

# $(g - 2)_\mu$ : recent improvements and preliminary update



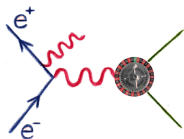
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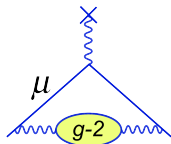


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# Overview

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# Where we left off... [HLMNT(11), J. Phys. G38 (2011), 085003]

<b>QED</b> contribution	11 658 471.808 (0.015) $\times 10^{-10}$	Kinoshita & Nio, Aoyama et al
<b>EW</b> contribution	15.4 (0.2) $\times 10^{-10}$	Czarnecki et al
<b>Hadronic</b> contribution		
<b>LO</b> hadronic	694.9 (4.3) $\times 10^{-10}$	HLMNT11
<b>NLO</b> hadronic	-9.8 (0.1) $\times 10^{-10}$	HLMNT11
<b>light-by-light</b>	10.5 (2.6) $\times 10^{-10}$	Prades, de Rafael & Vainshtein
<b>Theory TOTAL</b>	11 659 182.8 (4.9) $\times 10^{-10}$	
<b>Experiment</b>	11 659 208.9 (6.3) $\times 10^{-10}$	world avg
<b>Exp - Theory</b>	26.1 (8.0) $\times 10^{-10}$	3.3 $\sigma$ discrepancy

(Numbers taken from HLMNT11, arXiv:1105.3149)

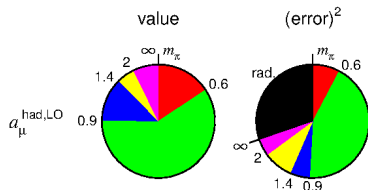
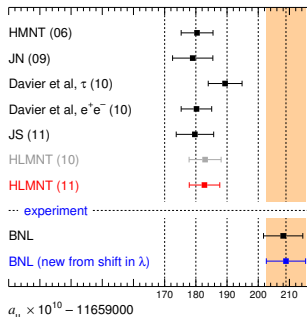
→  $a_\mu^{\text{had, LOVP}}$  still dominating uncertainty

→ Eagerly awaiting **new hadronic cross section data**

→ ...and **new results from Fermilab**

⇒ **Hints at increased discrepancy,**

$$\delta a_\mu = a_\mu^{\text{exp}} - a_\mu^{\text{SM}}$$



# Problems concerning data combination

## ⇒ All data is different

Does data include ISR, FSR, VP correction in  $\delta_{\text{rad}}$ ?

What forms do the uncertainties take: point-to-point, overall systematic percentage, covariance matrix...?

## ⇒ How do we re-bin/cluster data?

What cluster sizes do we allow?

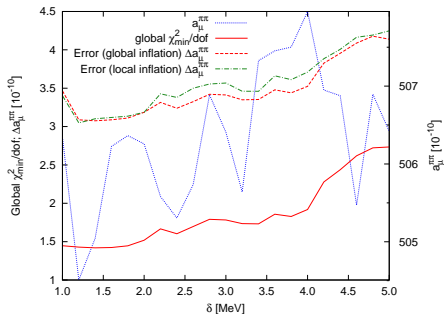
How do we cluster points at resonance peaks as opposed to in a continuum region?

## ⇒ How do we fit data?

What function/method is the best?

How do we construct a trustworthy representation of the uncertainties as an overall uncertainty on  $a_\mu$ ?

How do we achieve this whilst avoiding potential bias through error propagation?



# Existing/new data

⇒ Do we discard old, imprecise data?

Some older data sets ( $\sim 80$ s) are questionably inaccurate compared to newer sets?

Can we justify removing this data?

⇒ Differences between direct scan and radiative return data

More and more radiative return measurements...

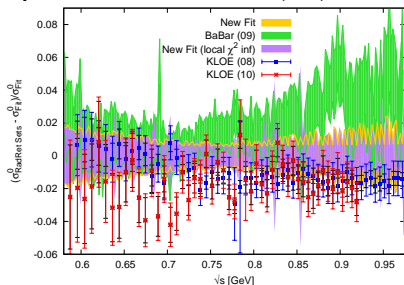
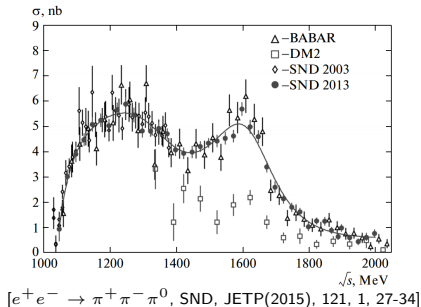
How do they compare?

⇒ BaBar problem

Why do BaBar measurements sit higher than all other data?

Causing tension between sets...

No longer just  $\pi^+\pi^-$

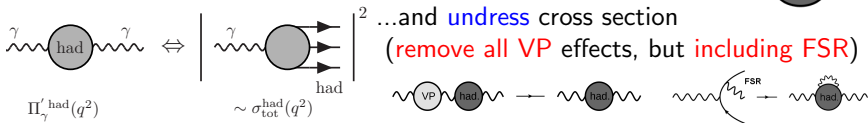
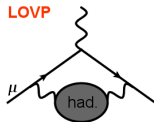


# Calculating $a_\mu^{\text{had, LOVP}}$

Need to **precisely** determine  $a_\mu^{\text{had, LOVP}} < 2 \text{ GeV}$

→ pQCD not useful

→ Use **analyticity** and the **optical theorem**



...as input into **dispersion relation**:

$$a_\mu^{\text{had, LOVP}} = \frac{1}{4\pi^3} \int_{s_{\text{th}}}^{\infty} ds \frac{\sigma_{\text{had}}^0(s)}{\sigma_{\text{pt}}(s)} \frac{4\pi\alpha^2}{3s} K(s)$$

- Sum  $\sim 30$  **exclusive channels** (isospin relate missing channels)
- Combine data from **many experiments**, different energy bins, uncertainties, correlations...
- Re-binning through **adaptive clustering** and minimise  $\chi^2$ -function

## Data combination: step-by-step

- Apply appropriate **radiative corrections** to data: VP, FSR etc.
- **Re-bin** data into energy-dependent clusters (**piecewise constant  $R$** )
- Determine cluster properties: energy, cross-section value, uncertainty...

$$R_m = \left[ \sum_k \sum_{i=1}^{N(k,m)} \frac{R_i^{(k,m)}}{(\text{d}\tilde{R}_i^{(k,m)})^2} \right] \left[ \sum_k \sum_{i=1}^{N(k,m)} \frac{1}{(\text{d}\tilde{R}_i^{(k,m)})^2} \right]^{-1}$$

- Construct uncertainty **covariance matrices** for all data
- **Fit data** according to  $\chi^2$  minimum
  - Determine **local  $\chi^2$**  per cluster and **global  $\chi_{\text{min}}^2/\text{d.o.f.}$**
- **Integrate** over fitted clusters as input into dispersion integral to determine  $a_\mu \pm \delta a_\mu$
- **Inflate error**,  $\delta a_\mu$ , according to local cluster  $\chi^2$

# The $f_k$ method

- Fit data and minimise a **non-linear  $\chi^2$ -function** [HLMNT, 2012]  
 → **Two fitting parameters**: initial cluster values  $R_m$ , normalisation factor  $f_k$  of each experiment  $k$

$$\chi^2(R_m, f_k) = \sum_{k=1}^{N_{exp}} \left\{ \left( \frac{1 - f_k}{df_k} \right)^2 + \left[ \sum_{m=1}^{N_{clu}} \sum_{i=1}^{N^{(k,m)}} \left( \frac{R_i^{(k,m)} - f_k R_m}{dR_i^{(k,m)}} \right)^2 \right]_{\text{w/o cov. mat}} \right. \\ \left. + \left[ \sum_{m=1}^{N_{clu}} \sum_{n=1}^{N_{clu}} \sum_{i=1}^{N^{(k,m)}} \sum_{j=1}^{N^{(k,n)}} (R_i^{(k,m)} - f_k R_m) C^{-1}(m_i, n_j) (R_j^{(k,n)} - f_k R_n) \right] \right\}$$

The  $f_k$ 's are **multiplicative** re-normalisation factors for the data **which vary as the  $\chi^2$ -function is minimised**

→ **Penalty Trick Method**



# What is a biased result?

→ *D'Agostini Bias* = Fit favours more precise measurement

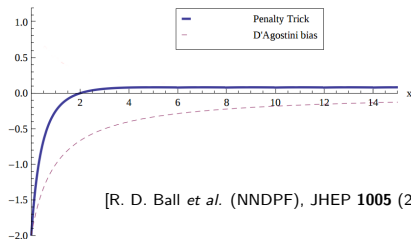
[D'Agostini, Nucl. Instrum. Meth. A **346** (1994) 306]

⇒ How do we include a global normalisation uncertainty whilst avoiding a D'Agostini bias? →  $f_k$  (penalty trick) method

$$\chi^2(R_m, f_k) = \sum_{k=1}^{N_{exp}} \left\{ \left( \frac{1 - f_k}{df_k} \right)^2 \right\} + \left[ \sum_{m=1}^{N_{clu}} \sum_{i=1}^{N^{(k,m)}} \left( \frac{R_i^{(k,m)} - f_k R_m}{dR_i^{(k,m)}} \right)^2 \right]_{w/o \text{ cov. mat}}$$

PENALTY TERM  
NORMALISATION FACTOR  
FIT PARAMETERS

Is the  $f_k$  method truly free from bias?



# The $f_k$ method: a toy example

Consider **two** measurements:  $R_i^{(k,m)}$  and  $R_j^{(l,m)}$  with **equal (uncorrelated) errors**...

$$\rightarrow df_k = df_l \equiv df \quad ; \quad dR_i^{(k,m)} = dR_j^{(l,m)} \equiv dR$$

$$\text{Unbiased solution} \rightarrow R_m = \bar{R}_m = \frac{1}{2}(R_i^{(k,m)} + R_j^{(l,m)})$$

Minimising w.r.t  $R_m$  and  $f_k$  and substituting, we find:  $R_m = \bar{R}_m(1 + \beta_m)$

$\Rightarrow \beta_m$  is the **bias contribution** to the cluster centre  $R_m$

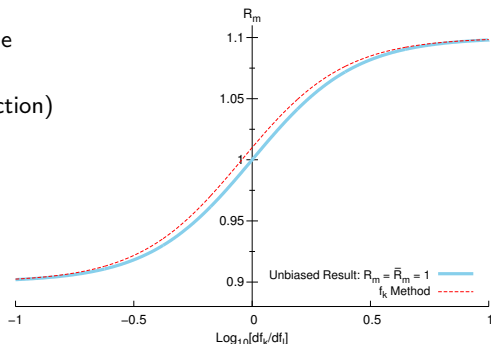
(due to **non-linear nature** of  $\chi^2$  function)

$\rightarrow$  Now, let

$$R_i^{(k,m)} = 0.9 \text{ and } R_j^{(l,m)} = 1.1$$

$$dR_i^{(k,m)} = dR_j^{(l,m)} \equiv dR = 0$$

$$\Rightarrow R_m = 1.0050125 \neq \bar{R}_m$$



# The $R_m^I$ method

Following → NNPDF Collaboration, R.D.Ball *et al.* [JHEP **1005** (2010) 075]

→ Benayoun *et al.* [Eur. Phys. J. C **75** (2015) no.12, 613]

⇒ Claim: **fix covariance matrices** during fitting

$$\mathbf{C}(i^{(k,m)}, j^{(l,n)}) = \mathbf{C}^{\text{stat}}(i^{(k,m)}, j^{(l,n)}) + \mathbf{C}^{\text{sys}\%}(i^{(k,m)}, j^{(l,n)}) R_m^0 R_n^0$$

⇒ ...to prevent error propagation skewing theory value

$$\chi^2 = \sum_{m,n=1}^{N_{clu}} \sum_{i,j=1}^{N^{(k;m,n)}} (R_i^{(k,m)} - R_m) \mathbf{C}^{-1}(i^{(k,m)}, j^{(l,n)}) (R_j^{(k,n)} - R_n)$$

→ Allows for inclusion of  
full correlated normalisation uncertainties  
without renormalisation factors

→ Iterative routine forces convergence to  
unbiased solution

→ Observe one (maximum two) iteration  
to convergence

→ Benayoun *et al.* predict reduction in  
mean value and improved uncertainty

Channel	$A = m$	$A = M_0$	Exp. Value
$\pi^+\pi^-$	$495.06 \pm 1.43$	$494.59 \pm 0.89$	$492.98 \pm 3.38$
$\pi^0\gamma$	$4.53 \pm 0.04$	$4.54 \pm 0.04$	$3.67 \pm 0.11$
$\eta\gamma$	$0.64 \pm 0.01$	$0.64 \pm 0.01$	$0.56 \pm 0.02$
$\pi^+\pi^-\pi^0$	$40.83 \pm 0.57$	$40.84 \pm 0.57$	$43.54 \pm 1.29$
$K_L K_S$	$11.56 \pm 0.08$	$11.53 \pm 0.08$	$12.21 \pm 0.33$
$K^+ K^-$	$16.79 \pm 0.20$	$16.90 \pm 0.20$	$17.72 \pm 0.52$
Total	$569.41 \pm 1.55$	$569.04 \pm 1.08$	$570.68 \pm 3.67$

# The $R_m^I$ method: a toy example

Consider **two** measurements:  $R_i^{(k,m)}$  and  $R_j^{(l,m)}$  with equal (uncorrelated) errors...

$$\rightarrow \chi^2(R_m) = \sum_{i=1}^{N^{(k,m)}} \frac{(R_i^{(k,m)} - R_m)^2}{(dR_{i; \text{stat}}^{(k,m)})^2 + (dR_{i; \text{sys}\%}^{(k,m)})^2 (R_m^0)^2}$$

Minimising w.r.t  $R_m$  and **reintroduce**  $\rightarrow df_k = df_l \equiv df$  ;  $dR_i^{(k,m)} = dR_j^{(l,m)} \equiv dR$

$$R_m = \left[ \frac{R_i^{(k,m)} + R_j^{(l,m)}}{(dR_{\text{stat}})^2 + (dR_{\text{sys}})^2 (R_m^0)^2} \right] \left[ \frac{2}{(dR_{\text{stat}})^2 + (dR_{\text{sys}})^2 (R_m^0)^2} \right]^{-1}$$

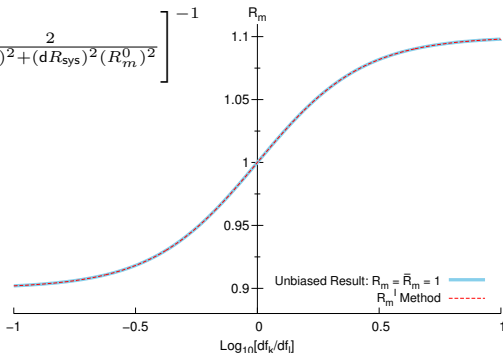
$$= \frac{1}{2} (R_i^{(k,m)} + R_j^{(l,m)})$$

$\rightarrow$  Again, let

$$R_i^{(k,m)} = 0.9 \text{ and } R_j^{(l,m)} = 1.1$$

$$dR_i^{(k,m)} = dR_j^{(l,m)} \equiv dR = 0$$

$$\Rightarrow R_m = 1 = \bar{R}_m$$



# What effect will this have on our results?

Use dominant  $\pi^+\pi^-$  channel as example of change to  $a_\mu^{\text{had, LOVP}}$

→ Consider **real cluster** with **one measurement**:

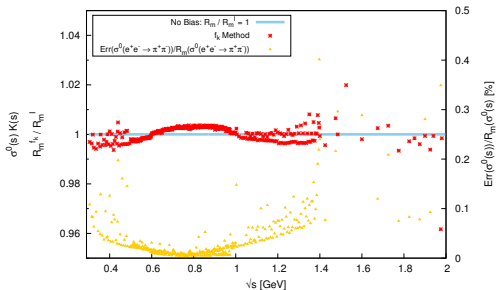
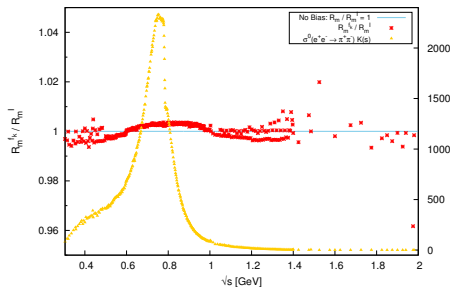
$$E_{cm}(\text{GeV}) = 0.4600, \sigma^0(\text{nb}) = 123.6418 \pm 22.4855 \Rightarrow R_m^0 = 123.6418$$

Experimental data includes **three more measurements in different clusters**.

⇒ **Normalisation uncertainty provides weighting** to cluster value through **fixed covariance matrix correlations**

$f_k$  method →  $R_m^{f_k} = 130.81$

;  $R_m^I$  method →  $R_m^I = 130.97$



# Are previous results reliable?

## Did the $f_k$ method incur a bias?

Compare  $f_k$  method and  $R_m^I$  method with **only multiplicative normalisation uncertainties**.

→ If we see **differences** in mean value, then **bias previously influenced the fit**.

→ **Previous results unreliable**

→ If we see **no differences** in mean value, then **bias did not influence fit** (any change comes from improved treatment of systematics)

→ **Previous results reliable**

*Example -  $\pi^+\pi^-$*

Set 1 - CMD-2(06) (0.7% Systematic Uncertainty), Set 2 - CMD-2(06) (0.8% Systematic Uncertainty), Set 3 - SND(04) (1.3% Systematic Uncertainty)

From 0.37 → 0.97 GeV

Fit Method:	$f_k$ Method		$R_m^I$ Method		
Channel	$a_\mu$	$\chi_{\min}^2/\text{d.o.f.}$	$a_\mu$	$\chi_{\min}^2/\text{d.o.f.}$	Difference
$\pi^+\pi^-$	$481.42 \pm 4.26$	1.10	$481.42 \pm 4.05$	1.02	0.00

# What new data is available to us since 2011 analysis?

Since HLMNT(11)...

- $\pi^0\gamma$ : SND (2016) (1 direct scan)
- $\pi^+\pi^-$ : KLOE (2012), BESIII (2015) (2 radiative return)
- $\pi^+\pi^-\pi^0$ : SND (2015) (1 direct scan)
- $\pi^+\pi^-\pi^+\pi^-$ : BaBar (2012) (1 radiative return)
- $\pi^+\pi^-\pi^+\pi^-\pi^+\pi^-$ : CMD-3 (2013) (1 direct scan)
- $K^+K^-$ : BaBar (2013) (1 radiative return)
- $K_S^0K_L^0$ : BaBar (2014), CMD-3 (2016) (1 radiative return, 1 direct scan)
- $K^+K^-\pi^+\pi^-$ : CMD-3 (2015) (1 direct scan)
- $\eta\gamma$ : SND (2014) (1 direct scan)
- $\eta\pi^+\pi^-$ : SND (2015) (1 direct scan)
- $\omega\pi^0$ : SND (2015) (1 direct scan)
- $p\bar{p}$ : CMD-3 (2015) (1 direct scan)
- $n\bar{n}$ : SND (2014) (1 direct scan)

**No new data for**  $\pi^+\pi^-\pi^0\pi^0$ ,  $\pi^+\pi^-\pi^+\pi^-\pi^0$ ,  $\pi^+\pi^-\pi^0\pi^0\pi^0$ ,  $K_S^0K^\pm\pi^\mp$ ,  $K^*K\pi$ ,  $K^+K^-\pi^0\pi^0$ ,  $K^+K^-\pi^+\pi^-\pi^0$ ,  $\eta\pi^+\pi^-\pi^+\pi^-$ ,  $\eta\omega$ ,  $\omega\pi^+\pi^-$ ,  $\eta\phi$ ,  $\phi\pi^0$ ,  $\phi\pi^+\pi^-$

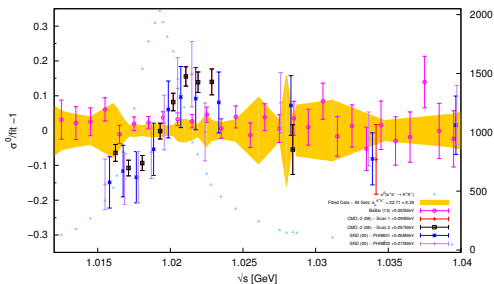
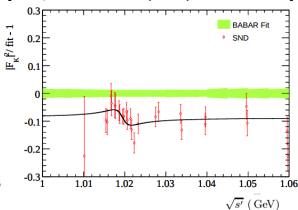
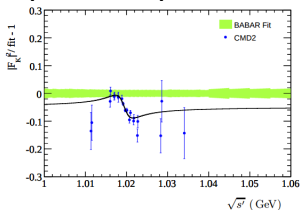
$e^+e^- \rightarrow K^+K^- : m_\phi$  calibration uncertainty study

→ Tension in fitted  $\Phi$  mass between BaBar, CMD-2 and SND

→ Study to shift energies of data points by differences in the  $\phi$ -mass calibration to potentially reduce tension

→ Allow variation in  $\delta m_\phi$

[ J. P. Lees *et al.* [BaBar Collaboration], Phys. Rev. D **88** (2013) no.3, 032013]



→ Marginal improvement in error and  $\chi^2/\text{d.o.f}$  but still not resolved

⇒ BaBar puzzle still remains

$$a_\mu^{K^+K^-} (\leq 2.0\text{GeV}) = 22.75 \pm 0.27$$

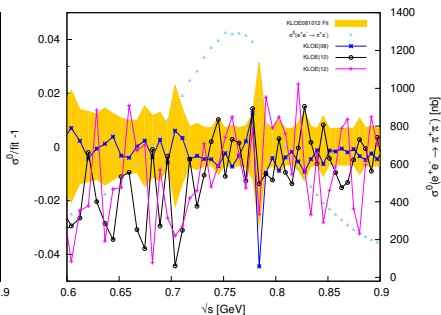
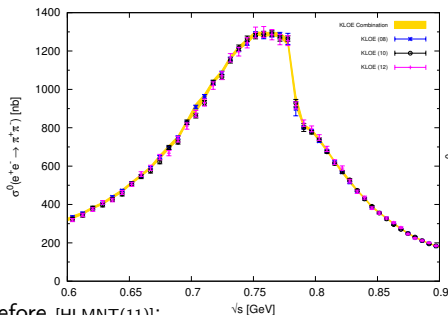
$$\chi_{\min}^2/\text{d.o.f.} = 2.09$$



# $e^+e^- \rightarrow \pi^+\pi^-$ : inclusion of KLOE(12) & KLOE combination matrix

[[KLOE Collaboration], Phys. Lett. B 720 (2013) 336; arXiv:1406.4639 [hep-ph]]

Combination of KLOE(08), KLOE(10) & KLOE(12) radiative return data



Before [HLMNT(11)]:

$$a_{\mu}^{\pi^+\pi^-} (0.305 \leq \sqrt{s} \leq 2.0 \text{ GeV}) = (505.77 \pm 3.09) \times 10^{-10} \quad ; \quad \chi^2_{\min}/\text{d.o.f.} = 1.35$$

After:

$$a_{\mu}^{\pi^+\pi^-} (0.305 \leq \sqrt{s} \leq 2.0 \text{ GeV}) = (505.05 \pm 2.19) \times 10^{-10} \quad ; \quad \chi^2_{\min}/\text{d.o.f.} = 1.37$$

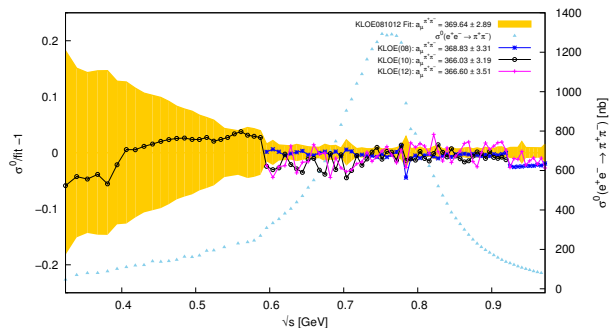
# KLOE combination issues

Individual KLOE measurements give (in range  $0.6 \leq \sqrt{s} \leq 0.9$  GeV):

		FJ03-VP	→	VP_HLMNT_v2.2
KLOE(08)	⇒	$368.83 \pm 3.31$	→	$368.27 \pm 3.31$
KLOE(10)	⇒	$366.03 \pm 3.19$	→	$365.48 \pm 3.19$
KLOE(12)	⇒	$366.60 \pm 3.51$		

Combining using  $R_m^I$  method:

KLOE(081012)	⇒	$369.94 \pm 2.89$	→	$369.64 \pm 2.89$
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⇒ KLOE(08) & KLOE(12)  
correlated systematics  
above  $\sim 0.9$  GeV drag fit up

⇒ Removing these  
correlated systematics gives:

$$a_\mu^{\pi^+\pi^-} = 367.70 \pm 2.70$$

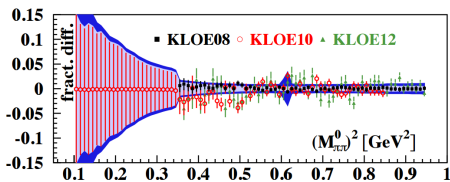
⇒ We believe  $R_m^I$   
method is free from bias  
BUT if so... what is the  
cause?

# KLOE combination: method comparison

⇒ KLOE preliminary combination [1406.4639; Acta Phys. Polon. B **46** (2015) 45] uses... “the best linear unbiased estimator (BLUE) method... (whereby) the BLUE values are constructed using only the covariance matrix with statistical uncertainties.”

→ Systematic uncertainties added separately to error contribution

⇒ Equivalent to minimising  $\chi^2$  function??????



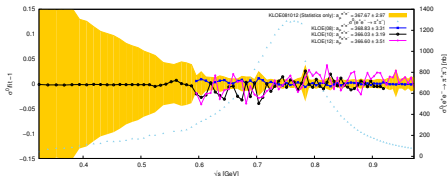
Courtesy of S. Müller

→ Using BLUE method...

$$a_{\mu}^{\pi^+\pi^-} (0.324 \leq \sqrt{s} \leq 0.975 \text{ GeV}) = 487.8 \pm 5.7$$

$$a_{\mu}^{\pi^+\pi^-} (0.592 \leq \sqrt{s} \leq 0.922 \text{ GeV}) = 378.1 \pm 2.8$$

⇒ Open for discussion...



→ Using  $R_m^I$  method

(restricted to **only statistical errors only in the  $\chi^2$** )

$$a_{\mu}^{\pi^+\pi^-} (0.324 \leq \sqrt{s} \leq 0.975 \text{ GeV}) = 486.7 \pm 5.2$$

$$a_{\mu}^{\pi^+\pi^-} (0.592 \leq \sqrt{s} \leq 0.922 \text{ GeV}) = 378.1 \pm 3.0$$

$$a_{\mu}^{\pi^+\pi^-} (0.6 \leq \sqrt{s} \leq 0.9 \text{ GeV}) = 367.67 \pm 2.97$$

# KLOE combination: $\pi^+\pi^-$ implications

⇒ Full combination of  $\pi^+\pi^-$   
data **remains stable**

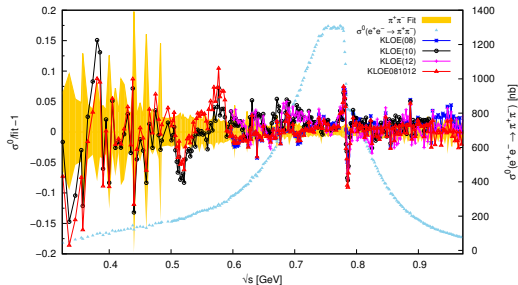
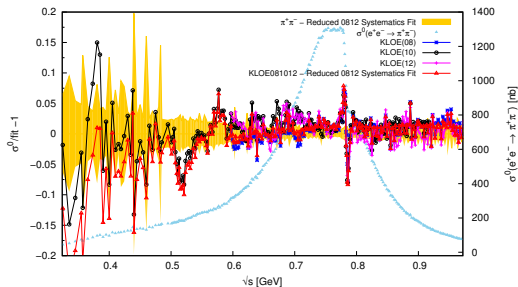
→ Removing KLOE(08) & KLOE(12)  
correlated systematics  $> 0.9$  GeV:

$$a_\mu^{\pi^+\pi^-} (\leq 2.0\text{GeV}) = 503.71 \pm 2.10$$

→ Calculating a full  $\pi^+\pi^-$   
combination with **all KLOE**  
systematics:

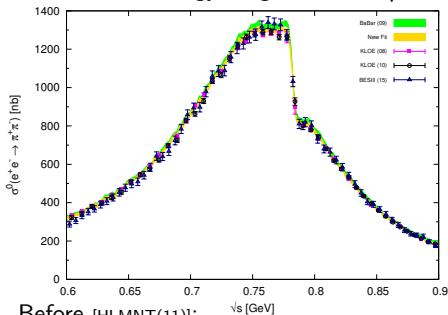
$$a_\mu^{\pi^+\pi^-} (\leq 2.0\text{GeV}) = 504.00 \pm 2.12$$

⇒ **Effect is less than 14% of the error**



$e^+e^- \rightarrow \pi^+\pi^-$ : inclusion of BESIII data [M. Ablikim *et al.* [BESIII

Collaboration], Phys. Lett. B 753 (2016) 629]

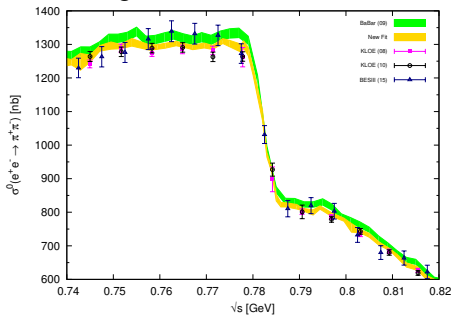
Measured in energy range of  $0.6 \leq \sqrt{s} \leq 0.9$  GeV using the radiative return method

Before [HLMNT(11)]:

$$a_{\mu}^{\pi^+\pi^-} (0.305 \leq \sqrt{s} \leq 2.0 \text{ GeV}) = (505.77 \pm 3.09) \times 10^{-10} \quad ; \quad \chi_{\min}^2/\text{d.o.f.} = 1.35$$

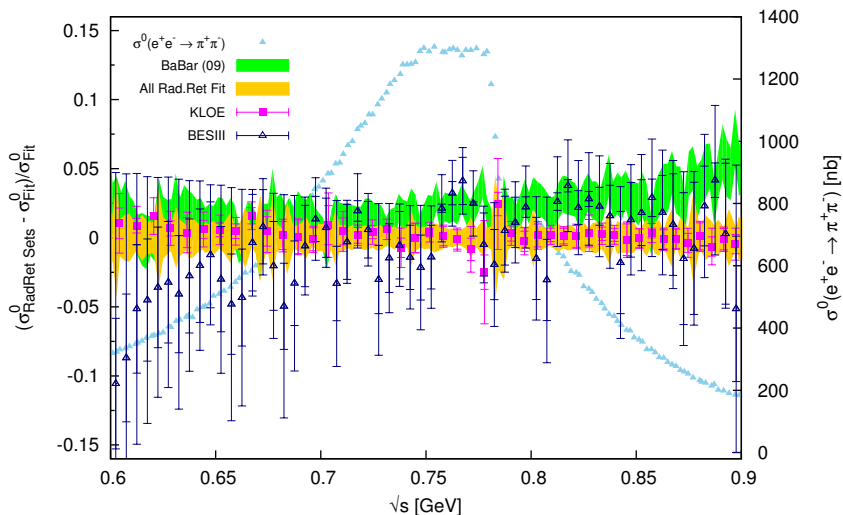
After (excluding new KLOE data):

$$a_{\mu}^{\pi^+\pi^-} (0.305 \leq \sqrt{s} \leq 2.0 \text{ GeV}) = (504.15 \pm 2.17) \times 10^{-10} \quad ; \quad \chi_{\min}^2/\text{d.o.f.} = 1.26$$



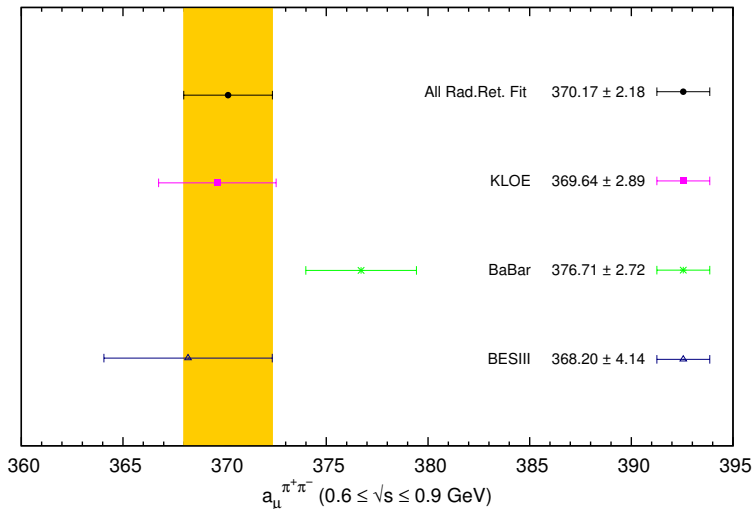
$e^+e^- \rightarrow \pi^+\pi^-$ : radiative return comparison

⇒ Using **full method** with **complete** radiative return data...



$e^+e^- \rightarrow \pi^+\pi^-$ : radiative return comparison

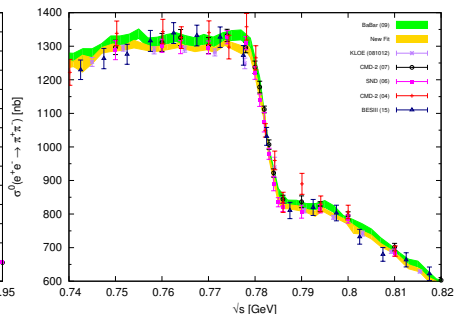
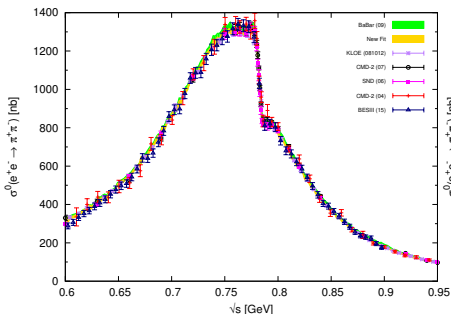
⇒ Using **full method** with **complete** radiative return data...



$a_\mu^{\pi^+\pi^-}$ : preliminary result

$\pi^+\pi^-$  channel contributes  $\sim 73\%$  to  $a_\mu^{\text{had}}$ , LOVP

→ Channel now combines **999 data points** from **25 experimental measurements** (5 radiative return measurements with full uncertainty covariance matrices)



⇒ Full combination yields...  $a_\mu^{\pi^+\pi^-} (0.6 \leq \sqrt{s} \leq 0.9 \text{ GeV}) = 370.88 \pm 1.52$

$a_\mu^{\pi^+\pi^-} (0.305 \leq \sqrt{s} \leq 2.0 \text{ GeV}) = (504.00 \pm 2.12) \times 10^{-10}$  ;  $\chi^2_{\text{min}}/\text{d.o.f.} = 1.28$



# Results comparison

Comparative (**PRELIMINARY**) results after introducing new fitting method and including new data in leading channels (in units  $10^{-10}$ ):

Channel	HLMNT(11)		This work		Difference
	$a_\mu(\leq 2\text{GeV})$	$\chi_{\min}^2/\text{d.o.f.}$	$a_\mu(\leq 2\text{GeV})$	$\chi_{\min}^2/\text{d.o.f.}$	
$\pi^+\pi^-$	$505.77 \pm 3.09$	1.39	$504.00 \pm 2.12$	1.28	-1.77
$\pi^+\pi^-\pi^0$	$47.51 \pm 0.98$	3.04	$47.68 \pm 0.97$	2.51	+0.17
$\pi^+\pi^-\pi^0\pi^0$	$20.73 \pm 1.28$	1.29	$20.42 \pm 1.11$	0.72	-0.31
$\pi^+\pi^-\pi^+\pi^-$	$14.73 \pm 0.48$	1.81	$14.91 \pm 0.47$	1.61	+0.18
$K^+K^-$	$22.12 \pm 0.41$	1.95	$22.75 \pm 0.27$	2.09	+0.63
$K_s^0 K_l^0$	$13.33 \pm 0.17$	1.10	$13.09 \pm 0.13$	0.73	-0.24
$\pi^+\pi^-\pi^+\pi^-\pi^0$	$1.42 \pm 0.09$	1.21	$1.39 \pm 0.09$	1.05	-0.03
$\pi^+\pi^-\pi^+\pi^-\pi^+\pi^-$	$0.30 \pm 0.01$	1.67	$0.28 \pm 0.01$	1.49	-0.02
Total:	$625.91 \pm 3.55$		$624.52 \pm 2.63$		-1.39

- Changes in almost each channel due to new method/new data
- Improved error estimate and goodness-of-fit
- Much more statistically reliable and trustworthy fitting method
- Reduction in overall mean value would mean increased  $g-2$  discrepancy,  $\delta a_\mu$
- Room for improvement in clustering, new data etc.  $\Rightarrow$  Results expected to change

# Outlook and the future...

- More data = **improved precision**
- Looking forward to **new measurements** from KLOE, CMD-3, SND, BaBar, BESIII (& Belle)
- AND **new experimental measurements** from Fermilab and J-PARC
- ⇒ As it stands (from **ONLY** the results in this talk)...

LO hadronic	694.9 (4.3)	→	693.5 (3.6) this work
Theory total	11659182.8 (4.9)	→	11659181.4 (4.4) this work
Experiment			11659208.9 (6.5) world avg
Exp - Theory	26.1 (8.0)	→	27.48 (7.8)
$\Delta a_\mu$	3.3 $\sigma$	→	3.5 $\sigma$

# Summary

- Data has to be treated with intense care and scrutiny in combination
- $f_k$  method shown to exhibit potential bias...
- $R_m^I$  method avoids bias by fixing covariance matrix through iterative fit
- Method allows for more stable, robust and trustworthy treatment of uncertainties (covariance matrices)
- $K^+K^-$  data tension unresolved through energy shift study  
⇒ BaBar problem
- Unresolved issues regarding combination of KLOE data and systematic uncertainties  
⇒ Full  $\pi^+\pi^-$  combination remains stable
- Lower mean value and reduced uncertainty for  $\pi^+\pi^-$  and  $a_\mu^{\text{had, LOVP}}$  overall
- Results indicate slight increase to  $(g-2)_\mu$  discrepancy

## Thank You

# Extra Slides

# An iterated fit

Iterating the fit ensures an unbiased solution  $\rightarrow$  Forces the fit to converge to an unbiased result

$$\chi_1^2(R_m) = \sum_{m=1}^{N_{clu}} \sum_{n=1}^{N_{clu}} \sum_{i=1}^{N^{(k,m)}} \sum_{j=1}^{N^{(k,n)}} (R_i^{(k,m)} - R_m) \mathbf{C}_0(i^{(k,m)}, j^{(l,n)}) (R_j^{(k,n)} - R_n)$$

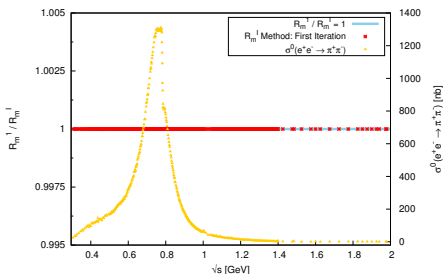
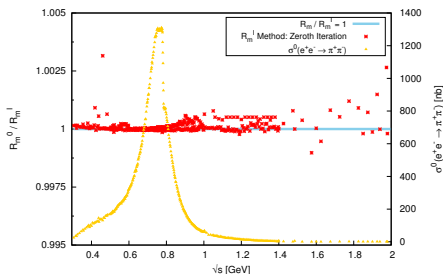
$$\mathbf{C}_0(i^{(k,m)}, j^{(l,n)}) = \mathbf{C}^{\text{stat}}(i^{(k,m)}, j^{(l,n)}) + \mathbf{C}^{\text{sys}\%}(i^{(k,m)}, j^{(l,n)}) R_m^0 R_n^0$$

$\Rightarrow$  Feed the fitted  $R_m$  values into the next iteration...

$$\chi_2^2(R_m) = \sum_{m=1}^{N_{clu}} \sum_{n=1}^{N_{clu}} \sum_{i=1}^{N^{(k,m)}} \sum_{j=1}^{N^{(k,n)}} (R_i^{(k,m)} - R_m) \mathbf{C}_1(i^{(k,m)}, j^{(l,n)}) (R_j^{(k,n)} - R_n)$$

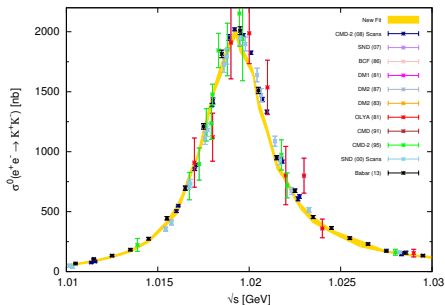
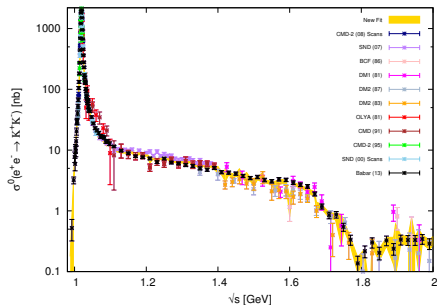
$$\mathbf{C}_1(i^{(k,m)}, j^{(l,n)}) = \mathbf{C}^{\text{stat}}(i^{(k,m)}, j^{(l,n)}) + \mathbf{C}^{\text{sys}\%}(i^{(k,m)}, j^{(l,n)}) R_m^1 R_n^1$$

Repeat until fit converges and returns final fitted values for clusters,  $R_m = R_m^I$



# $e^+e^- \rightarrow K^+K^-$ : inclusion of BaBar data [ J. P. Lees *et al.* [BaBar

Collaboration], Phys. Rev. D **88** (2013) no.3, 032013]



Before [HLMNT(11)]:

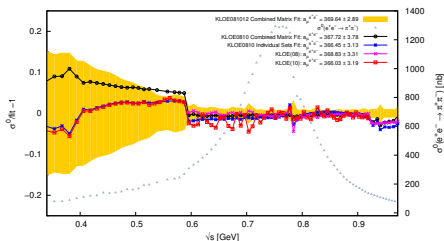
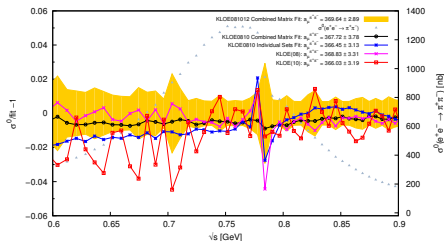
$$a_{\mu}^{\pi^+\pi^-}(\sqrt{s} \leq 2.0 \text{ GeV}) = (22.12 \pm 0.41) \times 10^{-10} \quad ; \quad \chi_{\min}^2/\text{d.o.f.} = 1.95$$

After:

$$a_{\mu}^{\pi^+\pi^-}(\sqrt{s} \leq 2.0 \text{ GeV}) = (22.75 \pm 0.27) \times 10^{-10} \quad ; \quad \chi_{\min}^2/\text{d.o.f.} = 2.09$$

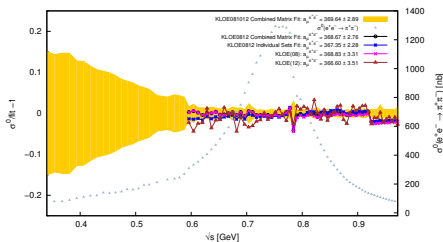
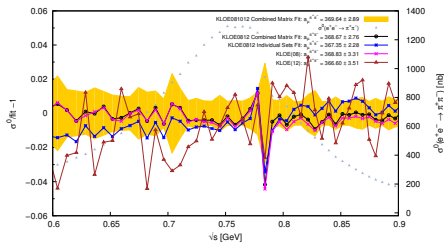
## KLOE Combination: KLOE(0810)

KLOE Sets	Individual Sets		Combination Matrix		
	$a_\mu$	$\chi^2_{\min}/\text{d.o.f.}$	$a_\mu$	$\chi^2_{\min}/\text{d.o.f.}$	Difference
KLOE(0810)	$366.45 \pm 3.13$	7.65	$367.72 \pm 3.78$	7.74	+1.27
KLOE(0812)	$367.35 \pm 2.28$	1.01	$368.67 \pm 2.76$	1.29	+1.32
KLOE(1012)	$366.91 \pm 2.29$	3.06	$366.49 \pm 2.27$	2.99	-0.42
KLOE(081012)	$367.34 \pm 2.06$	1.21	$369.64 \pm 2.89$	1.27	+2.30



## KLOE Combination: KLOE(0812)

KLOE Sets	Individual Sets		Combination Matrix		Difference
	$a_\mu$	$\chi^2_{\min}/\text{d.o.f.}$	$a_\mu$	$\chi^2_{\min}/\text{d.o.f.}$	
KLOE(0810)	$366.45 \pm 3.13$	7.65	$367.72 \pm 3.78$	7.74	+1.27
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## KLOE Combination: KLOE(1012)

KLOE Sets	Individual Sets		Combination Matrix		Difference
	$a_\mu$	$\chi^2_{\min}/\text{d.o.f.}$	$a_\mu$	$\chi^2_{\min}/\text{d.o.f.}$	
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