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e Cultura



Perspectives of direct Detection of supersymmetric Dark Matter in the MSSM and NMSSM



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based on : ArXive: 1703.01255

The MSSM and NMSSM

Parameters

MSSM

$$\begin{aligned}m_0 \\ m_{1/2} \\ A_0 \\ \mu \\ \tan \beta = \frac{v_u}{v_d}\end{aligned}$$

+

NMSSM

$$\lambda H_u H_d S + \frac{\kappa}{3} S^3 + \text{soft terms}$$

$$\lambda$$

$$\kappa$$

$$A_\lambda$$

$$A_\kappa$$

$$\mu_{eff} = \lambda < s >$$

Higgs mass

$$M_H^2 \approx M_Z^2 \cos^2 2\beta + \Delta_{\tilde{t}} + \lambda^2 v^2 \sin^2 2\beta - \frac{\lambda^2}{\kappa^2} (\lambda - \kappa \sin 2\beta)^2$$

Tree level (MSSM)

Rad Corr

NMSSM new terms

Choice of parameters

Full set of parameters (9) $m_0, m_{1/2}, A_0, \tan \beta, \lambda, \kappa, A_\lambda, A_\kappa, \mu_{eff}$

Reduced set of parameters (3) m_{H_1} or $m_{H_2}, m_{A_1}, m_{A_2} \approx m_{H_3} \approx m_{H^\pm}$

To fulfill all constraints including:

- The light Higgs mass of 125 GeV with correct couplings,
- Dark Matter abundance,
- LHC limits, etc

One gets Two scenarios:

I Large lambda and kappa and small tan beta

II Small lambda and kappa and larger tan beta

NMSSMTools 4.6.0
U. Ellwanger et al,
micrOMEGAs3.1
A. Pukhov et al.

- For both scenarios one can have either $H_1=H_{SM}$ or $H_2=H_{SM}$
- In both scenarios the turning point for either a singlino or higgsino-dominated LSP is around $2\kappa/\lambda = 1$

Dark Matter Content

Neutralino mass matrix

MSSM

$$\mathcal{M} = \begin{pmatrix} M_1 & 0 & -\frac{g_1 v_d}{\sqrt{2}} & \frac{g_1 v_u}{\sqrt{2}} \\ 0 & M_2 & \frac{g_2 v_d}{\sqrt{2}} & -\frac{g_2 v_u}{\sqrt{2}} \\ -\frac{g_1 v_d}{\sqrt{2}} & \frac{g_2 v_d}{\sqrt{2}} & 0 & -\mu \\ \frac{g_1 v_u}{\sqrt{2}} & -\frac{g_2 v_u}{\sqrt{2}} & -\mu & 0 \end{pmatrix}$$

NMSSM

$$\mathcal{M}_0 = \begin{pmatrix} M_1 & 0 & -\frac{g_1 v_d}{\sqrt{2}} & \frac{g_1 v_u}{\sqrt{2}} & 0 \\ 0 & M_2 & \frac{g_2 v_d}{\sqrt{2}} & -\frac{g_2 v_u}{\sqrt{2}} & 0 \\ -\frac{g_1 v_d}{\sqrt{2}} & \frac{g_2 v_d}{\sqrt{2}} & 0 & -\mu_{\text{eff}} & -\lambda v_u \\ \frac{g_1 v_u}{\sqrt{2}} & -\frac{g_2 v_u}{\sqrt{2}} & -\mu_{\text{eff}} & 0 & -\lambda v_d \\ 0 & 0 & -\lambda v_u & -\lambda v_d & 2\kappa s \end{pmatrix}$$

LSP=Dark Matter

$$\tilde{\chi}_1^0 = N_{1,1} |\tilde{B}\rangle + N_{1,2} |\tilde{W}^0\rangle + N_{1,3} |\tilde{H}_1^0\rangle + N_{1,4} |\tilde{H}_2^0\rangle + N_{1,5} |\tilde{S}\rangle$$

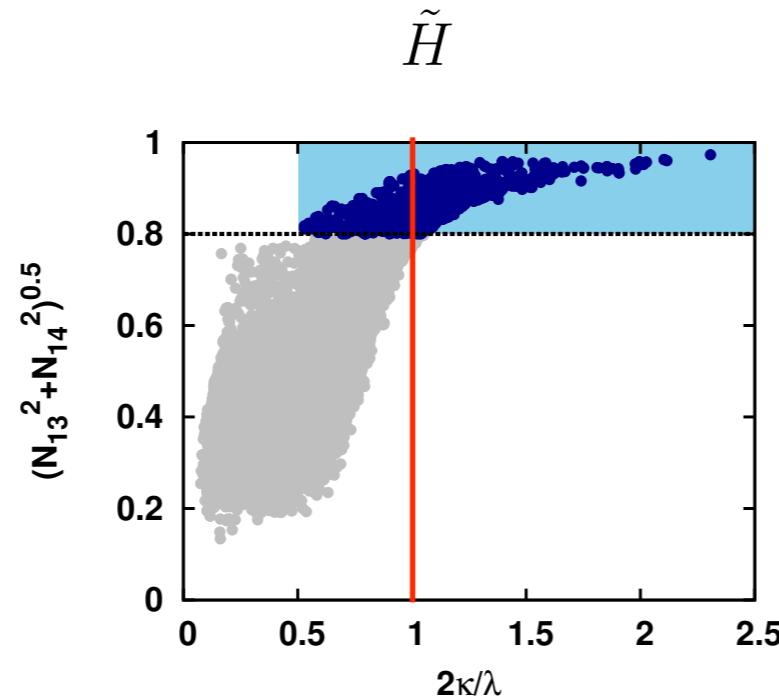
gaugino

higgsino

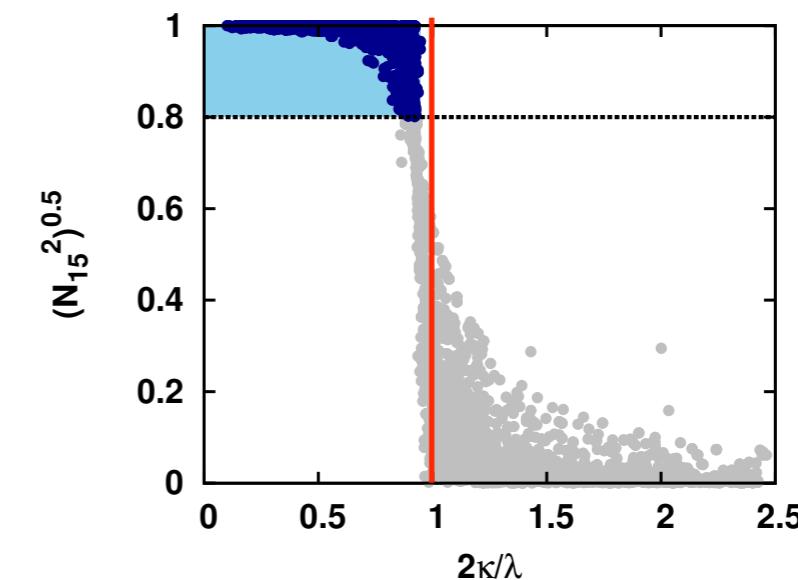
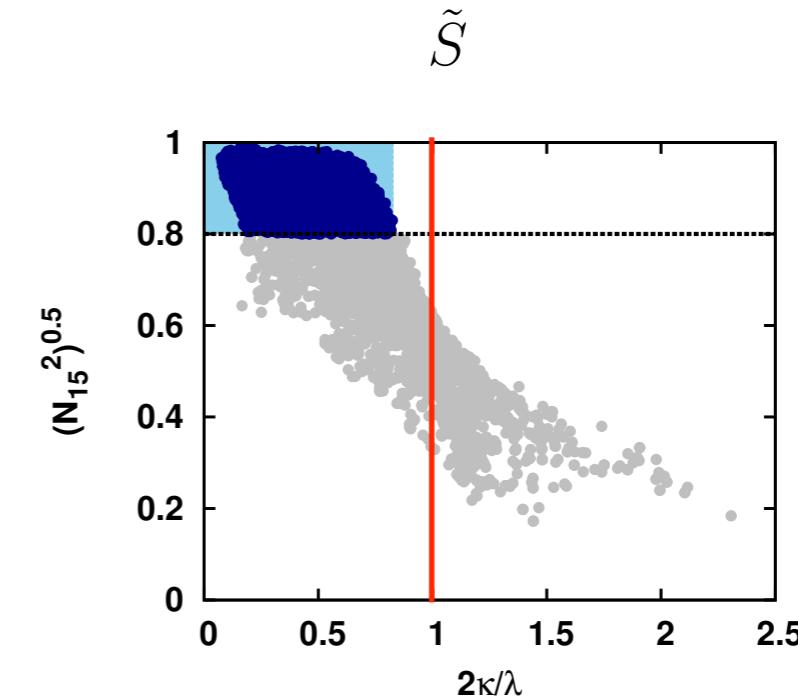
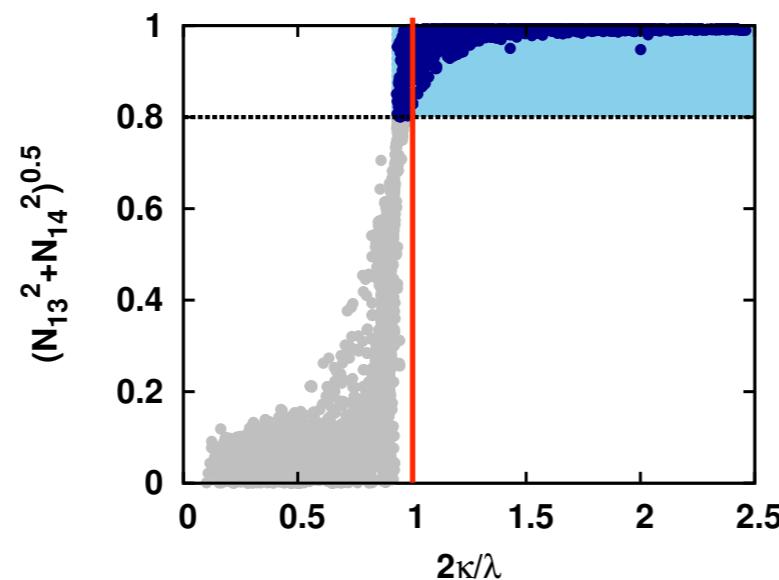
singlino

LSP Content in NMSSM

Scenario I
(large λ, κ ,
small $\tan \beta$)



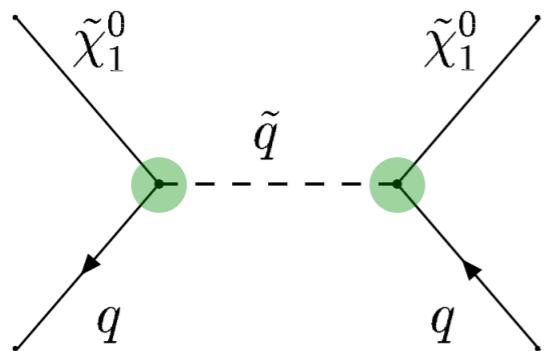
Scenario II
(small λ, κ ,
large $\tan \beta$)



Elastic WIMP-Nucleon Scattering

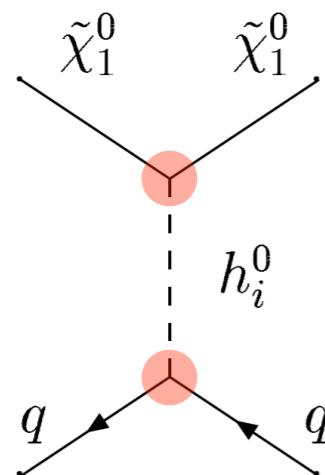
Spin-independent cross section

- Contributions from **squark-** and **Higgs-**exchanging diagrams:



Squark-exchange

$$\sigma_{\tilde{\chi}_1^0 - p} \propto \frac{m_r^2}{4\pi} \left(\frac{g'^2 \sin \theta}{m_{\tilde{q}}^2 - m_{\tilde{\chi}_1^0}^2} \right)^2 |N_{11}|^4$$



Higgs-exchange

It is the leading contribution, and increases when

$$\sigma_{\tilde{\chi}_1^0 - p} \propto \frac{m_r^2}{4\pi} \frac{\lambda_q^2}{m_h^4} |N_{13, 14}|^2 (g' N_{11} - g N_{12})^2$$

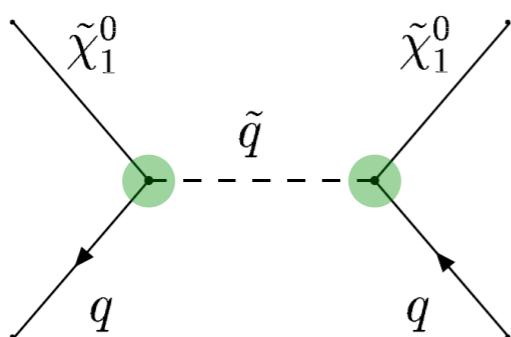
- The **Higgsino components** of the neutralino increase $\mu \downarrow$
- The **Higgs masses** decrease $m_h, m_{H^0}, m_{A^0} \downarrow$

Elastic WIMP-Nucleon Scattering

Spin-dependent cross section

Detectability

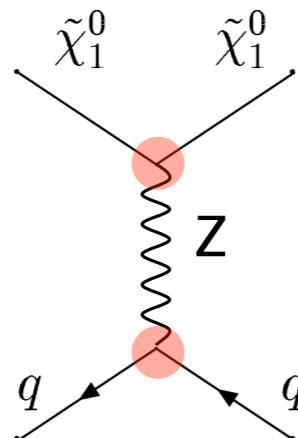
- Contributions from **squark-** and **Z-exchanging** diagrams:



Squark-exchange

$$\alpha_{2i}^{\tilde{q}} = \frac{1}{4(m_{1i}^2 - m_\chi^2)} [|Y_i|^2 + |X_i|^2] + \frac{1}{4(m_{2i}^2 - m_\chi^2)} [|V_i|^2 + |W_i|^2]$$

- Typically very small unless $m_q \sim m_\chi$



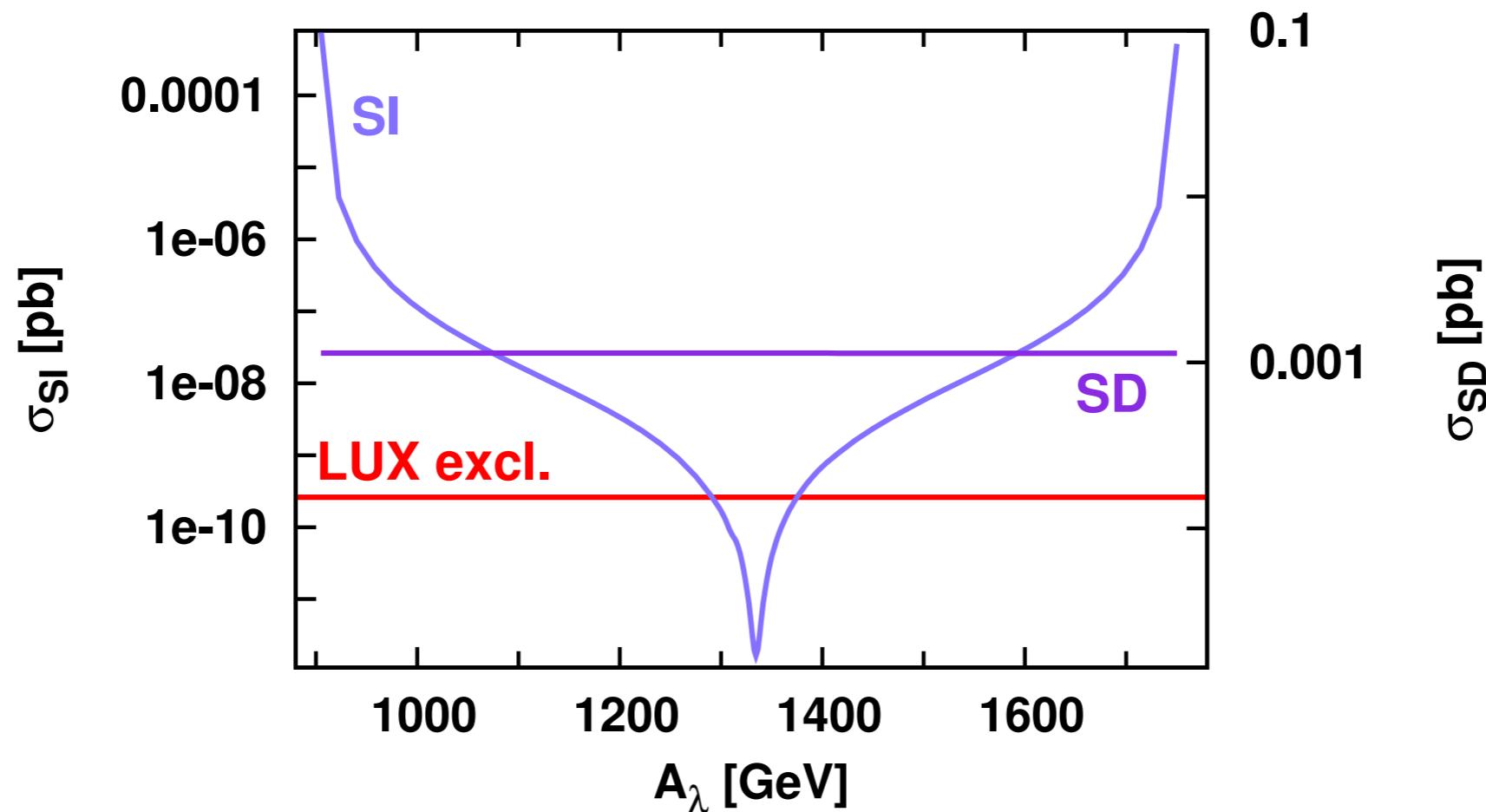
Z-exchange

$$\alpha_{2i}^Z = -\frac{g^2}{4m_Z^2 \cos^2 \theta_W} [|N_{13}|^2 - |N_{14}|^2] \frac{T_{3i}}{2}$$

Leading contribution but has an upper bound: $\sigma \leq 6.2 \times 10^{-2} \text{ pb}$

- It also increases with the neutralino **Higgsino components**: $\mu \downarrow$

SI versus SD X-sections



Negative interference of two Higgses for SI x-section $\sigma_{SI} \propto N_{13}^2 - N_{14}^2$

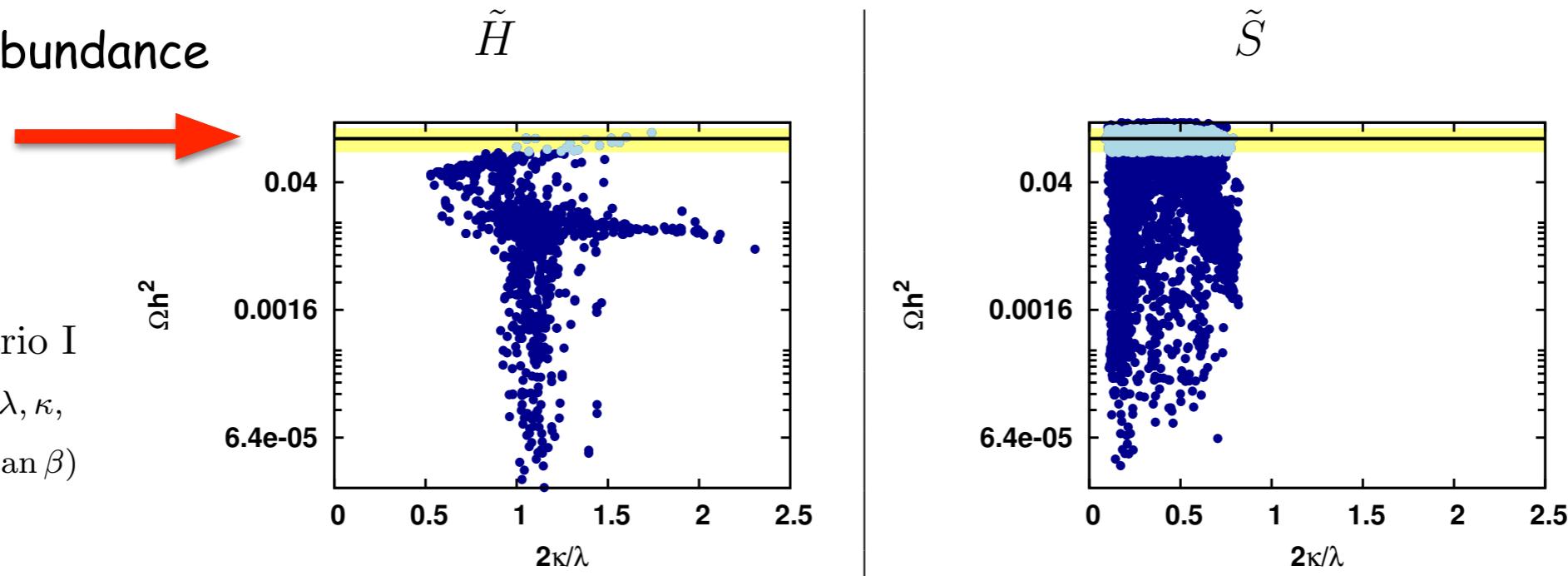
versus additive SD x-section

$$\sigma_{SD} \propto N_{13}^2 + N_{14}^2$$

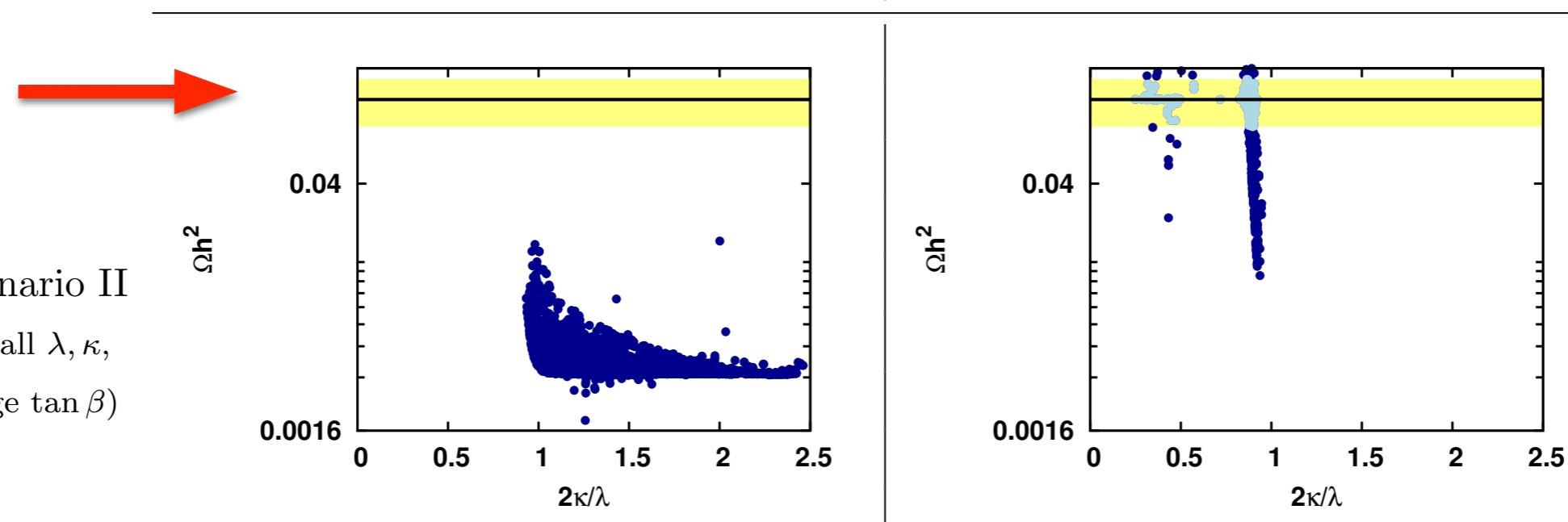
Relic Density Abundance

Exp. relic abundance

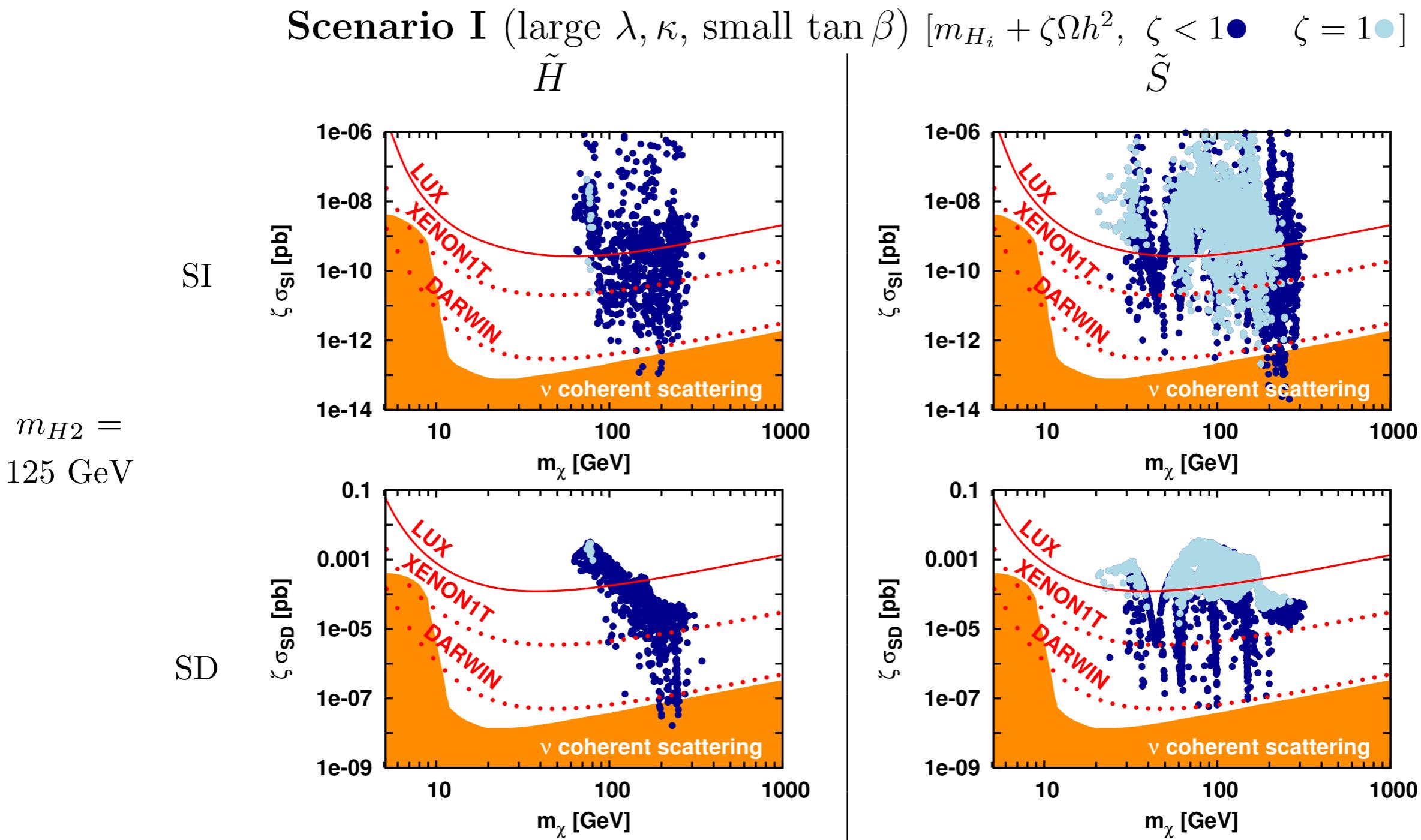
Scenario I
(large λ, κ ,
small $\tan \beta$)



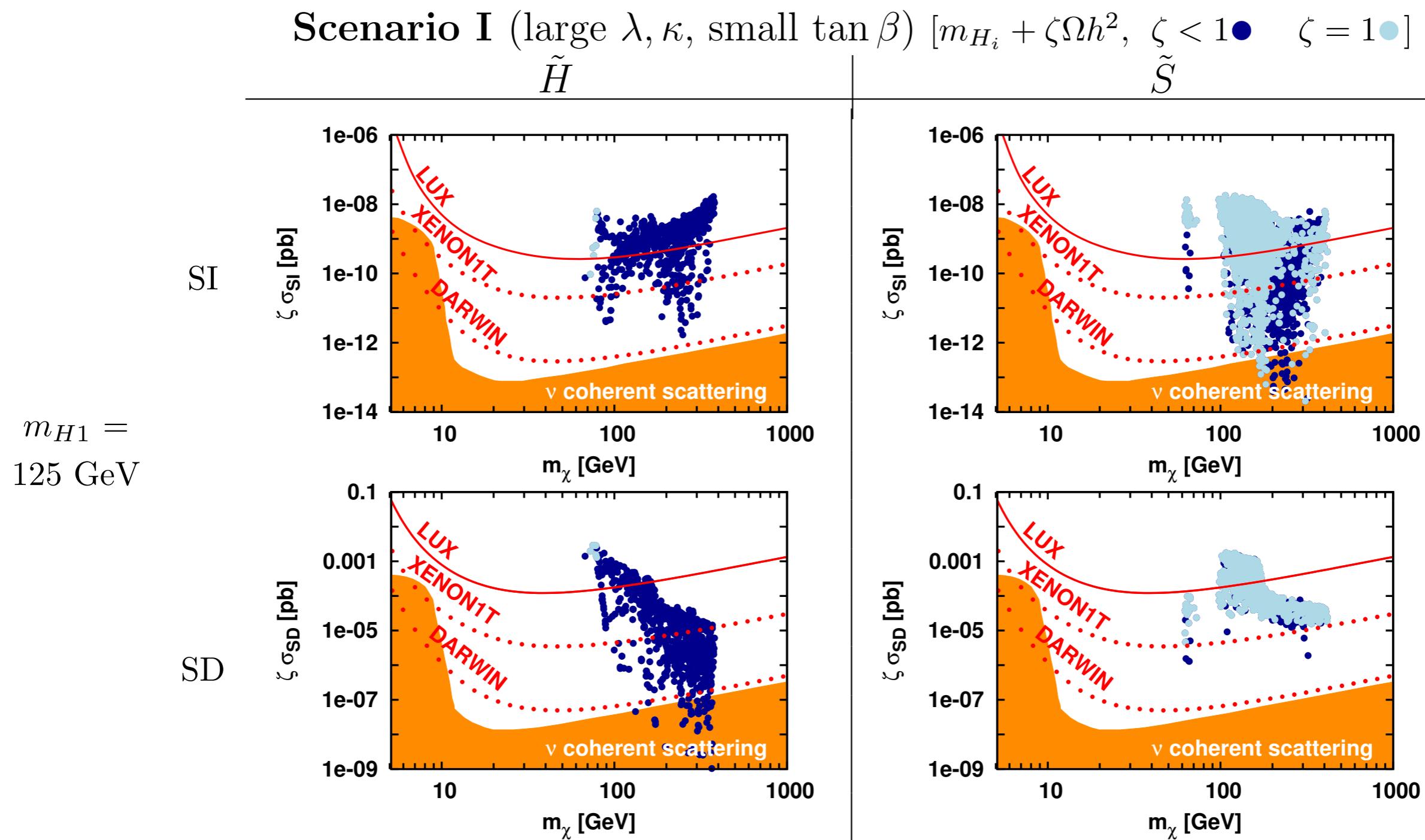
Scenario II
(small λ, κ ,
large $\tan \beta$)



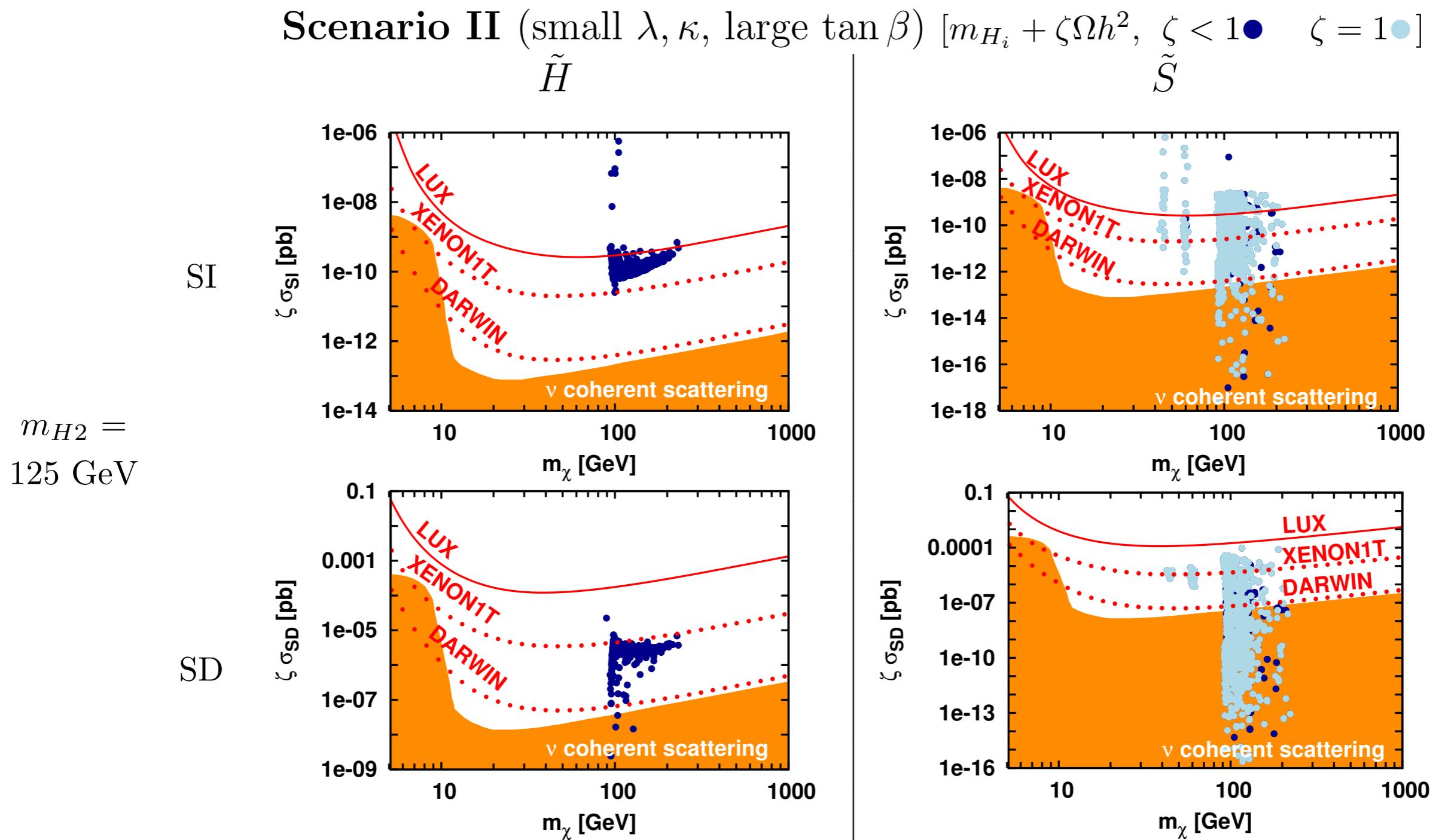
Direct Detection of Dark Matter



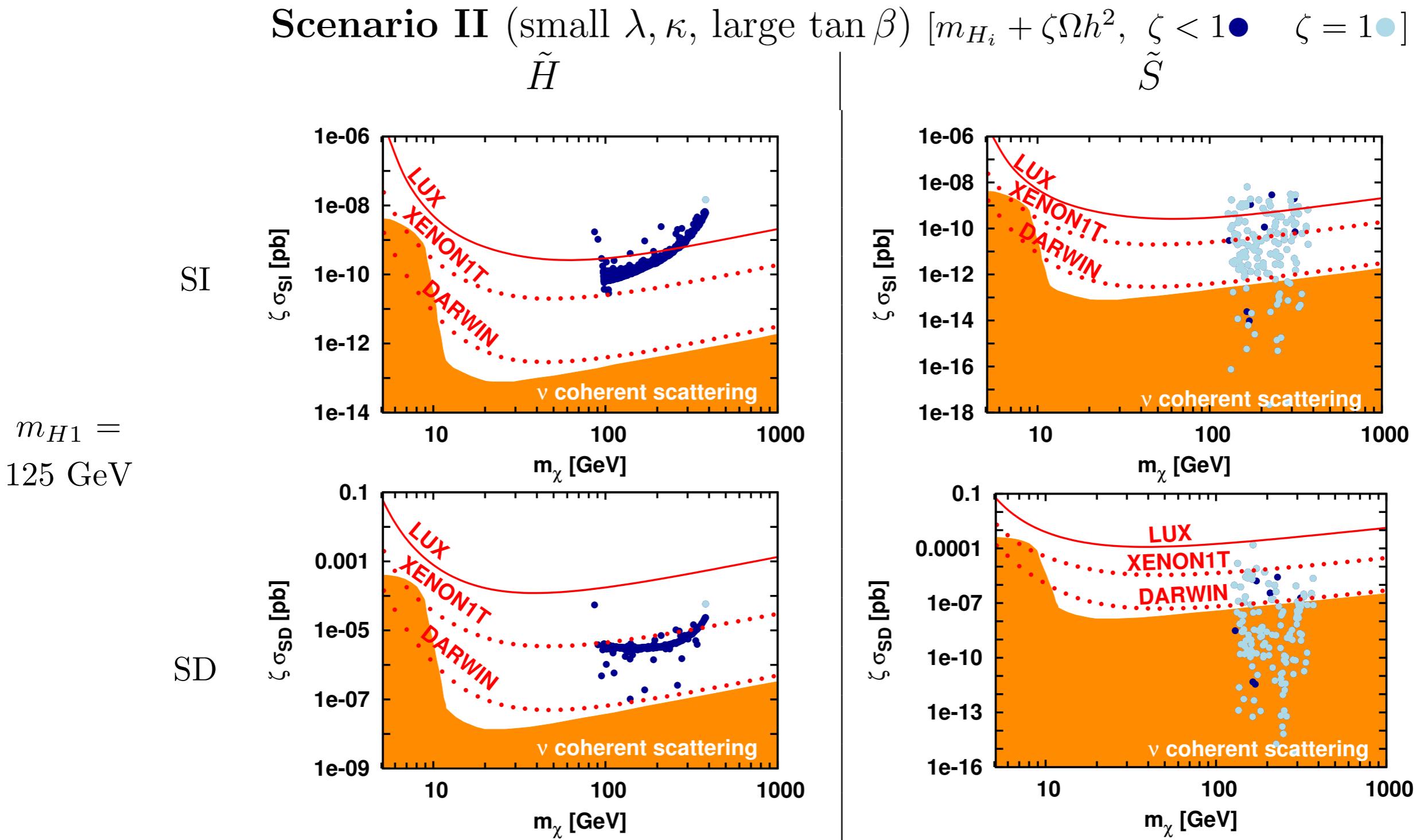
Direct Detection of Dark Matter



Direct Detection of Dark Matter

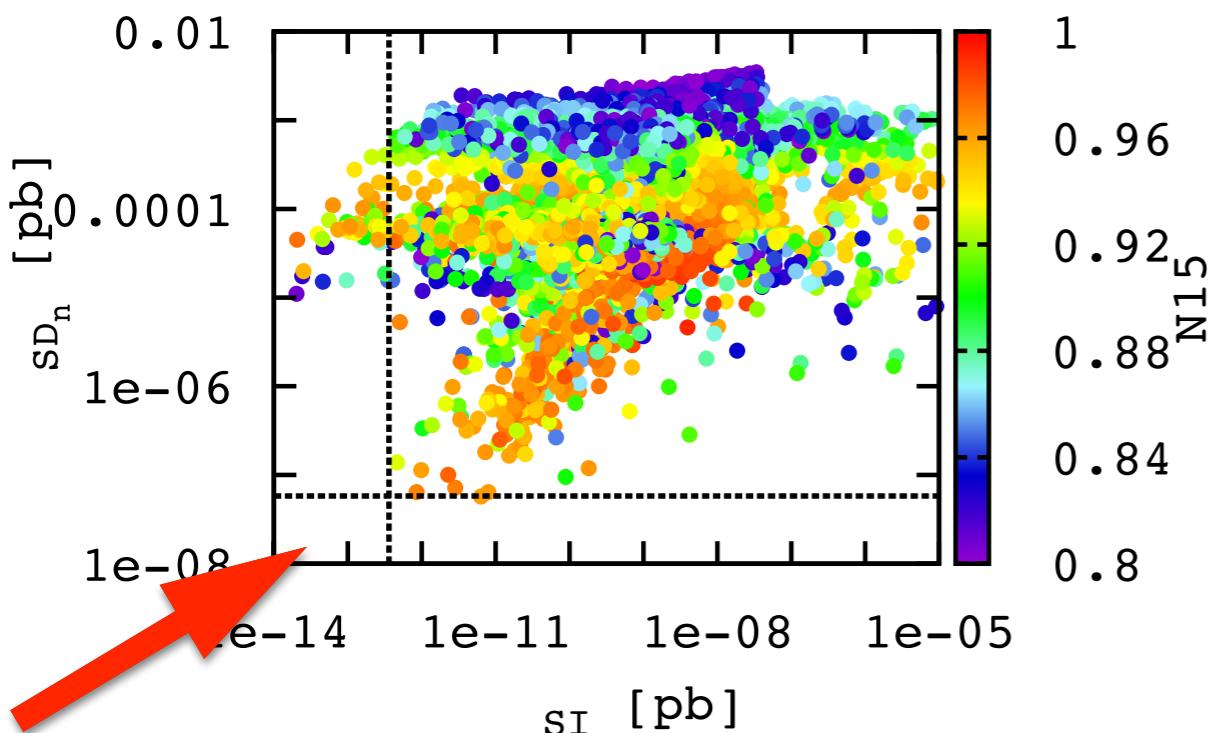


Direct Detection of Dark Matter



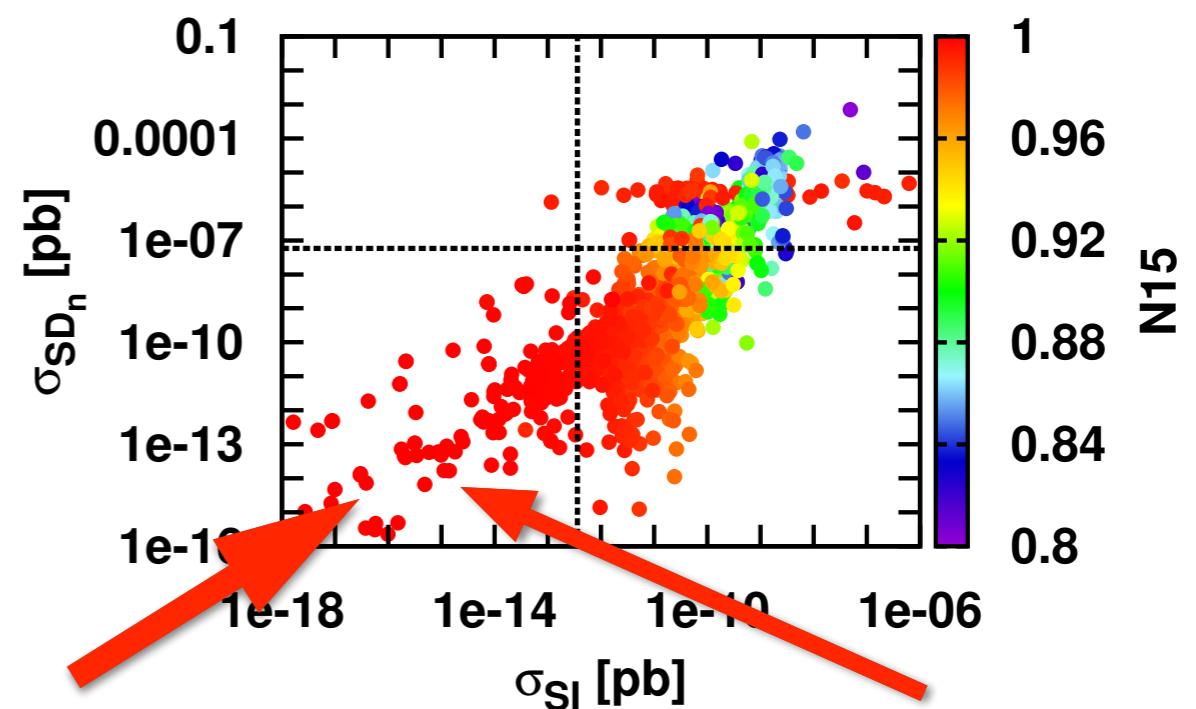
Experiment Reach in the Nearest Future

Scenario I



Will not be covered by experiment

Scenario II



Will not be covered by experiment

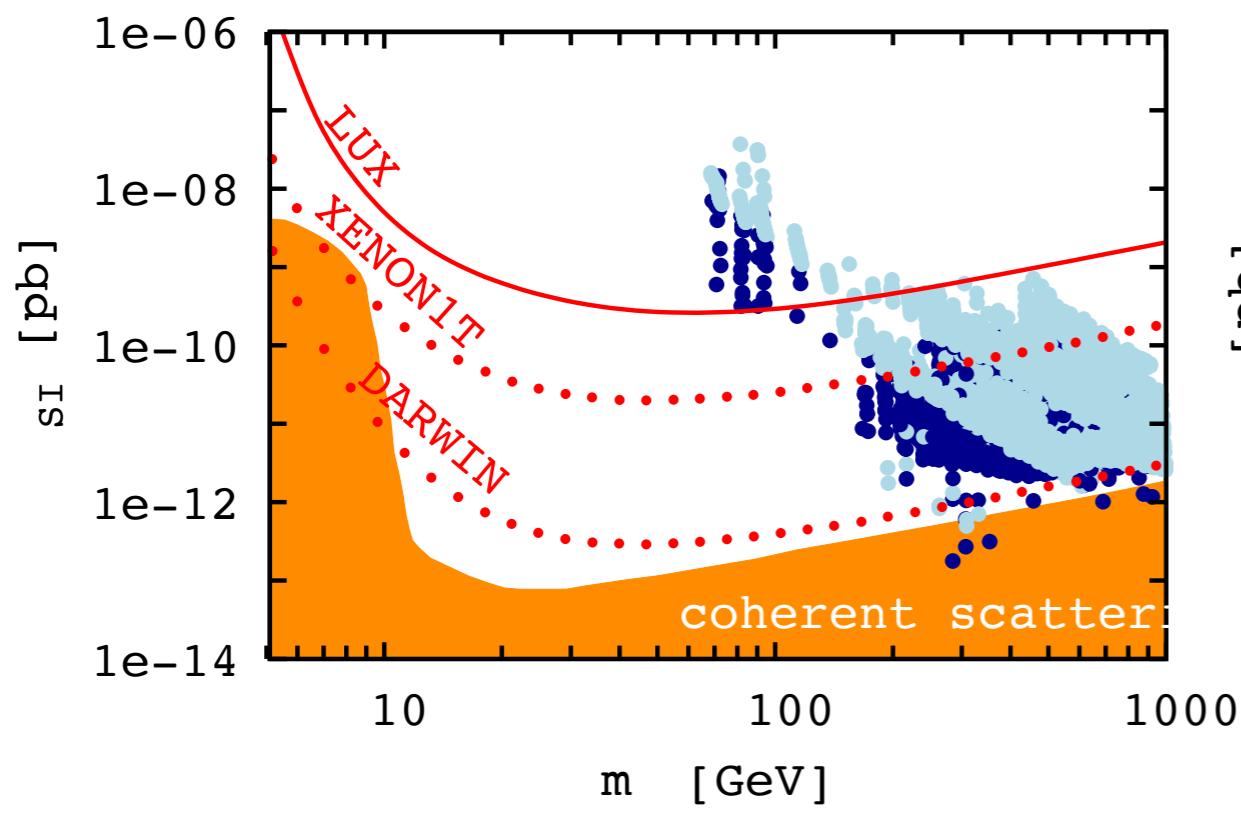
Almost pure singlino

Conclusion: in the framework of the NMSSM

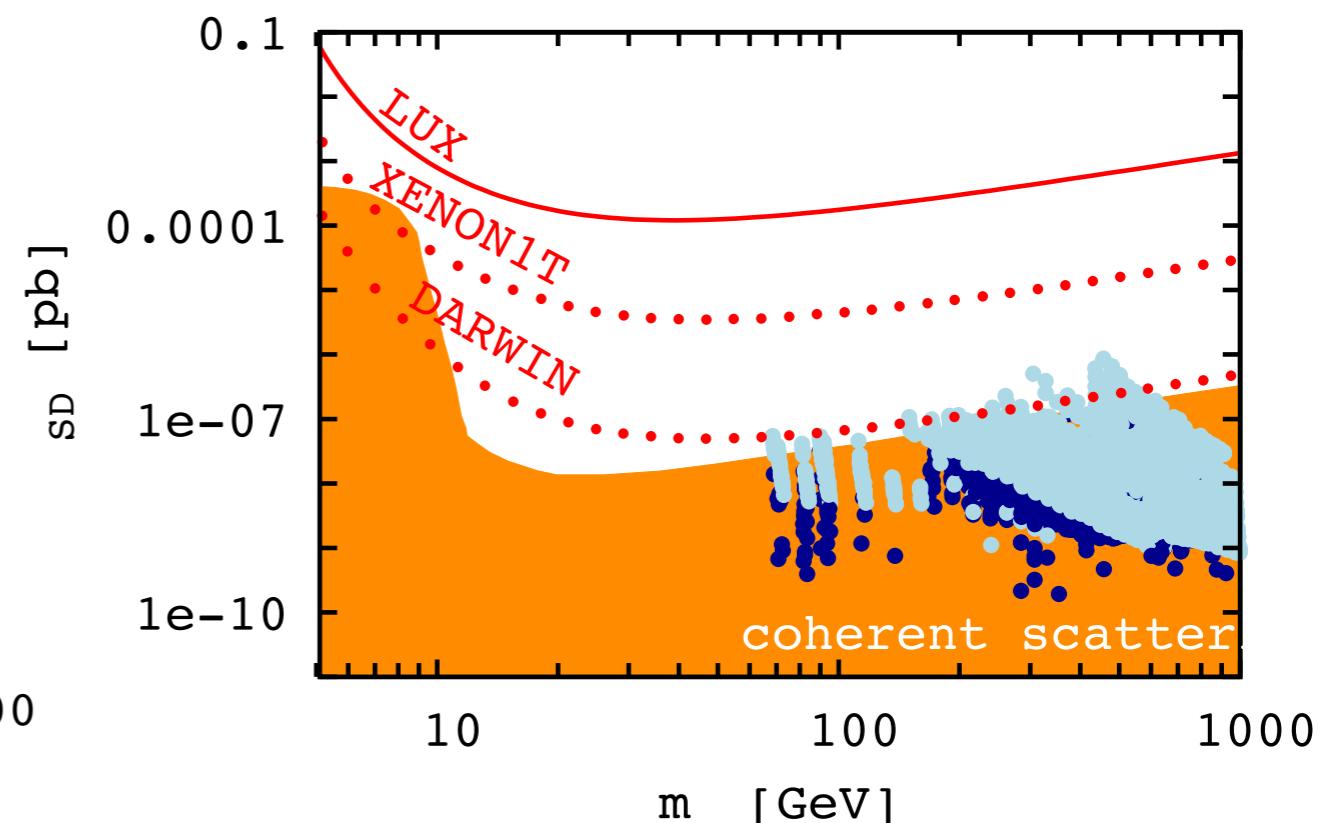
- for Scenario I future experiments will cover all allowed range
- for Scenario II future experiments will cover almost all allowed range except for small part of pure singlino LSP

Experiment Reach in the Nearest Future

CMSSM



LSP \approx Bino



Conclusion: in the framework of the MSSM
the SI searches will cover all allowed region,
the SD searches will not be essential

Conclusions

- In case of MSSM the future direct searches for DM will cover all allowed region of x-sections up to neutrino floor.
- In case of NMSSM and higgsino dominated LSP the future searches will also cover the whole range.
- In case of NMSSM and singlino dominated LSP the future searches for scenario I will cover the whole range and for scenarios II the small domain might remain which corresponds to almost $>90\%$ singlino DM.
- The SD dependent searches do not add significant information to SI searches in all cases.