Direct reactions of weakly-bound nuclei within a one dimensional model

Laura Moschini

Andrea Vitturi & Antonio Moro

Outline

Introduction
The model
Results
Conclusions

Introduction

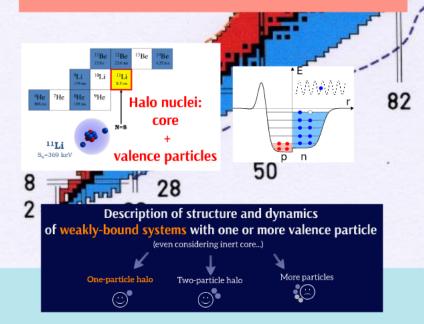
How to simplify the problem?

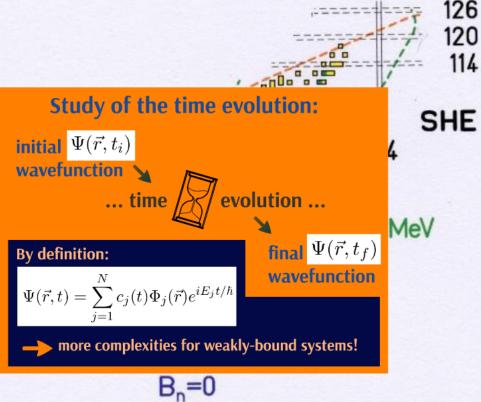
- ⚠ huge basis → truncations
- ⚠ continuum discretization

=>model-dependent approximations Let's move to one dimension!

$$\Psi(x,t) = \sum_{j=1}^{N} c_j(t) \Phi_j(x) e^{iE_j t/\hbar}$$

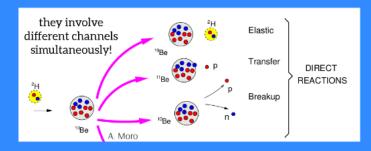
we can follow both the time evolutions and understand the limitations of approximations



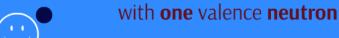


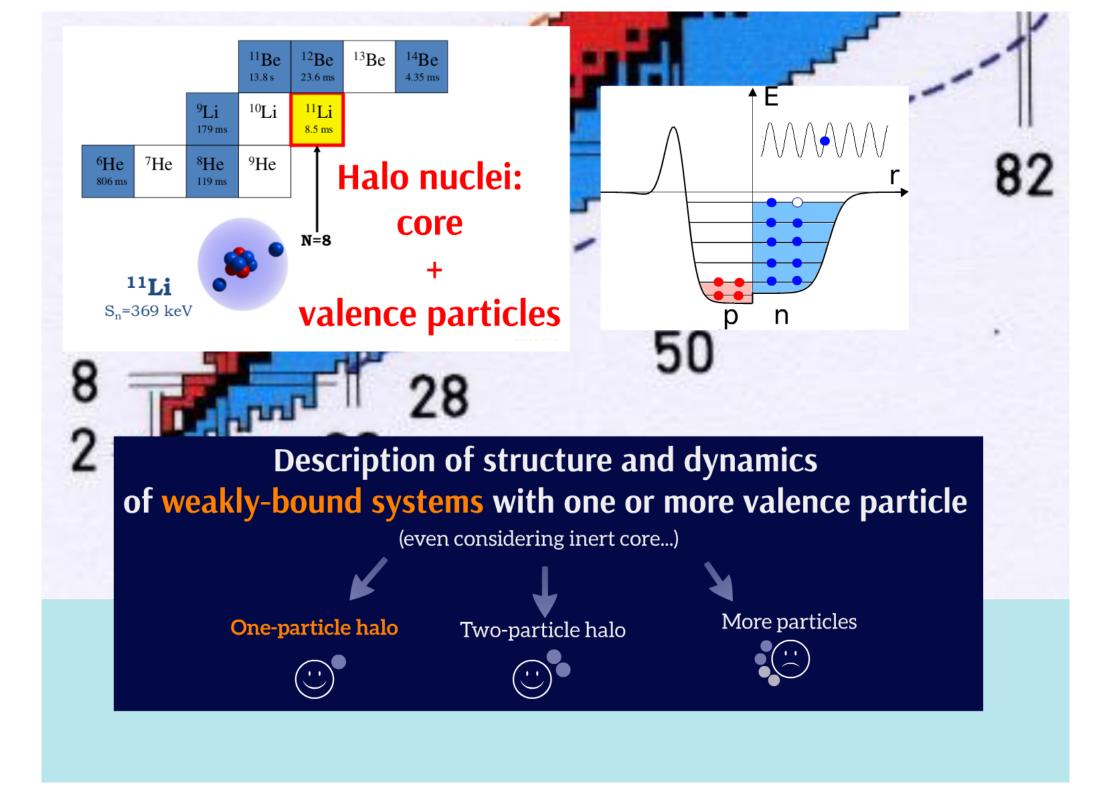
D_n-0

Study of **direct reactions** applied to

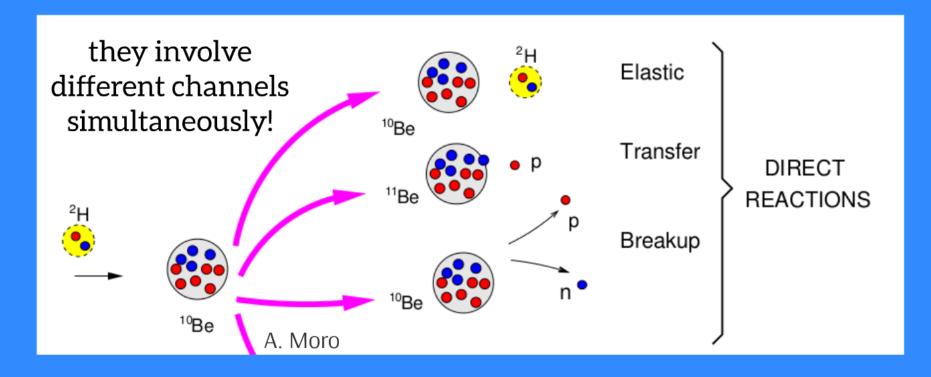


weakly-bound systems





Study of direct reactions applied to



weakly-bound systems
with one valence neutron

Study of the time evolution:

initial $\Psi(\vec{r},t_i)$ wavefunction

... time



evolution ...

By definition:

$$\Psi(\vec{r},t) = \sum_{j=1}^{N} c_j(t) \Phi_j(\vec{r}) e^{iE_j t/\hbar}$$

final $\Psi(\vec{r}, t_f)$ wavefunction

more complexities for weakly-bound systems!

How to simplify the problem?

- huge basis -> truncations
- _____ continuum discretization

=>model-dependent approximations Let's move to one dimension!

$$\Psi(x,t) = \sum_{j=1}^{N} c_j(t) \Phi_j(x) e^{iE_j t/\hbar}$$

we can follow both the time evolutions and understand the limitations of approximations

Introduction

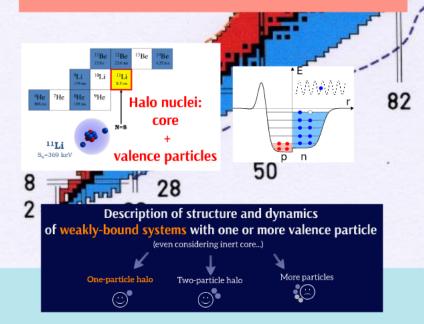
How to simplify the problem?

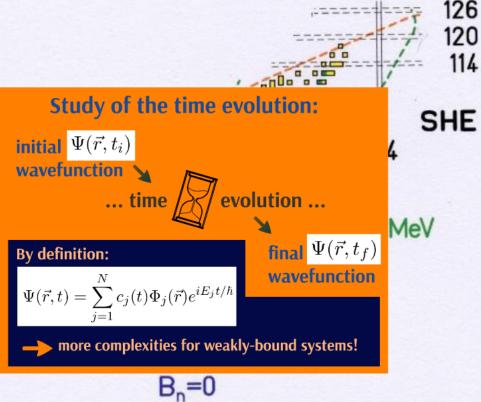
- ⚠ huge basis → truncations
- ⚠ continuum discretization

=>model-dependent approximations Let's move to one dimension!

$$\Psi(x,t) = \sum_{j=1}^{N} c_j(t) \Phi_j(x) e^{iE_j t/\hbar}$$

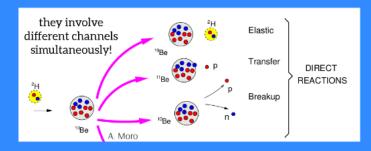
we can follow both the time evolutions and understand the limitations of approximations



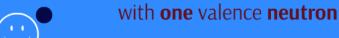


D_n-0

Study of **direct reactions** applied to



weakly-bound systems



Direct reactions of weakly-bound nuclei within a one dimensional model

Laura Moschini

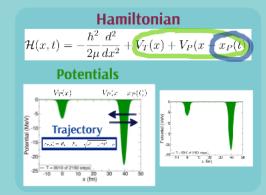
Andrea Vitturi & Antonio Moro

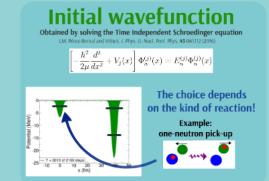
Outline

Introduction
The model
Results
Conclusions

The model

Problem definition and initial conditions





Time evolution

$$\Psi(x,t) = \sum_{j=1}^{N_T} c_j^T(t) \Phi_j^T(x) e^{iE_j^T t/\hbar} + \sum_{j=1}^{N_P} c_j^P(t) \Phi_j^P(x) e^{iE_j^P t/\hbar}$$

Exact model

Solution of time dependent Schroedinger equation

Vitturi and LM. J. Phys.: Conf. Ser. 590, 012007 (2015) itturi, LM, Hagino and Moro, AIP Conf. Proc. 1681, 060001 (2015)

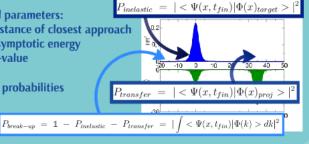
 $i\hbar \frac{d}{dt}\Psi(x,t) = \mathcal{H}(x,t)\Psi(x,t)$

 $\Psi(x, t + \Delta t) = \left(1 + \frac{i\Delta t}{2\hbar} \mathcal{H}\right)^{-1} \left(1 - \frac{i\Delta t}{2\hbar} \mathcal{H}\right) \Psi(x, t)$ Bonche, Koonin and Negele, Phys. Rev. C 13, 1226 (1976)

Initial parameters:

- distance of closest approach
- · asymptotic energy
- Q-value

Final probabilities



Coupled-channels

Initial condition:

$$c_j^P(t=-\infty)=0$$

$$c_j^T(t=-\infty)=\delta_{i,j}$$

 $i\hbar\frac{\partial c_j^T}{\partial t} = \sum c_k^T \langle \omega_j^T | V^P | \Psi_k^T \rangle + \sum c_k^P \langle \omega_j^T | V^T | \Psi_k^P \rangle$

$$i\hbar\frac{\partial c_j^P}{\partial t} = \sum c_k^T \langle \omega_j^P | V^P | \Psi_k^T \rangle + \sum c_k^P \langle \omega_j^P | V^T | \Psi_k^P \rangle$$

- i-th arget state is the initial state Esbensen, Broglia and Winther, Ann. Phys. 146, 149-173 (1983)
- inclusion of target AND projectile bases
- dual bases associated with the two wells to solve orthonormal problem
- time-dependent functions based on overlaps between target and projectile states

$$\langle \Psi_m^I | \omega_n^J \rangle = \delta_{I,J} \delta_{n,m}$$

New feature: continuum inclusion!!!

Final probabilities

$$P_j^{(T,P)}(t_f) = |c_j^{(T,P)}(t_f)|^2$$

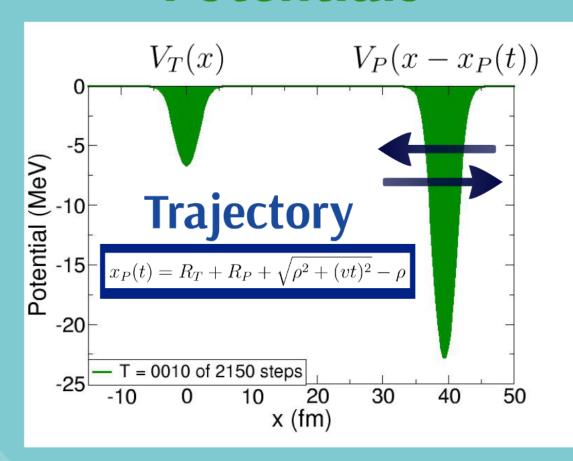
non-orthogonality of basis states => tot probability is not conserved during collision

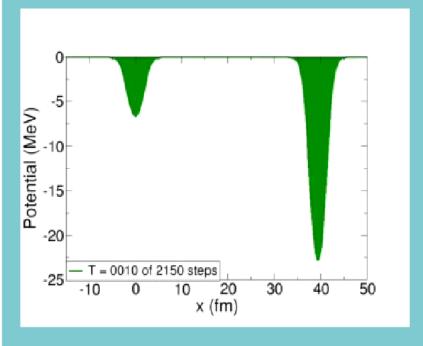
To be calculated after the collision when overlaps are zero

Hamiltonian

$$\mathcal{H}(x,t) = -\frac{\hbar^2}{2\mu} \frac{d^2}{dx^2} + V_T(x) + V_P(x - x_P(t))$$

Potentials



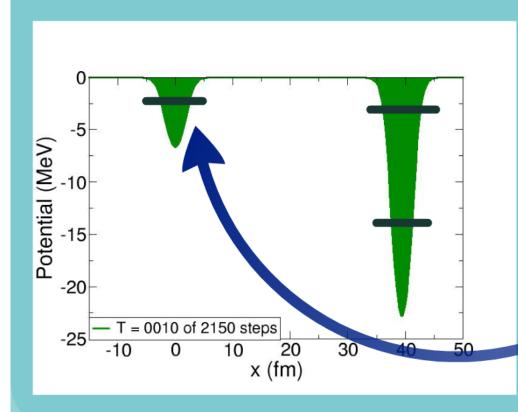


Initial wavefunction

Obtained by solving the Time Independent Schroedinger equation

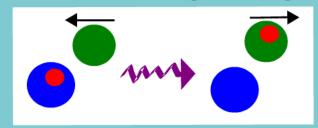
LM, Pérez-Bernal and Vitturi, J. Phys. G: Nucl. Part. Phys. 43 045112 (2016)

$$\left[-\frac{\hbar^2}{2\mu} \frac{d^2}{dx^2} + V_j(x) \right] \Phi_n^{(j)}(x) = E_n^{(j)} \Phi_n^{(j)}(x)$$



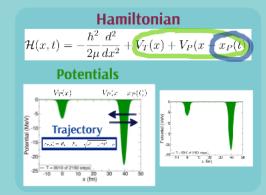
The choice depends on the kind of reaction!

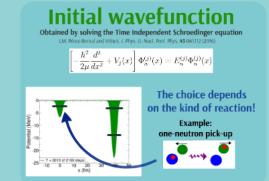
Example: one-neutron pick-up



The model

Problem definition and initial conditions





Time evolution

$$\Psi(x,t) = \sum_{j=1}^{N_T} c_j^T(t) \Phi_j^T(x) e^{iE_j^T t/\hbar} + \sum_{j=1}^{N_P} c_j^P(t) \Phi_j^P(x) e^{iE_j^P t/\hbar}$$

Exact model

Solution of time dependent Schroedinger equation

Vitturi and LM. J. Phys.: Conf. Ser. 590, 012007 (2015) itturi, LM, Hagino and Moro, AIP Conf. Proc. 1681, 060001 (2015)

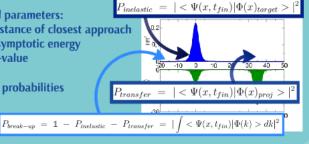
 $i\hbar \frac{d}{dt}\Psi(x,t) = \mathcal{H}(x,t)\Psi(x,t)$

 $\Psi(x, t + \Delta t) = \left(1 + \frac{i\Delta t}{2\hbar} \mathcal{H}\right)^{-1} \left(1 - \frac{i\Delta t}{2\hbar} \mathcal{H}\right) \Psi(x, t)$ Bonche, Koonin and Negele, Phys. Rev. C 13, 1226 (1976)

Initial parameters:

- distance of closest approach
- · asymptotic energy
- Q-value

Final probabilities



Coupled-channels

Initial condition:

$$c_j^P(t=-\infty)=0$$

$$c_j^T(t=-\infty)=\delta_{i,j}$$

 $i\hbar\frac{\partial c_j^T}{\partial t} = \sum c_k^T \langle \omega_j^T | V^P | \Psi_k^T \rangle + \sum c_k^P \langle \omega_j^T | V^T | \Psi_k^P \rangle$

$$i\hbar\frac{\partial c_j^P}{\partial t} = \sum c_k^T \langle \omega_j^P | V^P | \Psi_k^T \rangle + \sum c_k^P \langle \omega_j^P | V^T | \Psi_k^P \rangle$$

- i-th arget state is the initial state Esbensen, Broglia and Winther, Ann. Phys. 146, 149-173 (1983)
- inclusion of target AND projectile bases
- dual bases associated with the two wells to solve orthonormal problem
- time-dependent functions based on overlaps between target and projectile states

$$\langle \Psi_m^I | \omega_n^J \rangle = \delta_{I,J} \delta_{n,m}$$

New feature: continuum inclusion!!!

Final probabilities

$$P_j^{(T,P)}(t_f) = |c_j^{(T,P)}(t_f)|^2$$

non-orthogonality of basis states => tot probability is not conserved during collision

To be calculated after the collision when overlaps are zero

Exact model

Solution of time dependent Schroedinger equation

Vitturi and LM, *J. Phys.: Conf. Ser.* **590**, 012007 (2015) Vitturi, LM, Hagino and Moro, *AIP Conf. Proc.* **1681**, 060001 (2015)

$$i\hbar \frac{d}{dt}\Psi(x,t) = \mathcal{H}(x,t)\Psi(x,t)$$

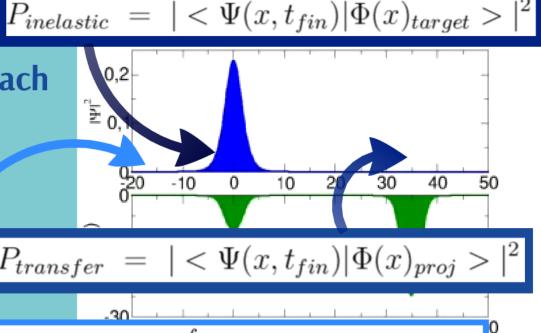
$$\Psi(x, t + \Delta t) = \left(1 + \frac{i\Delta t}{2\hbar} \mathcal{H}\right)^{-1} \left(1 - \frac{i\Delta t}{2\hbar} \mathcal{H}\right) \Psi(x, t)$$

Bonche, Koonin and Negele, Phys. Rev. C 13, 1226 (1976)

Initial parameters:

- distance of closest approach
- asymptotic energy
- Q-value

Final probabilities



$$P_{break-up} = 1 - P_{inelastic} - P_{transfer} = |\int \langle \Psi(x, t_{fin}) | \Phi(k) \rangle dk|^2$$

Coupled-channels

Initial condition:

$$c_j^P(t=-\infty)=0$$

$$c_j^T(t=-\infty)=\delta_{i,j}$$

i-th arget state is the initial state

$$i\hbar \frac{\partial c_j^T}{\partial t} = \sum c_k^T \langle \omega_j^T | V^P | \Psi_k^T \rangle + \sum c_k^P \langle \omega_j^T | V^T | \Psi_k^P \rangle$$

$$i\hbar \frac{\partial c_j^P}{\partial t} = \sum_{k} c_k^T \langle \omega_j^P | V^P | \Psi_k^T \rangle + \sum_{k} c_k^P \langle \omega_j^P | V^T | \Psi_k^P \rangle$$

Esbensen, Broglia and Winther, Ann. Phys. 146, 149-173 (1983)

- inclusion of target AND projectile bases
- dual bases associated with the two wells to solve orthonormal problem
- time-dependent functions based on overlaps between target and projectile states

$$\langle \Psi_m^I | \omega_n^J \rangle = \delta_{I,J} \delta_{n,m}$$

New feature: continuum inclusion!!!

Final probabilities

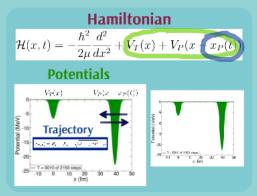
$$P_j^{(T,P)}(t_f) = |c_j^{(T,P)}(t_f)|^2$$

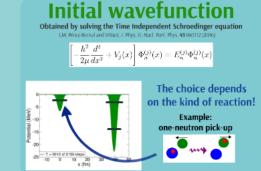
non-orthogonality of basis states => tot probability is not conserved during collision

To be calculated after the collision when overlaps are zero

The model

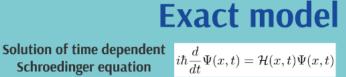
Problem definition and initial conditions

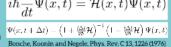


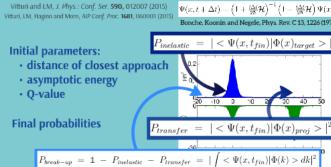


Time evolution

$$\Psi(x,t) = \sum_{j=1}^{N_T} c_j^T(t) \Phi_j^T(x) e^{iE_j^T t/\hbar} + \sum_{j=1}^{N_P} c_j^P(t) \Phi_j^P(x) e^{iE_j^P t/\hbar}$$







Coupled-channels

Initial condition:

$$c_j^P(t=-\infty)=0$$

$$c_j^T(t=-\infty)=\delta_{i,j}$$

$$\begin{split} &i\hbar\frac{\partial c_j^I}{\partial t} = \sum c_k^T \langle \omega_j^T | V^P | \Psi_k^T \rangle + \sum c_k^P \langle \omega_j^T | V^T | \Psi_k^P \rangle \\ &i\hbar\frac{\partial c_j^P}{\partial t} = \sum c_k^T \langle \omega_j^P | V^P | \Psi_k^T \rangle + \sum c_k^P \langle \omega_j^P | V^T | \Psi_k^P \rangle \end{split}$$

i-th arget state is the initial state Esbensen, Broglia and Winther, Ann. Phys. 146, 149-173 (1983)

- inclusion of target AND projectile bases
- dual bases associated with the two wells to solve orthonormal problem
- time-dependent functions based on overlaps between target and projectile states

$$\langle \Psi_m^I | \omega_n^J \rangle = \delta_{I,J} \delta_{n,m}$$

New feature: continuum inclusion!!!

Final probabilities

$$P_j^{(T,P)}(t_f) = |c_j^{(T,P)}(t_f)|^2$$

non-orthogonality of basis states => tot probability is not conserved during collision

To be calculated after the collision when overlaps are zero

Direct reactions of weakly-bound nuclei within a one dimensional model

Laura Moschini

Andrea Vitturi & Antonio Moro

Outline

Introduction
The model
Results
Conclusions

Results

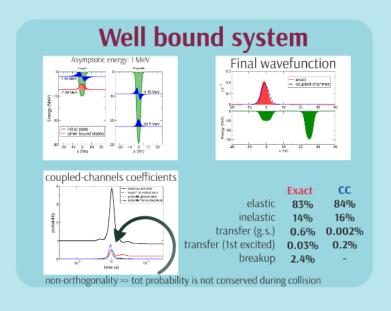
The same equation can be solved within

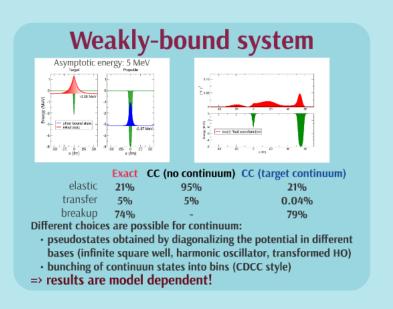
- the "exact" numerical method
- the coupled-channels formalism

The goal is to check

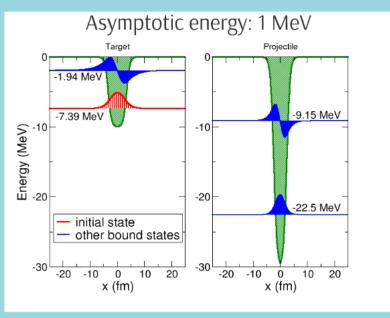
- the validity of the approximation
- the necessary truncation of the two bases
- · continuum discretization

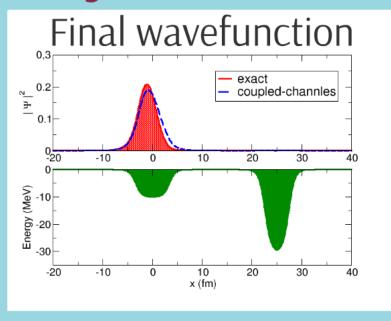
This is particular relevant in the case of weakly-bound systems

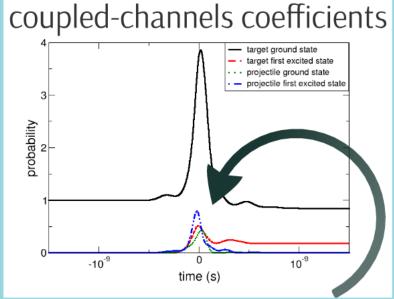




Well bound system



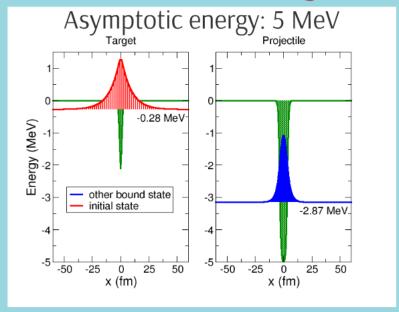


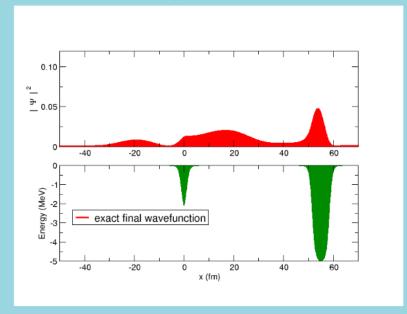


	Exact	
elastic	83%	84%
inelastic	14%	16%
transfer (g.s.)	0.6%	0.002%
ransfer (1st excited)	0.03%	0.2%
breakup	2.4%	-

non-orthogonality => tot probability is not conserved during collision

Weakly-bound system





Exact CC (no continuum) CC (target continuum)

 elastic
 21%

 transfer
 5%

 5%
 0.04%

 breakup
 74%

 79%

Different choices are possible for continuum:

- pseudostates obtained by diagonalizing the potential in different bases (infinite square well, harmonic oscillator, transformed HO)
- bunching of continuun states into bins (CDCC style)
- => results are model dependent!

Results

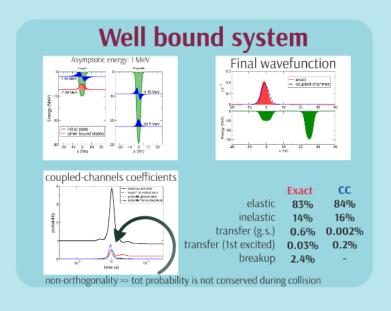
The same equation can be solved within

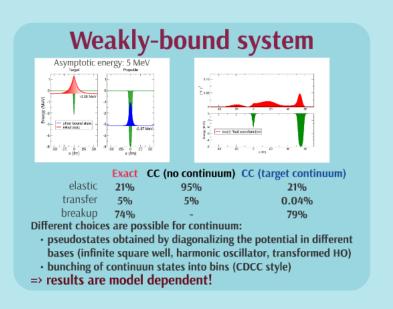
- the "exact" numerical method
- the coupled-channels formalism

The goal is to check

- the validity of the approximation
- the necessary truncation of the two bases
- · continuum discretization

This is particular relevant in the case of weakly-bound systems





Direct reactions of weakly-bound nuclei within a one dimensional model

Laura Moschini

Andrea Vitturi & Antonio Moro

Outline

Introduction
The model
Results
Conclusions

Conclusions

- We developed a one dimensional model to solve direct reactions involving nuclei with one valence neutron
- We deal with the problem by solving numerically the time-dependent Schroedinger equation and by applying the coupled-channels formalism
- We understand that for weakly-bound systems it is mandatory to include the continuum and that its definition is model dependent