

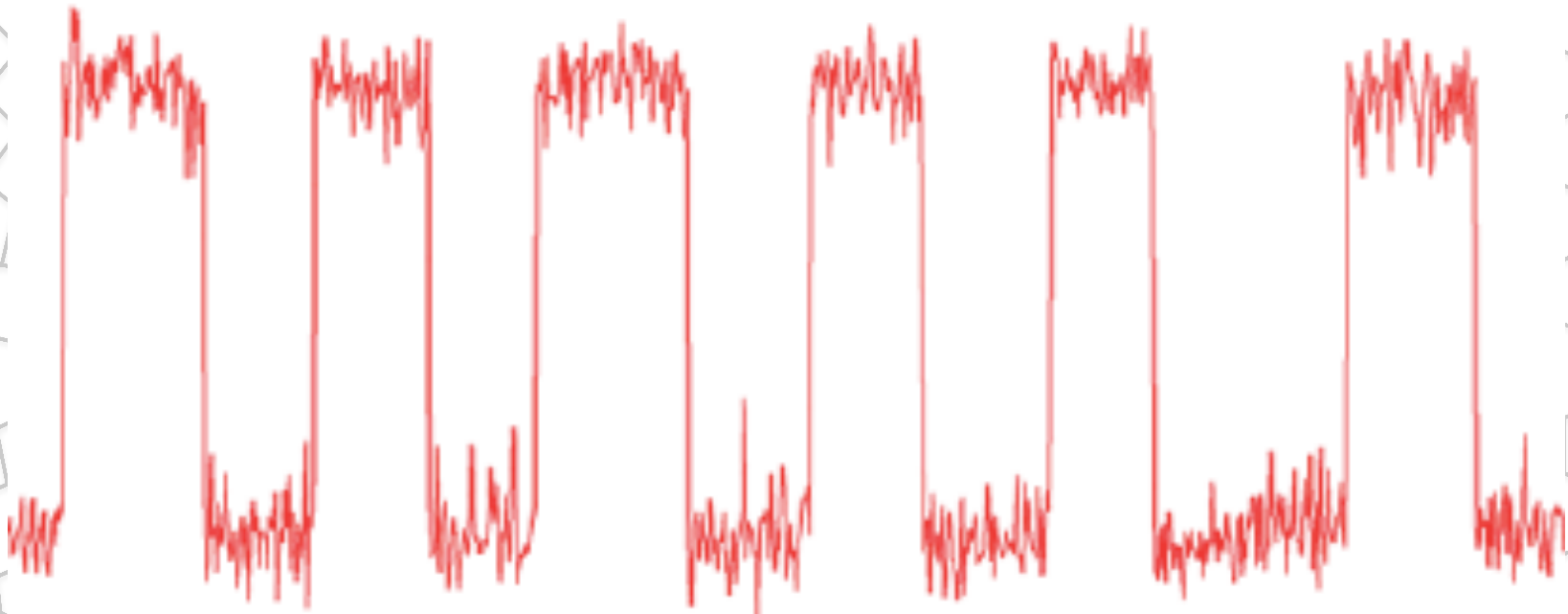
Plastic flow and stochastic resonance in soft glassy materials

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Not liquids, nor solids

"a **solid** has a stable definite shape"

"a **liquid** conforms to the shape of its container"

(from Wikipedia)



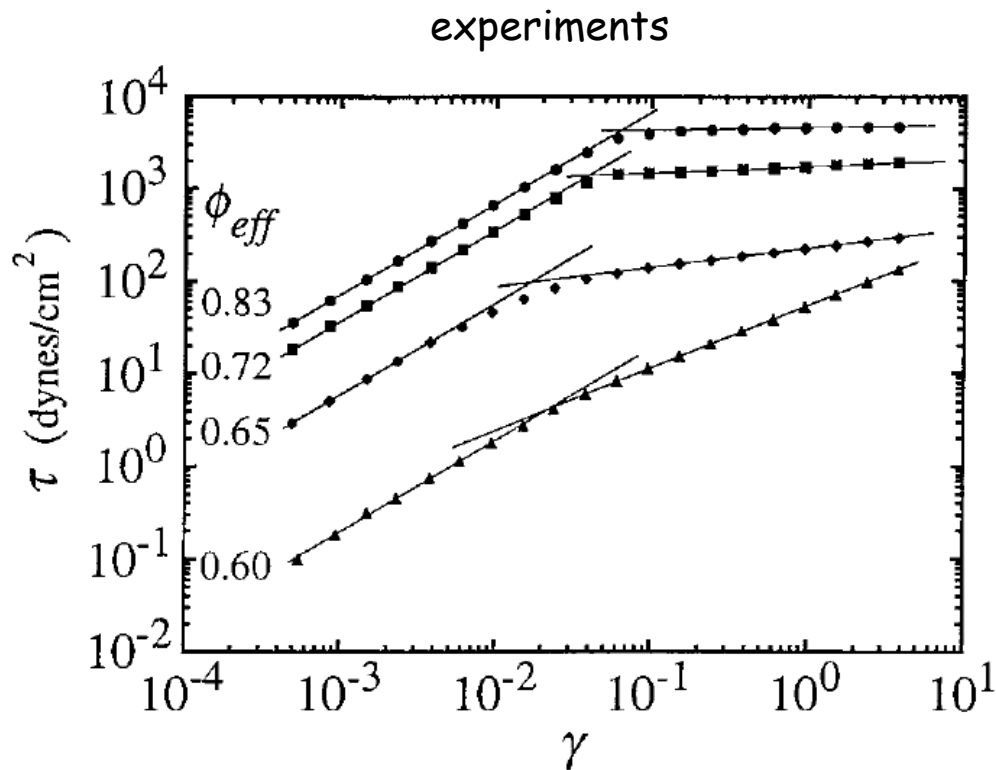
?!?

Yielding and non-Newtonian rheology

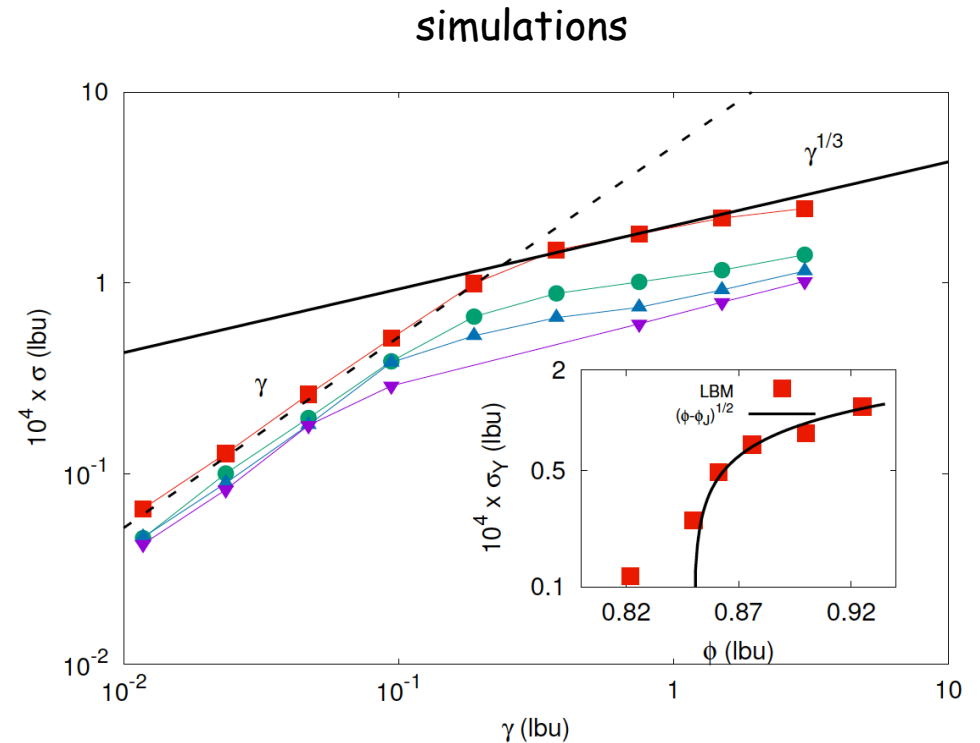
Stress-strain curves for dense microemulsions/dry foams

$$\sigma = \sigma_Y + K\dot{\gamma}^a \quad \text{Herschel-Bulkley relation}$$

σ_Y Yield stress



(TG Mason et al, J. Colloid Interface Sci. **179**, 439 (1996))



(AS et al, Europhys. Lett **114**, 64003 (2016))

Complex fluids

Emulsions, foams, gels, polymers, pastes...



Composition

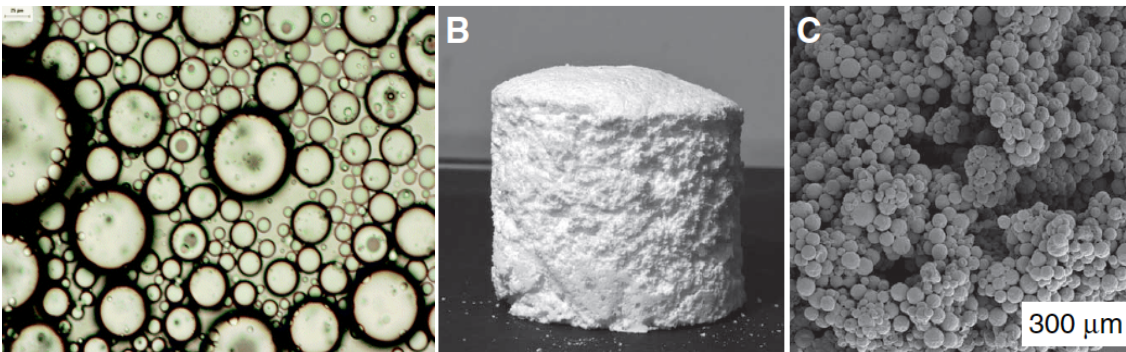
- Multiple phases
- Multiple components
- (Often **both**)

Structural complexity

- Dynamical arrest
- **Structural disorder**
- Heterogeneities

Rheological complexity

- Yield stress
- Viscoelasticity
- **Non-local plastic activity**



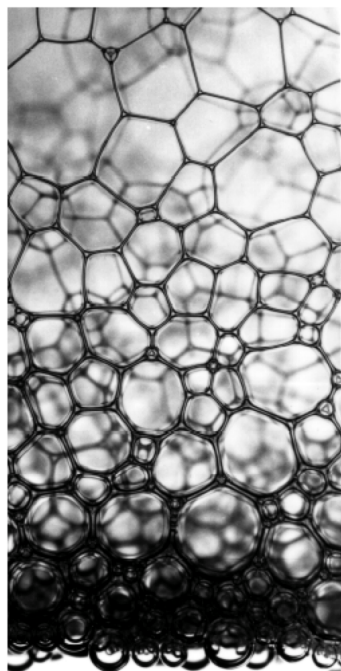
Foams and emulsions

gas (dispersed) + liquid (continuous)

liquid (dispersed) + liquid (continuous)

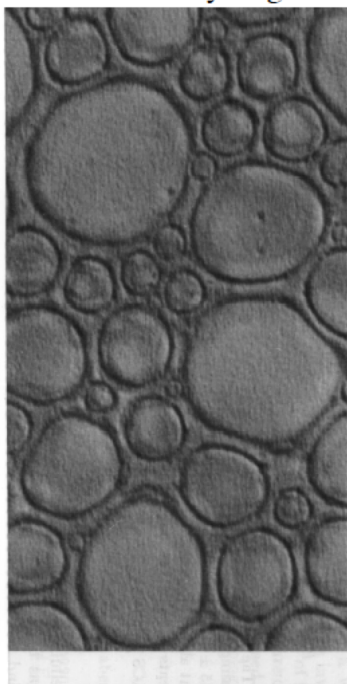
surfactant-stabilised films

5% SDS in water

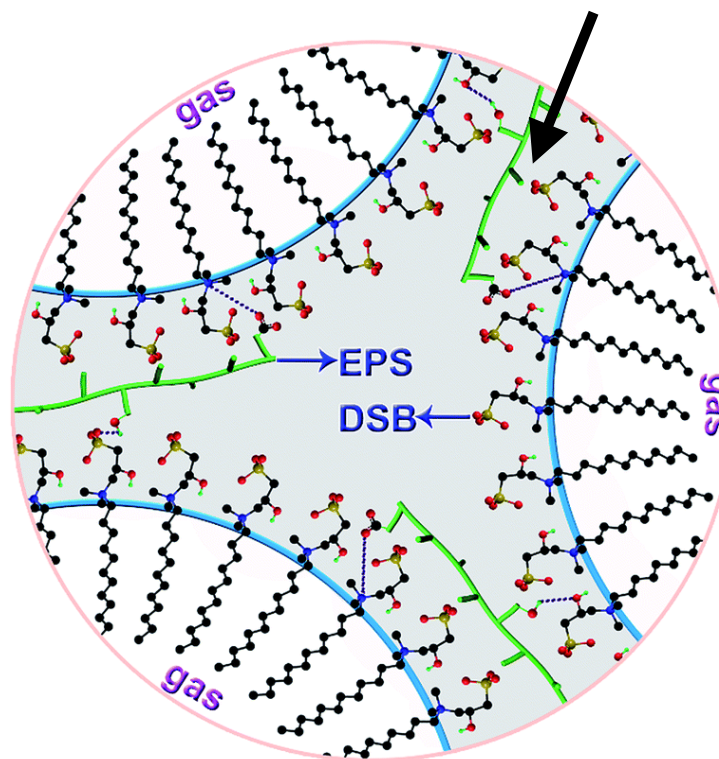


3 mm bubbles

Gillette Foamy Regular



30 μ m bubbles



(Deng et al, *RSC Adv.* **5**, 61868 (2015))



"Effective" modelling of surfactants. What is their the mechanical 'translation'?

Lattice Boltzmann Method (LBM): basics

Lattice Boltzmann Equation (LBE)

$$f_i^{(\alpha)}(\mathbf{r} + \mathbf{c}_i \Delta t, t + \Delta t) - f_i^{(\alpha)}(\mathbf{r}, t) = \Lambda_{ij}^{(\alpha)} \left(f_j^{(\alpha)} - f_j^{(\alpha)(eq)} \right)$$

FREE-STREAMING

COLLISIONS

(relaxation towards a local equilibrium)

$$f_i^{(\alpha)(eq)} = w_i \rho^{(\alpha)} \left[1 + \frac{\mathbf{u} \cdot \mathbf{c}_i}{c_s^2} + \frac{\mathbf{u} \mathbf{u} : (\mathbf{c}_i \mathbf{c}_i - c_s^2 \mathbf{1})}{c_s^4} \right]$$

local equilibrium
(low Mach number expansion of Maxwellian equilibrium)

$\alpha = A, B$
 $i = 0, 2, \dots, 8$

Hydrodynamic fields

density

$$\rho^{(\alpha)}(\mathbf{r}, t) = \sum_i f_i^{(\alpha)}(\mathbf{r}, t)$$

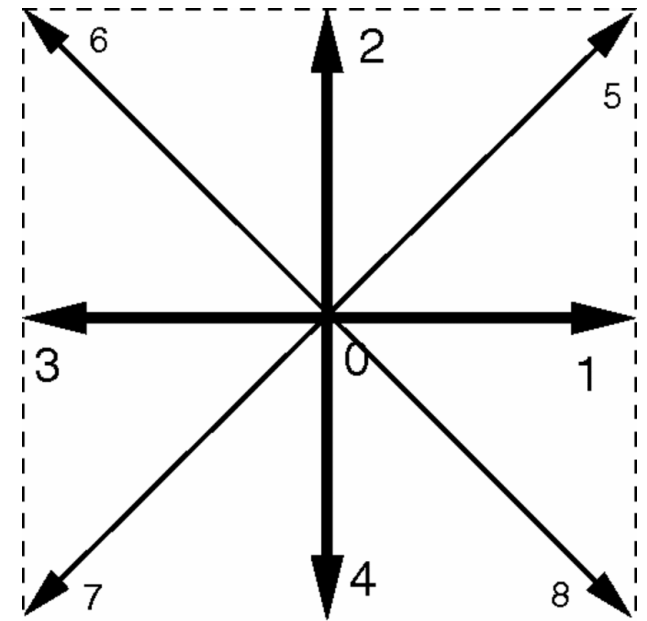
momentum

$$\rho(\mathbf{r}, t) \mathbf{u}(\mathbf{r}, t) = \sum_{i, \alpha} f_i^{(\alpha)}(\mathbf{r}, t) \mathbf{c}_i$$

Force enters through the local equilibrium

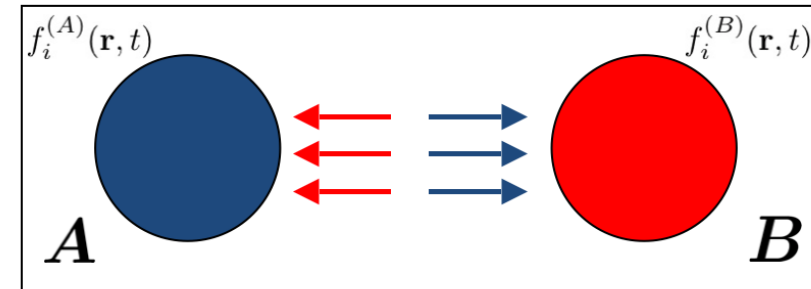
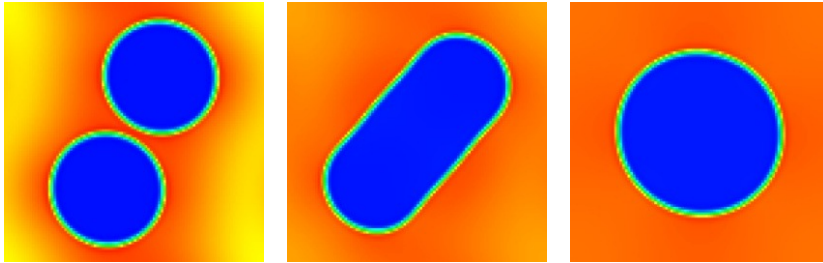
$$f_i^{(\alpha)(eq)}(\rho^{(\alpha)}, \mathbf{u}) \rightarrow f_i^{(\alpha)(eq)}(\rho^{(\alpha)}, \mathbf{u} + \tau_{LB} \mathbf{F}^{(\alpha)} / \rho^{(\alpha)})$$

lattice speeds

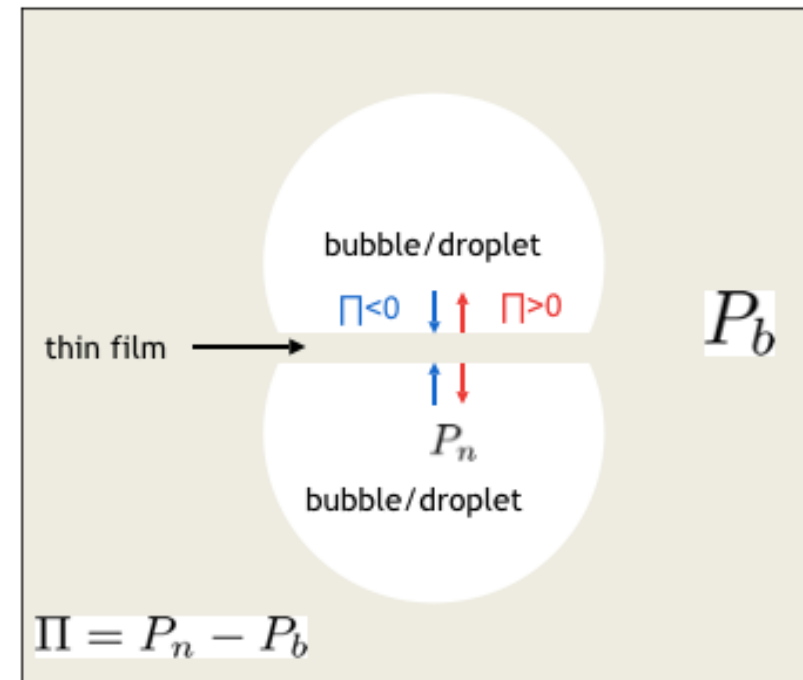
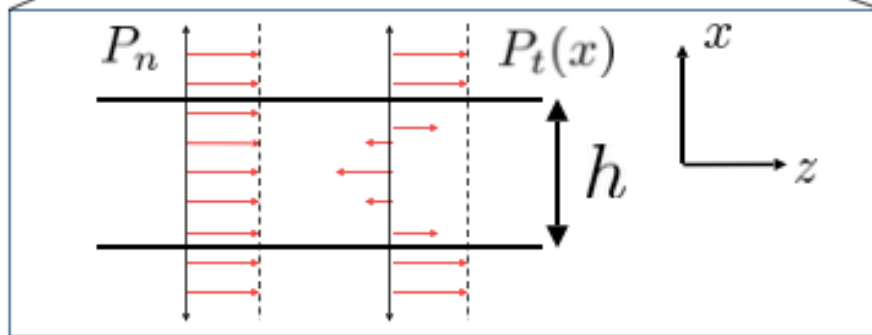
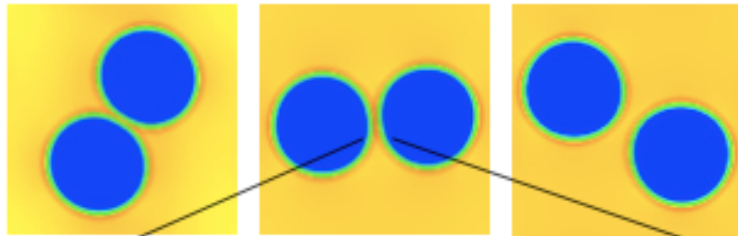


Modelling the surface tension and disjoining pressure

Surface tension (phase separation)

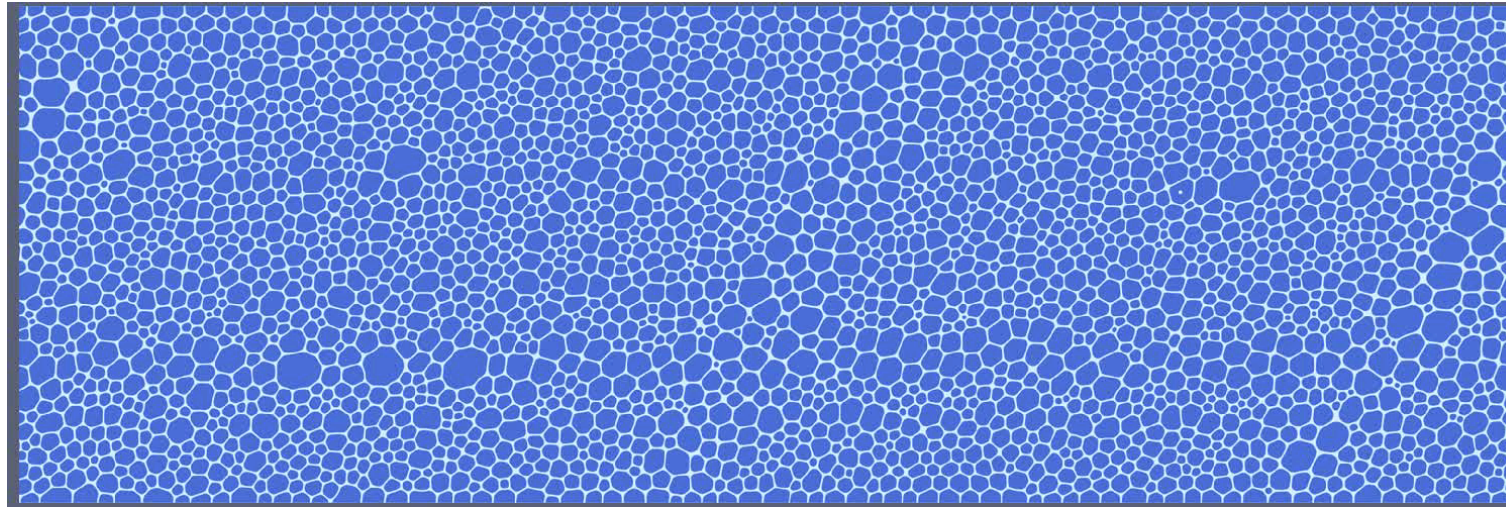


Disjoining pressure (stabilisation of thin films)



NUMERICAL "FOAM"

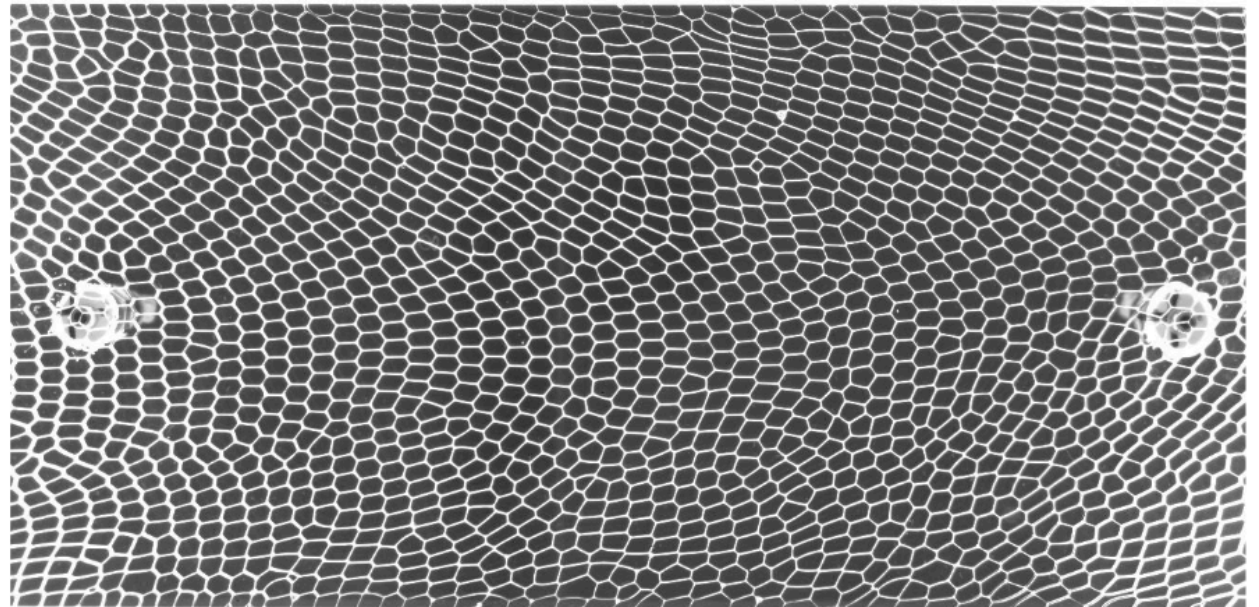
$+U(t)$



$-U(t)$

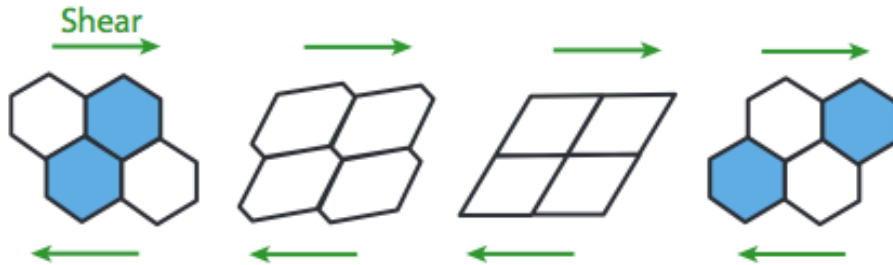
REAL ONE

$$U(t) = U_P \sin(\omega t)$$

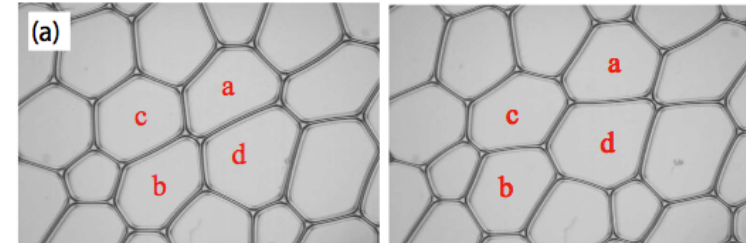


Plastic events and non-local stress relaxation

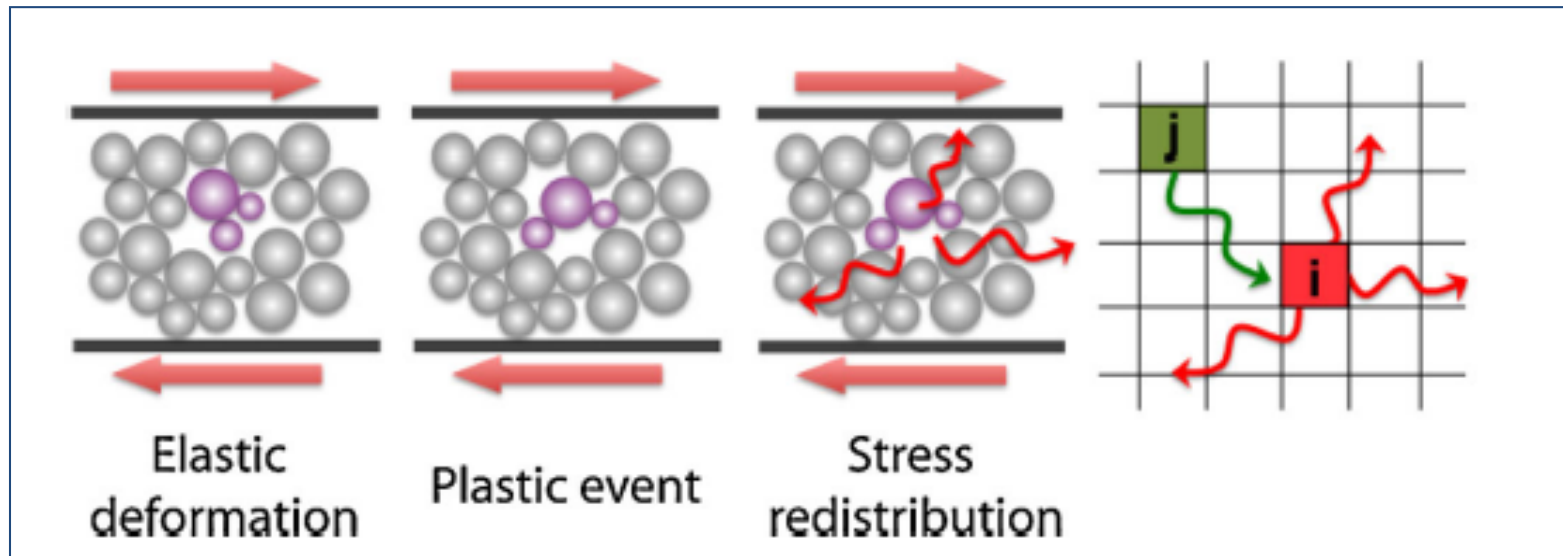
Plastic rearrangements: **T1 event**



(S. Cohen-Addad et al, Annu. Rev. Fluid Mech. **45**, 241 (2013))



(A. Kabla et al, J. Fluid Mech. **587**, 45 (2007))



(L. Bocquet et al, Phys. Rev. Lett. **103**, 036001 (2009))

Theory of Soft-Glassy Rheology

Equation for the probability distribution of stresses

$$\frac{\partial}{\partial t} P = -\dot{\gamma} \frac{\partial}{\partial l} P - \Gamma_0 e^{-(E - \frac{1}{2}kl^2)/x} P + \Gamma(t)\rho(E)\delta(l)$$

structural disorder



"complex" energy landscape
(distribution of energy wells)

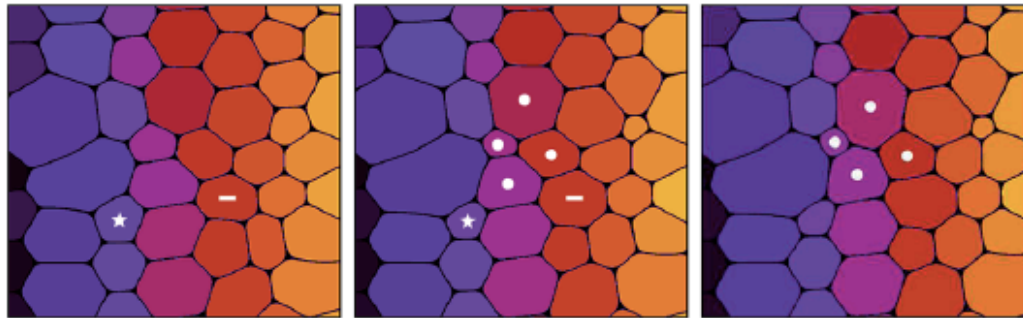
**(non-local) coupling among
plastic rearrangements**



effective "noise temperature"

Does such conjecture make any sense?

"Ageing" and rearrangements



(a) $t=t_0$

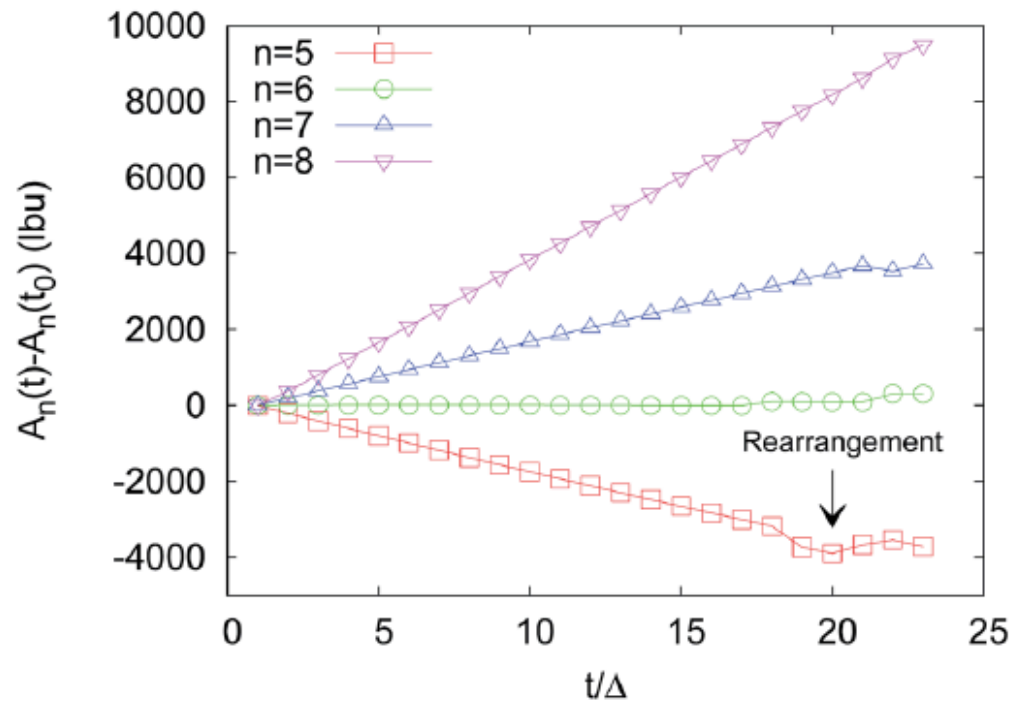
(b) $t=t_0+20\Delta$

(c) $t=t_0+23\Delta$

von Neumann's law

$$\frac{dA_n(t)}{dt} = K(n - 6)$$

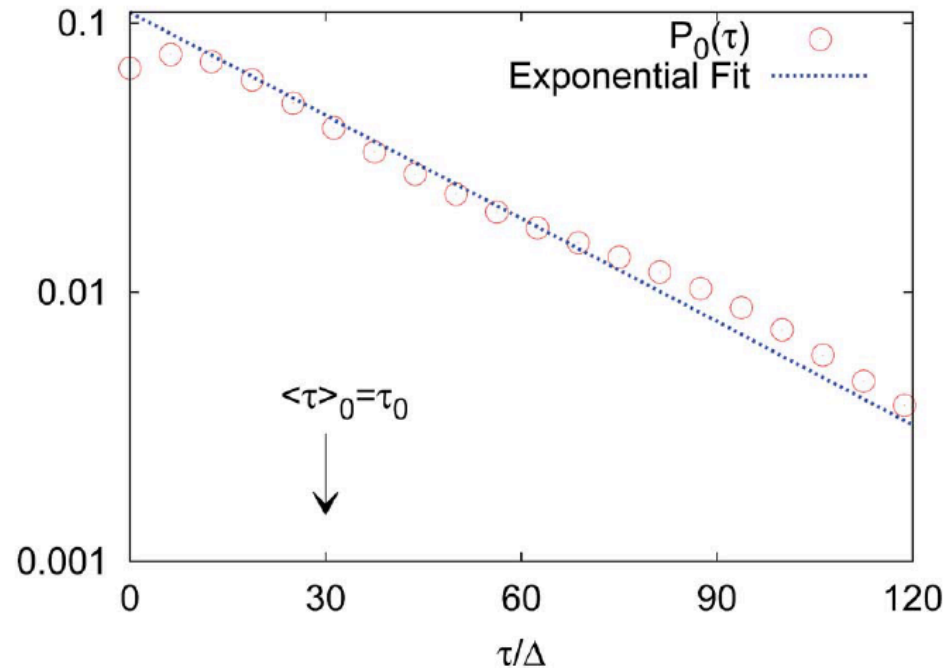
(variation of area of bubbles with n neighbours)



coarsening as a source of effective
"mechanical noise"

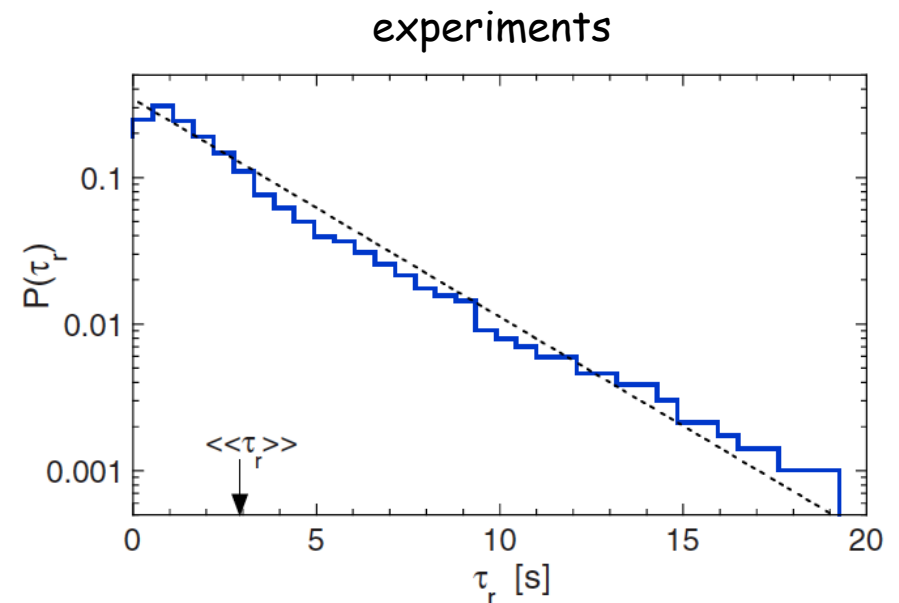
Statistics of rearrangements during coarsening

Probability density function of time lapses
between successive plastic rearrangements



$$P(\tau) \sim e^{-\tau/\tau_0}$$

τ_0 **internal** characteristic time-scale
(**NO STRAIN APPLIED!**)



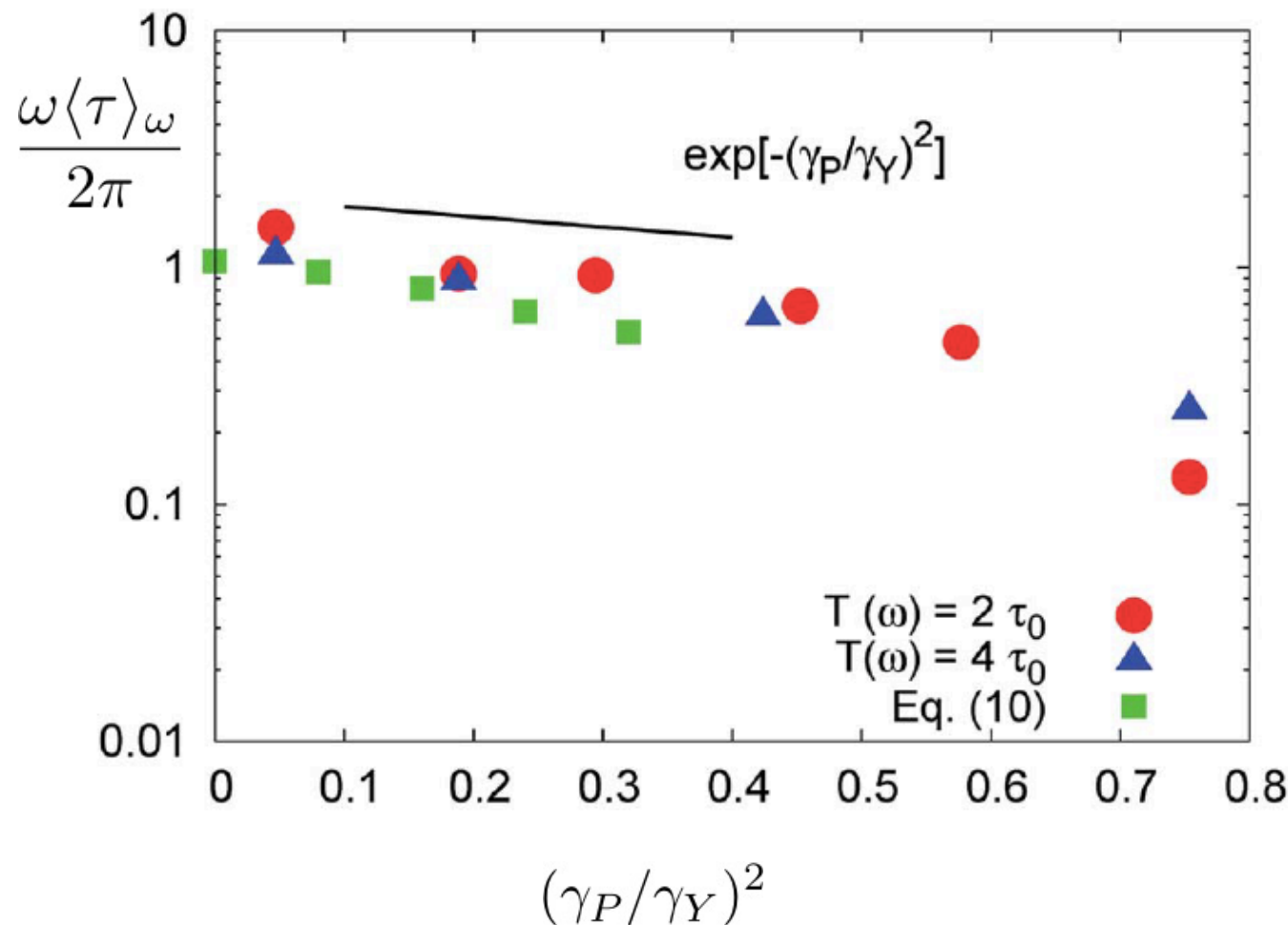
Mean time lapse under oscillatory strain

oscillatory strain

$$\gamma(t) = \gamma_P \sin(\omega t)$$

below yield conditions!

$$\gamma_P < \gamma_Y$$



Energy input $\sim E \frac{\gamma_P^2}{2}$

Energy barrier $\sim E \frac{\gamma_Y^2}{2}$



hints for the activated processes scenario...

Stochastic Resonance (SR) in a nutshell

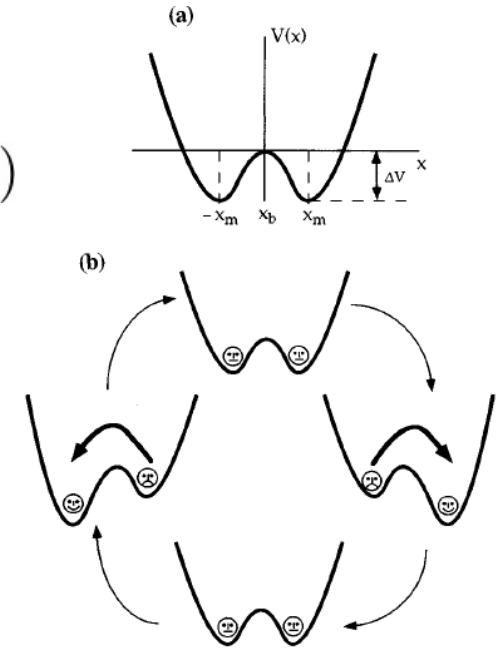
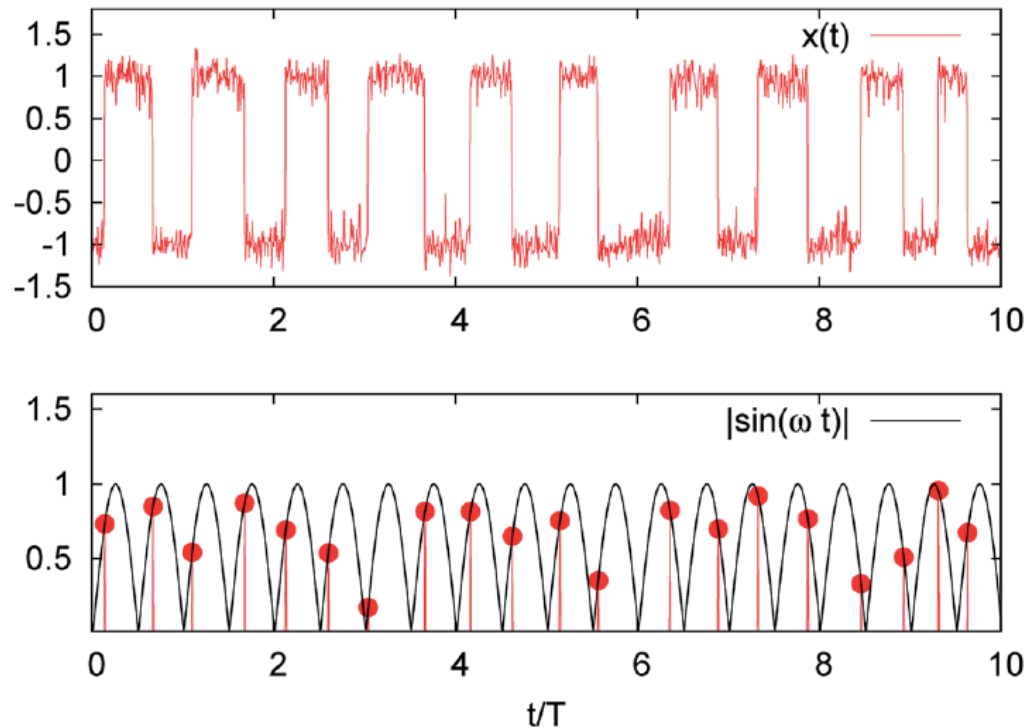
overdamped of a Brownian particle + periodic forcing

$$\dot{x} = x(1 - x^2) + A \sin(\omega t) + \xi(t) \quad \langle \xi(t)\xi(t') \rangle = \varepsilon \delta(t - t')$$

mean hopping time

$$A = 0 \quad \longrightarrow \quad \tau_x = \frac{\pi}{\sqrt{2}} \exp\left(\frac{2\Delta V}{\varepsilon}\right)$$

$$A \neq 0$$



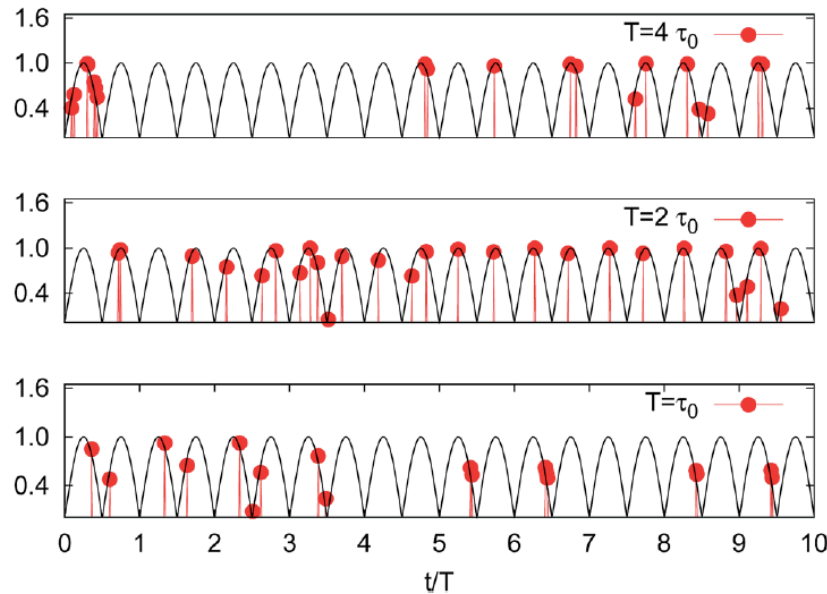
(L. Gammaitoni et al, Rev. Mod. Phys. **70**, 223-287 (1998))

matching condition

$$\frac{\pi}{\omega} = \tau_x$$

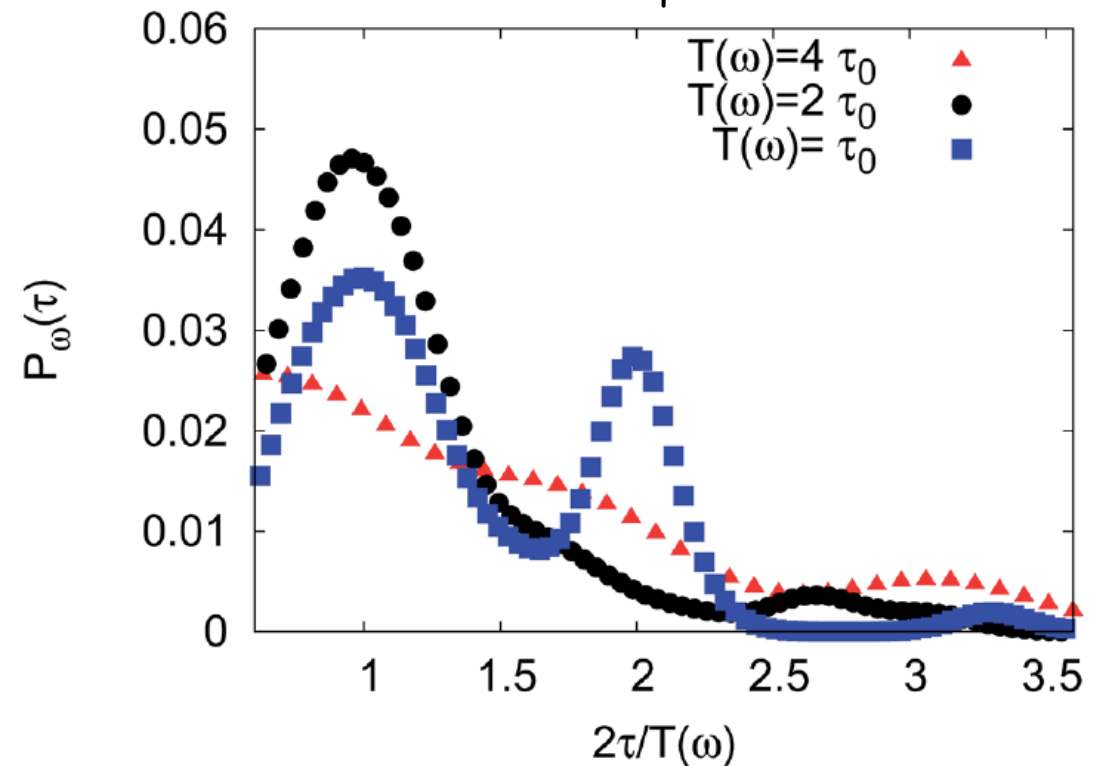
SR signatures in the plastic dynamics of Soft-Glassy Materials

$|\sin(\omega t)|$ and occurrence of plastic events



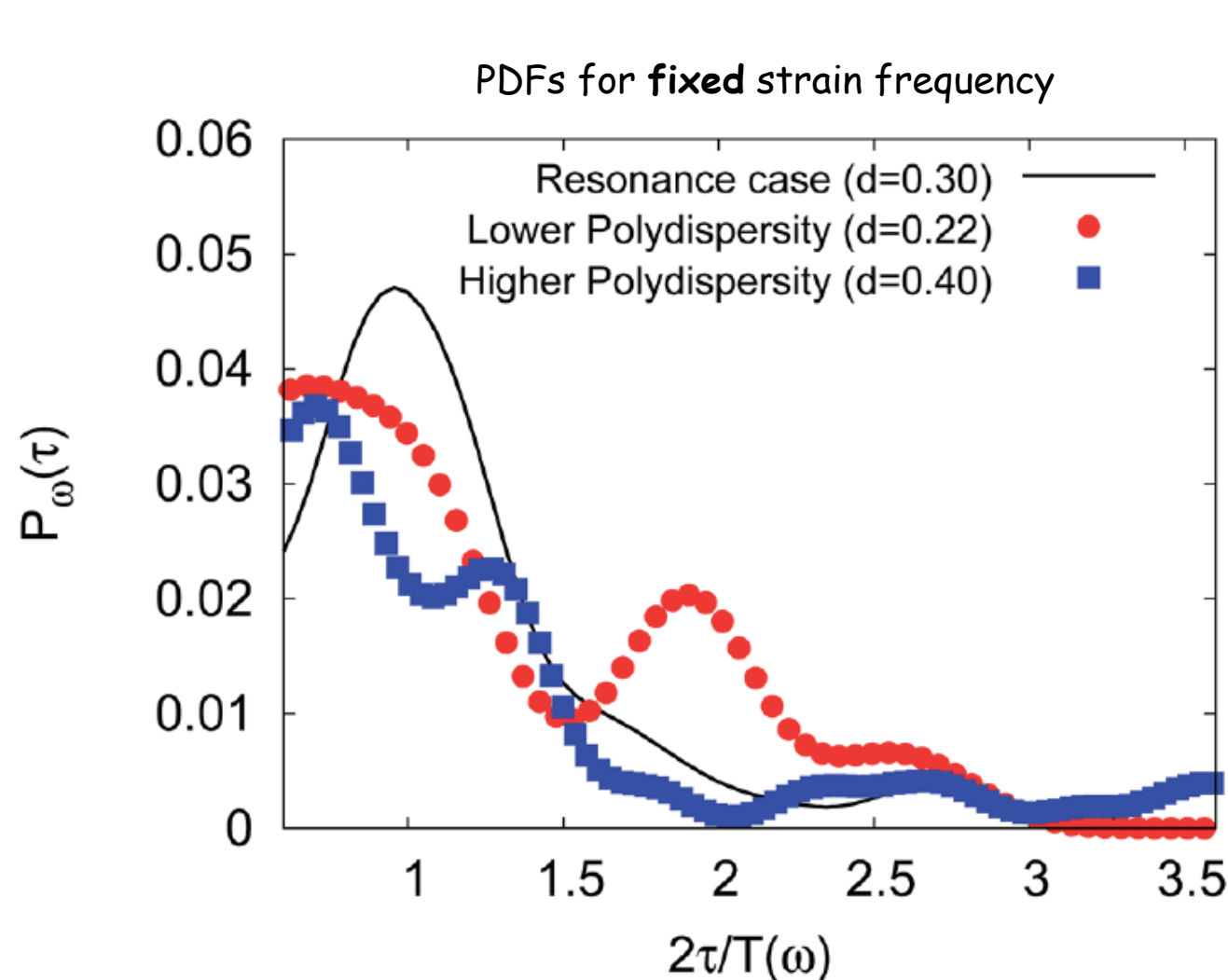
matching
frequency

PDFs of time lapses between
successive plastic events



how to “change” the intrinsic noise?

Polydispersity: a control on the intrinsic effective noise?



$$\text{polydispersity} \\ d = \frac{\langle \mathcal{A}^2 \rangle^{1/2}}{\langle \mathcal{A} \rangle}$$



changing polydispersity drives the system out of the matching condition!



polydispersity \sim noise

Take home messages...

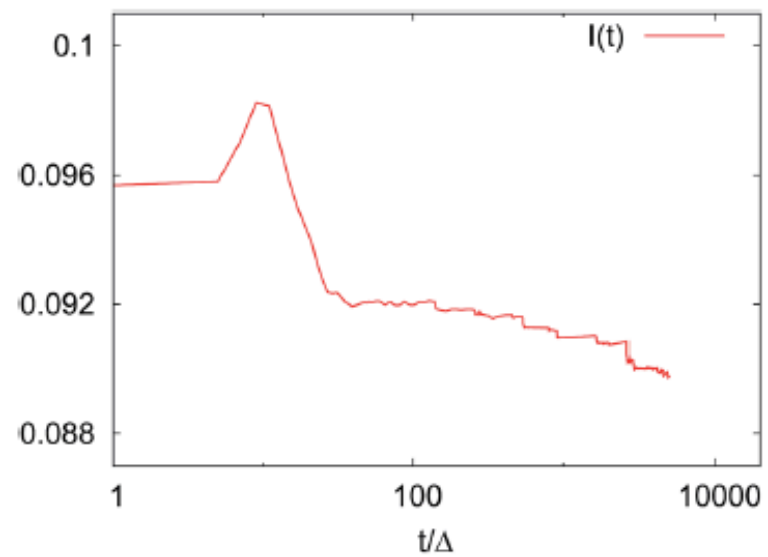
- **Intrinsic** characteristic **time-scale** of plastic activity induced by **coarsening**
- **Plastic rearrangements** can be regarded as **activated processes** induced by a “noise”
- **Signatures of stochastic resonance** from simulations of oscillatory strain
- **Intrinsic noise** correlates to **spatial disorder**, quantified by **polydispersity**

...and perspectives

- Quantification of the intrinsic noise from non-equilibrium fluctuation-dissipation theorems

CREDITS: R. Benzi and M. Sbragaglia (UniToV)
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(a)

$$I(t) = \frac{1}{L^2} \int |\nabla \phi(r)|^2 dr$$

