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R. Davis in 80ies: "Models of neutrino masses are so numerous that they can compose the Dark Matter in the Universe"

This can not be true: with present number of models the Universe would be overclosed many times contradicting observations

Davis's remark on v models- DM connection has new turn now:

- joint models of neutrino masses and dark matter
- understanding neutrino properties can steam from the Dark sector of theory



#### What is the problem?

No theory of flavor, in general No theory of quark masses and mixing

No physics BSM has been discovered yet at LHC ... Especially no SUSY → UV completion issue.

What is the hope?

Neutrinos are the key

#### Less ambitious:

It the neutrino which shed the light on all these problems Understand at least the difference of neutrino mass and mixing from quark mixing and masses



### **3v - paradigm**

#### All well established/confirmed results fit well a framework

- three neutrinos
  with interactions described by the standard model
  with masses and mixing
- negligible feedback of neutrino mass generation on the standard model







S. Weinberg

Violation of universality, unitarity

or maybe: hHLv<sub>R</sub>H

That's all? Will we learn more?





Different lines of developments depend on



Existence of neutrino states with large mixing with active neutrinos like LSND neutrinos

Can be clarified soon

#### Dirac or Majorana

Naturally small Dirac masses → require additional symmetry, new fermions, scalars

E.Ma , 2016

May take some (a lot of) time









Strong perturbation of 3v pattern:  $m_{\alpha\beta}^{ind} \sim m_4 U_{\alpha4} U_{\beta4} \sim \sqrt{\Delta m_{32}^2}$ 

 $m_{\alpha\beta}$   $m_{4} \circ_{\alpha4} \circ_{\beta4}$   $\gamma \Delta m_{32}$ 

Effect of possible sterile neutrinos can be neglected if  $m_{\alpha\beta}^{ind} \ll \frac{1}{2} \sqrt{\Delta m_{21}^2} \sim 3 \ 10^{-3} \ eV$  $|U_{\alpha4}|^2 < 10^{-3} \ (1 \ eV/m_4)$ 

#### Large flavor mixing from steriles

Mass matrix

 $\begin{array}{c} v_{e} \\ v_{\mu} \\ v_{\tau} \\ v_{\tau} \end{array} \left( \begin{array}{cccc} m_{ee} & m_{e\mu} & m_{e\tau} & m_{eS} \\ m_{\mu\mu} & m_{\mu\tau} & m_{\muS} \\ m_{\tau\tau} & m_{\tauS} \end{array} \right) \\ m_{\tau S} \\ m_{\tau S} \end{array} \right)$ no contribution from S to  $\beta\beta_{0\nu}$  decay, but S do contribute to oscillations eV scale seesaw  $m_a + m_{ind}$  $m_{\nu}$  =  $\mathbf{m}_{a} = \begin{pmatrix} 0.2 & 0.4 & 0.4 \\ \dots & 2.8 & 2.0 \\ \dots & \dots & 3.0 \end{pmatrix} \mathbf{10^{-2} eV} \qquad \mathbf{m}_{ind} = \frac{\mathbf{m}_{SS}}{1 eV} \begin{pmatrix} 2.0 & 2.0 & 4.5 \\ \dots & 2.0 & 4.5 \\ \dots & \dots & 10.0 \end{pmatrix} \mathbf{10^{-2} eV}$ 

produce dominant Enhance lepton mixing  $\mu\tau$ -block with small determinant  $m_{eS} m_{\mu S} m_{\tau S} m_{\sigma \tau S}$  may have Generate TBM mixing certain symmetry

### Nature of Neutro Mass

is the neutrino mass of the same origins as masses of other particles?





Similar to cosmological constant

#### Smallness:

Suppression wrt. the EW scale Why there is no usual scale Dirac masses?

No RH component → Dirac mass can not be formed

symmetry

See-saw or multisinglet mechanisms - suppression only -finite contribution negligible

see-saws type-I does both things simultaneously:

incomplete suppression

#### Finite value

Mechanisms unrelated to suppression of usual Dirac masses

Seesaw type II Radiative mechanisms



Similarly for Dirac neutrinos

Hard mass related to the EW scale small effective coupling small induced VEV formed by large VEV's (seesaw II)

Soft mass

VEV created at small scales melting at T ~ VEV

MAVAN

Environment dependent masses; relic neutrinos

Gravitationally induced mass Melting couplings

#### **Soft couplings and small VEV's**



Neutrino mass generation through the condensate (crossed blue circles) via non-perturbative interaction (green circle). Small neutrino masses from gravitational O-term

G. Dvali and L. Funcke, Phys.Rev. D93 (2016) no.11, 113002 arXiv:1602.03191 [hep-ph]

No  $\beta\beta_{0\nu}$  decay due to large  $q^2$  the vertex does not exist

 $\beta\beta_{0\nu}$  decay - unique process where neutrinos are highly virtual

Certain generic features independent on specific scenario can be considered on phenomenological level

#### **Probing Nature of neutrino mass**

Determination of masses, mass squared differences from processes at different conditions

Searches for dependence of mass on external variables:

Vacuum – media with different densities, fields Solar – KamLAND: ∆m<sub>21</sub><sup>2</sup> 2-3 mixing: T2K – NOvA

Energies (in medium, or if Lorentz is violated)







M. Shaposhnikov, et al

Everything below EW scale → small Yukawa couplings

L BAU	R 0.1 - few GeV split ~ few kev	<ul> <li>Small neutrino mass</li> <li>lepton asymmetry via oscillations</li> <li>can be produced in B-decays (BR ~ 10<sup>-10</sup>) etc, SHiP</li> </ul>		
WDM	3 - 10 kev	Decouples from generation of neutrino mass, RHN?	n –	radiative decays → 3.5 keV line? Extensions of model G.K Karananas,
		Unnatural Seesaw -small Dirac Yukawas	Co	nstrained GUT's? Flavor structure, mixing?

Left-right models

With inverse seesaw

Natural realization of low scale seesaw If low scale – small Dirac Yukawa couplings Usual coupling – inverse seesaw



$$M_L = h_L V_L \quad V_L \iff V_R$$
 is required

$$m_{D} = h < \Phi > M_{R} = h_{R}V_{R}$$

$$\begin{pmatrix} 0 & m_{D}^{T} & 0 \\ m_{D} & 0 & M_{R}^{T} \\ 0 & M_{R} & \mu \end{pmatrix} V_{L} \\V_{R} \\S$$

$$m_{v}^{IS} = m_{D}^{T} M_{R}^{-1T} \mu M_{R}^{-1} m_{D}$$

$$m_{D} = m_{q}$$

$$\Rightarrow \text{ embedding in } SO_{10}$$

$$B. Dev, R Mohapatra$$
flavor symmetry in  $\mu$ 

### **v- mass and Higgs physics**

Correction to  $\lambda$  - 4 point

coupling - vacuum stability

H

Η

bottom -up

Correction to Higgs mass



Upper bound on mass M<sub>R</sub> < 10<sup>7</sup> GeV → leptogenesis ? → cancellation (a kind of SUSY)

F. Vissani ... J Elias-Miro et al, R Volkas, et al, M. Fabbrichesi ... Other contributions from particles associated to neutrino mass generation, e.g. Higgs triplets

Η

C. Bonila et al, 1506.04031

Higgs as composite state of neutrinos



New strong int. Generate 4 fermionic coupling

Recent: J. Krog, C. T. Hill 1506.02843



I. Brivio, M. Trott, 1703.10924 [hep-ph] Whole Higgs potential is generated by the neutrino corrections

Both Higgs mass term and quartic coupling (absent at tree level) are generated by neutrino loops

RH neutrino masses is the origin of the EW scale ?

 $M_R = 10^7 - 10^9 GeV$ h = 10<sup>-6</sup> - 10<sup>-4.5</sup>

Dirac Yukawa coupling



#### Symmetry behind beside the second



#### **Mixing and symmetry**

#### Real or accidental?

P. F. Harrison, D. H. Perkins, W. G. Scott L. Wolfenstein

$$U_{tbm} = U_{23}(\pi/4) U_{12}$$
  
sin<sup>2</sup> $\theta_{12} = 1/3$  0.30 - 0.31

Accidental, numerology, useful for bookkeeping

Tri-bimaximal mixing

Accidental symmetry (still useful)

 $U_{tbm} = \begin{bmatrix} \sqrt{2/3} & \sqrt{1/3} & 0 \\ -\sqrt{1/6} & \sqrt{1/3} & -\sqrt{1/2} & 0.62 \\ -\sqrt{1/6} & \sqrt{1/3} & \sqrt{1/2} & 0.62 \end{bmatrix}$ 

There is no relation of mixing with masses (mass ratios)

#### Not accidental

Lowest order approximation which corresponds to weakly broken (flavor) symmetry of the Lagrangian

with some other physics and structures associated

flavons other new particles



Sum rules... But in most of situations – just accidental, rather than follow from symmetries



Always possible, the key is that relations are simple and can be consequences of simple symmetries

#### **Relations and symmetries**

 For TBM:
  $m_{12} = -m_{13}$   $m_{22} = m_{33}$   $m_{11} + m_{13} = m_{22} + m_{23}$ 
 $S_4$ 

For Cabibbo mixing: 2x2 matrix

$$m_{12} = \frac{\sin \theta_C}{1 - 2 \sin^2 \theta_C} (m_{11} - m_{22})$$

Relation between matrix elements which leads to Cabibbo mixing independently of values of matrix elements

symmetry which produces the relation dihedral D14 *C. Hagedorn, 1204.0715* 

Difficult to reconcile with required lepton symmetry

Mixing appears as a result of different ways of the flavor symmetry breaking in the neutrino and charged lepton (Yukawa) sectors.

**Residual symmetries approa** 

No connection of masses and mixing



A<sub>4</sub> S<sub>4</sub> T<sup>'</sup>

Residual symmetries of the mass matrices

E. Ma,

C. S. Lam

Generic symmetries which do not depend on values of masses

CP-transformations can be added

Discrete finite groups Flavons to break symmetries

#### **Flavons**



Realized for arbitrary values of neutrino and charged lepton mass

- for Majorana neutrinos

in the mass basis

 $m = diag(m_1, m_2, m_3)$ 

 $G_{v}$ 

S<sub>2</sub> = diag (- 1, 1, -1)

 $S_i^2 = I \qquad Z_2 \times Z_2$ 

The Klein group

$$\begin{split} m_{l} &= diag(m_{e}, m_{\mu}, m_{\tau}) \\ \hline G_{l} \\ T &= diag(e^{i\phi_{e}}, e^{i\phi_{\mu}}, e^{i\phi_{\tau}}) \\ \phi_{\alpha} &= 2\pi k_{\alpha}/m \\ \hline T^{m} &= I \\ \Sigma \phi_{\alpha} &= 0 \end{split}$$

### Intrinsic = residual symmetries

If intrinsic symmetries are residual symmetries of the unique symmetry group (follow from breaking of unique group) → bounds on elements of mixing matrix

$$(U_{PMNS} S_i U_{PMNS}^+ T)^p = I p - integer$$

Symmetry group condition

D. Hernandez, A.S. 1204.0445

For each i the equation gives two relations between mixingparameters

Two such equations for i = 1,2 fix the mixing matrix completely  $\rightarrow$  TBM





 $G_v = Z_2$ 

Two relations

D. Hernandez, A Y S. 1304.7738 [hep-ph]

 $\begin{array}{c}
\sin^2 \theta_{12} \\
0.6 \\
0.5 \\
0.4 \\
0.3 \\
\hline \\
0.01 \\
0.02 \\
0.03 \\
0.04 \\
0.04 \\
0.05 \\
\sin^2 \theta_{13}
\end{array}$ 



In the residual symmetries approach

for column of the mixing matrix:

$$|U_{\beta i}|^{2} = |U_{\gamma i}|^{2}$$
$$|U_{\alpha i}|^{2} = \frac{1+\alpha}{4 \sin^{2} (\pi k/m)}$$

k, m, p integers which determine symmetry group

 $k_{\alpha} = 0$ 



Schemes with nonzero 1-3 mixing can be obtained

### **CP-violation, phase**

Certain values of  $\delta$  follow from the flavor symmetries without additional assumptions

Specific transformations which can be responsible (control) value of the phase

Non-trivial CP phase : from generalized CP transformations when also flavor is changed

Usually maximal CP violation is predicted



### Constrained lepton Unification



#### **Quark and Lepton Mixing**

Patterns of mixing are strongly different

Non connected

Different mechanism of generation of masses of quarks and neutrinos

e.g. in seesaw type-II

In general:

Fully Partially connected connected by symmetry: **QLC** -relations difficult to realize  $\theta_{12}^{\dagger} \sim \pi/2 - \theta_{12}^{q} = \theta_{c}$  $\theta_{23}^{|} \sim \pi/2 - \theta_{23}^{|}$ Predicted from QLC  $\theta_{13}^{\dagger} \sim \frac{1}{\sqrt{2}} \theta_C$  $U_{PMNS} = U_{CKM}^+ U_X$  $U_{CKM} = V_{CKM}$ 

Bi-maximal or TBM matrix



Normal mass ordering

 $\lambda = \sin \theta_{C}$ 



Dependence of 1-3 mixing on 2-3 mixing for different values of the phase  $\alpha$ . Allowed regions from the global fit NuFIT 2015

Allowed values of parameters of  $U_X$ Best fit value:  $\theta_{23}^{\times} = 42^{\circ}$ 

RGE effect from maximal mixing value at high scale





Quarks and leptons know about each other, Q L unification, GUT or/and Common flavor symmetries



Some additional physics is involved in the lepton sector which explains smallness of neutrino mass and difference of the quark and lepton mixing patterns



Can be naturally realized in the seesaw type I

#### **Neutrinos and Unification**

Another indication - smallness on neutrino mass

The simplest connection:

See-saw type-I

$$\begin{array}{l} \begin{array}{l} RH - neutrinos \\ V_{EW}^{2} \\ M_{R} \sim \frac{V_{EW}^{2}}{m_{v}} = 10^{8} - 10^{14} \ GeV \end{array} \end{array}$$

$$\begin{array}{l} M_{3R} \sim M_{GUT} = 10^{16} \ GeV \\ still \ possible \end{array}$$

Double seesaw connection to the Planck scale

#### Unification and difference of quark and lepton mixing patterns?

#### **Seesaw and PMNS-CKM relation**

$$\begin{array}{ccc}
\nu & N \\
\nu & 0 & m_D \\
N & m_D^T & M_R
\end{array}$$

Dirac mass terms m<sub>D</sub> = Y<H> N have large Majorana masses M<sub>R</sub> >> m<sub>D</sub> P. Minkowski H. Fritsch M. Gell-Mann, T. Yanagida P. Ramond, R. Slansky S. L. Glashow R.N. Mohapatra, G. Senjanovic

Diagonalization:

if

$$m_{v} = -m_{D} (M_{R})^{-1} m_{D}^{T}$$

$$V_{CKM^{+}}$$

$$m_{D} = m_{D}^{q}$$
should give  $U_{x}$ 

$$U_{PMNS} = V_{CKM^{+}} U_{x}$$

### **More than usual see-saw?**

**Scale of see-saw** 

$$\mathbf{M}_{\mathsf{R}} = -\mathbf{m}_{\mathsf{D}}^{\mathsf{T}} \frac{1}{\mathbf{m}_{\mathsf{v}}} \mathbf{m}_{\mathsf{D}}$$

q - 1 similarity:  $m_D \sim m_q \sim m_1$ for one third generations  $M_R \sim 2 \ 10^{14} \ GeV$ 



Flavor structure

Difficult to reproduce

Can be explained in the framework of double seesaw



R.N. Mohapatra J. Valle

Three additional singlets S which couple with RH neutrinos

$$\begin{pmatrix} 0 & \mathbf{m}_{\mathsf{D}}^{\mathsf{T}} & \mathbf{0} \\ \mathbf{m}_{\mathsf{D}} & \mathbf{0} & \mathbf{M}_{\mathsf{D}}^{\mathsf{T}} \\ \mathbf{0} & \mathbf{M}_{\mathsf{D}} & \mathbf{M}_{\mathsf{S}} \end{pmatrix} \begin{bmatrix} v \\ v^{\mathsf{c}} \\ \mathsf{S} \end{bmatrix}$$

 $M_s >> M_D$  $M_s$  - scale of B-L violation

RH neutrinos get mass via see-saw

$$\mathbf{M}_{\mathsf{R}} = \mathbf{M}_{\mathsf{D}}^{\mathsf{T}} \mathbf{M}_{\mathsf{S}}^{-1} \mathbf{M}_{\mathsf{D}}$$

This explains

1. strong mass hierarchy 
$$M_D \sim m_D$$
 and  $M_S$  has no strong hierarchy

- 2. intermediate scale of masses if  $M_S \sim M_{Pl}$ ,  $M_D \sim M_{GU}$
- 3. Flavor structure:

$$\implies m_v = m_D^T M_D^{-1T} M_S M_D^{-1} m_D$$

A.Y.S M. Lindner, M.A. Schmidt A.Y.S

# In the provide the





Neutrino are special

via the portal:

Neutrino mass - seesaw Large lepton mixing Non Standard Interactions

SM is well protected

Singlet of symmetry group of hidden sector

Connection to the Higgs portal: H<sup>+</sup>H



#### **Realization scheme**

Patrick Ludl, A.S arXiv:1507.03494 [hep-ph]





### **Tests and problems**

 $\delta_{CP}^{\prime} \sim \delta_{CP}^{\prime} q$ 

Normal mass hierarchy, first 2-3 octant

Flavons, new fermions, new higgses are at GUT - Planck scale

Nothing should be observed at LHC which is responsible for neutrino masses

Proton decay

New elements related to CKM physics Very strong hierarchy of masses of RH neutrinos: Leptogenesis with second RH neutrino





### Connections

Neutrinos

g-2)<sub>u</sub>

Any discovery in these fields can have impact on neutrinos

LIC Observables Higgs physics Dark matter 

Anomalies n B-decays

Lepton universality

#### Axions Dark energy

Model dependent, not unique



#### Is the (hot) component of the DM

Mechanism of generation of small neutrino masses is related to DM



Neutrino portal connects DM and neutrinos

DM particles participate (appear in loops) in generation of neutrino mass

The same symmetry is responsible for smallness of neutrino mass and stability of the DM

### FIN CONCLUSION











After more than 40 years of theoretical studies, thousands of papers written we are not far from the beginning: "ground zero" determined by experimental measurements

Big temptation to present such a talk as collection of jokes, if not one point

Enormous efforts in determination of matrix elements, cross-sections, systematics, backgrounds...

And all this is to measure neutrino parameters

Determination of neutrino parameters is not the end of story

We measure neutrino parameters to establish

the underlying physics. In spite of scepticism searches for true theory of mass and mixing is the must



Probably correct elements of the theory of neutrino mass and mixing are already among numerous mechanisms, schemes models the goal is to identify them

Still something important can be is missed

Feeling is that the theory may not be simple and The progress may not be easy

We discuss matrix elements, cross-sections, systematics, backgrounds....

to measure neutrino parameters

We measure neutrino parameters to establish the underlying physics. In spite of scepticism searches for true theory of mass and mixing must continue

## $\begin{array}{l} \mbox{Principles and Codes} \\ \mbox{Input} \\ \mbox{Tr} \ \mbox{OFT} \ \mbox{Principles} \end{array}$

Gauge interactions, extended gauge group

Relevant experiment al data

Additional symmetries: discrete, continuous, local global

Spontaneous violation of symmetries

New fields: fermions, bosons in various representations of symmetry groups Beyond QFStringy mechanisms, selection rules, etc.

Computer code for model building



Viable Models



quark sector relation Still possible 2-3 mixing is close to maximal but 2-3 mass splitting is large. Complete degeneracy is disfavored by cosmology

> Simple symmetries → degeneracy, massless states

 $\theta \sim \sqrt{\frac{m_2}{m_3}}$ 



Deviations - consequences of symmetry (complicated groups)  $\rightarrow$  "direct"

Deviations - violation of (simple) symmetries  $\rightarrow$  "semi-direct"

"Sum rules"

Ref. Nothing fundamental model dependent

Deviations related to mass ratios?

 $Z_2 \times Z_2$  - TBM  $Z_2$  - only one column in the mixing matrix is fixed, e.g. TBM<sub>1</sub>

#### **Quark and Lepton Mixing**

No immediate relations,

equalities rent mechanism of generation of masses of quarks and neutrinos

e.g. in seesaw type-II

Still some relations can be obtained within GUT since the same 126 contributes to quark masses

$$\theta_{12}^{\dagger} \sim \pi/2 - \theta_{12}^{\dagger} = \theta_{c}$$
  
$$\theta_{23}^{\dagger} \sim \pi/2 - \theta_{23}^{\dagger}$$

 $\theta_{13}^{\dagger} \sim \frac{1}{\sqrt{2}} \theta_C$ 

QLC -relations

Predicted from QLC

Other quark mixing angles can be involved But they give small corrections to these relations