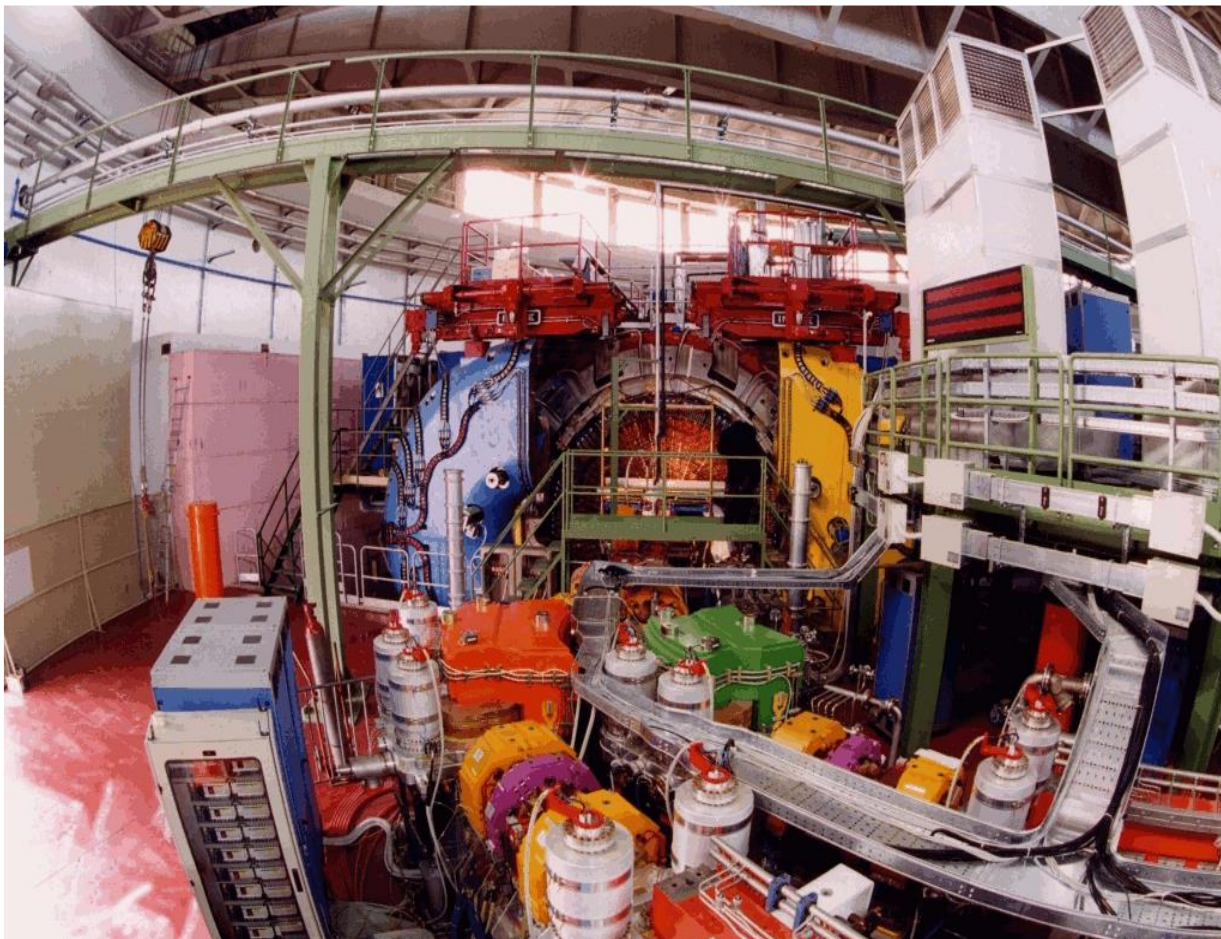


KLOE-2 Workshop on e^+e^- collision physics at 1 GeV

26-28 October 2016

INFN - Laboratori Nazionali di Frascati, Italy



DISCRETE SYMMETRIES WITH NEUTRAL MESONS

José Bernabéu

IFIC-Univ. Valencia

IFIC
INSTITUT DE FÍSICA
CORPUSCULAR

EXCELENCIA
SEVERO
OCHOA

VNIVERSITAT
ID VALÈNCIA

CSIC
CONSEJO SUPERIOR DE INVESTIGACIONES CIENTÍFICAS

FUNDAMENTAL ROLE OF SYMMETRY BREAKING

➤ Gauge Symmetry Breaking → Origin of Mass (1964-2012)

➤ Parity PV → Fields of definite transformation properties under Gauge Group: CHIRAL FIELDS

➡ Standard Model (1955 → 1957 → 1962 → 1967 → 1973)

➤ CPV → 3 Families of Elementary Fermions ↔ Mixing

➡ Flavour Physics (1964-1973-2001)

➤ TRV → Antiunitary ↔ $i \rightleftharpoons f$ → “impossible” (?)
for decaying particles ...

BYPASS (1999→2012-?) $\left\{ \begin{array}{l} \text{Decay as Filtering of parent state ONLY} \\ \text{Quantum Entanglement} \rightarrow \text{transfer to the orthogonal partner} \end{array} \right.$

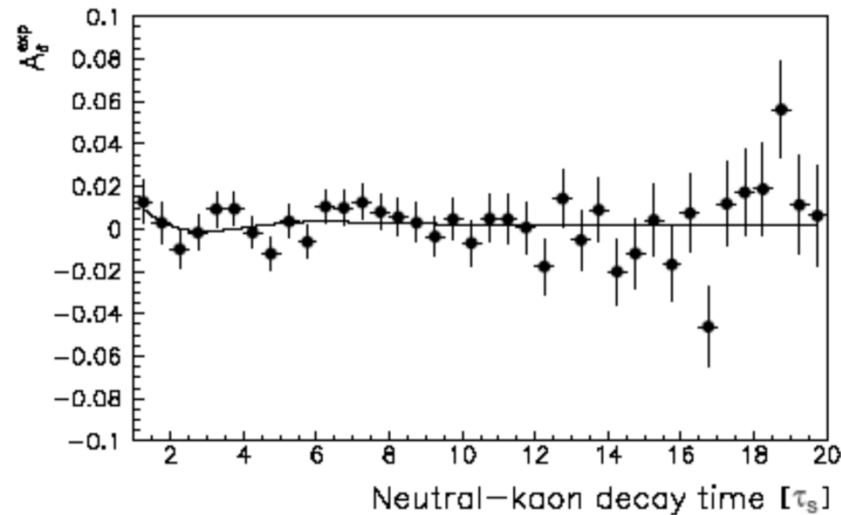
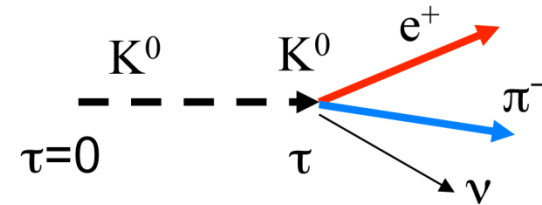
➤ CPTV ? ➡ - Beyond QFT paradigm
- Nothing in QM forbids CPTV

OUTLINE

- CPTV in SemiLeptonic Decays of Neutral Kaons
- Genuine T, CP, CPT Separate Asymmetries
in B_d -Transitions
- Direct Measurements of TRV
in Neutral Kaon Transitions at KLOE-2
- Direct Tests of CPT in Neutral Kaon Transitions at KLOE-2
- The ω -EFFECT
- Outlook

CPT test at CPLEAR

Test of **CPT** in the time evolution of neutral kaons using the semileptonic asymmetry



$$\left\{ \begin{array}{l} A_{\delta}(\tau) = \frac{\bar{R}_{+}(\tau) - \alpha R_{-}(\tau)}{\bar{R}_{+}(\tau) + \alpha R_{-}(\tau)} + \frac{\bar{R}_{-}(\tau) - \alpha R_{+}(\tau)}{\bar{R}_{-}(\tau) + \alpha R_{+}(\tau)} \\ R_{+(-)}(\tau) = R \left(K^0_{t=0} \rightarrow (e^{+(-)} \pi^{-(+)} \nu)_{t=\tau} \right) \\ \bar{R}_{- (+)}(\tau) = R \left(\bar{K}^0_{t=0} \rightarrow (e^{-(+)} \pi^{+(-)} \nu)_{t=\tau} \right) \\ \alpha = 1 + 4\Re \varepsilon_L \end{array} \right.$$

$$A_{\delta}(\tau \gg \tau_s) = 8\Re \delta$$

$$\Re \delta = (0.30 \pm 0.33 \pm 0.06) \times 10^{-3}$$

$A_S - A_L$ at KLOE – 2

Semileptonic decays of neutral kaons

$$|K_S\rangle = \frac{1}{\sqrt{2(1 + |\epsilon_S|^2)}} \left((1 + \epsilon_S) |K^0\rangle + (1 - \epsilon_S) |\bar{K}^0\rangle \right)$$

$$|K_L\rangle = \frac{1}{\sqrt{2(1 + |\epsilon_L|^2)}} \left((1 + \epsilon_L) |K^0\rangle - (1 - \epsilon_L) |\bar{K}^0\rangle \right)$$

$$\epsilon_{S/L} = \epsilon_K \pm \delta_K$$

parameter describing \mathcal{CP} violation

parameter describing \mathcal{CPT} violation

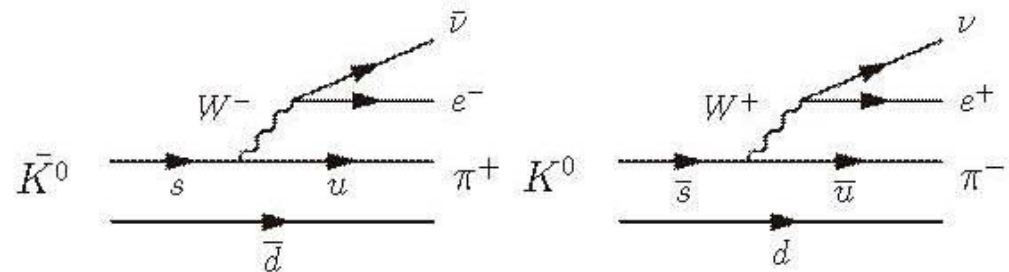
Possible semileptonic decays:

$$K^0 \rightarrow \pi^- e^+ \bar{\nu}$$

$$\bar{K}^0 \rightarrow \pi^+ e^- \nu$$

$$K^0 \rightarrow \pi^+ e^- \nu$$

$$\bar{K}^0 \rightarrow \pi^- e^+ \bar{\nu}$$



Two decays are allowed according to elementary Quarks ($\Delta S = \Delta Q$ rule)

$A_S - A_L$ at KLOE – 2

We can parametrize semileptonic amplitudes in the following way:

$$a + b = \langle \pi^- e^+ \nu | H_{weak} | K^0 \rangle = (if \mathcal{CP}) = a^* + b^*,$$

$$a^* - b^* = \langle \pi^+ e^- \bar{\nu} | H_{weak} | \bar{K}^0 \rangle = (if \mathcal{CP}) = a + b$$

and then we obtain following relations between symmetries and semileptonic amplitudes:

	\mathcal{CP}	\mathcal{T}	\mathcal{CPT}
a	$Im = 0$	$Im = 0$	
b	$Re = 0$	$Im = 0$	$= 0$

; CPTV in Decay Amplitude

$$\rightarrow y = -b/a$$

Charge Asymmetry

$$A_{S,L} = \frac{\Gamma(K_{S,L} \rightarrow \pi^- e^+ \nu) - \Gamma(K_{S,L} \rightarrow \pi^+ e^- \bar{\nu})}{\Gamma(K_{S,L} \rightarrow \pi^- e^+ \nu) + \Gamma(K_{S,L} \rightarrow \pi^+ e^- \bar{\nu})} = 2 [Re(\epsilon_K) \pm Re(\delta_K) - Re(y)]$$

$$A_L = (3.332 \pm 0.058_{\text{stat}} \pm 0.047_{\text{syst}}) \cdot 10^{-3} \text{ [KTeV Collaboration, PRL 88 (2002)]}$$

$$A_S = (1.5 \pm 9.6_{\text{stat}} \pm 2.9_{\text{syst}}) \cdot 10^{-3} \text{ [KLOE, PL B 636 (2006) 173-182]}$$

SYMMETRIES IN THE LAWS OF PHYSICS

- In Quantum Mechanics, there is an operator U_T implementing the T-symmetry acting on the states of the physical system, such that

$$U_T \vec{r} U_T^\dagger = \vec{r}, \quad U_T \vec{p} U_T^\dagger = -\vec{p}, \quad U_T \vec{s} U_T^\dagger = -\vec{s}$$

By considering the commutator $[r_j, p_K] = i\hbar\delta_{jK}I$

the operator U_T must be ANTI-UNITARY:

UNITARY- for conserving probabilities, ANTI- for complex conjugation

ANTIUNITARITY introduces many intriguing subtleties:

$$S_{i \rightarrow f} \xrightarrow{T} S_{U_T f \rightarrow U_T i}$$

T - Violation means Asymmetry under

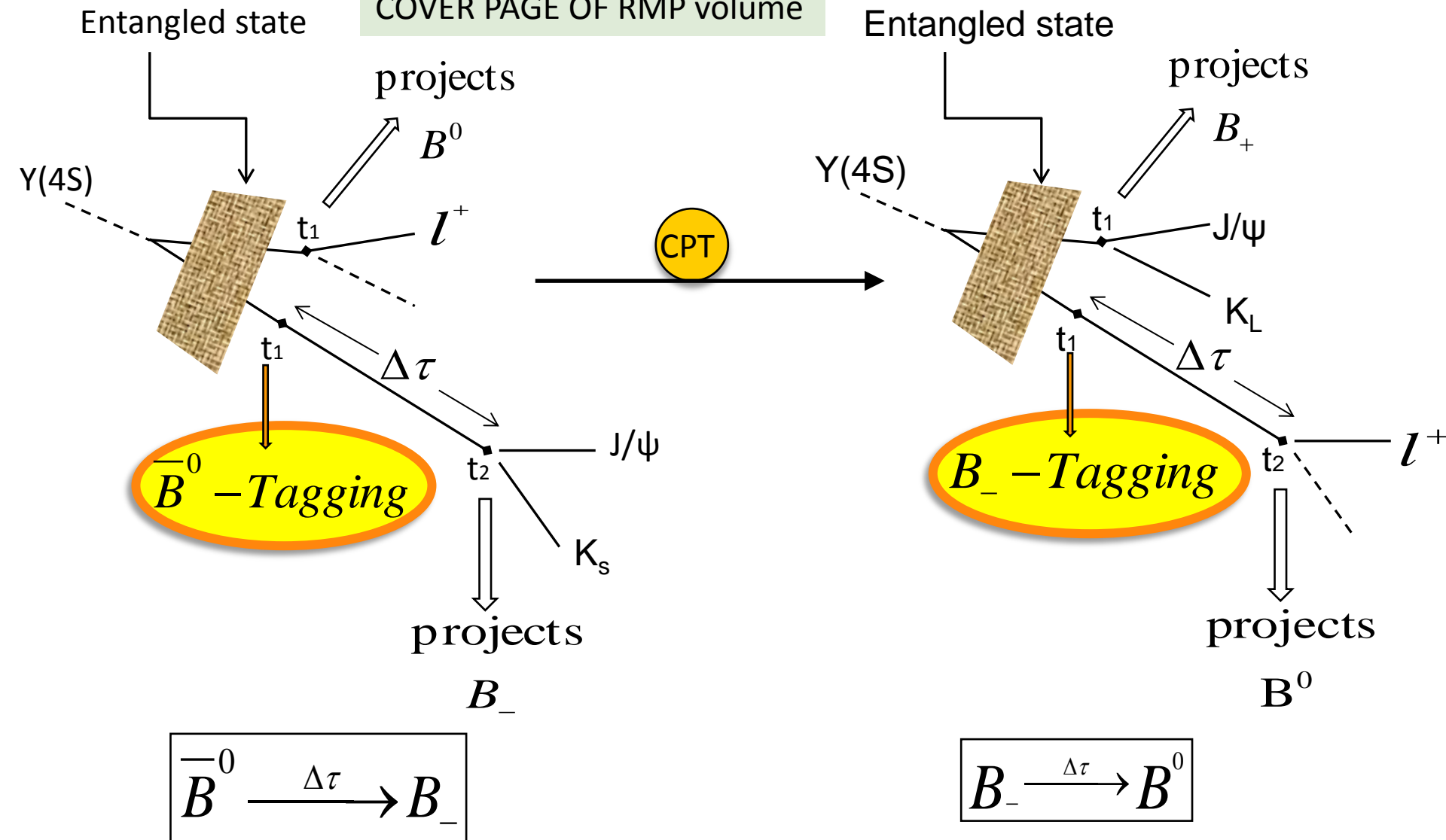
Interchange in \longleftrightarrow out states

- Similarly for ANTIUNITARY CPT which needs not only in \longleftrightarrow out, but also $i, f \rightarrow \bar{f}, \bar{i}$, in transitions.

WHAT IS CPT-TRANSFORMATION EXPERIMENTALLY ?

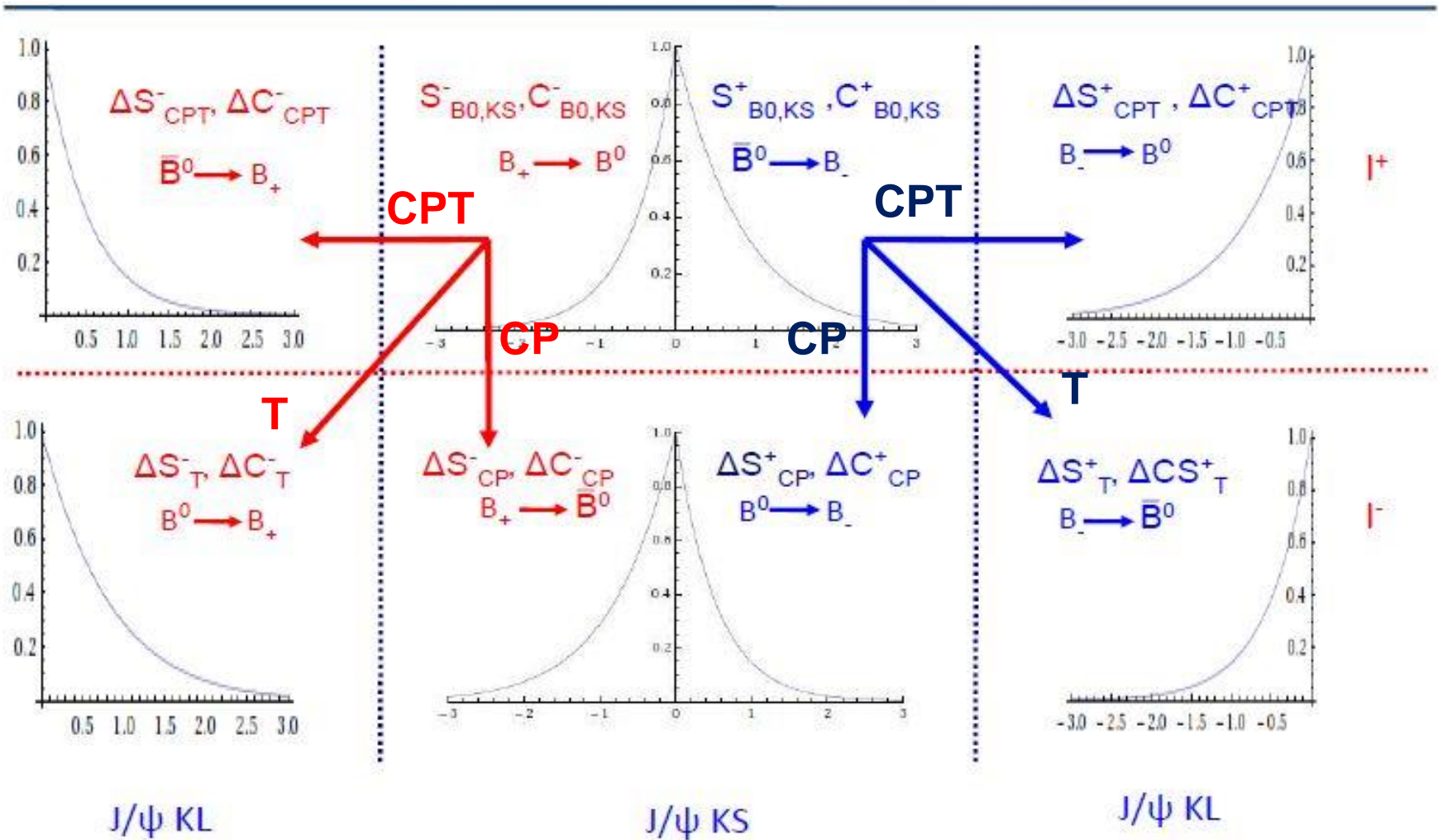
The problem is in the preparation and filtering of the appropriate initial and final meson states for a CPT-test in TRANSITIONS

COVER PAGE OF RMP volume



$\Delta S^\pm, \Delta C^\pm$ ASYMMETRY PARAMETERS

$$I_i(\Delta t) \sim e^{-\Gamma\Delta t} \{ C_i \cos(\Delta m\Delta t) + S_i \sin(\Delta m\Delta t) + C'_i \cosh(\Delta\Gamma\Delta t) + S'_i \sinh(\Delta\Gamma\Delta t) \}$$



GENUINE T, CP, CPT ASYMMETRIES

J.B., F. Botella, M. Nebot, JHEP 1606 (2016) 100

➤ 3 different Observables ΔC_h , ΔC_c , ΔS_c for each symmetry

9 Asymmetry parameters with different information content

$$\Delta S_c^T = -0.687 \pm 0.020 ; \Delta S_c^{cP} = -0.680 \pm 0.021$$

Impressive separate evidence of TRV, CPV

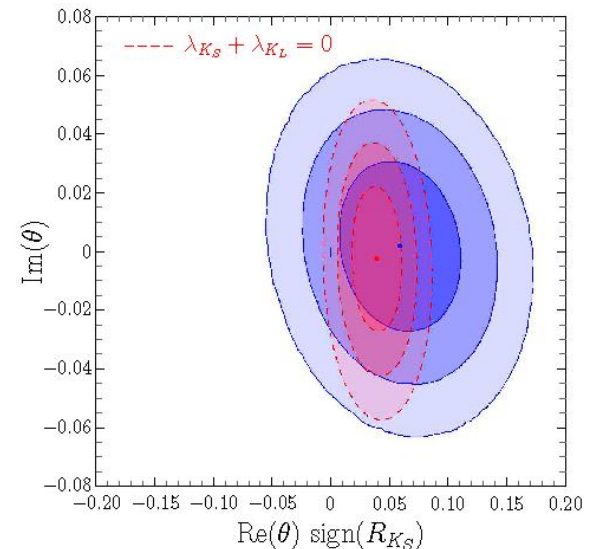
“Intriguing” 2σ - effect for CPTV $\Re\delta$

Analysis assuming perfect ENTANGLEMENT

➔ The two Time-Ordered Decays **f, g** satisfy

$$C_h(f,g) = C_h(g,f) ; C_c(f,g) = C_c(g,f) ; S_c(f,g) = -S_c(g,f)$$

➤ Weakening Entanglement → CPTV ω -effect violates this relation, in progress

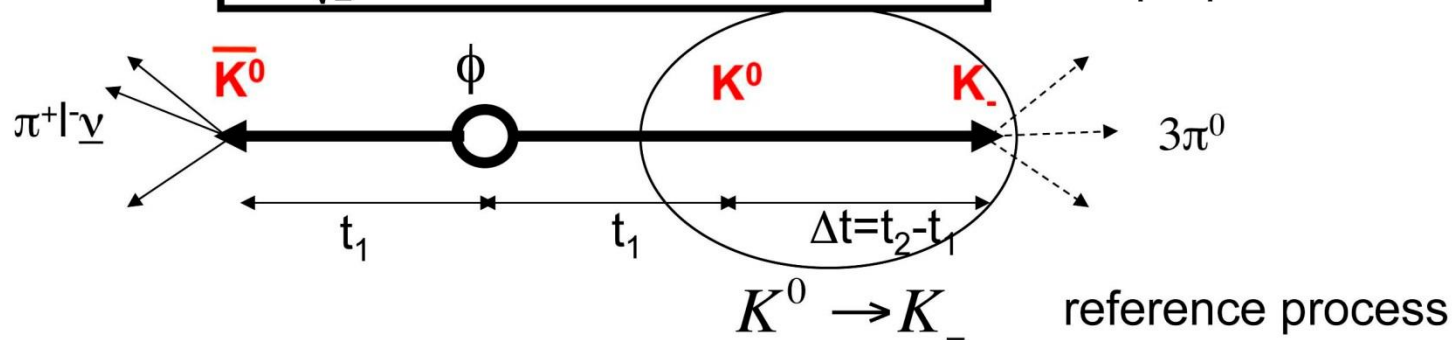


TIME REVERSAL VIOLATION IN NEUTRAL KAON TRANSITIONS

EPR correlations at a ϕ -Factory can be exploited to study T-conjugated TRANSITIONS between K^0, \bar{K}^0 and the orthogonal K_+, K_- states filtered by CP-eigenstate decay products

$$\begin{aligned}
 |i\rangle &= \frac{1}{\sqrt{2}} \left[|K^0(\vec{p})\rangle |\bar{K}^0(-\vec{p})\rangle - |\bar{K}^0(\vec{p})\rangle |K^0(-\vec{p})\rangle \right] \\
 &= \frac{1}{\sqrt{2}} \left[|K_+(\vec{p})\rangle |K_-(-\vec{p})\rangle - |K_-(-\vec{p})\rangle |K_+(\vec{p})\rangle \right]
 \end{aligned}$$

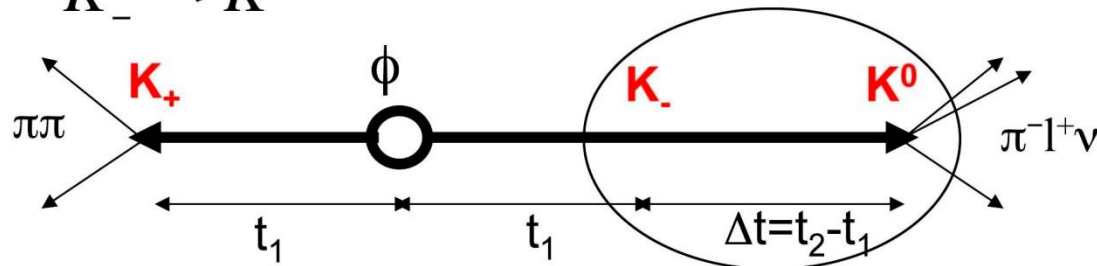
- decay as filtering measurement
- entanglement -> preparation of state



Note: CP and CPT conjugated process

$$\bar{K}^0 \rightarrow K_- \quad K_- \rightarrow \bar{K}^0$$

$K_- \rightarrow K^0$ T-conjugated process



TIME REVERSAL VIOLATION IN NEUTRAL KAON TRANSITIONS

T symmetry test

Reference		T -conjugate	
Transition	Decay products	Transition	Decay products
$\bar{K}^0 \rightarrow K_-$	$(\ell^+, \pi^0 \pi^0 \pi^0)$	$K_- \rightarrow \bar{K}^0$	$(\pi^0 \pi^0 \pi^0, \ell^-)$
$K_+ \rightarrow K^0$	$(\pi^0 \pi^0 \pi^0, \ell^+)$	$K^0 \rightarrow K_+$	$(\ell^-, \pi \pi)$
$\bar{K}^0 \rightarrow K_+$	$(\ell^+, \pi \pi)$	$K_+ \rightarrow \bar{K}^0$	$(\pi^0 \pi^0 \pi^0, \ell^-)$
$K_- \rightarrow K^0$	$(\pi \pi, \ell^+)$	$K^0 \rightarrow K_-$	$(\ell^-, \pi \pi)$

One can define the following ratios of probabilities:

$$\begin{aligned}
 R_1(\Delta t) &= P [K^0(0) \rightarrow K_+(\Delta t)] / P [K_+(0) \rightarrow K^0(\Delta t)] \\
 R_2(\Delta t) &= P [K^0(0) \rightarrow K_-(\Delta t)] / P [K_-(0) \rightarrow K^0(\Delta t)] \\
 R_3(\Delta t) &= P [\bar{K}^0(0) \rightarrow K_+(\Delta t)] / P [K_+(0) \rightarrow \bar{K}^0(\Delta t)] \\
 R_4(\Delta t) &= P [\bar{K}^0(0) \rightarrow K_-(\Delta t)] / P [K_-(0) \rightarrow \bar{K}^0(\Delta t)] .
 \end{aligned}$$

Any deviation from $R_i=1$ constitutes a violation of T-symmetry

TIME REVERSAL VIOLATION IN NEUTRAL KAON TRANSITIONS

$$R_1^{\text{exp}}(\Delta t) \equiv \frac{I(\ell^-, \pi\pi; \Delta t)}{I(3\pi^0, \ell^+; \Delta t)} = R_1(\Delta t) \times \frac{C(\ell^-, \pi\pi)}{C(3\pi^0, \ell^+)}$$

$$R_2^{\text{exp}}(\Delta t) \equiv \frac{I(\ell^-, 3\pi^0; \Delta t)}{I(\pi\pi, \ell^+; \Delta t)} = R_2(\Delta t) \times \frac{C(\ell^-, 3\pi^0)}{C(\pi\pi, \ell^+)}$$

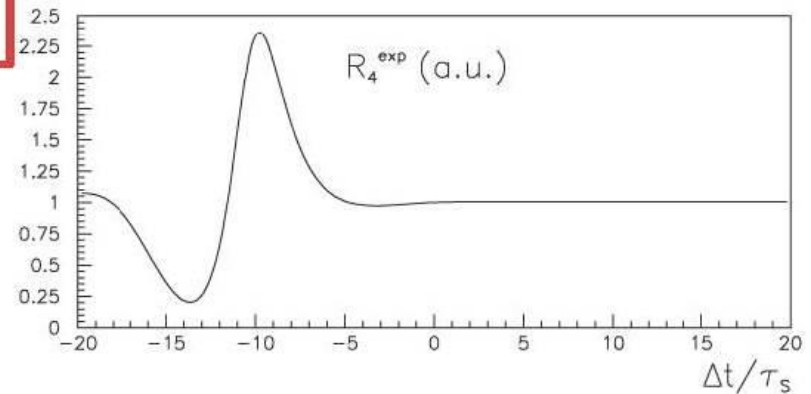
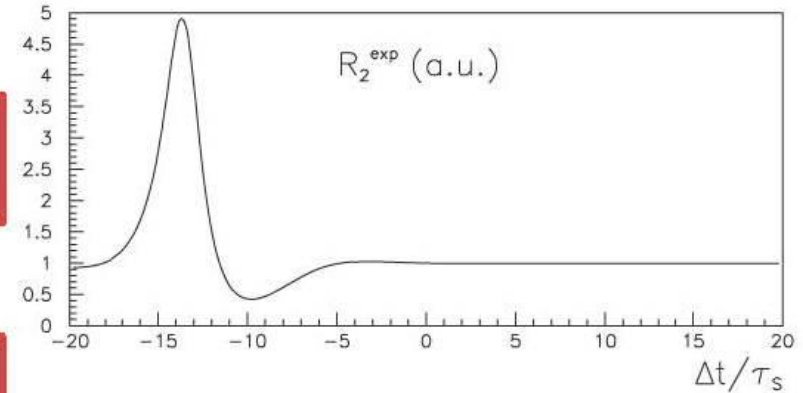
$$R_3^{\text{exp}}(\Delta t) \equiv \frac{I(\ell^+, \pi\pi; \Delta t)}{I(3\pi^0, \ell^-; \Delta t)} = R_3(\Delta t) \times \frac{C(\ell^+, \pi\pi)}{C(3\pi^0, \ell^-)}$$

$$R_4^{\text{exp}}(\Delta t) \equiv \frac{I(\ell^+, 3\pi^0; \Delta t)}{I(\pi\pi, \ell^-; \Delta t)} = R_4(\Delta t) \times \frac{C(\ell^+, 3\pi^0)}{C(\pi\pi, \ell^-)}$$

In practice two measurable ratios with $\Delta t < 0$ or > 0

$$R_2^{\text{exp}}(-\Delta t) = \frac{1}{R_3^{\text{exp}}(\Delta t)} = \frac{1}{R_3(\Delta t)} \times \frac{C(3\pi^0, \ell^-)}{C(\ell^+, \pi\pi)},$$

$$R_4^{\text{exp}}(-\Delta t) = \frac{1}{R_1^{\text{exp}}(\Delta t)} = \frac{1}{R_1(\Delta t)} \times \frac{C(3\pi^0, \ell^+)}{C(\ell^-, \pi\pi)}.$$



T test could be feasible at
KLOE-2 @ DAPHNE
with $L = O(10 \text{ fb}^{-1})$

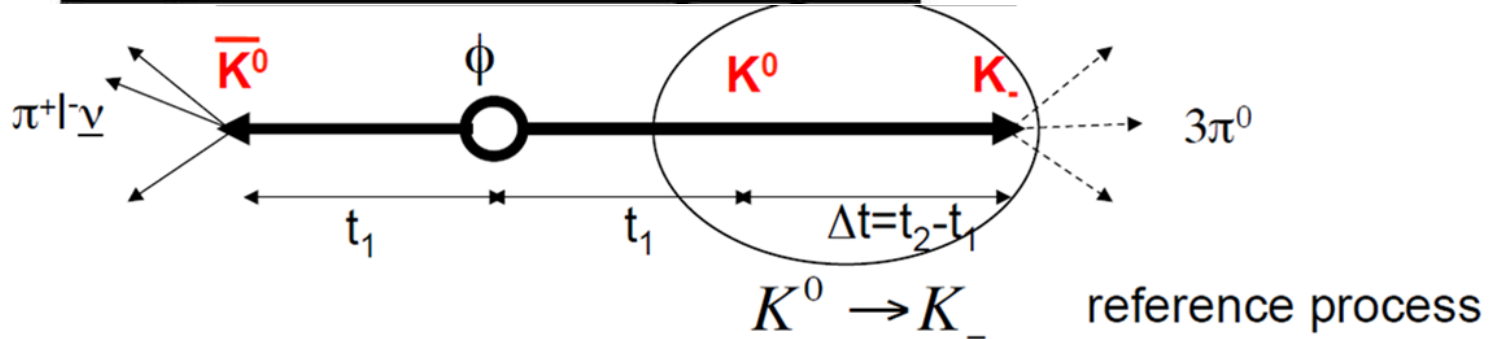
DIRECT TEST OF CPT IN NEUTRAL KAON TRANSITIONS

EPR correlations at a ϕ -Factory can be exploited to study CPT-conjugated Transitions involving Flavour $K^0 - \bar{K}^0$ and the filtered K^+ and K^- from CP-eigenstate Decay Products

$$|i\rangle = \frac{1}{\sqrt{2}} \left[|K^0(\vec{p})\rangle |\bar{K}^0(-\vec{p})\rangle - |\bar{K}^0(\vec{p})\rangle |K^0(-\vec{p})\rangle \right]$$

$$= \frac{1}{\sqrt{2}} \left[|K_+(\vec{p})\rangle |K_-(-\vec{p})\rangle - |K_-(-\vec{p})\rangle |K_+(\vec{p})\rangle \right]$$

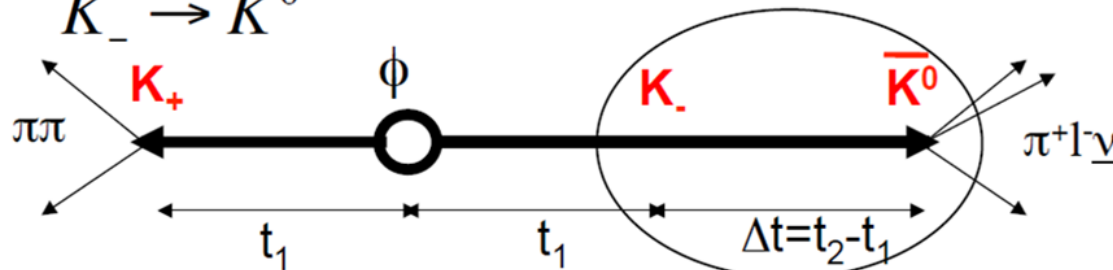
- decay as filtering measurement
- entanglement \rightarrow preparation of state



Note: CP and T conjugated process

$$\bar{K}^0 \rightarrow K_- \quad K_- \rightarrow K^0$$

$$K_- \rightarrow \bar{K}^0 \quad \text{CPT-conjugated process}$$



DIRECT TEST OF CPT IN NEUTRAL KAON TRANSITIONS

CPT symmetry test

Reference		\mathcal{CPT} -conjugate	
Transition	Decay products	Transition	Decay products
$K^0 \rightarrow K_+$	$(\ell^-, \pi\pi)$	$K_+ \rightarrow \bar{K}^0$	$(3\pi^0, \ell^-)$
$K^0 \rightarrow K_-$	$(\ell^-, 3\pi^0)$	$K_- \rightarrow \bar{K}^0$	$(\pi\pi, \ell^-)$
$\bar{K}^0 \rightarrow K_+$	$(\ell^+, \pi\pi)$	$K_+ \rightarrow K^0$	$(3\pi^0, \ell^+)$
$\bar{K}^0 \rightarrow K_-$	$(\ell^+, 3\pi^0)$	$K_- \rightarrow K^0$	$(\pi\pi, \ell^+)$

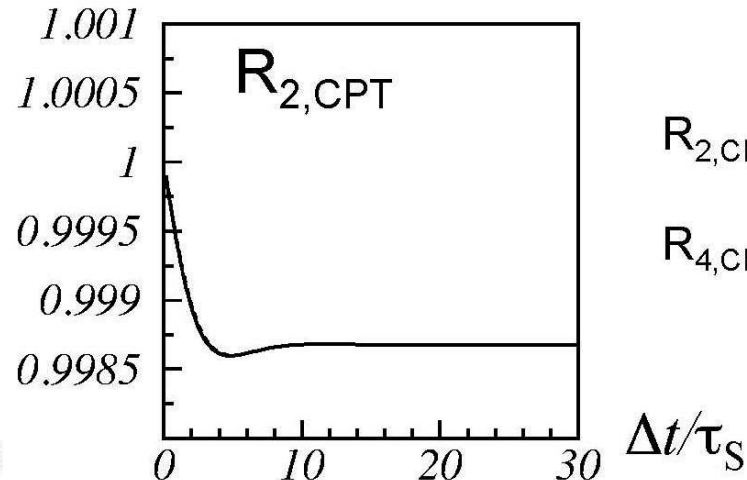
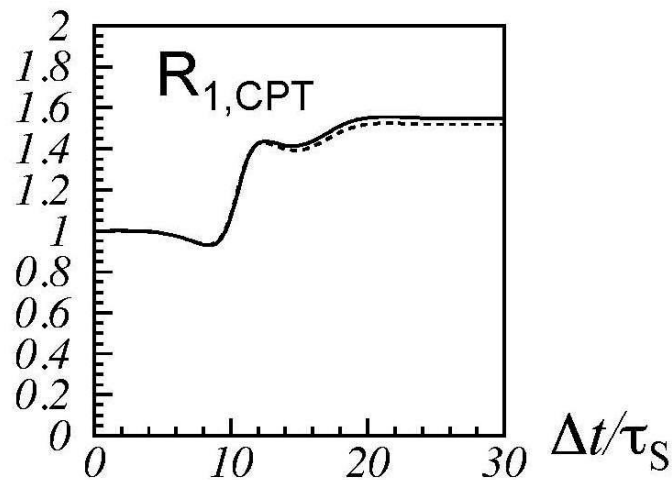
One can define the following ratios of probabilities:

$$\begin{aligned}
 R_{1,\mathcal{CPT}}(\Delta t) &= P [K_+(0) \rightarrow \bar{K}^0(\Delta t)] / P [K^0(0) \rightarrow K_+(\Delta t)] \\
 R_{2,\mathcal{CPT}}(\Delta t) &= P [K^0(0) \rightarrow K_-(\Delta t)] / P [K_-(0) \rightarrow \bar{K}^0(\Delta t)] \\
 R_{3,\mathcal{CPT}}(\Delta t) &= P [K_+(0) \rightarrow K^0(\Delta t)] / P [\bar{K}^0(0) \rightarrow K_+(\Delta t)] \\
 R_{4,\mathcal{CPT}}(\Delta t) &= P [\bar{K}^0(0) \rightarrow K_-(\Delta t)] / P [K_-(0) \rightarrow K^0(\Delta t)]
 \end{aligned}$$

Any deviation from $R_{i,\mathcal{CPT}}=1$ constitutes a violation of CPT-symmetry

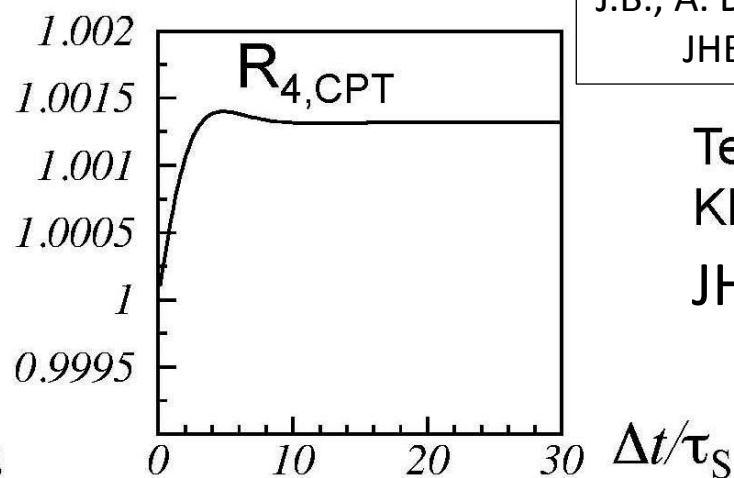
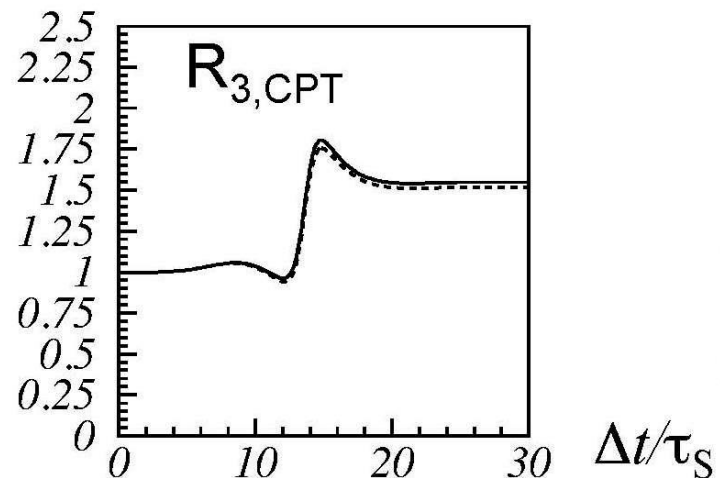
DIRECT TEST OF CPT IN NEUTRAL KAON TRANSITIONS

for visualization purposes, plots with $\text{Re}(\delta)=3.3 \cdot 10^{-4}$ $\text{Im}(\delta)=1.6 \cdot 10^{-5}$ (---- $\text{Im}(\delta)=0$)



$$R_{2,\text{CPT}}(\Delta t \gg \tau_S) = 1 - 4\text{Re}(\delta)$$

$$R_{4,\text{CPT}}(\Delta t \gg \tau_S) = 1 + 4\text{Re}(\delta)$$



J.B., A. Di Domenico, P. Villanueva,
JHEP 1510 (2015)139

Test feasible at
KLOE/KLOE-2
JHEP 2015

THE ω - EFFECT

In presence of decoherence and CPT violation induced by quantum gravity (CPT operator “ill-defined”) the definition of the particle-antiparticle states could be modified. This in turn could induce a breakdown of the correlations imposed by Bose statistics (EPR correlations) to the kaon state:

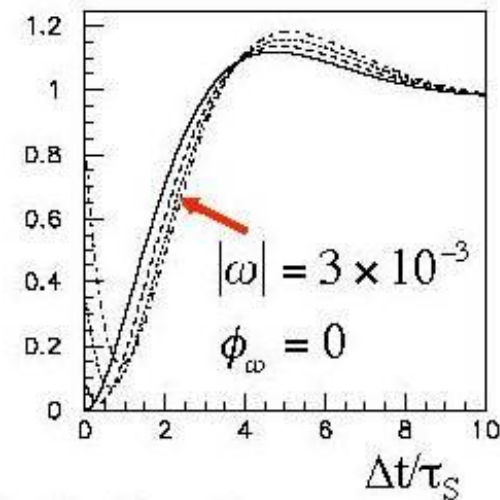
[JB, Mavromatos, Papavassiliou, PRL 92(2004) 131601]

$$|i\rangle \propto (|K^0\rangle|\bar{K}^0\rangle - |\bar{K}^0\rangle|K^0\rangle) + \omega(|K^0\rangle|\bar{K}^0\rangle + |\bar{K}^0\rangle|K^0\rangle) \\ \propto (|K_S\rangle|K_L\rangle - |K_L\rangle|K_S\rangle) + \omega(|K_S\rangle|K_S\rangle - |K_L\rangle|K_L\rangle)$$

at most one expects:

$$|\omega|^2 = O\left(\frac{E^2/M_{PLANCK}}{\Delta\Gamma}\right) \approx 10^{-5} \Rightarrow |\omega| \sim 10^{-3}$$

$I(\pi^+\pi^-, \pi^+\pi^-; \Delta t)$ (a.u.)



In some microscopic models of space-time foam arising from non-critical string theory:

[JB, Mavromatos, Sarkar, PRD 74(2006) 045014]

$$|\omega| \sim 10^{-4} \div 10^{-5}$$

The maximum sensitivity to ω is expected for $f_1=f_2=\pi^+\pi^-$

All CPTV effects induced by QG ($\alpha, \beta, \gamma, \omega$) could be simultaneously disentangled.

MEASUREMENT OF ω - EFFECT

Fit of $I(\pi^+\pi^-, \pi^+\pi^-; \Delta t, \omega)$:

- Analysed data: 1.5 fb^{-1}

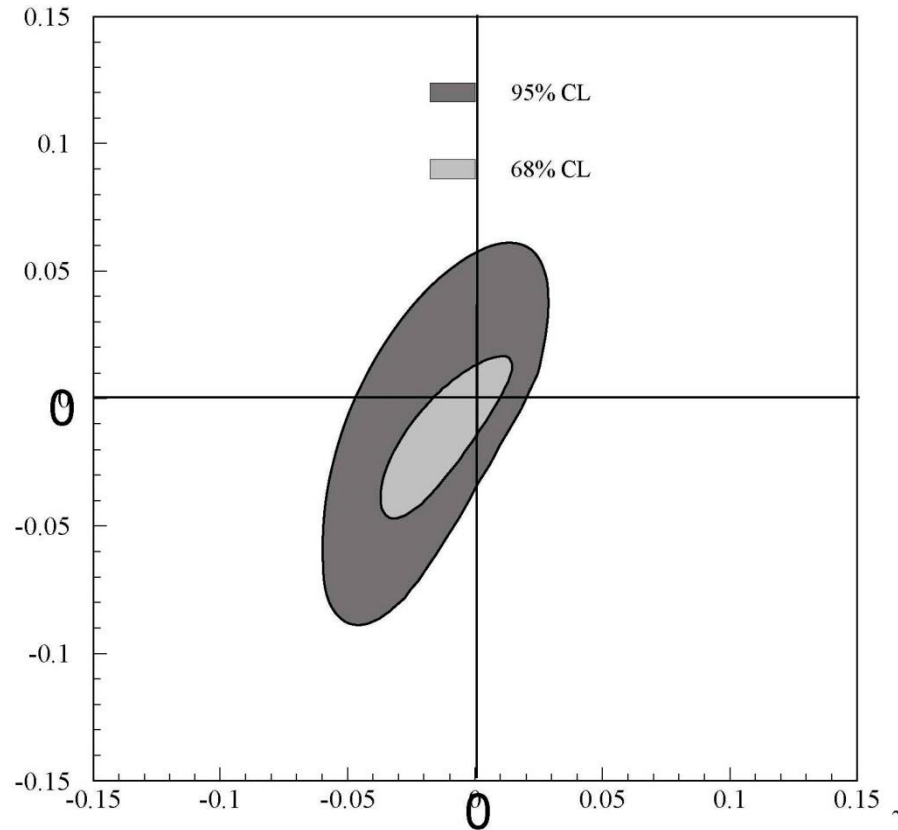
KLOE result: [PLB 642\(2006\) 315](#)
[Found. Phys. 40 \(2010\) 852](#)

$$\Re \omega = \left(-1.6_{-2.1}^{+3.0} \text{STAT} \pm 0.4_{\text{SYST}} \right) \times 10^{-4}$$

$$\Im \omega = \left(-1.7_{-3.0}^{+3.3} \text{STAT} \pm 1.2_{\text{SYST}} \right) \times 10^{-4}$$

$$|\omega| < 1.0 \times 10^{-3} \quad \text{at } 95\% \text{ C.L.}$$

$\Im \omega \times 10^{-2}$



$\Re \omega \times 10^{-2}$

In the B system [Alvarez, Bernabeu, Nebot JHEP 0611, 087]:

$$-0.0084 \leq \Re \omega \leq 0.0100 \quad \text{at } 95\% \text{ C.L.} ; \Re \omega \text{ and } \Im \omega \text{ in progress}$$

OUTLOOK

- Fundamental Role of Discrete Symmetry Breaking → Importance of Direct Asymmetries for Separate T, CP, CPT in Transitions → **Need of Entanglement for Neutral Mesons**
- DAPHNE → K_S decays → Charge Asymmetry in Semileptonic Decays of K_S at KLOE-2 → **Separation of CPTV in the Mass Matrix of $K^0 - \bar{K}^0$**
- Flavour-CP Transitions in Entangled $B^0 - \bar{B}^0$ have demonstrated **Genuine Separate Asymmetries for T and CP with high statistical significance** and compatibility with CPT invariance (with a 2σ tension)
- KLOE-2 is able to accomplish a **complete Program of Genuine Separate Asymmetries for T, CP and CPT in Flavour-CP Transitions for $K^0 - \bar{K}^0$** → Possible fake effects controlled in the same experiment
- The best way to study the **ω -effect** weakening Entanglement due to ill-defined CPT → **CPV ($\pi^+ \pi^-$, $\pi^+ \pi^-$) Correlated Decay at KLOE-2** → Distinguished signature from CPTV in the Effective Hamiltonian for $K^0 - \bar{K}^0$

**THANK YOU
VERY MUCH FOR
YOUR ATTENTION**

BACK-UP

CAN TR BE TESTED IN UNSTABLE SYSTEMS?

THE FACTS

- Taking as Reference $K^0 \rightarrow \bar{K}^0$ and calling (X,Y) the observed decays at times t_1 and t_2 , with $\Delta t \equiv t_2 - t_1 > 0$, the CP, T and CPT transformed transitions are

Transition	$K^0 \rightarrow \bar{K}^0$	$\bar{K}^0 \rightarrow K^0$	$\bar{K}^0 \rightarrow K^0$	$K^0 \rightarrow \bar{K}^0$	$K^0 \rightarrow \bar{K}^0$
(X,Y)	(l ⁻ , l ⁻)	(l ⁺ , l ⁺)	(l ⁺ , l ⁺)	(l ⁻ , l ⁻)	(l ⁻ , l ⁻)
Transformation	Reference	CP	T	CPT	Δt

➡ No way to separate T and CP if T were defined.

- T-operator is not defined for **decaying** states: its time reverse is not a physical state.
- The Kabir asymmetry NEEDS the interference of CP mixing with the “initial state interaction” to generate the effect, directly proportional to $\Delta\Gamma$.

The decay plays an essential role

- The time evolutions of $K^0 \rightarrow \bar{K}^0$ and $\bar{K}^0 \rightarrow K^0$ are equal, the asymmetry is time independent.

- In the WW approach, the entire effect comes from the overlap of non-orthogonal K_L , K_S states. If the **stationary** states were orthogonal ➡ no asymmetry.

- L. Wolfenstein: “it is not as direct a test of TRV as one might like”.

