

A self-consistent equation of motion multiphonon method for even and odd mass nuclei

G. De Gregorio

•Università di Napoli Federico II

•INFN



In Collaboration with

N. Lo Iudice, F. Andreozzi, A. Porrino

Univ. Federico II & INFN Napoli

F. Knapp

Charles University (Prague)

P. Veselý

NPI Řež (Prague)

Semiclassical

$$\{\alpha_\lambda, \pi_\lambda\} \rightarrow \{O_\lambda, O_\lambda^\dagger\}$$

EofM

$$[H, O_\lambda^\dagger] = \hbar\omega_\lambda O_\lambda^\dagger$$

**Collective
modes**

Microscopic

TDA mapping

$$O_\lambda^\dagger = \sum_{ph} c_{ph}(\lambda) a_p^\dagger a_h$$

EoM

$$[H, O_\lambda^\dagger] |> = \hbar\omega_\lambda O_\lambda^\dagger |>$$

RPA mapping

$$O_\lambda^\dagger = \sum_{ph} [X_{ph}(\lambda) a_p^\dagger a_h - Y_{ph}(\lambda) a_h^\dagger a_p]$$

EoM

$$[H, O_\lambda^\dagger] |0> = \hbar\omega_\lambda O_\lambda^\dagger |0>$$

$|0> \equiv$ correlated g.s

Beyond mean field: *Adopted Methods*

Non relativistic

- qp-Phonon

P. F. Bortignon et al. (Milano group)

- 2nd RPA

R. Roth et al.

D. Gambacurta et al.

Phenomenological

- QPM (*Soloviev School (Dubna)*)

Relativistic: RTBA

E. Litvinova, P. Ring, D. Vretenar.....

Our proposal: EMPM

*D. Bianco, F. Knapp, N. Lo Iudice, F. Andreozzi, A. Porrino, Phys. Rev. C **85**, 014313 (2012).*

Equation of motion phonon model (EMPM)

Eigenvalue problem

$$\mathbf{H} |\Psi_v\rangle = E_v |\Psi_v\rangle$$

$$|\Psi_v\rangle \in \mathcal{H} = \sum_n \oplus \mathcal{H}_n \quad \mathcal{H}_n \in |n; \beta\rangle \equiv \text{n-phonon basis states}$$

An obvious, but unmanageable, multiphonon basis

$$|\lambda_1, \dots, \lambda_i, \dots, \lambda_n\rangle = \mathbf{O}_{\lambda_1}^\dagger \dots \mathbf{O}_{\lambda_i}^\dagger \dots \mathbf{O}_{\lambda_n}^\dagger |0\rangle$$

$$\mathbf{O}_\lambda^\dagger = \sum_{ph} c_{ph}^\lambda \mathbf{a}_p^\dagger \mathbf{a}_h$$

A viable route

$$|\alpha_n\rangle = \sum_{\lambda\alpha} C_{\lambda\alpha}^\beta \mathbf{O}_\lambda^\dagger |\alpha_{n-1}\rangle$$

Construction of $|\mathbf{n}; \beta\rangle$: EoM

Assuming $|\alpha_{n-1}\rangle$ known, we solve the **Eq. of Motion**

$$\langle \alpha_n | [\mathbf{H}, \mathbf{O}^\dagger_\lambda] | \alpha_{n-1} \rangle = (\mathbf{E}_\beta^{(n)} - \mathbf{E}_\alpha^{(n-1)}) \langle \alpha_n | \mathbf{O}^\dagger_\lambda | \alpha_{n-1} \rangle$$


$$\mathcal{A} \mathcal{X} = \mathcal{E} \mathcal{X}$$


$$\mathcal{X}$$

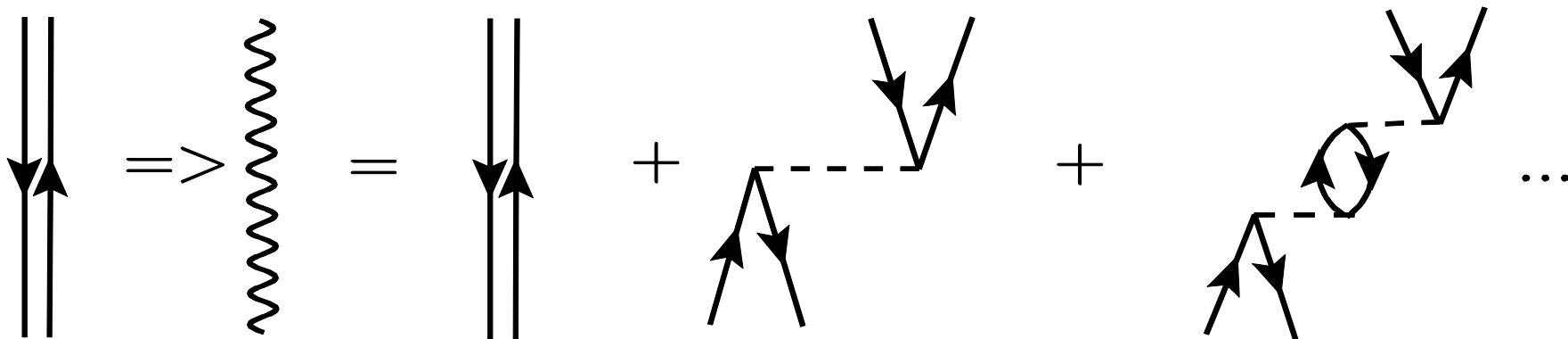
Where

$$\mathcal{A}_{(\mu\alpha)(\nu\gamma)} = (\mathbf{E}_\mu + \mathbf{E}_\alpha) \delta_{\alpha\gamma} \delta_{\mu\nu} + \mathcal{V}_{(\mu\alpha)(\nu\gamma)}$$

TDA matrix

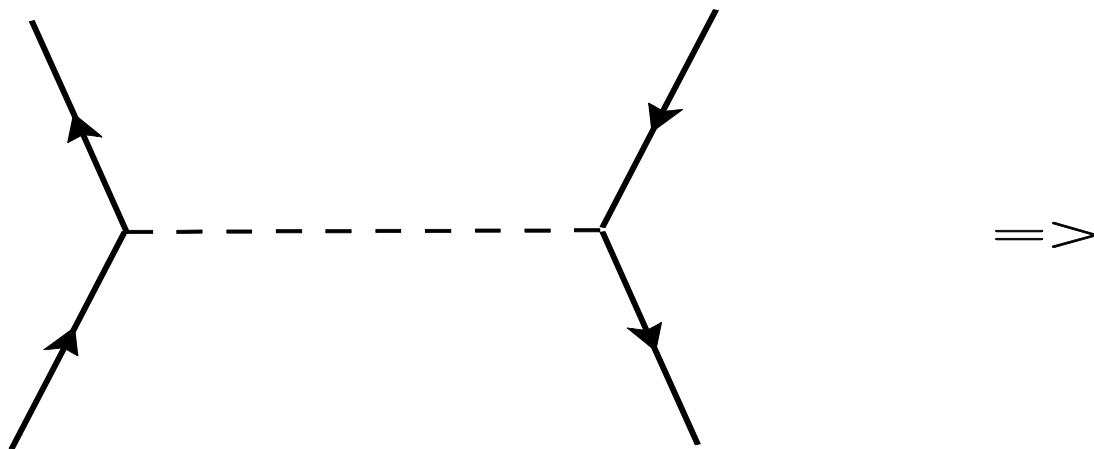
$$\mathbf{A}_{(ph)(p'h')} = (\epsilon_p - \epsilon_h) \delta_{pp'} \delta_{hh'} + \mathbf{V}_{ph'hp'}$$

From **p-h** to **TDA** to **EMPM**

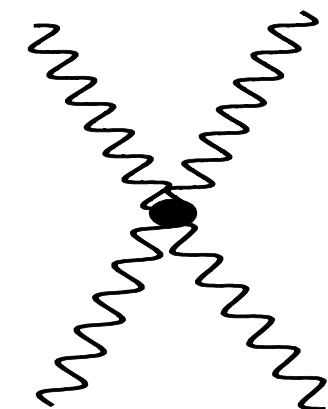


$$\varepsilon_p - \varepsilon_h \Rightarrow \mathbf{E}_\mu$$

$$|\mathbf{ph}\rangle \Rightarrow |\boldsymbol{\mu}\rangle = \sum_{\mathbf{ph}} \mathbf{c}_{\mathbf{ph}}^\mu \mathbf{a}_p^\dagger \mathbf{a}_h | \rangle$$



$$\mathbf{V}_{\mathbf{ph}'\mathbf{hp}'}$$



$$\mathcal{V}_{(\mu\nu)(\mu'\nu')}$$

Construction of $|\mathbf{n}; \beta\rangle$: EoM

Problem

$$\mathcal{A} \mathcal{X} = \mathcal{E} \mathcal{X}$$

is **not** a true Eigenvalue Eq.!

$\{\mathbf{O}^\dagger_\lambda |\mathbf{n}-1, \alpha\rangle\}$ form a **non-orthogonal redundant** basis



$\mathcal{X} = \langle \mathbf{n}, \beta | \mathbf{O}^\dagger_\lambda |\mathbf{n}-1, \alpha\rangle$ is **not** a **true** expansion coefficient

Recipe for solving the problem

1° step

$$|\alpha_n\rangle = \sum_{\lambda\alpha} C_{\lambda\alpha}^\beta \mathbf{O}_{\lambda}^\dagger |\alpha_{n-1}\rangle \}$$

$$|\lambda\alpha\rangle = \mathbf{O}_{\lambda}^\dagger |\alpha_n\rangle$$

$$(X = \langle \alpha_n | \lambda \alpha_{n-1} \rangle)$$

$$\mathcal{D} \equiv \{ \langle \lambda' \alpha_n' | \lambda \alpha_n \rangle \}$$

$$\mathcal{A} X = \mathcal{E} X$$

$$X = \mathcal{D} C$$

$$[\mathcal{H} - \mathcal{E}\mathcal{D}] C = \mathbf{0}$$

where

$$\mathcal{H} = \mathcal{A}\mathcal{D}$$

But \mathcal{D} is **singular** !

2° conclusive step: Choleski

$$[\mathcal{H} - \mathcal{E}\mathcal{D}] \mathbf{C} = \mathbf{0}$$

Cholesky



$$\mathcal{D} \rightarrow \mathbf{D}$$

$$[\mathbf{H} - \mathbf{E}] \mathbf{C} = \mathbf{0}$$



$$|n, \beta\rangle = \sum_{\lambda\alpha} C_{\lambda\alpha}^{\beta} \mathbf{O}_{\lambda}^{\dagger} |\alpha\rangle$$

$$\langle n', \alpha | n; \beta \rangle = \delta_{nn'} \delta_{\alpha\beta}$$

$$\mathcal{H} = \mathcal{A}\mathcal{D}$$

$$\mathcal{D} \equiv \{ \langle \lambda' \alpha' | \lambda \alpha \rangle \}$$

$$\mathbf{H} = \mathbf{D}^{-1} \mathcal{A}\mathcal{D}$$

Iterative Generation of n-phonon states

$$\mathbf{A} \mathbf{c} = \hbar \boldsymbol{\omega} \mathbf{c}$$



$$|\lambda\rangle = \sum_{\text{ph}} \mathbf{c}_{\text{ph}}^{\lambda} \mathbf{a}_{\text{p}}^{\dagger} \mathbf{a}_{\text{h}} | \rangle$$



$$\mathbf{H} \mathbf{C} = \mathbf{E} \mathbf{C}$$



$$|\alpha_n\rangle = \sum_{\lambda\alpha(n)} \mathbf{C}_{\lambda\alpha(n)}^{\alpha(n)} \mathbf{O}_{\lambda}^{\dagger} | \alpha_{n-1} \rangle$$

- **No approximations except for truncation!**
- **Pauli** principle fully accounted for:
- **No redundant** states!
- $|\alpha_n\rangle$ form an **orthonormal** basis

$$n = 2, 3, \dots$$

Eigenvalue problem in Multiphonon basis $\{|\alpha_n\rangle\} \equiv \{|\alpha_0\rangle, |\alpha_1\rangle, \dots\}$

$$\sum_{n', \beta(n')} [(\mathbf{E}_{\alpha(n)} - \epsilon_v) \delta_{nn'} \delta_{\alpha(n) \beta(n')} + \mathbf{V}_{\alpha(n) \beta(n')}] \mathbf{C}_{\beta(n')}^v = \mathbf{0}$$



$$|\Psi_v\rangle = \sum_{\alpha(n)} \mathbf{C}_{\alpha(n)}^v |\alpha_n\rangle$$

where

$$|\alpha_n\rangle = \sum_{\lambda \alpha(n-1)} \mathbf{C}_{\lambda \alpha(n-1)}^\beta \mathbf{O}_{\lambda}^\dagger |\alpha_{n-1}\rangle$$

EMPM in q-p scheme

One proceeds as in p-h scheme

$$\langle \alpha_n | [\mathbf{H}, \mathbf{O}_\lambda^\dagger] | \alpha_{n-1} \rangle = (E_\beta^{(n)} - E_\alpha^{(n-1)}) \langle \alpha_n | \mathbf{O}_\lambda^\dagger | \alpha_{n-1} \rangle$$

namely with the replacement

$$\mathbf{O}_\lambda^\dagger = \sum_{ph} c_{ph}(\lambda) \mathbf{a}_p^\dagger \mathbf{a}_h \quad \Rightarrow \quad \mathbf{O}_\lambda^\dagger = \sum_{r \leq s} c_{rs}(\lambda) \alpha_r^\dagger \alpha_s$$

Implementation

1° STEP : Intrinsic Hamiltonian

$$H = T_{\text{int}} + V_{\text{NN}}$$

where

$$T_{\text{int}} = \frac{1}{2m} \sum_i p_i^2 - \frac{P^2}{2M_{\text{cm}}} \quad V_{\text{NN}} = V_{\chi} = \text{NNLO}_{\text{opt}}$$

2° STEP : HF(B) Self-consistent basis

3° STEP : Construction of TDA phonons

(free of spurious admixtures induced by CM and particle number violation)

$$O_{\lambda}^{\dagger} = \sum_{\text{ph}} c_{\text{ph}}(\lambda) a_p^{\dagger} a_h$$

$$O_{\lambda}^{\dagger} = \sum_{\underline{r} \leq \underline{s}} c_{rs}(\lambda) \alpha_r^{\dagger} \alpha_s^{\dagger}$$

Implementation

4° STEP : Generation of Multiphonon basis

$$\{|\alpha_n\rangle\} \equiv \{|\alpha_0\rangle, |\alpha_1\rangle, |\alpha_2\rangle, \dots\}$$

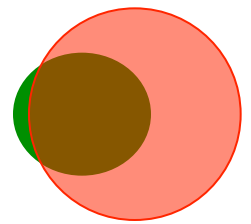
5° STEP : Solution of the eigenvalue problem in Multiphonon basis

$$|\Psi_\nu\rangle = \sum_{n\alpha} C_\alpha^{(\nu)} |n; \alpha\rangle$$

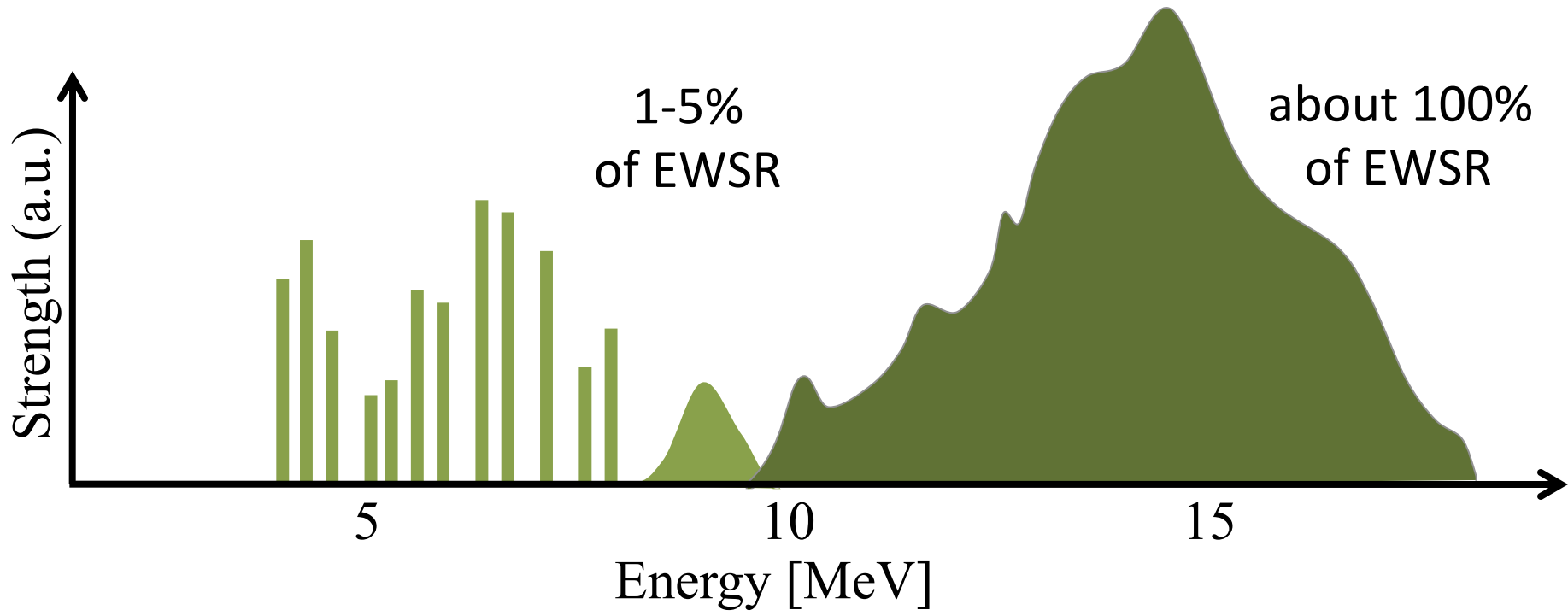
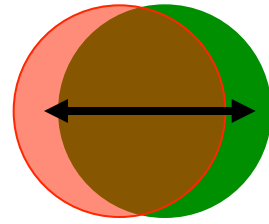
6° STEP : Transition Probabilities and cross section

Applications: Selfconsistent study of Dipole response

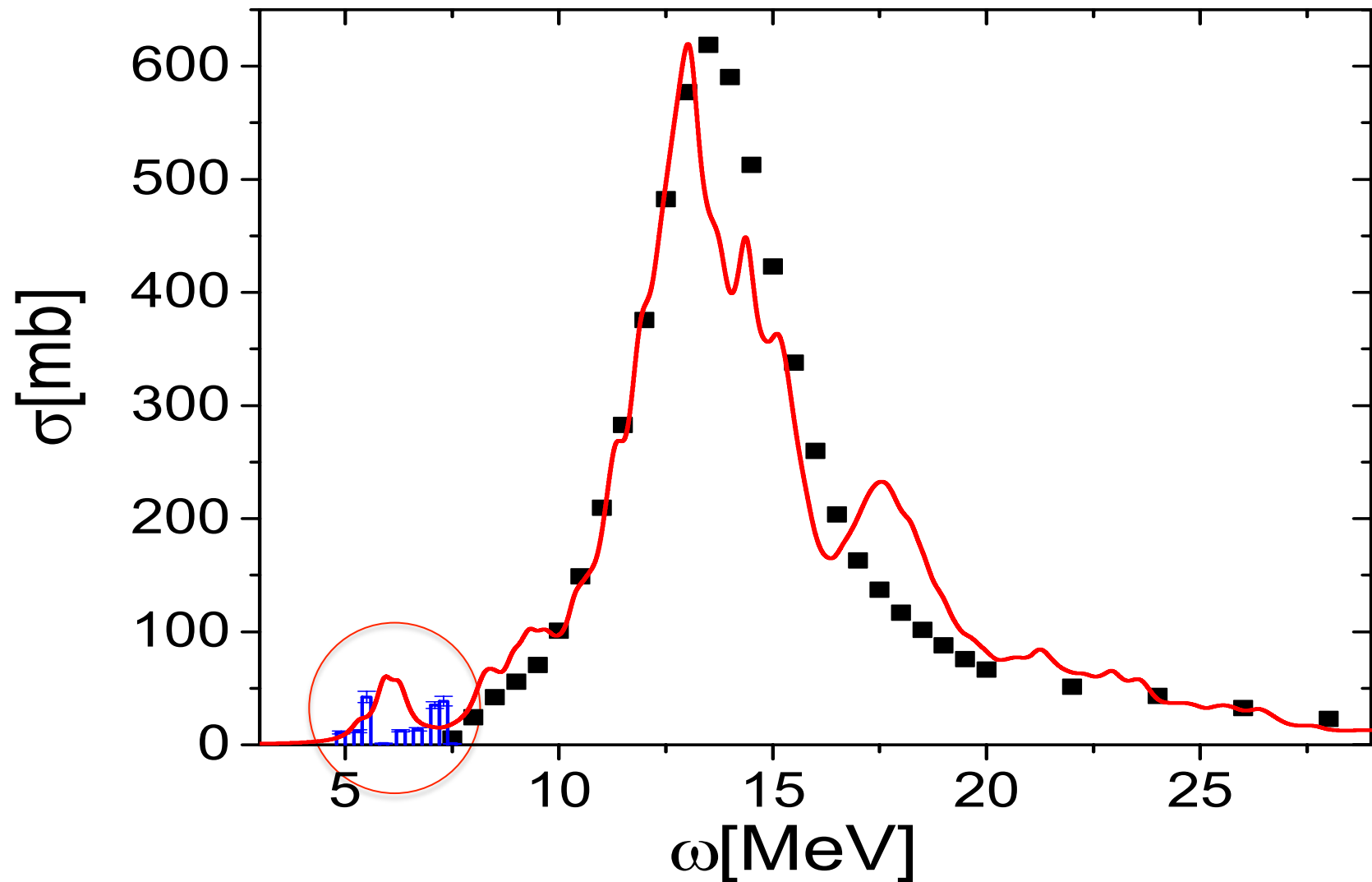
PDR



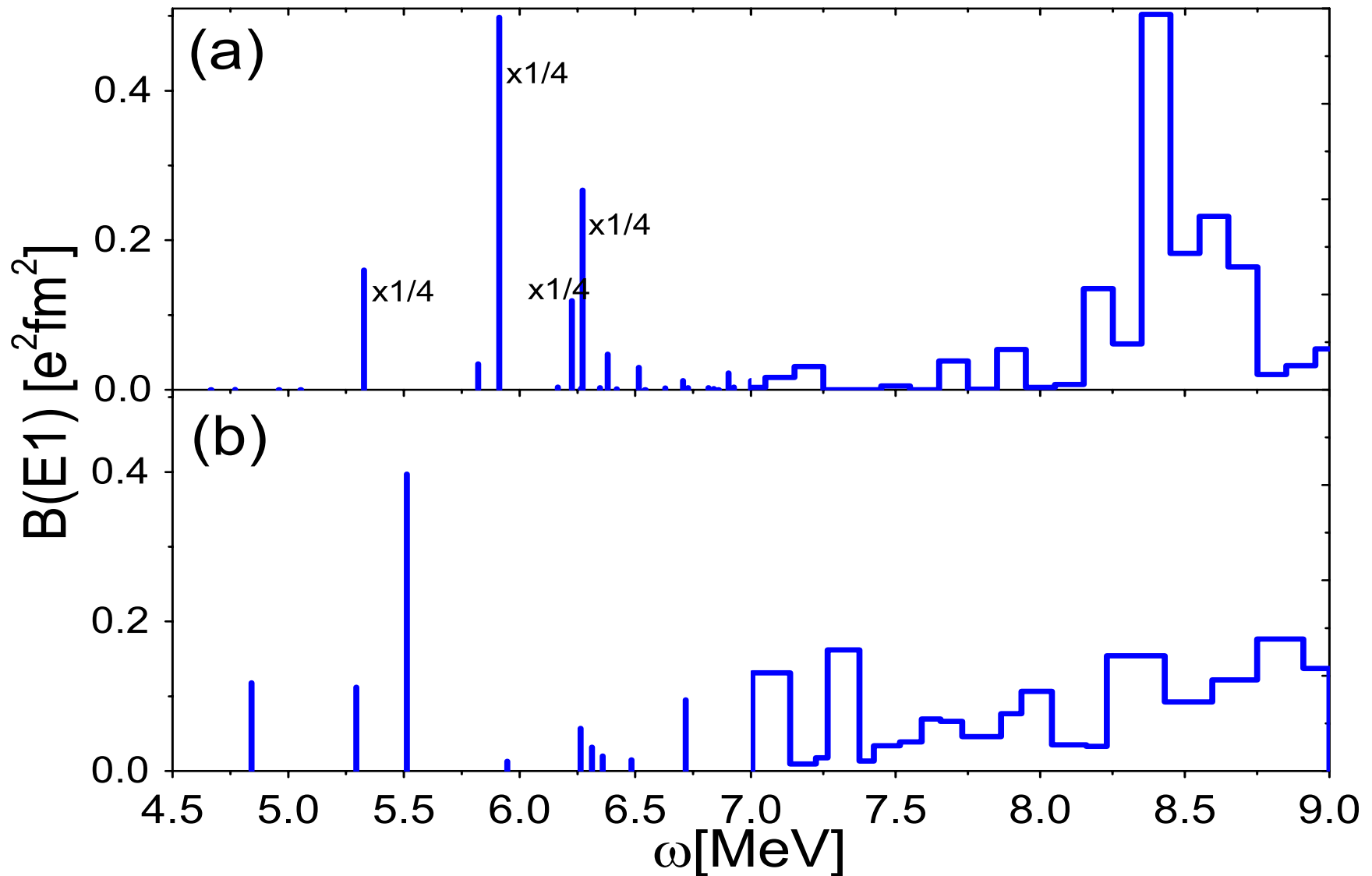
GDR



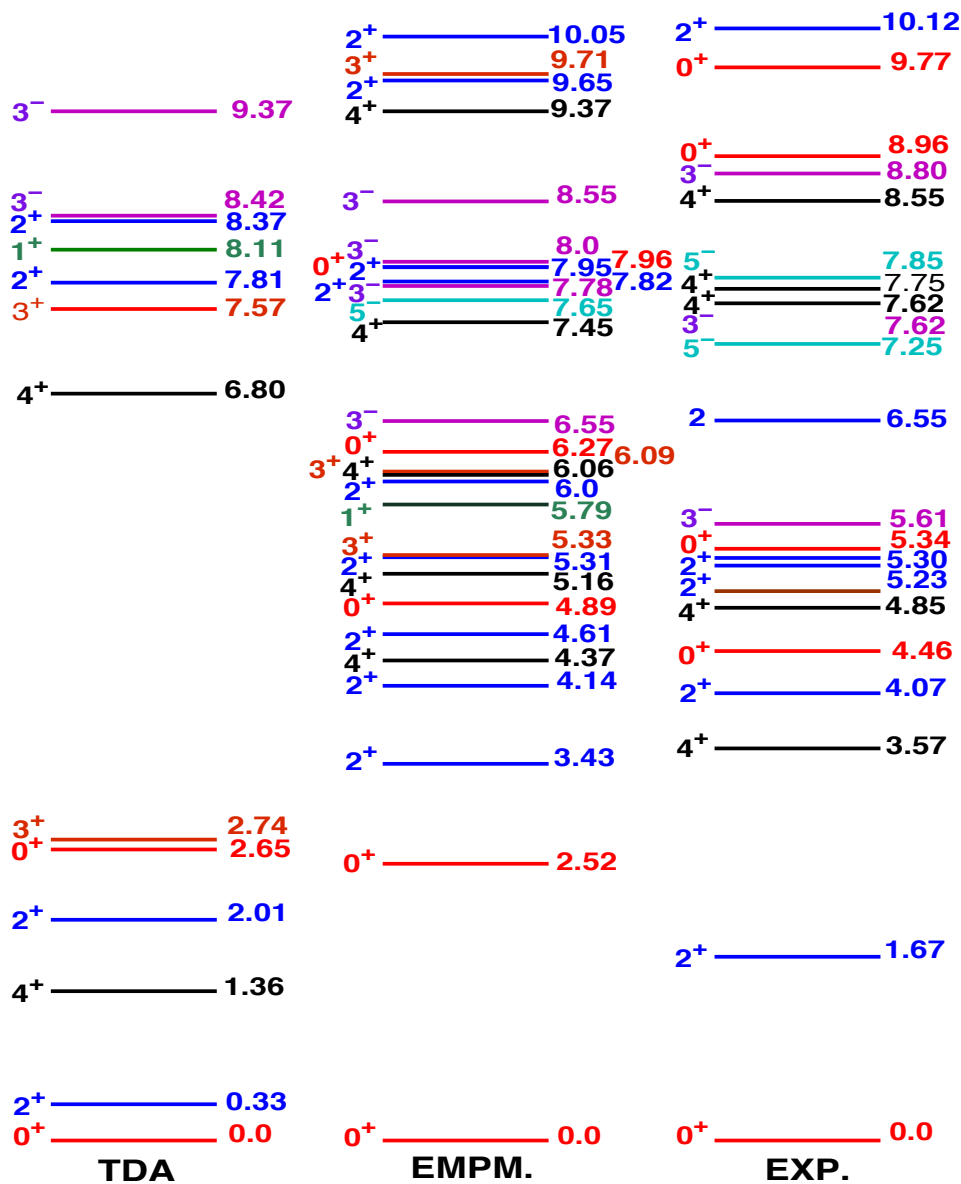
EMPM in p-h scheme: Application to neutron rich ^{208}Pb



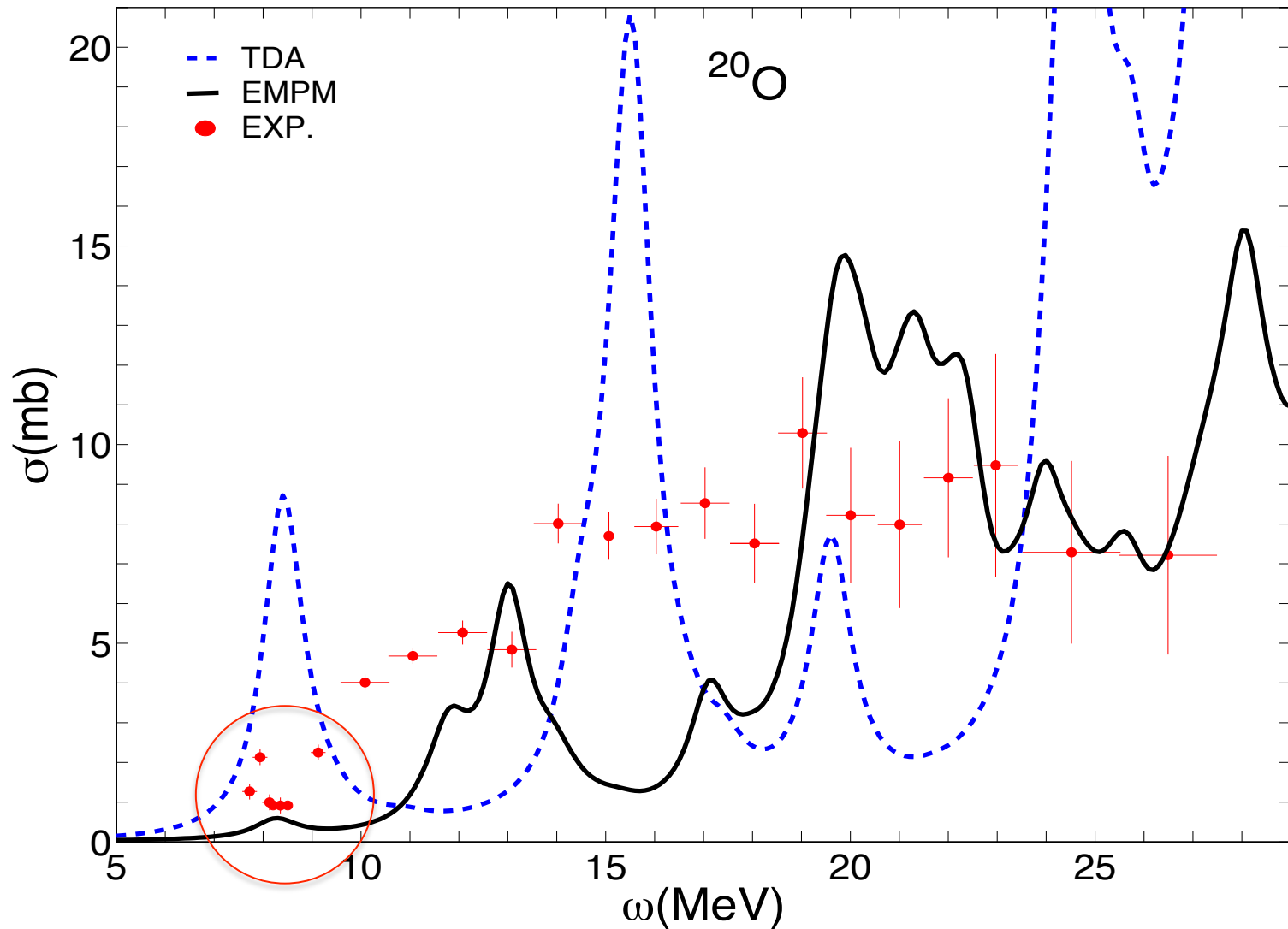
Low-lying Dipole strength in ^{208}Pb



EMPM in q-p scheme: Application to neutron rich ^{20}O



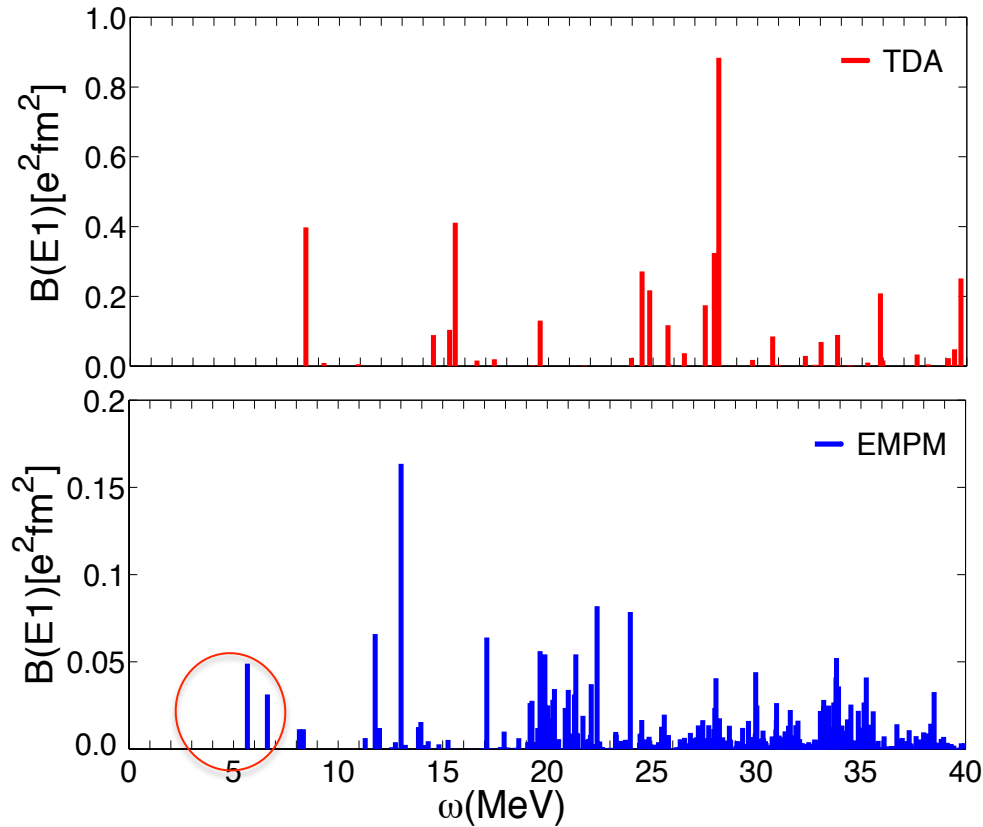
Cross section in ^{20}O



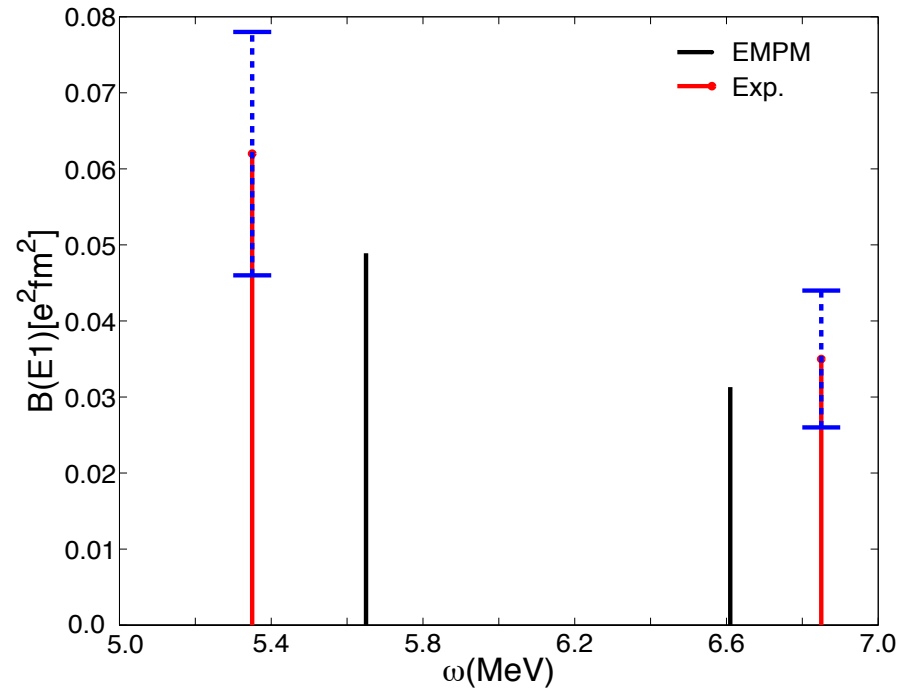
G. De Gregorio, F. Knapp, N. Lo Iudice, P. Vesely
Phys. Rev. C 044314 (2016)

E. Tryggestad, et al., Phys. Rev. C
67, 064309 (2003).

Dipole strength in ^{20}O



E. Tryggestad, et al., Phys Rev. C 67, 064309 (2003).



G. De Gregorio, F. Knapp, N. Lo Iudice, P. Vesely
Phys. Rev. C 044314 (2016)

Even Nuclei: Summary

EMPM

- Extends Mean Field without **approximations**, except for truncation, for any potential
- Includes **multiphonon states explicitly**

Results

2-phonon configurations are **needed** for

- describing the **fine** structure of the GDR and PDR
 - obtaining a **complete** low-energy spectrum

EMPM: Odd nuclei

$$|\lambda \alpha\rangle = O^\dagger_\lambda |\alpha\rangle$$



$$|p \alpha\rangle = a^\dagger_p |\alpha\rangle$$

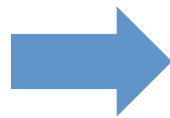
$$\langle \alpha_n | [H, O^\dagger_\lambda] | \alpha_{n-1} \rangle$$



$$\langle v_n | [H, a^\dagger_p] | \alpha_n \rangle$$

$$X = DC$$

$$AX = EX$$



$$[H - ED]C = 0$$

$$D \equiv \{ \langle p \alpha | p' \alpha' \rangle \}$$

Cholesky

$$D \rightarrow D$$

EMPM: Odd nuclei

$$[\mathbf{H} - \mathbf{E}] \mathbf{C} = \mathbf{0}$$

$$\mathbf{H} = \mathbf{D}^{-1} \mathcal{A} \mathcal{D}$$

$$|\mathbf{v}_n\rangle = \sum_{p\alpha} C_{p\alpha}^{\mathbf{v}} a_p^\dagger |\alpha\rangle$$

$$\langle \mathbf{v}' | \mathbf{v} \rangle = \delta_{\mathbf{v}\mathbf{v}'}$$

$$|\Psi_\mu\rangle = \sum_i C_{\mathbf{v}_i}^{(\mu)} |\mathbf{v}_i\rangle$$

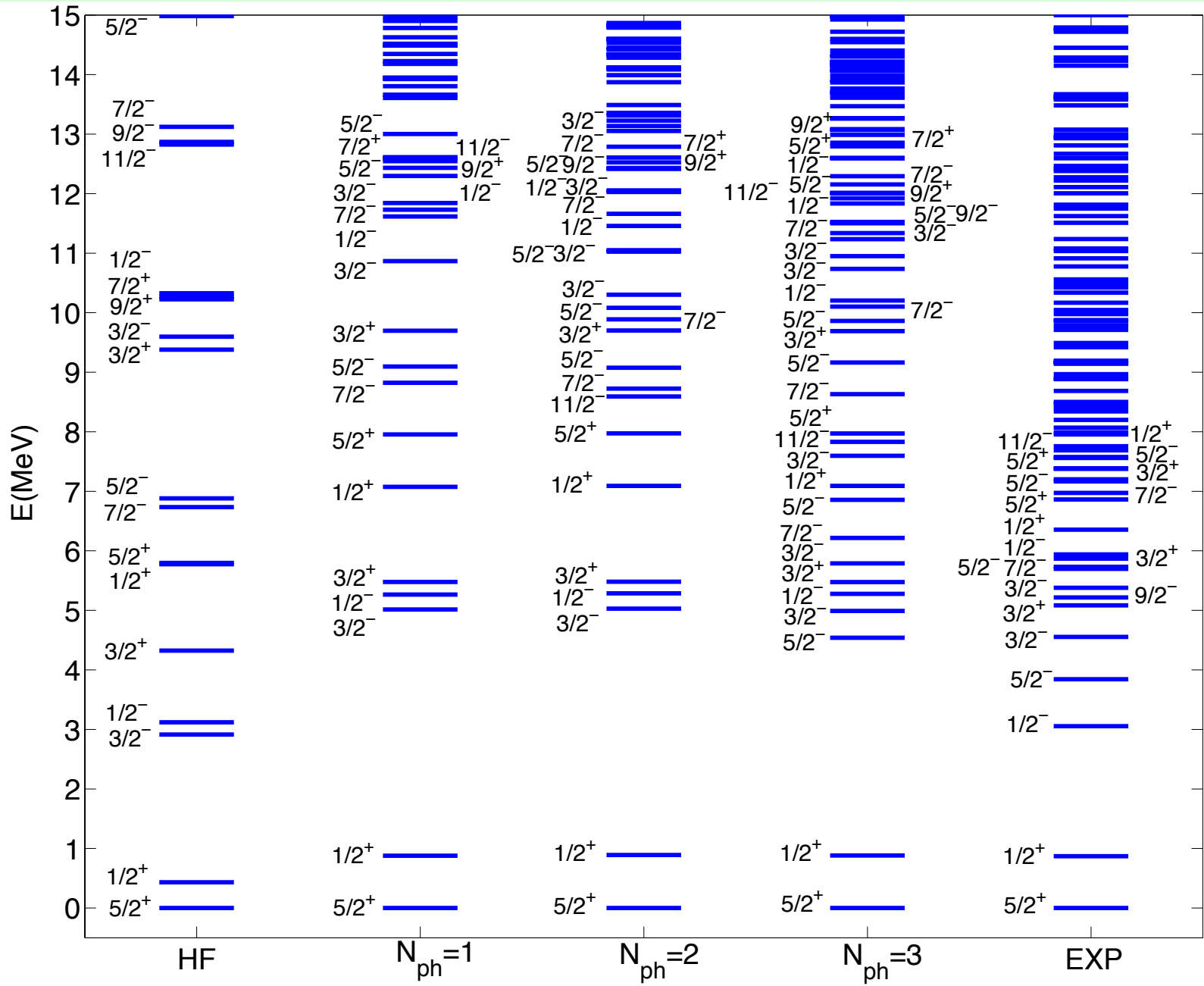
- **No approximations except for truncation !**
- **Pauli** principle fully accounted for:
- **No redundant** states!
- $|\mathbf{v}_n\rangle$ form an **orthonormal** basis

Implementation

Potential: NNLO_{opt}

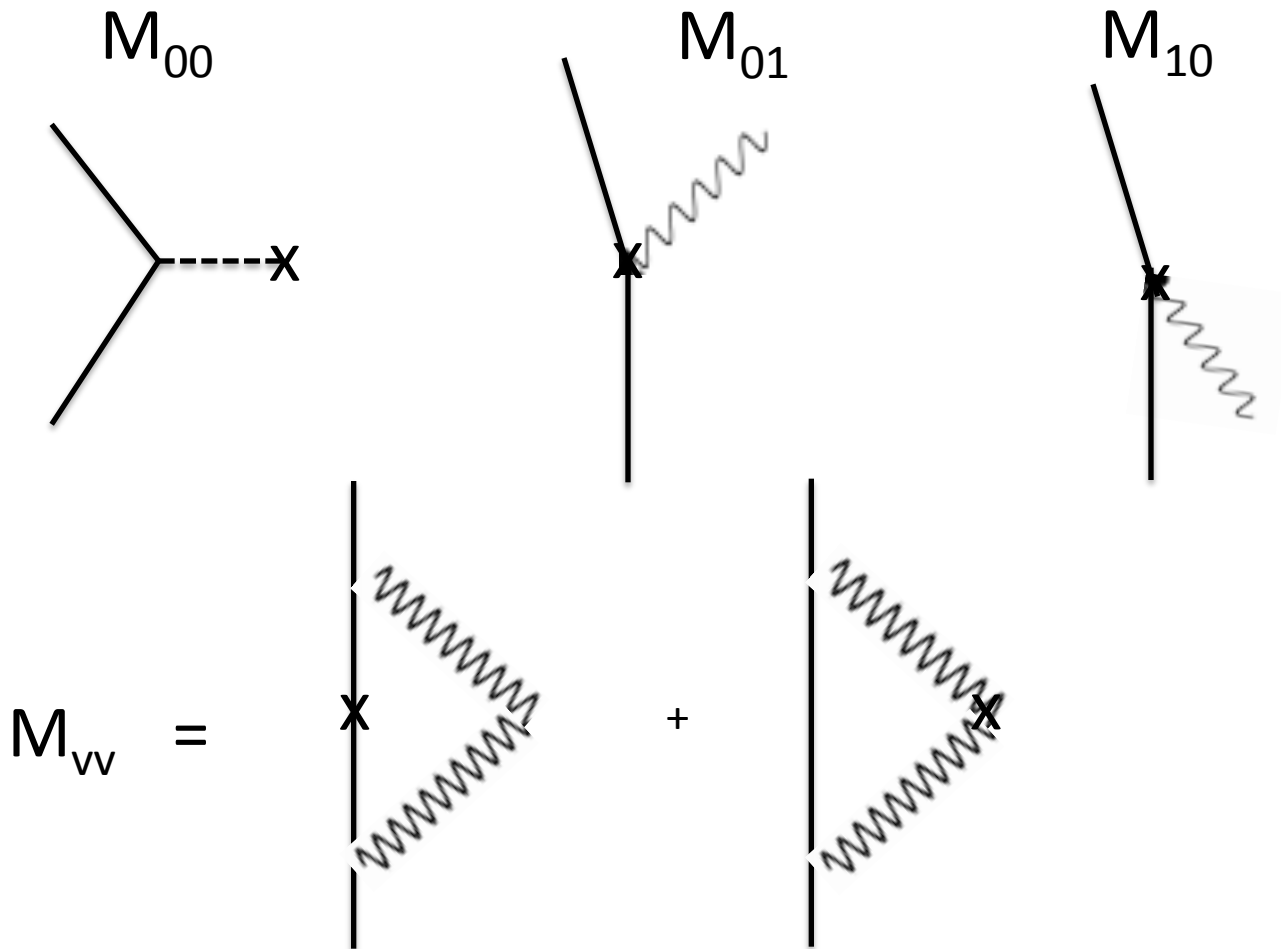
- Perform **HF**
- Construct **TDA** phonons (**free** of **CM** spurious admixtures)
- Generate the multiphonon basis $\{|\alpha_n\rangle\} \equiv \{|\alpha_0\rangle, |\alpha_1\rangle, |\alpha_2\rangle, |\alpha_3\rangle\}$
- Generate the **orthonormal** particle-phonon basis $\{|\mathbf{v}_n\rangle\}$
- Full diagonalization

^{17}O spectrum

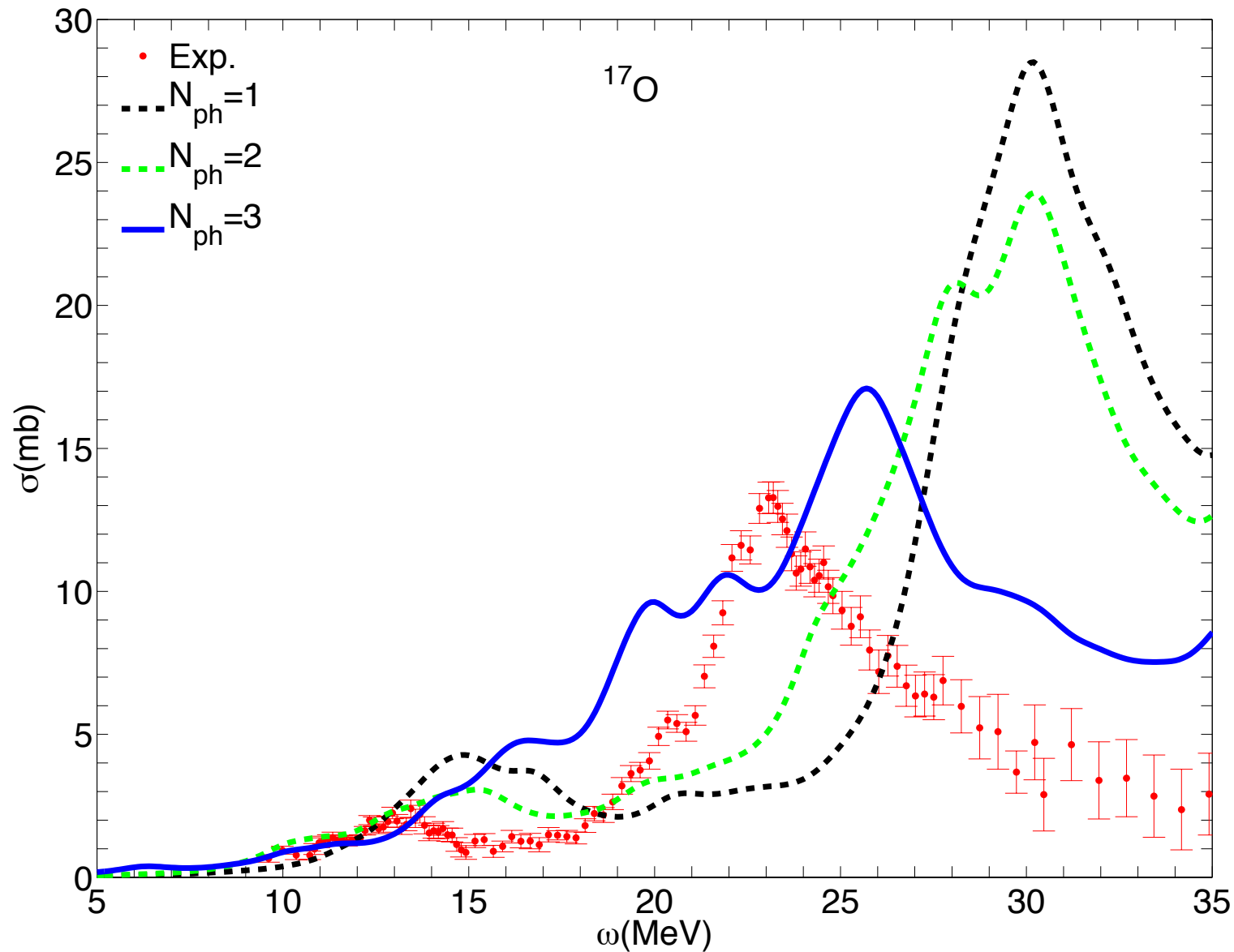


Transitions Amplitudes

$$\langle \Psi_{\Omega_f} \| M(\lambda) \| \Psi_{\Omega_i} \rangle = M_{00}(\lambda) + M_{01}(\lambda) + M_{10}(\lambda) + M_{11}(\lambda) + M_{12}(\lambda) + M_{21}(\lambda) + M_{22}(\lambda)$$

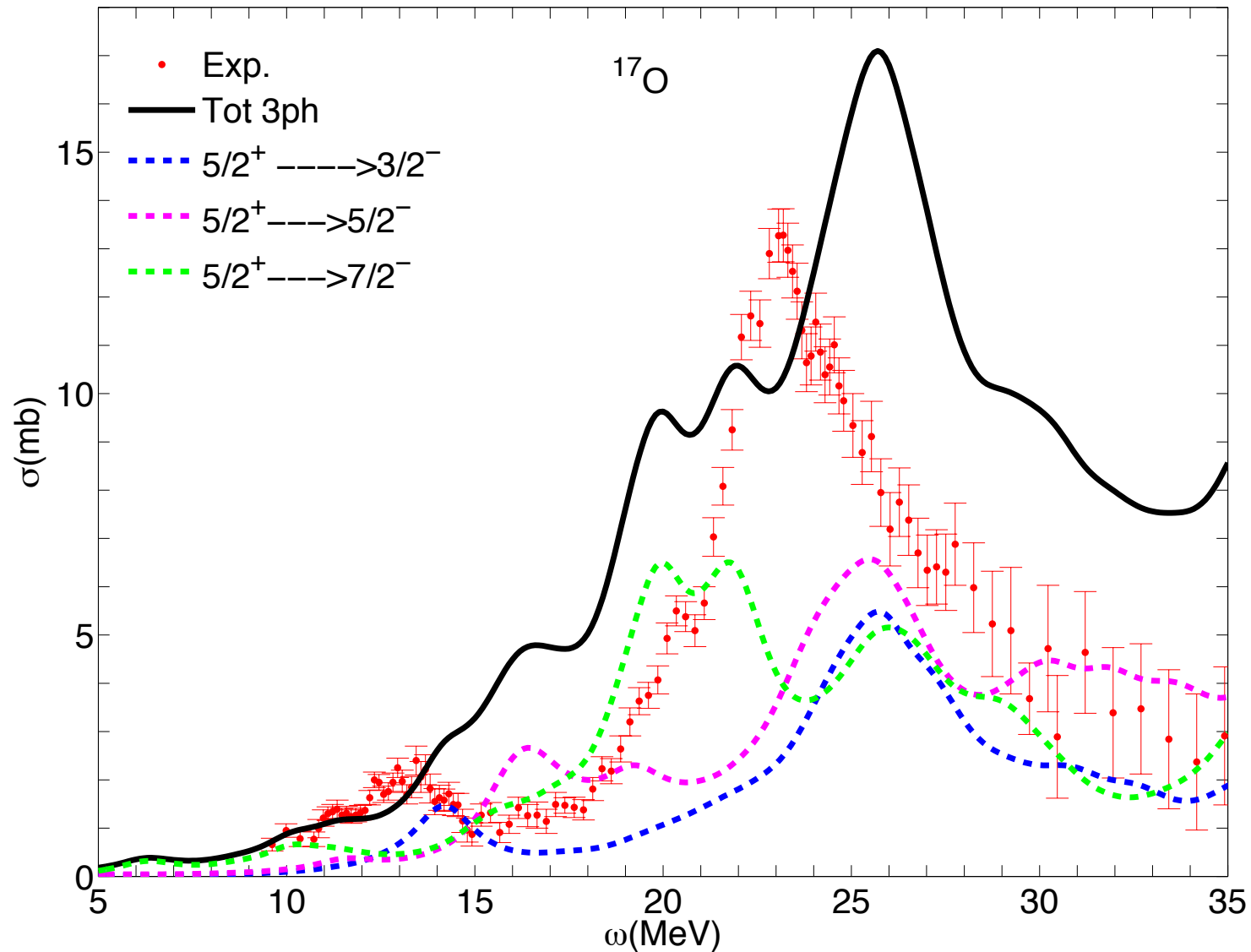


Cross section



J. W. Jury, et al., Phys Rev. C **21**, 503 (1980).

Cross section: Preliminary



Mean Values

^{17}O	HF	$N_{\text{ph}}=1$	$N_{\text{ph}}=2$	$N_{\text{ph}}=3$	Exp
Q(barn)	0	-0.00825	-0.00830	-0.00834	-0.025 ⁽¹⁾
μ (nm)	-1.9130	-1.8376	-1.834	-1.833	-1.893 ⁽¹⁾
log ft	3.2944	3.3891	3.3905	3.3913	3.358 ⁽²⁾

¹ N.P. Stone, Atomic Data and Nuclear Data Tables **90** (2005) 75–176

²D.R. Tilley Nuclear Physics A **636** (1998) 249-364

Odd nuclei : Concluding remarks

- One- and two-phonon states enhance greatly the density of levels consistently with experiments.
- Three-phonon states play an important role, their (strong) coupling to particle-phonon states $|v_1\rangle$
 - a) shift downwards the energies of $|v_1\rangle$
 - b) improve the cross section
 - c) increase the density of levels in low-energy part of the spectrum

Thank you