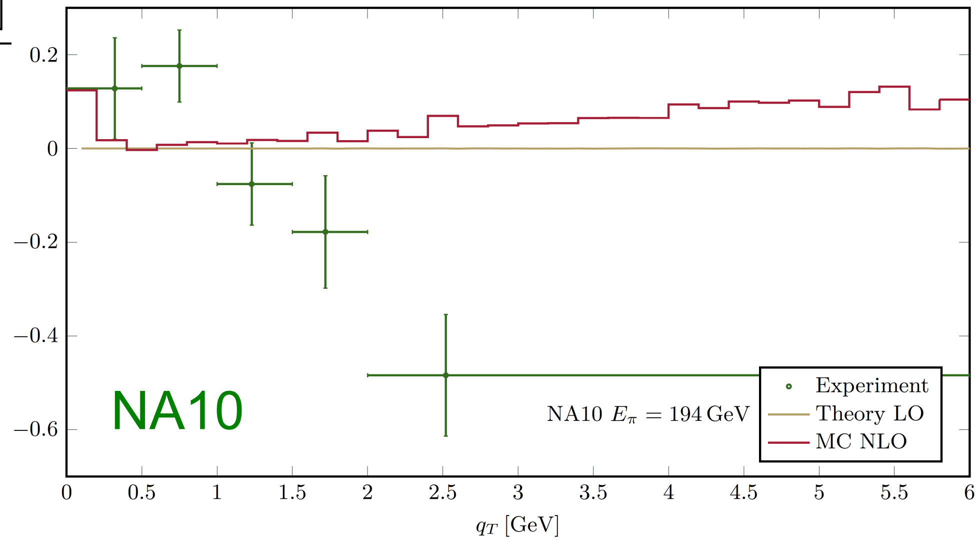


- note: positivity constraint $\lambda \leq 1$ Lam, Tung '78

Lam-Tung $1 - \lambda - 2\nu$

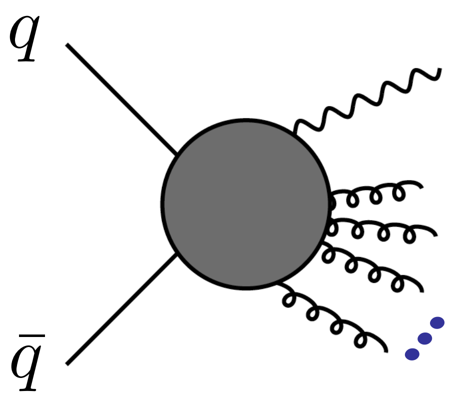






Resummation

- region $q_T \ll Q$:



$$\widehat{W}_T^{(k)} \propto \alpha_s^k \left(\frac{\log^{2k-1}(q_T^2/Q^2)}{q_T^2} \right)_+ + \ldots$$

"q_T logarithms"

- all-order resummation very well understood for W_T
  Collins-Soper-Sterman formalism

- 1-1 correspondence to TMD evolution

  Collins, Mert Aybat, Rogers, Qiu; Echevarria, Melis, Scimemi, d'Alesio; Kang, Prokudin, Sun, Yuan;…

- specifically structure functions at LO (CS frame):

$$\widehat{W}_T^{(1)} = -C_F \frac{\alpha_s}{2\pi} \frac{2\log(q_T^2/Q^2)+3}{q_T^2} + \ldots$$

$$\widehat{W}_L^{(1)} = 2\,\widehat{W}_{\Delta\Delta}^{(1)} = -C_F \frac{\alpha_s}{2\pi} \big(2\log(q_T^2/Q^2)+3\big) + \ldots$$

Boer, WV

- leading logs same to all orders in $\widehat{W}_T, \widehat{W}_L, \widehat{W}_{\Delta\Delta}$

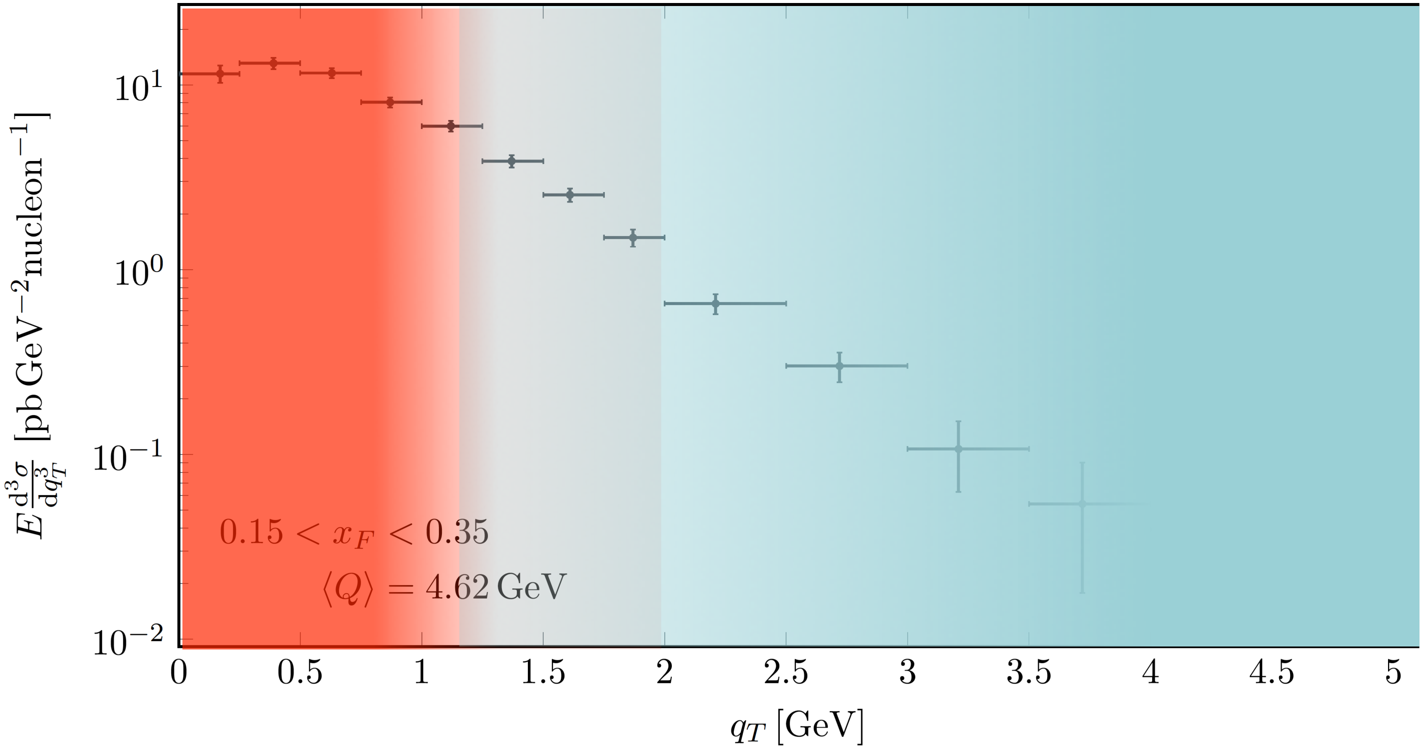
Berger, Qiu, Rodriguez

- differences at next-to-leading log Boer, WV

- at present, resummation for angular coeff. not fully understood
vital for TMD phenomenology!

- insights from TMD evolution might help Prokudin @ INT2017

- however, need to understand full q_T spectrum:

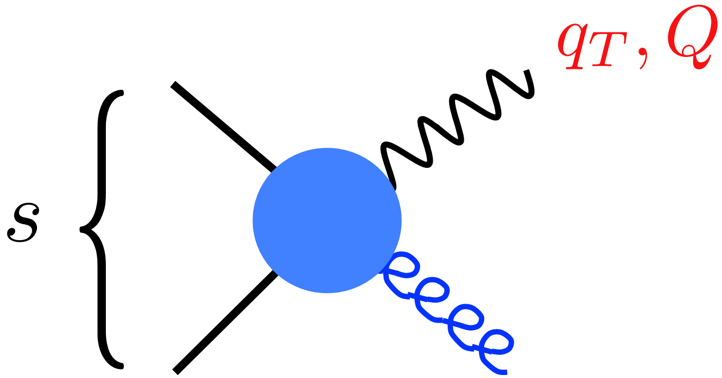


TMD regime

collinear regime

- crucial for matching ("Y term")
- develops different set of large logs

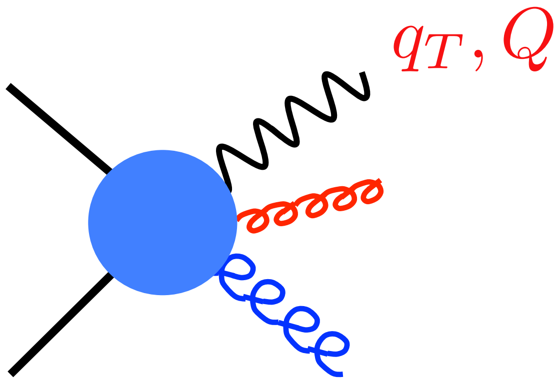
- LO :



$$\sqrt{s} \geq q_T + \sqrt{Q^2 + q_T^2}$$

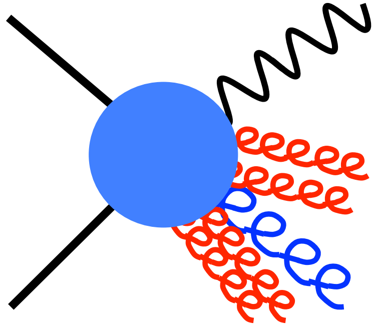
$$y_T \equiv \frac{q_T + \sqrt{q_T^2 + Q^2}}{\sqrt{s}} \leq 1$$

- NLO :



$$\frac{d\hat{\sigma}^{\mathrm{NLO}}}{dq_T} \propto \alpha_s \left[\mathcal{A} \log^2(1-y_T^2) + \mathcal{B} \log(1-y_T^2) + \mathcal{C} \right]$$

- N^{k}LO :



$$\frac{d\hat{\sigma}^{\mathrm{N}^k\mathrm{LO}}}{dq_T} \propto \alpha_s^k \log^{2k}(1-y_T^2) + \ldots$$

- threshold logarithms

de Florian, Kulesza, WV; Kidonakis, Gonsalves; …

- threshold resummation:

$$\hat{\sigma}_{q\bar{q}\to\gamma^* g}(N) \equiv \int_0^1 dy_T^2 (y_T^2)^{N-1} \frac{d\hat{\sigma}_{q\bar{q}\to\gamma^* g}}{dp_T}$$

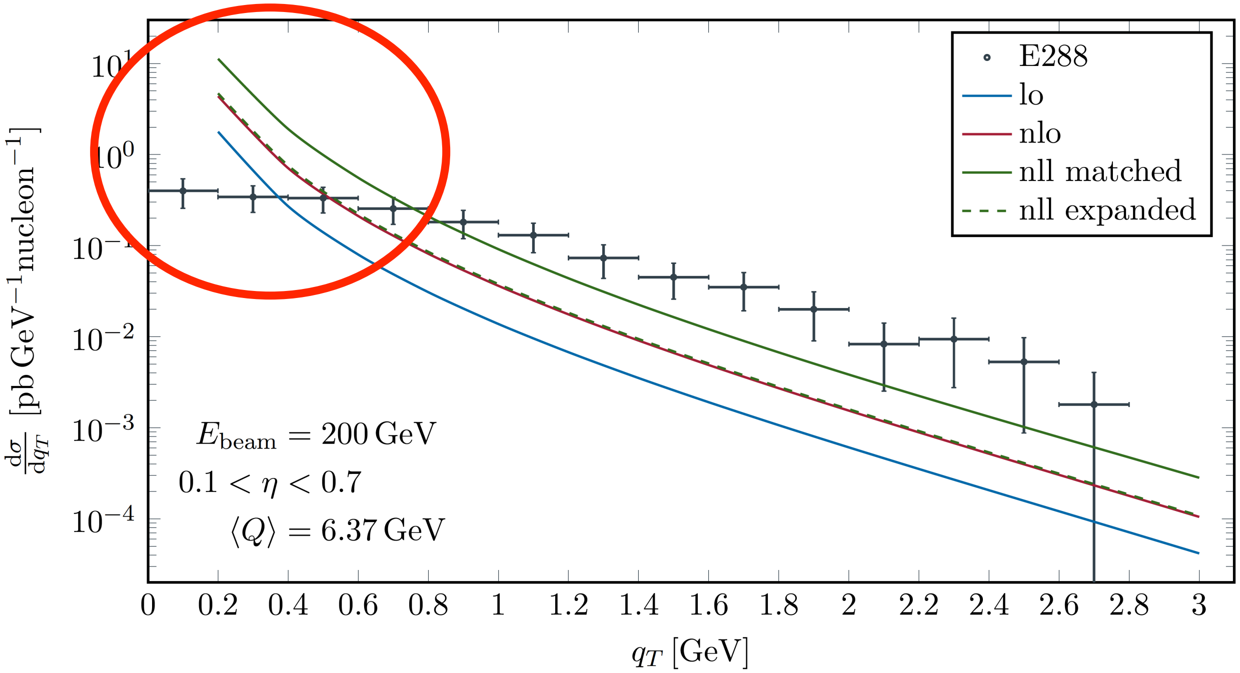
$$\hat{\sigma}^{(\mathrm{res})}_{q\bar{q}\to\gamma^* g}(N) = C_{q\bar{q}\to\gamma^* g}\, \underbrace{\Delta_N^q\, \Delta_N^{\bar{q}}\, J_N^g \Delta_N^{(\mathrm{int})}}_{}\, \hat{\sigma}^{\mathrm{LO}}_{q\bar{q}\to\gamma^* g}(N)$$

hard virtual corrections

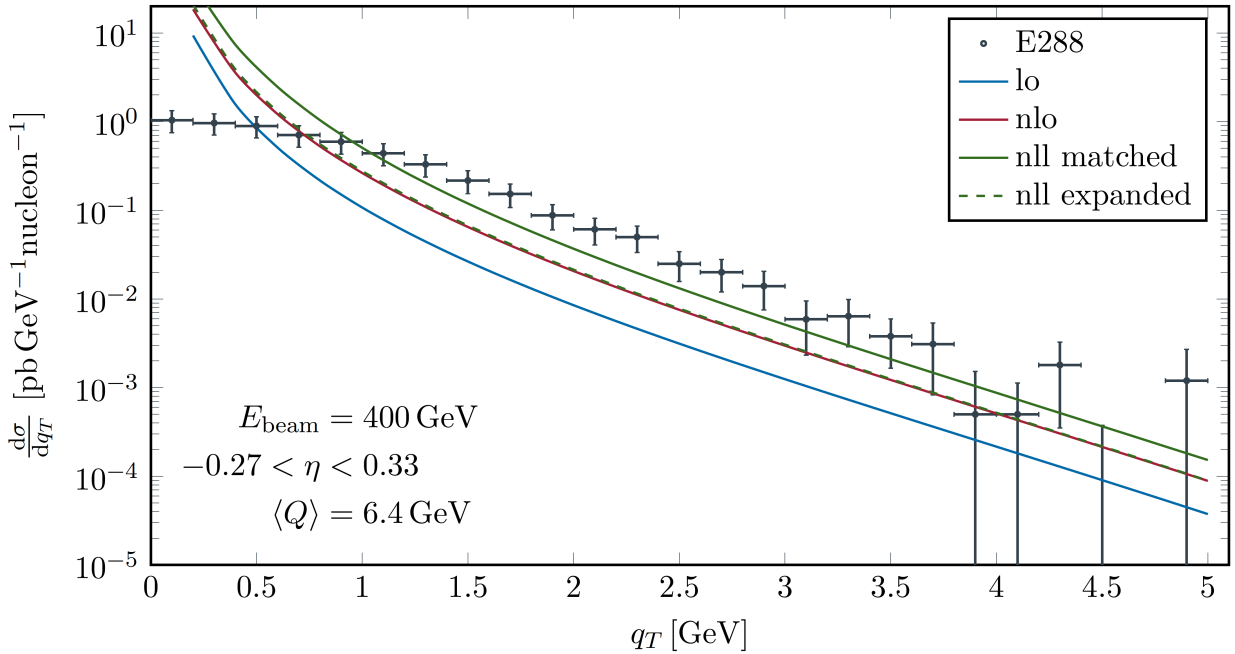
LO cross sec.

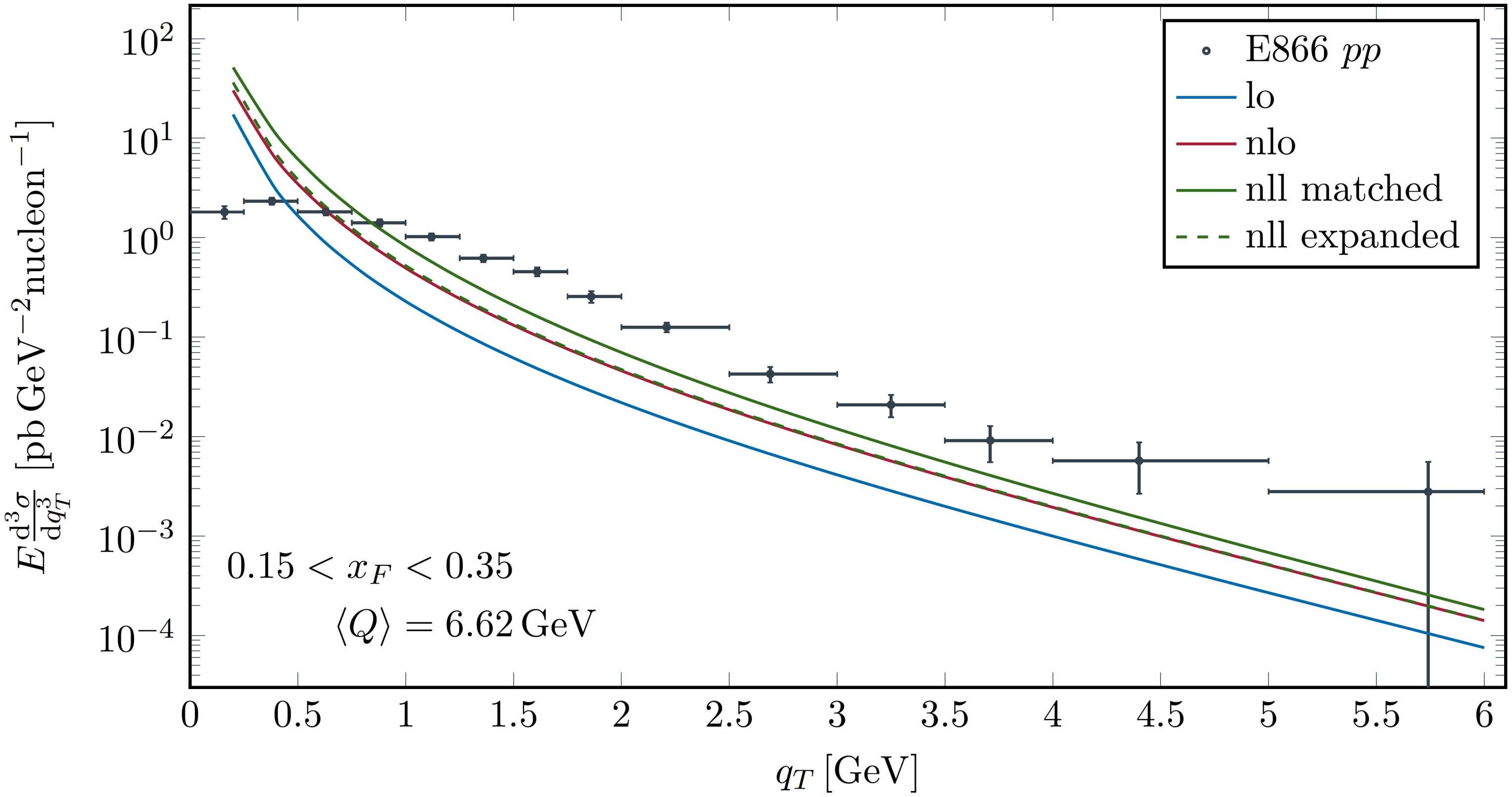
soft/coll. gluons

$$\ln \Delta_N^q = \int_0^1 dz \frac{z^{N-1}-1}{1-z} \int_{\mu_F^2}^{(1-z)^2 Q^2} \frac{dq^2}{q^2} A_q(\alpha_s(q^2))$$

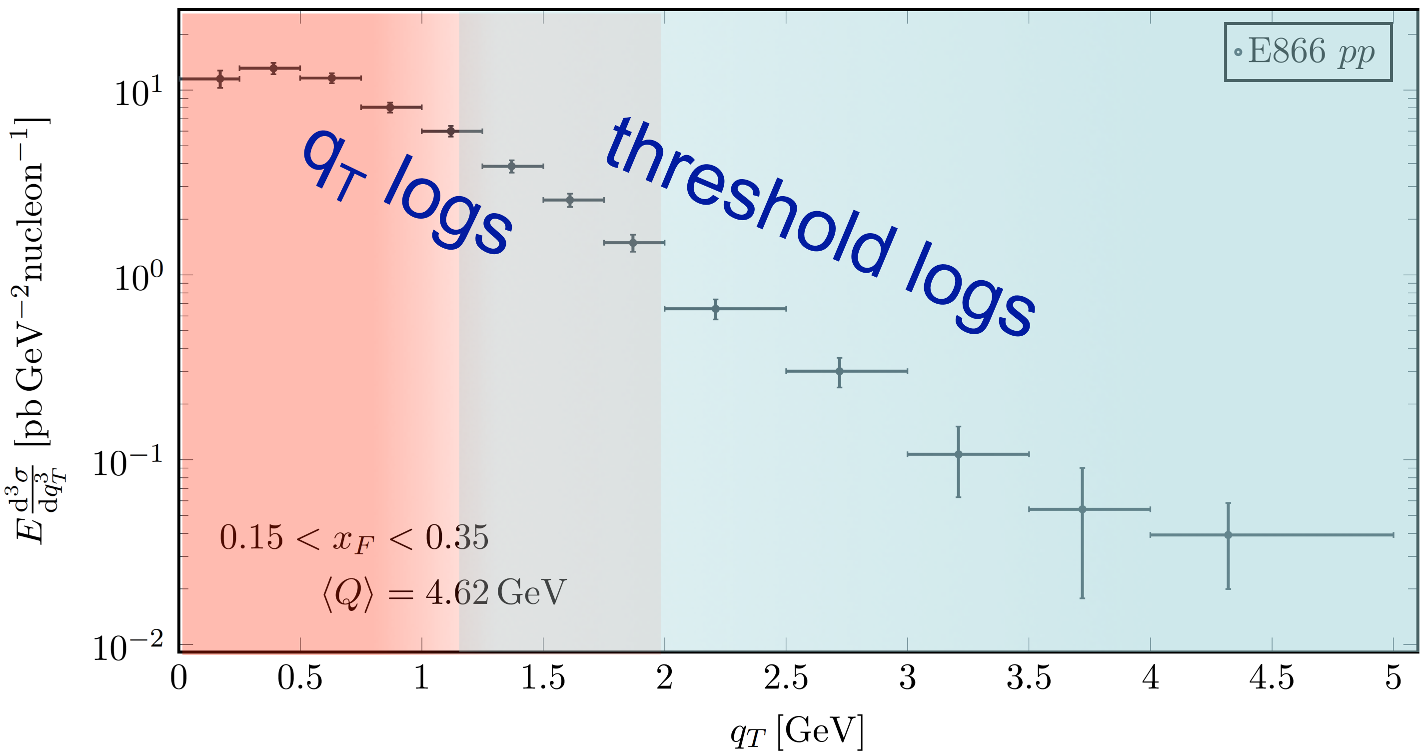


Lambertsen, Steiglechner, WV





- note, threshold resummation for angular dependences not known



- ultimately, will need joint resummation of both types of logs

- framework exists:

Laenen, Sterman, WV ’00
Lustermans, Waalewihn, Zeune ‘16
Muselli, Forte, Ridolfi ‘17

- e.g. inclusive Drell-Yan:

$$\hat{\sigma}^{(\mathrm{res})} \propto \exp\left[2\int_0^{Q^2}\frac{dk_\perp^2}{k_\perp^2}A_q(\alpha_s(k_\perp))\left[J_0(bk_\perp)\,K_0\left(\frac{2Nk_\perp}{Q}\right)+\ln\left(\frac{\bar{N}k_\perp}{Q}\right)\right]\right]$$

- "jointly resummed" cross section:

Laenen,Sterman,WV
Kulesza,Sterman,WV

$N \gg bQ$: threshold logs (e.g. b=0)

$bQ \gg N$: q_T logs

- Drell-Yan at high q_T: Muselli, Forte, Ridolfi

- for pure q_T resummation case:

$$\sigma^{\mathrm{eik}}(N,b) \approx \exp\left[-2\int_0^{Q^2}\frac{dk_\perp^2}{k_\perp^2}A_q(\alpha_s(k_\perp))\left(J_0(bk_\perp)-1\right)\ln\left(\frac{\bar{N}k_\perp}{Q}\right)\right]$$

- btw: shows how to match TMD and collinear framework

Conclusions and outlook:

- "dispel myth" that pQCD cannot describe λ, μ, ν
  overall reasonable description
- *not* meant to argue that there are no effects beyond fixed-order pQCD
- presently not a really good understanding of Drell-Yan cross section at high q_T
- serious studies of Boer-Mulders etc. should include pQCD radiative effects (including resummation, probably joint resummation)
- clear call for better data in fixed-target regime