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Newton
Fund



THE ROYAL
SOCIETY

Electron Capture processes in intermediate mass stars

Andrea Idini¹

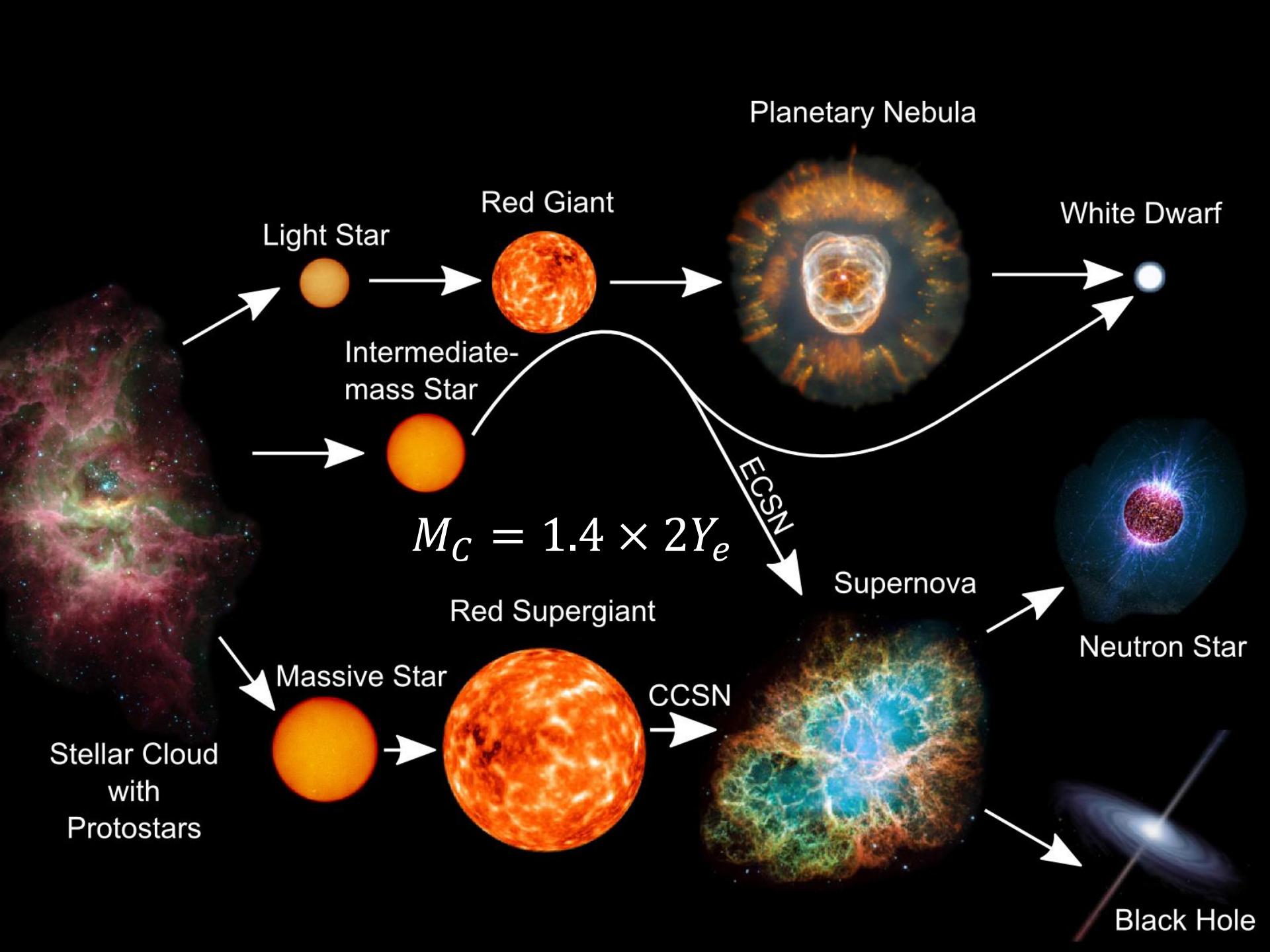
B. A. Brown², K. Langanke^{3,4}

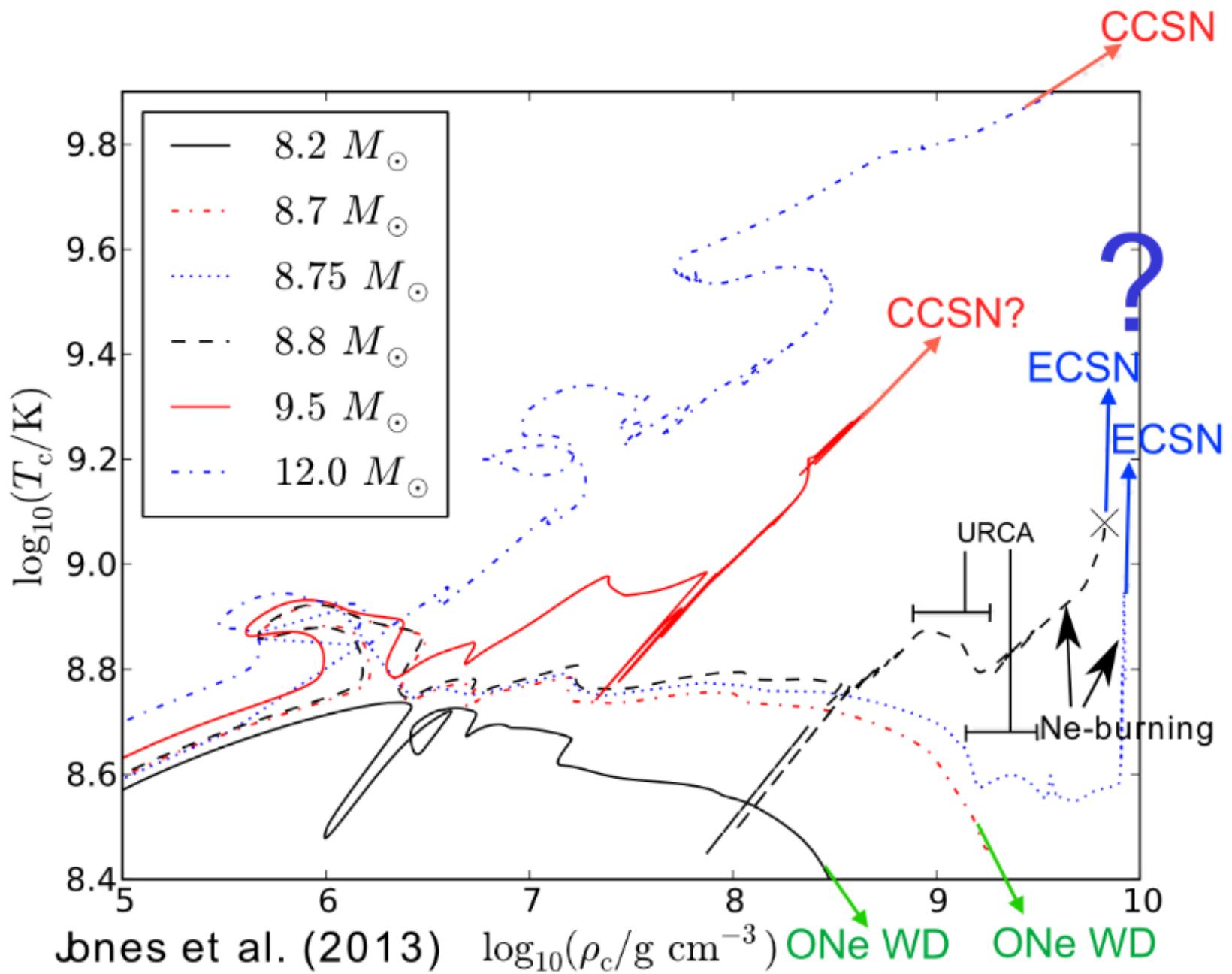
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^{20}Ne ($Q_{ec} = -7.5$ MeV)

^{24}Mg ($Q_{ec} = -3.4$ MeV)

Urca Cooling

e^- conserving

odd A
 $Q_2 > Q_1$

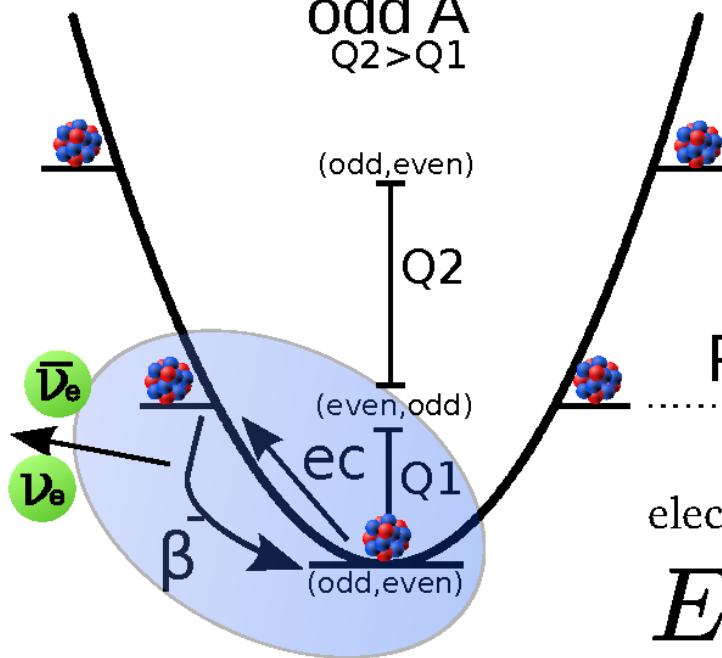
(odd,even)
Q2

(even,odd)

(odd,even)

Pairing

electron fermi energy:
 $E_f \propto \rho^{1/3}$



Heating

e^- decreasing
even A

$Q_1 > Q_2$

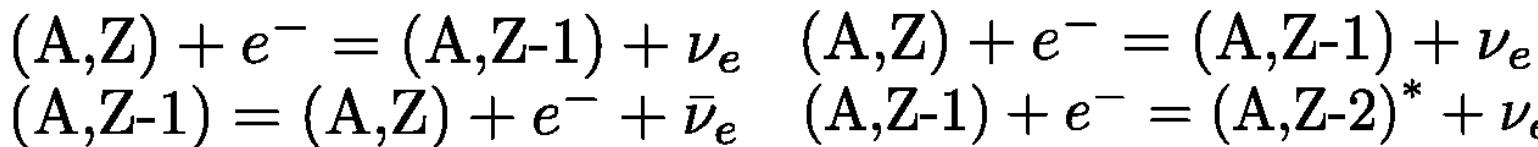
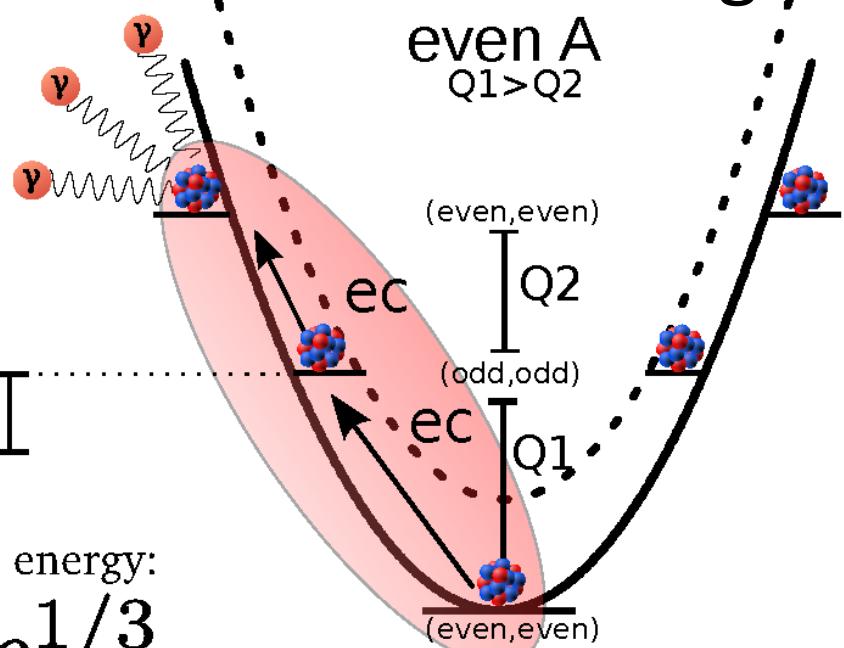
(even,even)

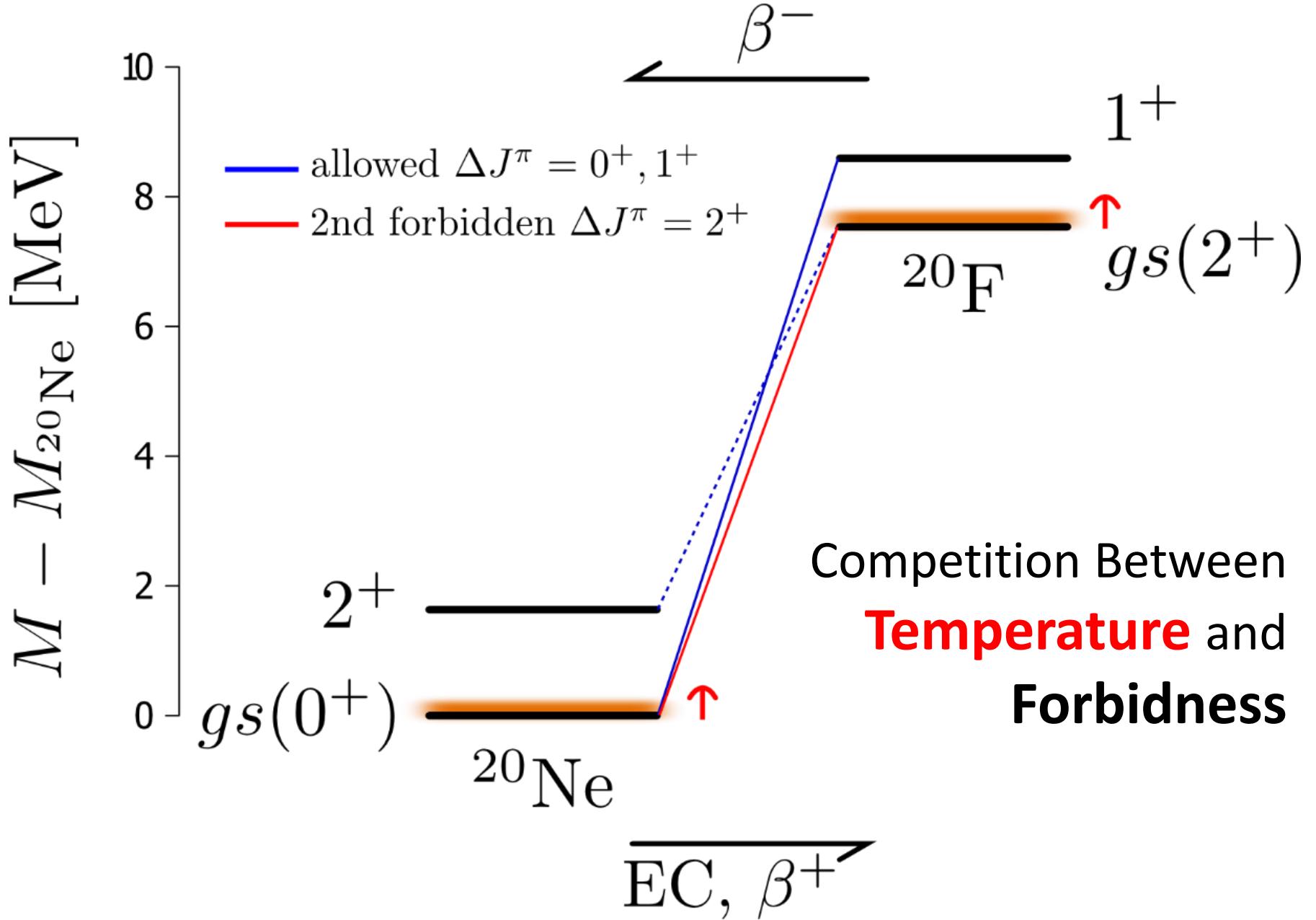
Q2

(odd,odd)

Q1

(even,even)

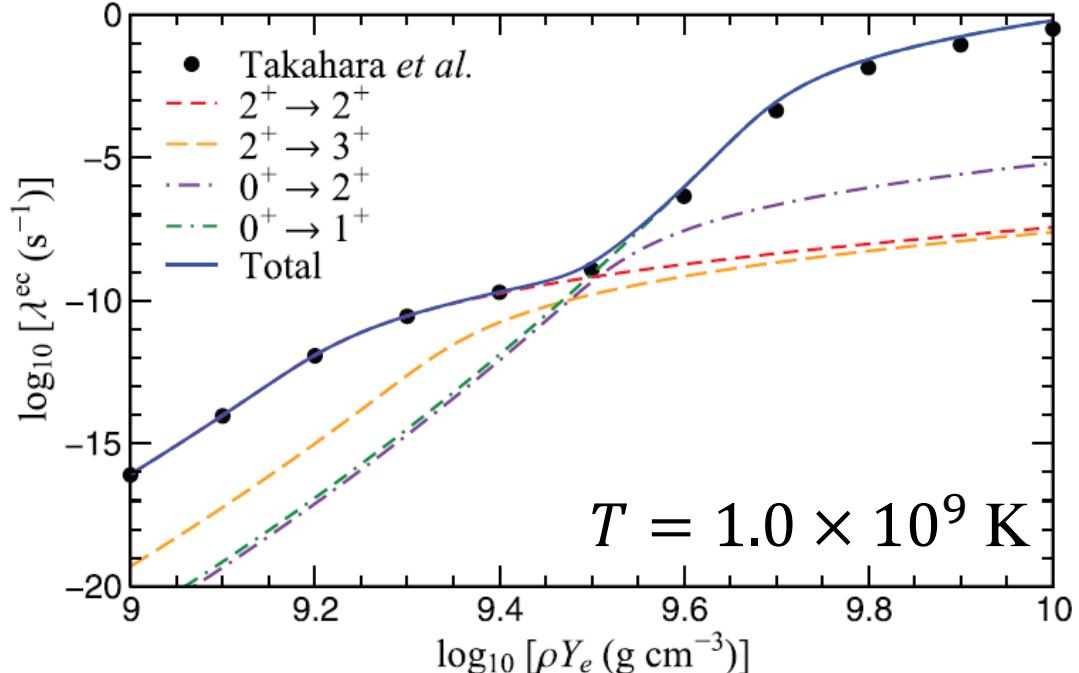
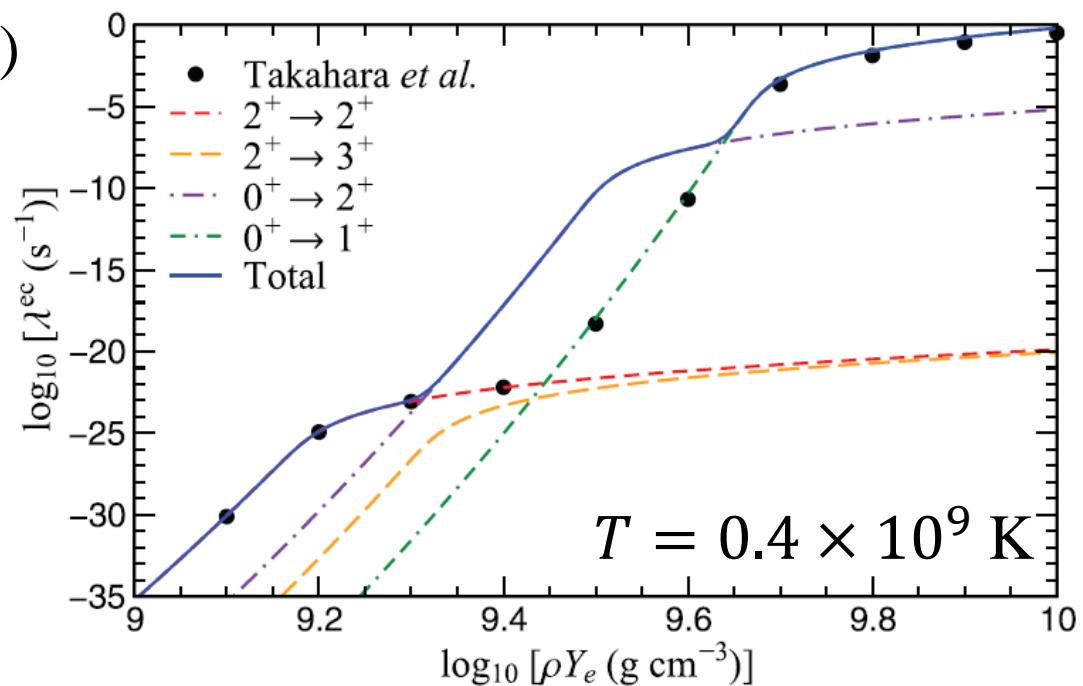




$BR \sim 1 \times 10^{-5}$
(exp. upper limit)

NPA 475, 1 (1987)

Estimate needed!
(possibly experiment)



$$\lambda^{ec} = \frac{\ln 2}{K} \int_{|Q|}^{+\infty} E_\nu^2 C(E) Ep f(E, kT, \mu) F(Z, E) dE$$

Shape Factor
(Nuclear Physics)
Fermi Distribution
(Statistical Description)

Kinematic **Fermi Function**

$$C(E) \propto \sum | \langle f | H_\beta | i \rangle |^2 \quad |f, i\rangle \text{ USD-B Shell Model}$$

Combination of ↓

$$\vec{r}^a, \vec{p}^b, \sigma, \tau$$

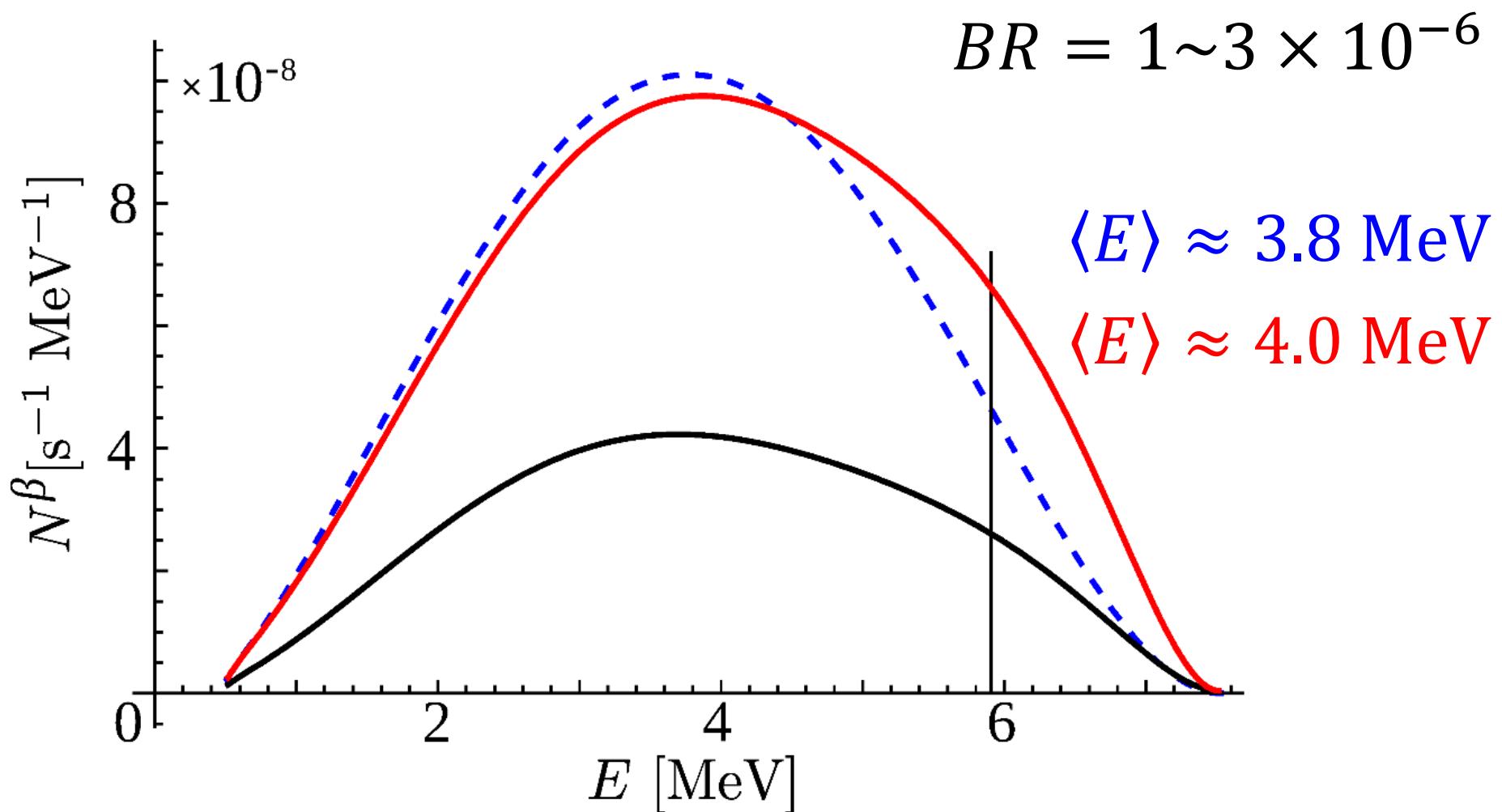
$C = \text{const}$ (allowed)

$C(E)$ (forbidden)

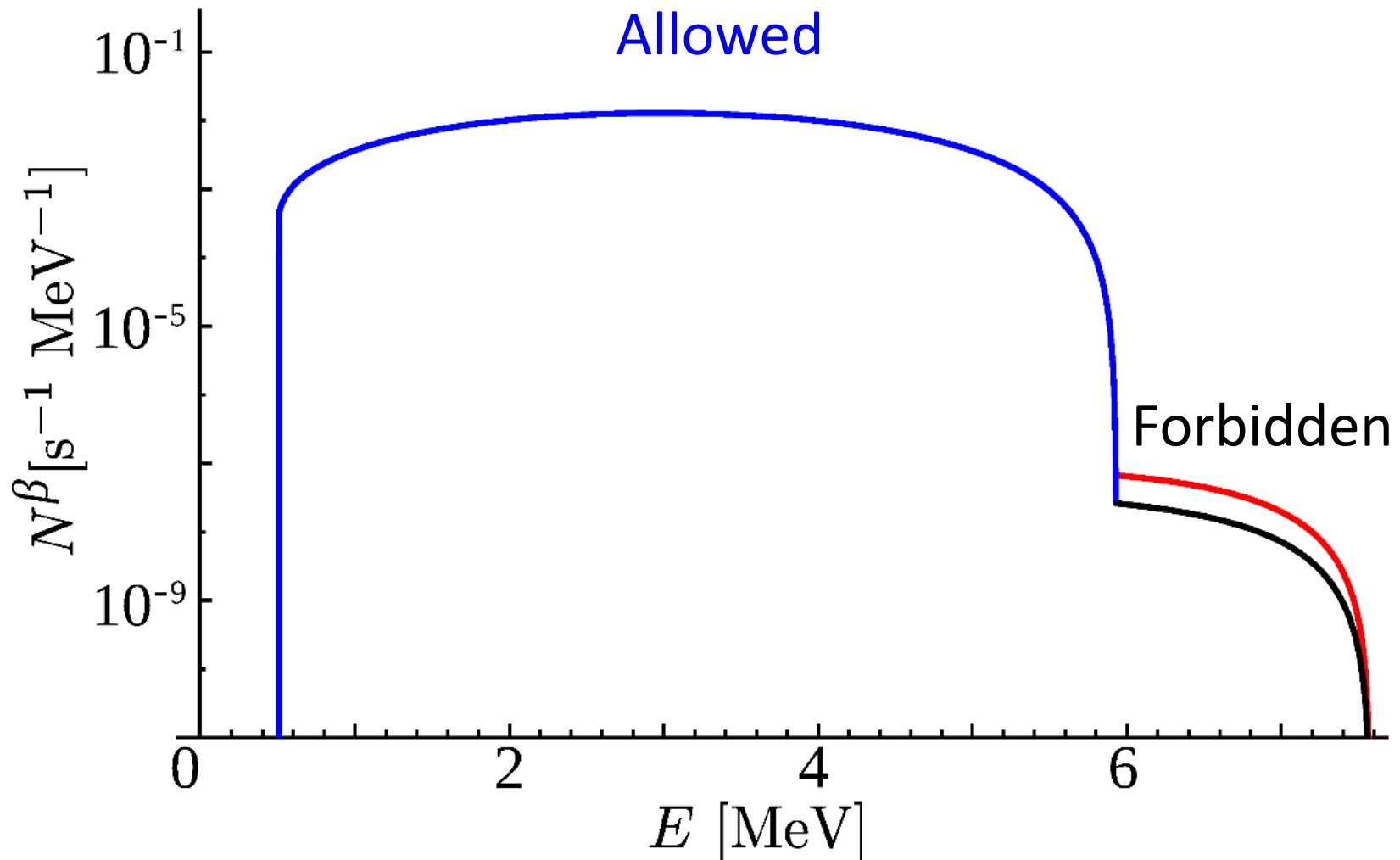
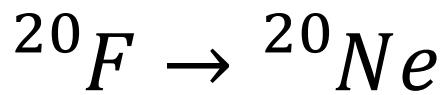
$$C \sim E^4 \quad (2^\circ \text{ forbidden})$$

$$\lambda^{\beta^-} = \frac{\ln 2}{K} \int_1^Q \underbrace{E_\nu^2 E p \, C(E) F(Z, E)}_{C_F} dE$$

$^{20}F(gs(2^+)) \rightarrow ^{20}Ne(gs(0^+))$ β -spectrum



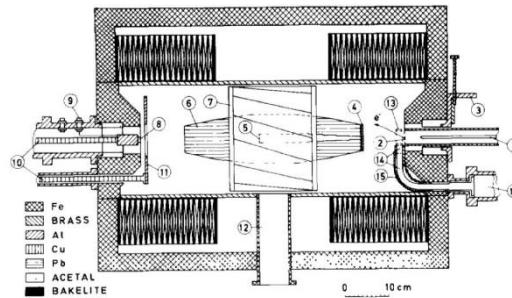
Total β -spectrum



MARA Spectrometer Jyväskylä

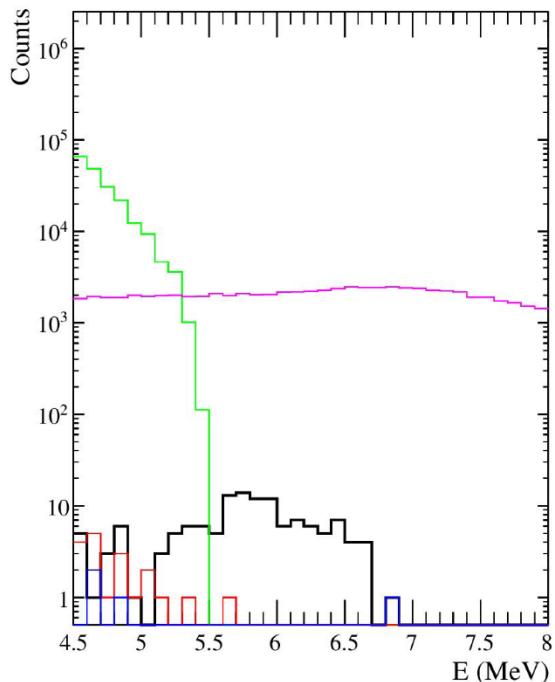
GEANT4 simulation

(1 week @ 40 kHz)

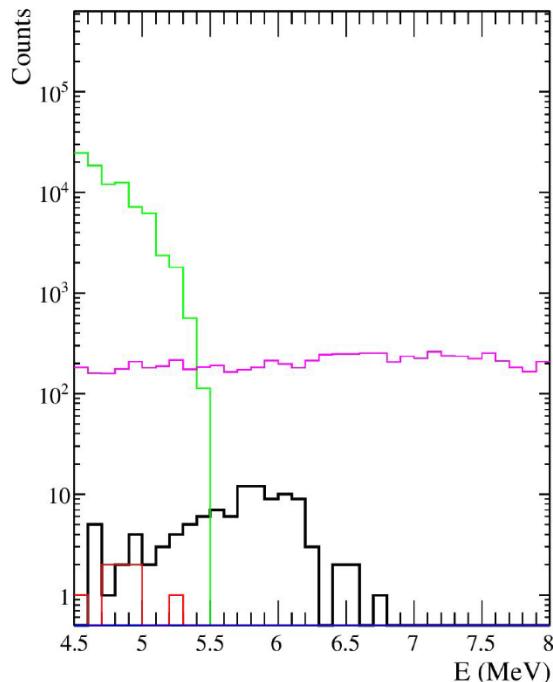


- $2^+ \rightarrow 0^+$ ($\text{br}=1 \times 10^{-6}$)
- $2^+ \rightarrow 2^+$
- $\beta\gamma$ summing
- $\beta\beta$ pile-up ($\tau=20\text{ns}$)
- CR muons

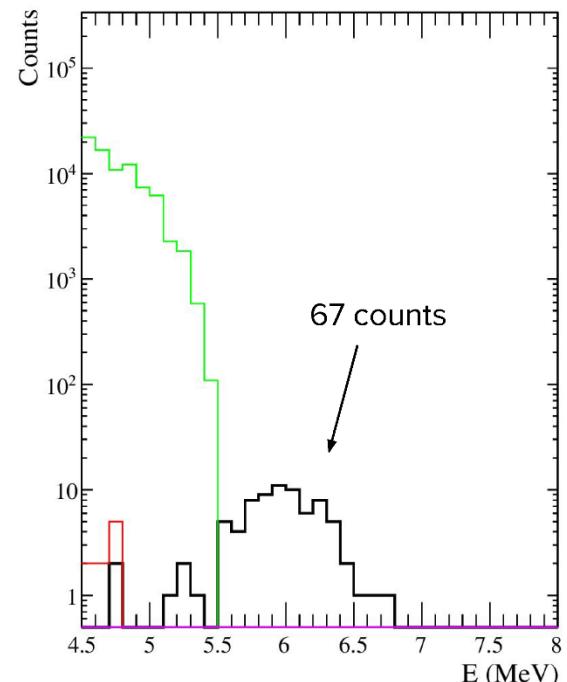
Single plastic detector



ΔE -E plastic telescope

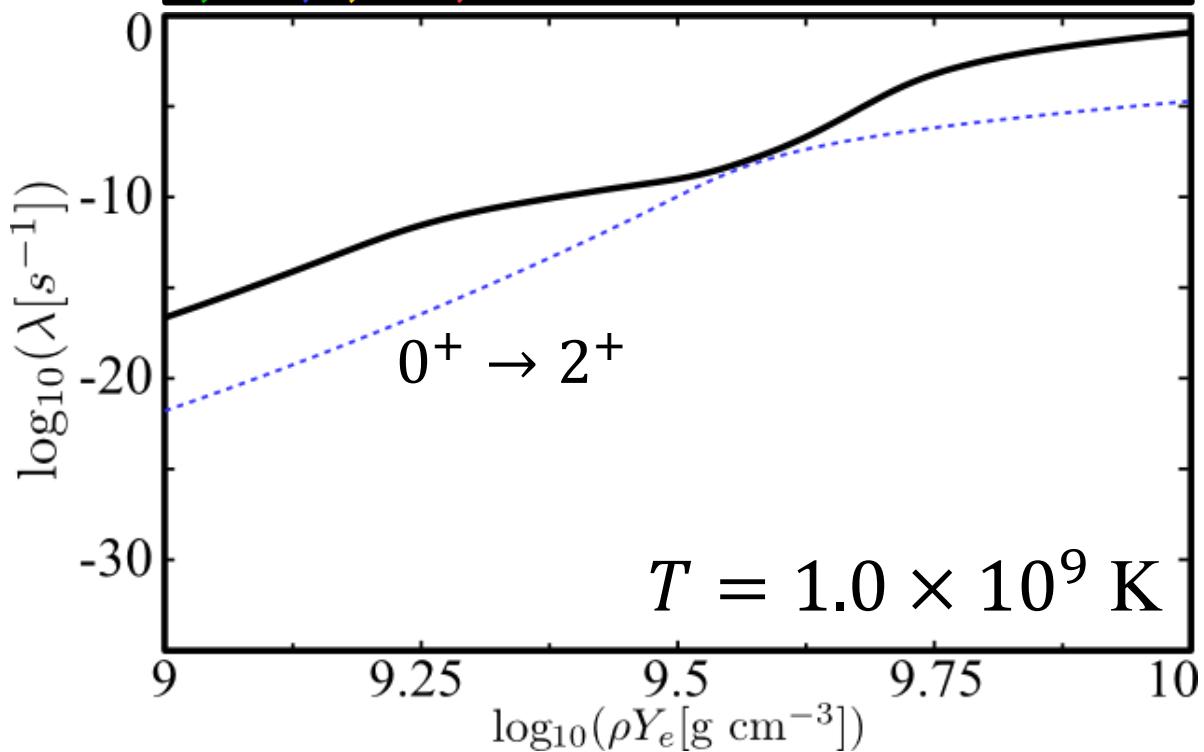
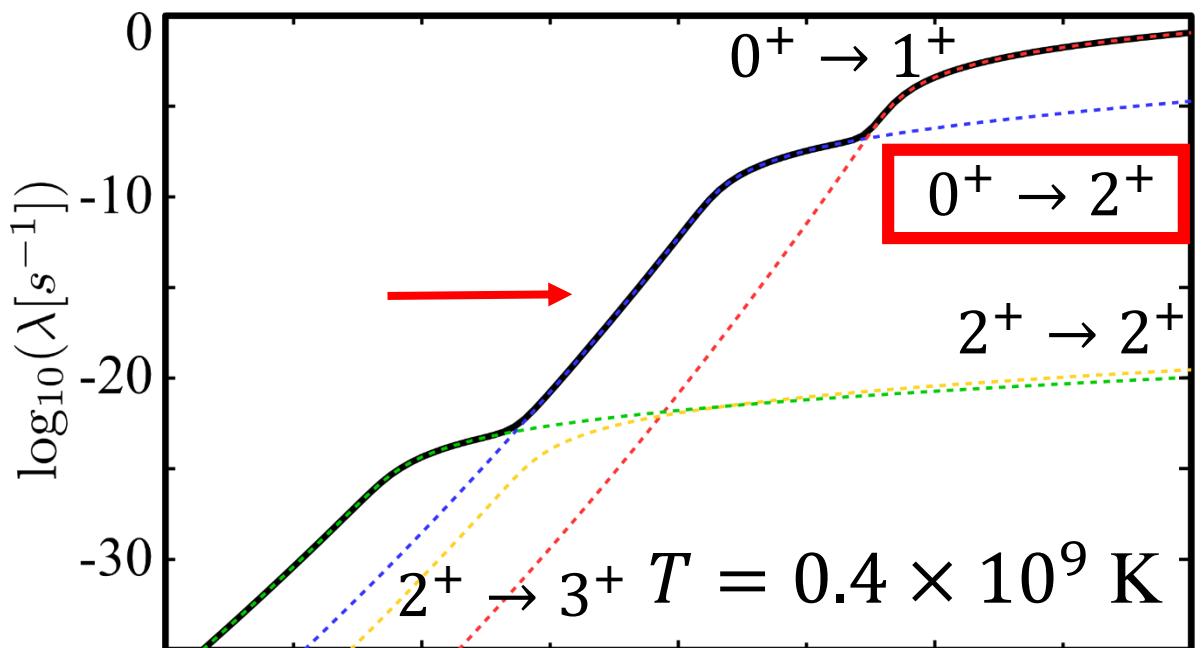


ΔE -E plastic telescope + μ veto



e^- capture

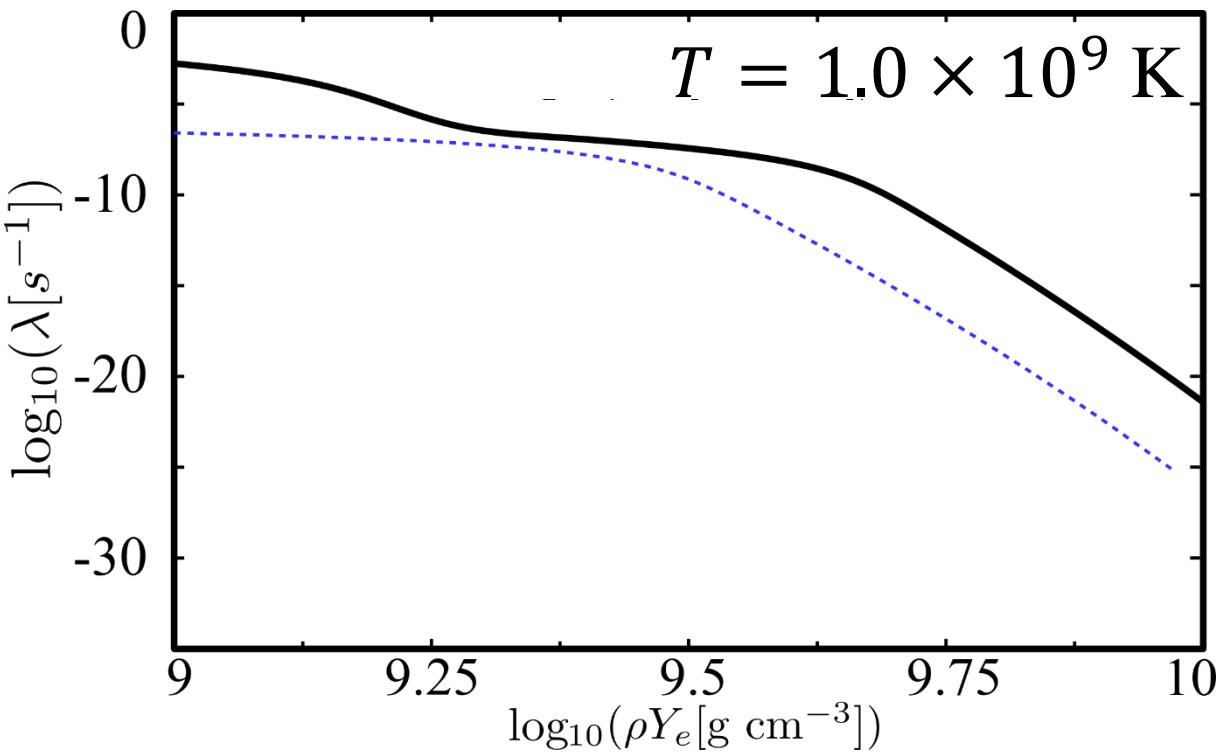
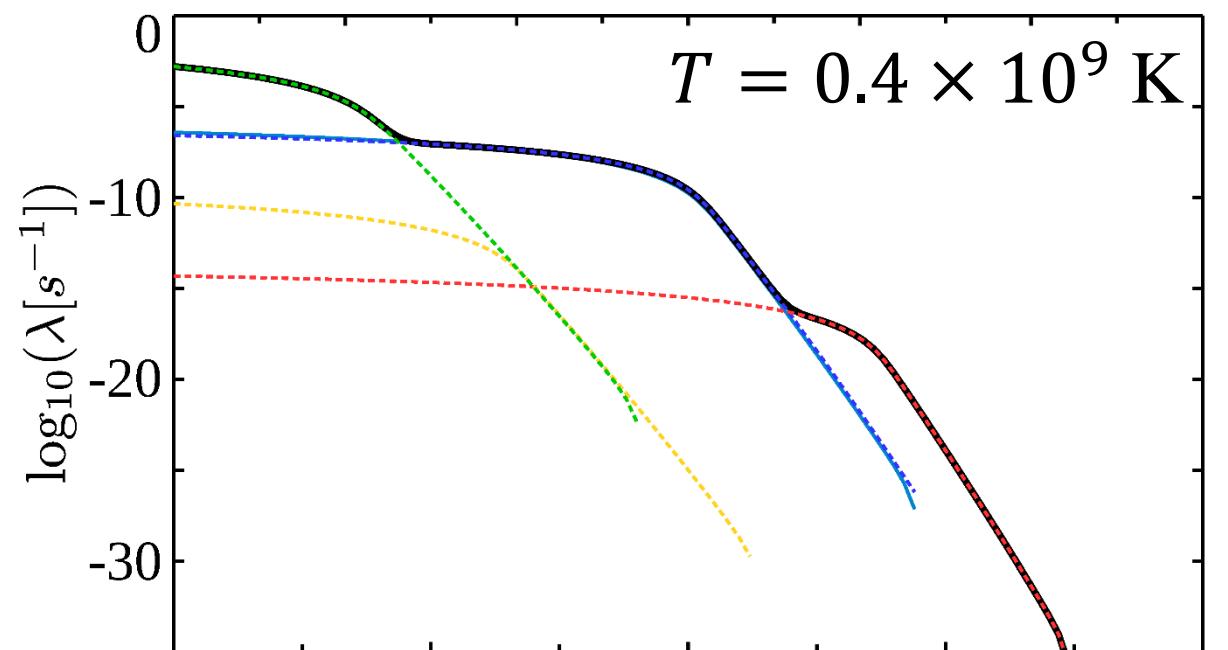
*In Astrophysical
Environment*

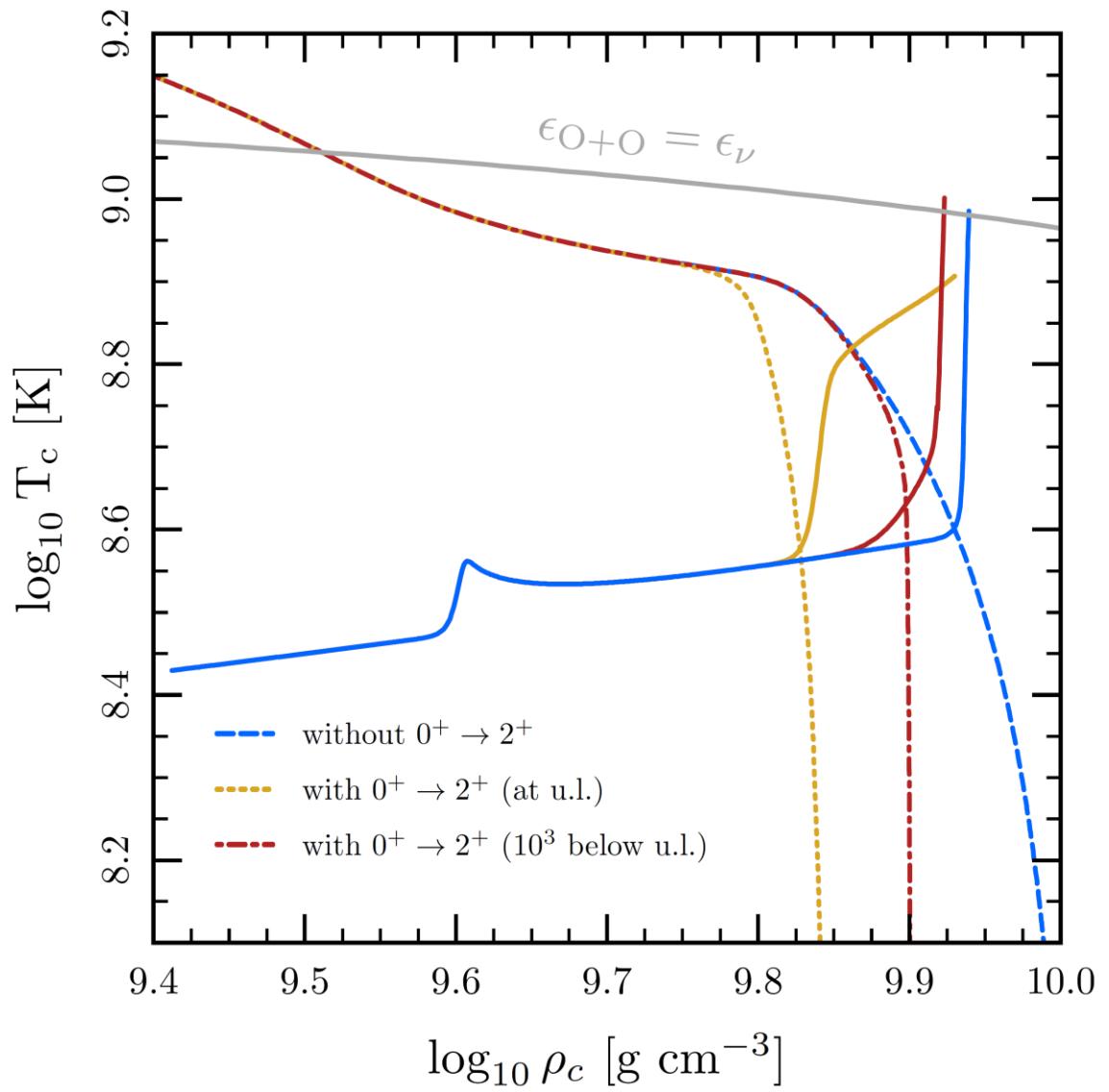


β^- decay

*In Astrophysical
Environment*

$0^+ \rightarrow 2^+$
dominant until
 $T \sim 0.7 \times 10^9$ K





Intermediate mass stars are on the verge of electron capture supernovae or disrupting. Forbidden transitions play an unexpected central role in their fate.

Thank you
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For higher forbidden transitions, we have

$$\begin{aligned} M_L(k_x, k_\nu^{(1)}) = & K_L(p_x R)^{k_x-1} (q_x R)^{k_\nu^{(1)}-1} \left\{ -[(2L+1)/L]^{1/2} {}^V F_{LL-11}^0 + (2k_x+1)^{-1/2} \alpha Z {}^V F_{LL0}^0(k_x, 1, 1, 1) \right. \\ & + [(2k_x+1)^{-1}(2k_\nu^{(1)}+1)^{-1} q_x R] {}^V F_{LL0}^0 - (2k_x+1)^{-1} \alpha Z [(L+1)/L]^{1/2} {}^A F_{LL1}^0(k_x, 1, 1, 1) \\ & \left. - [(2k_x+1)^{-1} W_x R + (2k_\nu^{(1)}+1)^{-1} q_x R][(L+1)/L]^{1/2} {}^A F_{LL1}^0 \right\}, \end{aligned} \quad (2.106a)$$

$$m_L(k_x, k_\nu^{(1)}) = K_L(p_x R)^{k_x-1} (q_x R)^{k_\nu^{(1)}-1} (2k_x+1)^{-1} R \left\{ {}^V F_{LL0}^0 - [(L+1)/L]^{1/2} {}^A F_{LL1}^0 \right\}, \quad (2.106b)$$

$$M_L(k_x, k_\nu^{(2)}) = -\tilde{K}_L(p_x R)^{k_x-1} (q_x R)^{k_\nu^{(2)}-1} (L+1)^{1/2} [(2k_x-1)(2k_\nu^{(2)}-1)]^{-1/2} \left\{ {}^V F_{LL0}^0 + (k_x - k_\nu^{(2)})(L+1)^{-1} [(L+1)/L]^{1/2} {}^A F_{LL1}^0 \right\}, \quad (2.106c)$$

$$M_{L+1}(k_x, k_\nu^{(2)}) = -\tilde{K}_L(p_x R)^{k_x-1} (q_x R)^{k_\nu^{(2)}-1} {}^A F_{(L+1)L1}^0. \quad (2.106d)$$

$$\begin{aligned}
{}^V \mathfrak{M}_{KK_0}^N(k_x, m, n, \rho) = & \sqrt{2}(2J_i + 1)^{-1/2} \left\{ G_{KK_0}(\kappa_f, \kappa_i) \int_0^\infty g_f(r, \kappa_f) \left(\frac{r}{R}\right)^{K+2N} I(k_x, m, n, \rho; r) g_i(r, \kappa_i) r^2 dr \right. \\
& \left. + \text{sign}(\kappa_f) \text{sign}(\kappa_i) G_{KK_0}(-\kappa_f, -\kappa_i) \int_0^\infty f_f(r, \kappa_f) \left(\frac{r}{R}\right)^{K+2N} I(k_x, m, n, \rho; r) f_i(r, \kappa_i) r^2 dr \right\},
\end{aligned} \tag{2.84a}$$

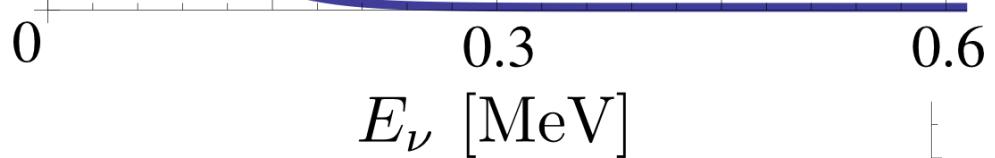
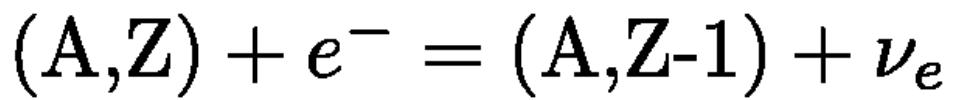
$$\begin{aligned}
{}^A \mathfrak{M}_{KL_1}^N(k_x, m, n, \rho) = & \sqrt{2}(2J_i + 1)^{-1/2} \left\{ G_{KL_1}(\kappa_f, \kappa_i) \int_0^\infty g_f(r, \kappa_f) \left(\frac{r}{R}\right)^{L+2N} I(k_x, m, n, \rho; r) g_i(r, \kappa_i) r^2 dr \right. \\
& \left. + \text{sign}(\kappa_f) \text{sign}(\kappa_i) G_{KL_1}(-\kappa_f, -\kappa_i) \int_0^\infty f_f(r, \kappa_f) \left(\frac{r}{R}\right)^{L+2N} I(k_x, m, n, \rho; r) f_i(r, \kappa_i) r^2 dr \right\},
\end{aligned} \tag{2.84b}$$

$$\begin{aligned}
{}^A \mathfrak{M}_{KK_0}^N(k_x, m, n, \rho) = & \sqrt{2}(2J_i + 1)^{-1/2} \left\{ \text{sign}(\kappa_i) G_{KK_0}(\kappa_f, -\kappa_i) \int_0^\infty g_f(r, \kappa_f) \left(\frac{r}{R}\right)^{K+2N} I(k_x, m, n, \rho; r) f_i(r, \kappa_i) r^2 dr \right. \\
& \left. + \text{sign}(\kappa_f) G_{KK_0}(-\kappa_f, \kappa_i) \int_0^\infty f_f(r, \kappa_f) \left(\frac{r}{R}\right)^{K+2N} I(k_x, m, n, \rho; r) g_i(r, \kappa_i) r^2 dr \right\},
\end{aligned} \tag{2.84c}$$

$$\begin{aligned}
{}^V \mathfrak{M}_{KL_1}^N(k_x, m, n, \rho) = & \sqrt{2}(2J_i + 1)^{-1/2} \left\{ \text{sign}(\kappa_i) G_{KL_1}(\kappa_f, -\kappa_i) \int_0^\infty g_f(r, \kappa_f) \left(\frac{r}{R}\right)^{L+2N} I(k_x, m, n, \rho; r) f_i(r, \kappa_i) r^2 dr \right. \\
& \left. + \text{sign}(\kappa_f) G_{KL_1}(-\kappa_f, \kappa_i) \int_0^\infty f_f(r, \kappa_f) \left(\frac{r}{R}\right)^{L+2N} I(k_x, m, n, \rho; r) g_i(r, \kappa_i) r^2 dr \right\}.
\end{aligned} \tag{2.84d}$$

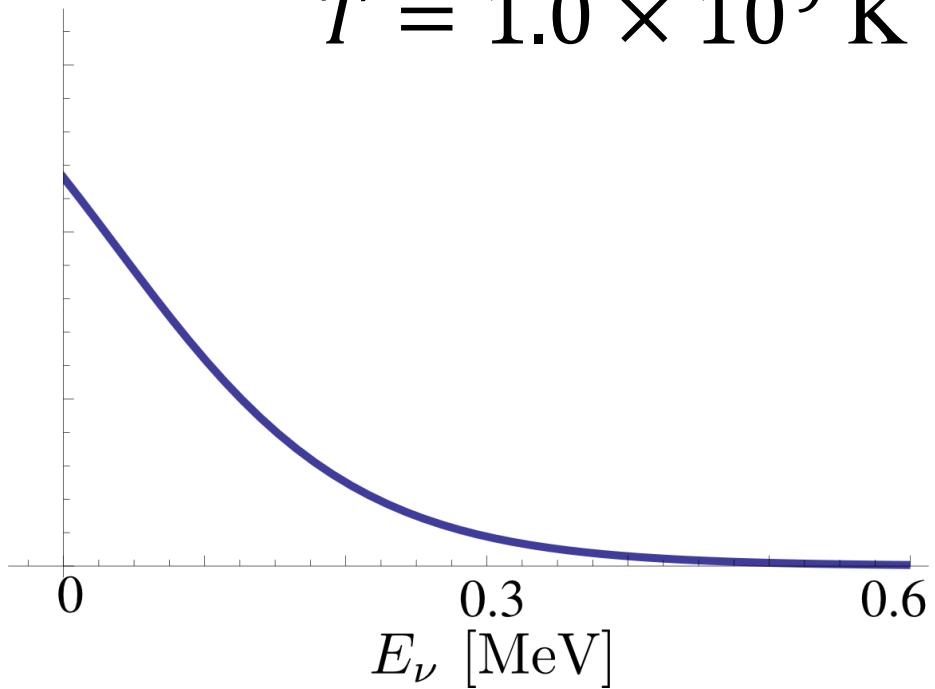
$T = 0.4 \times 10^9$ K

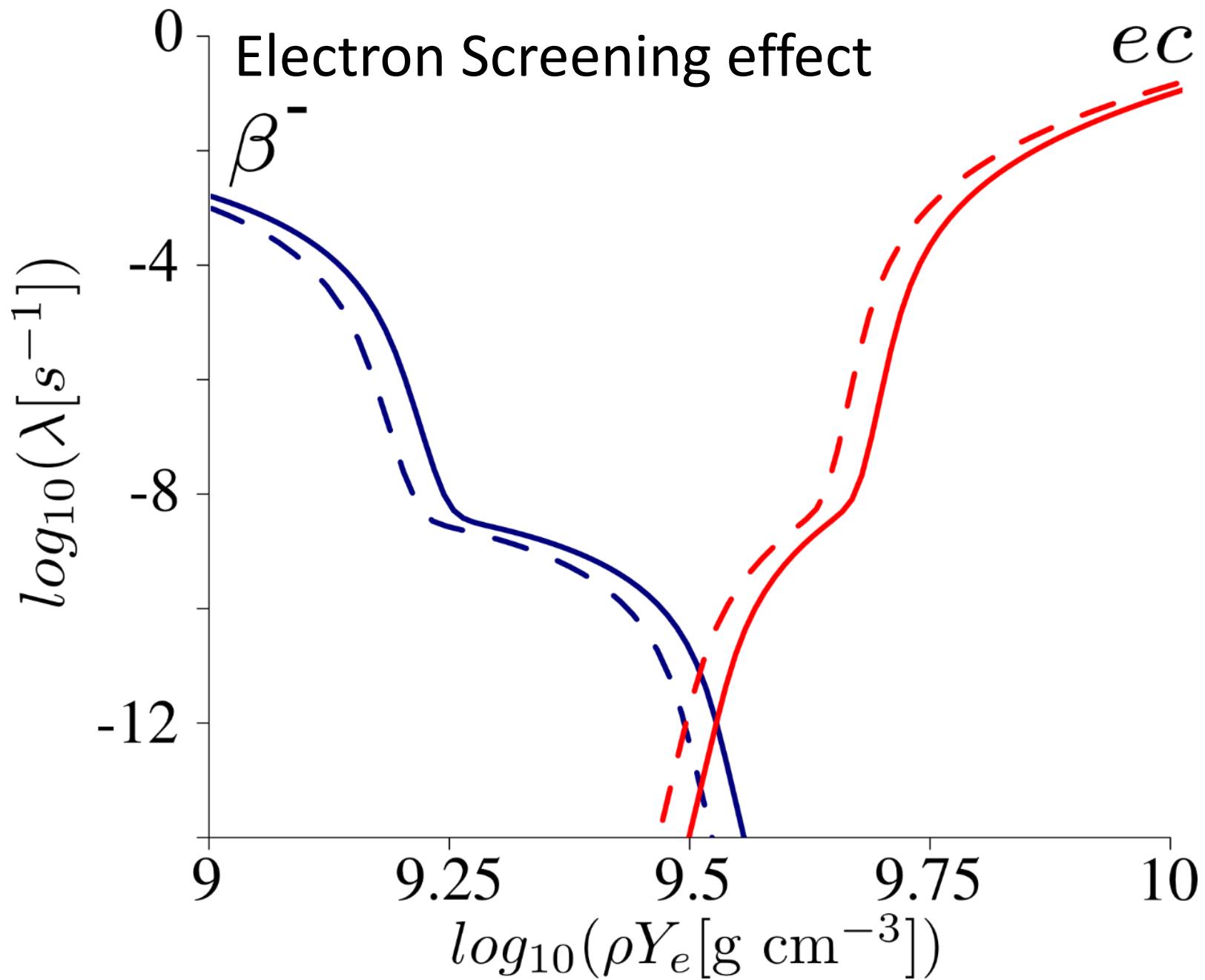
Neutrino EC-Spectrum

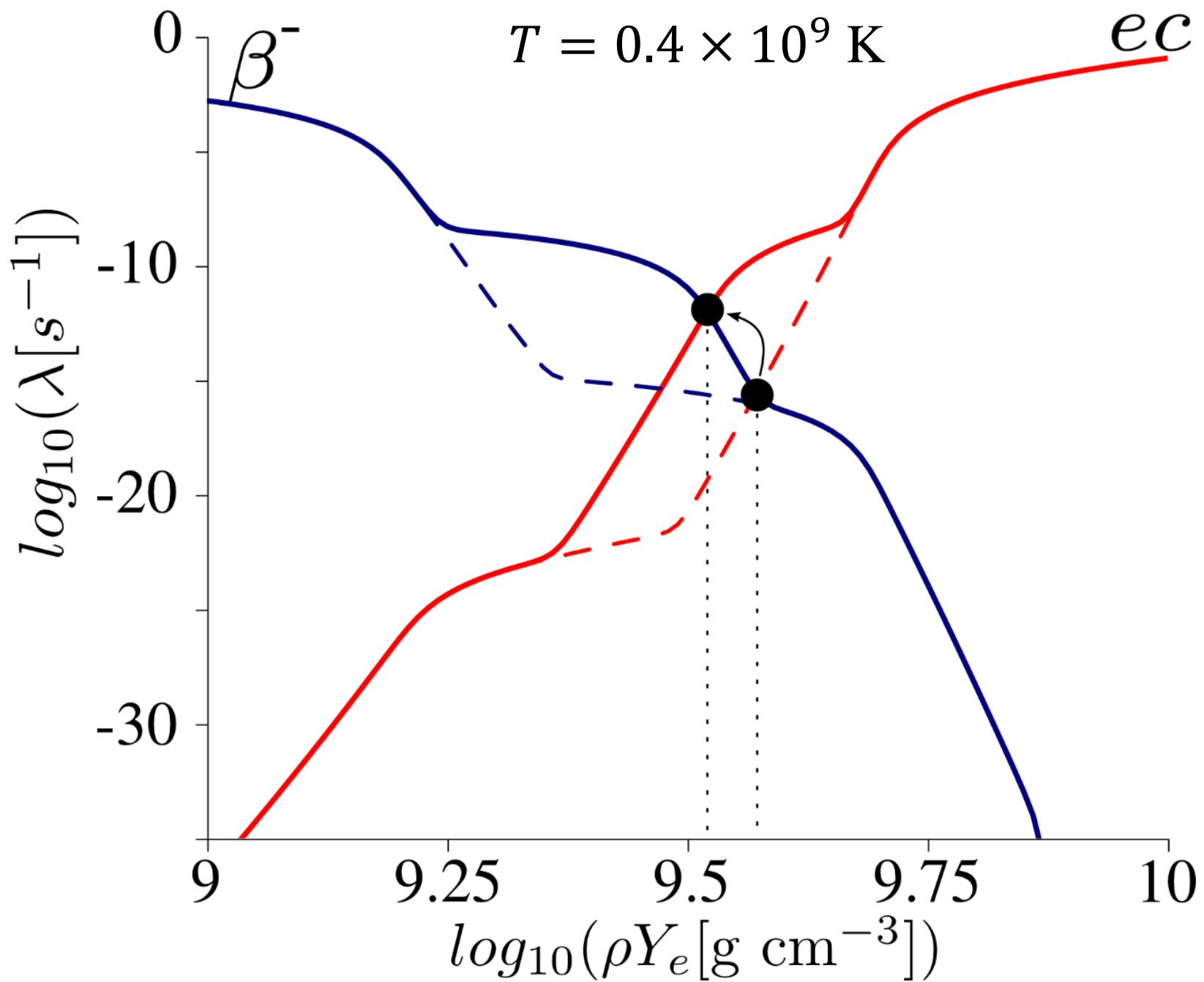


$$\rho Y_e = 10^{9.5} g/cm^3$$

$T = 1.0 \times 10^9$ K







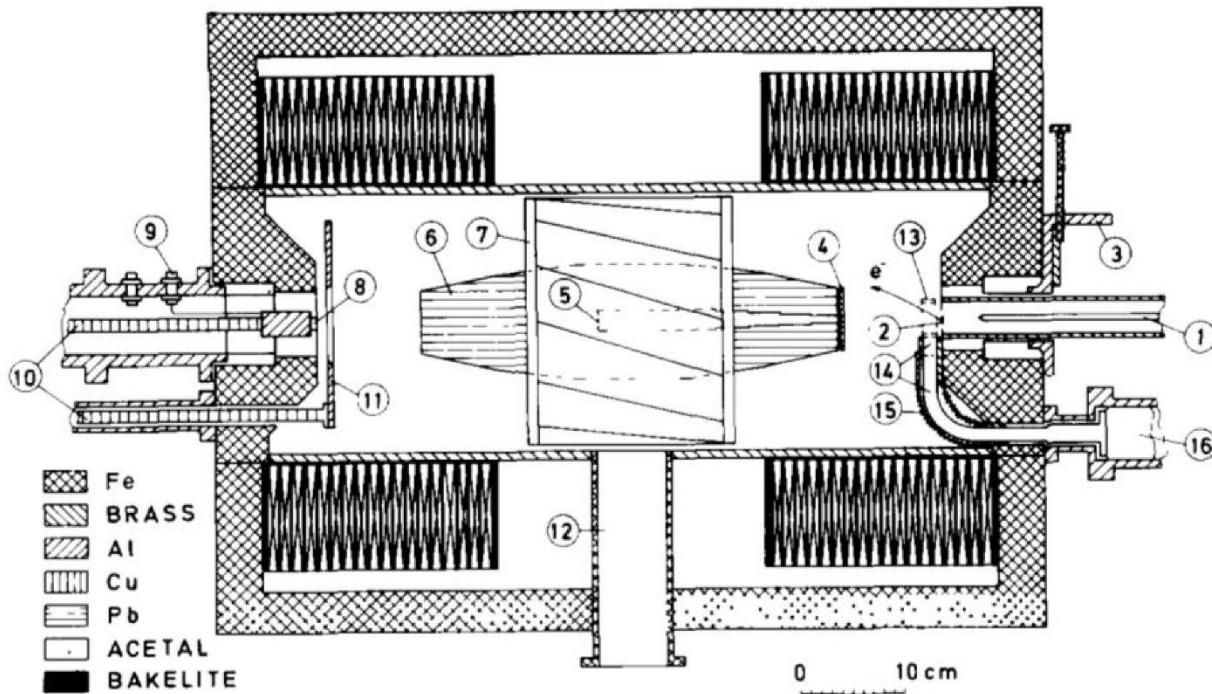


Fig. 1. Magnetic + semiconductor detector conversion-electron spectrometer based on an "old" Siegbahn–Släts magnet used as the electron transporter: (1) beam, (2) target, (3) target-changing system, (4) collimator and current measurement, (5) Faraday cup, (6) lead shield, (7) antiproton baffle, (8) detector, (9) feedthrough, (10) coldfingers, (11) coldtrap, (12) to pump, (13) cylindrical plastic scintillator, (14) light guide, (15) bellows reflector, (16) PM tube.