

Studying nuclear matter under extreme conditions with the ALICE experiment at the LHC

Roberto Preghenella Istituto Nazionale di Fisica Nucleare CERN

Aperitivo Scientifico
Department of Physics and Astronomy, University of Bologna
20 May 2016

Heavy-ion physics

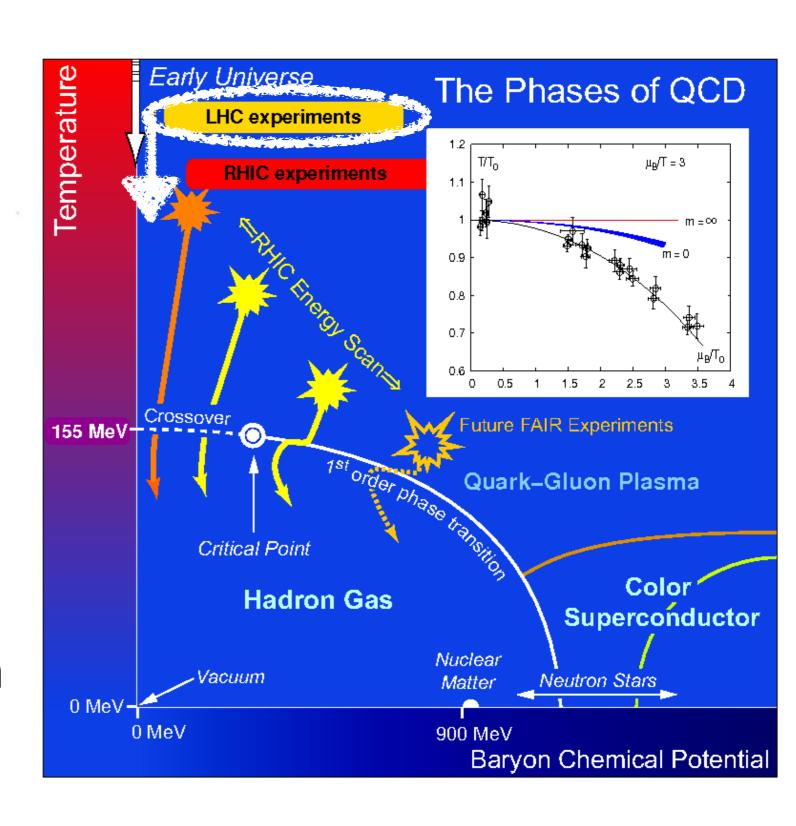
nuclear matter under extreme conditions
high temperature and energy-density

expected to undergo a phase-transition

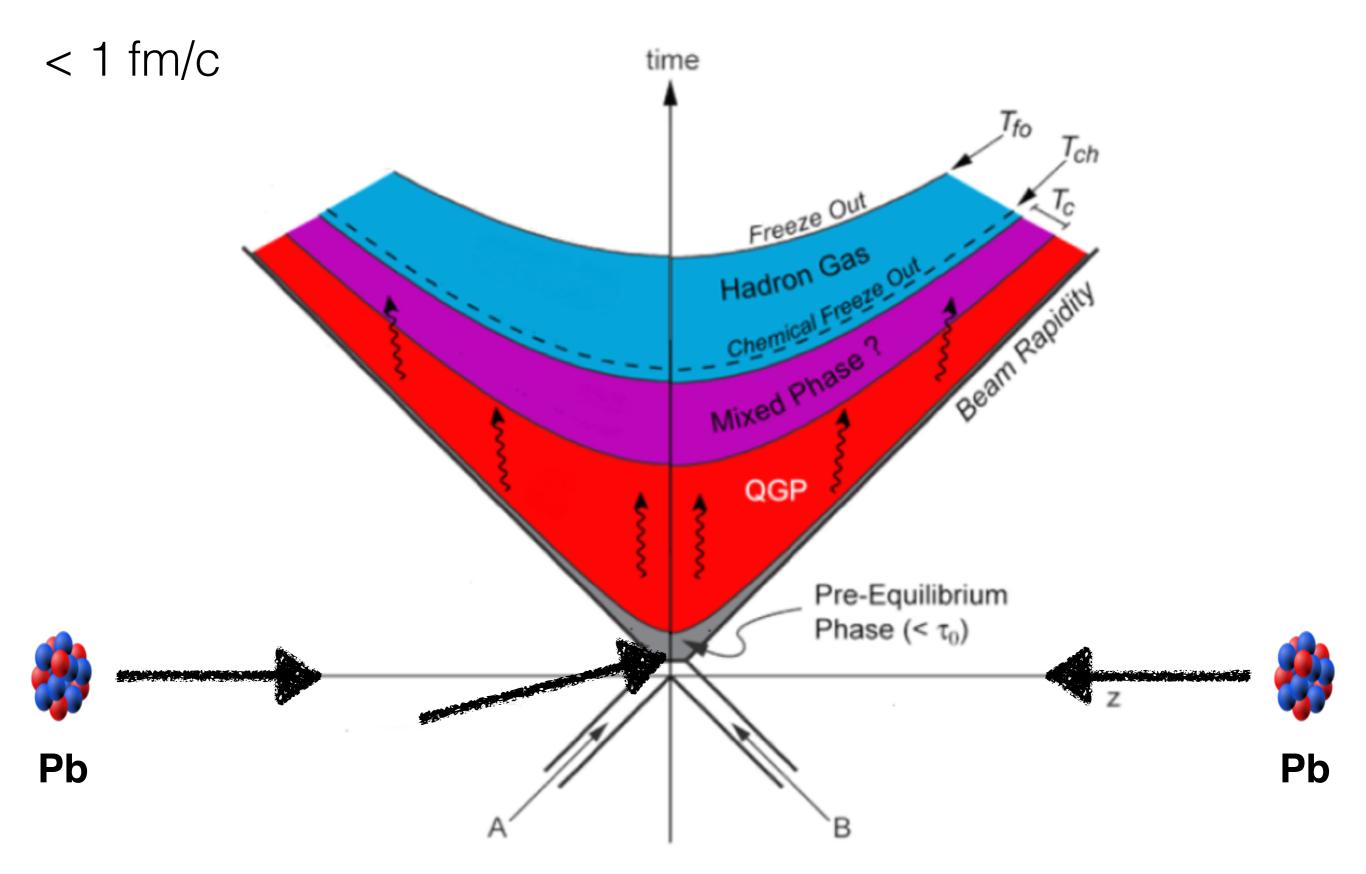
hadronic matter

Quark-Gluon Plasma (QGP)

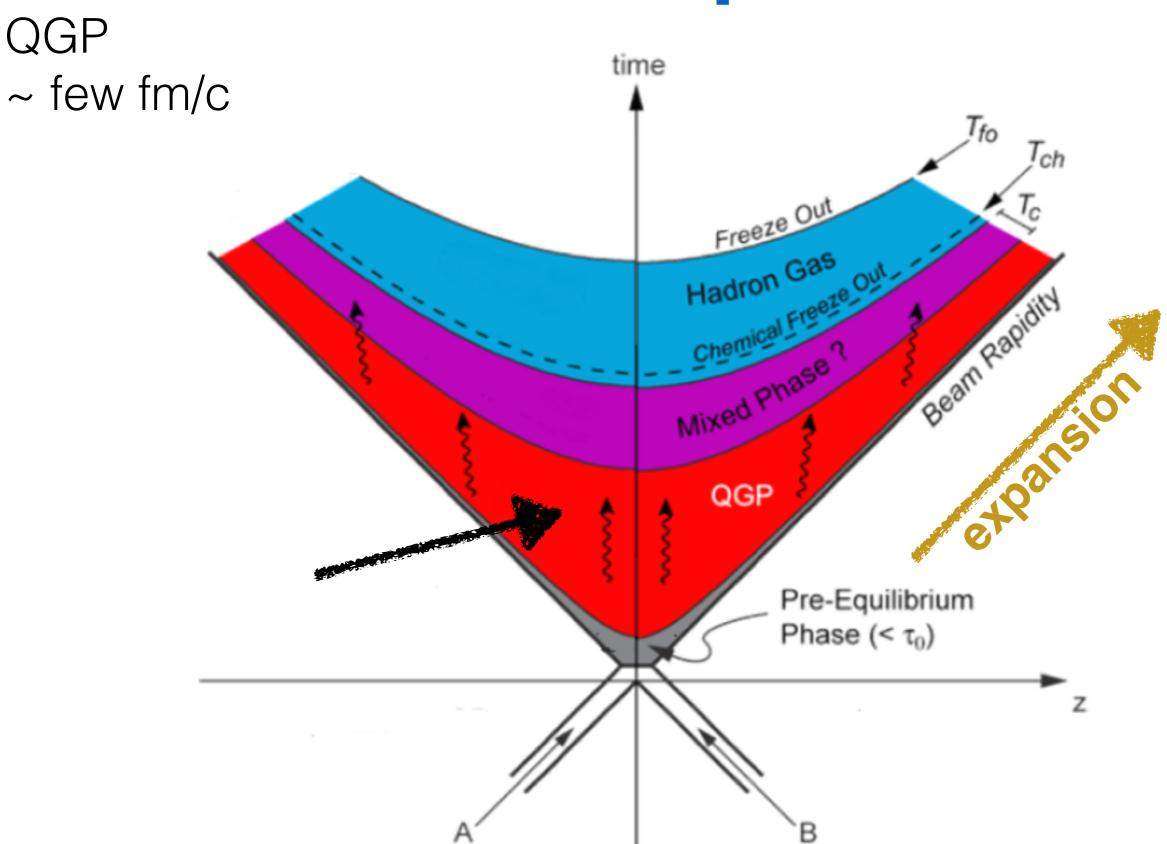
study the phase diagram and the properties of hot QCD matter



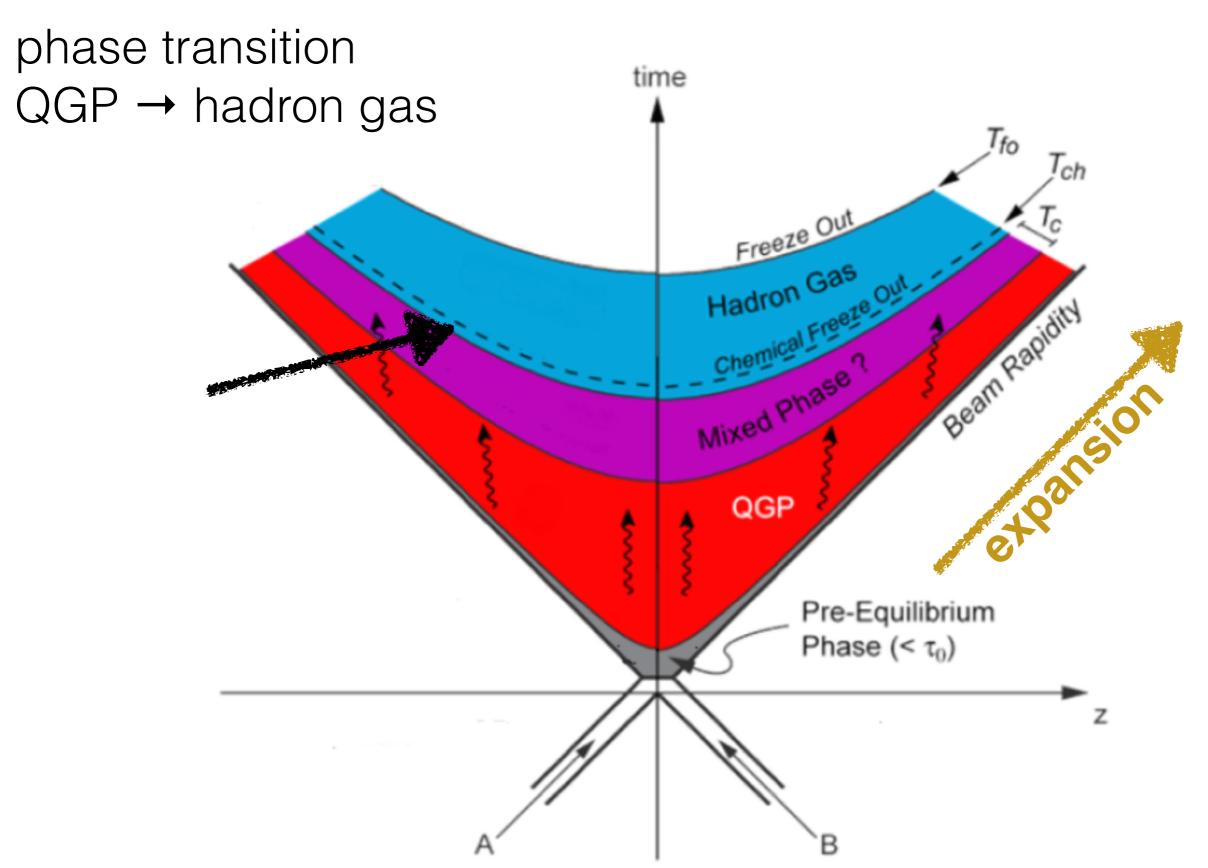
Hard scattering + thermalisation



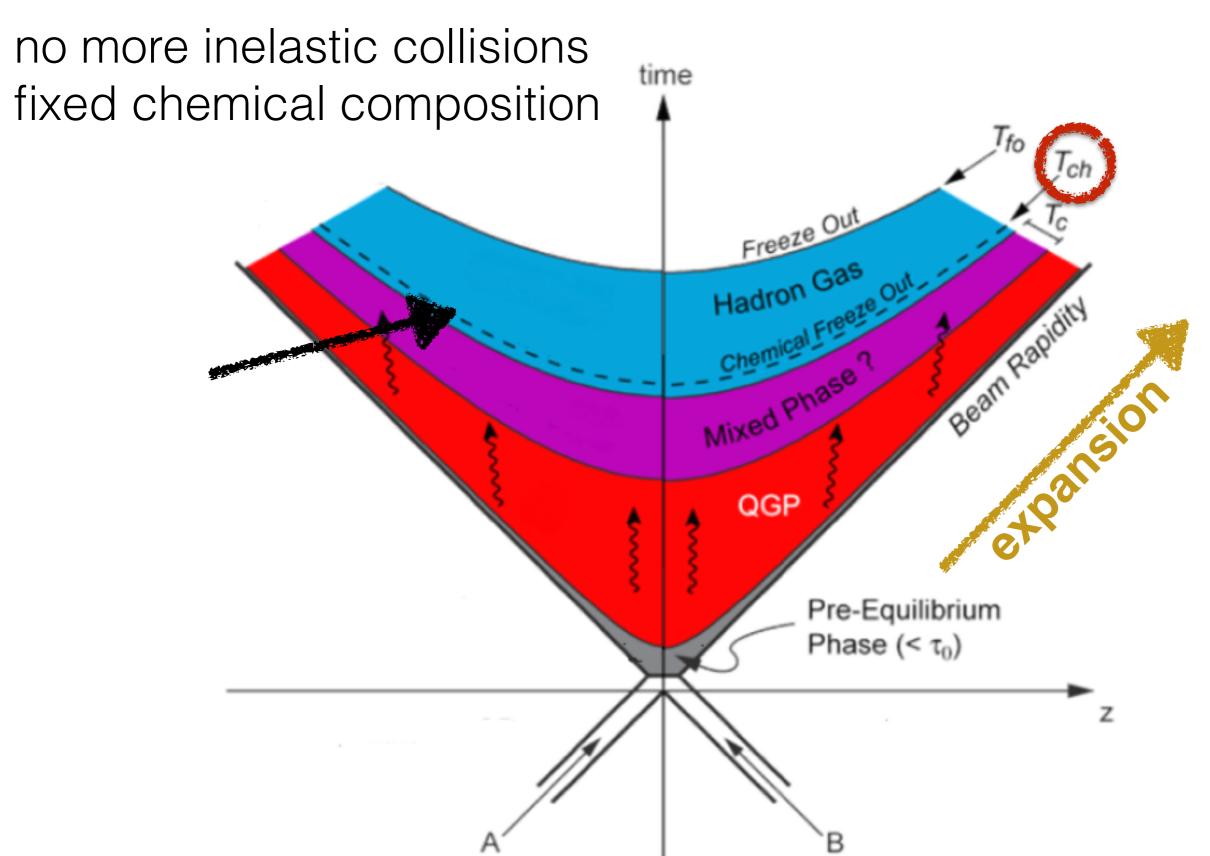
Partonic phase



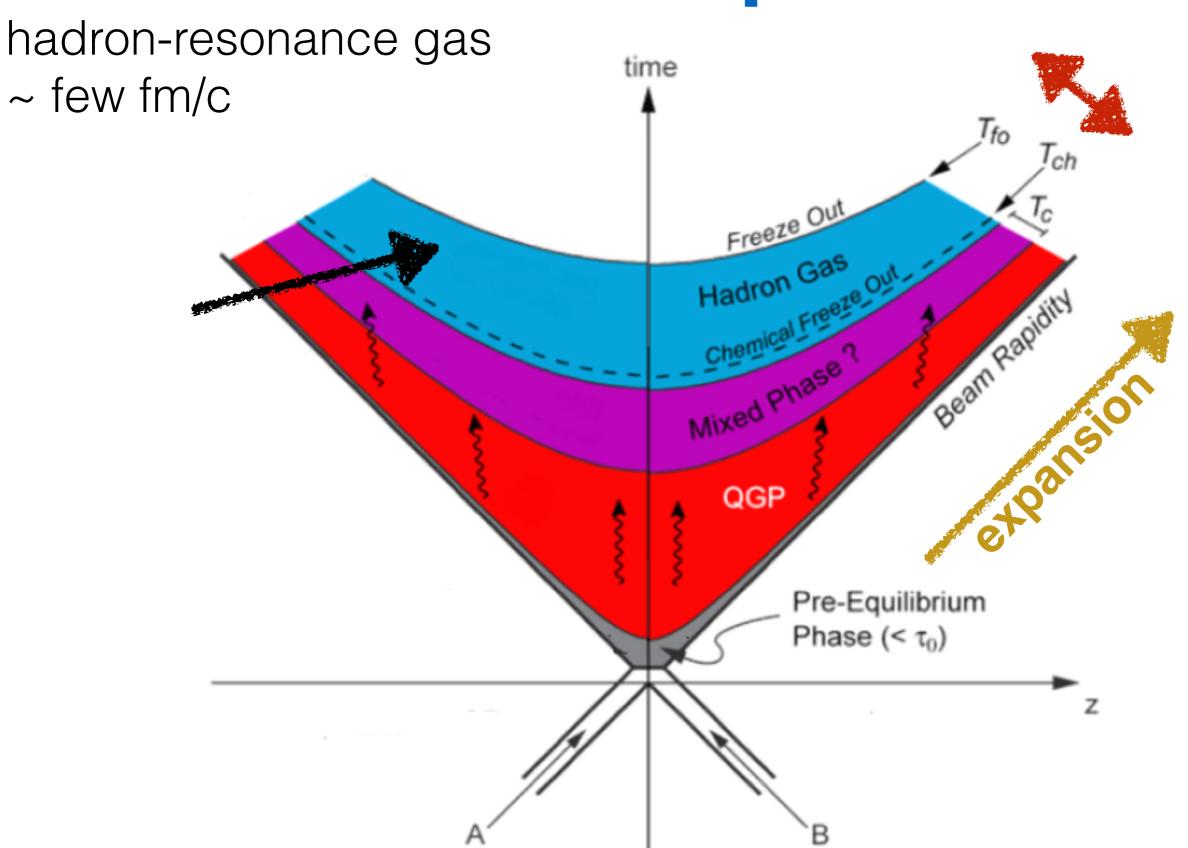
Hadronisation



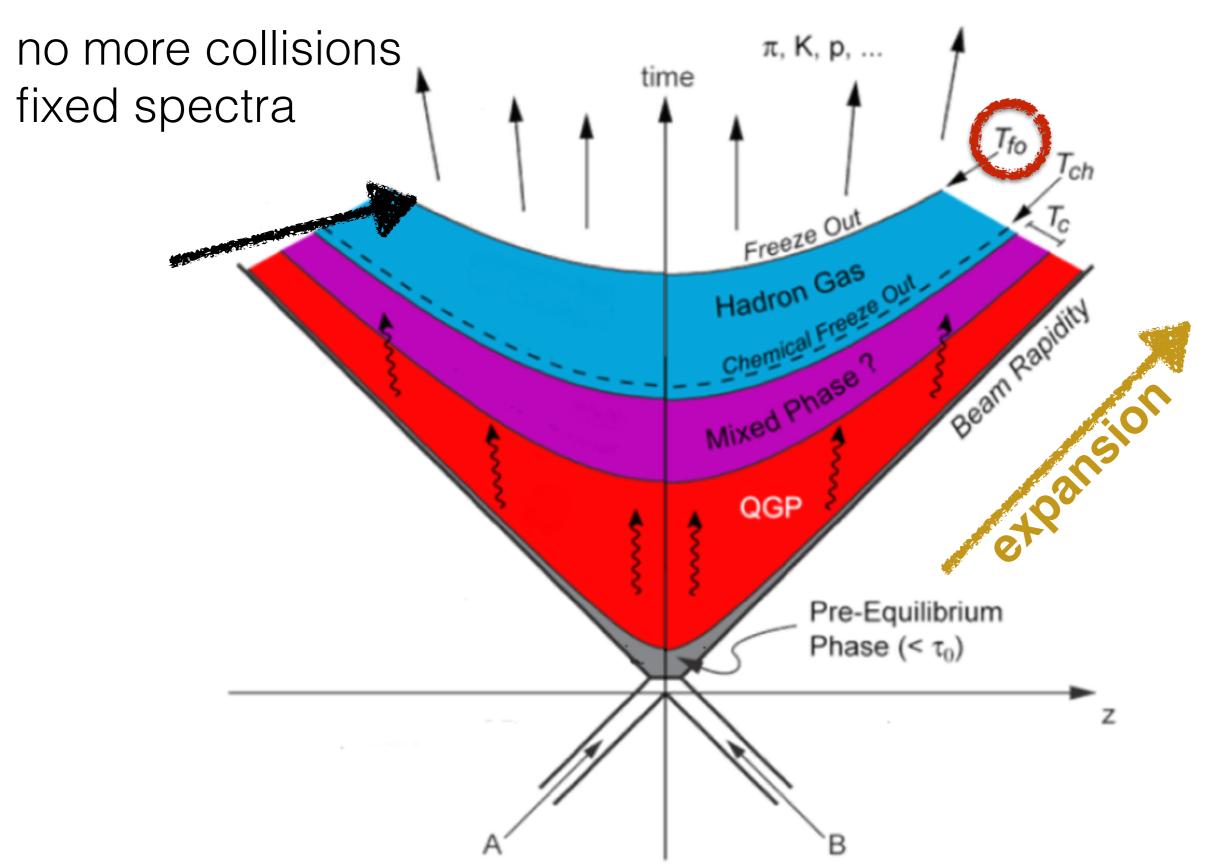
Chemical freeze-out



Hadronic phase



Kinetic freeze-out



Heavy-ion collisions at the LHC



The ALICE detector

a dedicated heavy-ion experiment at the LHC

designed to cope with

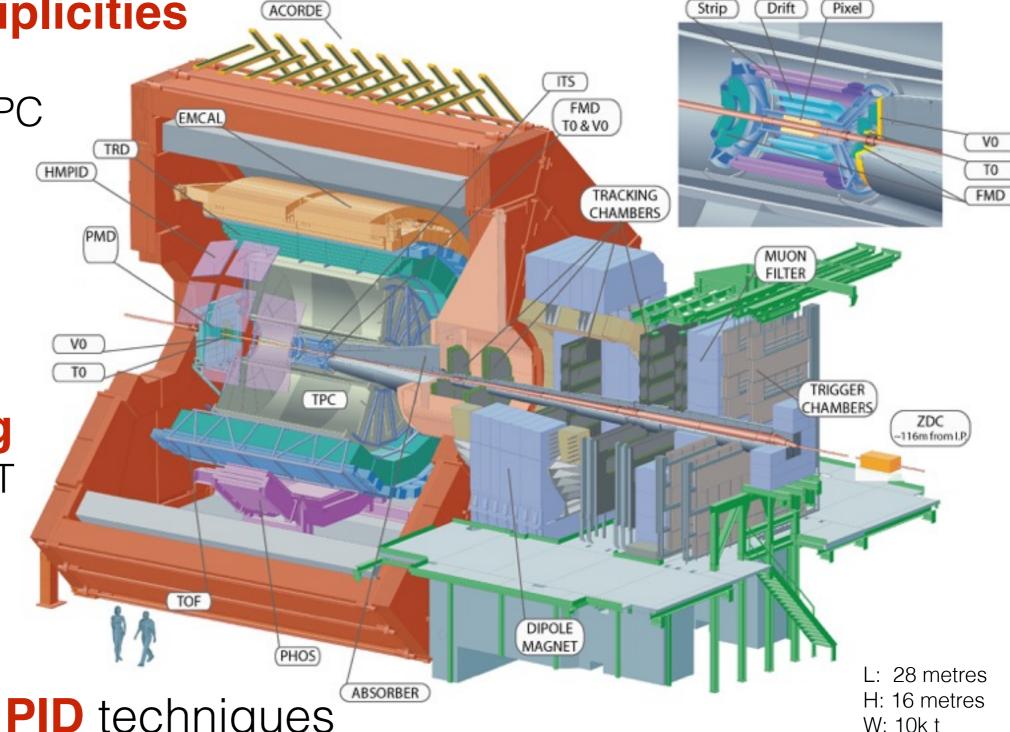
very high multiplicities

 $dN_{ch}/d\eta \le 8000$

3D tracking with TPC

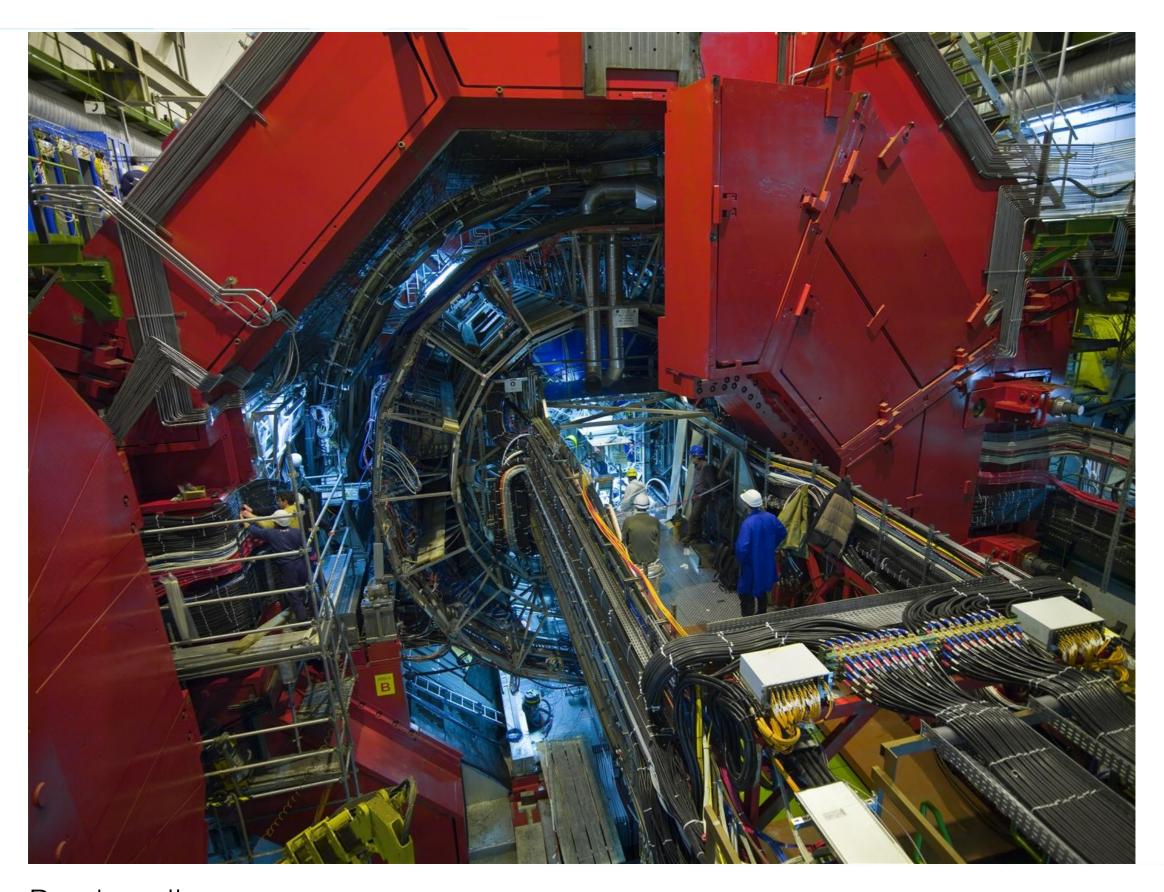
low-p_T tracking

moderate B = 0.5 Tthin materials

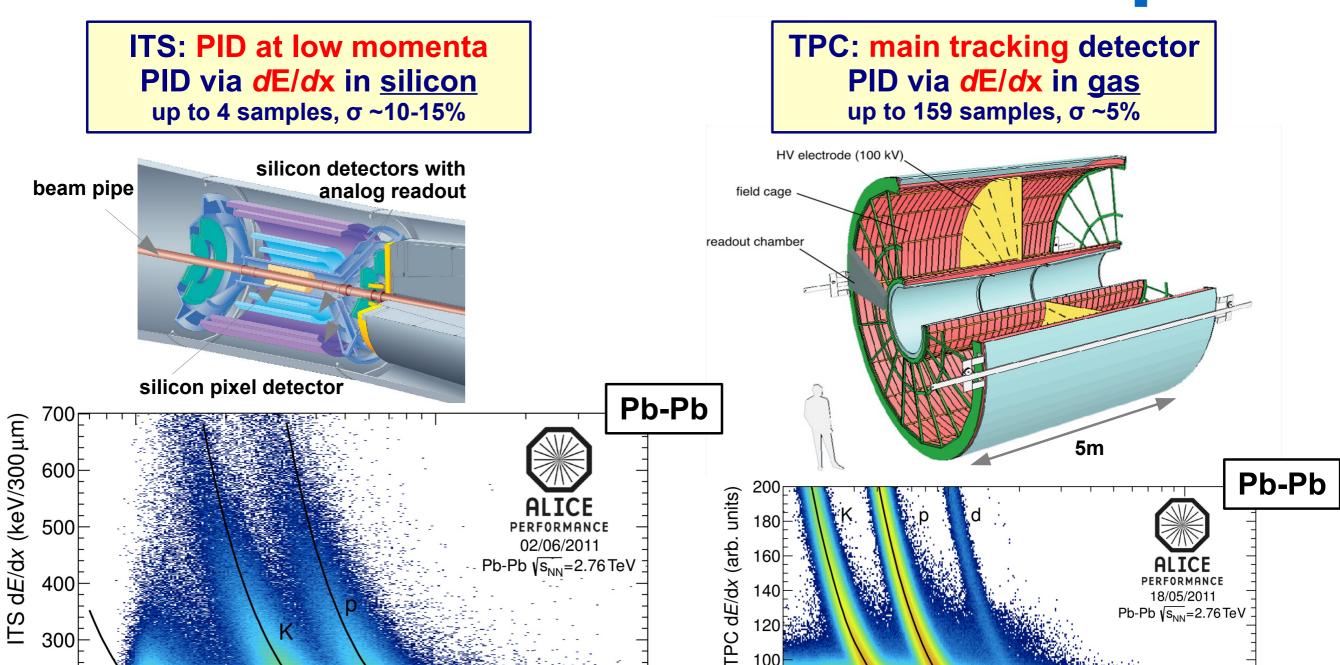


uses all known PID techniques

The ALICE detector



Particle-ID: dE/dx technique



80

20

ALI-PERF-3849

p (GeV/c)

0.2

0.3

O.07 0.1

0.3 0.4

200

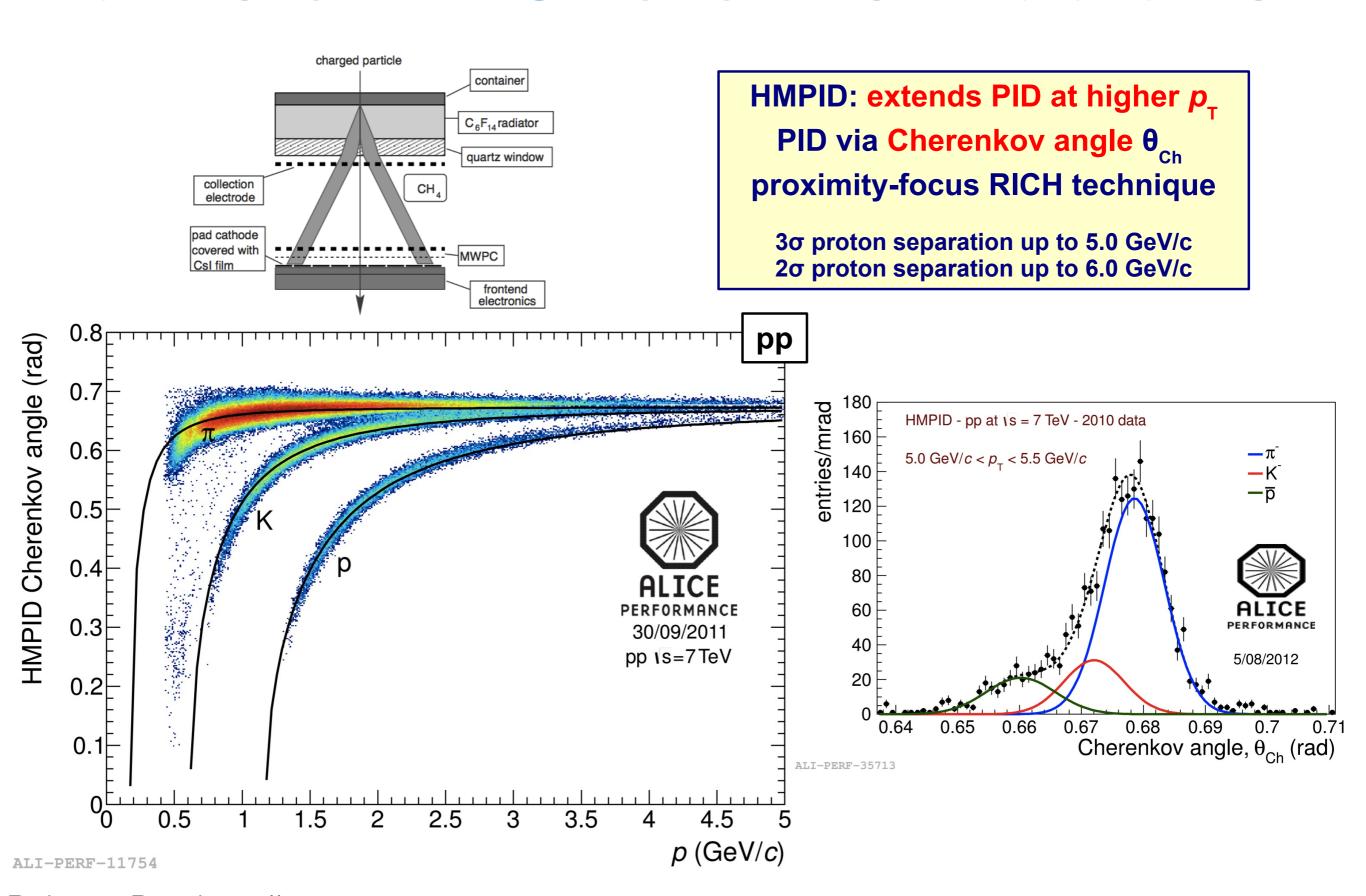
100

ALI-PERF-8369

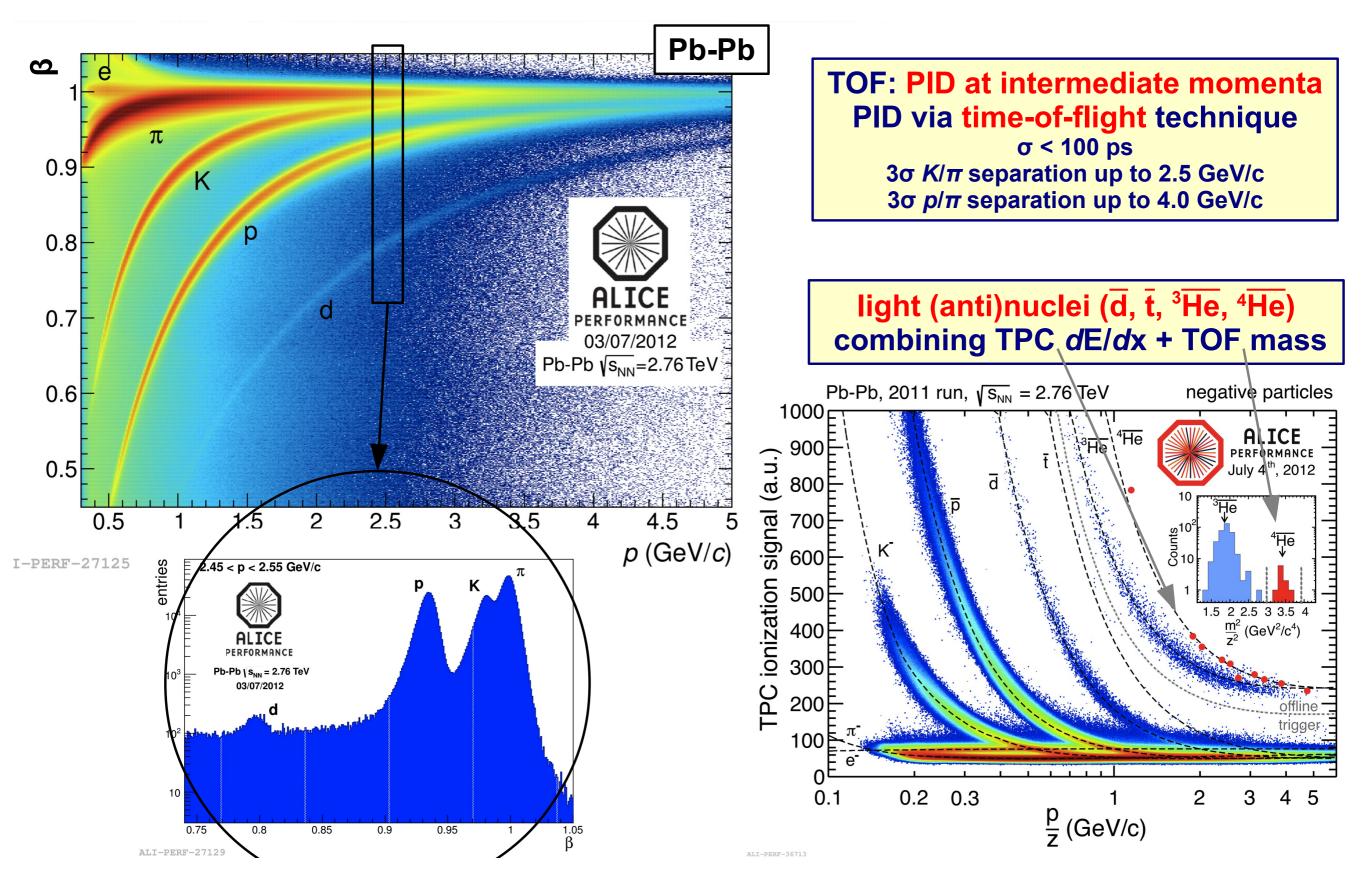
5 6 7 8 9 1 0

p (GeV/c)

Particle-ID: Cherenkov radiation



Particle-ID: time-of-flight technique

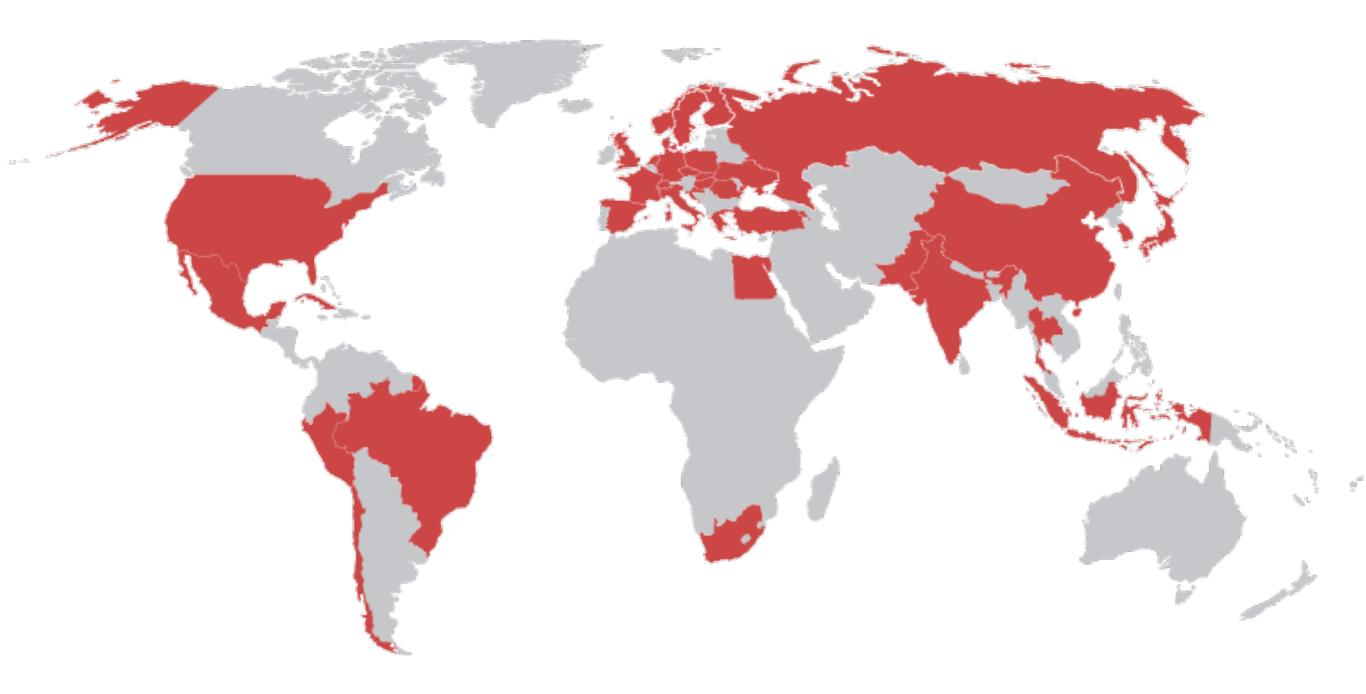


Me building the time-of-flight detector



ALICE Collaboration

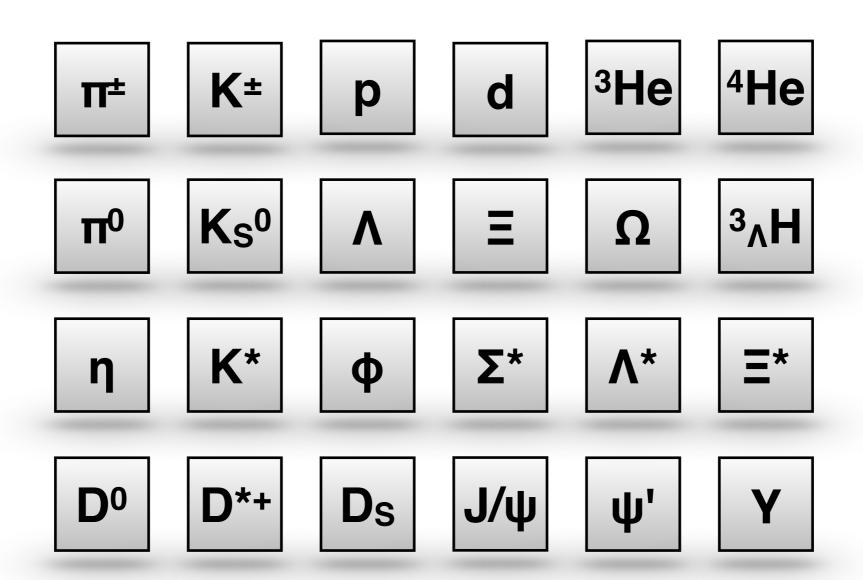
37 countries, 154 institutes, over 1500 members



goal is to study QCD phase transition and QGP properties

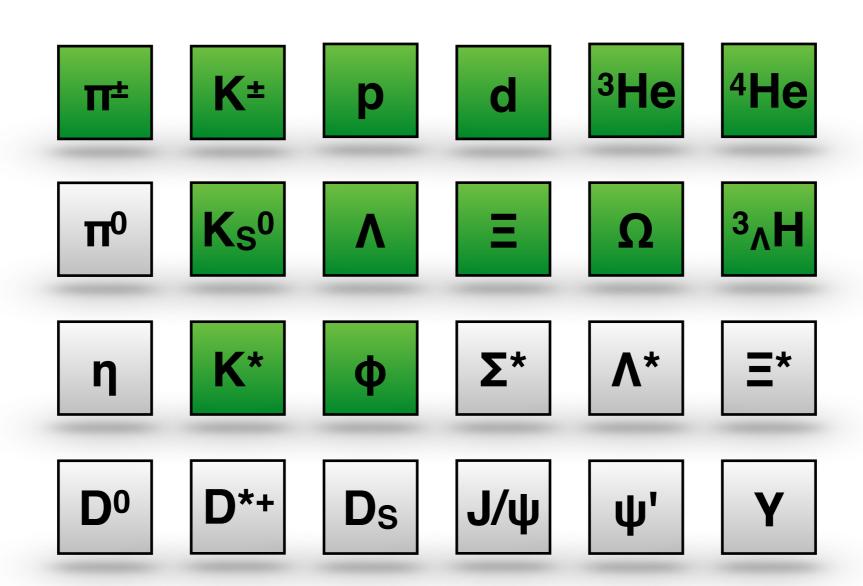
The particle zoo

ALICE has measured the production of a large number of particles, resonances and nuclei and anti-particles/nuclei



The particle zoo

ALICE has measured the production of a large number of particles, resonances and nuclei and anti-particles/nuclei



in the next slides, the focus will be on these particles

Soft heavy-ion physics at the LHC

what I will discuss in the following slides

- nucleus-nucleus collisions
 - → produce **hot nuclear matter**: QGP
 - → investigate QCD phase transition / diagram
 - → thermodynamics and collectivity
 - space-time evolution of the fireball
- proton-nucleus collisions
 - → control experiment
 - → disentangle cold / hot nuclear matter effects
 - → surprising features in high-multiplicity events

is far to be a comprehensive summary of soft heavy-ion physics

Soft heavy-ion physics at the LHC

p-Pb discussion has a similar flow, but not for today

material is in backup for whoever is interested

- nucleus-nucleus collisions
 - → produce **hot nuclear matter**: QGP
 - → investigate QCD phase transition / diagram
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 - → space-time **evolution** of the fireball
- proton-nucleus collisions
 - → control experiment
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is far to be a comprehensive summary of soft heavy-ion physics

Particle production in nucleus-nucleus collisions





Jet suppression

J. D. BJORKEN
Fermi National Accelerator Laboratory
P.O. Box 500, Batavia, Illinois 60510

FERMILAB-Pub-82/59-THY August, 1982

Energy Loss of Energetic Partons in Quark-Gluon Plasma: Possible Extinction of High $\rm p_T$ Jets in Hadron-Hadron Collisions.

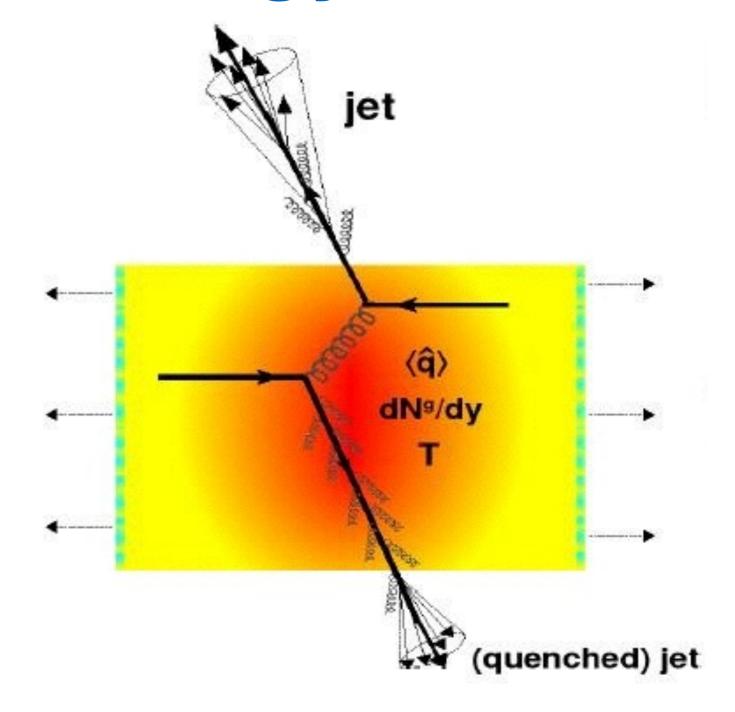
High energy quarks and gluons propagating through quark-gluon plasma suffer differential energy loss via elastic scattering from quanta in the plasma. This mechanism is very similar in structure to ionization loss of charged particles in ordinary matter. The dE/dx is roughly proportional to the square of the plasma temperature. For hadron-hadron collisions with high associated multiplicity and with transverse energy $dE_{\mathbf{r}}/dy$ in excess of 10 GeV per unit rapidity, it is possible that quark-gluon plasma is produced in the collision. If so, a produced secondary high-p, quark or gluon might lose tens of GeV of its initial transverse momentum while plowing through quark-gluon plasma

In-medium energy loss

partons produced in high Q² processes lose energy while traversing the medium

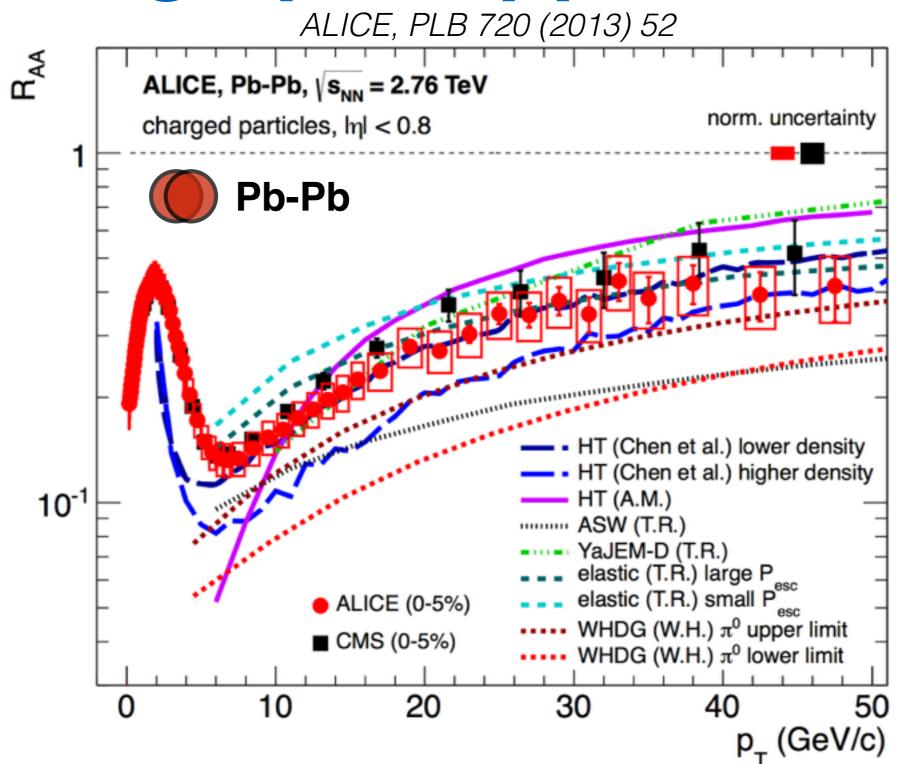
modification (suppression)
of high- p_T production
observable: nuclear
modification factor

$$R_{AA} = \frac{dN^{AA}/dp_T}{N_{coll}dN^{pp}/dp_T}$$



 $R_{AA} = 1$ for hard-processes in the absence of nuclear effects confirmed in Pb-Pb collisions at LHC (direct- γ , Z⁰ and W[±])

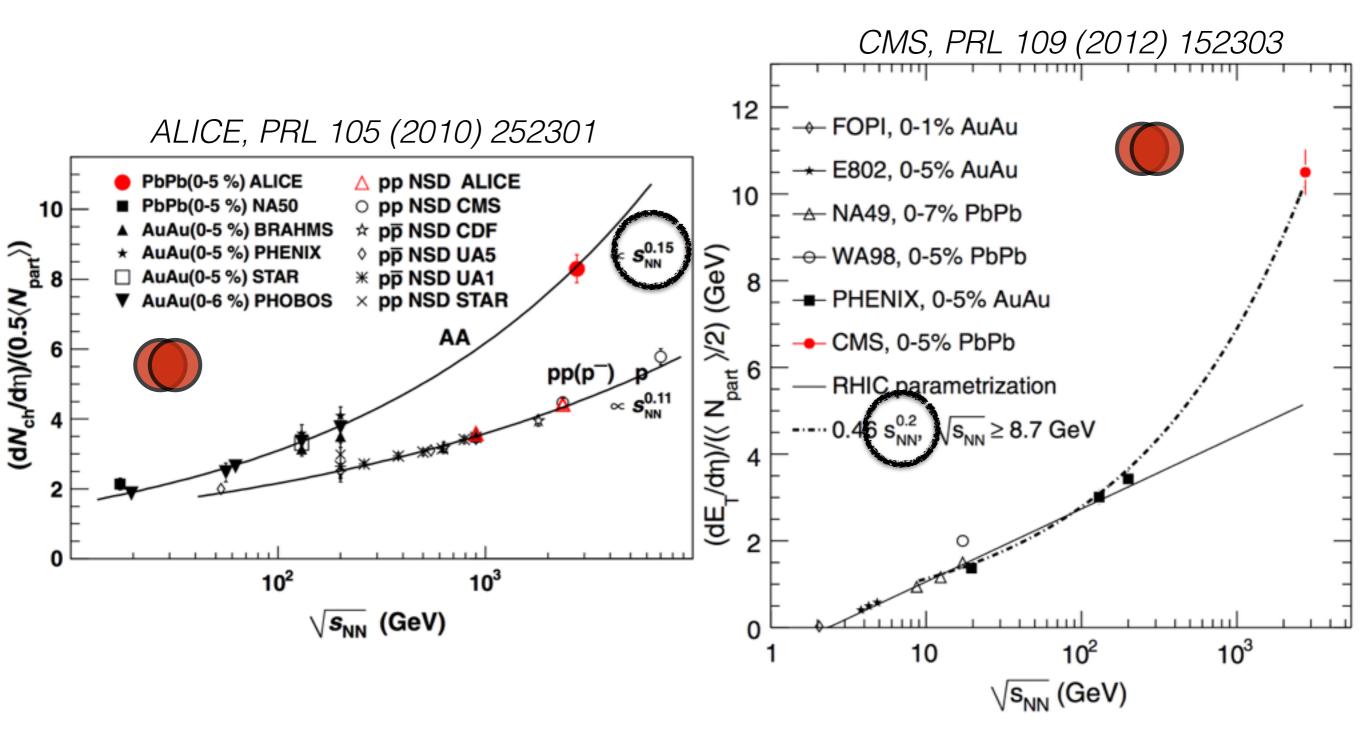
High-pt suppression



hadron production strongly modified in Pb-Pb collisions large suppression in a wide $p_{\rm T}$ range Roberto Preghenella

24

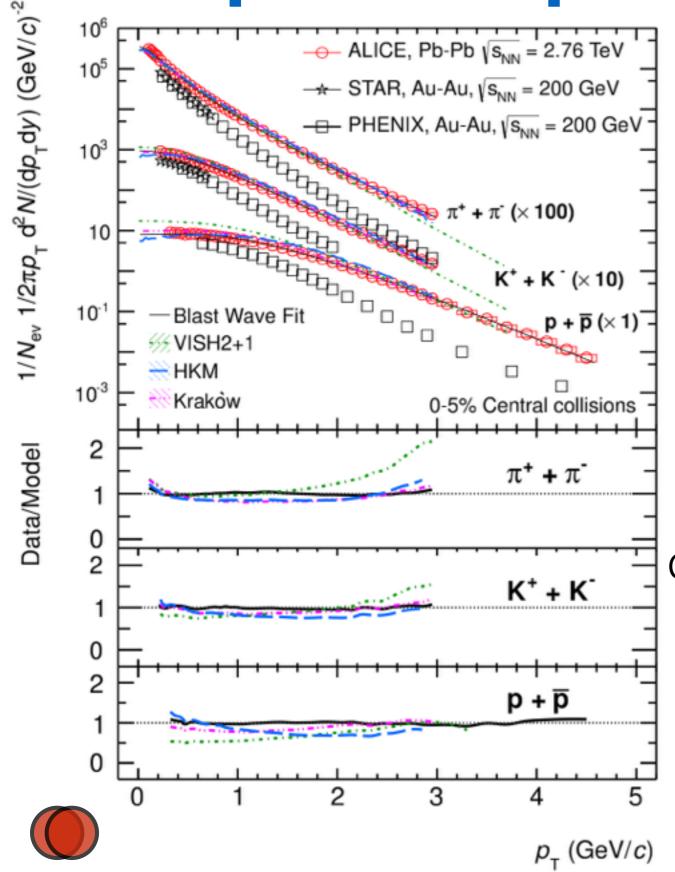
Multiplicity and transverse energy



(E_T) grows faster with energy than the multiplicity

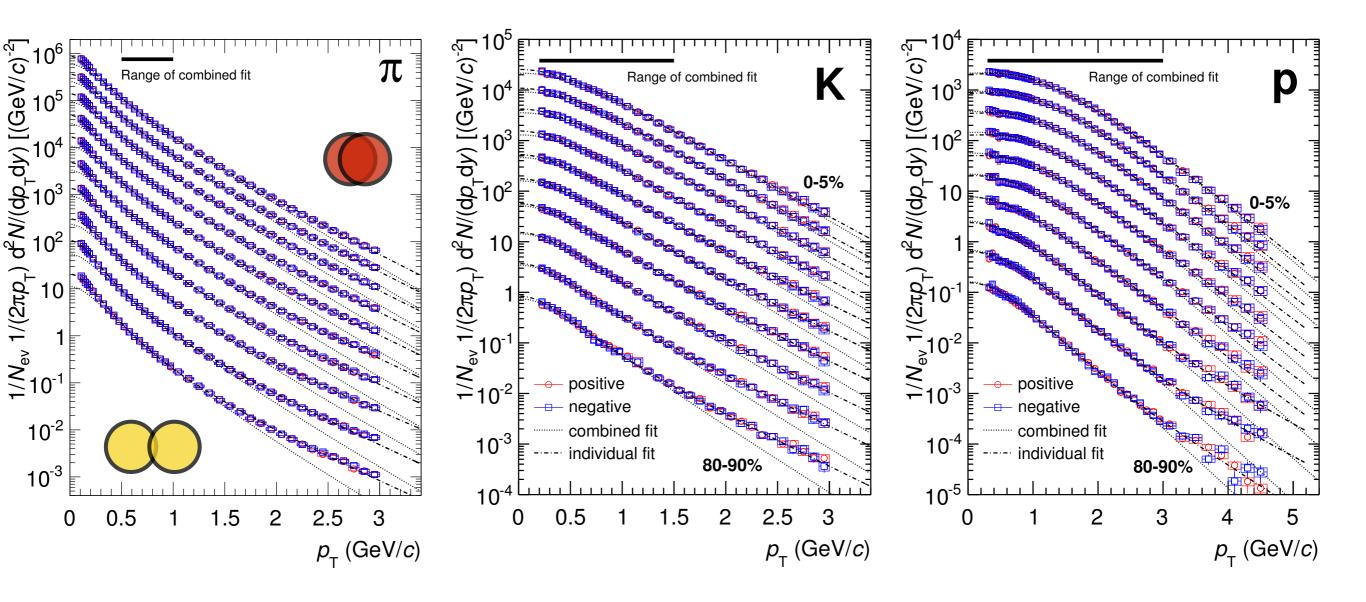
significant increase of $\langle E_T \rangle$ per particle compared to lower-energy data

Bulk particle production in Pb-Pb



transverse momentum
spectra in central Pb-Pb
collisions at the LHC are
significantly harder than in
central Au-Au collisions at RHIC

Bulk particle production in Pb-Pb

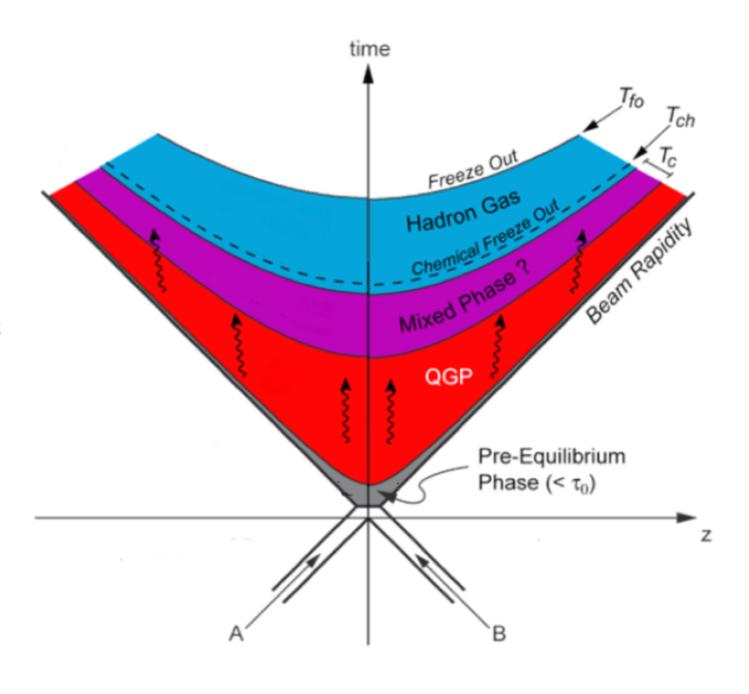


clear evolution of particle spectra → hardening with centrality more pronounced for protons than for pions mass ordering as expected from collective hydro expansion

Collective phenomena

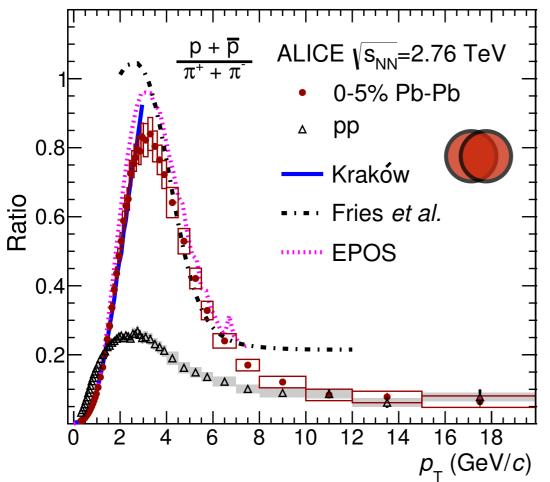
bulk matter created in high-energy heavy-ion collisions can be described in terms of hydrodynamics

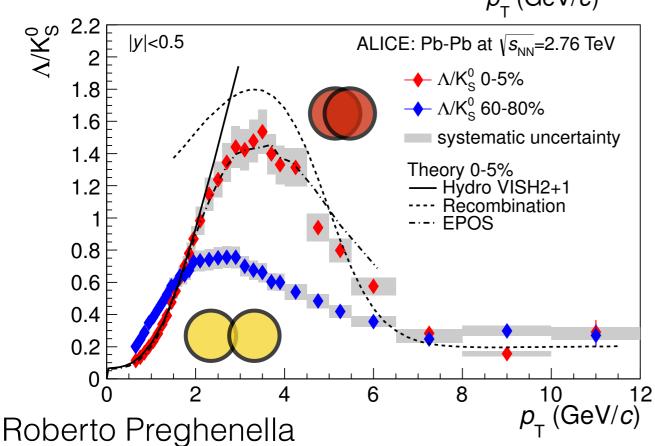
- initial hot and dense partonic matter rapidly expands
- collective flow develops and the system cools down
- phase transition to hadron gas when T_{critical} is reached resulting in



- dependence of the shape of the p_T distribution on the particle mass
- azimuthal anisotropic flow patterns (initial spatial anisotropy)

Baryon-meson enhancement in Pb-Pb





hydro model works fine for $p_T < 2 \text{ GeV}$ but deviates for higher p_T Song, PLB 658 (2008) 279

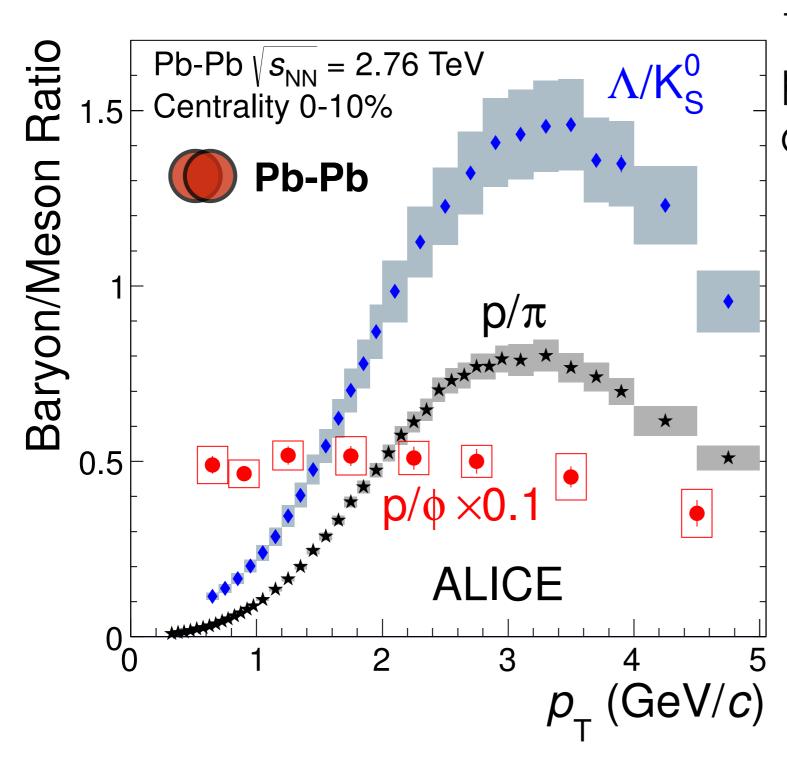
recombination approximately reproduces shape but overestimates effect

Fries, Ann.Rev.Nucl.Part.Sci. 58 (2008) 177

EPOS provides **good description** of data *Werner, PRL 109 (2012) 102301*

ALICE, PRL 111 (2013) 222301 ALICE, PLB 728 (2014) 25

p/ф spectra ratio in Pb-Pb



test baryon enhancement:

p: 938 MeV/c²

999

φ: 1018 MeV/c²

 $q\bar{q}$

spectral shapes are
very similar if particles
have similar mass
p/φ ratio is constant

the data seems to indicate that mass is the main parameter driving particle spectra

(as foreseen by hydro)

Collective phenomena

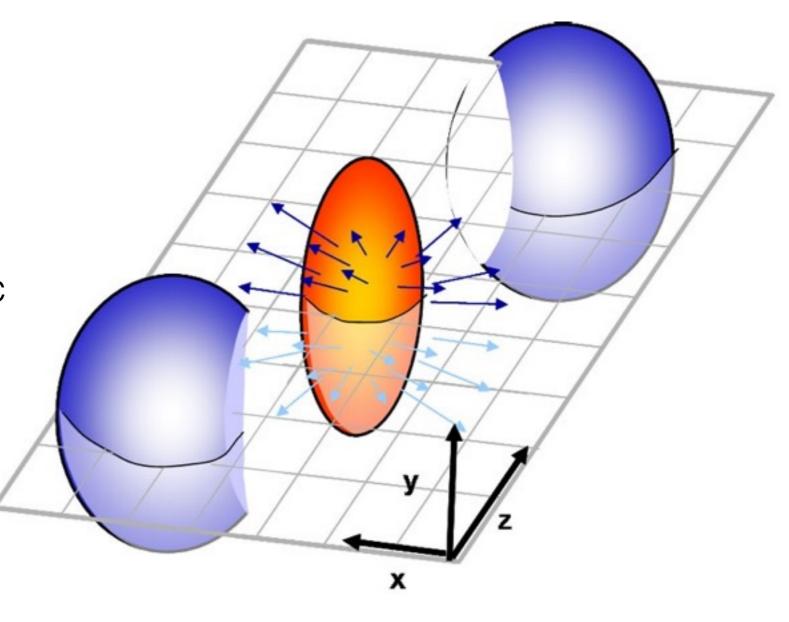
Bulk matter created in high-energy heavy-ion collisions can be described in terms of hydrodynamics

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 collective flow develops and the system cools down

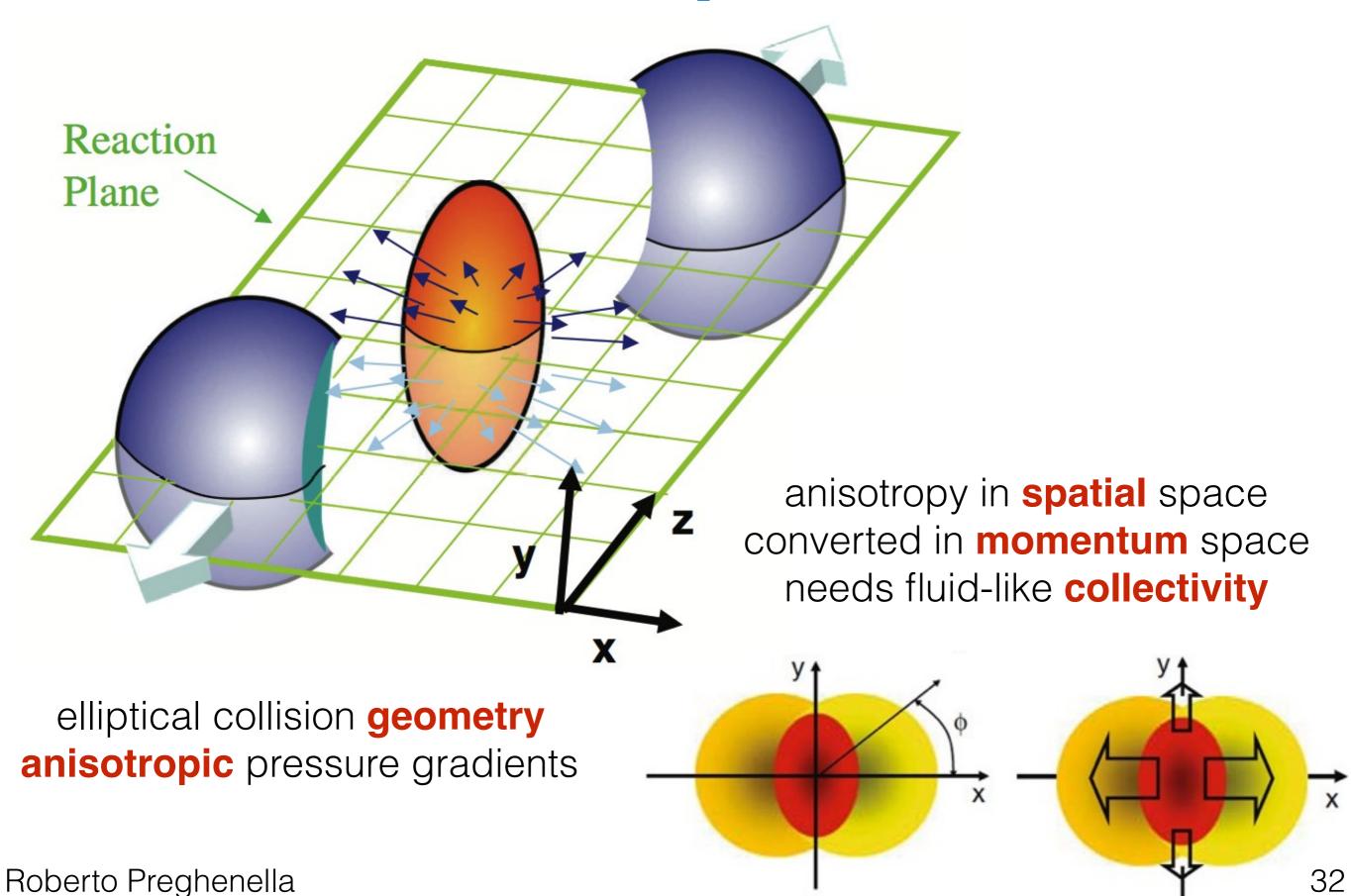
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resulting in

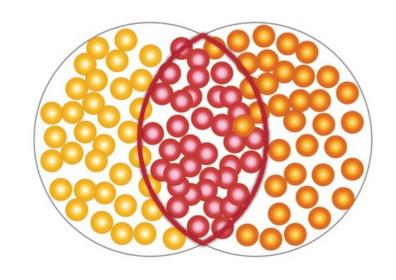


- dependence of the shape of the p_T distribution on the particle mass
- azimuthal anisotropic flow patterns (initial spatial anisotropy)

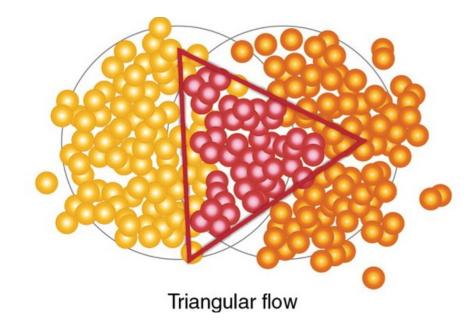
Anisotropic flow



Anisotropic flow



Elliptic flow

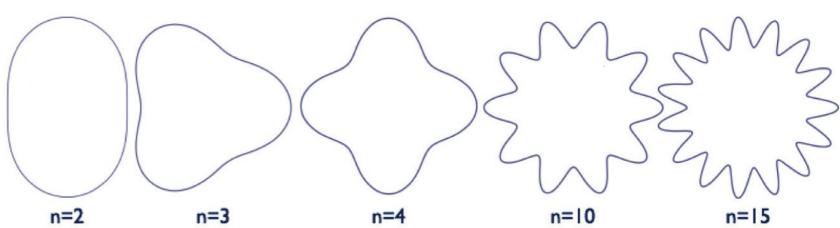


anisotropic momentum distributions

dependence can be decomposed in **Fourier series**

$$\frac{d^3N}{d^3p} = \frac{1}{2\pi} \frac{d^2N}{p_t dp_t dy} \left(1 + \sum_{n=1}^{\infty} 2v_n \cos[n(\phi - \Psi_R)] \right)$$

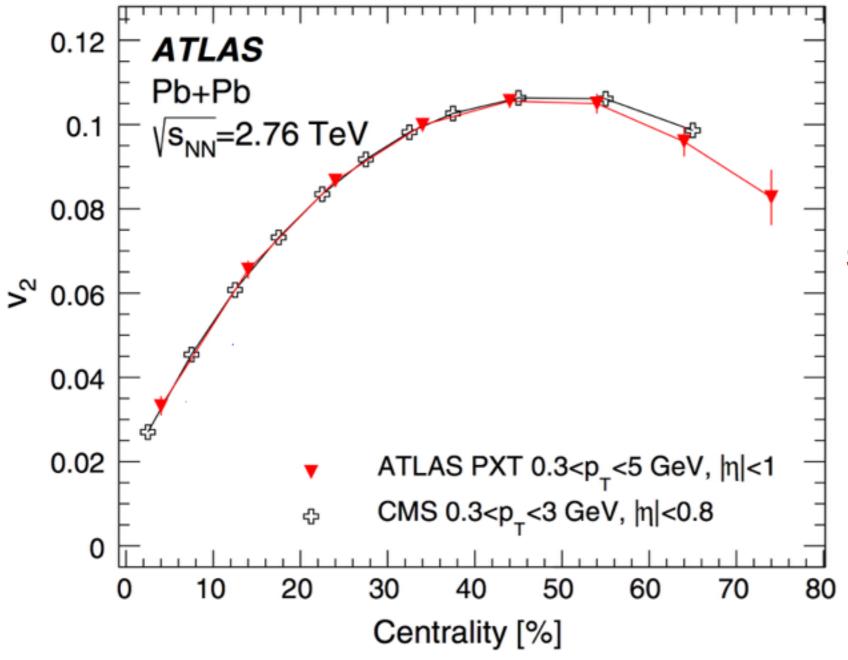
magnitude characterised by v_n coefficients

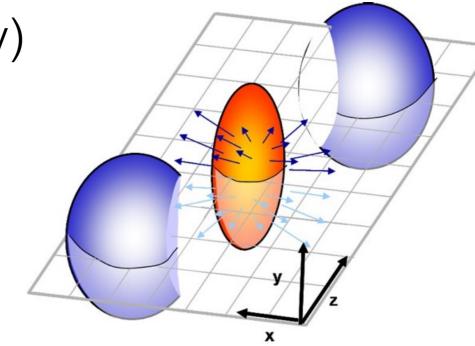


Collective anisotropic flow

spatial anisotropy (collisions geometry)

→ anisotropy in momentum space: V2



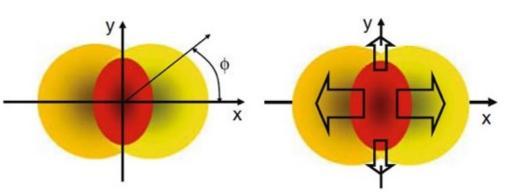


strong collective flow

persists at the LHC

flow is driven by

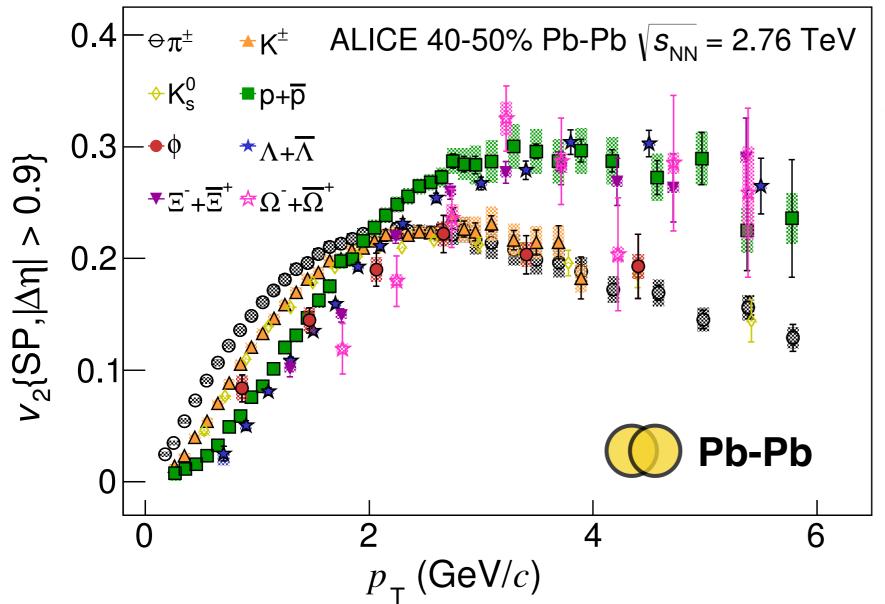
initial-state geometry

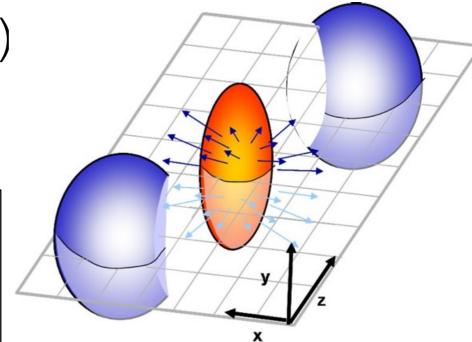


Collective anisotropic flow

spatial anisotropy (collisions geometry)

→ anisotropy in momentum space: V2





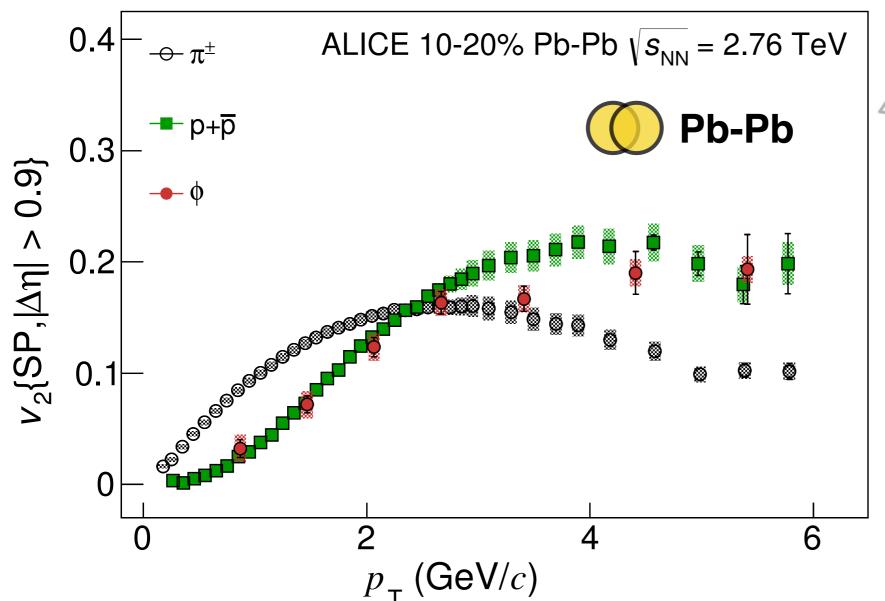
 v_2 measured for π^{\pm} , K^{\pm} , K^0_S , p, φ , Λ , Ξ , Ω

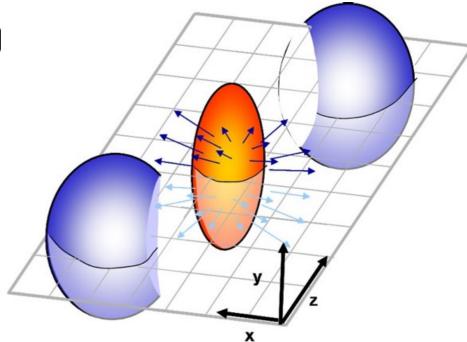
mass ordering attributed to common radial expansion velocity

Collective anisotropic flow

spatial anisotropy (collisions geometry)

→ anisotropy in momentum space: V2



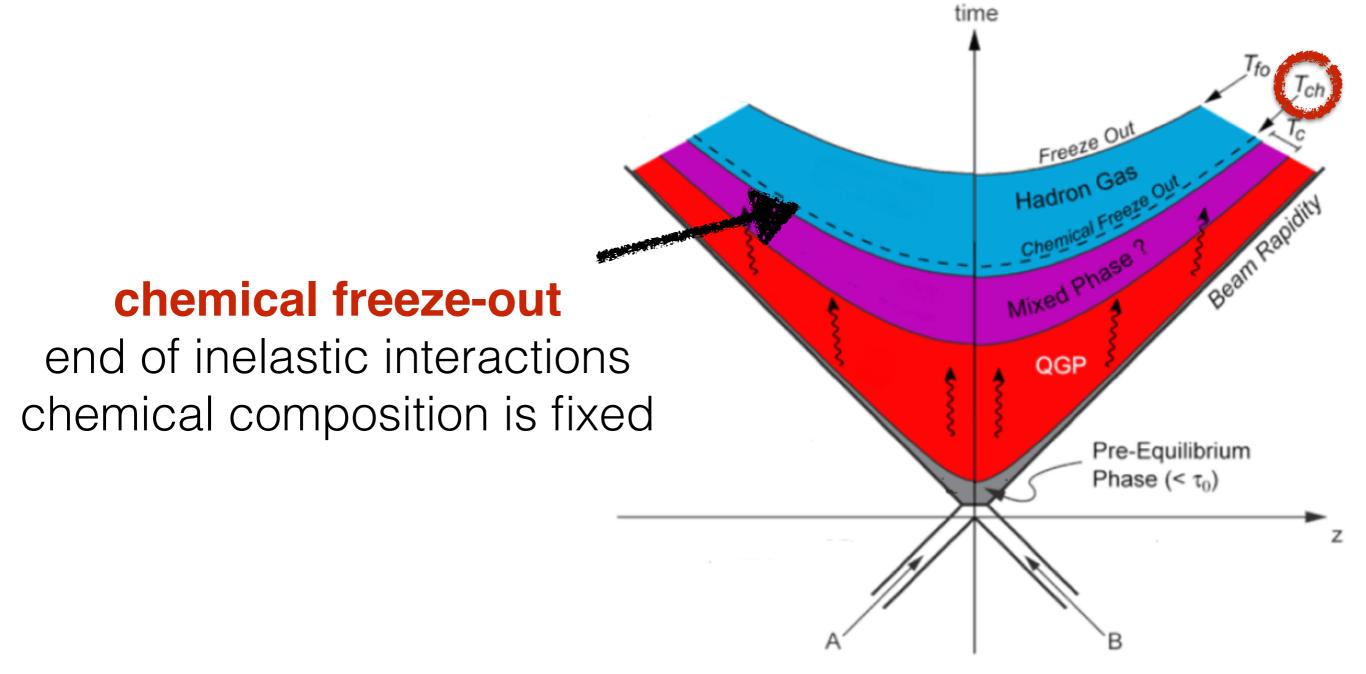


φ meson behaves like a proton

mass drives v_2 and spectra, not number of constituent quarks

Thermal model of hadron production

Chemical equilibrium achieved during or very shortly after phase transition



results of an analysis of the measured abundances allow one to get the **thermodynamic variables** (T, µ) at freeze-out

Thermal model of hadron production

Chemical equilibrium achieved during or very shortly after phase transition abundance described by Bose-Einstein or Fermi-Dirac distributions of an ideal relativistic quantum gas

$$n_j = \frac{g_j}{2\pi^2} \int_0^\infty p^2 dp (\exp\{[E_j(p) - \mu_j]/T\} \pm 1)^{-1}$$

 $E_j^2 = M_j^2 + \vec{p_j}^2$

n = particle density (N / V)

M = hadron mass

T = temperature

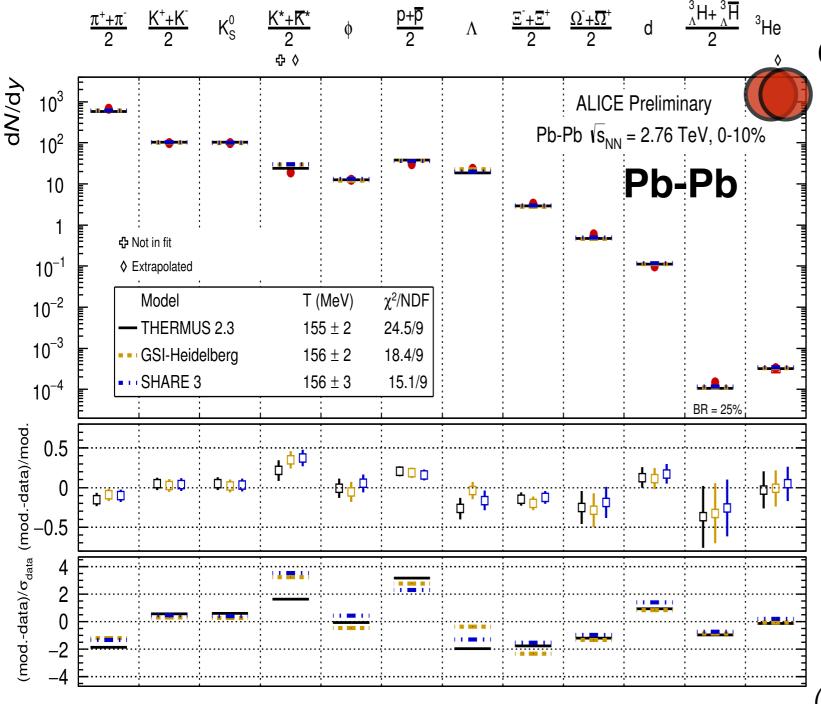
μ = chemical potential dE/dN

results of an analysis of the measured abundances allow on to set the thermodynamic variables (T, μ) at chemical freeze-out

Thermal model of hadron production

describe hadron yields as produced in chemical equilibrium

Andronic et al., NPA 772 (2006) 167



dN/dy of particle species well described in Pb-Pb x²/ndf ~ 2

same conclusion from different implementations single temperature

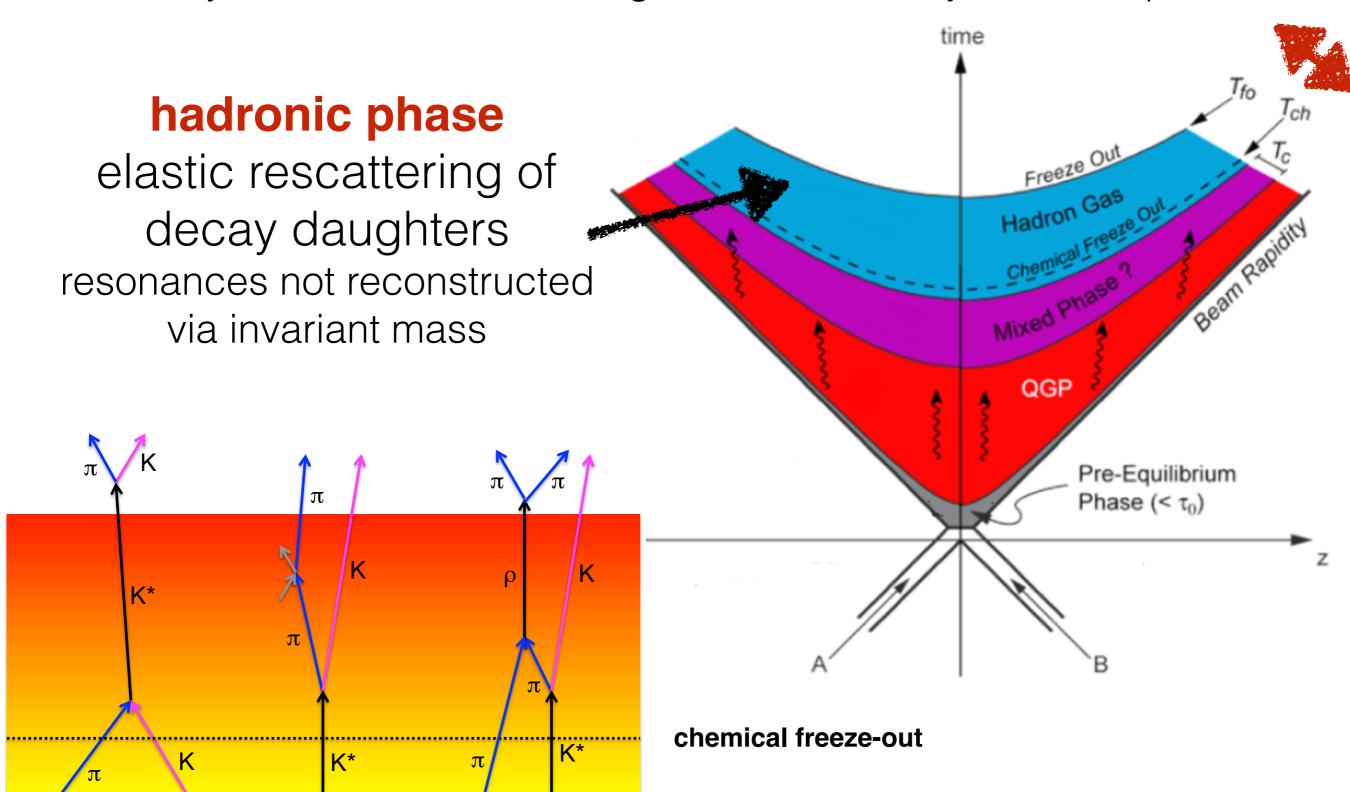
 $T_{\rm ch} \sim 156 \; {\rm MeV}$

deviations for K* and p hint at final-state interactions other mechanisms under investigation

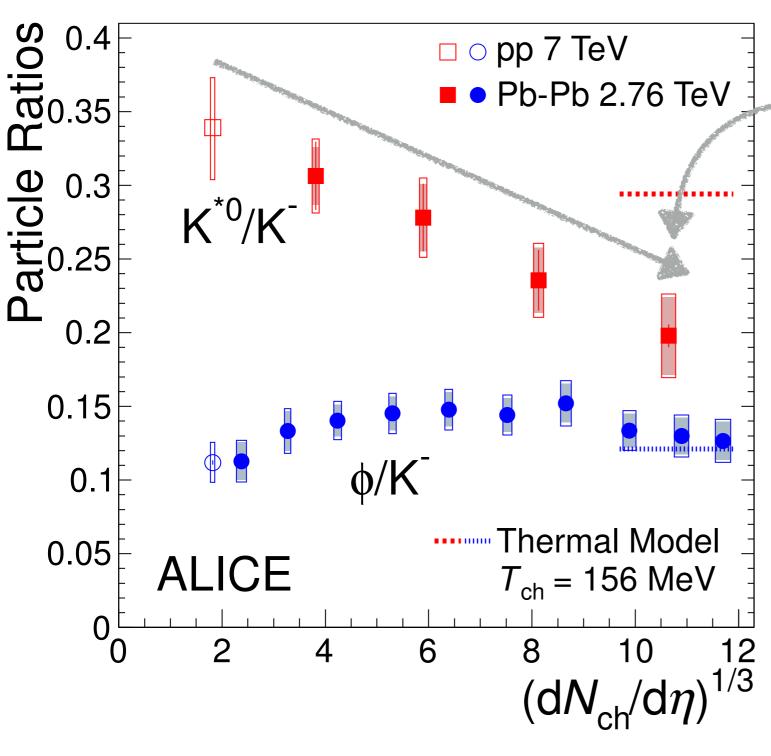
(flavour hierarchy, non-equilibrium, ...)

Interactions in the hadronic phase

measured yields of resonances might be modified by hadronic processed



K* suppression



K*/K shows clear suppression going from pp and peripheral Pb-Pb collisions to central Pb-Pb

not observed in φ/K

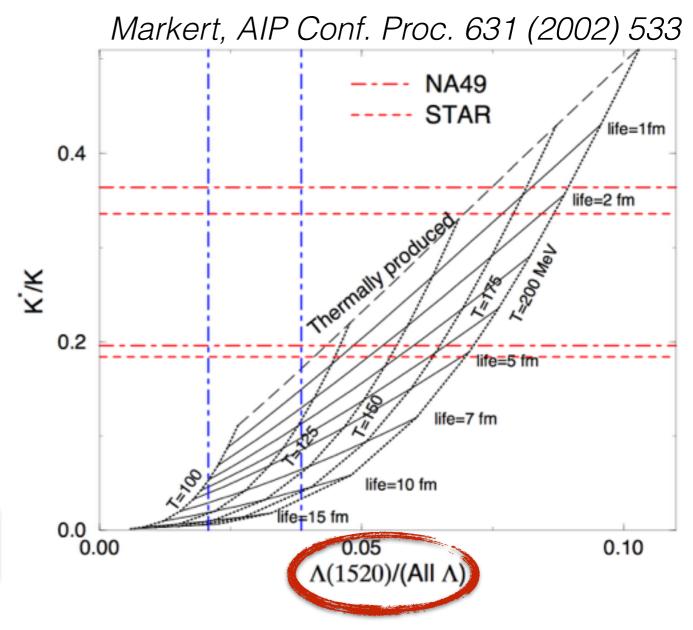
most favoured explanation re-scattering of the decay daughters with final-state hadronic medium τ_{K^*} (~4 fm/c) $\ll \tau_{\Phi}$



Properties of hadronic phase

- Model of Torrieri, Rafelski, et al.
 predicts particle ratios as functions
 of chemical freeze-out temperature
 and lifetime of hadronic phase
- Model Predictions:

Torrieri/Rafelski* no re-scattering $T_{ch} = 156 \text{ MeV}$ Prediction: $K^{*0}/K^{-} = 0.35$ Torrieri/Rafelski* no re-scattering measured K^{*0}/K^{-} Torrieri/Rafelski* $T_{ch} = 120\pm7 \text{ MeV}$ Torrieri/Rafelski* Prediction: $T_{ch} = 120\pm7 \text{ MeV}$ Torrieri/Rafelski* Prediction: Lifetime $\geq 2 \text{ fm/c}$ $T_{ch} = 156 \text{ MeV}$



important to measure other resonances to determine the upper limit for the lifetime

Particle production in proton-nucleus collisions





Summary

detailed study of the properties of hot QCD matter with nucleus-nucleus collisions at the LHC

signatures of thermalisation, final-state effects and collectivity

particle production evolves with increasing system size

baryon and K* suppression, strangeness and deuteron enhancement central Pb-Pb well described by GC thermal models, $T_{ch} = 156 \text{ MeV}$

bulk particle production in proton-nucleus shows nucleus-nucleus features and signatures of collectivity

non-zero elliptic flow, mass-dependence of p_T spectra and v_2 enhanced production of strange and multi-strange hadrons interesting! need more investigation on small systems

many more results and a bright future

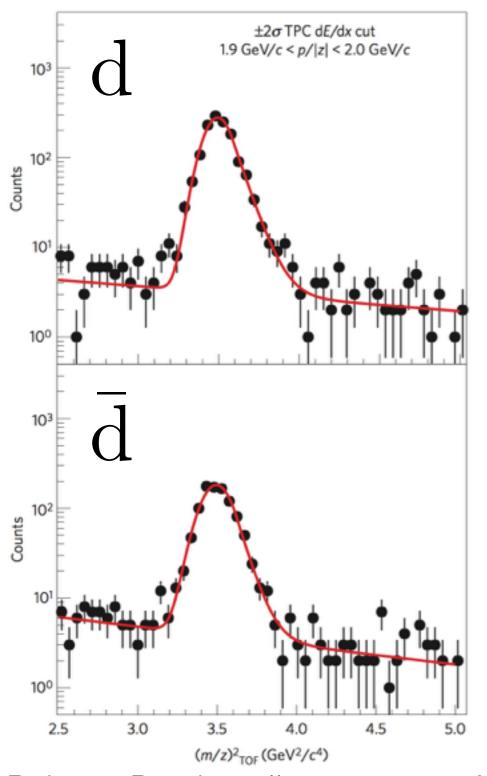
new data and more ideas for LHC Run-2

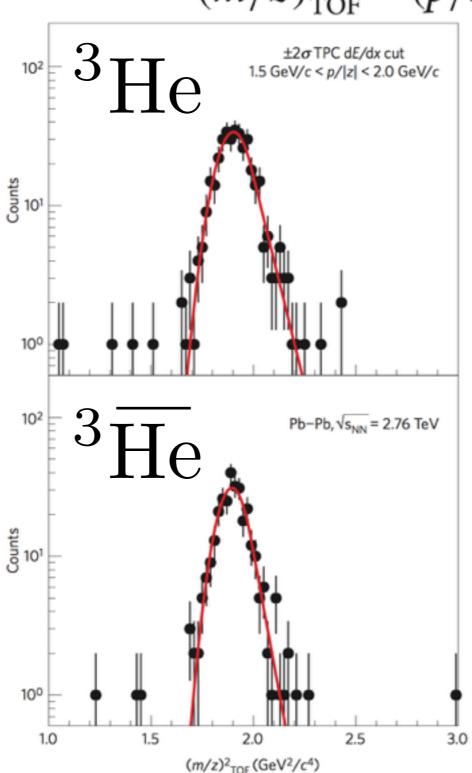
More than just Heavy-ion physics

CPT invariance in nuclear systems

precision measurement of nuclei mass with time-of-flight

$$(m/z)_{\text{TOF}}^2 = (p/z)^2 [(t_{\text{TOF}}/L)^2 - 1/c^2]$$





makes use of

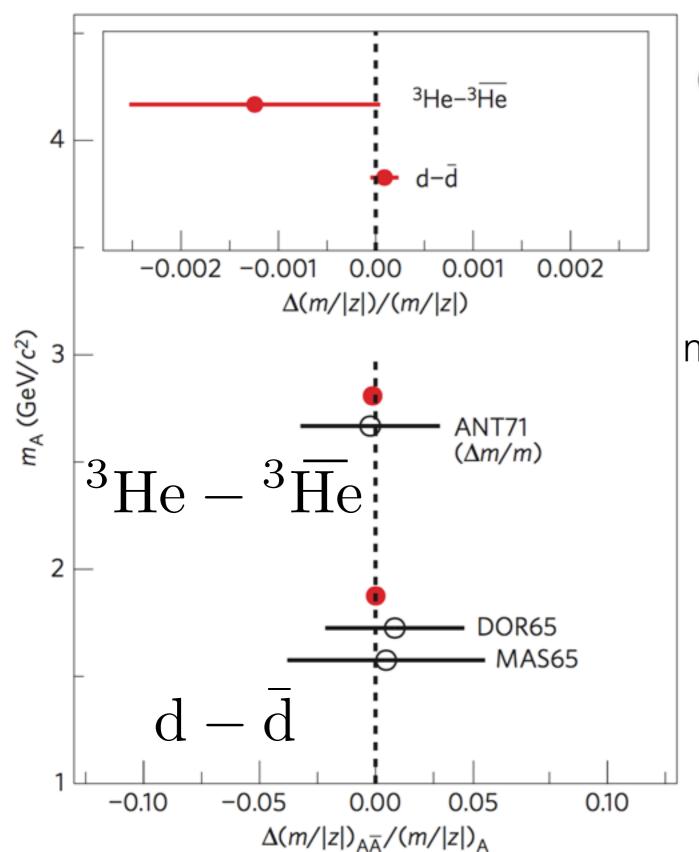
heavy-ion collisions as an efficient source of nuclei and anti-nuclei

combined with

high-precision tracking and identification capabilities of ALICE

Nature Physics 11 (2015) 811

CPT invariance in nuclear systems



$$(m/z)_{\text{TOF}}^2 = (p/z)^2 [(t_{\text{TOF}}/L)^2 - 1/c^2]$$

measuring mass differences

rather than absolute values

→ <u>reduced uncertainties</u> momentum, time-of-flight, track length

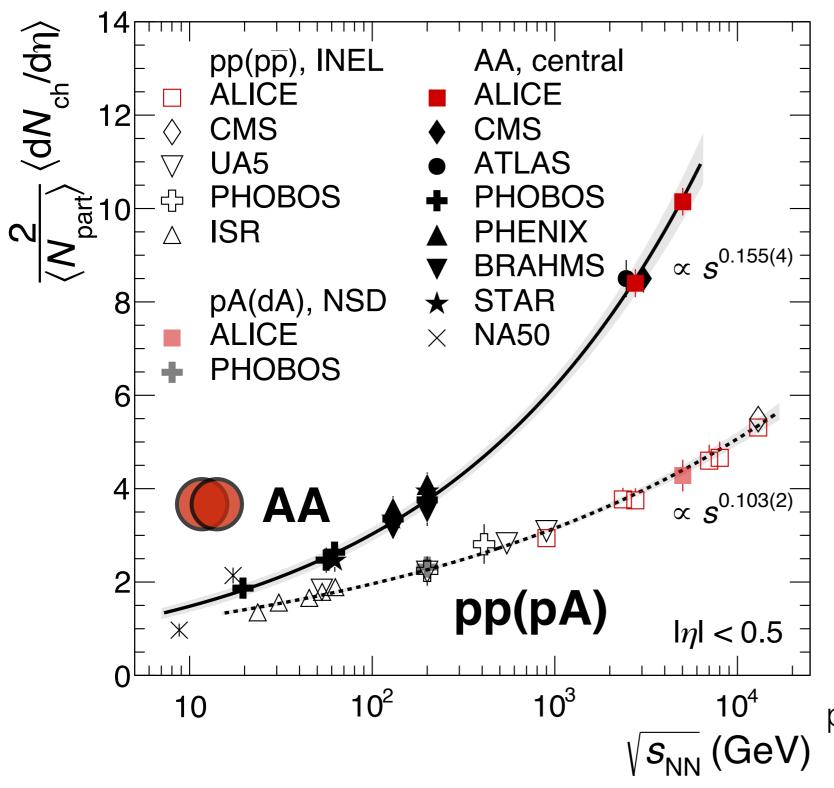
these results are

the highest precision direct measurement of the mass difference of nuclei/anti-nuclei improved by one to two orders of magnitude wrt. previous measurements (dating back to 1965 and 1971)

First results from LHC Run-2

Charged particles in Pb-Pb@5.02 TeV

centre-of-mass energy dependence



charged-particle multiplicity density

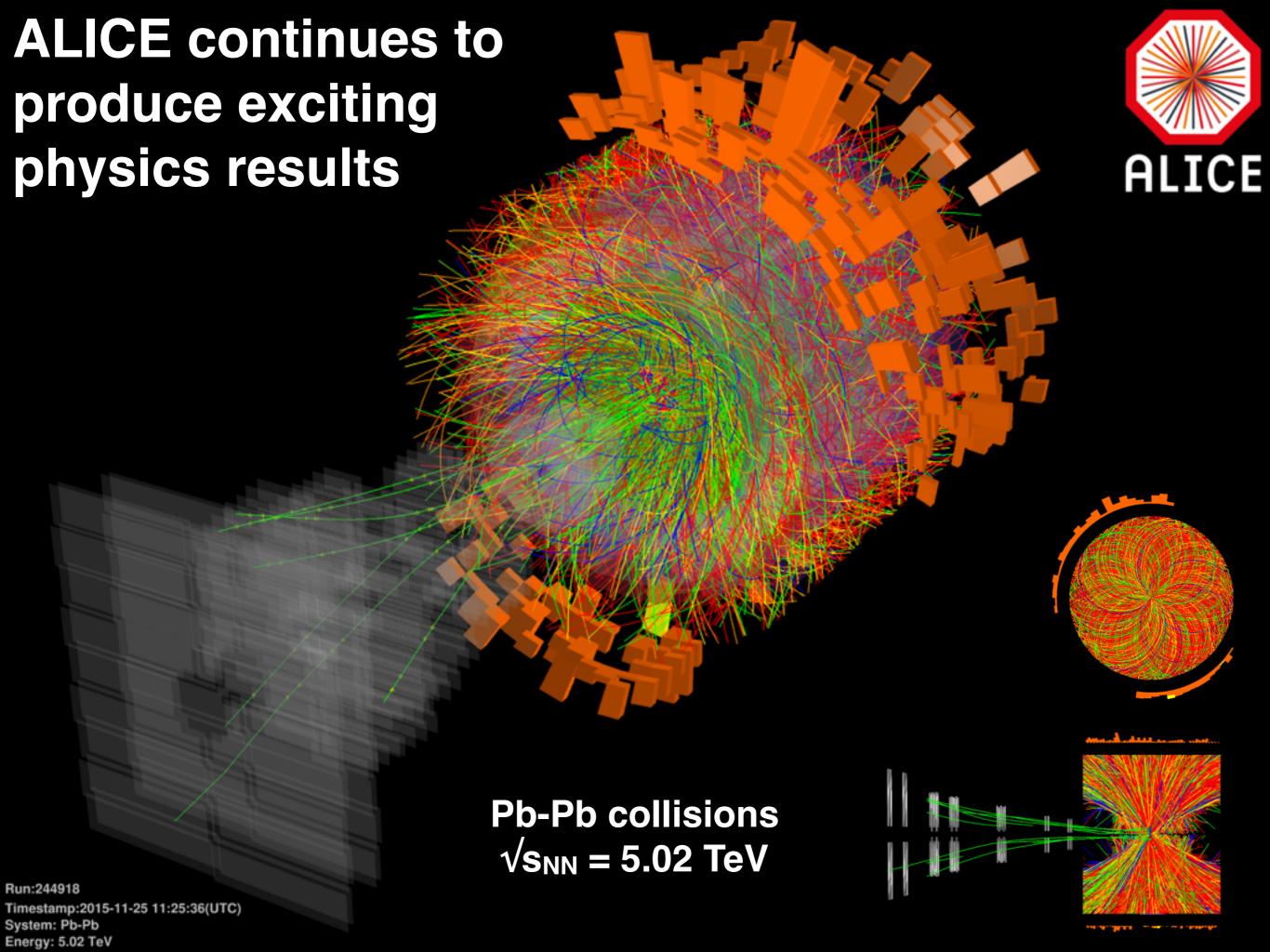
at mid-rapidity, |η| < 0.5 reaches a value of 1943 ± 56 in most central collisions

much stronger √s dependence than pp

2.4x larger charged-particle

multiplicity than p-Pb

at same energy
scaled by the average number of participating nucleon pairs \(N_{part} \)/2



Particle production in proton-nucleus collisions



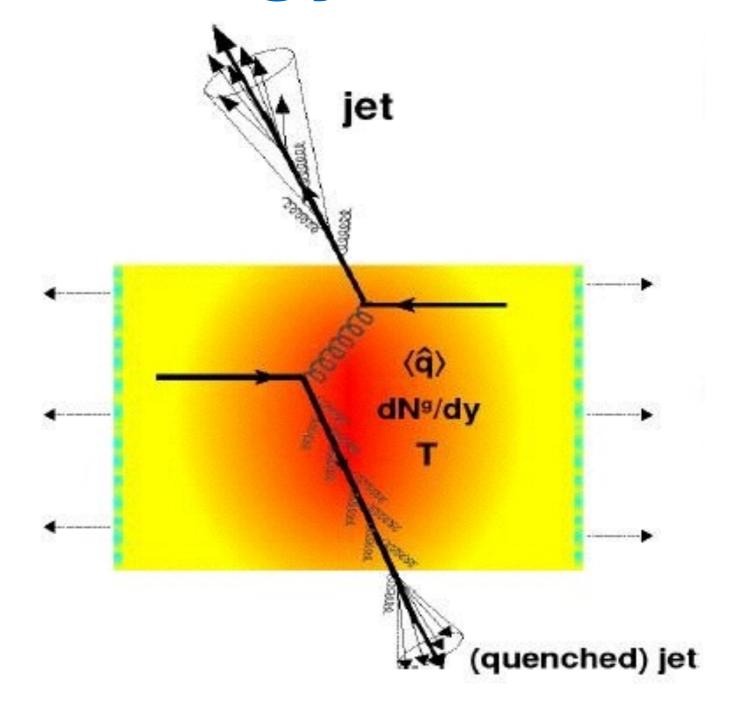


In-medium energy loss

partons produced in high Q² processes lose energy while traversing the medium

modification (suppression)
of high- p_T production
observable: nuclear
modification factor

$$R_{AA} = \frac{dN^{AA}/dp_T}{N_{coll}dN^{pp}/dp_T}$$



 $R_{AA} = 1$ for hard-processes in the absence of nuclear effects confirmed in Pb-Pb collisions at LHC (direct- γ , Z⁰ and W[±])

No nuclear modification in p-Pb

ALICE, EPJC 72 (2012) 1945 ALICE, Pb-Pb, $\sqrt{s_{NN}} = 2.76 \text{ TeV}$ charged particles, lnl < 0.8 Pb-Pb HT (Chen et al.) lower density HT (Chen et al.) higher density 10⁻¹ HT (A.M.) ASW (T.R.) YaJEM-D (T.R.) elastic (T.R.) large P esc elastic (T.R.) small P esc ALICE (0-5%) WHDG (W.H.) π⁰ upper limit CMS (0-5%) WHDG (W.H.) π⁰ lower limit p_{_} (GeV/c)

ALICE, PLB 720 (2013) 52

ALICE, PRL 110 (2013) 082302 ALICE, charged particles p+Pb $\sqrt{s_{NN}}$ = 5.02 TeV, NSD, $|\eta_{cms}|$ < 0.3 Pb+Pb $\sqrt{s_{NN}}$ = 2.76 TeV, 0%-5% central, | η | < 0.8 Pb+Pb $\sqrt{s_{NN}}$ = 2.76 TeV, 70%-80% central, $|\eta| < 0.8$ 1.2 0.8 0.6 0.4 0.2 p_{_} (GeV/c)

charged particle spectra **strongly modified in Pb-Pb** collisions in a wide p_T range

p-Pb confirms that it comes from a final-state effect parton in-medium energy loss

R_{pPb} at intermediate p_T

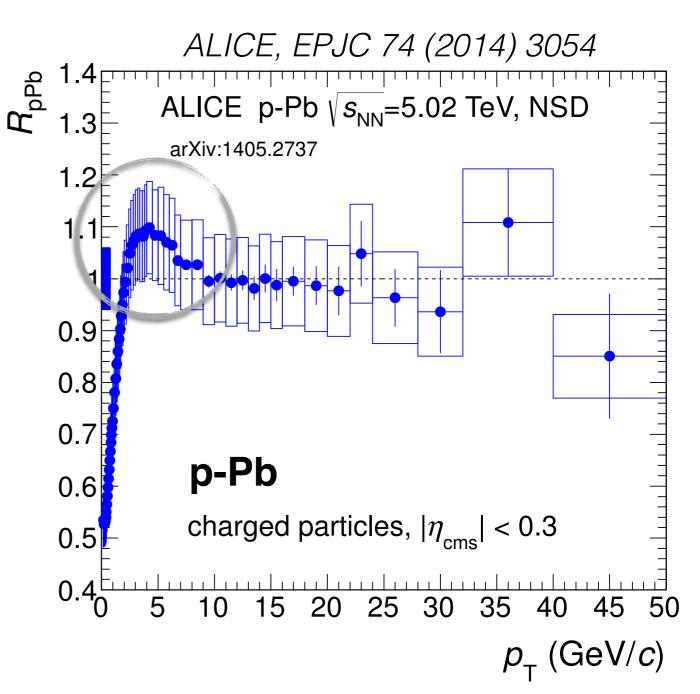
the data indicate a small enhancement at mid-p_T stronger enhancement is seen at lower energies

**Cronin, PRD 11 (1975) 3105*

traditional explanations of **Cronin enhancement**

multiple soft scatterings in the initial state prior to the hard scattering

Accardi, arXiv:hep-ph/0212148

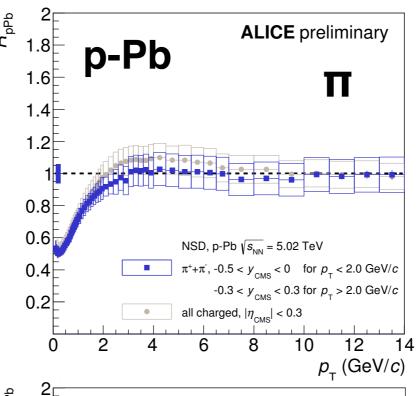


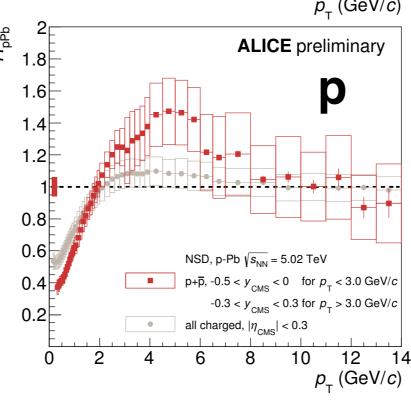
Identified particle RpPb

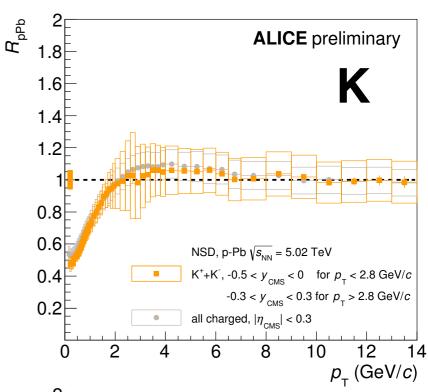
pions and kaons consistent with no modification at mid- p_T

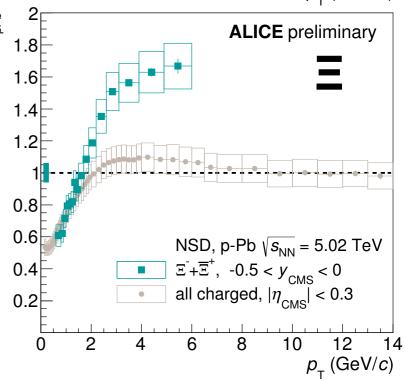
rather pronounced peak for **protons**

even stronger enhancement for cascades







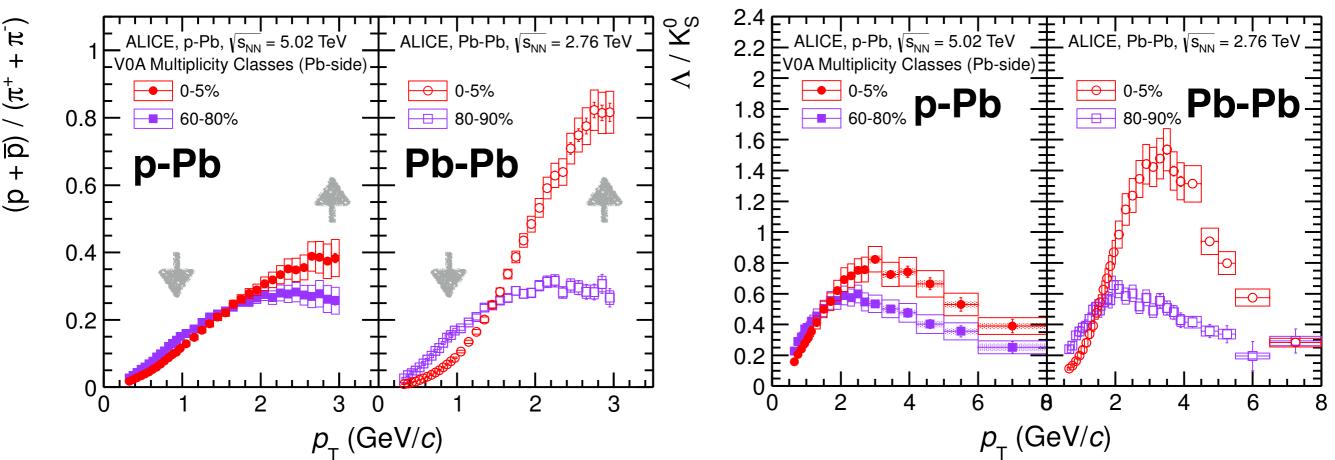


particle species dependence suggests final state effects

recombination, collective flow, ...

Baryon enhancement

ALICE, PLB 728 (2014) 25



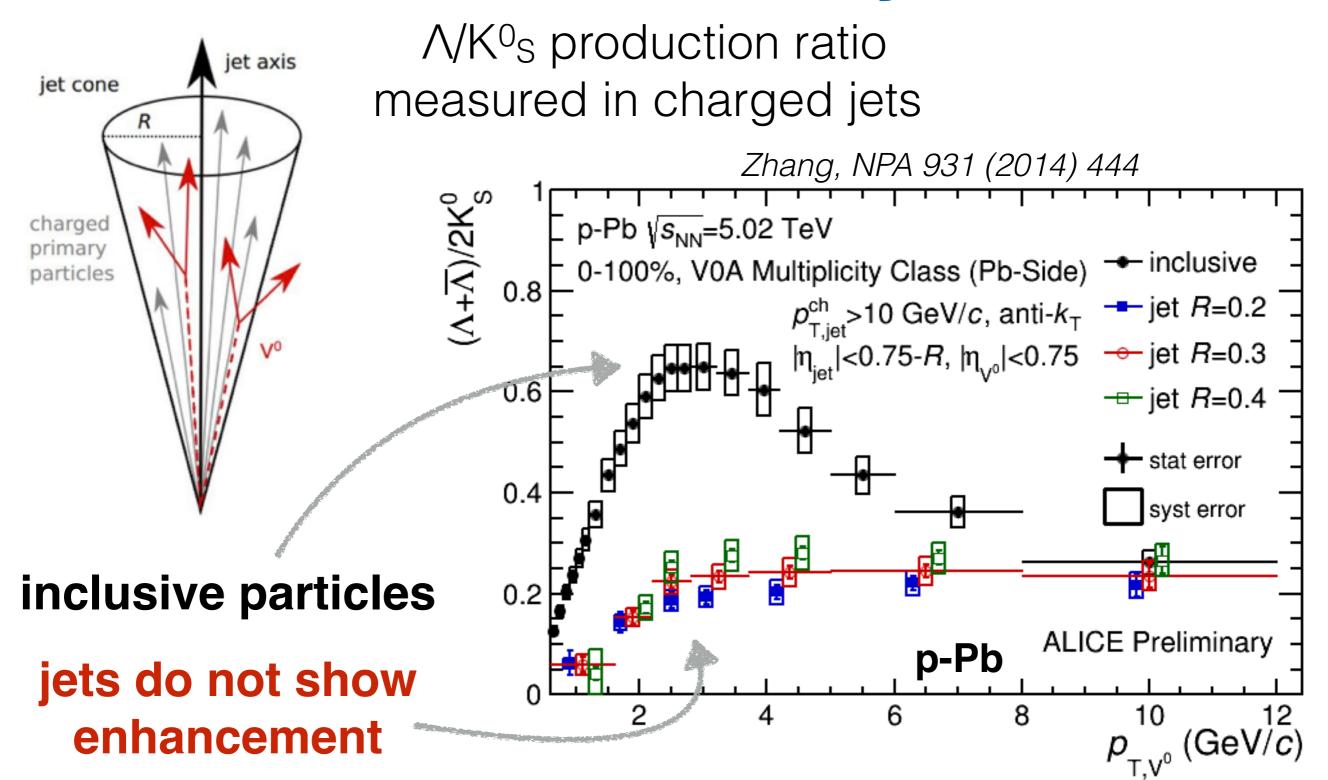
Significant centrality/multiplicity dependence of the ratios

enhancement at mid- p_T with increasing multiplicity corresponding depletion in the low- p_T region

Reminiscent of A-A observations

commonly understood in terms of collective flow / quark recombination

Where are the extra baryons from?

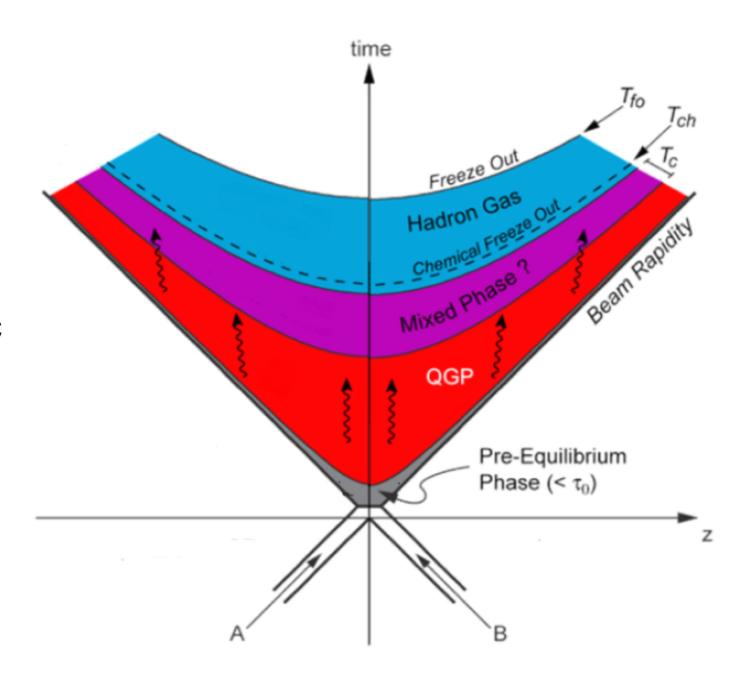


the extra baryons are not coming from jets

Collective phenomena

bulk matter created in high-energy heavy-ion collisions can be described in terms of hydrodynamics

- initial hot and dense partonic matter rapidly expands
- collective flow develops and the system cools down
- phase transition to hadron gas when T_{critical} is reached resulting in



- dependence of the shape of the p_T distribution on the particle mass
- azimuthal anisotropic flow patterns (initial spatial anisotropy)

are there final state dense matter effects in p-Pb?

Bulk π, K, p production in p-Pb

Blast-Wave

hydro-motivated fit thermal sources expanding with common velocity

EPOS LHC

full event generator with hydro evolution

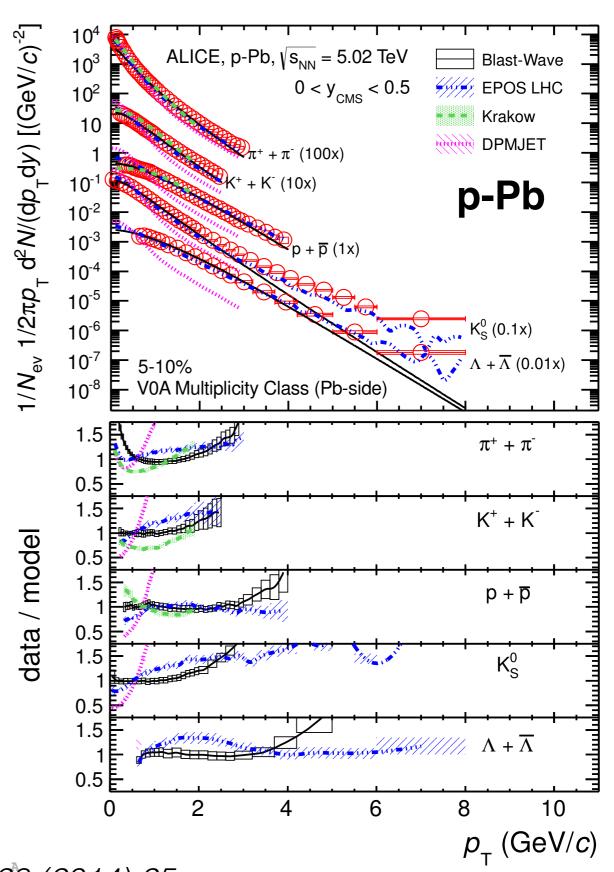
Krakow

3+1 viscous hydro

DPMJET

pQCD based

Models including hydrodynamics do a better job describing the data



Collective phenomena

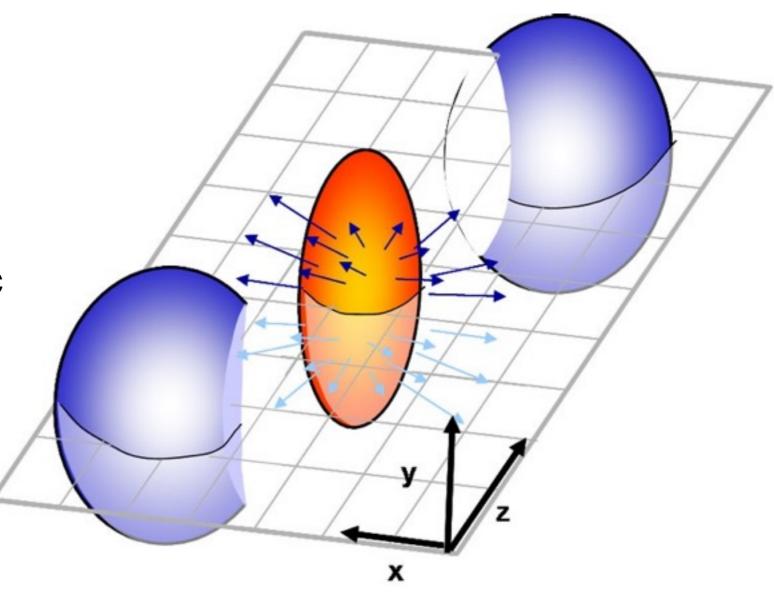
bulk matter created in high-energy heavy-ion collisions can be described in terms of hydrodynamics

 initial hot and dense partonic matter rapidly expands

 collective flow develops and the system cools down

• phase transition to hadron gas when T_{critical} is reached

resulting in

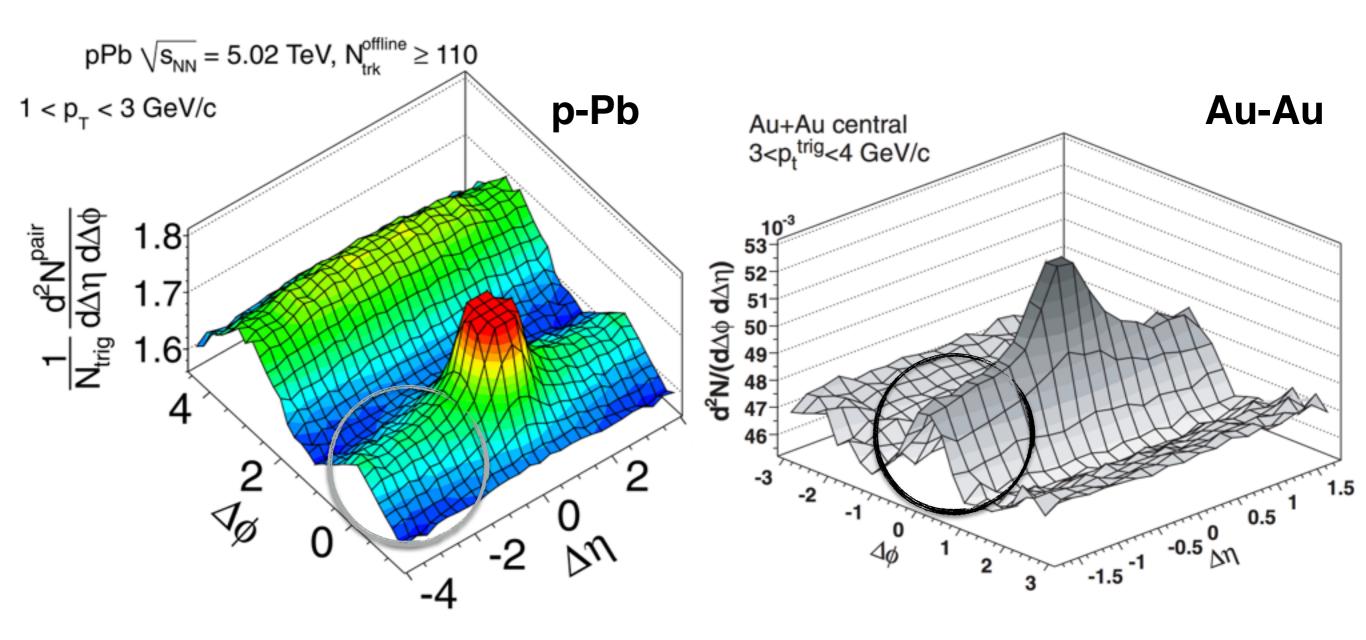


- dependence of the shape of the p_T distribution on the particle mass
- azimuthal anisotropic flow patterns (initial spatial anisotropy)

are there final state dense matter effects in p-Pb?

The ridge

long-range (2 < $|\Delta\eta|$ < 4), near-side ($\Delta\phi \approx 0$) resembles the ridge-like correlation seen in A-A collisions interpreted as consequence of hydrodynamic flow

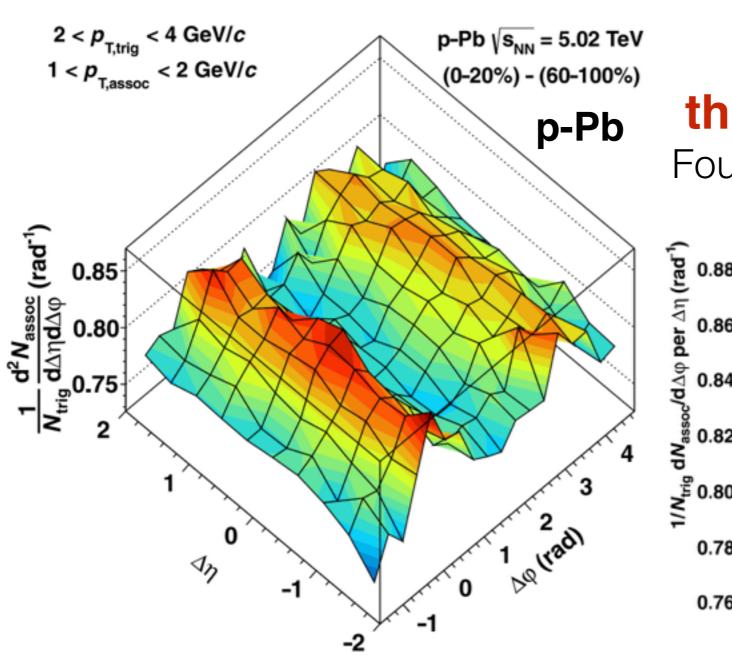


CMS, PLB 718 (2013) 795

STAR, PRC 80 (2010) 064912

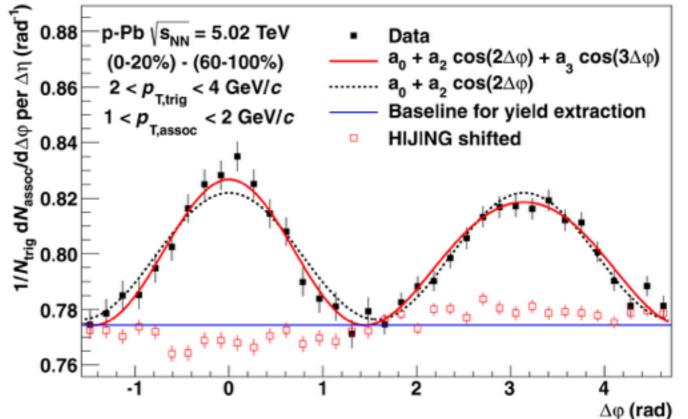
The double ridge

the ridge in p-Pb events triggered further investigations jet contribution removed by subtracting low-multiplicity events a double ridge structure was revealed



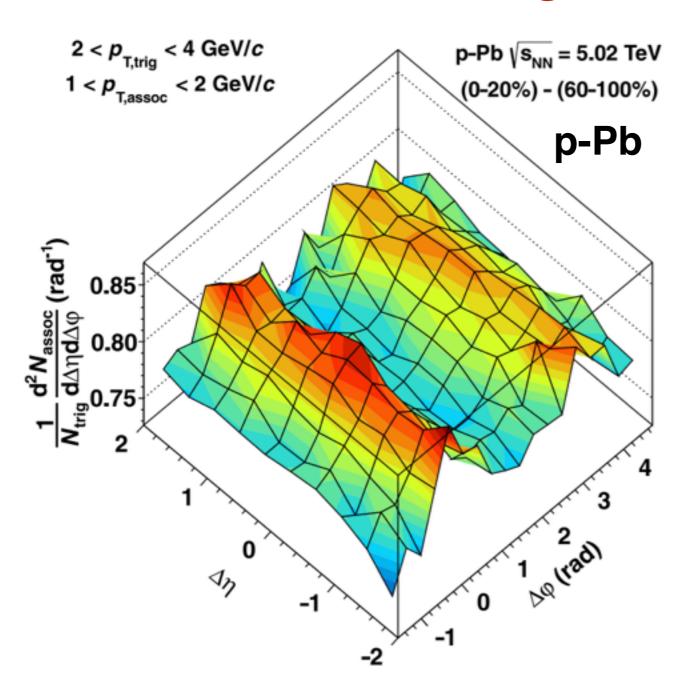
this looks so much like flow

Fourier decomposition of $\Delta \varphi$: V_2 , V_3 , ...



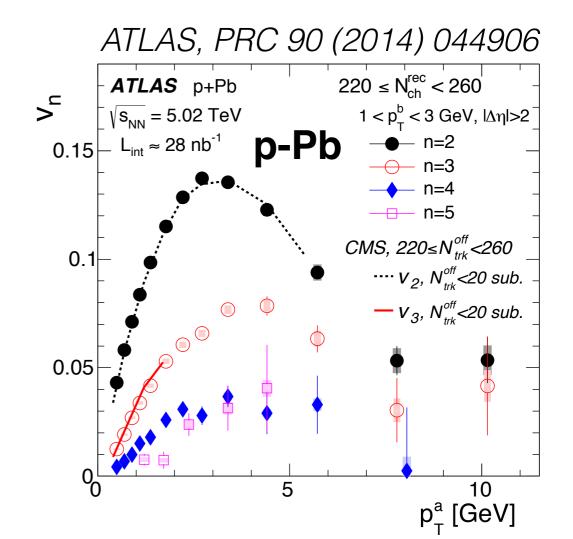
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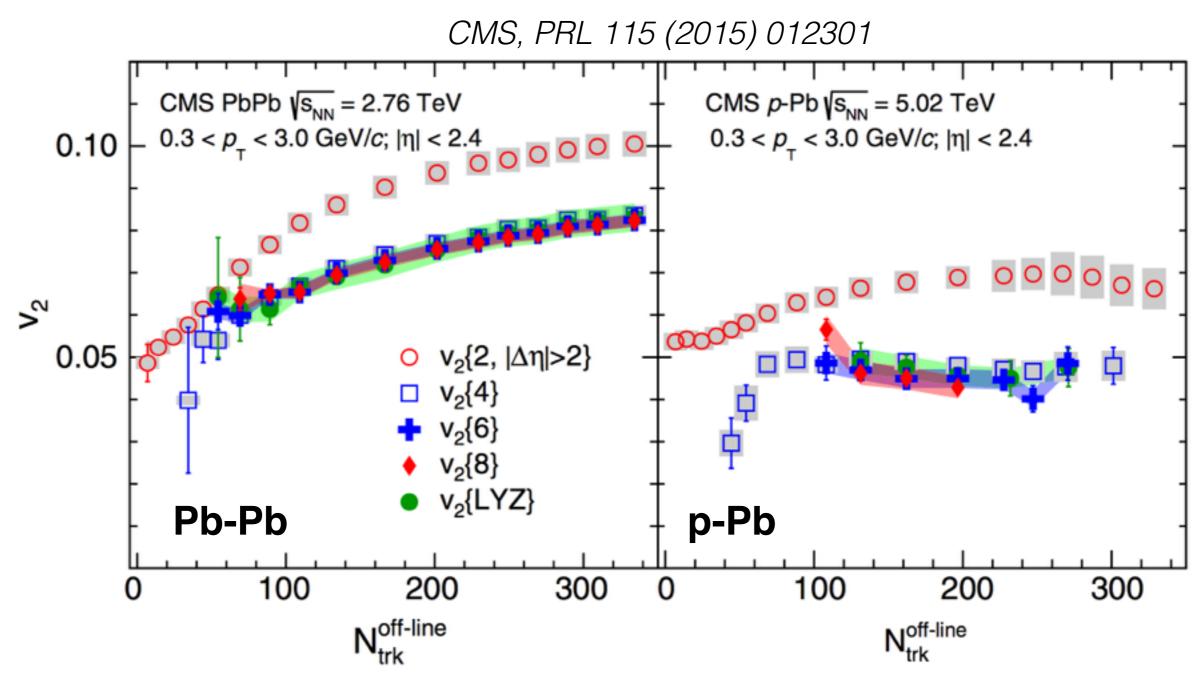


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Fourier decomposition of $\Delta \varphi$: V_2 , V_3 , ...

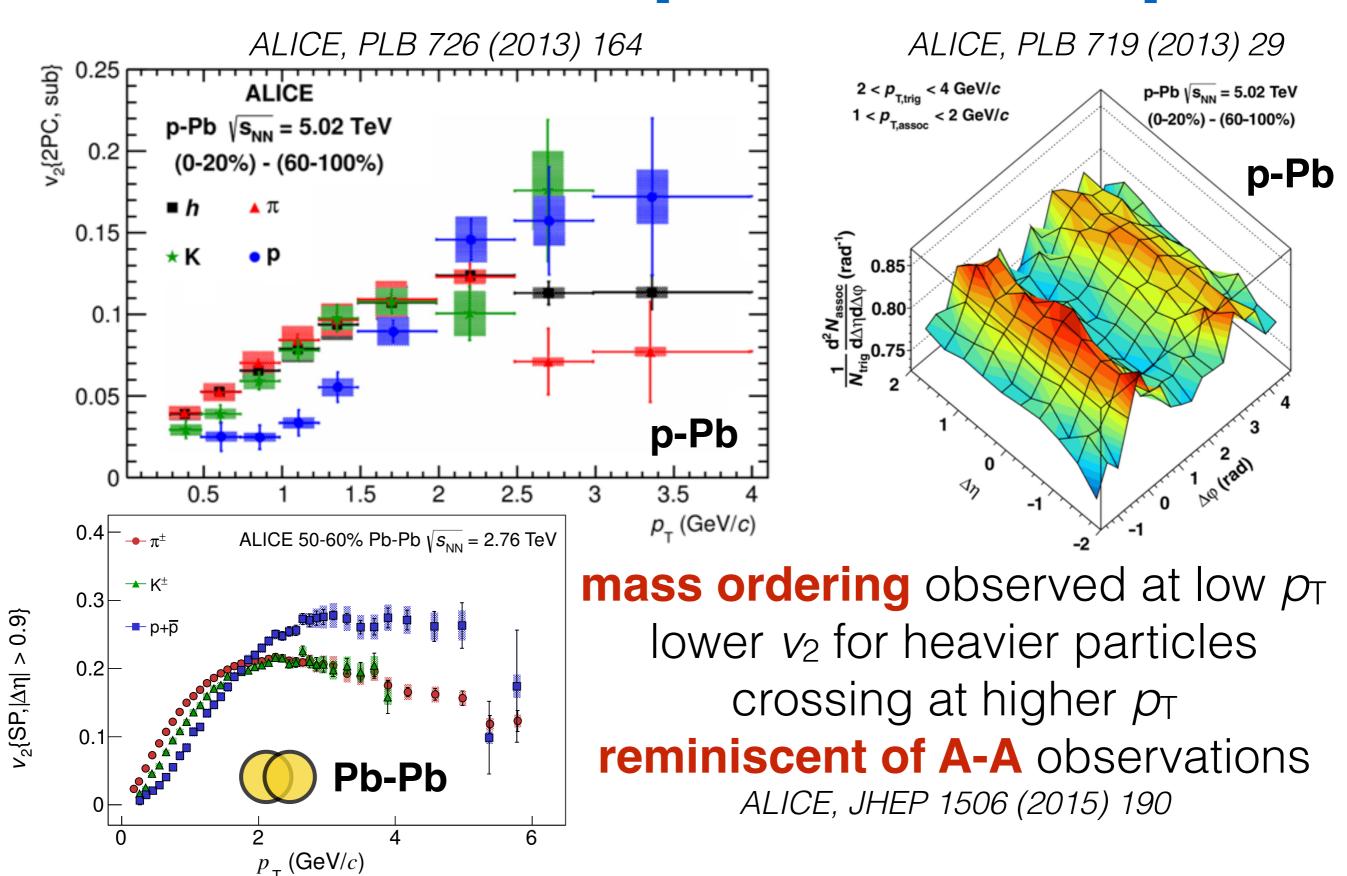


True collective effect

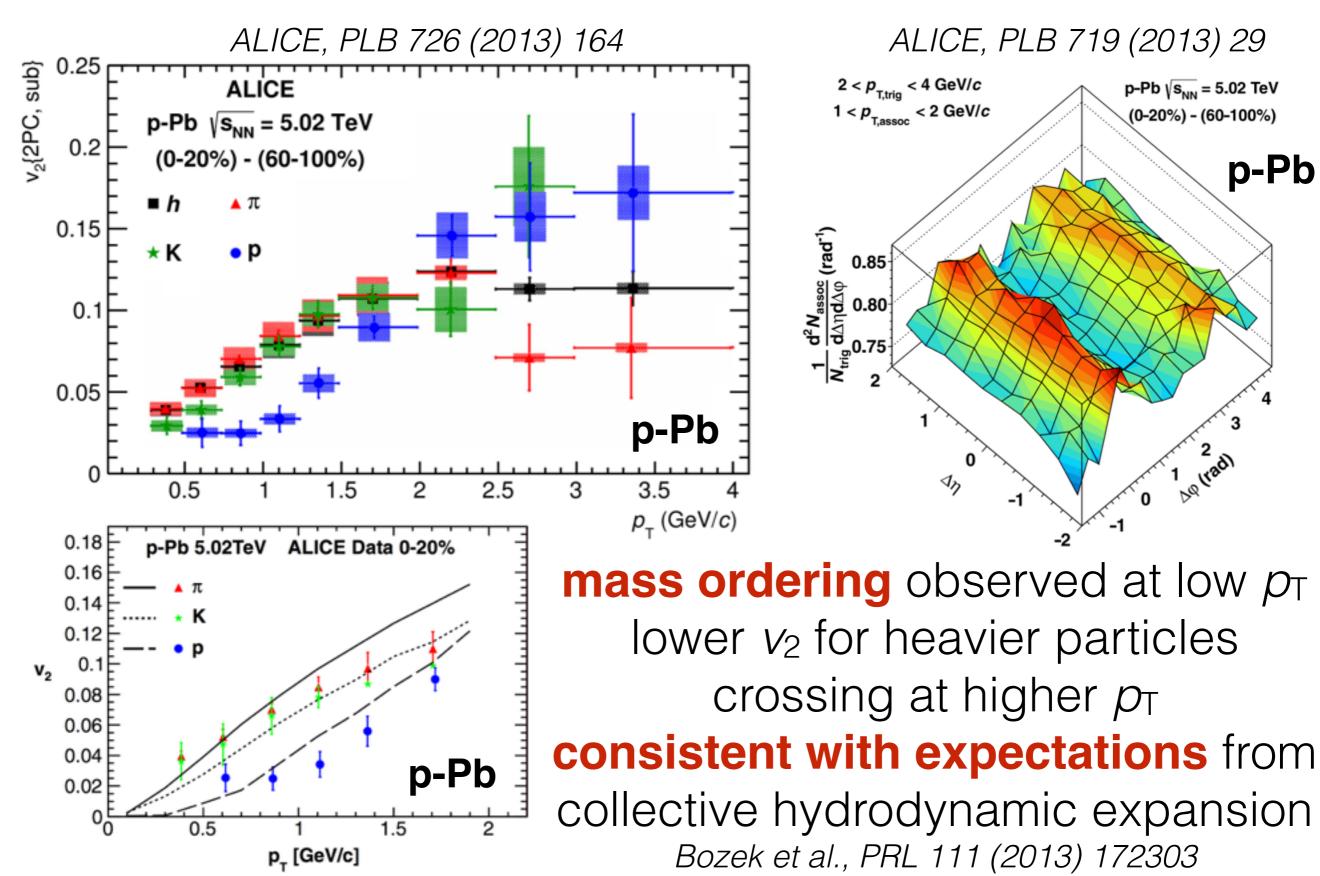


 v_2 stays large when computed with multi-particles $v_2\{4\} = v_2\{6\} = v_2\{8\} = v_2\{LYZ\}$ have different sensitivity to non-flow effects there is true collectivity in p-Pb

v₂ of identified particles in p-Pb



v₂ of identified particles in p-Pb



Strangeness enhancement

one of the first proposed QGP signatures

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PHYSICAL REVIEW LETTERS

19 APRIL 1982

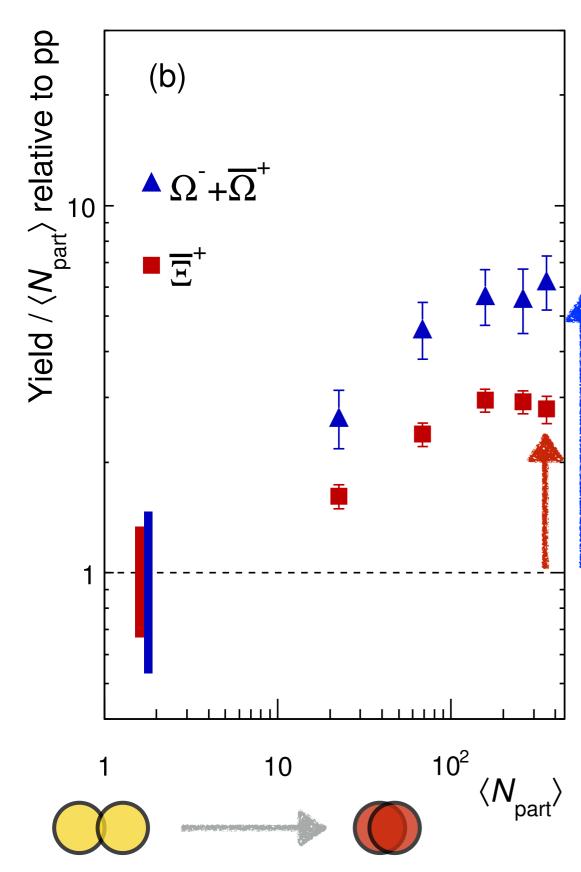
Strangeness Production in the Quark-Gluon Plasma

Johann Rafelski and Berndt Müller

Institut für Theoretische Physik, Johann Wolfgang Goethe-Universität, D-6000 Frankfurt am Main, Germany (Received 11 January 1982)

We thus conclude that strangeness abundance saturates in sufficiently excited quark-gluon plasma (T > 160 MeV, $E > 1 \text{ GeV/fm}^3$), allowing us to utilize enhanced abundances of rare, strange hadrons ($\overline{\Lambda}$, $\overline{\Omega}$, etc.) as indicators for the formation of the plasma state in nuclear collisions.

Strangeness production in Pb-Pb



strangeness enhancement

one of the first proposed QGP signatures Rafelski, PRL 48 (1982) 1066

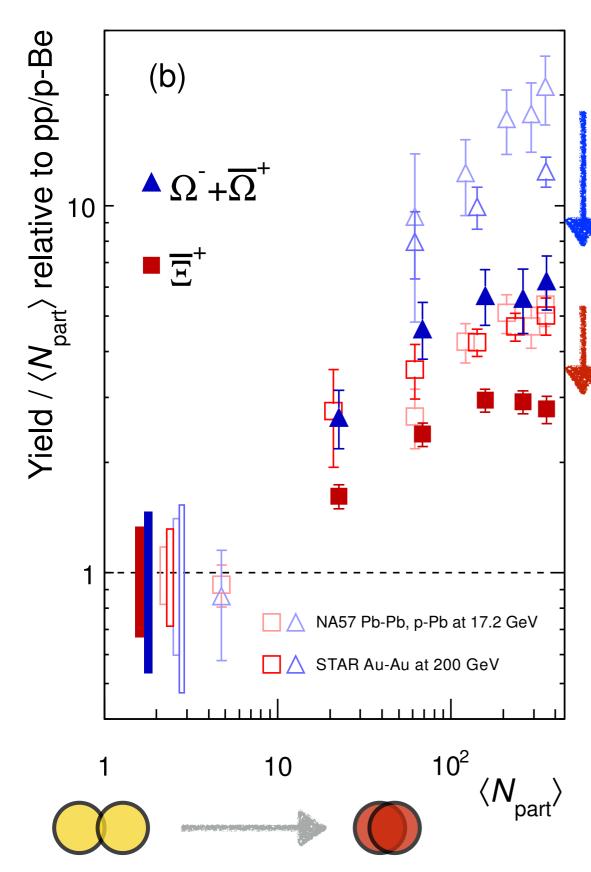
$$E = \frac{2}{\langle N_{part}^{PbPb} \rangle} \frac{(dN/dy)^{PbPb}}{(dN/dy)^{pp}}$$

strangeness-content hierarchy

Ξ (dss) enhanced

 Ω (sss) more enhanced

Strangeness production in Pb-Pb



strangeness enhancement

one of the first proposed QGP signatures Rafelski, PRL 48 (1982) 1066

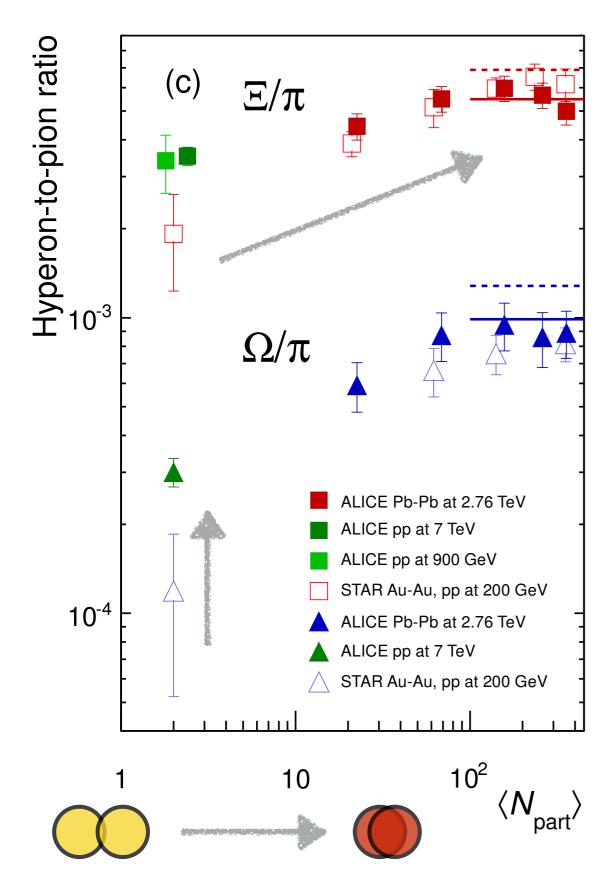
$$E = \frac{2}{\langle N_{part}^{PbPb} \rangle} \frac{(dN/dy)^{PbPb}}{(dN/dy)^{pp}}$$

strangeness-content hierarchy

 Ξ (dss) enhanced Ω (sss) more enhanced

decreasing trend with increasing √s (from SPS to LHC) progressive removal of canonical suppression in pp

Strangeness production in Pb-Pb



strangeness enhancement

one of the first proposed QGP signatures Rafelski, PRL 48 (1982) 1066

relative production of strangeness in pp collisions is larger at LHC

clear increase of strangeness production from pp to Pb-Pb

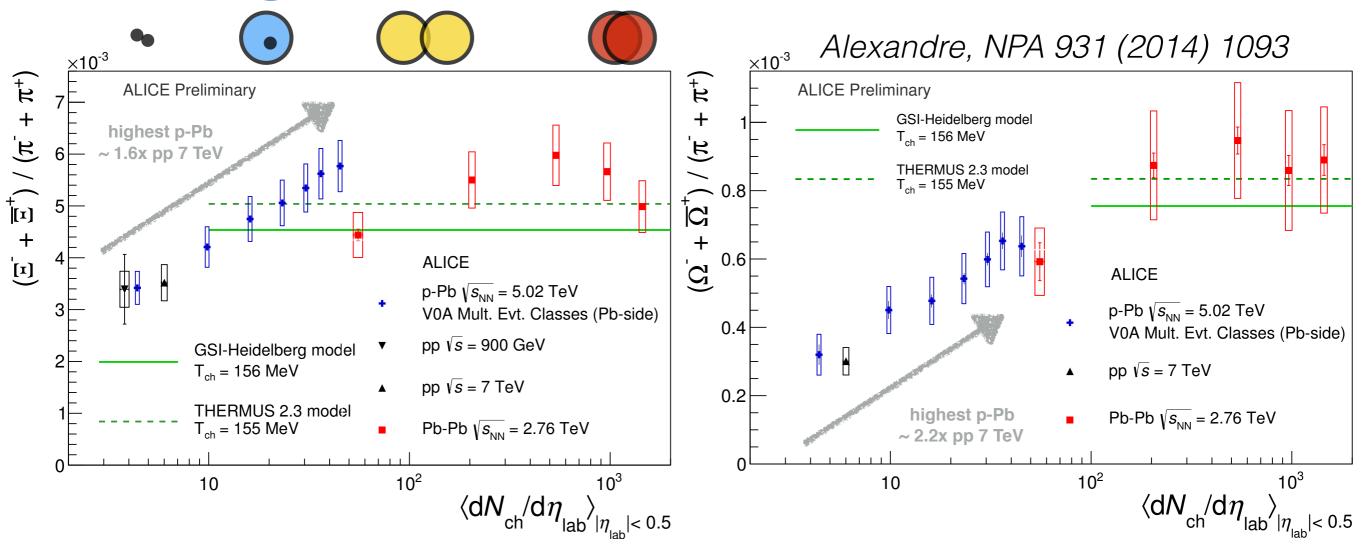
saturation of ratios for $N_{part} > 150$

match predictions from Grand Canonical thermal models

GSI-Heidelberg: $T_{ch} = 164 \text{ MeV}$

THERMUS: $T_{ch} = 170 \text{ MeV}$

Strangeness production in p-Pb



 Ξ/π and Ω/π ratios in p-Pb increase with increasing $\langle N_{ch} \rangle$

low-multiplicity

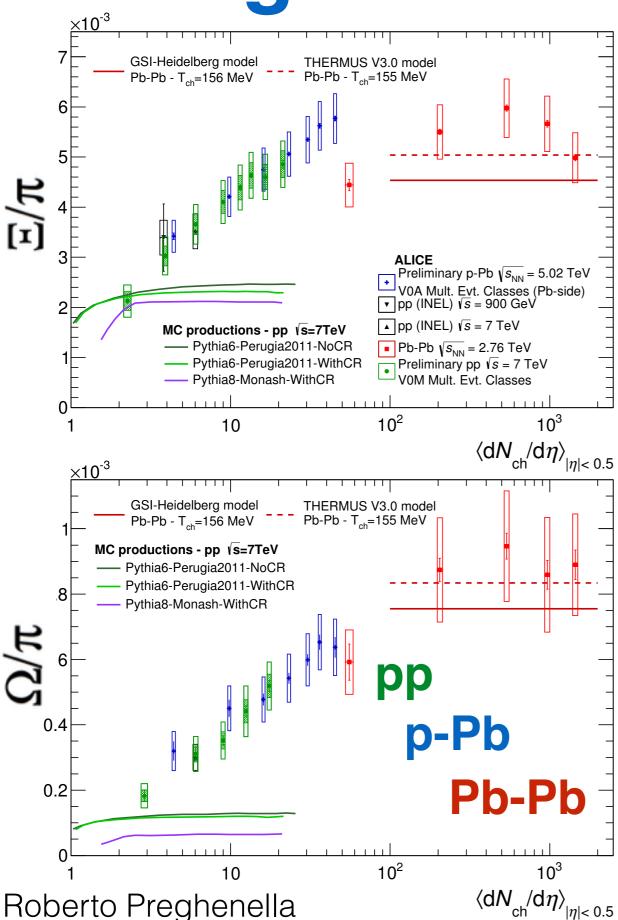
 Ξ and $\Omega \rightarrow$ consistent with pp

high-multiplicity

Ξ → compatible with central Pb-Pb

 $\Omega \rightarrow$ compatible with peripheral Pb-Pb

Strangeness enhancement in pp



also measured in pp and p-Pb collisions as a function of charged-particle multiplicity

first observation of enhanced production of strange particles in pp and p-Pb collisions

ratios to pions <u>reach values</u> <u>measured in Pb-Pb</u> collisions PYTHIA cannot reproduce the data