

# 3D Parton Distributions: Path to the LHC

November 29<sup>th</sup> - December 2<sup>nd</sup>, 2016 - LNF, Frascati, Italy

Studies of TMDs at  hermes

# 21 years ago



ELSEVIER

Nuclear Physics B 461 (1996) 197–237

NUCLEAR  
PHYSICS B

## The complete tree-level result up to order $1/Q$ for polarized deep-inelastic leptonproduction

P.J. Mulders<sup>a,b</sup>, R.D. Tangerman<sup>a</sup>

<sup>a</sup> *National Institute for Nuclear Physics and High-Energy Physics (NIKHEF), P.O. Box 41882, NL-1009 DB Amsterdam, The Netherlands*

<sup>b</sup> *Department of Physics and Astronomy, Free University, De Boelelaan 1081, NL-1081 HV Amsterdam, The Netherlands*

Received 18 October 1995; accepted 1 December 1995

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### Abstract

We present the results of the tree-level calculation of deep-inelastic leptonproduction, including polarization of target hadron and produced hadron. We also discuss the dependence on transverse momenta of the quarks, which leads to azimuthal asymmetries for the produced hadrons.

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- use semi-inclusive DIS for

- accessing the full momentum structure

Abstract

- parton polarimetry

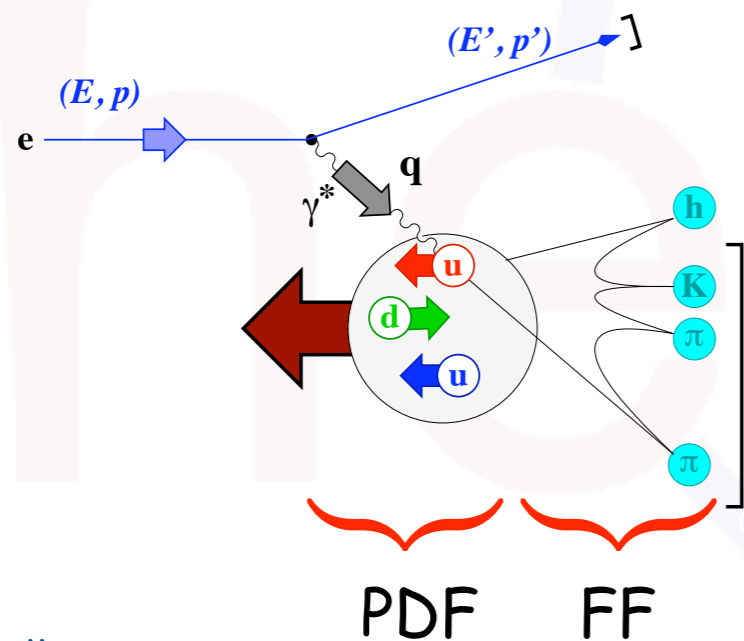
We present the complete tree-level calculation of deep-inelastic leptonproduction, including polarization of target hadron and produced hadron. We also discuss the dependence on transverse momenta of the quarks, which leads to azimuthal asymmetries for the produced hadrons.

# probing TMDs in semi-inclusive DIS

quark pol.

|              |   |                |                     |
|--------------|---|----------------|---------------------|
|              | U | L              | T                   |
| nucleon pol. | U |                | $h_1^\perp$         |
|              | L | $g_{1L}$       | $h_{1L}^\perp$      |
|              | T | $f_{1T}^\perp$ | $h_1, h_{1T}^\perp$ |

in SIDIS\*) couple PDFs to:



\*) semi-inclusive DIS with unpolarized final state

# probing TMDs in semi-inclusive DIS

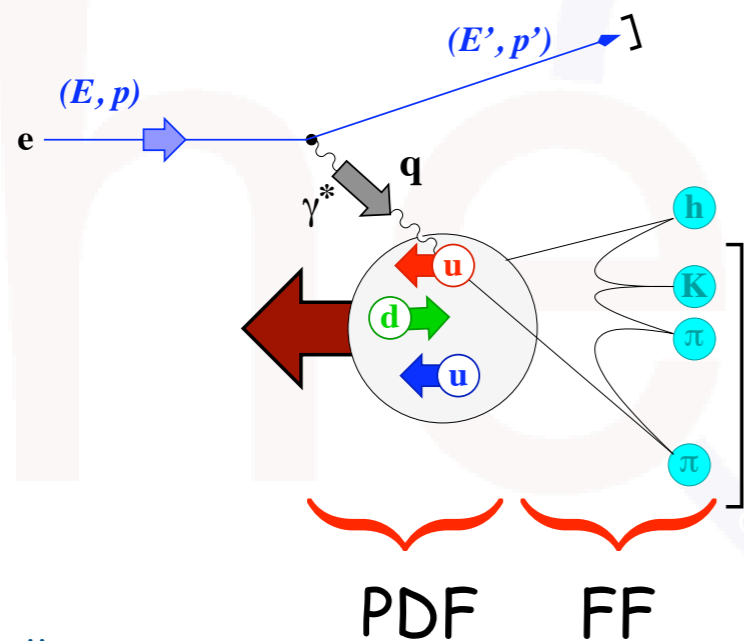
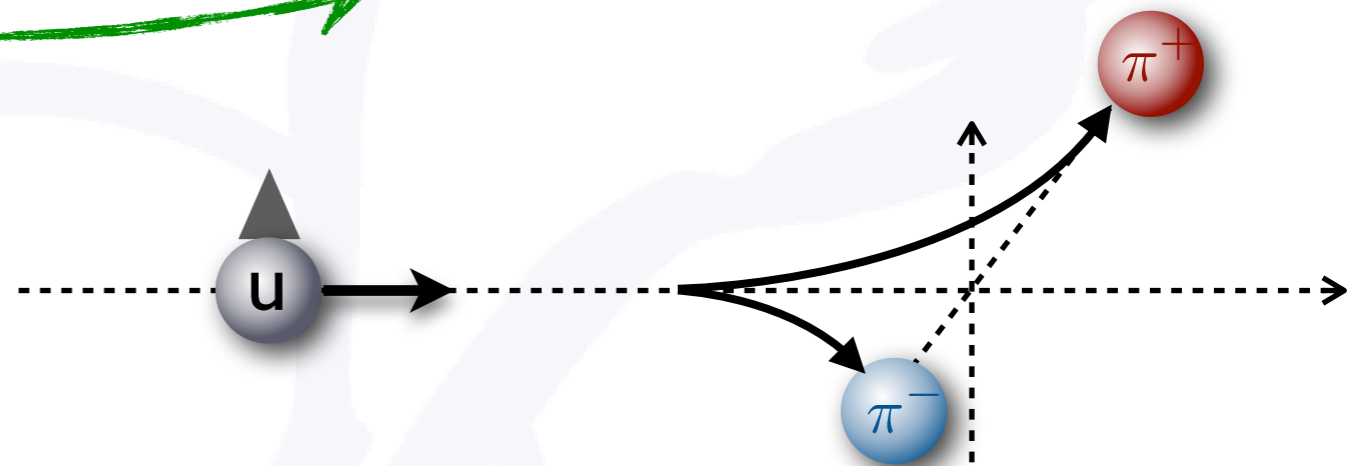
quark pol.

|   | U              | L        | T                   |
|---|----------------|----------|---------------------|
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| L |                | $g_{1L}$ | $h_{1L}^\perp$      |
| T | $f_{1T}^\perp$ | $g_{1T}$ | $h_1, h_{1T}^\perp$ |

nucleon pol.

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Collins FF:  $H_1^{\perp, q \rightarrow h}$



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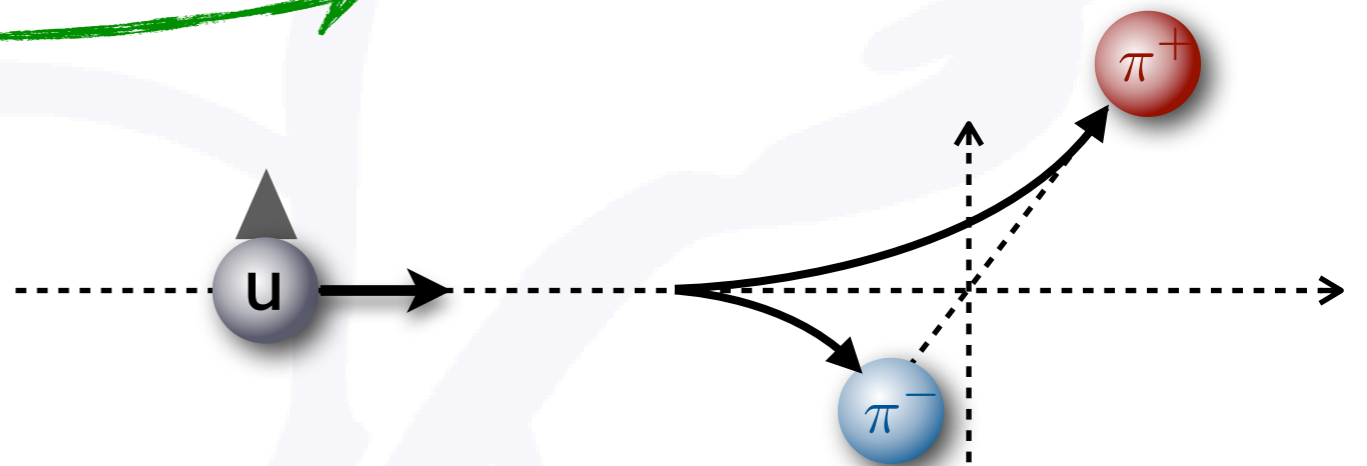
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quark pol.

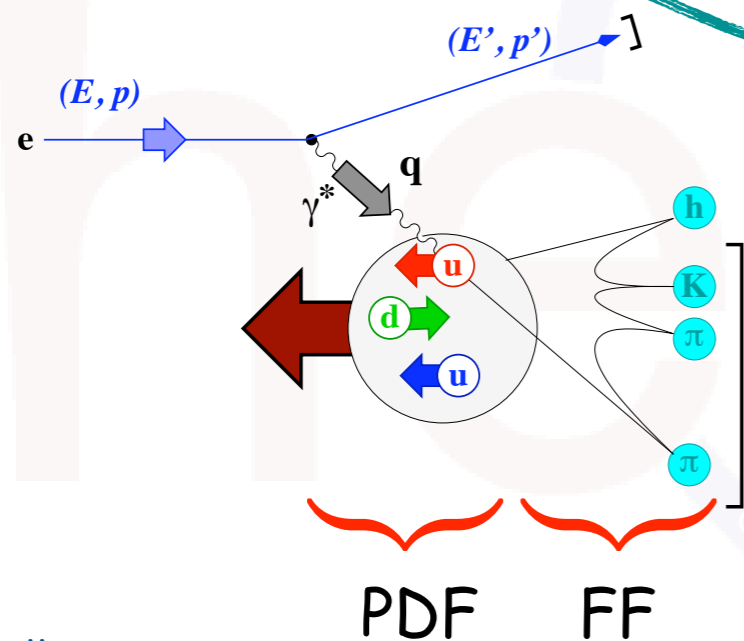
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ordinary FF:  $D_1^{q \rightarrow h}$



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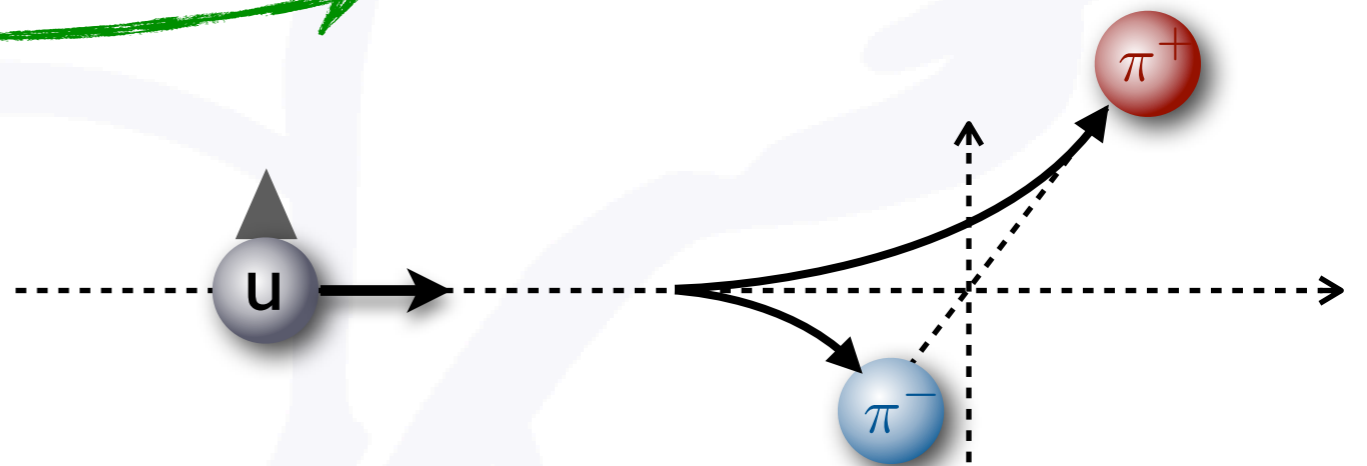
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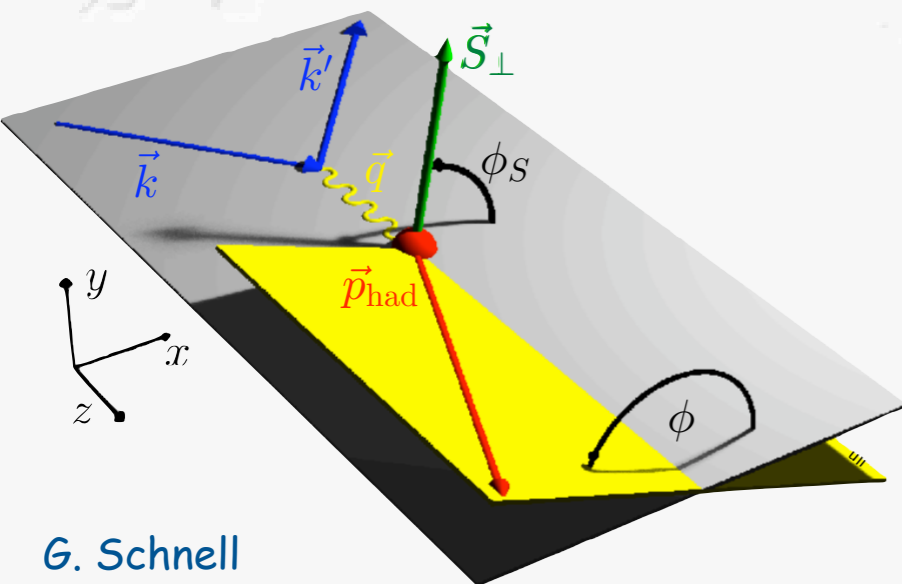
**gives rise to characteristic azimuthal dependences**

\*) semi-inclusive DIS with unpolarized final state

# one-hadron production ( $ep \rightarrow ehX$ )

$$\begin{aligned}
 d\sigma = & d\sigma_{UU}^0 + \cos 2\phi d\sigma_{UU}^1 + \frac{1}{Q} \cos \phi d\sigma_{UU}^2 + \lambda_e \frac{1}{Q} \sin \phi d\sigma_{LU}^3 \\
 & + S_L \left\{ \sin 2\phi d\sigma_{UL}^4 + \frac{1}{Q} \sin \phi d\sigma_{UL}^5 + \lambda_e \left[ d\sigma_{LL}^6 + \frac{1}{Q} \cos \phi d\sigma_{LL}^7 \right] \right\} \\
 & + S_T \left\{ \sin(\phi - \phi_S) d\sigma_{UT}^8 + \sin(\phi + \phi_S) d\sigma_{UT}^9 + \sin(3\phi - \phi_S) d\sigma_{UT}^{10} \frac{1}{Q} \right. \\
 & \quad \left. + \frac{1}{Q} (\sin(2\phi - \phi_S) d\sigma_{UT}^{11} + \sin \phi_S d\sigma_{UT}^{12}) \right. \\
 & \quad \left. + \lambda_e \left[ \cos(\phi - \phi_S) d\sigma_{LT}^{13} + \frac{1}{Q} (\cos \phi_S d\sigma_{LT}^{14} + \cos(2\phi - \phi_S) d\sigma_{LT}^{15}) \right] \right\}
 \end{aligned}$$

$\sigma_{XY}$   
 ↙ ↘  
**Beam Target**  
**Polarization**



Mulders and Tangerman, Nucl. Phys. B 461 (1996) 197

Boer and Mulders, Phys. Rev. D 57 (1998) 5780

Bacchetta et al., Phys. Lett. B 595 (2004) 309

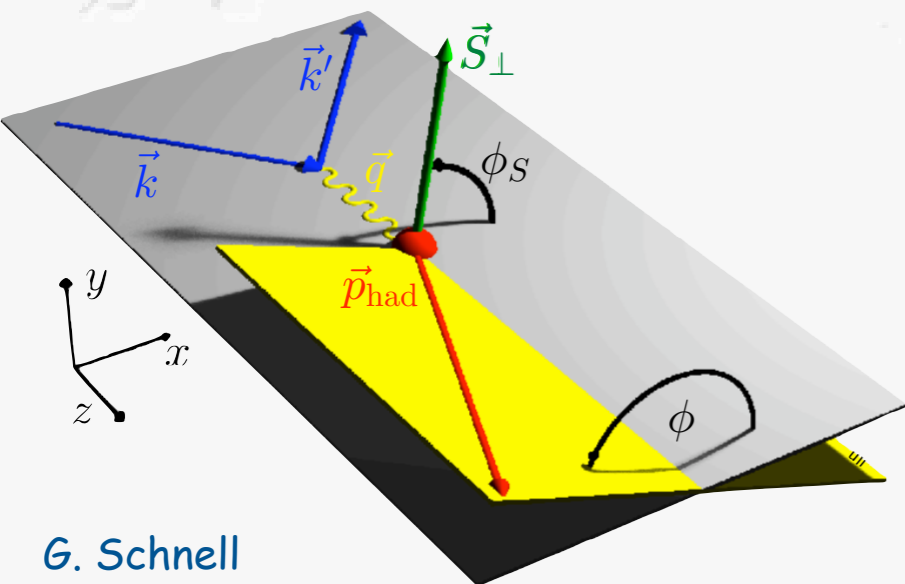
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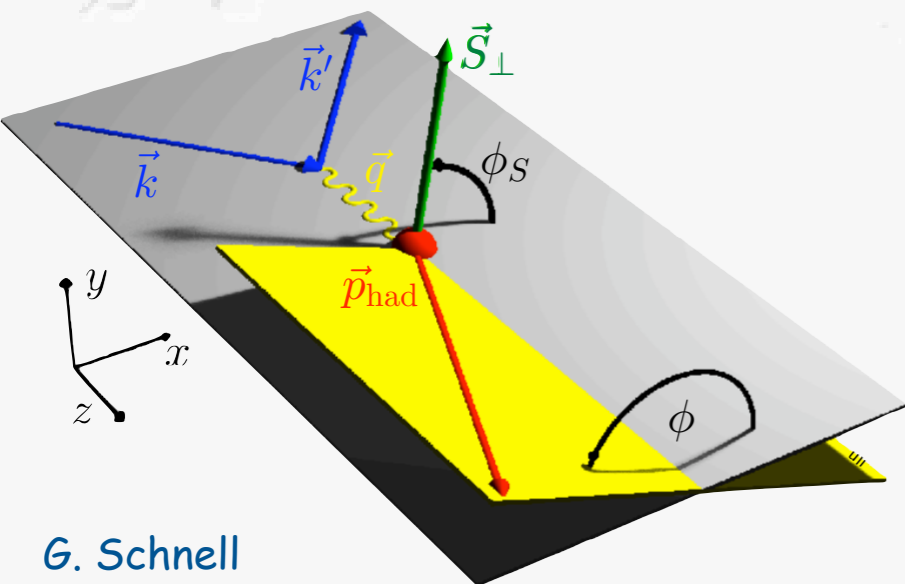
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# The HERMES Experiment

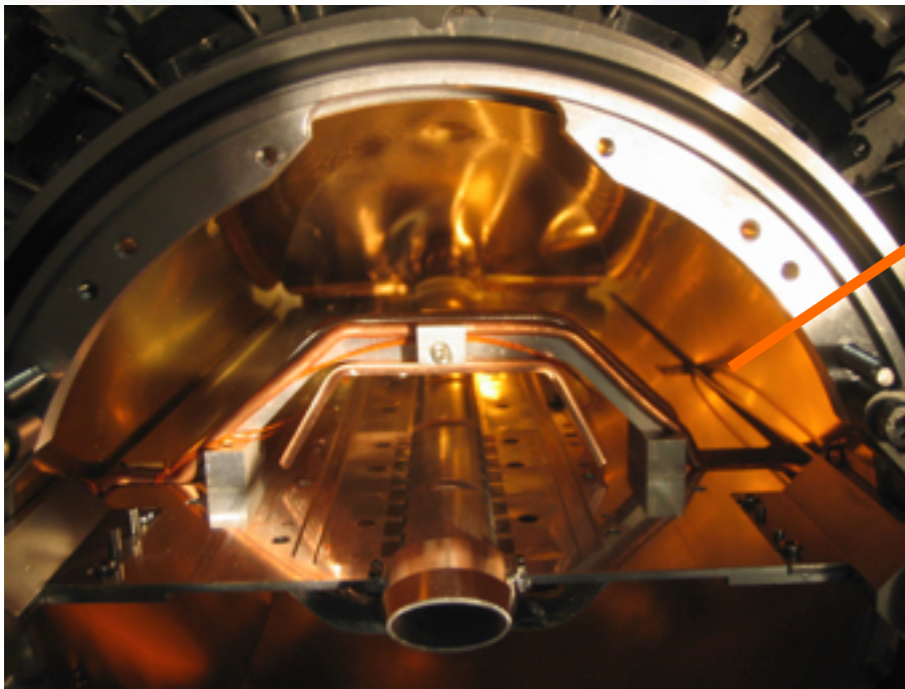
- 27.6 GeV HERA  $e^+/e^-$  beam



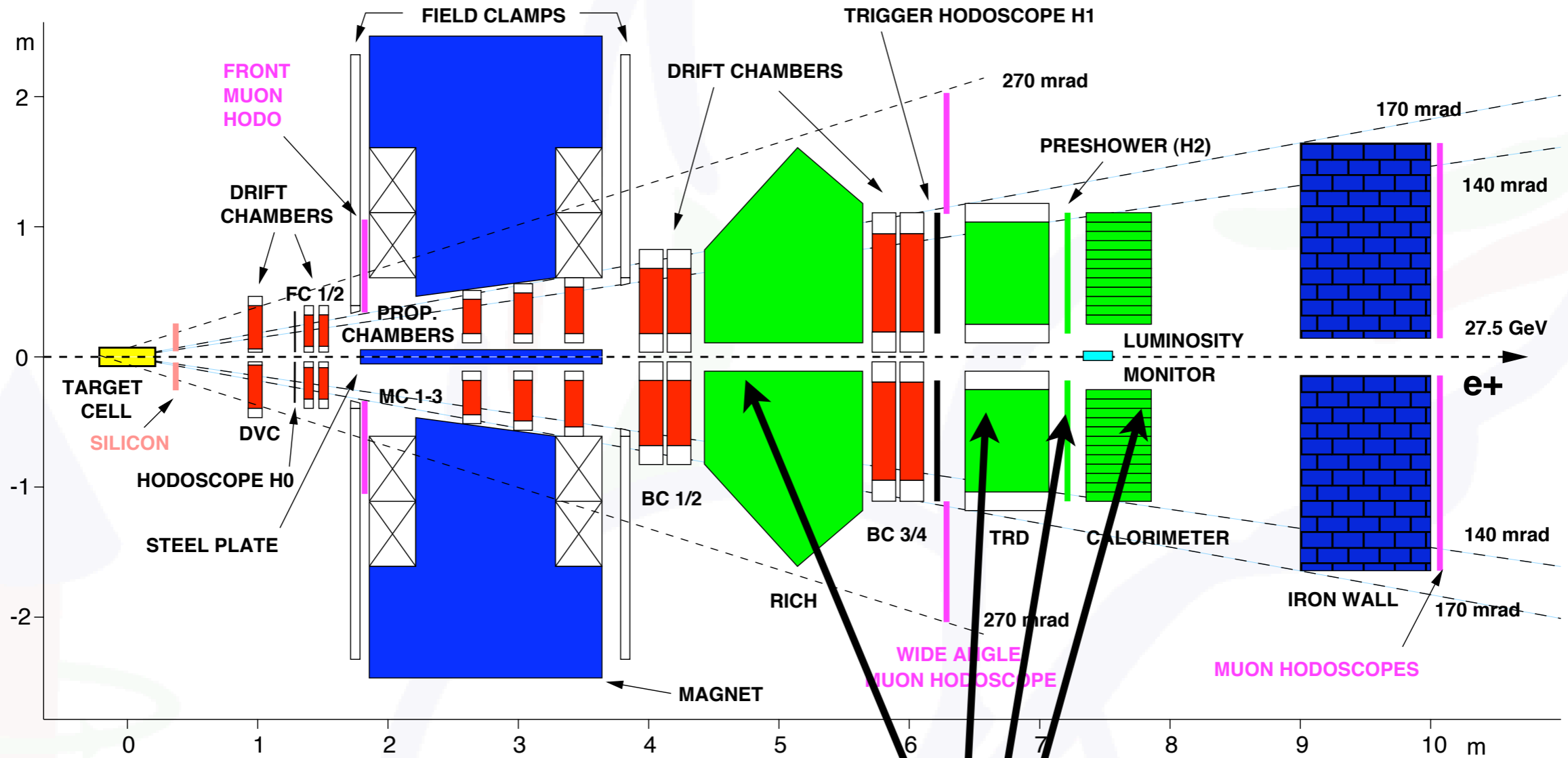
- longitudinally polarized

# The HERMES Experiment

- pure gas targets
- internal to lepton ring
- unpolarized ( $^1\text{H}$  ... Xe)
- long. polarized:  $^1\text{H}$ ,  $^2\text{H}$ ,  $^3\text{He}$
- transversely polarized:  $^1\text{H}$



# HERMES schematically



- pure gas targets internal to HERA 27.6 GeV lepton ring
- unpolarized ( $^1\text{H}$  ... Xe)
- long. polarized:  $^1\text{H}$ ,  $^2\text{H}$ ,  $^3\text{He}$
- transversely polarized:  $^1\text{H}$

Particle ID detectors allow for

- lepton/hadron separation
- RICH: pion/kaon/proton discrimination  $2\text{GeV} < p < 15\text{GeV}$

# hadron multiplicities in DIS

$$\frac{d^5\sigma}{dx dy dz d\phi_h dP_{h\perp}^2} \propto \left(1 + \frac{\gamma^2}{2x}\right) \{F_{UU,T} + \epsilon F_{UU,L} + \sqrt{2\epsilon(1-\epsilon)} F_{UU}^{\cos\phi_h} \cos\phi_h + \epsilon F_{UU}^{\cos 2\phi_h} \cos 2\phi_h\}$$

$$F_{XY,Z} = F_{XY,Z}(x, y, z, P_{h\perp})$$

target polarization  $\downarrow$   
 beam polarization  $\uparrow$     virtual-photon polarization  $\uparrow$

[see, e.g., Bacchetta et al., JHEP 0702 (2007) 093]

$$\gamma = \frac{2Mx}{Q}$$

$$\epsilon = \frac{1 - y - \frac{1}{4}\gamma^2 y^2}{1 - y + \frac{1}{2}y^2 + \frac{1}{4}\gamma^2 y^2}$$



# hadron multiplicities in DIS

hadron multiplicity:  
normalize to inclusive DIS  
cross section

$$\frac{d^2 \sigma^{\text{incl. DIS}}}{dxdy} \propto F_T + \epsilon F_L$$

$$\frac{d^4 \mathcal{M}^h(x, y, z, P_{h\perp}^2)}{dxdydzdP_{h\perp}^2} \propto \left(1 + \frac{\gamma^2}{2x}\right) \frac{F_{UU,T} + \epsilon F_{UU,L}}{F_T + \epsilon F_L}$$

$$\approx \frac{\sum_q e_q^2 f_1^q(x, p_T^2) \otimes D_1^{q \rightarrow h}(z, K_T^2)}{\sum_q e_q^2 f_1^q(x)}$$

$$\frac{d^5 \sigma}{dxdydzd\phi_h dP_{h\perp}^2} \propto \left(1 + \frac{\gamma^2}{2x}\right) \left\{ F_{UU,T} + \epsilon F_{UU,L} + \sqrt{2\epsilon(1-\epsilon)} F_{UU}^{\cos \phi_h} \cos \phi_h + \epsilon F_{UU}^{\cos 2\phi_h} \cos 2\phi_h \right\}$$

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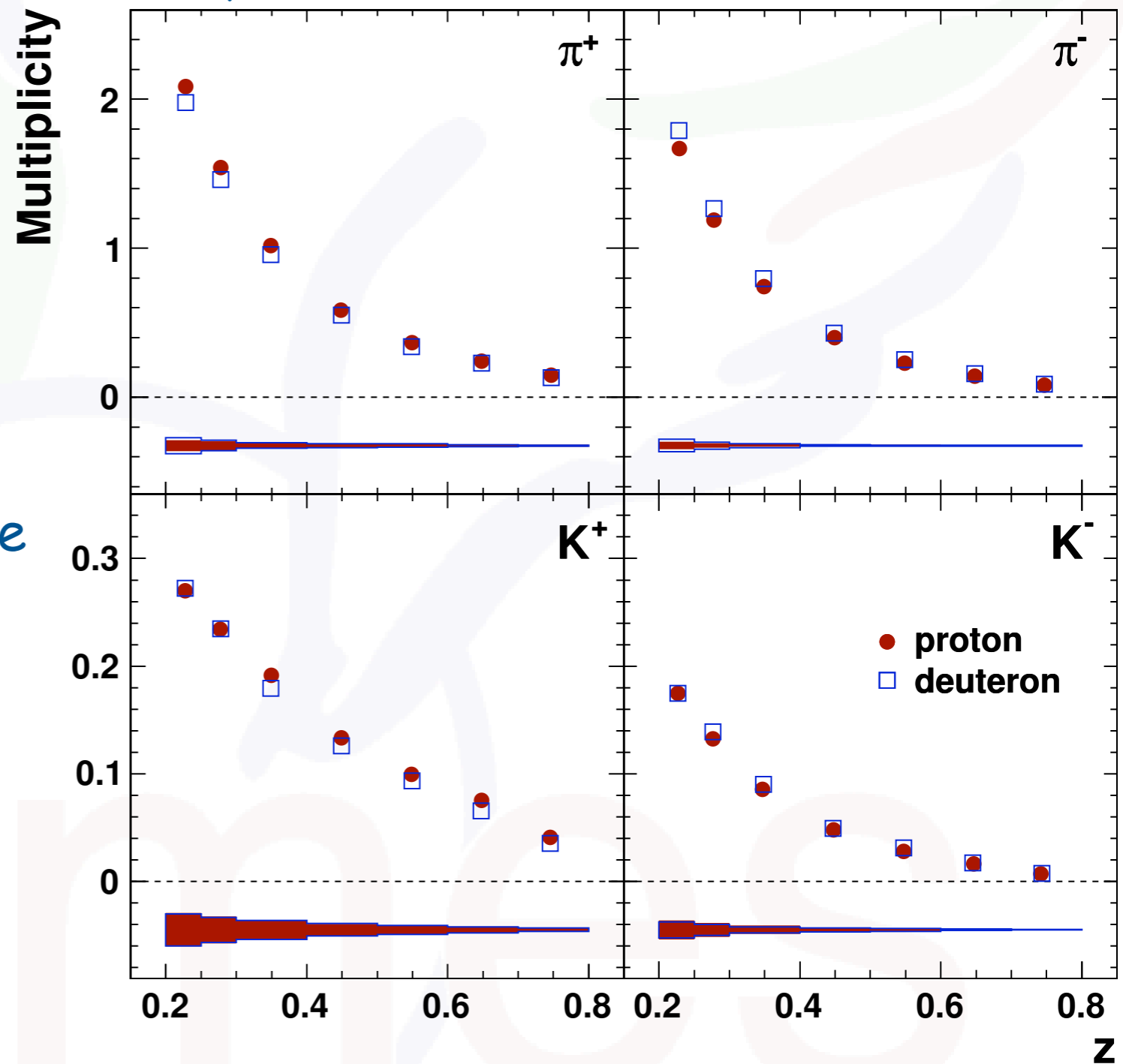
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# multiplicities @ HERMES

- extensive data set on pure proton and deuteron targets for identified charged mesons <http://www-hermes.desy.de/multiplicities>
- extracted in a multi-dimensional unfolding procedure

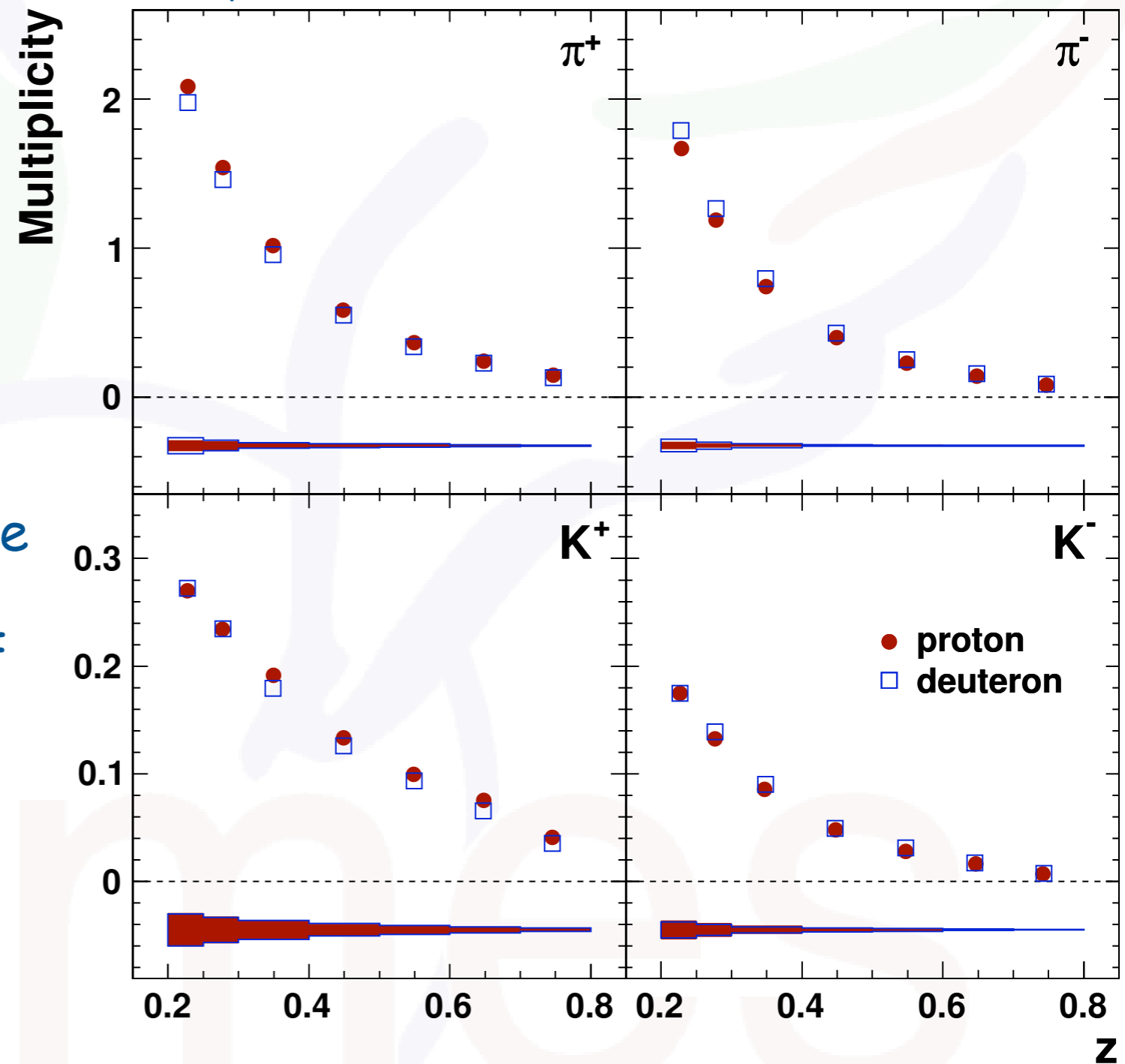
[Airapetian et al., PRD 87 (2013) 074029]



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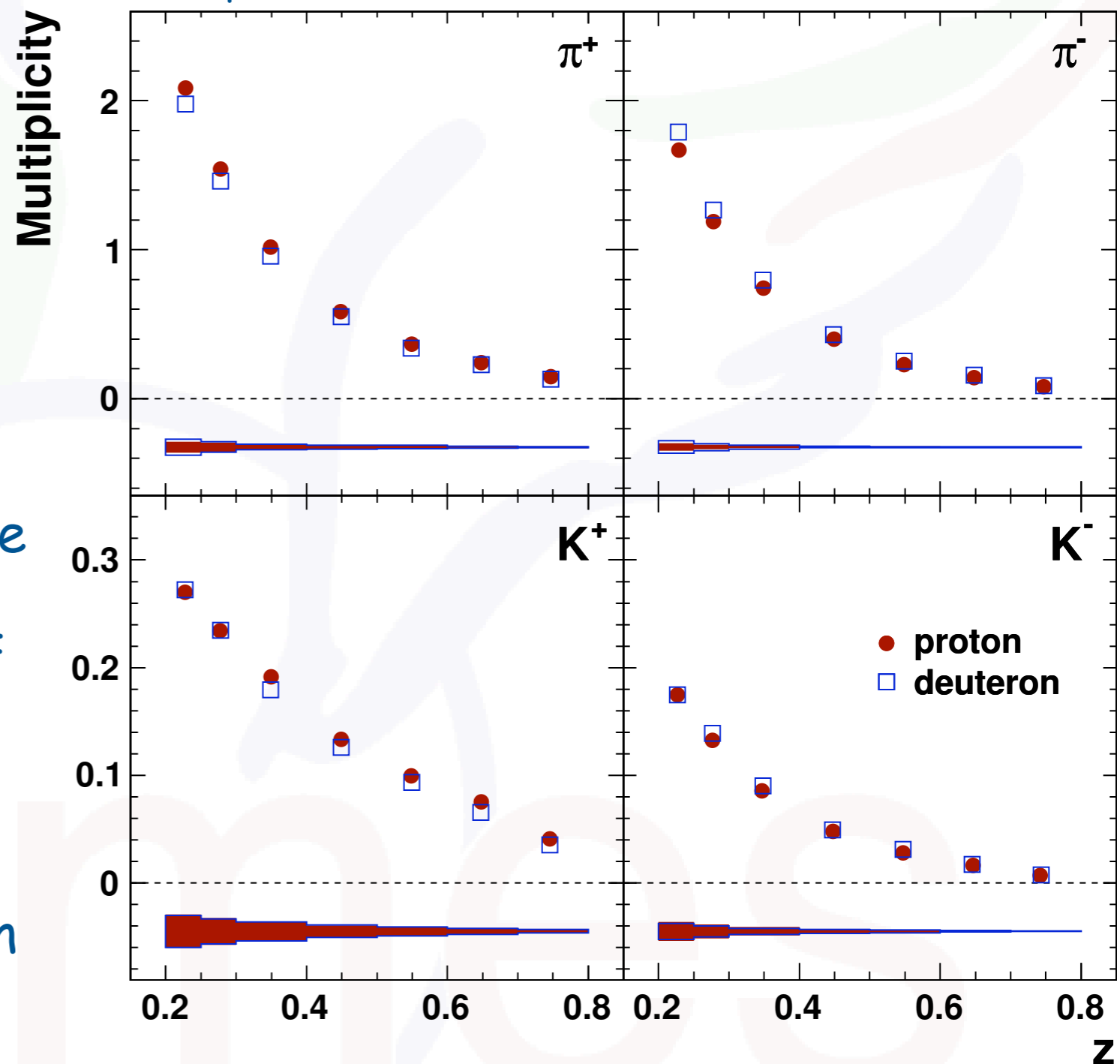
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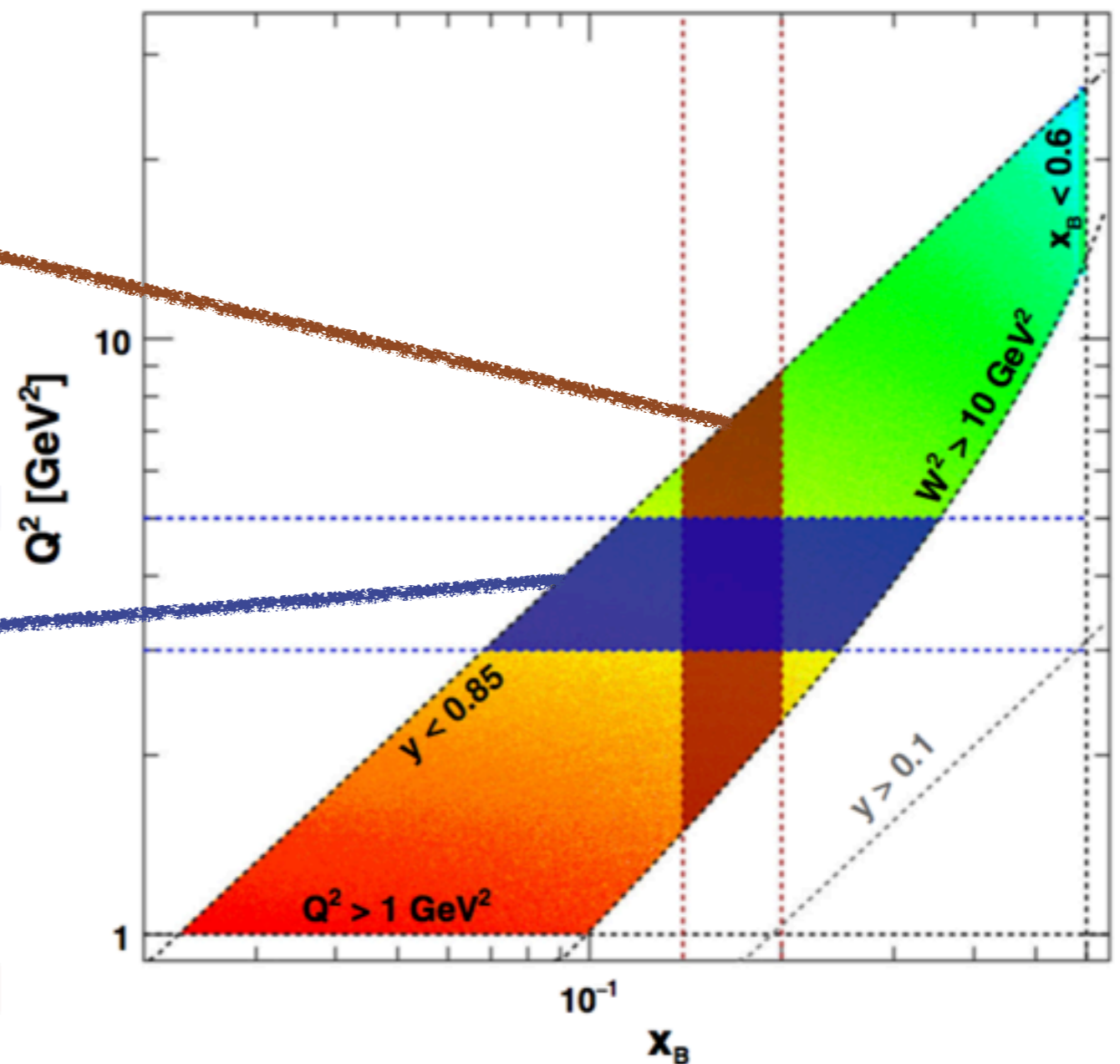
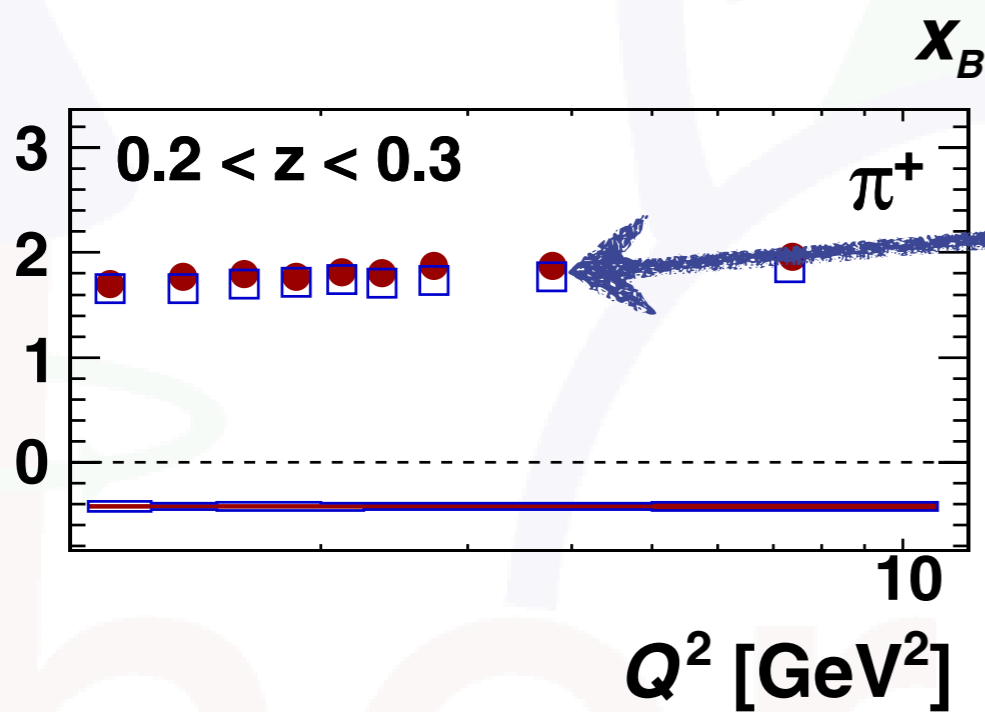
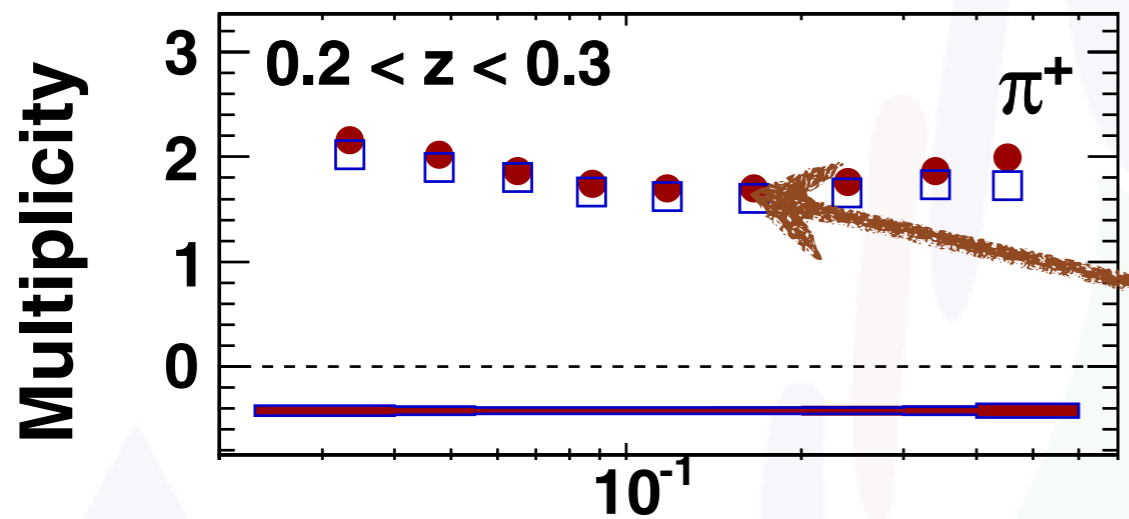
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- access to flavor dependence of fragmentation through different mesons and targets
- input to fragmentation function analyses

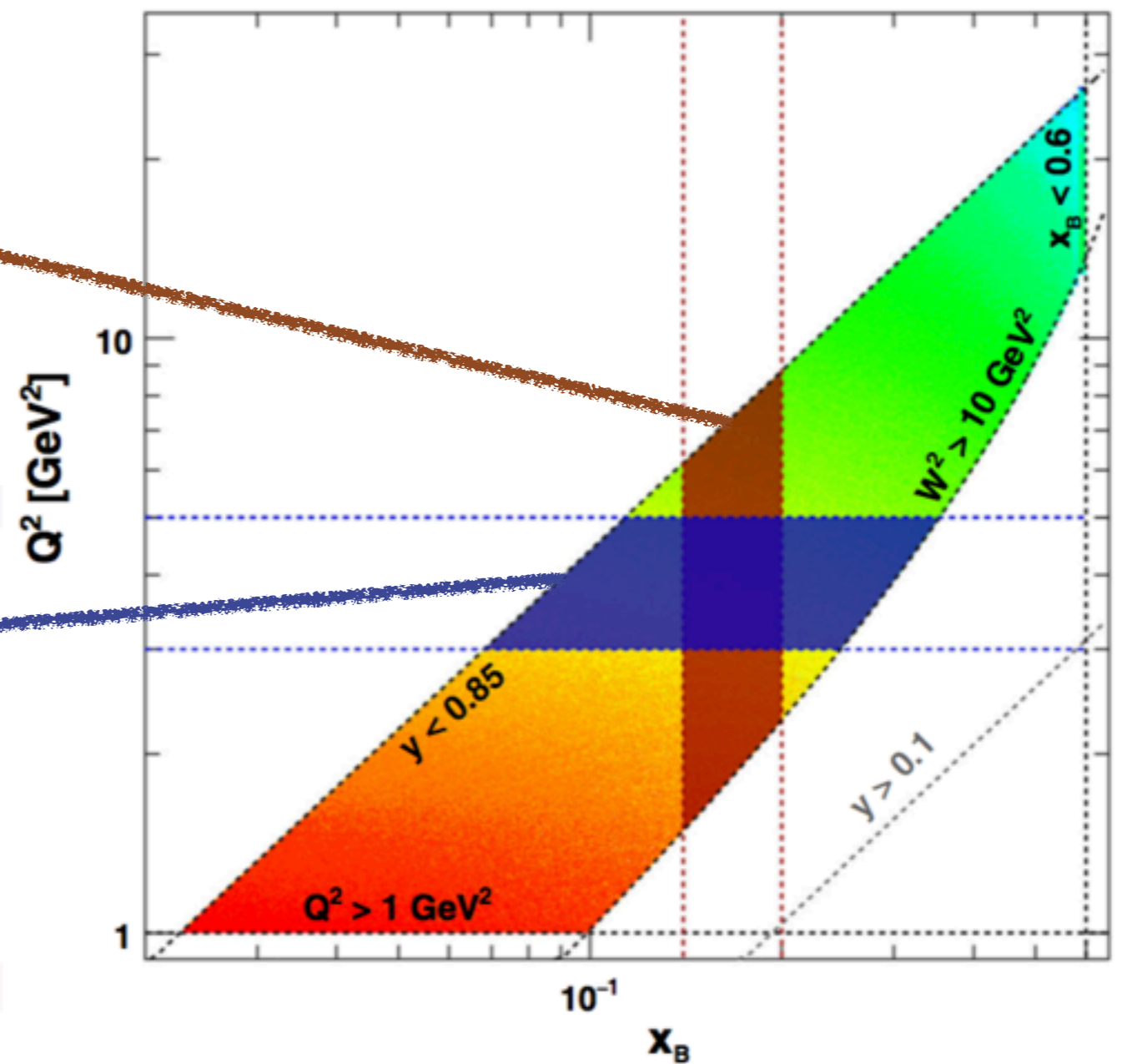
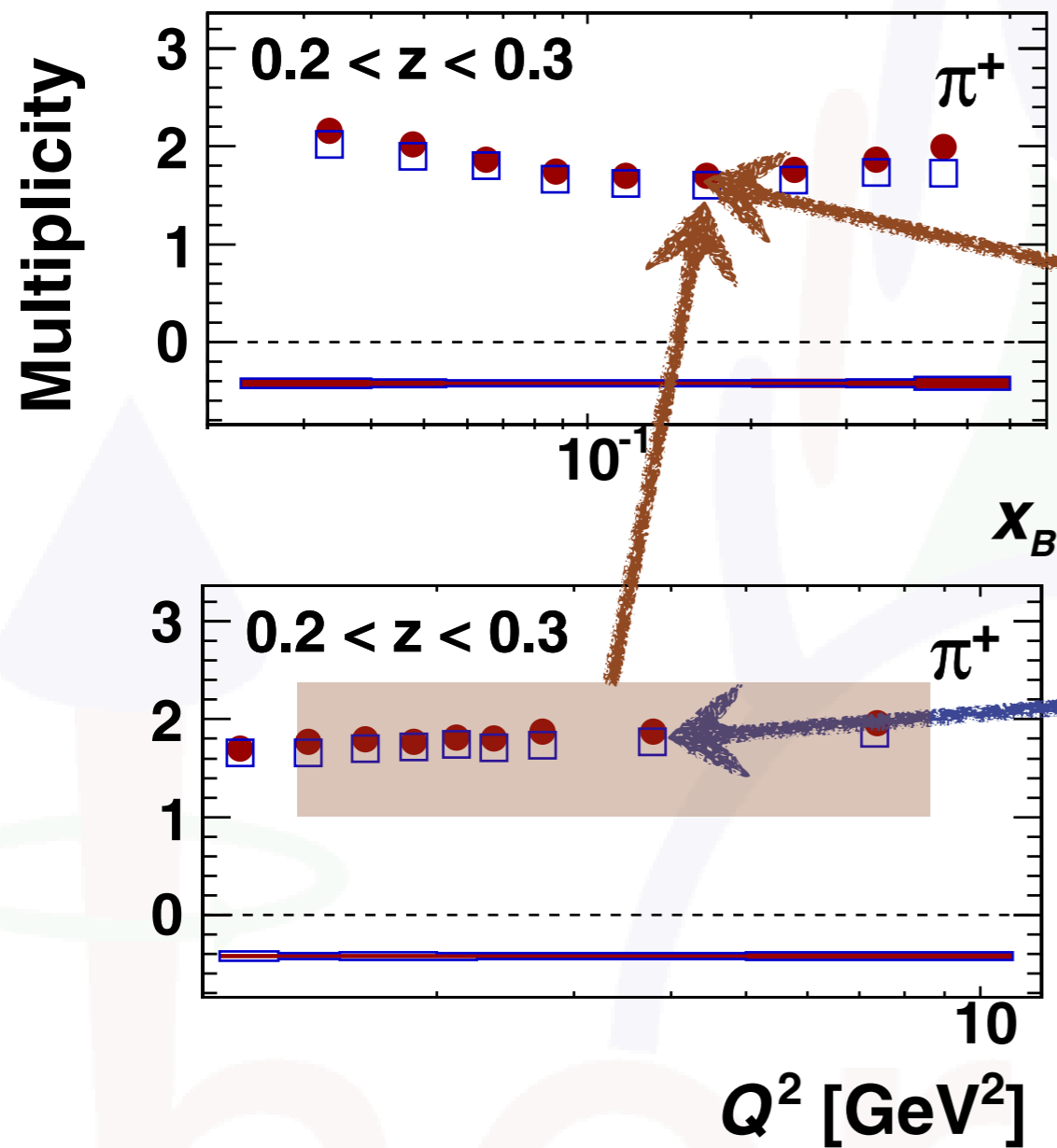
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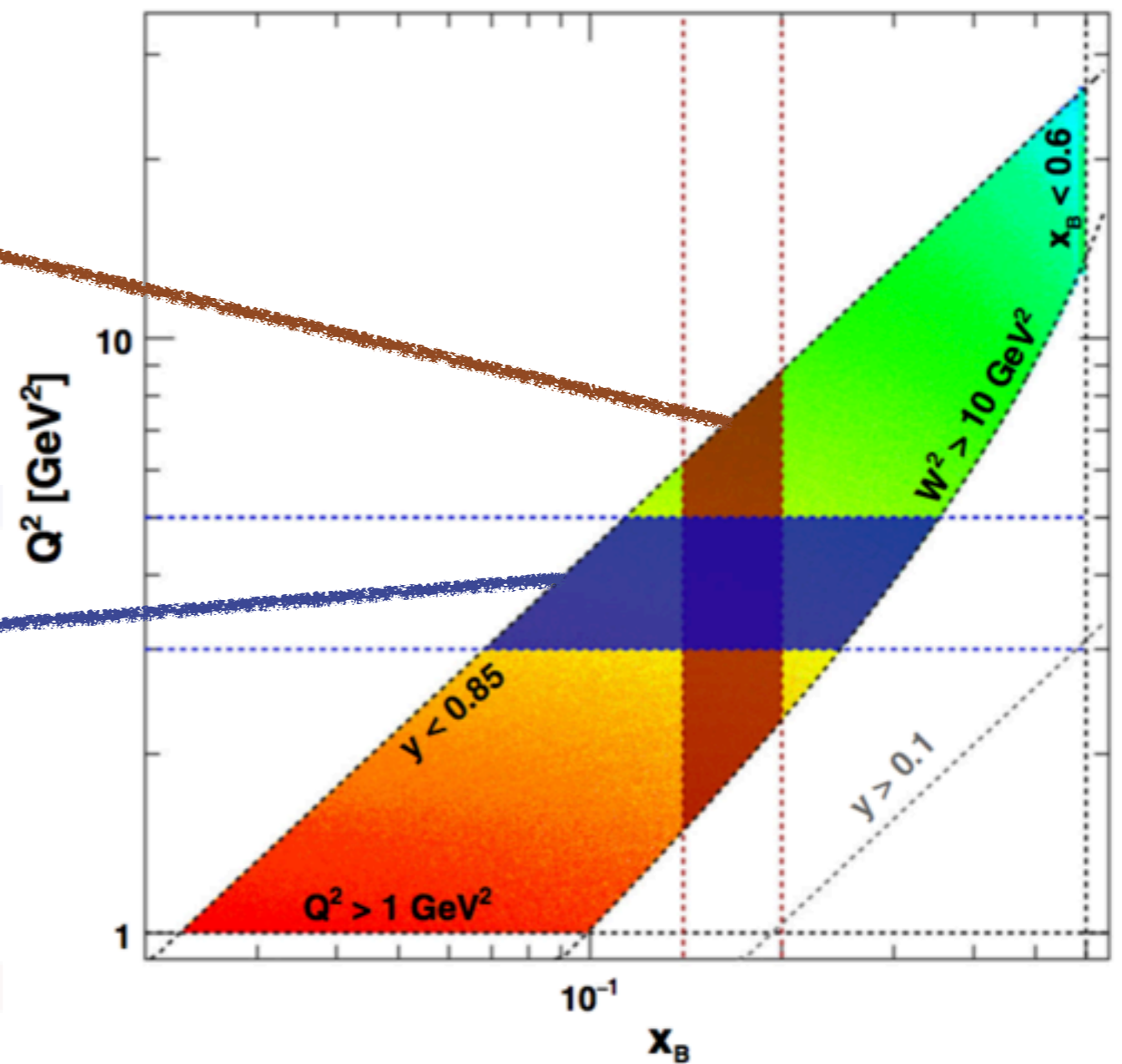
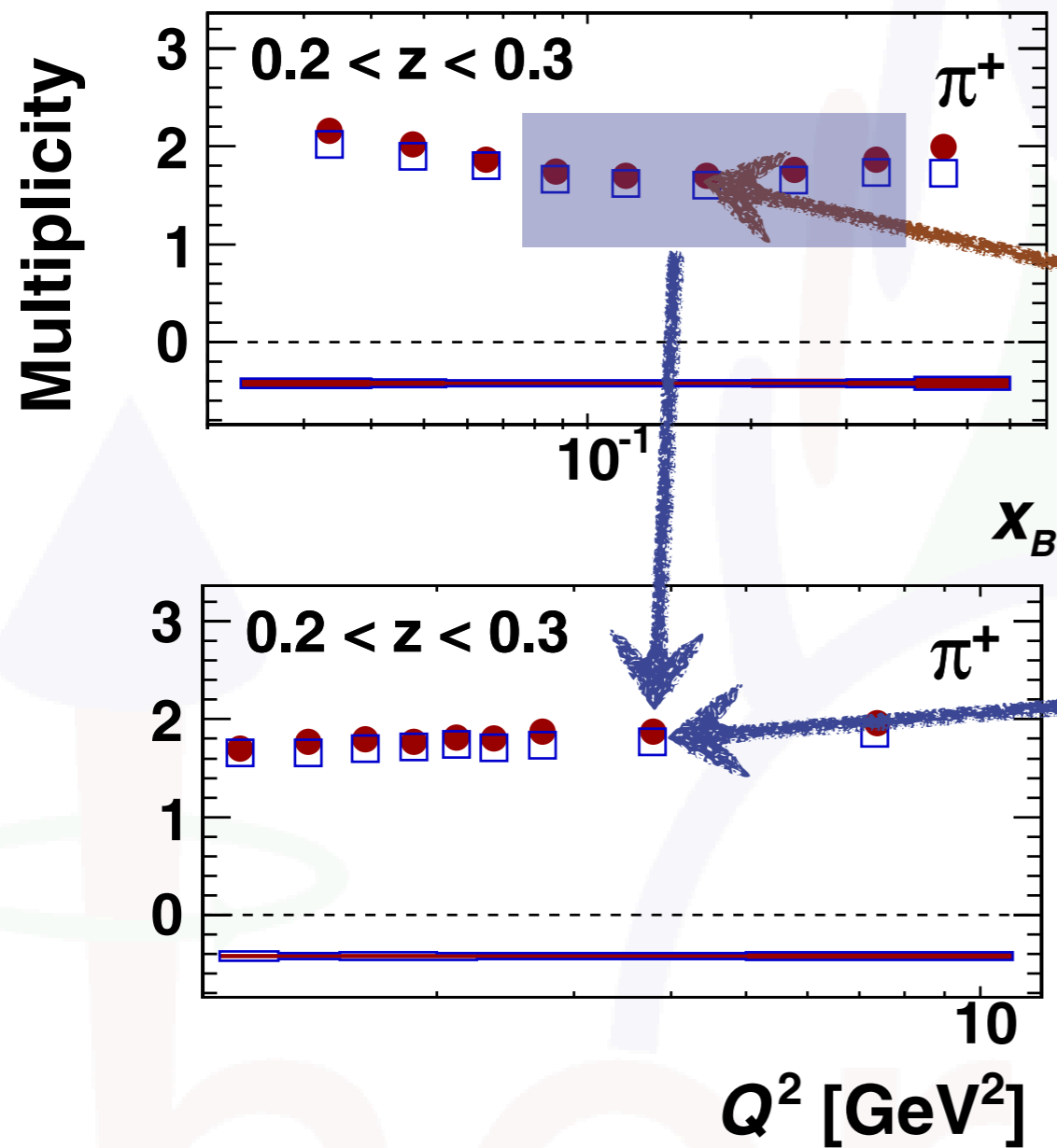
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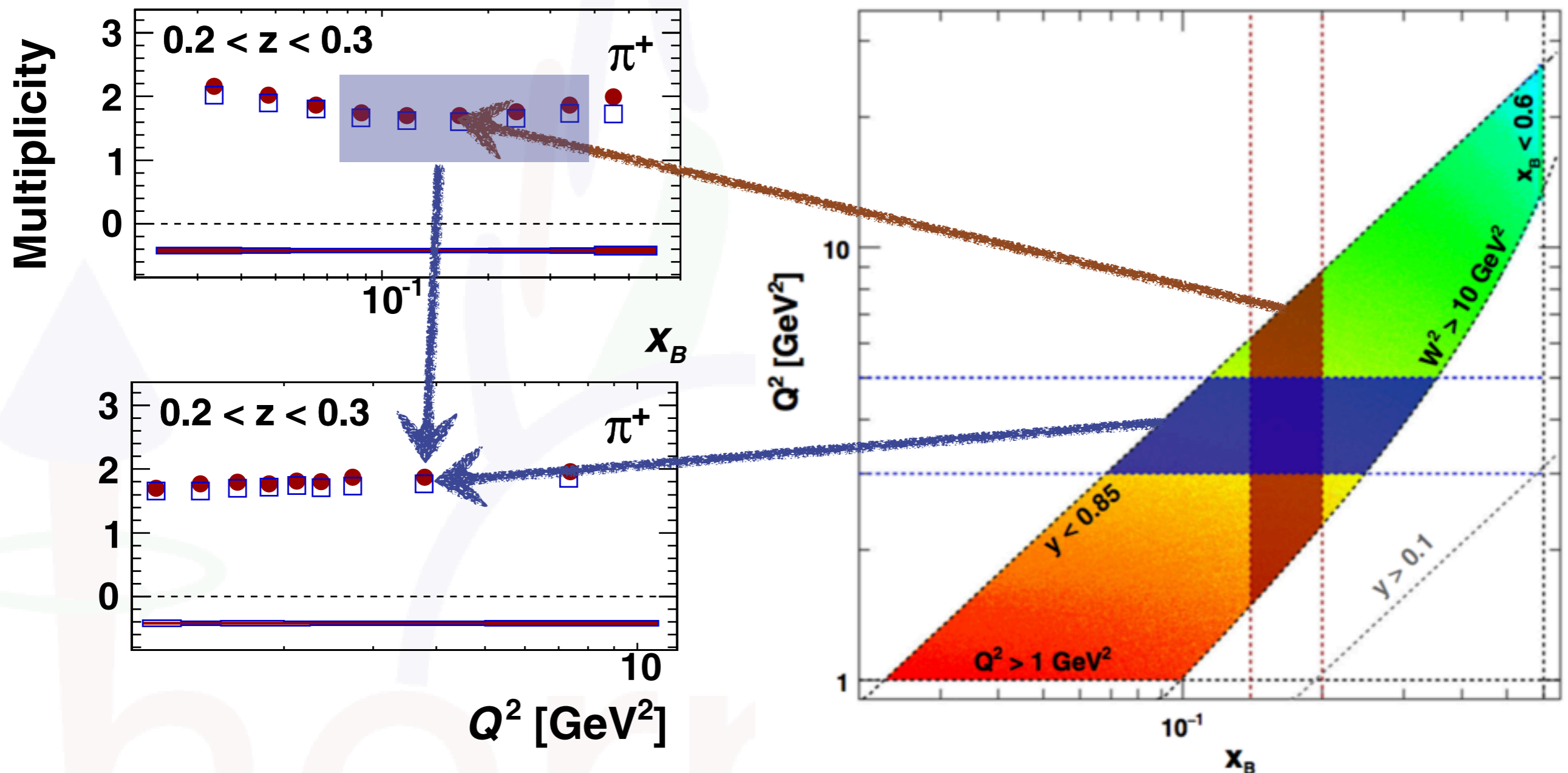
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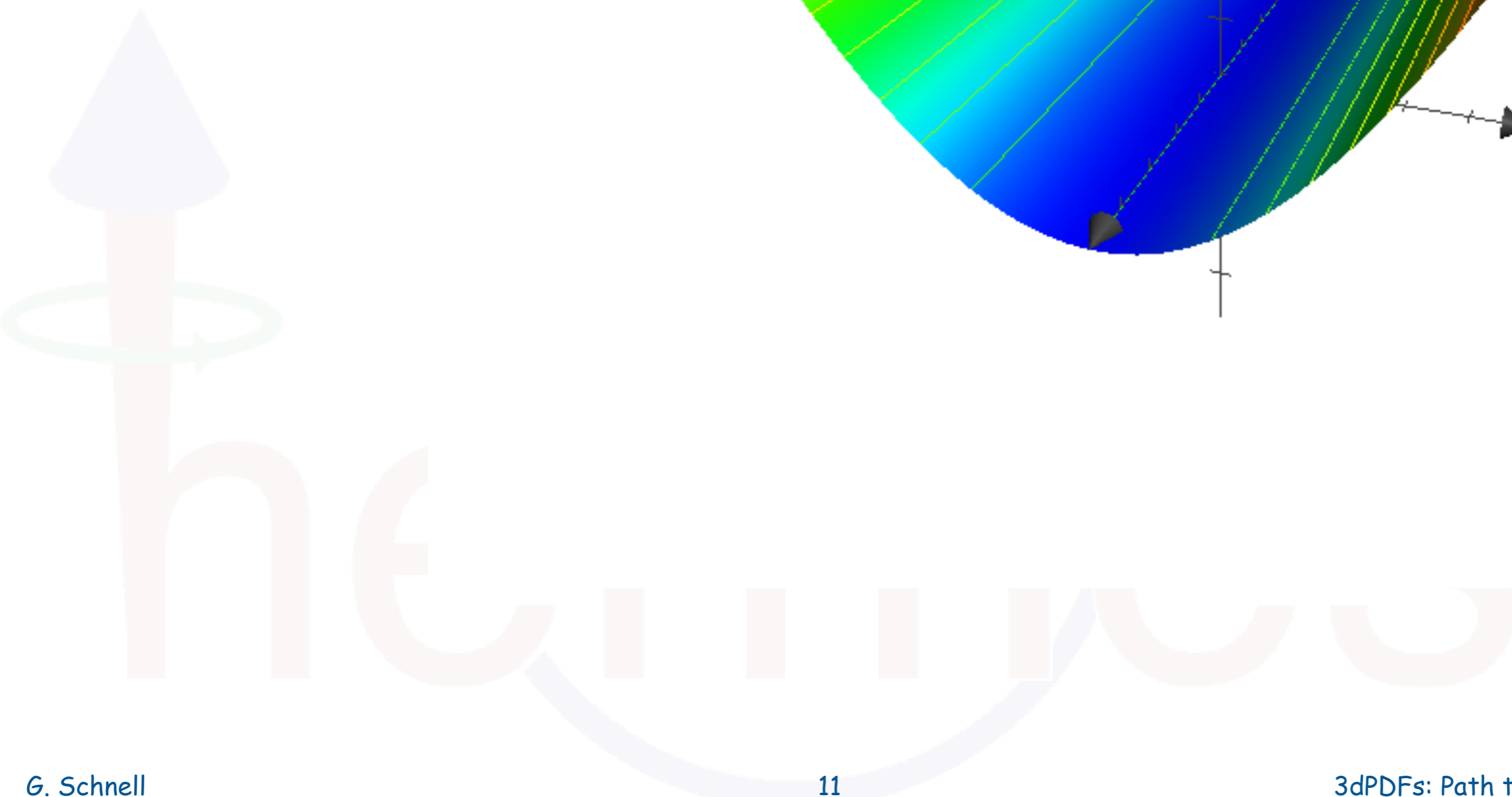
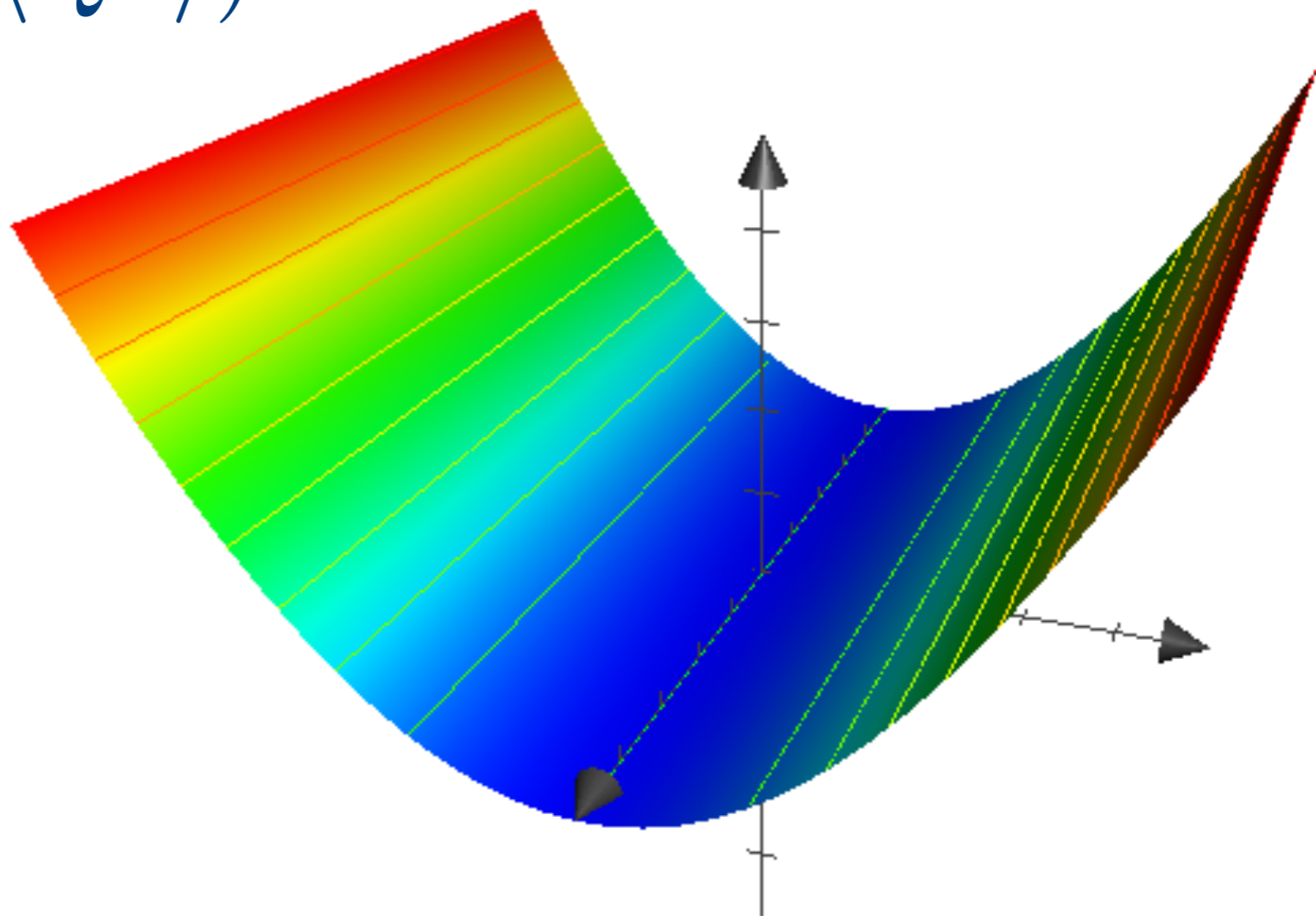
$$\langle \mathcal{M}(Q^2) \rangle_{Q^2} \neq \mathcal{M}(\langle Q^2 \rangle)$$



- even though having similar average kinematics, multiplicities in the two projections are different

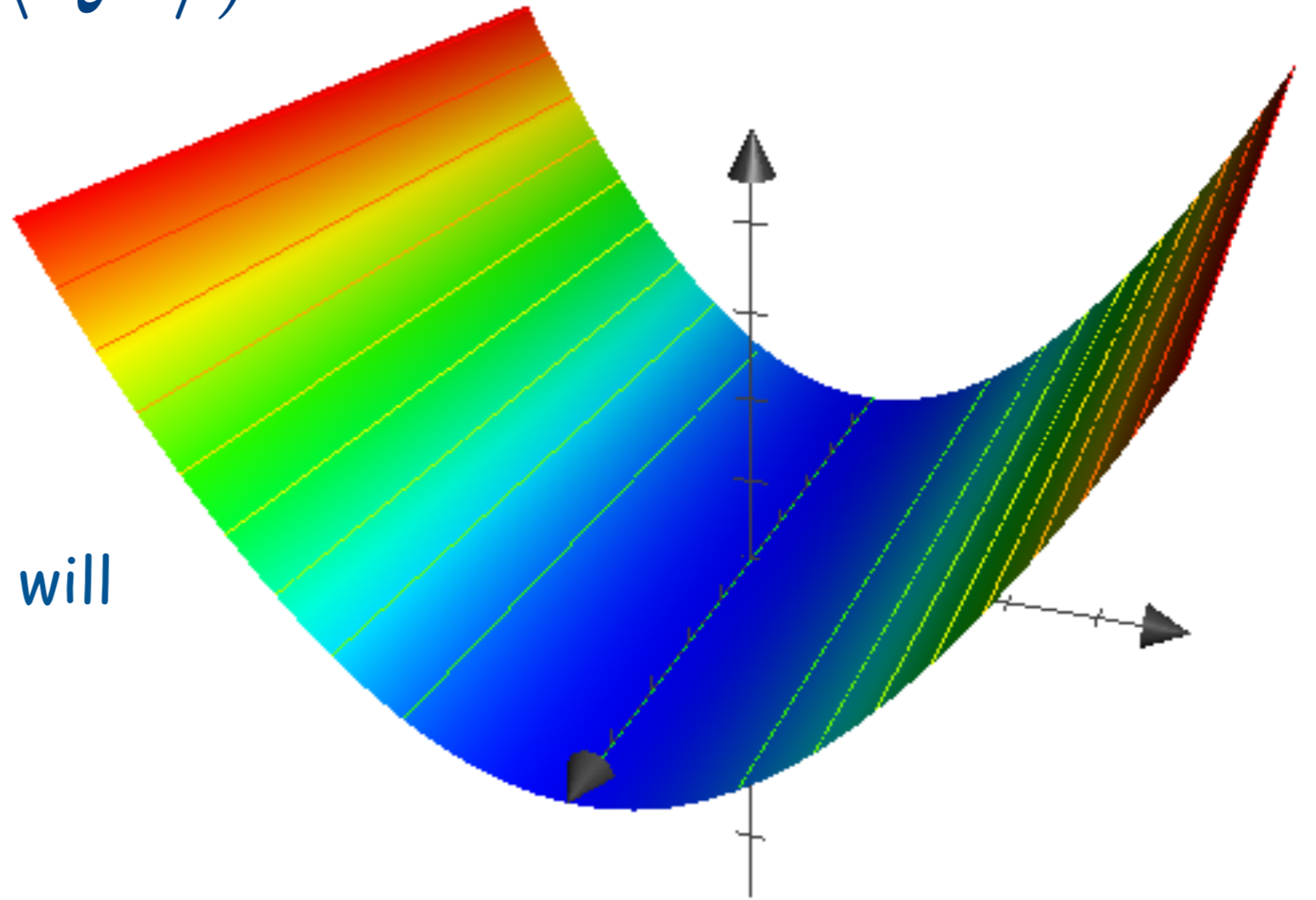


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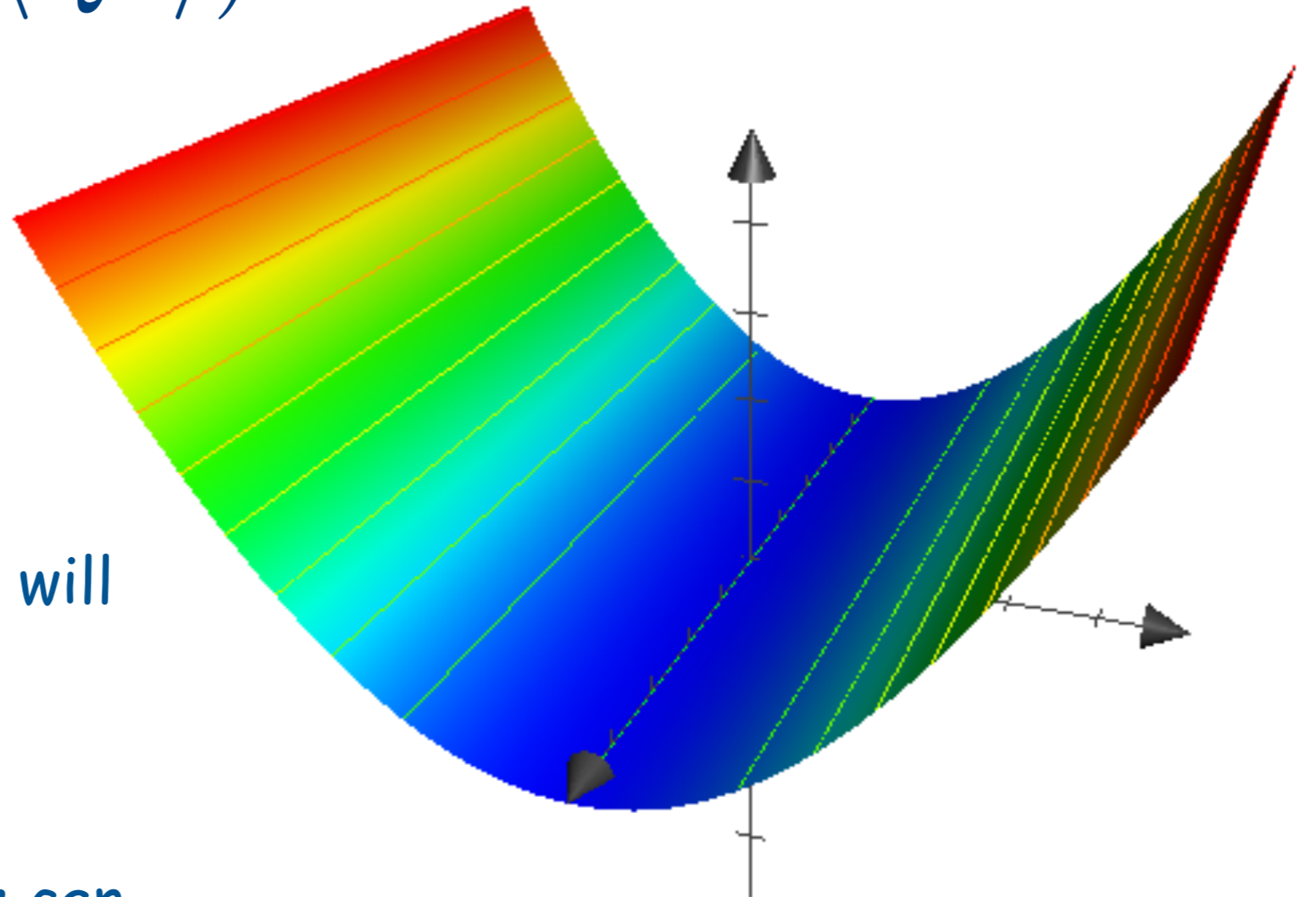
$$\langle \mathcal{M}(Q^2) \rangle_{Q^2} \neq \mathcal{M}(\langle Q^2 \rangle)$$

- the average along the valley will be smaller than the average along the gradient



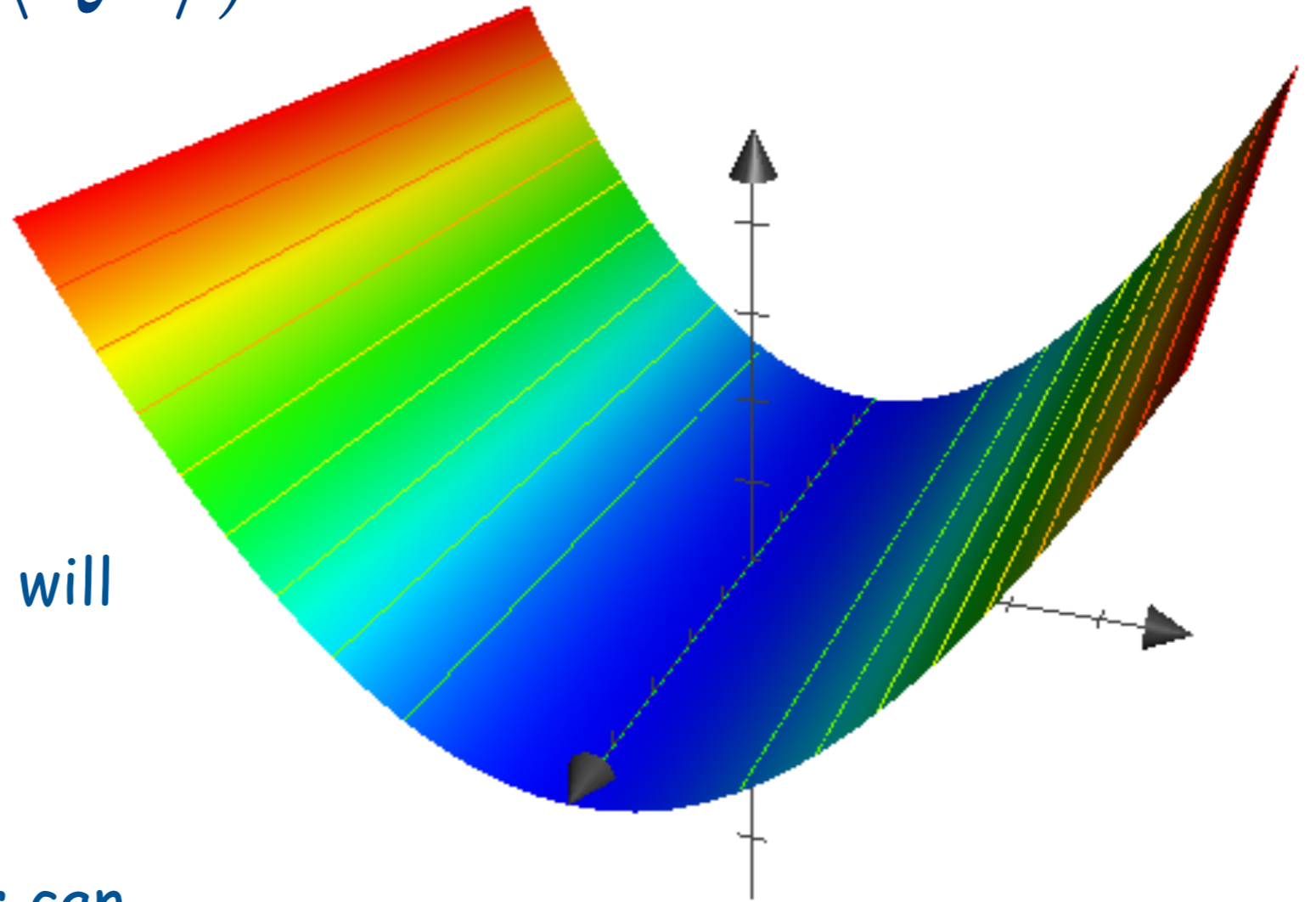
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- still the average kinematics can be the same



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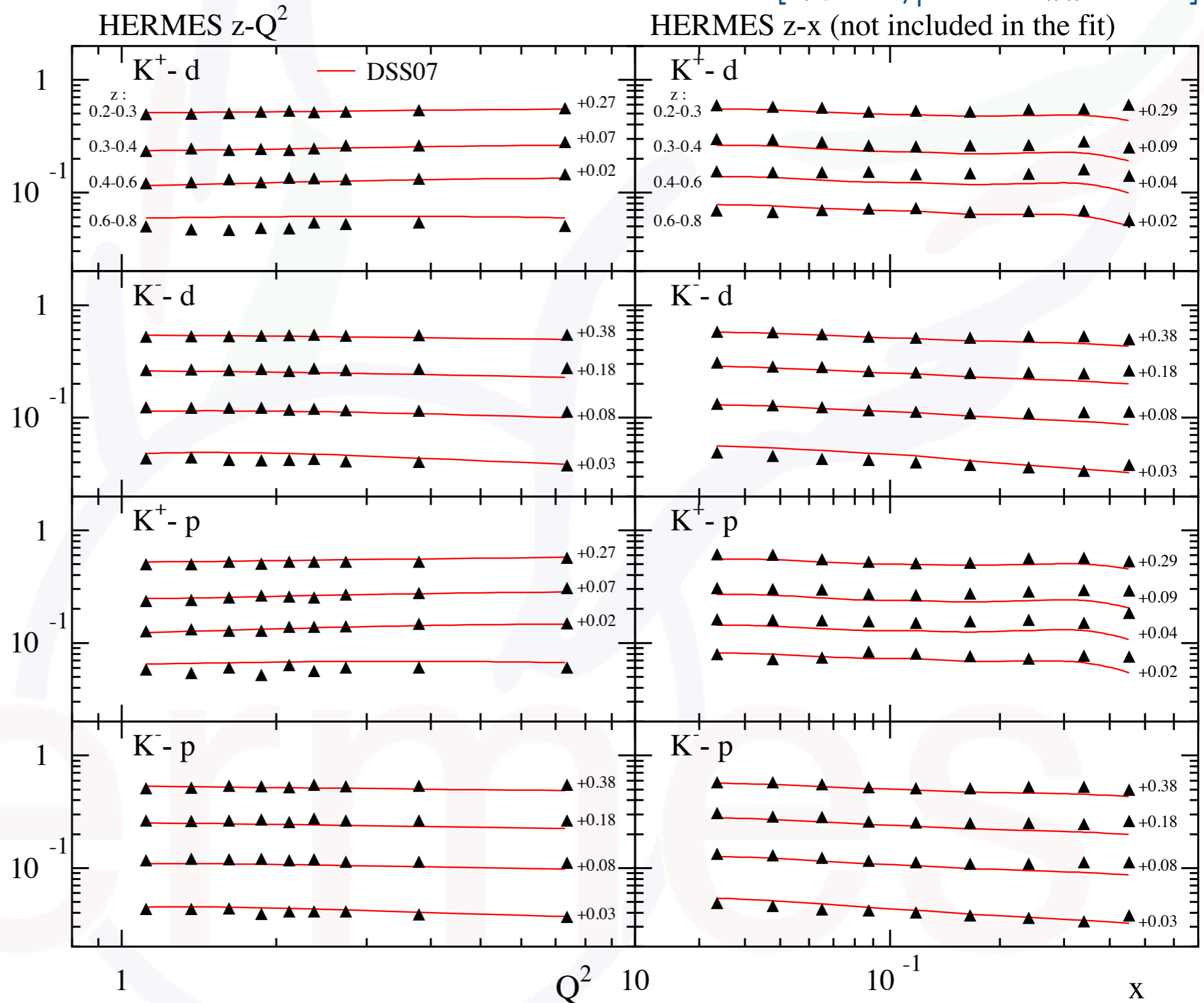


take-away message: integrate your cross sections over the kinematic ranges dictated by the experiment (and do not simply evaluate it at the average kinematics)

# integrating vs. using average kinematics

[R. Sassot, private communication]

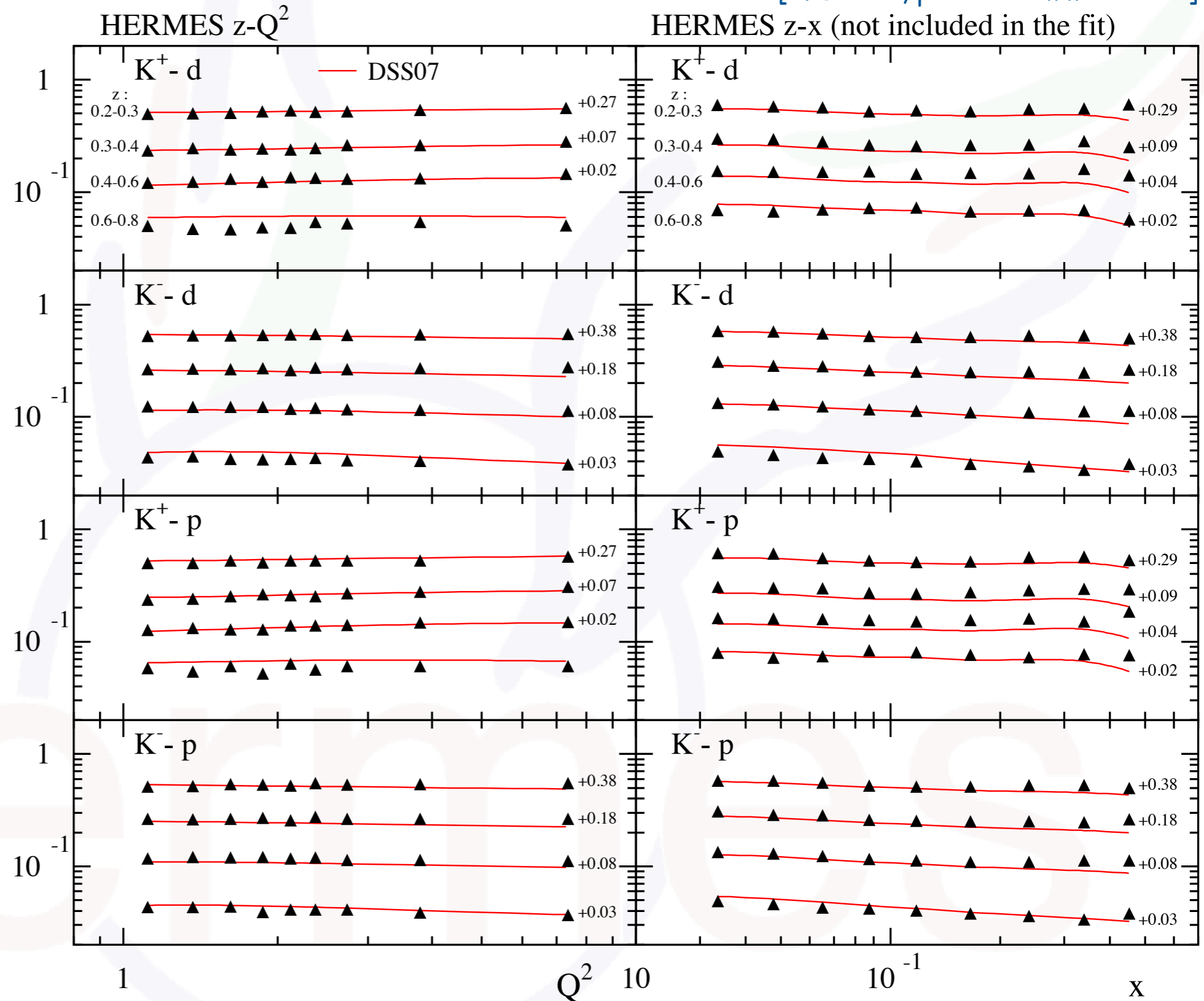
● (by now old)  
DSS07 FF fit to  
 $z$ - $Q^2$  projection



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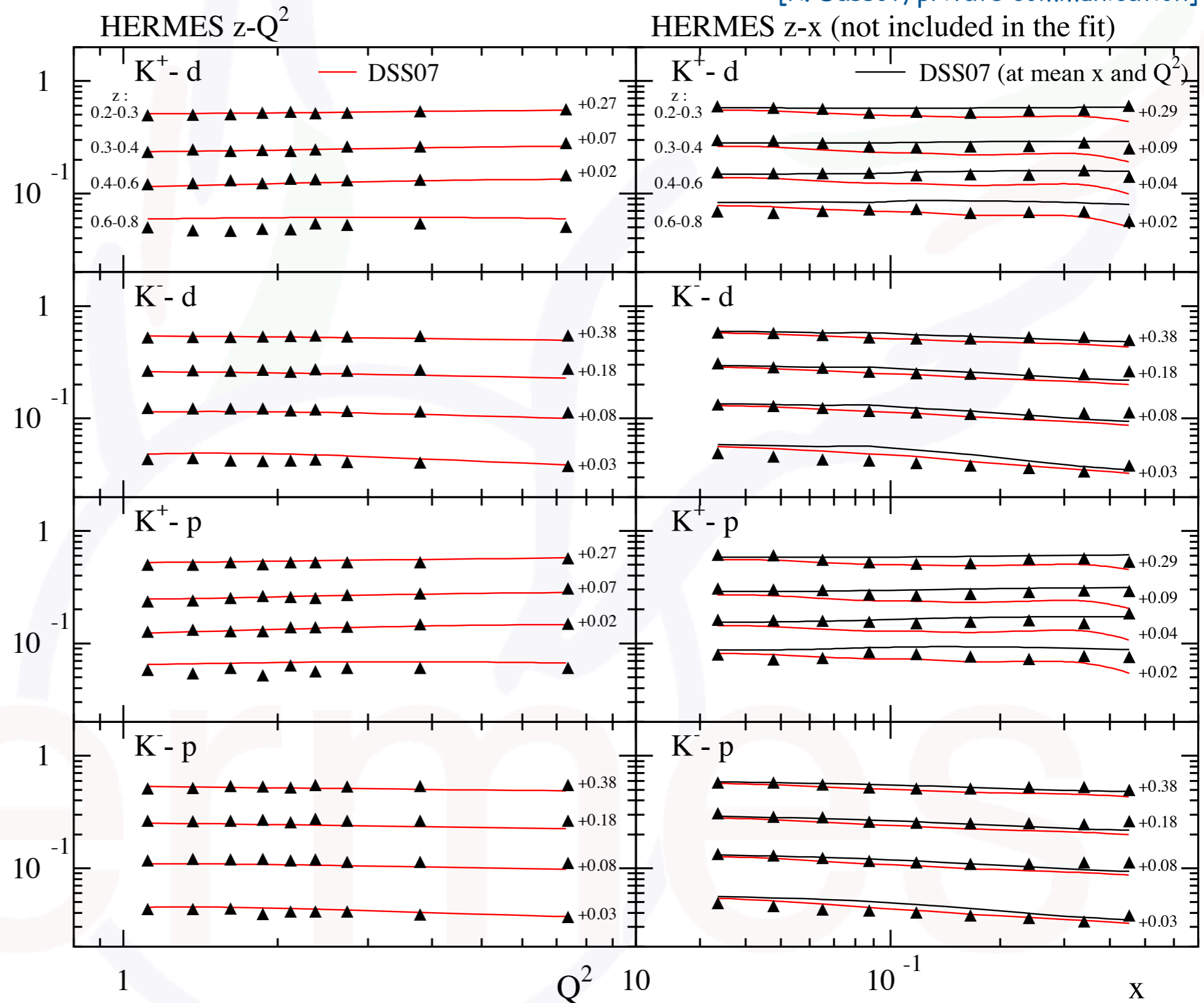
- (by now old) DSS07 FF fit to  $z$ - $Q^2$  projection
- $z$ - $x$  "prediction" reasonable well when using integration over phase-space limits (red lines)



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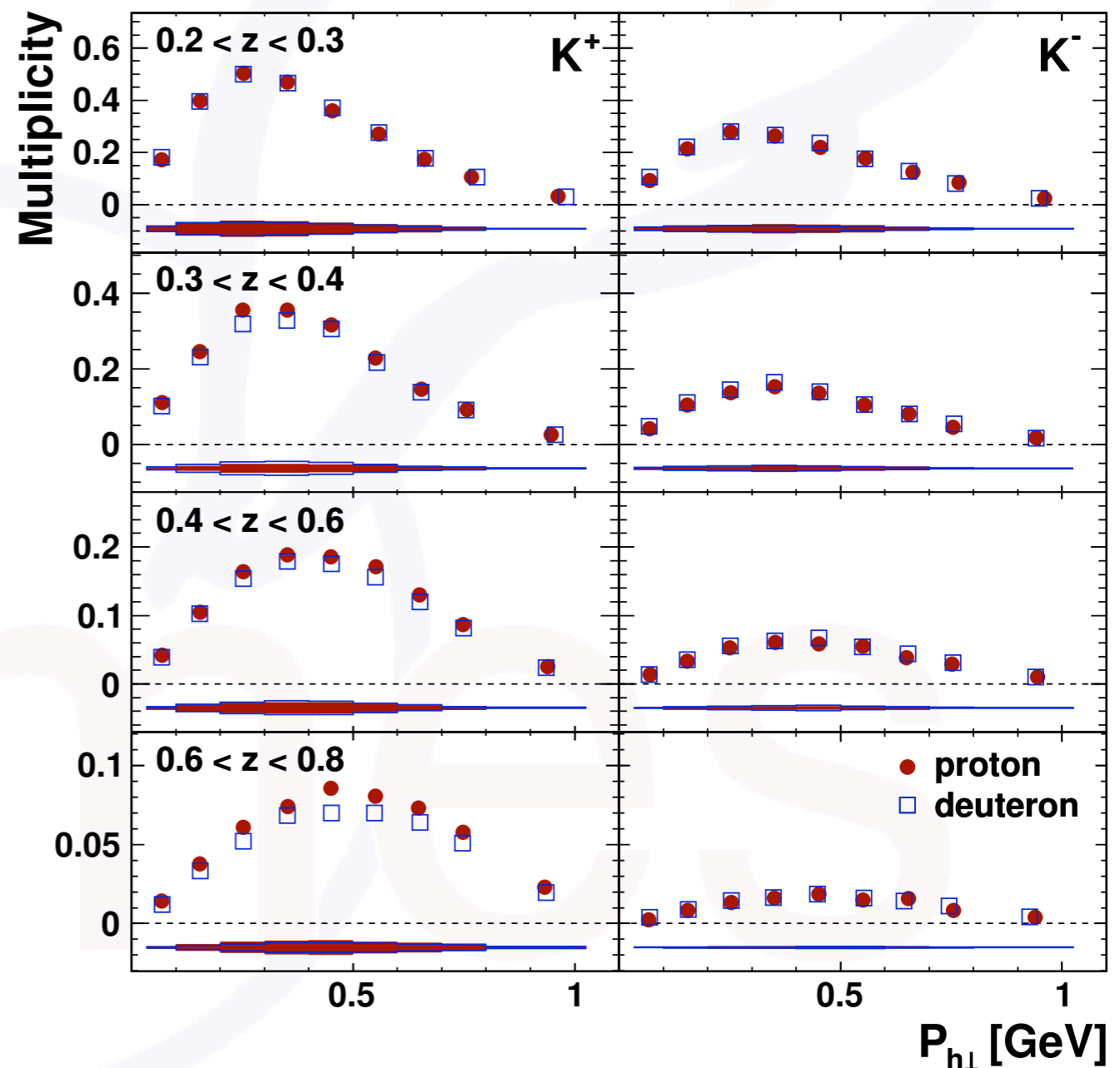
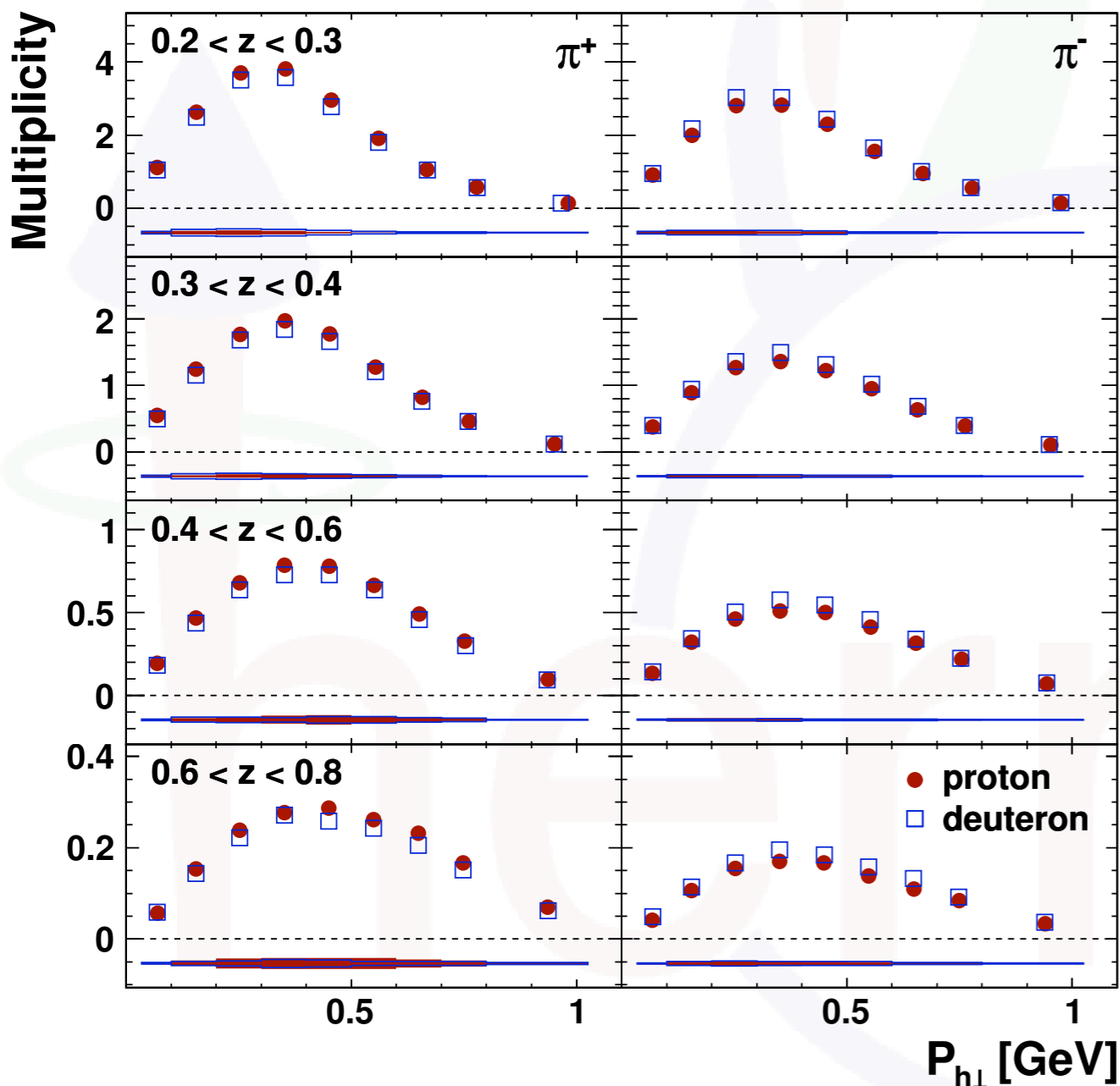
- (by now old) DSS07 FF fit to  $z$ - $Q^2$  projection
- $z$ - $x$  "prediction" reasonable well when using integration over phase-space limits (red lines)
- significant changes when using average kinematics



# transverse momentum dependence

- multi-dimensional analysis allows going beyond collinear factorization
- flavor information on transverse momenta via target variation and hadron ID

[Airapetian et al., PRD 87 (2013) 074029]



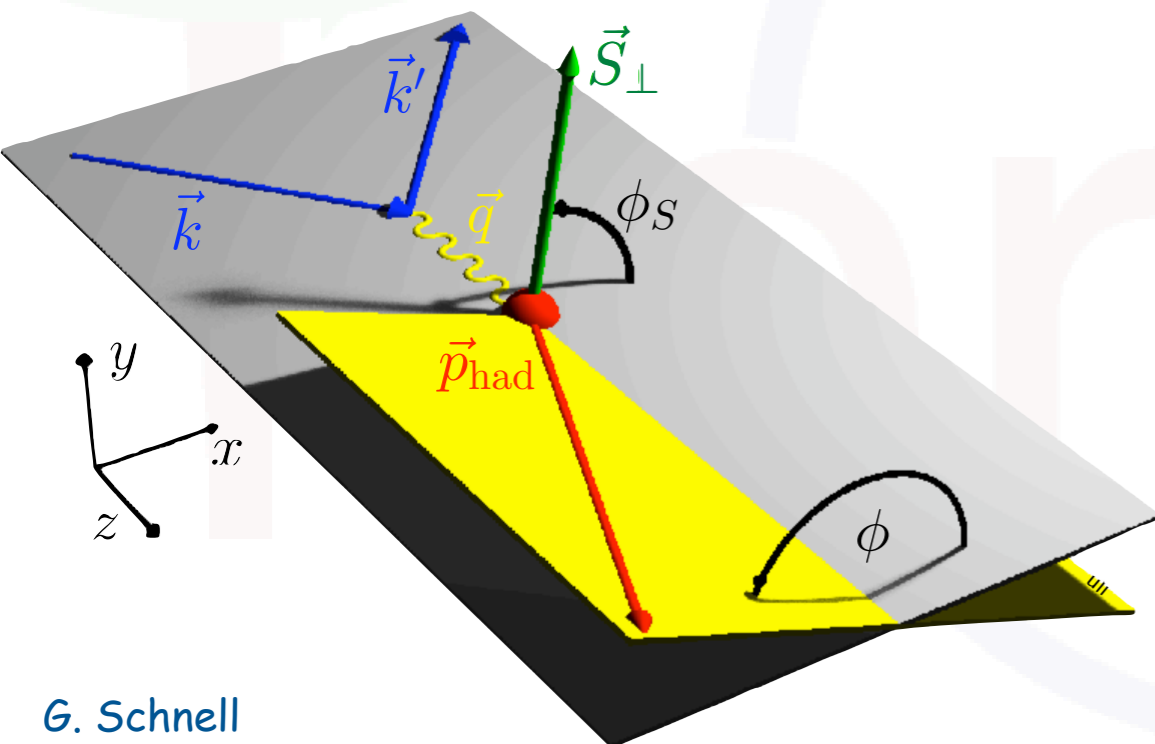


chiral-odd distributions

# transversely polarized quarks?

|   | U              | L        | T                   |
|---|----------------|----------|---------------------|
| U | $f_1$          |          | $h_1^\perp$         |
| L |                | $g_{1L}$ | $h_{1L}^\perp$      |
| T | $f_{1T}^\perp$ | $g_{1T}$ | $h_1, h_{1T}^\perp$ |

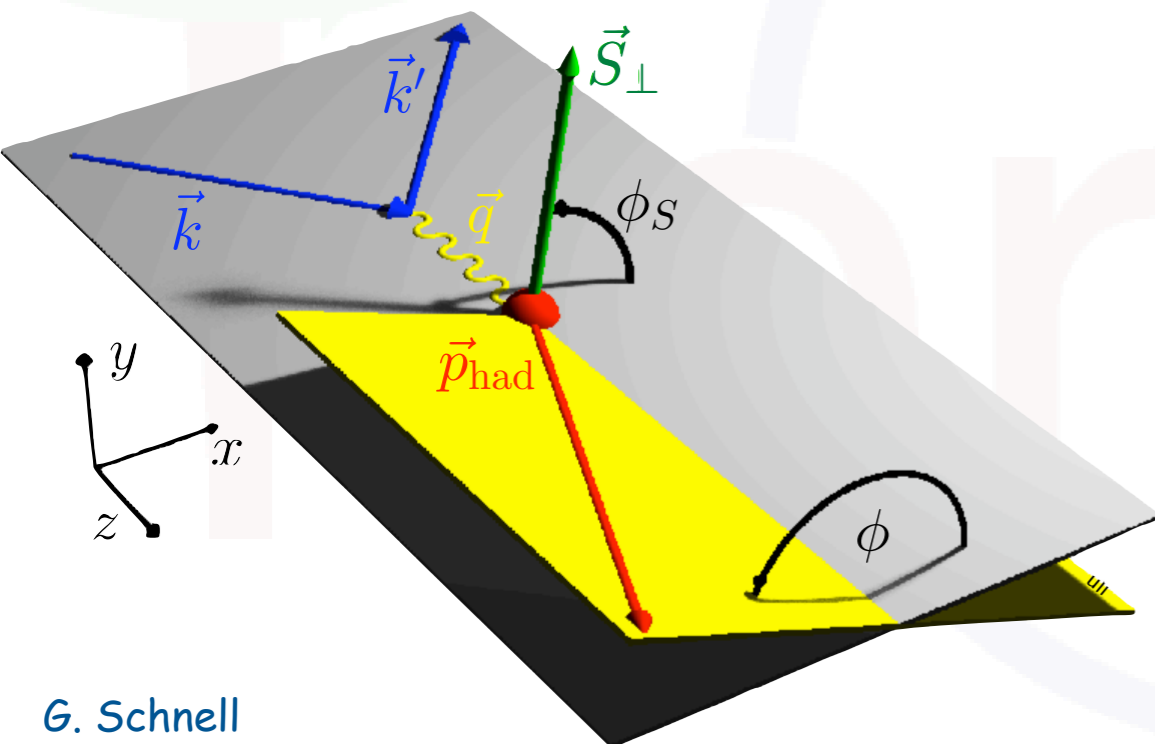
- look at characteristic azimuthal dependence of single-hadron lepto-production cross section



# transversely polarized quarks?

|   | U              | L        | T                   |
|---|----------------|----------|---------------------|
| U | $f_1$          |          | $h_1^\perp$         |
| L |                | $g_{1L}$ | $h_{1L}^\perp$      |
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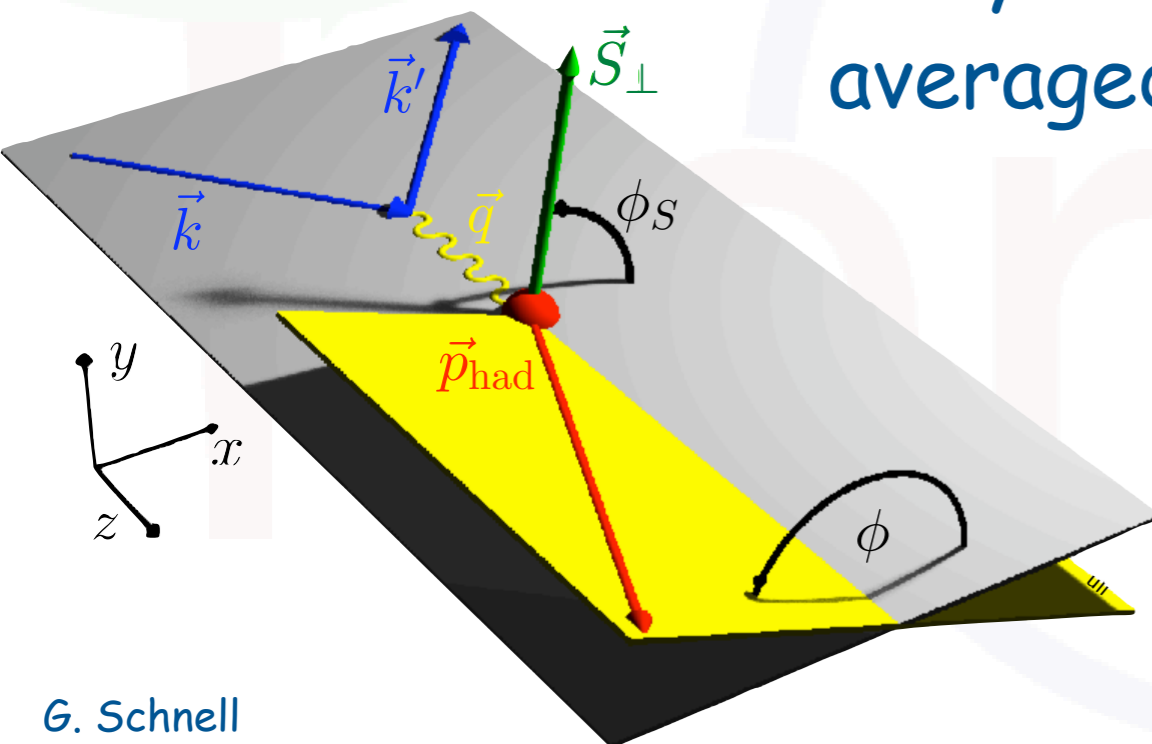
- look at characteristic azimuthal dependence of single-hadron lepto-production cross section
- in practice reverse nucleon-polarization orientation and form spin asymmetries



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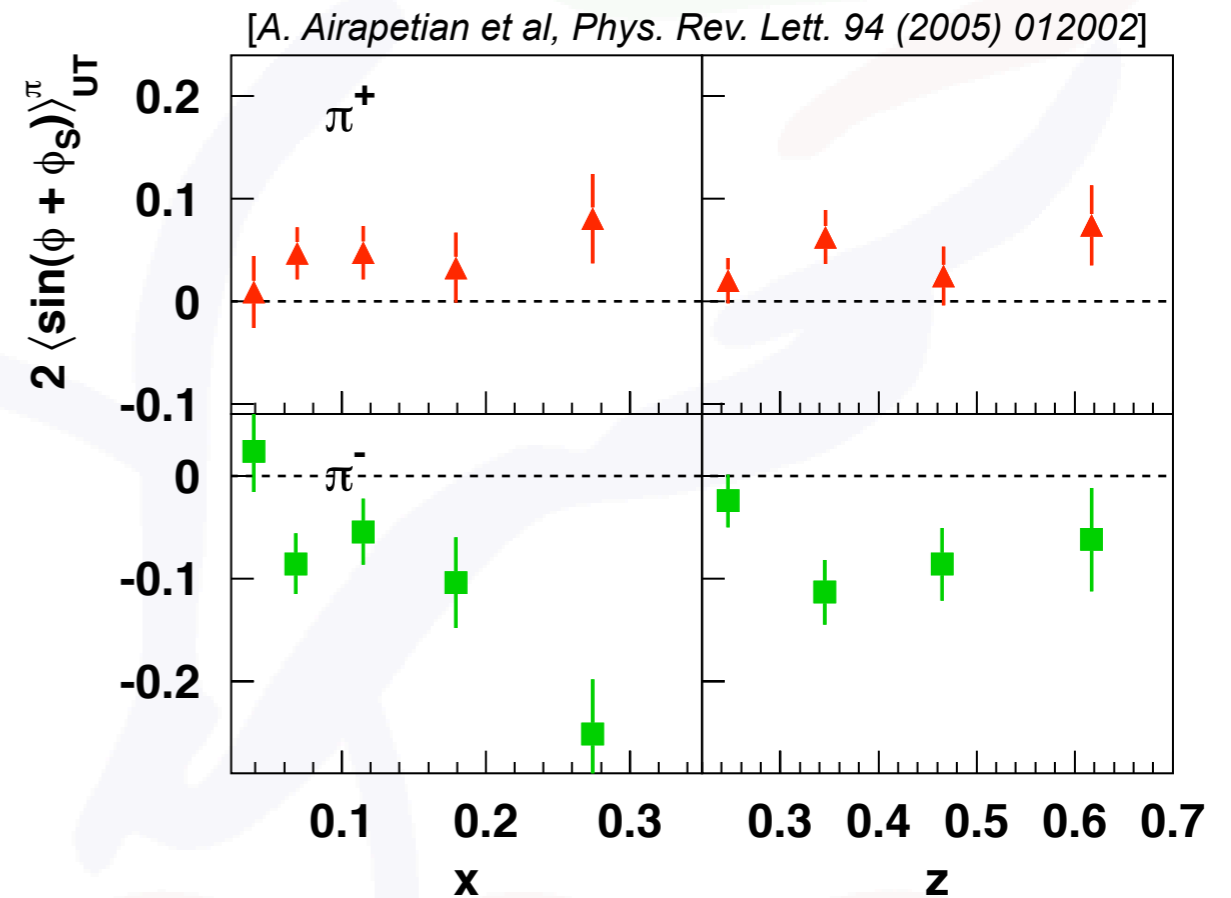
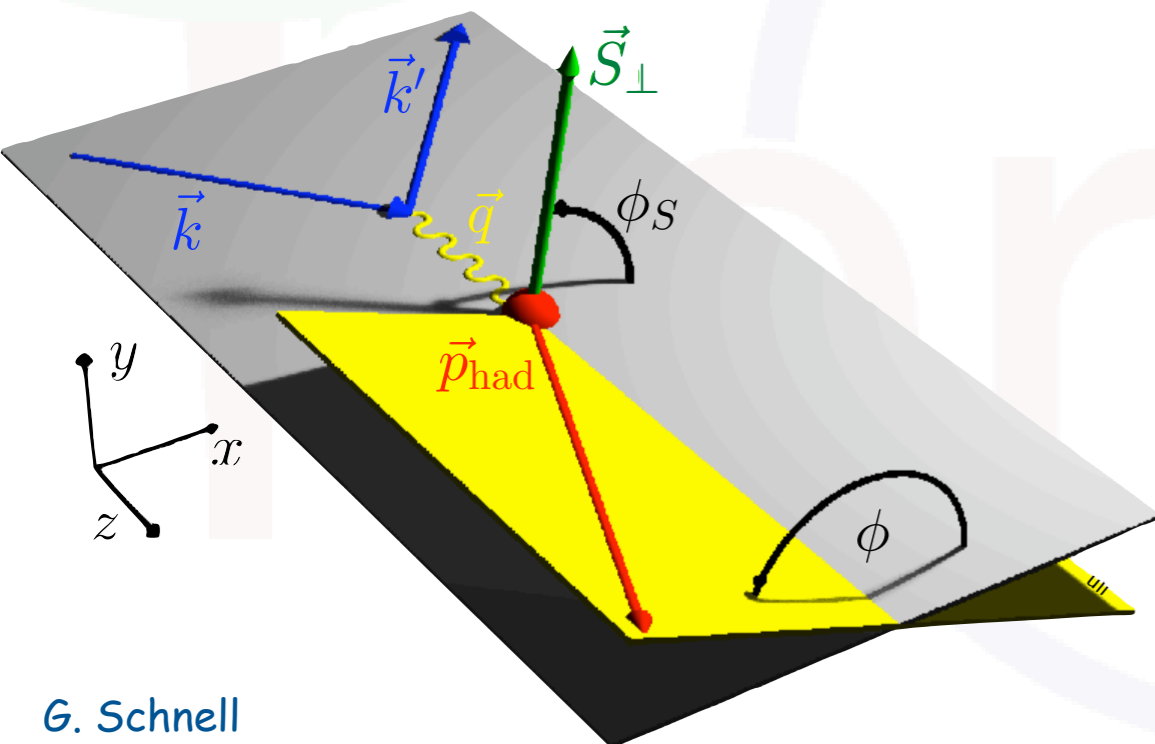
- look at characteristic azimuthal dependence of single-hadron lepto-production cross section
- in practice reverse nucleon-polarization orientation and form spin asymmetries
- many of the systematics of polarization-averaged observables cancel (e.g., luminosity)



# transversely polarized quarks?

|   |                |          |                     |
|---|----------------|----------|---------------------|
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- transverse polarization of quarks leads to large effects!



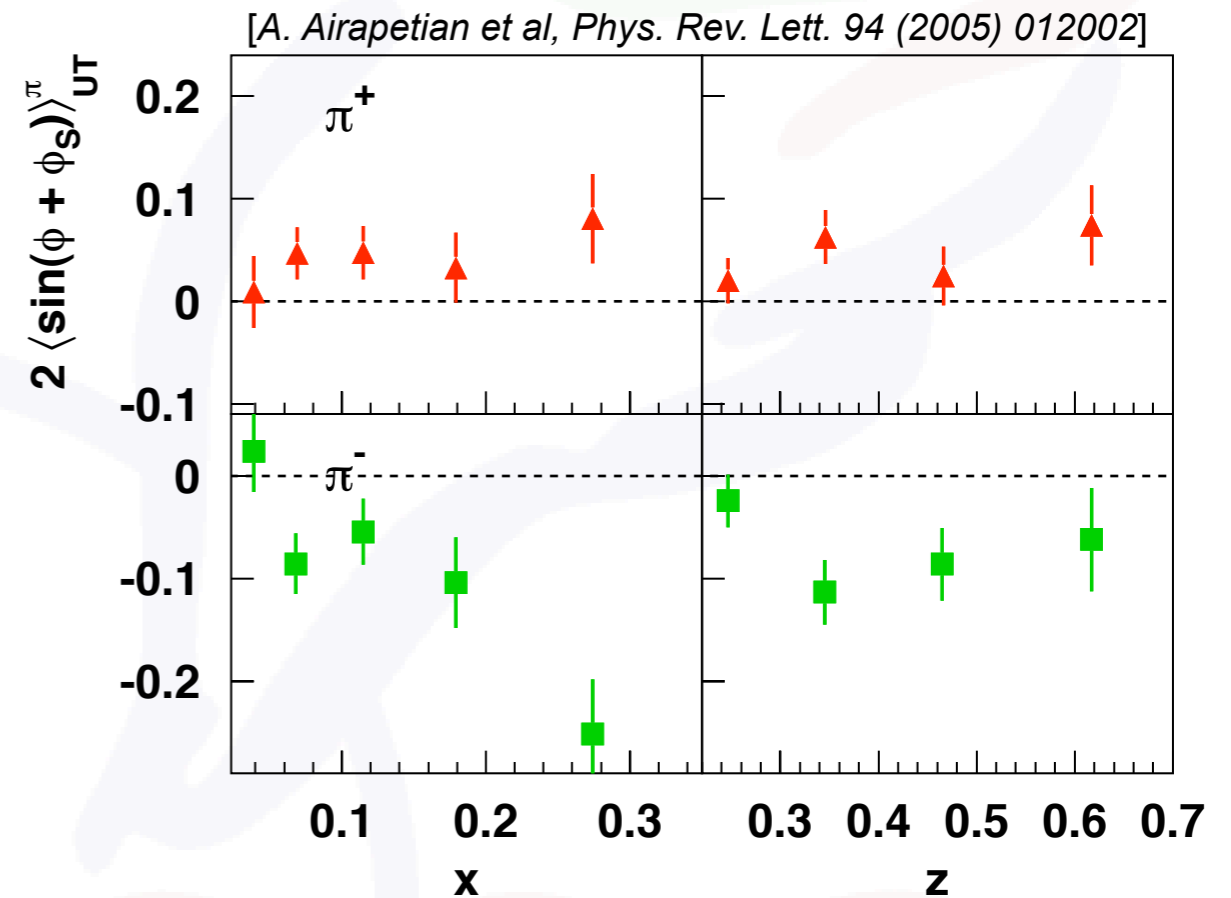
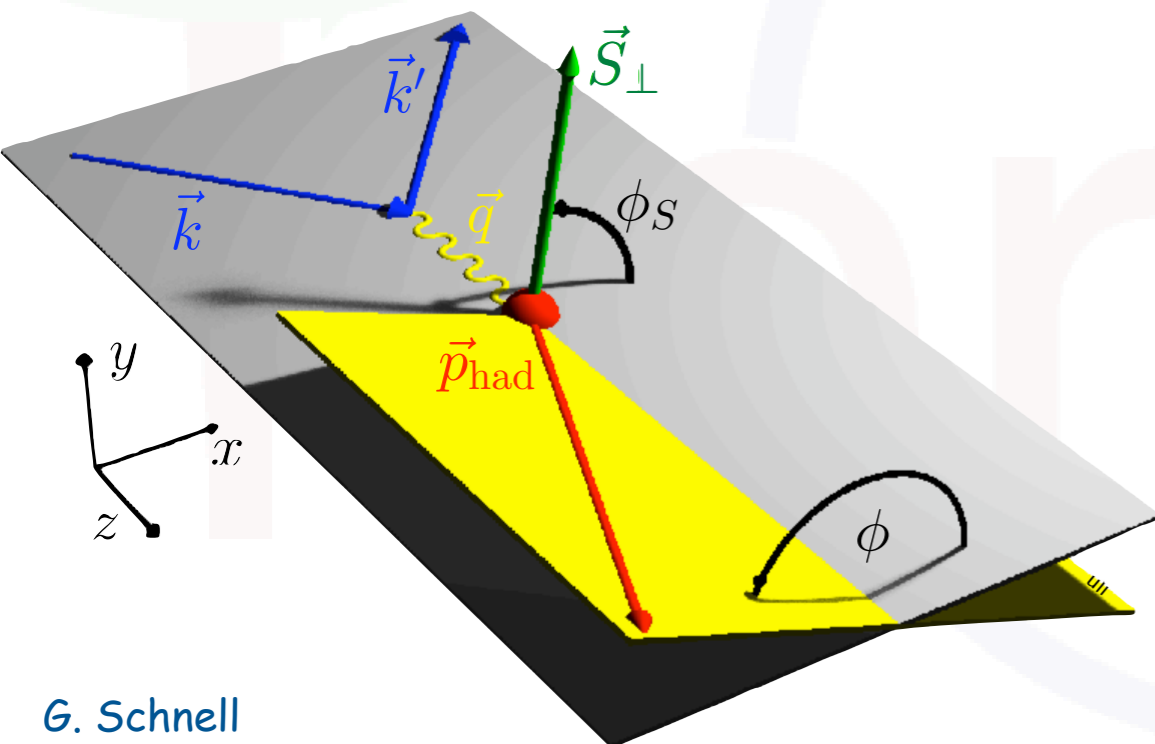
2005: First evidence from HERMES  
SIDIS on proton

Non-zero transversity  
Non-zero Collins function

# transversely polarized quarks?

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- transverse polarization of quarks leads to large effects!
- opposite in sign for charged pions



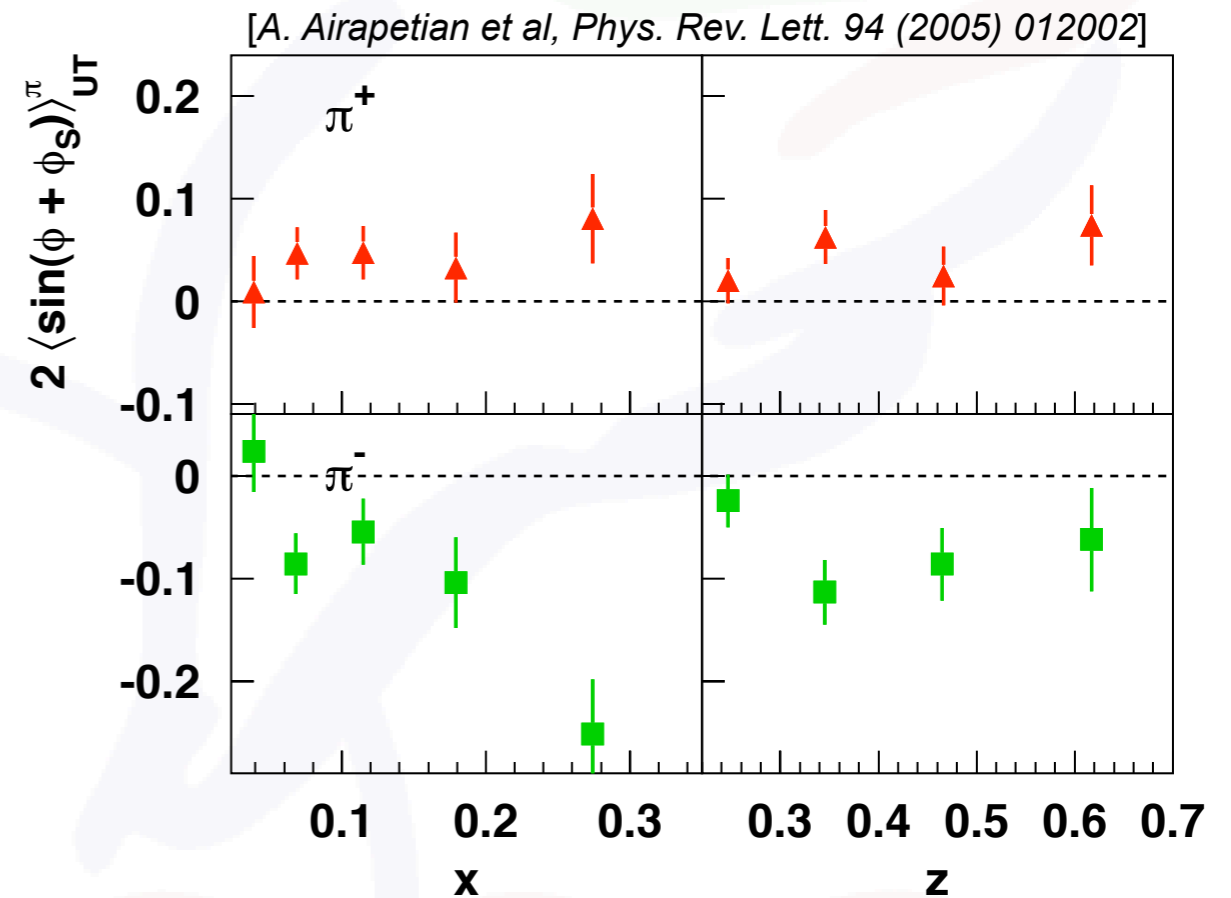
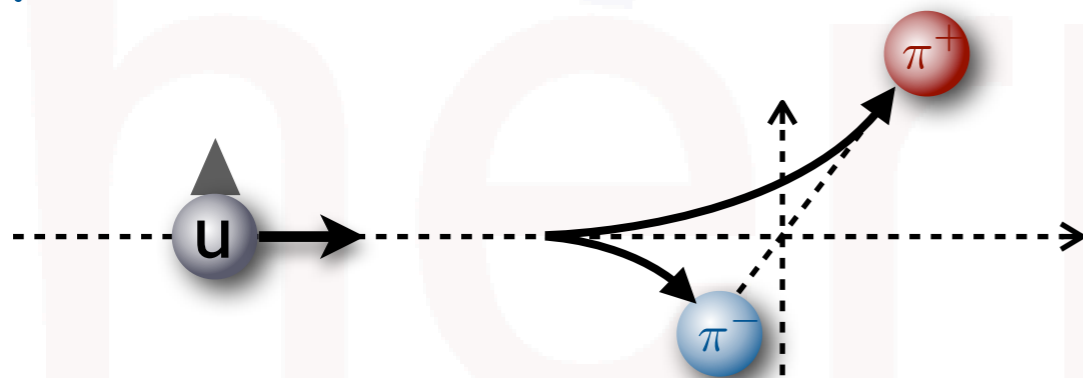
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# transversely polarized quarks?

|   |                |          |                     |
|---|----------------|----------|---------------------|
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- transverse polarization of quarks leads to large effects!
- opposite in sign for charged pions
- disfavored Collins FF large and opposite in sign to favored one



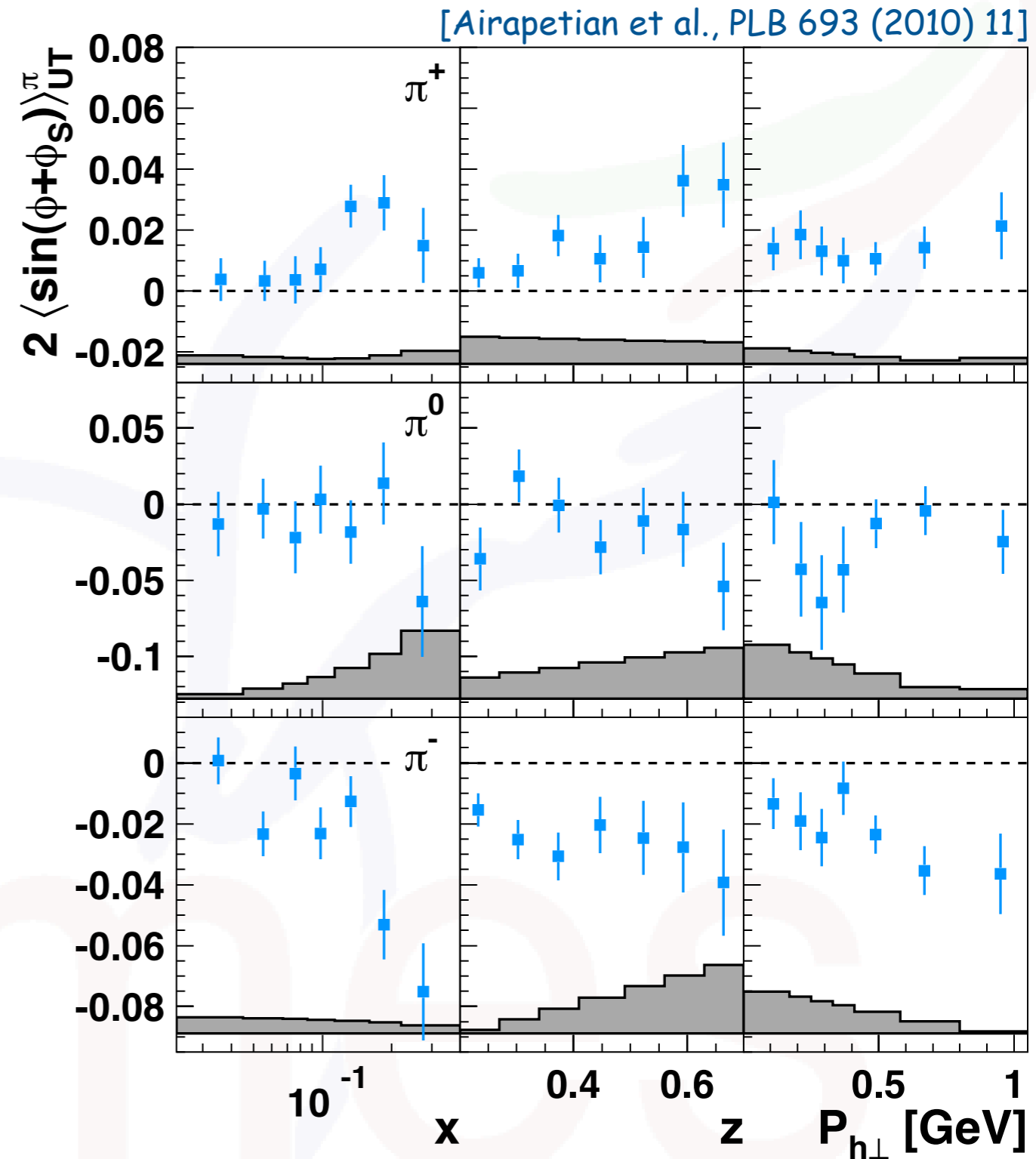
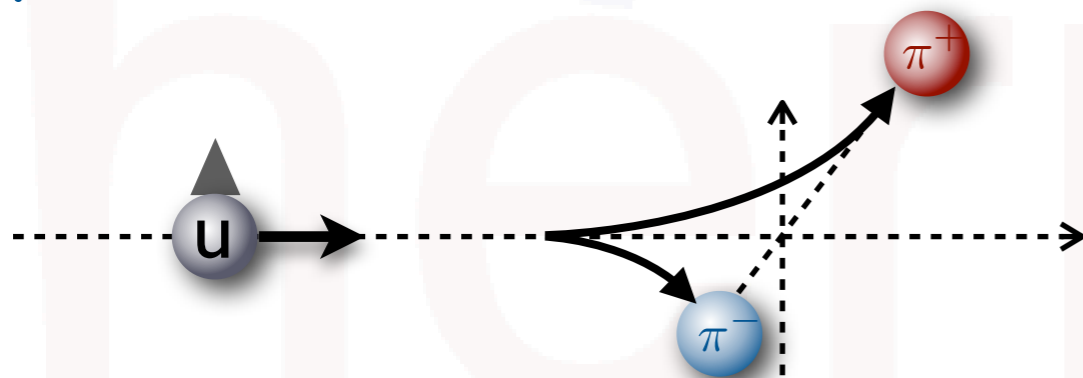
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# transversely polarized quarks?

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- transverse polarization of quarks leads to large effects!
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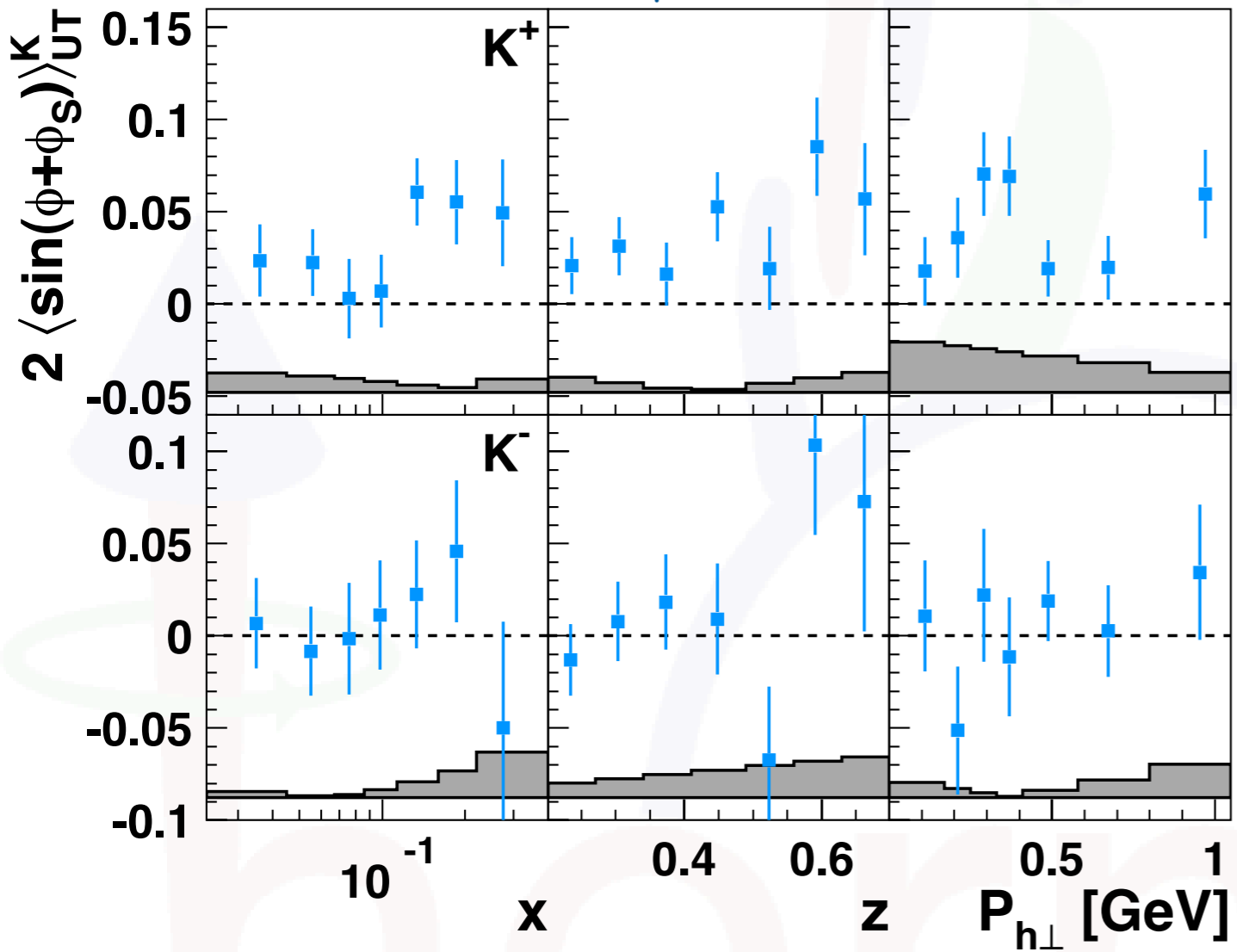




# Collins effect for kaons and (anti) protons

|   |                |          |                     |
|---|----------------|----------|---------------------|
|   | U              | L        | T                   |
| U | $f_1$          |          | $h_1^\perp$         |
| L |                | $g_{1L}$ | $h_{1L}^\perp$      |
| T | $f_{1T}^\perp$ | $g_{1T}$ | $h_1, h_{1T}^\perp$ |

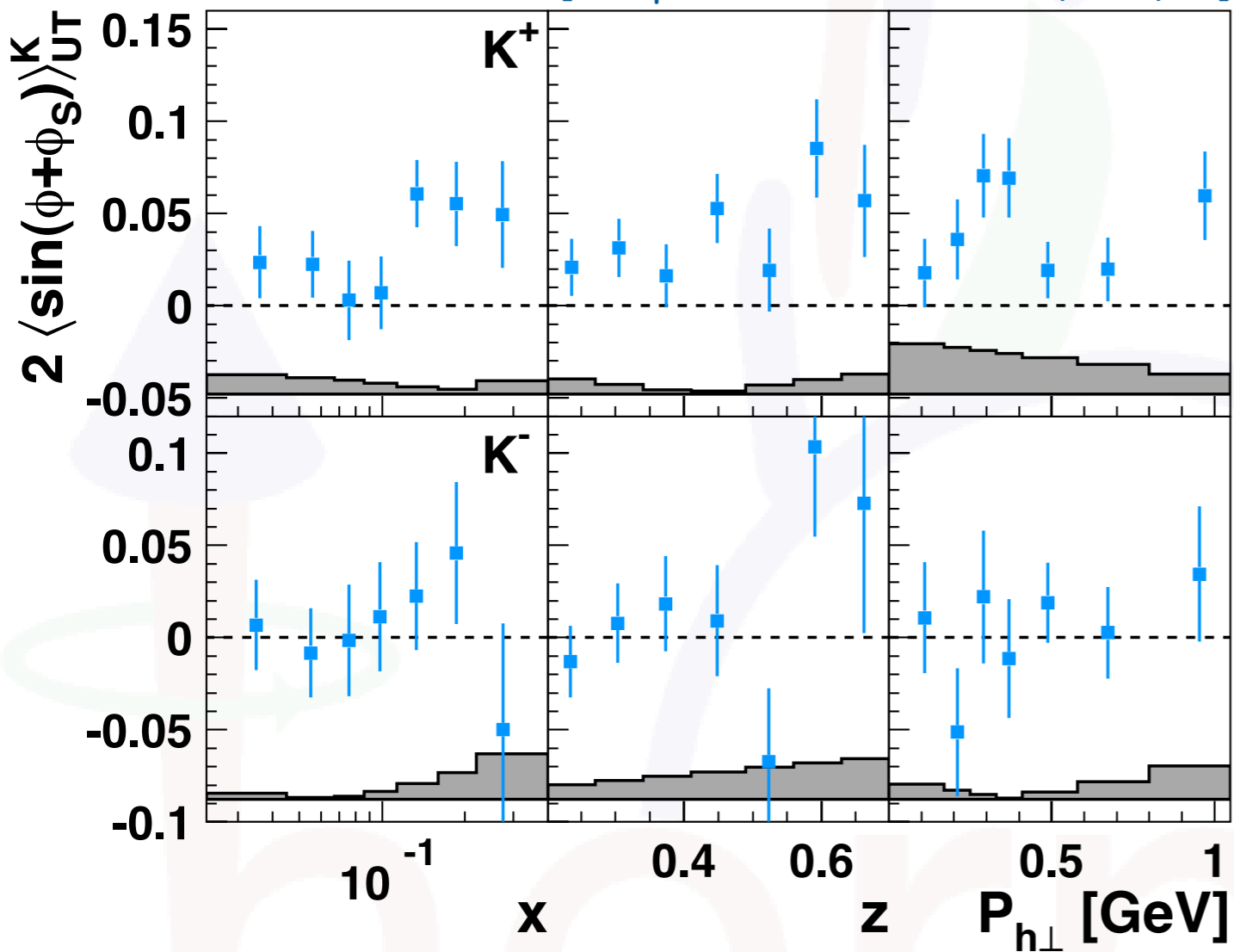
[Airapetian et al., PLB 693 (2010) 11]



# Collins effect for kaons and (anti) protons

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|   | U              | L        | T                   |
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[Airapetian et al., PLB 693 (2010) 11]

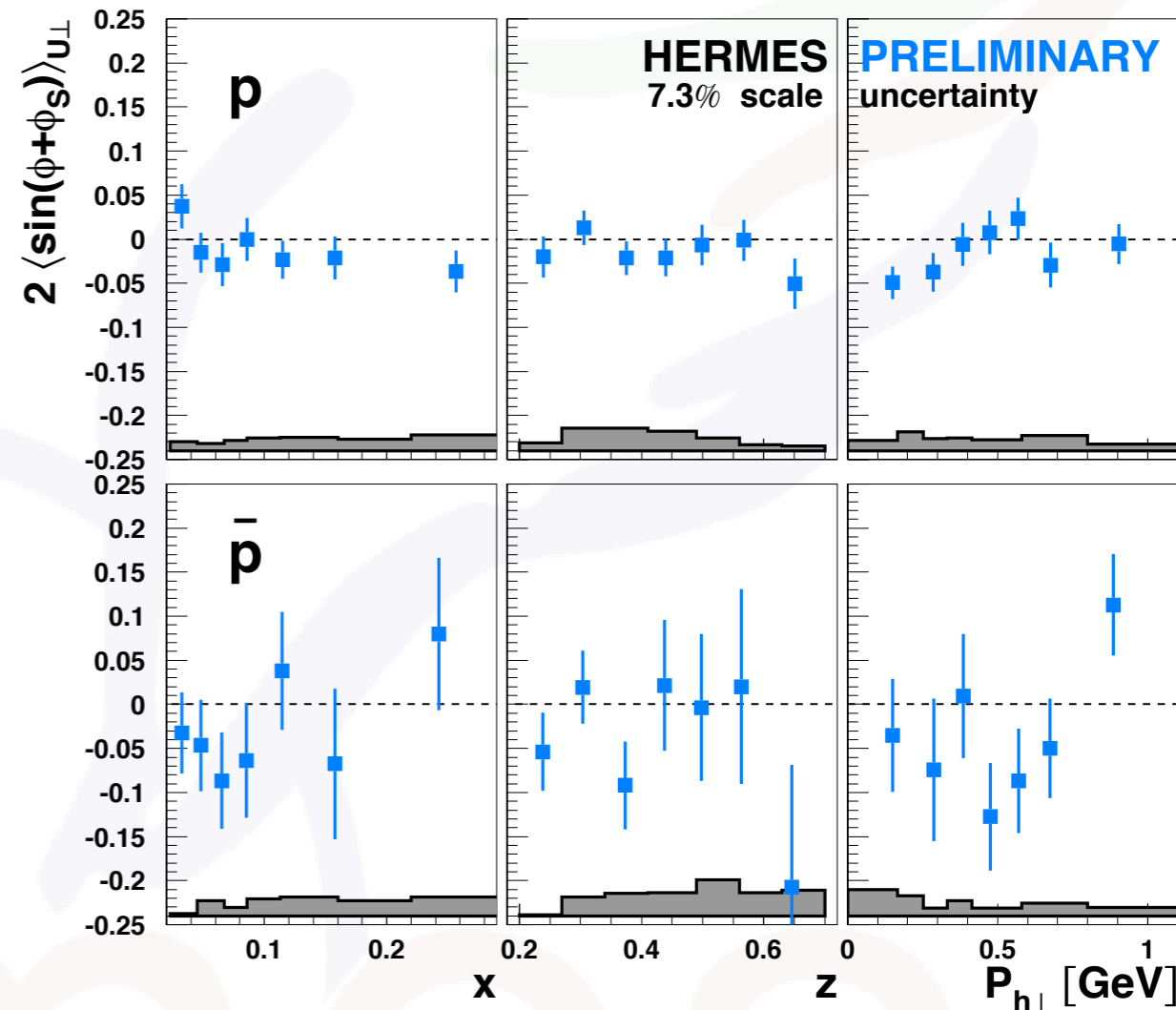
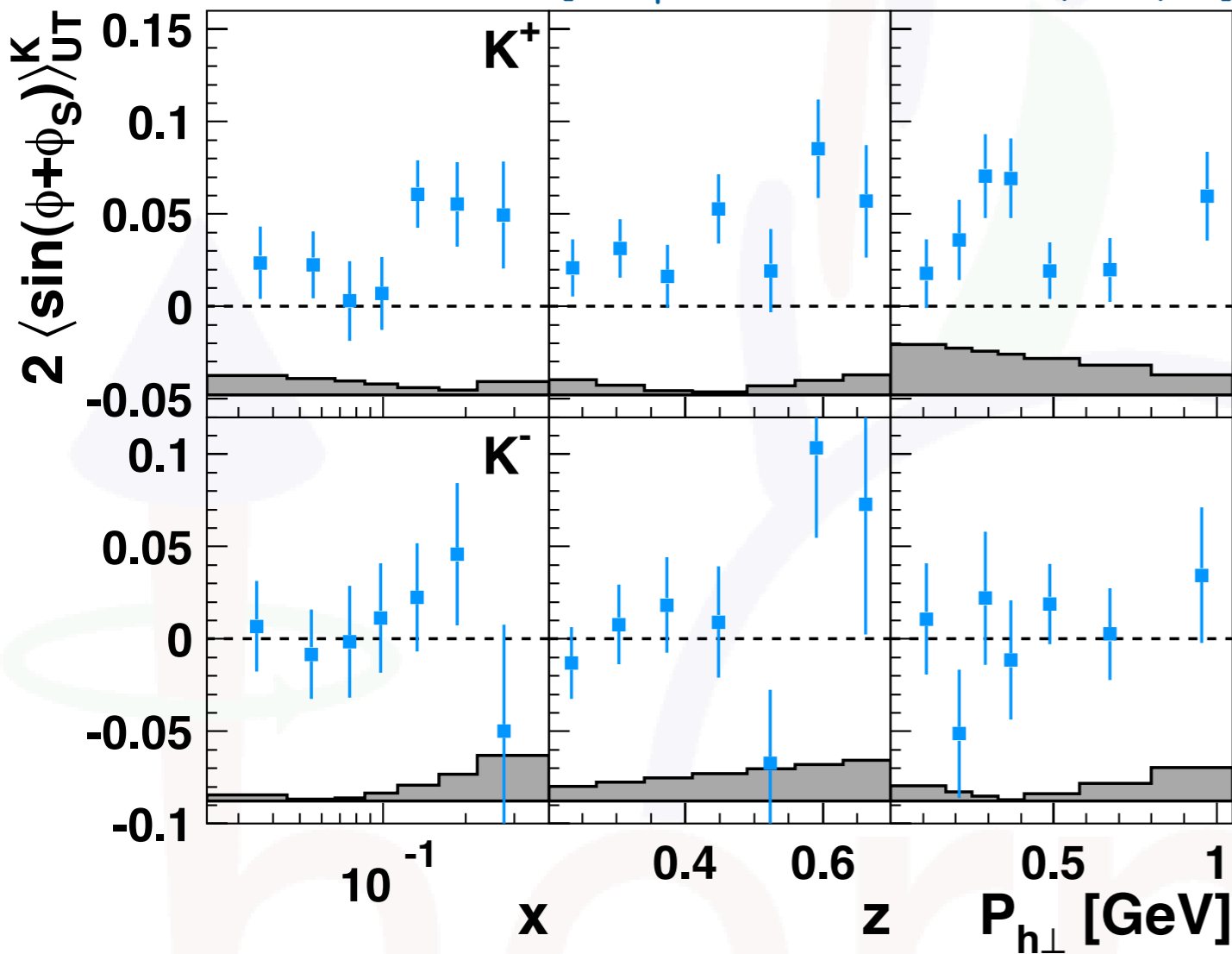


- positive Collins SSA amplitude for positive kaons

# Collins effect for kaons and (anti) protons

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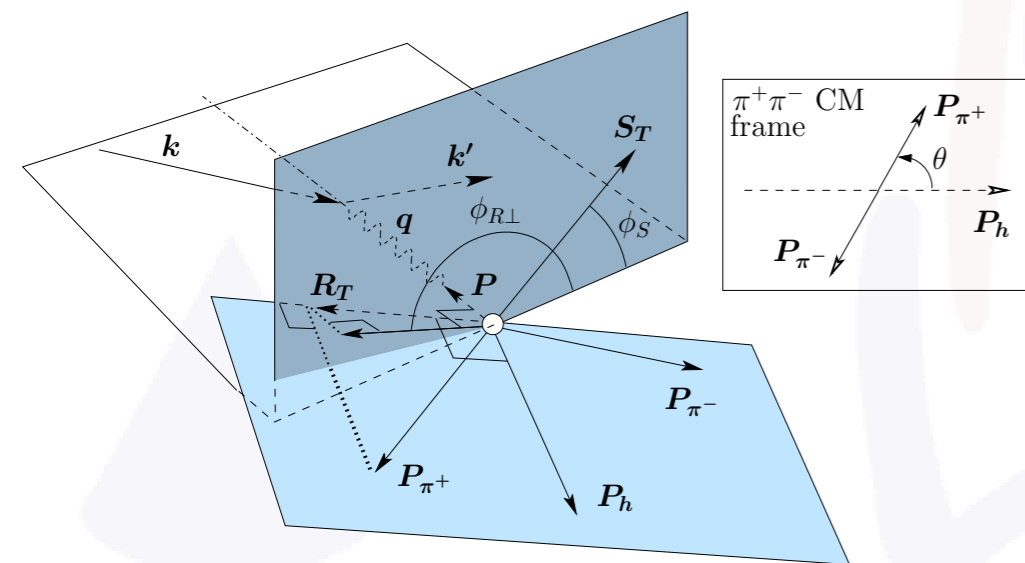
[Airapetian et al., PLB 693 (2010) 11]



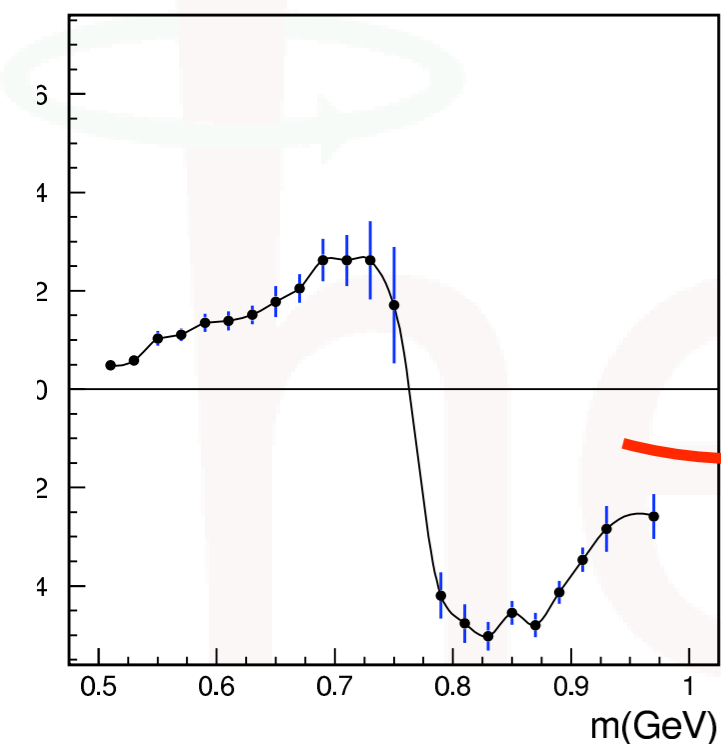
- positive Collins SSA amplitude for positive kaons
  - consistent with zero for negative kaons and (anti)protons
- ➔ vanishing sea-quark transversity and baryon Collins effect?

# Transversity through 2-hadron fragmentation

|   | U              | L        | T                   |
|---|----------------|----------|---------------------|
| U | $f_1$          |          | $h_1^\perp$         |
| L |                | $g_{1L}$ | $h_{1L}^\perp$      |
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$$A_{UT} \sim \sin(\phi_{R\perp} + \phi_S) \sin\theta h_1 H_1^{\triangleleft}$$



Jaffe et al. [hep-ph/9709322]:

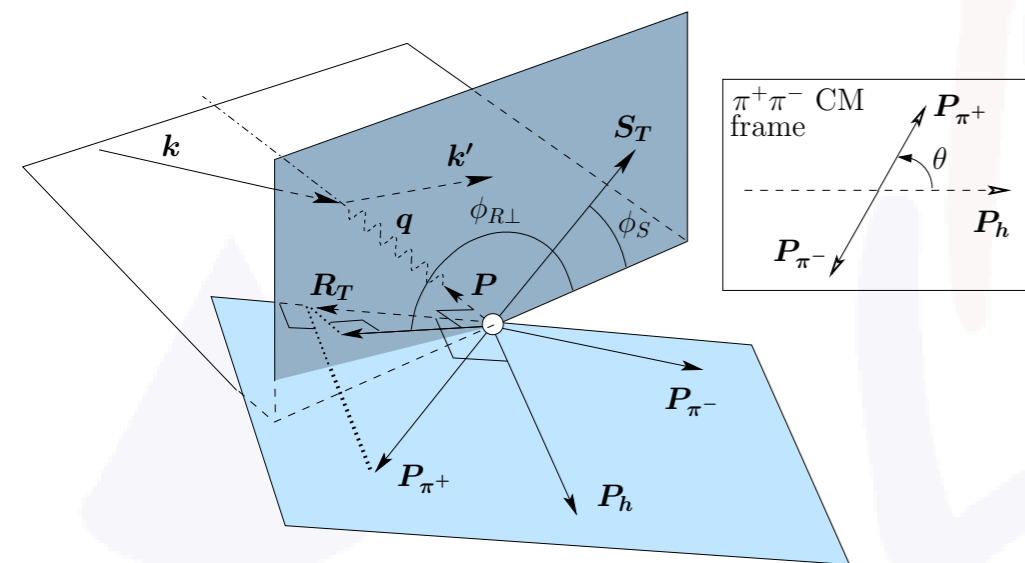
$$H_1^{\triangleleft, sp}(z, M_{\pi\pi}^2) = \frac{\sin\delta_0 \sin\delta_1 \sin(\delta_0 - \delta_1) H_1^{\triangleleft, sp'}(z)}{\delta_0 (\delta_1) \rightarrow \text{S(P)-wave phase shifts}}$$

$$= \mathcal{P}(M_{\pi\pi}^2) H_1^{\triangleleft, sp'}(z)$$

$\Rightarrow A_{UT}$  might depend strongly on  $M_{\pi\pi}$

# Transversity through 2-hadron fragmentation

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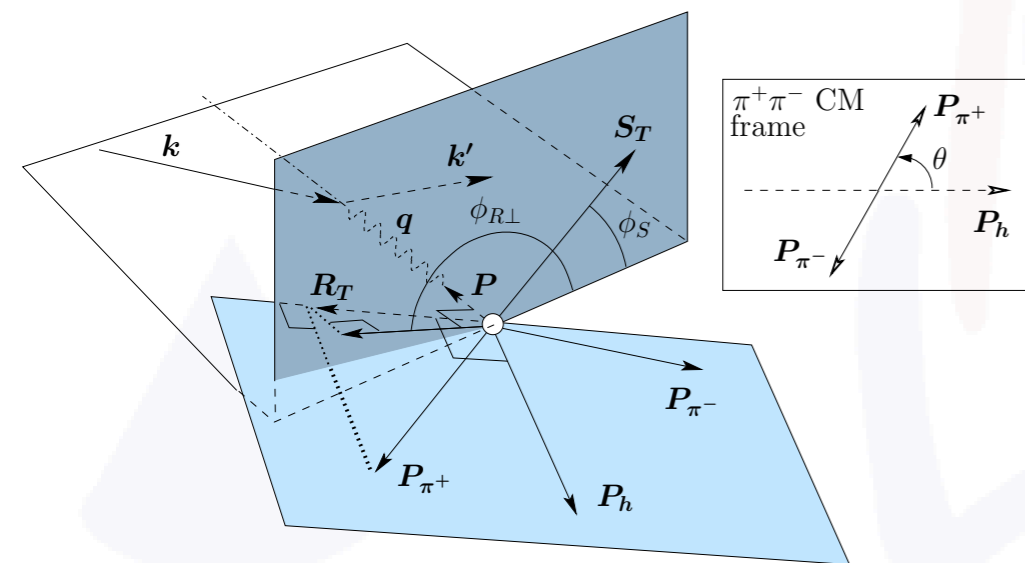


$$A_{UT} \sim \sin(\phi_{R\perp} + \phi_S) \sin\theta h_1 H_1^{\triangleleft}$$

- not only strong invariant-mass dependence, experimental challenges also because of
- transverse-momentum dependence
- theta dependence

# Transversity through 2-hadron fragmentation

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| U | $f_1$          |          | $h_1^\perp$         |
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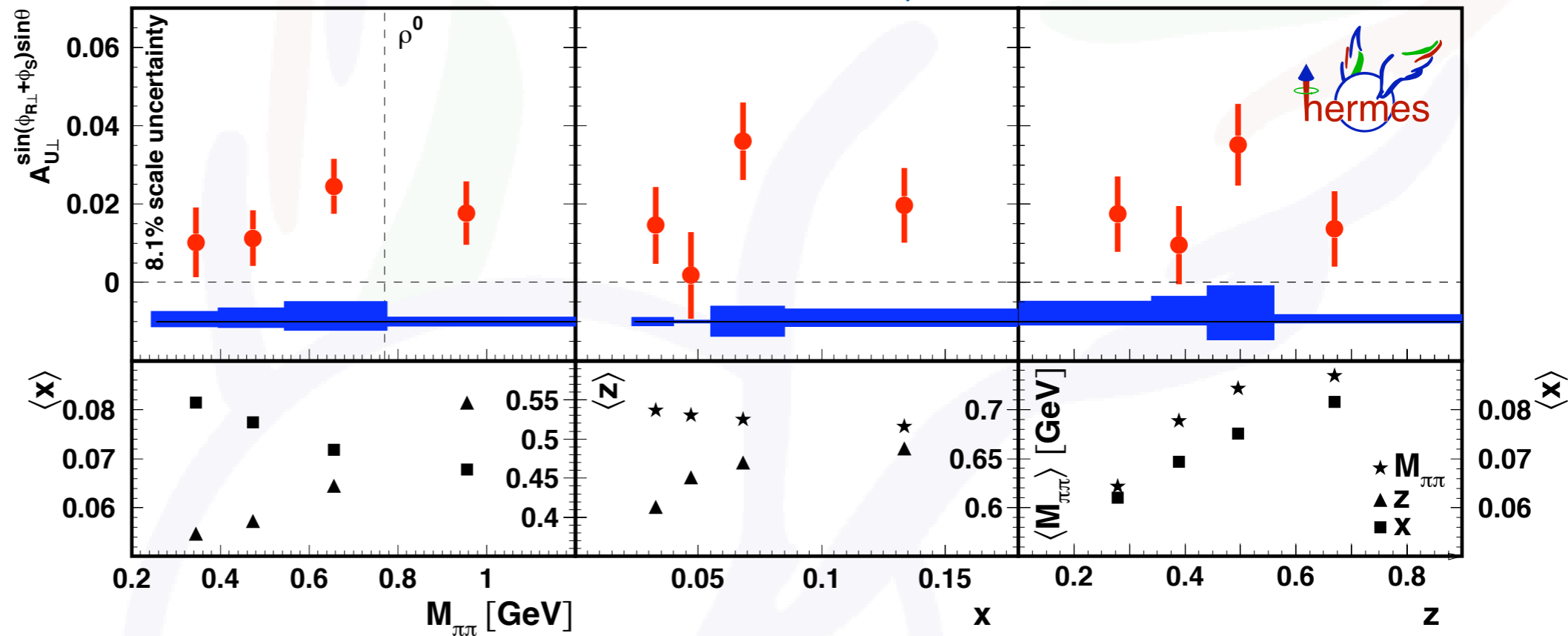
$$A_{UT} \sim \sin(\phi_{R\perp} + \phi_S) \sin\theta h_1 H_1^{\triangleleft}$$

- not only strong invariant-mass dependence, experimental challenges also because of
- transverse-momentum dependence
- theta dependence
- 9 vs. 6 (for single hadrons) dependences, too many to analyze simultaneously (at least with presently available data)

# Transversity through 2-hadron fragmentation

|   |                |          |                     |
|---|----------------|----------|---------------------|
|   | U              | L        | T                   |
| U | $f_1$          |          | $h_1^\perp$         |
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[A. Airapetian et al., JHEP 06 (2008) 017]



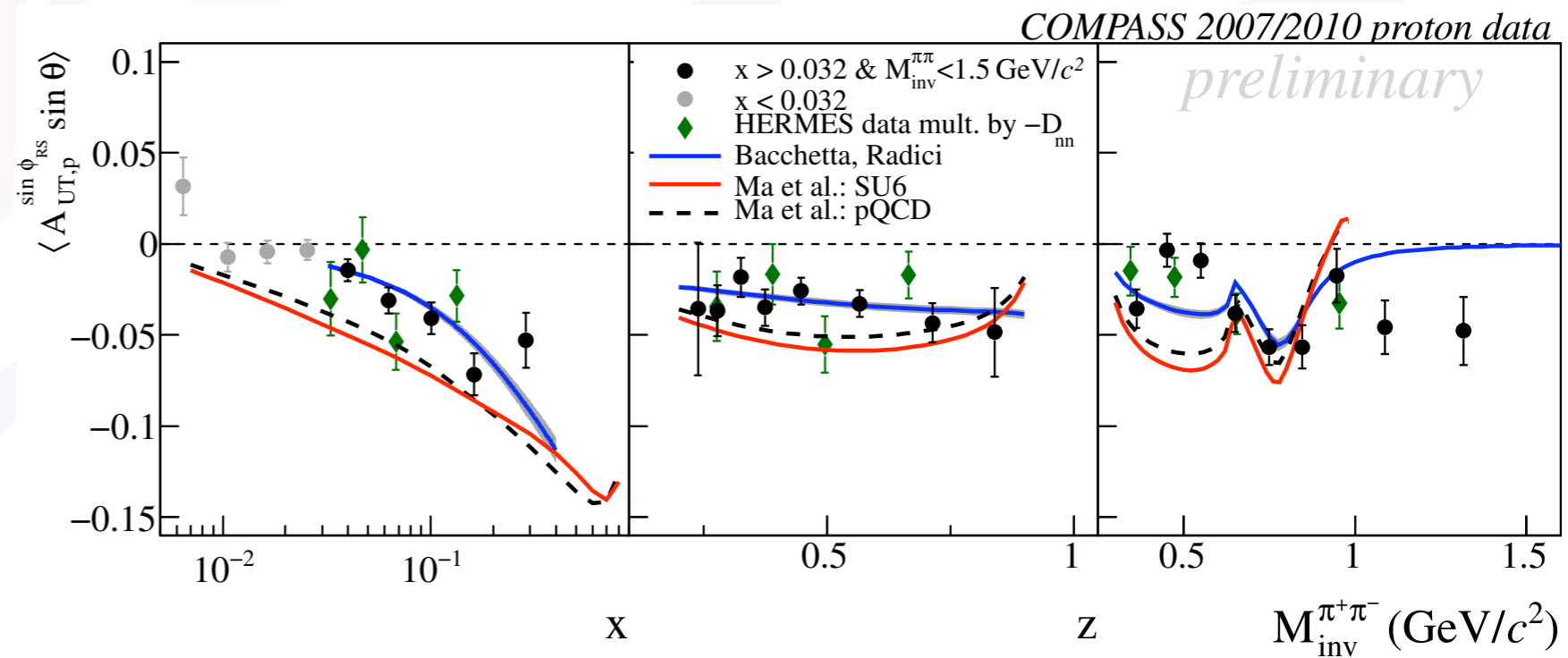
- systematics include
  - incomplete integration over transverse momentum (negligible)
  - contribution from higher partial waves in (unpolarized) denominator
  - integration over other variables, e.g.,  $A(\langle \text{kin.} \rangle) \neq \langle A(\text{kin.}) \rangle$

# Transversity through 2-hadron fragmentation

|   |                |          |                     |
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|   | U              | L        | T                   |
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- HERMES, COMPASS: for comparison scaled HERMES data by depolarization factor and changed sign

[A. Airapetian et al., JHEP 06 (2008) 017]  
 COMPASS 2007: [C. Adolph et al., Phys. Lett. B713 (2012) 10]  
 COMPASS 2010: [C. Braun et al., Nuovo Cimento C 035 (2012) 02]



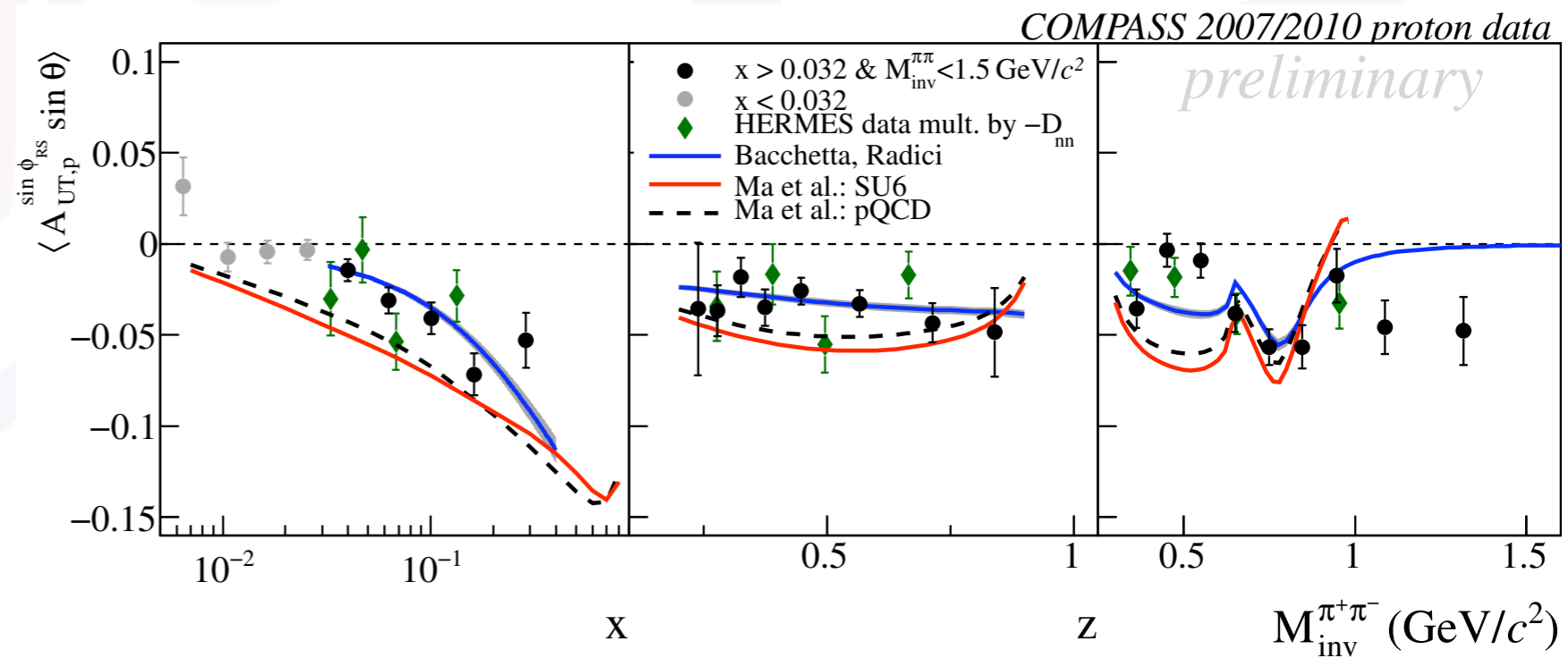


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- $^2\text{H}$  results consistent with zero

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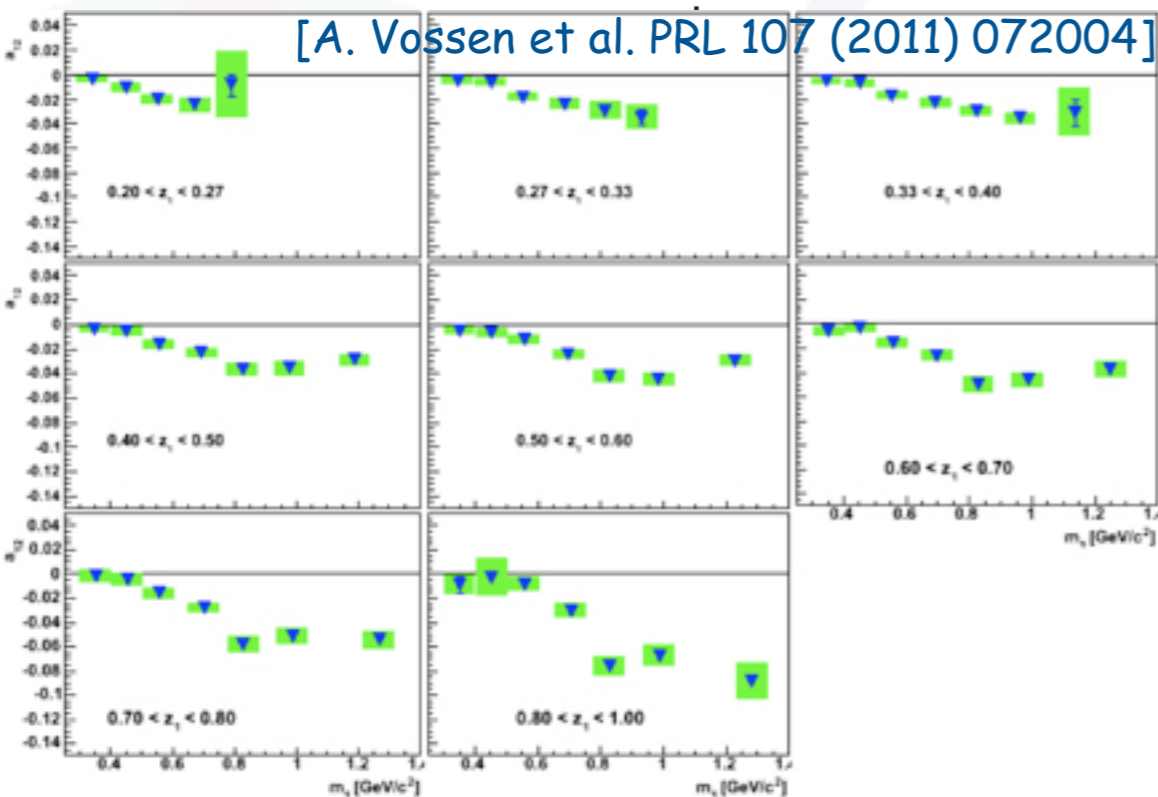
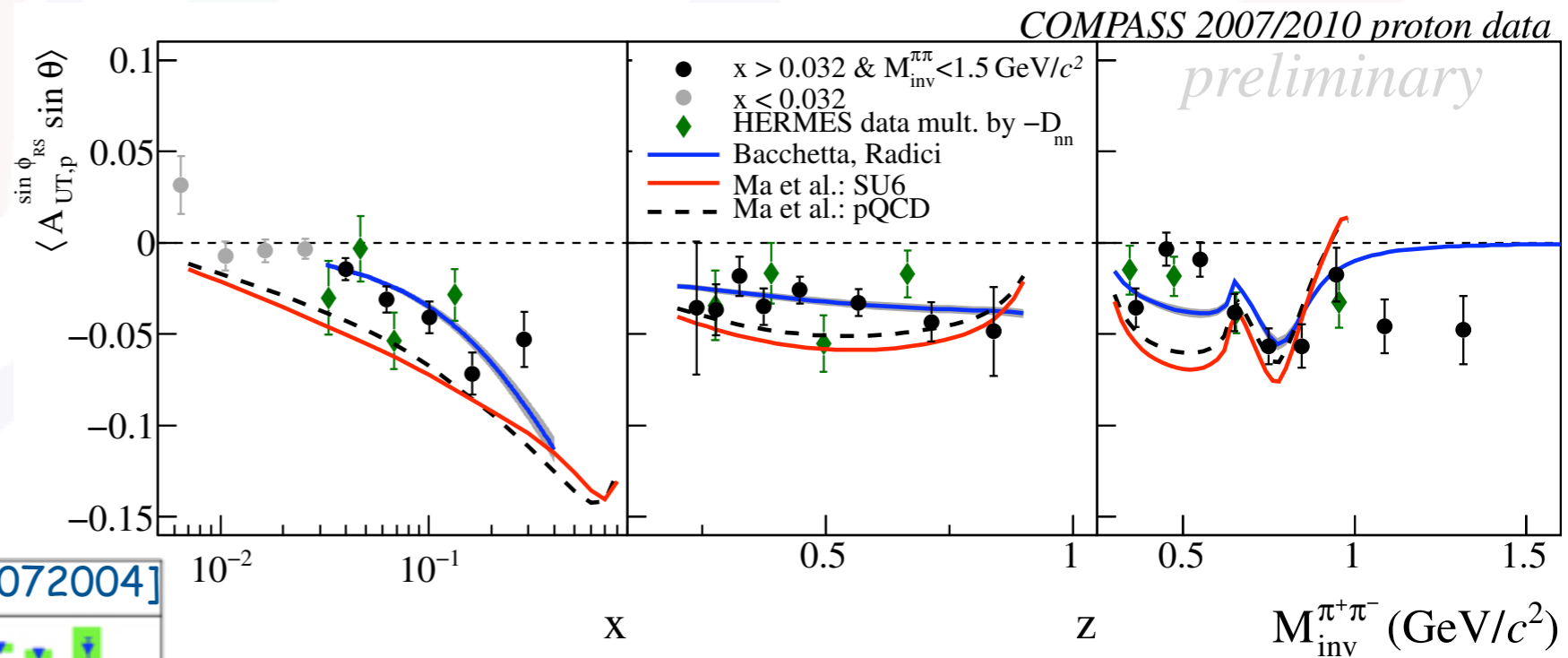


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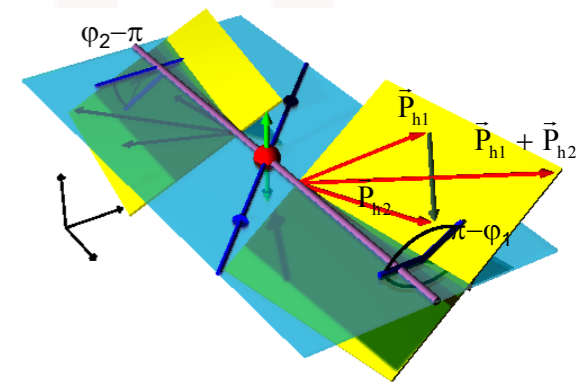
|   |                |          |                     |
|---|----------------|----------|---------------------|
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- data from  $e^+e^-$  by BELLE

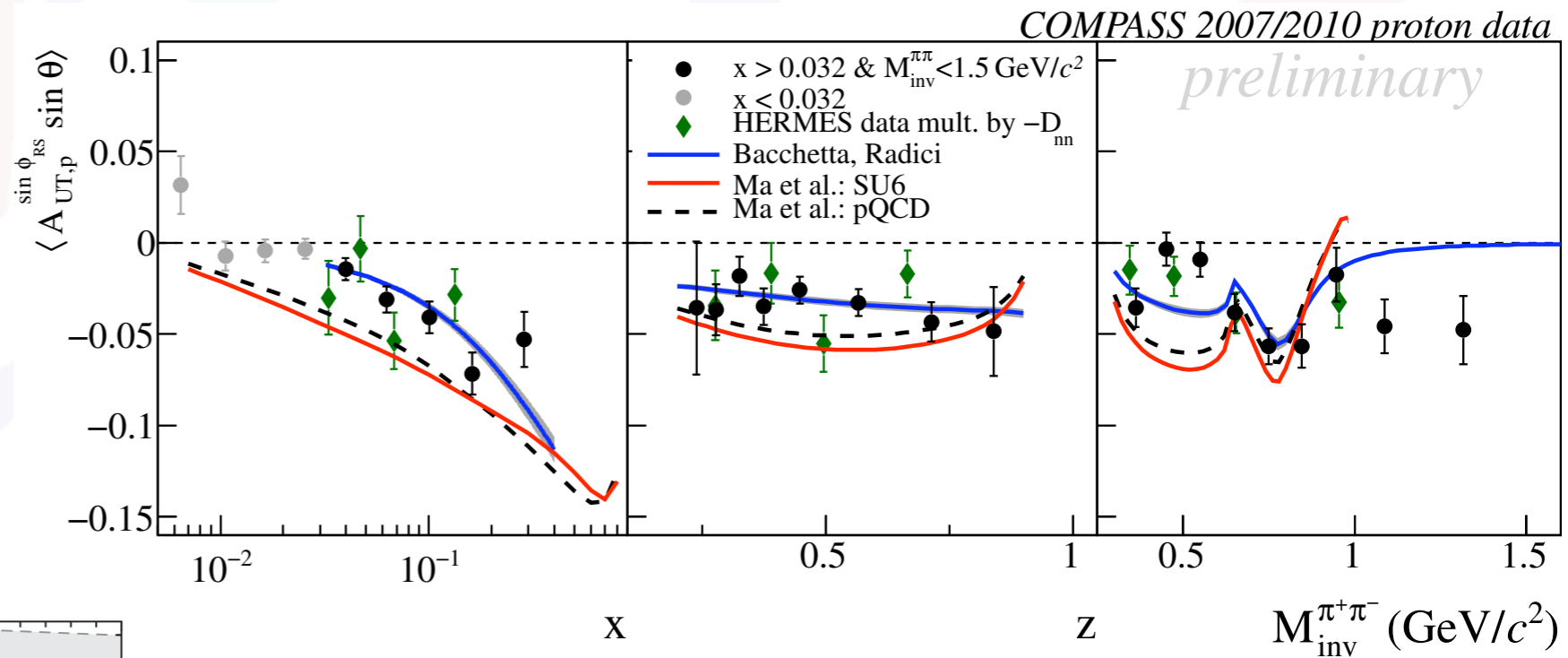


# Transversity through 2-hadron fragmentation

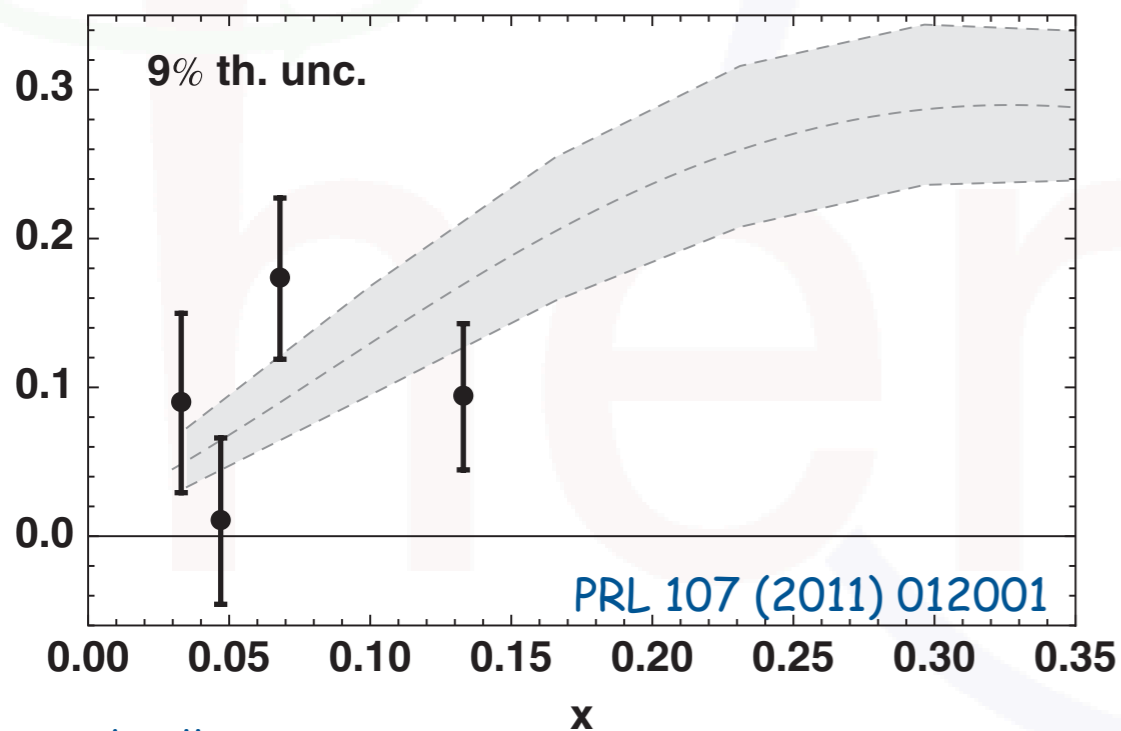
|   |                |          |                     |
|---|----------------|----------|---------------------|
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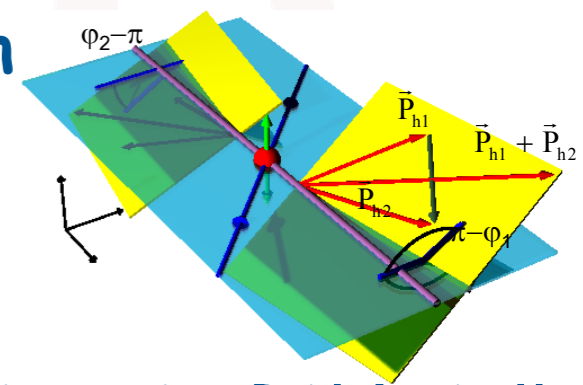
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$$x h_1^{u_v}(x) - x h_1^{d_v}(x)/4$$



- data from  $e^+e^-$  by BELLE allow first (collinear) extraction of transversity (compared to Anselmino et al.)

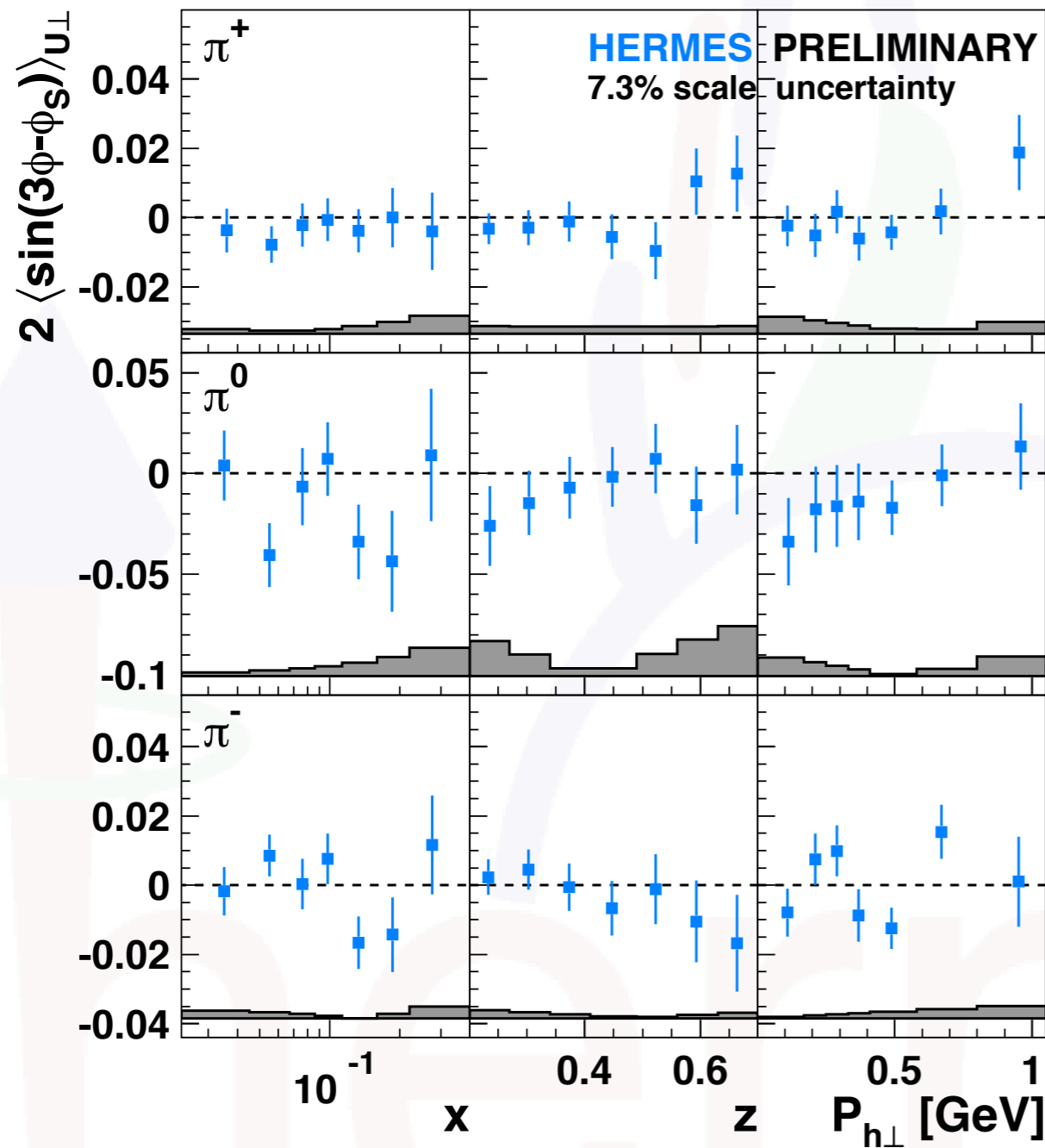


- updated analysis exists, not part of this talk

Transversity's friends

# Pretzelosity?

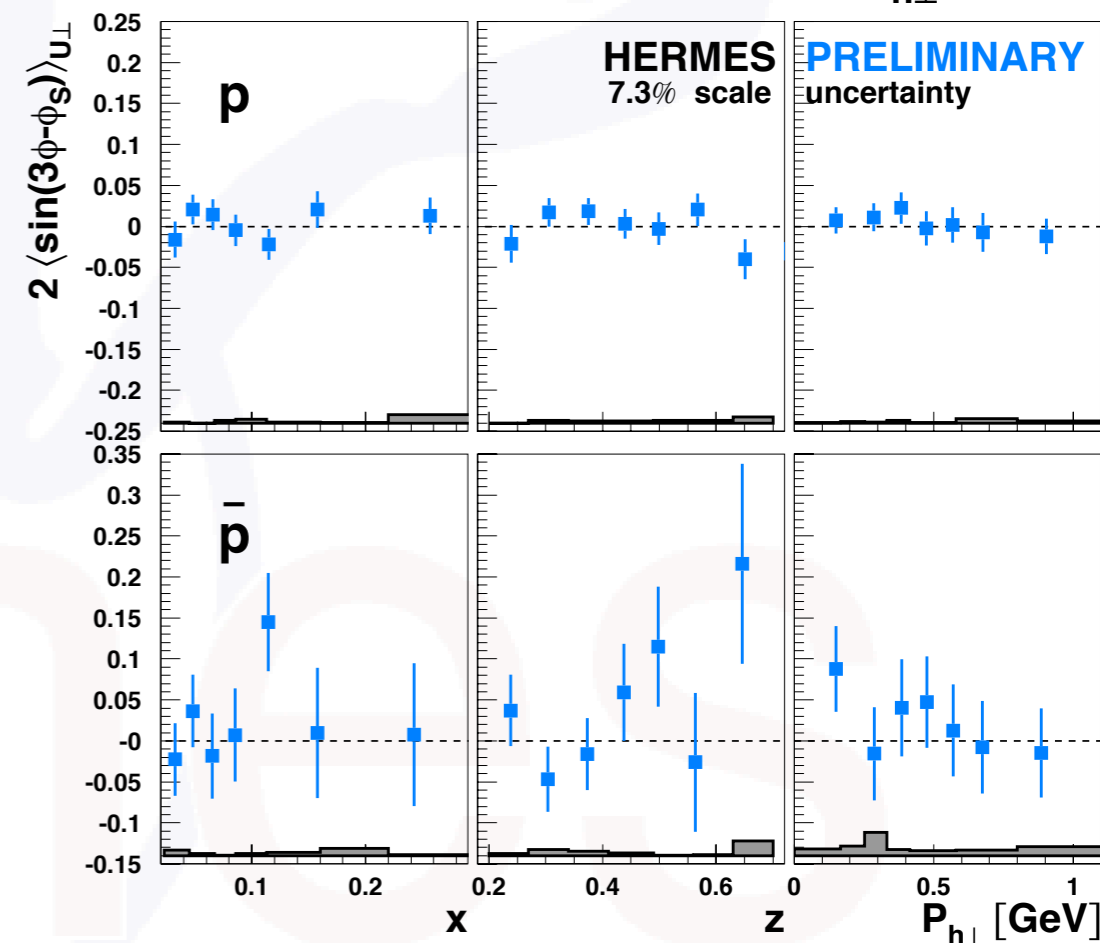
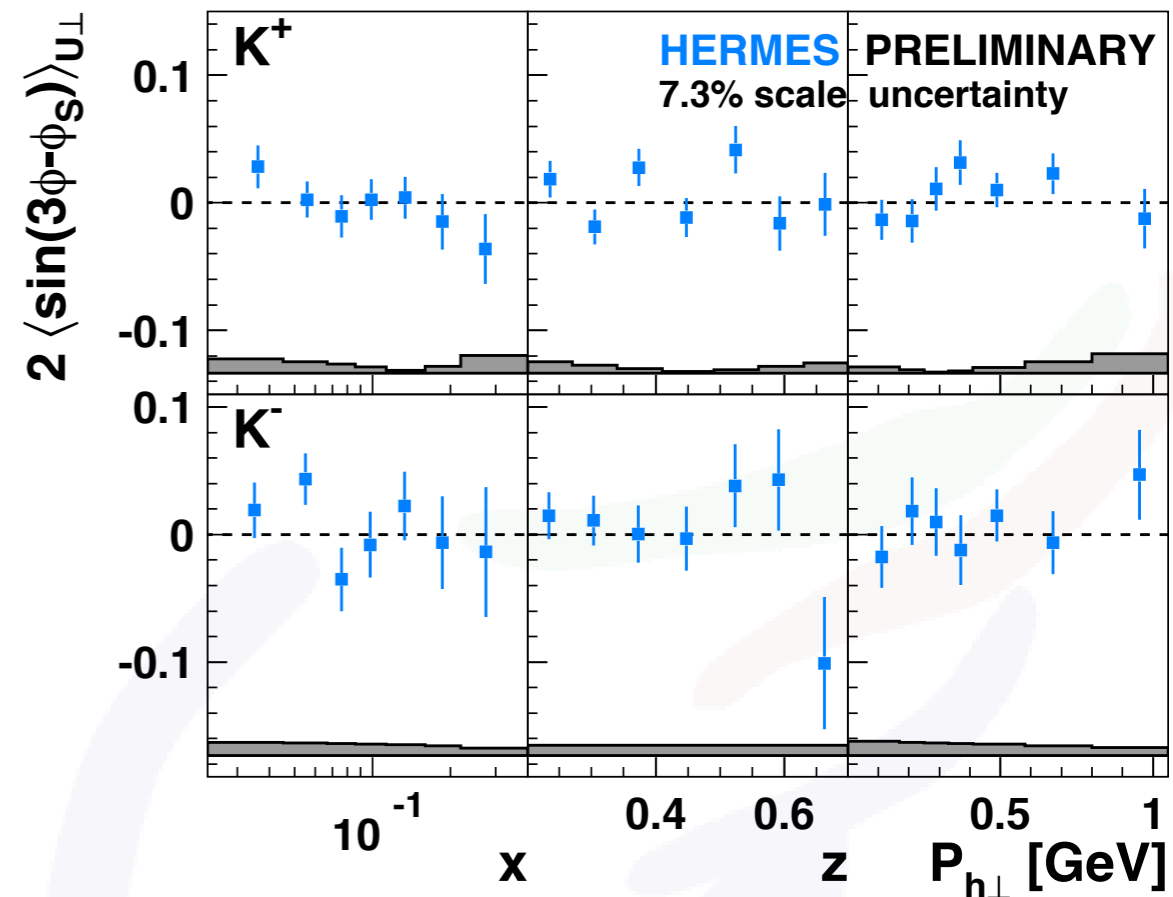
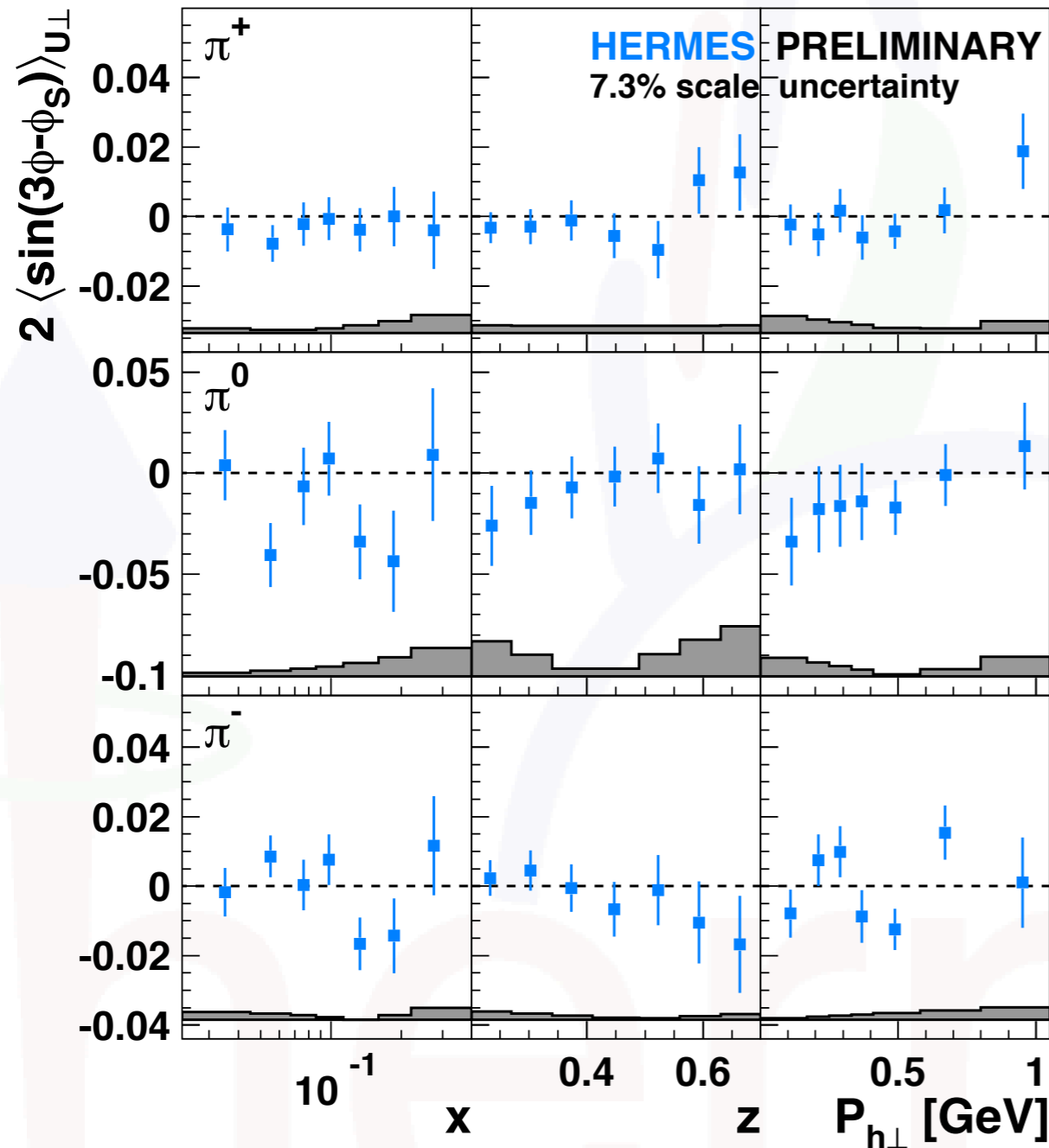
|   |                |          |                     |
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- consistent with zero; but suppressed by two powers of  $P_{h\perp}$  (compared to, e.g., transversity ⊗ Collins)

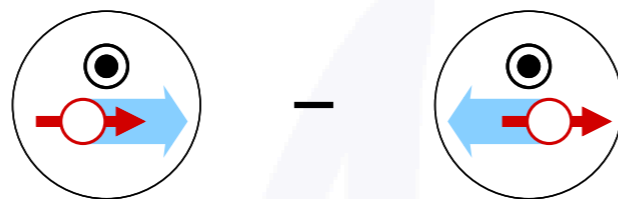
# Pretzelosity?

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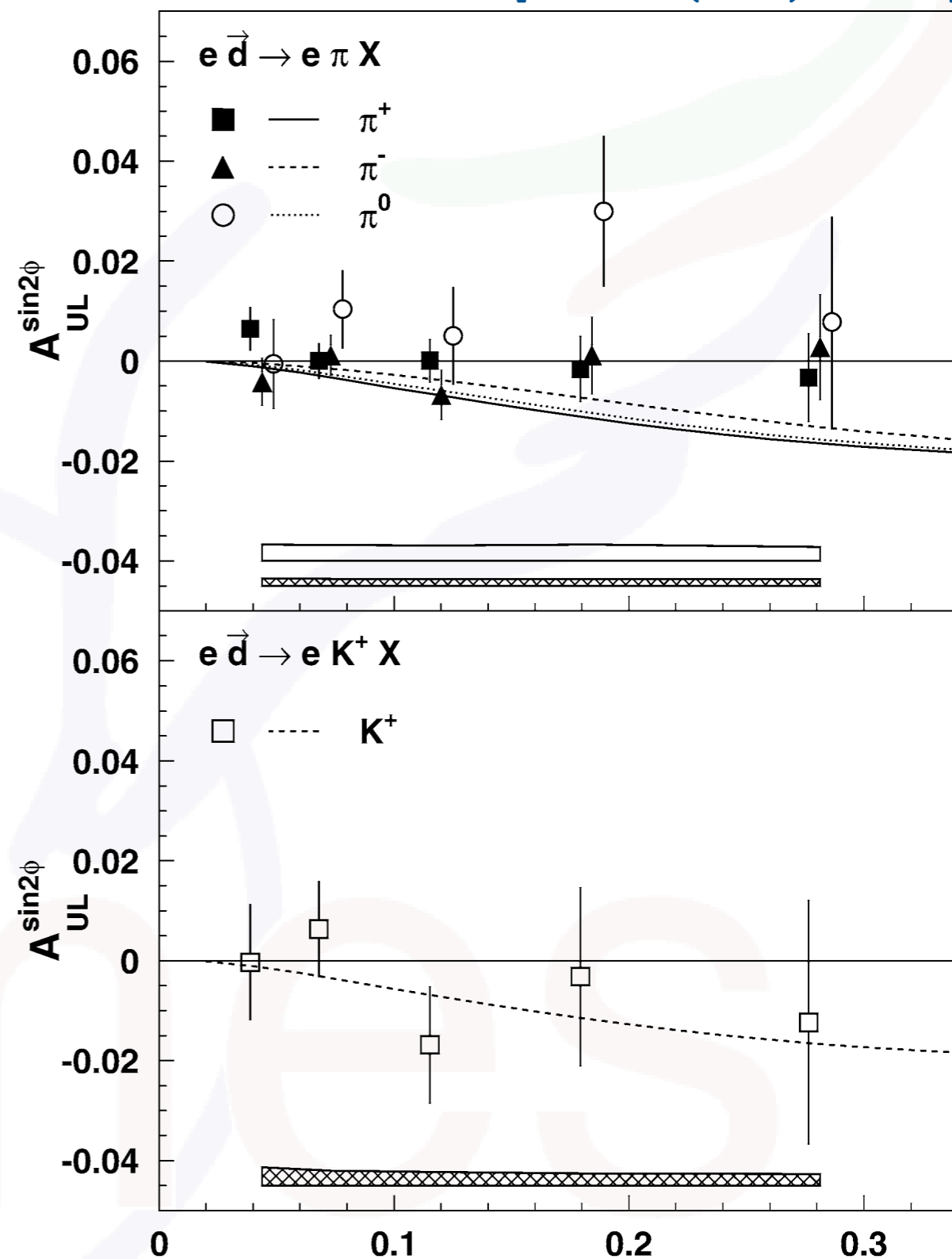
|   |                |          |                     |
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# Worm-Gear I

[PLB 562 (2003) 182-192]

- again: chiral-odd
- consistent with zero both for proton and deuteron

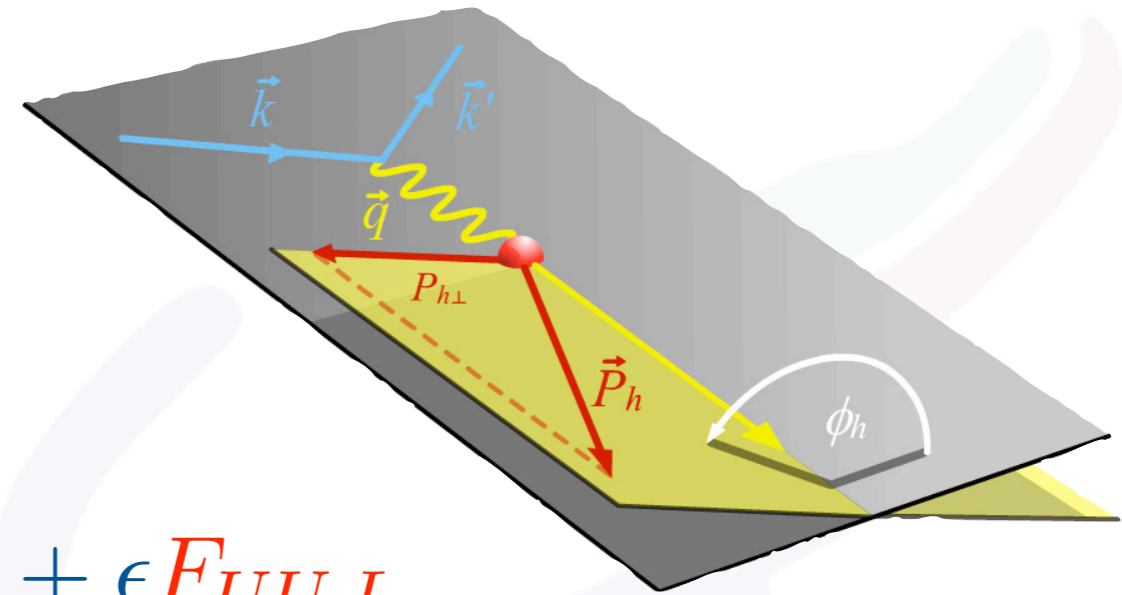


| Meson   | Deuterium target             | Proton target [2,3]          |
|---------|------------------------------|------------------------------|
| $\pi^+$ | $0.004 \pm 0.002 \pm 0.002$  | $-0.002 \pm 0.005 \pm 0.003$ |
| $\pi^0$ | $0.009 \pm 0.005 \pm 0.003$  | $0.006 \pm 0.007 \pm 0.003$  |
| $\pi^-$ | $0.001 \pm 0.003 \pm 0.002$  | $-0.005 \pm 0.006 \pm 0.005$ |
| $K^+$   | $-0.005 \pm 0.006 \pm 0.003$ | —                            |

$A_{UL}^{\sin 2\phi}$

[PLB 562 (2003) 182-192]

# cross section without polarization



$$\frac{d^5\sigma}{dx dy dz d\phi_h dP_{h\perp}^2} \propto \left(1 + \frac{\gamma^2}{2x}\right) \{F_{UU,T} + \epsilon F_{UU,L}$$

$$+ \sqrt{2\epsilon(1-\epsilon)} F_{UU}^{\cos\phi_h} \cos\phi_h + \epsilon F_{UU}^{\cos 2\phi_h} \cos 2\phi_h\}$$

$$F_{XY,Z} = F_{XY,Z}(x, y, z, P_{h\perp})$$

target polarization  
 ↓  
 beam polarization    virtual-photon polarization  
 ↑                    ↑

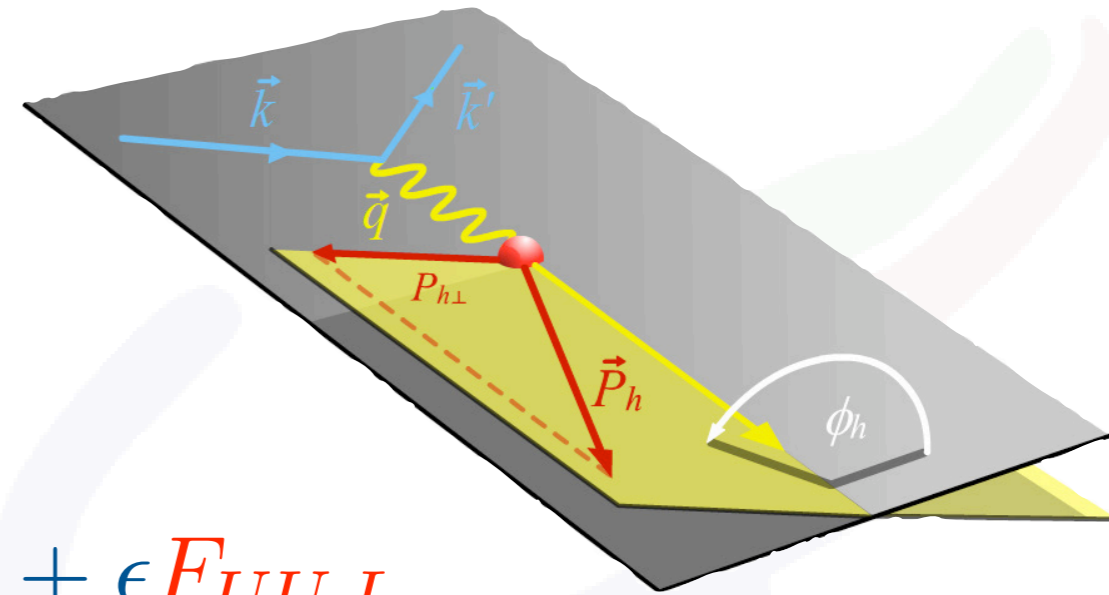
$$\gamma = \frac{2Mx}{Q}$$

$$\epsilon = \frac{1 - y - \frac{1}{4}\gamma^2 y^2}{1 - y + \frac{1}{2}y^2 + \frac{1}{4}\gamma^2 y^2}$$

[see, e.g., Bacchetta et al., JHEP 0702 (2007) 093]



# cross section without polarization



$$\frac{d^5\sigma}{dx dy dz d\phi_h dP_{h\perp}^2} \propto \left(1 + \frac{\gamma^2}{2x}\right) \{F_{UU,T} + \epsilon F_{UU,L} + \sqrt{2\epsilon(1-\epsilon)} F_{UU}^{\cos\phi_h} \cos\phi_h + \epsilon F_{UU}^{\cos 2\phi_h} \cos 2\phi_h\}$$

*leading twist*  
 $F_{UU}^{\cos 2\phi_h} \propto C \left[ \frac{2(\hat{P}_{h\perp} \cdot \vec{k}_T)(\hat{P}_{h\perp} \cdot \vec{p}_T) - \vec{k}_T \cdot \vec{p}_T}{MM_h} h_1^\perp H_1^\perp \right]$

*next to leading twist*  
 $F_{UU}^{\cos\phi_h} \propto \frac{2M}{Q} C \left[ \frac{\hat{P}_{h\perp} \cdot \vec{p}_T}{M_h} x h_1^\perp H_1^\perp - \frac{\hat{P}_{h\perp} \cdot \vec{k}_T}{M} x f_1 D_1 + \dots \right]$

BOER-MULDERS EFFECT (points to the blue box in the leading twist equation)

CAHN EFFECT (points to the red box in the next to leading twist equation)

Interaction dependent terms neglected (points to the ellipsis in the next to leading twist equation)

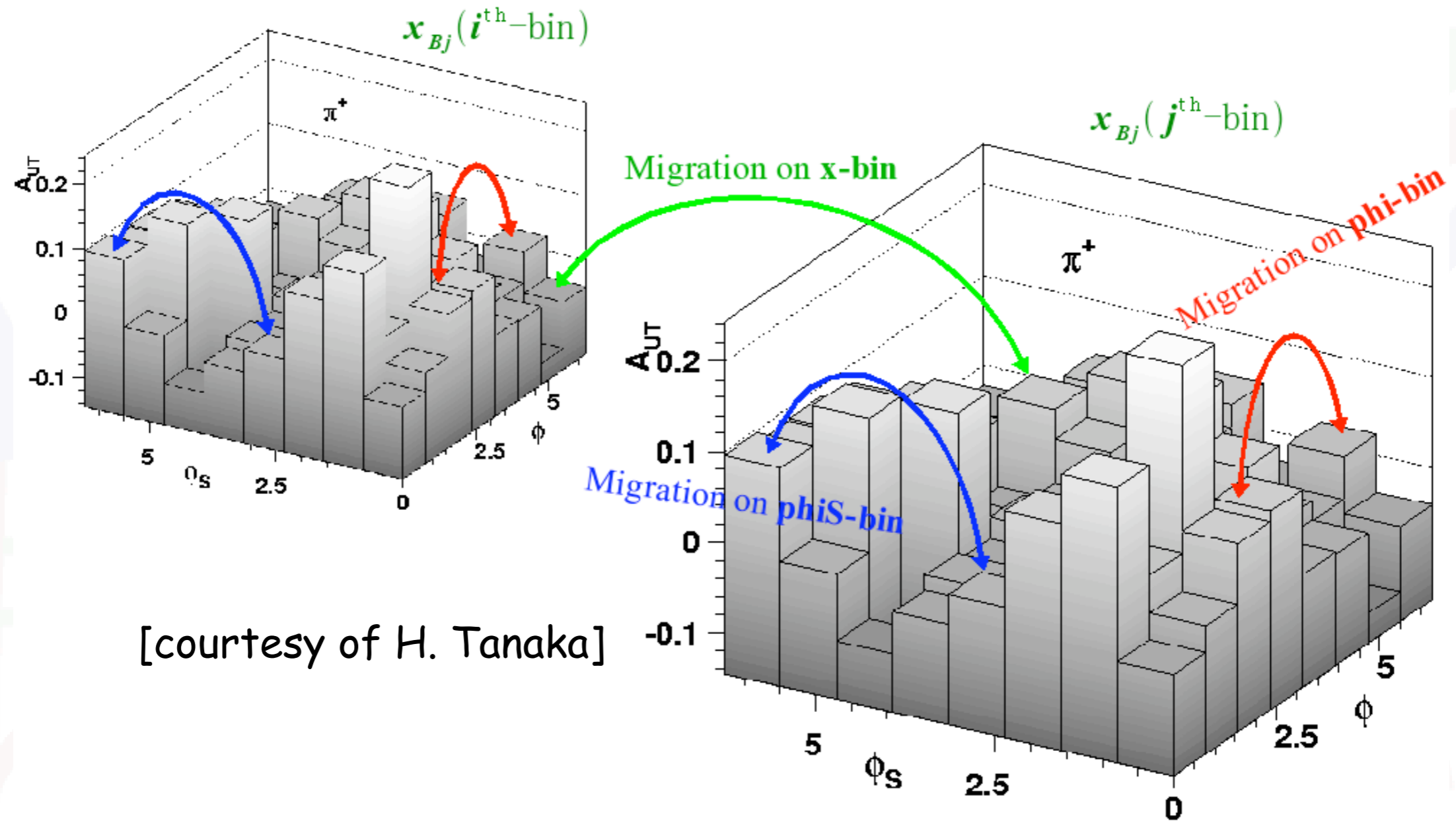
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$$\epsilon = \frac{1 - y - \frac{1}{4}\gamma^2 y^2}{1 - y + \frac{1}{2}y^2 + \frac{1}{4}\gamma^2 y^2}$$

[see, e.g., Bacchetta et al., JHEP 0702 (2007) 093]

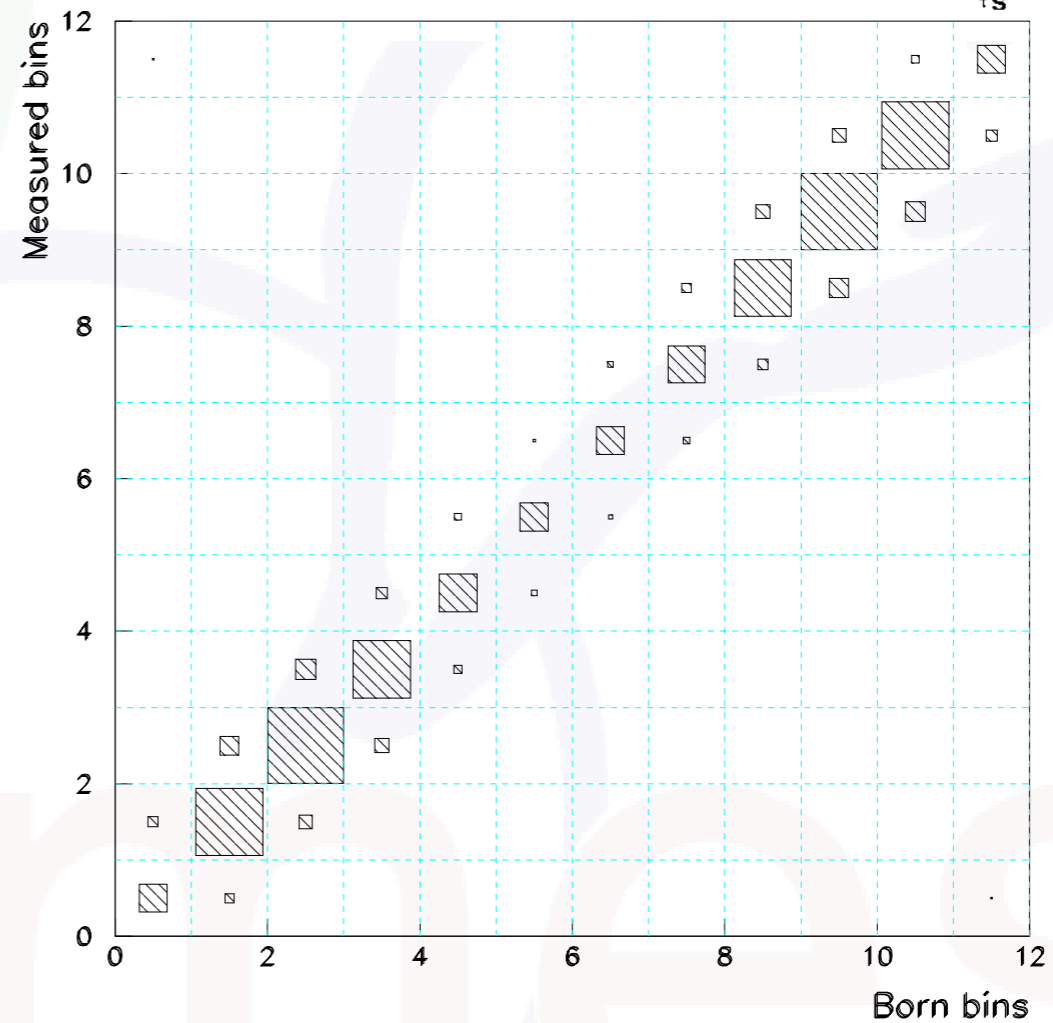
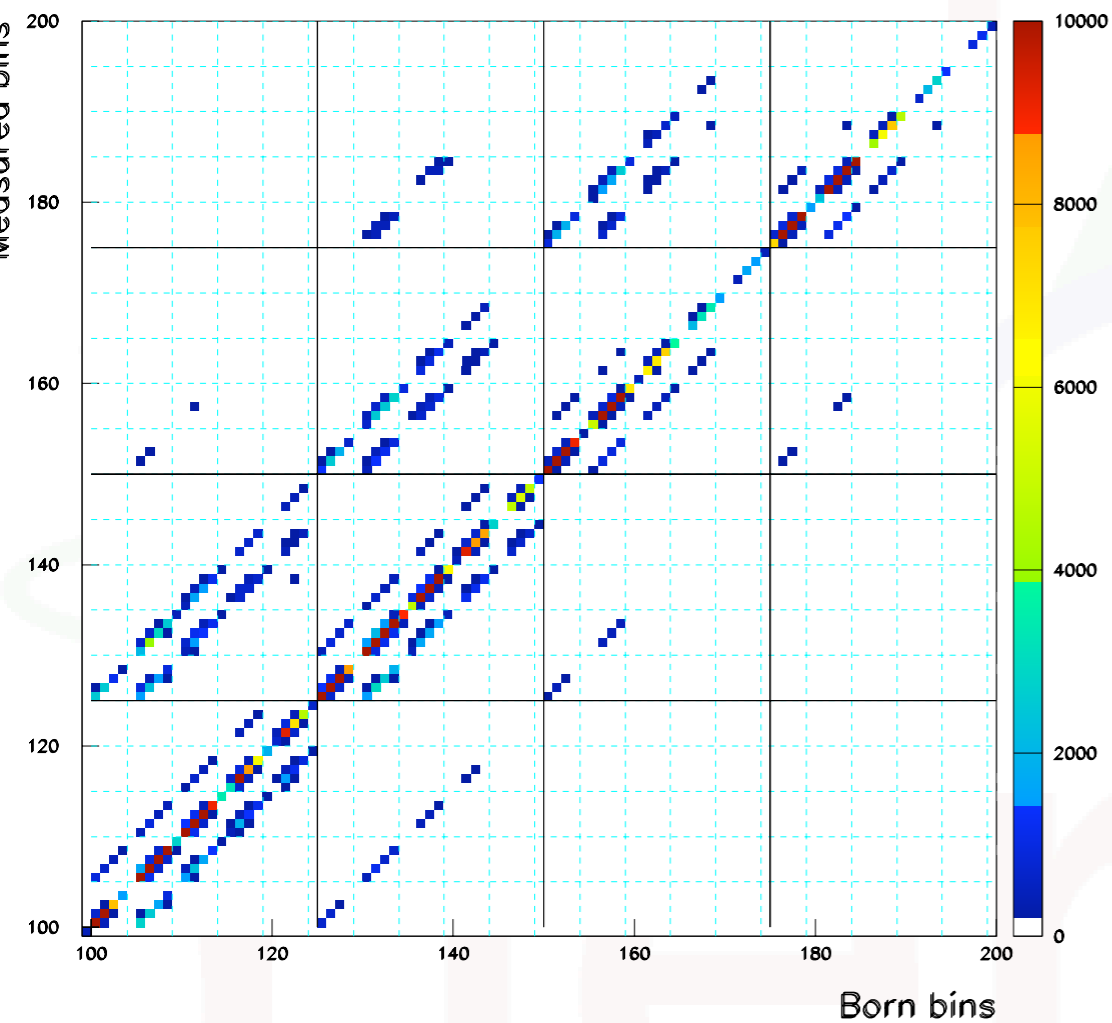
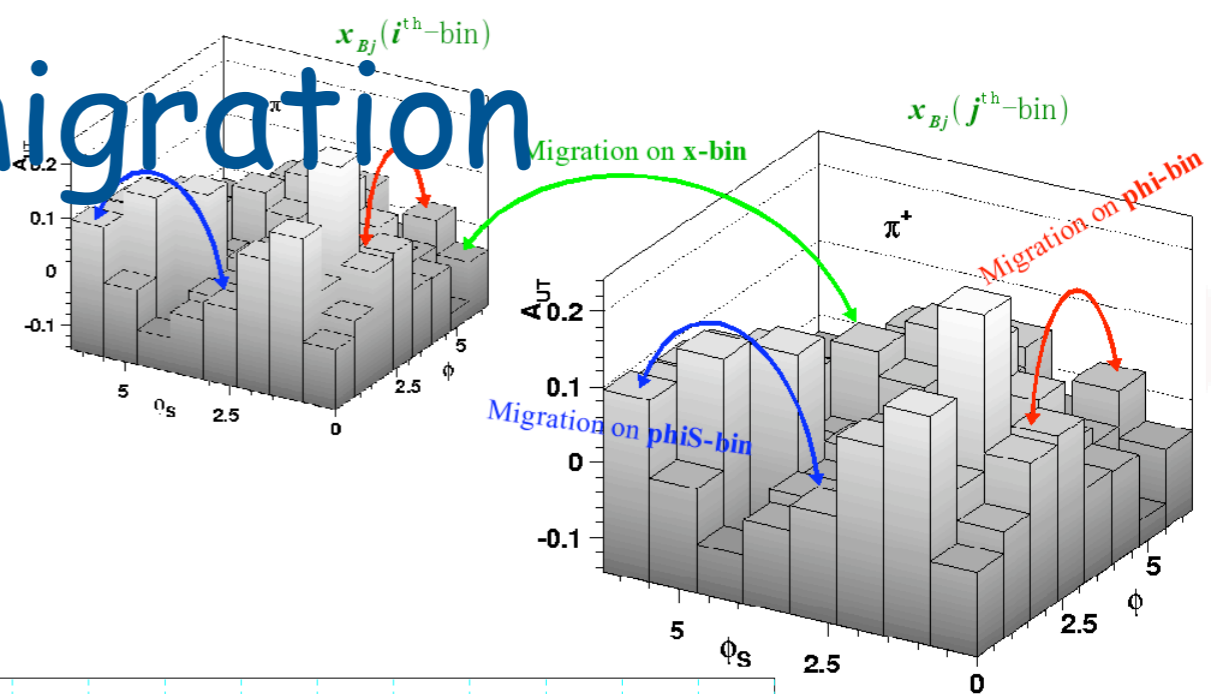
(Implicit sum over quark flavours)

# extraction I - event migration



[courtesy of H. Tanaka]

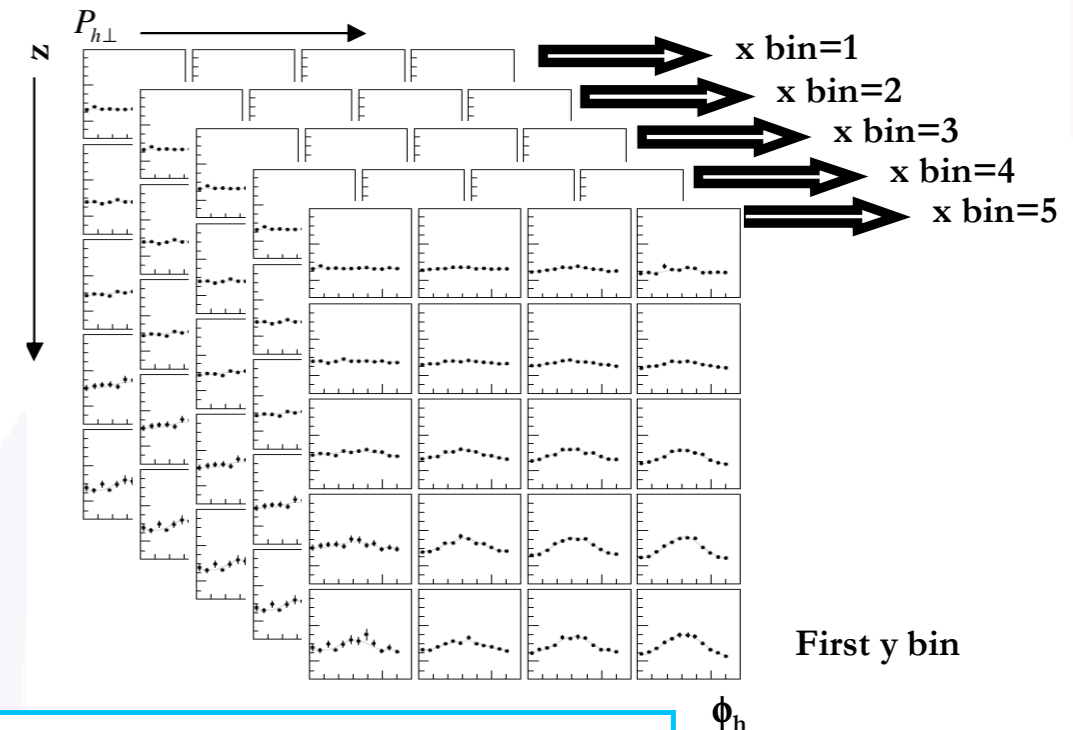
# extraction I - event migration



- migration correlates yields in different bins
- can't be corrected properly in bin-by-bin approach

# extraction II - unfolding

- Fully differential analysis in  $(x, y, z, P_{h\perp}, \phi)$
- Multi-dimensional unfolding: correction for finite acceptance, QED radiation, kinematic smearing, detector resolution



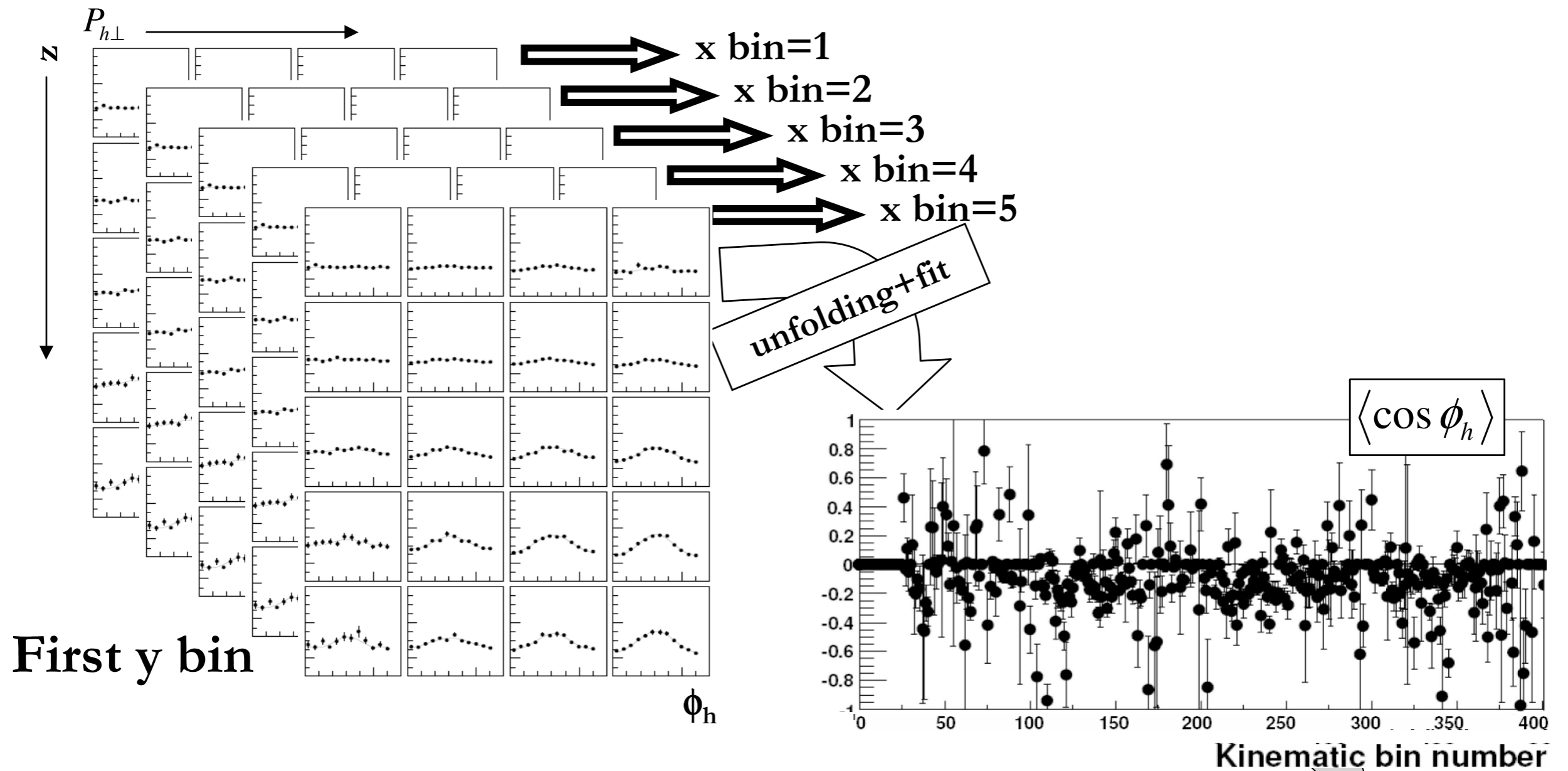
probability that an event generated with a certain kinematics is measured with a different kinematics

$$n_{EXP} = S n_{BORN} + n_{Bg}$$

$$n_{BORN} = S^{-1} [n_{EXP} - n_{Bg}]$$

includes the events smeared into the acceptance

# extraction III - projecting



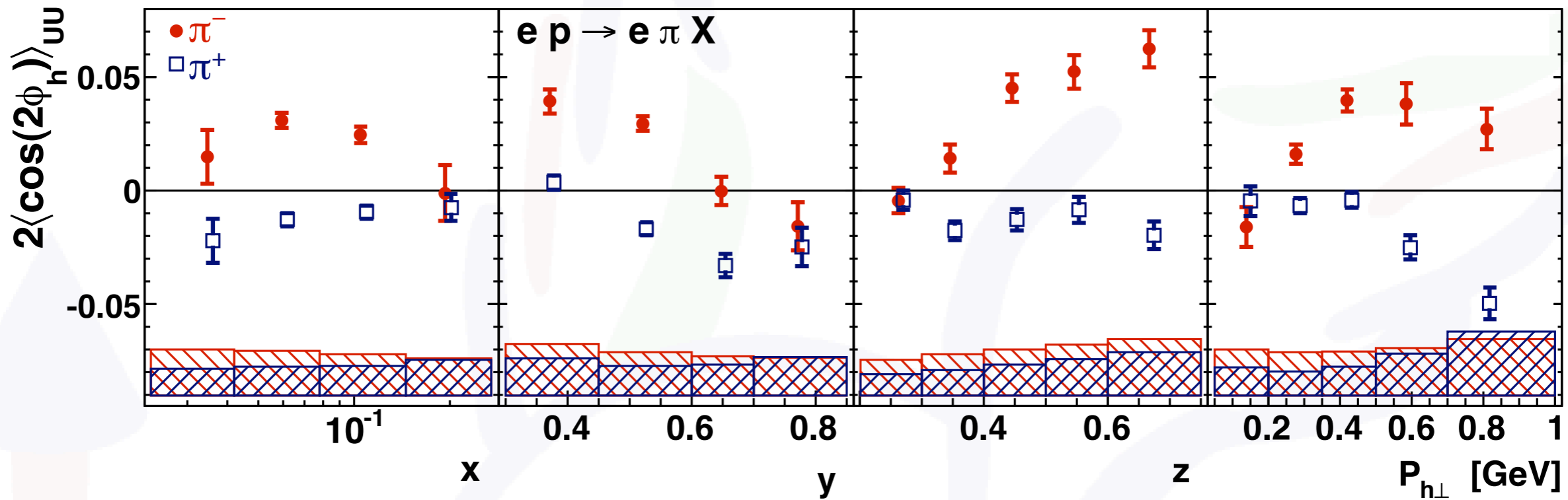
$$\langle \cos \phi \rangle(x_b) \approx \frac{\int_{0.3}^{0.85} dy \int_{0.2}^{0.75} dz \int_{0.05}^{0.75} dP_{h\perp}^2 \sigma^{4\pi}(\omega_{x_i=x_b}) \langle \cos \phi \rangle_{x_i=x_b}}{\int_{0.3}^{0.85} dy \int_{0.2}^{0.75} dz \int_{0.05}^{0.75} dP_{h\perp}^2 \sigma^{4\pi}(\omega_{x_i=x_b})}$$

|   |                |          |                     |
|---|----------------|----------|---------------------|
|   | U              | L        | T                   |
| U | $f_1$          |          | $h_1^\perp$         |
| L |                | $g_{1L}$ | $h_{1L}^\perp$      |
| T | $f_{1T}^\perp$ | $g_{1T}$ | $h_1, h_{1T}^\perp$ |



# signs of Boer-Mulders

[Airapetian et al., PRD 87 (2013) 012010]

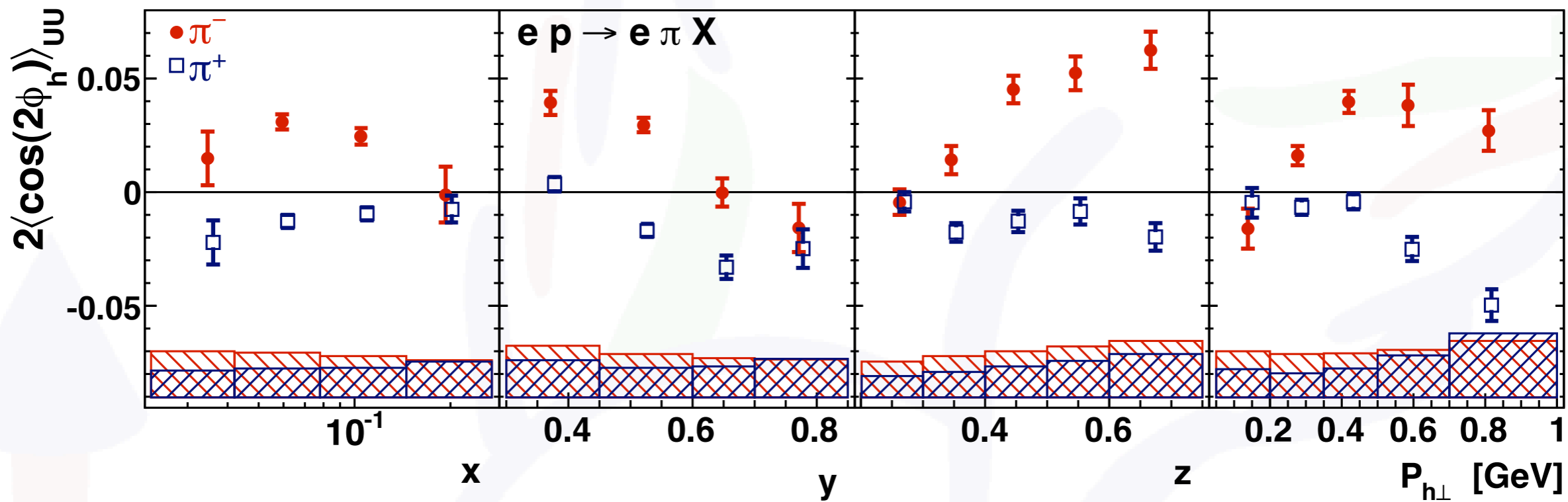


|   |                |          |                     |
|---|----------------|----------|---------------------|
|   | U              | L        | T                   |
| U | $f_1$          |          | $h_1^\perp$         |
| L |                | $g_{1L}$ | $h_{1L}^\perp$      |
| T | $f_{1T}^\perp$ | $g_{1T}$ | $h_1, h_{1T}^\perp$ |



# signs of Boer-Mulders

[Airapetian et al., PRD 87 (2013) 012010]



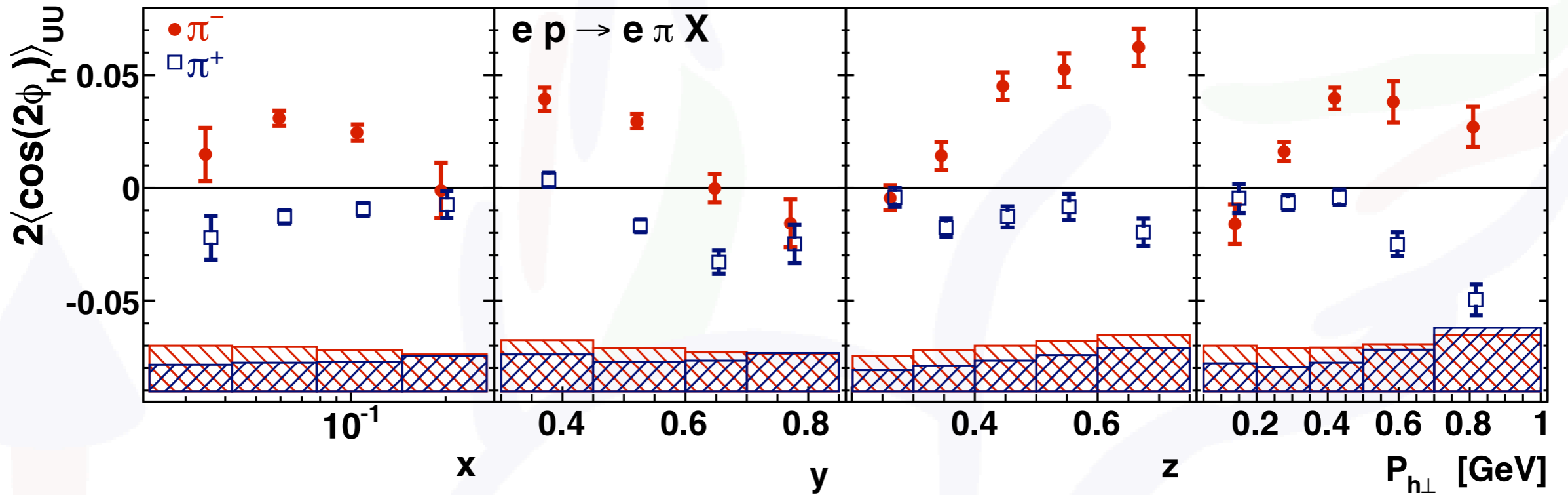
● modulations are not zero!

# signs of Boer-Mulders



|   | U              | L        | T                   |
|---|----------------|----------|---------------------|
| U | $f_1$          |          | $h_1^\perp$         |
| L |                | $g_{1L}$ | $h_{1L}^\perp$      |
| T | $f_{1T}^\perp$ | $g_{1T}$ | $h_1, h_{1T}^\perp$ |

[Airapetian et al., PRD 87 (2013) 012010]



- modulations are not zero!
- opposite sign for charged pions with larger magnitude for  $\pi^-$

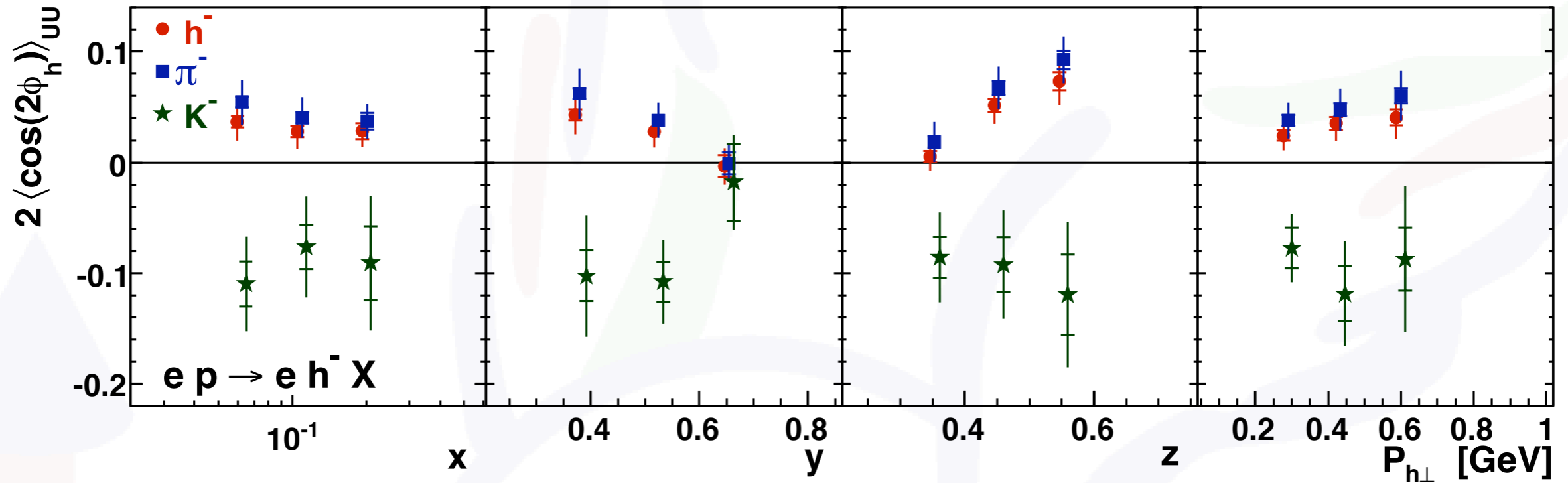


# signs of Boer-Mulders



|   | U              | L        | T                   |
|---|----------------|----------|---------------------|
| U | $f_1$          |          | $h_1^\perp$         |
| L |                | $g_{1L}$ | $h_{1L}^\perp$      |
| T | $f_{1T}^\perp$ | $g_{1T}$ | $h_1, h_{1T}^\perp$ |

[Airapetian et al., PRD 87 (2013) 012010]



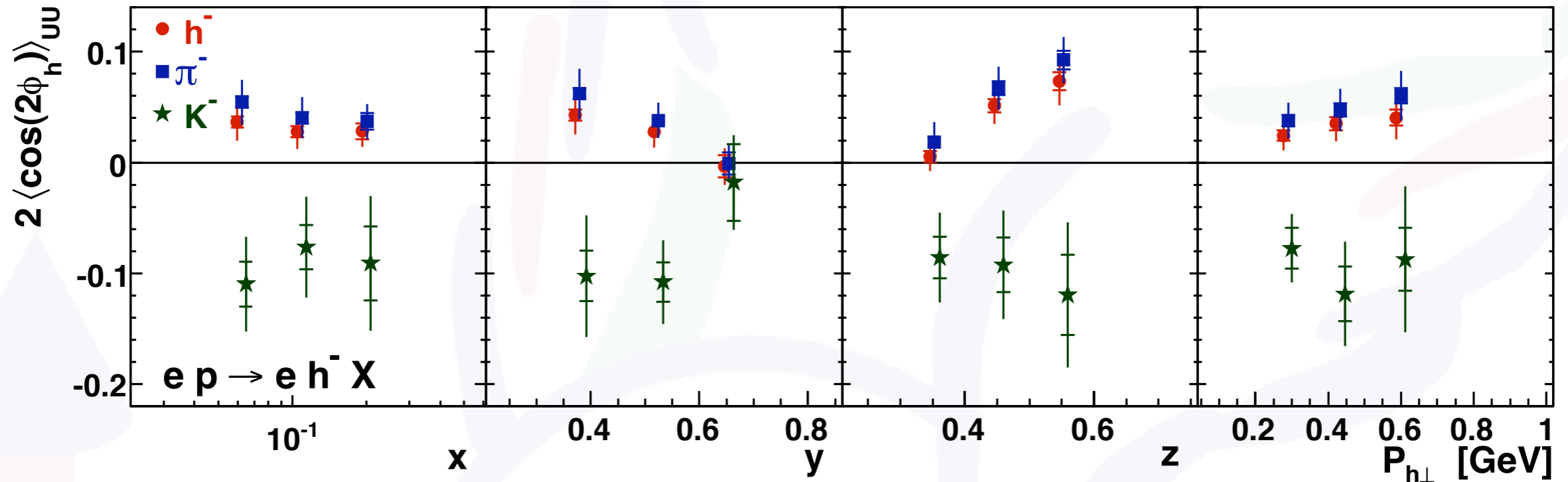
- modulations are not zero!
- opposite sign for charged pions with larger magnitude for  $\pi^-$
- intriguing behavior for kaons

# signs of Boer-Mulders



|   | U              | L        | T                   |
|---|----------------|----------|---------------------|
| U | $f_1$          |          | $h_1^\perp$         |
| L |                | $g_{1L}$ | $h_{1L}^\perp$      |
| T | $f_{1T}^\perp$ | $g_{1T}$ | $h_1, h_{1T}^\perp$ |

[Airapetian et al., PRD 87 (2013) 012010]



- modulations are not zero!
- opposite sign for charged pions with larger magnitude for  $\pi^-$
- intriguing behavior for kaons
- available in multidimensional binning, e.g., before projecting:  
<http://www-hermes.desy.de/cosnphi/>

# Cahn effect?

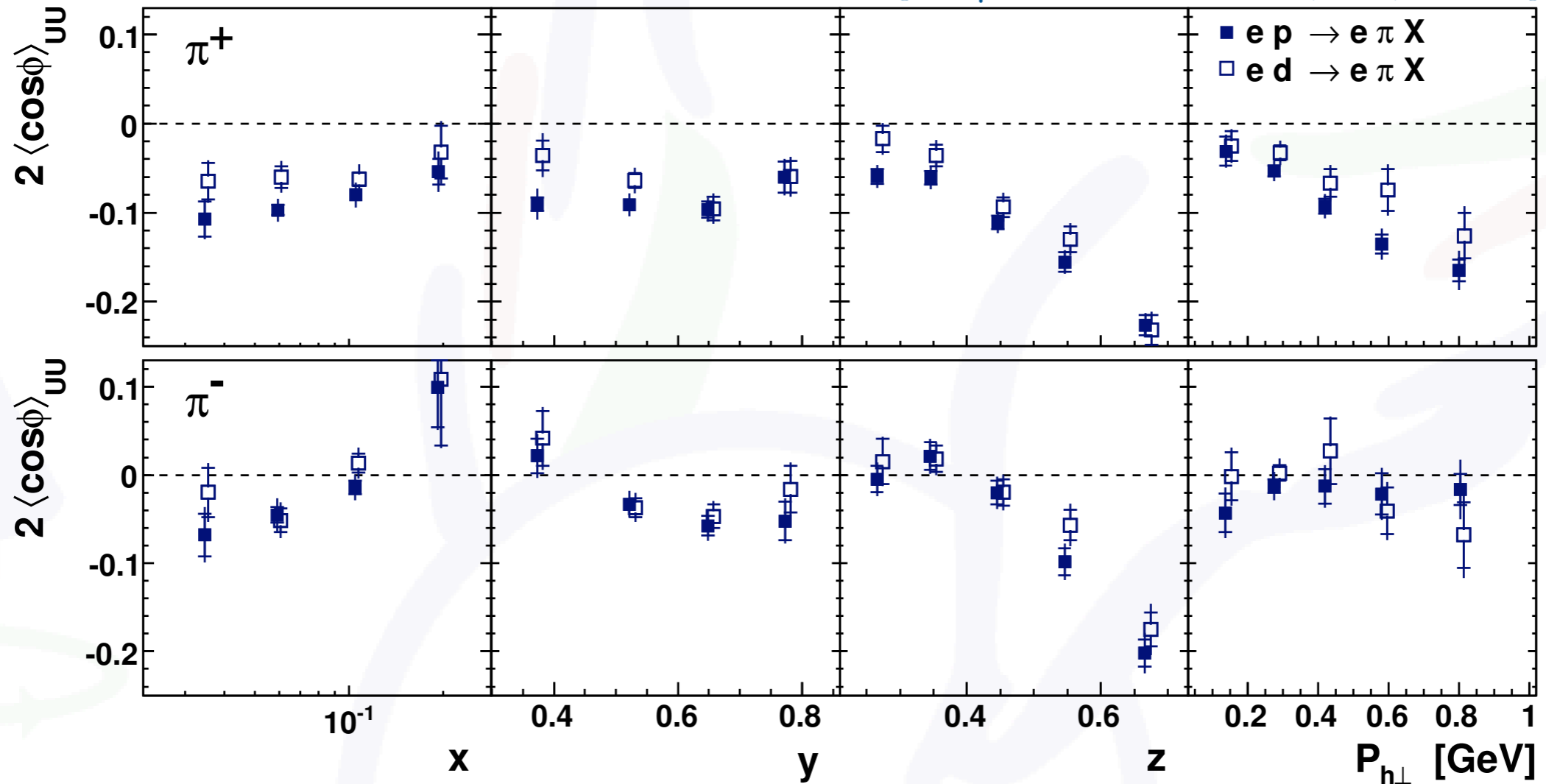
next to leading twist

$$F_{UU}^{\cos\phi_h} \propto \frac{2M}{Q} C$$

$$\left[ -\frac{\hat{P}_{h\perp} \cdot \vec{p}_T}{M_h} x h_1^\perp H_1^\perp - \frac{\hat{P}_{h\perp} \cdot \vec{k}_T}{M} x f_1 D_1 + \dots \right]$$

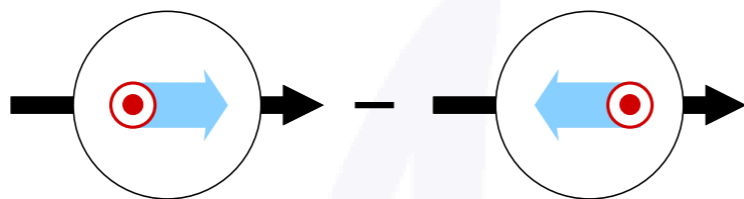
BOER-MULDERS EFFECT (blue arrow pointing to the first term)  
 CAHN EFFECT (red arrow pointing to the second term)  
 Interaction dependent terms neglected (black arrow pointing to the ellipsis)

[Airapetian et al., PRD 87 (2013) 012010]



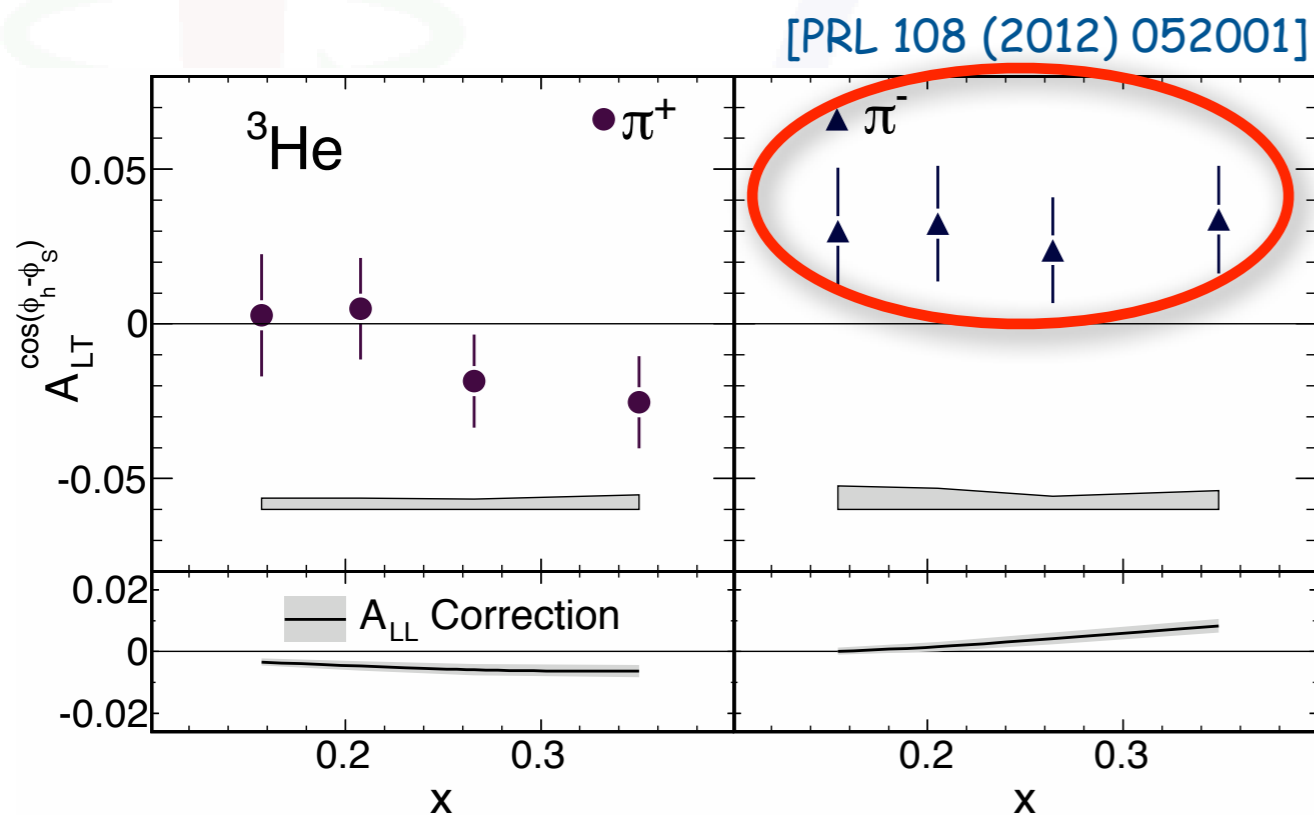
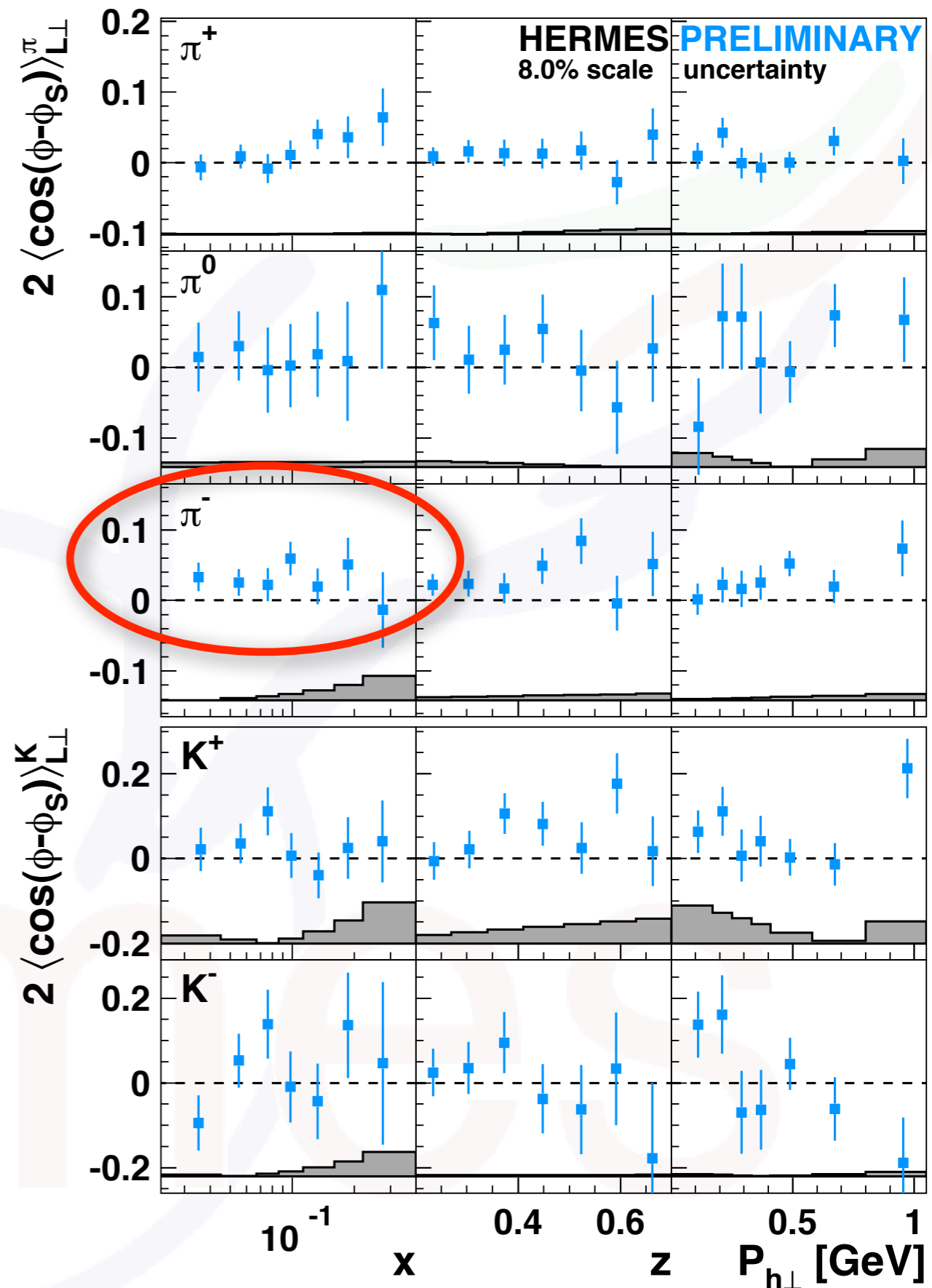
- no dependence on hadron charge **expected** for Cahn effect
- ➔ flavor dependence of transverse momentum
- ➔ sign of Boer-Mulders in  $\cos\phi$  modulation (indeed, overall pattern resembles B-M modulations)
- ➔ additional "genuine" twist-3 contributions?

|   |                |          |                     |
|---|----------------|----------|---------------------|
|   | U              | L        | T                   |
| U | $f_1$          |          | $h_1^\perp$         |
| L |                | $g_{1L}$ | $h_{1L}^\perp$      |
| T | $f_{1T}^\perp$ | $g_{1T}$ | $h_1, h_{1T}^\perp$ |

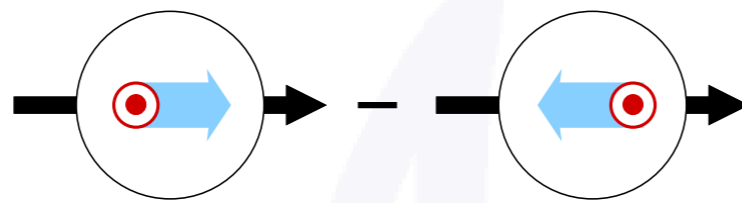


# Worm-Gear

- chiral even
- first direct evidence for worm-gear  $g_{1T}$  on
  - $^3\text{He}$  target at JLab
  - H target at HERMES

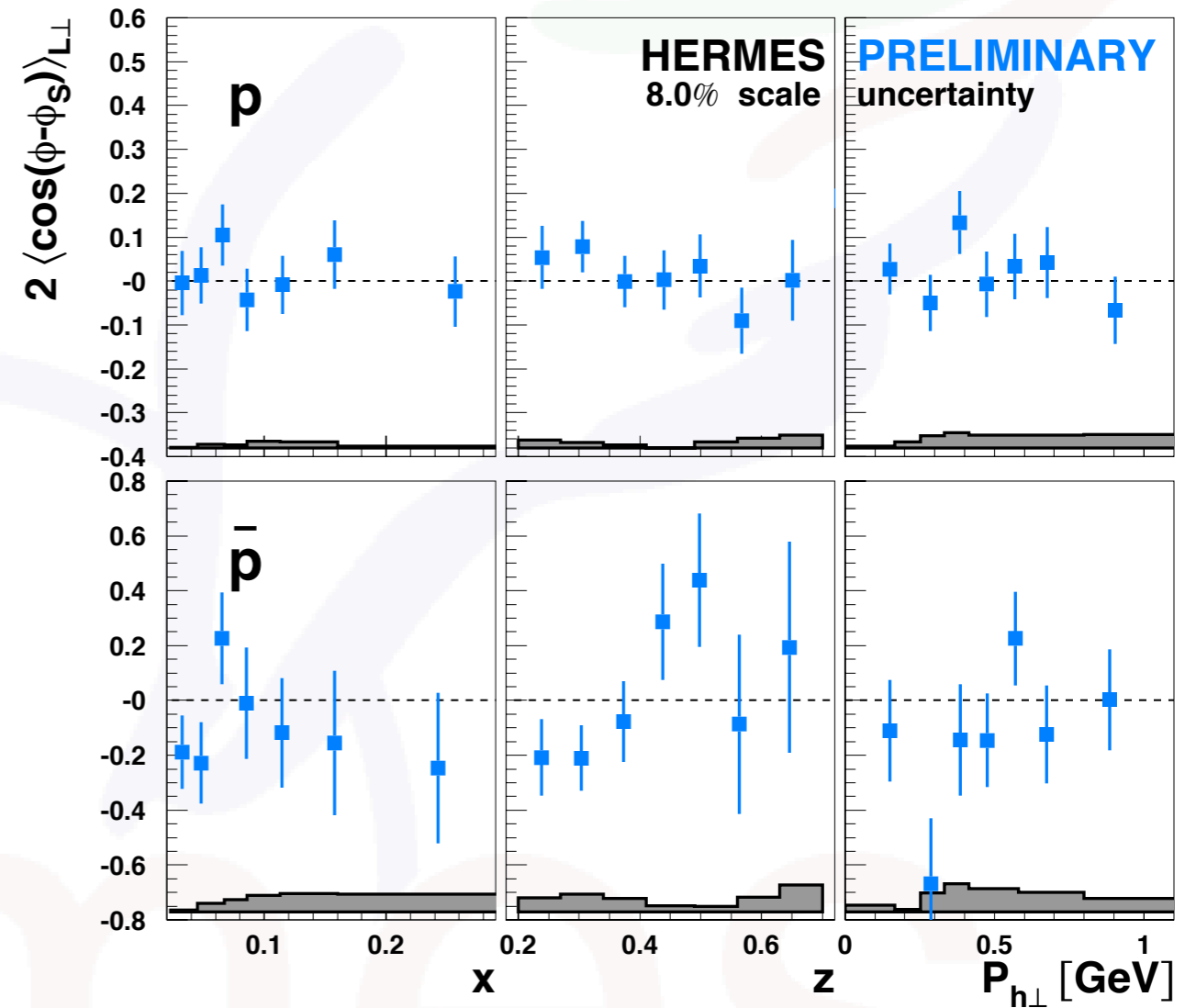


# Worm-Gear



|   | U              | L        | T                   |
|---|----------------|----------|---------------------|
| U | $f_1$          |          | $h_1^\perp$         |
| L |                | $g_{1L}$ | $h_{1L}^\perp$      |
| T | $f_{1T}^\perp$ | $g_{1T}$ | $h_1, h_{1T}^\perp$ |

- chiral even
- first direct evidence for worm-gear  $g_{1T}$  on
  - $^3\text{He}$  target at JLab
  - H target at HERMES
- results for protons and anti-protons consistent with zero



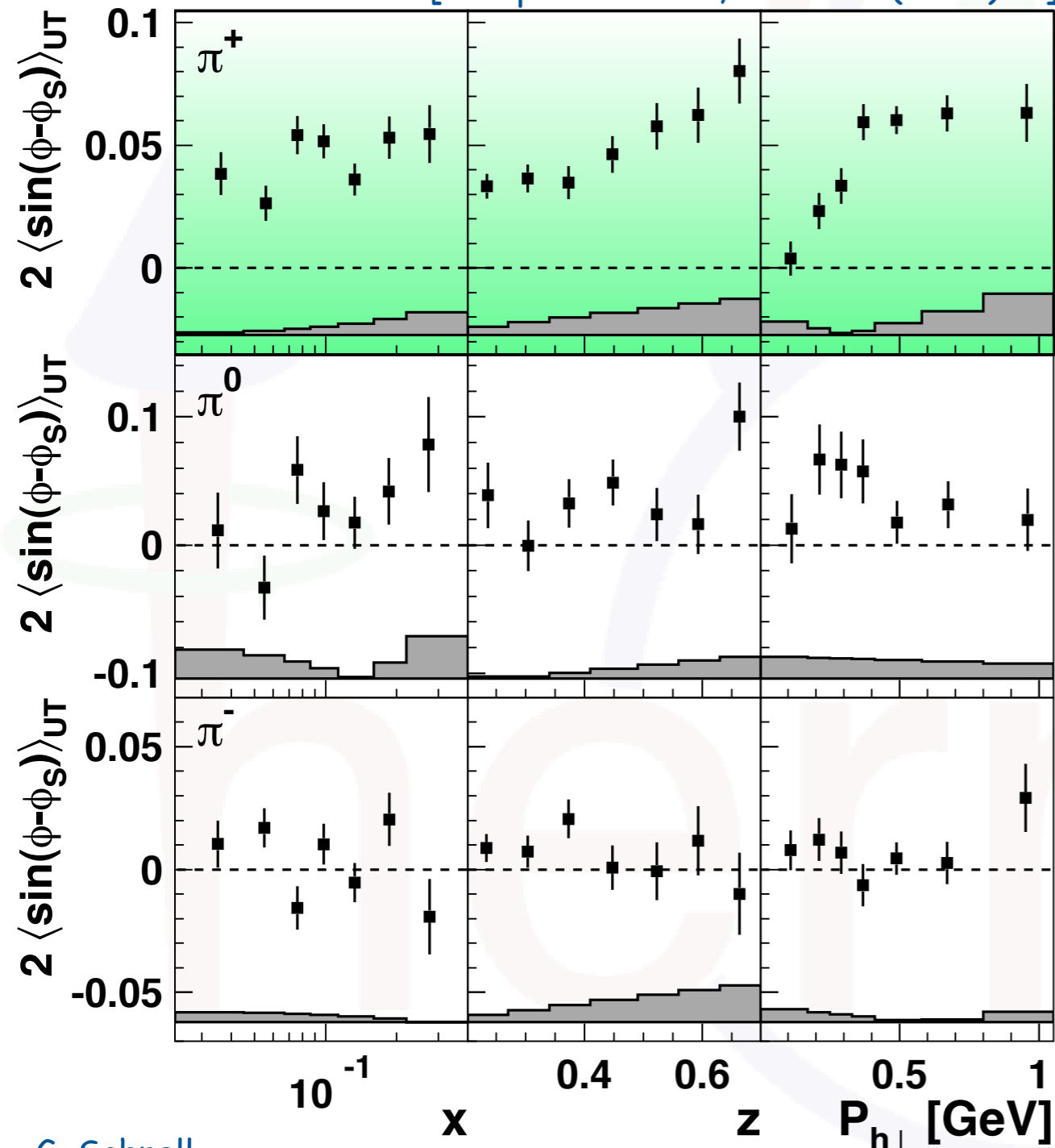
# Sivers amplitudes for pions

|   | U              | L        | T                   |
|---|----------------|----------|---------------------|
| U | $f_1$          |          | $h_1^\perp$         |
| L |                | $g_{1L}$ | $h_{1L}^\perp$      |
| T | $f_{1T}^\perp$ | $g_{1T}$ | $h_1, h_{1T}^\perp$ |

$$2\langle \sin(\phi - \phi_S) \rangle_{UT} =$$

$$= - \frac{\sum_q e_q^2 f_{1T}^{\perp,q}(x, p_T^2) \otimes_{\mathcal{W}} D_1^q(z, k_T^2)}{\sum_q e_q^2 f_1^q(x, p_T^2) \otimes D_1^q(z, k_T^2)}$$

[Airapetian et al., PLB 693 (2010) 11]



$\pi^+$  dominated by u-quark scattering:

$$\simeq - \frac{f_{1T}^{\perp,u}(x, p_T^2) \otimes_{\mathcal{W}} D_1^{u \rightarrow \pi^+}(z, k_T^2)}{f_1^u(x, p_T^2) \otimes D_1^{u \rightarrow \pi^+}(z, k_T^2)}$$

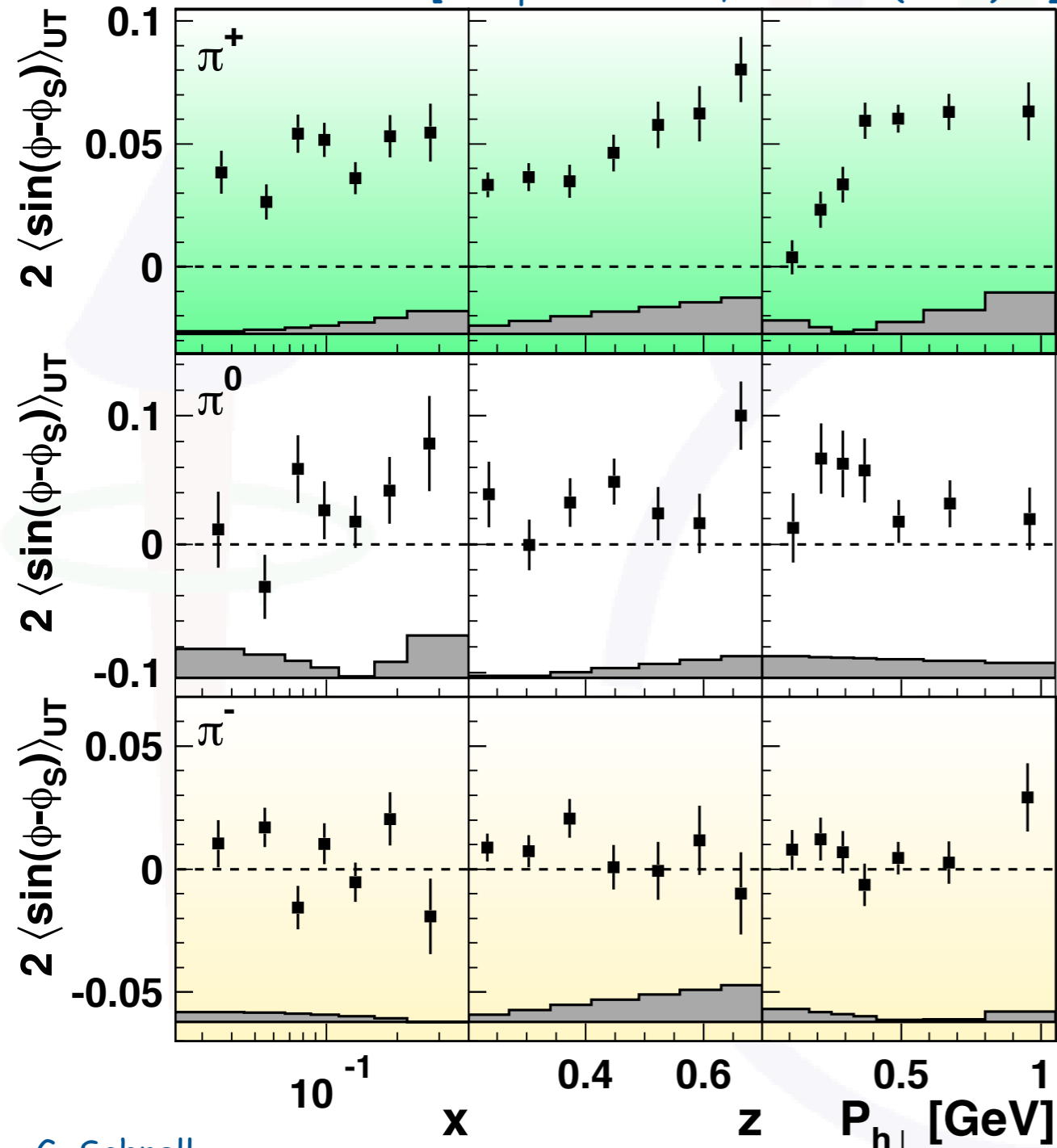
➡ u-quark Sivers DF < 0

# Sivers amplitudes for pions

|   | U              | L        | T                   |
|---|----------------|----------|---------------------|
| U | $f_1$          |          | $h_1^\perp$         |
| L |                | $g_{1L}$ | $h_{1L}^\perp$      |
| T | $f_{1T}^\perp$ | $g_{1T}$ | $h_1, h_{1T}^\perp$ |

$$2\langle \sin(\phi - \phi_S) \rangle_{UT} = - \frac{\sum_q e_q^2 f_{1T}^{\perp,q}(x, p_T^2) \otimes_{\mathcal{W}} D_1^q(z, k_T^2)}{\sum_q e_q^2 f_1^q(x, p_T^2) \otimes D_1^q(z, k_T^2)}$$

[Airapetian et al., PLB 693 (2010) 11]



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➡ u-quark Sivers DF < 0

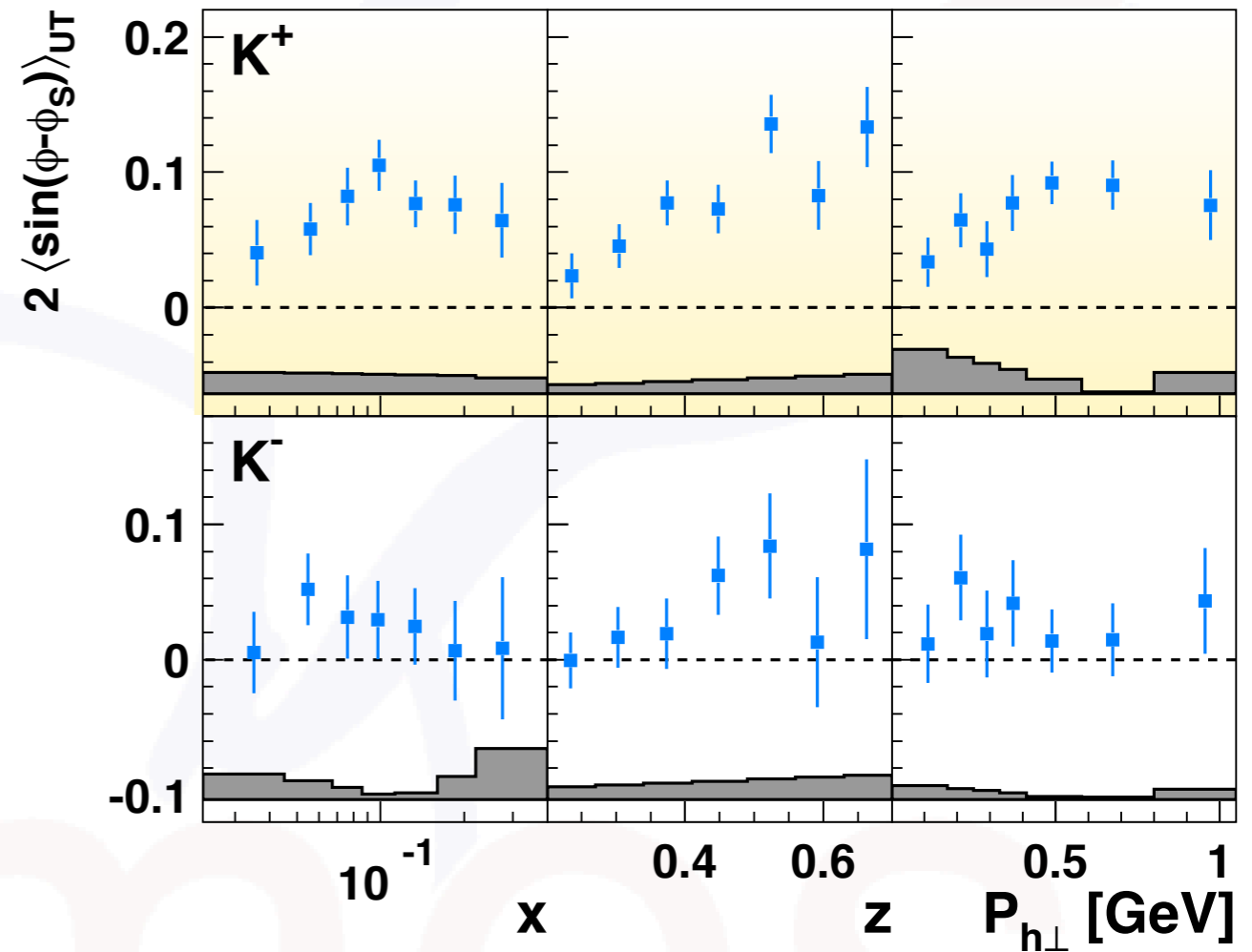
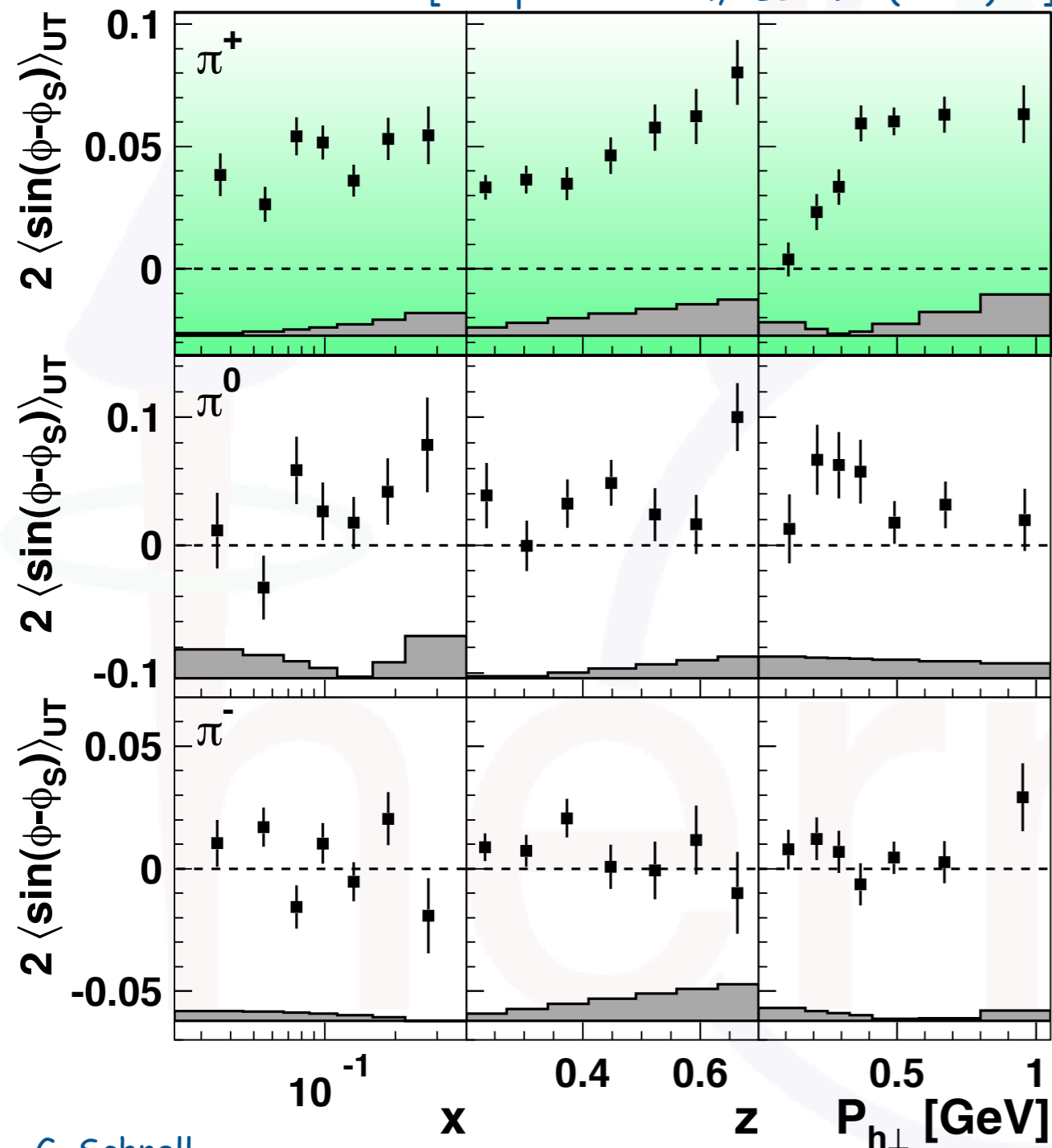
➡ d-quark Sivers DF > 0  
(cancellation for  $\pi^-$ )

# Sivers amplitudes for mesons

|   |                |          |                     |
|---|----------------|----------|---------------------|
|   | U              | L        | T                   |
| U | $f_1$          |          | $h_1^\perp$         |
| L |                | $g_{1L}$ | $h_{1L}^\perp$      |
| T | $f_{1T}^\perp$ | $g_{1T}$ | $h_1, h_{1T}^\perp$ |

$$2\langle \sin(\phi - \phi_S) \rangle_{UT} = - \frac{\sum_q e_q^2 f_{1T}^{\perp,q}(x, p_T^2) \otimes_{\mathcal{W}} D_1^q(z, k_T^2)}{\sum_q e_q^2 f_1^q(x, p_T^2) \otimes D_1^q(z, k_T^2)}$$

[Airapetian et al., PLB 693 (2010) 11]



➡ larger amplitudes for positive kaons vs. pions

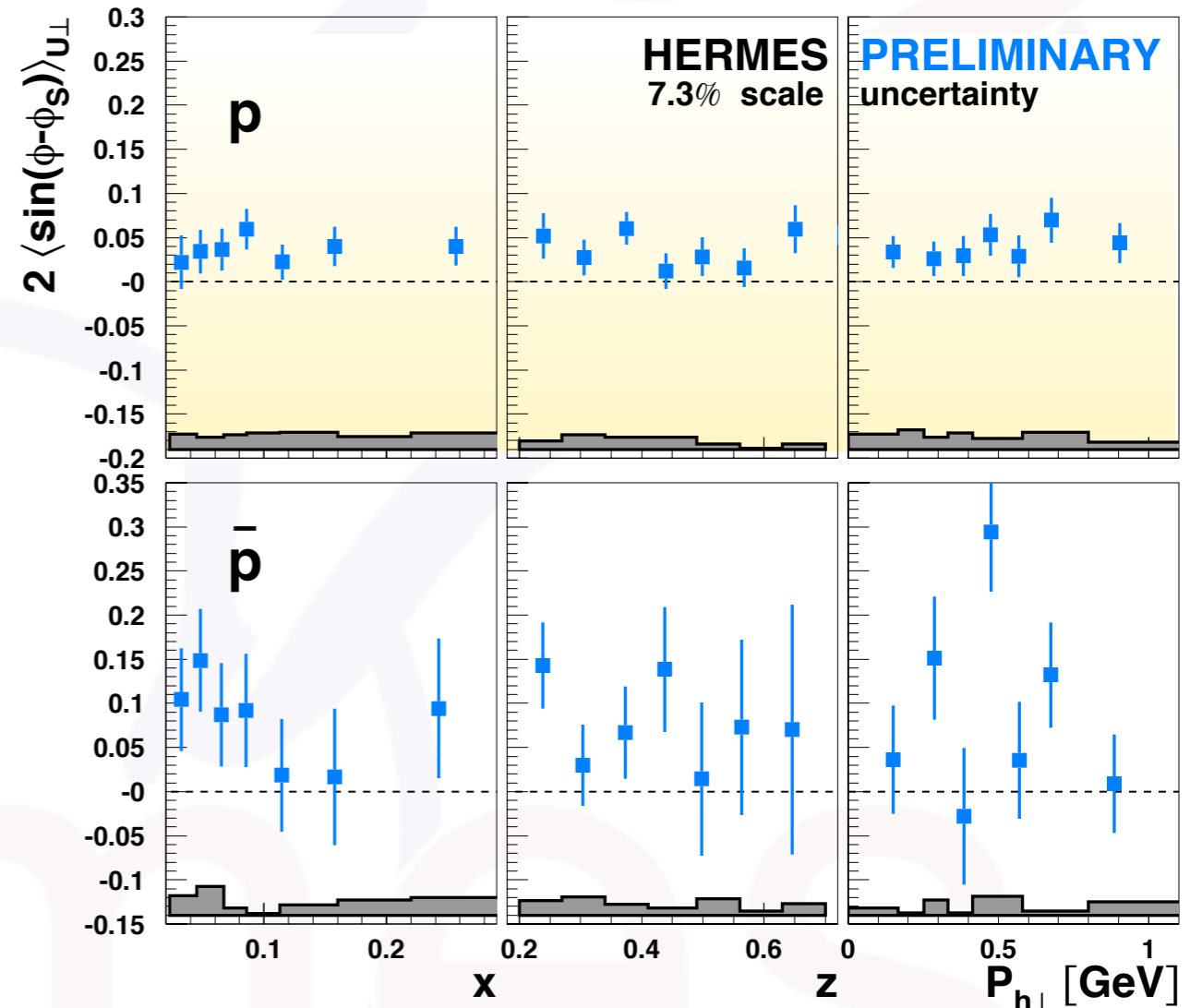
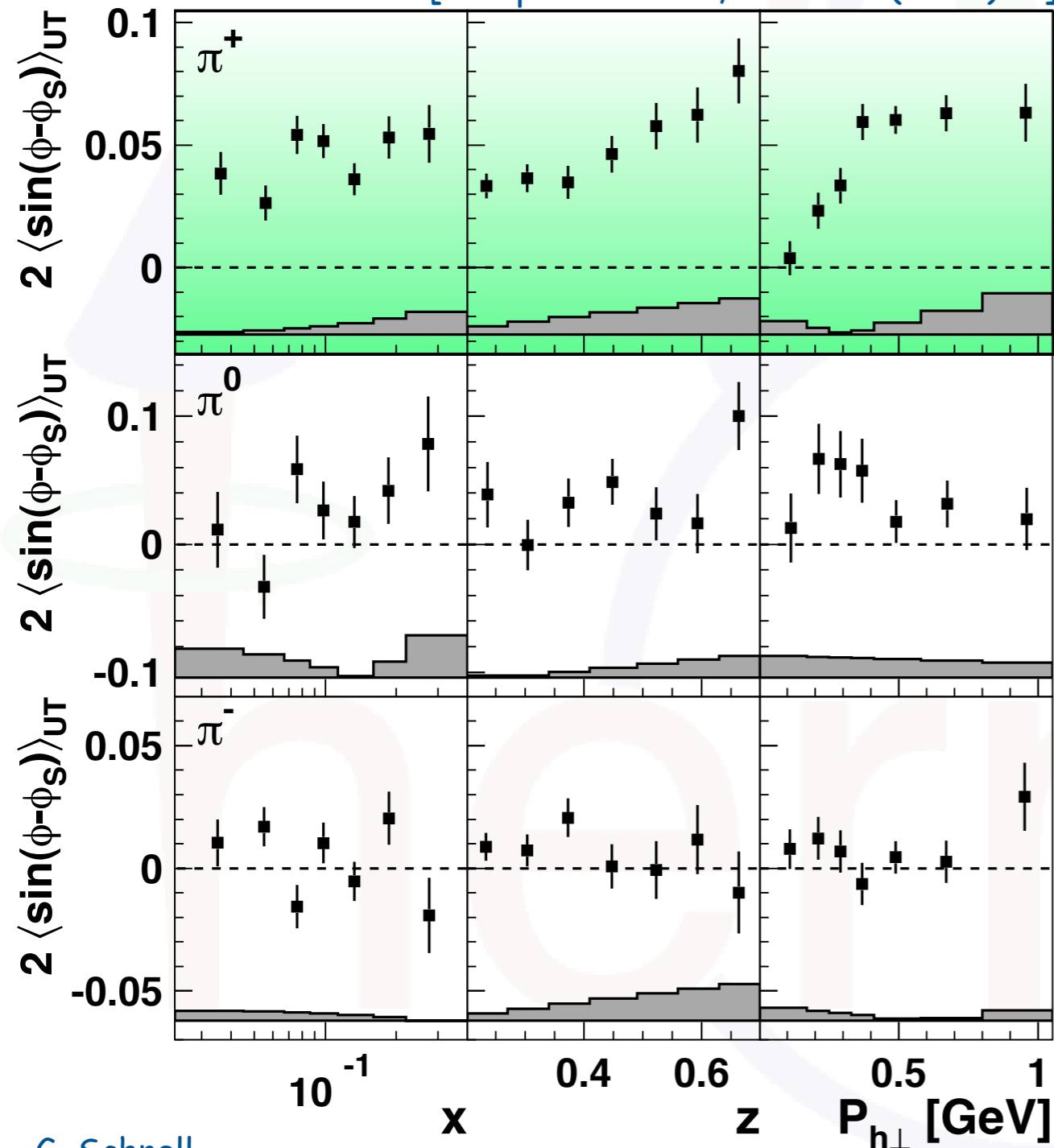


# Sivers amplitudes for baryons

|   |                |          |                     |
|---|----------------|----------|---------------------|
|   | U              | L        | T                   |
| U | $f_1$          |          | $h_1^\perp$         |
| L |                | $g_{1L}$ | $h_{1L}^\perp$      |
| T | $f_{1T}^\perp$ | $g_{1T}$ | $h_1, h_{1T}^\perp$ |

$$2\langle \sin(\phi - \phi_S) \rangle_{UT} = - \frac{\sum_q e_q^2 f_{1T}^{\perp,q}(x, p_T^2) \otimes_{\mathcal{W}} D_1^q(z, k_T^2)}{\sum_q e_q^2 f_1^q(x, p_T^2) \otimes D_1^q(z, k_T^2)}$$

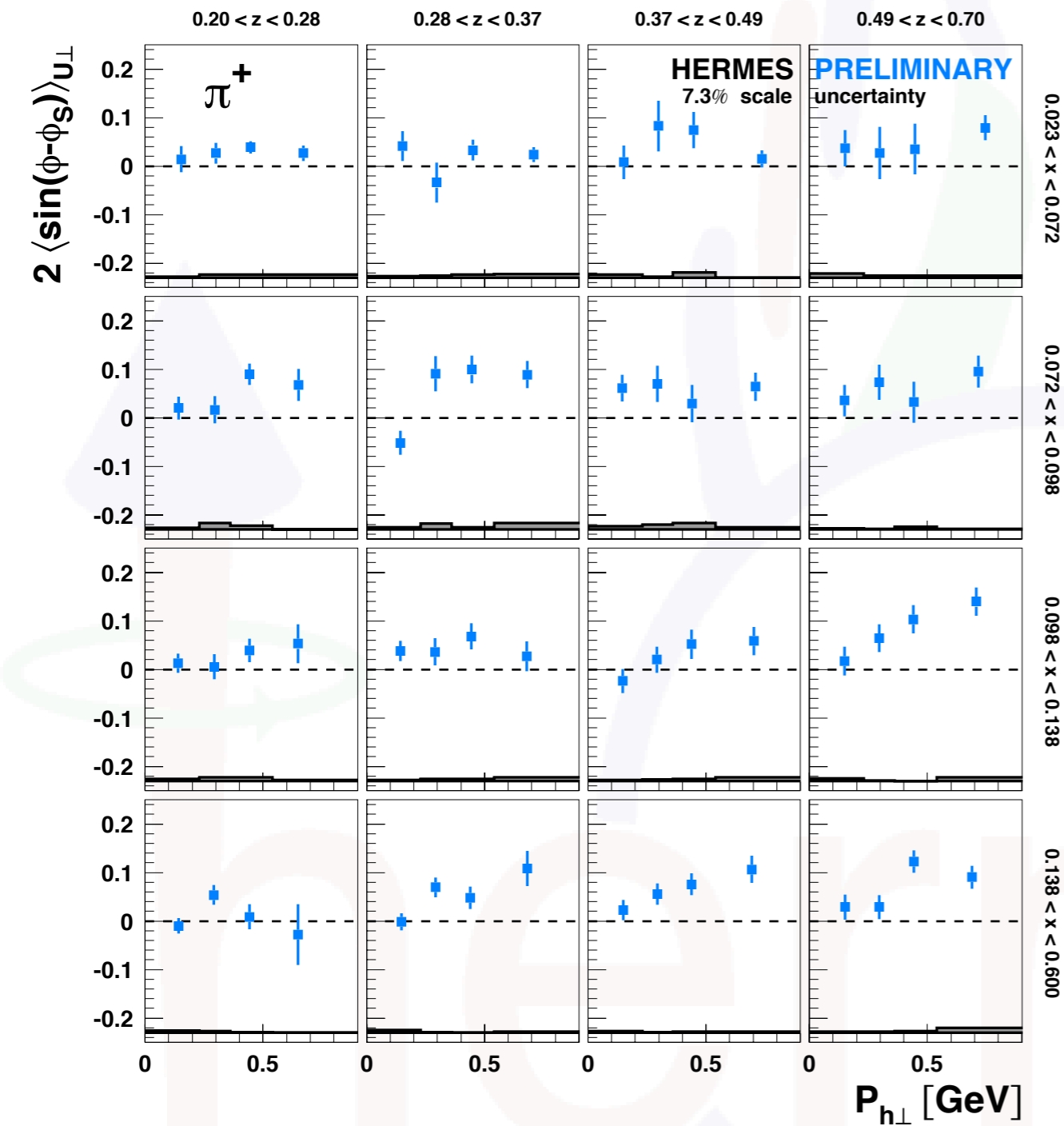
[Airapetian et al., PLB 693 (2010) 11]



similar amplitudes for positive pions and protons  $\rightarrow$  u-quark dominance (and not a FF effect)?

# Sivers amplitudes - 3d binning

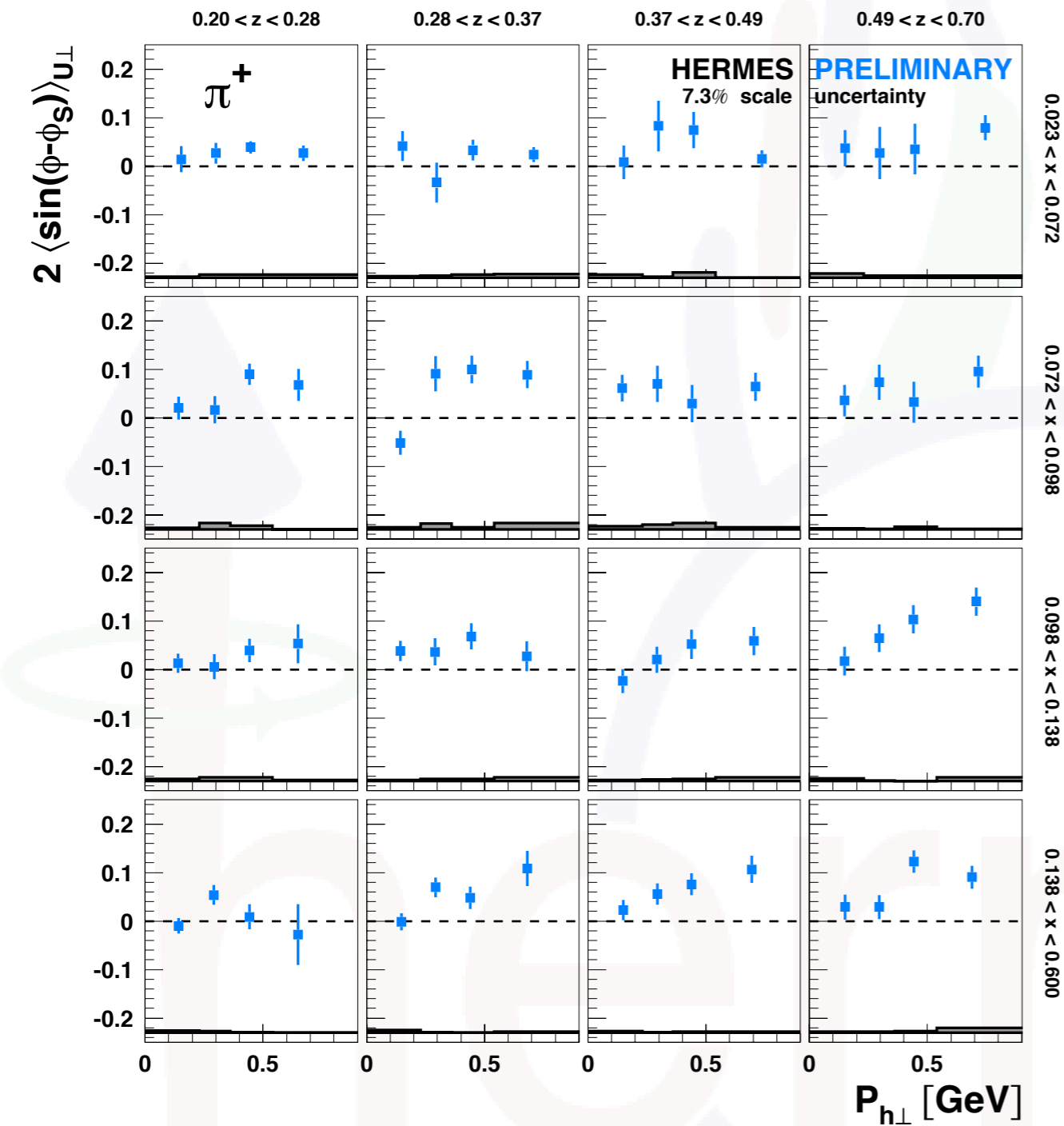
|   |                |          |                     |
|---|----------------|----------|---------------------|
|   | U              | L        | T                   |
| U | $f_1$          |          | $h_1^\perp$         |
| L |                | $g_{1L}$ | $h_{1L}^\perp$      |
| T | $f_{1T}^\perp$ | $g_{1T}$ | $h_1, h_{1T}^\perp$ |



- 3d analysis: 4x4x4 bins in  $(x, z, P_{h_\perp})$

# Sivers amplitudes - 3d binning

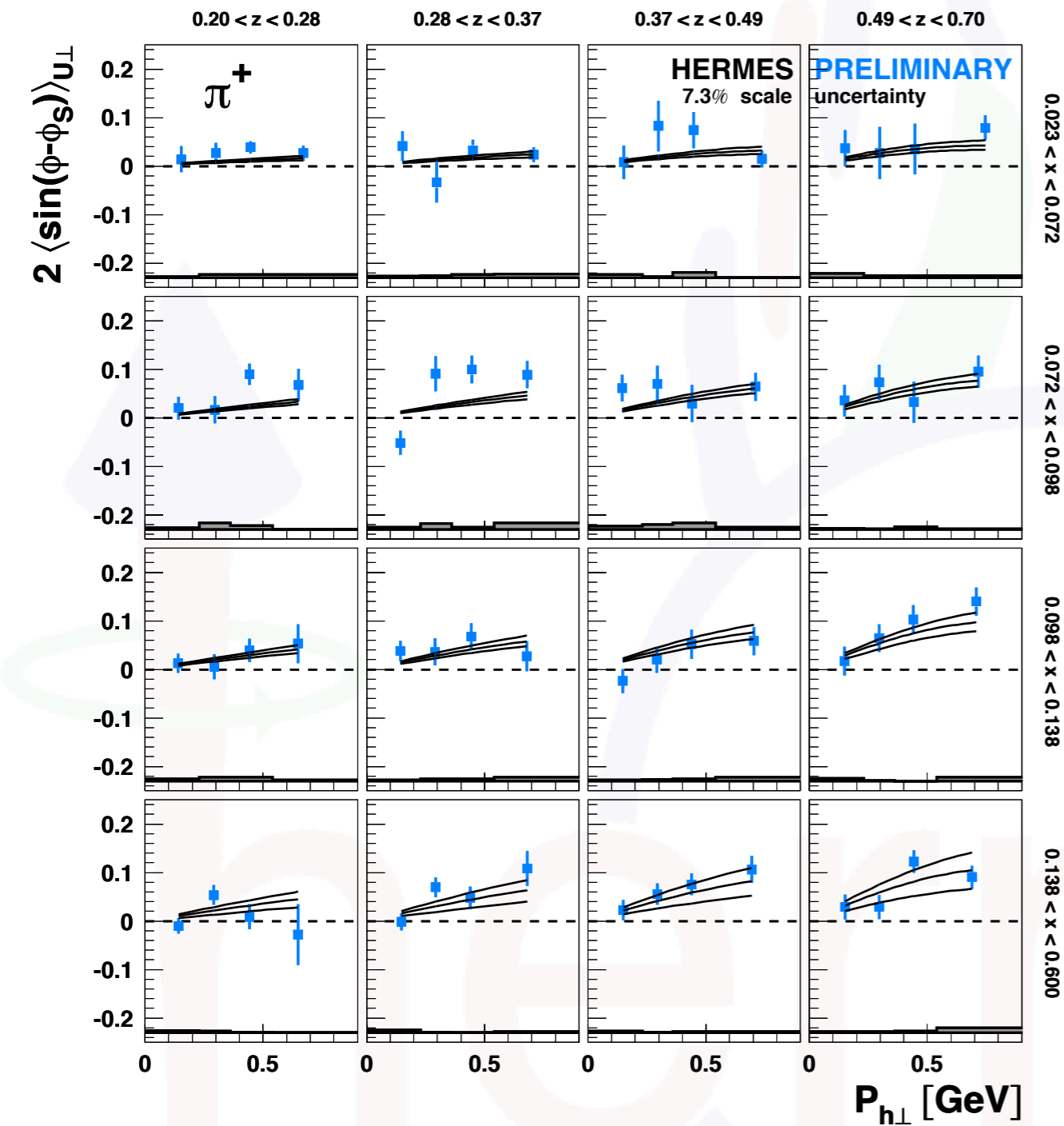
|   |                |          |                     |
|---|----------------|----------|---------------------|
|   | U              | L        | T                   |
| U | $f_1$          |          | $h_1^\perp$         |
| L |                | $g_{1L}$ | $h_{1L}^\perp$      |
| T | $f_{1T}^\perp$ | $g_{1T}$ | $h_1, h_{1T}^\perp$ |



- 3d analysis: 4x4x4 bins in  $(x, z, P_{h_\perp})$
- disentangle correlations
- isolate phase-space region with strong signal strength

# Sivers amplitudes - 3d binning

|   |                |          |                     |
|---|----------------|----------|---------------------|
|   | U              | L        | T                   |
| U | $f_1$          |          | $h_1^\perp$         |
| L |                | $g_{1L}$ | $h_{1L}^\perp$      |
| T | $f_{1T}^\perp$ | $g_{1T}$ | $h_1, h_{1T}^\perp$ |

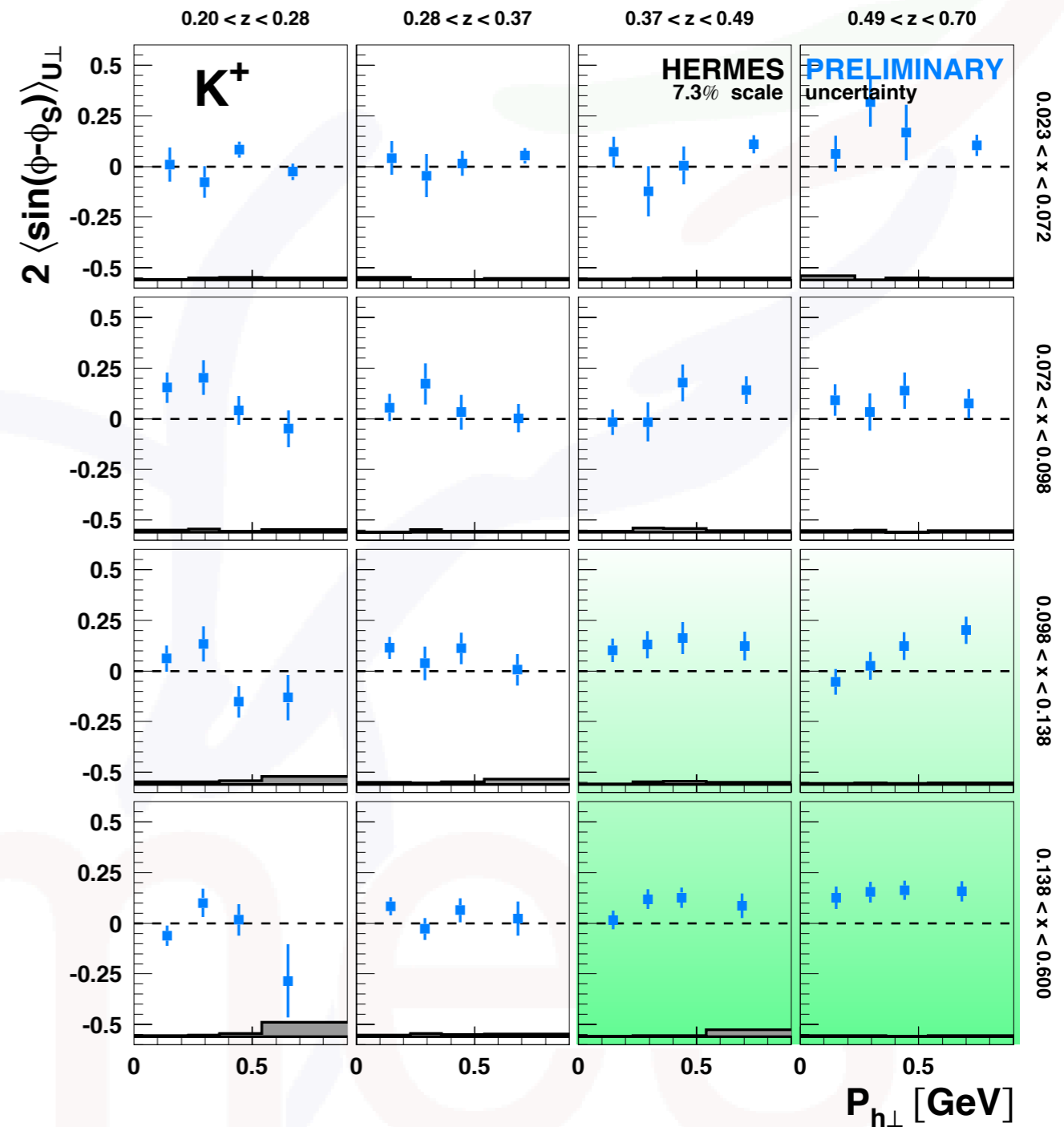


- 3d analysis: 4x4x4 bins in  $(x, z, P_{h\perp})$
- disentangle correlations
- isolate phase-space region with strong signal strength
- allows more detailed comparison with calculations (e.g., [10.1103/PhysRevD.86.014028](https://arxiv.org/abs/10.1103/PhysRevD.86.014028) fit - courtesy M. Boglione)

# Sivers amplitudes - 3d binning

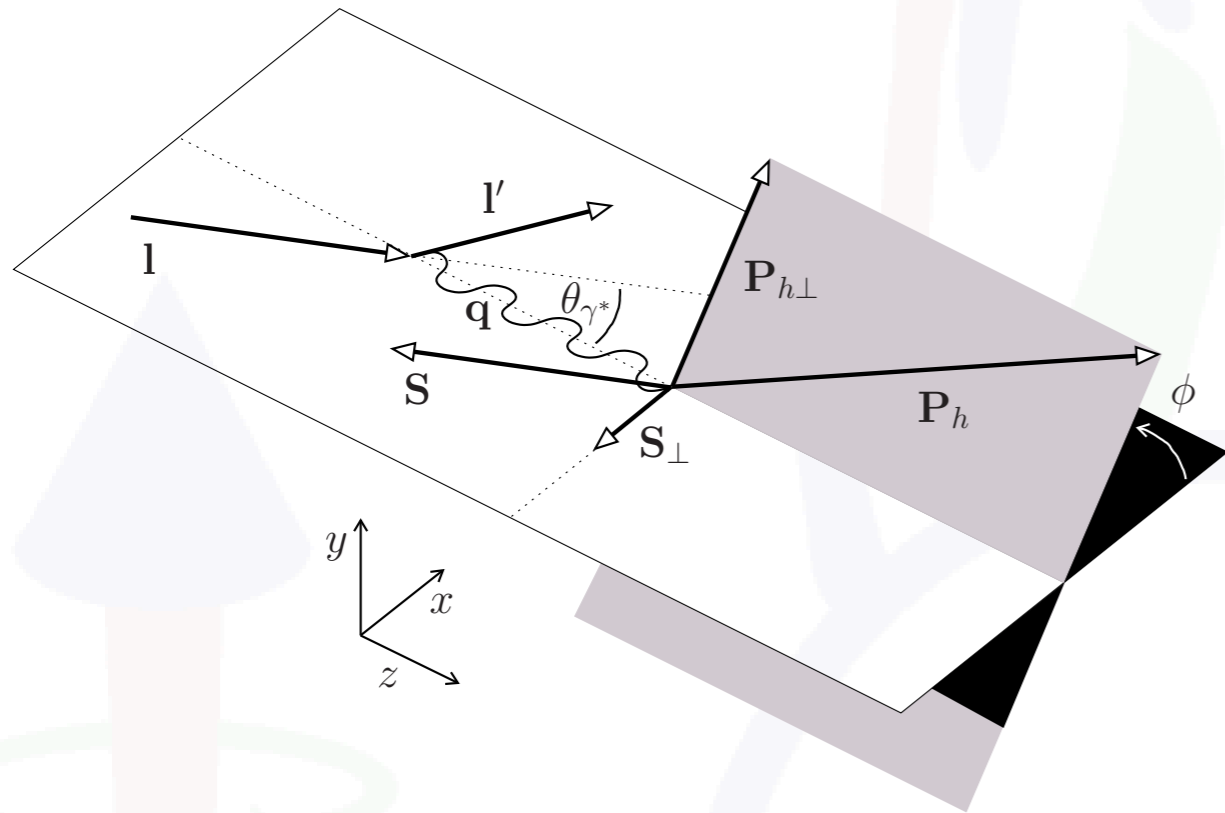
|   |                |          |                     |
|---|----------------|----------|---------------------|
|   | U              | L        | T                   |
| U | $f_1$          |          | $h_1^\perp$         |
| L |                | $g_{1L}$ | $h_{1L}^\perp$      |
| T | $f_{1T}^\perp$ | $g_{1T}$ | $h_1, h_{1T}^\perp$ |

- large  $K^+$  amplitudes  $O(20\%)$  seen at large values of  $(x, z)$
- region of purest "u-quark probe"



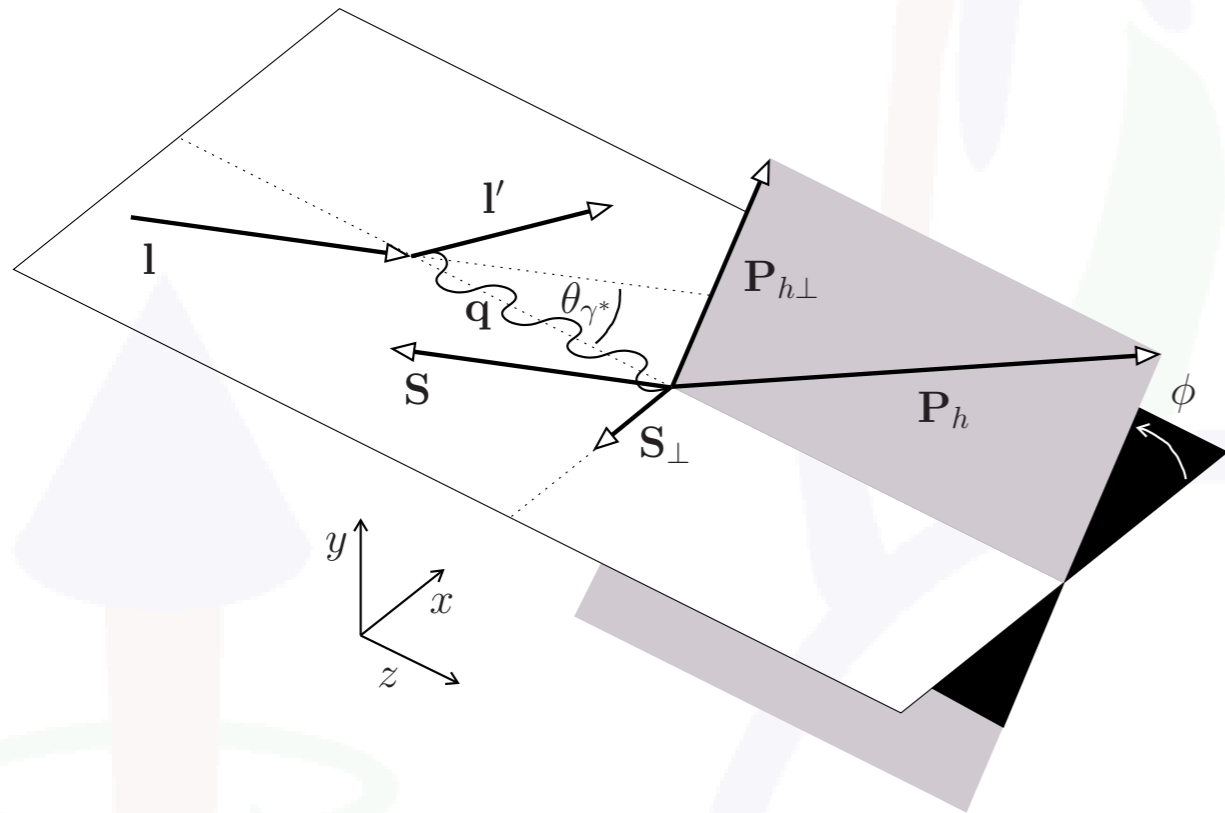
subleading twist

# Subleading twist I - $\langle \sin(\phi) \rangle_{UL}$



- in experiments: target polarized w.r.t. beam direction  
[Diehl&Sapeta EPJC41 (2005)]
- small transverse component w.r.t. virtual-photon direction when longitudinally polarized
- mixing of transverse and longitudinal target-spin asymmetries

# Subleading twist I - $\langle \sin(\phi) \rangle_{UL}$



- in experiments: target polarized w.r.t. beam direction  
[Diehl&Sapeta EPJC41 (2005)]
- small transverse component w.r.t. virtual-photon direction when longitudinally polarized
- mixing of transverse and longitudinal target-spin asymmetries

$$\begin{pmatrix} \langle \sin \phi \rangle_{UL}^I \\ \langle \sin(\phi - \phi_S) \rangle_{UT}^I \\ \langle \sin(\phi + \phi_S) \rangle_{UT}^I \end{pmatrix} = \begin{pmatrix} \cos \theta_{\gamma^*} & -\sin \theta_{\gamma^*} & -\sin \theta_{\gamma^*} \\ \frac{1}{2} \sin \theta_{\gamma^*} & \cos \theta_{\gamma^*} & 0 \\ \frac{1}{2} \sin \theta_{\gamma^*} & 0 & \cos \theta_{\gamma^*} \end{pmatrix} \begin{pmatrix} \langle \sin \phi \rangle_{UL}^q \\ \langle \sin(\phi - \phi_S) \rangle_{UT} \\ \langle \sin(\phi + \phi_S) \rangle_{UT} \end{pmatrix}$$

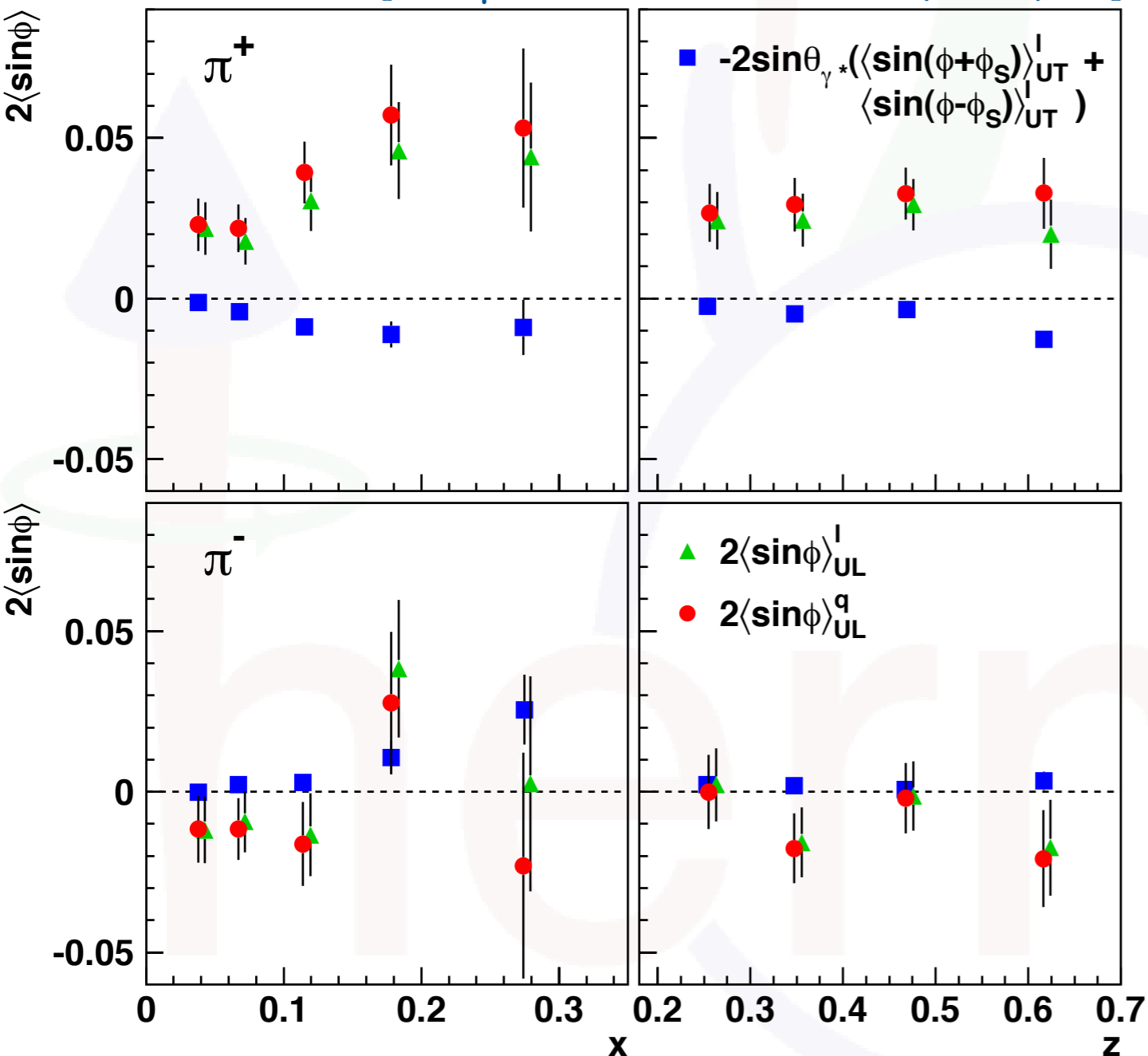
( $\cos \theta_{\gamma^*} \simeq 1$ ,  $\sin \theta_{\gamma^*}$  up to 15% at HERMES energies)



# Subleading twist I - $\langle \sin(\phi) \rangle_{UL}$

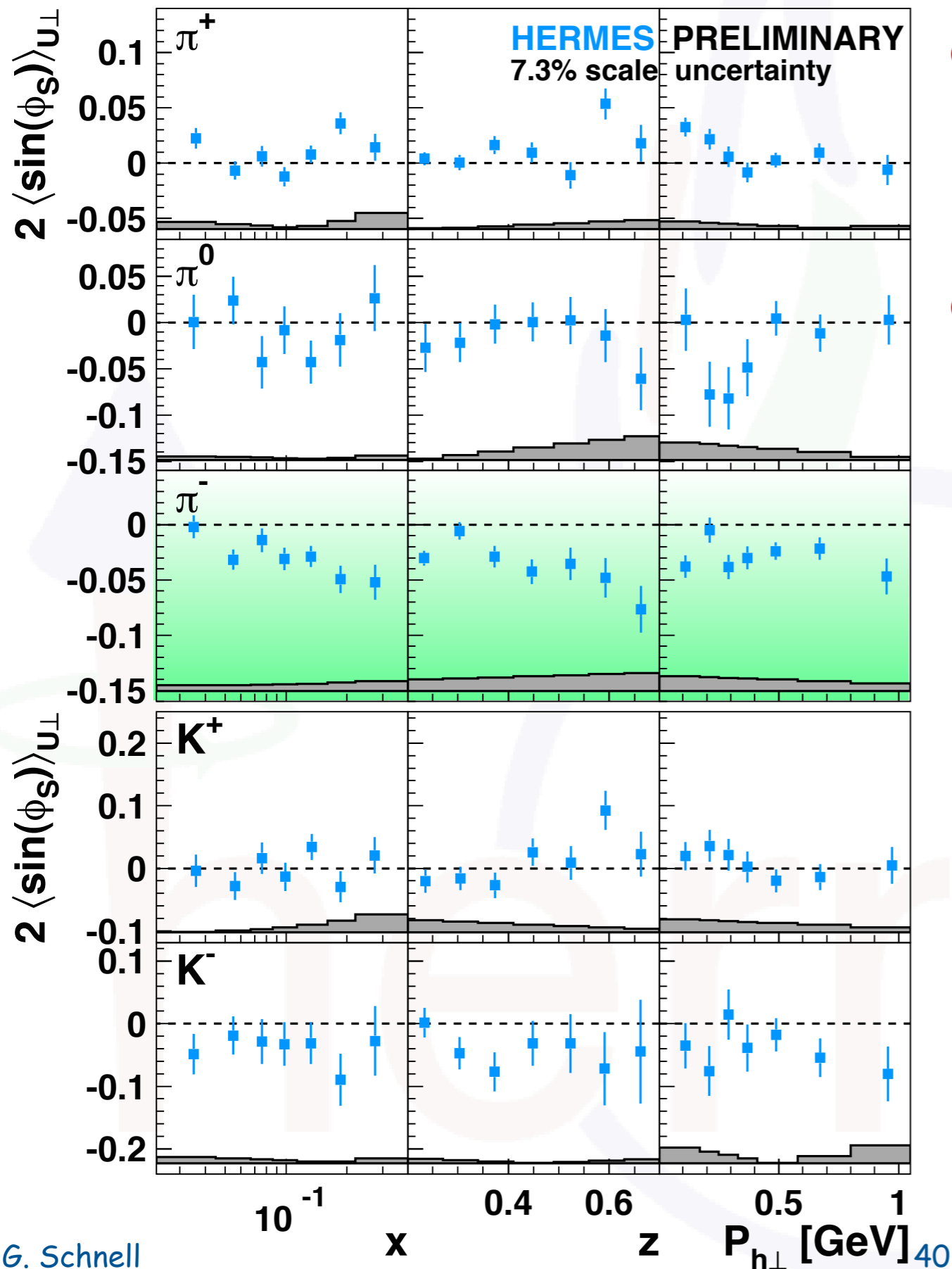
$$\langle \sin \phi \rangle_{UL}^q = \langle \sin \phi \rangle_{UL}^I + \sin \theta_{\gamma^*} \left( \langle \sin(\phi + \phi_S) \rangle_{UT}^I + \langle \sin(\phi - \phi_S) \rangle_{UT}^I \right)$$

[Airapetian et al., PLB 622 (2005) 14]



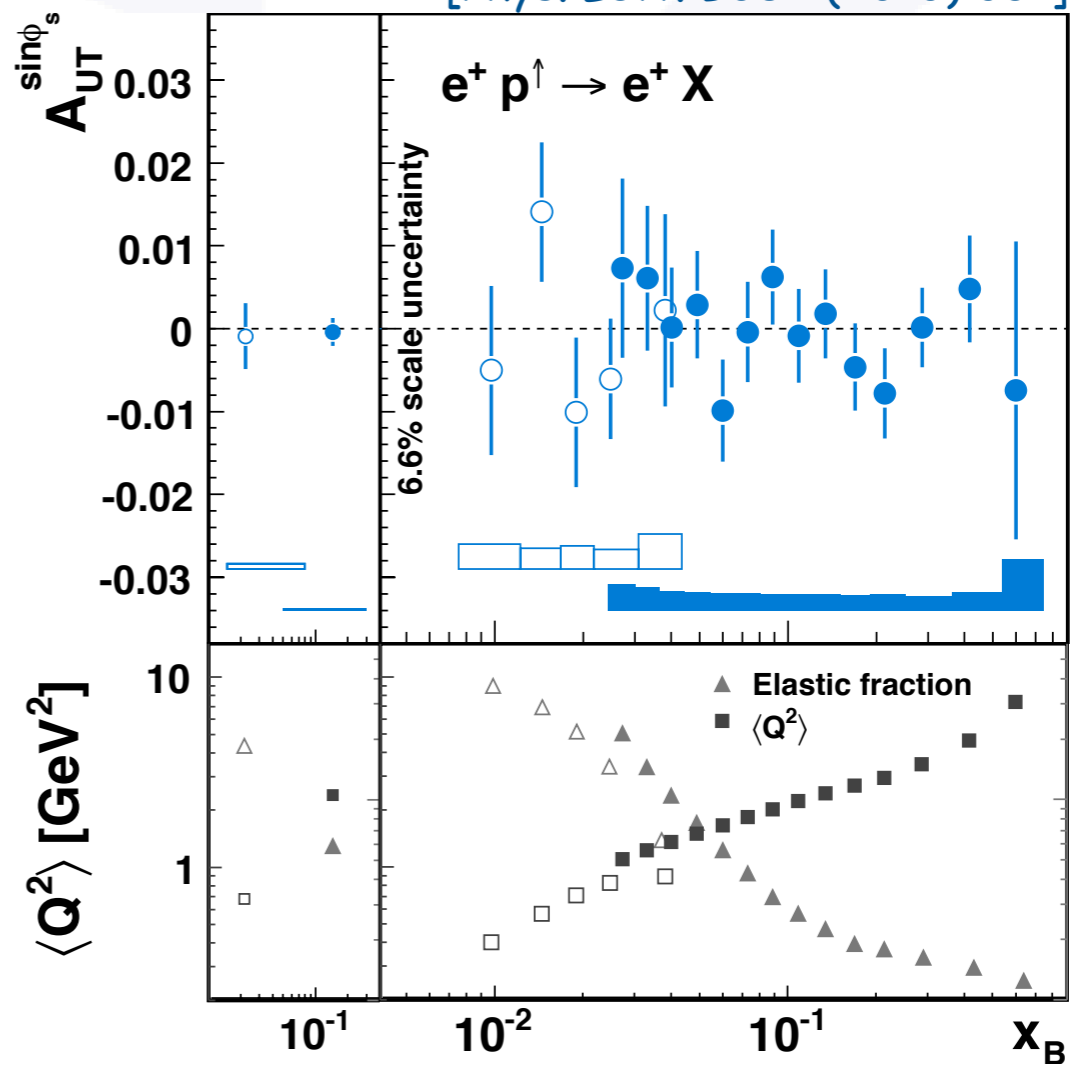
- experimental  $A_{UL}$  dominated by twist-3 contribution
- correction for  $A_{UT}$  contribution increases purely longitudinal asymmetry for positive pions
- consistent with zero for  $\pi^-$

# Subleading twist II - $\langle \sin(\phi_s) \rangle_{UT}$

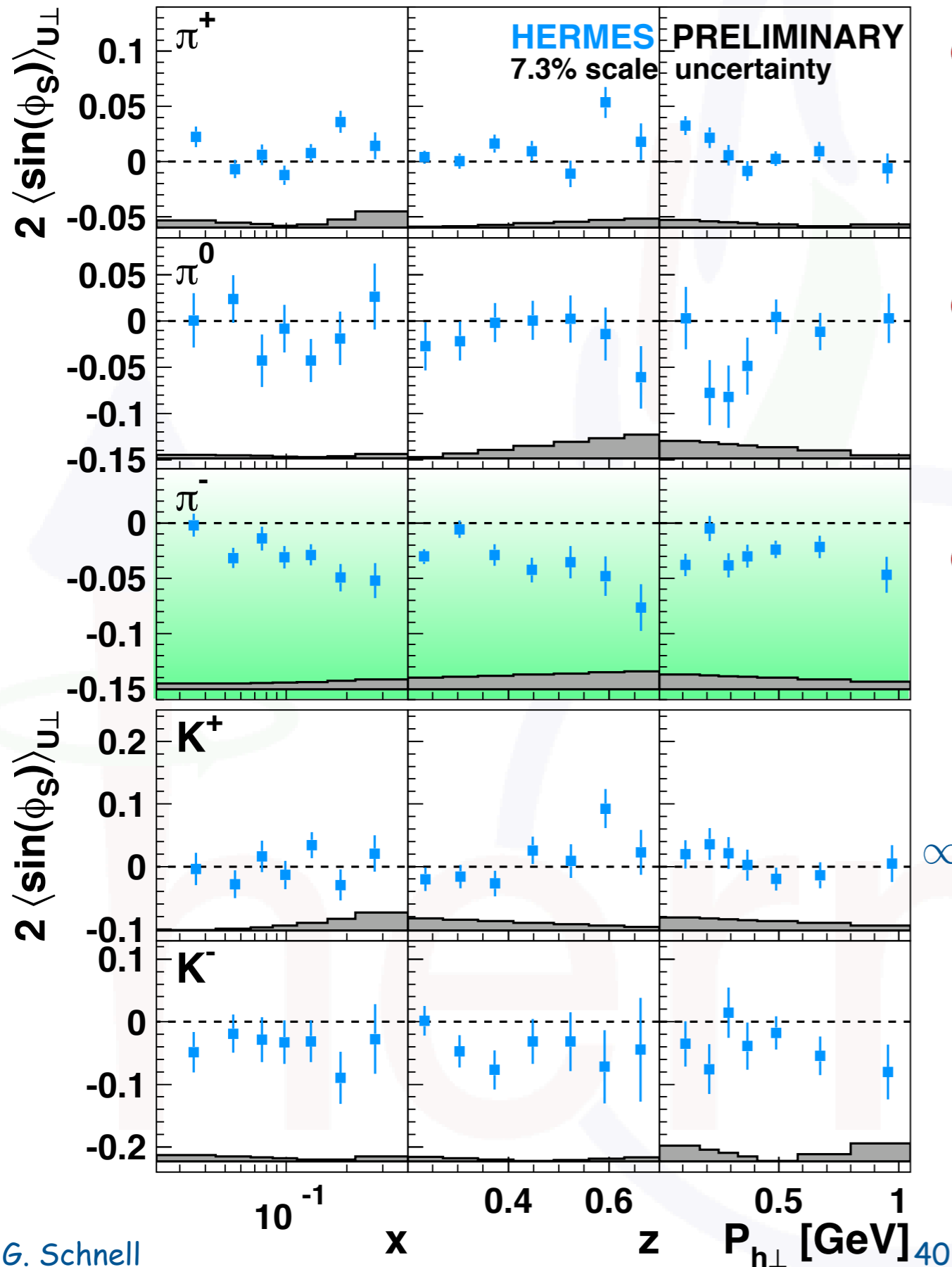


- significant non-zero signal observed for negatively charged mesons
- vanishes in inclusive limit, e.g. after integration over  $P_{h\perp}$  and  $z$ , and summation over all hadrons

[Phys. Lett. B682 (2010) 351]



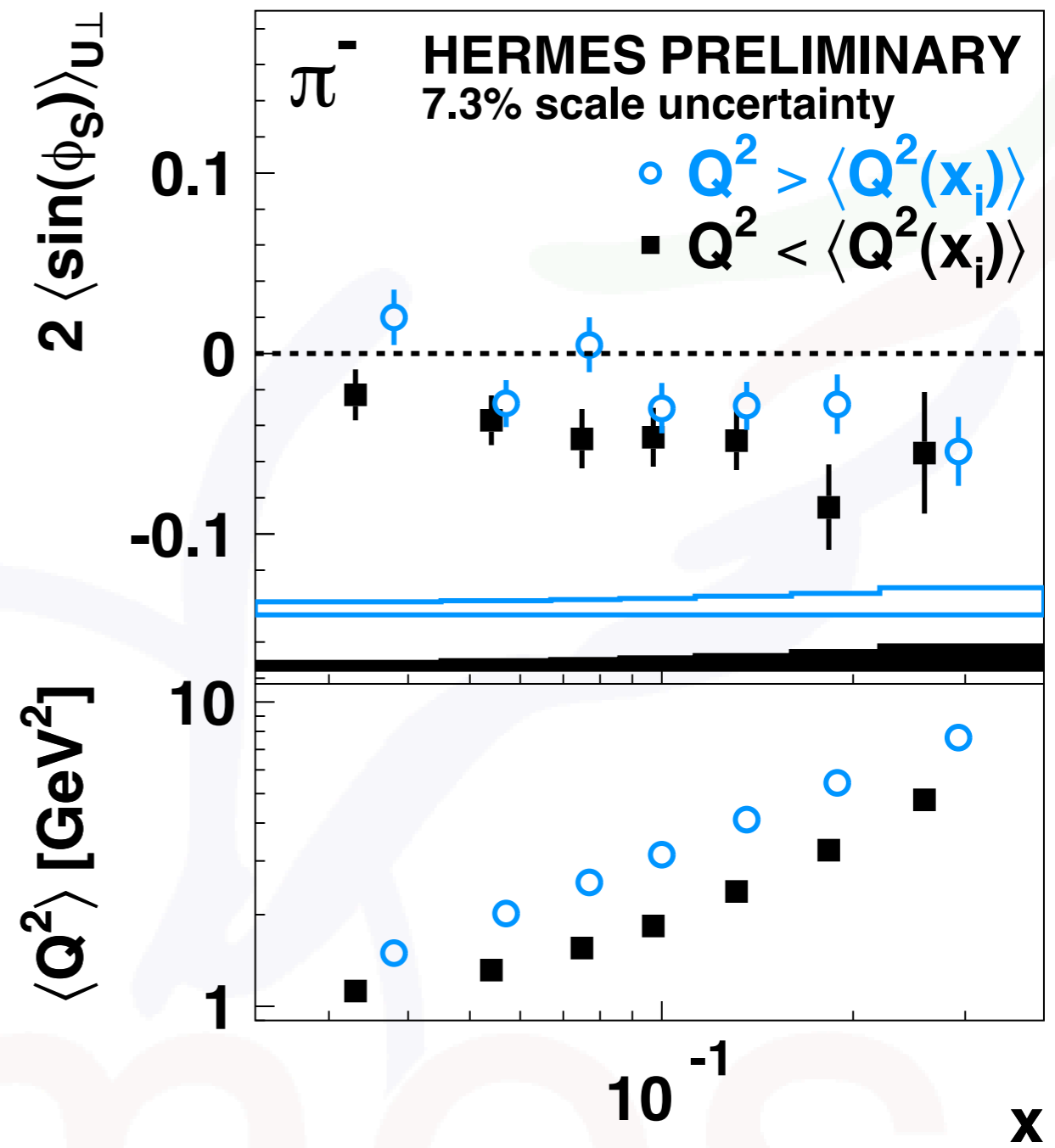
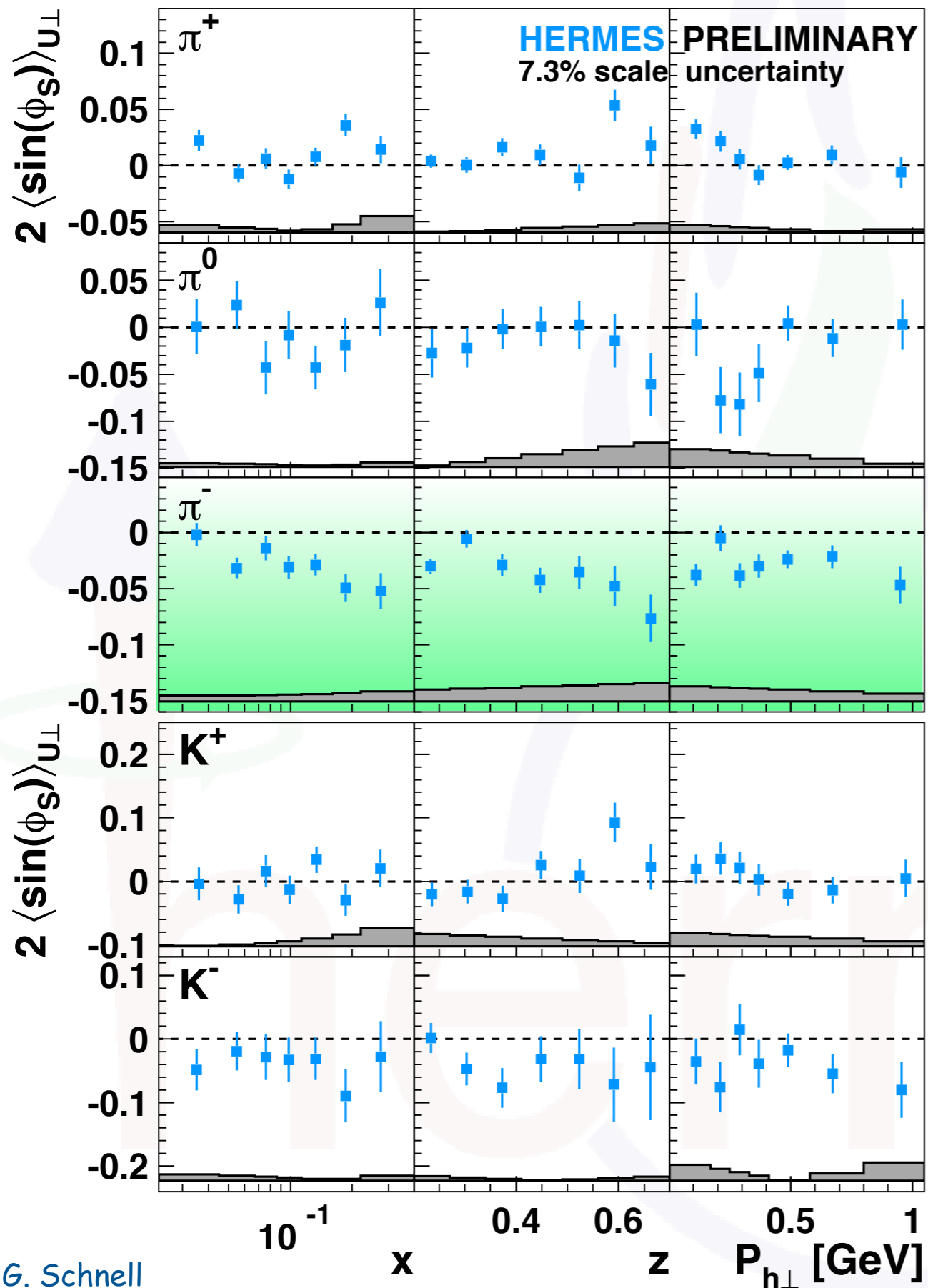
# Subleading twist II - $\langle \sin(\phi_s) \rangle_{U\perp}$



- significant non-zero signal observed for negatively charged mesons
- vanishes in inclusive limit, e.g. after integration over  $P_{h\perp}$  and  $z$ , and summation over all hadrons
- various terms related to transversity, worm-gear, Sivers etc.:

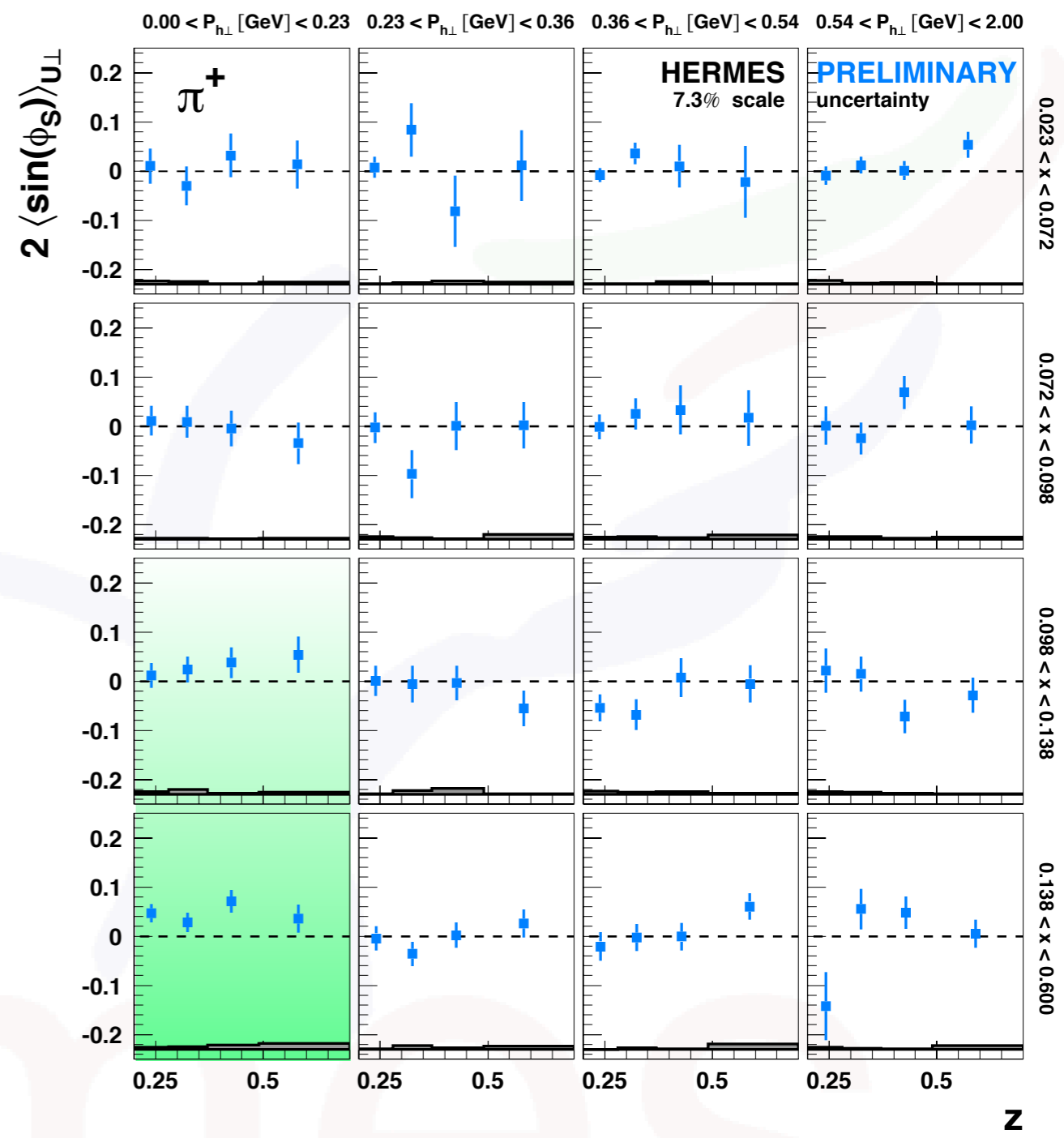
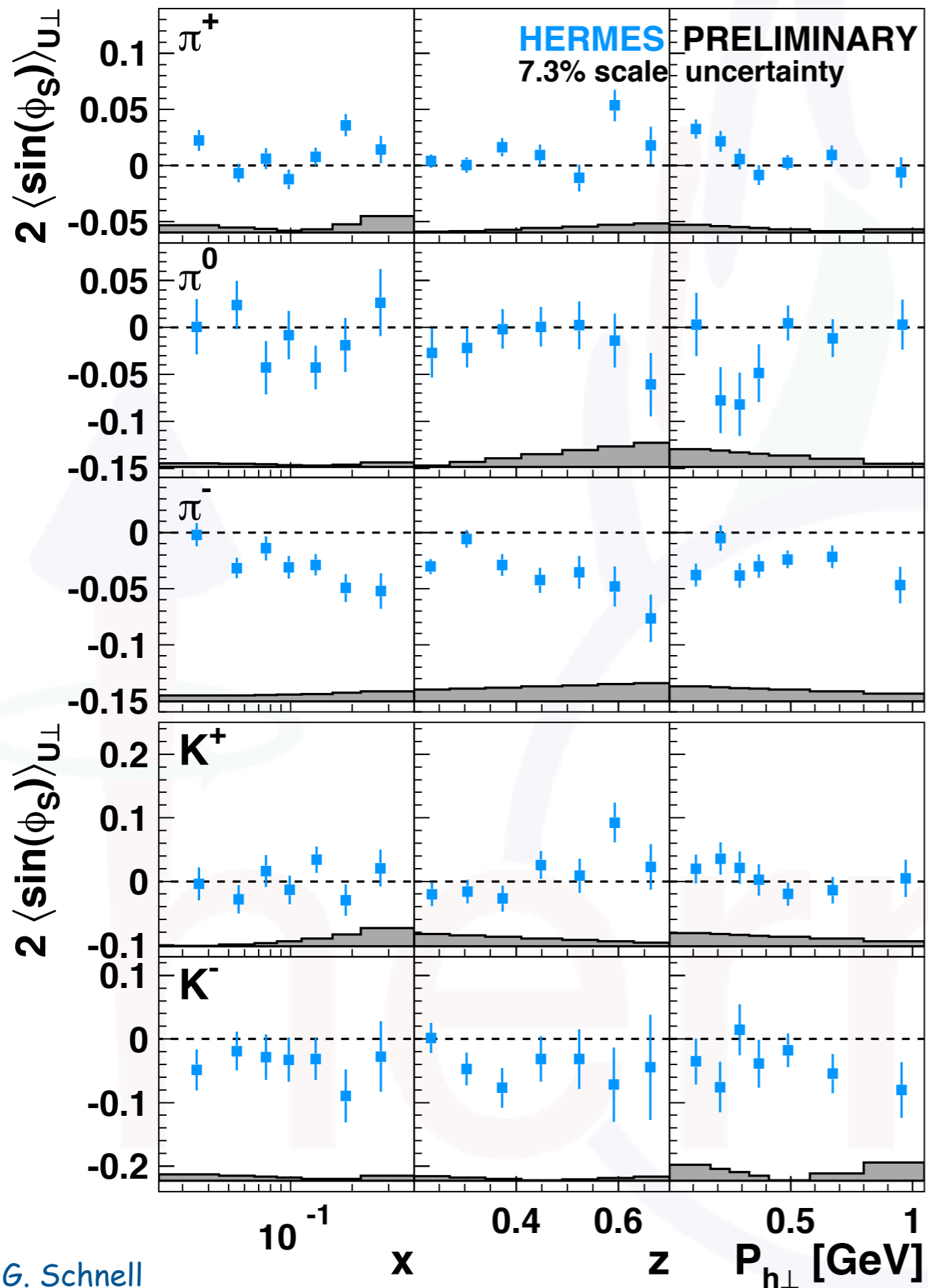
$$\propto \left( x f_{T1}^{\perp} D_1 - \frac{M_h}{M} h_1 \frac{\tilde{H}}{z} \right) - \mathcal{W}(p_T, k_T, P_{h\perp}) \left[ \left( x h_T H_1^{\perp} + \frac{M_h}{M} g_{1T} \frac{\tilde{G}^{\perp}}{z} \right) - \left( x h_T^{\perp} H_1^{\perp} - \frac{M_h}{M} f_{1T}^{\perp} \frac{\tilde{D}^{\perp}}{z} \right) \right]$$

# Subleading twist II - $\langle \sin(\phi_S) \rangle_{UT}$



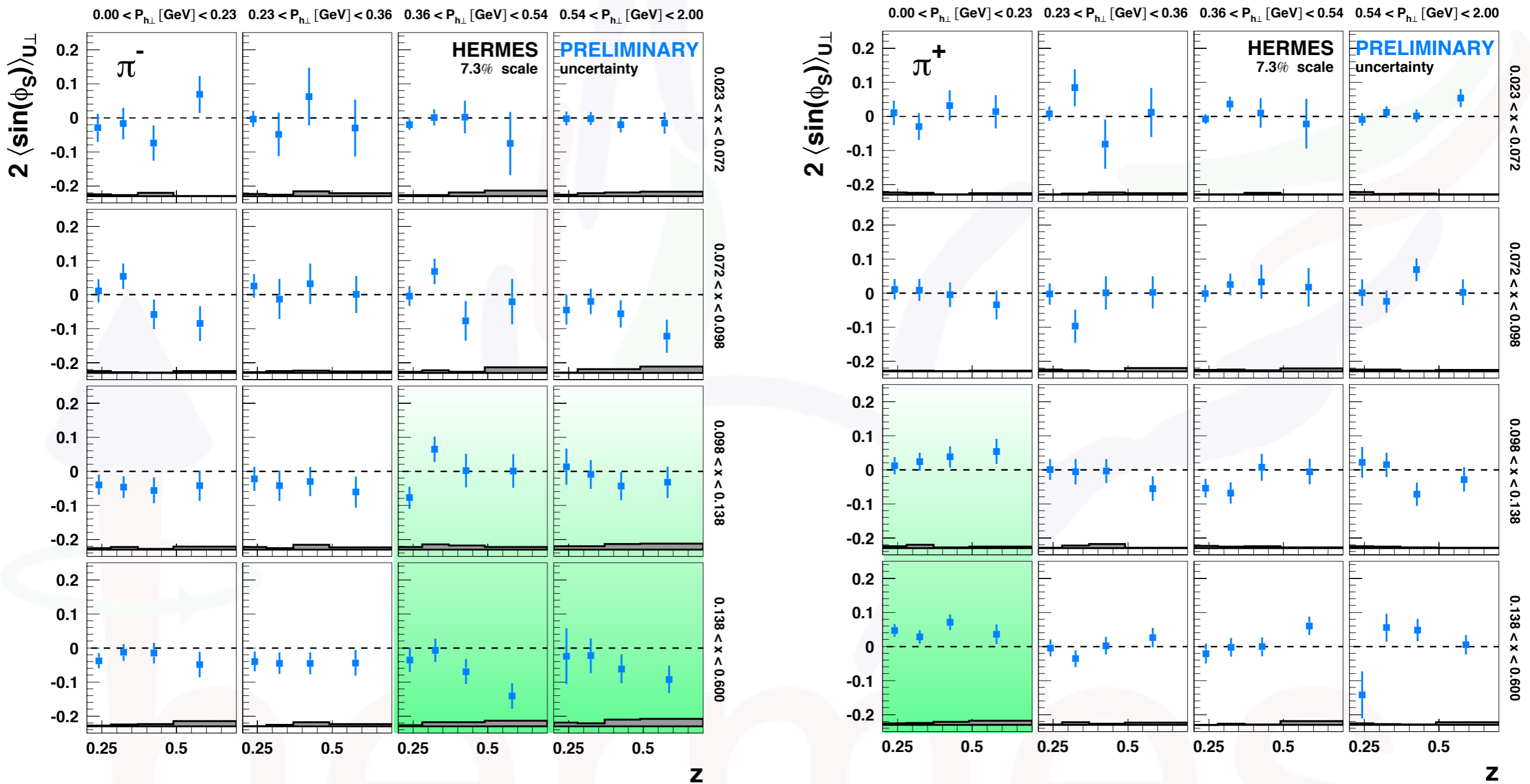
● hint of  $Q^2$  dependence seen in signal for negative pions

# Subleading twist II - $\langle \sin(\phi_S) \rangle_{UT}$



● positive amplitudes at low  $P_{h\perp}$   
also for positive pions

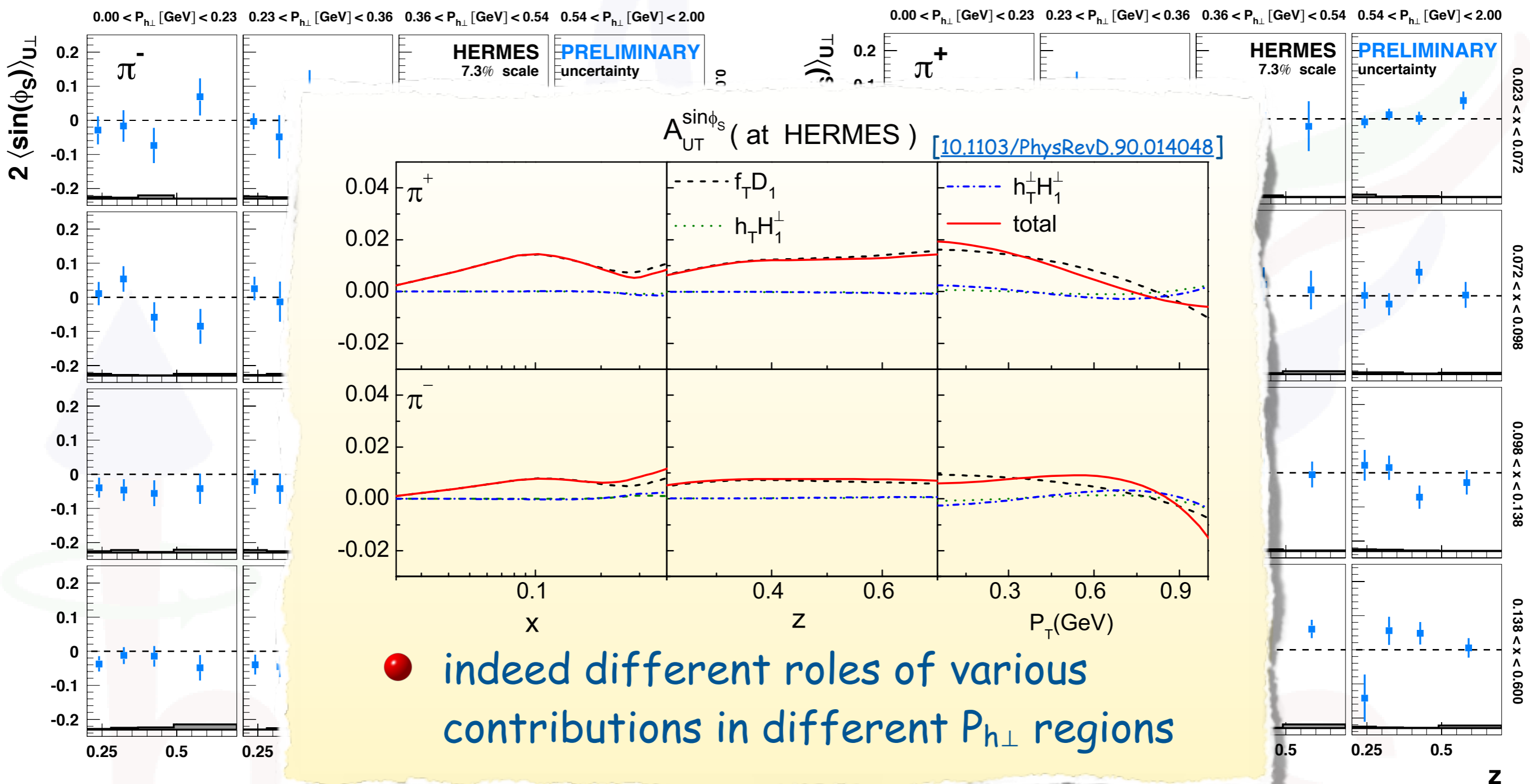
# Subleading twist II - $\langle \sin(\phi_S) \rangle_{UT}$



● nonzero amplitudes mainly at large  $P_{h\perp}$  in case of negative pions

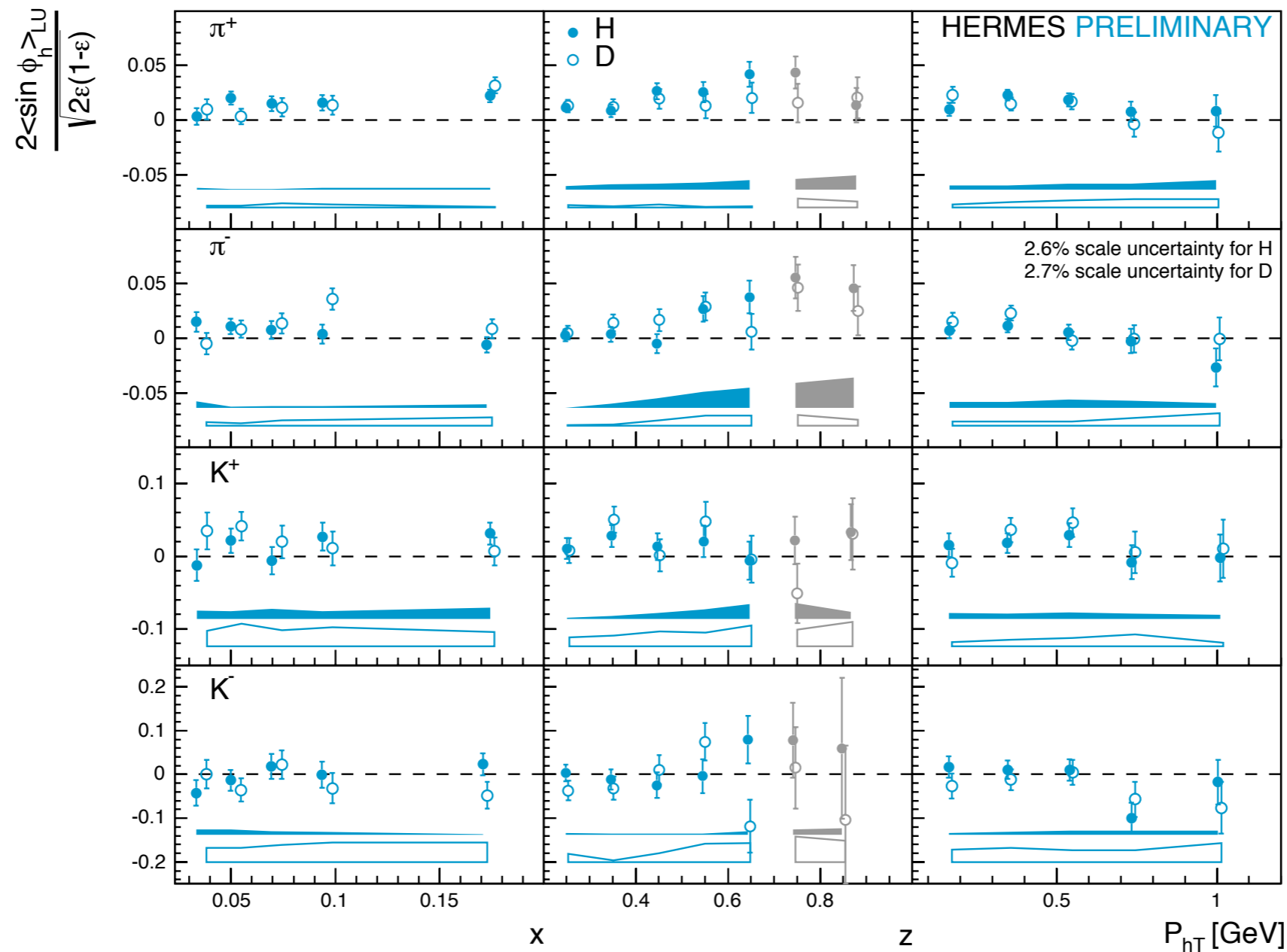
● positive amplitudes at low  $P_{h\perp}$  also for positive pions

# Subleading twist II - $\langle \sin(\phi_S) \rangle_{UT}$



# Subleading twist III - $\langle \sin(\phi) \rangle_{LU}$

$$\frac{M_h}{Mz} h_1^\perp E \oplus xg^\perp D_1 \oplus \frac{M_h}{Mz} f_1 G^\perp \oplus \text{red } eH_1^\perp$$

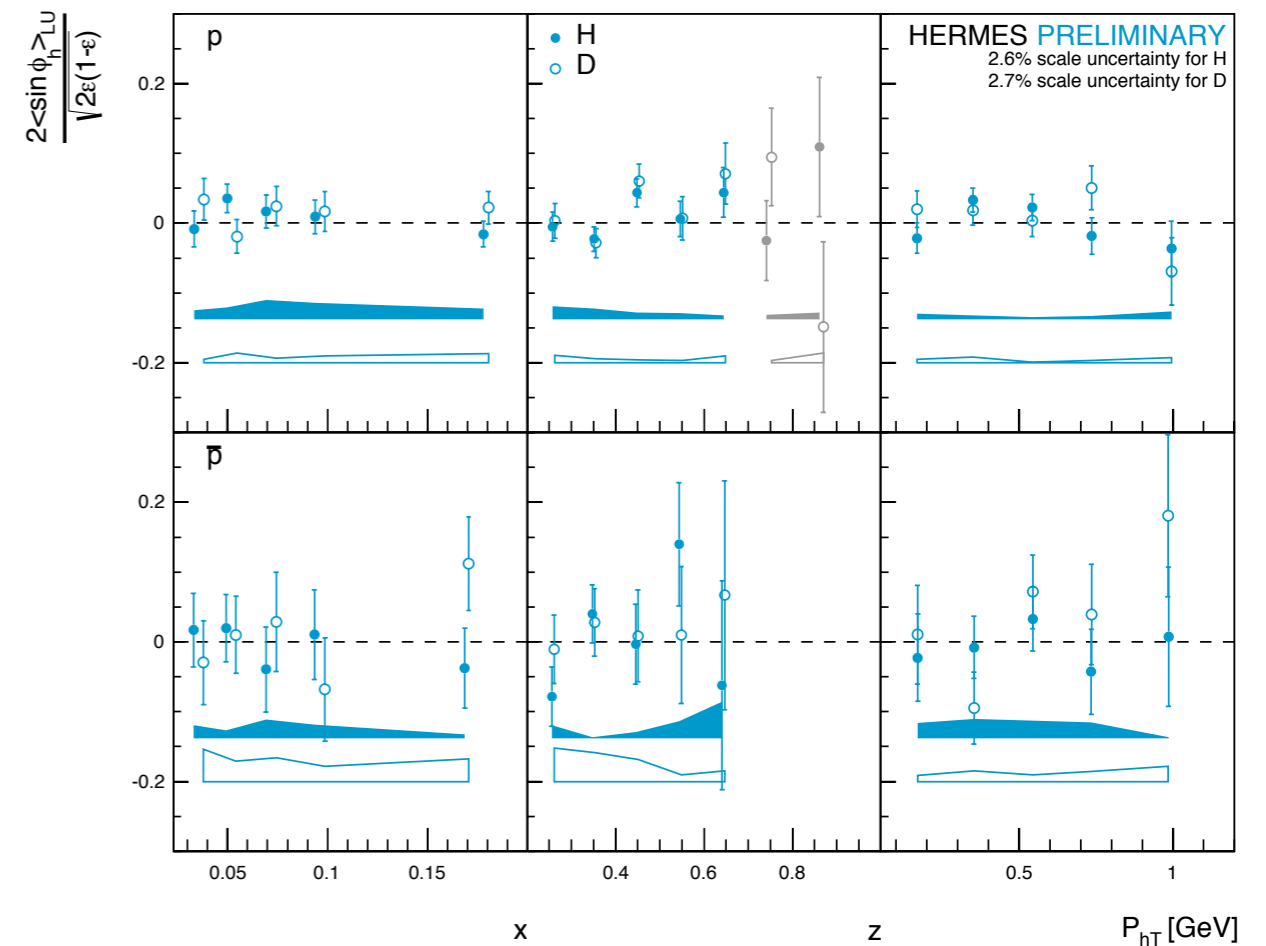
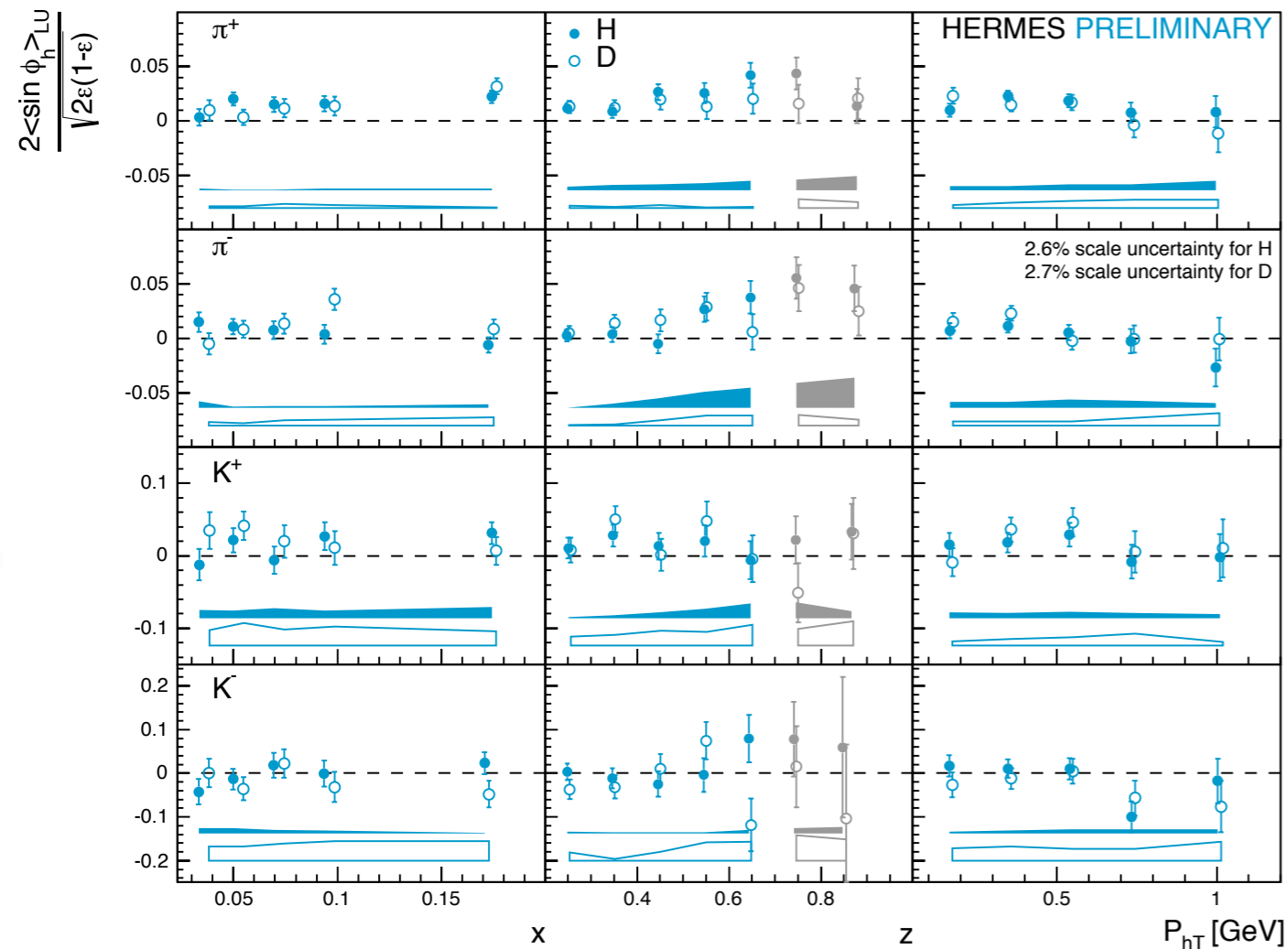


- significant positive amplitudes for (in particular positive) pions



# Subleading twist III - $\langle \sin(\phi) \rangle_{LU}$

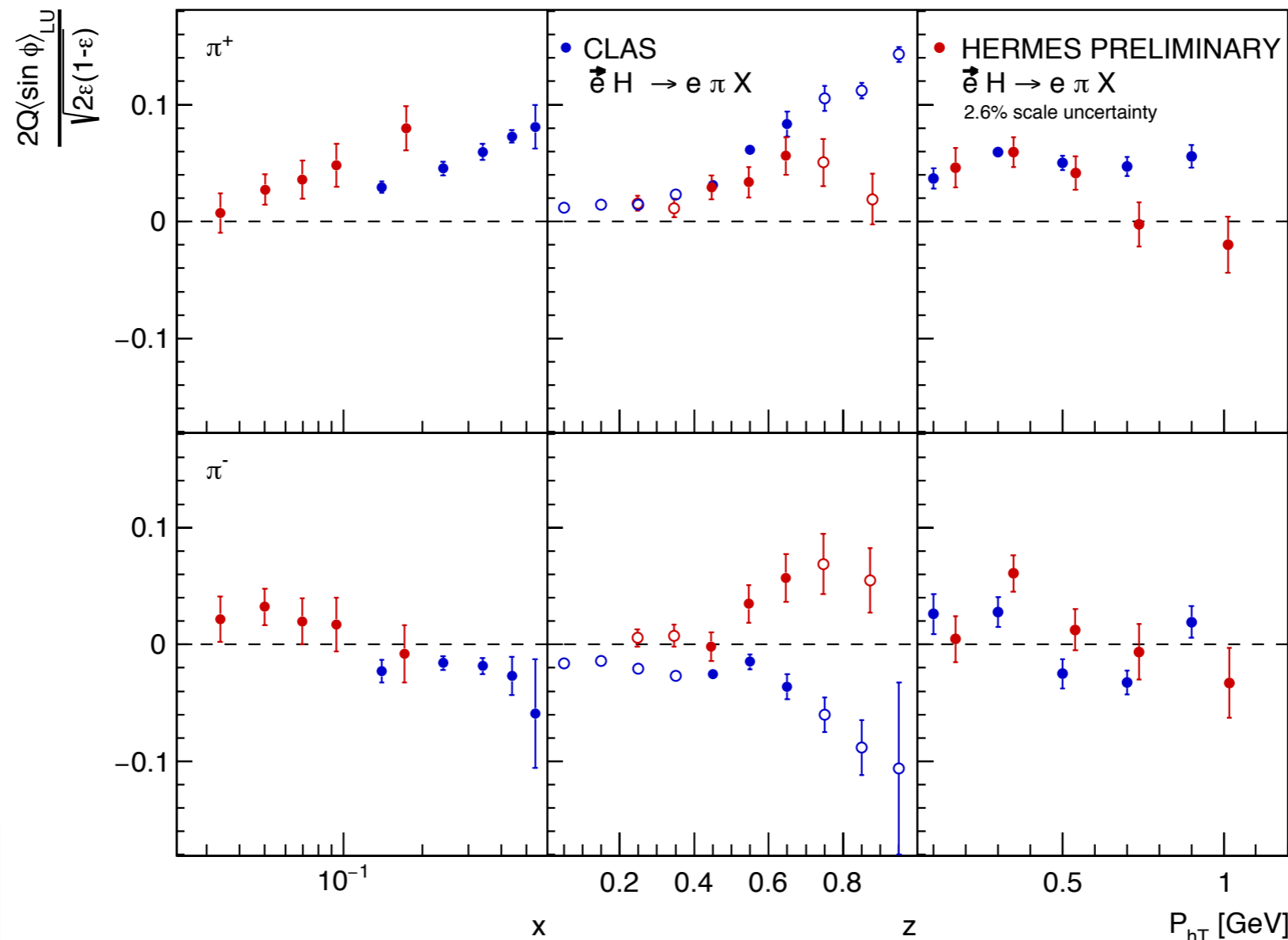
$$\frac{M_h}{M_z} h_1^\perp E \oplus x g^\perp D_1 \oplus \frac{M_h}{M_z} f_1 G^\perp \oplus x e H_1^\perp$$



- mostly consistent w/ zero for other hadrons (except maybe  $K^+$ )

# Subleading twist III - $\langle \sin(\phi) \rangle_{LU}$

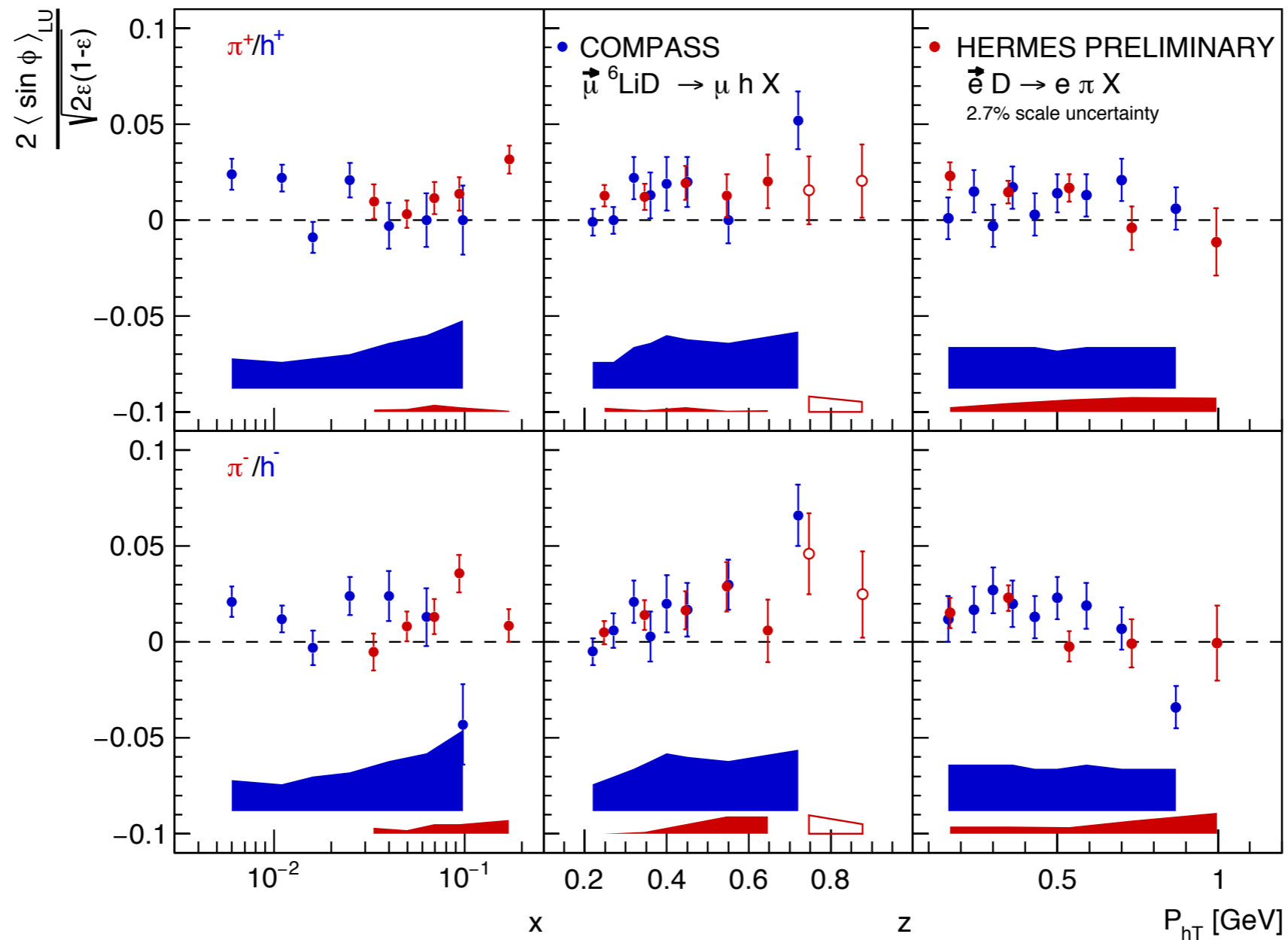
$$\frac{M_h}{M_z} h_1^\perp E \oplus x g^\perp D_1 \oplus \frac{M_h}{M_z} f_1 G^\perp \oplus x e H_1^\perp$$



- opposite behavior at HERMES/CLAS of negative pions in z projection due to different x-range probed
- CLAS more sensitive to  $e(x)$ Collins term due to higher x probed?

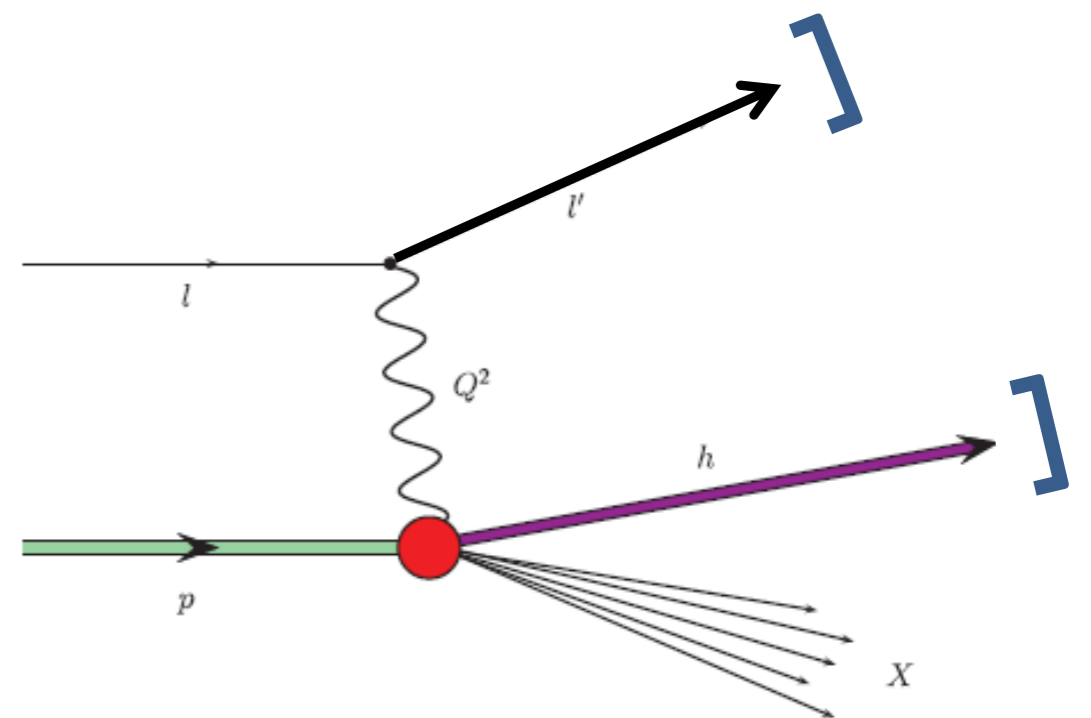
# Subleading twist III - $\langle \sin(\phi) \rangle_{LU}$

$$\frac{M_h}{M_z} h_1^\perp E \oplus x g^\perp D_1 \oplus \frac{M_h}{M_z} f_1 G^\perp \oplus x e H_1^\perp$$

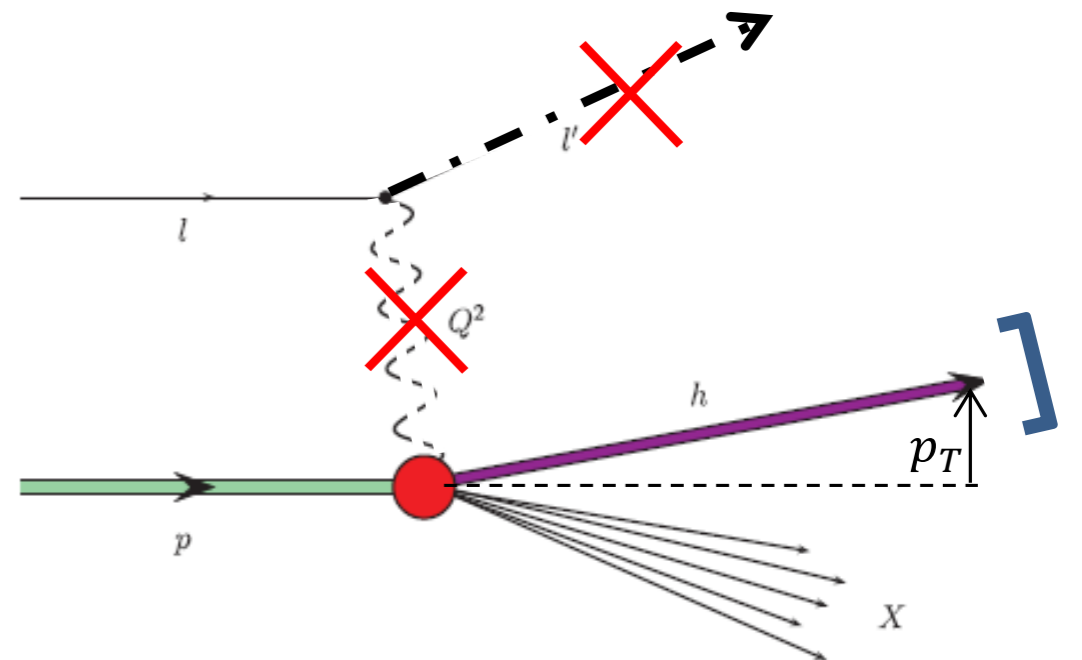


- consistent behavior for charged pions / hadrons at HERMES / COMPASS for isoscalar targets

# Semi-inclusive hadrons



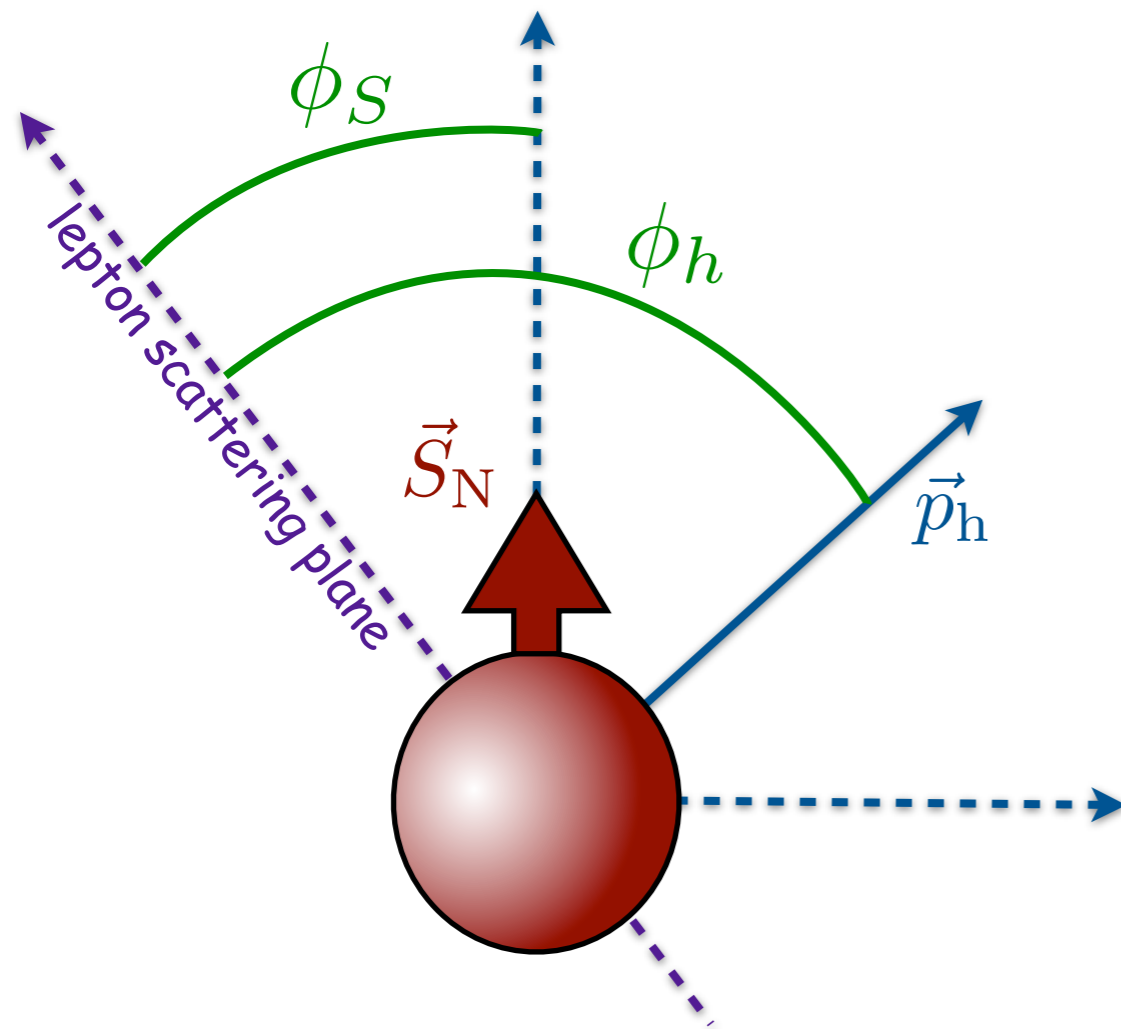
# ~~S~~emi-inclusive hadrons



[click here if \(likely\) out of time](#)

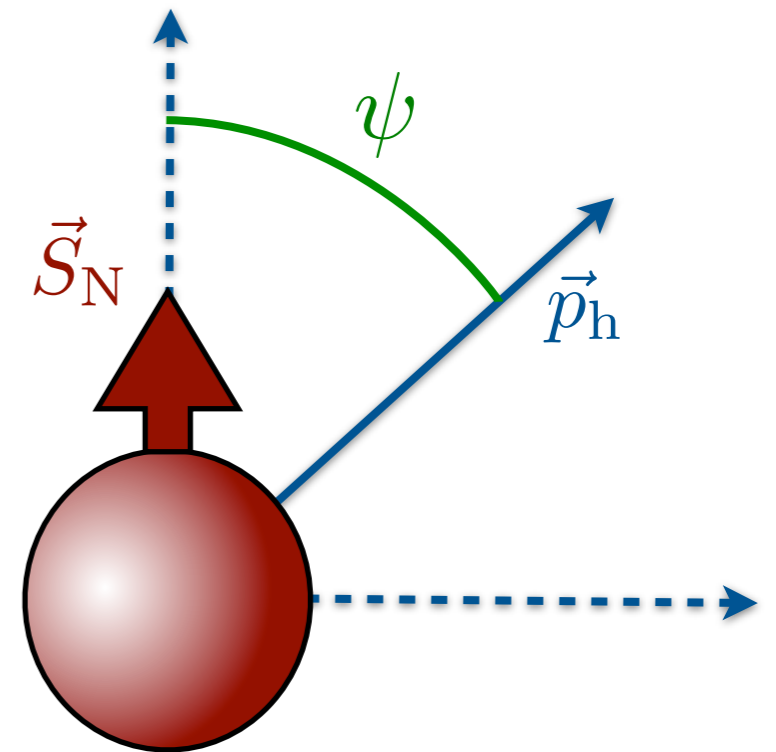
# Inclusive hadron electro-production

$$ep^{\uparrow} \rightarrow ehX$$



virtual photon going  
into the page

$$ep^{\uparrow} \rightarrow hX$$



lepton beam going  
into the page

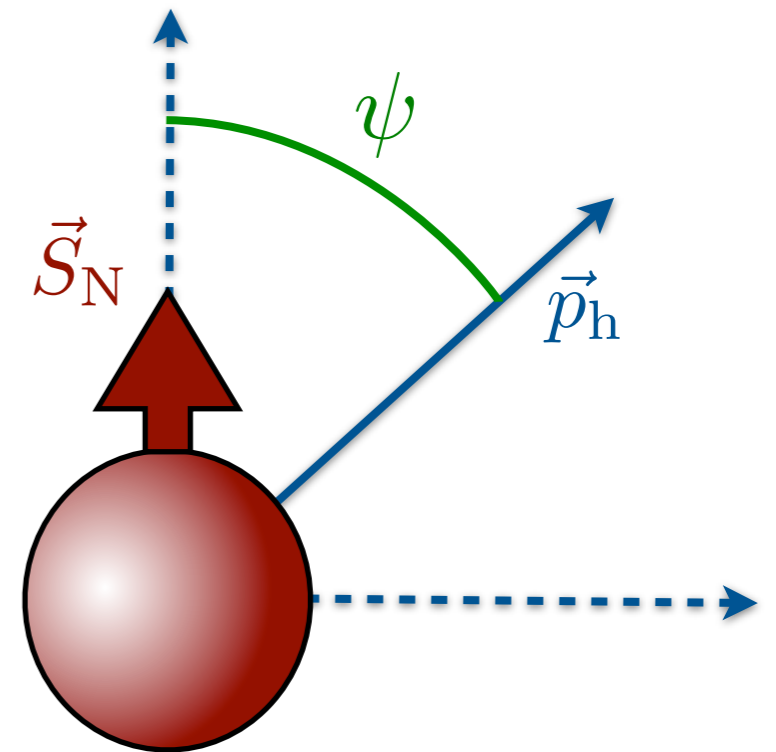
$$\psi \simeq \phi_h - \phi_s$$

→ "Sivers angle"

# Inclusive hadron electro-production

- scattered lepton undetected  
↳ lepton kinematics unknown

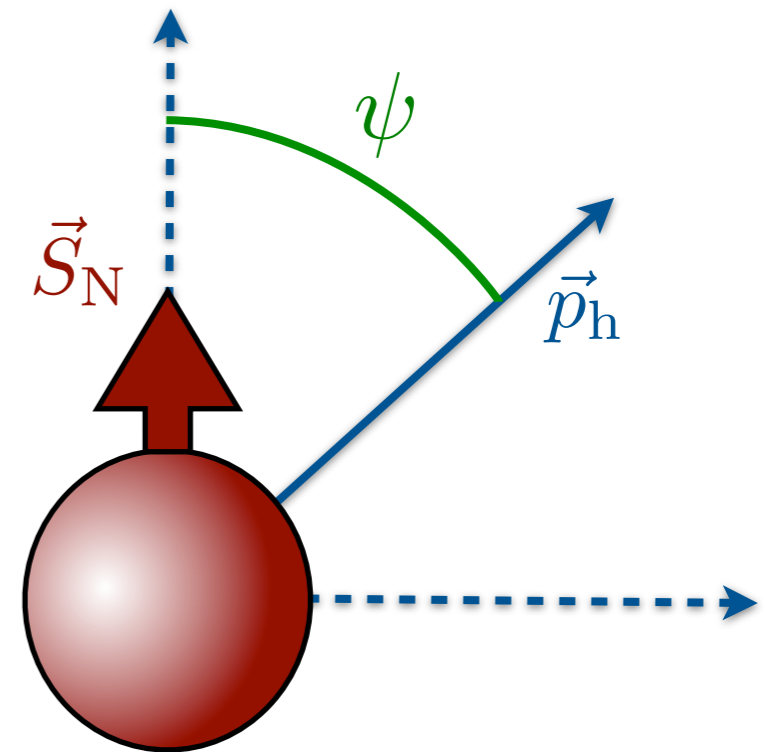
$$ep^{\uparrow} \rightarrow hX$$



# Inclusive hadron electro-production

- scattered lepton undetected  
↳ lepton kinematics unknown
- dominated by quasi-real photo-production (low  $Q^2$ )  
↳ hadronic component of photon relevant?

$$ep^{\uparrow} \rightarrow hX$$

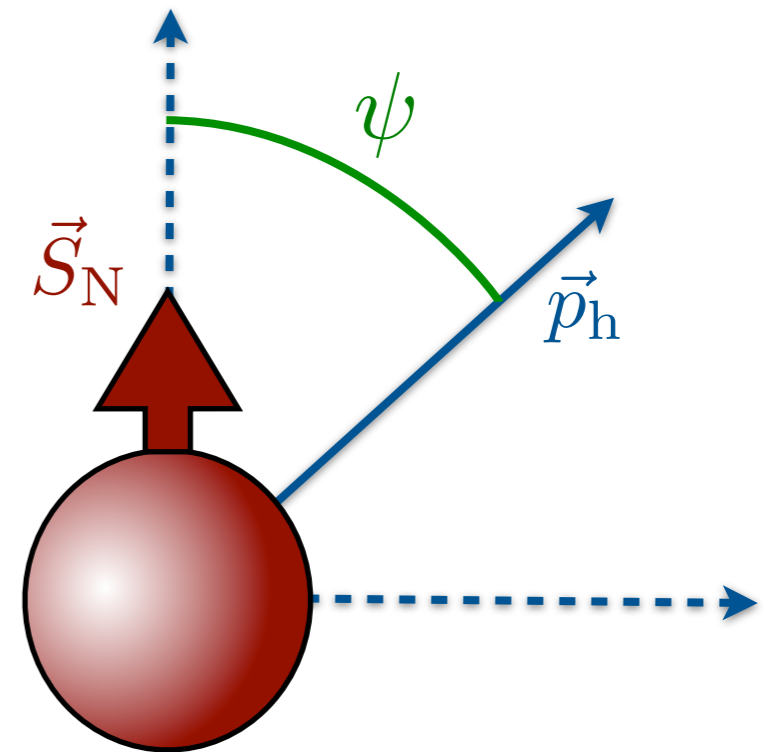




# Inclusive hadron electro-production

- scattered lepton undetected  
 ↳ lepton kinematics unknown
- dominated by quasi-real photo-production (low  $Q^2$ )  
 ↳ hadronic component of photon relevant?
- cross section proportional to  $S_N (\mathbf{k} \times \mathbf{p}_h) \sim \sin\psi$

$$ep^{\uparrow} \rightarrow hX$$



$$A_{UT}(P_T, x_F, \psi) = A_{UT}^{\sin\psi}(P_T, x_F) \sin\psi$$

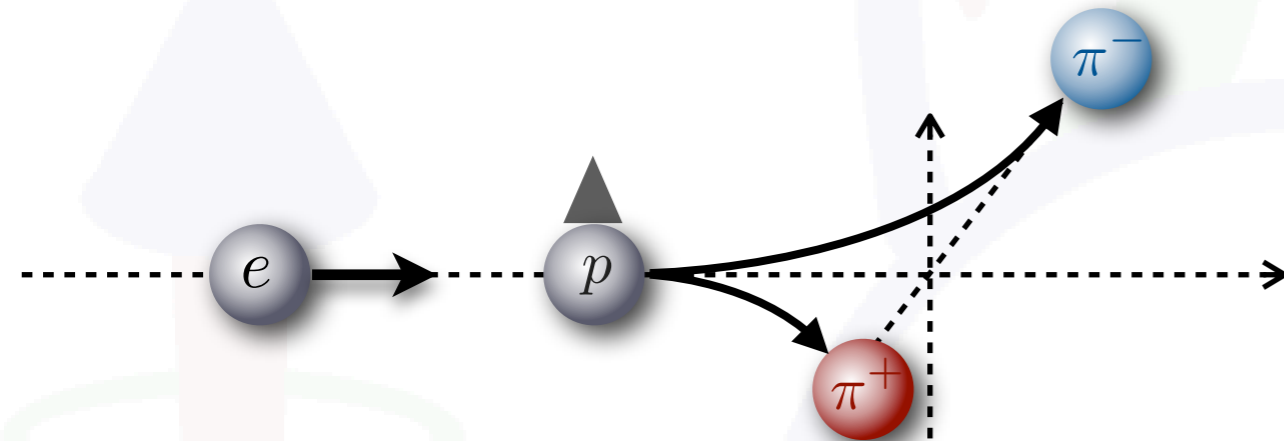
$$A_N \equiv \frac{\int_{\pi}^{2\pi} d\psi \sigma_{UT} \sin\psi - \int_0^{\pi} d\psi \sigma_{UT} \sin\psi}{\int_0^{2\pi} d\psi \sigma_{UU}}$$

$$= -\frac{2}{\pi} A_{UT}^{\sin\psi}$$

# 1D dependences of $A_{UT} \sin\psi$ amplitude

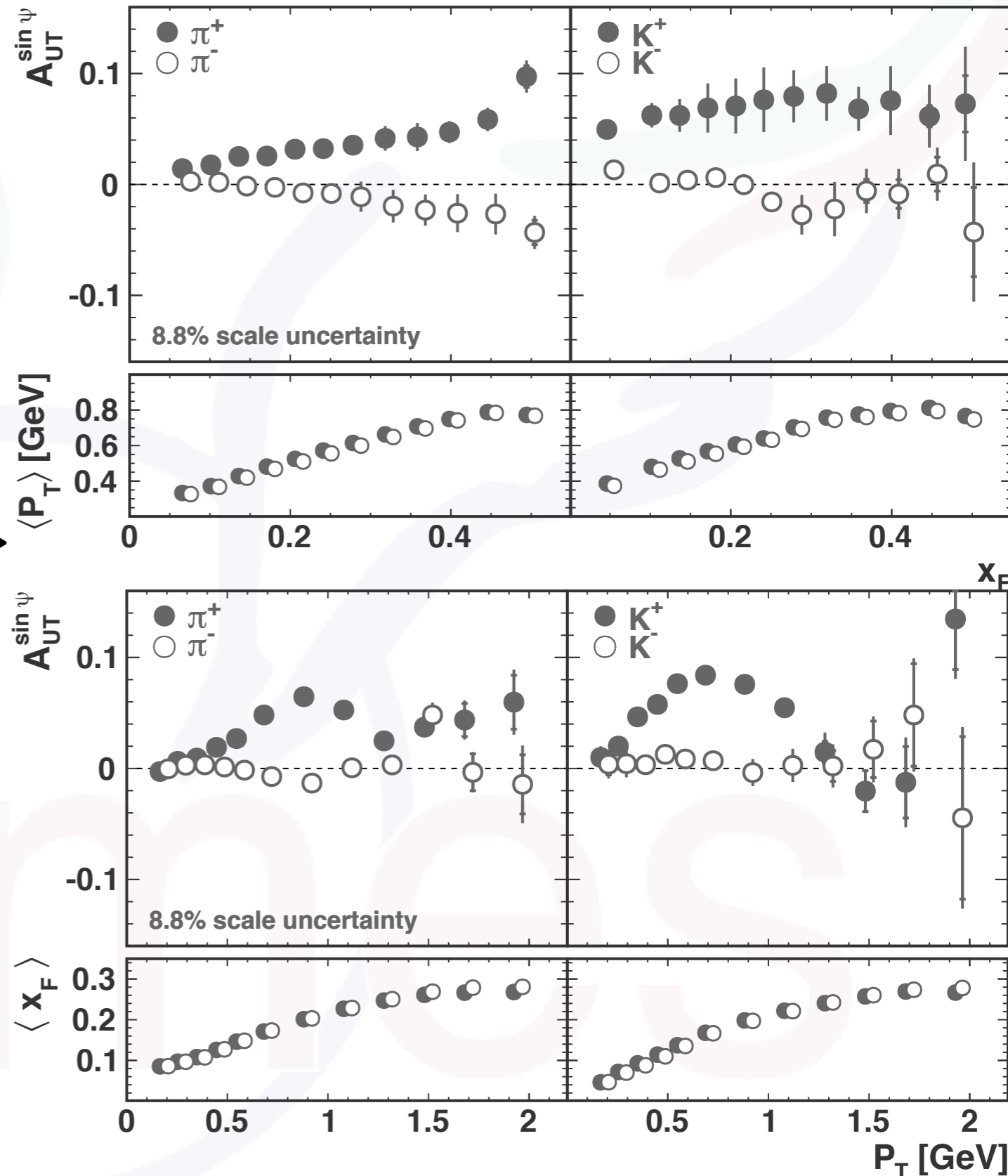
- clear left-right asymmetries for pions and positive kaons

- increasing with  $x_F$  (as in pp)



- initially increasing with  $P_T$  with a fall-off at larger  $P_T$

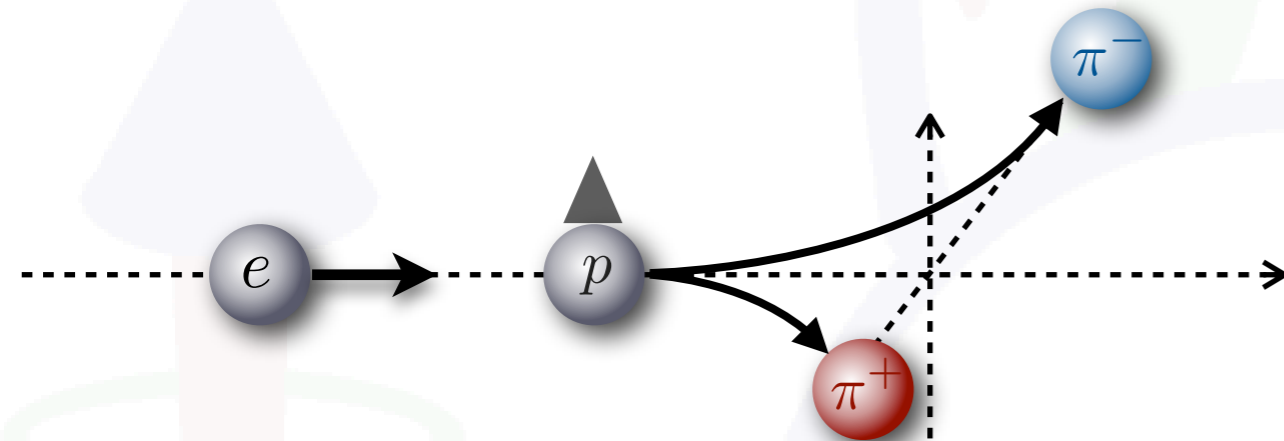
[Airapetian et al., Phys. Lett. B 728, 183-190 (2014)]



# 1D dependences of $A_{UT} \sin\psi$ amplitude

- clear left-right asymmetries for pions and positive kaons

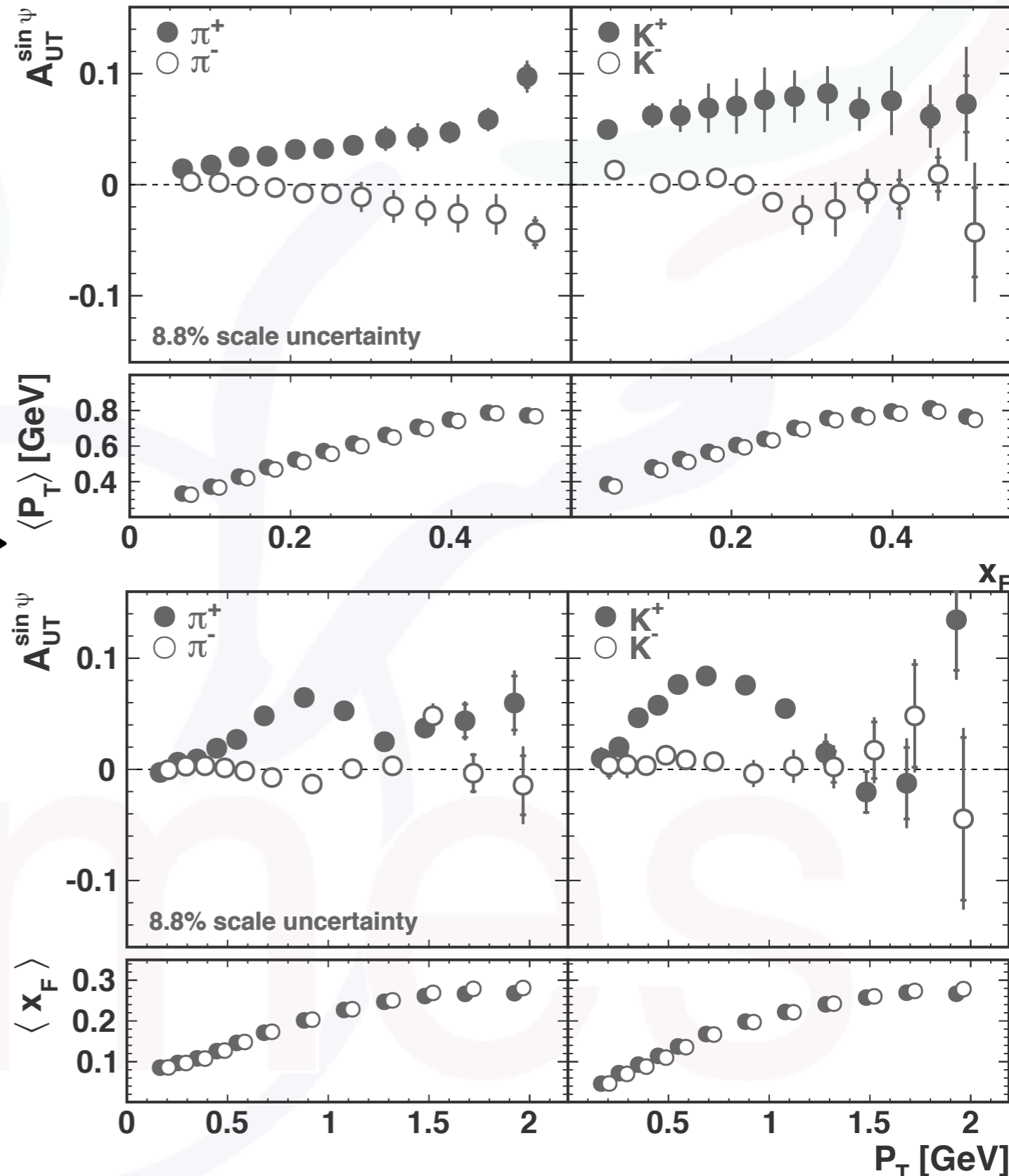
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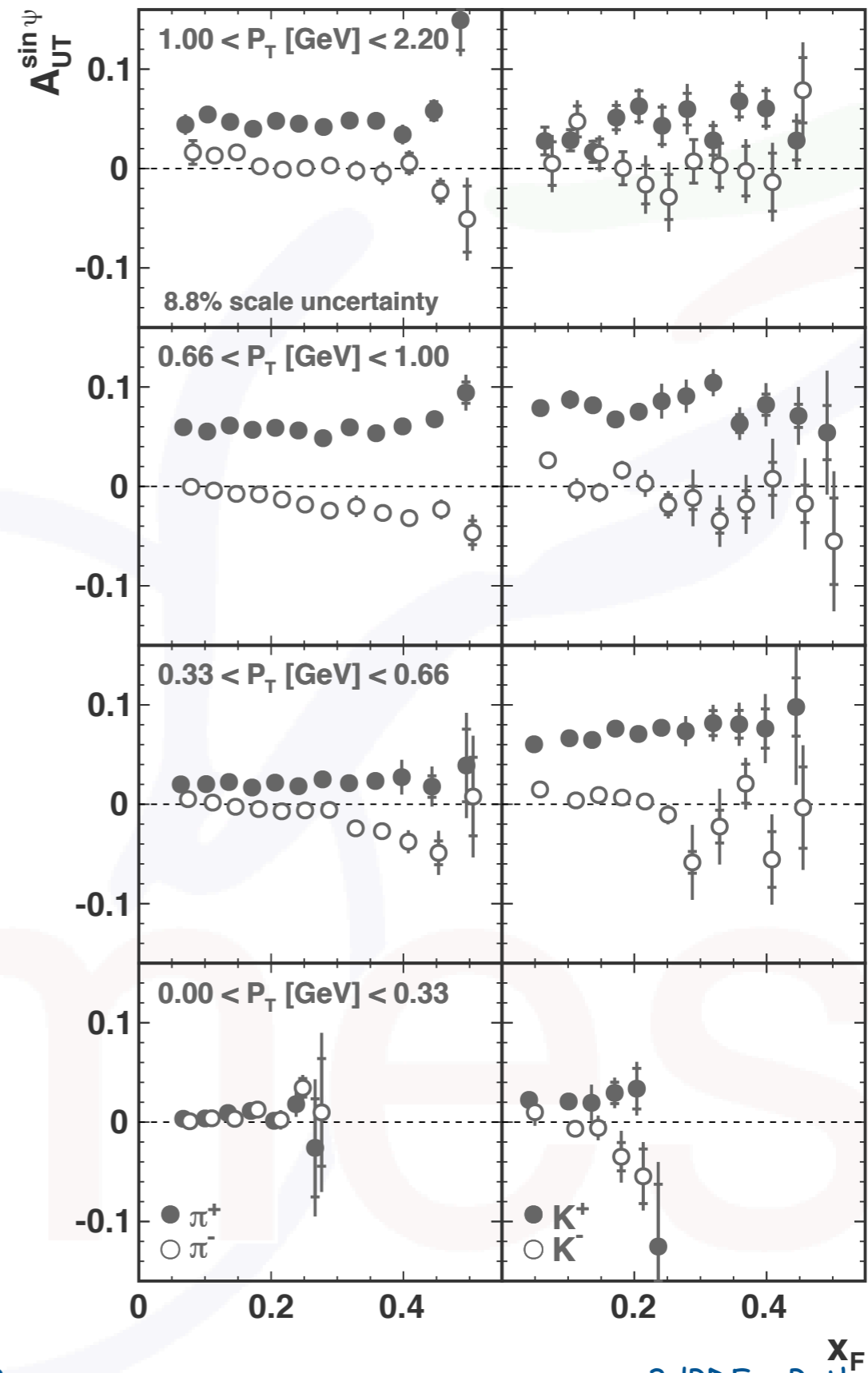
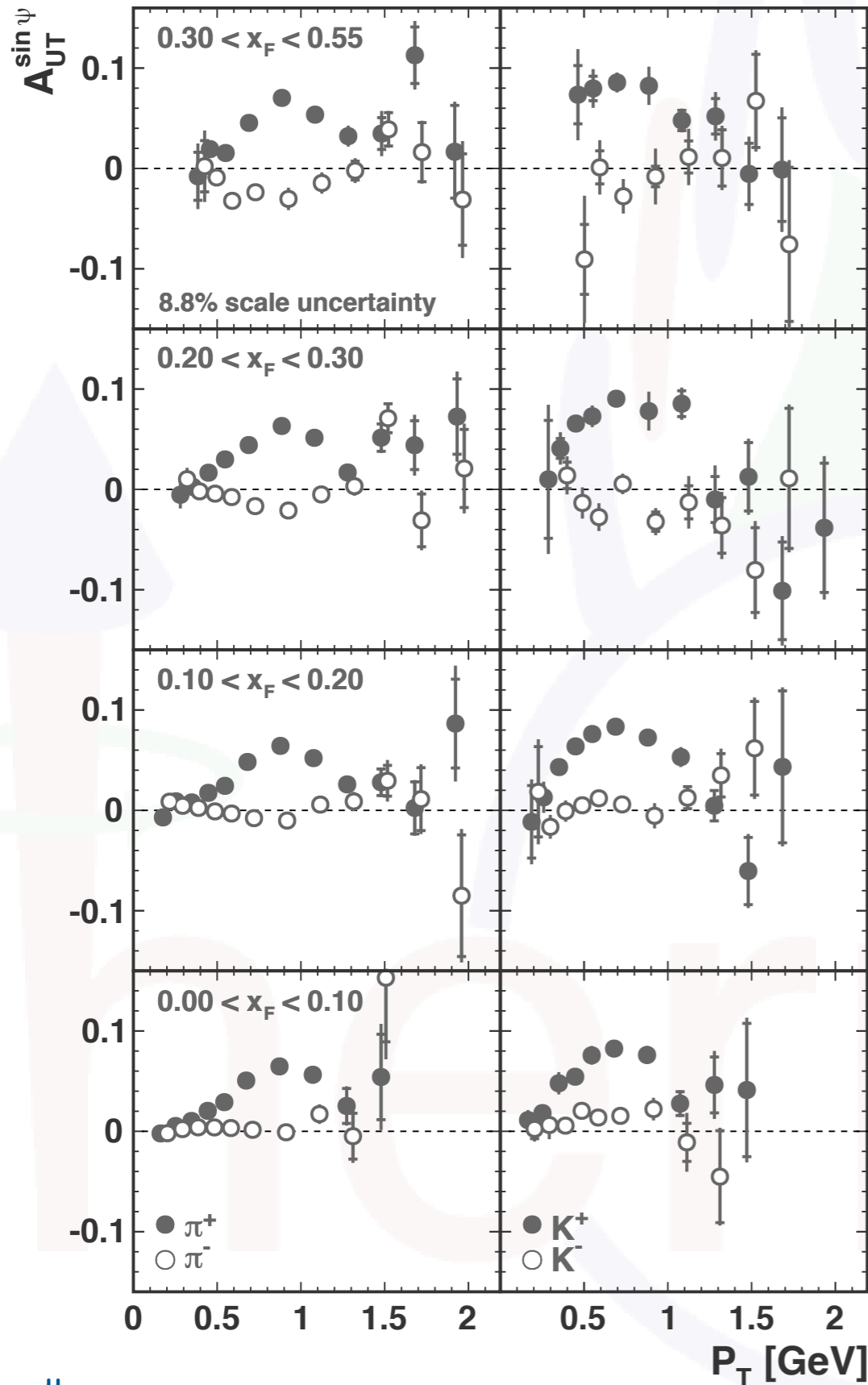
- $x_F$  and  $P_T$  correlated  
 ➔ look at 2D dependences

[Airapetian et al., Phys. Lett. B 728, 183-190 (2014)]



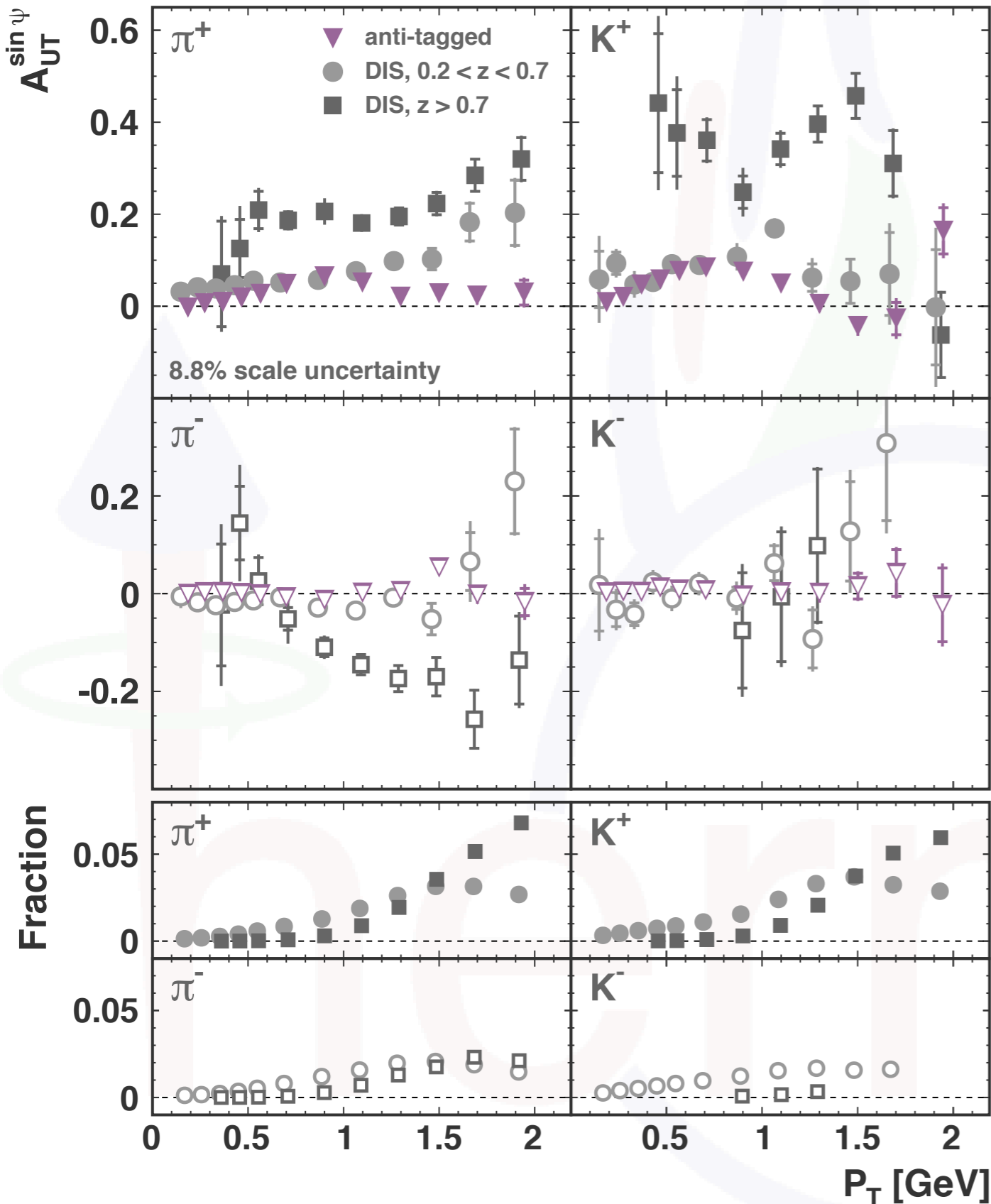
# Inclusive hadrons: 2D dependences

[Airapetian et al., Phys. Lett. B 728, 183-190 (2014)]



# Asymmetries of subprocesses

[Airapetian et al., Phys. Lett. B 728, 183-190 (2014)]



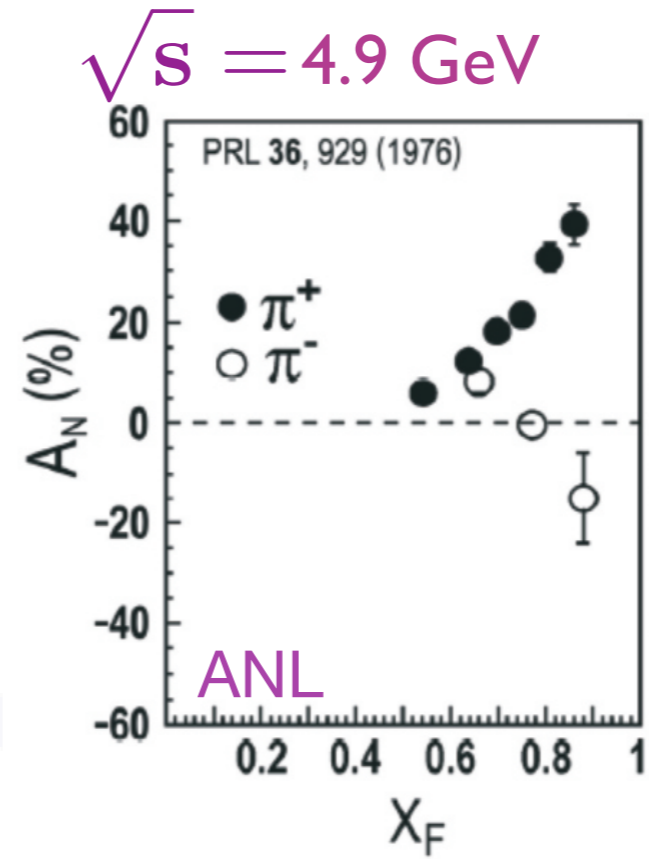
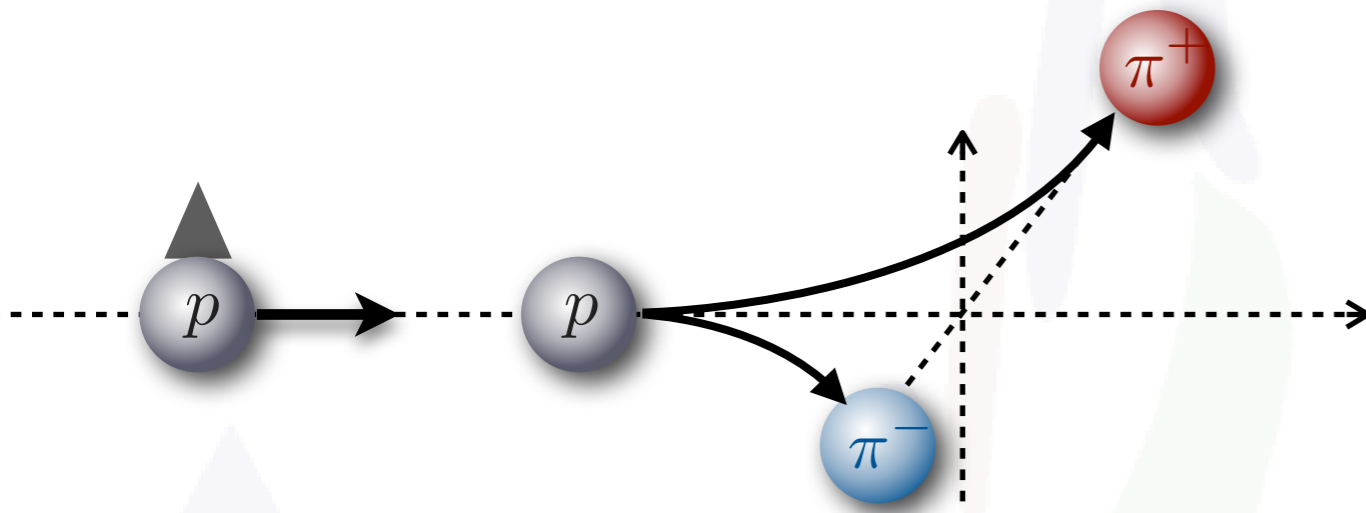
“anti-tagged”  
no lepton in  
acceptance

DIS  
 $0.2 < z < 0.7$

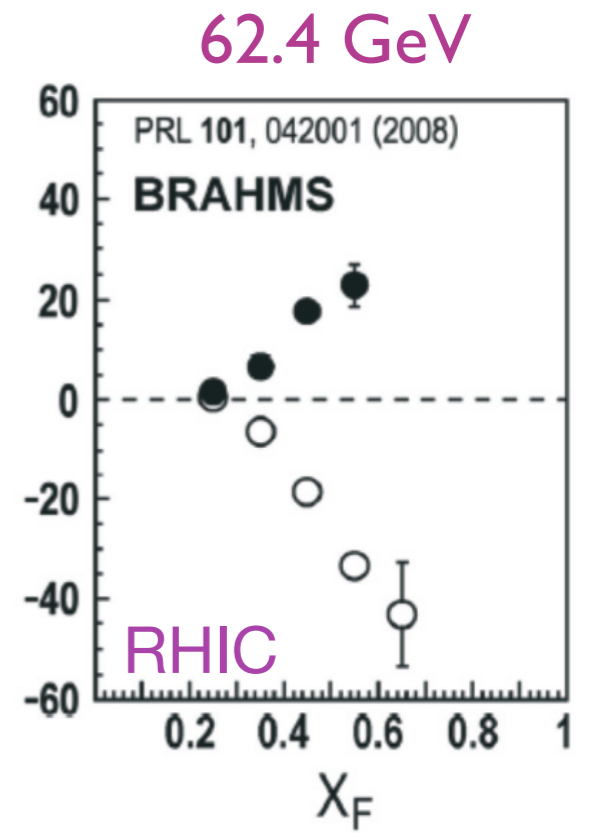
DIS  $z > 0.7$

- at large  $P_T$  significant contribution from DIS events ( $Q^2 > 1$ )
- asymmetries increase with larger  $z$
- large asymmetries also for  $\pi^-$  in case of  $z > 0.7$

# the other inclusive SSA

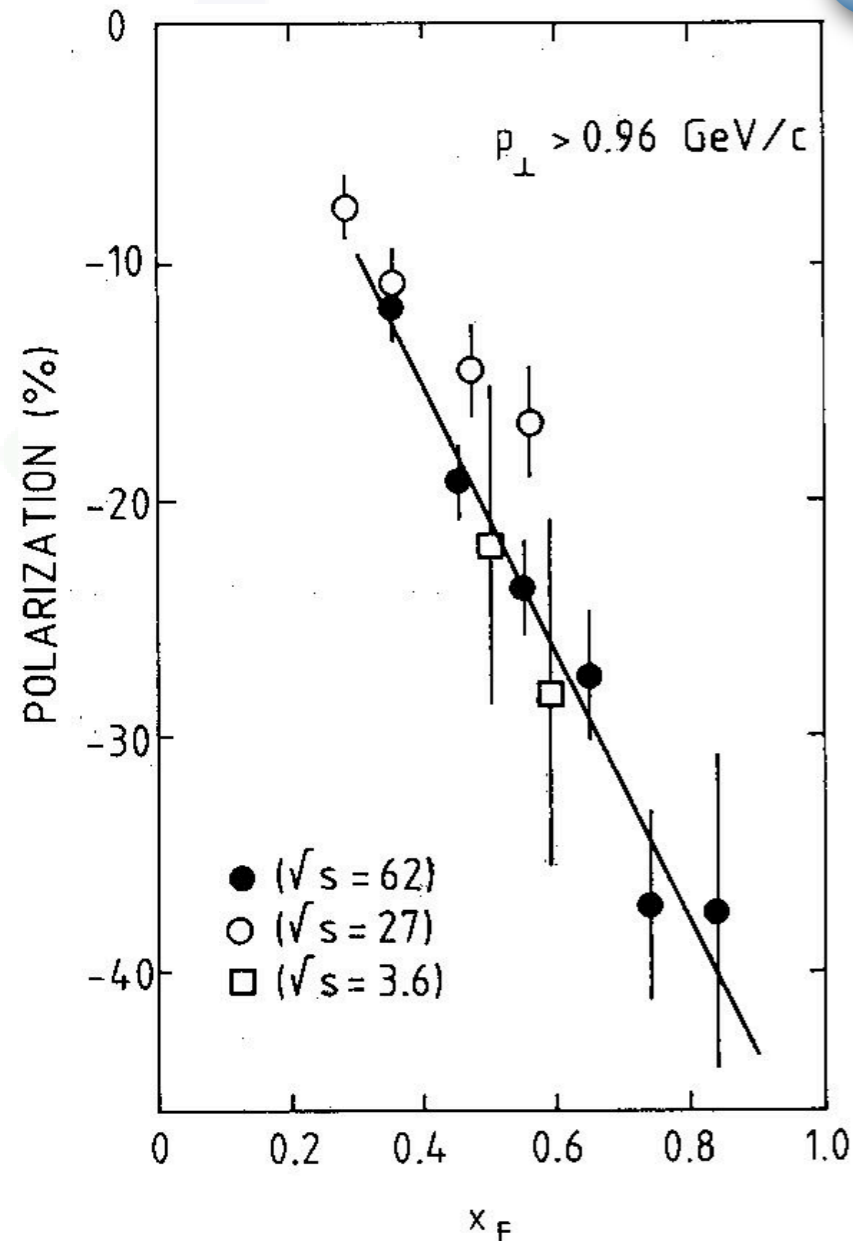
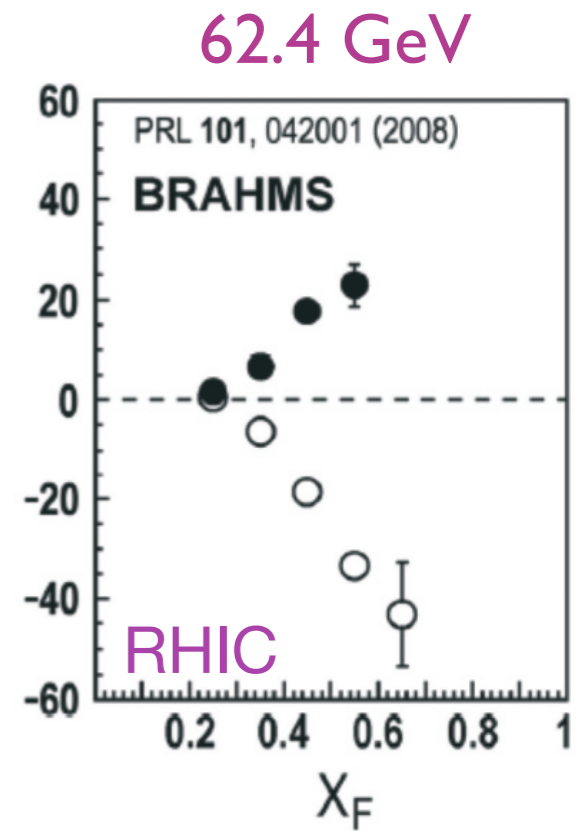
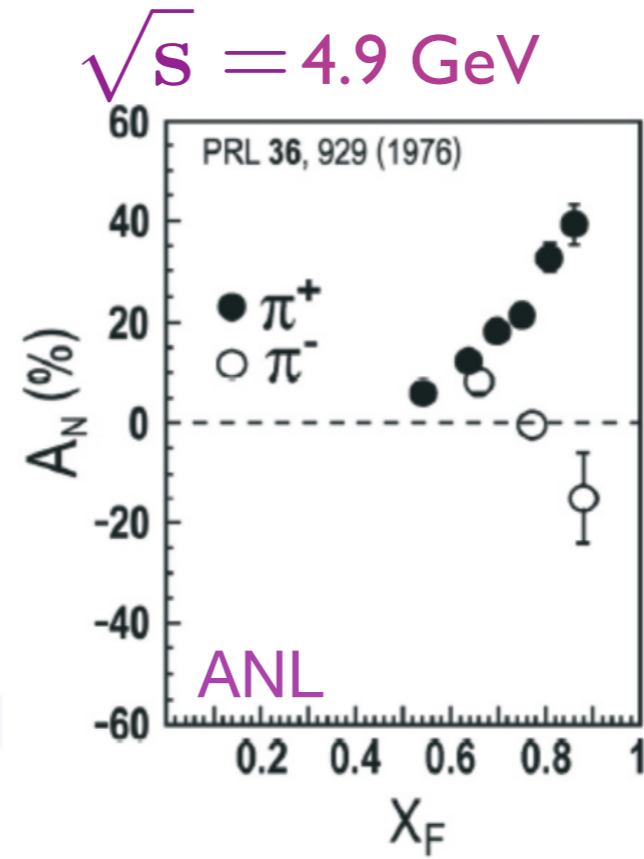
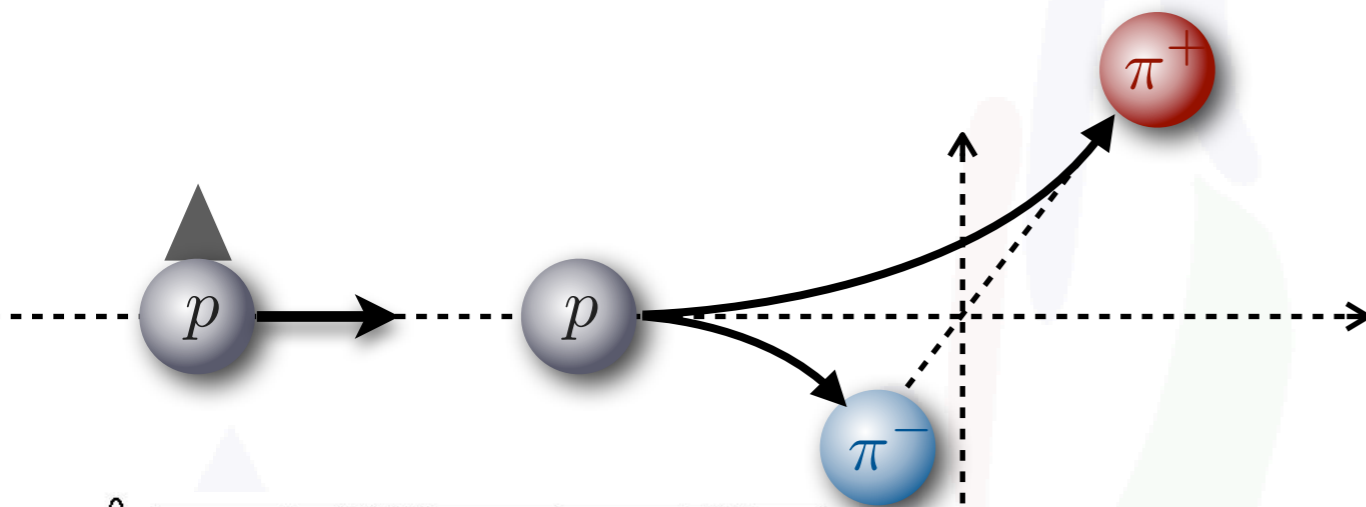


1976



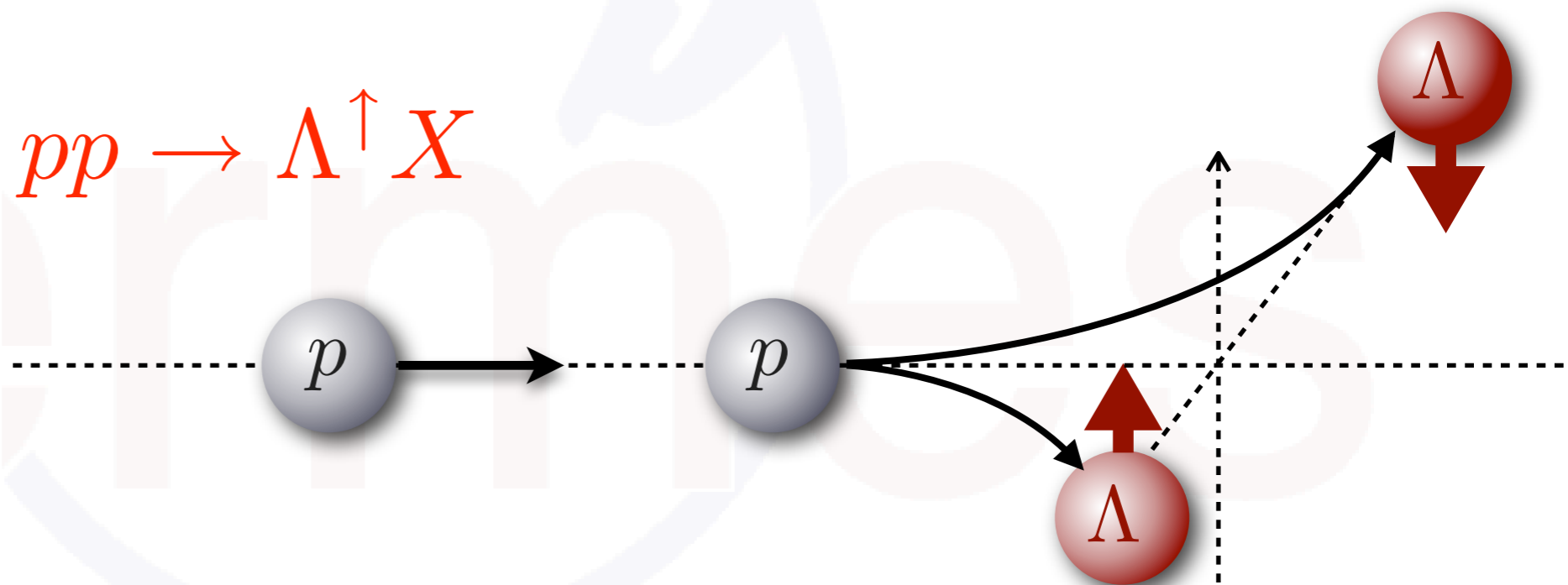
2008

# the other inclusive SSA



1976

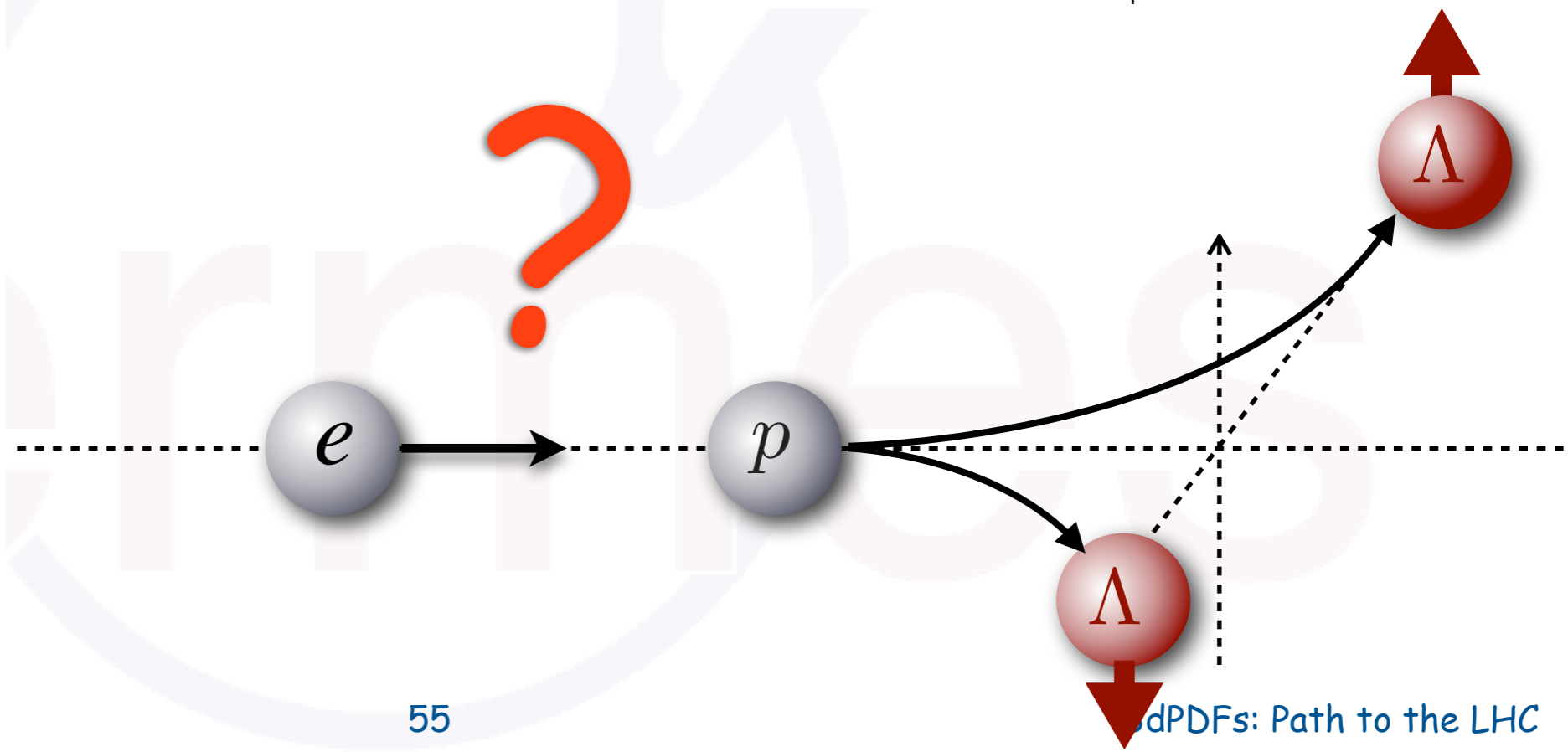
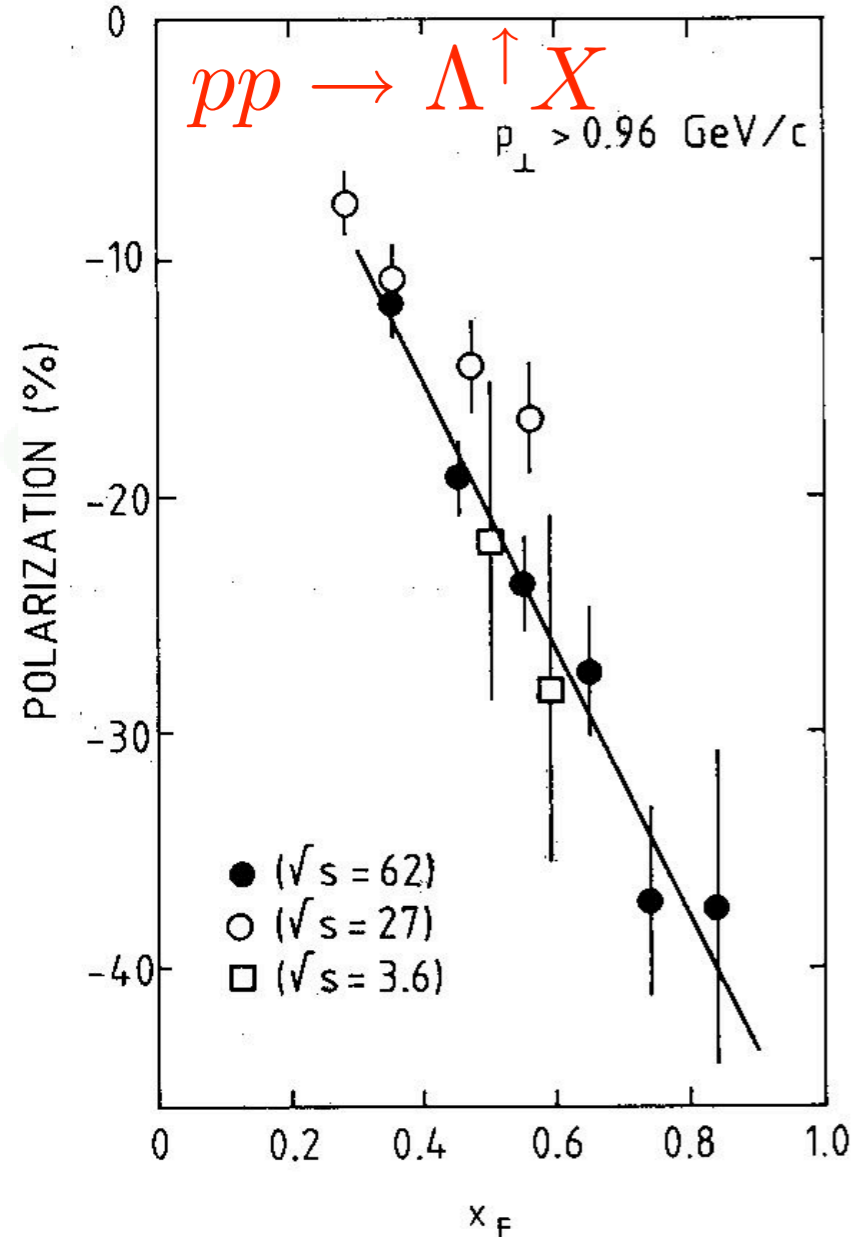
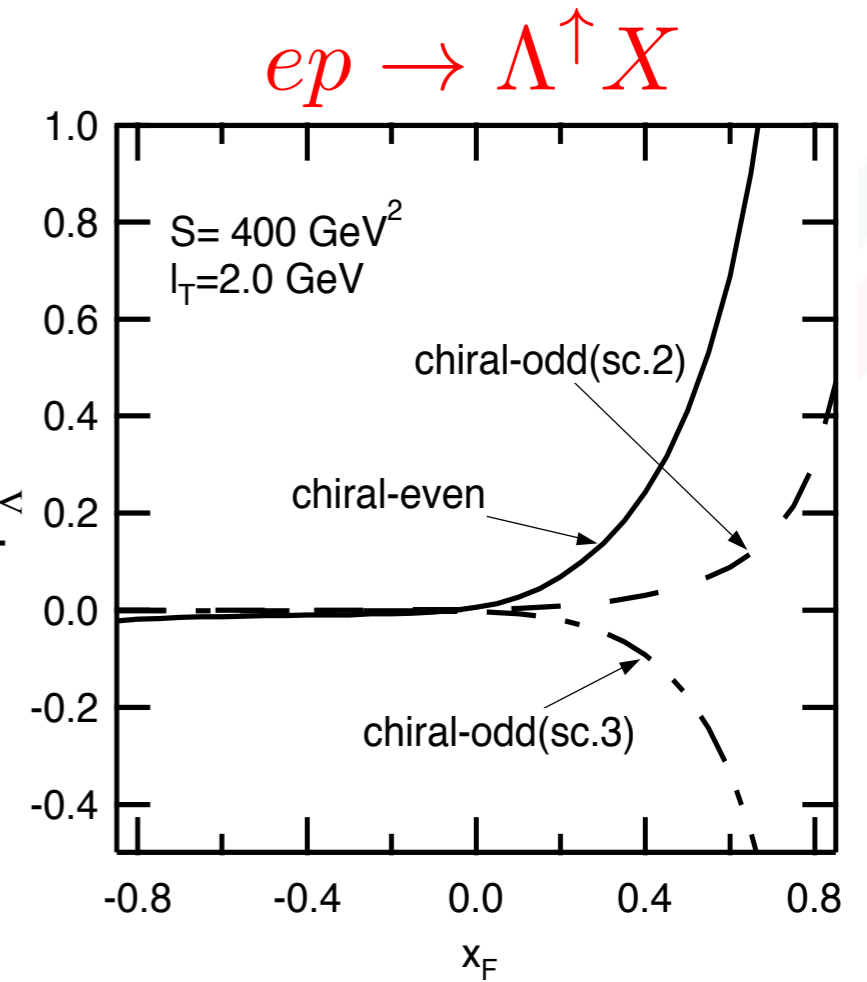
2008



# the other inclusive SSA

in SIDIS (large  $Q^2$ ) proportional to polarizing FF  $D_{1T}^\perp$  (naive T-odd, chiral-even

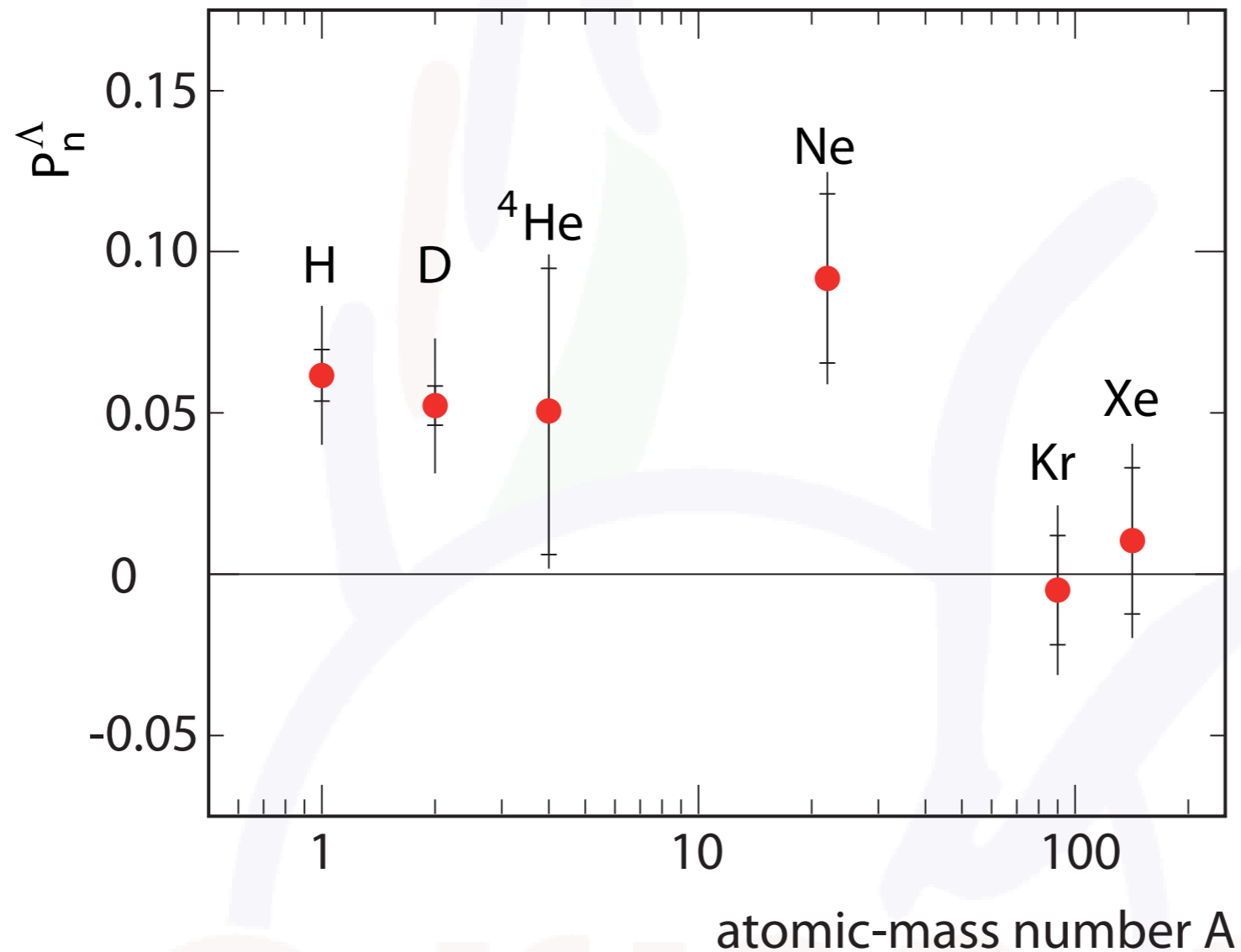
in twist-3 factorization opposite sign to pp





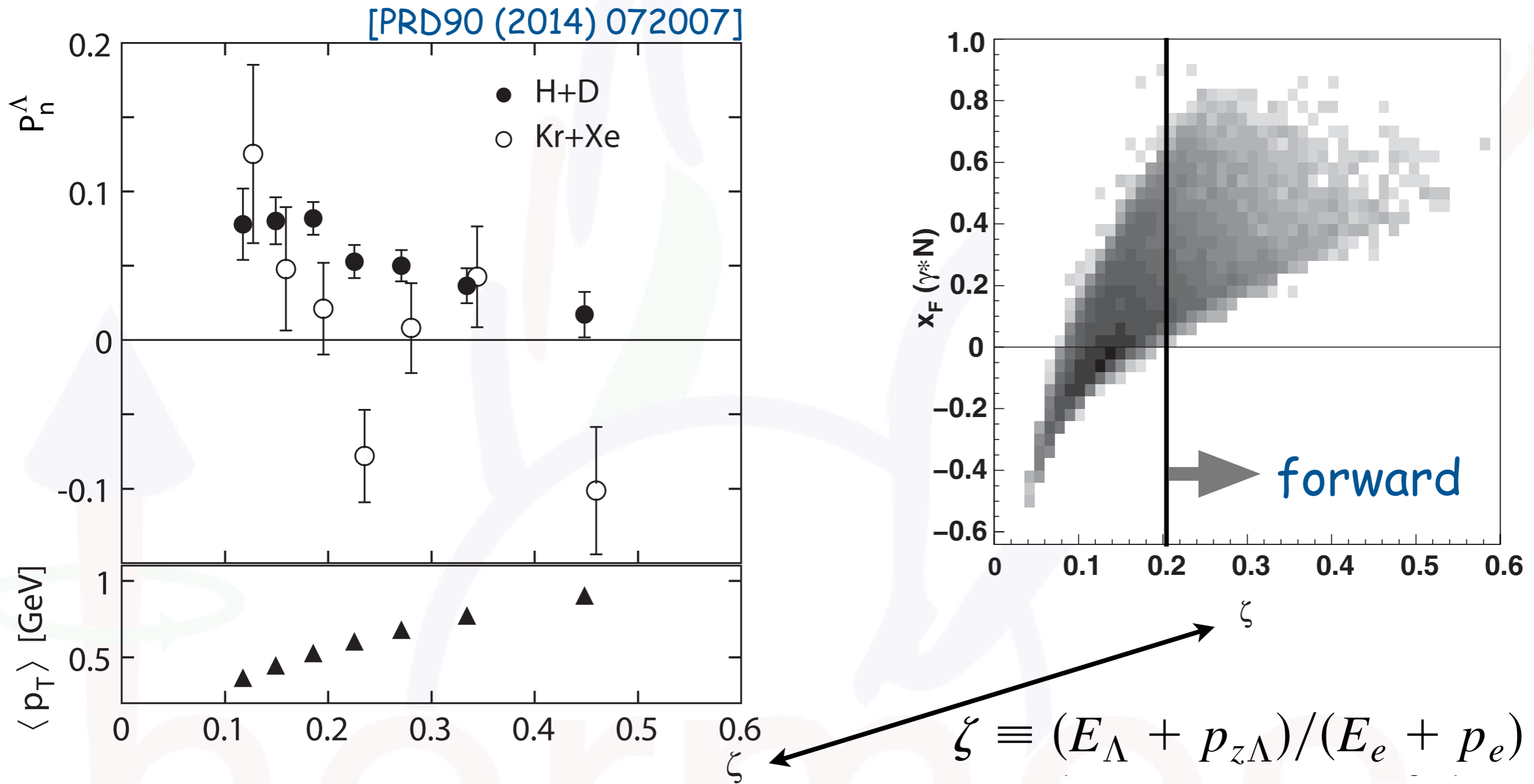
# the other inclusive SSA

[PRD90 (2014) 072007]



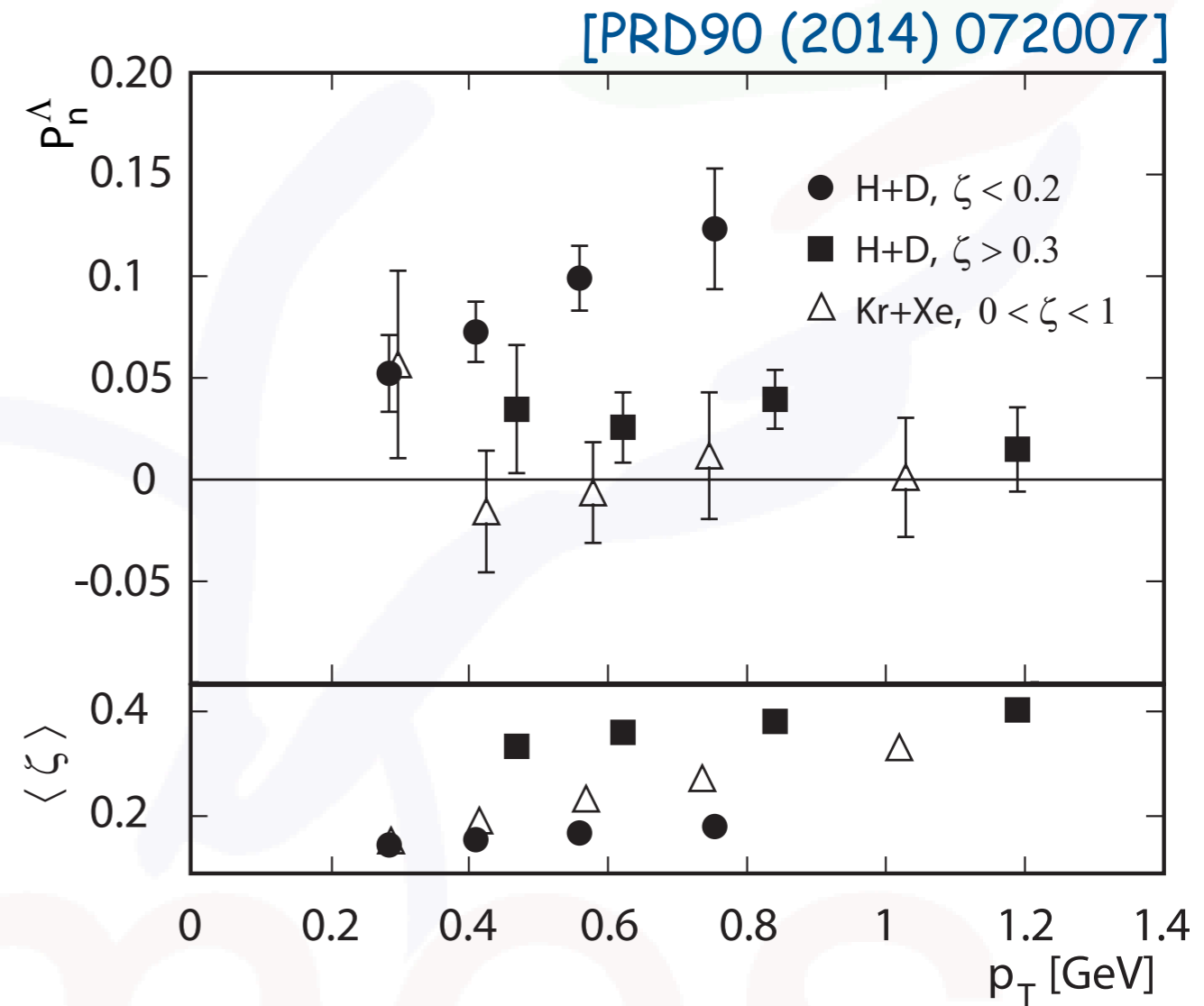
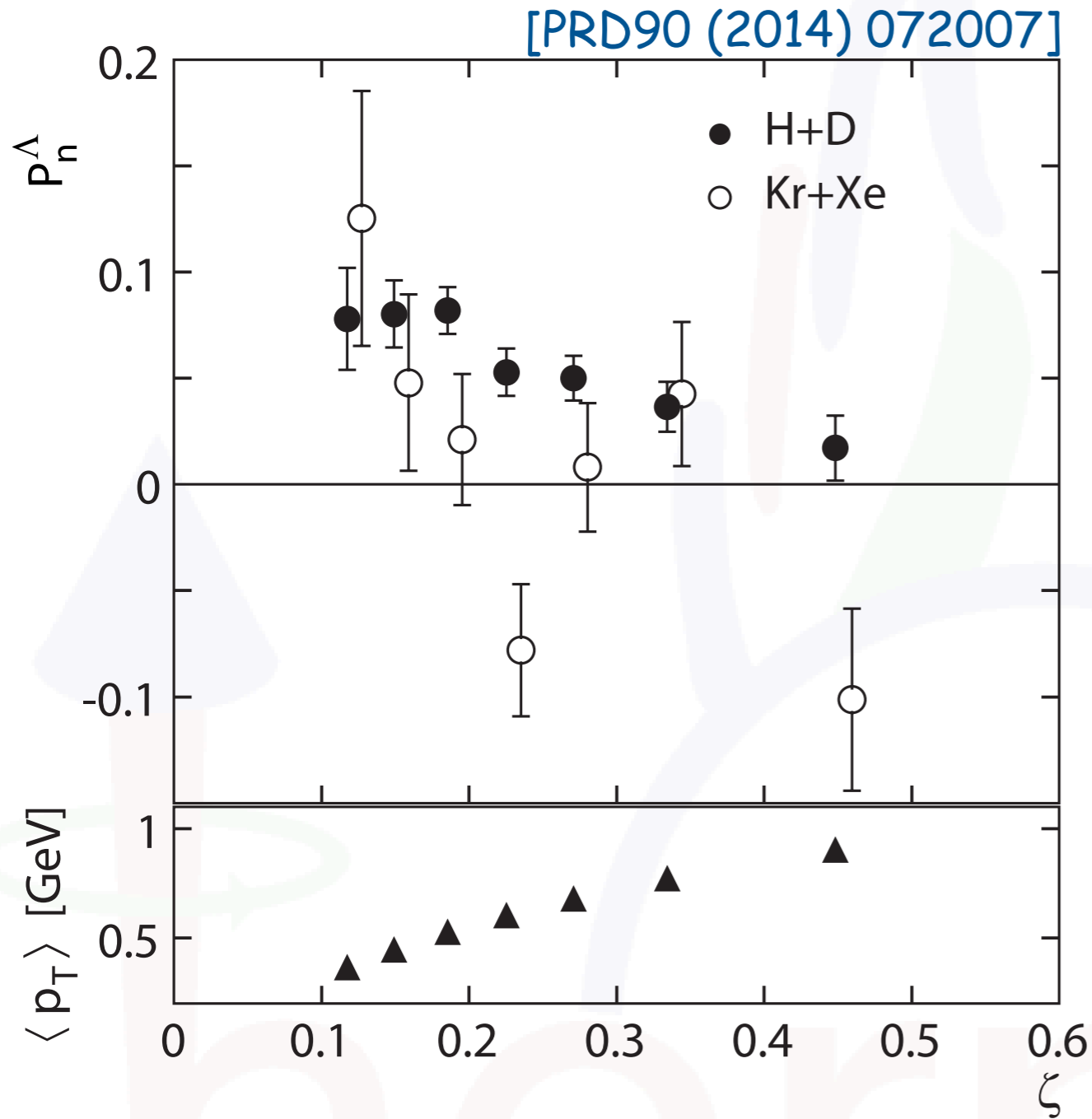
- clearly positive for light target nuclei
- consistent with zero for heavy targets

# the other inclusive SSA



- larger in backward direction w.r.t. incoming lepton
- consistent with  $x_F$  dependence of twist-3 calculation (opposite sign conventions for  $x_F$ !)

# the other inclusive SSA



- larger in backward direction w.r.t. incoming lepton
- distinct  $p_T$  dependences in forward and backward directions: rising with  $p_T$  in backward direction as in pp

# conclusions before the summary

- HERMES conceived almost 3 decades ago in order to solve the "spin crisis"
- measure precisely the quark-spin and somewhat the gluon spin contribution to the proton spin
- no orbital angular momentum on the menu
- no real transverse-spin physics
- up to  $g_2$  and the Burkhardt-Cottingham S.R. ...  
... and that mainly to have a more precise  $g_1$  measurement

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... and that mainly to have a more precise  $g_1$  measurement
- thanks also to the "believers" in the Frascati group, HERMES has published a wealth of transverse-spin results, among others, HERMES' most cited publications

# conclusions before the summary

- HERMES conceived almost 3 decades ago in order to address the "spin crisis"
- measure precisely the quark-spin contribution to the proton spin, and to determine that the gluon spin contribution is non-zero
- no orbital angular momentum measurements
- no real transverse spin measurements
  - up to 2000, at the DESY-Berlin synchrotron, by the HERMES collaboration, led by the group of R. Arndt-Cotttingham S.R. ...
  - ... and the desire to have a more precise  $g_1$  measurement
- the "believers" in the Frascati group, HERMES published a wealth of transverse-spin results, among which, HERMES' most cited publications

always be open to trying out new paths

