Probing transverse nucleon structure at high momentum transfer Trento / ECT* - 18-22/April/2016



Proton Form Factors

(Space-Like, EM p FF)

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Outlook

- Short EMFF overview and p-FF status
- JLab opportunity at high Q²
- SBS GEp experiment and expected results

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Nucleon, it is a mystery!

- Mass origin
 - Confinement / QCD
 - Internal dynamics
- Spin origin
 - Orbital Angular Momentum
 - Gluon Contribution
- (Proton) radius
- QED-QCD reaction mechanisms

Toward a unified picture of nucleon structure



Form Factors and GPDs



see M. Diehl and P. Kroll EPJ C (2013) 73:2397 model dependent attempt

«We note that the **electromagnetic form factors** provide indirect constraints on GPDs at **high values of t**, which will conceivably never be accessible in hard exclusive scattering processes.»

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FF «traditional» extraction method

- Rosenbluth separation in Born approximation (Z α <<1) $\frac{d\sigma_r}{d\Omega} = \varepsilon \ G_E^2(Q^2) + \tau \ G_M^2(Q^2), \qquad \tau = Q^2/(4M^2)$ $\varepsilon = \left[1 + 2(1+\tau)\tan^2(\theta_e/2)\right]^{-1}$
 - Can determine both absolute values of G_{E} and G_{M}
 - τ kinematically suppress G_E at high Q^2 and G_M at low Q^2
 - Measure cross sections for different ε, then make linear fits
 - Different "kinematic" approaches



The power of spin

- Akhiezer et al. Sov. Phys. JETP 6, 588 (1958), Sov. PHys. Dokl. 14 (1968), ...<
 Form factor accuracy can be improved by measuring interference term G_EG_M by means of beam helicity asymmetry with polarized target or recoil polarimetry
- About 30 years to get the needed technologies:
 - Polarized e beam (high current 100 µA and polarization>70%) and beam polarimetry (~3%)
 - Highly polarized targets
 - Efficient recoil polarimeters (reasonably high analyzing power)

FF «modern» extraction methods

• Recoil polarization: $H(\vec{e}, e', \vec{p}), {}^{2}H(\vec{e}, e', \vec{n})H \quad \frac{P_{t}}{P_{l}} \tan(\frac{\theta}{2}) \propto -\frac{G_{E}^{p}}{G_{M}^{p}}$

 P_t , P_l = trans. and long. polarization of the recoil proton

• Beam-Target polarization asymmetry:

$$\overrightarrow{H}(\vec{e}, e'p) \qquad A = \frac{N^+ - N^-}{N^+ + N^-} \approx \frac{G_E}{G_M}$$
³He($\vec{e}, e'n$)

many systematics (theory and exp.) cancel in ratio

Proton G_E/G_M – an «unexpected» discrepancy



$$rac{d\sigma}{d\Omega} \propto G_{Ep}^2 + rac{ au}{arepsilon} G_{Mp}^2$$

Rosenbluth Separation: assume single photon approximation

Prior to JLab/2000, expectations were that proton G_E/G_M fairly constant with Q^2

$$\mu \frac{G_{Ep}}{G_{Mp}} = -\mu \frac{P_t}{P_l} \frac{(E_{beam} + E_e)}{2M_p} \tan \frac{\vartheta_e}{2}$$

Polarization transfer from the incident electron to the scattered proton

At JLab, new class of experiments show proton G_E/G_M decreasing linearly with Q^2

$$R_{p} = \mu_{p} \frac{G_{E}(Q^{2})}{G_{M}(Q^{2})} \approx 1 - \underbrace{0.13 (Q^{2} - 0.29)}_{Pol. \ Transfer \ Discr.}$$

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 Q^2 / (GeV/c²)

• Description of the reaction mechanism is wrong (?)

• G_{E} and G_{M} do not scale the same way

TPE makes FF extraction model dependent

Cross section up to 2γ exchange approximation (nucleon model dependent):

$$d\sigma_{1\gamma+2\gamma} = C \left\{ |\tilde{G}_{M}|^{2} + \frac{\epsilon}{\tau} \left[|\tilde{G}_{E}|^{2} + 2|\tilde{G}_{M}|^{2} \left(\tau + \frac{|\tilde{G}_{E}|}{|\tilde{G}_{M}|}\right) Y_{2\gamma} \right] \right\}$$
Rosenbluth at 1_Y approx. Not negligible at high Q²

$$(R_{Rosenbluth}^{exp})^{2} \doteq \frac{|\tilde{G}_{E}|^{2}}{|\tilde{G}_{M}|^{2}} + 2 \left(\tau + \frac{|\tilde{G}_{E}|}{|\tilde{G}_{M}|}\right) Y_{2\gamma}$$
D. Adikaram et al. (CLAS12 TPE)
PRL 114 062003 2015
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PRL 114 062003 2015
D. Adikaram

FF at small Q²: the nucleon E,M radius

Nucleon size (radius) derived from Form Factors (at low Q²)

 $\begin{aligned} G_E^{p,n}(Q^2) &= G_E^{p,n}(0) - \frac{1}{6}Q^2 \langle r^2 \rangle_E^{p,n} + \cdots \Rightarrow \left\langle r^2 \rangle_E^{p,n} = -6 \left. \frac{d \, G_E^{p,n}(Q^2)}{dQ^2} \right|_{Q^2=0} \\ G_M^{p,n}(Q^2)/\mu &= 1 - \frac{1}{6}Q^2 \langle r^2 \rangle_M^{p,n} + \cdots \end{aligned}$ (similar for the magnetic radius)



2010 (Pohl et al.): Large disagreement between the proton charge radius from ep scattering/spectroscopy and µp Lamb shift



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Different variations of the Rosenbluth approach give consistent results

 \rightarrow exp. systematic under control

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"Non Rosenbluth" measurements



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Proton F_2/F_1 scaling – pQCD prediction



- Quark Helicity conservation
- Counting rules:
 - 1/Q² for gluon line
 - 1/Q² for helicity flip
- Photon absorbed by one «collinear» quark which interacts with the other two quarks by 2×gluon exchange

$$\rightarrow$$
 F₁ ~ 1/Q⁴, F₂ ~ 1/Q⁶

Modified pQCD: Belitsky et al. & other models include quark orbital angular momentum and gluon polarization effects to explain F_2/F_1 at arge O² 20/April/2016 (ETC*)

Naive Summary - Why and where are FFs interesting

- Fundamental properties of the nucleon
- Many fundamental models able to calculate them
- Large Q²:
 - Distinguish important models for GE/GM
 - Scaling behavior (flavor form factors components)
 - Constrain GPD at high x (valence quark dominate)
 - Smaller effective mass of the quarks
 - Toward pQCD dominated regime
- Very Low Q²:
 - Pion cluod effects
 - Precise nucleon size estimation

SBS original motivations (2007) - Proton G_E/G_M



- VMD (lachello, Lomon, Bijker), generally good description of all FFs
- Relativistic CQM (Miller, Gross, ...) spin dependent quark density
- Lattice QCD, starts to give prediction
- Dyson-Schwinger, dressed quarks, diquark correlation, ...
 - pQCD-based:
 G_E/G_M→const, Q²→ ∞
- GPD-based: direct connection to quark
 OAM, FF's constraint GPD's

Most of them agree with current data but diverge at higher, unexplored, Q²

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New SuperBigbite Spectrometer (SBS) in Hall A



SBS specifications (fully configured)



EMFF measurements at high Q² by pol. methods

Method:	Polarization Transfer: $p(\vec{e}, e'\vec{p})$	Target Perp. Polarization: $p^{\uparrow}(\vec{e}, e'p)$		
Measure (one photon approx.)	$\frac{P_t}{P_l} \tan\left(\frac{\theta}{2}\right) \propto \frac{G_E^p}{G_M^p}$ P _t , P _l : trans. and long. polarization of the recoil proton	$A = \frac{N^+ - N^-}{N^+ + N^-} \sim \frac{G_E^p}{G_M^p}$ N ⁺ and N ⁻ : events with opposite transverse target polarization		
Many systematics effects (theory and exp.) cancel in ratio				
Figure of Merit (stat.) Ω : acceptace L: Luminosity σ : elastic xsec ~ E^2/Q^{12} Pb: beam polarization	$\begin{array}{l} \Omega \mathrel{\ L} \sigma \mathrel{\ P_b^2} \mathrel{\ } \epsilon \mathrel{\ } A_y^2 \\ \sim \displaystyle \frac{\Omega \mathrel{\ } \mathrel{\ } L \mathrel{\ } \epsilon}{Q^{16}} \\ A_y: \text{ polarimeter analyzing power} \\ \epsilon: \text{ polarimeter efficiency} \end{array}$	$\begin{split} \Omega \ L \ \sigma \ P_b^2 \ P_T^2 \\ \sim & \frac{\Omega \ L}{Q^{12}} \ P_T^2 \\ P_T : \text{Target polarization} \end{split}$		
At Q ² ~10 GeV ² expected: $FoM_{pol_trans} \sim 10 \times FoM_{targ_pol}$ (target polarization cannot tolerate large L)				
 Challenges at high Q²: <u>need to maximize</u> (coincidence) acceptance (solid angle) luminosity polarimetry efficiency beam polarization (having the needed beam energy) keeping costs at «affordable» level 				
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Proton G_E/G_M at large Q^2 by polarization transfer



Absolute error < 0.1 Beam time = 60 days

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Large Luminosity ⇒ Large Background

Tracker	Area of interest	Hit Rate,
	for tracking, cm^2	kHz/cm ²
First	0.20 x 18	400
Second	$2\pi \ 0.35^2$	130
Third	π 4.8 ²	64
BigCal	$\pi \ 1.2^2$	173

 $L\sim 10^{39}\ /cm^2/s$

- Must be supported by the detectors
 ⇒ GEM technology
- Must be handled by the trigger:
 - spatial and time correlation between electron and proton elastically scattered
 - «high» energy threshold in segmented CALO's



GEp5: Proton Polarimeter (PP)

$$\frac{G_E}{G_M} = -\frac{P_t}{P_l} \frac{E+E'}{2M} \tan \frac{\theta}{2} \left[1 + (\text{few \%})_{2\gamma}\right]$$

Use azimuthal asymmetry of the proton scattering off matter induced by spin-orbit coupling



Polarimeter only measures components of proton spin that are **transverse** to the proton's momentum direction

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Number of scattered protons: $f^{\pm}(\vartheta,\varphi) = \frac{e^{\rho\rho}(\vartheta,\varphi)}{2\pi} \left[1 \pm A_{\gamma} \left(P_{\chi}^{\rho\rho} \sin\varphi + P_{\gamma}^{\rho\rho} \cos\varphi \right) \right]$ where \pm refers to electron beam helicity $A \doteq \frac{f^+ - f^-}{f^+ + f^-} = A_Y \left(P_X^{pp} \sin \varphi + P_V^{pp} \cos \varphi \right) = A_Y \cos(\phi - \delta)$ $\tan \delta \doteq \frac{P_x^{pp}}{P_x^{pp}}$ A (a.u.) 0.02 0.01 0 -0.01-0.0290 180 270 360 ϕ (degrees) $\sigma_{P_{X,Y}^{pp}} \sim \sqrt{2}/(A \cdot P_e \cdot \sqrt{N}) \implies \text{Maximize } P_e$

N=number of scattered proton, $\overline{P_e}$ beam polarization

• Require: Dipole magnet to precess P₁ at target to P_v^{pp}

GEp5 Projected results





- •Settle the question of a zero crossing of GEp/GMp
- Constrain GPDs at high x, large t



8.8

11.0

8.0

12.0

10

30

0.032

0.074

New software tools may improve analysis

from C. Fanelli (ECT* 2015)

A Markov chain Monte Carlo has been used to extract the polarization transfers from the Likelihood:

$$\mathcal{L}(\mathbf{P}) = \prod_{i=1}^{N_{evt}} \frac{1}{2\pi} \left[1 + (a_1 + h\epsilon^{(i)}A_y^{(i)} \sum_{j=x,z} S_{yj}^{(i)}P_j + A_y^{(i)}S_{yy}^{(i)}P_y) \cos \varphi^{(i)} + (b_1 - h\epsilon^{(i)}A_y^{(i)} \sum_{j=x,z} S_{xj}^{(i)}P_j - A_y^{(i)}S_{xy}^{(i)}P_y) \sin \varphi^{(i)} + (b_1 - h\epsilon^{(i)}A_y^{(i)} \sum_{j=x,z} S_{xj}^{(i)}P_j - A_y^{(i)}S_{xy}^{(i)}P_y) \sin \varphi^{(i)} + (b_1 - h\epsilon^{(i)}A_y^{(i)} \sum_{j=x,z} S_{xj}^{(i)}P_j - A_y^{(i)}S_{xy}^{(i)}P_y) \sin \varphi^{(i)} + (b_1 - h\epsilon^{(i)}A_y^{(i)} \sum_{j=x,z} S_{xj}^{(i)}P_j - A_y^{(i)}S_{xy}^{(i)}P_y) \sin \varphi^{(i)} + (b_1 - h\epsilon^{(i)}A_y^{(i)} \sum_{j=x,z} S_{xj}^{(i)}P_j - A_y^{(i)}S_{xy}^{(i)}P_y) \sin \varphi^{(i)} + (b_1 - h\epsilon^{(i)}A_y^{(i)} \sum_{j=x,z} S_{xj}^{(i)}P_j - A_y^{(i)}S_{xy}^{(i)}P_y) \sin \varphi^{(i)} + (b_1 - h\epsilon^{(i)}A_y^{(i)} \sum_{j=x,z} S_{xj}^{(i)}P_j - A_y^{(i)}S_{xy}^{(i)}P_y) \sin \varphi^{(i)} + (b_1 - h\epsilon^{(i)}A_y^{(i)} \sum_{j=x,z} S_{xj}^{(i)}P_j - A_y^{(i)}S_{xy}^{(i)}P_y) \sin \varphi^{(i)} + (b_1 - h\epsilon^{(i)}A_y^{(i)} \sum_{j=x,z} S_{xj}^{(i)}P_j - A_y^{(i)}S_{xy}^{(i)}P_y) \sin \varphi^{(i)} + (b_1 - h\epsilon^{(i)}A_y^{(i)} \sum_{j=x,z} S_{xj}^{(i)}P_j - A_y^{(i)}S_{xy}^{(i)}P_y) \sin \varphi^{(i)} + (b_1 - h\epsilon^{(i)}A_y^{(i)} \sum_{j=x,z} S_{xj}^{(i)}P_j - (b_1 - h\epsilon^{(i)}A_y^{(i)}) + (b_1 - h\epsilon^{(i)}A_y^{(i)} \sum_{j=x,z} S_{xj}^{(i)}P_j - (b_1 - h\epsilon^{(i)}A_y^{(i)}) + (b_1 - h\epsilon^{(i)}A_y^{(i)} \sum_{j=x,z} S_{xj}^{(i)}P_j - (b_1 - h\epsilon^{(i)}A_y^{(i)}) + (b_1 - h\epsilon^{(i)}A_y^{(i)} \sum_{j=x,z} S_{xj}^{(i)}P_j - (b_1 - h\epsilon^{(i)}A_y^{(i)}) + (b_1 - h\epsilon^{(i)}A_y^{(i)} + (b_1 - h\epsilon^{(i)}A_y^{(i)}) + ($$

$$+a_2\cos 2\varphi^{(i)}+b_2\sin 2\varphi^{(i)}+\cdots]$$



New software tools may improve analysis

- Likelihood function from azimuthal asymmetry, event by event observed quantities
- Bayesian MCMC (uniform priors distribution)
- Estimate posterior distribution functions of polarization transfer
- Pro's
 - global picture
 - search max. no approximation
 - ad hoc priors could improve uncertainties

Similar approaches in: N. Sato/ECT* 2016 for PDFs extraction Super Rosenbluth Separation - J. C. Bernauer 2010

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GMp in JLab/HallA (E12-07-108) at high momentum transfer

Measure proton elastic cross section at $Q^2 \sim 7 - 17.5$ GeV² with high statistical precision Keep systematics under control (beam energy, target density, scattering angle)



Kinematics at smaller ϵ respect to existing SLAC data: GEp contributions smaller TPE still affecting the GMp extraction, but not the elastic cross section measurements

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... but higher QED corrections are coupled to nucleon internal structure: the probe gets correlated to the investigated sample!

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Nucleon, still a mystery!

- Mass origin
 - Confinement / QCD
 - Internal dynamics
- Spin origin
 - Orbital Angular Momentum
 - Gluon Contribution
- (Proton) radius

sample image show (actual product may vary)

QED-QCD reaction mechanisms

Current experimental programs on FF, GPD, TMD, QCD spectroscopy ... (thanks to technological progresses in Lumi, Polarization, Coherent Light production ...) & theoretical developments

are going to offer real chance to shed light on this mystery