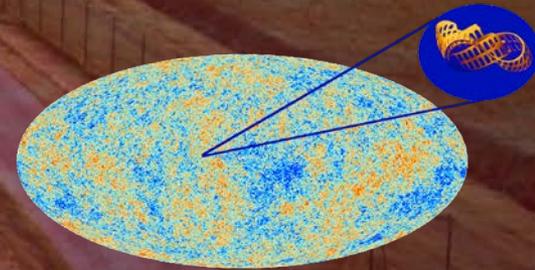


The search for a stochastic background of gravitational waves

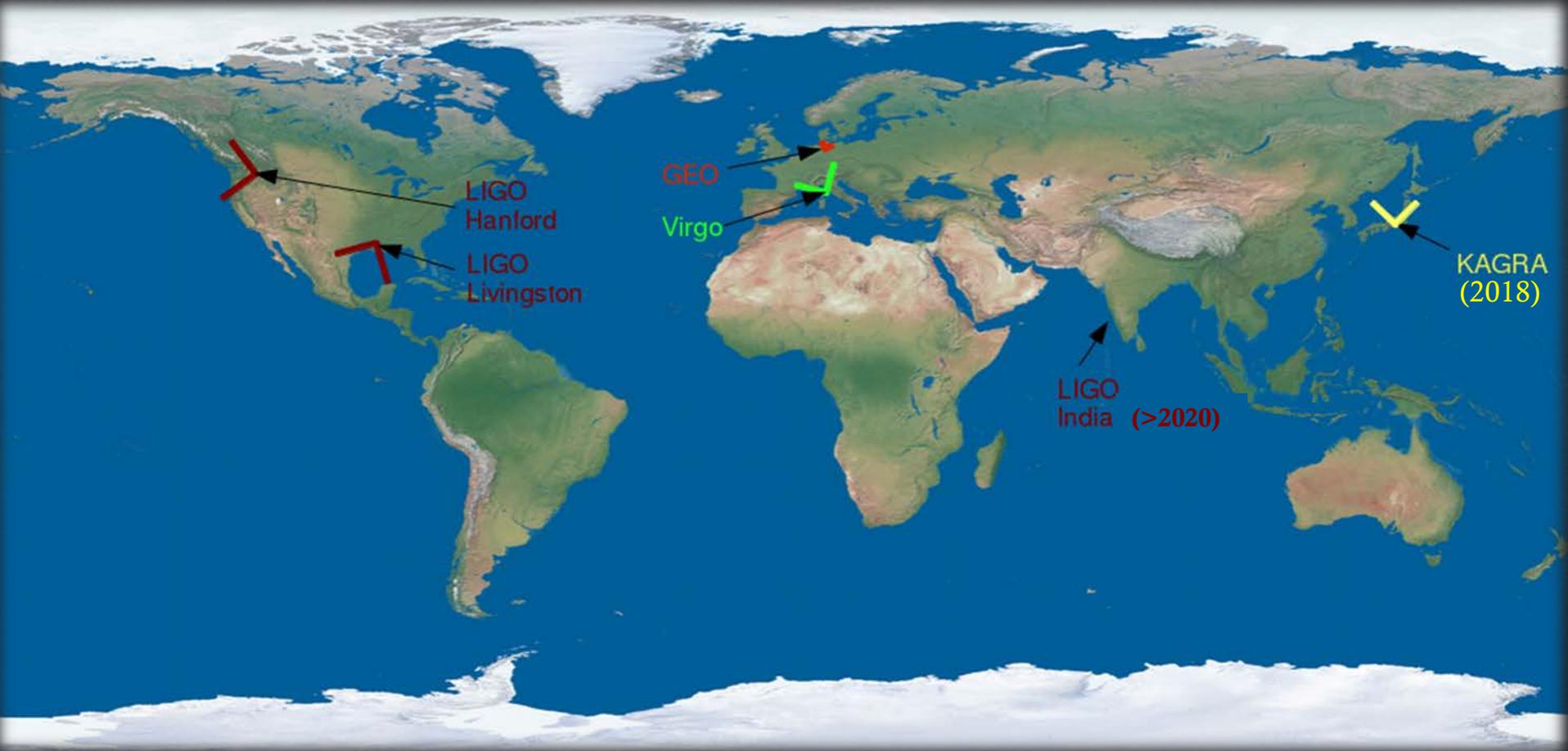


Workshop «String Theory and Inflation»
20 September 2016

Physics Department, University of Rome Tor Vergata

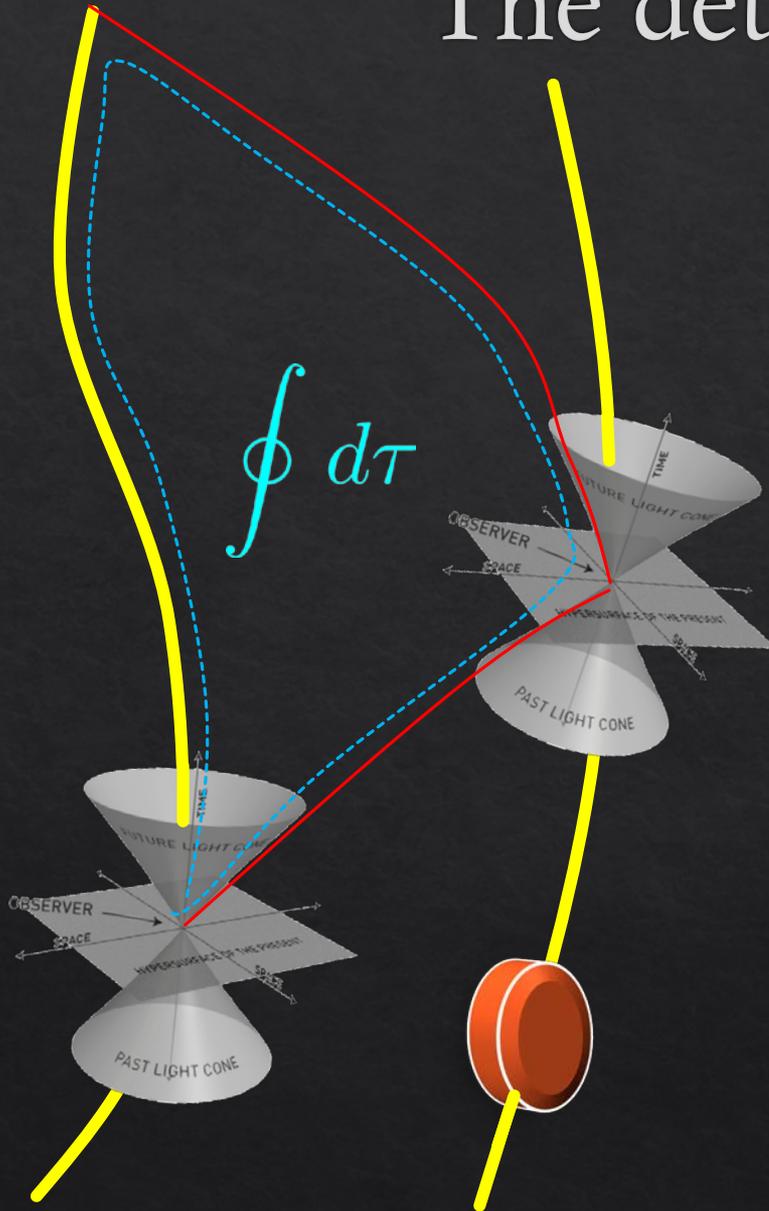
Giancarlo Cella – INFN Pisa & Virgo collaboration

The Gravitational Wave Detectors Network



- ◇ Currently LIGO-H & LIGO-V operative (first scientific run ended, now stopped until the end of the year for upgrade)
- ◇ Virgo will join at the next scientific run.
- ◇ This will be important to improve sensitivity, coverage and estimation of parameters

The detector principle



◇ Description can be coordinate dependent

◇ Physical observable is not

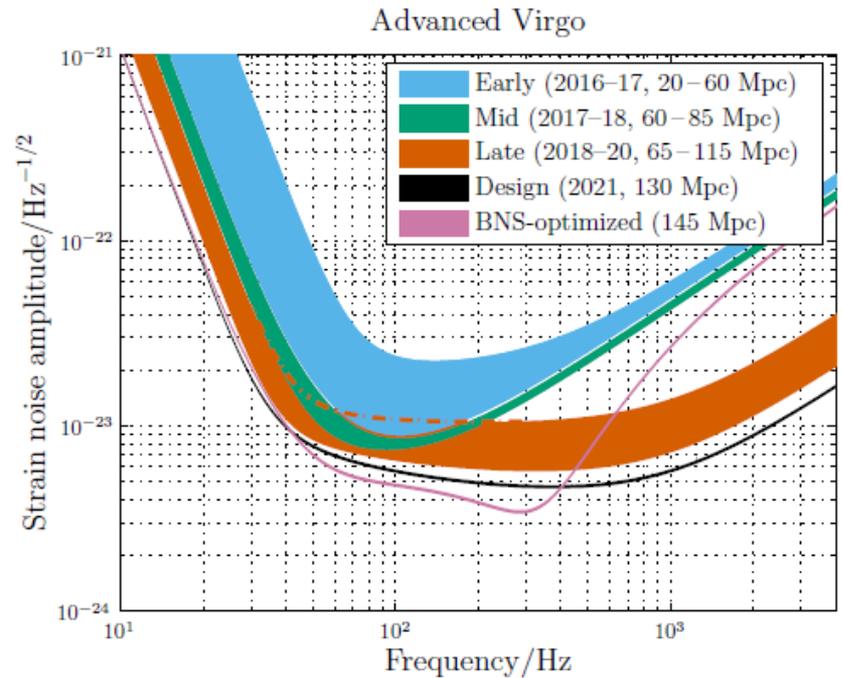
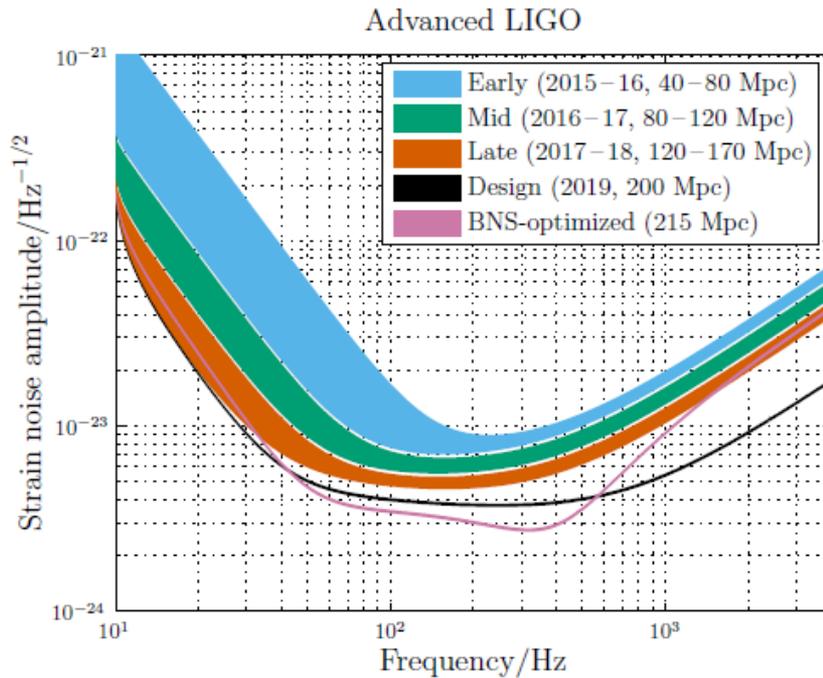
◇ Intuitive picture (when $\lambda_{GW} \gg FL$): tidal force

$$F_i = \frac{1}{2} m \frac{d^2 h_{ij}^{TT}}{dt^2} L_j$$

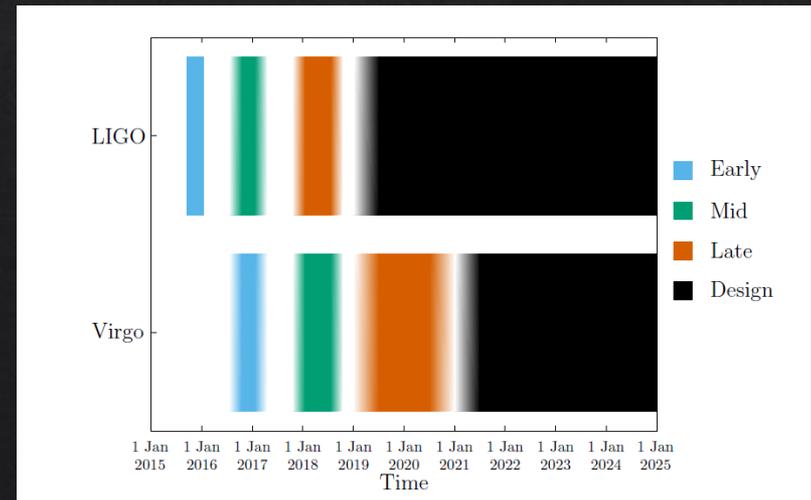
on the end mirrors.



The observing scenario



- ◇ Plan: a series of scientific runs with intermediate commissioning interruptions
- ◇ Sensitivity will increase in steps toward the design one
- ◇ Quite successful until now....



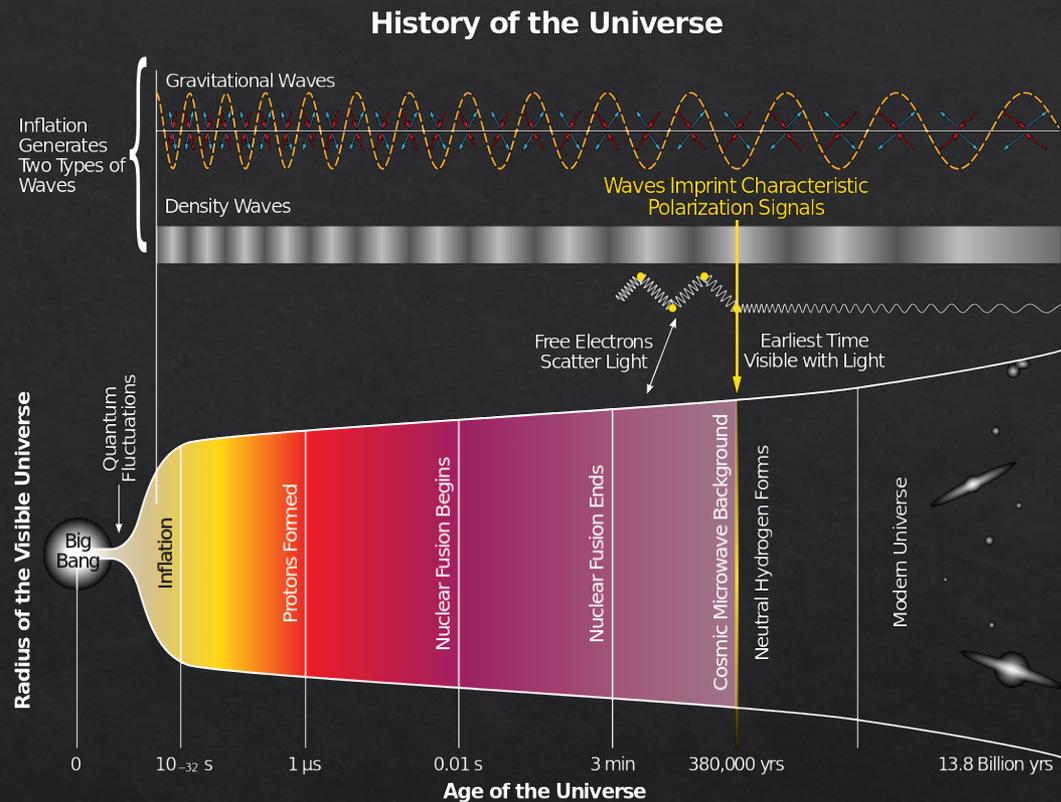
Stochastic background in a nutshell

A stochastic background can be seen as

- a GW field which evolves from an initially random configuration
- the result of a superposition of many uncorrelated and unresolved sources

Two different kinds:

- **Cosmological:**
 - signature of the early Universe: coupling of gravitational field is small!
 - *inflation, cosmic strings, phase transitions...*
- **Astrophysical:**
 - sources since the beginning of stellar activity
 - *compact binaries, supernovae, rotating NSs, core-collapse to NSs or BHs, supermassive BHs...*



Typical «first approximations» :

- 1) Gaussian, because sum of many contributions
- 2) Stationary, because physical time scales much larger than observational ones
- 3) Isotropic (at least for cosmological backgrounds)
- 4) Unpolarized

If these are true, SB is completely described by its power spectrum

Description (simplest model)

1. Correlation between GW modes

$$\langle h_A^*(f, \hat{\Omega}, \psi) h_B(f', \hat{\Omega}', \psi') \rangle = \delta_{AB} \delta(f-f') \frac{\delta^2(\hat{\Omega}, \hat{\Omega}') \delta(\psi - \psi')}{4\pi} \frac{1}{2\pi} S_h(f)$$

2. Connection with GW energy density

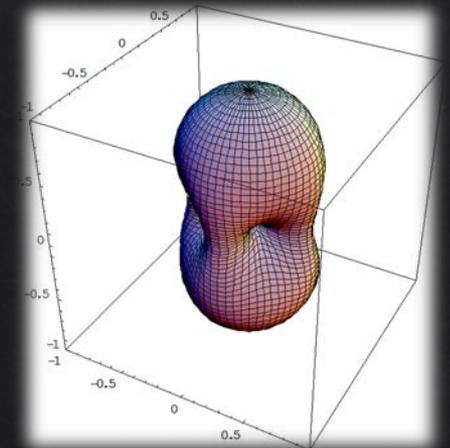
$$h_0^2 \Omega_{gw}(f) = \frac{1}{\rho_c} \frac{d\rho_{gw}}{d \log f} = \frac{4\pi^2 h_0^2}{3H_0^2} f^3 S_h(f)$$

3. Strain at the detector: sum over modes

$$h_{ij}(t, \vec{r}) = \sum_{P=+, \times} \int_{S^2} d\hat{\Omega} \varepsilon_{ij}^P(\hat{\Omega}) \int_{-\infty}^{\infty} df \tilde{h}_P(f, \hat{\Omega}) e^{i2\pi f(t - \hat{\Omega} \cdot \vec{r})}$$

4. Signal at the detector: projection on the detector's tensor

$$h(t, r) = D^{ij} h^{ij}(t, r)$$



How it is possible to (directly) detect it

◊ We have a vector-valued Gaussian stochastic process

$$\diamond \underline{x}_i = \left(x_i^{\text{Virgo}}, x_i^{\text{Hanford}}, \dots \right) \quad dP = \mathcal{N} \prod_f \exp \left(-\frac{1}{2} \tilde{\underline{x}}_f^+ C^{-1}(f) \tilde{\underline{x}}_f \right) d\tilde{\underline{x}}_f$$

$$\underline{x}_i = \underline{h}_i + \underline{n}_i \quad \langle \underline{n}_i \otimes \underline{n}_j \rangle = I c_{ij}$$

◊ We must discriminate between two hypothesis:

$$\mathcal{H}_0 \quad C = C_N = \begin{pmatrix} S_N^{\text{Virgo}} & 0 & \dots & 0 \\ 0 & S_N^{\text{Hanford}} & \dots & 0 \\ \vdots & \vdots & \ddots & 0 \\ 0 & 0 & 0 & S_N^{\text{Kagra}} \end{pmatrix}$$

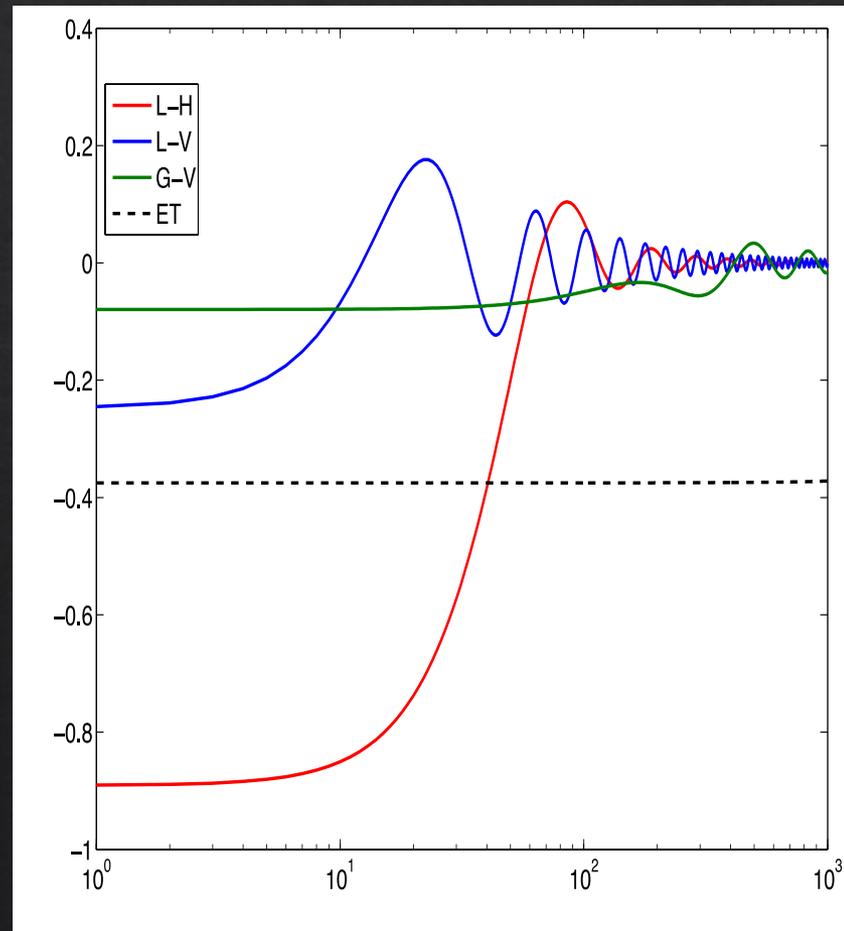
$$\mathcal{H}_1 \quad C = C_N + S_h \begin{pmatrix} \gamma_{V,V} & \gamma_{V,H} & \dots & \gamma_{V,K} \\ \gamma_{V,H} & \gamma_{H,H} & \dots & \gamma_{H,K} \\ \vdots & \vdots & \ddots & \dots \\ \gamma_{V,K} & \gamma_{H,K} & \dots & \gamma_{K,K} \end{pmatrix}$$

◊ Optimal statistic (two detectors):

$$Y = \lambda \int df \frac{\tilde{x}_1^*(f) \tilde{x}_2(f) \gamma_{12}(f) S_h(f)}{S_{n,1}(f) S_{n,2}(f)}$$

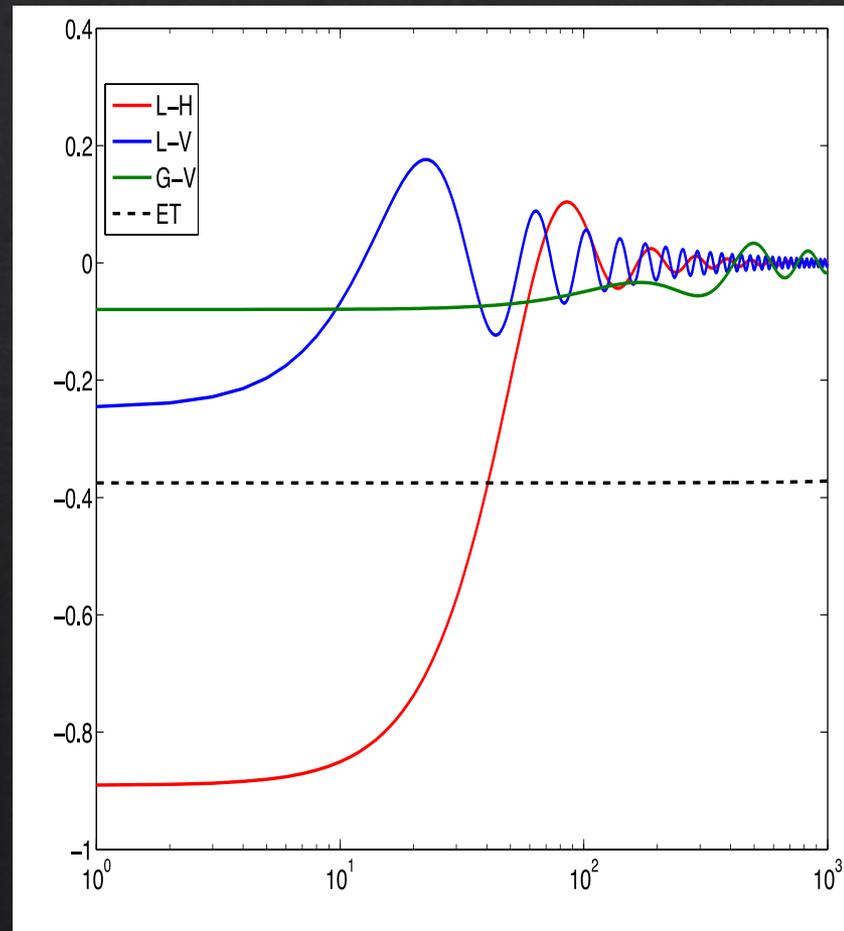
Overlap reduction function (a.k.a. coherence)

$$\text{SNR}_Y^2 := \frac{\mu_Y^2}{\sigma_Y^2} = 2T \int_0^\infty S_h^2(f) \frac{\gamma_{12}^2(f)}{S_{n,1}(f)S_{n,2}(f)} df$$



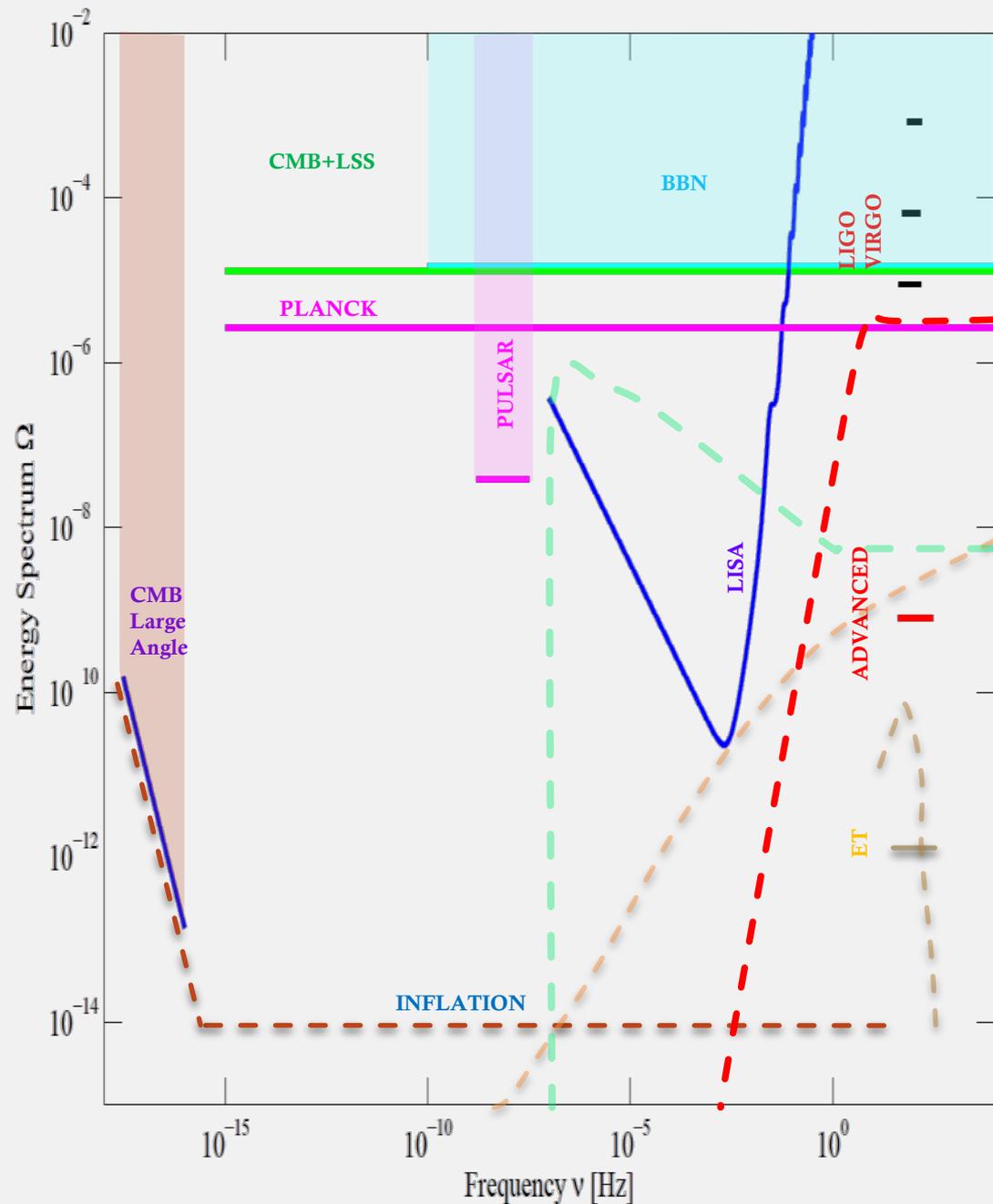
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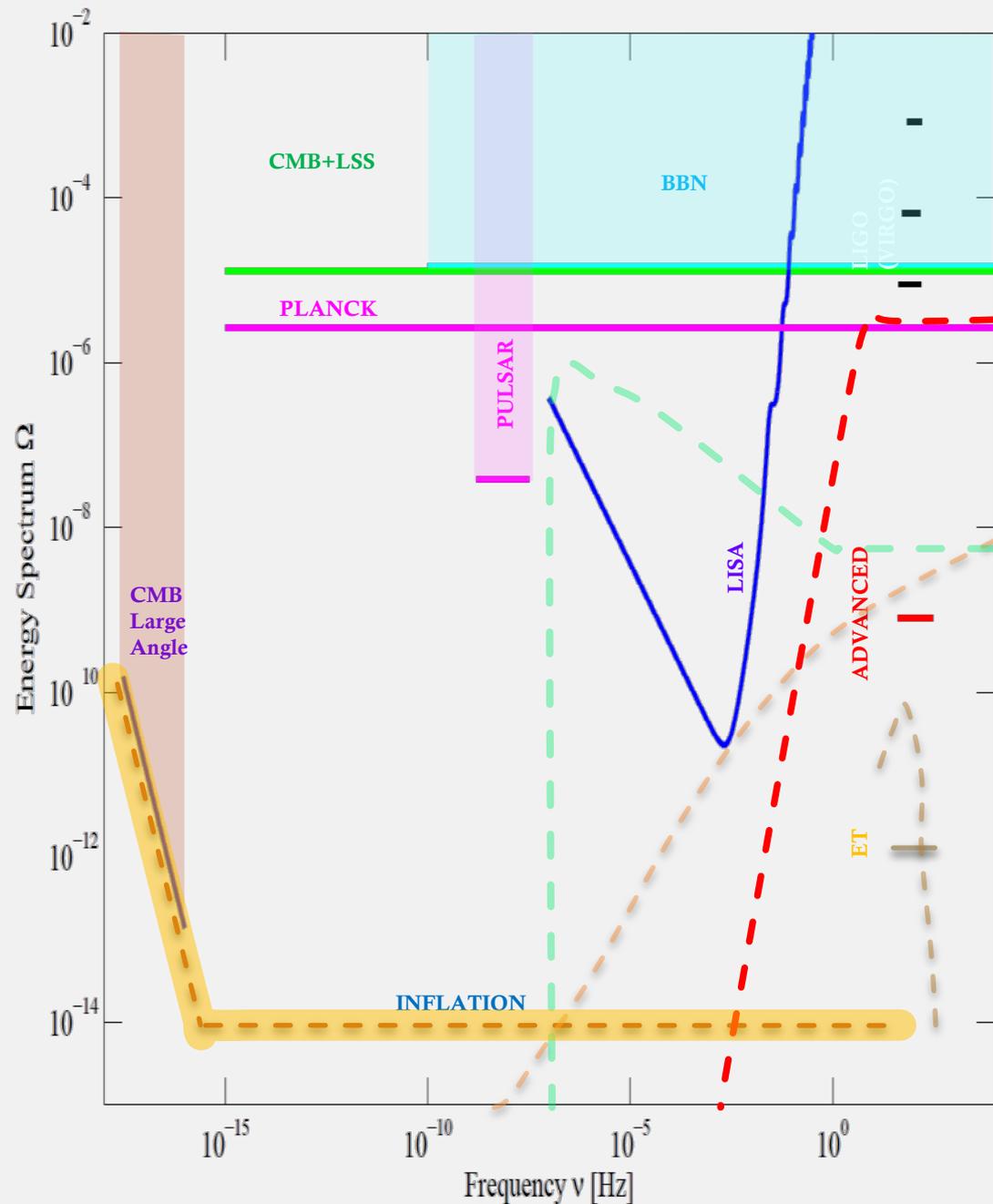
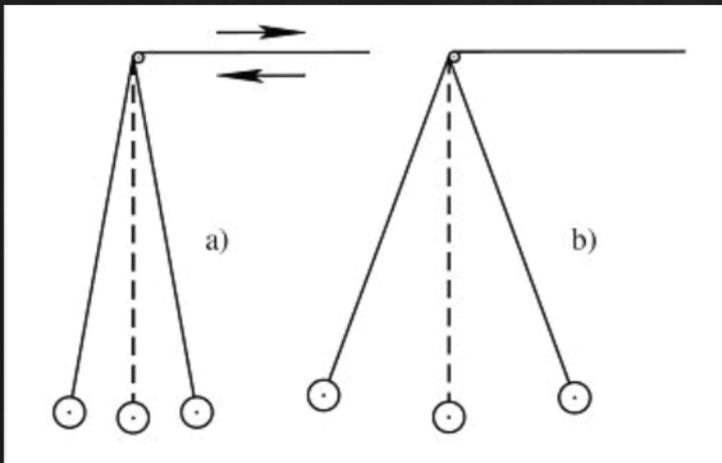
Upper limits & expected sensitivities

- **Big-Bang Nucleosynthesis model and observations** constrain the total GW energy at the time of BBN (integral bound)
- Similar bound from **CMB observations**
- Too much GW gives too much **large angle anisotropy** by the Sachs Wolfe effect
- Signal from **millisecond pulsars** works as a (big) interferometer:



Inflation

- Parametric amplification of vacuum quantum fluctuations
- Standard inflationary models are weakly dependent on frequency
- Tight bound: CMB large angle anisotropy
- Out of reach of advanced detectors by 5 orders of magnitude

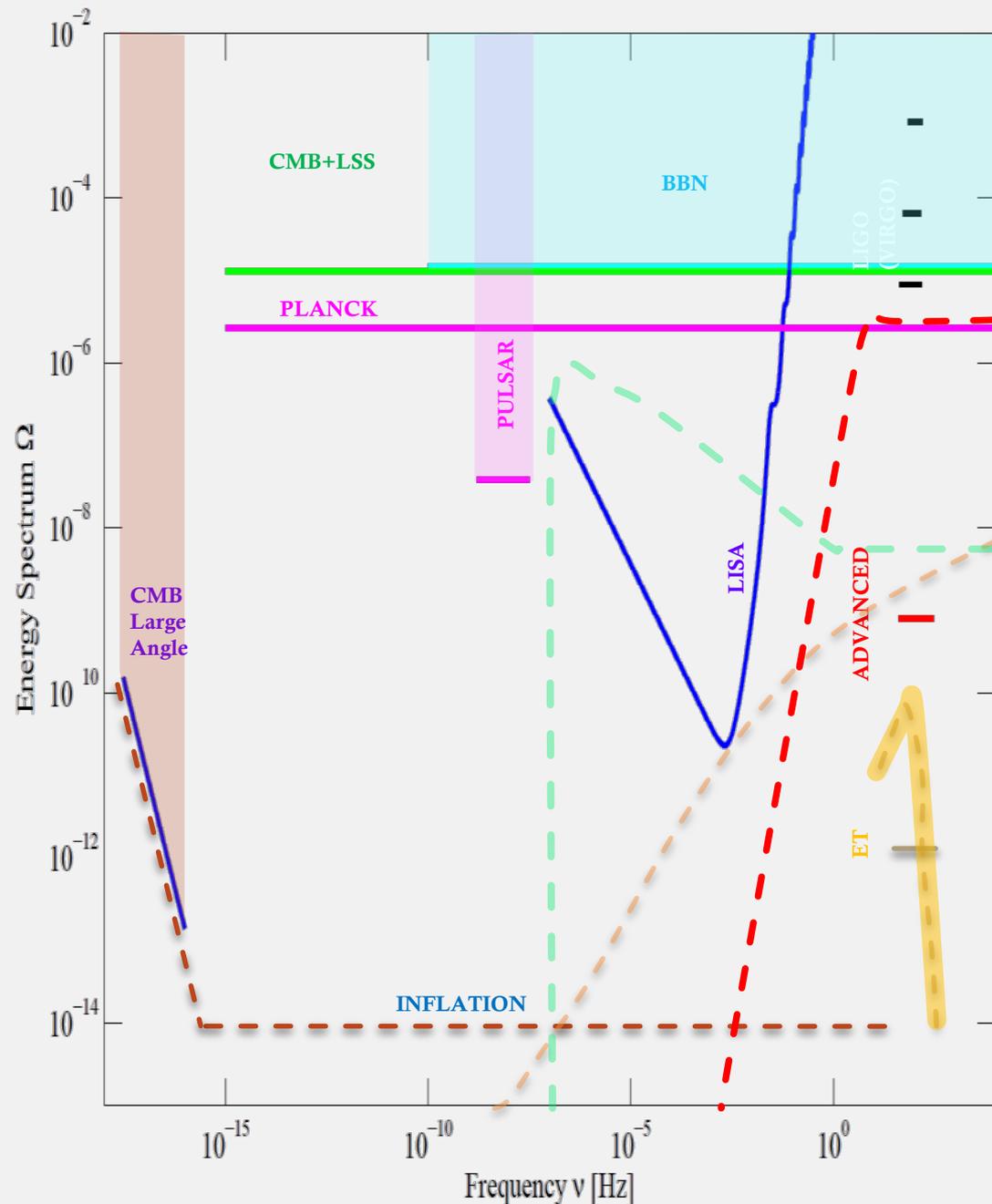


Resonant preheating



- During a resonant preheating phase at the end of the inflation, inflaton energy can be transferred efficiently to other particles
- Can produce a significant GW background
- Spectrum peak depends on energy scale (here 10 GeV, higher frequencies for a larger scale)

Easter & Lim, JCAP 0604, 010 (2006)
Easter et al., PLR 99, 221301 (2007)
Easter, Nucl. Phys. Proc. Suppl. 194, 33 (2009)

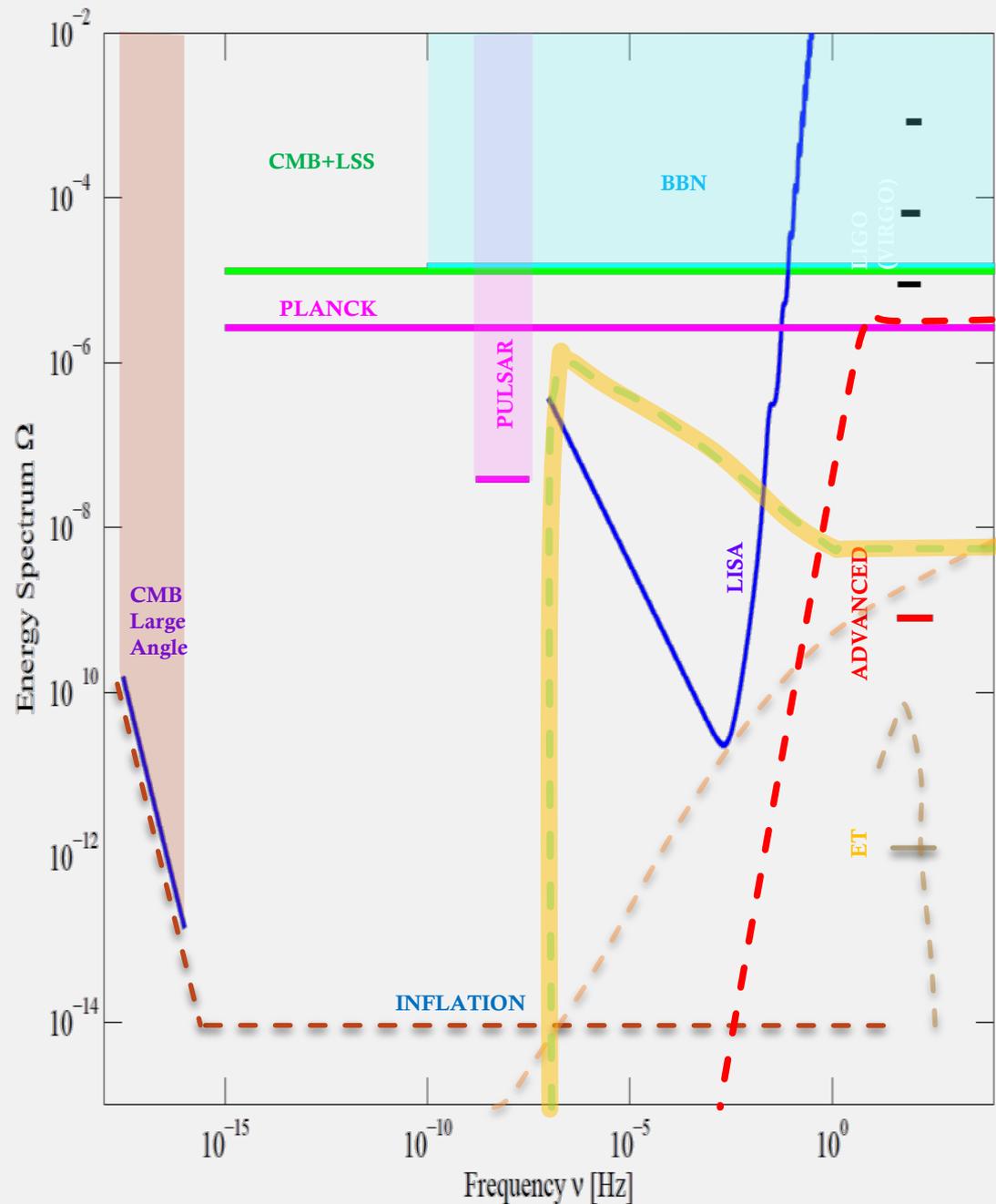


Cosmic (super)string models

- Dynamical network:
 - Strings entering in the horizon
 - Interconnection: loops generation
 - Radiation (GW and other fields): loop destruction
- Most efficient emission mechanisms: cusps and kinks
- Integrating over the whole universe leads to a GW background
- Large parameter space



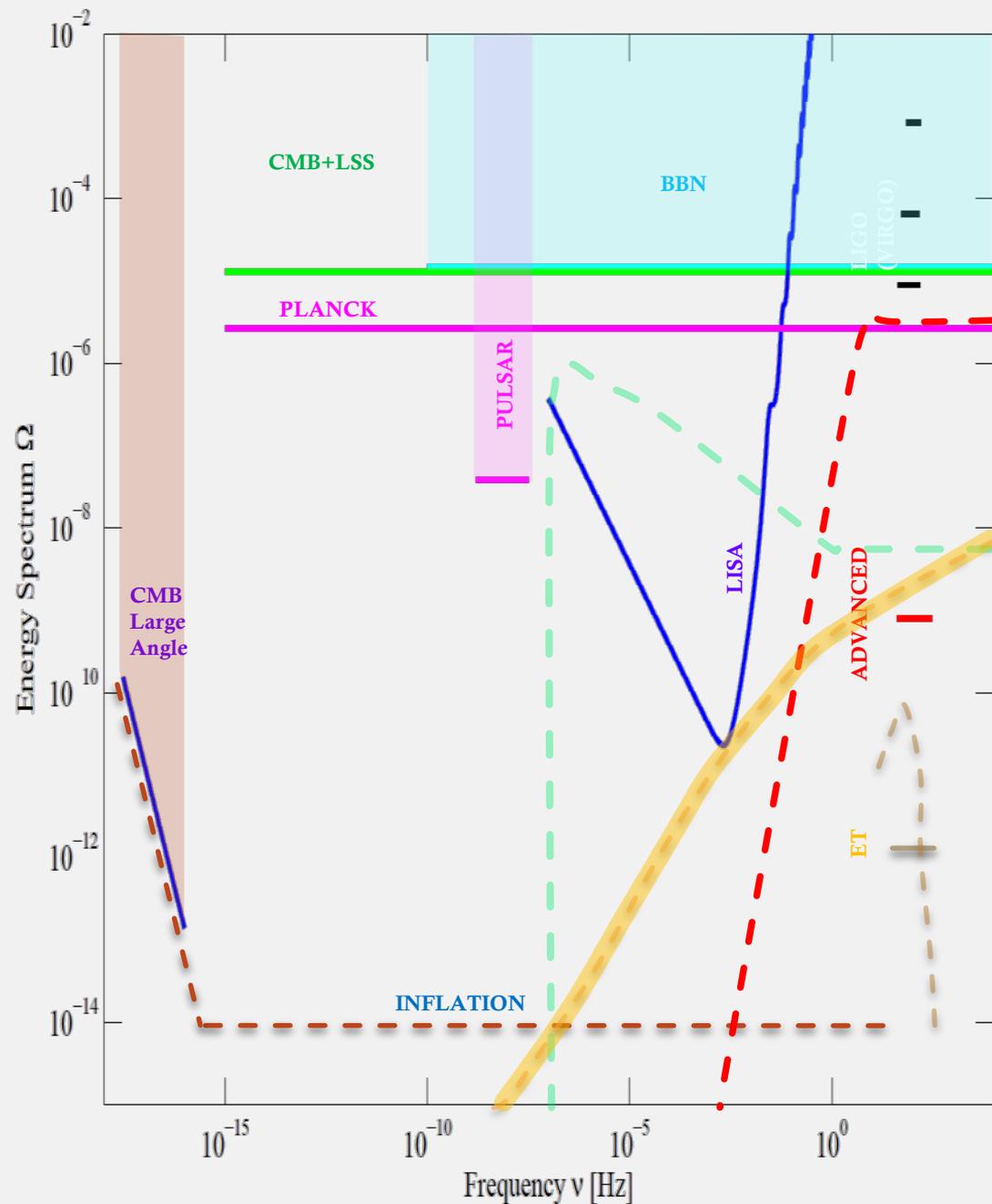
Damour & Vilenkin, PRL 85, 3761 (2000)
 Siemens et al., PRL 98, 111101 (2007)
 Olmez et al., PRD 81, 104028 (2010)



Axion-based inflation models

- Models include axion-gauge couplings
- Gauge backreaction on the inflaton extends inflation
- The late inflationary phase increases GW production at high frequencies

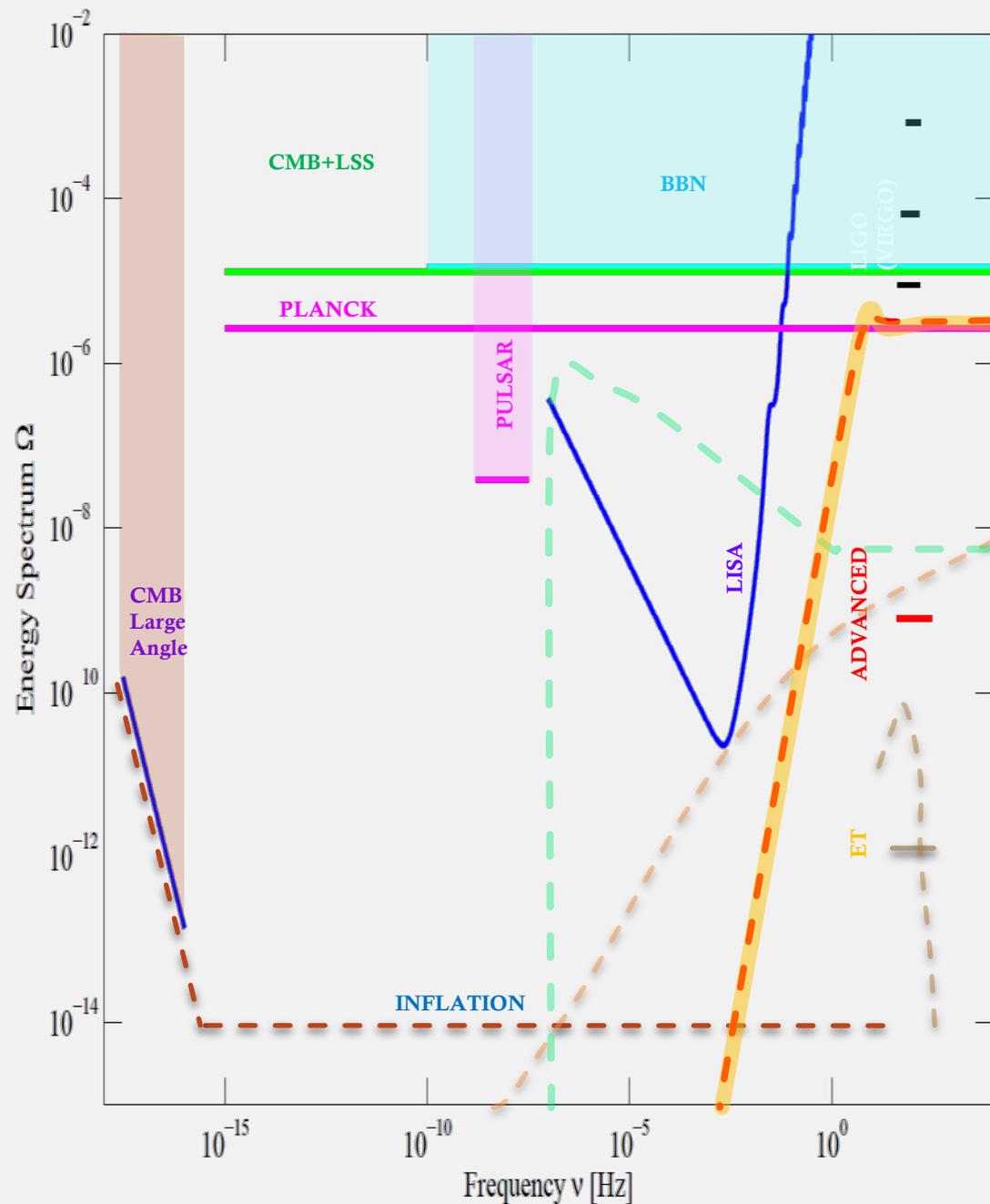
Barnaby, Pajer and Peloso - Phys. Rev. D 85, 023525



Pre-BB models

- Alternative cosmologies
- Evade the CMB large angle anisotropy bound
- Evade the BBN-CMB integral bound
- Can be significant at Virgo/LIGO frequencies

Gasperini & Veneziano, Phys. Rep. 373, 1 (2003)
Buonanno et al., PRD 55, 3330 (1997)



Isotropic upper limits

$$\Omega_{GW}(f) = \Omega_{\alpha} \left(\frac{f}{f_{ref}} \right)^{\alpha}$$

Best current upper limit

Frequency (Hz)	f_{ref} (Hz)	α	Ω_{α}	95% C.L. upper limit	Previous limits
41.5–169.25	...	0	$(-1.8 \pm 4.3) \times 10^{-6}$	5.6×10^{-6}	7.7×10^{-6}
170–600	...	0	$(9.6 \pm 4.3) \times 10^{-5}$	1.8×10^{-4}	...
600–1000	900	3	0.026 ± 0.052	0.14	0.35
1000–1726	1300	3	-0.077 ± 0.53	1.0	...

PRL 113, 231101 (2014)

Colocated interferometers

Band (Hz)	95% C.L. UL ($\times 10^{-3}$)
460–1000	0.77
460–537	1.11
537–628	2.12
628–733	1.18
733–856	2.53
856–1000	2.61

O_1 analysis
(advanced detectors)
is in progress.

Phys. Rev. D 91, 022003 (2015)

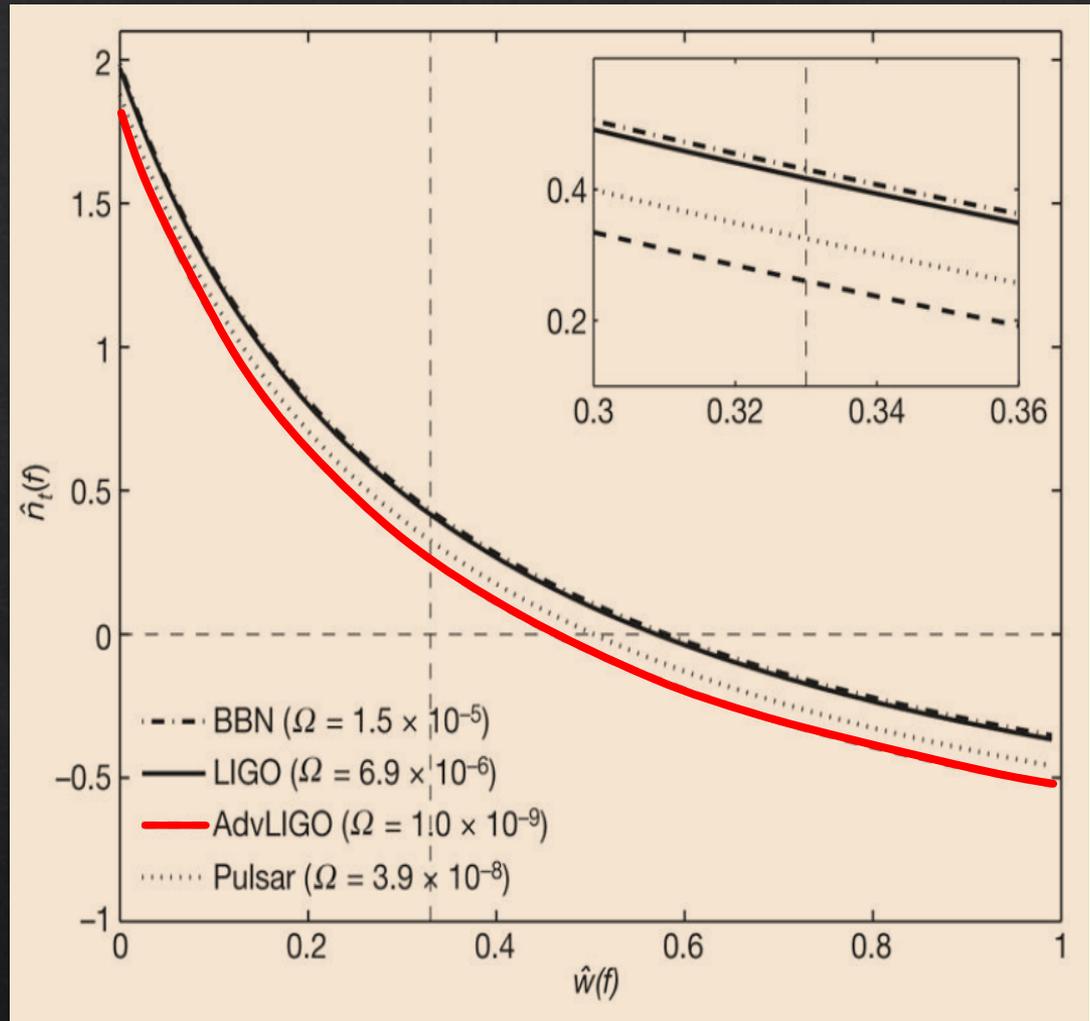
Constraint on early universe state equation

Spectrum
parameterization:

$$\Omega_{GW}(f) = A f^{\hat{\alpha}(f)} f^{\hat{n}(f)} r$$

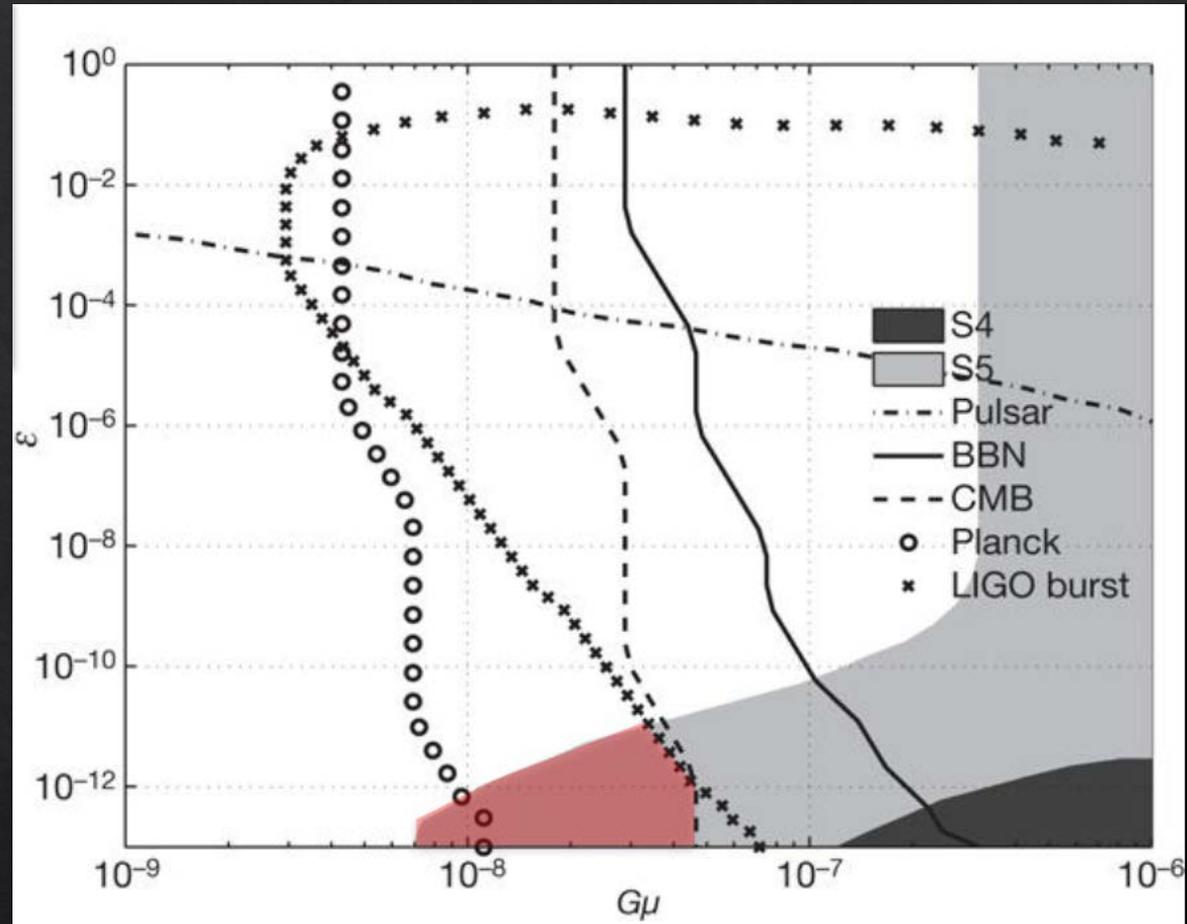
$$\hat{\alpha}(f) = 2 \frac{3\hat{w}(f) - 1}{3\hat{w}(f) + 1}$$

- $\hat{n}_t(f)$ effective tensor tilt parameter
- r ratio of tensor and scalar perturbation amplitudes (here $r = 0.1$)
- $\hat{w}(f)$ equation of state parameter



Constraint on cosmic strings models

- Network of cosmic strings parameterized by:
- String tension μ
- Reconnection probability p
- Loop size (parameterized by ε)
- $G\mu < 10^{-6}$ (CMB observations)
- ε unconstrained
- $10^{-4} < p < 1$ ($p = 10^{-3}$ here)



Region excluded: entire plane will be probed by advanced detectors.

LSC/Virgo Collaboration, Nature 460, 990-994 (2009)
PRD 80, 062002 (2009)
PRL 98, 111101 (2007)

Constraints on pre-Big-Bang models

Spectrum:

$$\Omega_{GW}(f) \propto f^3$$

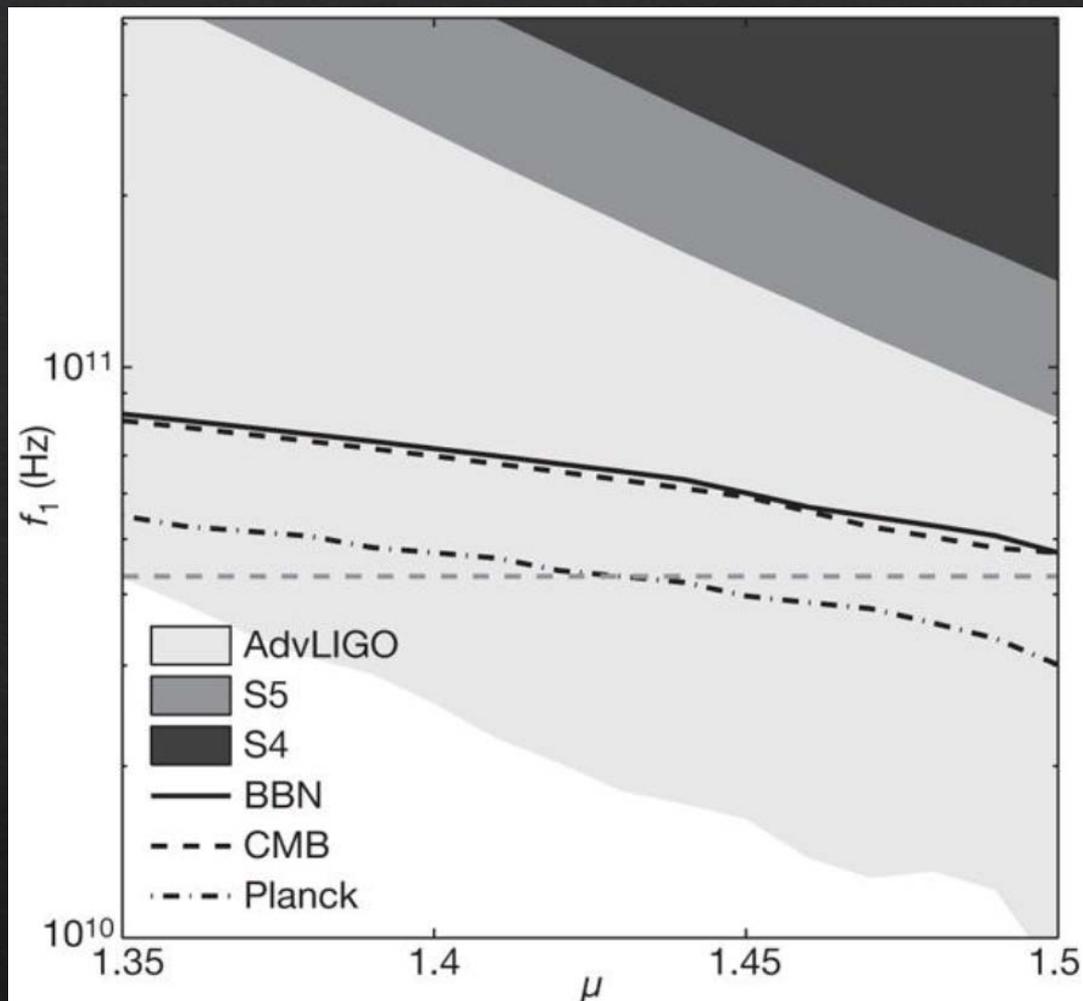
below f_s , $f_s = 30$ Hz here

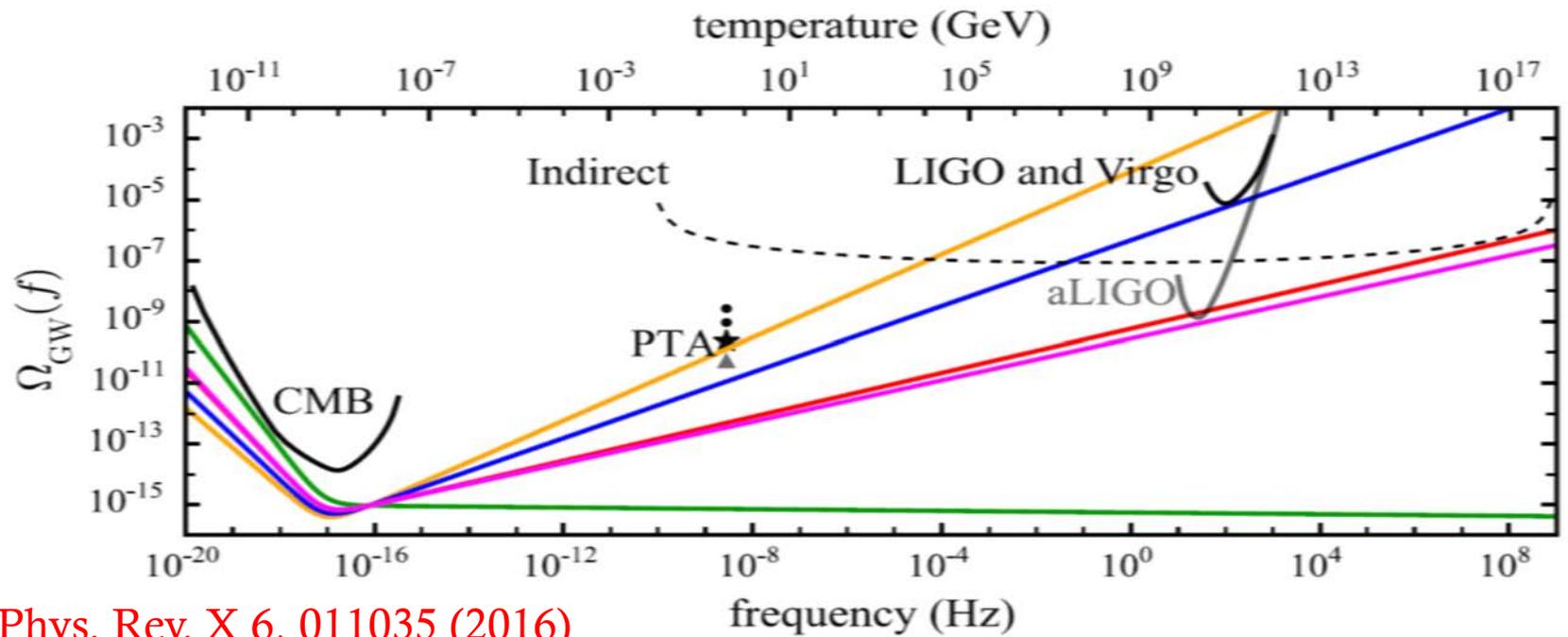
$$\Omega_{GW}(f) \propto f^{3-2\mu}$$

above f_s , $f_s = 30$ Hz here

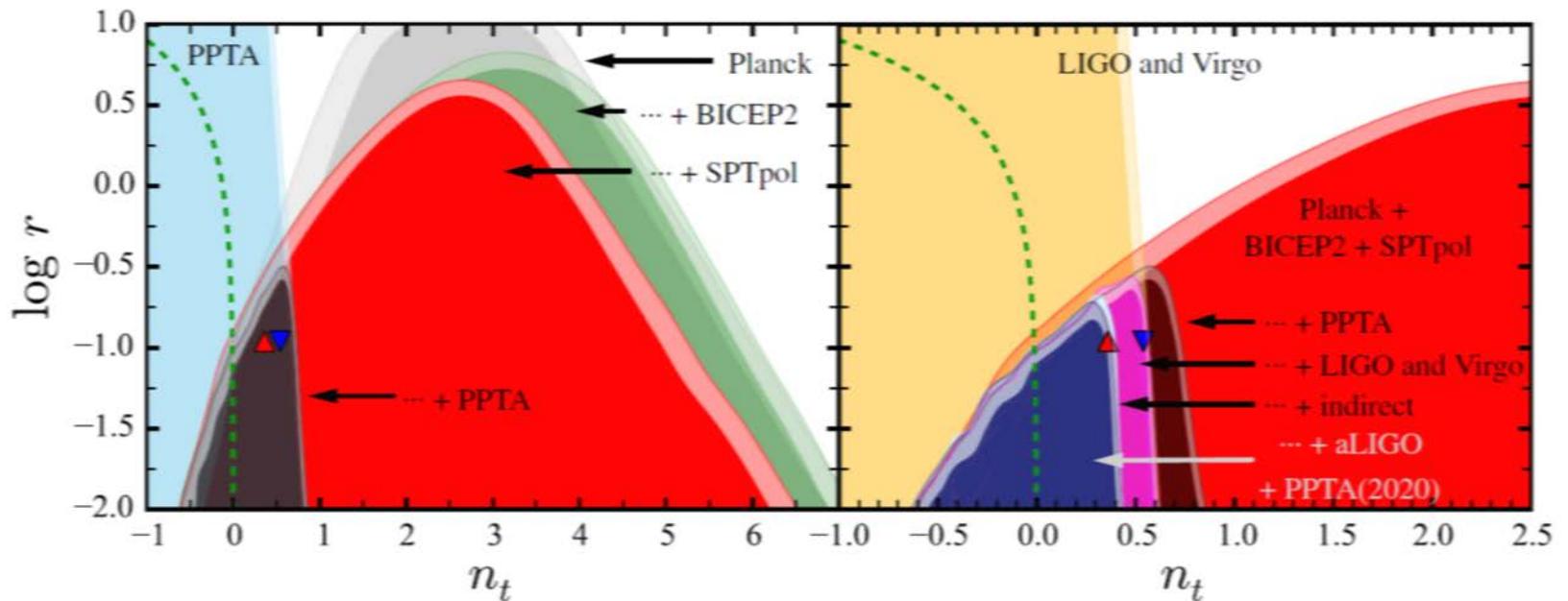
f_1 : cut off frequency

(a factor 10 from 4.3×10^{10})





Phys. Rev. X 6, 011035 (2016)



Astrophysical Stochastic Background

Core collapse supernovae

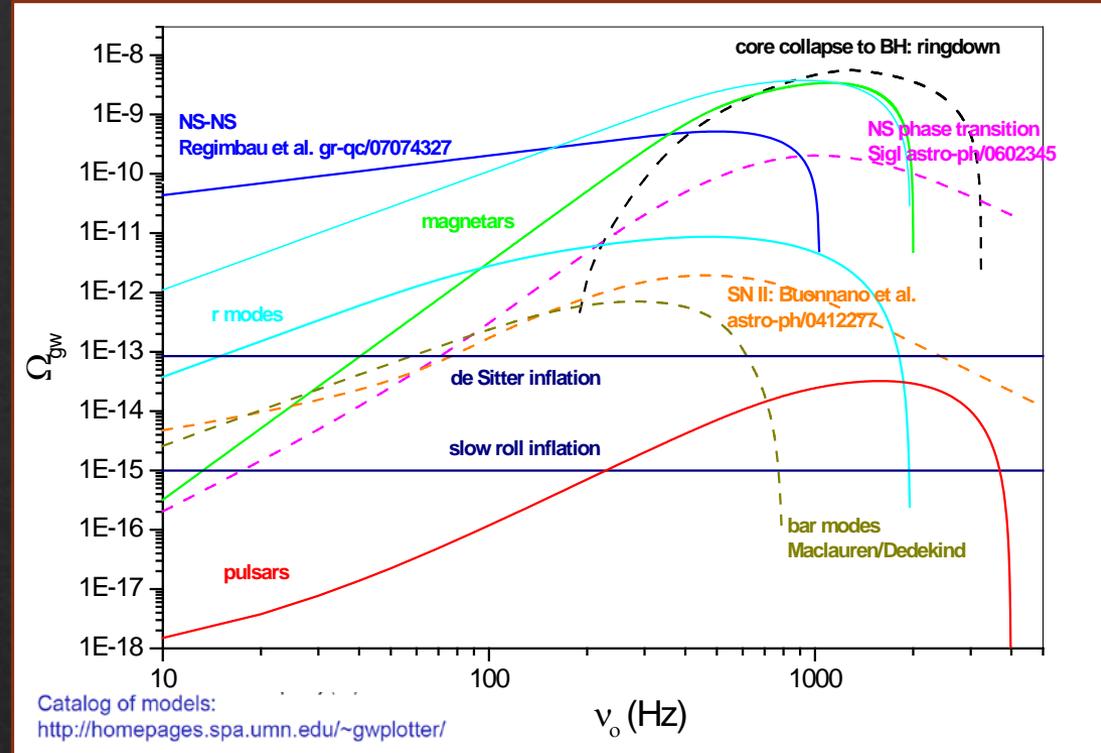
- **Neutron star formation:** *Blair & Ju 1996, Coward et al. 2001-02, Howell et al. 2004, Buonanno et al. 2005*
- **Stellar Black Hole formation:** *Ferrari et al. 1999, de Araujo et al. 2000-04*

Neutron stars

- **tri-axial emission:** *Regimbau & de F. Pacheco 2001-06*
- **bar or r-modes:** *Owen et al. 1998, Ferrari et al. 1999, Regimbau 2001*
- **phase transitions:** *Sigl 2006*

$$\Omega_{gw}(f) = \int P(\underline{\theta}) \Omega_{gw}(f, \underline{\theta}) d\underline{\theta}$$

$$\Omega_{gw}(f, \underline{\theta}) = \frac{f}{\rho_c} \int \frac{dR}{dz} (z, \underline{\theta}) \frac{1}{4\pi r^2(z)} \frac{dE_{gw}}{df} (\underline{\theta}, f(1+z)) dz$$



Stellar Compact Binaries

- **near coalescence (NS, BH):** *Regimbau et al. 2006-07, Coward et al. 2005 (BNS), Howell et al. 2007 (BBH)*
- **low frequency inspiral phase:** *Ferrari et al. 2002, Farmer & Phinney 2002, Cooray 2004 (WD-NS)*

Capture of compact objects by SMBHs :

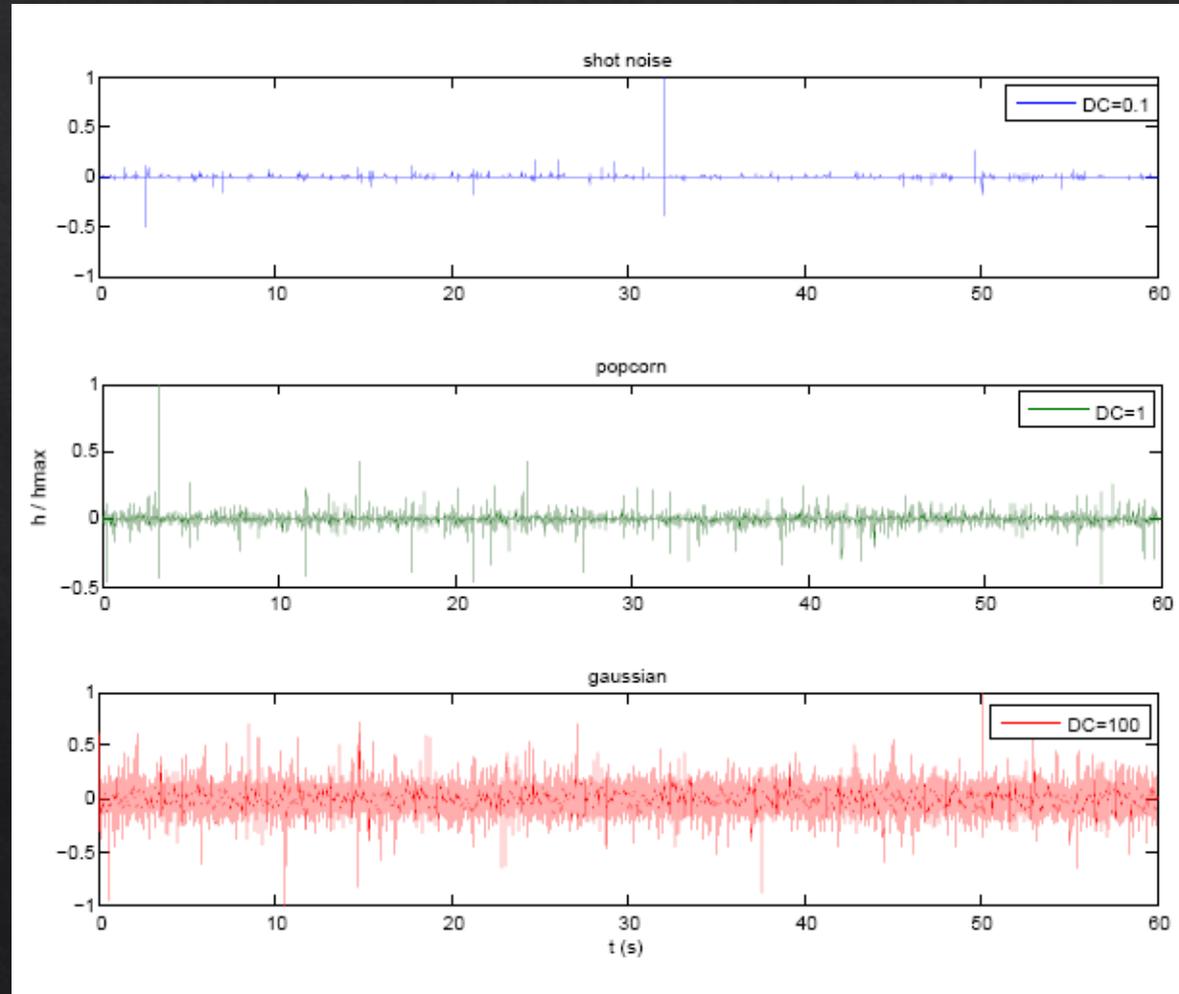
Barack & Cutler 2004

Astrophysical Stochastic Background

$$D(z) = \int_0^z \tau (1 + z') \frac{dR}{dz'} dz'$$

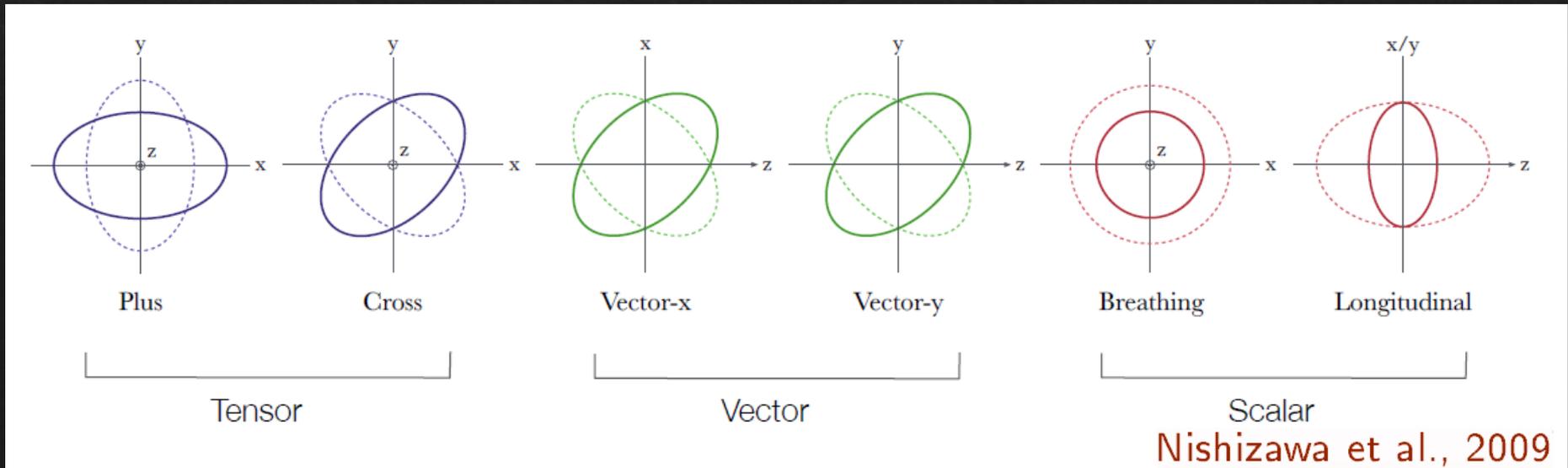
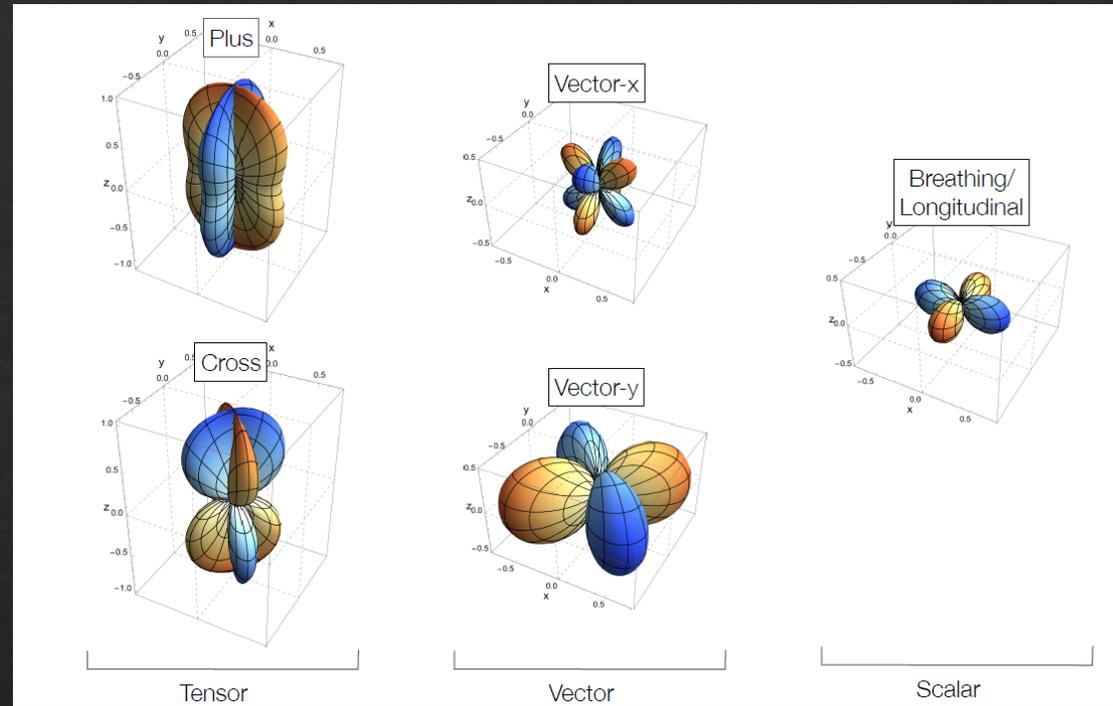
Duty cycle: ratio between the observed typical duration of the event and the average time between events.

- $D \ll 1$: resolved sources
 - Burst data analysis, optimal filtering
- $D \simeq 1$: «popcorn noise»
 - Maximum likelihood statistic (Drasco et al. 2003)
 - Probability event horizon (Coward et al. 2005)
- $D \gg 1$: Gaussian stochastic background
 - Cross correlation statistic (isotropic/anisotropic)

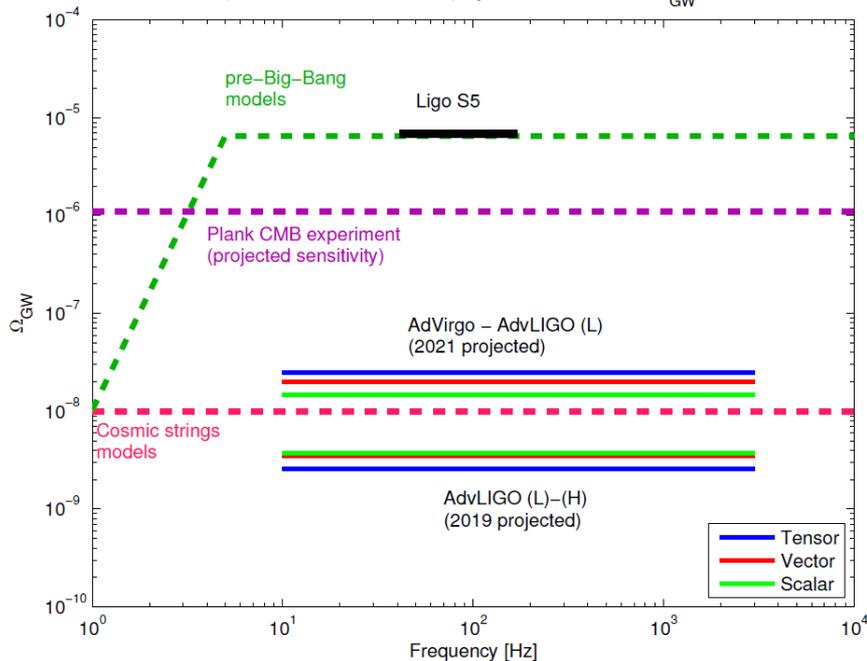


Non standard polarizations

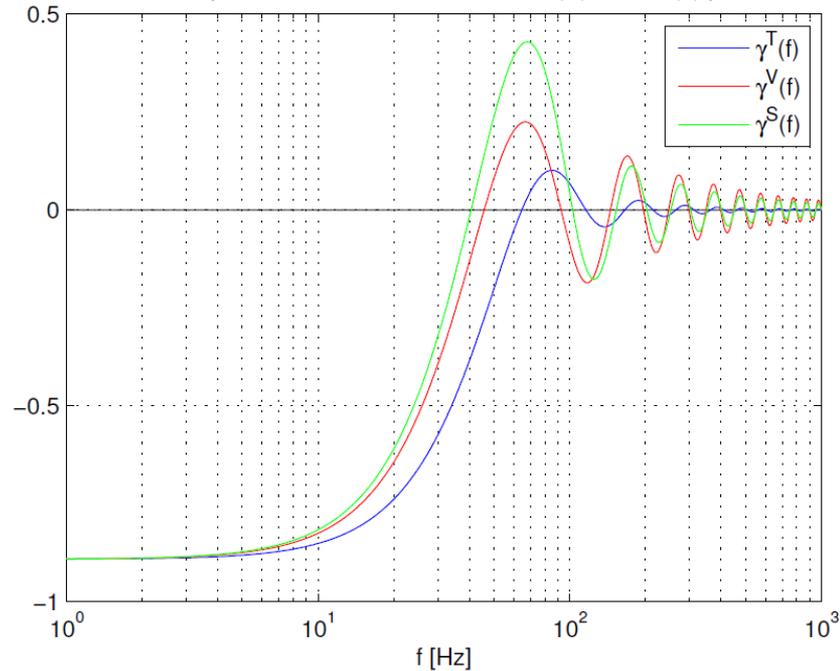
- Looking at alternative theories of gravity
- Indirect constraints on scalar modes
 - Binary pulsars
 - WMAP
- Need many detector pairs to disentangle different contributions



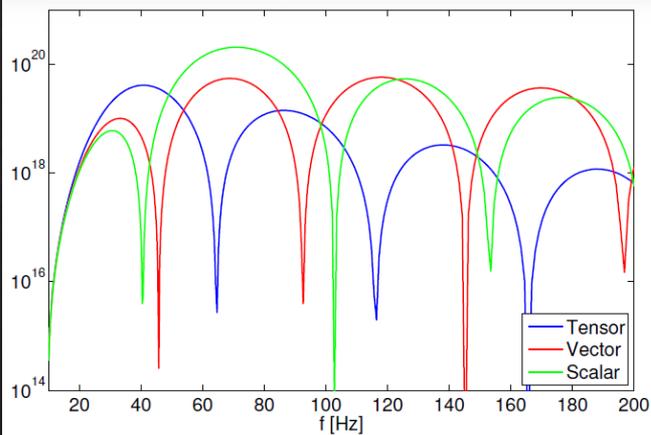
Comparison of different SGWB projections and models: $\Omega_{\text{GW}} \propto \text{const.}$



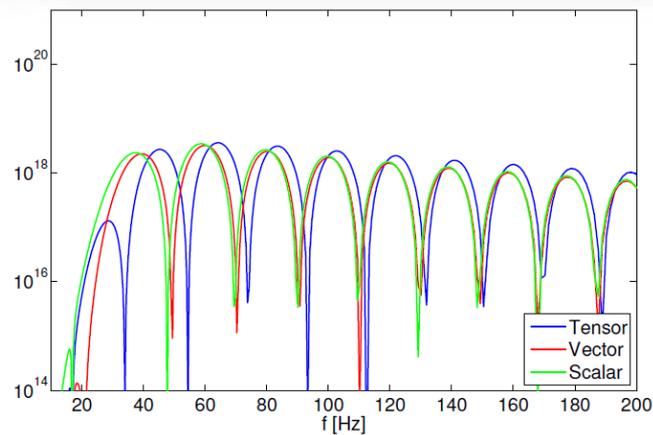
Overlap reduction function for the LIGO(H) – LIGO(L) pair



SNR integrand LIGO H – LIGO L

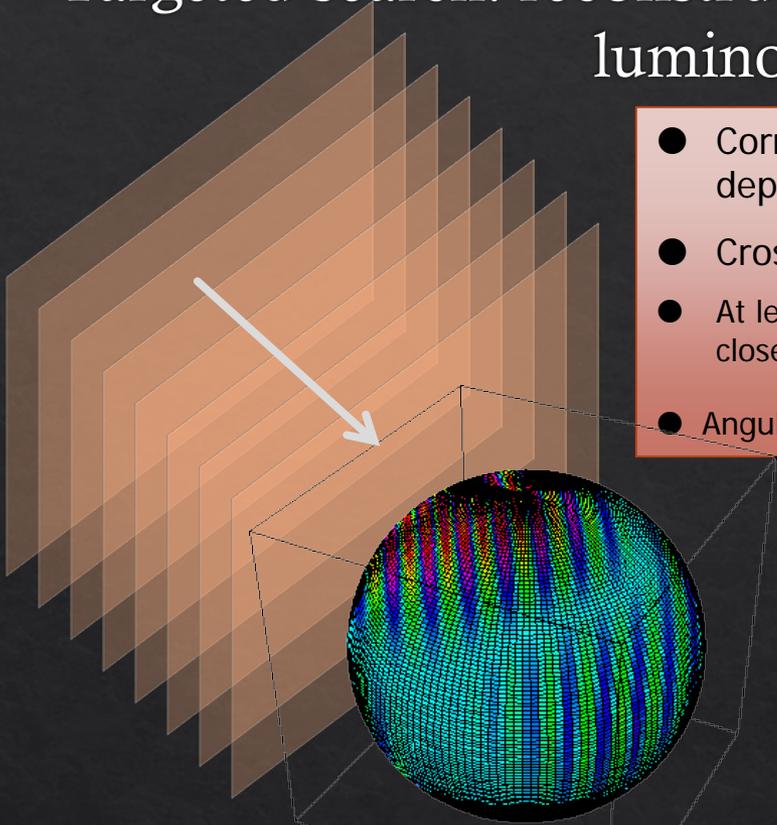


SNR integrand LIGO L – Virgo

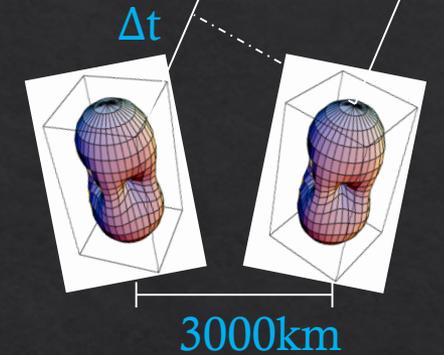


$$\Omega(f) = \Omega_T \left(\frac{f}{f_0} \right)^{\alpha_T} + \Omega_V \left(\frac{f}{f_0} \right)^{\alpha_V} + \Omega_S \left(\frac{f}{f_0} \right)^{\alpha_S}$$

Targeted search: reconstruct a map of the gravitational wave luminosity in the sky



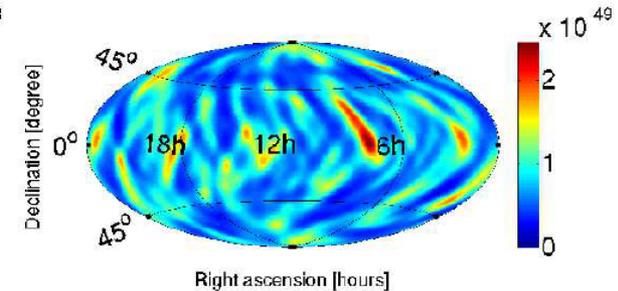
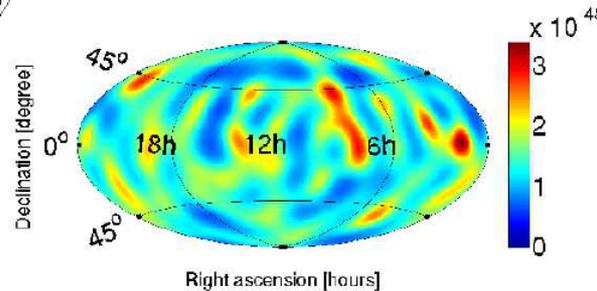
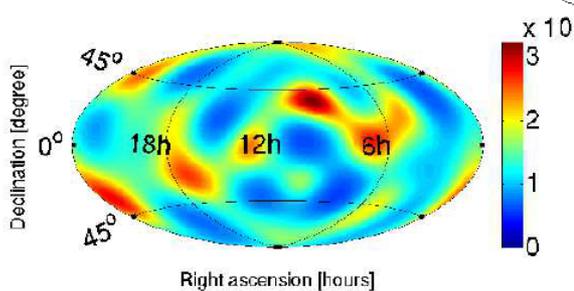
- Correct the direction-dependent modulation
- Cross-Correlate
- At least 3 detectors needed to close the inverse problem.
- Angular resolution limited by λ/D



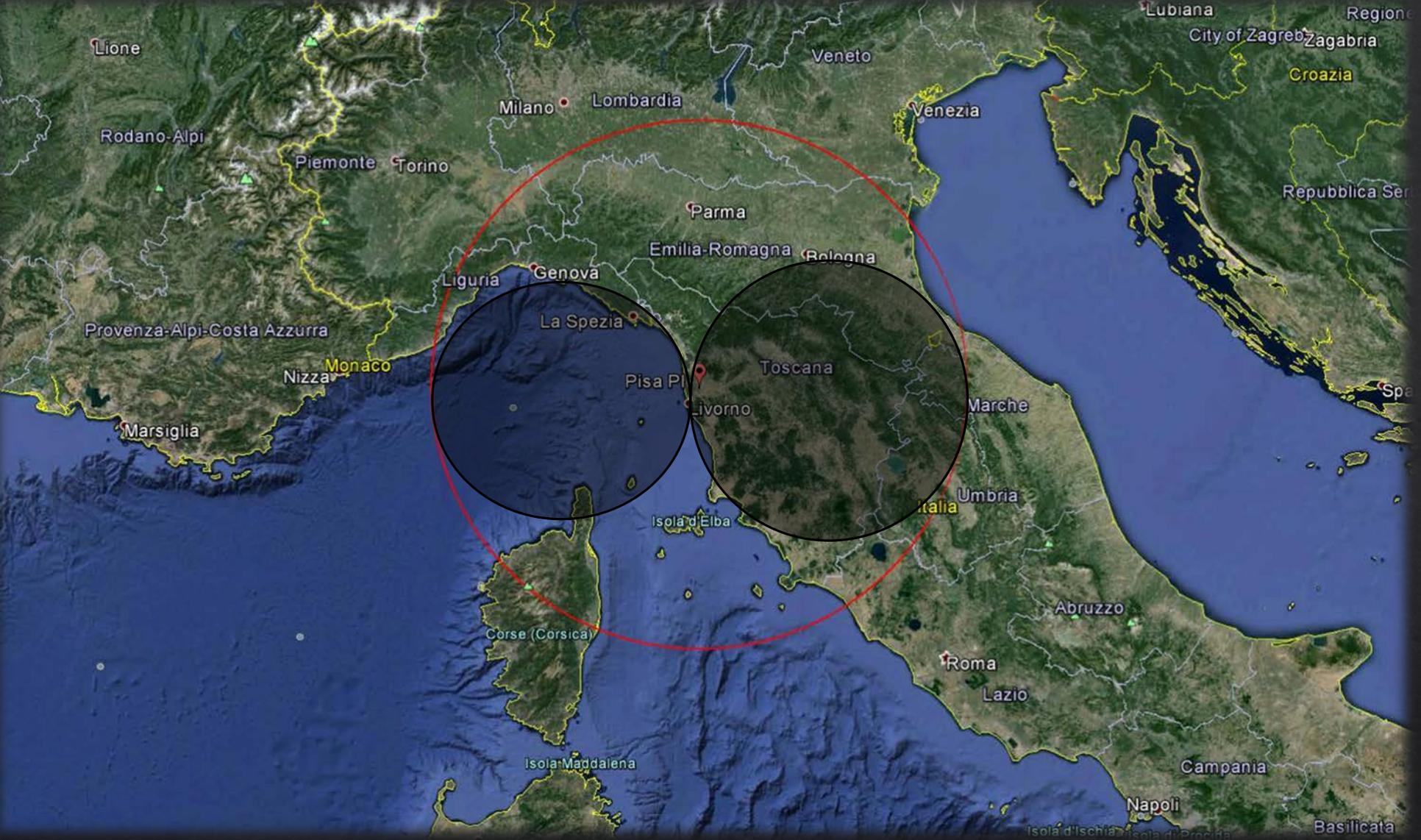
$$\Omega_{gw}(f) = \frac{2\pi^2}{3H_0^2} f^2 S_h(f) \int_{S^2} d\hat{\Omega} \mathcal{P}(\hat{\Omega})$$

$$\mathcal{P}(\vec{\Omega}) \equiv \int \eta(\hat{\Omega}') \delta^2(\hat{\Omega}, \hat{\Omega}') d\hat{\Omega}' \quad \text{«radiometer» search}$$

$$\mathcal{P}(\vec{\Omega}) \equiv \sum_{\ell, m} c_{\ell m} Y_{\ell m}(\hat{\Omega}) \quad \text{«SA» search}$$

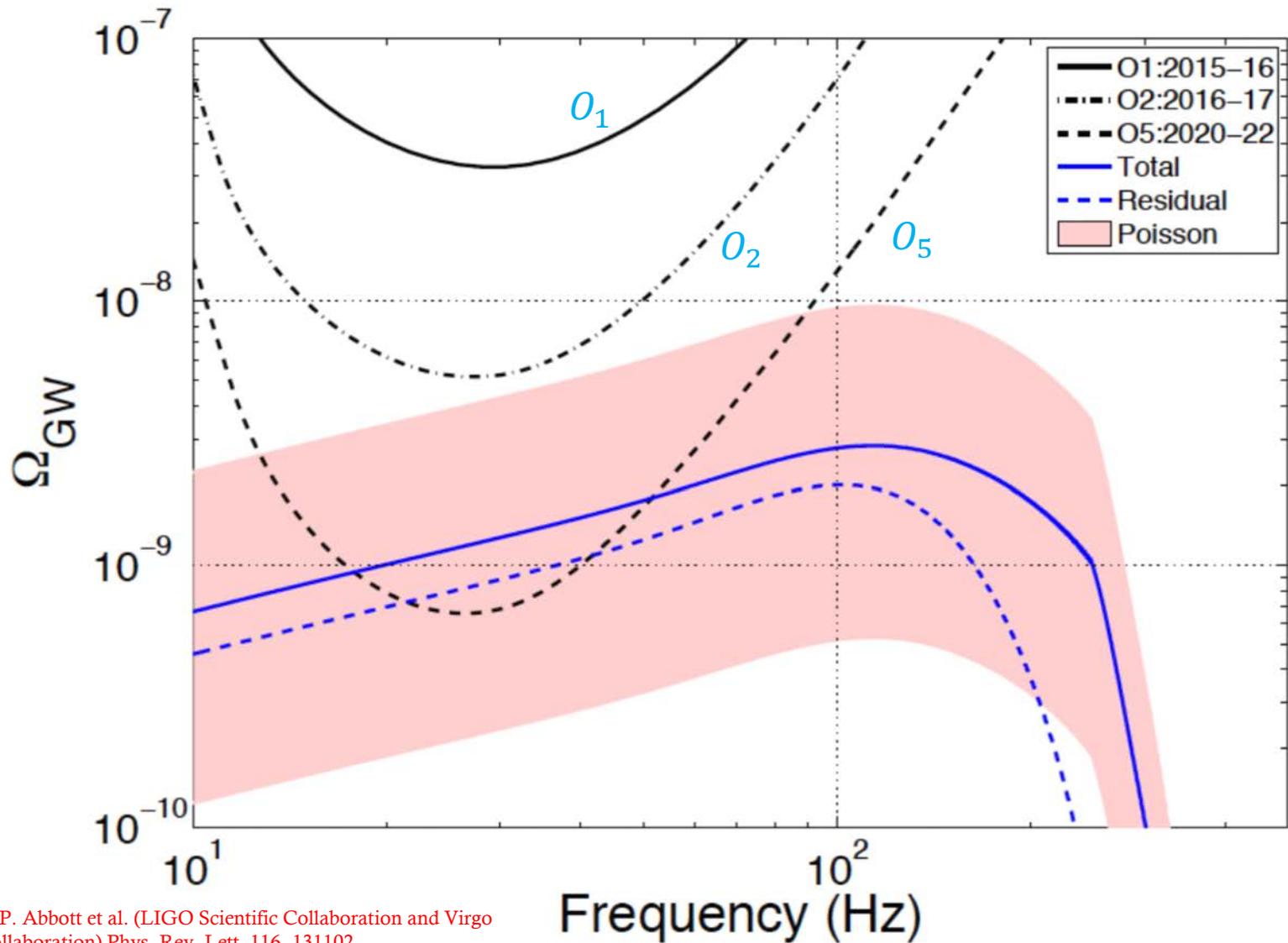




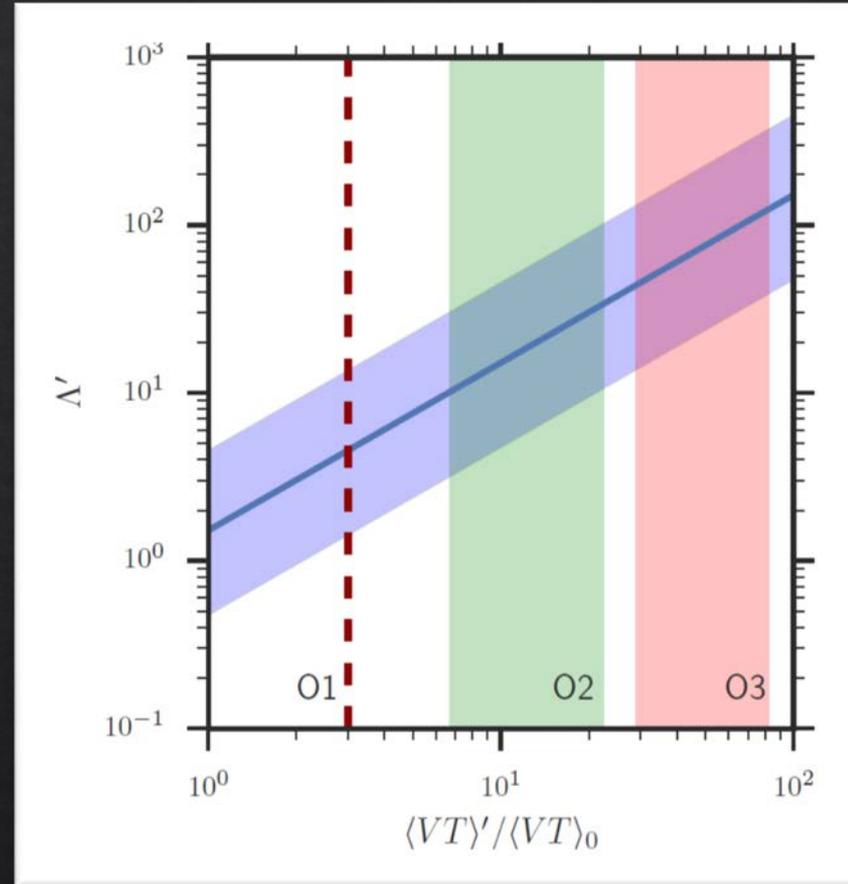
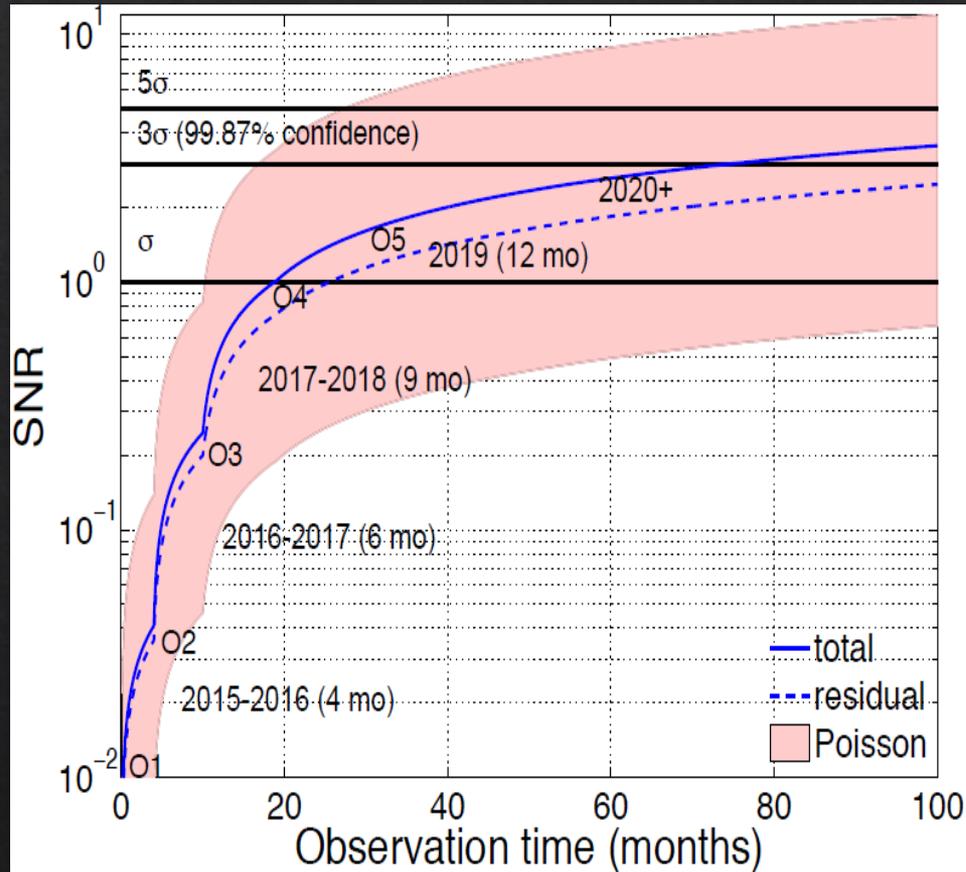


$$\left(\frac{v}{c}\right)^2 \lesssim 0.36 \quad d \gtrsim 180\text{km} \quad R_s \sim 90\text{km}$$

Implications of LIGO first detection

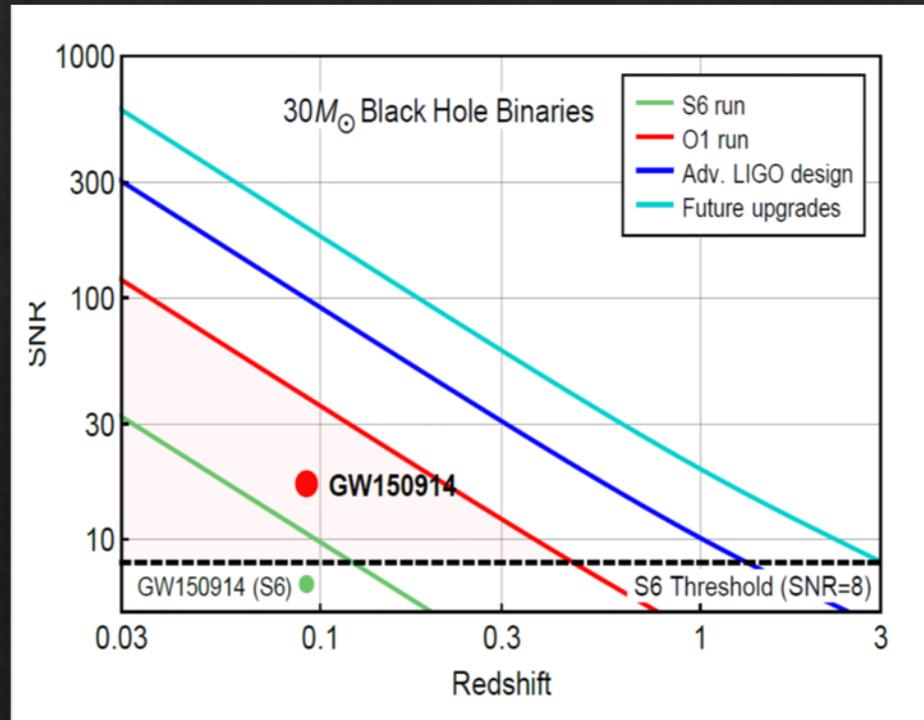


Perspectives



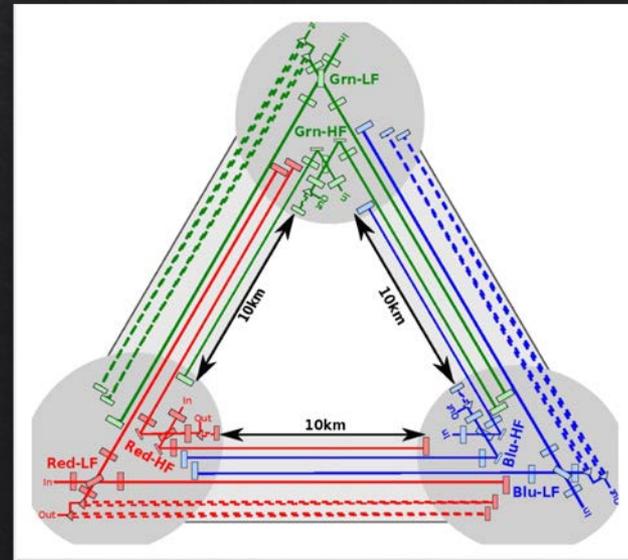
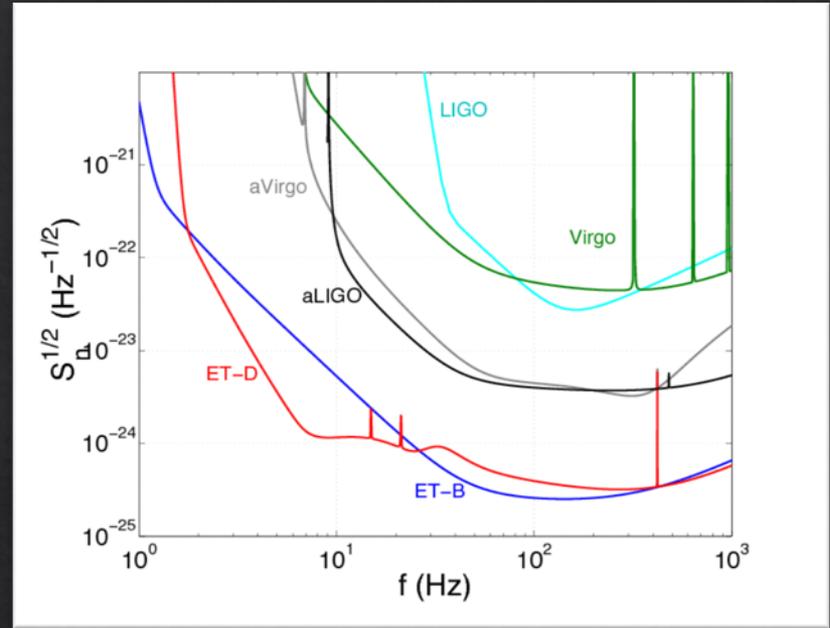
Expectations for future runs

The median value and 90% credible interval for the expected number of highly significant events (FARs <1/century) as a function of surveyed time-volume in an observation



Signal-to-noise ratio vs z
for 30 M_{\odot} BBH

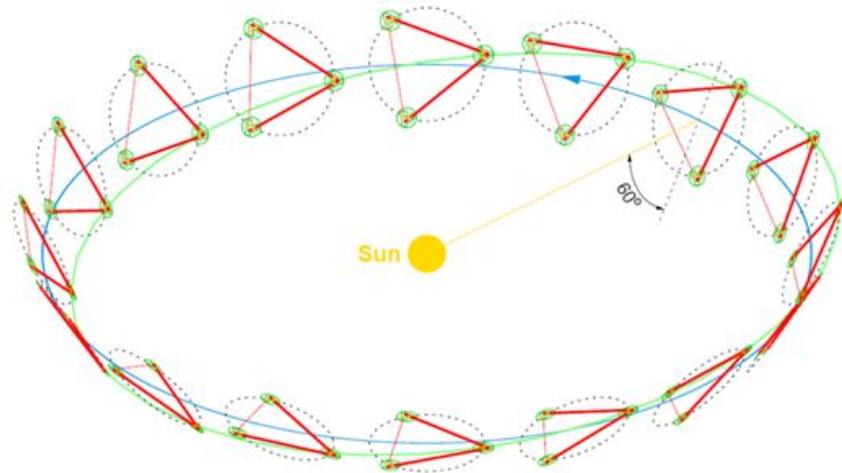
What next? The third generation detectors



What next? Space detectors

eLISA Space Based GW Detector

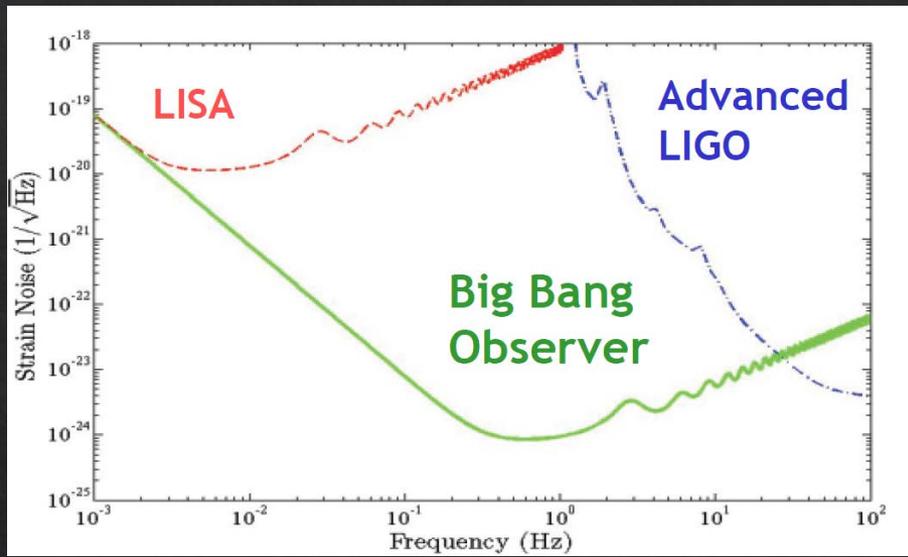
- Laser Interferometer in Space Antenna, LISA, provides unique capabilities
 - Immune to seismic noise
 - Long baseline provides 0.001 - 1Hz GW spectrum sensitivity needed for observing massive black hole mergers
- Multiple identical or similar detectors to improve detection confidence



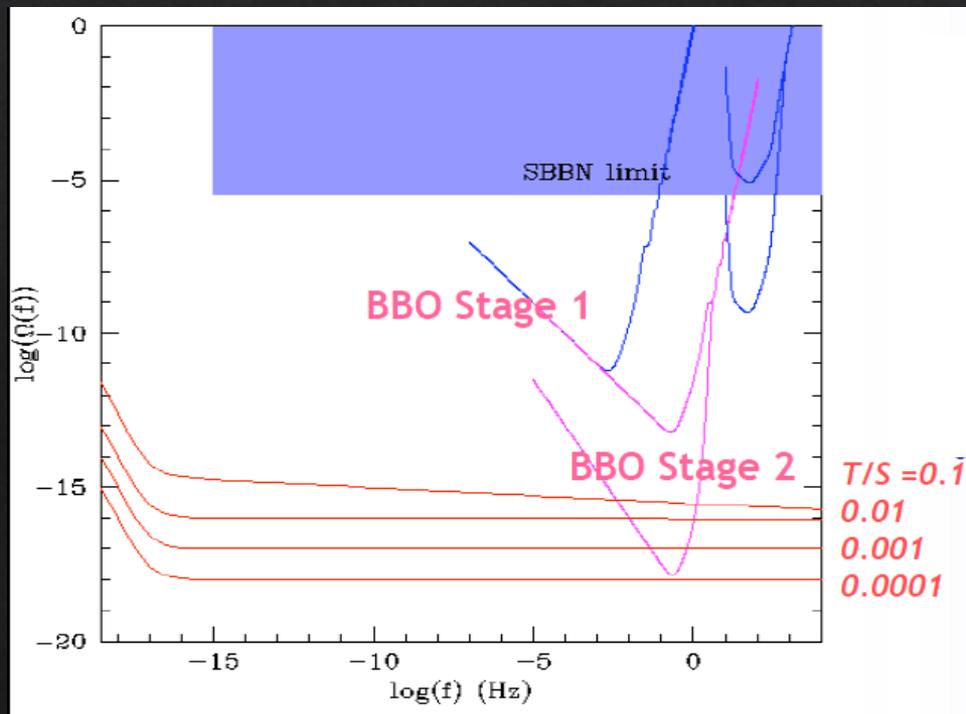
LISA: a mission to detect and observe gravitational waves, O Jennrich, in Gravitational Wave and Particle Astrophysics, Proc SPIE v5500

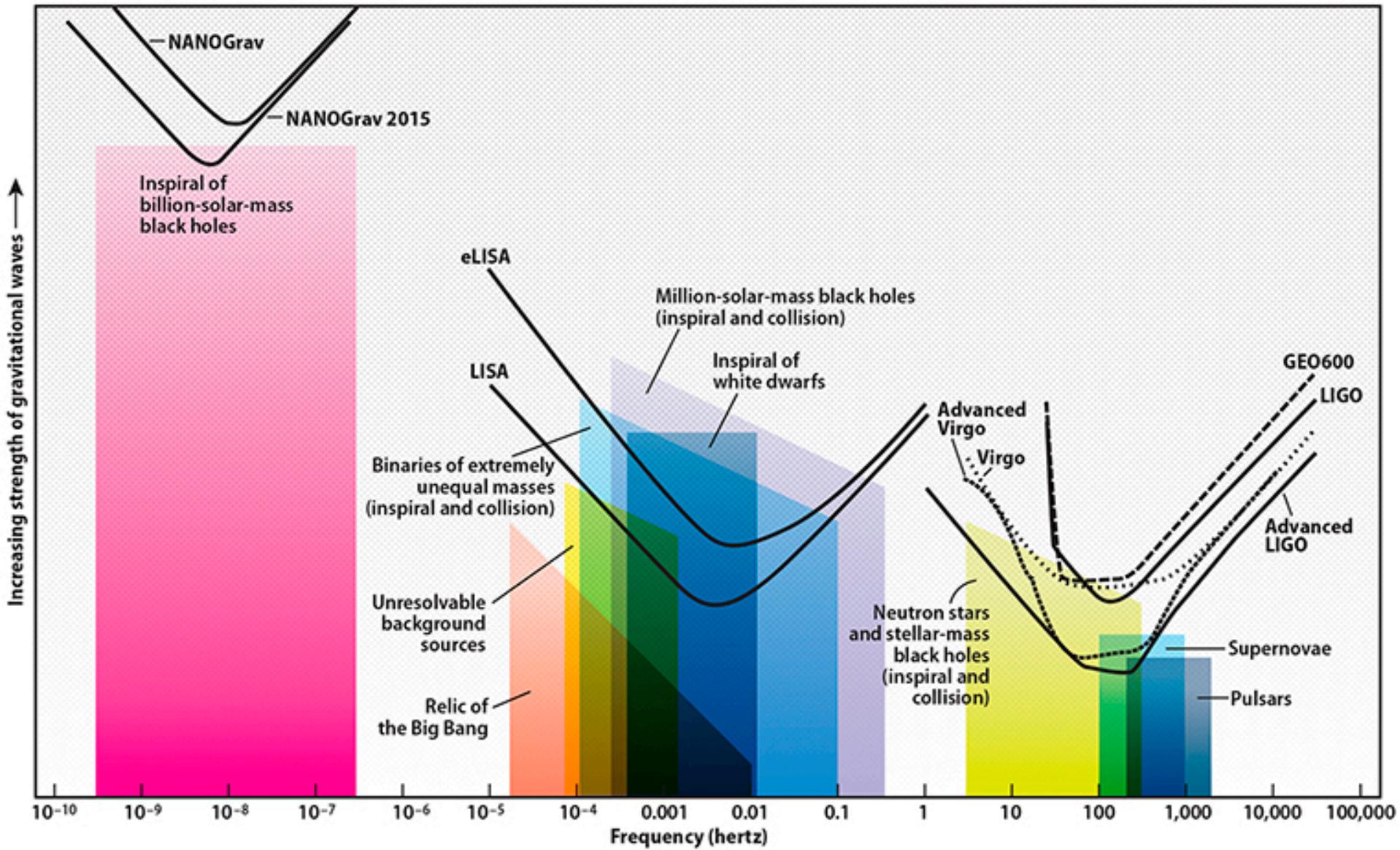
Big Bang Observer

- ◇ Fill the gap between (e)LISA & earthbound detectors
- ◇ Designed to detect SB produced by inflation
- ◇ Space based
- ◇ Shorter arm than (e)LISA



- Higher laser power
- Improved acceleration noise
- Two phases





Thank you for your attention...