## **The On-Shell Story**

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# the most complicated calculation ...



=2π











#### summary

- $\checkmark$  the matter covered here is less general than advertised
- I will mostly cover asymmetries and branching fractions of exclusive decays (honorary mention of some inclusive channels too...)
- $\checkmark$  my survey of literature will be limited
- $\checkmark$  my survey of numerical results will be limited
- ✓ I will shamelessly promote HEPfit and how we can extract hadronic uncertainties from data
- ✓ I will also promote flavio and send the bill to the developer later
- ✓ I will discuss mostly SM with honorary mentions of extension to BSM
- ✓ I do not qualify as an expert in QCD, yet...

#### some details...

$$Q_{7} = \frac{e}{8\pi^{2}} \left[ m_{b} \bar{s} \sigma_{\mu\nu} (1+\gamma_{5}) b + m_{s} \bar{s} \sigma_{\mu\nu} (1-\gamma_{5}) b \right] F^{\mu\nu} \equiv Q_{7}^{L} + \frac{m_{s}}{m_{b}} Q_{7}^{R}$$
$$a_{7L}^{c,u} = C_{7} + O(\alpha_{s}, 1/m_{b}) \qquad a_{7R}^{c} = \frac{m_{s}}{m_{b}} C_{7} + O\left(\frac{1}{m_{b}}, \frac{\alpha_{s}}{m_{b}}\right)$$

 $\checkmark$  the leading order in expansion is factorized and a pure short distance contribution

- $\checkmark$  the contributions at  $1/m_b$  do not factorize and are commonly classified as power correction estimated using various techniques
- ✓ the right handed contribution from the SM is helicity suppressed but is the essential ingredient for the non-zero value of certain observables like the mixing induced photon polarization asymmetries.
- $\checkmark$  a priori, there is no reason why this suppression should be manifest in new dynamics

$$\bar{\mathcal{A}}_{L(R)} \equiv \frac{G_F}{\sqrt{2}} \left( \lambda_u a_{7L(R)}^u + \lambda_c a_{7L(R)}^c \right) \langle \bar{K}^* \gamma_{L(R)} | Q_7^{L(R)} | \bar{B} \rangle$$
$$\boxed{\lambda_u \sim \lambda^4} \left[ \lambda_c \sim \lambda^2 \right]$$
$$\langle \bar{K}^*(p,\eta) \gamma_{L(R)}(q,e) | Q_7^{L(R)} | \bar{B} \rangle \equiv -\frac{e}{2\pi^2} m_b T_1^{B \to K^*}(0) S_{L(R)}$$

### in addition to branching fractions...

$$A_{\rm CP}(B_q(t) \to V\gamma) = \frac{\Gamma(\bar{B}_q(t) \to \bar{V}\gamma) - \Gamma(B_q(t) \to V\gamma)}{\Gamma(\bar{B}_q(t) \to \bar{V}\gamma) + \Gamma(B_q(t) \to V\gamma)} = \frac{S(B_q \to V\gamma)\sin(\Delta M_q t) + A_{\rm CP}(B_q \to V\gamma)\cos(\Delta M_q t)}{\cosh(y_q t/\tau_{B_q}) - A_{\Delta\Gamma}(B_q \to V\gamma)\sinh(y_q t/\tau_{B_q})}$$

$$\overline{\mathrm{BR}}(B_s \to \phi \gamma) = \left[\frac{1 - A_{\Delta \Gamma}(B_s \to \phi \gamma) y_s}{1 - y_s^2}\right] \mathrm{BR}(B_s \to \phi \gamma)$$

$$\begin{split} A_{CP} &= \frac{\Gamma(\bar{B}^{0}(t) \to \bar{K}^{*0}\gamma) - \Gamma(B^{0}(t) \to K^{*0}\gamma)}{\Gamma(\bar{B}^{0}(t) \to \bar{K}^{*0}\gamma) + \Gamma(B^{0}(t) \to K^{*0}\gamma)} = S\sin(\Delta m_{B}t) - C\cos(\Delta m_{B}t) \\ \bar{A}_{L(R)} &= \mathcal{A}(\bar{B}^{0} \to \bar{K}^{*0}\gamma_{L(R)}) \qquad \mathcal{A}_{L(R)} = \mathcal{A}(B^{0} \to K^{*0}\gamma_{L(R)}) \\ S &= \frac{2\operatorname{Im}\left(\frac{q}{p}(\mathcal{A}_{L}^{*}\bar{\mathcal{A}}_{L} + \mathcal{A}_{R}^{*}\bar{\mathcal{A}}_{R})\right)}{|\mathcal{A}_{L}|^{2} + |\mathcal{A}_{R}|^{2} + |\bar{\mathcal{A}}_{L}|^{2} + |\bar{\mathcal{A}}_{R}|^{2}} \\ S^{\mathrm{SM},s_{R}} &= -\sin(2\beta) \frac{m_{s}}{m_{b}} (2 + O(\alpha_{s})) \qquad C = \frac{|\mathcal{A}_{L}|^{2} + |\mathcal{A}_{R}|^{2} - |\bar{\mathcal{A}}_{L}|^{2} - |\bar{\mathcal{A}}_{R}|^{2}}{|\mathcal{A}_{L}|^{2} + |\bar{\mathcal{A}}_{R}|^{2} + |\bar{\mathcal{A}}_{L}|^{2} + |\bar{\mathcal{A}}_{R}|^{2}} \\ \frac{q}{p} &= \sqrt{\frac{M_{12}^{*}}{M_{12}}} = e^{-2i\beta} \qquad S^{\mathrm{SM},m_{s}} = -2\sin 2\beta \frac{m_{s}}{m_{b}} = -0.0249 \pm 0.0016 \end{split}$$

### the charm loop

#### computation done with QCDF

A. Ali and A. Y. Parkhomenko, Eur. Phys. J. C 23 (2002) 89 [arXiv:hep-ph/0105302];
M. Beneke, T. Feldmann and D. Seidel, Eur. Phys. J. C 41 (2005) 173 [arXiv:hep-ph/0412400];

T. Becher, R. J. Hill and M. Neubert, Phys. Rev. D  $\mathbf{72}$  (2005) 094017 [arXiv:hep-ph/0503263].

S. W. Bosch and G. Buchalla, Nucl. Phys. B **621** (2002) 459 [arXiv:hep-ph/0106081] and JHEP **0501** (2005) 035 [arXiv:hep-ph/0408231].

#### computation done with QCD sum rules

P. Ball and R. Zwicky, Phys. Lett. B642, 478 (2006), arXiv:hep-ph/0609037 [hep-ph].
A. Khodjamirian, R. Ruckl, G. Stoll, and D. Wyler, Phys. Lett. B402, 167 (1997), arXiv:hep-ph/9702318 [hep-ph].
M. B. Voloshin, Phys. Lett. B397, 275 (1997), arXiv:hep-ph/9612483 [hep-ph].

$$S^{\text{SM},s_R} = -0.027 \pm 0.006(m_{s,b}) \pm 0.001(\sin(2\beta)) \quad S^{\text{SM},\text{soft gluons}} = -2\sin(2\beta) \left( -\frac{C_2}{C_7} \frac{L - \tilde{L}}{36m_b m_c^2 T_1^{B \to K^*}(0)} \right) = 0.005 \pm 0.015 \pm 0.0015 \pm 0.00$$

 $S^{\text{SM}} = S^{\text{SM}, s_R} + S^{\text{SM}, \text{soft gluons}} = -0.022 \pm 0.015^{+0}_{-0.01}$ 

#### computation done with pQCD sum rules on the light cone

M. Matsumori and A. I. Sanda, Phys. Rev. D73, 114022 (2006), arXiv:hep-ph/0512175 [hep-ph].

$$S^{\rm SM}_{K^*\gamma} = -(3.5\pm1.7) imes 10^{-2}$$

#### estimate using SCET

B. Grinstein and D. Pirjol, Phys. Rev. **D73**, 014013 (2006), arXiv:hep-ph/0510104 [hep-ph].

B. Grinstein, Y. Grossman, Z. Ligeti, and D. Pirjol, Phys. Rev. D71, 011504 (2005), arXiv:hep-ph/0412019 [hep-ph]

$$S_{K^*\gamma} \sim 10\%$$



### leading QCDF corrections

#### Compare to $C_7 = 0.3$ in SM



68%
95%
99%

		$B^0 \to K^* \gamma$	$B^+ \to K^* \gamma$	$B_s \to \phi \gamma$
Vertex correct	tions	-(7.8)	$\pm 1.0) - (1.1 =$	$\pm 0.3)i$
Spectator scatteri	$\log Q_{1-6}$	-0.7 - 1.3i	-0.7-1.3i	-0.7 - 1.7i
Spectator scatter	$ring Q_8$	-0.3	-0.3	-0.4
Weak annihila	tion	-0.4	+0.9	-0.5

A. Paul, D.M. Straub, "Constraints on new physics from radiative B decays", JHEP 1704 (2017) 027 [arXiv:1608.02556].

Contributions to  $\Delta C_7$  in units of  $10^{-2}$  for the three decays.

#### sub-leading non-factorizable corrections

Compare to  $C_7 = 0.3$  in SM



### the form factors

#### LCSR:

$T_1(0) =$	$= 0.282 \pm 0.00$	0.031
$T_1(0) =$	= $0.309 \pm$	0.027

#### LCSR + Lattice:

 $T_1(0) = 0.312 \pm 0.027$  $T_1(0) = 0.299 \pm 0.012$  for  $B \to K^* \gamma$ , for  $B_s \to \phi \gamma$ ,

for  $B \to K^* \gamma$ , for  $B_s \to \phi \gamma$ .

#### fit to experimental data:

$T_1(0) = 0.316^{+0.016}_{-0.015}$	for $B \to K^* \gamma$ ,
$T_1(0) = 0.280^{+0.020}_{-0.022}$	for $B_s \to \phi \gamma$ .



A. Paul, D.M. Straub, "Constraints on new physics from radiative B decays", JHEP 1704 (2017) 027 [arXiv:1608.02556].

#### some experimental numbers

 $\begin{array}{ll} 10^{4} \times {\rm BR}(B \to X_{s}\gamma)_{E_{\gamma}>1.6\,{\rm GeV}} & 3.27 \pm 0.14 \\ 10^{5} \times {\rm BR}(B^{+} \to K^{*}\gamma) & 4.21 \pm 0.18 \\ 10^{5} \times {\rm BR}(B^{0} \to K^{*}\gamma) & 4.33 \pm 0.15 \\ 10^{5} \times \overline{{\rm BR}}(B_{s} \to \phi\gamma) & 3.5 \pm 0.4 \\ S(B^{0} \to K^{*}\gamma) & -0.16 \pm 0.22 \\ A_{\rm CP}(B^{0} \to K^{*}\gamma) & -0.002 \pm 0.015 \\ A_{\Delta\Gamma}(B_{s} \to \phi\gamma) & -1.0 \pm 0.5 \end{array}$ 

### error budget in the branching fractions





### fit to the branching fractions measurements



x-axis shows experimental mean +/- 5 s.d.

### error budget in the CP asymmetries





### and finally...







### questions for the future upgrade

- $\checkmark$  Can measurements of photon polarization be made at O(SM) or better?
- ✓ What can be done that Belle II cannot do, maybe baryonic modes?
- $\checkmark$  If NP is not yet found, what can be done to constrain parametrizations of long distance effects?

## Thank you...!!



To my Mother and Father, who showed me what I could do,

and to Ikaros, who showed me what I could not.

"To know what no one else does, what a pleasure it can be!"

– adopted from the words of

Eugene Wigner.

