



*SM&FT 2017 - The XVII Workshop on  
Statistical Mechanics and non-perturbative Field Theory*



Istituto Nazionale di Fisica Nucleare

**PAOLA RUGGIERO**

# ENTANGLEMENT MEASURES IN MANY-BODY QUANTUM SYSTEMS: THE NEGATIVITY SPECTRUM

with P. CALABRESE, V. ALBA

PRB 94, 195121 (2016)

December 13th 2017



## Black hole physics

Bekenstein-Hawking entropy

$$S_{BH} = \frac{A}{4G_N}$$

## Numerical algorithms

Based on MPS, tensor networks  
Ex: DMRG, iTEBD

## Phase transitions

- Probing quantum criticality
- Universal scaling in critical systems

## ENTANGLEMENT in MANY-BODY SYSTEMS

## Randomness

Entanglement is a good indicator of *localization*

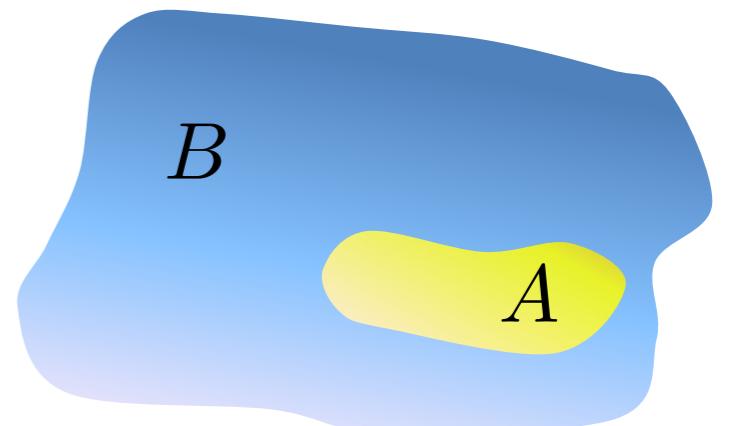
## Phases of matter

Entanglement useful to detect and characterize *topological order*

# HOW DO WE MEASURE ENTANGLEMENT?

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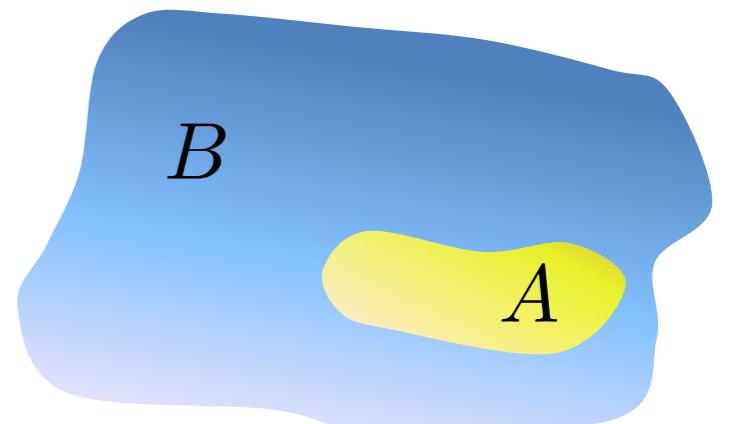


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- ▶ Entanglement entropy: Pure states

$$S_A = -\text{tr}(\rho_A \log \rho_A)$$

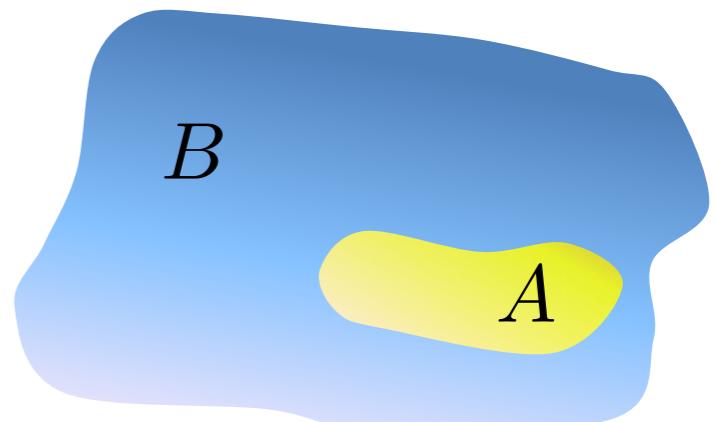


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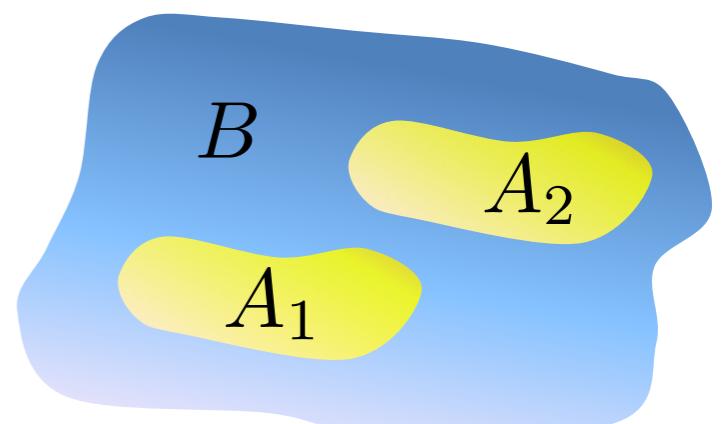


- ▶ Entanglement negativity: Mixed bipartite systems  $\longleftrightarrow$  Tripartite system

*Partial transpose:*  $\langle e_i, e_j | \rho_A^{T_2} | e_k e_l \rangle = \langle e_i, e_l | \rho_A | e_k e_j \rangle$

*Not positive definite!*

$$\mathcal{E}_{A_1, A_2} = \ln \text{Tr} |\rho_A^{T_2}|$$



# REPLICA APPROACH TO ENTANGLEMENT MEASURES

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**Entanglement entropy:**  $S_A = -\text{tr}(\rho_A \log \rho_A)$

$\text{Tr} \rho_A^n$

Evaluated for  $n \in \mathbb{N}$  and analytically continued

$$S_A = - \lim_{n \rightarrow 1} \frac{\partial}{\partial n} \text{Tr} \rho_A^n$$

[Calabrese, Cardy, 2004]

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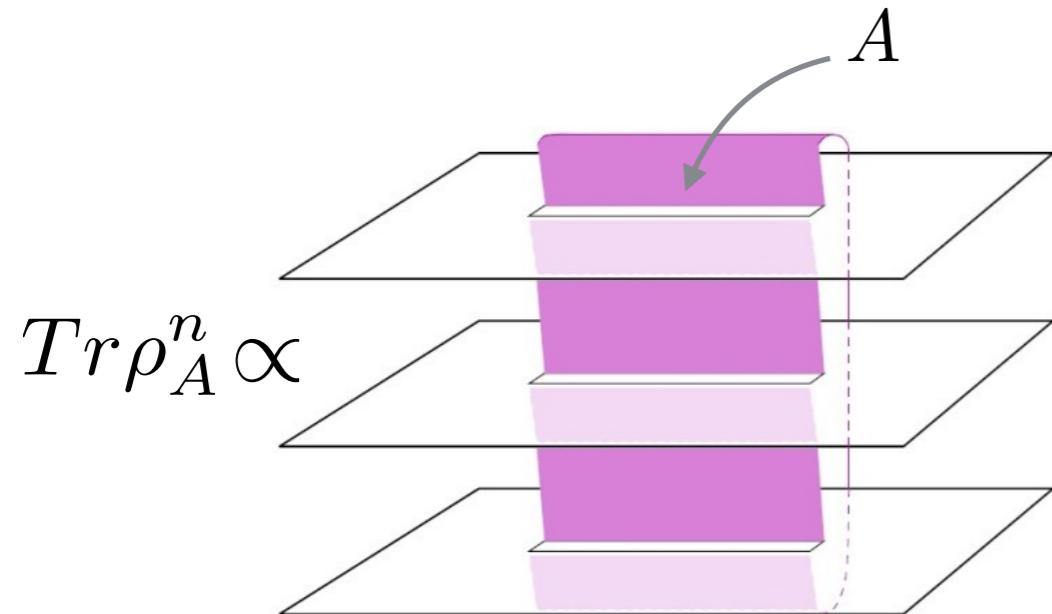


**Entanglement negativity:**  $\mathcal{E} \equiv \ln \text{Tr} |\rho^{T_2}|$

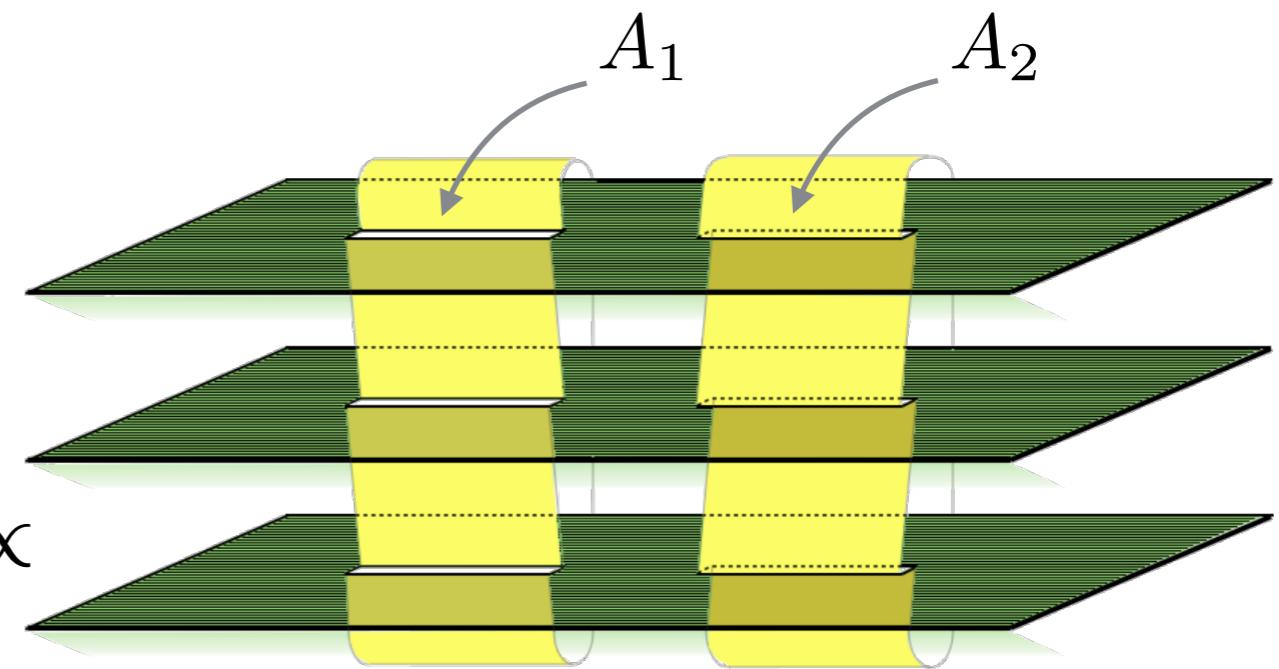
$$\begin{cases} \text{Tr}(\rho^{T_2})^{n_e} = \sum_{\lambda_i > 0} |\lambda_i|^{n_e} + \sum_{\lambda_i < 0} |\lambda_i|^{n_e} \\ \text{Tr}(\rho^{T_2})^{n_o} = \sum_{\lambda_i > 0} |\lambda_i|^{n_o} - \sum_{\lambda_i < 0} |\lambda_i|^{n_o} \end{cases} \rightarrow \mathcal{E} = \lim_{n_e \rightarrow 1} \ln \text{Tr}(\rho^{T_2})^{n_e}$$

[Calabrese, Cardy, Tonni, 2011]

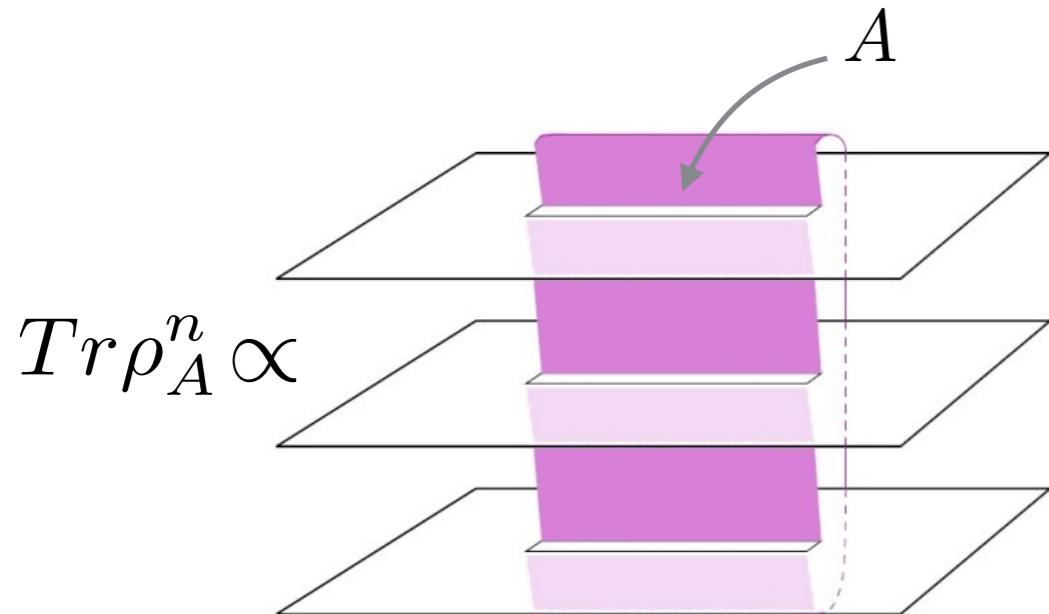
# QFT, PATH INTEGRAL FORMALISM



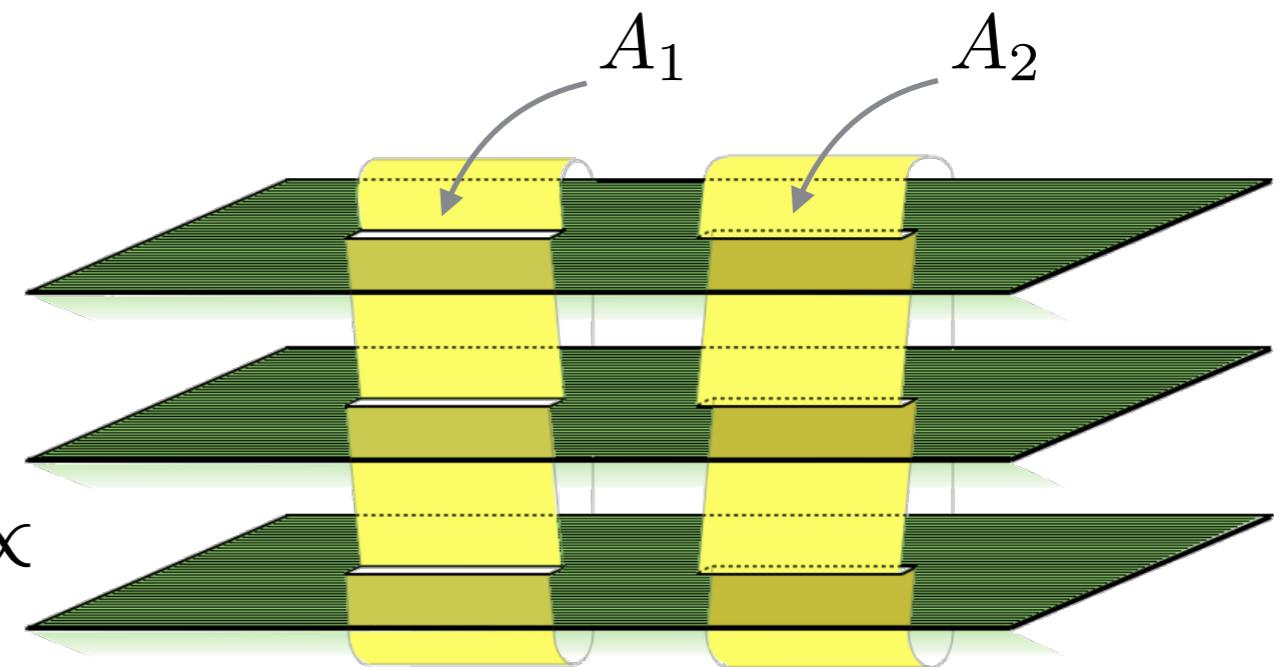
$$Tr(\rho_{A_1}^{T_2})^n \propto$$



# QFT, PATH INTEGRAL FORMALISM



$$Tr(\rho_A^{T_2})^n \propto$$



“BRANCH POINT” TWIST FIELDS:  $\mathcal{T}, \tilde{\mathcal{T}}$

$$tr \rho_A^n = \mathcal{N} \langle \mathcal{T}(u_1, 0) \tilde{\mathcal{T}}(v_1, 0) \cdots \mathcal{T}(u_m, 0) \tilde{\mathcal{T}}(v_m, 0) \rangle$$

$$tr(\rho_A^{T_2})^n = \mathcal{N} \langle \mathcal{T}(u_1, 0) \tilde{\mathcal{T}}(v_1, 0) \tilde{\mathcal{T}}(u_2, 0) \mathcal{T}(v_2, 0) \rangle$$

*Explicit results in CFTs!*

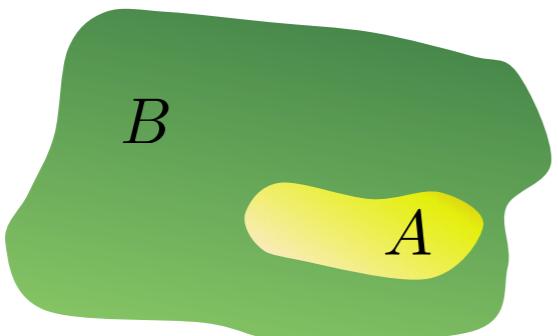
# NEGATIVITY SPECTRUM of ONE DIMENSIONAL CFTs

with V. Alba, P. Calabrese

*Phys. Rev. B 94, 195121 (2016)*

# ENTANGLEMENT & NEGATIVITY SPECTRUM

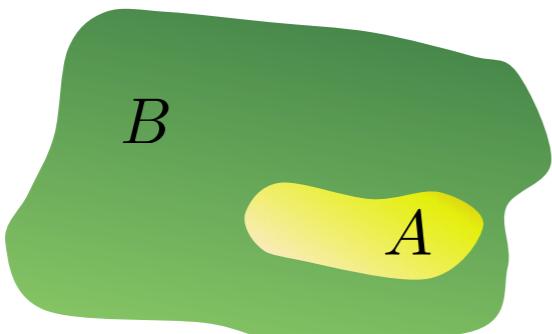
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***Reduced density matrix (RDM)***

$$\rho_A = \text{tr}_B \rho$$

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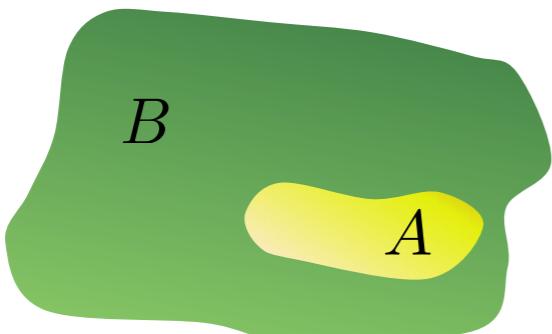
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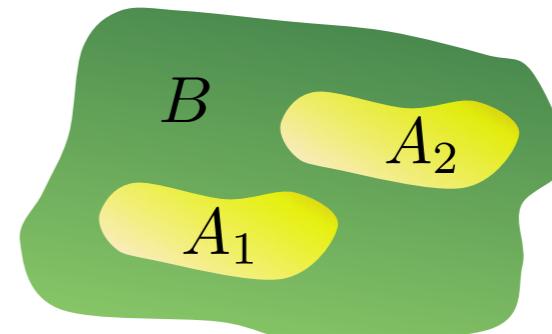
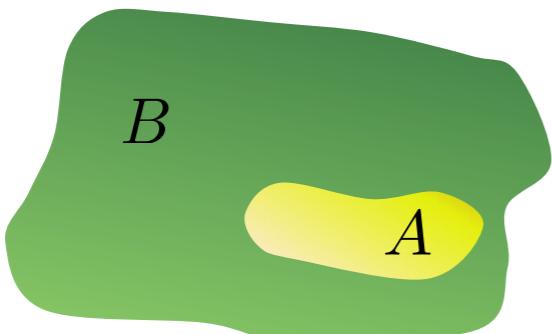
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► ENTANGLEMENT SPECTRUM:  $\{\mu_i\}$

[Li, Haldane, 2008]

- Crucial to probe topological phases of matter
- [Calabrese, Lefevre, 2008]: Analytic derivation of eigenvalues distribution

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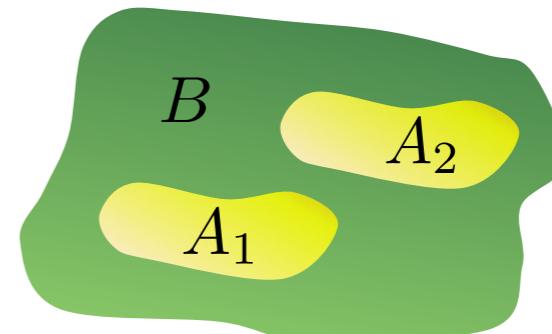
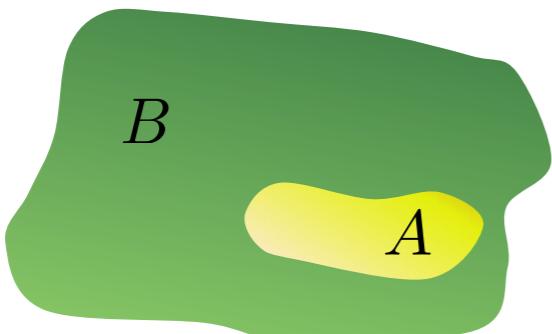
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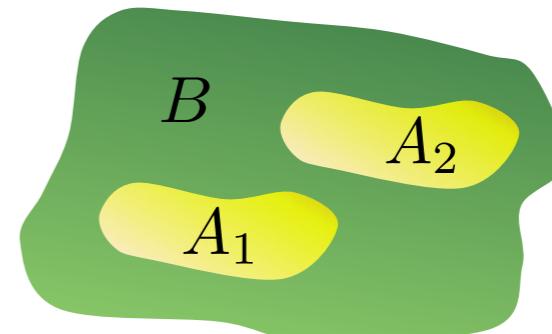
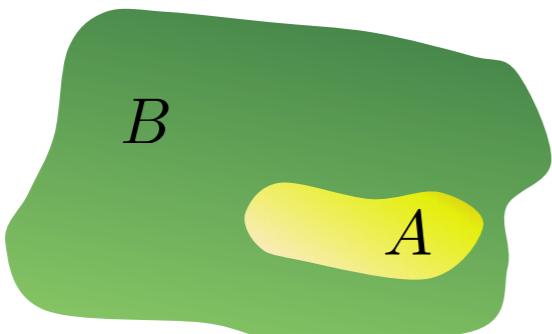
**Partial transpose of the RDM**

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► “NEGATIVITY SPECTRUM”:  $\{\lambda_i\}$

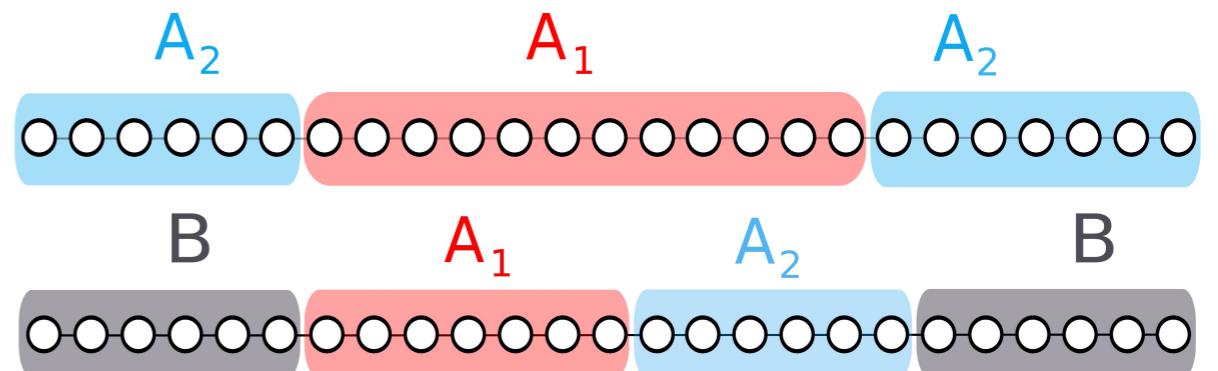
[Ruggiero, Alba, Calabrese, 2016]

- Can we spot new phases of matter? [Mbeng, Alba, Calabrese, 2016]
- Can we construct simpler entanglement measures for mixed states?

# NEGATIVITY SPECTRUM

## EIGENVALUES DISTRIBUTION

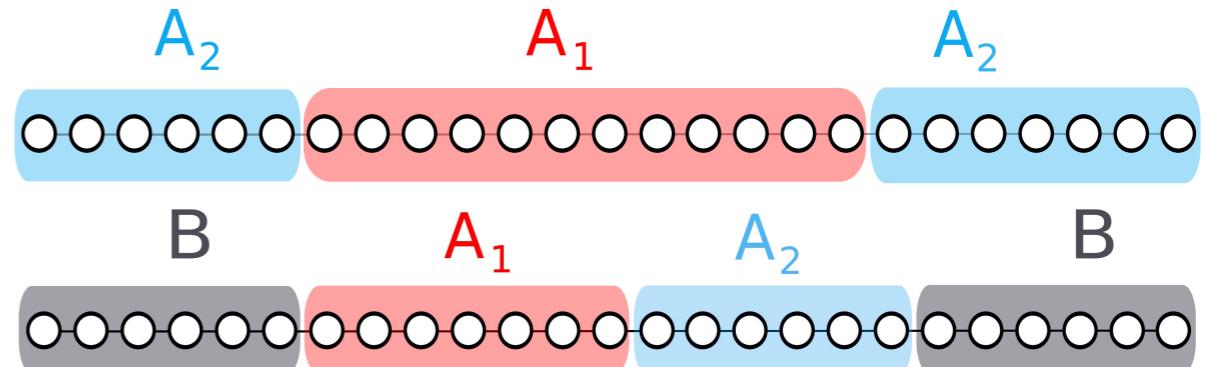
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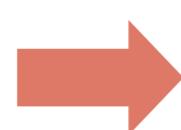
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The moment problem:

- $R_n^{T_2} \equiv Tr \left( \rho_A^{T_2} \right)^n = \sum_i \lambda_i^n = \int d\lambda \lambda^n P(\lambda)$  **(moments)**

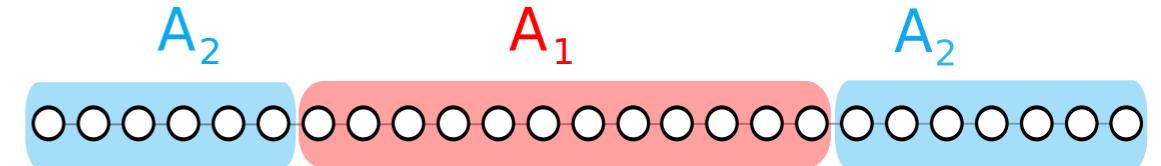
- $f(z) \equiv \frac{1}{\pi} \int d\lambda \frac{\lambda P(\lambda)}{z - \lambda} = \frac{1}{\pi} \sum_{n=1}^{\infty} R_n^{T_2} z^{-n}$  **(Stieltjes transform)**



$$P(\lambda) = \frac{1}{\lambda} \lim_{\epsilon \rightarrow 0} \text{Im} f(\lambda - i\epsilon)$$

# NEGATIVITY SPECTRUM in CFTs

Pure state:



$$R_n^{T_2} = \begin{cases} Tr \rho_{A_1}^{n_o} \\ (Tr \rho_{A_1}^{n_e/2})^2 \end{cases} \sim \begin{cases} e^{-b(n_o - \frac{1}{n_o})} & n_o \text{ odd} \\ e^{-b(n_e - \frac{4}{n_e})} & n_e \text{ even} \end{cases}$$

$$\xi \equiv \sqrt{b \ln(\lambda_M/\lambda)}$$

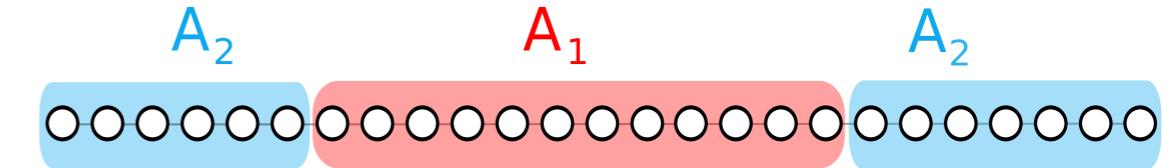
$$b \equiv -\lambda_M = \frac{c}{6} \ln \ell_1$$

$$P(\lambda) = \delta(\lambda_M - \lambda) + \frac{b\theta(\lambda_M - |\lambda|)}{|\lambda|\xi} \left[ \frac{\text{sgn}(\lambda)}{2} I_1(2\xi) + I_1(4\xi) \right]$$

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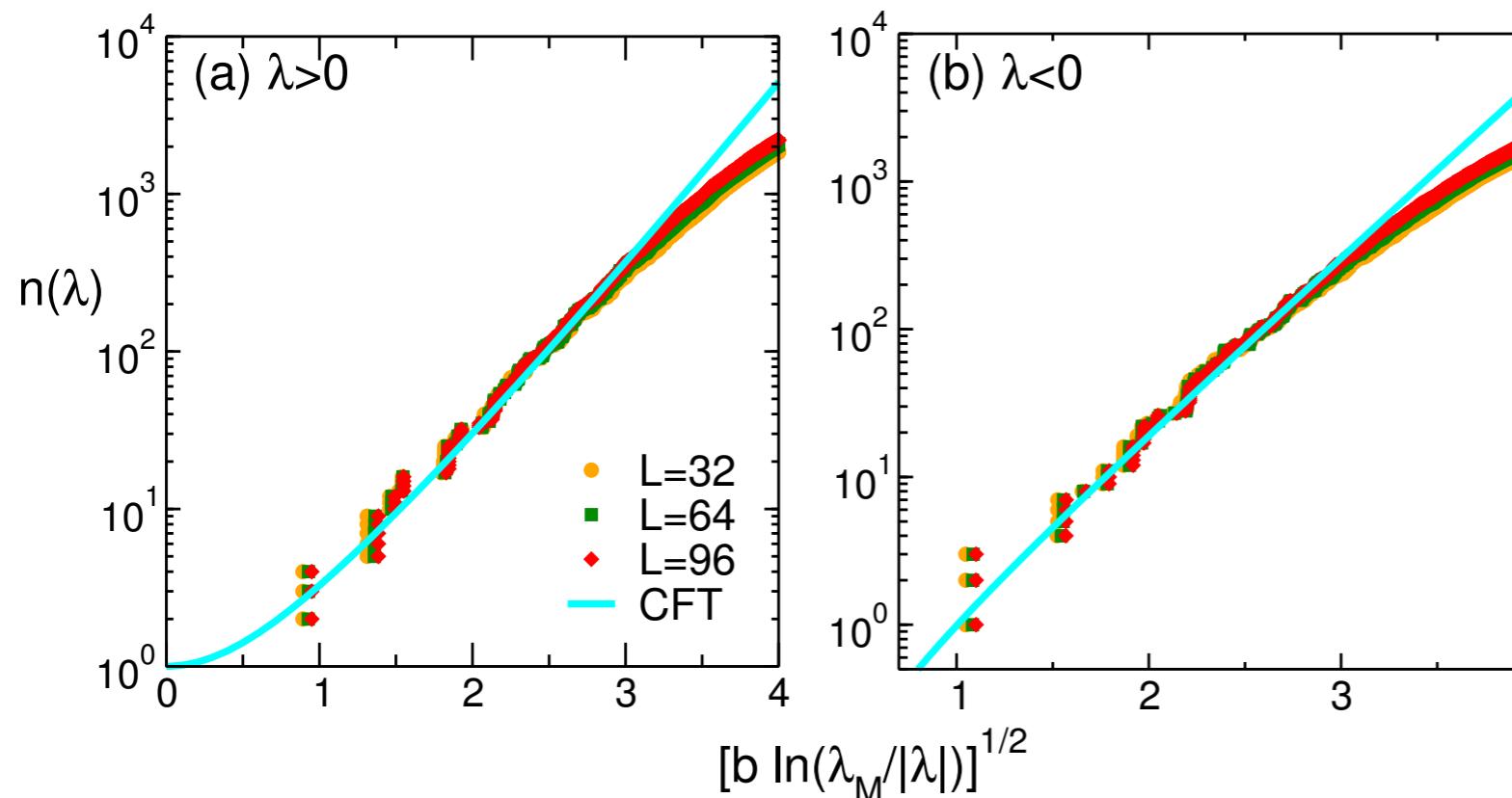
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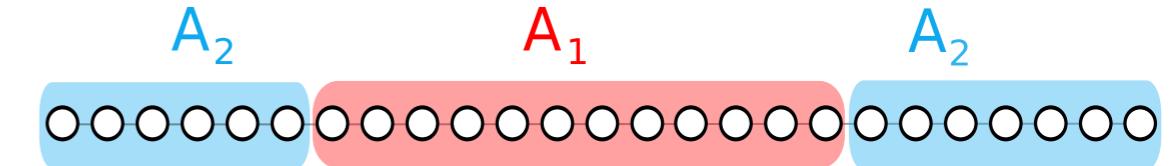
## NUMBER DISTRIBUTION FUNCTION

$$n(\lambda) \equiv \begin{cases} \int_{\lambda}^{\lambda_M} d\lambda P(\lambda) & \lambda > 0 \\ \int_{\lambda_m}^{\lambda} d\lambda P(\lambda) & \lambda < 0 \end{cases}$$

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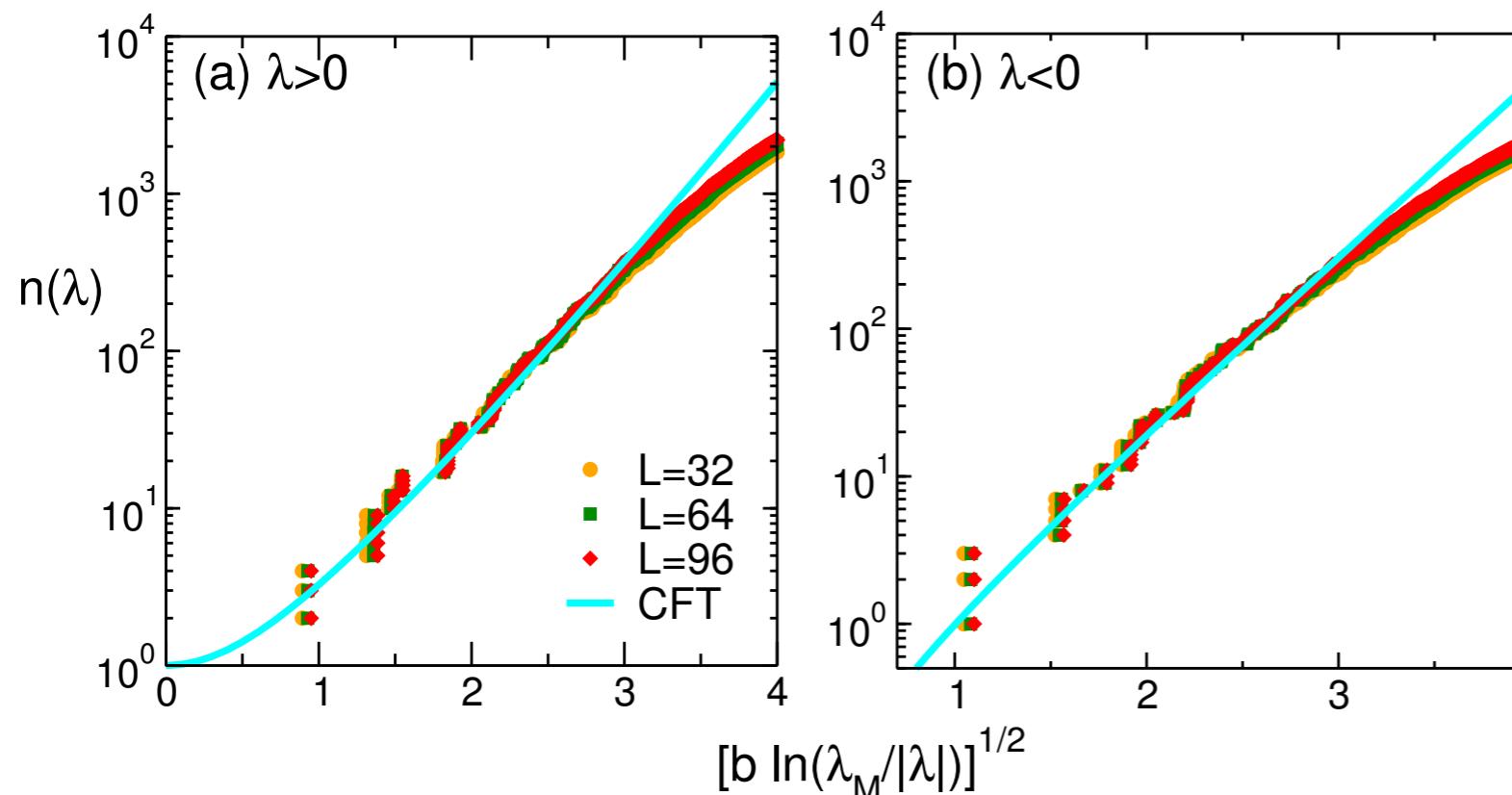
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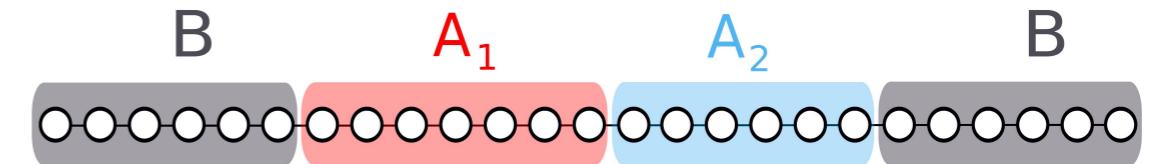
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**“Superuniversal”!**  
(in analogy with the entanglement spectrum)

# NEGATIVITY SPECTRUM in CFTs

Mixed state:

$$R_n^{T_2} \sim \begin{cases} (\ell_1 \ell_2)^{-\frac{c}{6}\left(\frac{n_e}{2} - \frac{2}{n_e}\right)} (\ell_1 + \ell_2)^{-\frac{c}{6}\left(\frac{n_e}{2} + \frac{1}{n_e}\right)} \\ (\ell_1 \ell_2 (\ell_1 + \ell_2))^{-\frac{c}{12}\left(n_0 - \frac{1}{n_0}\right)} \end{cases}$$



$$\tilde{\xi} \equiv \sqrt{\tilde{b} \ln(\lambda_M/\lambda)} \quad \omega \equiv \frac{\ell_2}{\ell_1}$$

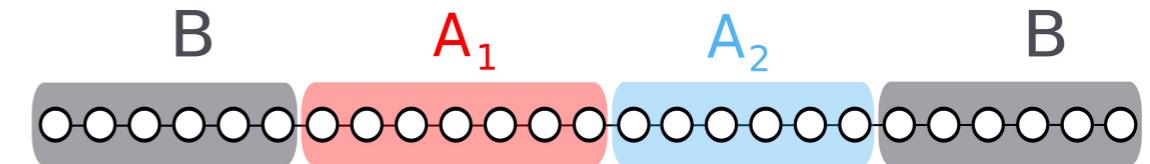
$$\tilde{b} \equiv 2b + \frac{c}{6} \ln \frac{\omega}{(1+\omega)^2}$$

$$P(\lambda) = \delta(\lambda_M - \lambda) + \frac{\theta(\lambda_M - |\lambda|)}{2|\lambda|} \left[ \frac{b}{\xi} I_1(2\xi) \operatorname{sgn}(\lambda) + \frac{\tilde{b}}{\tilde{\xi}} I_1(2\tilde{\xi}) \right]$$

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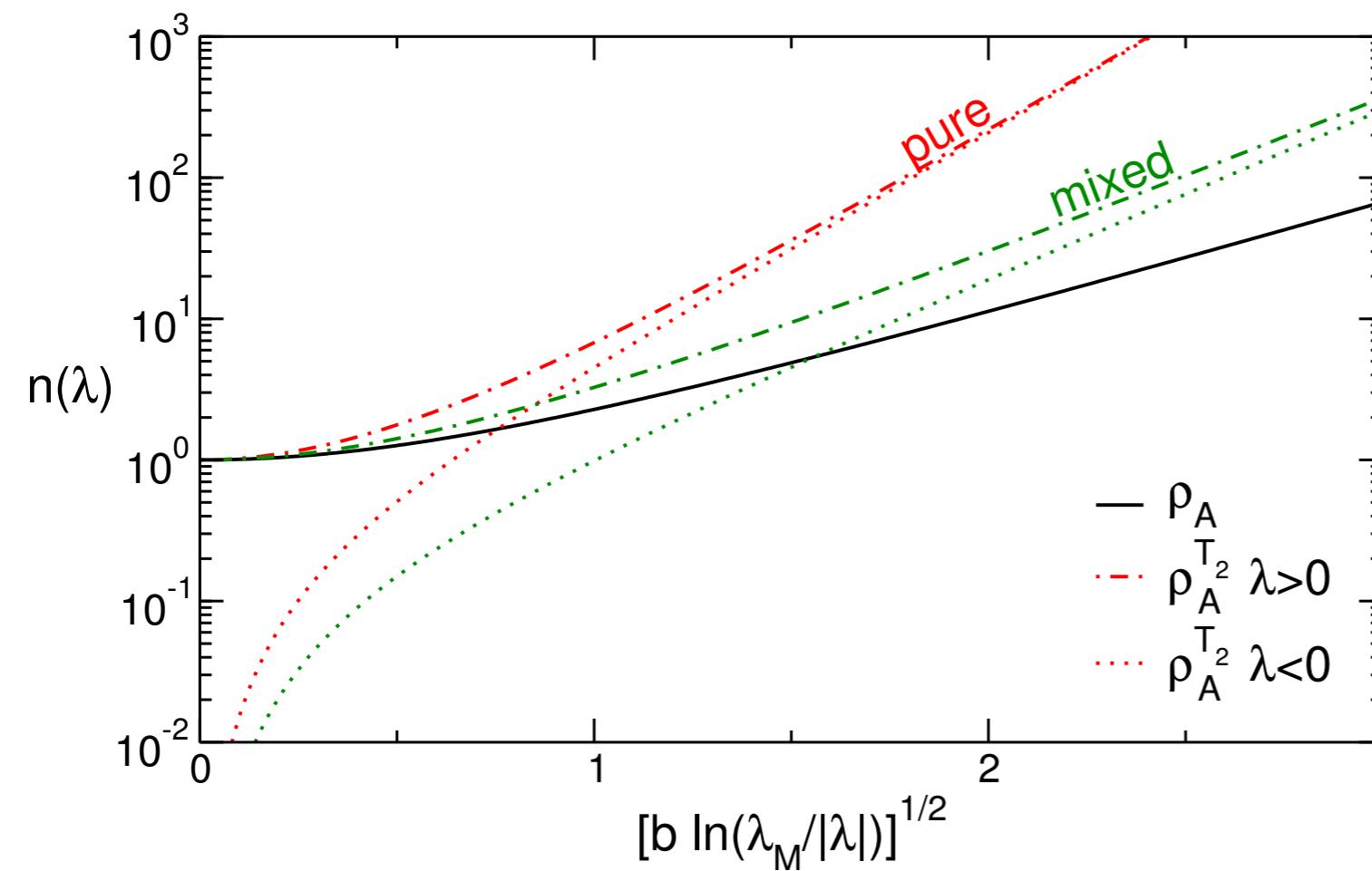
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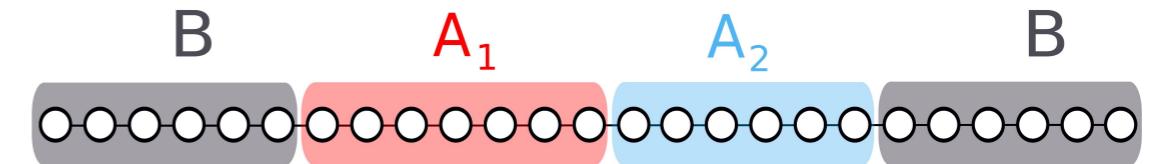
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$$n(\lambda; \xi, \tilde{\xi}) = n(\lambda; c, \omega)$$

# NEGATIVITY SPECTRUM in CFTs

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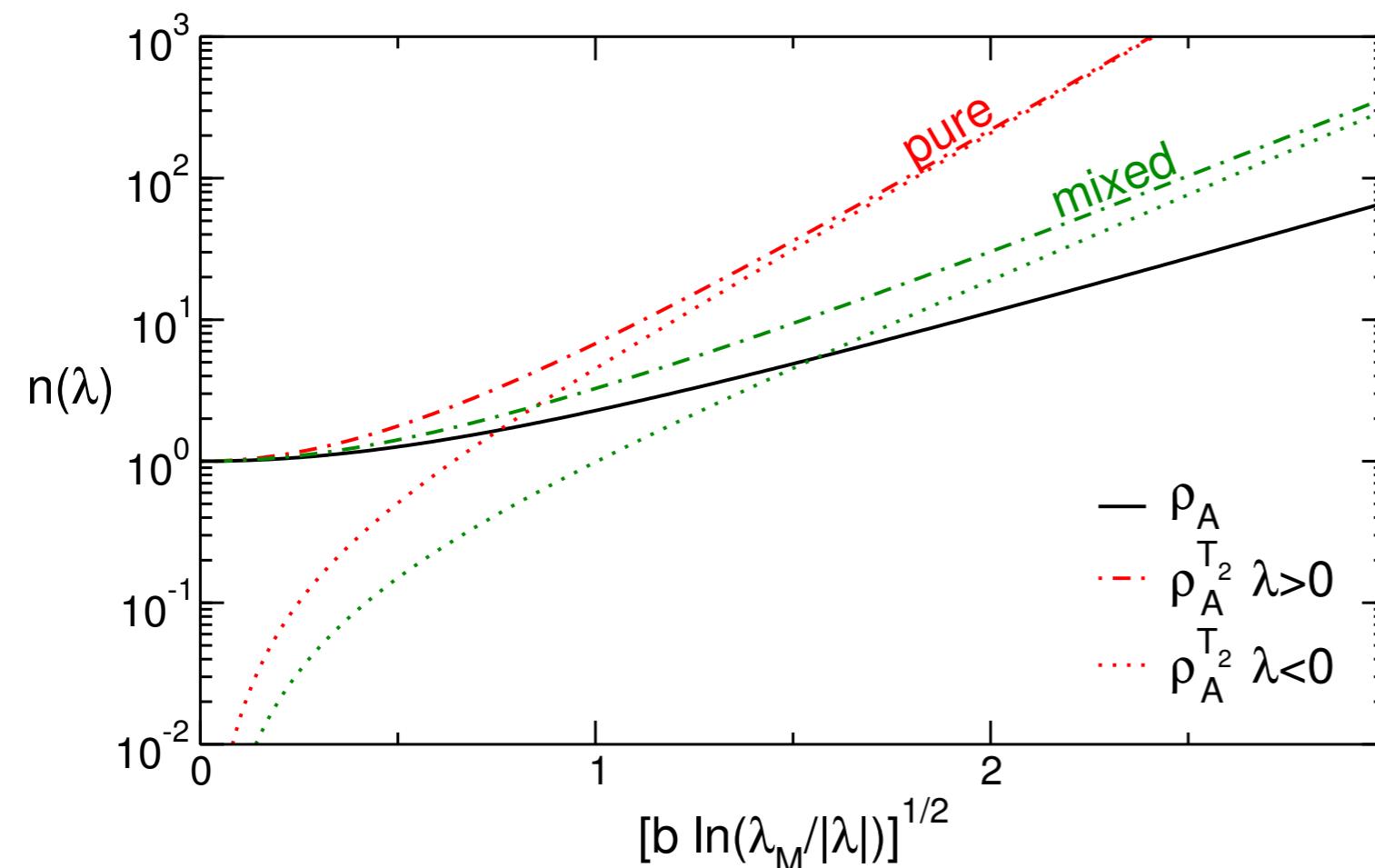
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## NUMBER DISTRIBUTION FUNCTION

$$n(\lambda; \xi, \tilde{\xi}) = n(\lambda; \cancel{\xi}, \cancel{\tilde{\xi}})$$

Limit  $\ell_1 \rightarrow \infty$  :

- “Superuniversal”
- Independent on the geometry of the tripartition

Thank you for the attention.

