

Spontaneous wave function collapse models: an introduction

Training school: Are spin-statistics connection and quantum theory exact? The endeavor for the theory beyond standard quantum mechanics

LNF-INFN Frascati, 19th – 21st December 2016

(Angelo Bassi – University of Trieste & INFN)

Motivation

It would seem that the theory is exclusively concerned about 'results of measurement', and has nothing to say about anything else. What exactly qualifies some physical systems to play the role of 'measurer'? Was the wavefunction of the world waiting to jump for thousands of millions of years until a single-celled living creature appeared? Or did it have to wait a little longer, for some better qualified system . . . with a PhD? If the theory is to apply to anything but highly idealised laboratory operations, **are we not obliged to admit that more or less 'measurement-like' processes are going on more or less all the time, more or less everywhere? Do we not have jumping then all the time?**

The GRW model is based on the acceptance of the fact that the Schrodinger dynamics, governing the evolution of the wavefunction, has to be modified by the inclusion of stochastic and nonlinear effects.

**Against
'measurement'**

JOHN BELL

Modify the Schrödinger equation

J.S.Bell

Speakable and Unspeakable in Quantum Mechanics

E.P. Wigner

in: Quantum Optics, Experimental gravity and Measurement theory, Plenum, NY (1983)

A.J. Leggett

Supplement Progr. Theor. Phys. 69, 80 (1980)

H.P. Stapp

In: Quantum Implications: Essay in Honor of David Bohm, Routledge & Kegan Paul, London (1987)

S. Weinberg

Phys. Rev. Lett. 62, 486 (1989).

R. Penrose

In: Quantum Concepts of Space and Time, Oxford U.P. (1985)

S.L. Adler

Quantum Theory as an emergent phenomenon, CUP (2009)

G.C. Ghirardi, A. Rimini, T. Weber

Phys. Rev. D 34, 470 (1986)

P. Pearle

Phys. Rev. A 39, 2277 (1989)

L. Diosi

L. Diosi, *Phys. Rev. A* 40, 1165 (1989)

The GRW model

Systems are described by the **wave function**, which evolves according to the following dynamics.

The wave function is subject to **random jumps**, independently for each “particle”:

$$\psi \longrightarrow \frac{\psi_{\mathbf{x}}^i}{\|\psi_{\mathbf{x}}^i\|}, \quad \psi_{\mathbf{x}}^i = \hat{L}_{\mathbf{x}}^i \psi, \quad \hat{L}_{\mathbf{x}}^i = \frac{1}{(\pi r_C)^{3/4}} \exp \left[-\frac{(\hat{\mathbf{q}}_i - \mathbf{x})^2}{2r_C^2} \right]$$

i = degree of freedom (“ i ’th particle” of a multi-particle system), \mathbf{x} = point in space, $\hat{\mathbf{q}}_i$ = position operator of the “ i ’th particle”.

The jumps are **random in time** (distributed according to a Poissonian distribution with mean frequency λ), **and in space** (with probability density $\|\psi_{\mathbf{x}}^i\|^2$)

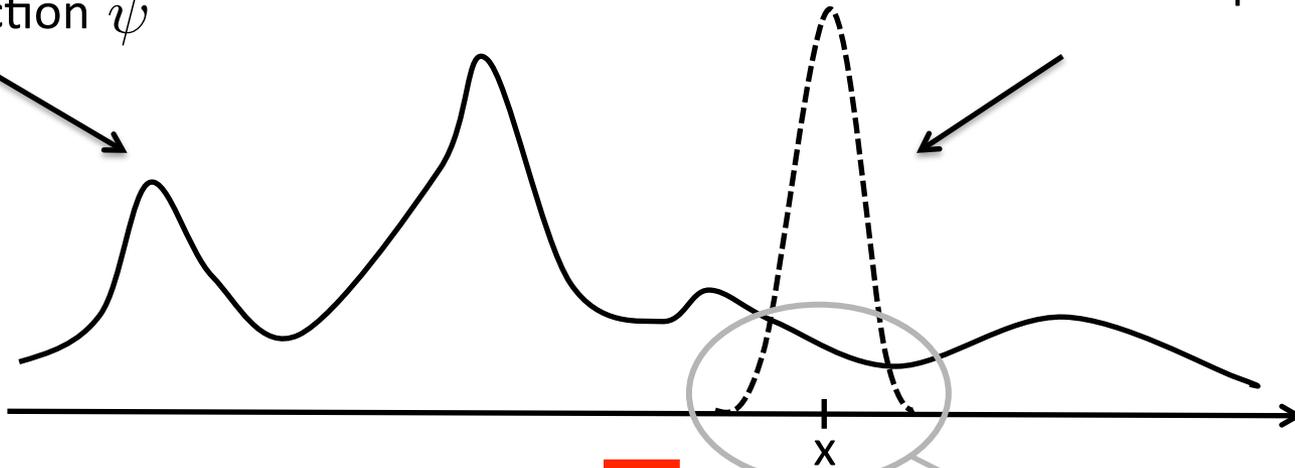
In between jumps, the wave function evolves according to the **Schrödinger equation**.

Note. Two phenomenological parameters: λ and r_C .

The jump

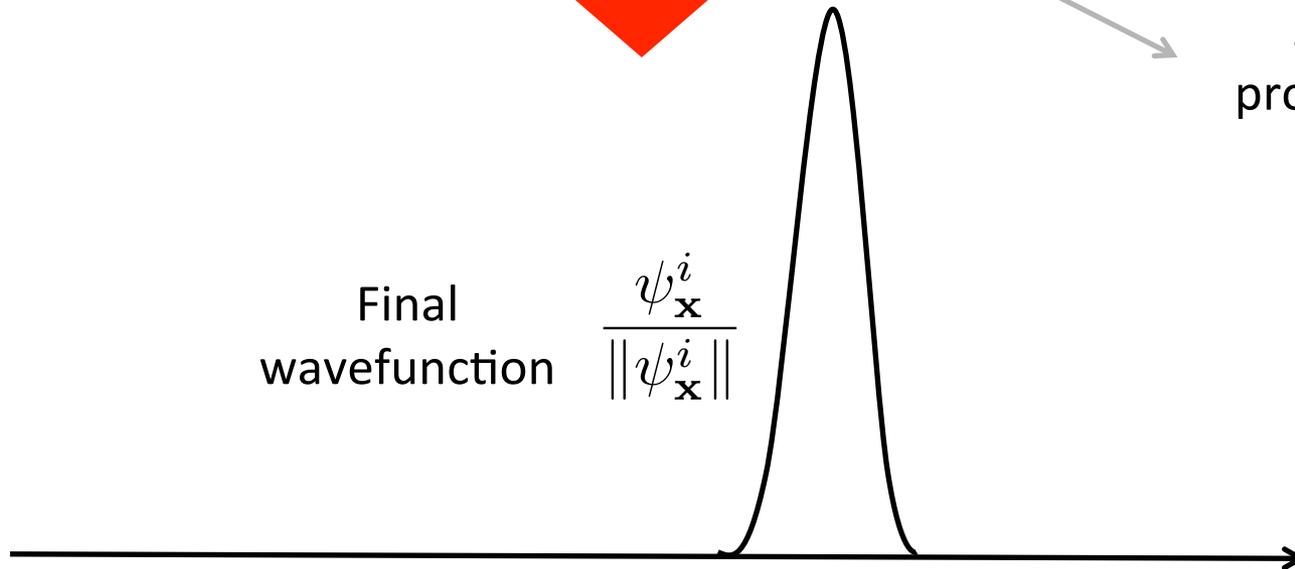
Initial
wavefunction ψ

Jump operator $\hat{L}_{\mathbf{x}}^i$



Jump
probability

Final
wavefunction $\frac{\psi_{\mathbf{x}}^i}{\|\psi_{\mathbf{x}}^i\|}$



Diffusion Process in Hilbert Space

L. Diosi, *Phys. Rev. A* 40, 1165 (1989)

$$d\psi_t = \left[-\frac{i}{\hbar} \hat{H} dt + \sqrt{\lambda} (\hat{q} - \langle \hat{q} \rangle_t) dW_t - \frac{\lambda}{2} (\hat{q} - \langle \hat{q} \rangle_t)^2 dt \right] \psi_t$$

$$\langle \hat{q} \rangle_t = \langle \psi_t | \hat{q} | \psi_t \rangle$$

nonlinearity

W_t = standard Wiener process

stochasticity

All of quantum (and classical) mechanics follows

(Mass-proportional) CSL model

P. Pearle, *Phys. Rev. A* 39, 2277 (1989). G.C. Ghirardi, P. Pearle and A. Rimini, *Phys. Rev. A* 42, 78 (1990)

$$\frac{d}{dt}|\psi_t\rangle = \left[-\frac{i}{\hbar}H + \frac{\sqrt{\gamma}}{m_0} \int d^3x (M(\mathbf{x}) - \langle M(\mathbf{x}) \rangle_t) dW_t(\mathbf{x}) - \frac{\gamma}{2m_0^2} \int \int d^3x d^3y G(\mathbf{x} - \mathbf{y}) (M(\mathbf{x}) - \langle M(\mathbf{x}) \rangle_t) (M(\mathbf{y}) - \langle M(\mathbf{y}) \rangle_t) \right] |\psi_t\rangle$$

$$M(\mathbf{x}) = ma^\dagger(\mathbf{x})a(\mathbf{x})$$

$$G(\mathbf{x}) = \frac{1}{(4\pi r_C^2)^{3/2}} \exp[-(\mathbf{x})^2/4r_C^2]$$

The operators are function of the space coordinate. **The collapse occurs in space.**

Two parameters

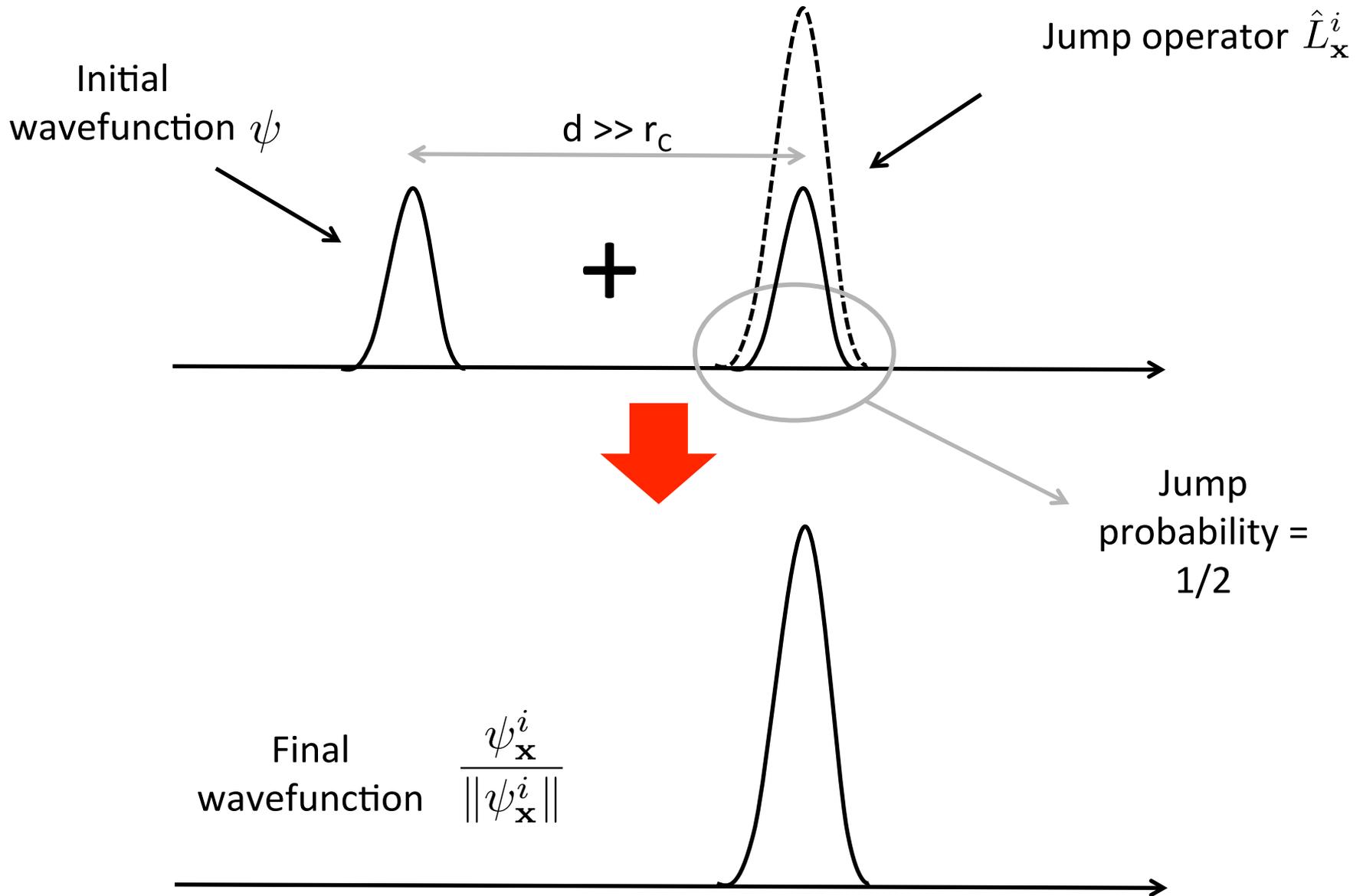
γ = collapse strength

r_C = localization resolution

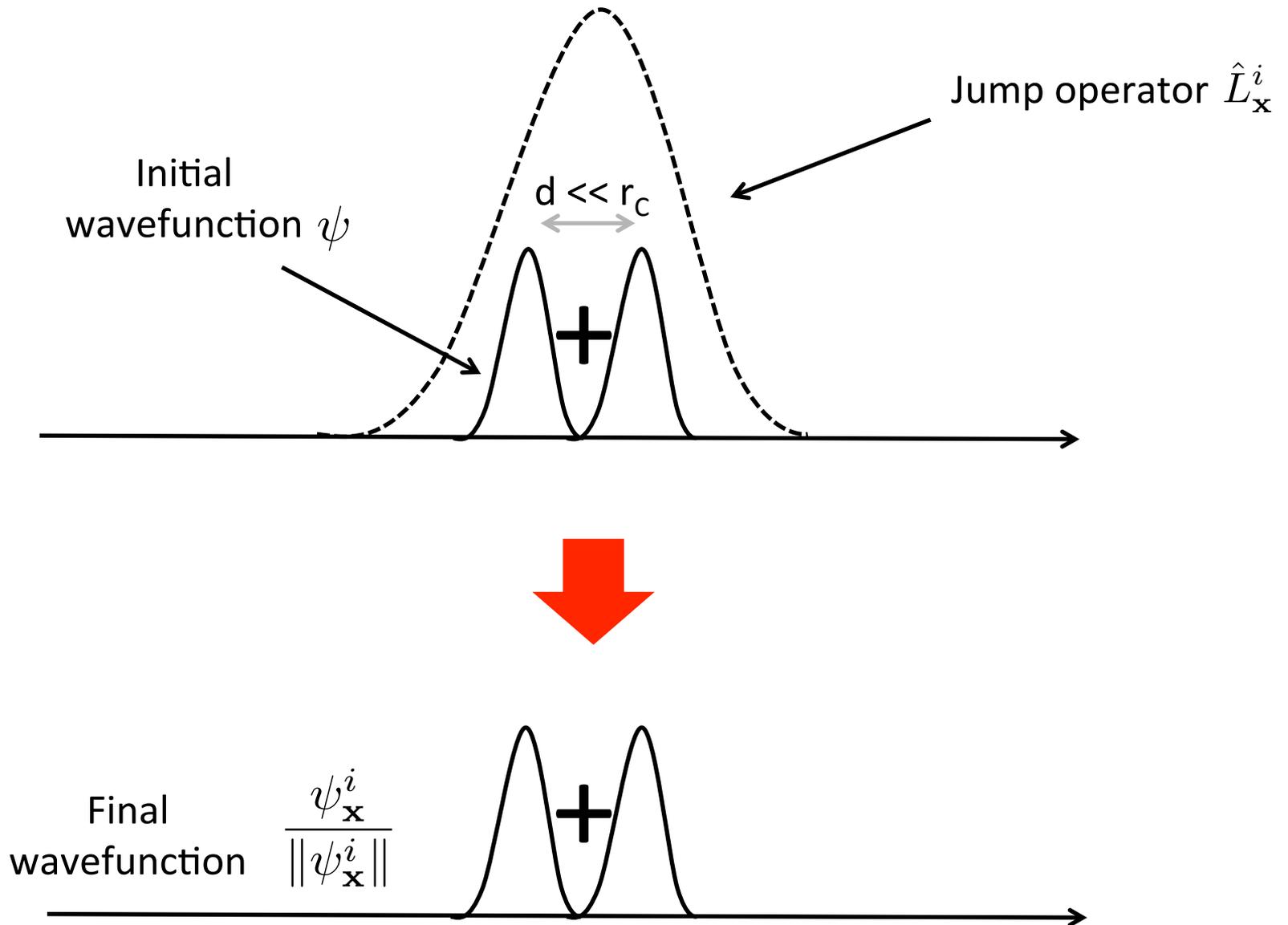


$$\lambda = \gamma / (4\pi r_C^2)^{3/2} = \text{collapse rate}$$

Example: "large" superposition



Example: "small" superposition

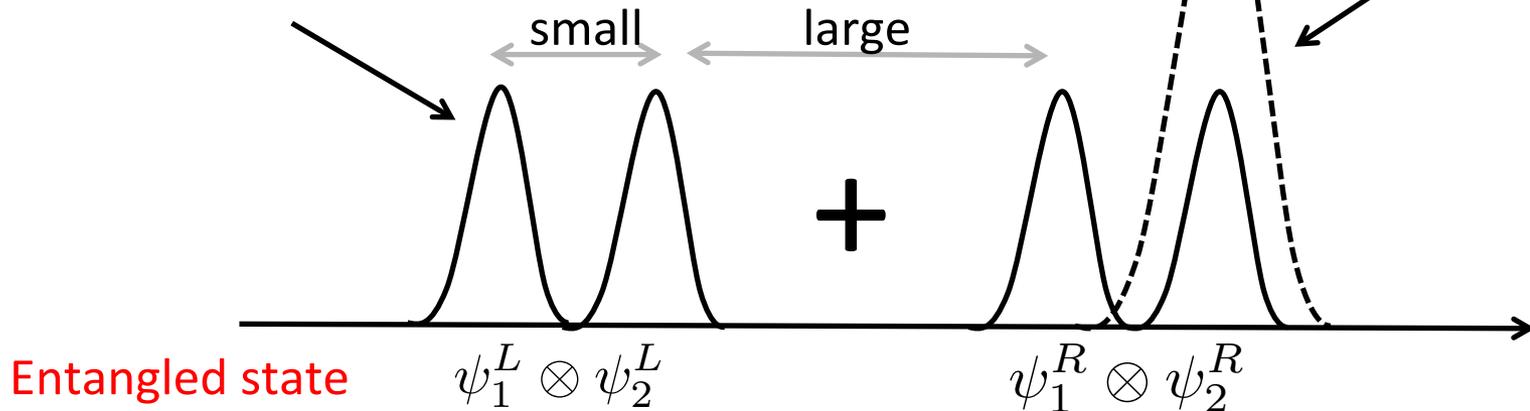


Amplification mechanism

Initial “2-particle” wavefunction

Rigid object: system left + system right

Jump operator
on “particle” 2



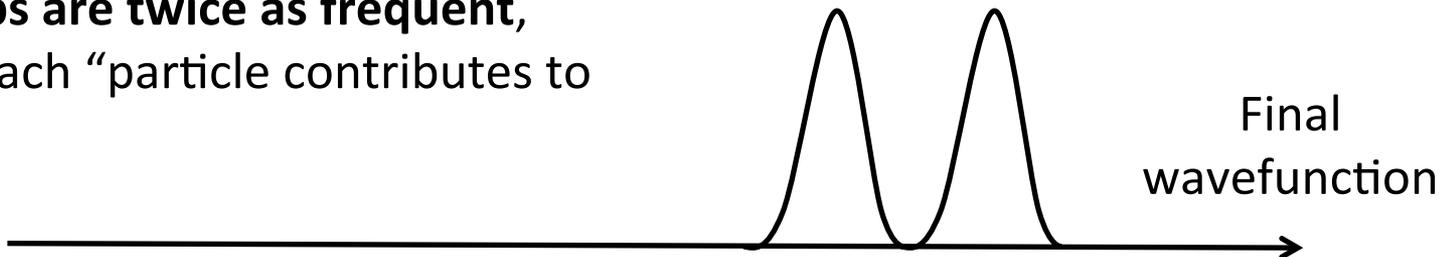
Entangled state

$$\psi_1^L \otimes \psi_2^L$$

$$\psi_1^R \otimes \psi_2^R$$



Such **jumps** are **twice as frequent**,
because each “particle contributes to
them

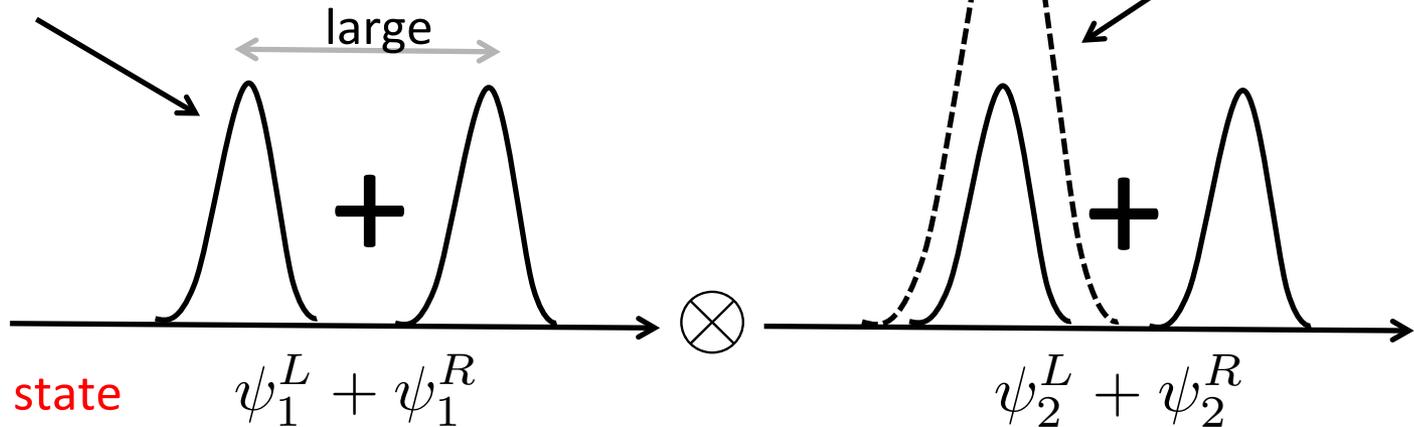


However

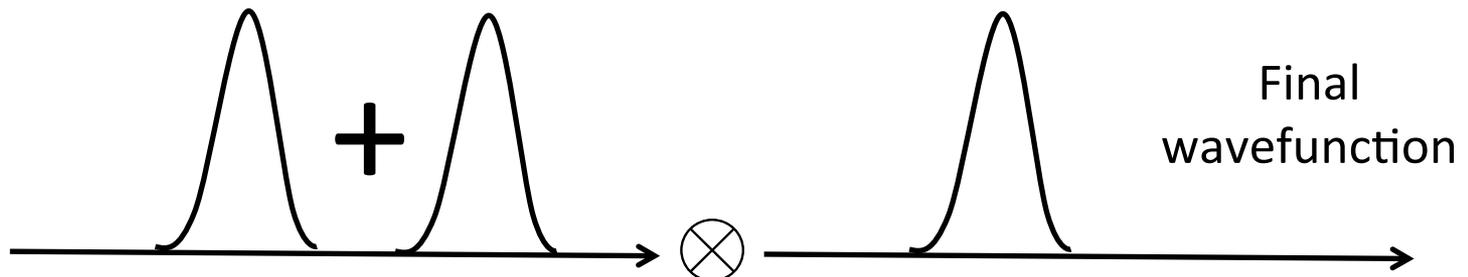
Initial "2-particle" wavefunction

Ideal gas: particles are independent

Jump operator
on "particle" 2



The jump on one particle did not affect
the state of the other particle!



Choice of the parameters (GRW)

$$\lambda = 10^{-16} \text{ s}^{-1}$$

For single isolated particles jumps almost never occur. However, for macroscopic objects

$$N \lambda = 10^{24} \times 10^{-16} \text{ s}^{-1} = 10^8 \text{ s}^{-1}$$

Macro-objects are almost instantly localised

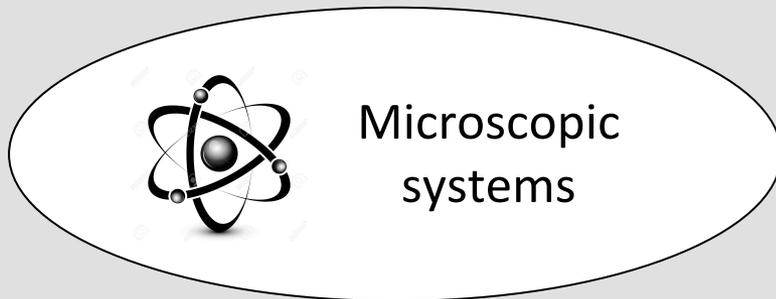
$$r_c = 10^{-7} \text{ m}$$

Mesoscopic distance. Microscopic superpositions are not affected. Macroscopic superpositions are.

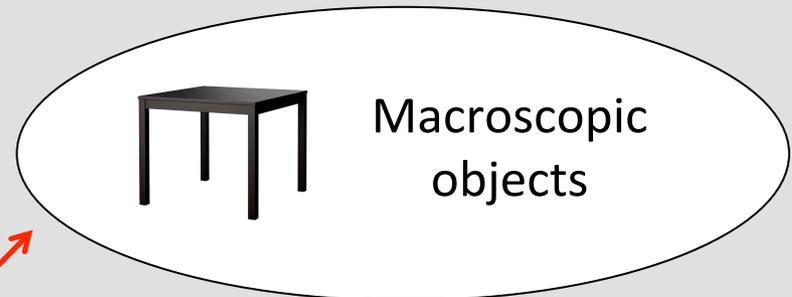
The overall picture

Stable. λ too small

Hilbert space



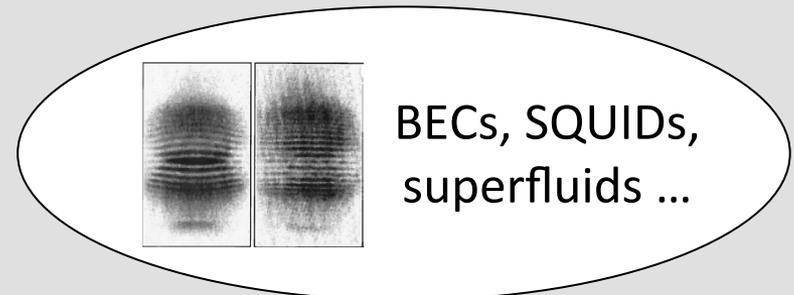
Stable. Already localized ($d \ll r_c$)



Unstable! $N\lambda$ large and $d \gg r_c$



Stable. No cat-like superposition



Ontology

```
graph TD; A[Ontology] --> B[Collapse models]; A --> C[GRW theory];
```

Collapse models

They are phenomenological models

The wave function is not fundamental in any way

Still looking for the underlying theory

No idea what the appropriate ontology is

GRW theory

Are they consistent theories of nature?

Yes (or perhaps) with an appropriate ontology

- Wave function ontology
- Mass density ontology
- Flash ontology
- Particle ontology

How to modify the Schrödinger equation?

The **no-faster-than-light condition** heavily constraints the possible ways to modify the Schrödinger equation.

In particular, it requires that **nonlinear terms** must always be accompanied by appropriate **stochastic terms**.

N. Gisin, *Hel. Phys. Acta* 62, 363 (1989). *Phys. Lett. A* 143, 1 (1990)

N. Gisin and M. Rigo, *Journ. Phys. A* 28, 7375 (1995)

J. Polcinski, *Phys. Rev. Lett.* 66, 397 (1991)

H.M. Wiseman and L. Diosi, *Chem. Phys.* 268, 91 (2001)

S.L. Adler, “Quantum Theory as an Emergent Phenomenon”, C.U.P. (2004)

A. Bassi, D. Dürr and G. Hinrichs, *Phys. Rev. Lett.* 111, 210401 (2013).

L. Diosi, *Phys. Rev. Lett.* 112, 108901 (2014)

M. Caiaffa, A. Smirne and A. Bassi, in preparation

Collapse models in space

REVIEW: A. Bassi and G.C. Ghirardi, *Phys. Rept.* 379, 257 (2003)

REVIEW: A. Bassi, K. Lochan, S. Satin, T.P. Singh and H. Ulbricht, *Rev. Mod. Phys.* 85, 471 (2013)

White noise models

All frequencies appear with the same weight

Colored noise models

The noise can have an arbitrary spectrum

Infinite temperature models

No dissipative effects

GRW / CSL

G.C. Ghirardi, A. Rimini, T. Weber, *Phys. Rev. D* 34, 470 (1986)

G.C. Ghirardi, P. Pearle, A. Rimini, *Phys. Rev. A* 42, 78 (1990)

QMUPL

L. Diosi, *Phys. Rev. A* 40, 1165 (1989)

DP

L. Diosi, *Phys. Rev. A* 40, 1165 (1989)

Non-Markovian CSL

P. Pearle, in *Perspective in Quantum Reality* (1996)

S.L. Adler & A. Bassi, *Journ. Phys. A* 41, 395308 (2008). arXiv: 0807.2846

Non-Markovian QMUPL

A. Bassi & L. Ferialdi, *Phys. Rev. Lett.* 103, 050403 (2009)

Finite temperature models

Dissipation and thermalization

Dissipative QMUPL

A. Bassi, E. Ippoliti and B. Vacchini, *J. Phys. A* 38, 8017 (2005).

Dissipative GRW & CSL

A. Smirne, B. Vacchini & A. Bassi *Phys. Rev. A* 90, 062135 (2014)

A. Smirne & A. Bassi *Nat. Sci. Rept.* 5, 12518 (2015)

Non-Markovian & dissipative QMUPL

L. Ferialdi, A. Bassi *Phys. Rev. Lett.* 108, 170404 (2012)

Three ontologies for GRW

Wave function ontology: there is only ψ . End of the story. Because macro-superpositions never occur, this is consistent with our physical worldview.

Plus: Very economical

Minus: One has to accept that the 3D world of particles we have experience of, is an “illusion”. Problem with tails.

Mass density ontology: “project” the wave function to 3D space, and create a mass density distribution:

$$m(\mathbf{x}, t) = \sum_{i=1}^N m_i \int d^{3N}q \delta^{(3)}(\mathbf{x} - \mathbf{q}_i) |\psi(q)|^2, \quad q = (\mathbf{q}_1, \dots, \mathbf{q}_N)$$

Flash ontology: we are made of flashes, the points around which jumps occur.

Plus: These ontologies makes direct reference to the world as we perceive it.

Minus: One still needs the wave function. A nomological entity?

Relativistic Collapse models

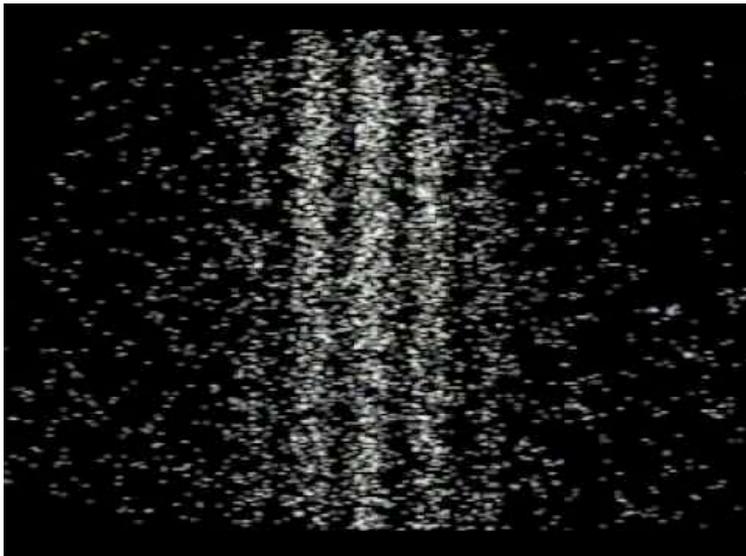
Several relativistic collapse models proposed. No one fully satisfactory so far.

Problem: they must be nonlocal, and that goes directly against relativity.

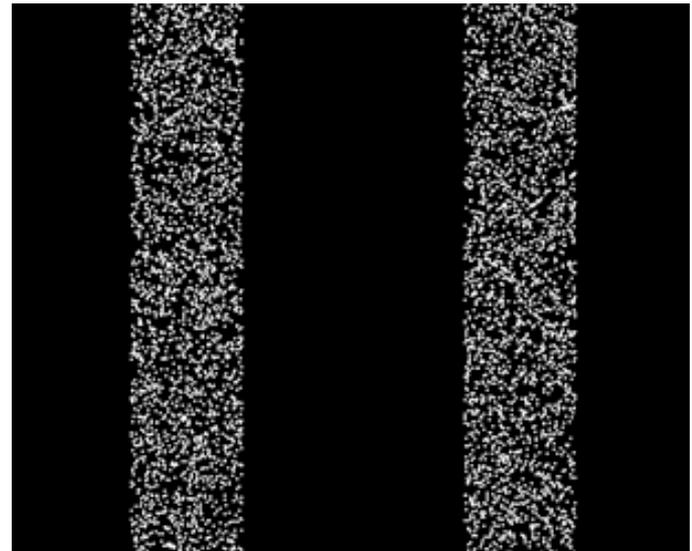
A proposal. Relativity is wrong. It is an approximation. There is a privileged reference frame (the cosmic one). That matches with the idea that the collapse is driven by a cosmological random field.

Experimental tests

The obvious way to test collapse models is with matter-wave interferometry



Prediction of quantum mechanics
(no environmental noise)



Prediction of collapse models
(no environmental noise)

Interferometric Experiments



Atom Interferometry

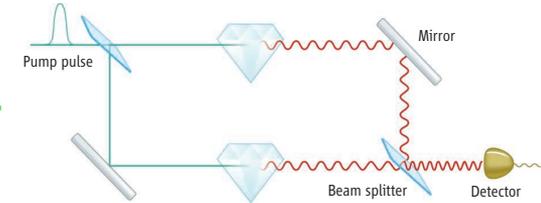
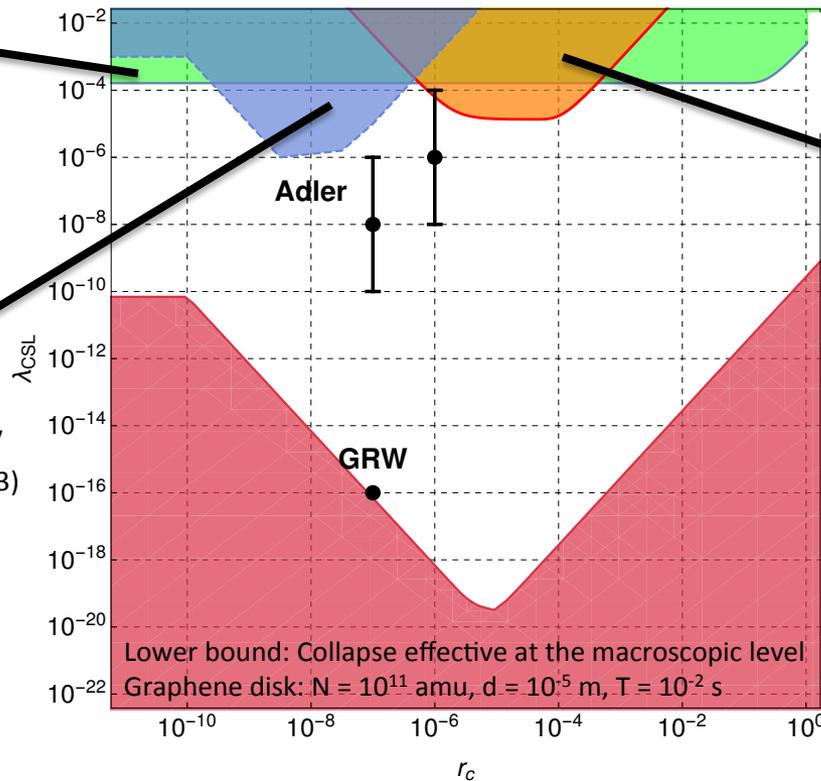
T. Kovachy *et al.*, Nature **528**, 530 (2015)

$M = 87 \text{ amu}$
 $d = 0.54 \text{ m}$
 $T = 1 \text{ s}$

Molecular Interferometry

S. Eibenberger *et al.* PCCP **15**, 14696 (2013)
 M. Toros *et al.*, ArXiv 1601.03672

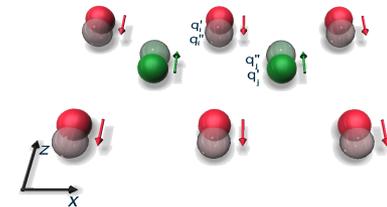
$M = 10^4 \text{ amu}$
 $d = 10^{-7} \text{ m}$
 $T = 10^{-3} \text{ s}$



Entangling Diamonds

K. C. Lee *et al.*, Science. **334**, 1253 (2011).
 S. Belli *et al.*, PRA **94**, 012108 (2016)

$M = 10^{16} \text{ amu}$
 $d = 10^{-11} \text{ m}$
 $T = 10^{-12} \text{ s}$



To improve interferometric tests, it will be necessary to go to micro-gravity environment in outer space. COST Action QTSpace. http://www.cost.eu/COST_Actions/ca/CA15220

Non interferometric tests

M. Bahrami, M. Paternostro, A. Bassi, H. Ulbricht, *Phys. Rev. Lett.* **112**, 210404 (2014). S. Nimmrichter, K. Hornberger, K. Hammerer, *Phys. Rev. Lett.* **113**, 020405 (2014). L. Diósi, *Phys. Rev. Lett.* **114**, 050403 (2015)

The collapse induces a **Brownian motion** on the system

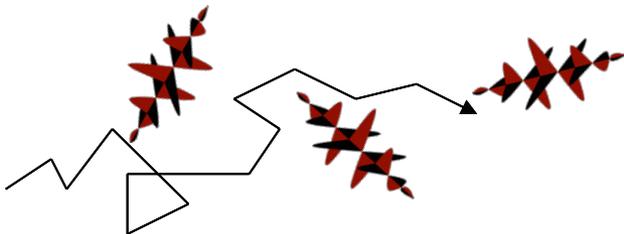


Charged free particle

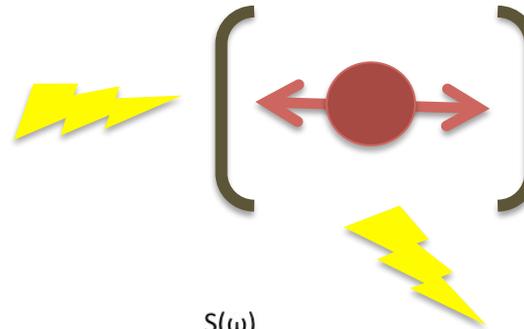
1. Quantum mechanics



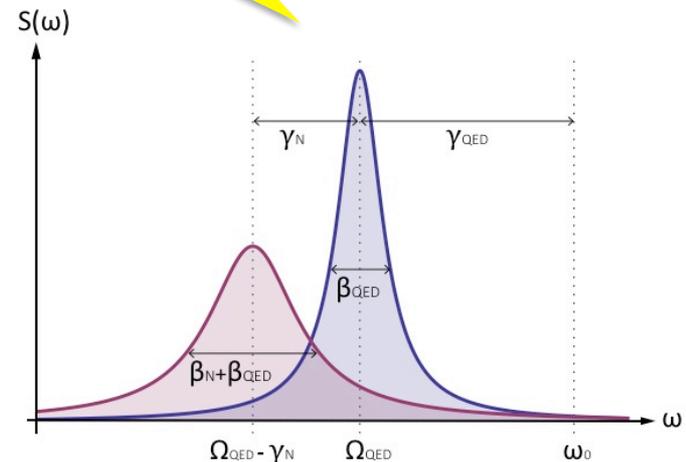
2. Collapse models



Spontaneous photon emission

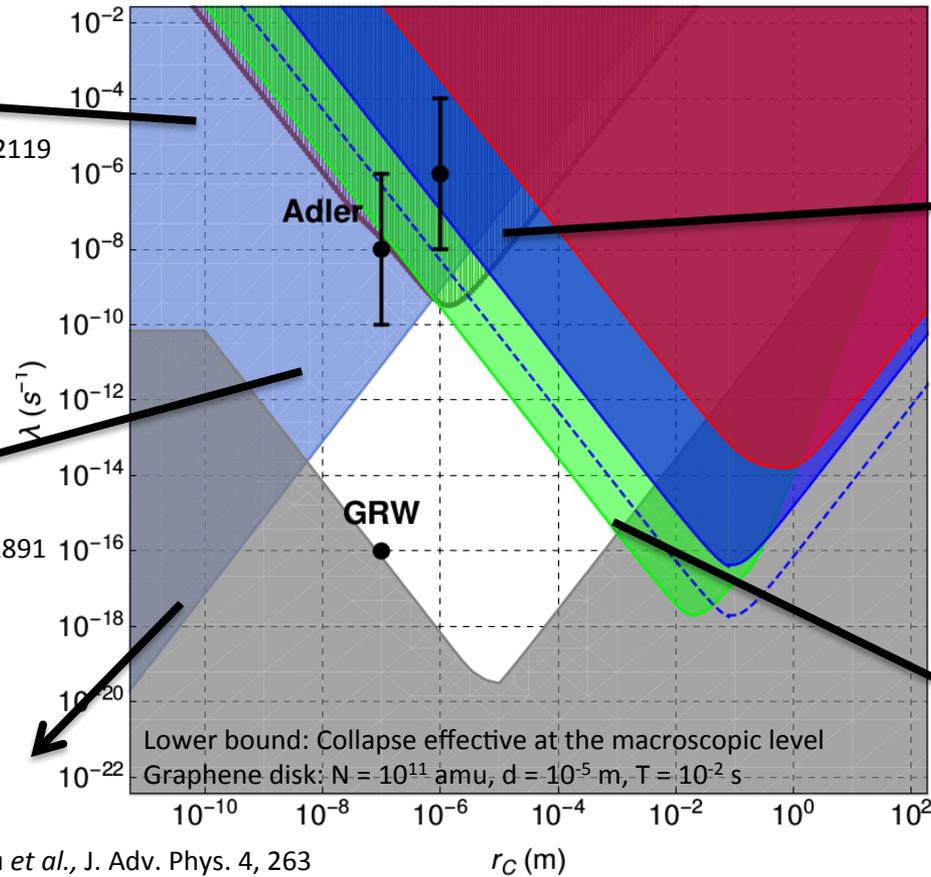
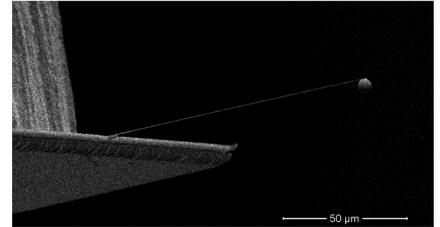


**Nano-particle
in a optical
cavity**



Extra shift and broadening

Non-Interferometric Experiments



Cold Atom Gas

F. Laloë *et al.*, *Phys. Rev. A* 90, 052119 (2014)

$M = 87$ amu
 $T = 30$ s

Cold Atom Gas

M. Bilardello *et al.*, *ArXiv* 1605.01891

$M = 87$ amu
 $T = 1$ s

X-Ray

C. Curceanu *et al.*, *J. Adv. Phys.* 4, 263 (2015).

$M = 1$ Kg
 $T =$ days / months

Cantilever

A. Vinante *et al.*, *Phys. Rev. Lett.* 116, 090402 (2016)

$M = 10^{14}$ amu
 $T = \infty$

Gravitational wave detectors

M. Carlesso *et al.*, *ArXiv* 1606.04581

$M = 10^{30}$ amu
 $T = \infty$

Acknowledgments

The Group (www.qmts.it)

- Postdocs: S. Donadi, F. Fassioli, A. Grossardt
- Ph.D. students: G. Gasbarri, M. Toros, M. Bilardello, M. Carlesso, S. Bacchi, L. Curcuraci
- Graduate students: A. Rampichini, T. Guaita, M.M. Marchese

Collaborations with M. Paternostro, H. Ulbricht, A. Vinante, C. Curceanu.



www.units.it

UNIVERSITÀ
DEGLI STUDI DI TRIESTE



www.cost.eu



Istituto Nazionale
di Fisica Nucleare

www.infn.it