

A natural Fermi scale

Majorana Lectures
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Riccardo Barbieri
SNS and INFN, Pisa

The Fine Tuning problem of the Fermi scale

$$\delta m_H^2 = \text{---} \circlearrowleft \text{SM} \text{---} + \text{---} \circlearrowleft \text{New} \text{---}$$

in isolation grows like Λ^2
where $\Lambda \equiv$ the "cutoff"
i.e. the highest scale at
which the theory is valid:
a signal of the sensitivity of m_H

knows about Λ_{NP}
i.e. any higher physical scale
to which the Higgs boson
is possibly coupled
 $M_{Pl}, M_{GUT} \dots$ you name it

$$\delta m_H^2 \approx \frac{\lambda^2}{16\pi^2} \Lambda_{NP}^2$$

The ElectroWeak scale determined by the Higgs mass

Why $G_F^{-1/2} \ll \Lambda_{NP}$?

Why $G_F^{-1/2} \ll G_N^{-1/2}$?

Another look at the FT problem

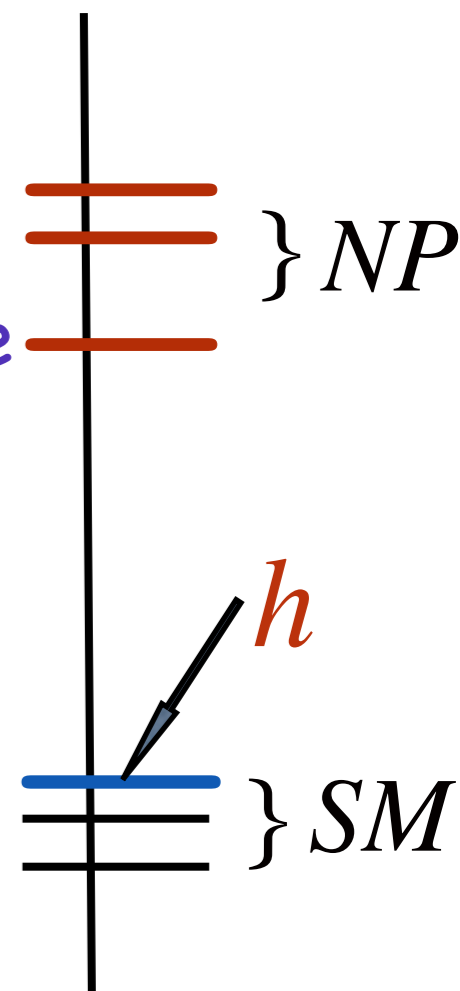
There are infinitely many theories at short distances, that give the same physics of the Standard Model, as long as their “low energy” spectrum coincides with the one of the SM

including the Higgs boson

We only know of approximate symmetries that can explain this

In all explicit examples, without unwarranted cancellations, new phenomena required at a scale Λ_{NP} not far from m_h

⇒ What are and where are they?

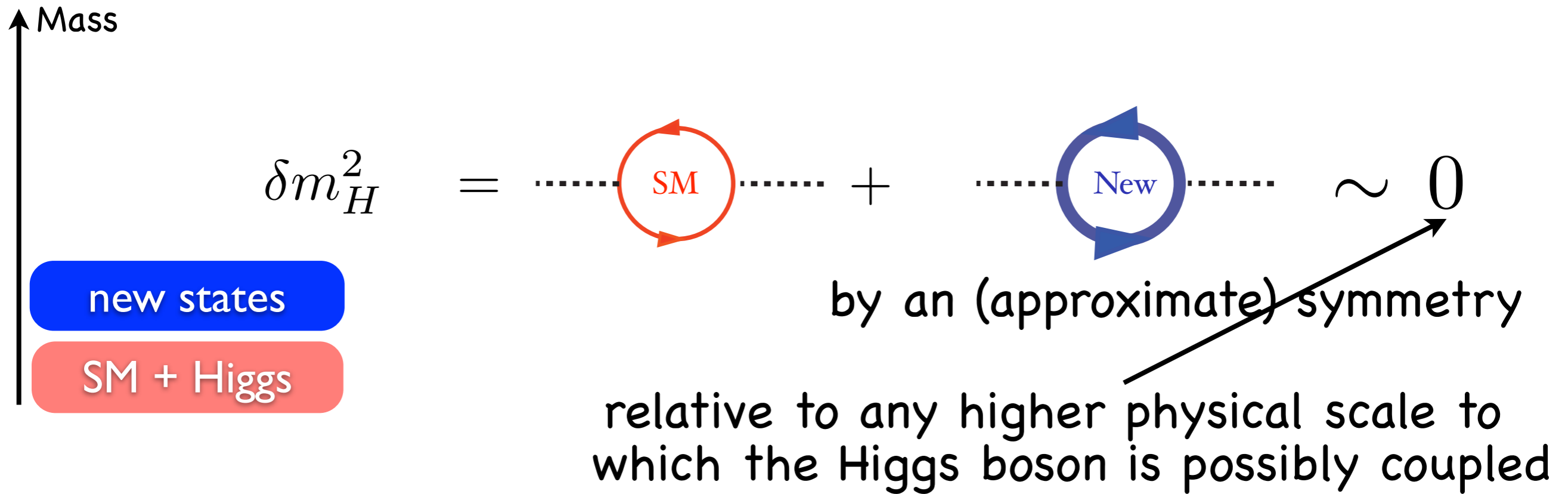


The (many) reactions to the FT problem

0. Ignore it and view the SM in isolation
(if no other short distance scale, what about gravity?)
- ➔ 1. Cure it by symmetries: SUSY, Higgs as PGB
2. A new strong interaction nearby
3. A new strong interaction not so nearby: quasi-CFT
4. Saturate the UV nearby: extra-dimensions around the corner
5. Warp space-time: RS
- ➔ 6. Accept it: anthropic selection, the multiverse,
the 10^{120} vacua of string theory

Anything else?

Natural theories



Symmetries crucial to the empirical success of particle physics so far
Have they exhausted their role?

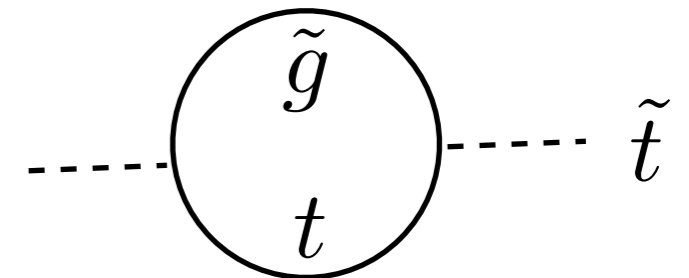
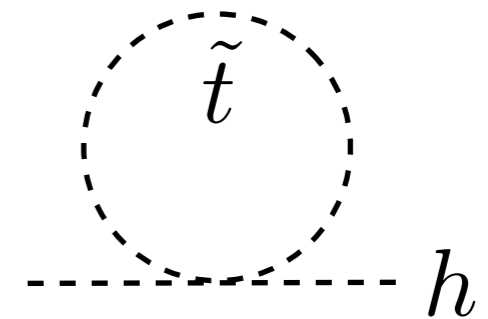
SUSY and its crucial configuration

The key equations:

$$\frac{m_h^2}{2} \approx -|\mu|^2 + m_u^2$$

$$\delta m_u^2 \approx -\frac{3y_t^2}{8\pi^2} (m_{\tilde{t}_L}^2 + m_{\tilde{t}_R}^2 + A_t^2) \log M/m_{\tilde{t}}$$

$$\delta m_{\tilde{t}}^2 \approx \frac{8\alpha_s}{3\pi} m_{\tilde{g}}^2 \log M/m_{\tilde{t}}$$



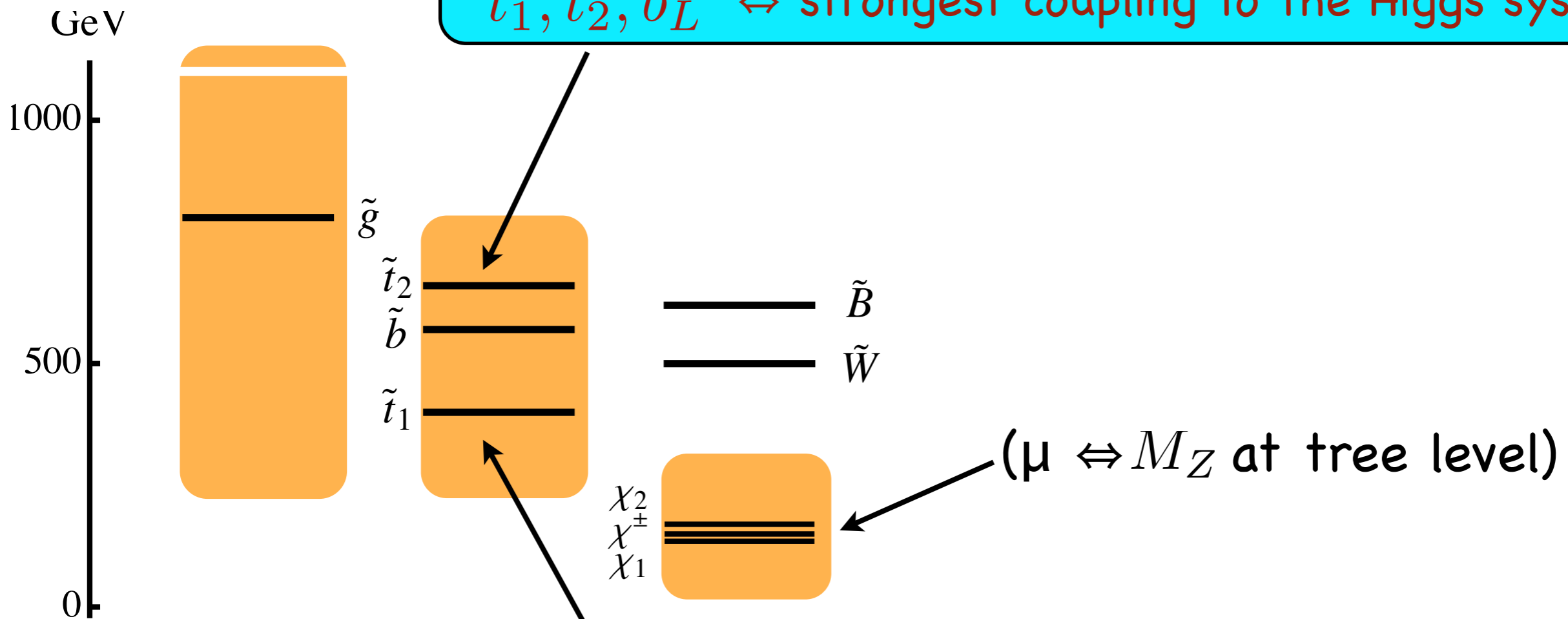
(to be made more precise in any given SB-mediation scheme)

All s-particles other than $\tilde{g}, \tilde{t}_L, \tilde{t}_R, \tilde{b}_L, \tilde{h}$ weakly constrained

The crucial configuration

“s-particles at their naturalness limit”

$\tilde{t}_1, \tilde{t}_2, \tilde{b}_L \Leftrightarrow$ strongest coupling to the Higgs system



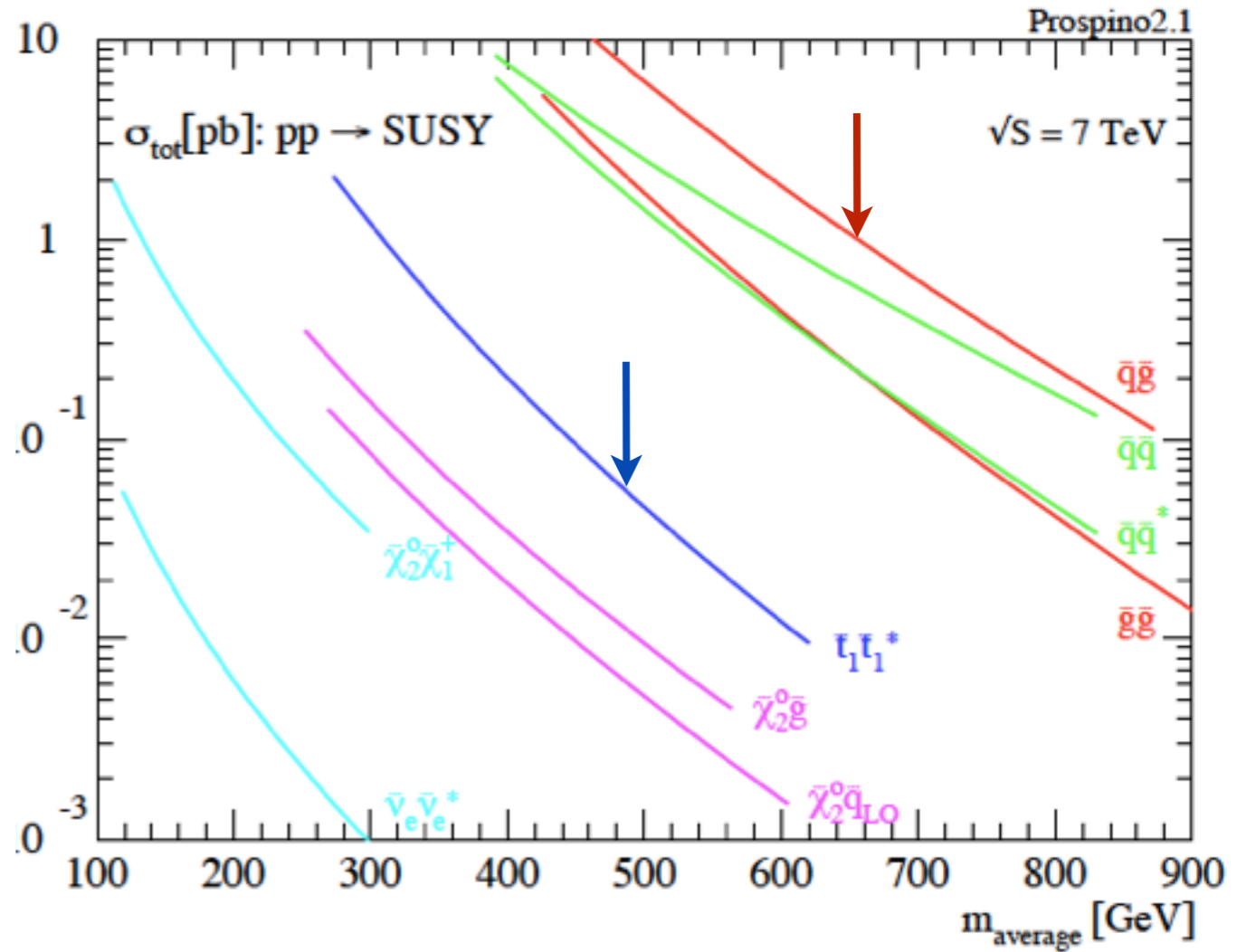
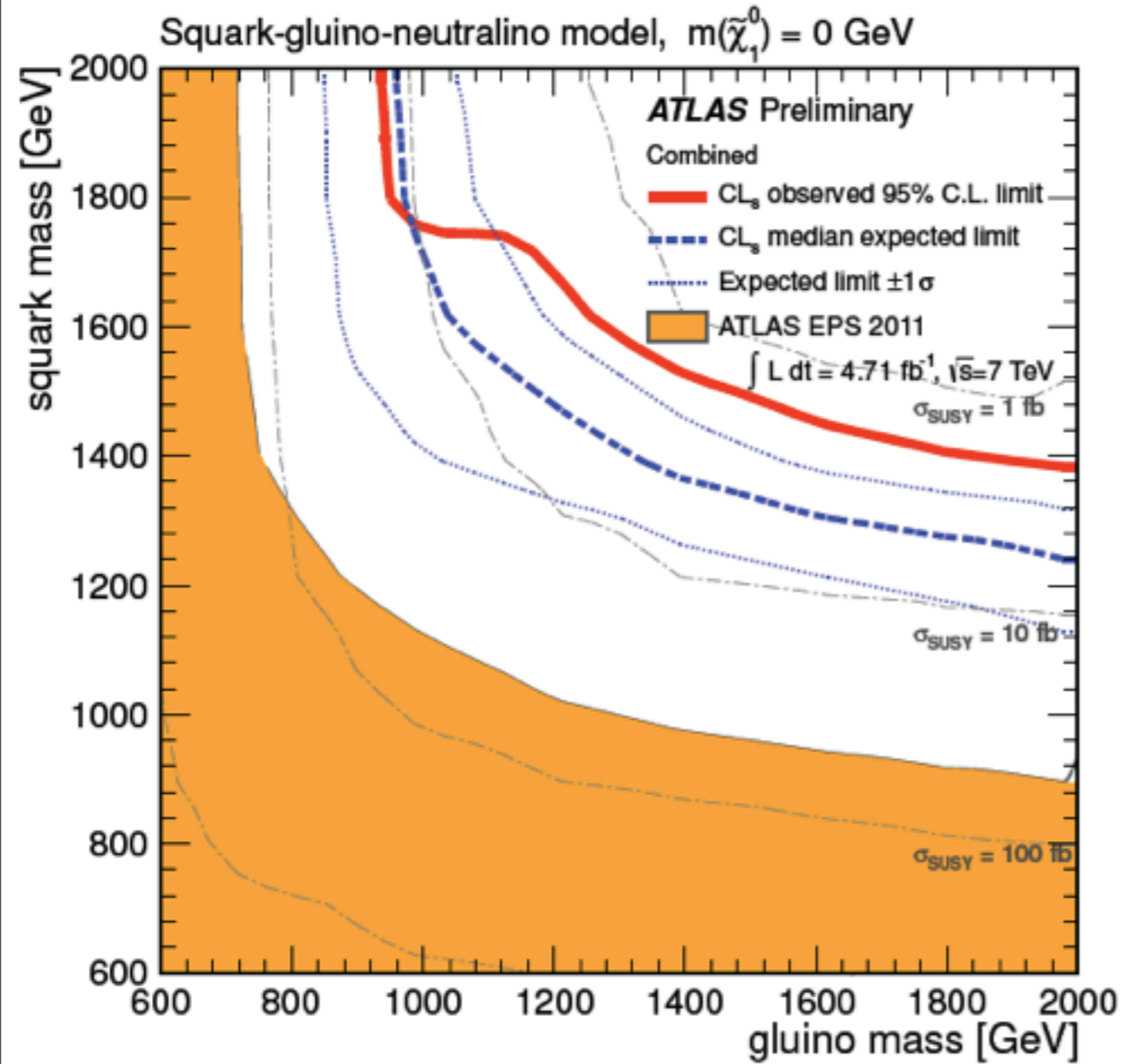
$(\mu \Leftrightarrow M_Z$ at tree level)

$\tilde{q}_1, \tilde{q}_2, \tilde{b}_R$ heavy enough ($\geq \tilde{g}$) to be ~ irrelevant

orange areas indicative and dependent on the Higgs mass

\tilde{B}, \tilde{W} not much constrained but expected below $m_{\tilde{g}}$

A first look at the LHC



$$\begin{aligned}
 g q &\rightarrow \tilde{g} \tilde{q} \\
 q q &\rightarrow \tilde{q} \tilde{q} \\
 q \bar{q} &\rightarrow \tilde{q} \tilde{q}^*
 \end{aligned}$$

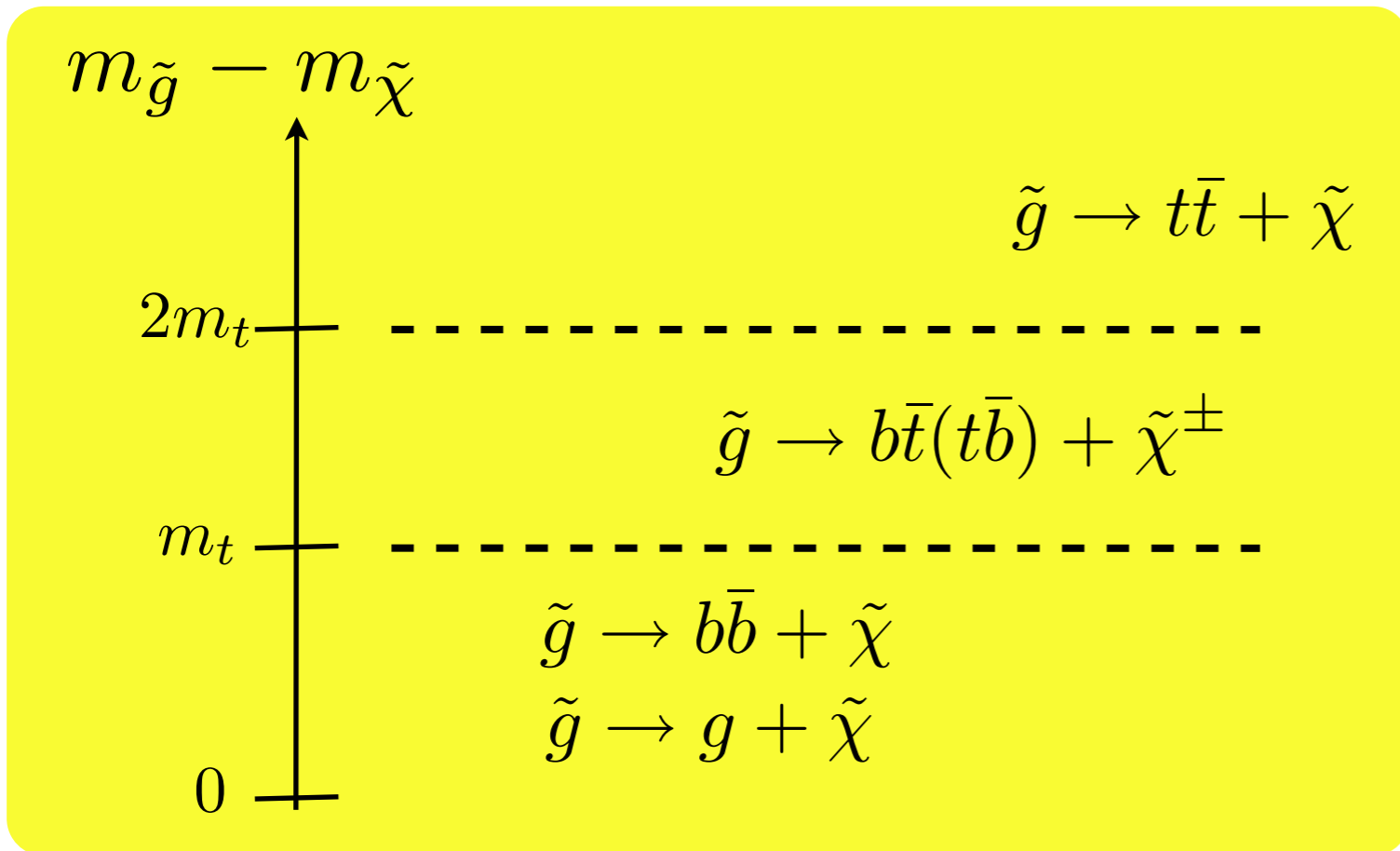
$$\Rightarrow m_{\tilde{g}}, m_{\tilde{q}_{1,2}} > 1 \div 1.5 \text{ TeV}$$

(for χ^0, \tilde{g} not too close in mass)

\tilde{t}, \tilde{b} unconstrained

A synthetic description of the LHC phenomenology

$$pp \rightarrow \tilde{g}\tilde{g} \quad \text{dominant over} \quad pp \rightarrow \tilde{t}\tilde{t}^* \quad (\tilde{b}\tilde{b}^*)$$



$$\begin{aligned}
 pp &\rightarrow \tilde{g}\tilde{g} \rightarrow t\bar{t}t\bar{t} + \chi\chi \\
 pp &\rightarrow \tilde{g}\tilde{g} \rightarrow t\bar{t}b\bar{b}(t\bar{t}t\bar{b}) + \chi\chi \\
 pp &\rightarrow \tilde{g}\tilde{g} \rightarrow t\bar{t}b\bar{b}(t\bar{t}b\bar{b}) + \chi\chi \\
 pp &\rightarrow \tilde{g}\tilde{g} \rightarrow t\bar{t}b\bar{b} + \chi\chi
 \end{aligned}$$

$$\chi = \chi^{\pm}, \chi_1, \chi_2$$

3 body final states either by cascade or direct
 ($m_{\tilde{t}}, m_{\tilde{b}}$ almost don't matter)

When phase space opens up, $\tilde{g} \rightarrow b\bar{b}\chi$ suppressed

If $\mu < M_1, M_2$ then χ^{\pm}, χ^0 close in mass

current bounds on $\tilde{g}, \tilde{t}, \tilde{b}$

ATLAS, $\int L dt \approx 2 fb^{-1}$

$\tilde{b}_1 \tilde{b}_1^*$ (MSSM)	$\tilde{b}_1 \rightarrow b \tilde{\chi}_1^0$	$m_{\tilde{b}_1} = 390 \text{ GeV} (m_{\tilde{\chi}_1^0} = 0)$	2 b -jets
$\tilde{b}_1 \tilde{b}_1^*$ (MSSM)	$\tilde{b}_1 \rightarrow b \tilde{\chi}_1^0$	$m_{\tilde{b}_1} = 350 \text{ GeV} (m_{\tilde{\chi}_1^0} = 120 \text{ GeV})$	2 b -jets
$\tilde{g} \tilde{g}, \tilde{b}_1 \tilde{b}_1^*$ (MSSM)	$\tilde{g} \rightarrow \tilde{b}_1 b, \tilde{b}_1 \rightarrow b \tilde{\chi}_1^0$	$m_{\tilde{g}} = 920 \text{ GeV} (m_{\tilde{b}_1} < 800 \text{ GeV})$	0 ℓ + b -jets
$\tilde{g} \tilde{g}$ (simpl. model)	$\tilde{g} \rightarrow \bar{b} \tilde{\chi}_1^0$	$m_{\tilde{g}} = 900 \text{ GeV} (m_{\tilde{\chi}_1^0} < 300 \text{ GeV})$	0 ℓ + b -jets
$\tilde{g} \tilde{g}, \tilde{t}_1 \tilde{t}_1^*$ (MSSM)	$\tilde{g} \rightarrow \tilde{t}_1 t, \tilde{t}_1 \rightarrow t \tilde{\chi}_1^0$	$m_{\tilde{g}} = 620 \text{ GeV} (m_{\tilde{t}_1} < 440 \text{ GeV})$	1 ℓ + b -jets
$\tilde{g} \tilde{g}, \tilde{t}_1 \tilde{t}_1^*$ (MSSM)	$\tilde{g} \rightarrow \tilde{t}_1 t, \tilde{t}_1 \rightarrow t \tilde{\chi}_1^0$	$m_{\tilde{g}} = 650 \text{ GeV} (m_{\tilde{t}_1} < 450 \text{ GeV})$	2 ℓ SS
$\tilde{g} \tilde{g}$ (simpl. model)	$\tilde{g} \rightarrow t \bar{t} \tilde{\chi}_1^0$	$m_{\tilde{g}} = 700 \text{ GeV} (m_{\tilde{\chi}_1^0} < 100 \text{ GeV})$	1 ℓ + b -jets
$\tilde{g} \tilde{g}$ (simpl. model)	$\tilde{g} \rightarrow t \bar{t} \tilde{\chi}_1^0$	$m_{\tilde{g}} = 650 \text{ GeV} (m_{\tilde{\chi}_1^0} < 215 \text{ GeV})$	2 ℓ SS
$\tilde{g} \tilde{g}$ (simpl. model)	$\tilde{g} \rightarrow tb + \tilde{\chi}_1^0$	$m_{\tilde{g}} = 710 \text{ GeV} (m_{\tilde{\chi}_1^0} < 100 \text{ GeV})$	1 ℓ + b -jets

My rough summary:

$$m_{\tilde{g}} \gtrsim 700 \text{ GeV} \quad m_{\tilde{b}} \gtrsim 350 \text{ GeV} \quad m_{\tilde{t}} \gtrsim 200 \text{ GeV}$$

(from Tevatron searches)

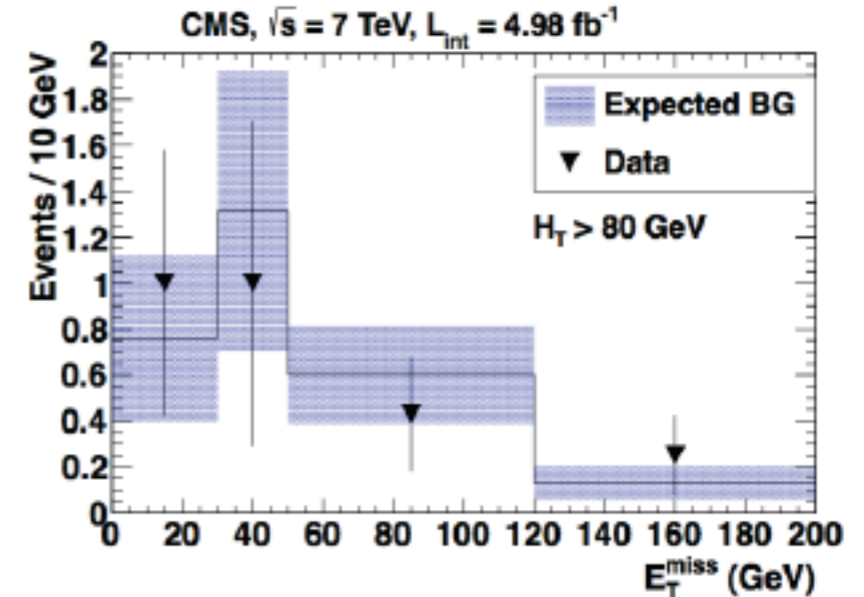
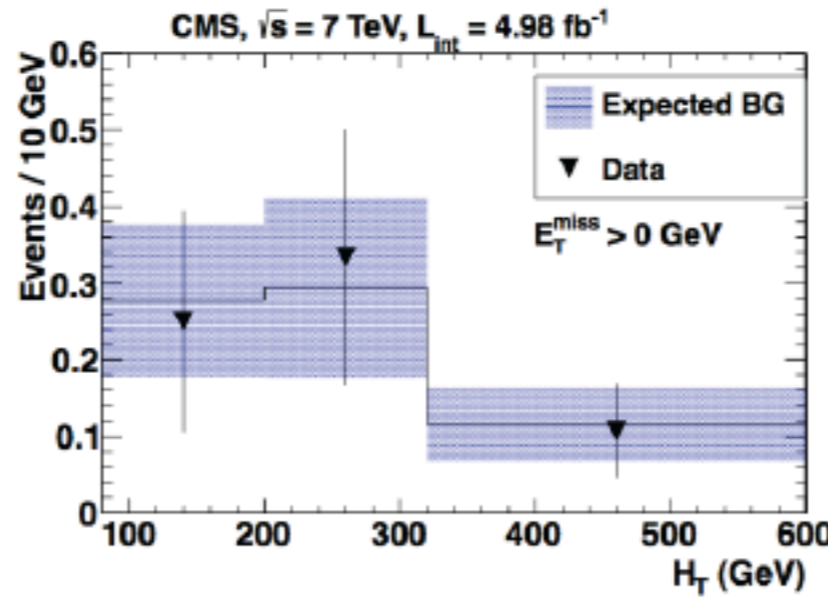
CMS in progress on 3d generation squarks

Kovalskyi

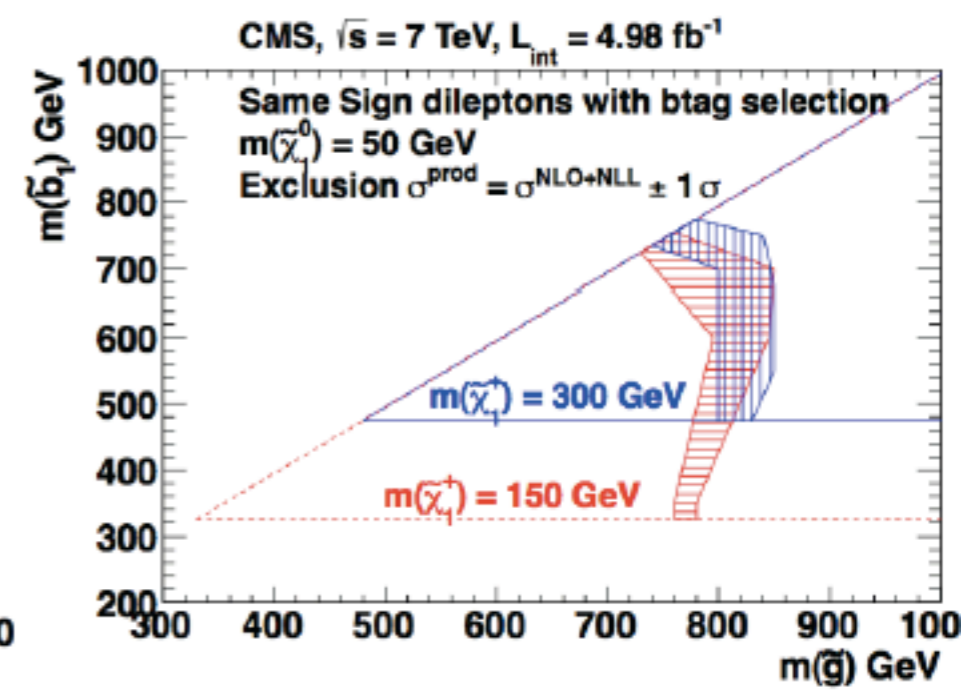
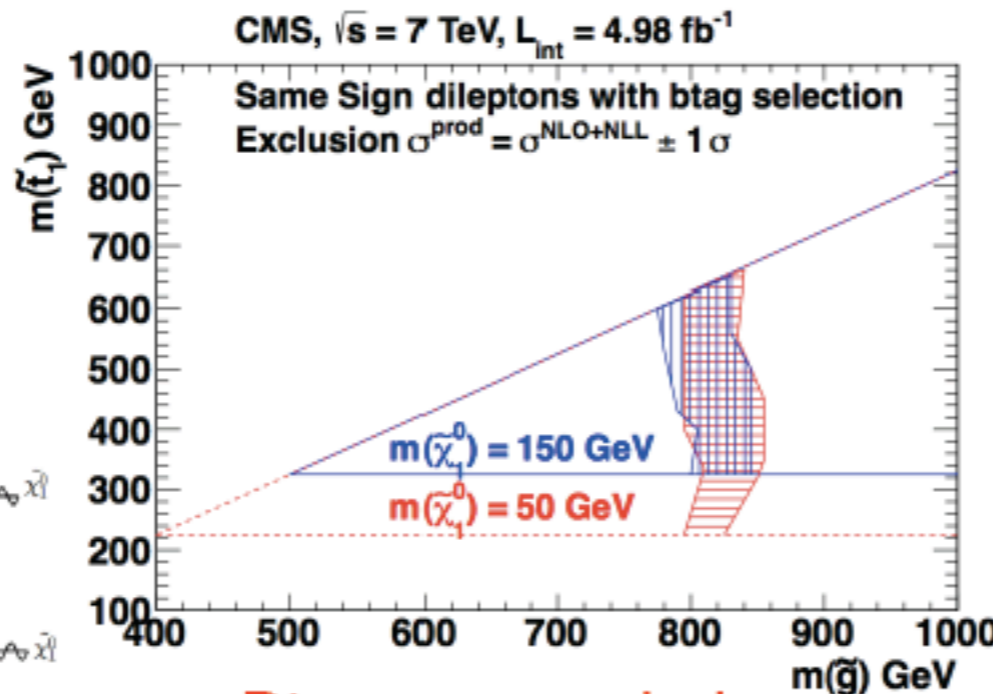
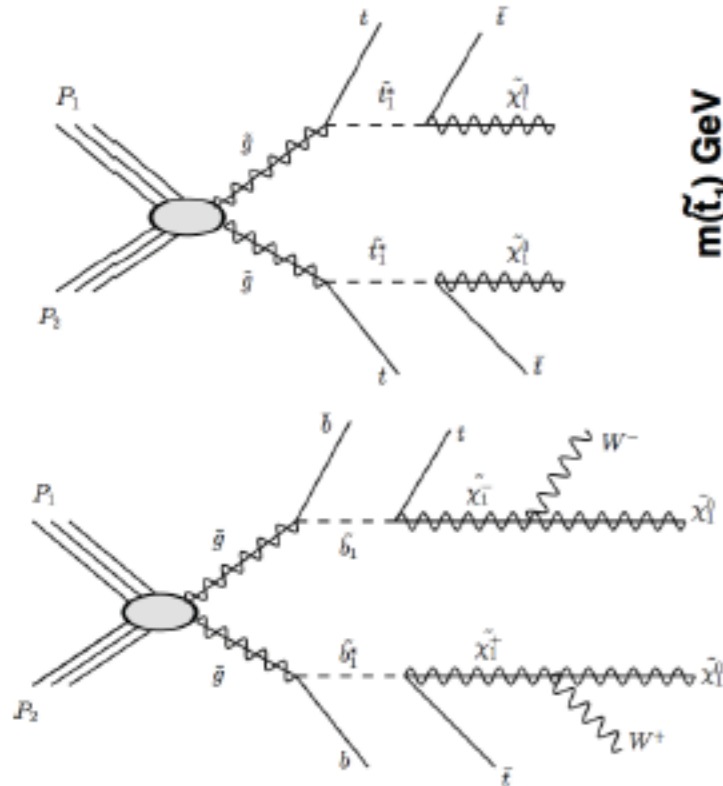
Same-sign leptons + MET + b-jets

For SUSY to remain natural with Higgs $\sim 120\text{GeV}$

- ▶ s-top/bottom $< 400\text{ GeV}$
- ▶ gluino $< 1500\text{ GeV}$



Gluon mitigated s-top and s-bottom production



Direct s-top and s-bottom production searches are in the pipeline

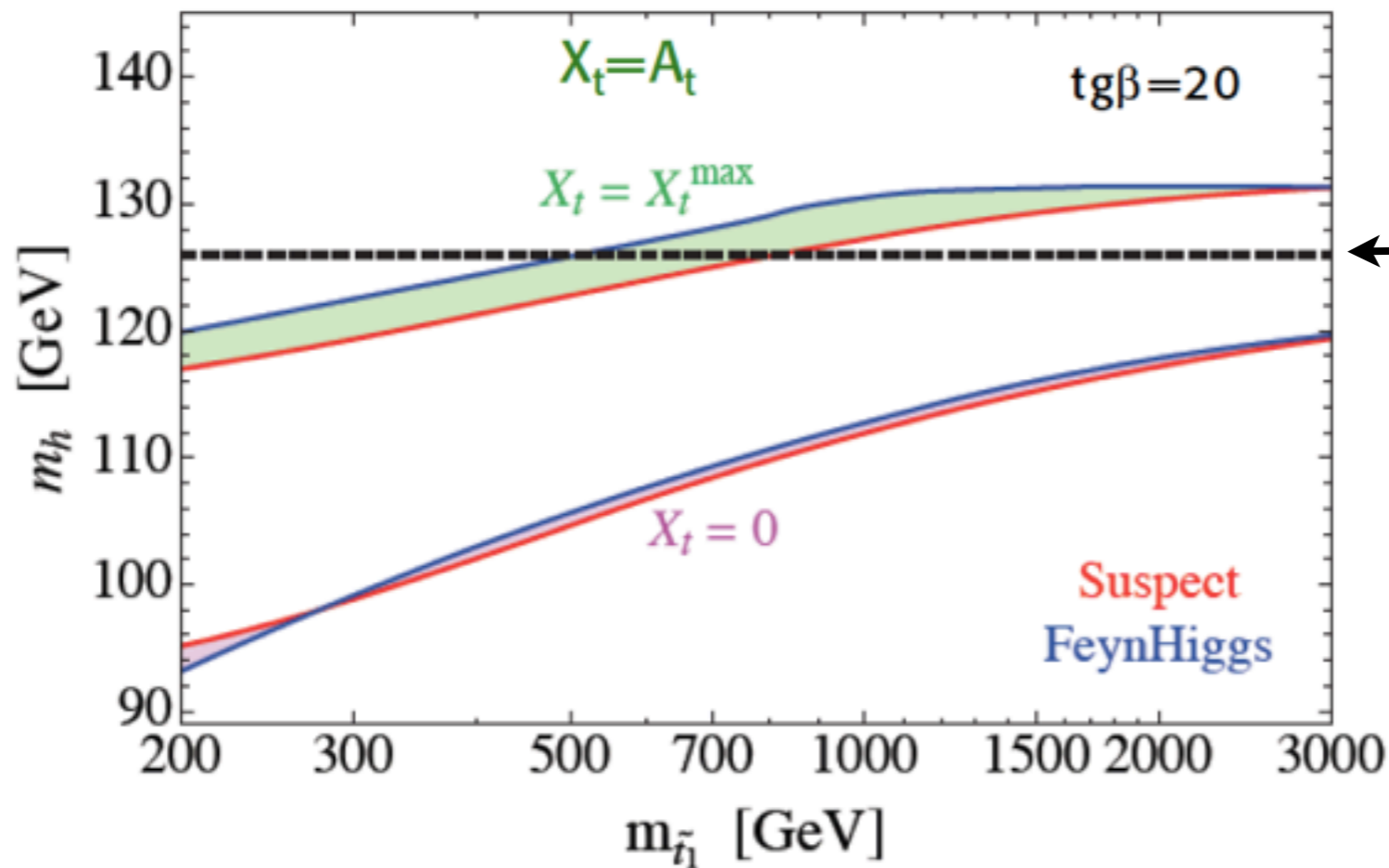
Where is the supersymmetric Higgs boson?

MSSM

$$m_h^2 \approx m_Z^2 \cos^2 2\beta + \frac{3}{(4\pi)^2} \frac{m_t^4}{v^2} \left[\ln \frac{m_{\tilde{t}}^2}{m_t^2} + \frac{X_t^2}{m_{\tilde{t}}^2} \left(1 - \frac{X_t^2}{12m_{\tilde{t}}^2} \right) \right]$$



MSSM Higgs Mass



putative LHC discovery

Hall, Piller, Ruderman 2011

MSSM expectation well OK with current allowed "SM range" but 125 GeV too high for naturally light stops

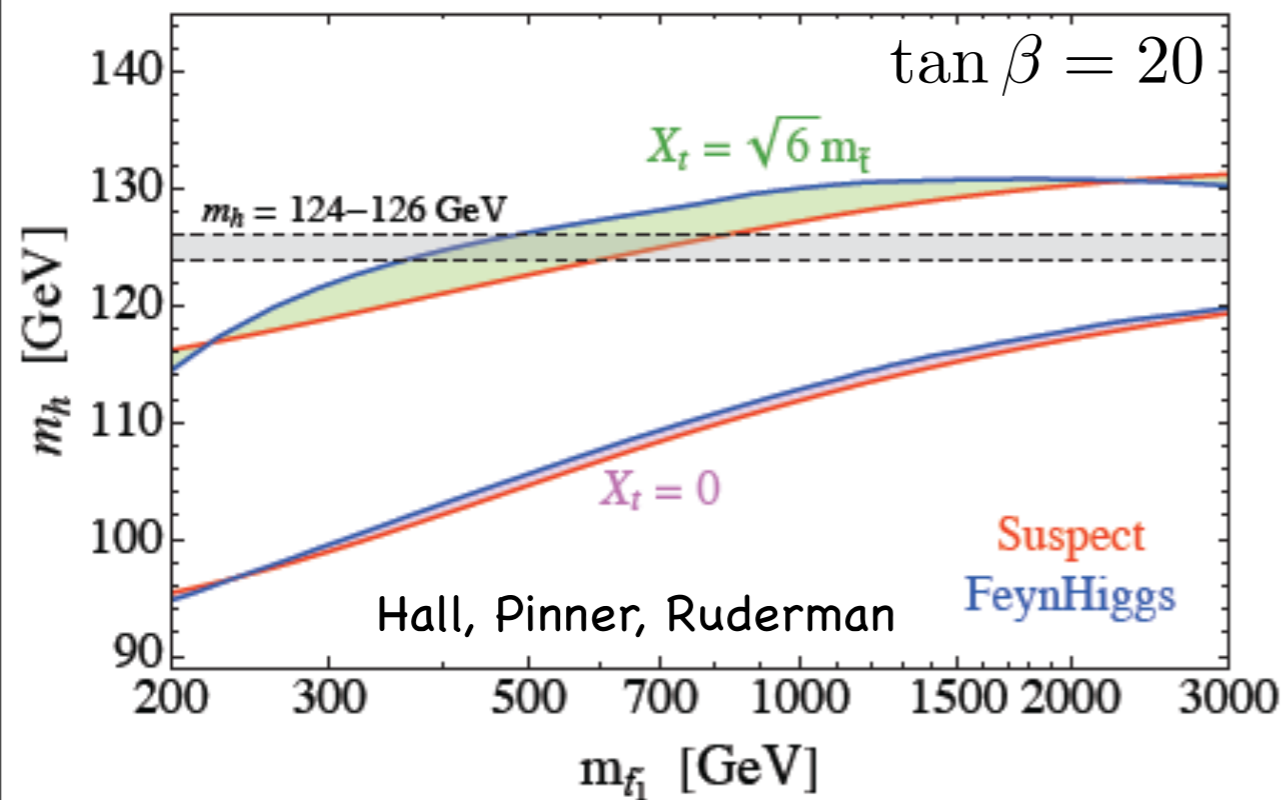
What about the Higgs mass?

The NMSSM as a possible way out

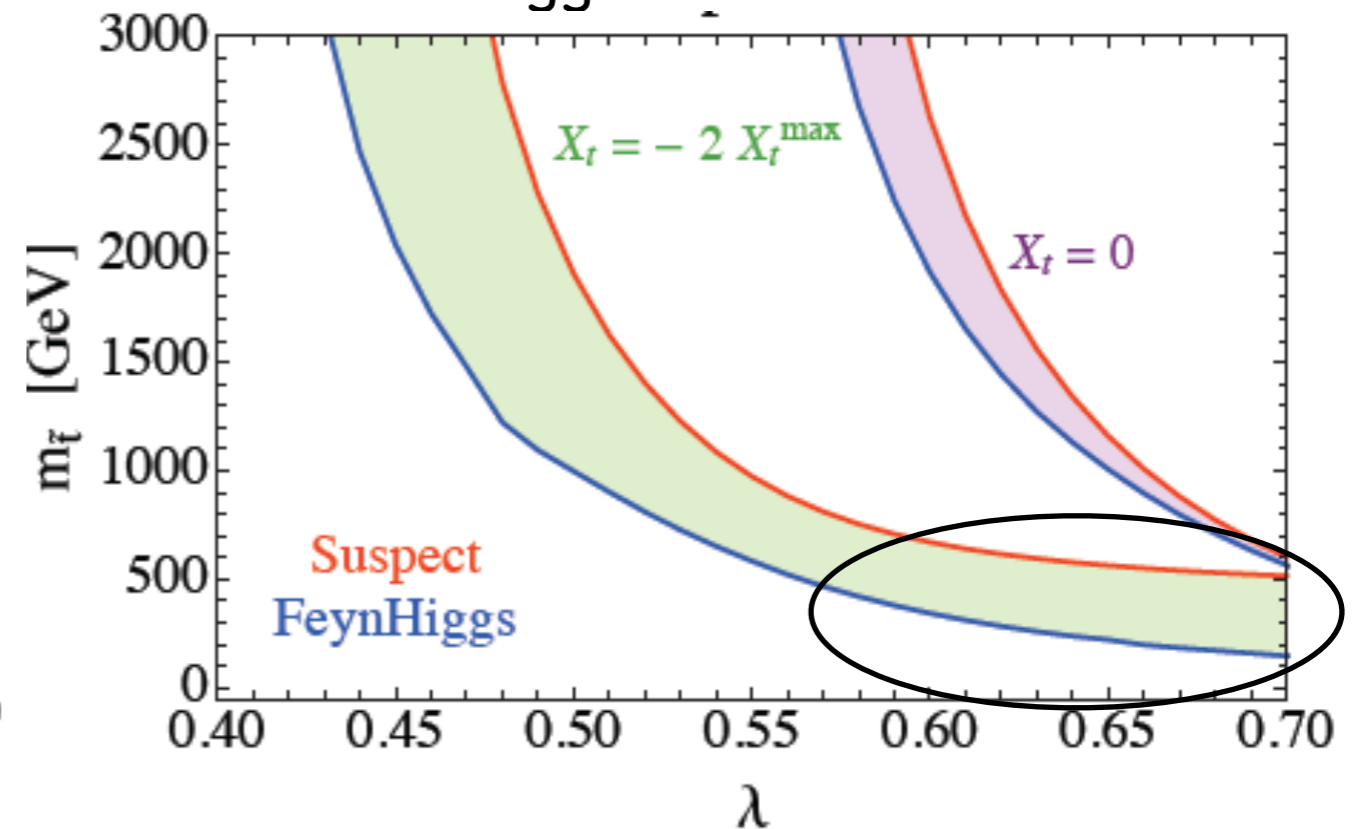
(since the 70's, Fayet, etc)

$$f = \mu H_1 H_2 \Rightarrow f = \lambda S H_1 H_2 \quad m_h^2 \leq m_Z^2 (\cos^2 2\beta + \frac{2\lambda^2}{g^2 + g'^2} \sin^2 2\beta) + \delta_t(m_{\tilde{t}_1}, X_t)$$

MSSM Higgs Mass



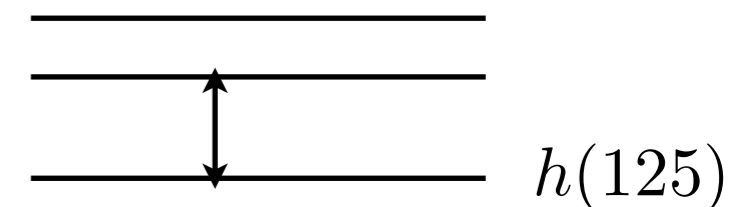
NMSSM Higgs Mass at 124-126 GeV



before mixing with the other 2 scalars

\Rightarrow 2 options:

$h(125)$



A pessimistic reaction

The SUSY scale, M_S , and the Fermi scale, $G_F^{-1/2}$,
not so tied together as we thought

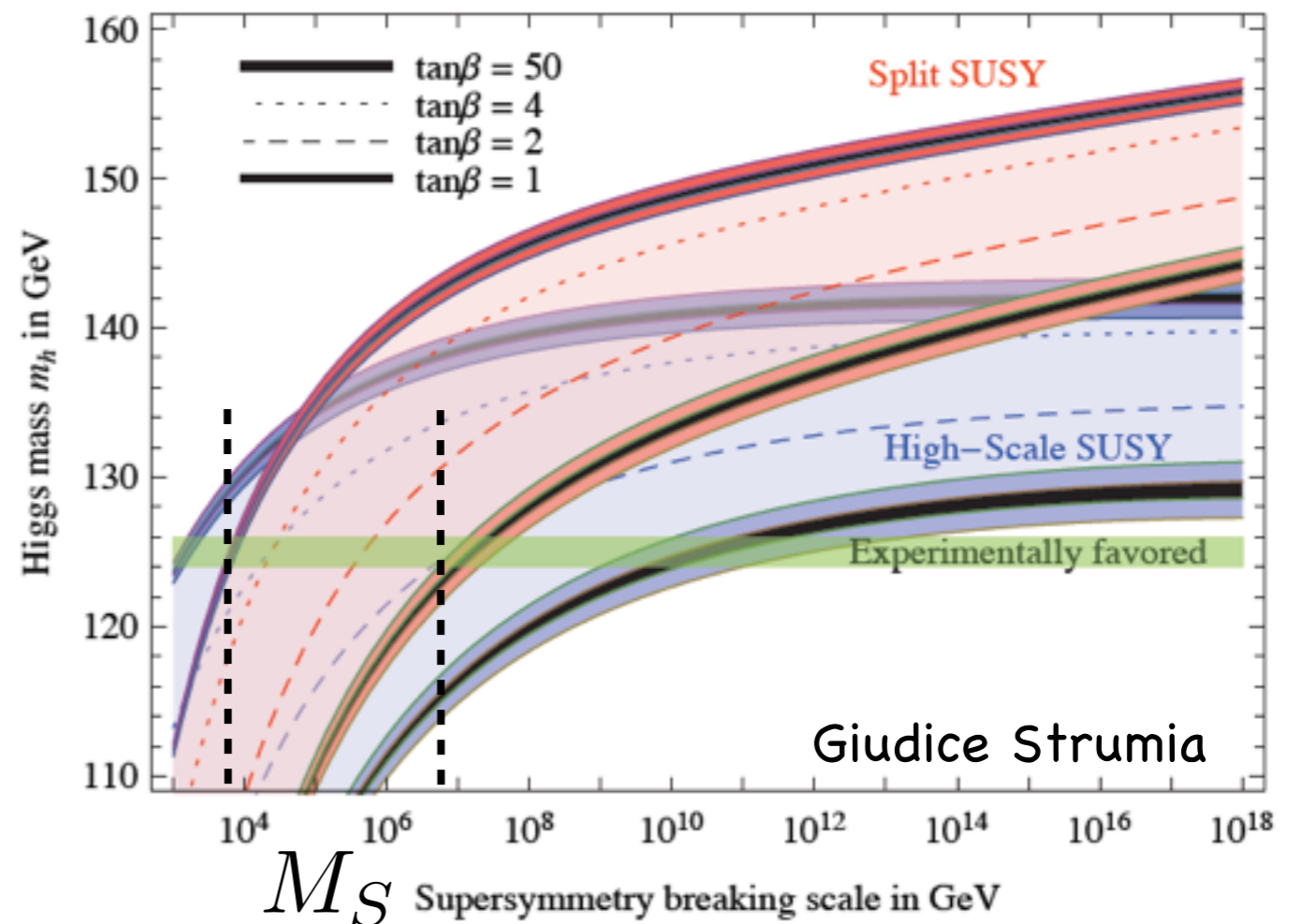
⇒ flavour physics, CPV as in SM

⇒ no SUSY at LHC so far

An “extreme” example:

Split-SUSY ≡
SUSY scalars at $\sim M_S$
SUSY fermions at $\sim G_F^{-1/2}$

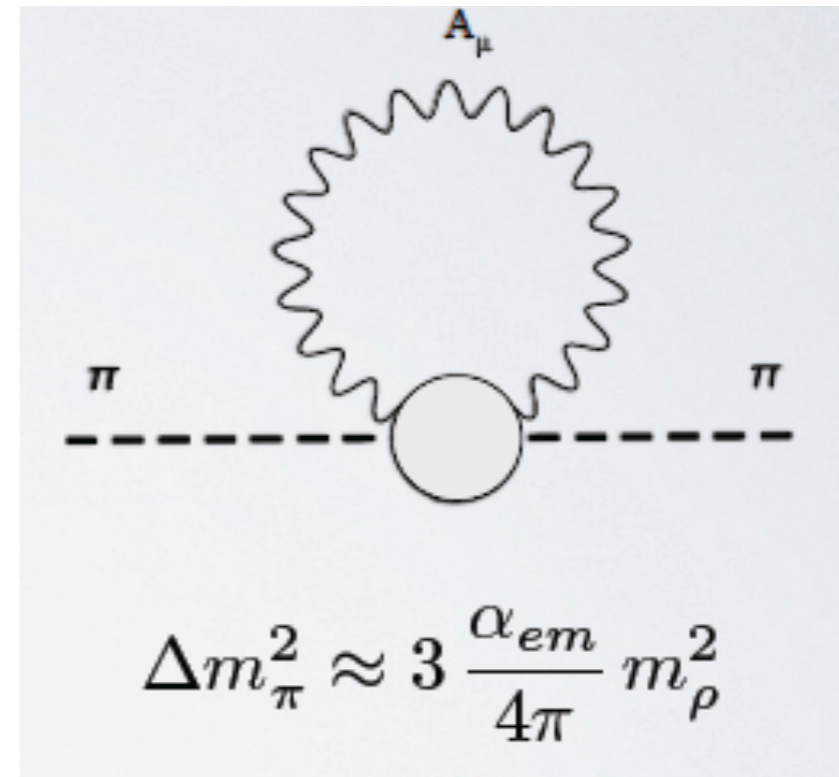
(“High-scale” SUSY ≡
all s-particles at M_S)



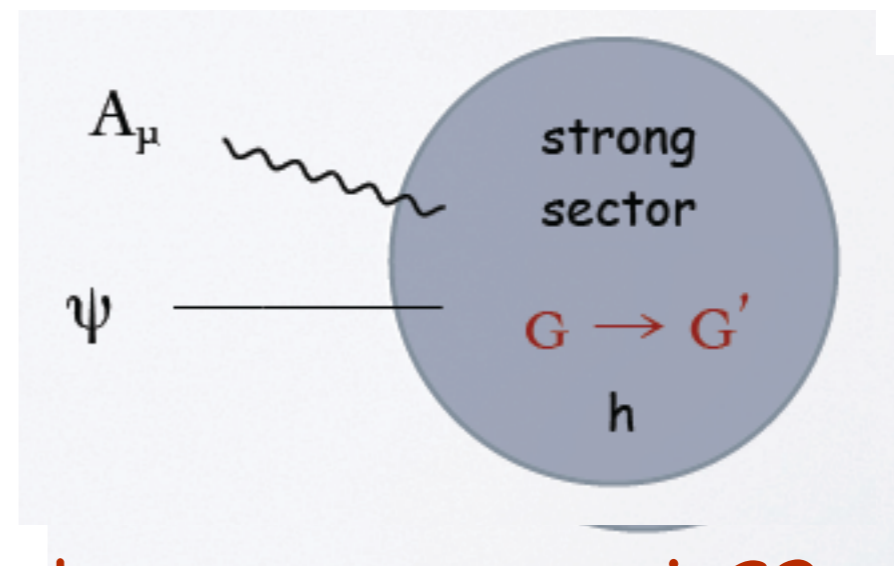
The Higgs boson as a PGB

The pion as an analogy:

$$SU(2)_L \times SU(2)_R \Rightarrow SU(2)_I$$
$$\Delta m_\pi^2 = m_{\pi^+}^2 - m_{\pi^0}^2$$



A new strong sector
at the TeV scale



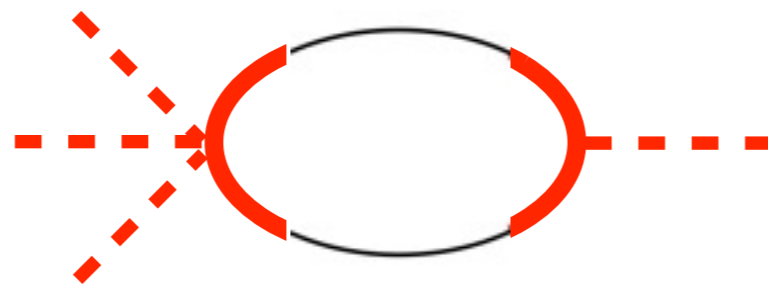
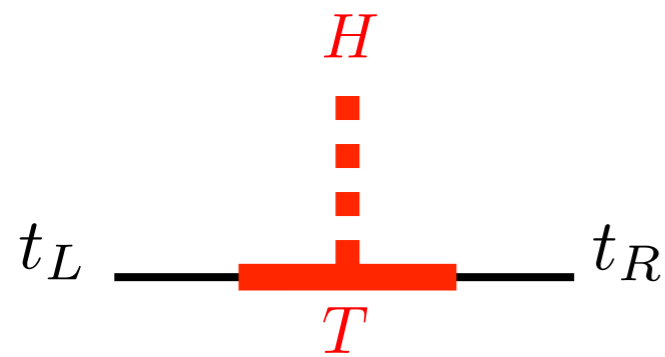
Like the pion in QCD, the Higgs boson as a quasi GB
of a spontaneously broken global symmetry

More in detail

$$\delta m_H^2 = \dots \text{SM} \dots + \dots \text{New} \dots \sim 0$$

Heavy "composite" fermions

Heavy "composite" fermion mass



$$m_h \sim m_t \frac{M_T}{\pi f}$$

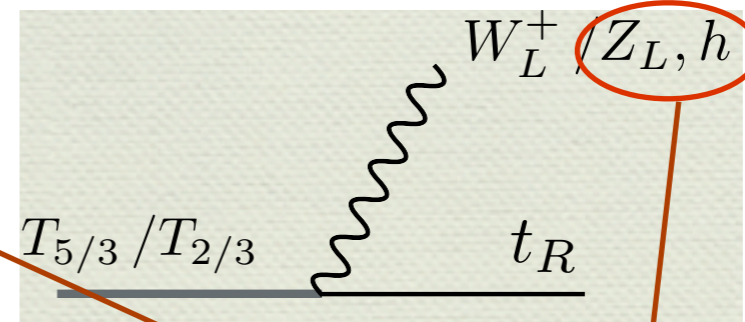
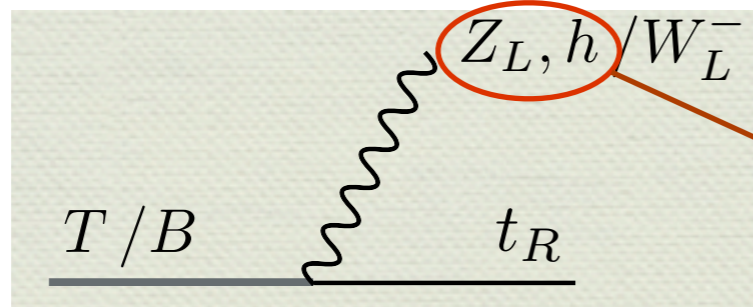
symmetry breaking scale

Most common:

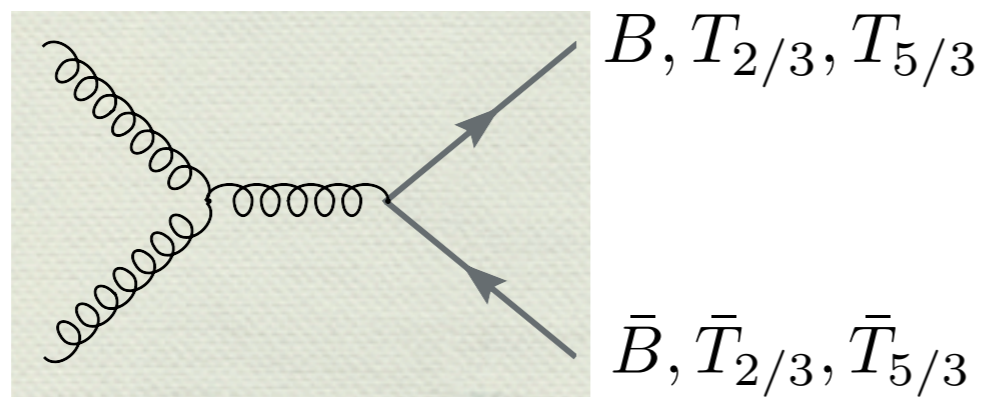
$$(2, 2)_{2/3} = \left[Q = \begin{pmatrix} T \\ B \end{pmatrix} \quad Q' = \begin{pmatrix} T_{5/3} \\ T_{2/3} \end{pmatrix} \right]$$

Phenomenology of the "composite" fermions

Heavy-light couplings

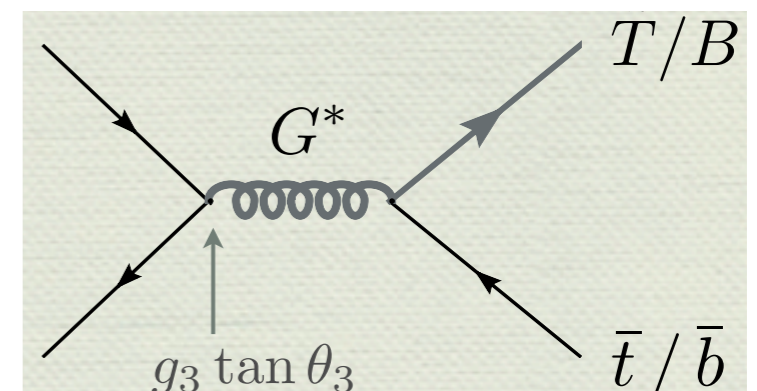
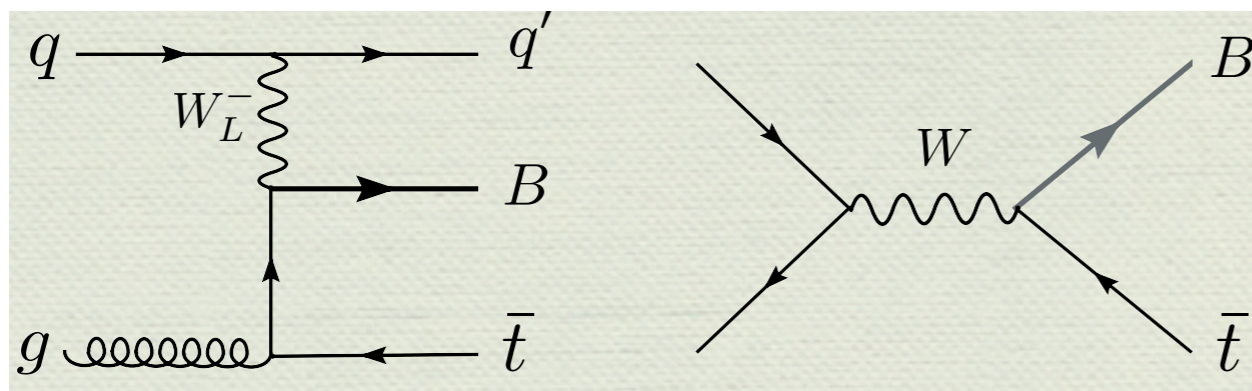


Pair production



Not like a sequential family

Single production



(Some of) the existing limits

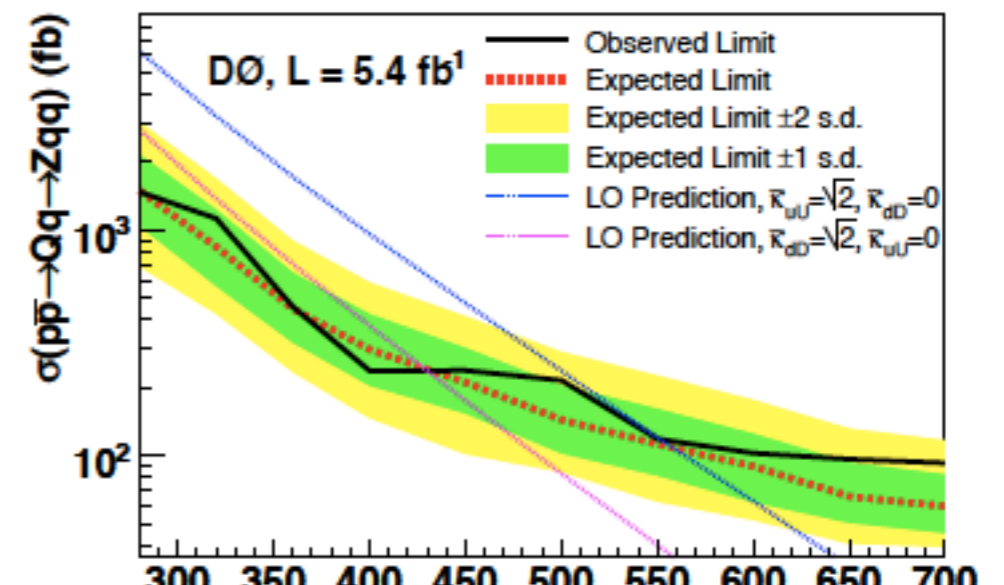
Contino, Servant 2008

Pair production

[CMS L=1.14 fb ⁻¹ PAS-EXO-11-050	$T\bar{T} \rightarrow WbW\bar{b} \rightarrow b\bar{b}l^+l^-\cancel{E}_T$	$m_T > 422 \text{ GeV}$
[CMS L=0.80 fb ⁻¹ PAS-EXO-11-051	$T\bar{T} \rightarrow WbW\bar{b} \rightarrow b3jl^\pm\cancel{E}_T$	$m_T > 450 \text{ GeV}$
[CMS L=191 pb ⁻¹ PAS-EXO-11-005	$T\bar{T} \rightarrow tZ\bar{t}Z \rightarrow (l^+l^-)l^\pm jj$	$m_T > 417 \text{ GeV}$
[CMS L=1.14 fb ⁻¹ PAS-EXO-11-036	$B\bar{B} \rightarrow WtW\bar{t} \rightarrow l^\pm l^\pm b3j\cancel{E}_T$ $\rightarrow lll b1j\cancel{E}_T$	$m_B > 495 \text{ GeV}$

Single production

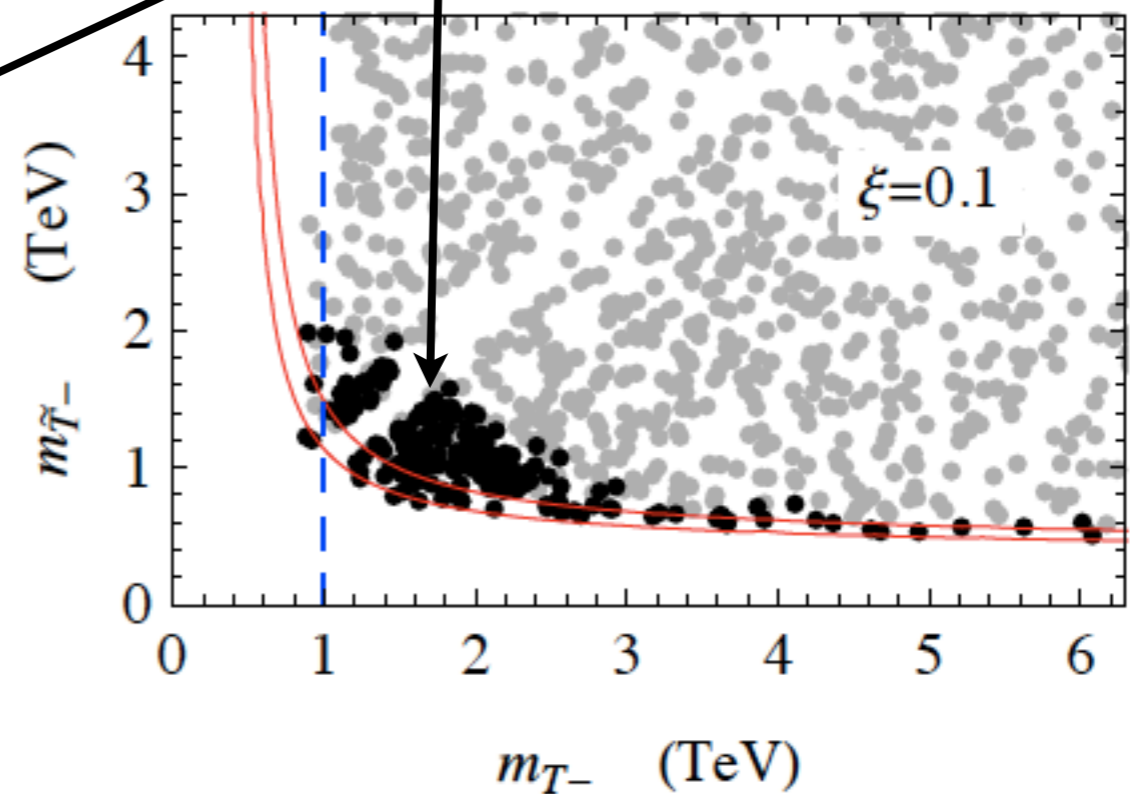
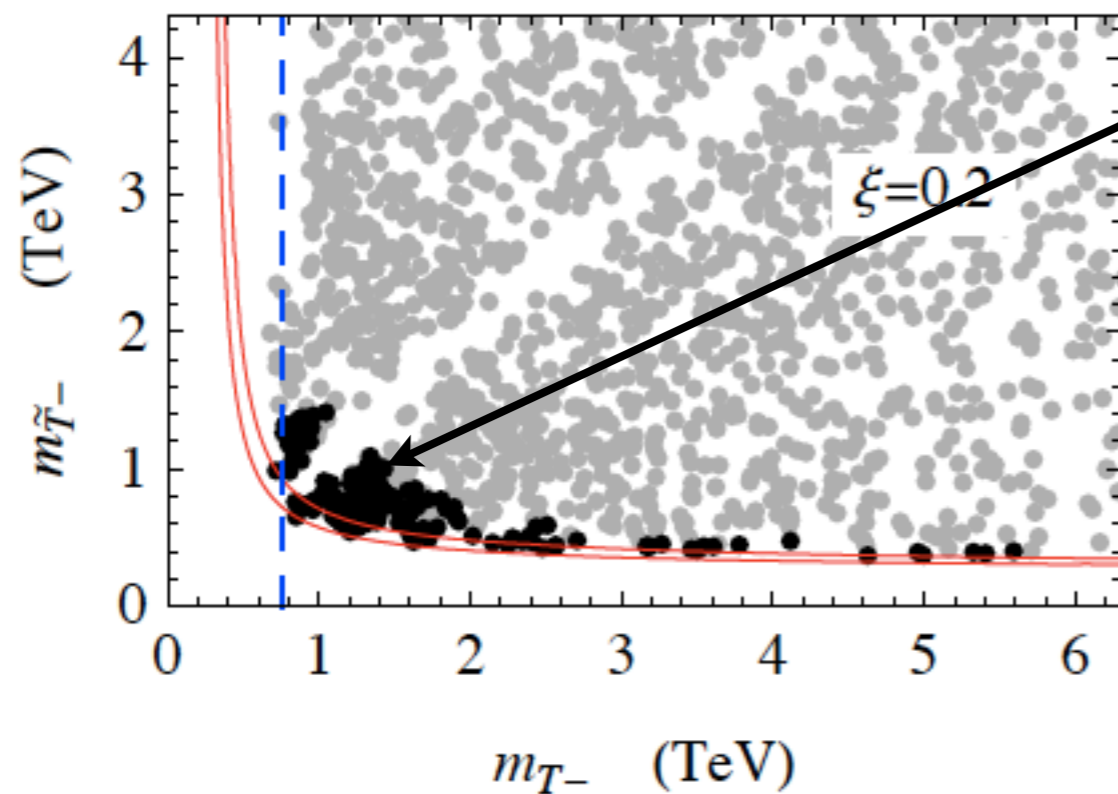
[D0 L=5.4 fb ⁻¹ arXiv:1010.1466	$Q\bar{q} \rightarrow Wq\bar{q} \rightarrow l^\pm jj\cancel{E}_T$ $\rightarrow Zq\bar{q} \rightarrow (l^+l^-)jj$
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Relatively light composite fermions preferred by currently "allowed" SM range for m_H

$$m_h \sim A \frac{m_T m_{\tilde{T}}}{\pi v} \xi, \quad A = O(1)$$

$$m_h \in (115, 130) \text{ GeV}$$



Matsedonskyi, Panico, Wulzer, 2012

T, \tilde{T} lightest fermionic partners of t_L, t_R

$\xi = \frac{v^2}{f^2}$ made small by fine tuning

Higgs boson branching ratios

$$\mathcal{L} = \mathcal{L}_{gauge}^{SM} + \frac{v^2}{4} \langle (D_\mu U)^\dagger (D_\mu U) \rangle \left(1 + a \frac{h}{v}\right) + \frac{v}{\sqrt{2}} \bar{Q}_{Li} U Q_{Ri} \left(1 + c \frac{h}{v}\right)$$

$$U(x) = e^{i\hat{\pi}(x)/v}, \quad \hat{\pi}(x) = \tau^a \pi^a$$

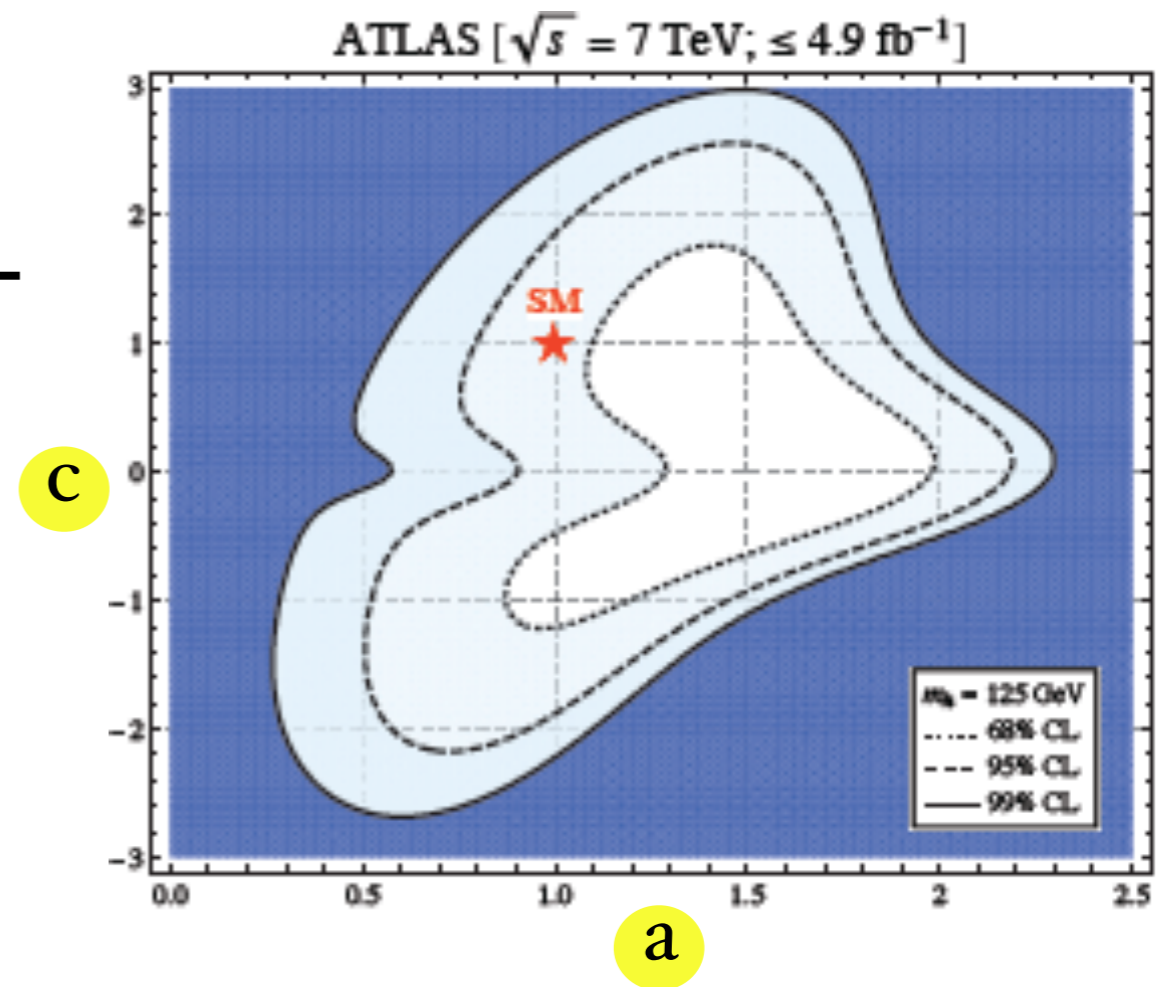
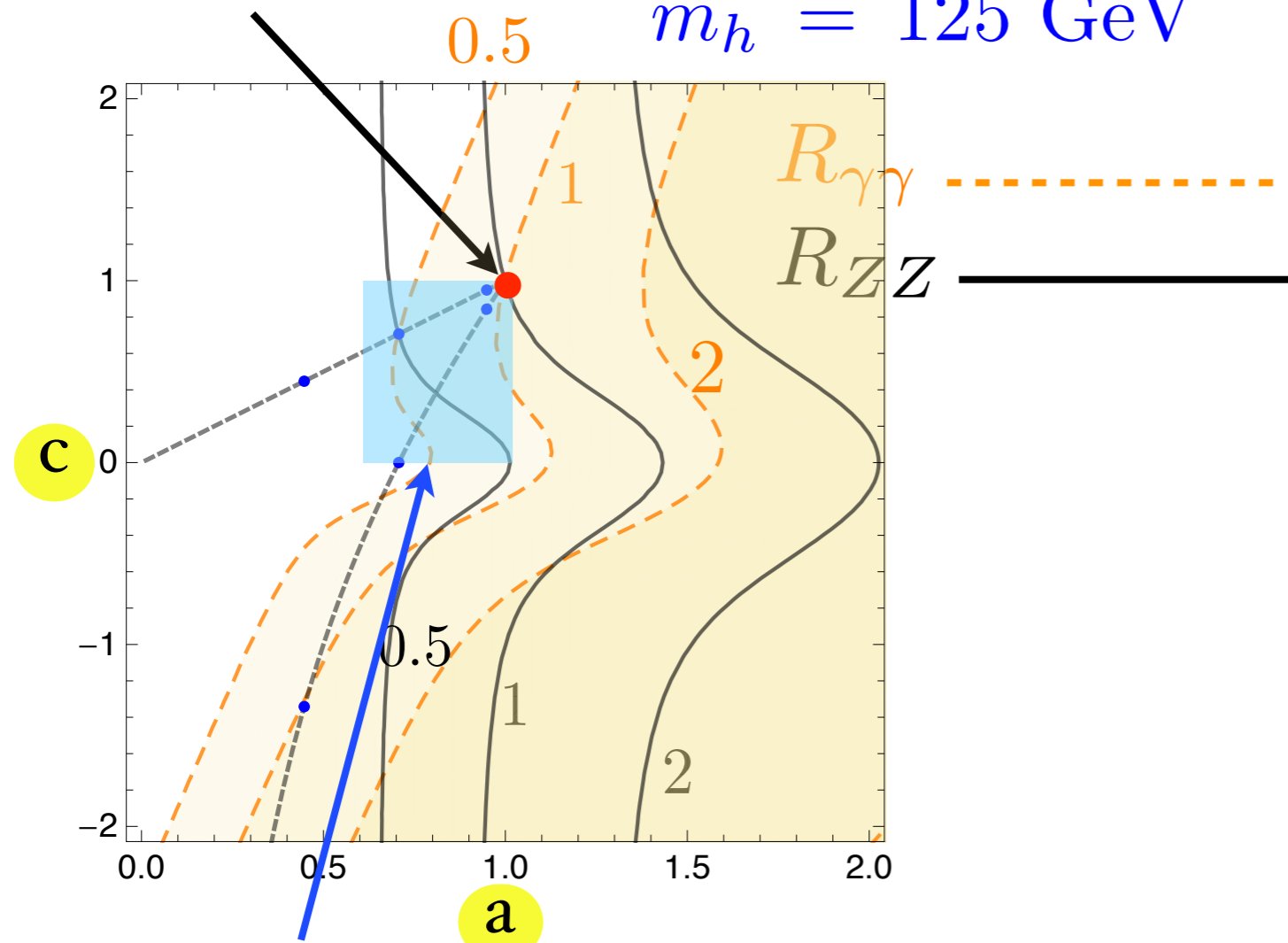
" π " = longitudinal W,Z

$$Q_{Ri} = \begin{pmatrix} \lambda_{ij}^u u_{Rj} \\ \lambda_{ij}^d d_{Rj} \end{pmatrix}$$

Standard Model: $a = c = 1$

$m_h = 125 \text{ GeV}$

Azatov, Contino, Galloway 2012



blu area preferred by EWPT

Summary

1. What underlies EWSB still unknown, as is the solution of the “little hierarchy” problem
2. To discover the Higgs boson and to know (some of) its properties has such far reaching implications that it pays to be patient a bit more to draw any conclusion
3. To discover (or to exclude natural) supersymmetry important to focus on $m_{\tilde{g}}, m_{\tilde{t}}, m_{\tilde{b}}, m_{\tilde{h}}$
4. Among “Exotica” the searches for the heavy fermions of the Higgs=PGB picture stand up