# Nonlinear equations from linear ones 

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## Matrix equations

- $G$ - a $n$-dimensional Lie group,
- $\mathfrak{g}$-its Lie algebra.
- Equation

$$
\frac{d}{d t} K(t)=M(t) K(t), \quad K(0)=I
$$

where

- $\mathbb{R} \ni t \mapsto M(t) \in \mathfrak{g}$ - a given curve in $\mathfrak{g}$,
- $K(t)$ - a curve in $G$, which is a solution to the equation.


## Motivation

In quantum mechanical applications the Schrödinger equation for an $n$-level system governed by a time-dependent Hamiltonian reads:

$$
\begin{equation*}
i \frac{d \psi}{d t}=H(t) \psi \tag{1}
\end{equation*}
$$

Writing the solution of with an initial condition $\psi(0)$ as
$\psi(t)=U(t) \psi(0)$ and substituting $H(t)=i M(t)$, we obtain:

$$
\frac{d}{d t} U(t)=M(t) U(t), \quad U(0)=I
$$

where $U(t) \in G=U(n)$.

This equations can be also treated as a classical control system on the Lie group $G$.

## Wei-Norman method

- $G$ - $n$-dimensional Lie group
- $\mathfrak{g}$ - its Lie algebra (simple, complex)
- $\mathbb{R} \ni t \mapsto M(t) \in \mathfrak{g}$ - a curve in $\mathfrak{g}$.
- $K(t)$ - a curve in $G$ given by the differential equation:

$$
\frac{d}{d t} K(t)=M(t) K(t), \quad K(0)=I
$$

- $X_{k}, k=1, \ldots, n$ is some basis in $\mathfrak{g}$, then:

$$
M(t)=\sum_{k=1}^{n} a_{k}(t) X_{k}
$$

- We look for the solution $K(t)$ in the form

$$
K(t)=\prod_{k=1}^{n} \exp \left(u_{k}(t) X_{k}\right)
$$

## Wei-Norman method

- Differentiating $\left({ }^{\prime}=d / d t\right)$ and commuting...

$$
K^{\prime}=\sum_{l=1}^{n} u_{l}^{\prime} \prod_{k<l} \operatorname{Ad}_{\exp \left(u_{k} X_{k}\right)} X_{l} K
$$

where Ad is the adjoint action of $G$ on $\mathfrak{g}$,

$$
\operatorname{Ad}_{g} X:=g X^{-1}, \quad g \in G, \quad X \in \mathfrak{g} .
$$

- Using

$$
\operatorname{Ad}_{\exp (f \cdot X)}=\exp \left(f \cdot \operatorname{ad}_{X}\right)
$$

where $\operatorname{ad}_{X}=[X, \cdot]$ is the adjoint action of $\mathfrak{g}$ on itself, we obtain:

- ... we get

$$
K^{\prime}=\sum_{l=1}^{n} u_{l}^{\prime} \prod_{k<l} \exp \left(u_{k} \operatorname{ad}_{x_{k}}\right) X_{l} K .
$$

## Wei-Norman method

- Comparing

$$
\frac{d}{d t} K(t)=\underbrace{\sum_{l=1}^{n} u_{l}^{\prime} \prod_{k<l} \exp \left(u_{k} \operatorname{ad}_{X_{k}}\right) \cdot X_{l}}_{M(t)=\sum_{k=1}^{n} a_{k}(t) X_{k}} K(t)
$$

- ... we obtain equations for the (unknown) coefficients $u_{j}$

$$
\mathbf{a}=A \mathbf{u}^{\prime}, \quad \mathbf{u}^{\prime}=A^{-1} \mathbf{a}
$$

where

$$
A_{j l}=A_{j l}^{(l)}, \quad A^{(l)}=\prod_{k<l} \exp \left(u_{k} \operatorname{ad}_{X_{k}}\right)
$$

It can be shown, that $A$ is invertible, at least locally

## Choice of basis

- A depends on the choice of (an ordered) basis in $\mathfrak{g}$
- Example: $\mathfrak{s l}(2, \mathbb{C})$ (J. Cariñena, J. Grabowski, G. Marmo, Lie-Scheffers Systems: A Geometric AApproach) for

$$
X_{1}=\frac{1}{2}\left[\begin{array}{cc}
1 & 0 \\
0 & -1
\end{array}\right], \quad X_{2}=\left[\begin{array}{ll}
0 & 1 \\
0 & 0
\end{array}\right], \quad X_{3}=\left[\begin{array}{ll}
0 & 0 \\
1 & 0
\end{array}\right]
$$

we get

$$
u_{1}^{\prime}=a_{1} e^{-u_{2}}-a_{3} u_{1}^{2} e^{u_{2}}, \quad u_{2}^{\prime}=a_{2}+2 a_{3} u_{1} e^{u_{2}}, \quad u_{3}^{\prime}=a_{3} \mathrm{e}^{u_{2}}
$$

whereas, for

$$
X_{1}=\left[\begin{array}{ll}
0 & 1 \\
0 & 0
\end{array}\right], \quad X_{2}=\frac{1}{2}\left[\begin{array}{cc}
1 & 0 \\
0 & -1
\end{array}\right], \quad X_{3}=\left[\begin{array}{ll}
0 & 0 \\
1 & 0
\end{array}\right]
$$

a much 'nicer' Riccati system

$$
u_{1}^{\prime}=a_{1}+a_{2} u_{1}+a_{3} u_{1}^{2}, \quad u_{2}^{\prime}=a_{2}+2 a_{3} u_{1}, \quad u_{3}^{\prime}=a_{3} \mathrm{e}^{u_{2}}
$$

- How to find such a 'canonical' ordered basis for an arbitrary Lie group?


## Example: spin in rotating magnetic field

- Time dependent Hamiltonian:

$$
H(t)=-\frac{B \cos (\omega t)}{2} \sigma_{x}-\frac{B \sin (\omega t)}{2} \sigma_{y}=-\frac{B}{2}\left[\begin{array}{cc}
0 & \exp (-i \omega t) \\
\exp (i \omega t) & 0
\end{array}\right], \quad U^{\prime}=-i H U
$$

- In the above 'nice' basis the equations read:

$$
u_{1}^{\prime}(t)=\frac{i B}{2}\left(\mathrm{e}^{-i \omega t}+\mathrm{e}^{i \omega t}\left(u_{1}(t)\right)^{2}\right), \quad u_{2}^{\prime}(t)=i B \mathrm{e}^{i \omega t} u_{1}(t), \quad u_{3}^{\prime}(t)=\frac{i B}{2} \mathrm{e}^{i \omega t} \mathrm{e}^{u_{2}(t)}
$$

- ... and are easily solved

$$
U=\left[\begin{array}{cc}
\left(\cos \left(\frac{\Omega t}{2}\right)+\frac{i \sin \left(\frac{\Omega t}{2}\right) \omega}{\Omega}\right) \mathrm{e}^{-i \frac{\omega t}{2}} & -\frac{B}{2 \Omega}\left(\mathrm{e}^{-i t \Omega}-1\right) \mathrm{e}^{i \frac{\Omega-\omega}{2} t} \\
\frac{B}{2 \Omega}\left(\mathrm{e}^{i t \Omega}-1\right) \mathrm{e}^{-i \frac{\Omega-\omega}{2} t} & \left(\cos \left(\frac{\Omega t}{2}\right)-\frac{i \sin \left(\frac{\Omega t}{2}\right) \omega}{\Omega}\right) \mathrm{e}^{i \frac{\omega t}{2}}
\end{array}\right], \quad \Omega=\sqrt{B^{2}+\omega^{2}}
$$

## Example: $\mathfrak{s l}(3, \mathbb{C})$

1. A system of two coupled Riccati equations:

$$
\begin{aligned}
& u_{1}^{\prime}=a_{1}+\left(2 a_{5}-a_{4}\right) u_{1}+a_{6} u_{2}-a_{8} u_{1}^{2}-a_{7} u_{1} u_{2}, \\
& u_{2}^{\prime}=a_{2}+a_{3} u_{1}+\left(a_{4}+a_{5}\right) u_{2}-a_{8} u_{1} u_{2}-a_{7} u_{2}^{2},
\end{aligned}
$$

2. A scalar Riccati equation for $u_{3}$ :

$$
u_{3}^{\prime}=\left(a_{3}-a_{8} u_{2}\right)+\left(2 a_{4}-a_{5}+a_{8} u_{1}-a_{7} u_{2}\right) u_{3}+\left(a_{7} u_{1}-a_{6}\right) u_{3}^{2},
$$

3. The rest

$$
\begin{aligned}
u_{4}^{\prime} & =a_{4}-a_{6} u_{3}+a_{7}\left(u_{1} u_{3}-u_{2}\right) \\
u_{5}^{\prime} & =a_{5}-a_{8} u_{1}-a_{7} u_{2} \\
u_{6}^{\prime} & =\left(a_{6}-a_{7} u_{1}\right) \mathrm{e}^{2 u_{4}-u_{5}}, \\
u_{7}^{\prime} & =\left(a_{7} u_{3}+a_{8}\right) u_{6} \mathrm{e}^{-u_{4}+2 u_{5}}+a_{7} \mathrm{e}^{u_{4}+u_{5}}, \\
u_{8}^{\prime} & =\left(a_{8}+a_{7} u_{3}\right) \mathrm{e}^{-u_{4}+2 u_{5}},
\end{aligned}
$$

which are solved by simple consecutive integrations, once solutions of the Riccati equations are known.

## Arbitrary simple Lie algebra $\mathfrak{g}$

- The construction of a 'canonical' basis can be done for all simple algebras (classical and exceptional) except the exceptional ones $G_{2}, F_{4}$, and $E_{8}$
- The resulting Wei-Norman equations split into several blocks of coupled Riccati and linear equations and 'trivial' ones (solvable by consecutive integrations)
- The construction hinges on two things
- A decomposition of a simple Lie algebra into a sum of commutative algebras (can be done for all simple Lie algebras)
- A particular property of the adjoint endomorhism (in all simple Lie algebras, except $G_{2}$, $F_{4}$, and $E_{8}$ )


## Decomposition

- The Cartan decomposition

$$
\mathfrak{g}=\bigoplus_{\alpha \in \Phi_{-}} \mathfrak{g}_{\alpha} \oplus \mathfrak{h} \oplus \bigoplus_{\alpha \in \Phi_{+}} \mathfrak{g}_{\alpha}
$$

where $\mathfrak{h}$ - the Cartan subalgebra, $\mathfrak{g}$ - the root spaces,

$$
\mathfrak{g}_{\alpha}:=\{X \in \mathfrak{g}:[H, X]=\alpha(H) X \forall H \in \mathfrak{h}\}
$$

and $\Phi_{ \pm}$- the set of positive (negative) roots.

- Decomposition into commutative subalgebras

For a basis $\left\{\alpha_{1}, \ldots, \alpha_{N}\right\}$ of the roots system consisting of positive simple roots

define
$\Phi_{k}:=\left\{\beta: \beta=\sum_{k}^{N} n_{i} \alpha_{i}, n_{i}>0\right\}, \quad \mathfrak{a}_{k}:=\operatorname{span}\left\{X_{\beta}: \beta \in \Phi_{k}\right\}, \quad \tilde{\mathfrak{a}}_{k}:=\operatorname{span}\left\{X_{-\beta}: \beta \in \Phi_{k}\right\}$
$\mathfrak{a}_{k}$ are commutative and

$$
\mathfrak{g}=\bigoplus_{k=1}^{N} \mathfrak{a}_{k} \oplus \mathfrak{h} \oplus \bigoplus_{j=1}^{n} \tilde{\mathfrak{a}}_{j}
$$

## Properties of the adjoint endomorphism

- $\mathfrak{g}$ - a simple Lie algebra not equal to $G_{2}, F_{4}, E_{8}$
- $\alpha-$ a root and $X_{\alpha} \in \mathrm{g}$ - the corresponding root vector


## Then

1. The image of $\left(\operatorname{ad}_{X_{\alpha}}\right)^{2}$ is equal to $\mathfrak{g}_{\alpha}$
2. $\left(\mathrm{ad}_{X_{\alpha}}\right)^{3}=0$.

- Remark: for $G_{2}, F_{4}$, and $E_{8}$, there are roots for which $\left(\operatorname{ad}_{X_{\alpha}}\right)^{3} \neq 0$, but for all roots we have $\left(\mathrm{ad}_{X_{\alpha}}\right)^{5}=0$
- Corollary: For $X \in \mathfrak{a}_{k}$ or $X \in \widetilde{\mathfrak{a}}_{k}$ the matrix of $\exp \left(\operatorname{ad}_{X}\right)$ is a quadratic polynomial in $\mathrm{ad}_{X}$ and it is block diagonal with respect to the decomposition

$$
\mathfrak{g}=\mathfrak{a}_{1} \oplus \ldots \oplus \mathfrak{a}_{k-1} \oplus \underbrace{\left(\mathfrak{b}_{k} \oplus \mathfrak{h} \oplus \widetilde{\mathfrak{b}}_{k}\right)}_{\text {one block }} \oplus \widetilde{\mathfrak{a}}_{k-1} \oplus \ldots \oplus \tilde{\mathfrak{a}}_{1}
$$

- The equation $\mathbf{u}^{\prime}=A^{-1}$ a separates into blocks:
- matrix $N \times N$ Riccati equation for $u_{1}, \ldots, u_{N}$, corresponding to $\mathfrak{a}_{1}$,
- matrix $(N-1) \times(N-1)$ Riccati equation for $u_{N+1}, \ldots, u_{2 N-1}$, corresponding to $\mathfrak{a}_{2}$,

- scalar Riccati equation for $u_{N(N+1) / 2}$, corresponding to $\mathfrak{a}_{N}$,
- the remaining equations which, are solved by consecutive integrations (provided the solutions of Riccati equations are known).


## Summary

- An algorithm for reducing the highly non linear system of Wei-Norman equations

$$
\sum_{k=1}^{n} a_{k} X_{k}=\sum_{l=1}^{n} u_{l}^{\prime} \prod_{k<l} \exp \left(u_{k} \operatorname{ad}_{X_{k}}\right) \cdot X_{l},
$$

for the parameters $u_{k}$ to a hierarchy of Riccati matrix equations and integrals.

- It is known that every system of Riccati equations is related to a Lie group action and solution of this system is equivalent to solution of the system of the form

$$
\frac{d}{d t} K(t)=M(t) K(t), \quad K(0)=I
$$

but not every system of this form is equivalent to Riccati equation system. Here we provide an explicit construction of a hierarchy of Riccati matrix equations equivalent to the equation on a Lie group for all classical groups (and the exceptional ones $E_{6}$ and $E_{7}$ )

- For the exceptional algebras $G_{2}, F_{4}$, a nd $E_{8}$ the nonlinearities are of 4-th order (c.f. the property of the adjoint endomorphism)


## References

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