



**NUMEN: outlook towards high beam intensities
experiments**

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NUclear MAtrix Elements for N eutrinoless double beta decay



JMEN collaboration

- LABORATORI NAZIONALI DEL SUD, CATANIA, ITALY
- ISTITUTO DI FISICA E ASTRONOMIA, UNIV. DI CATANIA, CATANIA, ITALY
- UNIVERSITÀ DI CATANIA, CATANIA, ITALY
- ISTITUTO STUDI DI ENNA "KORE", ENNA, ITALY
- UNIVERSITÀ DI GENOVA, GENOVA, ITALY
- UNIVERSITÀ DI TORINO, TORINO, ITALY
- INSTITUT FÜR THEORETISCHE PHYSIK, GIESSEN UNIVERSITY, GERMANY
- DEPARTMENT OF PHYSICS AND HINP, THE UNIV. OF IOANNINA, IOANNINA, GREECE
- INSTITUTO DE FISICA DA UNIVERSIDADE DE SAO PAULO, BRAZIL
- INSTITUTO DE FISICA DA UNIV. FEDERAL FLUMINENSE, NITEROI, BRAZIL
- IZMIR ANADOLU UNIVERSITY, ANTALYA, TURKEY



$$1 / T_{1/2}^{0\nu} (0^+ \rightarrow 0^+) = G_{01} \left| M^{\beta\beta 0\nu} \right|^2 \left| \langle \dots \right|$$

Emerging collaborations

- UNIVERSIDAD DE CIENCIAS NUCLEARES, UNAM, MEXICO
- UNIVERSIDAD DE COSTA RICA, SAN JOSE, COSTA RICA
- INSTITUT DE PHYSIQUE, UNIVERSITÉ HASSAN II – CASABLANCA, MOROCCO

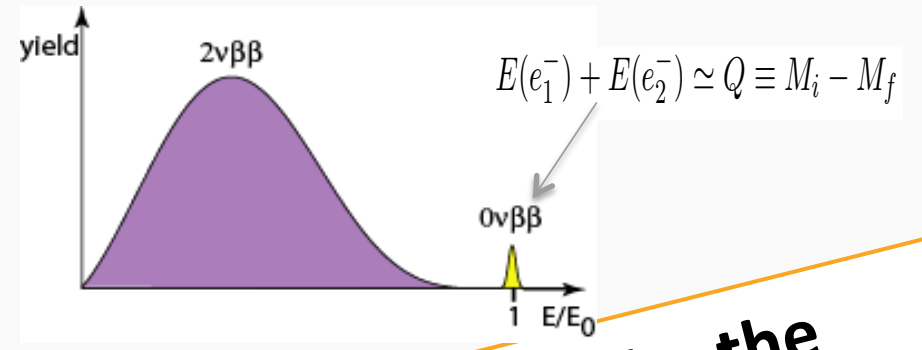
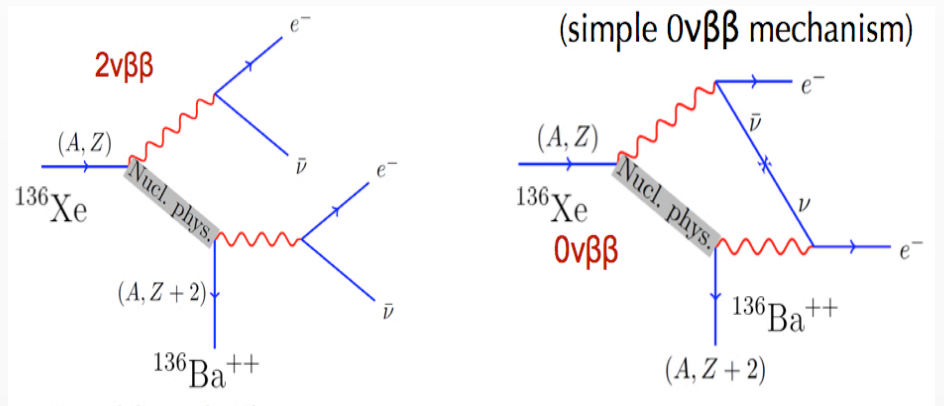
NATIONAL LABORATORY



from **neutrino oscillations** we know \longrightarrow **neutrino mass $\neq 0$**

What about the **m_ν absolute value** and the neutrino **nature** :
Dirac or Majorana particle?

Double β -decay



DBD : $(A, Z) \rightarrow (A, Z+2) + 2e^- + 2\nu_e$

$\approx 10^{19-21} \text{ y}$



The very rare $0\nu\beta\beta$ decay is potentially the best way to probe the Majorana or Dirac nature of neutrino and to extract its effective mass.

DBD : $(A, Z) \rightarrow (A, Z+2) + 2e^-$

- Respect the conservation of lepton number
- Neutrino has mass
- Neutrino is Majorana particle
- Violates the leptonic number conservation

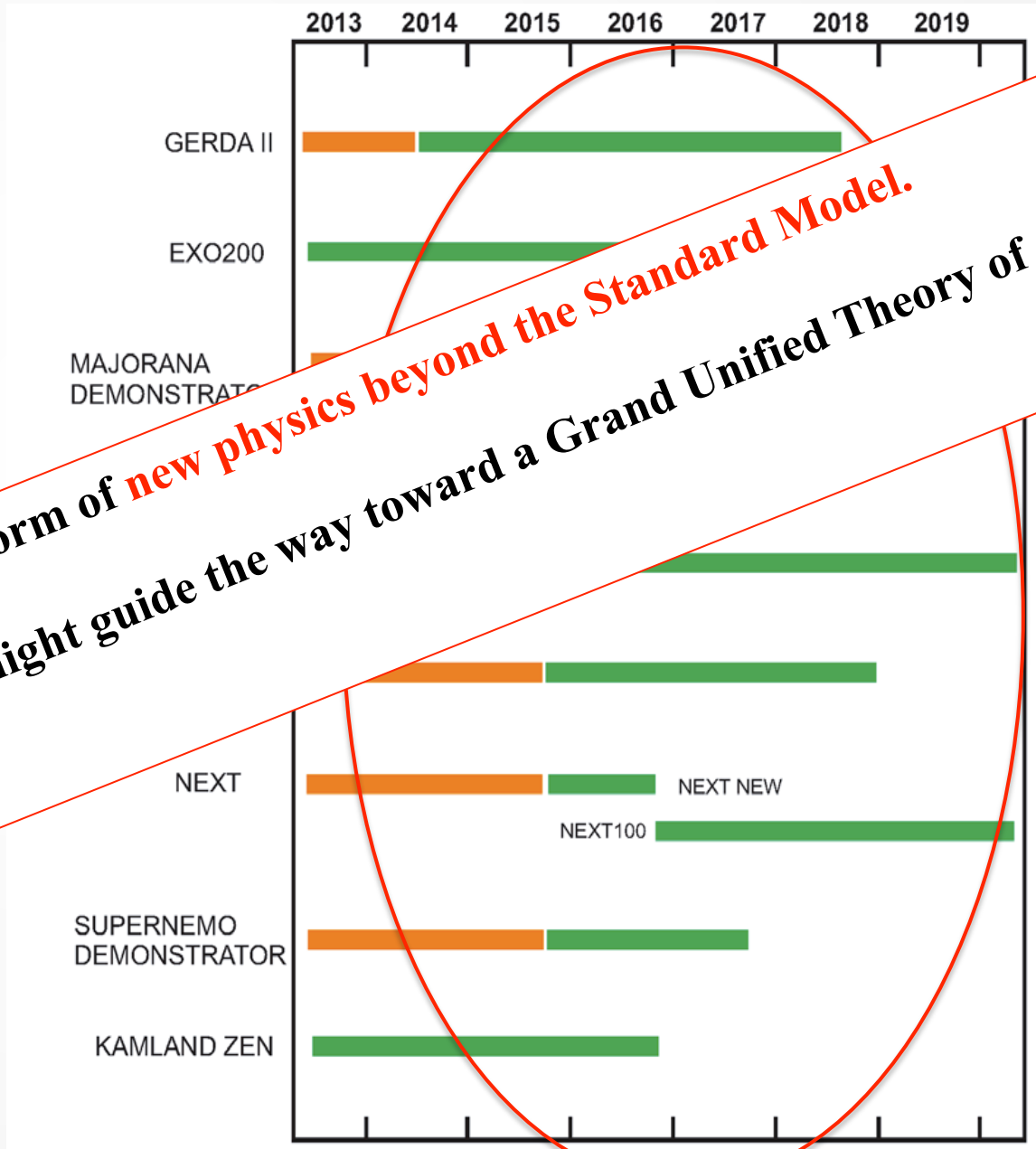
$^{82}\text{Se}, ^{100}\text{Mo}, ^{48}\text{Ca}, ^{76}\text{Ge}, \dots$

Search for $0\nu\beta\beta$ decay: a worldwide race



Experiment	Isotope	Lab	Status
GERDA II	^{76}Ge	LNGS	Phase I completed Migration to Phase II
EXO200	^{130}Te	LNGS	Data taking / Construction
MAJORANA DEMONSTRATOR	^{76}Ge	SURF	Construction
LEGEND-200	^{130}Te	SNOLAB	R&D / Construction
CUORE	^{82}Se (or others)	LSM	R&D / Construction
AMON	^{48}Ca	SNOLAB	R&D
PROSPECT	^{76}Ge	[Japan]	R&D
NUCLEON	^{100}Mo	[Korea]	R&D

NLDBD Report April 24, 2014



Neutrino mass explanations: all are based on some form of new physics beyond the Standard Model.

Measurement of the neutrino masses and mixing might guide the way toward a Grand Unified Theory of fundamental interactions.

The role of nuclear physics

For the $0\nu\beta\beta$ decay rate can be expressed as a product of independent factors, that also depend on a nuclear structure function containing physics beyond the Standard Model through the masses and the mixing coefficients of the neutrinos species :

$$1 / T_{1/2}^{0\nu} (0^+ \rightarrow 0^+) = G_{01} |M^{\beta\beta 0\nu}|^2 \left| \frac{\langle m_\nu \rangle}{m_e} \right|^2 \rightarrow \langle m_\nu \rangle = \sum_i |U_{ei}|^2 m_i$$

new physics inside !

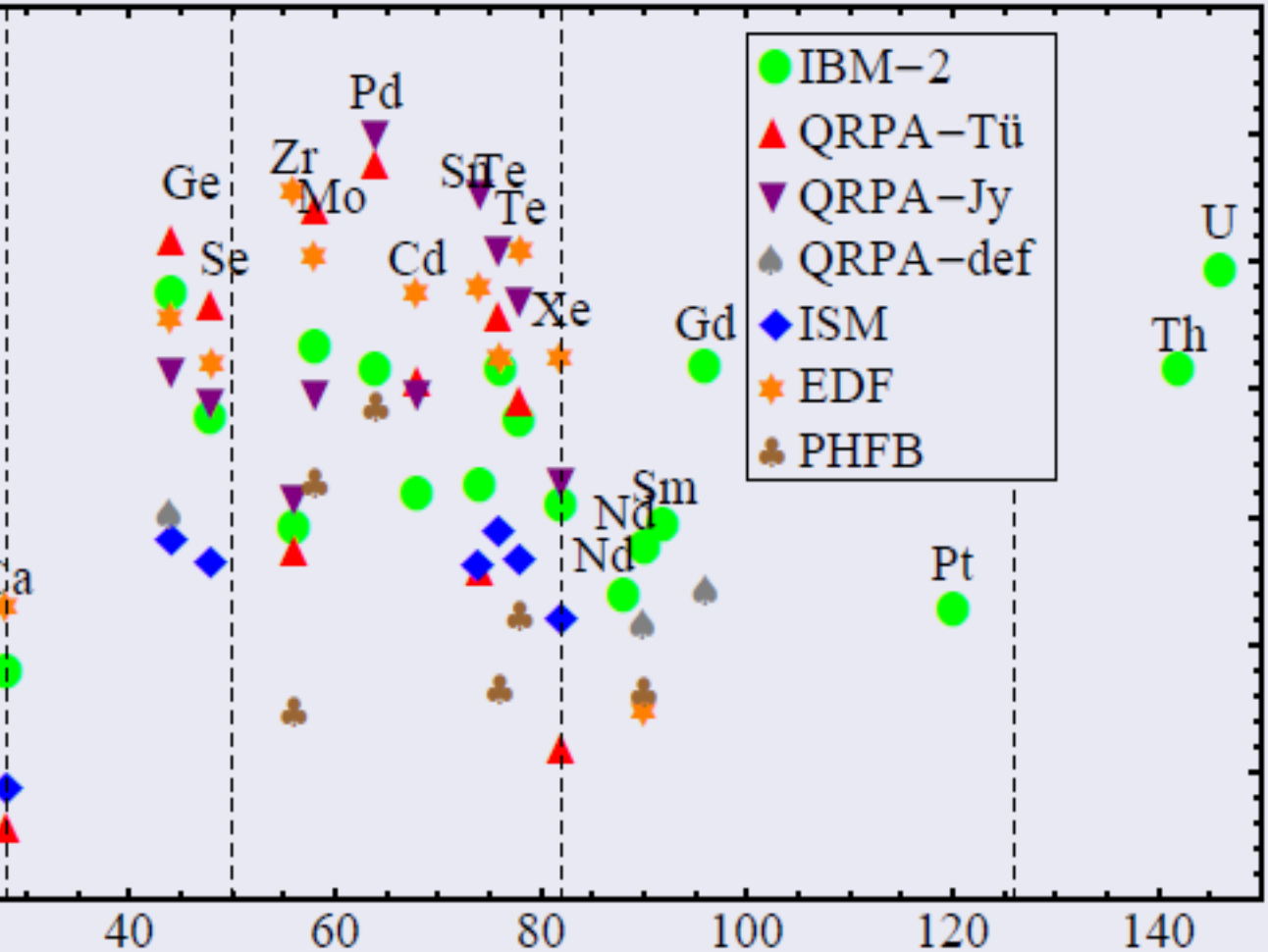
$$|M_\varepsilon^{\beta\beta 0\nu}|^2 = \left| \left\langle 0_f \left\| \hat{O}_\varepsilon^{\beta\beta 0\nu} \right\| 0_i \right\rangle \right|^2$$

Thus, if the $M^{\beta\beta 0\nu}$ nuclear matrix elements were known with sufficient precision, the neutrino mass could be established from $0\nu\beta\beta$ decay rate measurements

State of the art NME calculations

Evaluation of $\left| M_{\varepsilon}^{\beta\beta 0\nu} \right|^2 = \left| \left\langle \Psi_f \left| \hat{O}_{\varepsilon}^{\beta\beta 0\nu} \right| \Psi_i \right\rangle \right|^2$

$$M_{\varepsilon}^{(0\nu)} = M_{GT}^{(0\nu)} - \left(\frac{g_V}{g_A} \right)^2 M_F^{(0\nu)} + M_T^{(0\nu)}$$



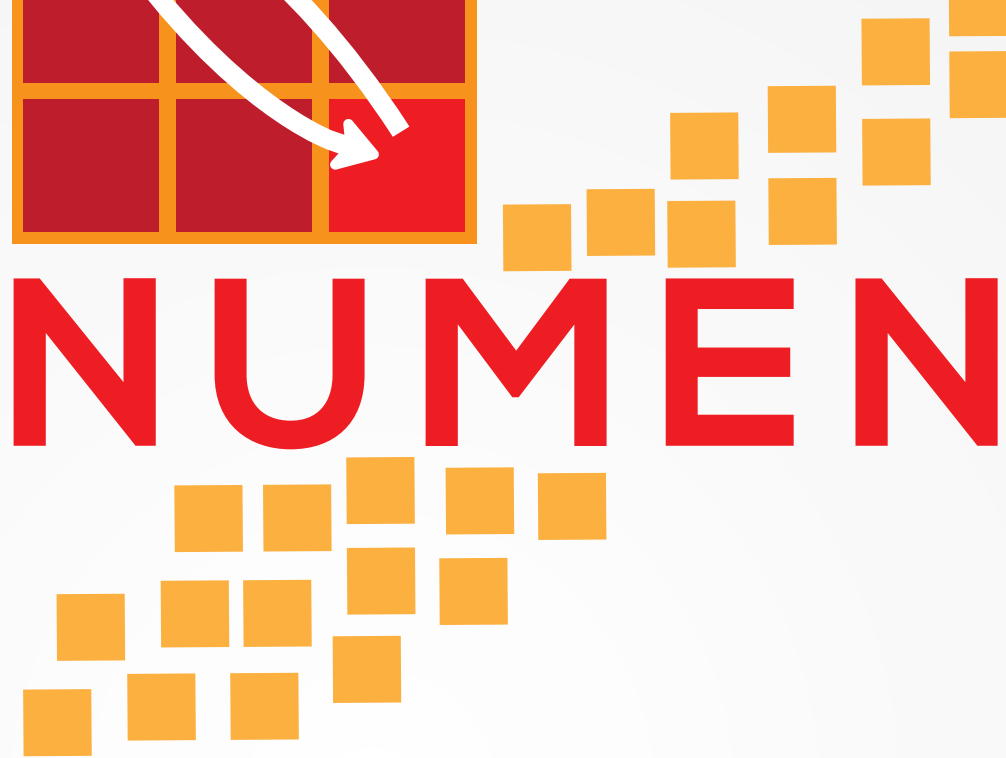
✓ **Calculations** (still sizeable uncertainties): QRPA, Large scale shell model, IBM

E. Caurier, et al., PRL 100 (2008) 052503
 N. L. Vaquero, et al., PRL 111 (2013) 142501
 J. Barea, PRC 87 (2013) 014315
 T. R. Rodriguez, PLB 719 (2013) 174
 F. Simkovic, PRC 77 (2008) 045503.
 F. Iachello et al. NPB 237-238 (2013) 21 - 23

✓ **Measurements** (still not conclusive $0\nu\beta\beta$):

- (π^+, π^-)
- single charge exchange
- electron capture
- transfer reactions ...

N. Auerbach, Ann. Of Phys. 192 (1989) 77
 S.J. Freeman and J.P. Schiffer JPG 39 (2012) 124
 D. Frekers, Prog. Part. Nucl. Phys. 64 (2010) 281
 J.P. Schiffer, et al. PRL 100 (2008) 112501



Qualitative is not enough...

**The challenge:
to access quantitative information !**

Double charge exchange reactions

Double charge exchange reactions are characterized by the transfer of isospin $\Delta T_Z = \pm 2$ with the mass number unchanged

Induced by strong interaction

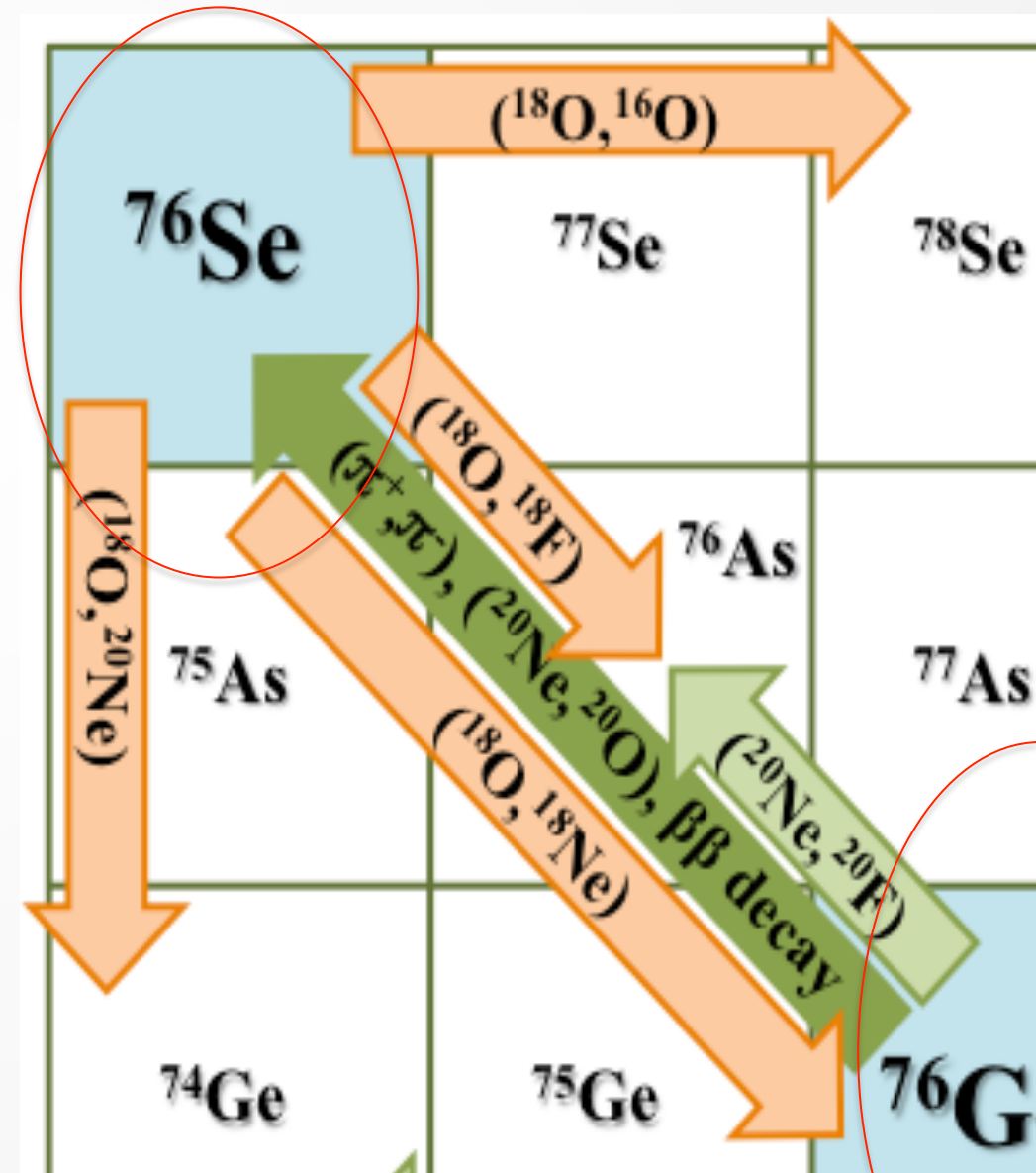
Sequential nucleon transfer mechanism
order:

► Brink's Kinematical matching conditions

D.M.Brink, et al., Phys. Lett. B 40 (1972) 37

Resonance exchange mechanism 2nd order

Possibility to go in both directions





Initial and final states: Parent/daughter states of the $0\nu\beta\beta$ are the same as those of target/residual nuclei in the DCE;

Spin-Isospin mathematical structure of the transition operator: **Fermi, Gamow-Teller rank-2 tensor together with higher L components** are present in both cases;

Connection between β -decay and Single Charge Exchange

Normalization of the charge exchange cross-section

$$B(\alpha) = \frac{1}{2J_i + 1} |M(\alpha)|^2$$

For single CEX:

$$\frac{d\sigma}{d\Omega}(q, \omega) = \hat{\sigma}_\alpha(E_p, A) F_\alpha(q, \omega) B_T(\alpha) B_P(\alpha)$$

$$\frac{B(GT)[(^3\text{He}, t); q = 0]}{B(GT)[\beta_{decay}]} = 1$$

➔ In the hypothesis of a surface localized process (for direct quasi elastic processes).

In the simple model one can assume that the DCE process is just a second order charge exchange, where projectile and target exchange two uncorrelated isovector virtual mesons.

Generalization

to DCE:

$$\frac{d\sigma^{DCE}}{d\Omega}(q, \omega) = \hat{\sigma}_\alpha^{DCE}(E_p, A) F_\alpha^{DCE}(q, \omega) B_T^{DCE}(\alpha) B_P^{DCE}(\alpha)$$

Superconducting Cyclotron

Superconducting Cyclotron in full operation since 1996.

Accelerate from Hydrogen to Helium.

...

Maximum nominal energy is 80 MeV/u.



MAGNEX magnetic spectrometer

Achieved resolution

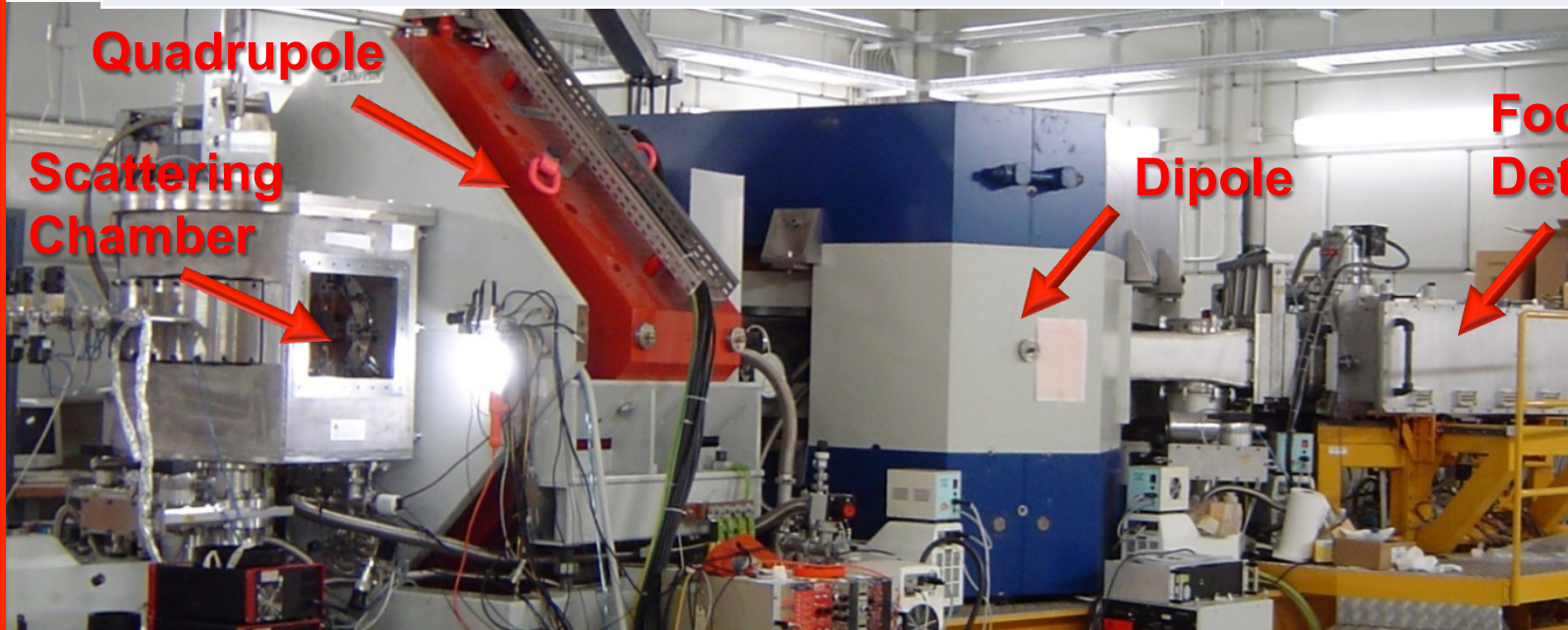
Energy $\Delta E/E \sim 1/1000$

Angle $\Delta\theta \sim 0.2^\circ$

Mass $\Delta m/m \sim 1/160$

F. Cappuzzello et al., *MAGNEX: an innovative large acceptance spectrometer for nuclear reaction studies*, in *Types, Uses and Safety* (Nova Publisher Inc., NY, 2011) pp. 1–63.

Optical characteristics	Measured value
Maximum magnetic rigidity	1.8 T m
Solid angle	50 msr
Momentum acceptance	-14.3%, +10%
Momentum dispersion for $k = -0.104$ (cm/%)	3.68



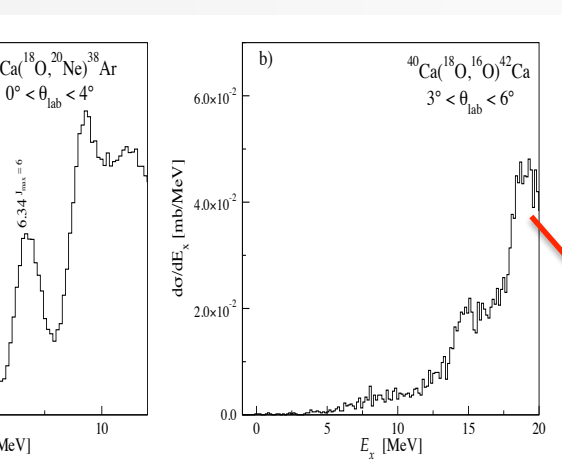
¹⁸O⁷⁺ beam from LNS Cyclotron at 270 MeV (10 pA)

⁴⁰Ca solid target of 300 μg/cm²

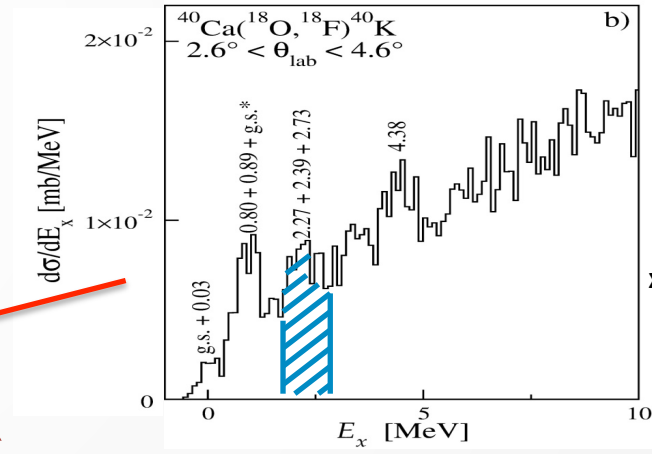
Ejectiles detected by the MAGNEX spectrometer

Angular setting

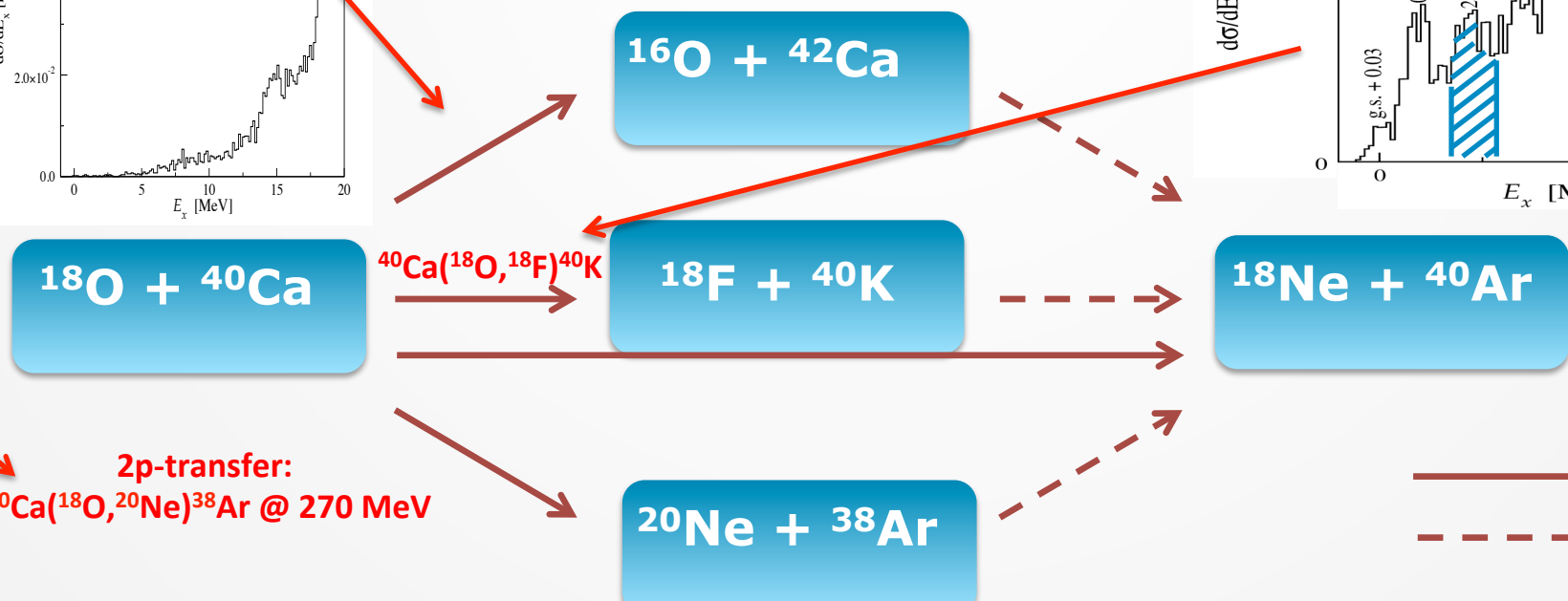
$$\theta_{opt} = 4^\circ \longrightarrow -2^\circ < \theta_{lab} < 10^\circ$$



2n-transfer:
(¹⁸O, ¹⁶O)⁴²Ca @ 270 MeV

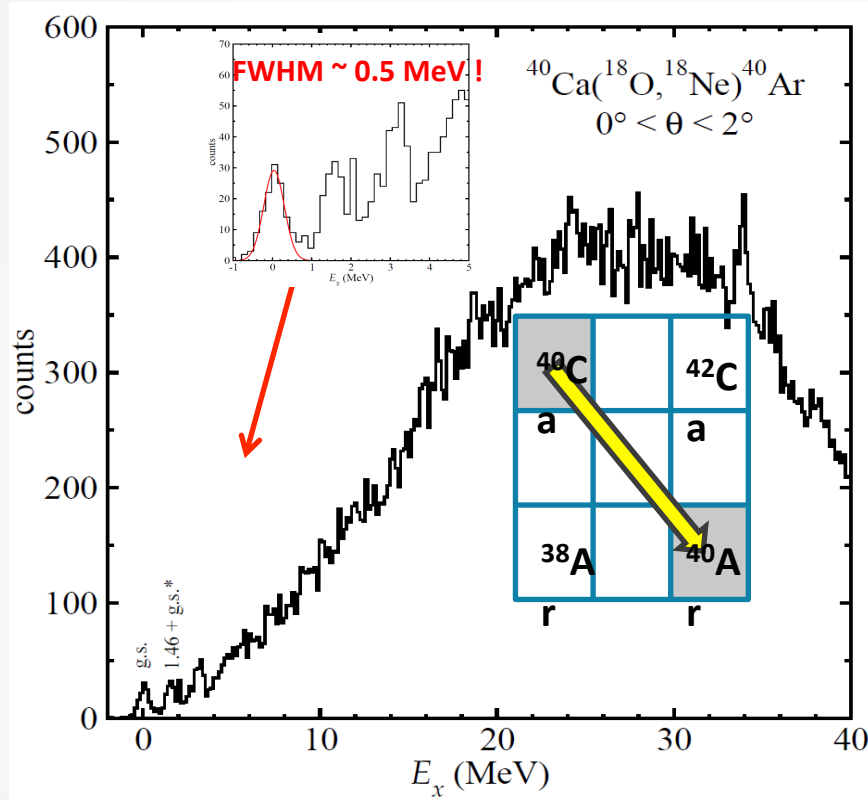


x-section (2M) ≈ 0.5
Extracted B(GT) from Y. F.



Measured
Not measured

Measured energy spectrum of ^{40}Ar at **very forward angles** with an energy resolution of FWHM ~ 0.5 MeV.



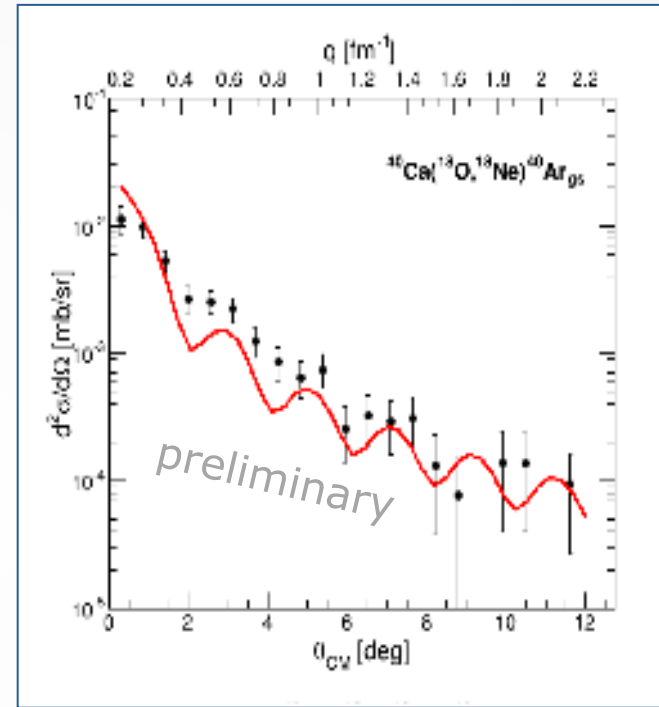
^{40}Ar 0^+ ground state is well separated from both the first excited ^{40}Ar 2^+ at 1.46 MeV and the ^{18}Ne excited state at 1.887 MeV

$$d\sigma^{\text{DCE}}/d\Omega = 11 \mu\text{b}/\text{sr} \quad \text{at } \theta_{\text{cm}} = 0^\circ$$

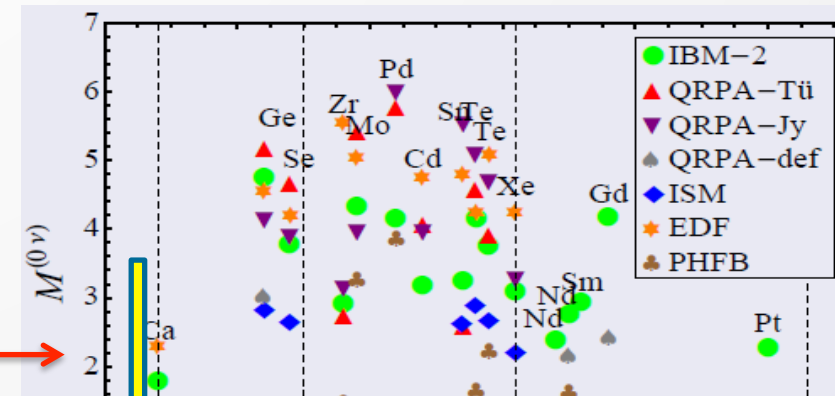
speculate: a comparison between ^{48}Ca and ^{40}Ca

$$|M^{0\nu\beta\beta}(^{40}\text{Ca})|^2 = 0.37 \pm 0.18$$

Differential cross-section of the transition $^{40}\text{Ca}_{\text{g.s.}}(^{18}\text{O}, ^{18}\text{Ne})^{40}\text{Ar}_{\text{g.s.}}$ @ 270 MeV



The position of the minima is well described by a Bessel function oscillation pattern is not expected in complex multistep transfer reactions





Experimental limits moving towards “hot-cases” :

Determination of nuclear matrix elements seems to be at
our reach... but :

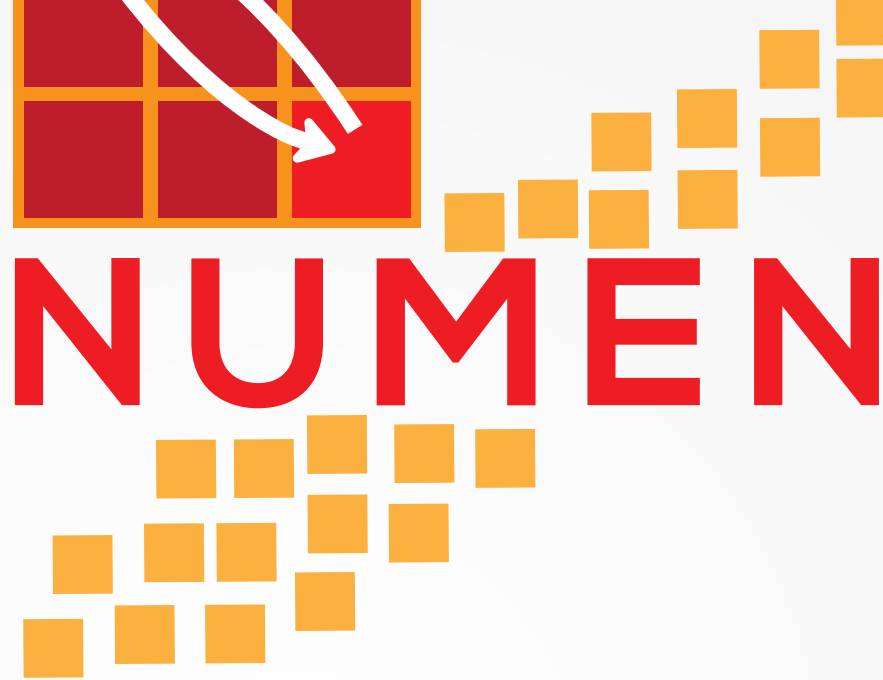
The ($^{18}\text{O}, ^{18}\text{Ne}$) reaction is particularly **advantageous**, but it is of $\beta^+\beta^+$ kind

None of the reactions of $\beta^-\beta^-$ kind looks like as favourable as the ($^{18}\text{O}, ^{18}\text{Ne}$)
($^{18}\text{Ne}, ^{18}\text{O}$) requires a radioactive beam
($^{20}\text{Ne}, ^{20}\text{O}$) or ($^{12}\text{C}, ^{12}\text{Be}$) have smaller $B(\text{GT})$

In some cases **gas target** will be necessary, e.g. ^{136}Xe or ^{130}Xe

In some cases the **energy resolution** is not enough to separate the g.s. from
the excited states in the final nucleus → Coincident **detection of γ -rays**

Much higher beam current is needed !



Present technology is not enough...

The challenge:

to measure at very high rates of heavy ions!

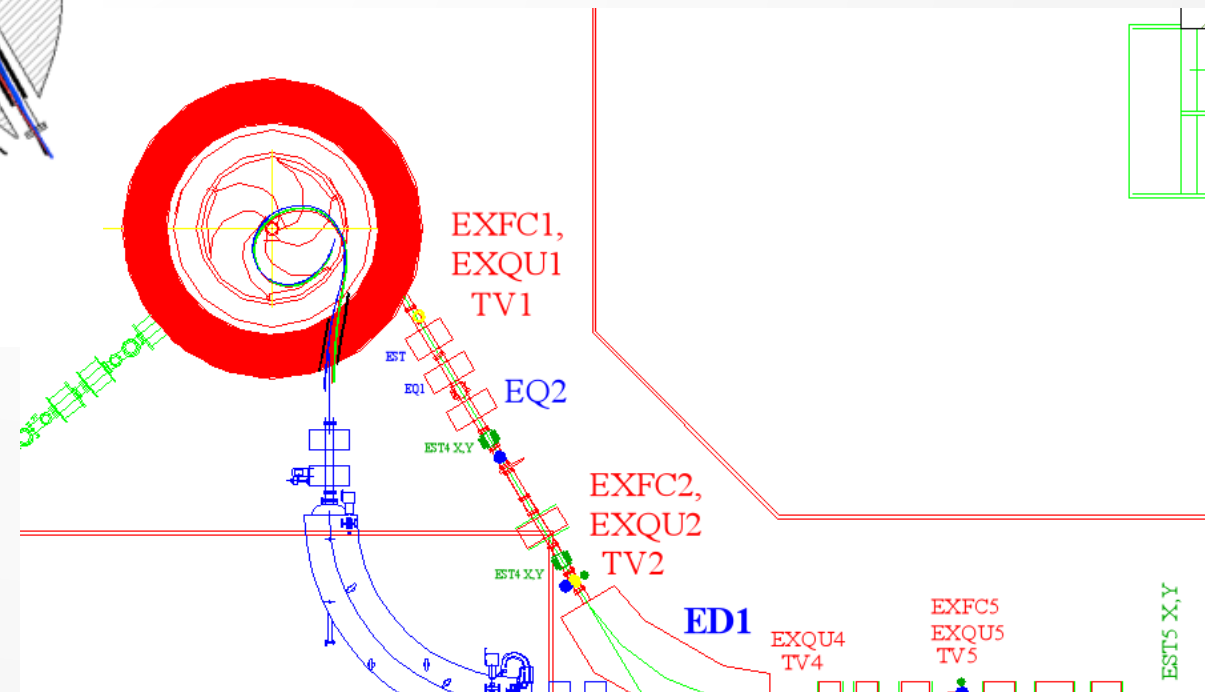
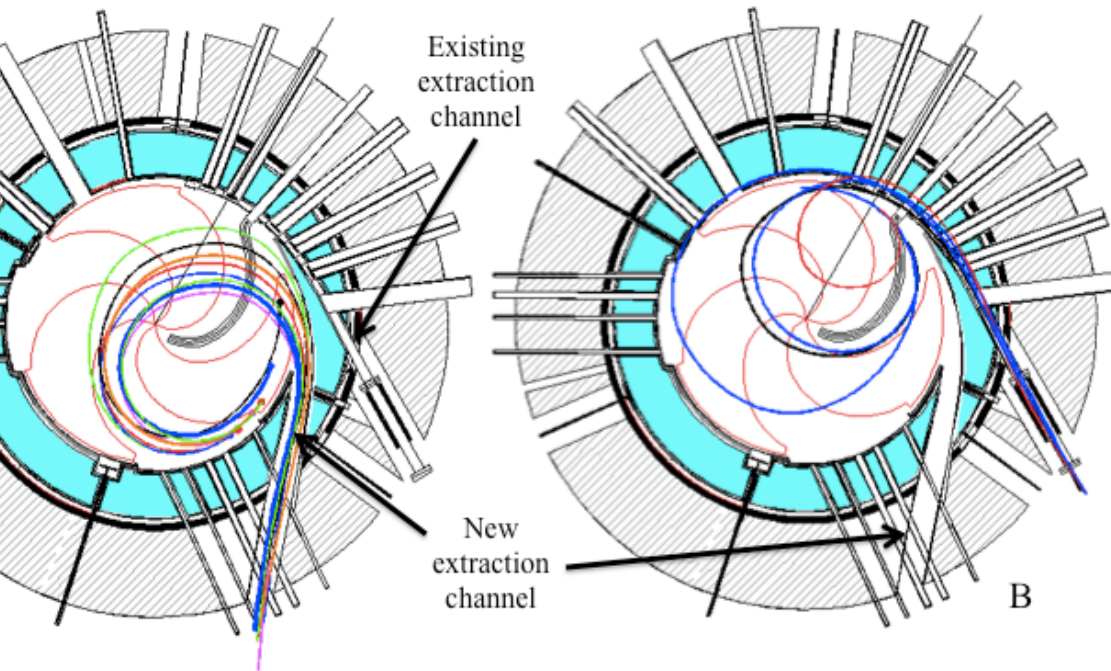
Upgraded set-up to work with two orders of magnitude more current than the present

Substantial change in the technologies used in **CS** and in the **MAGNEX** detector

Major upgrade of LNS facilities: The CS accelerator

The **CS** accelerator current (from 100 W to 5-10 kW);

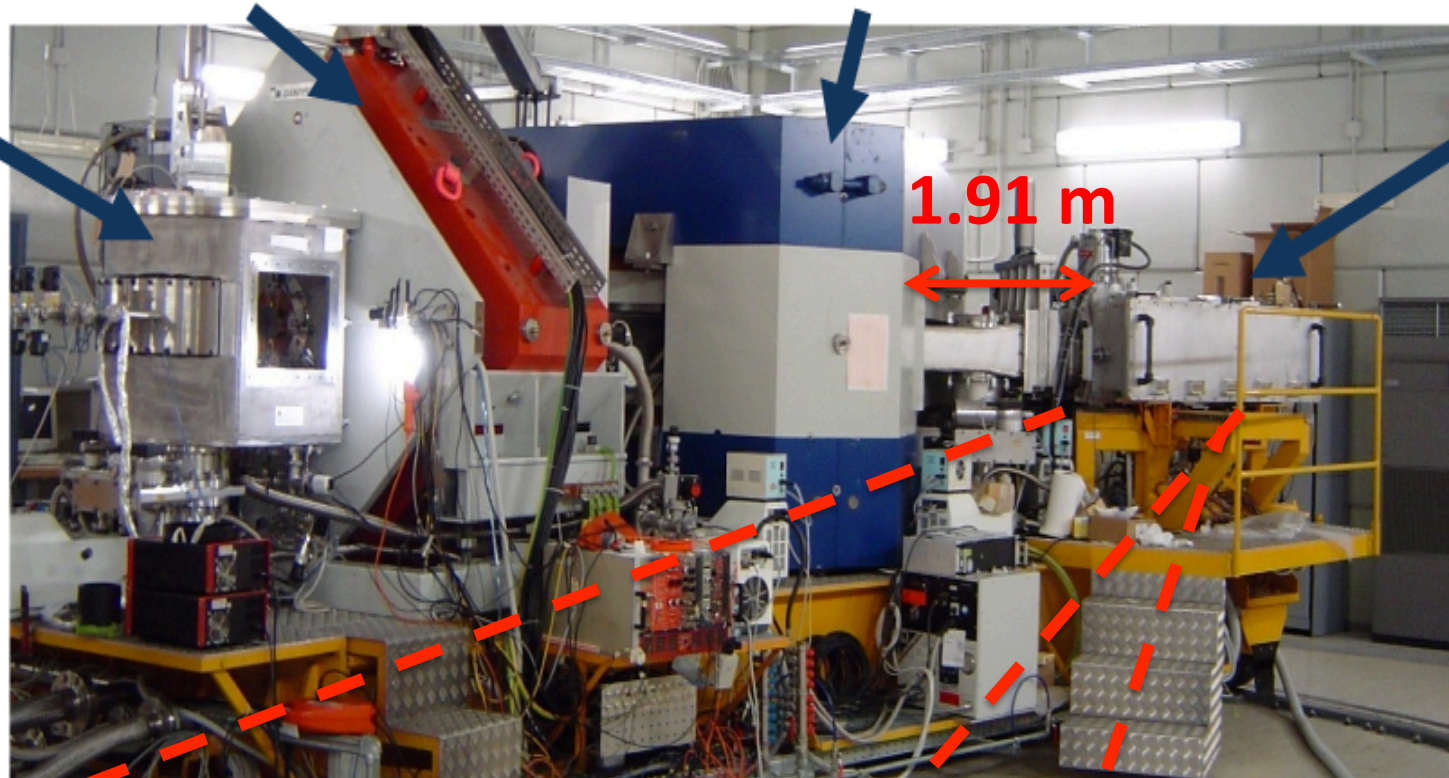
From electrostatic extraction to extraction by stripping



The **beam transport line**
transmission efficiency to
nearly 100%

Quadrupole

Dipole



1.91 m

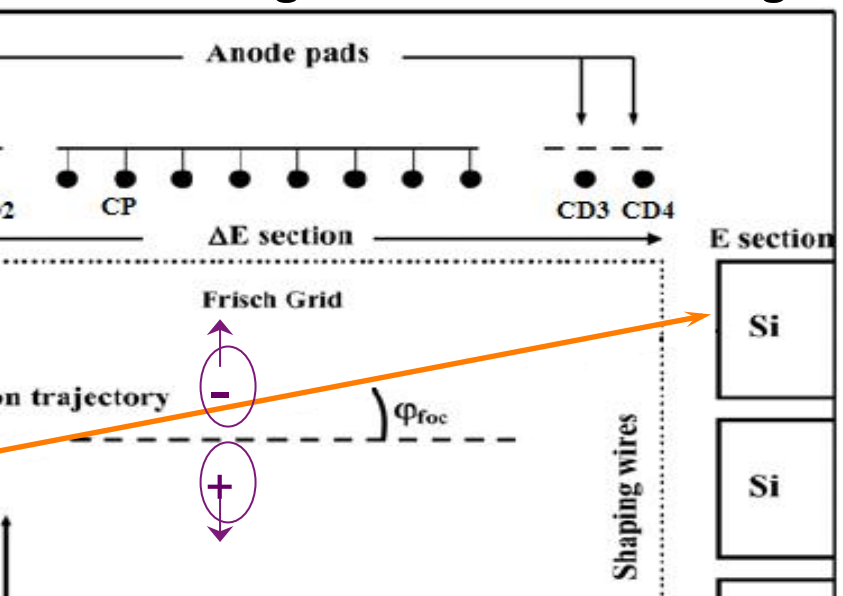
Focal Plane Detector

active area
140x20 cm²

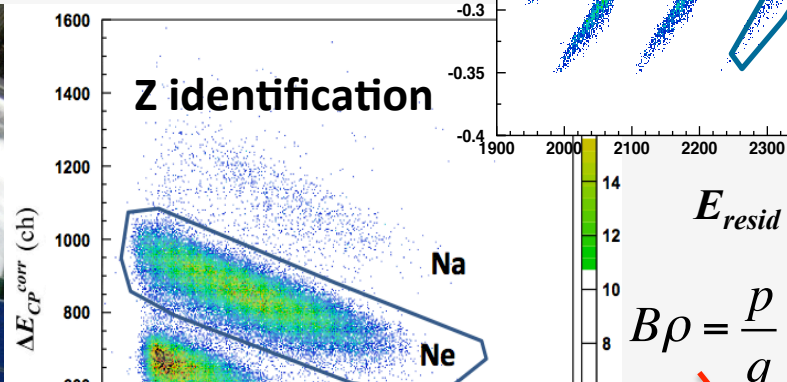
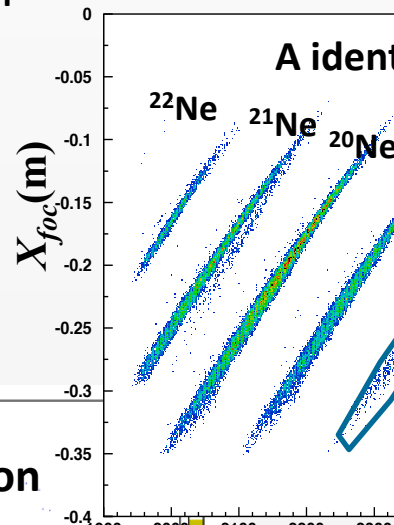
a **ionization drift chamber**, five independent proportional counter tubes, of which are position-sensitive and composed of **stopping silicon detectors**.

Pure isobutane pressure range: 5-100mbar; 600-800 Volt, wires 20 micron

Multiwire gas tracker and ΔE stage



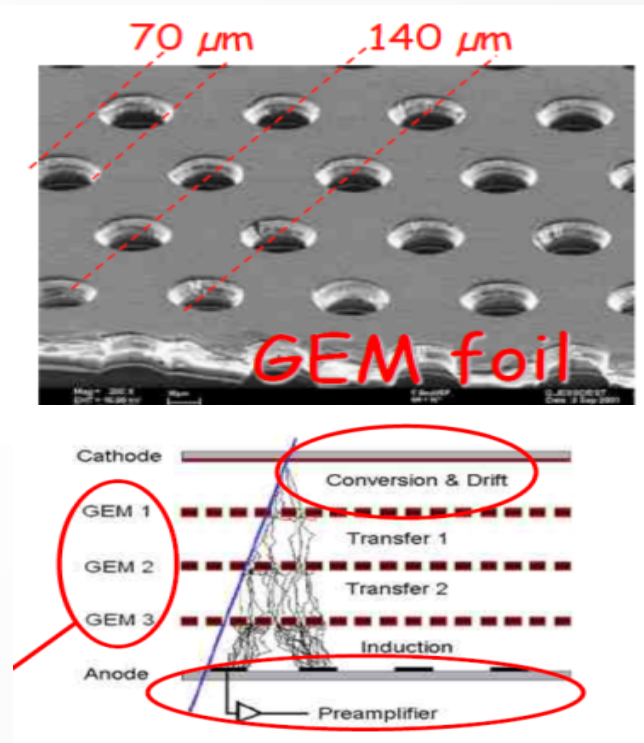
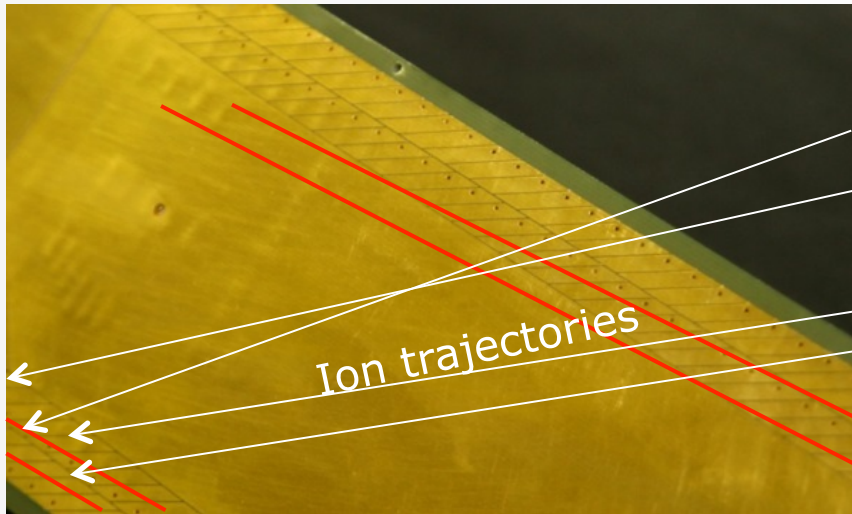
Wall of 60 stopping 7 X 5 cm² Silicon detectors surface covered 100 X 21 cm²



from ≈ 2 kHz \longrightarrow to several MHz

1. A new gas tracker

limited to ~ 5 kHz due to the slow drift of positive ions from the multiplication wires to the Frish grid



From multi-wire tracker

To micro-pattern electron multipliers tracker

➤ R&D key issue : GEM-based tracker at **low pressure and wide dynamic range**

from ≈ 2 kHz \longrightarrow to several MHz

A new particle identification wall

From

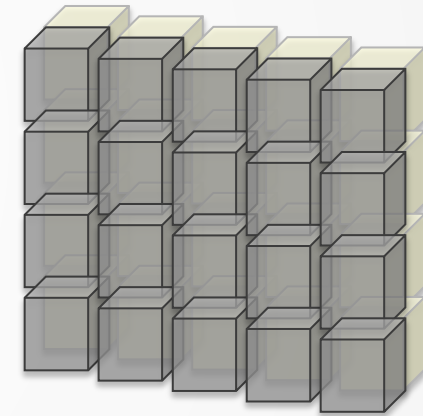
of 7×5 cm² silicon pad detectors



Double-hit probability at 100 kHz $> 30\%$
High cost for breaking with localized exposure
Pads tend to deteriorate at about 10^8 - 10^9 ions/cm²

To

A wall of 1×1 cm² telescopes of SiC detectors



- PID function decoupled from gas tracker
- Double-hit probability at 100 kHz $< 1\%$
- Much smaller maintenance cost
- First irradiation tests successfully performed in collaboration with CNR
- Radiation hard material (10^{13} ions/cm²)

Thickness of the telescope must be chosen in order to permit the detection of the ejecta over a wide dynamical range of incident energies (10 to 60 MeV/A)

$(^{18}\text{O}, ^{18}\text{Ne})$ DCE reaction:

In order to have a good signal in **the first stage** of the telescope :

$$\Delta E_1 \approx 100 \mu\text{m} \text{ for } 10 \text{ to } 30 \text{ MeV/A}$$

$$\Delta E_2 \approx 1000 \mu\text{m} \text{ for } 30 \text{ to } 60 \text{ MeV/A}$$

For the **second stage**

$$E \approx 1000 \mu\text{m} \text{ for } 10 \text{ to } 60 \text{ MeV/A}$$

Good solution would be **a triple stage telescope** to cover the whole energy range for the EN experimental campaign guarantying an unambiguous identification of the ion mass, atomic number and charge state.



A technological challenge !

To maintain the actual performances at high rates:

<i>NUMEN requirements</i>	
1 cm ² three-stage telescope	
thickness of ΔE_1 stage 100 μm	
thickness of ΔE_2 stage 1000 μm	
thickness of E stage 1000 μm	
high geometrical efficiency ($\sim 80\%$)	
good energy resolution (1%)	
good time resolution (1-2 ns)	
hard to radiation damage ($>10^{13}$ ions/cm ²)	
capacitance ~ 100 pF	
implementation of the pulse-shape analysis	

Thickness of the detector must be chosen in order to permit the detection of the ejectiles in a wide dynamical range of energies (10 to 60 MeV)

a high degree of segmentation is required in order to limit the probability of multiple hit events probability below 10% in the FPD.

To obtain accurate measurements of absolute cross sections

Diamond	GaN	4H SiC	Si
5.5	3.39	3.26	1.12
10^7	$4 \cdot 10^6$	$2.2 \cdot 10^6$	$3 \cdot 10^7$
1800	1000	800	1450
1200	30	115	450
$2.2 \cdot 10^7$		$2 \cdot 10^7$	$0.8 \cdot 10^7$

1

substitution of the present **Focal Plane Detector (FPD) gas tracker** with **a micro-patterned electron multipliers foils** for a tracker system;

2

substitution of the **wall of silicon pad** stopping detectors with a wall of telescopes based on **SiC-SiC detectors** ;

3

introduction of an array of detectors for measuring the **coincident γ -rays**;

4

enhancement of the maximum **magnetic rigidity**

The NUMEN goals

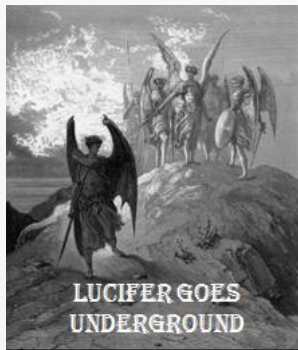
- The NUMEN Holy Graal: the unit cross section
- A new generation of DCE constrained $0\nu\beta\beta$ NME theoretical calculation
- The ratio of measured cross sections can give a model independent way to compare the sensitivity of different half-life experiments.



^{76}Ge



^{130}Te



^{82}Se



^{136}Xe



^{116}Cd

strong impact in future development of the field.



“...non è facile fare fisica molto innovativa se questa non è sorretta da tecnologie altrettanto innovative”