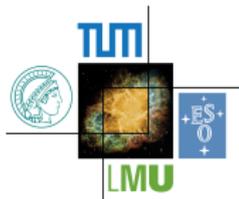


# Theory prospects with $b \rightarrow c(u)l\nu$ transitions

Martin Jung



**DFG** Deutsche  
Forschungsgemeinschaft

“Beyond the LHCb Phase-1 Upgrade”

La Biodola, Italy, 29th of May 2017

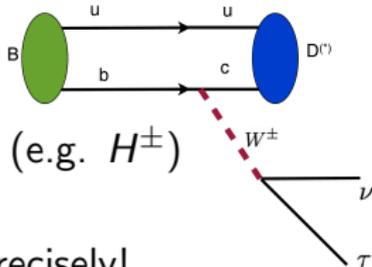
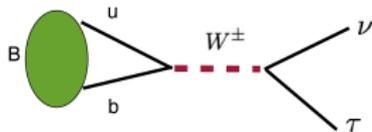
# Importance of (semi-)leptonic hadron decays

In the Standard Model:

- Determination of  $|V_{ij}|$  (7/9)

Beyond the Standard Model:

- Leptonic decays  $\sim m_f^2$ 
  - ➔ large relative NP influence possible (e.g.  $H^\pm$ )
- NP in semi-leptonic decays moderate
  - ➔ Need to understand the SM very precisely!
- NP: Relative to tree,  $\tau$  least constrained



Key advantages:

- Large rates
- Minimal hadronic input
  - ➔ This input is systematically improvable

Additionally: (almost) all flavour **anomalies** involve leptons

## Generalities

If  $R(D, D^*)$  are real, they will be established before 2nd upgrade

Consequently the objectives change:

- Differentiation between structures in  $b \rightarrow c\tau\nu$ 
  - ➡ Distributions in  $q^2$  + angles, polarization of  $\tau, D^*$
- Flavour structure on the lepton side ( $\rightarrow \mu$  vs.  $e$ )
  - ➡ improvements for electrons?
- Flavour structure on the quark side (e.g.  $b \rightarrow u$  vs.  $b \rightarrow c$ )
  - ➡ Possibilities in charm decays? (not part of this talk)

A lot of this is not yet done, insufficient data

Close collaboration of experiment and theory necessary

Objectives of this talk:

- Examples of challenging systematics (th + exp)
- Status of present tensions
- Identification of “clean” observables with differentiating power

## New systematics: BR measurements and isospin violation

Branching ratio measurements require normalization. . .

- $B$  factories: depends on  $\Upsilon \rightarrow B^+ B^-$  vs.  $B^0 \bar{B}^0$
- LHCb: normalization mode, usually obtained from  $B$  factories

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Assumptions entering this normalization:

- PDG: assumes  $r_{+0} \equiv \Gamma(\Upsilon \rightarrow B^+ B^-) / \Gamma(\Upsilon \rightarrow B^0 \bar{B}^0) \equiv 1$
- LHCb: (mostly) assumes  $f_u \equiv f_d$ , uses  $r_{+0}^{\text{HFAG}} = 1.058 \pm 0.024$

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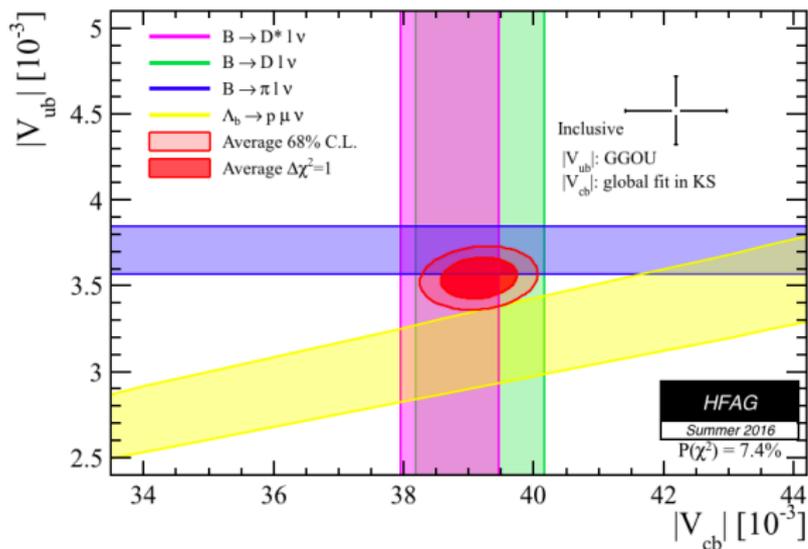
Both approaches problematic: [MJ'16 [1510.03423]]

- Potential large isospin violation in  $\Upsilon \rightarrow BB$  [Atwood/Marcano'90]
- Measurements in  $r_{+0}^{\text{HFAG}}$  assume isospin in exclusive decays
  - ➡ This is one thing we want to test!
- ➡ Avoiding this assumption yields  $r_{+0} = 1.027 \pm 0.037$
- Isospin asymmetries test NP with  $\Delta I = 1, 3/2$  (e.g.  $b \rightarrow s\bar{u}u$ )
  - ➡ Isospin asymmetry  $B \rightarrow J/\psi K$ :  $A_I = -0.009 \pm 0.024$

Affects every percent-level BR measurement  
 $B \rightarrow J/\psi K$  can be used to determine  $f_u/f_d$ !

# $|V_{xb}|$ : inclusive versus exclusive

Long-standing problem:



- Very hard to explain by NP [Crivellin/Pokorski'15] (but see [Colangelo/de Fazio'15])
- Likely experimental/theoretical systematics

# $|V_{xb}|$ : Recent developments

$V_{cb}$ :

Recent Belle  $B \rightarrow D, D^* \ell \nu$  analyses

Recent lattice results for  $B \rightarrow D$

[FNAL/MILC, HPQCD, RBC/UKQCD (ongoing)]

➡  $B \rightarrow D$  between incl. +  $B \rightarrow D^*$

New lattice result for  $B \rightarrow D^*$  [HPQCD]

➡  $V_{cb}^{\text{incl}}$  cv, compatible with old result

$B \rightarrow D^* \ell \nu$  re-analyses with CLN,

$|V_{cb}| = 39.3(1.0)10^{-2}$  [Bernlochner+'17]

+ BGL [Bigi+, Grinstein+'17] (Belle only),

$|V_{cb}| = 40.4(1.7)10^{-2}$

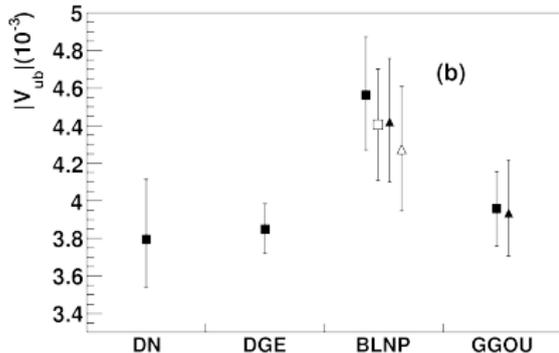
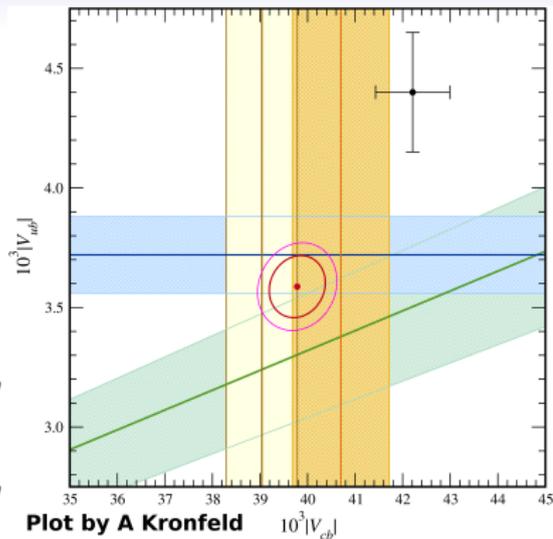
New BaBar analysis of  $V_{ub}$  incl.:

Dependence on theory treatment!

➡ GGOU  $2\sigma$  lower than WA

➡ Compatible w/ PDG exclusive avg

Hints towards resolution, not yet conclusive



## Prospects $b \rightarrow (u, c)(e, \mu) @ \text{LHCb}$

Potential unambiguous  $|V_{xb}|$  determination before phase-II upgrade

➡ Measuring  $b \rightarrow u, c\ell\nu$  not about this

Instead, model-independent determinations of NP contributions

- If FNU in  $b \rightarrow c$  is confirmed, expect “something” in  $b \rightarrow u$
- Also, with  $b \rightarrow c\tau\nu$  affected,  $\mu$  vs.  $e$  important to check
- Universality checks of right-handed currents interesting

$|V_{ub}/V_{cb}|$  from  $\Lambda_b$  important ingredient right now. . .

- Tests different NP combinations than mesonic modes
- Which observables are measurable?
- How much can we reduce the **systematics**?
- FFs need improvement, but not the main issue

$B_s \rightarrow K\ell\nu$  essentially probes the same physics as  $B \rightarrow \pi\ell\nu$

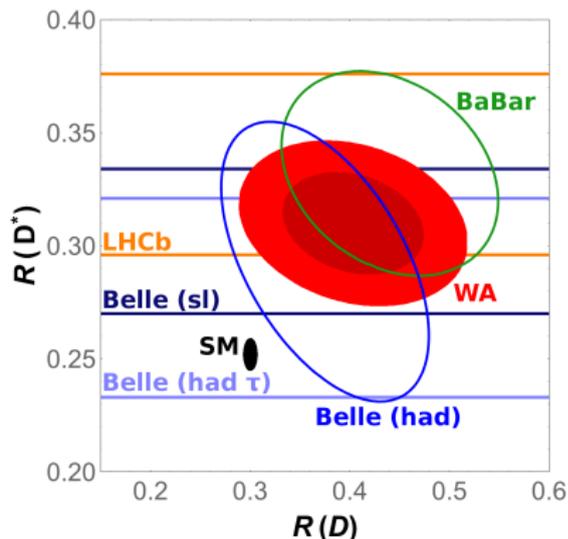
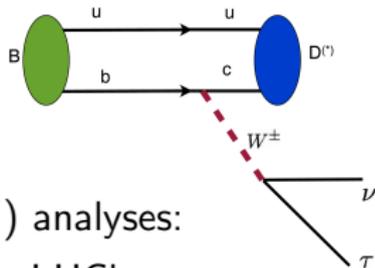
➡ direct competition with Belle II

$B \rightarrow p\ell\nu$  interesting new idea

➡ Challenging, qualitative theory progress required!

# Experimental Situation for $b \rightarrow c\tau\nu$ 2017

$$R(X) \equiv \frac{\text{Br}(B \rightarrow X\tau\nu)}{\text{Br}(B \rightarrow X\ell\nu)}$$



4 recent  $R(D^{(*)})$  analyses:

- $R(D^*)$  from LHCb [1506.08614]
- Belle update + new measurement (had./sl tag) [1507.03233,1603.06711],  $\tau$ -polarization +  $R(D^*)(\tau \rightarrow \text{had})$  [1608.06391]

↪ **4.0 $\sigma$  tension** [HFAG]

Further  $b \rightarrow c\tau\nu$  inputs:

- Differential rates from Belle, BaBar
- Total width of  $B_c$
- ( $b \rightarrow X_c\tau\nu$  by LEP)

contours: 68% CL  
filled: 95(68)% CL

## SM predictions [see also Zoltan's talk]

SI amplitude: kinematics  $\times$  FC coupling (SM: CKM)  $\times$  form factor

Strategy SM predictions:  $V_{cb}$  + leading FF cancels data + theoretical input from LQCD/HQET for FF ratios

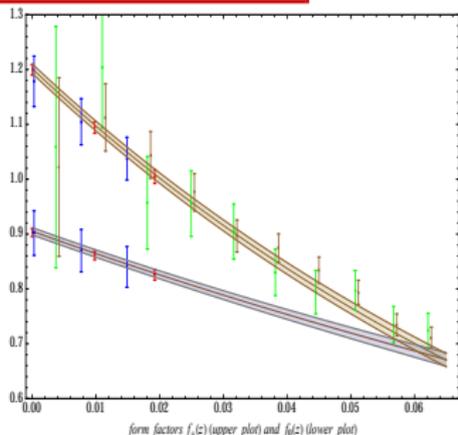
$B \rightarrow D$ : 2 form factors  $f_{+,0}$

- Data determines shape of  $f_+(q^2)$
- LQCD required for  $f_0$ : fit HPQCD + FNAL/MILC, use  $f_+(0) = f_0(0)$
- $\rightarrow R(D) = 0.301 \pm 0.003$  [Bigi/Gambino'16]

$B \rightarrow D^*$ : 4 form factors  $V, A_{0,1,2}$

- 3/4  $\rightarrow$  data (+HQET, unitarity  $\rightarrow$  CLN)
- HQET for  $A_0$  [Falk/Neubert], enhance uncertainty [Fajfer/Kamenik]
- $\rightarrow R(D^*) = 0.252 \pm 0.003$ , (0.257 from re-analysis [Bernlochner+'17])
- LQCD for non-maximal recoil underway

(Very) good control, effect too large to be in CLN relations



## NP in (semi-)leptonic decays

EFT for  $b \rightarrow c\tau\nu$  transitions (no light  $\nu_R$ , SM:  $C_{V_L} = 1$ ,  $C_{i \neq V_L} = 0$ ):

$$\mathcal{L}_{\text{eff}}^{b \rightarrow c\tau\nu} = -\frac{4G_F}{\sqrt{2}} V_{cb} \sum_j^5 C_j \mathcal{O}_j, \quad \text{with}$$

$$\mathcal{O}_{V_{L,R}} = (\bar{c}\gamma^\mu P_{L,R}b)\bar{\tau}\gamma_\mu\nu, \quad \mathcal{O}_{S_{L,R}} = (\bar{c}P_{L,R}b)\bar{\tau}\nu, \quad \mathcal{O}_T = (\bar{c}\sigma^{\mu\nu}P_Lb)\bar{\tau}\sigma_{\mu\nu}\nu.$$

NP models typically generate **subsets**; for a charged scalar:

NP couplings  $C_{S_{L,R}}$  (**complex**),  $C_{V_L} = C_{V_L}^{\text{SM}} = 1$ ,  $C_{V_R} = C_T = 0$

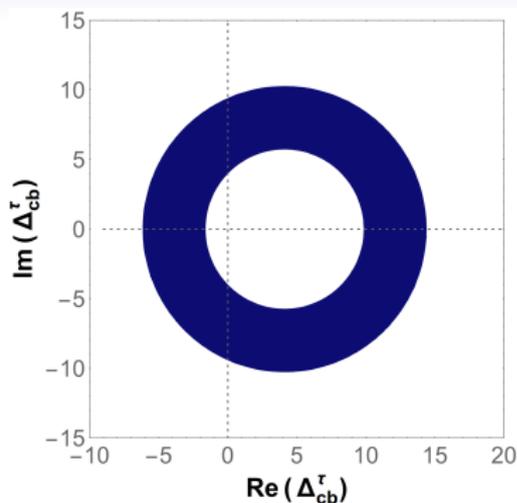
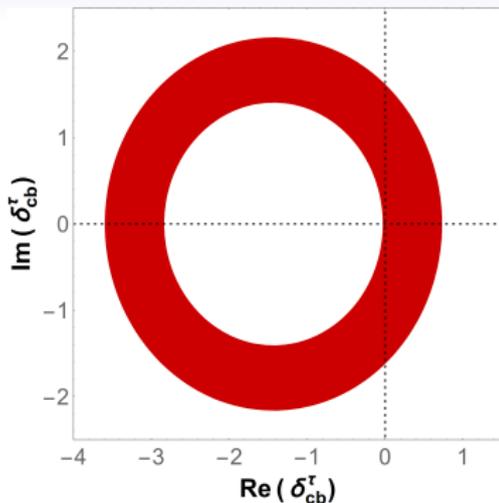
- Model-independent subclass as long as  $C_{S_{L,R}}$  general
- Phenomenologically  $C_{S_{L,R}}^{quqd^l} \sim m_{qu} m_l$  (e.g. Type III)

➡ Used to illustrate here, appearing combinations:

$$R(D) : \delta^{cbl} \equiv \frac{(C_{S_L} + C_{S_R})(m_B - m_D)^2}{m_l(\bar{m}_b - \bar{m}_c)} \quad R(D^*) : \Delta^{cbl} \equiv \frac{(C_{S_L} - C_{S_R})m_B^2}{m_l(\bar{m}_b + \bar{m}_c)}$$

Can trivially explain  $R(D^{(*)})$ ! Exclusion possible with **specific flavour structure** or **more  $b \rightarrow c\tau\nu$  observables!**

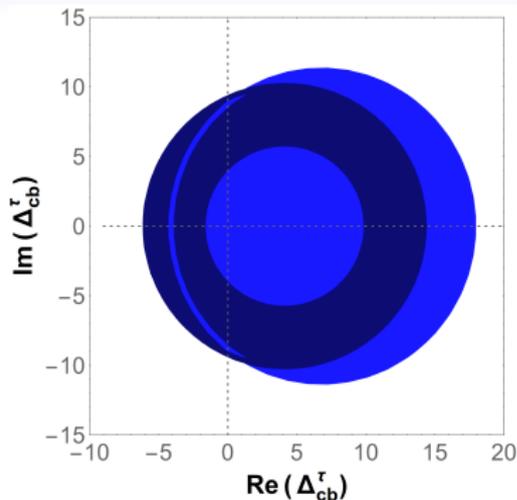
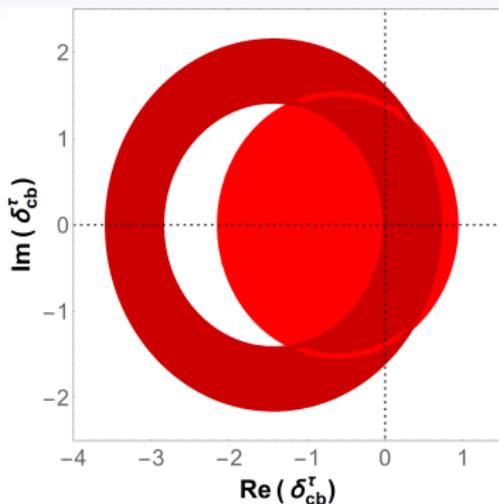
## $b \rightarrow c\tau\nu$ data and scalar NP [Celis/MJ/Li/Pich'17]



$R(D), R(D^*)$ :

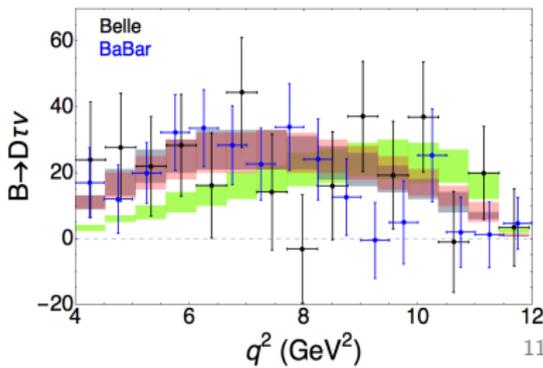
- $R(D)$  compatible with SM at  $\sim 2\sigma$
- Preferred scalar couplings from  $R(D^*)$  huge ( $|C_{S_L} - C_{S_R}| \sim 1 - 5$ )
- Can't go beyond circles with just  $R(D, D^*)$ !

# $b \rightarrow cTV$ data and scalar NP [Celis/MJ/Li/Pich'17]

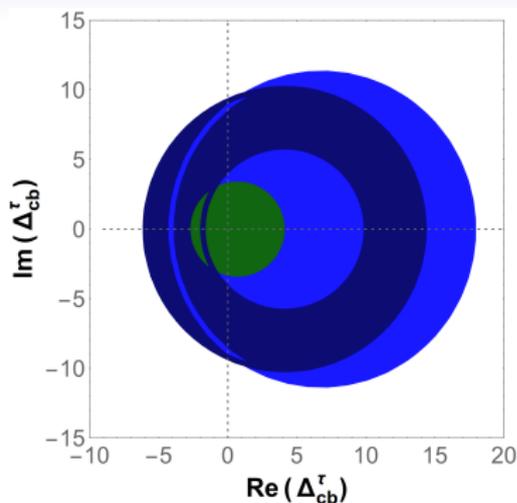
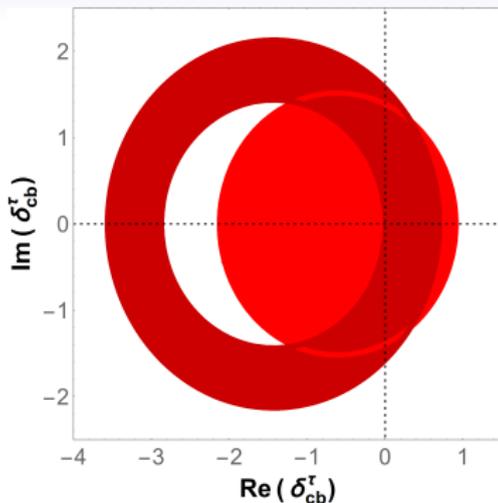


## Differential rates:

- compatible with SM and NP
- already now constraining, especially in  $B \rightarrow DTV$
- “theory-dependence” of data needs addressing [Bernlochner+'17]



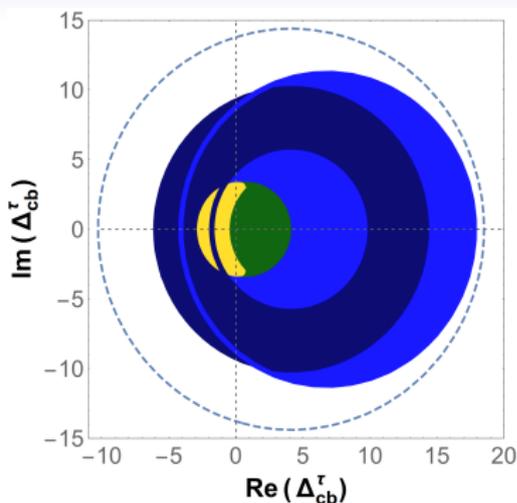
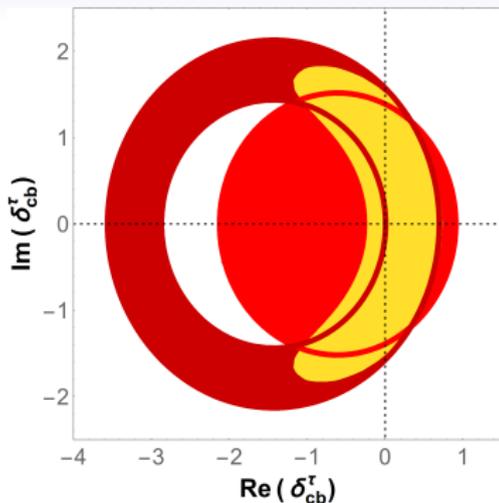
## $b \rightarrow c\tau\nu$ data and scalar NP [Celis/MJ/Li/Pich'17]



### Total width of $B_c$ :

- $B_c \rightarrow \tau\nu$  is an obvious  $b \rightarrow c\tau\nu$  transition
  - ➡ not measurable in foreseeable future
  - ➡ can oversaturate total width of  $B_c$ ! [X.Li+'16]
- Excludes second real solution in  $\Delta_{cb}^\tau$  plane (even scalar NP for  $R(D^*)$ ? [Alonso+'16])

## $b \rightarrow c\tau\nu$ data and scalar NP [Celis/MJ/Li/Pich'17]



### $\tau$ polarization:

- So far not constraining (shown:  $\Delta\chi^2 = 1$ )
- Differentiate NP models: with scalar NP [Celis/MJ/Li/Pich'13]

$$X_2^{D^{(*)}}(q^2) \equiv R_{D^{(*)}}(q^2) \left[ A_\lambda^{D^{(*)}}(q^2) + 1 \right] = X_{2,SM}^{D^{(*)}}(q^2)$$

Consistent explanation in 2HDMs possible, flavour structure?

## Differentiating models with $b \rightarrow c\tau\nu$ observables

Large  $R(D^*)$  possible with NP in  $V_L$  ( $\hat{R}(X) = R(X)/R(X)_{SM}$ ):

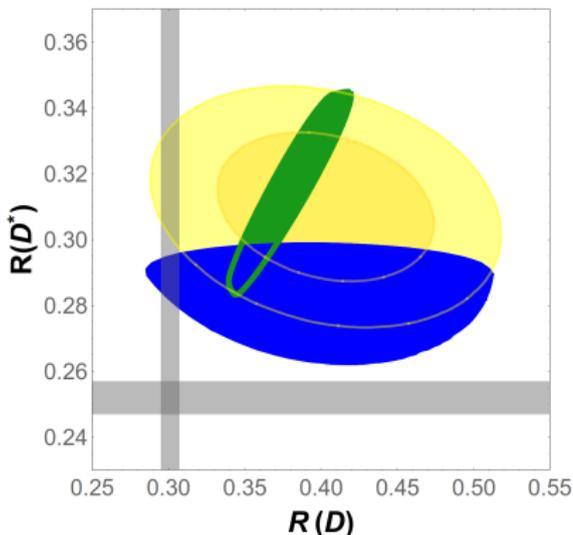
- trivial prediction:  $\hat{R}(D) = \hat{R}(D^*) = \hat{R}(\Lambda_c) = \dots \stackrel{exp}{\sim} 1.25$
- can be related to anomaly in  $B \rightarrow K^{(*)}\ell^+\ell^-$  modes
- $\hat{R}(X_c) = 0.99 \pm 0.10$  measured by LEP, oversaturation
- issues with  $\tau \rightarrow \mu\nu\nu$  [Feruglio+'16] and  $b\bar{b} \rightarrow X \rightarrow \tau^+\tau^-$  [Faroughy+'16]

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Fit results for the two scenarios for  $B \rightarrow D^{(*)}\tau\nu$ :

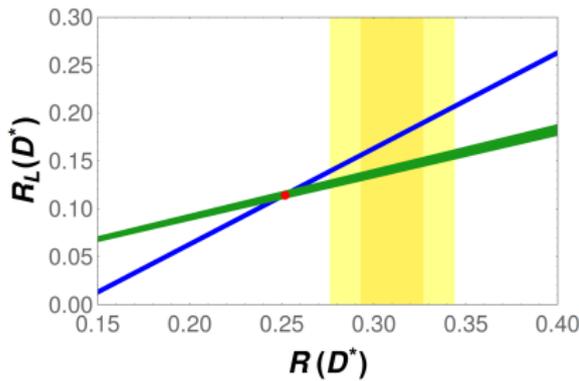
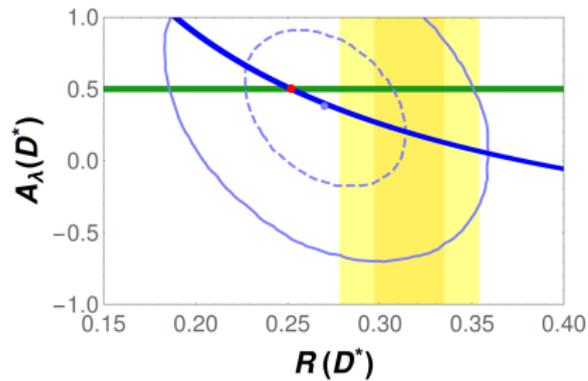


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Fit predictions for polarization-dependent  $B \rightarrow D^*\tau\nu$  observables:

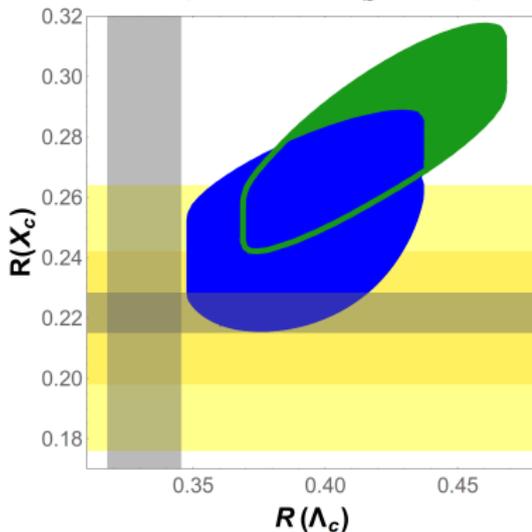


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- issues with  $\tau \rightarrow \mu\nu\nu$  [Feruglio+'16] and  $b\bar{b} \rightarrow X \rightarrow \tau^+\tau^-$  [Farouhy+'16]

Fit predictions for  $B \rightarrow X_{c\tau\nu}$  and  $\Lambda_b \rightarrow \Lambda_{c\tau\nu}$ :



## NP in $b \rightarrow u\tau\nu$ transitions

$b \rightarrow u\tau\nu$  less explored experimentally,  $|V_{ub}/V_{cb}|^2 \lesssim 1\%$ :

- $R(\tau) \equiv BR(B \rightarrow \tau\nu)/BR(B \rightarrow \pi\ell\nu)$  about  $1.8\sigma$  from SM
- $R(\pi)$  not significantly measured yet
- ➡ Data consistent with SM as well as sizable NP

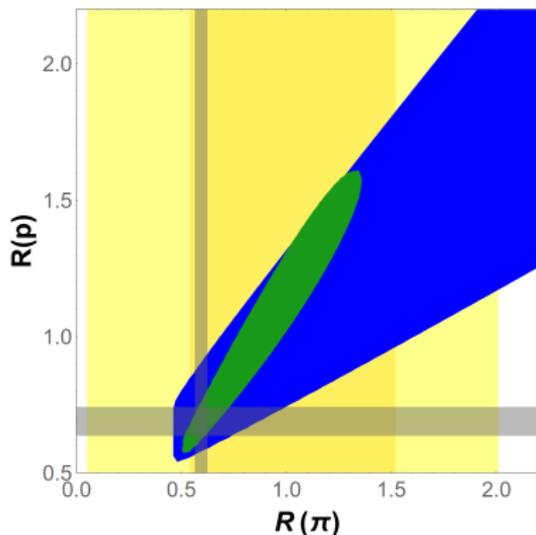
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Analyse  $b \rightarrow u\tau\nu$  individually:

•  $R(\tau)$  yields correlation between  $R(\pi)$  and  $R(\rho)$



More observables needed!  
 $\Lambda_b$  provides uncommon parameter combinations  
 $B_s \rightarrow K^{(*)}\tau\nu$  decays competitive? Detector requirements?  
Pionic final states possible?

# Conclusions

Excellent physics potential for LHCb beyond Run 4

- Present tensions:
  - $V_{xb}$  exclusive vs. inclusive: progress possible/probable
  - $b \rightarrow cTV$ : indications of lepton-non-universal NP
  - ➡ New measurements/observables constrain NP more severely
- Any BR measurement at the (few-)% level requires dealing with production asymmetries @  $B$  factories properly
- Should tensions be real, they're established by LS 3
  - ➡ Expect smaller deviations anyway (smaller  $R(D^*)$  would improve most NP interpretations)
  - ➡ Need to pin down precise structure of NP (Dirac, flavour)
  - ➡ Clean observables available to differentiate between different NP
  - ➡ Need for distributions + polarization measurements
- Chance to constrain  $b \rightarrow u$  transitions like  $b \rightarrow c$  now
  - ➡ Experimentally challenging, which detector changes could help?

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**Thank you for your attention!**

## Generic features and issues in 2HDMs

Charged Higgs possible as explanation of  $b \rightarrow c\tau\nu$  data...

However, typically expect  $\Delta R(D^*) < \Delta R(D)$

Generic feature: Relative influence larger in leptonic decays!

- No problem in  $b \rightarrow c\tau\nu$  since  $B_c \rightarrow \tau\nu$  won't be measured
- Large charm coupling required for  $R(D^*)$
- ➡ Embedding  $b \rightarrow c\tau\nu$  into a viable model complicated!
- ➡  $D_{d,s} \rightarrow \tau, \mu\nu$  kill typical flavour structures with  $C_{S_{L,R}} \sim m$
- ➡ Only fine-tuned models survive all (semi-)leptonic constraints

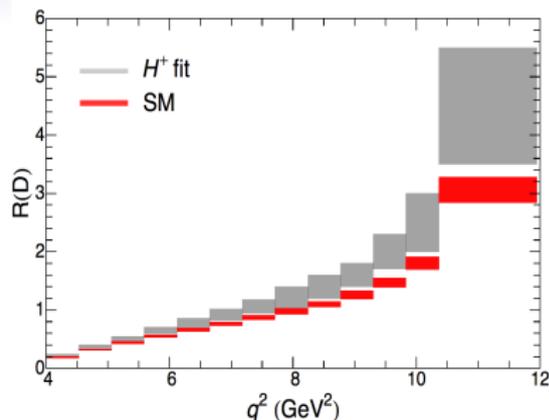
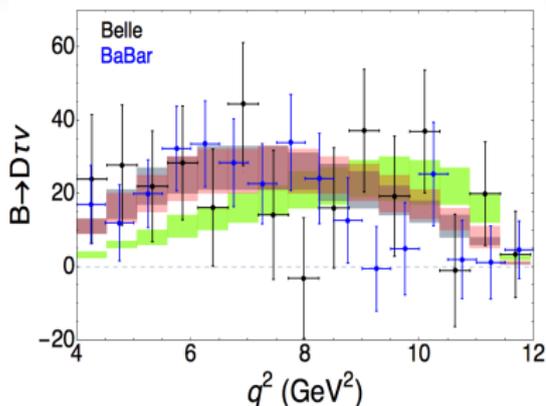
$b \rightarrow s\ell\ell$  very complicated to explain with scalar NP

➡ 2HDM alone tends to predict  $b \rightarrow s\ell\ell$  to be QCD-related

$b\bar{b} \rightarrow (H, A) \rightarrow \tau^+\tau^-$  poses a severe constraint [Faroughy+'16, Admir's talk]

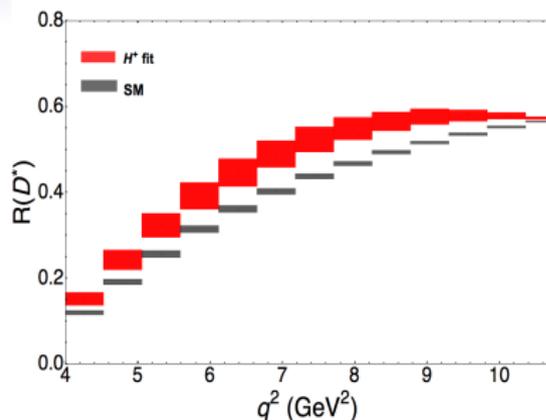
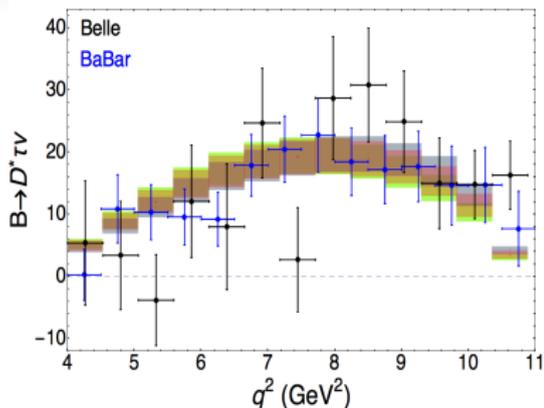
2HDMs strongly prefer a smaller value for  $R(D^*)$ !

# The differential distributions $d\Gamma(B \rightarrow D^{(*)}\tau\nu)/dq^2$



- Data stat. uncertainties only, BaBar rescaled
- Bands 68% CL (bins highly correlated):
  - Grey: NP fit including  $R(D)$
  - Red: SM fit (distributions only)
  - Green: Allowed by  $R(D)$ , excluded by distribution
- Need better experimental precision, ideally  $dR(D)/dq^2$
- Parts of NP parameter space clearly excluded

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 Red: SM fit (distributions only)  
 Green: Allowed by  $R(D^*)$ , excluded by distribution
- Need better experimental precision, ideally  $dR(D^*)/dq^2$
- Not very restrictive at the moment

# Implications of the Higgs EFT for Flavour: $q \rightarrow q' \ell \nu$

$b \rightarrow c \tau \nu$  transitions (SM:  $C_{V_L} = 1, C_{i \neq V_L} = 0$ ):

$$\mathcal{L}_{\text{eff}}^{b \rightarrow c \tau \nu} = -\frac{4G_F}{\sqrt{2}} V_{cb} \sum_j^5 C_j \mathcal{O}_j, \quad \text{with}$$

$$\mathcal{O}_{V_{L,R}} = (\bar{c} \gamma^\mu P_{L,R} b) \bar{\tau} \gamma_\mu \nu, \quad \mathcal{O}_{S_{L,R}} = (\bar{c} P_{L,R} b) \bar{\tau} \nu,$$

$$\mathcal{O}_T = (\bar{c} \sigma^{\mu\nu} P_L b) \bar{\tau} \sigma_{\mu\nu} \nu.$$

- All operators are independently present already in the linear EFT
- However: Relations between **different** transitions:  
 $C_{V_R}$  is **lepton-flavour universal** [see also Cirigliano+'09]  
Relations between charged- and neutral-current processes, e.g.  
 $\sum_{U=u,c,t} \lambda_{Us} C_{S_R}^{(U)} = -\frac{e^2}{8\pi^2} \lambda_{ts} C_S^{(d)}$  [see also Cirigliano+'12, Alonso+'15]
- These relations are again absent in the non-linear EFT

## Matching for $b \rightarrow c\ell\nu$ transitions

$$C_{V_L} = -\mathcal{N}_{CC} \left[ C_L + \frac{2}{v^2} c_{V5} + \frac{2V_{cb}}{v^2} c_{V7} \right],$$

$$C_{V_R} = -\mathcal{N}_{CC} \left[ \hat{C}_R + \frac{2}{v^2} c_{V6} \right],$$

$$C_{S_L} = -\mathcal{N}_{CC} (c'_{S1} + \hat{c}'_{S5}),$$

$$C_{S_R} = 2\mathcal{N}_{CC} (c_{LR4} + \hat{c}_{LR8}),$$

$$C_T = -\mathcal{N}_{CC} (c'_{S2} + \hat{c}'_{S6}),$$

where  $\mathcal{N}_{CC} = \frac{1}{2V_{cb}} \frac{v^2}{\Lambda^2}$ ,  $C_L = 2c_{LL2} - \hat{c}_{LL6} + \hat{c}_{LL7}$  and  $\hat{C}_R = -\frac{1}{2}\hat{c}_{Y4}$ .

## List of minimal $\chi^2$ values

Scenario	$\chi_{\min}^2$	# obs.	# pars.	central values ( $\delta_{cb}^\tau, \Delta_{cb}^\tau$ )
$R(D^{(*)})$ only				
SM	23.1	2	0	—
S1	0	2	4	$(0.2 + 0.7i, 10.0 - 6.3i)$
S1 real	0	2	2	$(0.4, -3.6)$
$g_L^{cb\tau}$	0	2	2	$g_L^{cb\tau} = -1.3 - 0.6i$
$g_R^{cb\tau}$	9.1	2	2	$g_R^{cb\tau} = 0.3 + 0.i$
$g_{V_L}$	0.2	2	1	$ g_{V_L}  = 1.12$
$R(D^{(*)}), d\Gamma/dq^2, \Gamma_{B_c}$				
SM	65.9	61	4	—
S1	49.2	61	8	$(0.4 + 0.i, -2.4 + 0.i)$
S1 real	49.2	61	6	$(0.4, -2.4)$
$g_L^{cb\tau}$	55.4	61	6	$g_L^{cb\tau} = -0.4 + 0.8i$
$g_R^{cb\tau}$	55.4	61	6	$g_R^{cb\tau} = 0.3 + 0.i$
$g_{V_L}$	42.4	61	5	$ g_{V_L}  = 1.12$
$R(D^{(*)}), d\Gamma/dq^2, \Gamma_{B_c}, R(X_c)$				
SM	65.9	62	4	—
S1	50.4	62	8	$(0.3 + 0.i, -2.4 + 0.i)$
S1 real	50.4	62	6	$(0.3, -2.4)$
$g_L^{cb\tau}$	55.4	62	6	$g_L^{cb\tau} = -0.4 - 0.8i$
$g_R^{cb\tau}$	56.1	62	6	$g_R^{cb\tau} = 0.2 + 0.i$
$g_{V_L}$	46.7	62	5	$ g_{V_L}  = 1.10$