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A new resonance at 750 GeV? Who ordered that?
 Main features and first interpretations
 Example of scenario: hMSSM+VLFs
 Implications: looking at the bright side of life
 Summary

A. Angelescu, G. Moreau, AD: 2HDMs/MSSM+VLFs, arXiv:1512.04921 Y. Mambrini, G. Arcadi, AD: Dark Matter issues, arXiv:1512.04913 J. Ellis, R. Godbole, J. Quevillon, AD: Collider Signatures, arXiv:1601.03696

Rome, 22/01/2016 A new resonance at 750 GeV? – Abdelhak Djouadi – p.1/26

It all started with a rumor on Resonances: and the smiles in LAL-Orsay ATLAS group... I was not the only one to guess apparently: CERN auditorium not empty on December 15!





It had a smell of December 2011, the other Higgstorical day....

A new resonance at 750 GeV?

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CERN Jamboree: LHC results at $\sqrt{\mathrm{s}}=$ 13 TeV and L= 3.2 fb $^{-1}$ or 2.6 fb $^{-1}$

ATLAS di-photon results: 3.9 σ local excess at 750 GeV (but only 2.3 σ after LEE?). Total width of about 45 GeV (but smaller width possible).



CMS di-photon results:

2.6 σ local excess at 760 GeV (but only 1.2 σ after LEE?). Signal larger with 8TeV data. Total width apparently small. (and analysis targets spin-2).



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And?

Experimentalists:

Too early to say anything! It is only three poor sigmas!



and if you insist a little bit:



So do your job and collect data (and leave the theorists enjoy!)

Poor theorists:

Waiting for 30 years for NP, starting to get desperate... something interesting appears.



Do your job and interpret data! (healthy exercise anyway...)

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Tsunami of theory papers trying to interpret the 750 GeV diphotons:

10 papers the very first day, 100 at the end of the year, about 170 papers as of today.. Nature article/Dorigo blog:



Florilège of explanations:

- cascading heavy quarks,
- collimated 2x2 photons,
- new gauge bosons Z'+X
- sgoldstinos and other SUSY,
- quirks, hidden valleys?
- statistical fluctuation...

But most papers are thinking

about a new heavy resonance:

- Dark matter mediators
- Technipions/Goldstones, ...
- Axions, Radions/Dilatons,...
- Gravitons or any spin 2...
- Higgs bosons...

and other possibilities...

I try some quick/basic interpretations...

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σ B [pb]

If resonance: obviously integer spin: the observation is made in $\mathbf{X} \rightarrow \gamma \gamma$: the Landau–Yang theorem



(orbital momentum conservation):
rules our case of spin-one particle.
(ways to evade that but curious...).
Either spin-zero or spin-two.
Spin-2 has democratic couplings:
(as in the case of KK gravitons eg):

should also appear in $\ell\ell$, jj, VV, Vh no sign of that in other searches.

Spin-zero is more likely.

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Does it come from gg or $q\bar{q}$? Well, to cope with 8 TeV data, it should better come from gg:

 $\mathcal{R}_{i} = rac{(\sigma_{s}^{i}/\sqrt{\sigma_{B}})_{^{13 \text{ TeV}}}}{(\sigma_{s}^{i}/\sqrt{\sigma_{B}})_{^{8 \text{ TeV}}}}$ $\mathcal{R}_{i=gg} \simeq 3 \, v.s. \mathcal{R}_{i=q\bar{q}} \simeq 1.7$ gg: still tension with 8 TeV data...



Prefers gg via heavy particles to light quarks: likely to be Higgs–like! It is a scalar or pseudoscalar Higgs boson: we baptize it $\Phi =$ H or A.

 Φ production cross section? Fit all data and make a χ^2 : ATLAS at 13 TeV run only, CMS at both 8 and 13 TeV runs, $\Rightarrow \sigma(\Phi) = 6 \pm 2$ fb

pretty large cross section!

 $\sigma_{13 \text{ TeV}}(\text{pp} \rightarrow \Phi \rightarrow \gamma\gamma) \text{ [fb]}$

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The $\Phi {
m gg}$ and $\Phi \gamma \gamma$ couplings should be induced by heavy fermion loops: $\left| \Gamma(\Phi \to \mathbf{g}\mathbf{g}, \gamma \gamma) \propto \mathbf{g}_{\mathbf{s}, \mathbf{w}}^{\mathbf{2}} \right| \sum_{\mathbf{F}} \lambda_{\Phi \mathbf{F} \mathbf{F}} / \mathbf{m}_{\mathbf{F}} \times \mathbf{A}_{\mathbf{1/2}}^{\Phi}(\tau_{\mathbf{F}}) \right|$ $\mathbf{A_{1/2}^{S}} = \mathbf{2} \left[\tau_{\mathbf{F}} + (\tau_{\mathbf{F}} - \mathbf{1}) \mathbf{f}(\tau_{\mathbf{F}}) \right] \tau_{\mathbf{F}}^{-2}$ $au_{f F}={f M^2_{f \Phi}}/4{f m^2_{f F}}$ $A_{1/2}^{P} = 2\tau_{F}^{-1}f(\tau_{F})$ $\mathbf{f}(\tau_{\mathbf{F}}) = \arcsin^2(\tau_{\mathbf{F}}^{-1/2})$ for $\tau_{\mathbf{F}} \ge 1$ $\mathrm{m_F}\gg\mathrm{M_\Phi}\Rightarrow rac{\mathrm{A_{1/2}^P}
ightarrow+2}{\mathrm{A_{1/2}^S}
ightarrow+rac{4}{2}}$ A/P $\mathbf{A}_{1/2}^{\Phi}(\tau_{\mathbf{F}})$ 4.5 3.5For big loop contributions, 3 $\operatorname{Re}(A_{1/2}^{A/P})$ we need (simultaneously?): 2.5 $\operatorname{Im}(\operatorname{A_{1/2}^{A/P}})$ big Yukawas, 2 big charge/color, 1.5 $\operatorname{Re}(A_{1/2}^{H/S})$ $\operatorname{Im}(A_{1/2}^{H/S})$ 1 • $\mathrm{m_F} \approx rac{1}{2} \mathrm{M}_{\Phi}$, 0.5many fermions... 0 0.3 0.53 50.110 $au_{\mathrm{F}} = \mathrm{M}_{\Phi}^2/4\mathrm{m}_{\mathrm{F}}^2$

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Narrow width (as in CMS?): Φ couples only via loops, also to WW, ZZ, $Z\gamma$ In addition, one has $m_F \gtrsim \frac{1}{2}M_{\Phi}$ so that there are no decays $\Phi \rightarrow f\bar{f}, F\bar{F}$ Effective Lagrangian approach with the field strengths and their duals:

$$\mathcal{L}_{eff}^{\mathbf{S/P}} = \frac{e^2}{4v} \mathbf{c}_{\Phi\gamma\gamma} \Phi \mathbf{F}_{\mu\nu} \mathbf{F}^{\mu\nu} / \tilde{\mathbf{F}}^{\mu\nu} + \frac{\mathbf{g}_s^2}{4v} \mathbf{c}_{\Phi gg} \Phi \mathbf{G}_{\mu\nu} \mathbf{G}^{\mu\nu} / \tilde{\mathbf{G}}^{\mu\nu}$$

$$BR(\Phi \to \gamma\gamma) = \frac{\Gamma(\Phi \to \gamma\gamma)}{\Gamma(\Phi \to \gamma\gamma) + \Gamma(\Phi \to gg)} \approx \frac{\Gamma(\Phi \to \gamma\gamma)}{\Gamma(\Phi \to gg)} \approx \frac{\mathbf{c}_{\Phi\gamma\gamma}^2}{\mathbf{c}_{\Phigg}^2} \frac{\alpha}{\mathbf{s}\alpha_s} \approx 10^{-2}$$
Only vector-like fermion loops,

discuss several possibilities:

 $\begin{array}{l} \mbox{model 1: an } e_{\mathbf{Q}} = \frac{2}{3} \; T_{\mathbf{R},\mathbf{L}} \; \mbox{singlet.} \\ \mbox{model 2: } e_{\mathbf{Q}} = \frac{2}{3}, -\frac{1}{3} \; (\mathbf{U},\mathbf{D})_{\mathbf{R},\mathbf{L}}, \\ \mbox{model 3: } \; (\mathbf{U},\mathbf{D})_{\mathbf{R},\mathbf{L}}, \; T_{\mathbf{R},\mathbf{L}}, \; B_{\mathbf{R},\mathbf{L}}, \\ \mbox{model 3: } \; (\mathbf{U},\mathbf{D})_{\mathbf{R},\mathbf{L}}, \; T_{\mathbf{R},\mathbf{L}}, \; B_{\mathbf{R},\mathbf{L}}, \\ \mbox{model 4: } \; (\mathbf{U},\mathbf{D})_{\mathbf{R},\mathbf{L}}, \; T_{\mathbf{R},\mathbf{L}}, \; B_{\mathbf{R},\mathbf{L}}, \\ \; (\mathbf{L}^1,\mathbf{L}^2)_{\mathbf{R},\mathbf{L}}, \; E_{\mathbf{R},\mathbf{L}} \end{array}$

LHC Φ xsection reproduced for perturbative $\lambda^2/4\pi < 1/2$ and not too large VLF masses...



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Large width scenario (as in ATLAS): Φ couples directly to heavy particles:

- the couplings to W and bosons: are all eaten by the SM-like h state, – only fermion couplings allowed: either tops, bottoms, or new ones... Again in the effective approach: $\mathcal{L}_1 = \mathcal{L}_{\mathbf{S}} + \mathbf{c_f} \mathbf{m_F} / \mathbf{v} imes \mathbf{\Phi} \mathbf{f} \mathbf{f}$ $\mathcal{L}_1 = \mathcal{L}_{\mathbf{P}} + \mathrm{i} \mathbf{c_f} \mathbf{m_F} / \mathbf{v} \mathbf{\Phi} \mathbf{\overline{f}} \gamma_5 \mathbf{f}$ with the SM vev $\mathbf{v} \approx$ 246 GeV; can fit $\sigma imes {
m BR}$ and $\Gamma_{oldsymbol{\Phi}}\!pprox\!50$ GeV for reasonable $c_{\Phi {\bf g} {\bf g}}, c_{\Phi \gamma \gamma}$ and $m_{F}.$



The best way to describe the large width possibility is the 2HDM/MSSM An example would be the hMSSM: AD, Maiani, Polosa, Quevillon, Riquer arXiv:1307.5205 [hep-ph] and arXiv:1502.05653 [hep-ph].

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In the MSSM: two Higgs doublets: $H_1 = \begin{pmatrix} H_1^0 \\ H_1^- \end{pmatrix}$ and $H_2 = \begin{pmatrix} H_2^+ \\ H_2^0 \end{pmatrix}$, After EWSB (which can be made radiative: more elegant than in the SM): Three dof to make W_L^\pm , $Z_L \Rightarrow$ 5 physical states left out: h, H, A, H^\pm Only two free parameters at tree-level: $tan\beta$, M_A but rad. cor. important: $M_h {\lesssim} M_Z |cos2\beta| + RC {\lesssim} 130 \; GeV \;, \; M_H {\approx} M_A {\approx} M_{H^\pm} {\lesssim} M_{EWSB}$ – Couplings of h, H to VV are suppressed; no AVV couplings (CP). – For $tan\beta \gg 1$: couplings to b (t) quarks enhanced (suppressed).



In the decoupling limit: MSSM reduces to SM but with a light SM Higgs. This decoupling limit occurs in 2HDM extension: alignment limit... a SM-like light and two CP-odd like heavy Higgses with cplg to t,b, τ only $\Rightarrow h \equiv H_{SM}$, $\Phi = H, A$

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 $\begin{array}{l} \label{eq:MSSM Higgs sector simple at tree level: only two basic inputs, } \tan\beta, M_A \\ \mbox{Radiative corrections make it complicated as } RC = f(M_S, X_t, X_b, \mu, M_i) \\ \mbox{ex: } M_h^2 \stackrel{M_A \gg M_Z}{\to} M_Z^2 |cos^2 2\beta| + \frac{3\bar{m}_t^4}{2\pi^2 v^2 \sin^2\beta} \left[\log \frac{M_S^2}{\bar{m}_t^2} + \frac{X_t^2}{M_S^2} \left(1 - \frac{X_t^2}{12\,M_S^2} \right) \right] \\ \mbox{Only information so far on the MSSM from the LHC } \Rightarrow \begin{cases} M_h = 125 \; {\rm GeV} \\ M_S \gtrsim 1 \; {\rm TeV} \end{cases} \end{cases}$

hMSSM: trade the value $M_h \!=\! 125~GeV$ against the radiative corrections. Back to tree-level: only two inputs for Higgs sector and no SUSY parameters:

$$\begin{split} M_{H}^{2} = \frac{(M_{A}^{2} + M_{Z}^{2} - M_{h}^{2})(M_{Z}^{2}c_{\beta}^{2} + M_{A}^{2}s_{\beta}^{2}) - M_{A}^{2}M_{Z}^{2}c_{2\beta}^{2}}{M_{Z}^{2}c_{\beta}^{2} + M_{A}^{2}s_{\beta}^{2} - M_{h}^{2}}\\ \alpha = -\arctan\left(\frac{(M_{Z}^{2} + M_{A}^{2})c_{\beta}s_{\beta}}{M_{Z}^{2}c_{\beta}^{2} + M_{A}^{2}s_{\beta}^{2} - M_{h}^{2}}\right) \end{split}$$

One also has the relation $M^2_{H^\pm}\simeq M^2_A+M^2_W$ which is more general...

Effective and "model-independent" approach of the MSSM Higgs sector:

- good: very simple, economical, .. and opens the possibility of low taneta
- bad: requires large $\mathbf{M_S}$ at low taneta (but apparently this is the case...?)
- ugly: needs large fine-tuning (but theory already fined-tuned anyway..).



Large width scenario (as in ATLAS): obtained from Φ -fermion couplings – couplings to massive gauge bosons all eaten by the SM–like 125 GeV h, – only couplings to fermions allowed: either tops, bottoms, or new ones...

 $\mathbf{g}_{\Phi tt} = \frac{\mathbf{m}_t}{\mathbf{v}} \cot \beta, \ \mathbf{g}_{\Phi bb} = \frac{\mathbf{m}_b}{\mathbf{v}} \tan \beta, \ \mathbf{g}_{\Phi \tau \tau} = \frac{\mathbf{m}_{\tau}}{\mathbf{v}} \tan \beta$

with $\tan\beta = v_2/v_1$ small $\tan\beta \approx 1$ or large $\tan\beta \approx m_t/m_b \approx 60$ $-\tan\beta \approx 1 : BR(\Phi \to t\bar{t}) \approx 1, BR(\gamma\gamma) \approx 10^{-5}, \ \Gamma_{\Phi} \approx 30$ GeV. $-\tan\beta \approx 60 : BR(\Phi \to b\bar{b}) \approx .9, BR(\gamma\gamma) \approx 10^{-7}, \ \Gamma_{\Phi} \approx 30$ GeV.



– $an\!eta$ pprox3–10: allow for light lepton (DM?) decays to get $\Gamma_{\Phi}pprox 30$ GeV.

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Large values $\tan\beta \gtrsim 30$: $\sigma(gg, b\bar{b} \rightarrow \Phi \rightarrow \tau \tau)$ too large ATLAS+CMS very sensitive \Rightarrow region totally excluded.



Low values $\tan\beta \lesssim 1$: $\sigma(gg \rightarrow \Phi \rightarrow t\bar{t})$ too large ATLAS+CMS searches sensitive \Rightarrow region being excluded.



NB: analysis valid for spin–1 (no interference with gg \rightarrow tt bkg) full Φ analysis in progress

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Unfortunately hMSSM or 2HDM without any new particle does not make it!

Rates for $gg \rightarrow \Phi \rightarrow \gamma\gamma$: $\sigma(H) = 0.85$ fb at 8 TeV $BR(H \rightarrow \gamma\gamma) \approx 6 \times 10^{-6}$ $\sigma(A) = 1.70$ fb at 8 TeV $BR(A \rightarrow \gamma\gamma) \approx 7 \times 10^{-6}$ $\sigma \times BR(H + A) \approx 10^{-2}$ fb We are short by a factor 500...

Include a bunch of VLFs:

- 3 families of 2 VLL doublets
- 3 doubly charged leptons
 one family of VLQ and VLL

(we set tan β =3 to reduce Γ_{Φ}) with usual Yukawa couplings optimal effect at $m_F = \frac{1}{2}M_{\Phi}$ (But watch out for light Higgs).



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VLFs will also contribute to SM–like Higgs gg and $\gamma\gamma$ loops!



Figure 1: Contours of constant $\sum \sigma(gg \to \Phi) \times \text{BR}(\Phi \to \gamma\gamma)$ and $\mu_{\gamma\gamma}$ in the $\{y_L^u, y_R^d\}$ plane, for MSSM (left) and type II 2HDM (right) including the $\mu_{\gamma\gamma} = 1.16 \pm 0.18 \pm 0.15$ constraint.

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4. Implications: Dark Matter

Φ resonance is ideal mediator for Dark Matter: case of fermion X — cosmological relic density Ωh^2 obtained by annihilation XXo Φ oSM.



Good prospects for direct/indirect detection in astrophysical experiments.



4. Implications: singlet resonance at colliders

Reproduce Φ resonance in pp: same prod. process $gg \to \Phi \gamma \gamma$ grows with the gluon luminosity



Ideal for HE-LHC, FCC-hh, SPPC 2 orders magnitude more at 100TeV check other WW,ZZ,Z γ final states.

Future e+e- HE linear colliders \neg can be turned into $\gamma\gamma$ colliders 80% energy and same luminosity $\Rightarrow \Phi$ production in $\gamma\gamma \rightarrow \Phi$



Ideal machine for a diphoton state: Measure precisely $\Phi\gamma\gamma$ coupling Check CP properties of resonance.

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4. Implications: doublet resonance at colliders

Many more processes if Φ is in a 2HDM/hMSSM like scenario; in pp:



4. Implications: doublet resonance at colliders

Many more processes if Φ is in a 2HDM/hMSSM like scenario; in e+e-:-



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4. Implications: vector–like fermions

The vector-like fermions can be produced in pair or singly at colliders:



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4. Implications: doublet resonance at colliders

First pair production of VLQs in pp and then single production via mixing:



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4. Implications: vector–like fermions

First pair production of VLQs in pp and then single production via mixing:

	Vector-like quark mass sensitivity				Vector-like lepton mass sensitivity			
model	$100 {\rm fb}^{-1}$	$300 \mathrm{fb}^{-1}$	$300 {\rm fb}^{-1}$	$20ab^{-1}$	$100 {\rm fb}^{-1}$	$300 {\rm fb}^{-1}$	$300 {\rm fb}^{-1}$	$20ab^{-1}$
	$13 { m TeV}$	$14 { m TeV}$	$33 { m TeV}$	$100 { m TeV}$	$13 { m TeV}$	$14 { m TeV}$	$33 { m TeV}$	$100 { m TeV}$
1	1.4	1.7	3.1	11.7		-		
2	1.5	1.8	3.4	12.7		-		
3	1.6	2.0	3.7	13.7		-		
4	1.6	2.0	3.7	13.7	0.56	0.73	1.7	5.3

Table 1: Prospective model sensitivities to massive vector-like quarks (left) and leptons (right) [with the particle masses in TeV] in the indicated pp collider and scenario from extrapolations of the present LHC searches.

4. Implications: doublet resonance at colliders

Pair production of VLLs in e+e- and then single production via mixing:



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5. Summary

And? Too early to conclude. But life suddenly became bright...

It is really a new resonance? Or is it simply a mirage?



If true then the future is bright! (bye-bye the multiverse ... and plenty of new physics!) But again we should hear the experimentalists and their usual :



and wait for the coming data. In summer we will now more (but until then we can enjoy!)

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